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Cosmological Naturalness

02 Aug 2022 - PASCOS, Heidelberg

Naturalness, 2022?

- The EFT “Standard Model” explains almost everything we see and we do not see
- Naturalness = dimensional analysis works. Already from G. Galilei $t \sim \sqrt{l/g}$. But:

$$m_{\text{Higgs}}, \rho_{\text{vacuum}}, \theta_{\text{QCD}} \lll \text{dimensional analysis}$$

- Concrete problem, not aesthetic, if SM as an EFT, with calculable parameters
- LHC has basically made 2.5 discoveries:
 - A. Higgs boson
 - B. unnaturalness of m_{Higgs}
 - C. quasi-criticality (if SM extrapolated)

Naturalness, 2022?

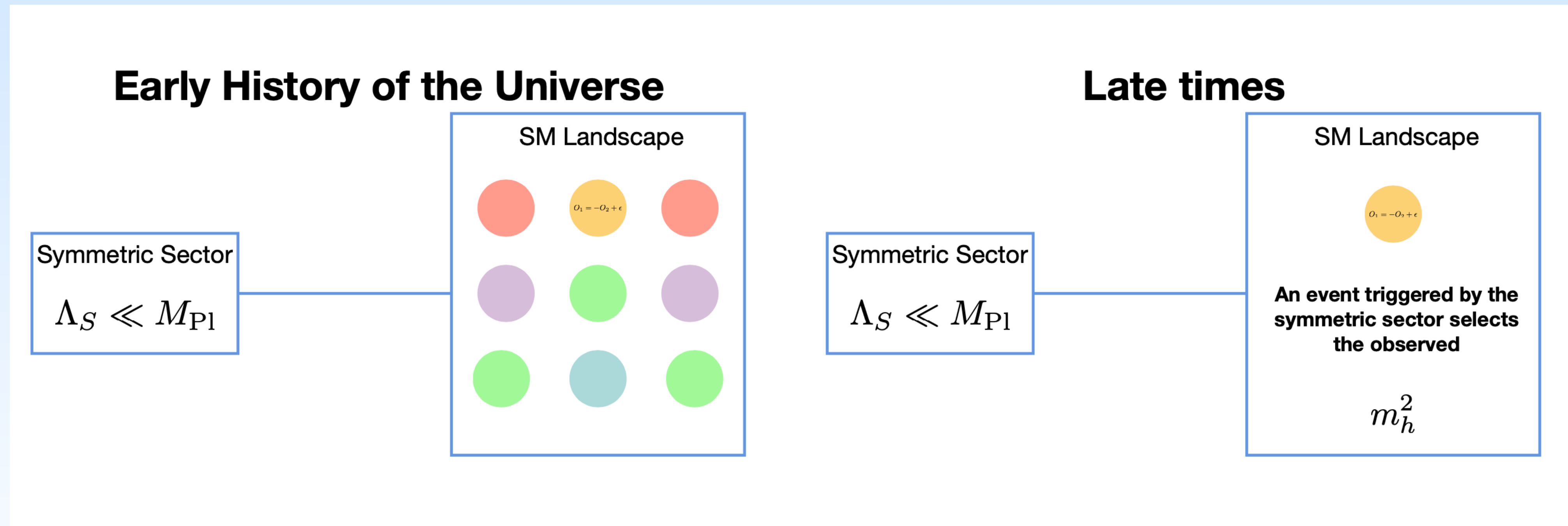
Essentially 3.5 possibilities:

- LHC will soon discover new physics related to m_{Higgs} (partial solution/non-solution)
- epicycles of low-scale SUSY or compositeness
- new frameworks in which m_{Higgs} is natural (ideally, not currently known)
- Nature is unnatural, but unnaturalness is selected dynamically (= cosmologically)

[..., Strumia, DT 2002.02463; D'Agnolo, DT 2106.04591; 2109.13249]

Cosmological Naturalness

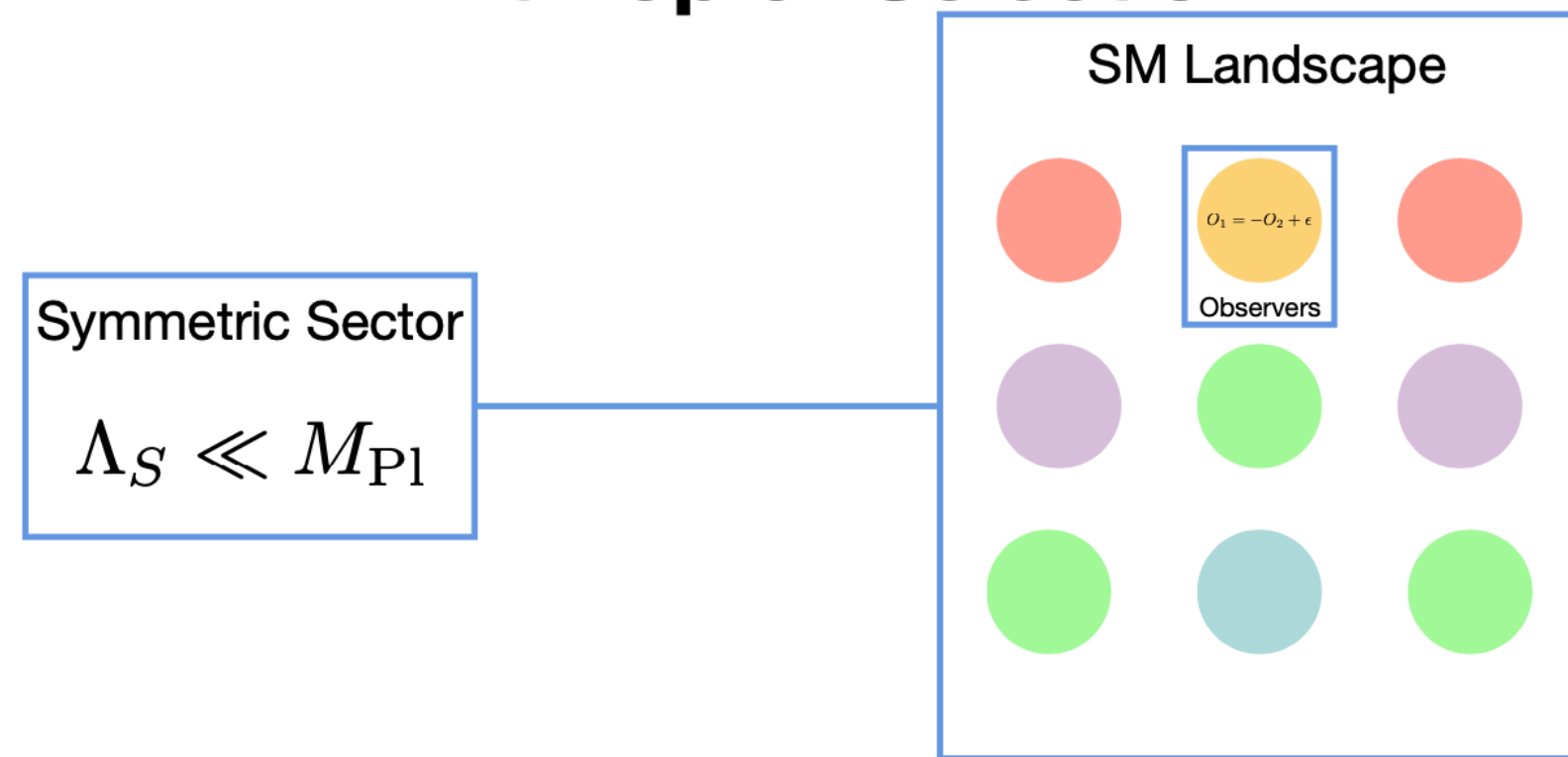
i.e. how cosmology can select a small Higgs mass



