

Global interpretation of LHC indications within the Georgi-Machacek Higgs model

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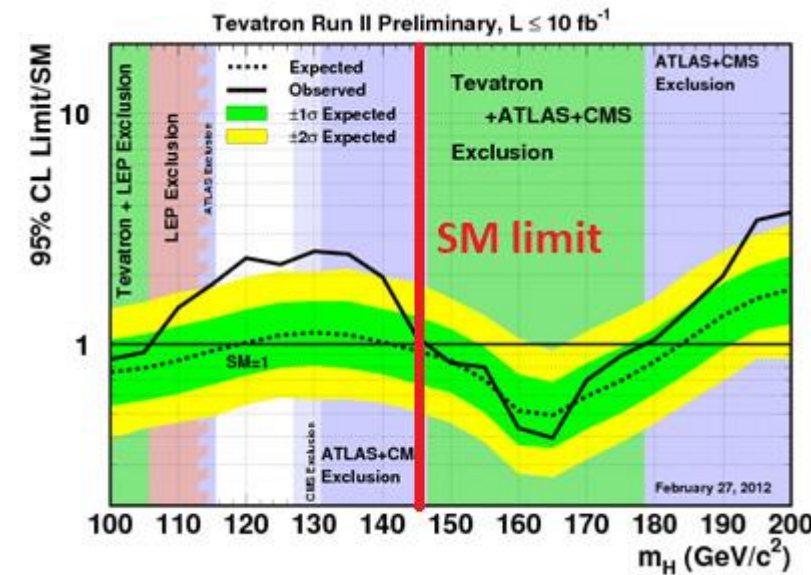
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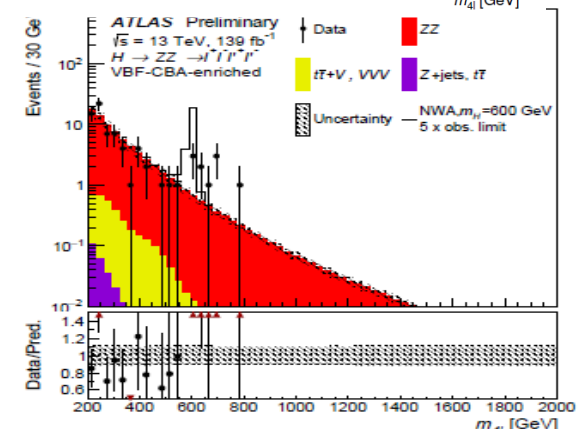
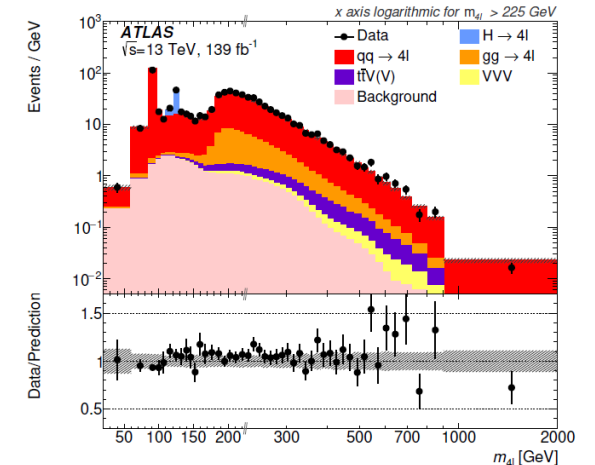
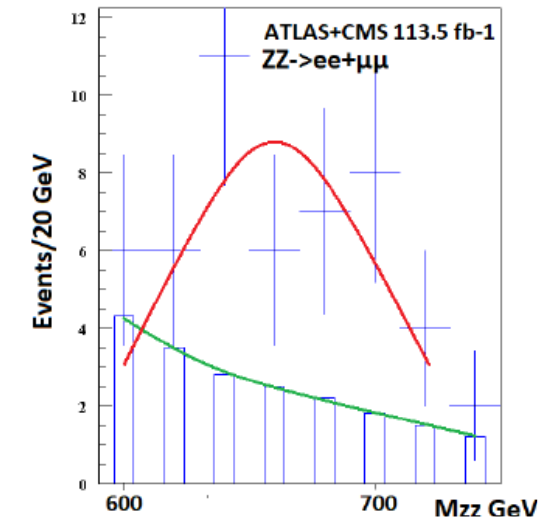
Introduction



- The SM Higgs discovery has relied on the following ingredients:
 - A well tested “**effective theory**”, the SM, that predicted the mass and the couplings of the Higgs particle
 - A > 5 s.d. evidence for 2 signals **coincidentally** observed by ATLAS and CMS for this particle (Tevatron started with a ~ 3 s.d. indication + $M_h < 145$ GeV from Mw)
- The following question therefore arises: will LHC meet a similar scenario ?
- Here I will restrict to searches for scalars resulting from **broken symmetries** which, in a large class of BSM models, are the lightest objects in similitude with **pions of QCD**, the **SM h** being one of them
- The dominant phenomenological BSM framework assumes that these scalars belong to doublets and singlets but this can also comprise triplets in an extension proposed in 1985 by **Georgi and Machacek** which is consistent with the **parameter ρ** being very close to 1
- I will show that while not yet reaching 5 s.d. – although not so far from it – there are several **indications** close to that level, which do not seem to match the two doublet models but can be accommodated within reasonable extensions of **GM**
- Some of these indications are not “far fetched” and belong to the gold plated channels of LHC which led to the Higgs discovery

1st indication : H->ZZ into 4 leptons

- The cleanest channel for discoveries
- From a combination of published histograms done in [1806.04529](#) with 113.5 fb⁻¹ from CMS (2/3) and ATLAS (1/3) one observes a peak at **~650 GeV** with s/b=42/14 ~4.3 s.d. local significance
- With 139 fb-1 ATLAS sees a much smaller effect but at the same mass [2103.01918](#)
- The VBF channel from ATLAS provides a better s/b, but with about 50% less efficiency
- With 139 fb-1, with **sequential cuts**, an excess is observed at the same mass, s/b=9/2 ~3 s.d., for **VBF->H(650)->ZZ** [ATLAS-CONF-2020-032](#)
- **Machine learning search** does not confirm (see T. Lagouri talk)
- The CMS analyses, inclusive+VBF, are not yet published
- These results call for a combination of both analyses before one can draw a valid conclusion
- Could stop here but...



Evidence for $VBF \rightarrow H(650) \rightarrow W+W- \rightarrow e\mu\nu\nu$

CMS PAS HIG-20-016

- 3 samples were analysed by CMS reaching 138 fb⁻¹
- Has a large top background even after b-jet vetoing
- The best evidence reaches 3.8 s.d. is for **$VBF \rightarrow H(650) \rightarrow e\mu\nu\nu$**
- This result does not suffer from the “look elsewhere” downgrading
- The VBF cross section ~ 160 fb
- Awaiting for the same analysis from ATLAS before reaching any quantitative interpretation of this resonance
- See [2205.12234](#) for an alternate interpretation of these signals

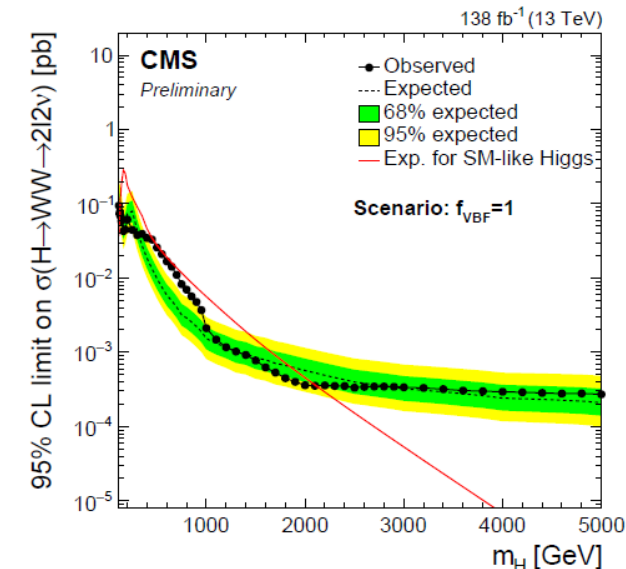
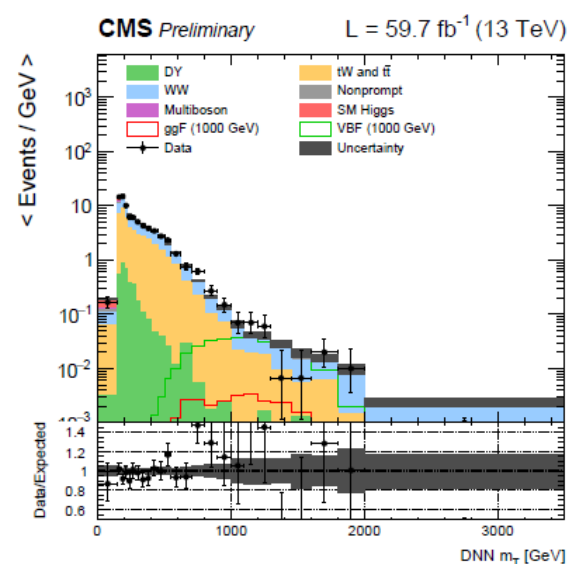
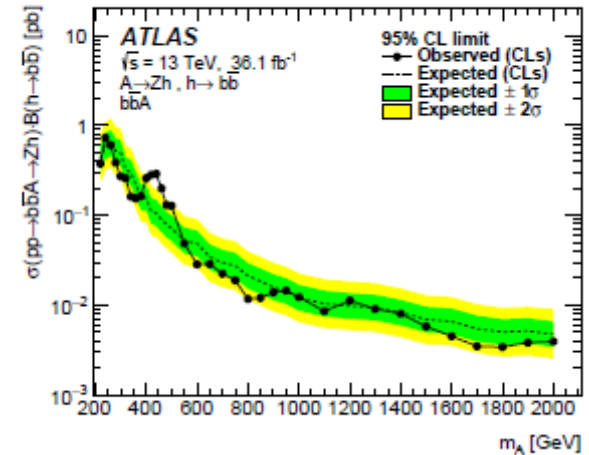
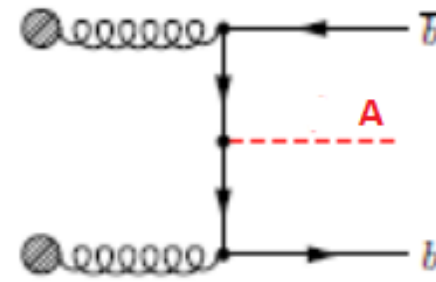


Table 3: Summary of the signal hypotheses with highest local significance for each f_{VBF} scenario. For each signal hypothesis the resonance mass, production cross sections, and the local and global significances are given.

Scenario	Mass [GeV]	ggF cross sec. [pb]	VBF cross sec. [pb]	Local signi. [σ]	Global signi. [σ]
SM f_{VBF}	800	0.16	0.057	3.2	1.7 ± 0.2
$f_{VBF} = 1$	650	0.0	0.16	3.8	2.6 ± 0.2
$f_{VBF} = 0$	950	0.19	0.0	2.6	0.4 ± 0.6
floating f_{VBF}	650	2.9×10^{-6}	0.16	3.8	2.4 ± 0.2

2d indication : a resonance at 400 GeV seen in various modes

- CMS sees 3.5 s.d. for $t\bar{t}$ at ~ 400 GeV $\Gamma/M=4\%$
- CMS took into account **interference** with the QCD background (major issue at LHC)
- Signals were observed by **ATLAS** in $\tau\tau$, $\tau\tau+b$ and in $hZ+b$ (CP-odd) at ~ 400 GeV
- $\tau\tau$, $\tau\tau+b$ not confirmed by **CMS-PAS-HIG-21-001**
- ATLAS signals at ~ 3 s.d. local, but giving **~ 6 s.d. global** when combined
- Caveat: $hZ+b$ not completed by ATLAS with full stat
- Does not fit with **MSSM** which predicts decoupling for $A \rightarrow hZ$ [1802.09122](https://arxiv.org/abs/1802.09122)
- **-> Try an other model !**



- <https://arxiv.org/abs/2001.04770>
- <https://arxiv.org/abs/2003.07112>

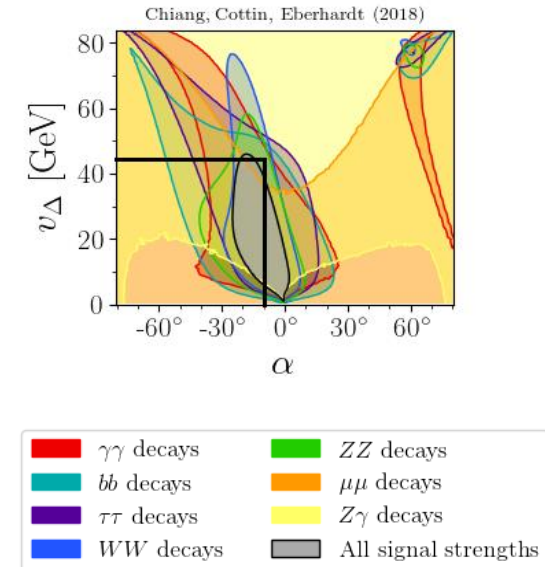
Reaction	Mass GeV	Nb of s.d.	Ref	ArXiv
$X(400) \rightarrow t\bar{t}$	400	3.5		1908.01115
$X(400) \rightarrow \tau\tau$	400	2.2		2002.12223
$X(400) \rightarrow \tau\tau+b$	400	2.7		2002.12223
$A(400) \rightarrow h(125)Z+b$	440	3.6		1712.06518
$X(400) + \text{high pt } e/\mu$	400	3		2002.11325

Giorgi-Machacek for pedestrians

- Allows $I=3/2$, H^{++} , without violating $\rho = M^2 w / Mz^2 \cos^2 \theta w = 1$ at tree level
- Is achieved by combining 1 isospin doublet (v_ϕ) + 2 triplets with the same vacuum expectations :

$$\rho = \frac{\tilde{v}_\phi^2 + 4\tilde{v}_\chi^2 + 4\tilde{v}_\xi^2}{\tilde{v}_\phi^2 + 8\tilde{v}_\chi^2} = \frac{v^2}{v^2 + 4(\tilde{v}_\chi^2 - \tilde{v}_\xi^2)} = 1 \text{ with } v_\chi = v_\xi$$

