

# Minimal SU(6) Gauge-Higgs Grand Unification or How to Unify Gauge Interactions with Their Breaking Sector

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Based on 2104.07366 [PRD] & work in progress,  
in collaboration with Andreas Bally, Simone Blasi, Florian Goertz, & Sascha Weber

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# Unification in Physics

**Unification**  $\longleftrightarrow$  deeper understanding of Nature:

- Maxwell  $\rightarrow$  **electromagnetism**
- Einstein  $\rightarrow$  **spacetime**
- Glashow, Salam, Weinberg  $\rightarrow$  **Standard Model** of **Electroweak (EW)** interactions
- Georgi, Glashow  $\rightarrow$  strong + EW interactions  $\rightarrow$   **$SU(5)$  grand unified theory (GUT)**
- Strong + EW interactions + Higgs sector?  $\rightarrow$  go to 5D (warped) spacetime!

GAUGE-HIGGS GRAND UNIFICATION

$$\Rightarrow A_M = (A_\mu, A_5), \{G_\mu^a, W_\mu^i, B_\mu\} \subset A_\mu, H \subset A_5$$

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# Hierarchy Problems: Gauge, Flavour, Doublet–Triplet Splitting

- UV model  $\rightarrow$  Higgs + heavy VL fermions:

$$-\mathcal{L}_{\text{UV}} \supset m_L^2 |H|^2 + M^2 (\bar{L}L + \bar{E}E) + (y \bar{L}H E + \text{h.c.})$$

- Match on an EFT (the SM is an EFT!):

$$-\mathcal{L}_{\text{EFT}} \supset m^2 |H|^2$$

- One-loop matching at  $\mu = M$  (avoid large logs):

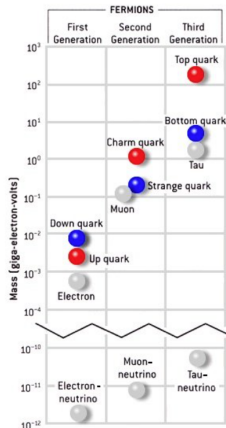
$$m^2(M) = m_L^2(M) + \frac{y^2}{16\pi^2} M^2$$

- But if  $\frac{y^2}{16\pi^2} M^2 \gg m^2 \Rightarrow$  very precise cancellation (fine tuning) to get  $m^2 \simeq (100 \text{ GeV})^2$ :

$$m_L^2 \simeq -\frac{y^2}{16\pi^2} M^2 \gg m^2$$

- Doublet–triplet splitting:  $\mathbf{5}_{SU(5)} \rightarrow H(1, 2)_{1/2} \oplus S_1(3, 1)_{-1/3}$

$\rightarrow m_h \simeq 100 \text{ GeV}$  but  $m_{S_1} \geq 10^{12} \text{ GeV}$  (proton decay)  $\Rightarrow$  "tree-level hierarchy problem"!



# Contents

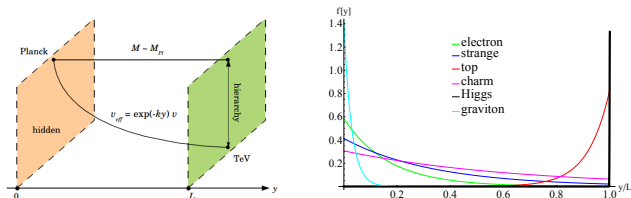
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- 2 Gauge–Higgs Grand Unified Theories (GHGUTs)
- 3 Minimal  $SU(6)$  GHGUT
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# The Randall-Sundrum (RS) Model [Randall, Sundrum '99]



- Setup  $\rightarrow$  two 4D worlds ("branes") situated at the endpoints of the ED:
  - $y = 0 \rightarrow$  Planck/UV/hidden brane;
  - $y = \pi R \equiv L \rightarrow$  TeV/IR/visible brane;
- Between the two branes  $\rightarrow$  **bulk** with anti-de-Sitter ( $AdS_5$ ) **warped geometry**:

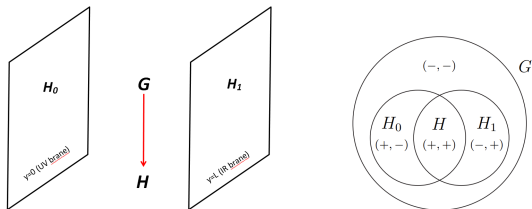
$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2 \Rightarrow g_{MN} = (e^{-2ky} \eta_{\mu\nu}, -1);$$

$$k \sim M_{Pl}, \quad kL \simeq \mathcal{O}(30) \Rightarrow ke^{-kL} = \mathcal{O}(1 \div 10 \text{ TeV});$$