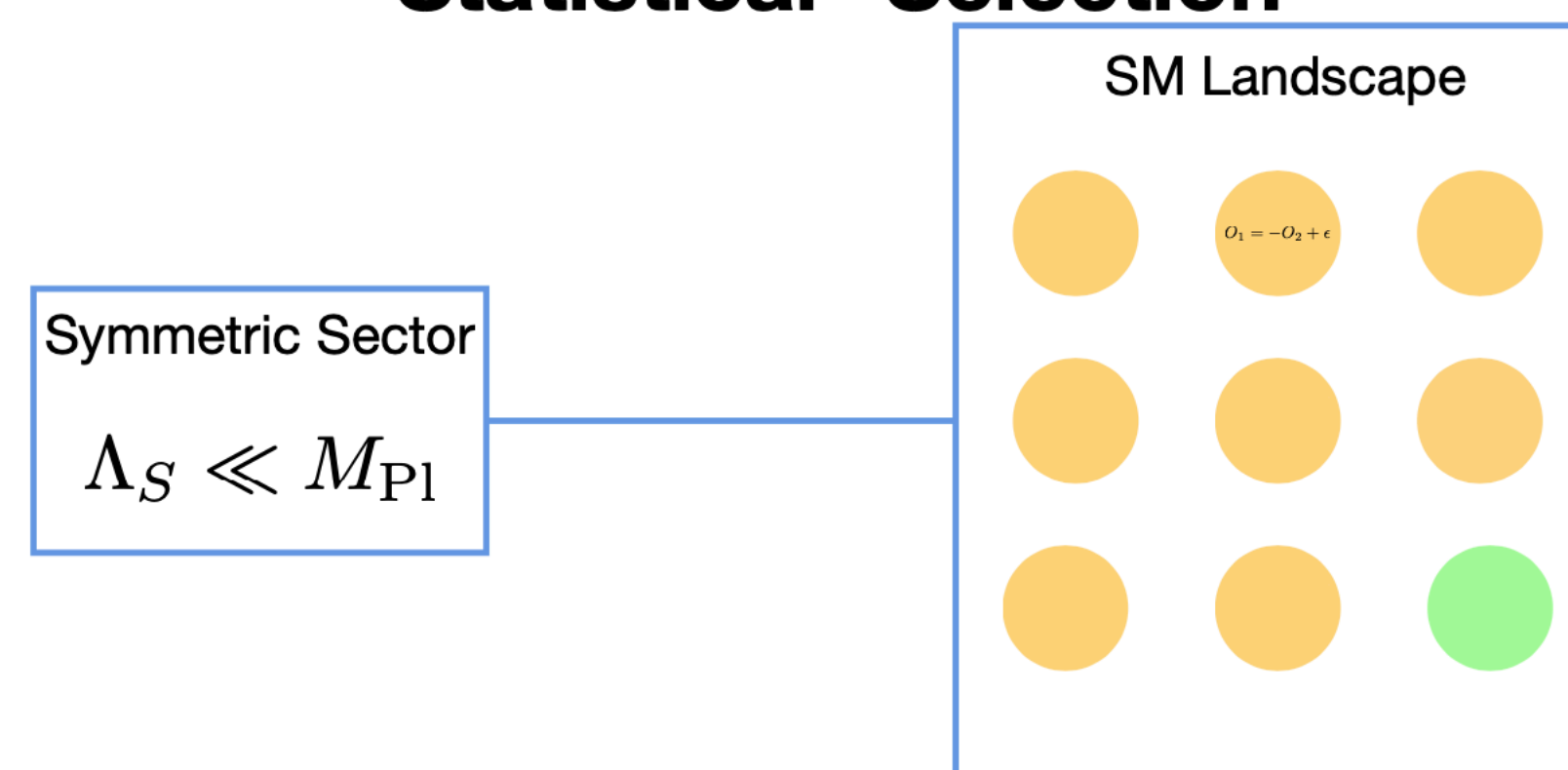
Cosmological Naturalness

i.e. how cosmology can select a small Higgs mass

“Anthropic” Selection



“Statistical” Selection



vacuum accumulation
[Dvali, Vilenkin]

light Higgs inflates most
[Geller, Hochberg, Kuflik; Cheung, Saraswat]

Self-Organized Localization
[Giudice, McCullough, You]

