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FCT
Fundação para a Ciência e a Tecnologia

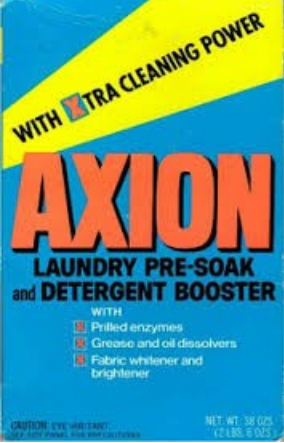
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Evaporating primordial black holes, the string axiverse, and hot dark radiation

Marco Calzà, John March-Russell, and João G. Rosa.

The goal

A new way to probe the total number of ALPs
with $m < \text{few MeV}$ through the spin distribution
of PBHs that are evaporating today!



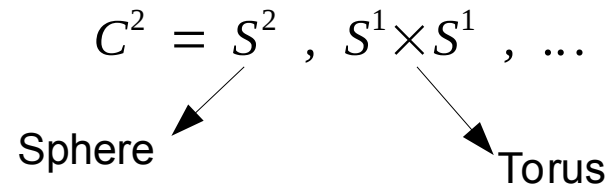
The String Axiverse

Scalar field with a shift symmetry in 4D
 No mass terms by perturbative effects
 Mass is generated by non-perturbative effects

String theory: $E \downarrow \rightarrow SM + \mathcal{G} \rightarrow 10D$ (4 s-t + 6 compactified)

Compactified extra dimensions are the key concept for the Axiverse!

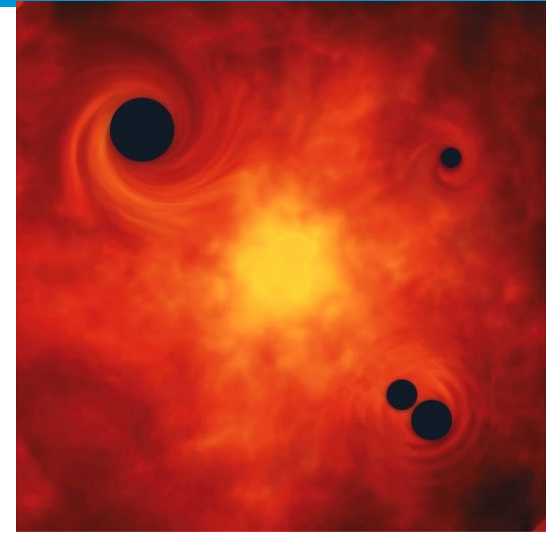
$$a(x) = \frac{1}{\mathcal{A}_{C^n}} \int_{C^n} B_n$$



6 extra d + many ways to compactify = $(N_a \sim [100-10^5])$

Primordial BH

- PBHs are BHs formed in the **early Universe**
- Through the gravitational collapse of **overdensities** in the **cosmic plasma**
- **Masses** can be several orders of magnitude **below the solar mass**



$M \sim 10^{12}$ kg evaporates enough to show changes in a_* in presence of many scalars. ($T > \text{few MeV}$)

Radiation Dominated Era \rightarrow small spin (percent level).

Early Matter Dominated Era \rightarrow Nearly extremal

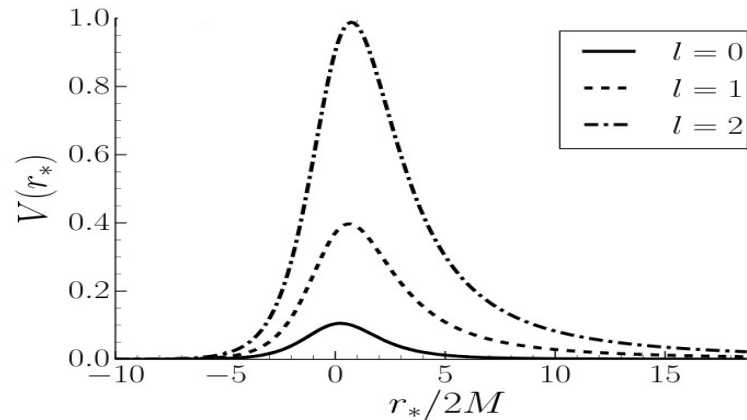
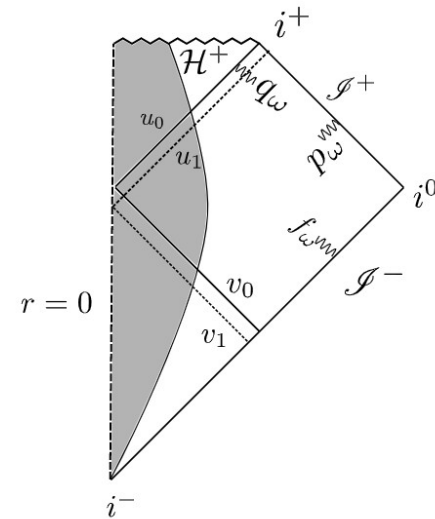
BH evaporation