- **Gauge hierarchy** problem solution?  $\rightarrow$  **Higgs localized** close to/on the **IR** brane

$$\Lambda_{EW} \sim ke^{-kL} \equiv m_{KK}$$

# Gauge–Higgs (Grand) Unification: General Aspects [Manton '79; Fairlie '79, Hosotani '83]

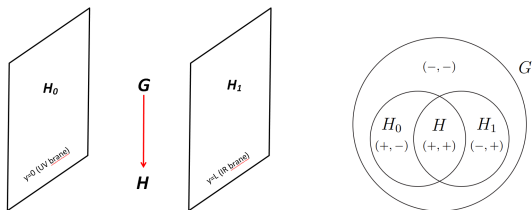


- Simple idea  $\rightarrow$  Higgs embedded in 5th component of the 5D gauge field,  $A_5$
- Gauge symmetry forbids Higgs potential at tree level  $\rightarrow$  Higgs naturally light!
- Bulk gauge group  $G \rightarrow$  broken on the branes through BCs:

$$y = 0 : G \rightarrow H_0, \quad y = L : G \rightarrow H_1$$

- At low energies  $\rightarrow$  unbroken gauge group  $H = H_0 \cap H_1$
- Higgs embedded into  $G/H_0 \cap G/H_1$

# Gauge–Higgs (Grand) Unification: General Aspects



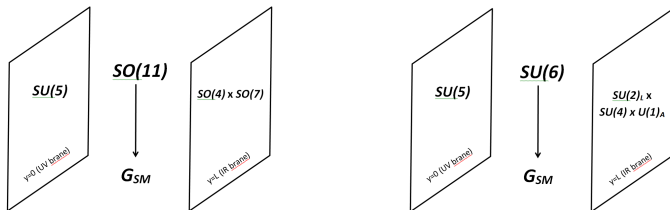
$$G \supset G_{GUT} = SU(5), SO(10), \dots$$

$$H \equiv H_0 \cap H_1 = G_{SM}$$

$$\text{Higgs} \supset G/H_0 \cap G/H_1$$

- Unification of gauge interactions with their breaking sector
- Warped geometry explains hierarchies  $\rightarrow$  gauge, flavour
- No doublet–triplet splitting problem
- Proton decay forbidden at all orders in perturbation theory!

## Existing GHGUT Models



- $SU(6)$  (mostly flat extra dimension)

SUSY models

Hall, Nomura, Smith arXiv:0107331

Burdman, Nomura arXiv:0210257

Haba, Hosotani, Kawamura, Yamashita arXiv:0401183

non- SUSY models

Lim, Maru arXiv:0706.1397

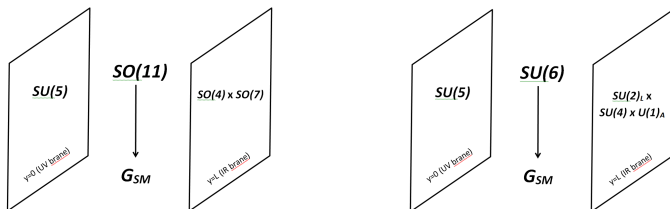
Maru, Yagatai arXiv:1903.08359 arXiv: 1911.03465

- $SO(11)$  (warped extra dimension)

Hosotani, Yamatsu arXiv: 1504.03817

Furui, Hosotani, Yamatsu arXiv: 1606.07222

## Existing GHGUT Models



Severe pheno challenges:

- too light exotic fermions due to large irreps of bigger symmetry
- difficult to obtain correct EWSB,  $m_h$  too low
- massless/degenerate SM fermions, no flavour

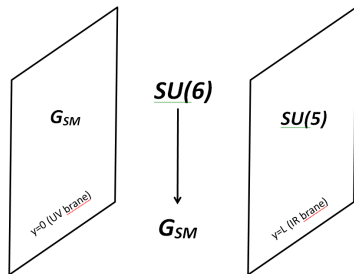
Solutions?