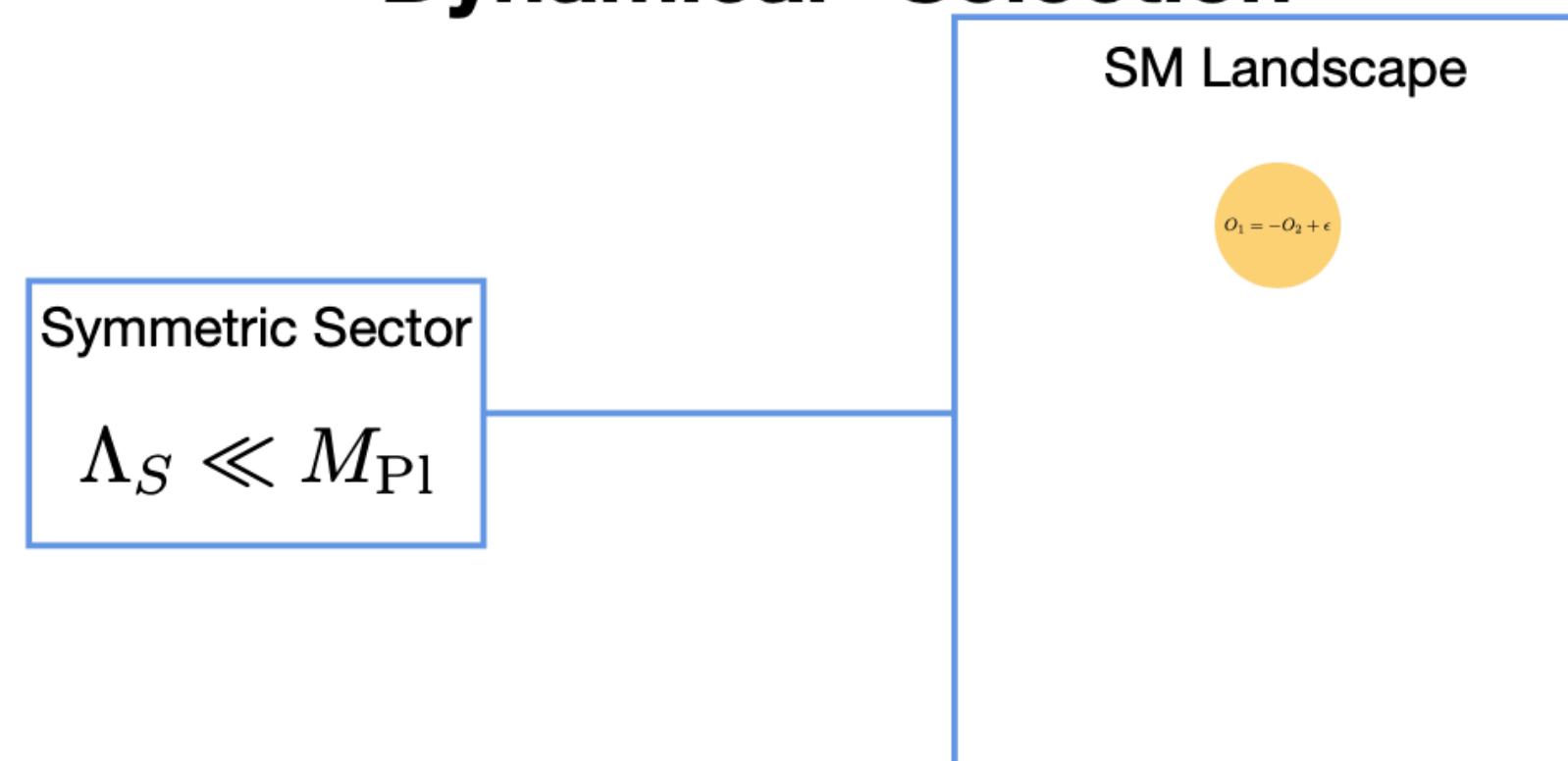
light Higgs from small CC
[Arvanitaki, Dimopoulos, Gorbenko, Huang, Van Tilburg]

relaxion
[Graham, Kaplan, Rajendran]

NNaturalness
[Arkani-Hamed, Cohen, D’Agnolo, Hook, Kim, Pinner]

Selfish Higgs
[Giudice, Kehagias, Riotto]

“Dynamical” Selection



Precarious Naturalness
[Strumia, DT]

crunching dilaton
[Csáki, D’Agnolo, Geller, Ismail]

Sliding Naturalness
[D’Agnolo, DT, 2106.04591, 2109.13249]

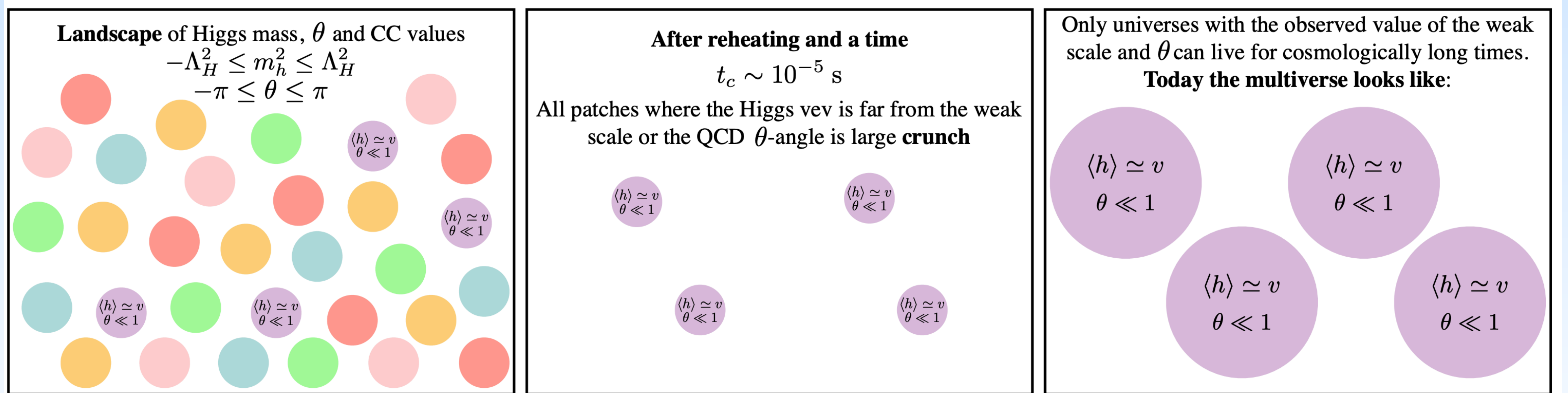
Sliding Naturalness

A novel way to select $0 < \langle h \rangle \lesssim O(100)$ GeV
and solve jointly the strong-CP problem,
explain DM, ...

