Searching for dark radiation at the LHC

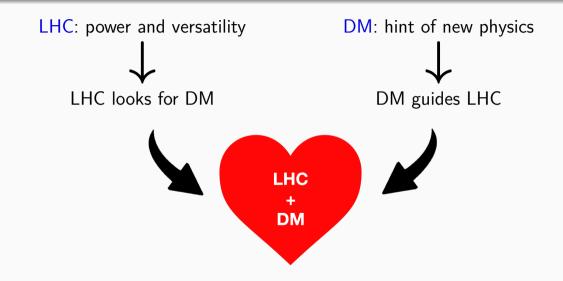
Based on 2204.01759

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A loving relationship...



Simple observation:

 $H(T_{\text{EW}}) \leftrightarrow \text{LHC}$ length

Interactions effective at the EW scale lead to macroscopic decay lengths!

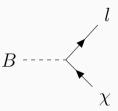
But $\Omega h_{\rm DM}^2$ is not compatible with that...

But there are other cosmological observations!

 $\Delta N_{\rm eff}$ in the near future: $\sigma_{\rm CMB-S4}=0.03$

Model

- New particle content: B, χ
- Massive B decays into light χ
- $B = (B_e, B_\mu, B_\tau)$ charged under SM



$$\mathcal{L}_{\mathsf{NP}} \supset B^T \cdot y_l \cdot (\bar{l}_R \chi) + h.c.$$
$$y_l = \begin{pmatrix} y & 0 & 0 \\ 0 & y & 0 \\ 0 & 0 & y \end{pmatrix} \quad \text{with } y \lesssim 10^{-6}$$

Calculating ΔN_{eff}

 $\Delta N_{\rm eff}$ is the extra radiation added on top of SM

$$\Delta N_{\text{eff}}(x) = \frac{\rho_{\chi}(x)}{\rho_{1\nu}(x)} = \frac{Z_{\chi}(x) s_0^{4/3}}{\frac{7}{8} \left(\frac{4}{11}\right)^{4/3} \rho_{\gamma,0}},$$
$$Z_{\chi}(x) \equiv \frac{\rho_{\chi}(x)}{s^{4/3}(x)}$$

Freeze-in via parent decay: pretty easy!

 $Z_{\chi}(x)$ can be derived by Boltzmann equation!

(1)

(2)

Calculating ΔN_{eff} (standard)

After integrating Boltzmann Equation:

$$\tilde{H}xs^{4/3}(x)\frac{\mathrm{d}Z_{\chi}}{\mathrm{d}x} = \frac{m_B^4\Gamma_B}{8\pi^2}\mathcal{I}(\cdot)$$
(3)

$$\mathcal{I} = \int \frac{\mathrm{d}^3 p_{\chi}}{2E_{\chi}} \int \frac{\mathrm{d}^3 p_{\ell}}{2E_{\ell}} \frac{2E_{\chi}}{E_B} \left(1 - \frac{f_{\chi}(E_{\chi})}{f_{\chi}^{\mathsf{eq}}(E_{\chi})} \right) f_B^{\mathsf{eq}}(E_B) \delta(E_B - E_{\chi} - E_{\ell})$$
(4)

Usual assumptions:

- B decays while non-relativistic
- Backreaction $\chi \operatorname{SM} \to B$ is negligible

Calculating ΔN_{eff} (refined calculation)

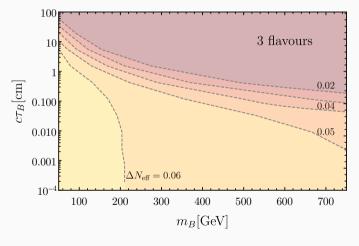
- We relax these assumptions to get a better determination of the parameter space!
- Relativistic treatment for DM already discussed in Arcadi et al.(<u>1906.07659</u>), De Romeri et al.(<u>2003.12606</u>)

$$\tilde{H}xs^{4/3}(x)\frac{\mathrm{d}Z_{\chi}}{\mathrm{d}x} = \frac{m_B^4\Gamma_B}{8\pi^2}\mathcal{I}(x, T_{\chi}, \mathsf{spins})$$

 ${\mathcal I}$ can be tabulated and is provided at the arXiv link.