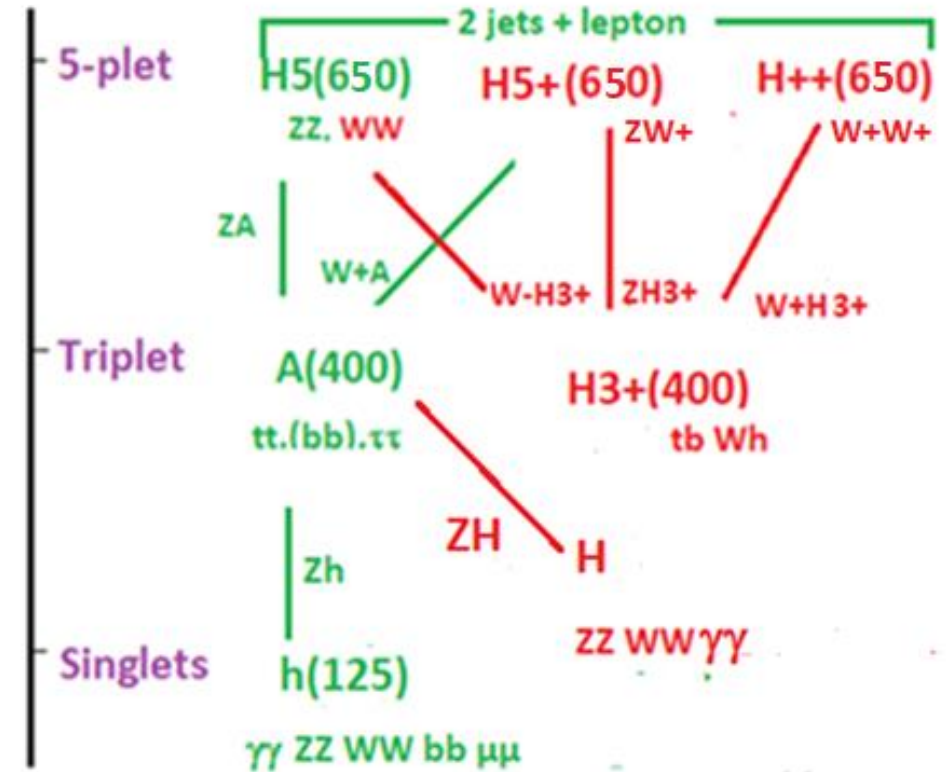
- Predicts a **Quintet** $H5^{++}$ $H5^+$ $H50$ $H5^-$ $H5^{--}$ **Fermiophobic**, only produced by **VBF** for which $H(650)$ could be a candidate
- + **Triplet Gaugephobic** $H3^+$ $H30$ (CP-odd) $H3^-$ -> **A(400)**
- Mass degeneracy inside multiplets usually assumed but unnecessary for $\rho=1$ [2111.14195](https://arxiv.org/abs/2111.14195)
- + **Singlets** **h(125)** and **H1** **mixing angle** α
- **Unitarity constraints** impose $m_H < 700$ GeV and fixing m_5 and m_3 , that $m_{H1} < 250$ GeV
- Allows $A(400) \rightarrow hZ$ but $A(400) \rightarrow H1Z \sim 25$ times larger if $m_{H1} \sim m_h$
- Couplings depend on 2 mixing angles constrained by LHC observations
- Tentative choice: $\sin \alpha \sim -0.15$ and $\sin \theta_H \sim 0.5$ ($v_\chi = 43$ GeV) to agree with PM



[1807.10660](https://arxiv.org/abs/1807.10660)

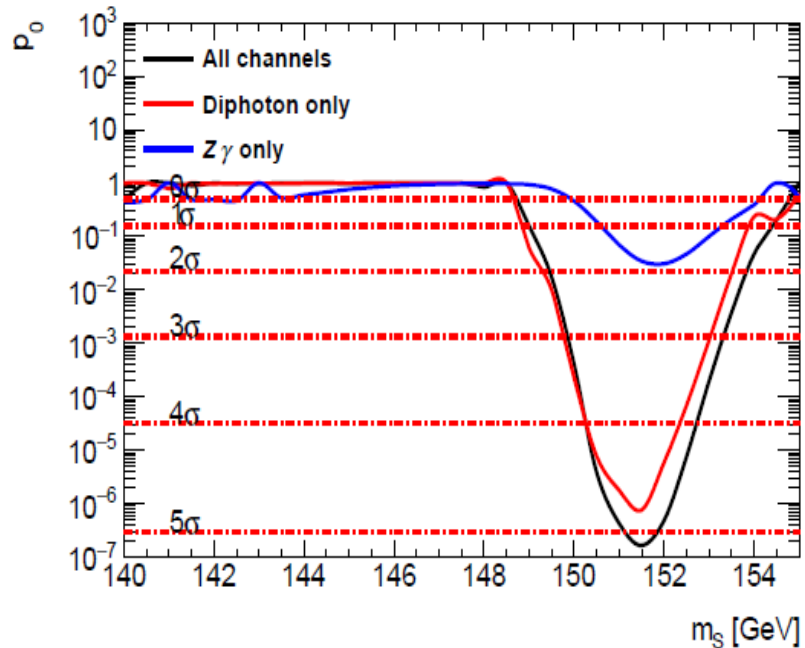
3d indication: cascades

- The GM picture suggests the presence of many **cascades** which allow various strategies of signature: like sign leptons, leptons + b jets, ttW etc...
- Cascades could explain the various topological anomalies recorded by **Budenbrock et al** [1901.05300](#) in ATLAS and CMS with an overall significance ~ 8 s.d.
- Similar analysis performed by CMS in [2202.08676](#)
- **2 additional scalars H(270) S(150)** were predicted by Budenbrock et al which tend to cascade into each other $H \rightarrow SS^*$ and $H \rightarrow Sh(125)$

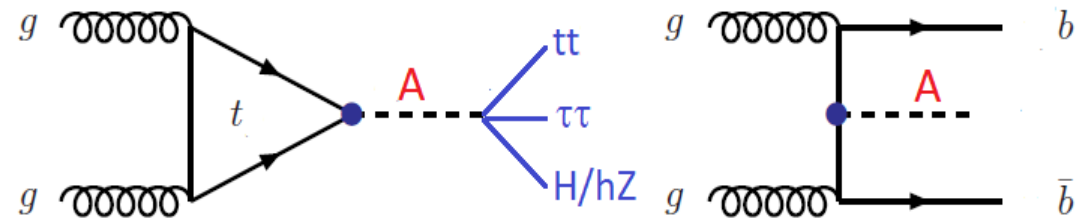


4th indication H(151)->2γ + E_{miss} + ...

- Following [1901.05300](#), signal observed by Crivellin et al. [2109.02650](#) asking ETMiss or presence of additional leptons and b jets (global **4.8 s.d.**)



- GM interpretation:



- 1st mechanism has $\sigma \sim 6000$ fb
- ETMiss comes from $Z \rightarrow \nu\nu$ **without invoking other particles**
- The second mechanism allows to understand Zh(125)+bjets observed by ATLAS

ZWW: a paradox ?

- GM predicts for $BR(A \rightarrow HZ) = 54\%$
- For H, in MeV: $\Gamma_{\text{tot}_{SM}} = 5$ $\Gamma(WW) = 4.5$ $\Gamma(ZZ) = 0.5$ $\Gamma(\gamma\gamma) = 0.025$
- $\Gamma_{BSM} = ?$ **Visible and invisible**
- Predicts $\delta\sigma(ZWW) \sim 3000$ fb if $\Gamma_{BSM} = 0$
- Excluded by CMS limit $\delta\sigma(ZWW) < 400$ fb [2006.11191](#)
- Would require $\Gamma_{BSM} \sim 40$ MeV to lower $BR(WW)$
- $BR_{\text{inv}} < 30\%$ deduced from h(125) favours $\Gamma_{BSM_{\text{vis}}}$
- If $H \rightarrow BSM$ visible, it is likely that this happens also for h(125)

H(151)->aa ?

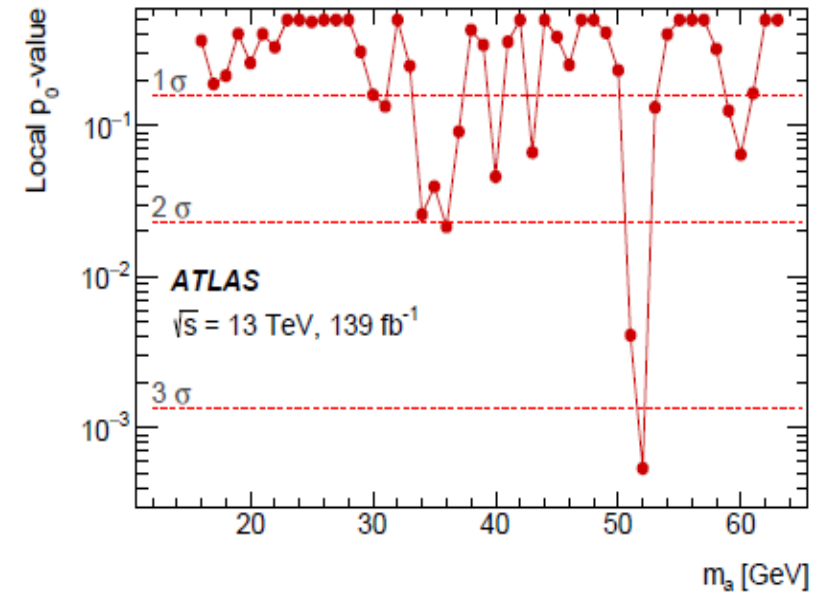
[2110.00313](#)

- ATLAS observes an evidence **3.3 σ local 1.7 σ global** for **h(125)->aa** scalar or pseudo-scalar $m_a=52$ GeV
- Such an object is well motivated (DM, EW phase transition, g-2...)
- $BR(h \rightarrow \mu\mu bb) \sim 2 \cdot 10^{-4}$ $\mu\mu bb/bbbb \sim 1/1300$ hence $\Gamma(h \rightarrow aa) \sim 1$ MeV
- Could be much larger for H in the GM model

$$h = \cos\alpha H_1 - \sin\alpha H_1' \quad H = \sin\alpha H_1 + \cos\alpha H_1'$$

assuming $a(52)$ couples to the H_1' component

$$\Gamma(H \rightarrow aa) = \Gamma(h \rightarrow aa) / \sin^2\alpha = 44 \text{ MeV !}$$



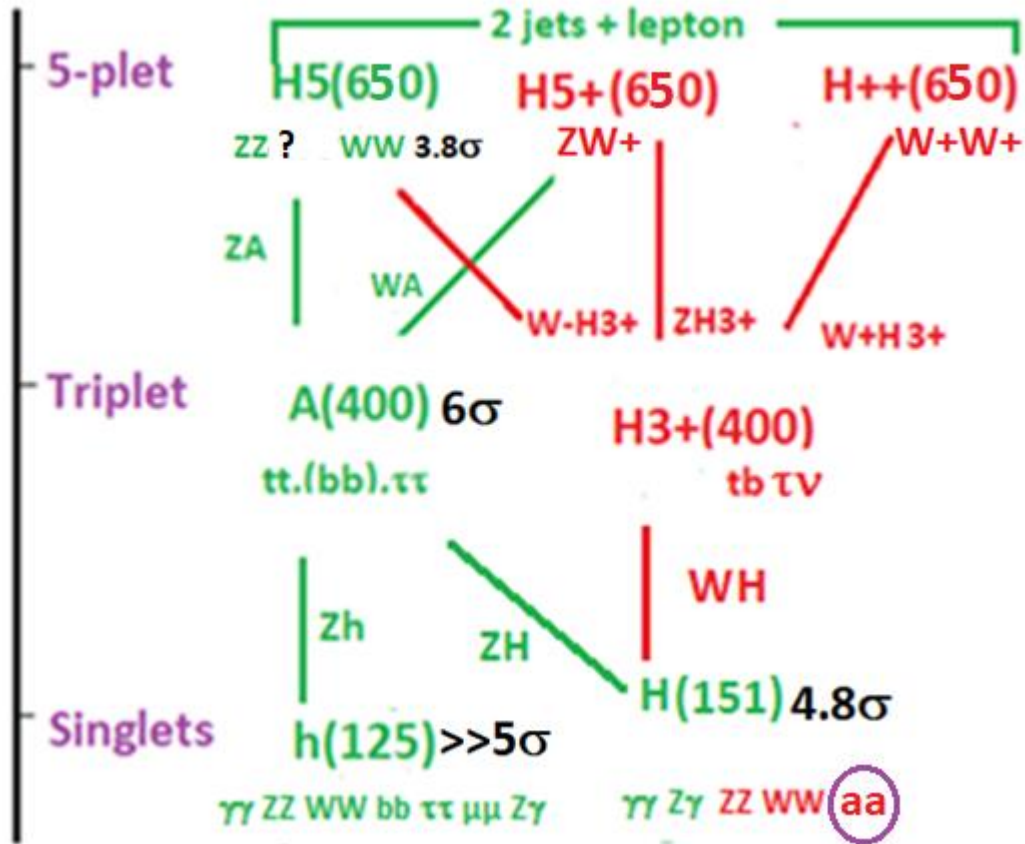
Final evaluation

Γ_{Htot} MeV	Γ_{WW} MeV	Γ_{ZZ} MeV	$\Gamma_{\gamma\gamma}$ MeV	Γ_{aa} MeV	Γ_{inv} MeV
5-50	4.5	0.5	0.025	45	<2-20

Γ_{Atot} GeV	Γ_{tt} GeV	Γ_{hZ} GeV	Γ_{HZ} GeV	Γ_A BSM GeV
16	6.4	0.4	8	?