R. T. D'Agnolo, D. Teresi, 2106.04591 PRL, 2109.13249 JHEP

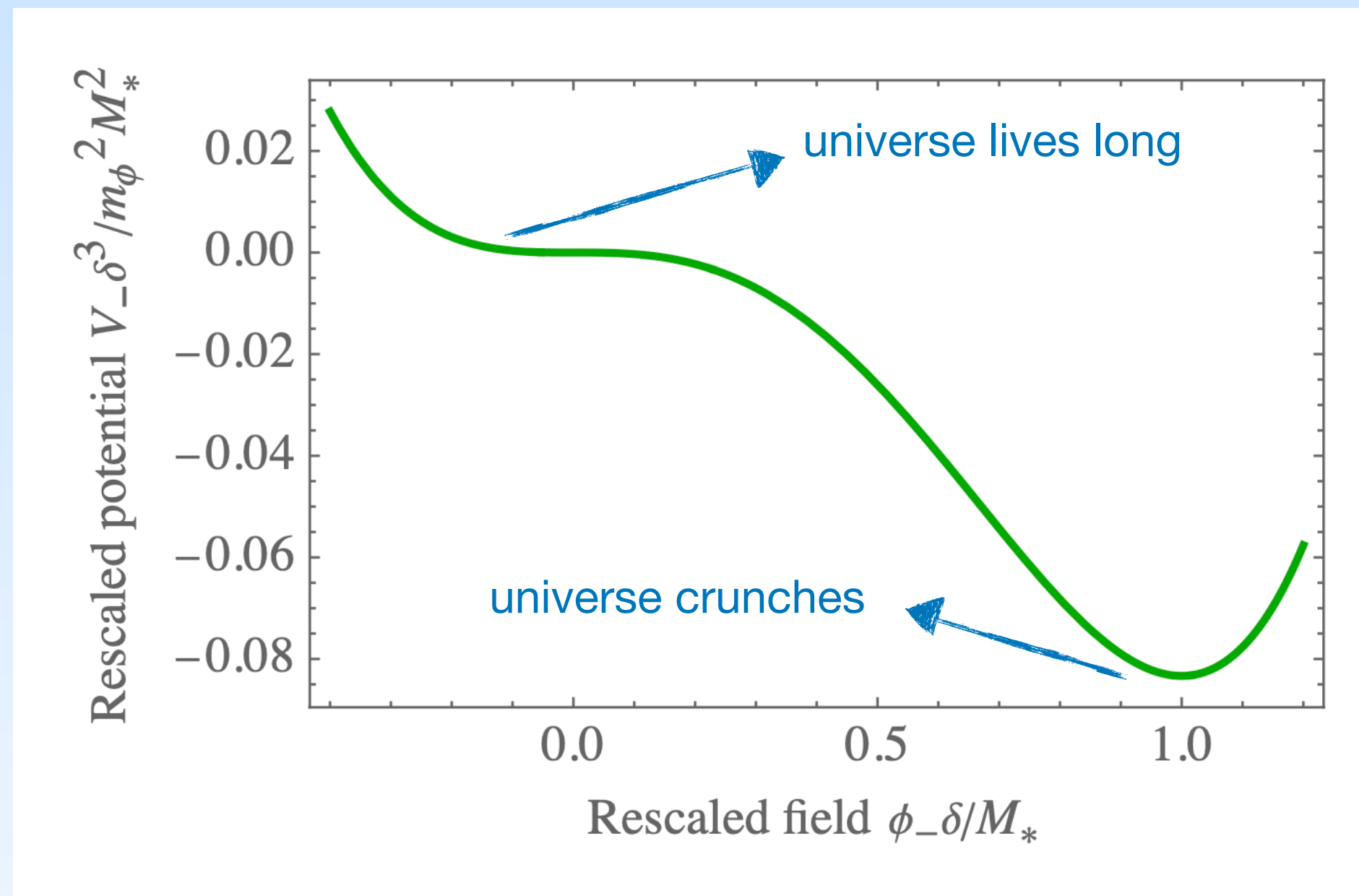
Sliding Naturalness

sketch of the mechanism



Sliding Naturalness

Two Scalars to Rule Them All