Spacetime before and after the formation of an horizon

(Hawking 1975)

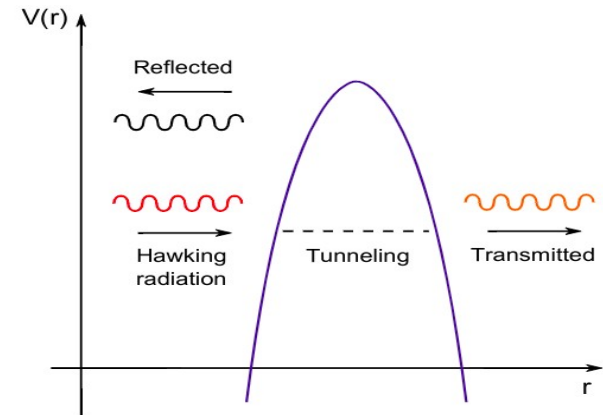
$$\text{In a (1+1)D s-t: } n_\omega = \frac{1}{\left(e^{\frac{2\omega\pi}{\kappa}} - 1\right)}, \quad T_H = \frac{\kappa}{2\pi}$$

$$\text{In a 4D s-t: } \nabla^\mu \nabla_\mu \Phi = 0 \Rightarrow \dots \Rightarrow \left(\frac{d^2}{dx^2} + \omega^2 - V(r) \right) \psi(r) = 0$$



BH geometry acts as a potential barrier that filters Hawking radiation.

$$n_\omega = \frac{\Gamma(\omega)}{\left(e^{\frac{\omega}{T_H}} - 1\right)}, \quad T_H = \frac{\kappa}{2\pi}$$



BH evaporation

Radial Teukolsky Equation:

$$\Delta^{-s} \frac{d}{dr} \left(\Delta^{s+1} \frac{dR}{dr} \right) + \left(\frac{K^2 - 2is(r-M)K}{\Delta} + 4is\omega r - \lambda \right) R = 0 \longrightarrow \Gamma_{l,m}^s(\omega)$$

$\forall \exists$ field & mode: $n_{l,m}^s(\omega) = \frac{\Gamma_{l,m}^s(\omega)}{\left(e^{\frac{2k\pi}{\kappa}} - 1 \right)}$ $\kappa = \frac{\sqrt{1-a_*^2}}{2} r_+$ $k = \omega - m\Omega_H$

Page ('74-'76)
Hiscock et al ('98)

Mass & angular momentum: $\begin{pmatrix} f_s \\ g_s \end{pmatrix} = \frac{1}{2\pi} \sum_{p,l,m} \int_0^\infty \frac{\Gamma_{l,m}^s(\omega)}{\left(e^{\frac{2k\pi}{\kappa}} \pm 1 \right)} \begin{pmatrix} x \\ m a_*^{-1} \end{pmatrix} dx \quad x = \omega M$