- Previous results allow to fix the widths of A & H
- CMS indicates $\Gamma_{Atot} \sim \mathbf{16 GeV}$, leaving few MeV for BSM
- **H(151)** decays are dominantly BSM, not necessarily aa but visible, from the LHC measurements of WW
- $BR(Z \rightarrow \nu\nu) \sigma(A \rightarrow HZ) BR(H \rightarrow 2\gamma) = 0.2 * 3000 (0.025 / \Gamma_{Htot})$
 0.42 ± 0.13 fb in Crivellin et al. suggesting $\Gamma_{Htot} \sim 40$ MeV
- $\sigma(4\ell) = \sigma(A \rightarrow HZ) BR(H \rightarrow ZZ) BR(ZZ \rightarrow 4\ell) = 3000 (0.5 / \Gamma_{Htot}) 0.07^2 = 0.18$ fb
compatible with < 0.28 fb in Crivellin et al.

The GM+a picture

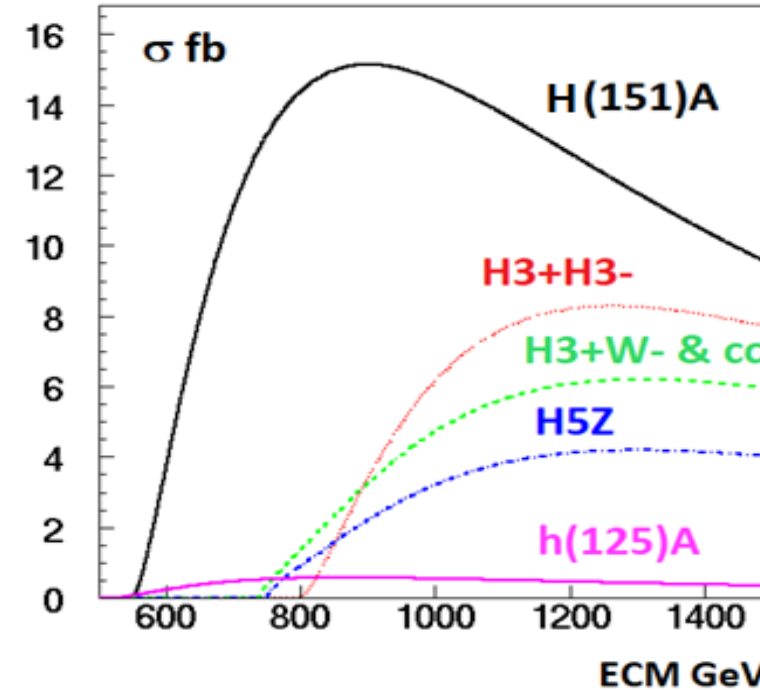
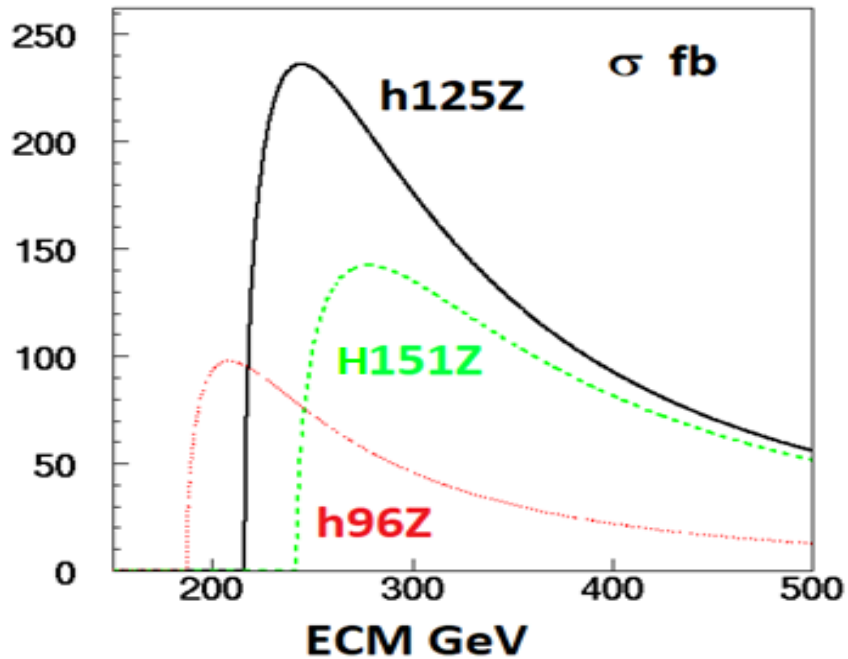


- GM predicts the mass and couplings of $H(151)$ with a BSM final state aa
- $A \rightarrow H(151)Z$ with $H(151) \rightarrow aa$ and $a \rightarrow bb$ with $\sigma \sim 5000 \text{ fb}$ offers a major new topology for discoveries
- This topology is **similar to $h(125)Z$ with $h \rightarrow bb$** , which could explain the apparent large cross section for “ hZ ” seen by ATLAS which could be a fake

GM cross sections in e^+e^-

- At 1 TeV

- At 250 GeV

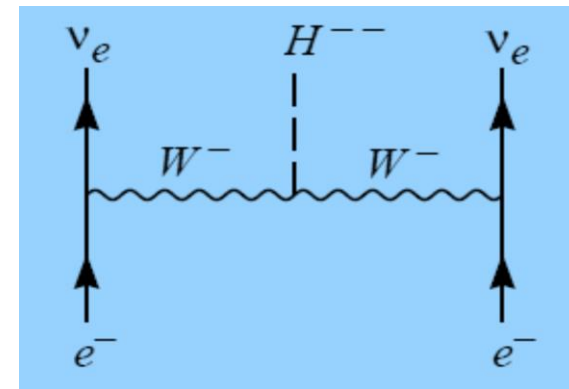
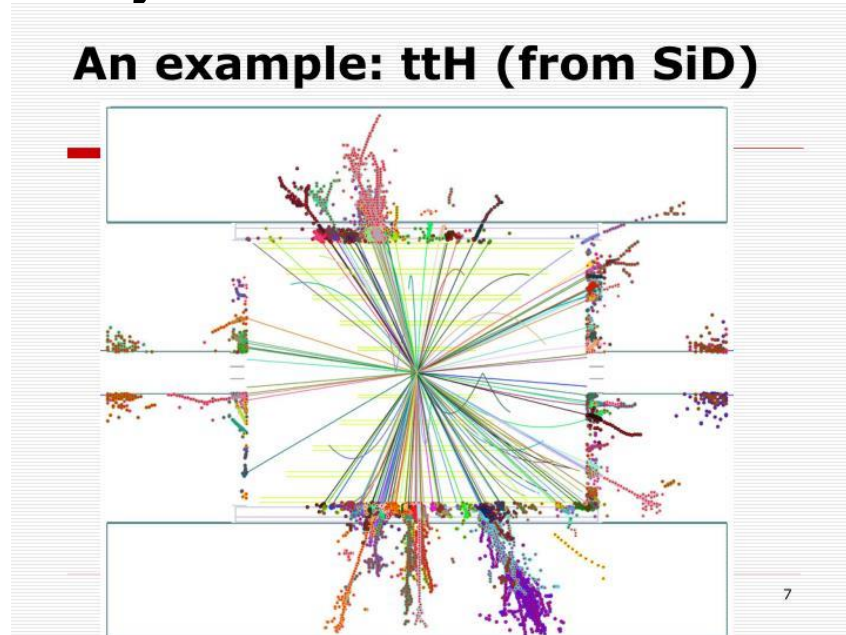


- Large xsections allowing very precise measurements

- Assumes mass degeneracy inside multiplets
- $H5Z$ is only a guess and needs confirmation
- Complex modes which requires highest \mathcal{L} and reconstruction efficiency

Prospects for a LC

- $ECM=1\text{ TeV}$ is sufficient to observe the full GM scalar spectrum
- Requires highest possible **luminosity**, $\sim 8000\text{ fb}^{-1}$ with ILC at 1 TeV [1903.01629](#)
- $\Omega > 0.9$ is needed for jet reconstruction and **b tagging** for these complex modes
 $\epsilon_{\text{eff}} \sim \Omega^n$ $n \sim 8$



- $>1.5\text{ TeV}$ to pair produce $H^{++} H^{--}$ in e^+e^-
 but 1 TeV enough in **VBF** $e^-e^- \rightarrow W^-W^- \nu\nu \rightarrow H^{--} \nu\nu$ \sim few fb for single production

Conclusions

- Evidence is building up at LHC for the presence of a rich scalar spectroscopy in the framework of the **GM model** which predict scalars lighter than 700 GeV therefore **fully testable at LHC**
- An indication coming from Crivellin et al. suggests that **H(151)** could have a **BSM dominant decay**, a scenario comforted by ATLAS observing $h(125) \rightarrow aa$, with $m_a = 52$ GeV
- There are a few concerns, e.g. about H(650) into WW and ZZ which would benefit from an official **combination of ATLAS and CMS analyses**
- These results potentially open good prospects for HL-LHC discoveries/confirmations, with the following message: **do not assume that BSM physics necessarily means masses > 1 TeV**
- All neutral scalars of GM have been observed while the absence of **H3+, H5+ and H5++** can still be understood (see Appendix)
- GM offers an entirely new landscape for e+e- colliders under discussion and motivates a **linear e+e- collider reaching no less than 1 TeV**
- Complex final states will have a **critical impact on the design of future LC detectors**

BBC: Large hadron collider: A revamp that could revolutionise physics



P. Richard CERN May 2022

APPENDIX

Missing slides (lack of time)

GM constraints on MH and BR(H(151)->2γ)

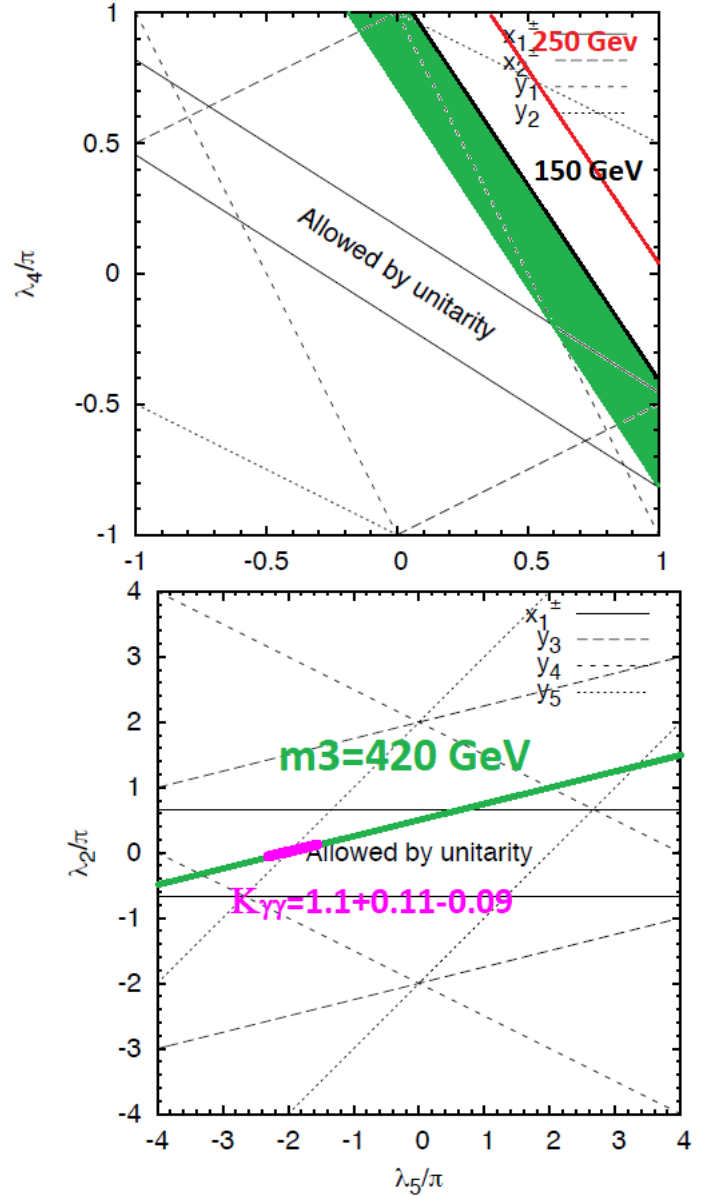
- The GM model gives strict **unitary limits** on the scalar masses and on the coefficient of the scalar potential:

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ -\phi^{+*} & \phi^0 \end{pmatrix},$$

$$X = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ -\chi^{+*} & \xi^0 & \chi^+ \\ \chi^{+++} & -\xi^{+*} & \chi^0 \end{pmatrix}$$

$$V(\Phi, X) = \frac{\mu_2^2}{2} \text{Tr}(\Phi^\dagger \Phi) + \frac{\mu_3^2}{2} \text{Tr}(X^\dagger X) + \lambda_1 [\text{Tr}(\Phi^\dagger \Phi)]^2 + \lambda_2 \text{Tr}(\Phi^\dagger \Phi) \text{Tr}(X^\dagger X) + \lambda_3 \text{Tr}(X^\dagger X X^\dagger X) + \lambda_4 [\text{Tr}(X^\dagger X)]^2 - \lambda_5 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^b) \text{Tr}(X^\dagger t^a X t^b) - M_1 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^b) (UXU^\dagger)_{ab} - M_2 \text{Tr}(X^\dagger t^a X t^b) (UXU^\dagger)_{ab}.$$

- A priori **m < 700 GeV for all scalars** [1404.2640](#)
- Knowing m3 and m5 one finds **mH < 250 MeV** for the missing scalar
- From LHC observations and **h(125)->2γ** one can extract all parameters and compute **Γ(H(151)->2γ) = 0.025 MeV** [1708.08753](#)