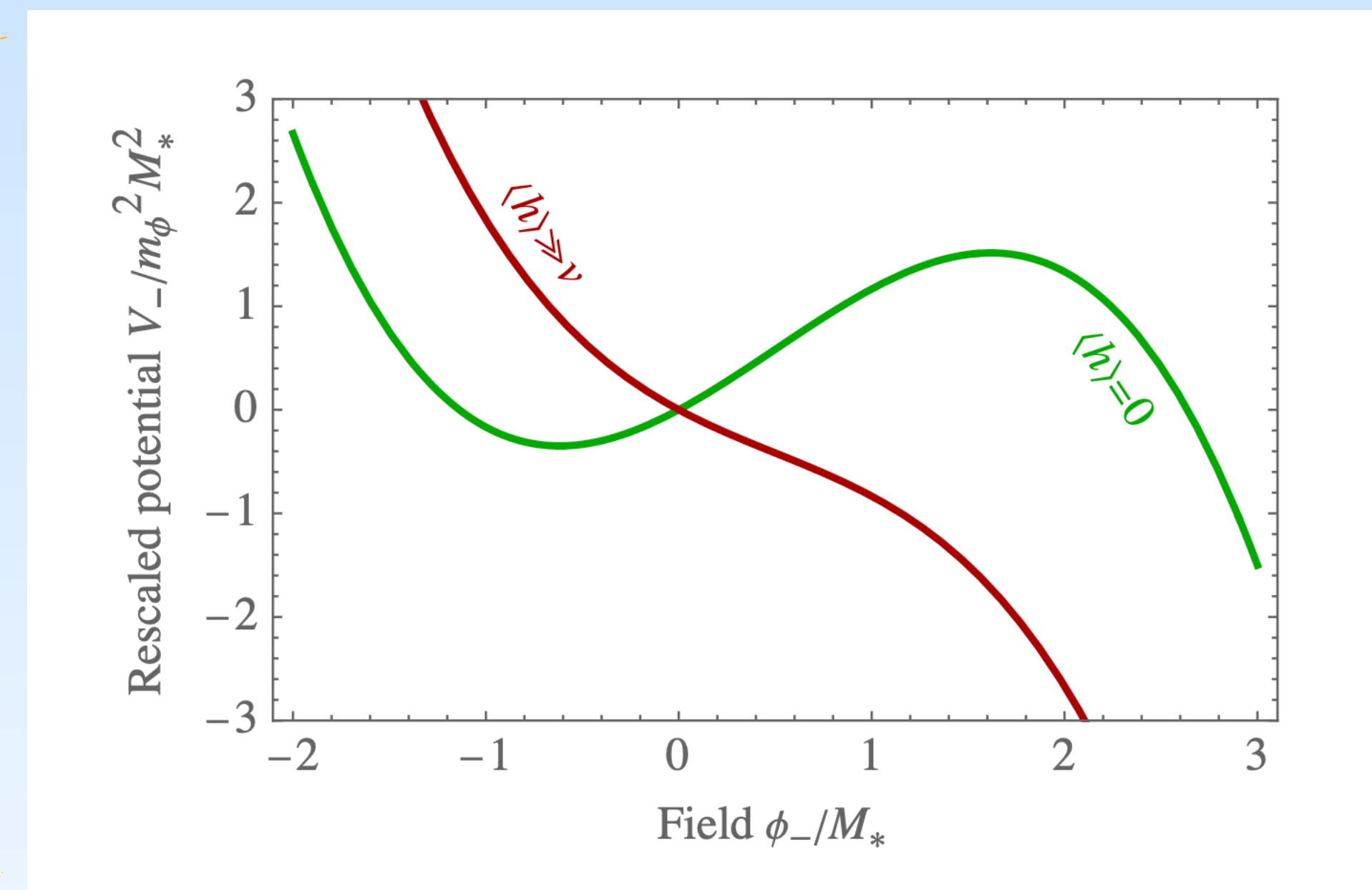
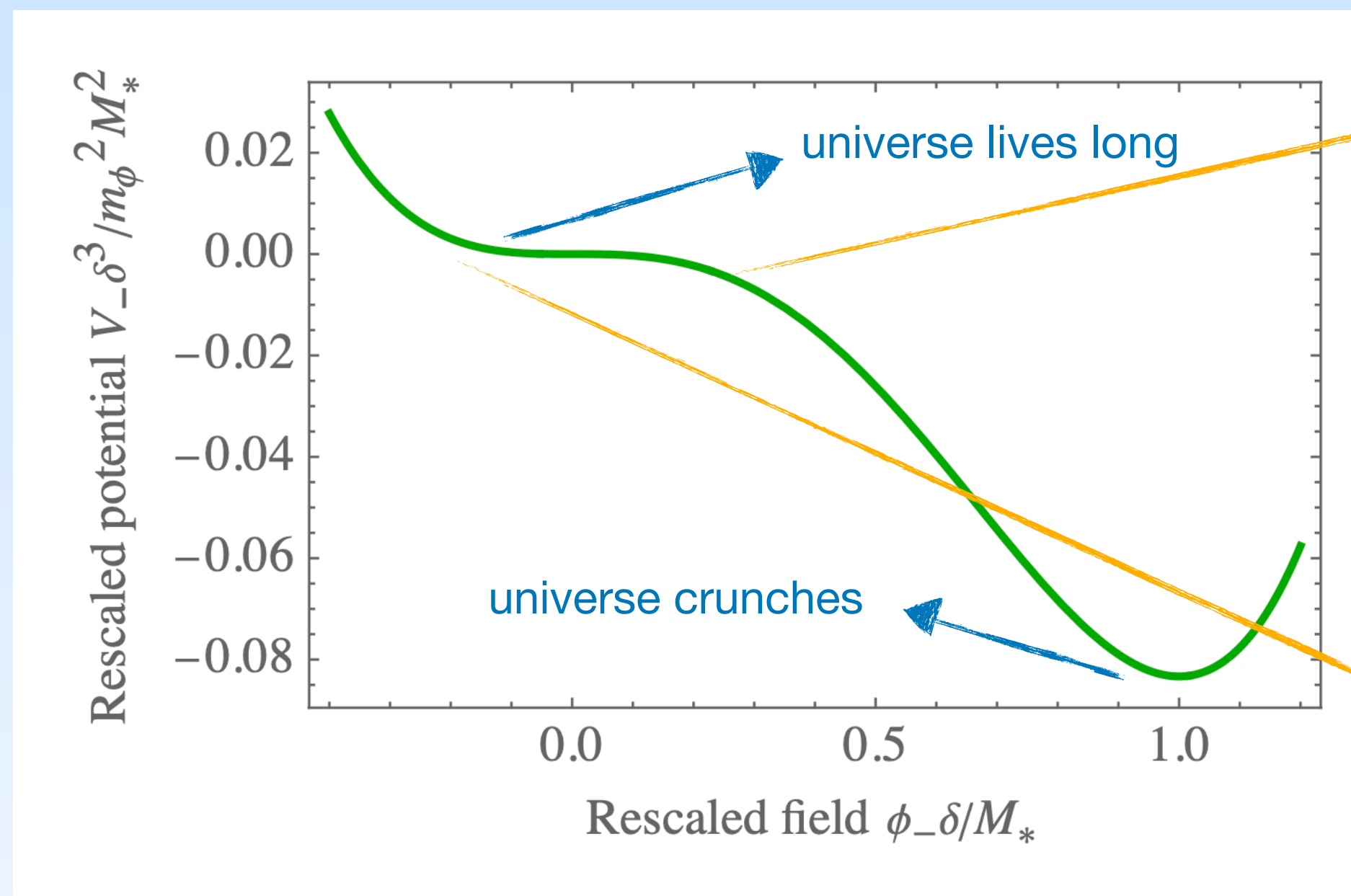


- At global minimum $V \sim -M_*^4$ too big to be compensated by a CC in the landscape
- At local minimum CC in the landscape can be tuned $V \approx 0$