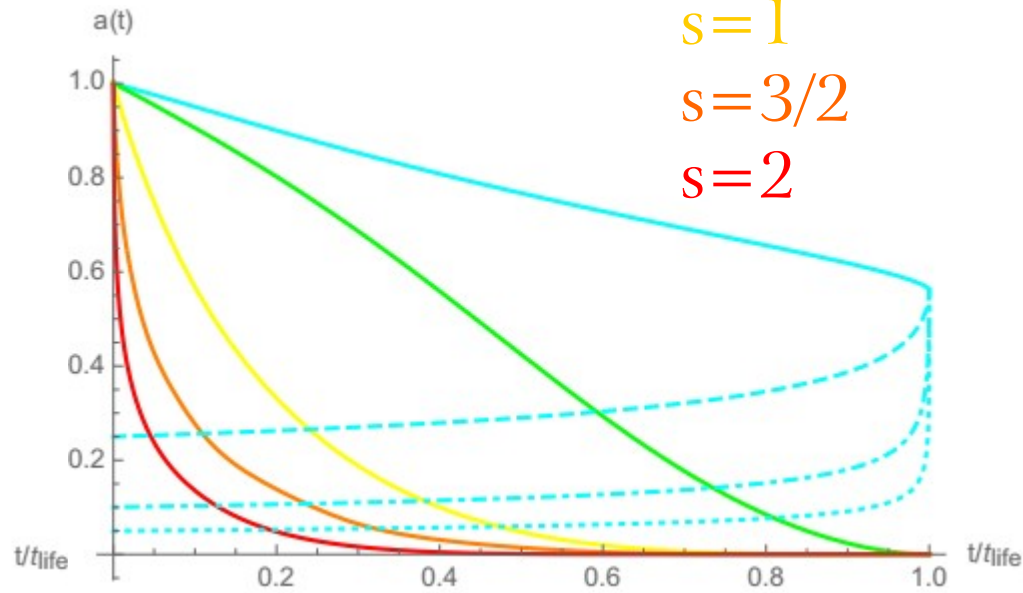
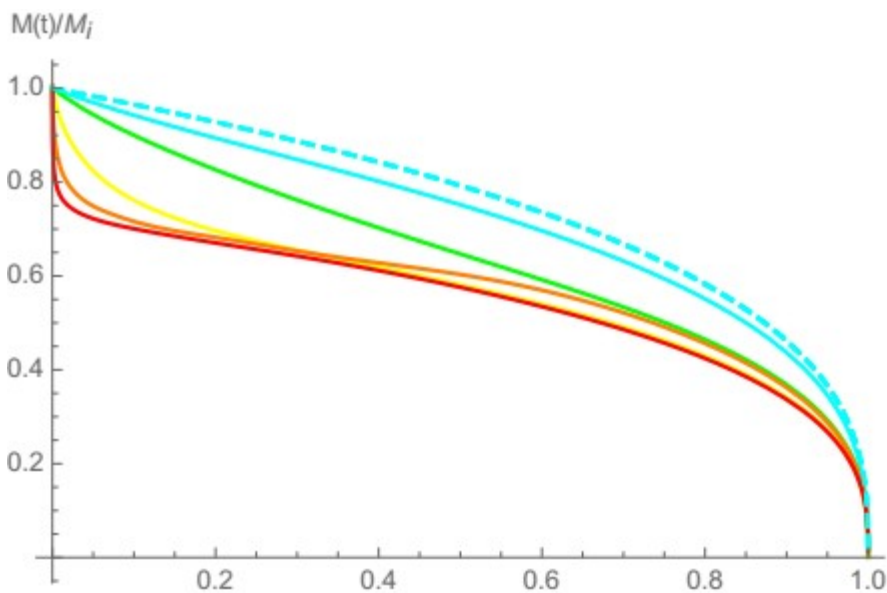
BH evaporation

$$y = -\ln[a_*] \quad z = -\ln[M/M_i] \quad \tau = -M_i^3 t$$

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$$\frac{dz}{dy} = \frac{1}{h} = \frac{f}{g-2f} \quad \frac{d\tau}{dy} = \frac{e^{3z}}{g-2f} \quad \frac{da_*}{dt} = \frac{-a_* h f}{M^3}$$

$s=0$
 $s=1/2$
 $s=1$
 $s=3/2$
 $s=2$



————— $a_{*i}=0.99$ - - - - $a_{*i}=0.25$ - · - · - $a_{*i}=0.1$ ······ $a_{*i}=0.05$