- go to 6D (Hosotani, Yamatsu, 1706.03503, 1710.04811)
- introduce new BSM 5D multiplets, additional brane fermions, ... (Maru, Yatagai, 1903.08359, 1911.03465)

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## Minimal SU(6) GHGUT [AA, Bally, Blasi, Goertz '21]



- Our approach  $\rightarrow$  relax breaking pattern (and  $SU(5)$  on IR brane, not on UV brane)
- No light exotics, realistic SM fermion masses + flavour
- Extra scalars besides the Higgs:  $\dim(SU(6)/SU(5)) = 35 - 24 = 11$  real scalar dof's

## Minimal SU(6) GHGUT: Gauge Sector

- Symmetry breaking pattern:

$$A_\mu^a T^a \sim \left( \begin{array}{ccc|cc|c} (++) & (++) & (++) & (--) & (--) & (--) \\ (++) & (++) & (++) & (--) & (--) & (--) \\ (++) & (++) & (++) & (--) & (--) & (--) \\ \hline (--) & (--) & (--) & (++) & (++) & (--) \\ (--) & (--) & (--) & (++) & (++) & (--) \\ \hline (--) & (--) & (--) & (--) & (--) & (--) \end{array} \right), \quad A_5^a T^a : (+ \leftrightarrow -)$$

- Branching rule for SU(6) adjoint:

$$\begin{aligned} \mathbf{35} &\rightarrow (\mathbf{8}, \mathbf{1})_0 \oplus (\mathbf{1}, \mathbf{3})_0 \oplus (\mathbf{1}, \mathbf{1})_0 \oplus [(\mathbf{3}, \mathbf{2})_{-5/6} \oplus (\mathbf{3}^*, \mathbf{2})_{5/6}] \\ &\oplus [(\mathbf{3}, \mathbf{1})_{-1/3} \oplus (\mathbf{3}^*, \mathbf{1})_{1/3}] \oplus [(\mathbf{1}, \mathbf{2})_{1/2} \oplus (\mathbf{1}, \mathbf{2})_{-1/2}] \oplus (\mathbf{1}, \mathbf{1})_0 \\ &\sim G_\mu^a \oplus W_\mu^i \oplus B_\mu \oplus V_\mu(\mathbf{3}, \mathbf{2})_{-5/6} \oplus S_1(\mathbf{3}, \mathbf{1})_{-1/3} \oplus H \oplus s \end{aligned}$$

- Extended pNGB sector  $\rightarrow$  Higgs + leptoquark (LQ) + singlet  
 $\Rightarrow$  no doublet-triplet splitting! ( $\mathbf{5}_{SU(5)} = [\text{LQ}, \text{Higgs}]$ )
- $\mathcal{O}(m_{KK})$  masses for (1st KK mode of) GUT gauge bosons  $X, Y \supset V_\mu(\mathbf{3}, \mathbf{2})_{-5/6}$   
 $\Rightarrow$  GUT without a desert! [Pomarol '00]



## Minimal SU(6) GHGUT: Fermion Sector

- Fermion content for each generation:

$$\begin{array}{l}
 20_L \rightarrow \left\{ \begin{array}{l} 10_L \rightarrow q'(3, 2)_{1/6}^{+, -} \oplus \tilde{u}'^c(3^*, 1)_{-2/3}^{+, -} \oplus e'^c(1, 1)_1^{+, -} \\ 10_L^* \rightarrow \tilde{q}^c(3^*, 2)_{-1/6}^{+, -} \oplus u(3, 1)_{2/3}^{-, -} \oplus \tilde{e}(1, 1)_{-1}^{+, -} \end{array} \right. \\
 15_L \rightarrow \left\{ \begin{array}{l} 10_L \rightarrow q(3, 2)_{1/6}^{+, +} \oplus \tilde{u}^c(3^*, 1)_{-2/3}^{-, +} \oplus e^c(1, 1)_1^{+, +} \\ 5_L \rightarrow d'(3, 1)_{-1/3}^{-, +} \oplus l'^c(1, 2)_{1/2}^{-, +} \end{array} \right. \\
 6_L \rightarrow \left\{ \begin{array}{l} 5_L \rightarrow d(3, 1)_{-1/3}^{-, -} \oplus l^c(1, 2)_{1/2}^{-, -} \\ 1_L \rightarrow \nu^c(1, 1)_0^{+, +} \end{array} \right. \\
 1_L \rightarrow \nu'^c(1, 1)_0^{+, -}
 \end{array}
 \qquad
 \begin{array}{l}
 20 \rightarrow 10 \oplus 10^* \\
 \quad \Downarrow M_{q/e}, M_{\tilde{u}} \\
 15 \rightarrow 10 \oplus 5 \\
 \quad \Downarrow M_{d/l} \\
 6 \rightarrow 5 \oplus 1 \\
 \quad \Downarrow M_{\nu} \\
 1 \rightarrow 1
 \end{array}$$