Sliding Naturalness

Two Scalars to Rule Them All

universe crunches if $\langle h \rangle \gg v$

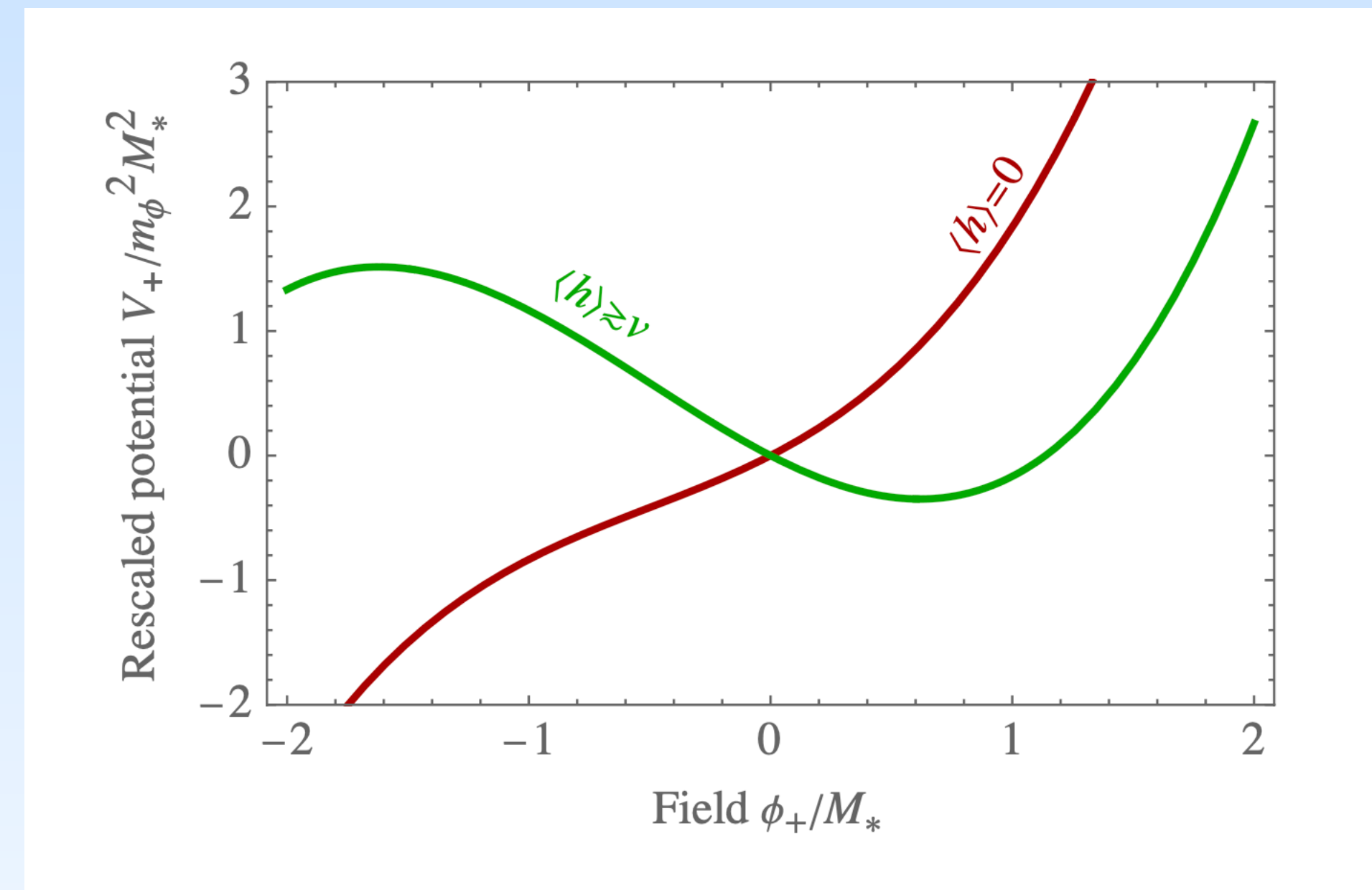
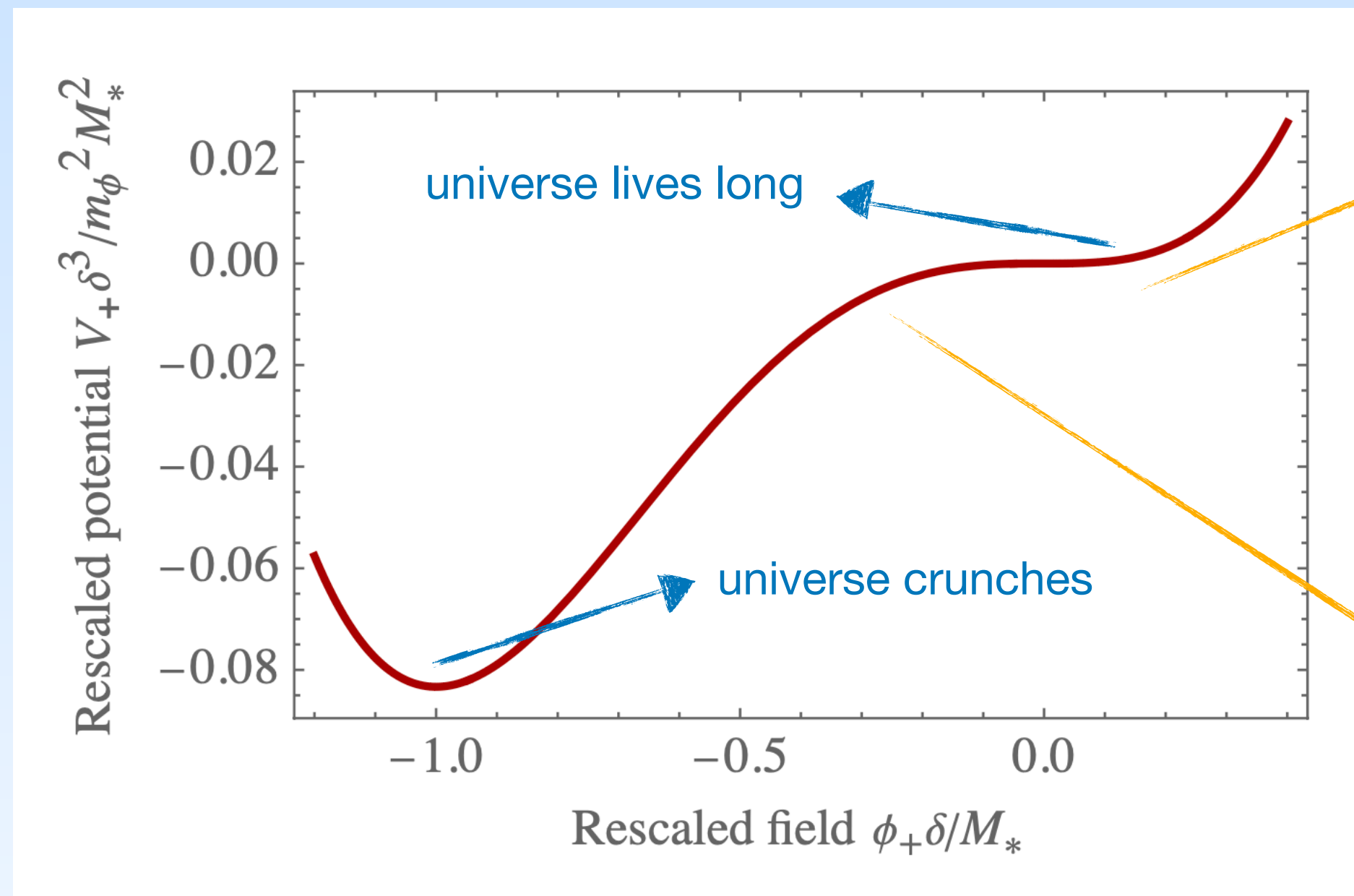


