

PASCOS 2022: **S.M.A.S.H.E.D.**
*Standard Model Axion Seesaw Higgs Inflation
Extended for Dirac Neutrinos*

based on **2207.08142**

Maximilian Berbig (BERBIG@PHYSIK.UNI-BONN.DE)

26.07.2022



- 1 Quick review of S.M.A.S.H.
- 2 Dirac Neutrinos
- 3 Axion pheno in S.M.A.S.H.E.D.
- 4 Summary

- **SM**: self explanatory...
- **Axion**:
 - strong CP via $U(1)_{PQ}$
 - cold **DM** candidate
- **Seesaw mechanism**:
 - heavy N for Majorana m_ν
 - baryogenesis via leptogenesis
- **Higgs-portal inflation**:
 - $T_{RH} \simeq 10^9$ GeV
 - complete therm. history

A single new scale

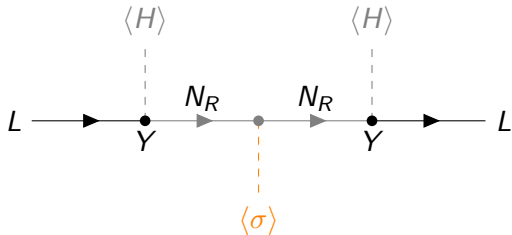
PQ breaking scale v_σ

= axion decay constant f_a

= Seesaw mass scale M

¹(Ballesteros, Redondo, Ringwald 2016)

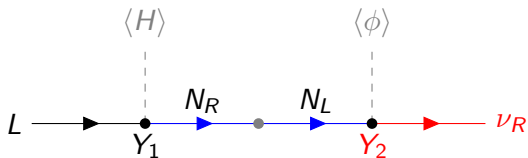
Type I Seesaw: Majorana vs. Dirac



(Yanagida 1979 et al.)

$$m_\nu^{\text{Maj.}} \sim \frac{v_H^2}{M} \sim \frac{v_H^2}{v_\sigma}$$

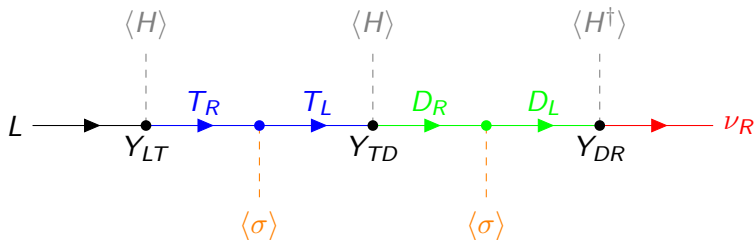
all N : SM singlets



(Roncadelli, Wyler 1983)

$$m_\nu^{\text{Dir.}} \sim \frac{v_H v_\phi}{M} \sim \begin{cases} \frac{v_H v_\phi}{v_\sigma} \\ \frac{v_H v_\sigma}{M} \end{cases}$$

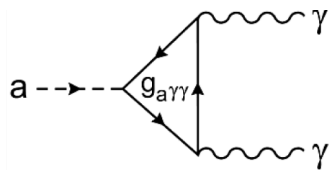
Dimension 6 Dirac Seesaw



$$T_{L,R} \sim (1, 3, 0), \quad D_{L,R} \sim (1, 2, 1/2), \quad M_T, M_D < 10^9 \text{ GeV}$$

$$m_\nu \sim Y_{LT} Y_{TD} Y_{DR} \frac{v_H^3}{M_T M_D} \sim \frac{v_H^3}{v_\sigma^2}$$

Axion Pheno in S.M.A.S.H.E.D. (1)



source: 2106.03424

3 gens. of T and D in loops:

- QCD:

$$N = \sum_f (\chi_{f_L} - \chi_{f_R}) t(f),$$

- QED:

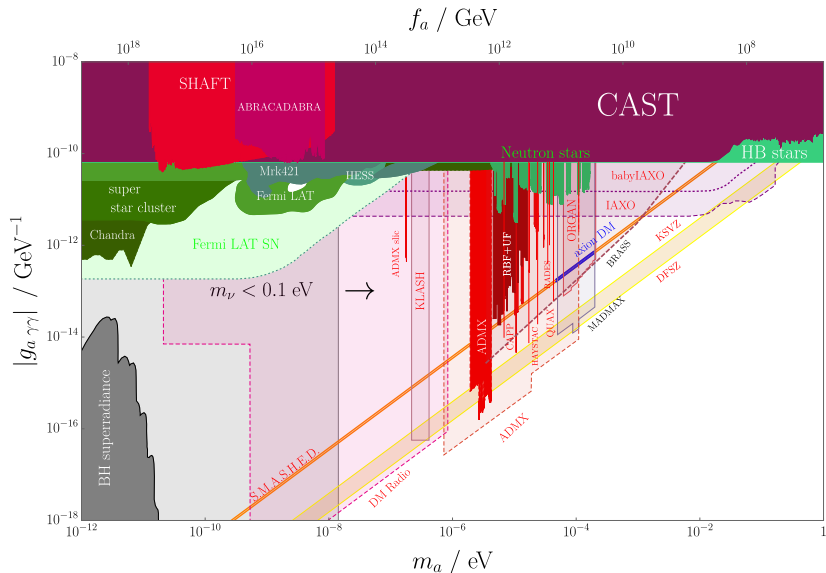
$$E = \sum_f (\chi_{f_L} - \chi_{f_R}) Q_{EM}(f)^2$$

Axion to photon coupling:

$$g_{a\gamma\gamma} = \frac{\alpha_{EM}}{2\pi f_a} \left(\frac{E}{N} - 1.92 \right)$$

- KSVZ, DFSZ: $E/N = 2/3, 8/3$
- S.M.A.S.H.E.D: $E/N = 18 + 2/3, 18 + 8/3$

Axion pheno in S.M.A.S.H.E.D.² (2)



²limits from C. O'hare at <https://cajohare.github.io/AxionLimits/>

Dirac Leptogenesis³ in S.M.A.S.H.E.D.