Higgs self coupling predictions

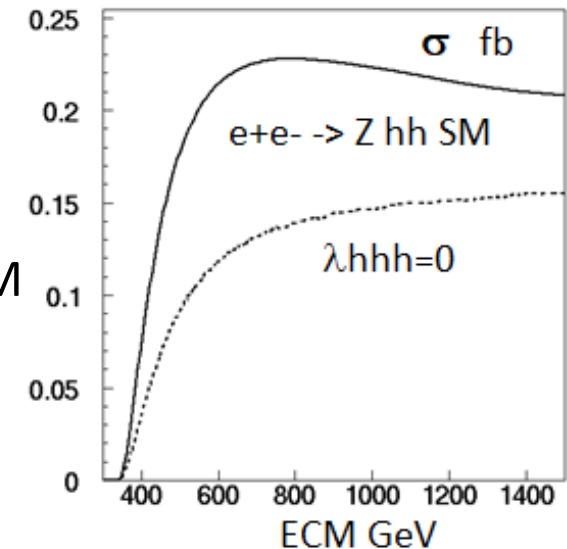
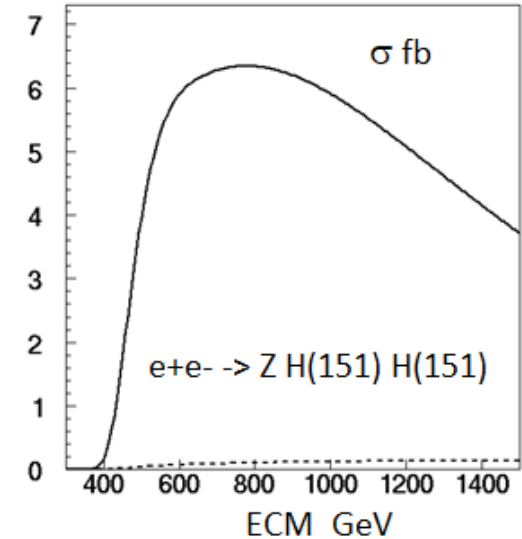
m3 GeV	m5 GeV	s_α	sH	λ_1/π	λ_2/π	λ_5/π	λ_3/π	λ_4/π	M1 GeV	M2 GeV
420	660	-0.15	0.5	0.02	0.01	-1.7+0.2-0.6	0.6	-0.2	1000	1000

- Fixing the GM parameters of the scalar potential from LHC observables, one can deduce the self couplings

$$\lambda_{HHH}/SM \sim 22 \quad \lambda_{hHH}/SM \sim 3.36 \quad \lambda_{hhH}/SM \sim 1.6 \quad \lambda_{hhh}/SM \sim 0.14$$

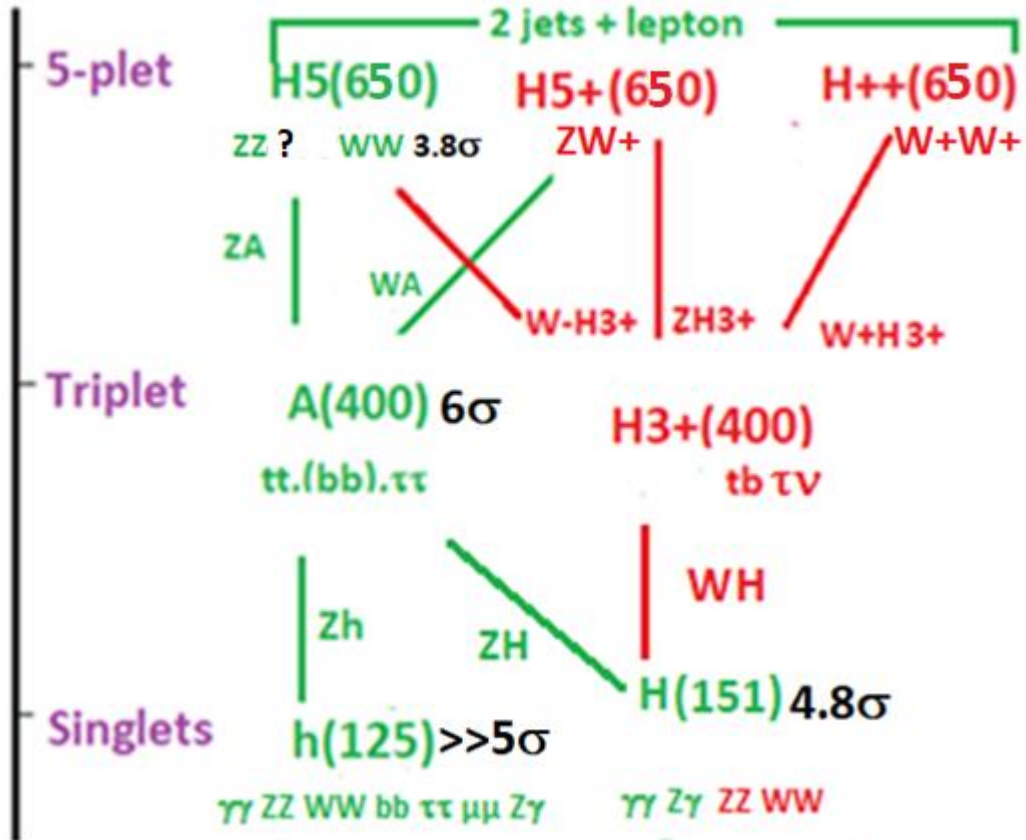
- $e^+e^- \rightarrow ZHH$ reaches ~ 4 fb at $\sqrt{s} \sim 600$ GeV ~ 27 times Zhh SM
- $e^+e^- \rightarrow Zhh$ almost unchanged /SM from the contribution of $H \rightarrow hh$
- Recalling that H should decay in $\sim 90\%$ of the cases into BSM modes, presumably into a pair of CP-odd light scalars $H \rightarrow aa$
- At LHC one expects that final states with HH originate either from a Higgsstrahlung, VBF or from the cascade $A \rightarrow H^* Z \rightarrow HHZ$
- The Z spectator can be used as a **tagger**
- These measurements are sensitive to the two parameters M1 and M2 of the GM potential recalling that

$$g_{HHH} = 24\lambda_1 s_\alpha^3 v_\phi + 6s_\alpha c_\alpha (\sqrt{3}s_\alpha v_\chi + c_\alpha v_\phi) (2\lambda_2 - \lambda_5) + 8\sqrt{3}c_\alpha^3 v_\chi (\lambda_3 + 3\lambda_4) - \frac{3\sqrt{3}}{2} M_1 s_\alpha^2 c_\alpha - 4\sqrt{3} M_2 c_\alpha^3,$$



The new GM+a picture

- GM predicts correctly the mass and couplings of H(151)
- All neutrals scalars are indicated
- Charged Higgs absent but **H3+** not searched into **WH(151)** which can dominate (also aW)



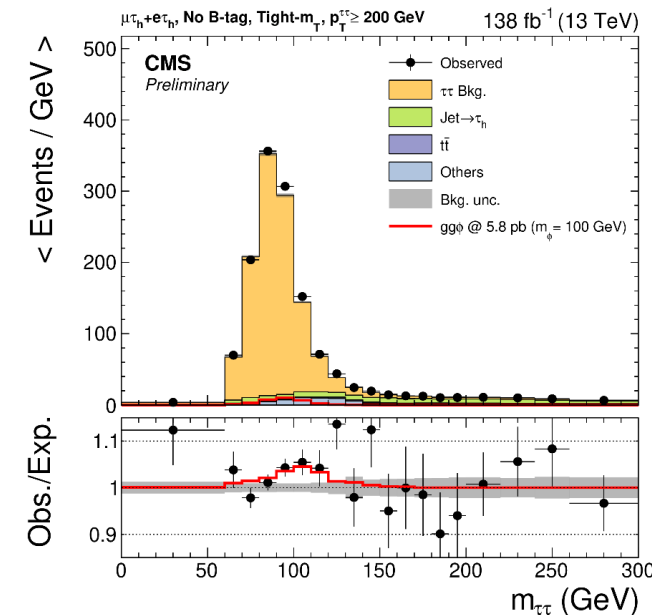
Type	coupling /SM, MSSM	$s\alpha=-0.15$ $sH=0.5$
h(125)WW/ZZ	$c\alpha cH-1.63s\alpha sH$	0.98
H(151)WW/ZZ	$s\alpha cH+1.63c\alpha sH$	0.68
h(125)tt,bb	$c\alpha/cH$	1.14
H(151)tt,bb	$s\alpha/cH$	0.17
Att,bb, $\tau\tau$	$\tan H$	0.58
H5WW, H5ZZ	$0.57sH,-1.15sH$	$0.27,-0.58$
H5AZ,H5H3+W-	$1.16cH$	1
H5+H3+Z,H5+AW+	cH	0.87
h(125)AZ,hH3+W-	$1.63(s\alpha cH+0.6c\alpha sH)$	0.28
H(151)AZ,HH3+W-	$1.63(c\alpha cH-0.6s\alpha sH)$	1.48
H5+W-Z,H5++W+W+	$-2sH,2.48sH$	$1.0,1.24$
H3+H3-Z	1	1

What about h(96) ?

(see also S. Heinemeir at this WS)

- $h(96) \rightarrow 2\gamma$ by CMS, not by ATLAS which however does not reach the same sensitivity
- Could it be the singlet H ?
- One would expect $A \rightarrow H(96)Z \sim 30 * A \rightarrow h(125)Z$, which does not seem to be the the case
- Could be an additional isosinglet embedded in an **EGM** or be the **Radion** expected within Randall Sundrum phenomenology [1712.06410](#)
- A 3.1 s.d. indication of $h(100) \rightarrow \tau\tau$ is seen by CMS after subtraction (see [2205.03187](#) for a CP-odd interpretation)
- Ideal for an e^+e^- collider $h(96) \rightarrow bb \tau\tau gg$

CMS-PAS-HIG-21-001



GM model issues

The GM model for advanced

- GM is constituted by one doublet ϕ and two triplets, one complex χ and one real ξ , with the same vacuum expectations to get $\rho=1$
- H1 and H1' have following composition

$$H_1^0 = \phi^{0,r},$$

$$H_1^{0r} = \sqrt{\frac{1}{3}}\xi^0 + \sqrt{\frac{2}{3}}\chi^{0,r}.$$

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}, \quad \chi = \begin{pmatrix} \chi^{++} \\ \chi^+ \\ \chi^{0*} \end{pmatrix}, \quad \xi = \begin{pmatrix} \xi^+ \\ \xi^0 \\ \xi^- \end{pmatrix}$$

$$Y=1/2 \quad T=1/2 \quad v\phi \quad Y=1 \quad T=1 \quad v\chi \quad Y=0 \quad T=1 \quad v\xi$$

$$\rho = \frac{\tilde{v}_\phi^2 + 4\tilde{v}_\chi^2 + 4\tilde{v}_\xi^2}{\tilde{v}_\phi^2 + 8\tilde{v}_\chi^2} = \frac{v^2}{v^2 + 4(\tilde{v}_\chi^2 - \tilde{v}_\xi^2)}.$$

- Only ϕ couples to fermions
- They form the following physical states, dominantly triplet

$$H_5^{++} = \chi^{++},$$

$$H_5^+ = \frac{(\chi^+ - \xi^+)}{\sqrt{2}},$$

$$H_5^0 = \sqrt{\frac{2}{3}}\xi^0 - \sqrt{\frac{1}{3}}\chi^{0,r},$$

$$H_3^+ = -s_H \phi^+ + c_H \frac{(\chi^+ + \xi^+)}{\sqrt{2}},$$

$$H_3^0 = -s_H \phi^{0,i} + c_H \chi^{0,i}.$$

- The physical states are

$$h = \cos \alpha H_1^0 - \sin \alpha H_1^{0r},$$

$$H = \sin \alpha H_1^0 + \cos \alpha H_1^{0r}.$$

- The mixing angle α has to be small to avoid altering the doublet properties of the SM h(125)
- E.g. $\sin \alpha = -0.15$ & $s_H = 0.5$, $v\phi = 213$ GeV for the doublet, $v\xi = v\chi = 43.5$ GeV for the triplets

A GM interpretation of the Mw excess?

- Recall that

$$\rho = \frac{\tilde{v}_\phi^2 + 4\tilde{v}_\chi^2 + 4\tilde{v}_\xi^2}{\tilde{v}_\phi^2 + 8\tilde{v}_\chi^2} = \frac{v^2}{v^2 + 4(\tilde{v}_\chi^2 - \tilde{v}_\xi^2)}$$

- From the CDF recent measurement of M_W which implies a deviation from 1 of this parameter one may wonder if GM could provide the simplest explanation ?
- See 3 GM references: [2204.12898](#) [2204.07844](#) [2204.05760](#)
- In [2204.12898](#) it is shown that this is indeed possible provided the scalar potential is modify to avoid generating unwanted effects (a bit technical, sorry)

Top-down approach to GM

- The GM model does not pronounce on the **origin of symmetry breaking** and is therefore **incomplete**
- In a **UV complete scenarios**, one assumes a symmetry breaking mechanism from extended symmetries to a group containing the SM $SU(2) \times U(1)$, accompanied by **pseudo Nambu-Goldstone bosons** in addition to the SM Higgs doublet which can result in GM
- A recent reference [1703.06064](#) shows how this mechanism can naturally generate **Higgs triplets**
- Some of these theories are **strongly interacting** and therefore constraints are extracted by lattice calculations
- In any case the GM model can be easily tested since it predicts scalars with masses below 700 GeV which can be ruled out at LHC