- All SM fermion masses ( $\nu$ 's also) reproduced with  $\mathcal{O}(1)$  fermion bulk mass parameters ( $c$ 's) and brane masses ( $M$ 's)
- No light exotics
- All brane masses allowed by symmetry are **necessary**:
  - +  $M_{q/e}, M_{d/l}, M_{\nu}$  to get correct SM fermion masses w/o light exotics
  - +  $M_{\tilde{u}}$  to get correct EWSB and Higgs mass

# Proton Decay: 4D GUT vs 5D GHGUT

- Ordinary 4D  $SU(5)$  GUT  $\rightarrow$  can write Yukawa term using  $\bar{\psi}$  and  $\psi^c$   
 $\Rightarrow q_L$  and  $u_R^c$  fit in one  $SU(5)$  multiplet and EWSB generates the  $u$  mass:

$$\mathbf{10}_L \rightarrow q_L(\mathbf{3}, \mathbf{2})_{1/6} \oplus u_R^c(\mathbf{3}^*, \mathbf{1})_{-2/3} \oplus e_R^c(\mathbf{1}, \mathbf{1})_1$$

$$\mathcal{L}_{4d} \supset i \bar{\mathbf{10}}_L \gamma^\mu D_\mu \mathbf{10}_L - y \bar{\mathbf{10}}_L \mathbf{10}_L^c \mathbf{5}_H + \text{h.c.} \supset g_{SU(5)} \bar{u}_R^c \gamma^\mu V_\mu q_L - y \bar{q}_L \tilde{H} u_R + \text{h.c.}$$

- GUT  $V_\mu$  boson has diquark couplings  $\Rightarrow$  proton decay!
- GHGUT  $\rightarrow$  Higgs included in covariant derivative  $\rightarrow$  only  $\bar{\psi} H \psi$  Yukawas  
 $\Rightarrow$  no EWSB mass for  $u$  if  $q_L$  and  $u_R^c$  in the same  $SU(5)$  multiplet!

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- GUT  $V_\mu$  boson and  $SU(5)$  Higgs partner do not have diquark couplings
- In fact, baryon no.  $\rightarrow$  accidental symmetry in GHGUT  
 $\Rightarrow$  proton stable at all orders in perturbation theory!

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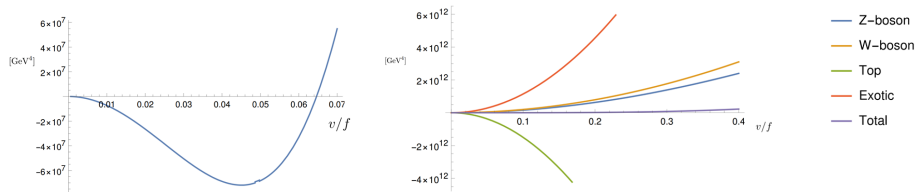
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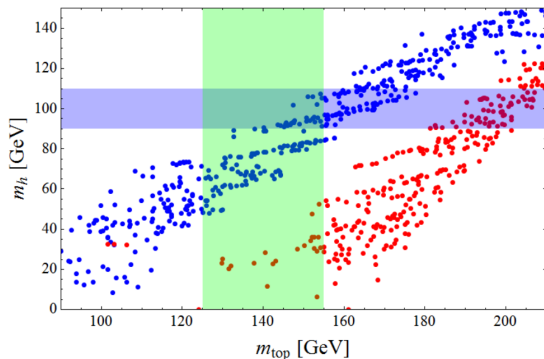
## pNGB Effective Potential

$$m_{KK} = 10 \text{ TeV}, \quad \frac{v_{SM}}{f} \simeq 0.05 \Rightarrow f \simeq 5 \text{ TeV}$$



- Exotic contribution crucial  $\rightarrow$  counterbalances top contribution
- $\mathcal{O}(0.1\%)$  fine tuning to get correct vev, but then  $m_h = \mathcal{O}(100)$  GeV w/o any extra tuning!

