

# Loop Quantum Cosmology with vector fields

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# Outline

- ▶ Motivation

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- ▶ GR with the vector field domination

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- ▶ Introduction to Loop Quantum Cosmology

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- ▶ Loop Quantum Cosmology and anisotropic Universe
- ▶ Summary

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- ▶ Large scale structure from vectors
- ▶ Embedding of vector fields in Loop Quantum Cosmology

# The Ashtekar formalism

Let us assume the perfect fluid with the vector field. Then

$$A_\mu = A(t)\delta^z_\mu, \quad ds^2 = dt^2 - a^2(dx^2 + dy^2) - b^2dz^2. \quad (1)$$

We construct the scalar constraint (Hamiltonian)

$$\mathcal{C} = -\frac{1}{\gamma^2 p_1 \sqrt{p_2}} (c_1^2 p_1^2 + 2c_1 c_2 p_1 p_2) + p_1 \sqrt{p_2} \rho, \quad (2)$$

where Ashtekar variables are defined by

$$p_1 = L_1 L_2 ab, \quad p_2 = L_1^2 a^2, \quad c_1 = \gamma L_1 \dot{a}, \quad c_2 = \gamma L_2 \dot{b}. \quad (3)$$

## Physical interpretation

$a^2$ ,  $ab$  - areas in anisotropic Universe

$c_1$ ,  $c_2$  - external curvature along  $x$ ,  $y$  and  $z$  axes

# Cosmology with vector fields

From Hamilton equations

$$\dot{p}_1 = \frac{\gamma}{2}\{p_1, \mathcal{C}\}, \dot{p}_2 = \gamma\{p_2, \mathcal{C}\}, \quad \dot{c}_1 = \frac{\gamma}{2}\{c_1, \mathcal{C}\}, \dot{c}_2 = \gamma\{c_2, \mathcal{C}\}, \quad (4)$$

one obtains Einstein's equations

$$H(H + 2\mathcal{H}) = \rho = \rho_f + \rho_A, \quad (5)$$

$$\frac{\ddot{a}}{a} + \frac{\ddot{b}}{b} + H\mathcal{H} = -p_{\perp} = -p_f - p_A, \quad (6)$$

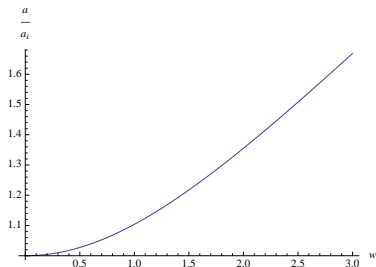
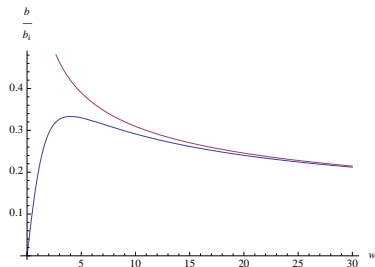
$$2\frac{\ddot{a}}{a} + H^2 = -p_{\parallel} = -p_f + p_A, \quad (7)$$

where  $H = \dot{a}/a$ ,  $\mathcal{H} = \dot{b}/b$  and

$$\rho_A = (\dot{A}^2 + m^2 A^2)/b^2, \quad p_A = (\dot{A}^2 - m^2 A^2)/b^2$$

# Massless vector field domination

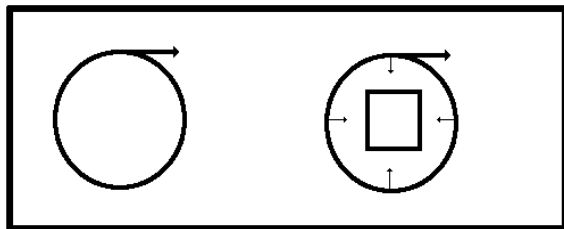
For  $m = 0$  one obtains  $b = E\dot{a}$ , where  $E = \text{const} > 0$ . The Big Bang has finite energy density  $\rho_I$