BH evaporation

A BH is **not** evaporating through **only one field!!!**

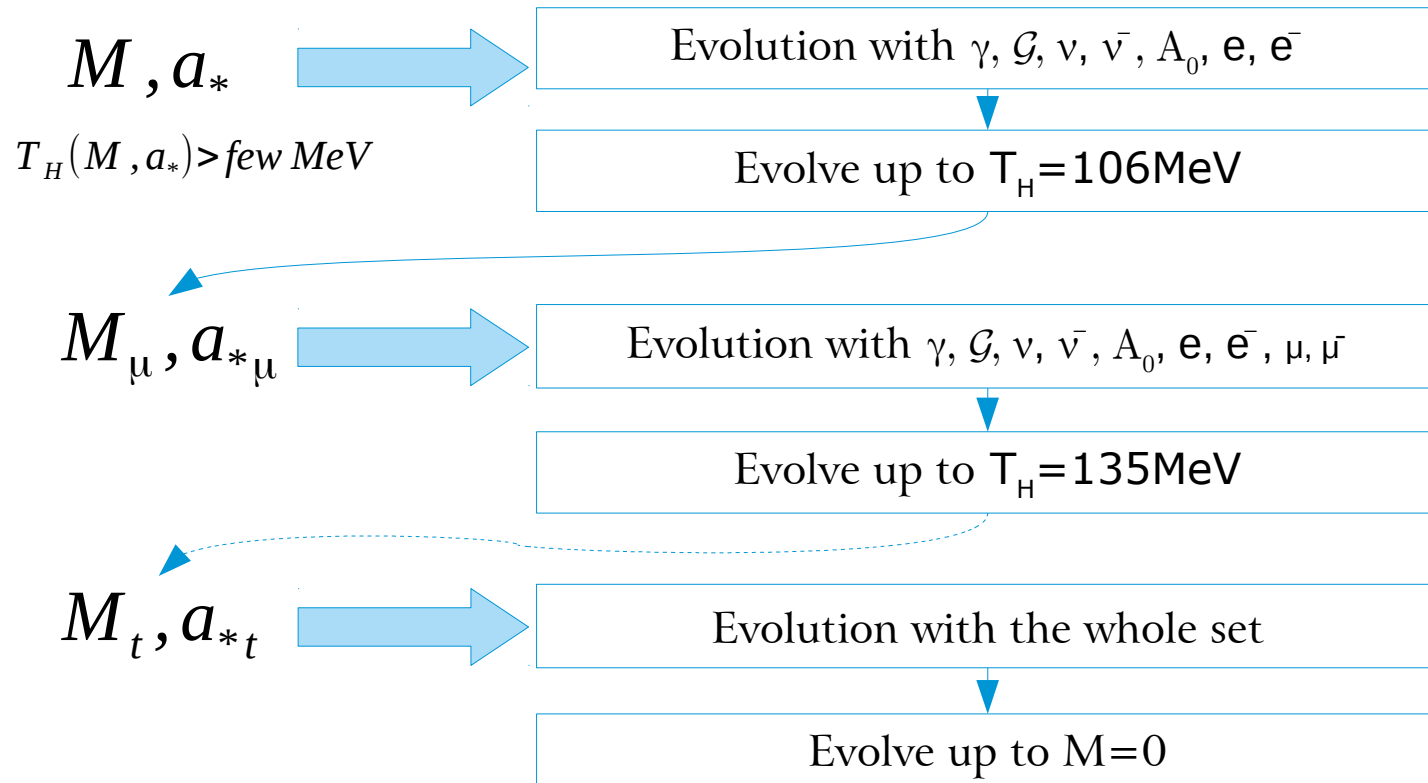
$$\begin{aligned} f_{tot} &= n_0 f_0 + n_{1/2} f_{1/2} + n_1 f_1 + n_{3/2} f_{3/2} + n_2 f_2 \\ g_{tot} &= n_0 g_0 + n_{1/2} g_{1/2} + n_1 g_1 + n_{3/2} g_{3/2} + n_2 g_2 \end{aligned} \quad \left. \vphantom{\begin{aligned} f_{tot} \\ g_{tot} \end{aligned}} \right\} \rightarrow h_{tot}$$