## Scalar Masses



$$m_{S_1}^2 \simeq \frac{3\zeta(3)(3g_s^2 - 2y_t^2)}{16\pi^2} m_{KK}^2 \simeq (0.2 m_{KK})^2 \simeq 2 \text{ TeV}, \quad m_s < 0.05 m_{KK} \simeq 500 \text{ GeV}$$

Bound on  $m_{KK}$  pushed higher by  $pp \rightarrow S_1^* S_1 \Rightarrow m_{KK} \gtrsim 7 - 8 \text{ TeV}$

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## Summary and Conclusions

- GHGUT → appealing model-building direction:
  - + unification of gauge interactions and EWSB
  - + GUTs w/o perturbative proton decay and doublet-triplet splitting problem
  - + solution to gauge and flavour hierarchies
  - + Higgs mass calculable → microscopic understanding of EWSB
  
- Viable  $SU(6)$  GUT by introducing more breaking on the boundaries:
  - + minimal fermionic content ⇒ correct SM spectrum + flavour hierarchies, no light exotics
  - + correct EWSB ⇒ correct  $m_h$  follows naturally (no extra tuning needed)
  - + new scalars →  $S_1(\mathbf{3}, \mathbf{1})_{-1/3}$  LQ w/ Yukawas predicted by  $SU(5)$  + scalar singlet
  - +  $m_{KK}$  pushed upwards by bounds on  $m_{S_1}$  ⇒ competitive with flavour constraints
  
- Lots of directions to explore!
  - baryogenesis, Higgs physics, running of gauge couplings (see Sascha Weber's talk), ...

THANK YOU FOR YOUR ATTENTION!

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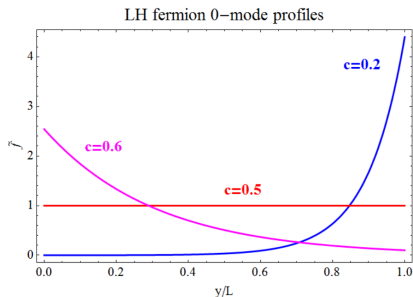
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# BACKUP

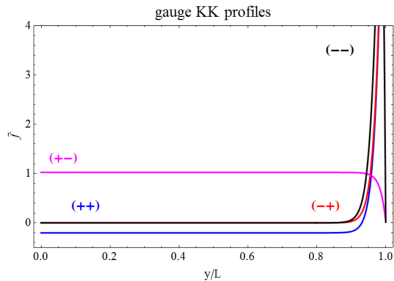
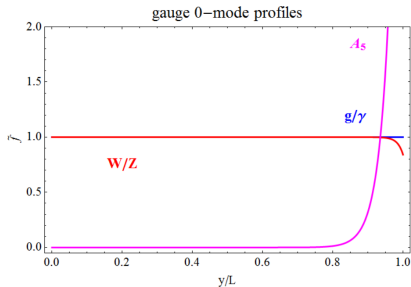
# Solving the Flavour Hierarchy Problem [Gherghetta, Pomarol '00; Grossman, Neubert '00]

- SM Yukawas  $\sim$  overlap of fermion profiles with IR-localized Higgs
- Flavour hierarchy explained geometrically ( $c \rightarrow -c$  for RH):



Exponential dependence on  $c$  of Yukawas  $\Rightarrow$   
**SM FLAVOUR HIERARCHY FROM  $\mathcal{O}(1)$  PARAMETERS!**

# Gauge Bosons in Warped Space [Gherghetta, Pomarol '00; Grossman, Neubert '00]

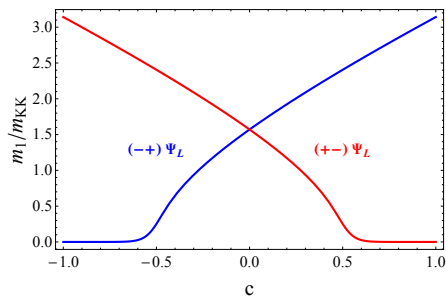


Masses of 1st KK modes of gauge bosons ( $m_{KK} = k e^{-kL}$ ):

$$m_{++} \simeq 2.45 m_{KK}, \quad m_{-+} \simeq 2.40 m_{KK}, \quad m_{+-} \simeq 0.25 m_{KK}, \quad m_{--} \simeq 3.83 m_{KK}$$

# Exotic Fermions from Warped Space

- $(-+)$  BCs forbid 0-mode, but 1st KK mass exponentially suppressed for certain values of  $c$   
 $\Rightarrow$  light exotics!