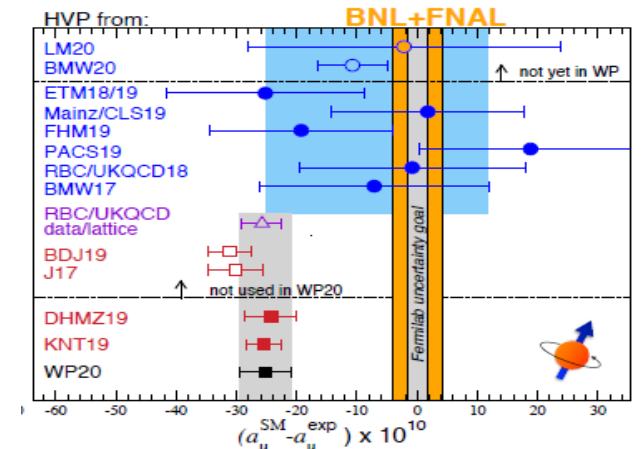
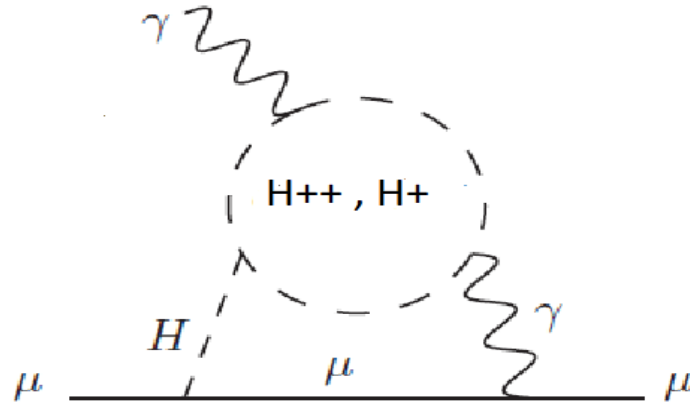
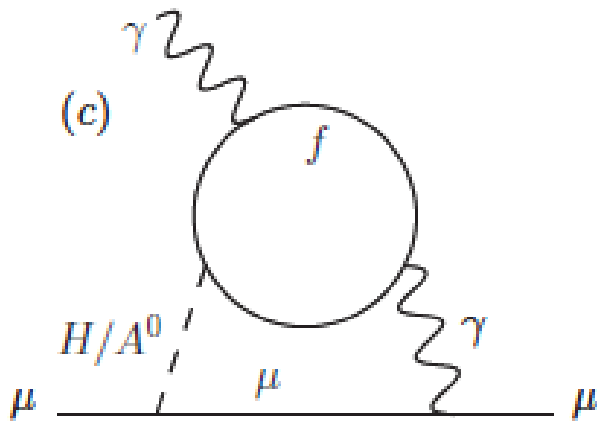
GM and LHC precision measurements on h(125)

- $\sigma(\text{gg} \rightarrow \text{h})$ **+28%** from htt in GM
- Γ_{htot} **+25%** due to $\text{h} \rightarrow \text{aa}$ and **+16%** from Γ_{bb}
- LHC measures $\sigma \Gamma_{\text{X}} / \Gamma_{\text{tot}}$ hence \sim **compensation** $\text{ggF} / \Gamma_{\text{tot}}$
apparent $\Gamma_{\gamma\gamma}$ $\Gamma_{\text{ww/zz}}$ -12% from true values
- Γ_{bb} comes from $\text{WWF} \rightarrow \text{h}$ and $\text{h} + \text{Z/W}$ with \sim compensation on
 Γ_{bb} -12% from true value
- LHC does not reveal GM due to these compensating effects
except for **VBF \rightarrow h \rightarrow WW*** where a **-40%** effect could become
visible with better accuracies
- [ATLAS-CONF-2021-014](#) $\mu_{\text{VBF}} = 1.0 \pm 0.20$
[CMS PAS HIG-20-013](#) $\mu_{\text{VBF}} = 0.70 \pm 0.26$
- Separation and high accuracy will be available in e+e-

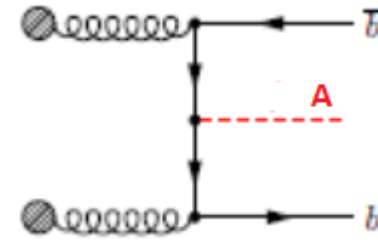
Type	coupling/SM	$s\alpha = -0.15$ $s_H = 0.5$
h(125)WW/ZZ	$c\alpha c_H - 1.63s\alpha s_H$	0.99
H(151)WW/ZZ	$s\alpha c_H + 1.63c\alpha s_H$	0.68
Att,bb, $\tau\tau$	$\tan H$	0.58
h(125)tt,bb	$c\alpha / c_H$	1.14
H(151)tt,bb	$s\alpha / c_H$	0.17
Zh(125)A	$1.63(s\alpha c_H + 0.6c\alpha s_H)$	0.28
ZH(151)A	$1.63(c\alpha c_H - 0.6s\alpha s_H)$	1.48
Wh(125)H3 ⁺	$1.63(s\alpha c_H + 0.6c\alpha s_H)$	0.28
WH(151)H3 ⁺	$1.63(c\alpha c_H - 0.6s\alpha s_H)$	1.48

What about $g-2$ in eGM+a(52) ?

- Usual wisdom is that extra scalars need to be **<100 GeV** to influence $g-2$
- The two loop contributions (**Barr Zee**) dominate and gives the right sign
- [2104.03275](#) with 2 doublets+1 triplet predicts a significant ($\geq 10^{-9}$) contribution to a_μ provided there is an enhanced $\mu\mu$ coupling ($\sim \times 100$)
- A quantitative prediction for the eGM+a model with 2 doublets + 2 triplets + a(52) seems promising

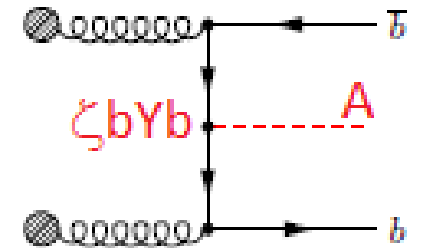
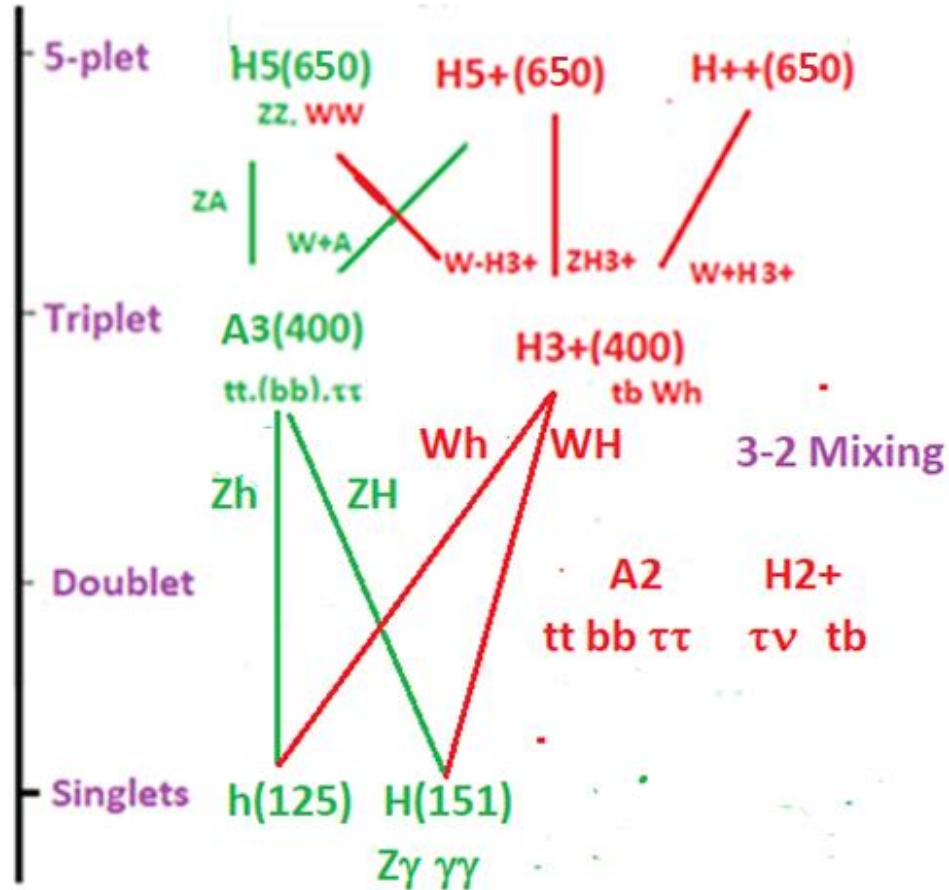


Need for extending GM



- Is GM satisfying the various observations ?
- The answer is NO, in particular for what concerns the **fermionic couplings** of $A(400)$ which tell us that **$Y_{t \sim SM}$** while **$Y_{b, \tau} \gg SM \text{ and GM}$**
- The remedy is to add an **extra doublet** and benefit from an enhancement of **$Y_{b, \tau} \sim \tan\beta \sim 20$** 'à la MSSM'. Too naïve since then **$Y_{t \sim 1/\tan\beta}$**
- The **Yukawa alignment mechanism** is a more general scheme sufficient to suppress **FCNC** and allowing an **independent tuning for u, d, ℓ**
- It assumes that both doublets couple to all fermions requiring **$Y_{2f} = \xi_f Y_{1f}$** where Y_{1f} and Y_{2f} are the Yukawa couplings to the two doublets ϕ_1 and ϕ_2 , and where ξ_f is an arbitrary constant which can be complex and differ for **u, d, ℓ**
- One can then have **$Y_{b, \tau} \gg SM$** even if $\tan\beta \sim 1$ and $Y_{t \sim SM}$ [0908.1554](https://arxiv.org/abs/0908.1554)

GM+2 doublets model



- Yukawa couplings
- A2HDS $\zeta_{\ell, d} \sim 20$ $\zeta_u \sim 0.7$

	Y_1^d	Y_1^u	Y_1^ℓ	Y_2^d	Y_2^u	Y_2^ℓ	ζ_u	ζ_d	ζ_ℓ
Type I	0	0	0	x	x	x	t_β^{-1}	t_β^{-1}	t_β^{-1}
Type II	x	0	x	0	x	0	t_β^{-1}	$-t_\beta$	$-t_\beta$
Type X	0	0	x	x	x	0	t_β^{-1}	t_β^{-1}	$-t_\beta$
Type Y	x	0	0	0	x	x	t_β^{-1}	$-t_\beta$	t_β^{-1}
A2HDS	x	x	x	x	x	x	$\frac{\xi_u - t_\beta}{1 + \xi_u t_\beta}$	$\frac{\xi_d - t_\beta}{1 + \xi_d t_\beta}$	$\frac{\xi_\ell - t_\beta}{1 + \xi_\ell t_\beta}$

Giving up mass degeneracy in the scalar sector

- **Mass degeneracy** inside H5 and H3 is not necessary to insure the constraint $\rho \sim 1$ as pointed out in [2111.14195](#)
- This aspect is very important for searches
- This aspect is critical for experimental searches and could explain the absence of charged Higgs scalars in ZW and W+W+ if $m_{H3} < m_A$, in which case $H5^{++} \rightarrow H3 + W+$ competes with W+W+

$$V(\Phi, X) = \frac{\mu_2^2}{2} \text{Tr}(\Phi^\dagger \Phi) + \frac{\mu_3^2}{2} \text{Tr}(X^\dagger X) + \lambda_1 [\text{Tr}(\Phi^\dagger \Phi)]^2 + \lambda_2 \text{Tr}(\Phi^\dagger \Phi) \text{Tr}(X^\dagger X) + \lambda_3 \text{Tr}(X^\dagger X X^\dagger X) + \lambda_4 [\text{Tr}(X^\dagger X)]^2 - \lambda_5 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^a) \text{Tr}(X^\dagger t^a X t^a) - M_1 \text{Tr}(\Phi^\dagger \tau^a \Phi \tau^a) (UXU^\dagger)_{ab} - M_2 \text{Tr}(X^\dagger t^a X t^a) (UXU^\dagger)_{ab}.$$

$$V(\phi, \chi, \xi) = \bar{\mu}_2^2 \phi^\dagger \phi + \bar{\mu}_3^2 \chi^\dagger \chi + \frac{\bar{\mu}_\xi^2}{2} \xi^\dagger \xi + \bar{\lambda}_1 (\phi^\dagger \phi)^2 + \bar{\lambda}_2 |\bar{\chi}^\dagger \chi|^2 + \bar{\lambda}_3 (\phi^\dagger \tau^a \phi) (\chi^\dagger t^a \chi) + [\bar{\lambda}_4 (\bar{\phi}^\dagger \tau^a \phi) (\chi^\dagger t^a \xi) + \text{h.c.}] + \bar{\lambda}_5 (\phi^\dagger \phi) (\chi^\dagger \chi) + \bar{\lambda}_6 (\phi^\dagger \phi) (\xi^\dagger \xi) + \bar{\lambda}_7 (\chi^\dagger \chi)^2 + \bar{\lambda}_8 (\xi^\dagger \xi)^2 + \bar{\lambda}_9 |\chi^\dagger \xi|^2 + \bar{\lambda}_{10} (\chi^\dagger \chi) (\xi^\dagger \xi) - \frac{1}{2} [M'_1 \phi^\dagger \Delta_2 \bar{\phi} + \text{h.c.}] + \frac{M_1}{\sqrt{2}} \phi^\dagger \Delta_0 \phi - 6M_2 \chi^\dagger \bar{\Delta}_0 \chi.$$

- In [2204.12898](#) it was shown that with this type of potential one can interpret the result of CDF II by assuming that $v_\xi \neq v_\chi$

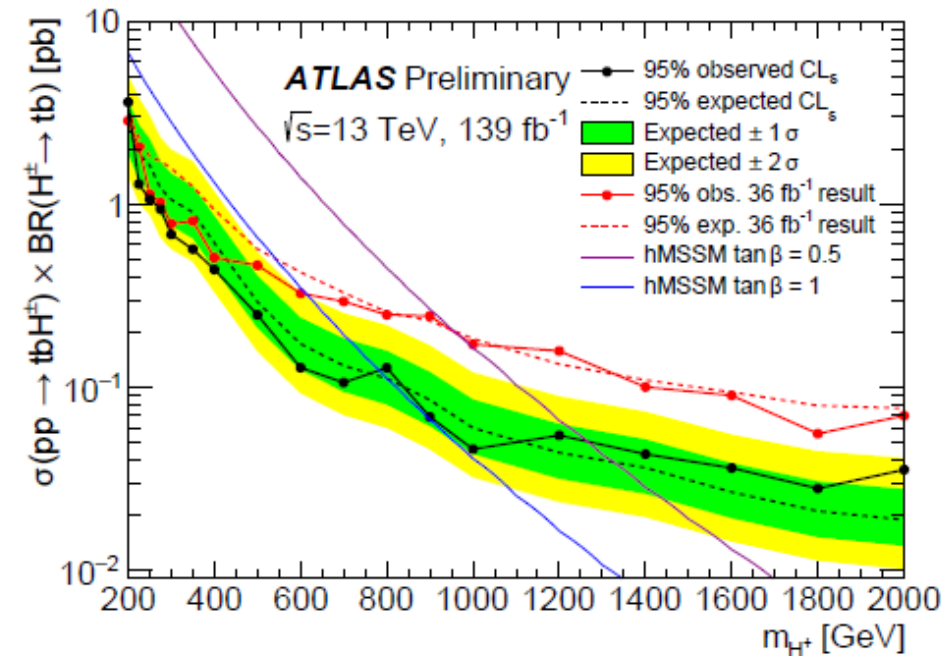
What is meant by eGM+a ?