Evaporating BH: $M \downarrow$ & $T_H \uparrow \rightarrow$ emitted particle set **changes!!!**

Particles emission with $m > T_H$ is exponentially **suppressed**

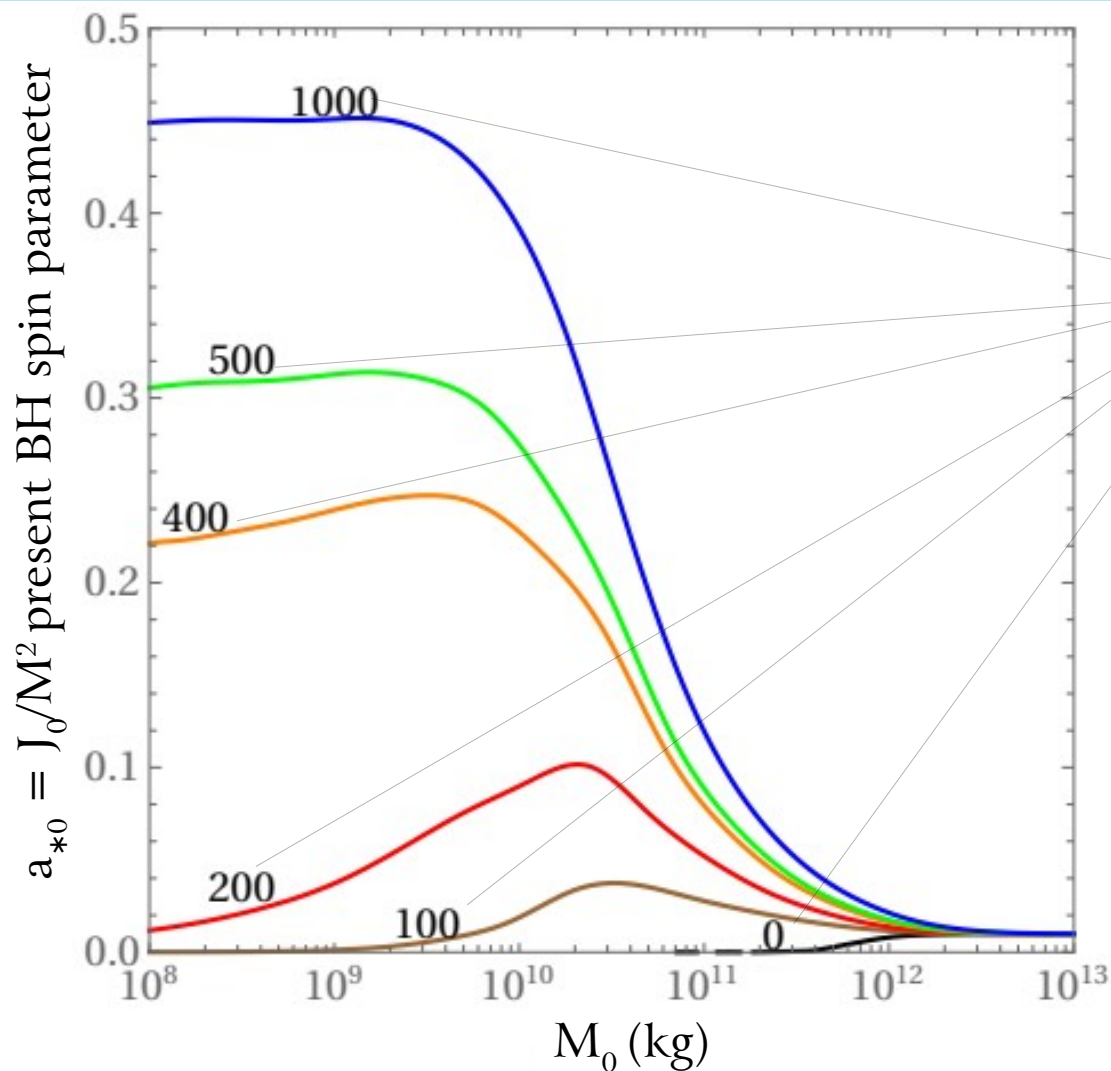
Approximation: particles are considered **massless** for $m < T_H$ and are **otherwise absent** from the emission spectrum.