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Sliding Naturalness

Two Scalars to Rule Them All

universe crunches if $\langle h \rangle \lesssim v$



- At global minimum $V \sim -M_*^4$ too big to be compensated by a CC in the landscape
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Sliding Naturalness

Two Scalars to Rule Them All

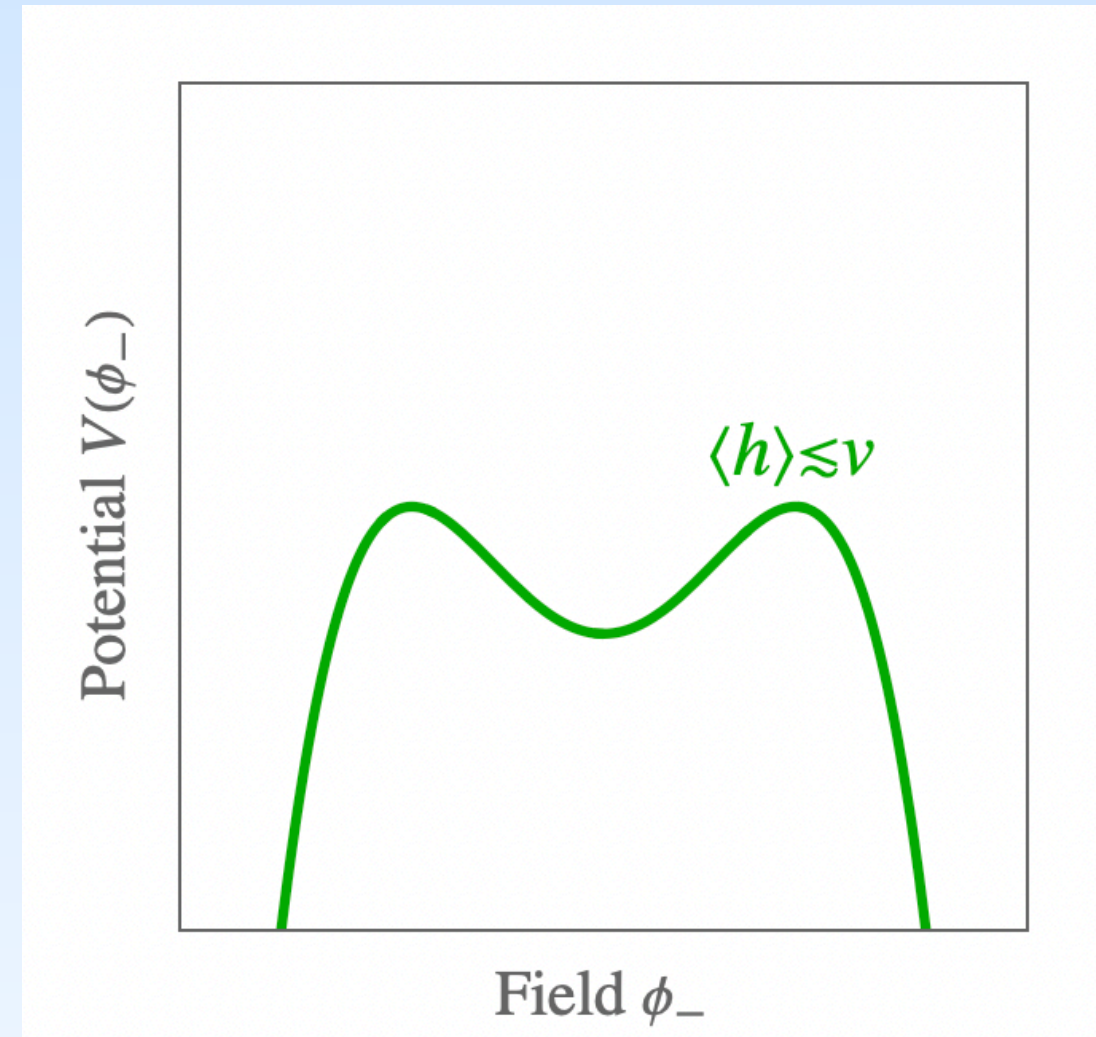
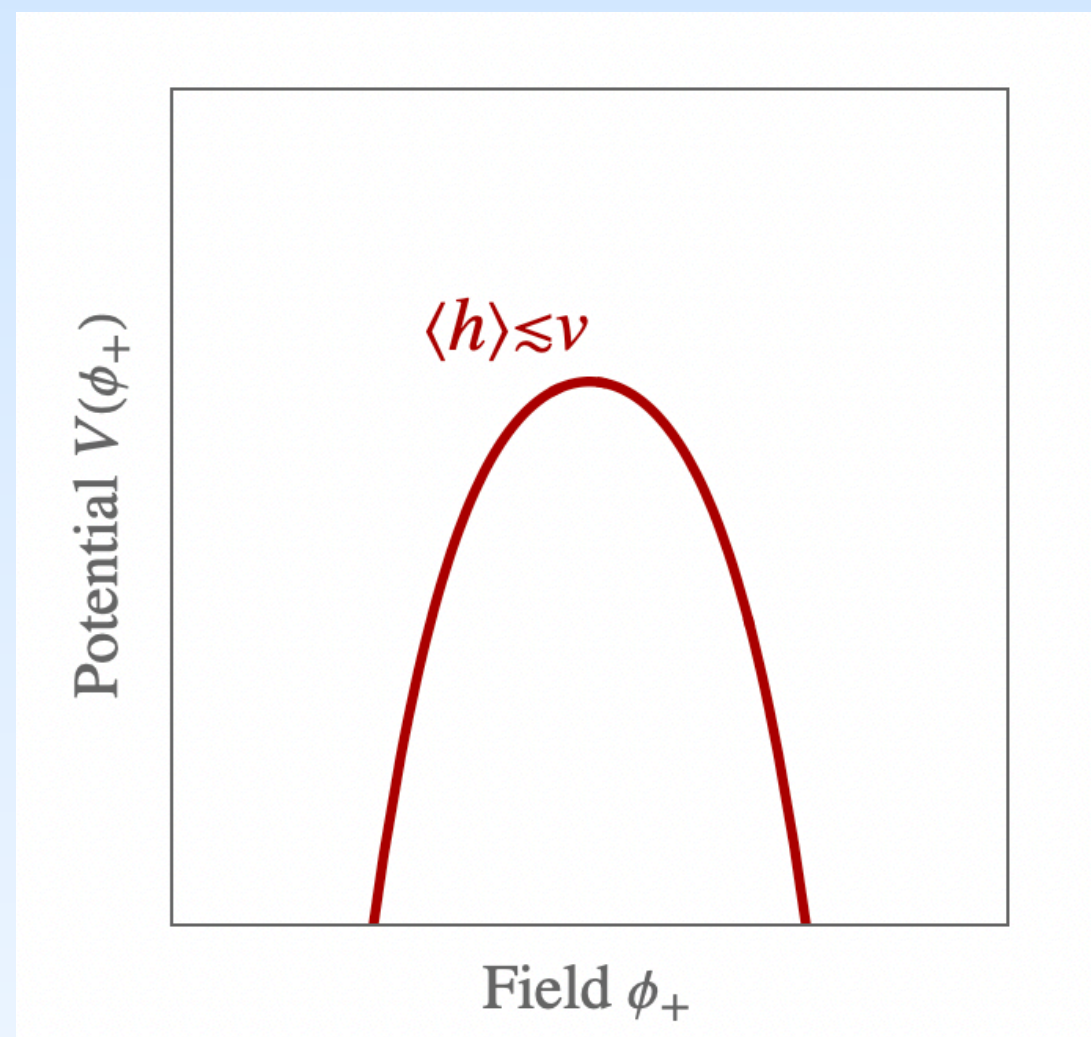
the only universes that live long
are those where