- LHC observations result into eGM+a: 2 doublets, 2 triplets, 1 CP-odd singlet a, 1 CP-even singlet h(96)
- With custodial symmetry breaking terms it can break **mass degeneracy** of 3-plets and 5-plets usual in GM [2111.14195](#)
- Lifting the triplet vacuum degeneracy it can provide the **ρ parameter** contribution needed for the CDF II M_W excess [2204.12898](#)
- The **additional doublet** allows to cope with A decays, using **Yukawa alignment** to explain fermion couplings
- The **additional singlet CP-odd a(52)** introduced to explain **H(151) decays** can be used in various ways:
 - With a large **$a\mu\mu$ coupling enhancement** it can explain the g-2 effect
 - [2202.12631](#) also assume that a(52) can decay into **DM** pairs
 - **EW phase transition** is helped by a(52) [0705.2425](#)
- The role of h(96), left out of eGM+a, is unclear unless one assumes that it is the isoscalar needed to stabilize the brane of RS
- eGM+a would have a very **complex scalar potential** with 12 isoscalars

LHC issues

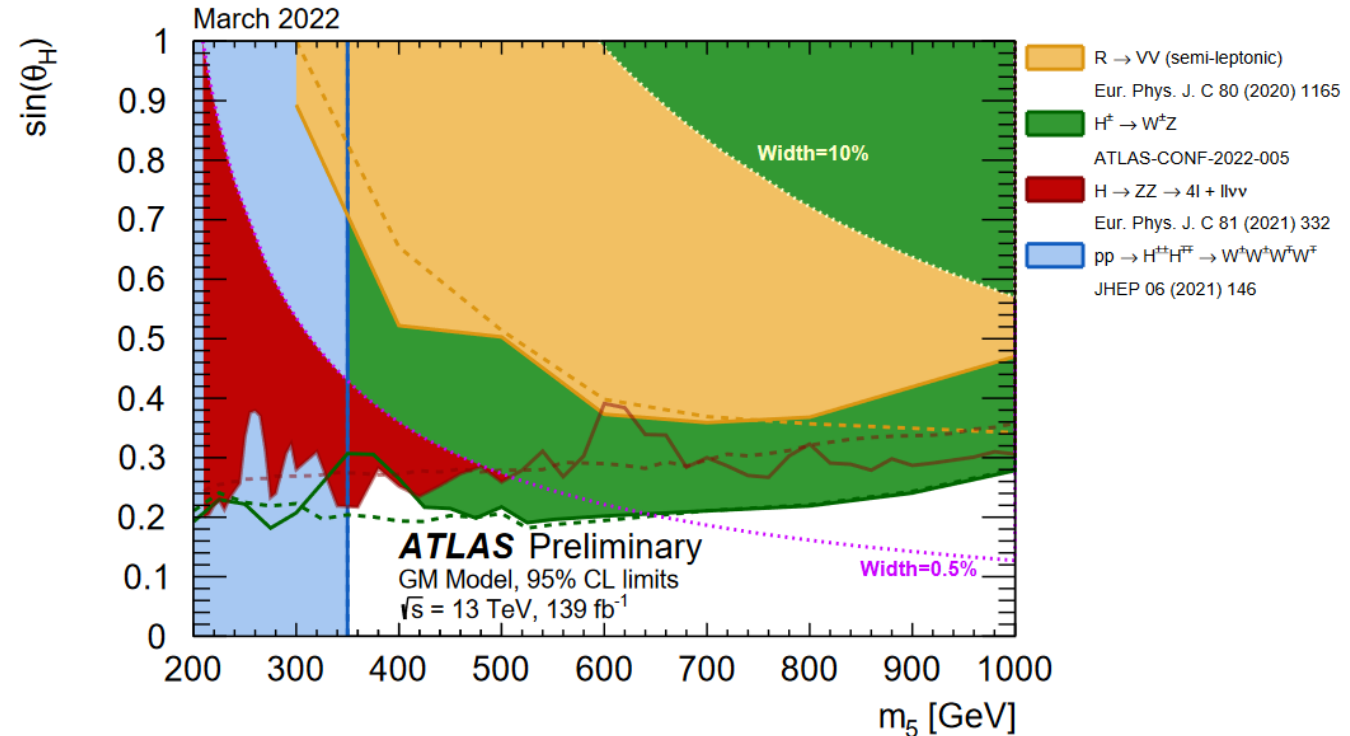
A worry

- How present indications, if real, can progress with increased luminosity ?
- There is a common tendency to expect **massive objects**, hence to optimize the various selections accordingly, in contrast to the strategy used to discover $h(125)$
- This may result in an absence of progress for relatively light scalars, a concern in the present scenario
- An example of this is shown in the search for $H \rightarrow tb$ where progress in luminosity is unclear for masses ~ 400 GeV which are relevant in our case

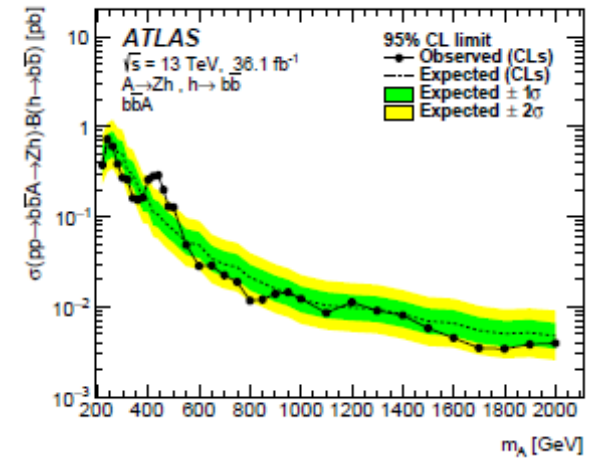


ATLAS GM summary plot

- ATL-PHYS-PUB-2022-008
- The present solution which assumes $\sin\theta_H=0.5$ and $m_5=650$ GeV seems incompatible with this summary plot of ATLAS
- Seems covered by $H \rightarrow W+Z$ not dominant when $ZH3+$ opens up which is the case
- An additional contribution of $H \rightarrow W+a$ competing with ZW is also ignored
- $H(650)$ relative width $\sim 0.5\%$ is invalid in the assumed scenario ($100/650=15\%$)

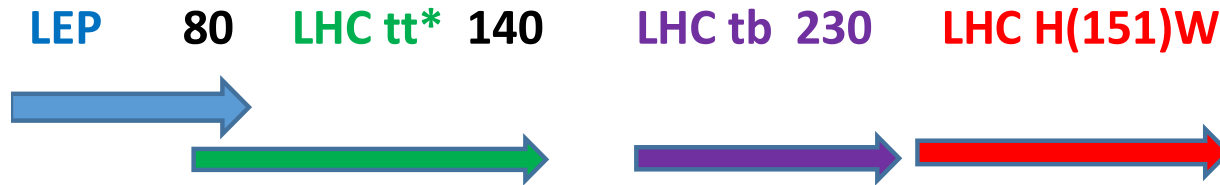
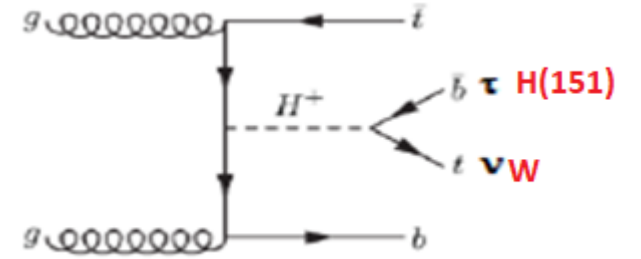


Why does ATLAS observe $A \rightarrow hZ$?



- The $A \rightarrow hZ$ GM coupling seems too small to explain the ATLAS observation
- $gg \rightarrow A \rightarrow H(151)Z$ has a cross section 20 to 30 times larger $\sim 3000 \text{ fb}$
- Assuming $H(151) \rightarrow a(52)a(52)$ with $a \rightarrow bb$ in $\sim 90\%$ of cases, which fraction of these events can fake hZ ?
- ATLAS requires two “small R-jets” and restricts the m_h range to 110-145 GeV, which implies that a large fraction of $H(151)Z$ events will be lost
- This can explain why the observed excess in ggF is only $\sim 200 \text{ fb}$, implying that $\sim 10\%$ of the HZ decays are accepted by this analysis
- For the bbA topology, the $4b$ final states could allow accidental selections of the $gg \rightarrow A$ contribution
- Given the ratio hZ/HZ , one can therefore not exclude that the observed signal is dominantly due to $A \rightarrow HZ$
- The reconstruction efficiency of the “ hZ ” channel depends on the experimental selections and may vary, perhaps explaining apparent contradictory results

What about H3+ ?

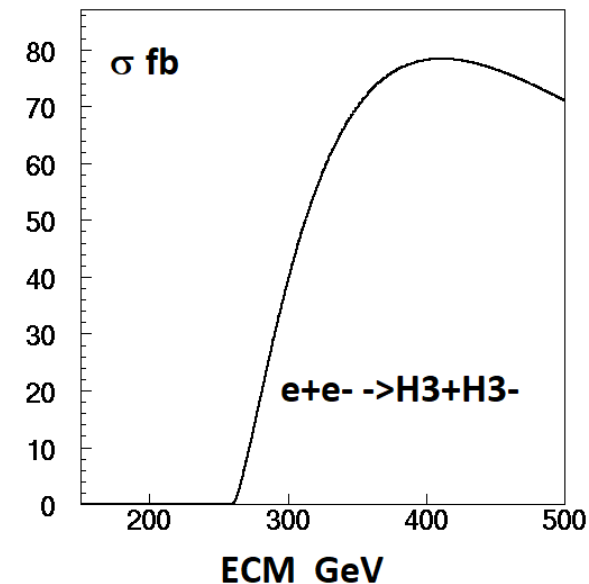
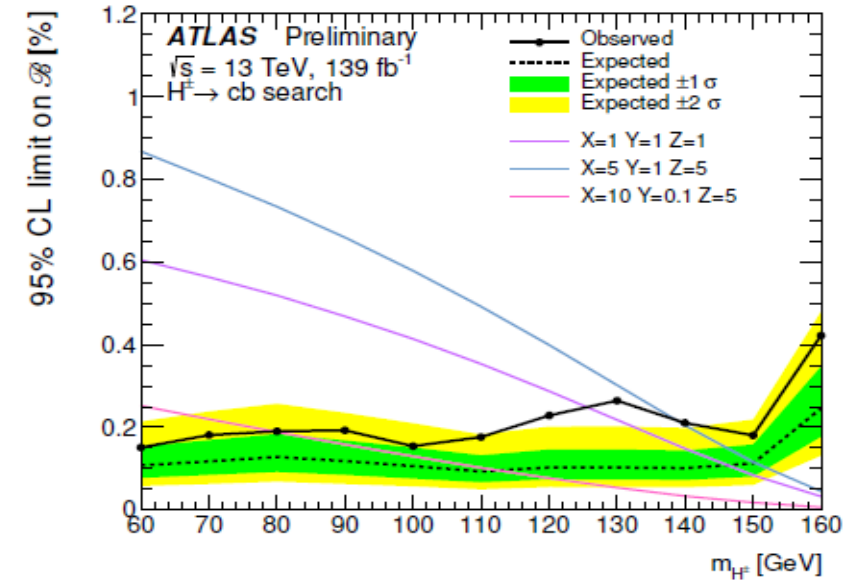


- $m_{H3+}=m_A=400$ GeV is predicted in GM however [2111.14195](#) tells us that mass degeneracy can be relaxed
- -> Full coverage is needed noting that there is a **blind zone** between 140 and the tb threshold
- The GM dominant decay mode $H3+ \rightarrow H(151)W+$ has not been searched for (hW+ subdominant)
- For $m_{H3+}=400$ GeV GM predicts $BR(HW)/BR(tb) \sim 1$
- An excess is reported on $H+ \rightarrow cb$ (next slide) with $m_{H+}=130$ GeV, consistent with the eGM model

ATL-PHYS-PUB-2020-006
23 March 2020

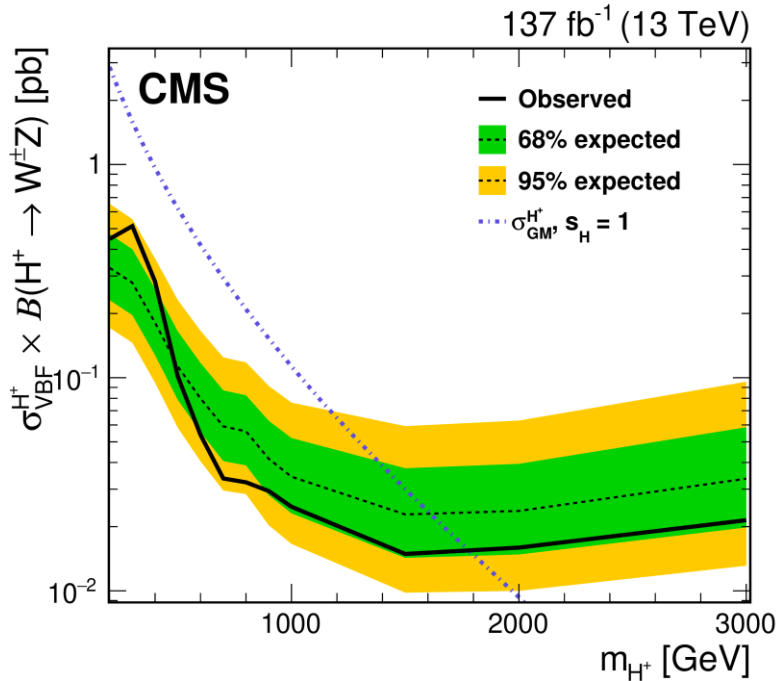
What if $m_{H^{3\pm}}=130$ GeV ?