Set up of our description



String Axiverse \rightarrow **N_a scalars field** in addition to the SM particles and the graviton.

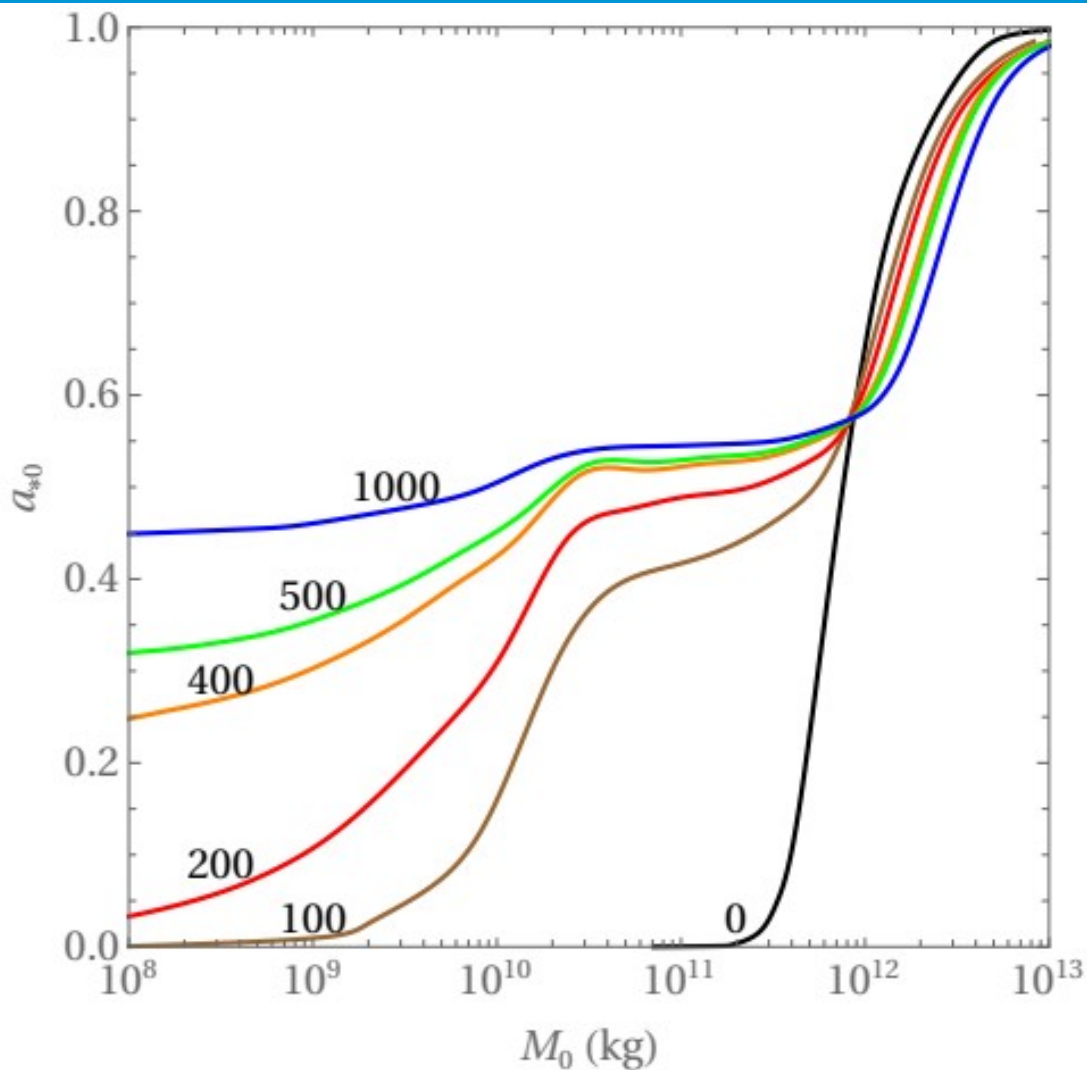
Axiverse fingerprint in PBHs evaporation