$$m_1^{(-+)} \sim e^{(c+\frac{1}{2})kL} m_{KK}, \quad c < -0.5$$

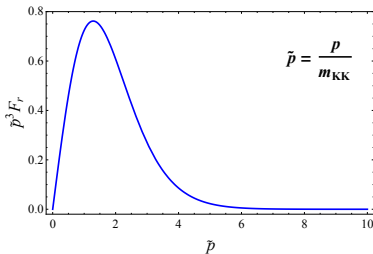
- Light exotics  $\rightarrow$  generic problem in GHGUTs b/c large fermion irreps

# Gauge–Higgs (Grand) Unification: Higgs Potential [Hosotani '83]

- $A_5$  appears through  $F_{\mu 5} \supset \partial_\mu A_5 \Rightarrow$  shift symmetry!
  - $\rightarrow$  explicitly broken by  $\bar{\Psi}\gamma_5 A_5 \Psi$  and  $f_{abc} f_{ade} \eta^{\mu\nu} A_\mu^a A_\nu^d A_5^b A_5^e$
- $\Rightarrow$  Higgs potential generated at one loop  $\rightarrow$  Hosotani mechanism
- Higgs  $\sim$  pseudo Nambu–Goldstone boson (pNGB) arising from  $SU(3) \rightarrow G_{EW}$ ;
- Typical contribution (of species  $r$ ) to the Higgs potential:

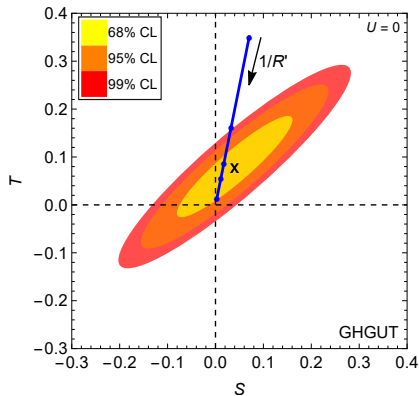
$$V_r(h) = \frac{N_r}{16\pi^2} \int_0^\infty dp p^3 \log \left[ 1 + F_r(p^2) \sin^2 \left( \frac{h}{f} \right) \right], \quad f = \frac{2}{g_{5D}^2 k} m_{KK} = \frac{m_{KK}}{g_*}$$

- Analogy with QCD:  $f \leftrightarrow f_\pi \simeq 130$  MeV,  $m_{KK} = g_* f \leftrightarrow m_\rho \simeq 770$  MeV



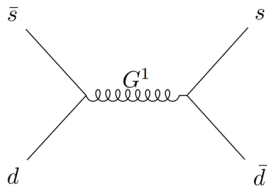
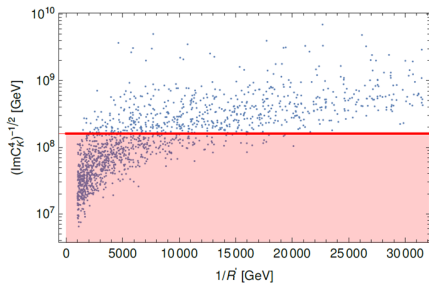
## Electroweak Precision Tests

$$\begin{cases} S = \frac{3\pi}{2} \frac{v^2}{m_{KK}^2} \simeq 0.003 \\ T = \frac{\alpha_{em}^{-1}}{10} \frac{v^2}{f^2} \simeq 0.03 \end{cases} \Rightarrow m_{KK} = g_* f \gtrsim 5 \text{ TeV at } 2\sigma$$



## Meson Mixing

## Meson Mixing



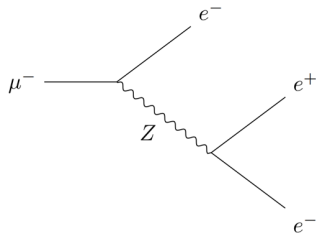
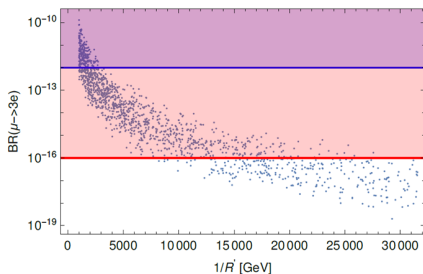
$$\mathcal{H} \supset C_4^K Q_4^{ds} \sim \left( \frac{g^{ds}}{m_{G^1}} \right)^2 (\bar{d}_L^\alpha s_{R,\alpha}) (\bar{d}_R^\beta s_{L,\beta})$$