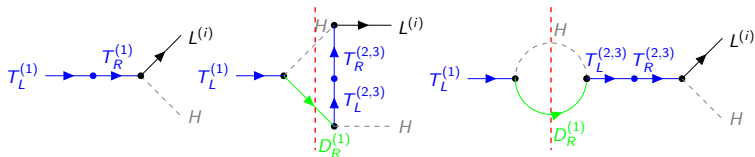
$$\frac{n_B - n_{\bar{B}}}{s} \simeq \kappa \cdot \varepsilon_L \cdot \underbrace{C_{sph}}_{\simeq 0.37} \cdot \underbrace{\frac{n_T + n_{\bar{T}}}{s}}_{\text{no. of } T \text{ and } \bar{T}} \stackrel{!}{=} 10^{-10}$$

- κ : efficiency from **Boltzmann equations**

- larger κ for **Dirac-T**

- ε_L : CP violating asymmetry from **decays** ↓

- larger due to **dim. 6** m_ν



³(Dick, Lindner, Ratz, Wright 1999)

Summary

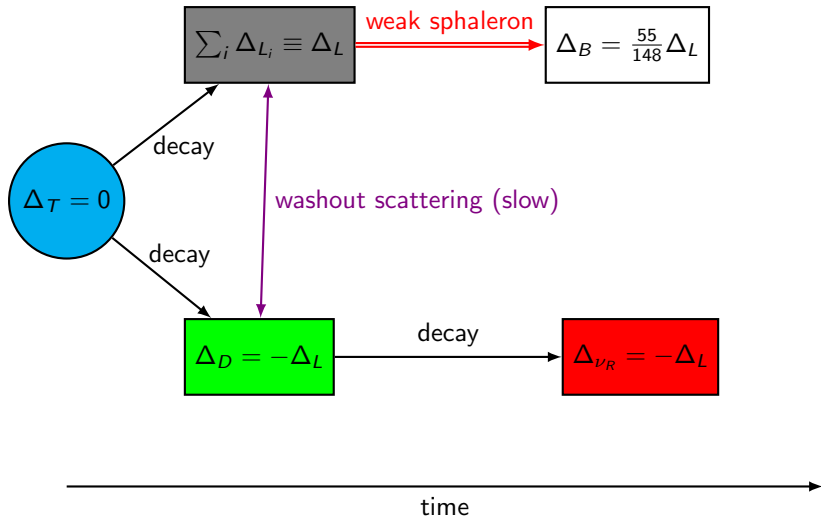
- extended S.M.A.S.H. for Dirac ν without spoiling anything (so far...)
- **dimension 6** operator for m_ν
- $|g_{a\gamma\gamma}|$ enhanced by **one order of magnitude**
- dark radiation $\Delta N_{\text{eff.}} \simeq 0.142 (3 \nu_R) + 0.027 (\text{axion}) \simeq 0.17$
- Dirac Leptogenesis with larger κ & ε_L

⇒ Don't dismiss Dirac neutrinos!

Appendix

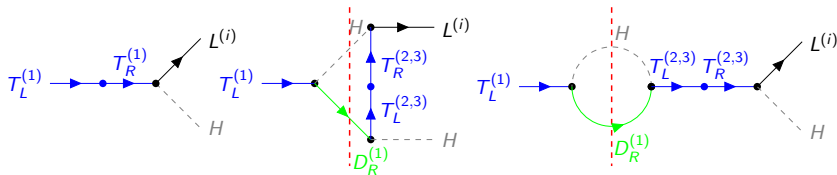
Here be dragons

Dirac Leptogenesis in S.M.A.S.H.E.D. (2)⁴



⁴ $\Delta_\psi \equiv (n_\psi - n_{\bar{\psi}})/s$ with s the entropy density

Dirac Leptogenesis in S.M.A.S.H.E.D. (3)



• Enhancement of ϵ_L :

- Type III Seesaw (*Davidson, Ibarra 2002*):

$$\epsilon_L < 3 \times 10^{-9} \cdot \left(\frac{M_T}{10^8 \text{ GeV}} \right) \cdot \left(\frac{m_\nu}{0.1 \text{ eV}} \right) \Rightarrow M_T \gtrsim 1.5 \times 10^{10} \text{ GeV}$$

- S.M.A.S.H.E.D.:

$$\epsilon_L < 3 \times 10^{-3} \cdot \left(\frac{M_T}{10^8 \text{ GeV}} \right) \cdot \left(\frac{m^{\text{eff.}}}{100 \text{ keV}} \right)$$

- $m_\nu = Y_{LT} Y_{TD} Y_{DR} v_H^3 / (M_T^{(2,3)} M_D^{(2,3)})$

- $m^{\text{eff.}} = Y_{LT} Y_{TD} v_H^2 / M_T^{(2,3)}$ missing $\sim Y_{DR} v_H / M_D^{(2,3)}$ b.c. no ν_R