Number of ALPs in addition to the SM + \mathcal{G}

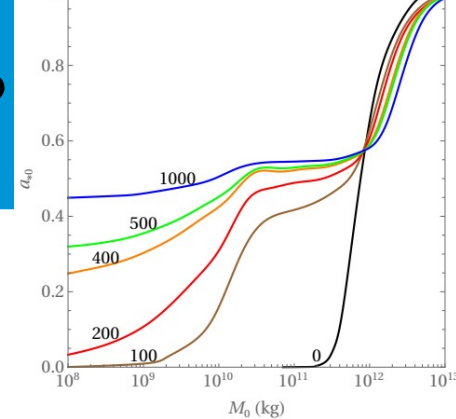
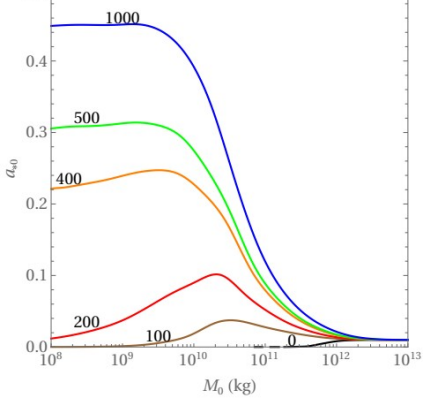
Present PBH spin, a_{*0} , as a function of their present mass, M_0 , for an initial population with spin $a_* = 0.01$ and varying mass. Curves labeled by number of light ALPs.

Axiverse fingerprint in PBHs evaporation



Present PBH spin, a_{*0} , as a function of their present mass, M_0 , for an initial population with spin $a_* = 0.99$ and varying mass. Curves labeled by number of light ALPs.

Why is this so interesting?



ALPs \rightarrow cosmological and astrophysical effects \rightarrow signatures of individual axions (mass ranges), not of the whole ‘string axiverse’.

The PBH spin distribution from evaporation process in the presence of many light scalar fields cannot, to our knowledge, be mimicked by other processes \rightarrow unique signature of an underlying theory with a large number of light scalars.

Hot Dark Radiation

Integrated flux of ALPs from a single PBH in the relevant PBH spin range

$$\sim 3 \times 10^{22} N_a (10^{10} / M) s^{-1}$$

- Reasonable N_a → Hawking luminosity is ALPs dominated
- Hot (10^{10} kg → $T_H \sim 1$ GeV)
- Dark to the SM
- Not red-shifted (It is now evaporating)
- Usual constraints (BBN, CMB...) do not apply → potentially the present ALPs hot ‘dark-radiation’ $\rho > \rho_{\text{CMB}}$.

Detection of background energetic dark axions is a striking signal for both axiverse physics and the existence of Hawking evaporating PBHs.

Evaporating primordial black holes, the string axiverse, and hot dark radiation

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(Dated: October 27, 2021)

We show that primordial black holes (PBHs) develop non-negligible spins through Hawking emission of the large number of axion-like particles generically present in string theory compactifications. This is because scalars can be emitted in the monopole mode ($l = 0$), where no angular momentum is removed from the BH, so a sufficiently large number of scalars can compensate for the spin-down produced by fermion, gauge boson, and graviton emission. The resulting characteristic spin distributions for 10^8 - 10^{12} kg PBHs could potentially be measured by future gamma-ray observatories, provided that the PBH abundance is not too small. This yields a unique probe of the total number of light scalars in the fundamental theory, independent of how weakly they interact with known matter. The present local energy density of hot, MeV-TeV, axions produced by this Hawking emission can possibly exceed ρ_{CMB} . Evaporation constraints on PBHs are also somewhat weakened.

Superstring theory is one of the leading candidates for a fundamental theory combining quantum gravity and

Light string axions can have a wide range of cosmological and astrophysical effects, e.g. steps in the matter

Thanks for your attention!!!



“That’s all Folks!”