$\mu \rightarrow eee$ 

Tree-level flavor violation:  $\mu^- \rightarrow e^- e^+ e^-$

SINDRUM (1988)  $\text{Br}(\mu \rightarrow 3e) < 10^{-12}$

Mu3e (2026(?))  $\text{Br}(\mu \rightarrow 3e) < 10^{-16}$



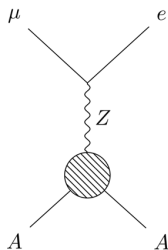
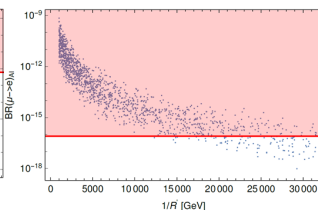
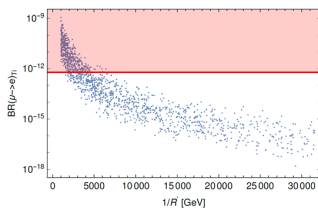


# $\mu \rightarrow e$ Conversion in Nuclei

Tree-level flavor violation:  $\mu \rightarrow e$  conversion

SINDRUM II (1998)  $\text{Br}(\mu \rightarrow e)_{\text{Ti}} < 6.1 \times 10^{-13}$

Mu2e (20??)  $\text{Br}(\mu \rightarrow e)_{\text{Al}} < 8 \times 10^{-17}$



$\mu \rightarrow e\gamma$ 

- Random scan over  $m_{KK}$ , brane mass matrices:

$$0.2 < \left| (M_{q/e})_{ij} \right|, \left| (M_{d/l})_{ij} \right| < 2$$

- LQ contribution dominates  $\mu \rightarrow e\gamma$ !
- $m_{KK} \gtrsim 7$  TeV; other onstraints less stringent

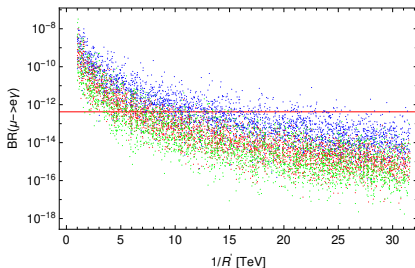
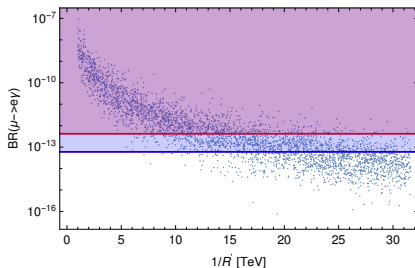
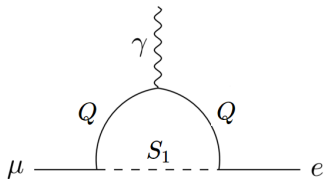


Figure: Left: current (red) and future (blue) constraints (MEG) on  $\mu \rightarrow e\gamma$ .

Right: breakdown of contributions to  $\mu \rightarrow e\gamma$ : LQ (blue), Higgs (red), Z (green).

## Scalar Leptoquark Yukawas

$$q_{L_i} = \begin{pmatrix} V^\dagger u_L \\ d_L \end{pmatrix}_i, \quad \ell_{L_i}^c = \begin{pmatrix} e_L^c \\ -U^* \nu_L^c \end{pmatrix}_i, \quad V \equiv V_{CKM}, \quad U \equiv U_{PMNS}, \quad W \equiv U_{d_R} U_{e_L}^T$$

- Higgs + scalar LQ come from the same  $SU(5)$  multiplet  $\Rightarrow$  Yukawas highly correlated:

$$-\mathcal{L}_{S_1} = S_1 \left\{ \bar{q}_{L_i} [y_d W]_{ij} \ell_{L_j}^c + \bar{d}_{R_i} [W U^* y_\nu]_{ij} \nu_{R_j}^c + \bar{u}_{R_i} [y_u V y_d^{-1} W y_e]_{ij} e_{R_j}^c \right\} + \text{h.c.}$$

- $S_1$  couples almost exclusively to  $t_R \tau_R^c$ :

$$[y_u V y_d^{-1} W y_e] \sim \begin{pmatrix} 10^{-6} & 10^{-4} & 10^{-4} \\ 10^{-4} & 10^{-3} & 10^{-2} \\ 10^{-3} & 10^{-2} & 1 \end{pmatrix}$$

$$[y_d W] \sim \begin{pmatrix} 10^{-5} & 10^{-6} & 10^{-6} \\ 10^{-4} & 10^{-4} & 10^{-4} \\ 10^{-3} & 10^{-2} & 10^{-2} \end{pmatrix}$$