- An excess is reported by ATLAS: a 3 sd deviation in $H^{\pm} \rightarrow cb$ with $m_{H^{3\pm}}=130$ GeV [EPS-HEP2021, 631](#)
- Main source is $gg \rightarrow tt^*$, $t^* \rightarrow H+b$ with a cross section of $\sim 10 \mu\text{b}$ but a huge background
- $\mathcal{B} \equiv \text{BR}(t \rightarrow Hb)\text{BR}(H \rightarrow cb) = 0.16 \pm 0.6\%$
- Inconsistent with standard GM but acceptable within eGM [2111.14195](#)
- For H^{++} $\text{BR}(H^{3+}W^+)/\text{BR}(W^+W^+) \sim (cH/sH)^2 \sim 3$ could explain why W^+W^+ is marginally observed at ~ 500 GeV by CMS
- $\text{BR}(H^{3+}H^{3+})/\text{BR}(W^+W^+)$ non negligible, depends on the details of the scalar potential
- $\Gamma(H^{50} \rightarrow H^{3+}W) = 17$ GeV hence $\Gamma_{H^5} = 130$ GeV, still acceptable
- $e^+e^- \rightarrow H^{3+}H^{3-}$, $H^{3\pm}W$ have large cross sections and good visibility

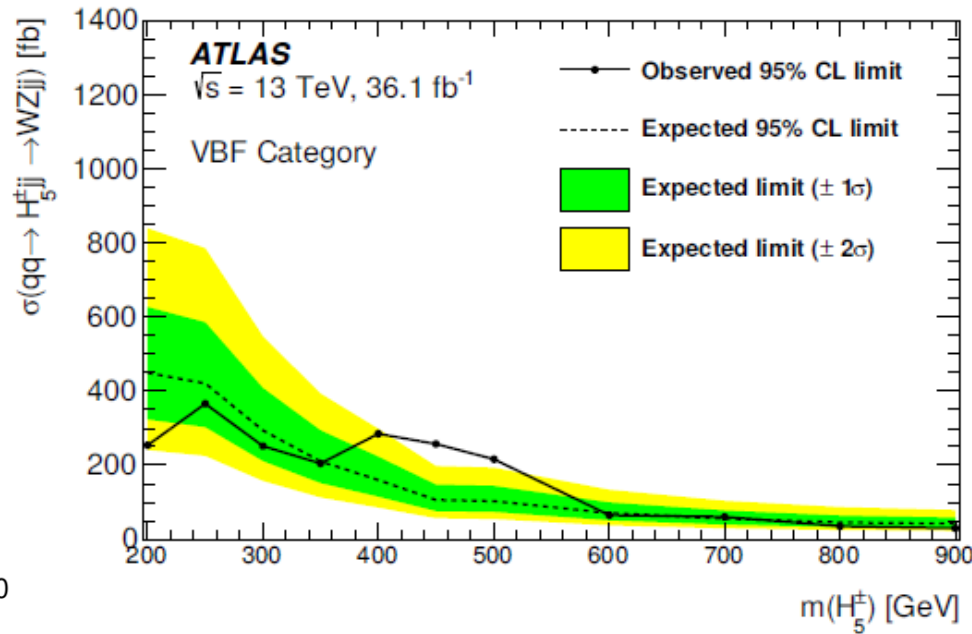


What about H5+ and H5++ ?

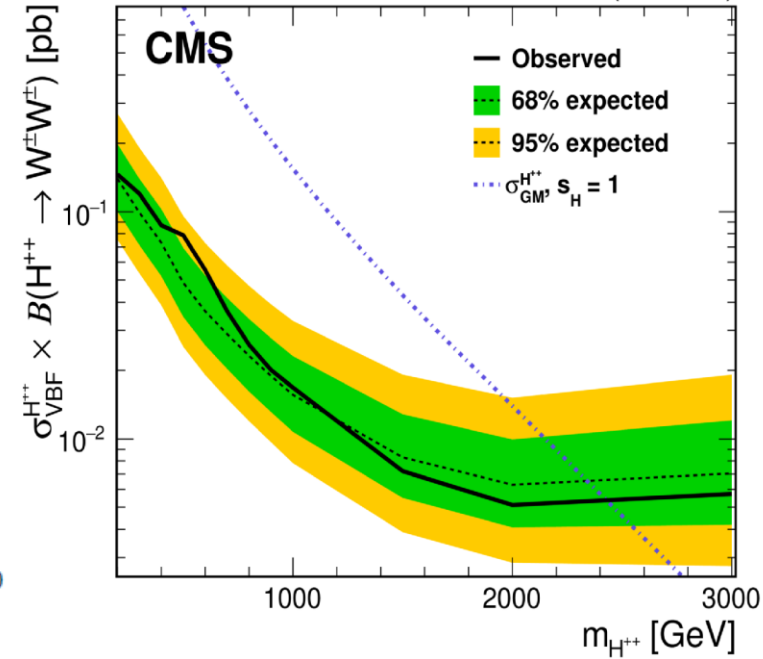
[2104.04762](#)



[1806.01532](#)



137 fb⁻¹ (13 TeV)



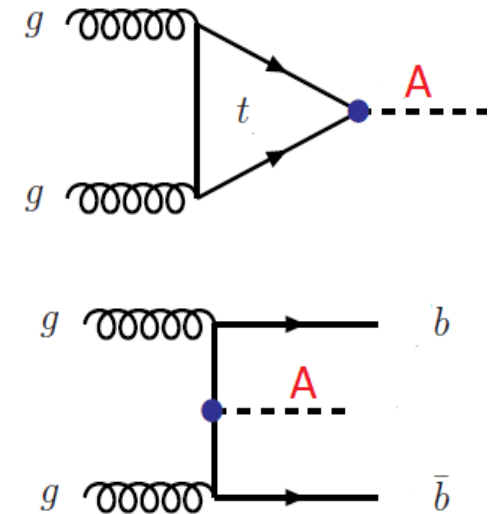
- CMS cross sections assume $s_H=1$ are divided by 4 for $s_H=0.5$
- If H3+ is light H3+Z and H3+W+ become dominant and these resonances become wide
- Coincident excess at $m_{H5+} \sim 450$ GeV for ATLAS (2.9 sd) & CMS while GM predicts 650 GeV
- Not excluded in eGM [2111.14195](#)

Confusing role of $a(52)$

- This particle can mix with H_{30} (and A_2 in eGM)
- One can have transitions like $H_{5(650)} \rightarrow Za$, $H_{5^+} \rightarrow W+a$, $A \rightarrow ah/H$, $H_{3^+} \rightarrow Wa$ suppressed by $\sin^2\theta$ but kinematically favoured
- Will compete with the dominant modes ZZ/WW , ZW , $H(151)Z$
- Could perhaps explain the absence of H_{5^+} and H_{3^+} signals
- $a(52)$ decays predominantly into bb (90%) and $\tau\tau$ (10%)
- One expects that $A \rightarrow H(151)Z$ with $H \rightarrow aa$ which can be misidentified as $A \rightarrow h(125)Z$ with $h(125) \rightarrow bb$
- These modes can be searched at LHC
- Many models assume that a connects to DM $a \rightarrow \chi\chi$

Prospects at HL-LHC

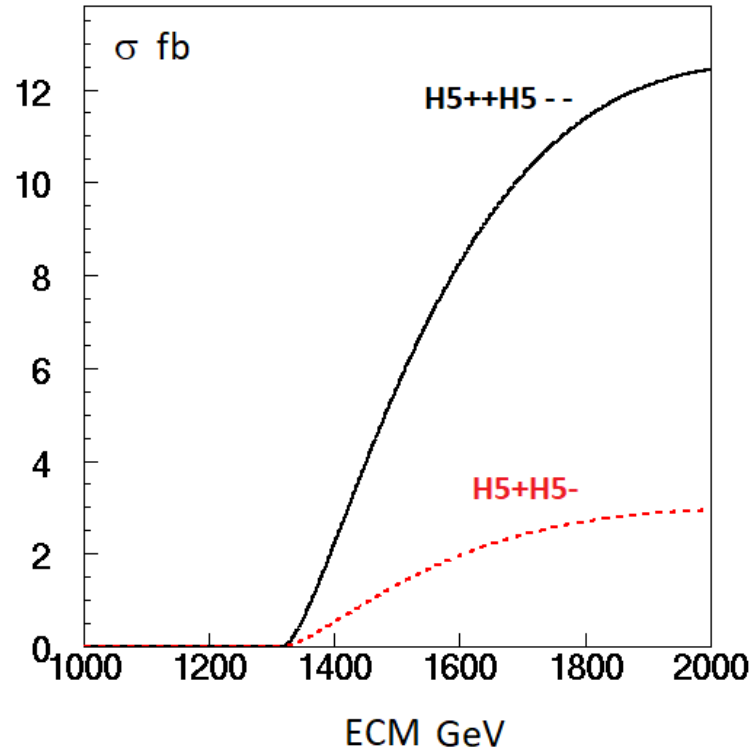
- eGM is a gold mine for LHC with a very rich spectrum of scalars with multiple signatures
- LHC has to:
- Confirm simplest topologies $H(650) \rightarrow ZZ \rightarrow 4\ell, WW \rightarrow \ell\ell\nu\nu$
 $H(151) \rightarrow 2\gamma + E_{\text{miss}}$
- Confirm $h(96) \rightarrow 2\gamma$
- Confirm $A \rightarrow tt$ & $A \rightarrow hZ, \tau\tau + b$
- Search for $A \rightarrow H(151)Z$ & $H_{3+} \rightarrow H(151)W$ with $H \rightarrow aa \rightarrow 4b$
- Confirm $H_{3+}(130) \rightarrow cb$
- Confirm $H_{5+-} \rightarrow W+Z$ $H_{5++-} \rightarrow W+W+$ with VBF selection
- Select VBF and try to reconstruct complex final states like $H_{5-} \rightarrow AZ \rightarrow ttZ$ and $H_{5+-} \rightarrow AW \rightarrow ttW+$
- Improve searches on $H_{++} \rightarrow W+W+$
- ...