# Loop Quantum Cosmology in FRW

In isotropic case  $a = b$ , so  $p_1 = p_2$  and  $c_1 = c_2$



The loop correction appears due to  $\Delta$  - the minimal value of the area operator. This changes  $c$  to

$$c \rightarrow \sin(\bar{\mu}c), \quad (8)$$

where  $\bar{\mu} = \sqrt{\Delta/p}$

Hamiltonian of matter fields and  $p$  remains unchanged

# LQC in anisotropic Universe

In diagonal Bianchi I model LQC changes  $c_i$  to

$$c_i \rightarrow \sin(\bar{\mu}_i c_i), \quad (9)$$

where  $\bar{\mu}_i$  are some generalisations of  $\bar{\mu}$ . Two ways of doing that:

## The old $\bar{\mu}_i$ scheme

Szulc '08:  $\bar{\mu}_i = \sqrt{\Delta/p_i} \Rightarrow$  problems with scaling

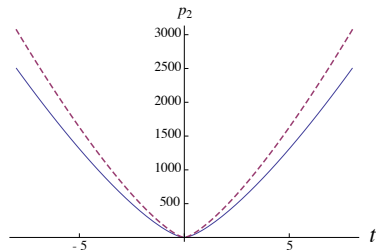
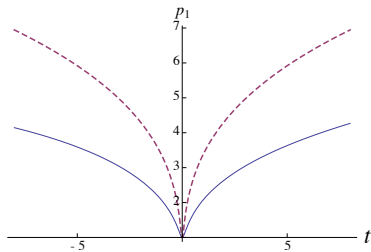
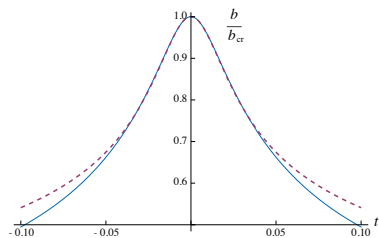
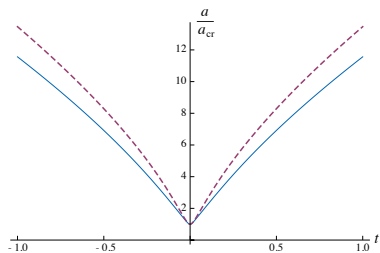
## The new $\bar{\mu}_i$ scheme

Ashtekar, Wilson-Ewing '09:  $\bar{\mu}_1 = \sqrt{\Delta/p_2}$ ,  $\bar{\mu}_2 = \sqrt{\Delta p_2}/p_1$

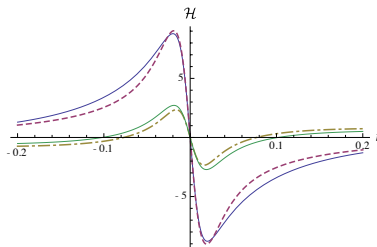
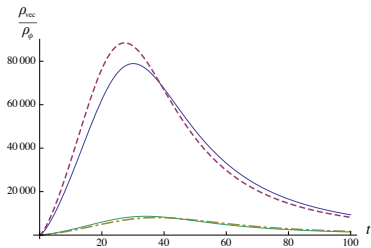
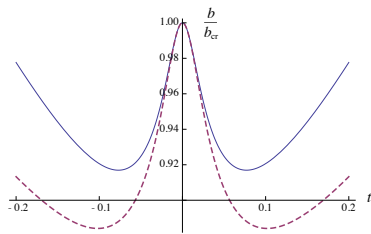
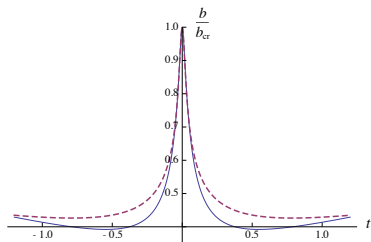
In both schemes  $p_i$  and  $\mathcal{C}_{mat}$  are not changed by LQC

And now... The most important thing

# LQC with massless vector field



# LQC with massive vector and scalar



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- ▶ Quasi Big Bounce solution for massless vectors
- ▶ Isotropic low energy limit for massive vectors
- ▶ In LQC with vector fields  $b(t)$  shrink for a while, but...
- ▶ everything is fine since areas always grow