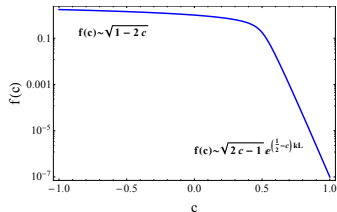
- Generates  $C_9^{\mu\mu} = -C_{10}^{\mu\mu}$  at one loop, but  $y_R^{t\mu} \sim 10^{-2}$  too small to explain  $b \rightarrow s\mu\mu$

## CKM and PMNS Matrices

- Quark mass matrices in flavour basis:

$$m_u \simeq \frac{g_* v}{2\sqrt{2}} f_{c15} M_{q/e}^\dagger f_{-c20}, \quad m_d \simeq -\frac{g_* v}{2\sqrt{2}} f_{c15} M_{d/l}^\dagger f_{-c6},$$

$$f_c = \text{diag}[f(c_1), f(c_2), f(c_3)], \quad f(c) = \sqrt{\frac{1-2c}{1-e^{(2c-1)kL}}}$$



- Unitary matrices diagonalizing  $m_{u,d}$  carry specific hierarchies [Csaki+ '08, Goertz+ '08]:

$$U_{uL}, U_{dL} \sim \begin{pmatrix} 1 & \frac{f(c_{15,1})}{f(c_{15,2})} & \frac{f(c_{15,1})}{f(c_{15,3})} \\ \frac{f(c_{15,1})}{f(c_{15,2})} & 1 & \frac{f(c_{15,2})}{f(c_{15,3})} \\ \frac{f(c_{15,1})}{f(c_{15,3})} & \frac{f(c_{15,2})}{f(c_{15,3})} & 1 \end{pmatrix}, \quad \begin{aligned} U_{uR} &\sim U_{uL} \Big|_{c_{15,i} \leftrightarrow -c_{20}} \\ U_{dR} &\sim U_{uL} \Big|_{c_{15,i} \leftrightarrow -c_6} \end{aligned}$$

- CKM inherits hierarchies from  $U_{uL}, U_{dL}$ :

$$\frac{f(c_{15,1})}{f(c_{15,2})} \sim \lambda, \quad \frac{f(c_{15,2})}{f(c_{15,3})} \sim \lambda^2 \Rightarrow V_{CKM} \equiv U_{uL} U_{dL}^\dagger \sim \begin{pmatrix} 1 & \lambda & \lambda^3 \\ \lambda & 1 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

## CKM and PMNS Matrices

- What about the PMNS matrix?

$$U_{PMNS} \equiv U_{eL} U_{\nu L}^\dagger, \quad U_{eL}, U_{\nu L} \sim \begin{pmatrix} 1 & \frac{f(-c_{6,1})}{f(-c_{6,2})} & \frac{f(-c_{6,1})}{f(-c_{6,3})} \\ \frac{f(-c_{6,1})}{f(-c_{6,2})} & 1 & \frac{f(-c_{6,2})}{f(-c_{6,3})} \\ \frac{f(-c_{6,1})}{f(-c_{6,3})} & \frac{f(-c_{6,2})}{f(-c_{6,3})} & 1 \end{pmatrix}$$

- All relevant  $c$ -parameters fixed [Csaki+ '08, Goertz+ '08]

$$\frac{m_c}{m_t} \sim \lambda^2 \frac{f(-c_{20,2})}{f(-c_{20,3})}, \quad \frac{m_u}{m_t} \sim \lambda^3 \frac{f(-c_{20,1})}{f(-c_{20,3})}, \quad \frac{m_b}{m_t} \sim \frac{f(-c_{6,3})}{f(-c_{20,3})},$$

$$\frac{m_s}{m_b} \sim \lambda^2 \frac{f(-c_{6,2})}{f(-c_{6,3})}, \quad \frac{m_d}{m_b} \sim \lambda^3 \frac{f(-c_{6,1})}{f(-c_{6,3})}$$

- $\Rightarrow$  no large hierarchies in  $U_{eL}, U_{\nu L}$ :

$$\frac{f(-c_{6,1})}{f(-c_{6,2})} \sim 0.25, \quad \frac{f(-c_{6,2})}{f(-c_{6,3})} \sim 0.5$$

- Fix fermion masses and  $V_{CKM} \Rightarrow$  no large hierarchy in  $U_{PMNS}$ !