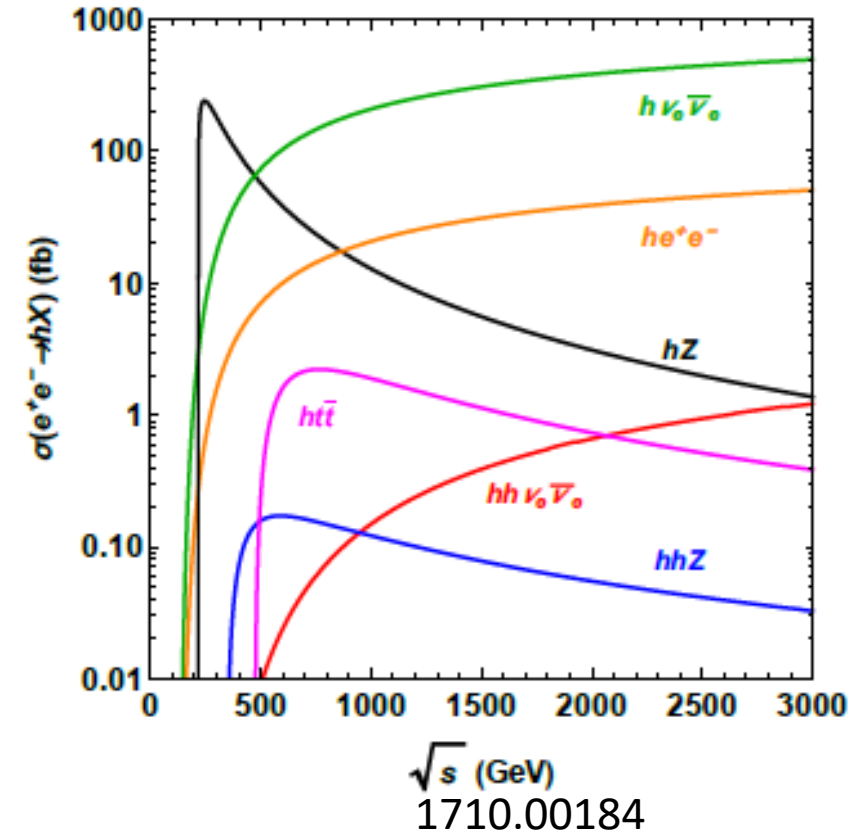
e^+e^- Colliders

High energy e+e- cross-sections

- H5++(650) and H5+(650) in GM
- SM h(125) cross-sections

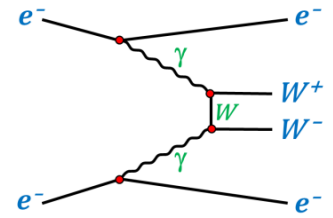


Could be lighter in eGM



\sqrt{s} (GeV)
1710.00184

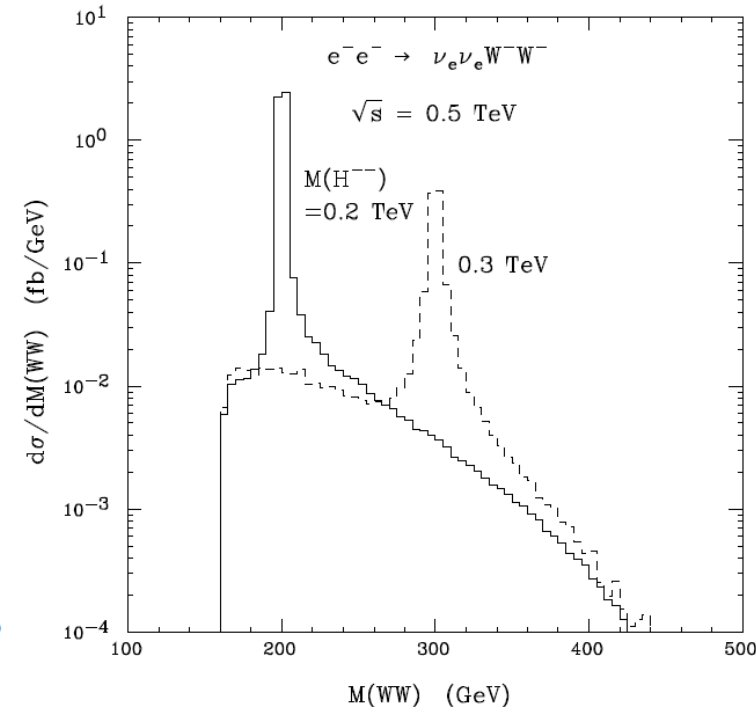
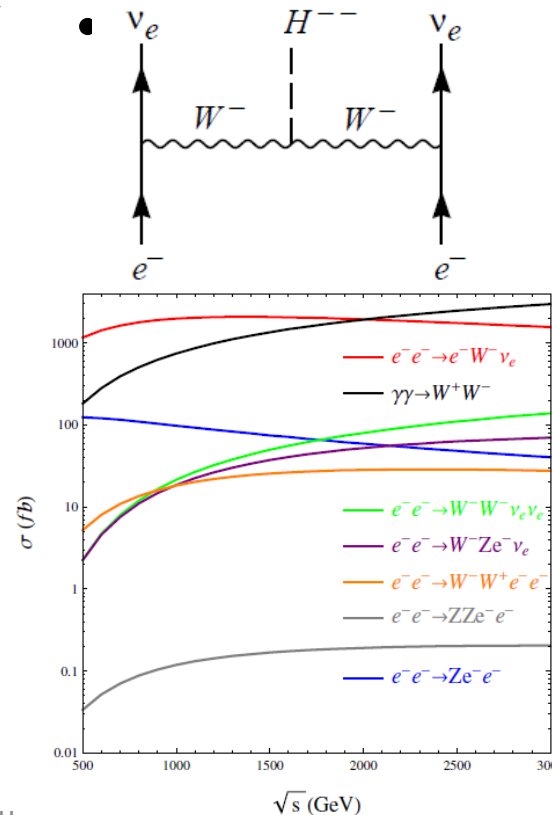
e-e- collider issues for H⁻ - (650)



- $\mathcal{L}_{e-e} \sim 70\% \mathcal{L}_{e+e-}$ seems feasible (Drutskoy at ILCX2021)
- $\sigma_{\text{VBF}}(e-e \rightarrow H^{--}(650)) = 3.5 \text{ fb}$ at 1 TeV
- With polarized beams, 80% for e⁻, luminosity gain 1.8² for W⁻W⁻
- Switch from e⁺e⁻ will require changing polarity of many magnets
- Telnov ERL scheme seems feasible
- Ref 9404335 $\sin\theta_H = 0.7$ (2 times smaller signal in my case)
- [2109.05855](#) ($\sin\theta_H = 0.1$)

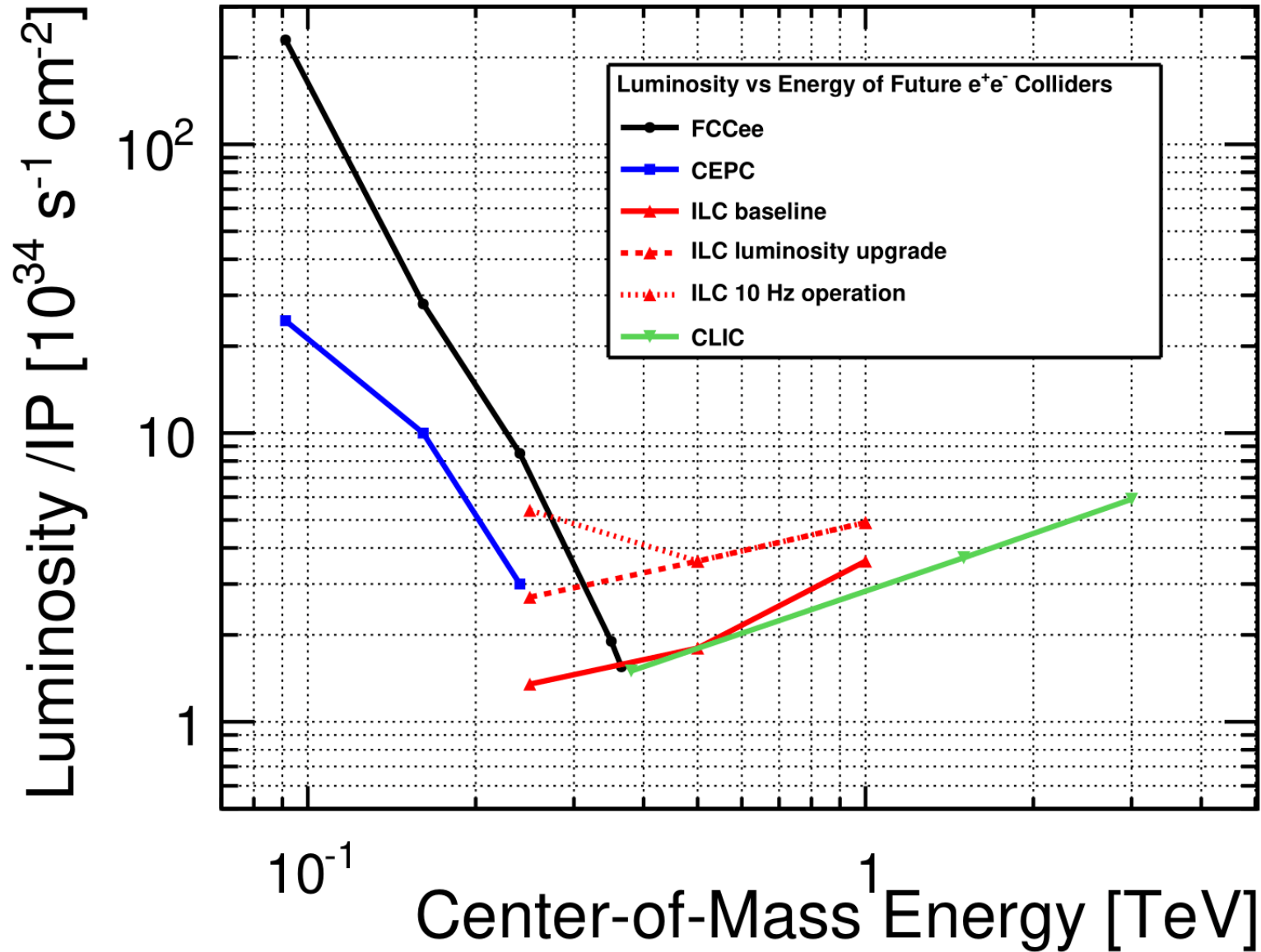
- **H⁻(650) → H³⁻(400)W⁻** with **H³⁺ → tb, HW** easily separated from W⁻W⁻ using **b-tagging** + multijet cuts

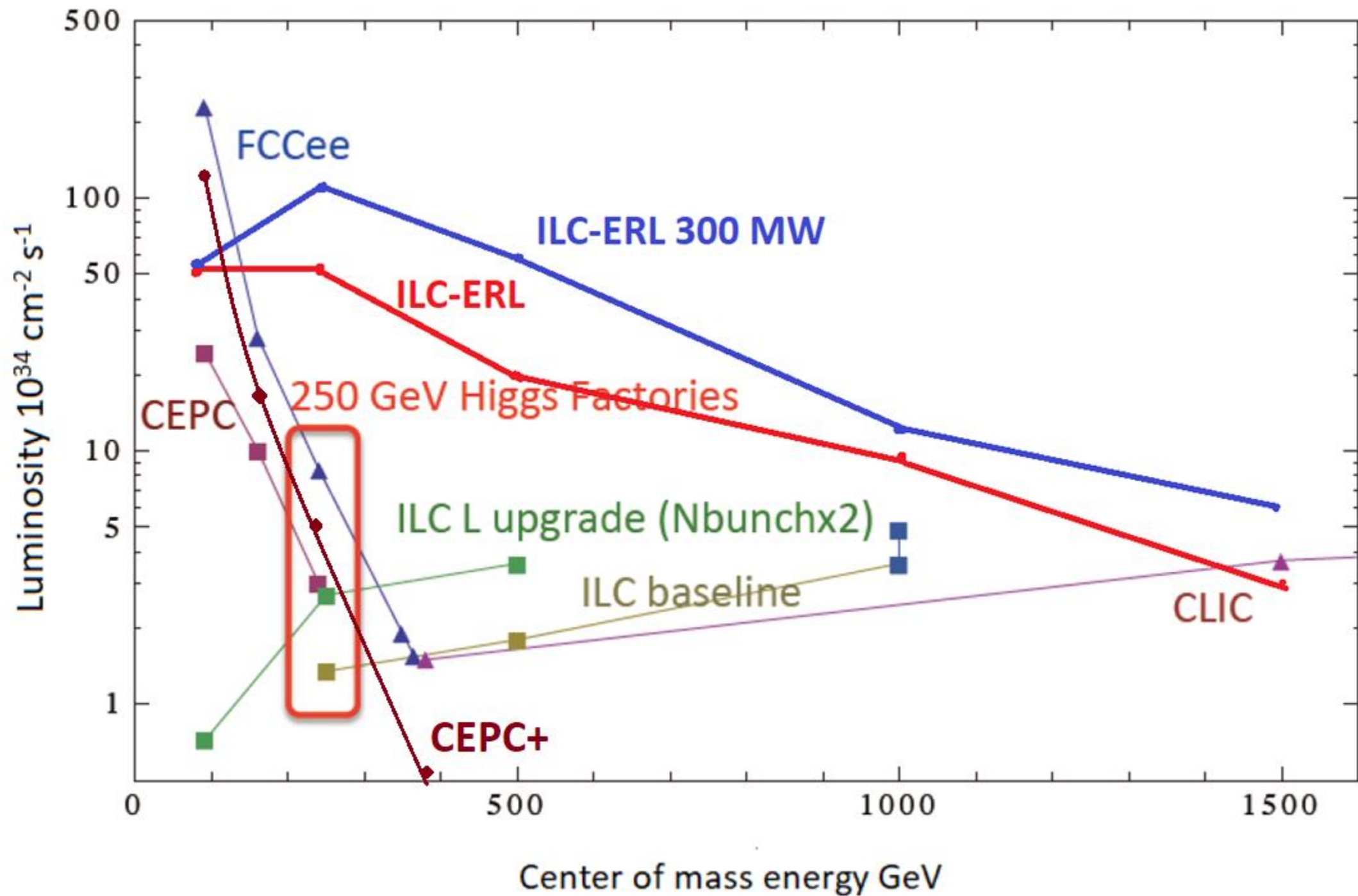
[9404335](#)



LUMINOSITY at 1 TeV

- In reference [1903.01629](#) a running scenario of ILC at **1 TeV collecting 8000 fb⁻¹** has been envisaged
- Beneficial for **Higgs self-coupling** measurement
- Discoveries at LHC would boost these studies at ILC and CLIC
- Convert ILC into an ERL [2105.11015](#) ?





Quantity	Symbol	Unit	Initial	\mathcal{L} Upgrade	Z pole	500	Jpgrades	1000
Centre of mass energy	\sqrt{s}	GeV	250	250	91.2	500	250	1000
Luminosity	\mathcal{L}	$10^{34} \text{cm}^{-2} \text{s}^{-1}$	1.35	2.7	0.21/0.41	1.8/3.6	5.4	5.1
Polarization for e^-/e^+	$P_-(P_+)$	%	80(30)	80(30)	80(30)	80(30)	80(30)	80(20)
Repetition frequency	f_{rep}	Hz	5	5	3.7	5	10	4
Bunches per pulse	n_{bunch}	1	1312	2625	1312/2625	1312/2625	2625	2450
Bunch population	N_e	10^{10}	2	2	2	2	2	1.74
Linac bunch interval	Δt_b	ns	554	366	554/366	554/366	366	366
Beam current in pulse	I_{pulse}	mA	5.8	8.8	5.8/8.8	5.8/8.8	8.8	7.6
Beam pulse duration	t_{pulse}	μs	727	961	727/961	727/961	961	897
Average beam power	P_{ave}	MW	5.3	10.5	1.42/2.84 ^{*)}	10.5/21	21	27.2
RMS bunch length	σ_z^*	mm	0.3	0.3	0.41	0.3	0.3	0.225
Norm. hor. emitt. at IP	$\gamma\epsilon_x$	μm	5	5	5	5	5	5
Norm. vert. emitt. at IP	$\gamma\epsilon_y$	nm	35	35	35	35	35	30
RMS hor. beam size at IP	σ_x^*	nm	516	516	1120	474	516	335
RMS vert. beam size at IP	σ_y^*	nm	7.7	7.7	14.6	5.9	7.7	2.7
Luminosity in top 1 %	$\mathcal{L}_{0.01}/\mathcal{L}$		73 %	73 %	99 %	58.3 %	73 %	44.5 %
Beamstrahlung energy loss	δ_{BS}		2.6 %	2.6 %	0.16 %	4.5 %	2.6 %	10.5 %
Site AC power	P_{site}	MW	111	138	94/115	173/215	198	300
Site length	L_{site}	km	20.5	20.5	20.5	31	31	40

Table 4.1: Summary table of the ILC accelerator parameters in the initial 250 GeV staged configuration and possible upgrades. A 500 GeV machine could also be operated at 250 GeV with 10 Hz repetition rate, bringing the maximum luminosity to $5.4 \cdot 10^{34} \text{cm}^{-2} \text{s}^{-1}$ [26]. *): For operation at the Z-pole additional beam power of 1.94/3.88 MW is necessary for positron production.