(5)

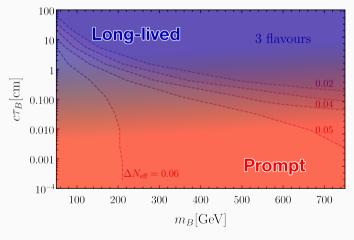
$\Delta N_{\rm eff}$ result and LHC parameter space



- Two model parameters: mass m_B , lifetime τ_B
- Cosmologically interesting

- CMB-S4 target:
$$\Delta N_{\rm eff} = 0.06$$

$\Delta N_{\rm eff}$ result and LHC parameter space



LHC consideration:

- needs prompt searches
- needs LLP searches

Prompt searches

Recast SUSY searches for $l l + \not\!\!\!E_T$ ATLAS (<u>1908.08215</u>), CMS (<u>2012.08600</u>)

Central question: how does the sensitivity change for macroscopic decay lengths?

Particles should decay promptly (i.e. before some Δx):

$$p(x < \Delta x) = 1 - \exp\left(-\frac{\Delta x}{\beta \gamma c \tau}\right) \approx \frac{\Delta x}{\beta \gamma c \tau},$$

Lifetime effect enters in the impact parameter cuts!

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(6)

These cuts are different for ATLAS and CMS!

ATLAS: $|d_0| < 3(5) \ \sigma(d_0)$ for $e^- \ (\mu^-)$, where $\sigma(d_0) \simeq 20 \ \mu$ m

CMS: $|d_0| < 0.5 \mathrm{mm}$ for e^- and μ^-

Of course there are also other differences (taken into account in DELPHES cards and analysis)

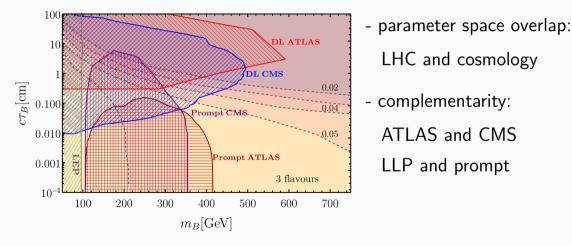
"Recast" SUSY searches for displaced leptons ATLAS (2011.07812), CMS (2110.04809)

Limits provided as $\sigma_{BB}(m_B, c\tau_B)$: can be applied directly to our model

Still different cuts for ATLAS and CMS on the impact parameter

ATLAS: $|d_0| \in [3mm, 300mm]$ **CMS**: $|d_0| \in [0.1mm, 100mm]$

Collider constraints on $\Delta N_{\rm eff}$

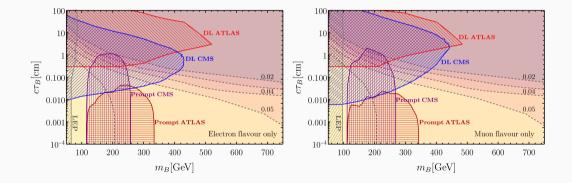


Conclusions

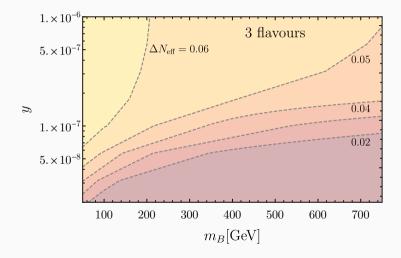
- $\bullet\,$ Calculations of $\Delta N_{\rm eff}$ has been improved to better determine the decay lengths
- ATLAS and CMS have different cuts which result in differences in parameter space probed
- The interesting parameter space lies at the boundary of prompt and long-lived searches → complementarity!

BACKUP

Single flavour case



Coupling parameter space



Prompt ATLAS:

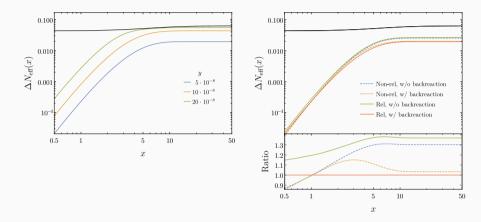
bins in m_{T2} , $e\mu$ as signal region $|z_0 \sin \theta| < 0.5$ mm

Prompt CMS:

bins in $p_T^{\rm miss}$, $e\mu$ as control region $|z_0|<1{\rm mm}$

LLP CMS does not provide limits on $\sigma(m_B, c\tau_B)$ for the single flavour scenario, so an approximation on the mass dependence is used.

Effect of approximations



Not portrayed here: magnitude of corrections depends sensitively on the parameters!

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