$$v \lesssim \langle h \rangle \lesssim v$$

The Lagrangian

$$V_+ = -\frac{m_+^2}{2}\phi_+^2 - \frac{m_+^2}{M_+^2}\phi_+^4$$

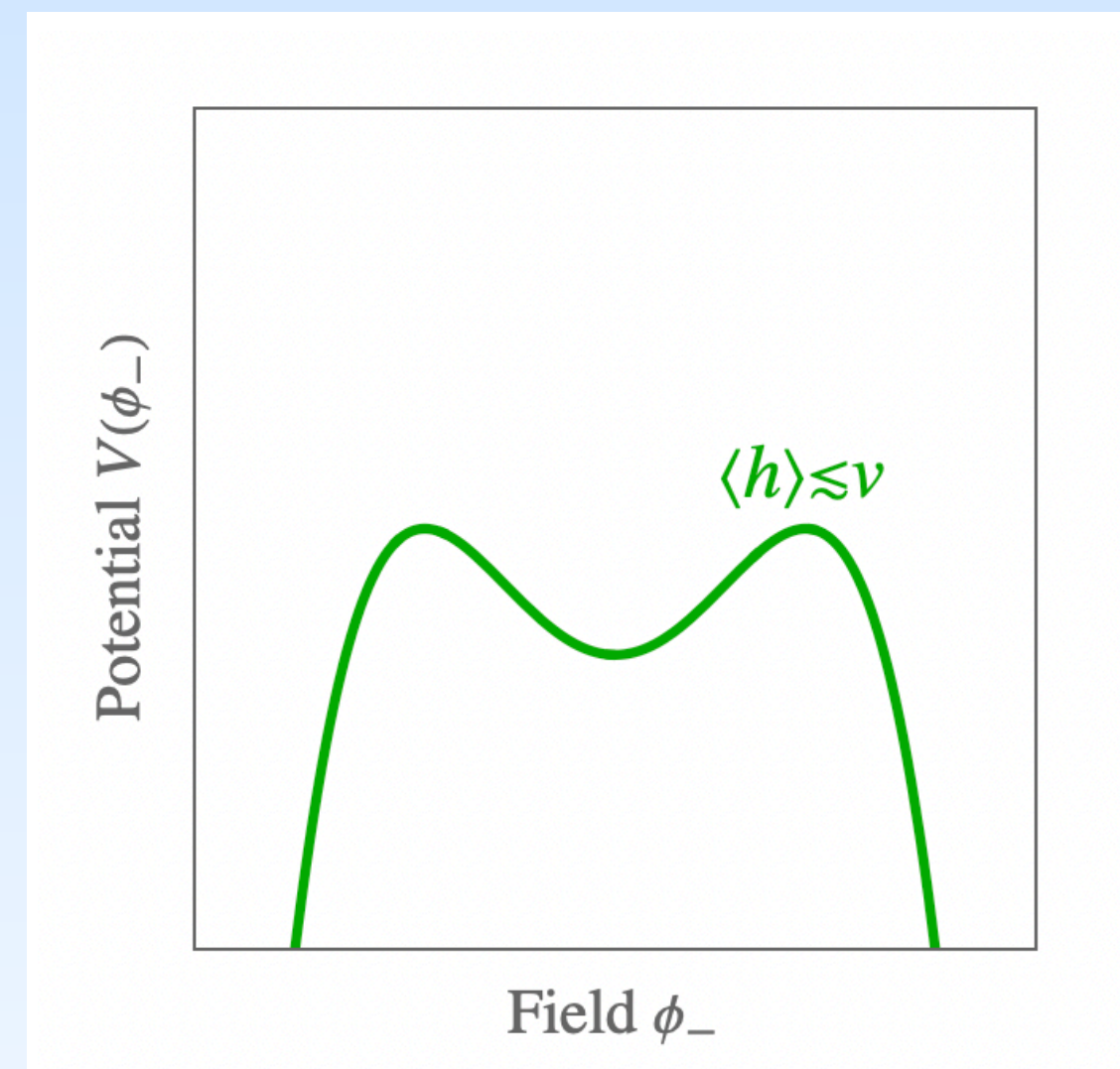
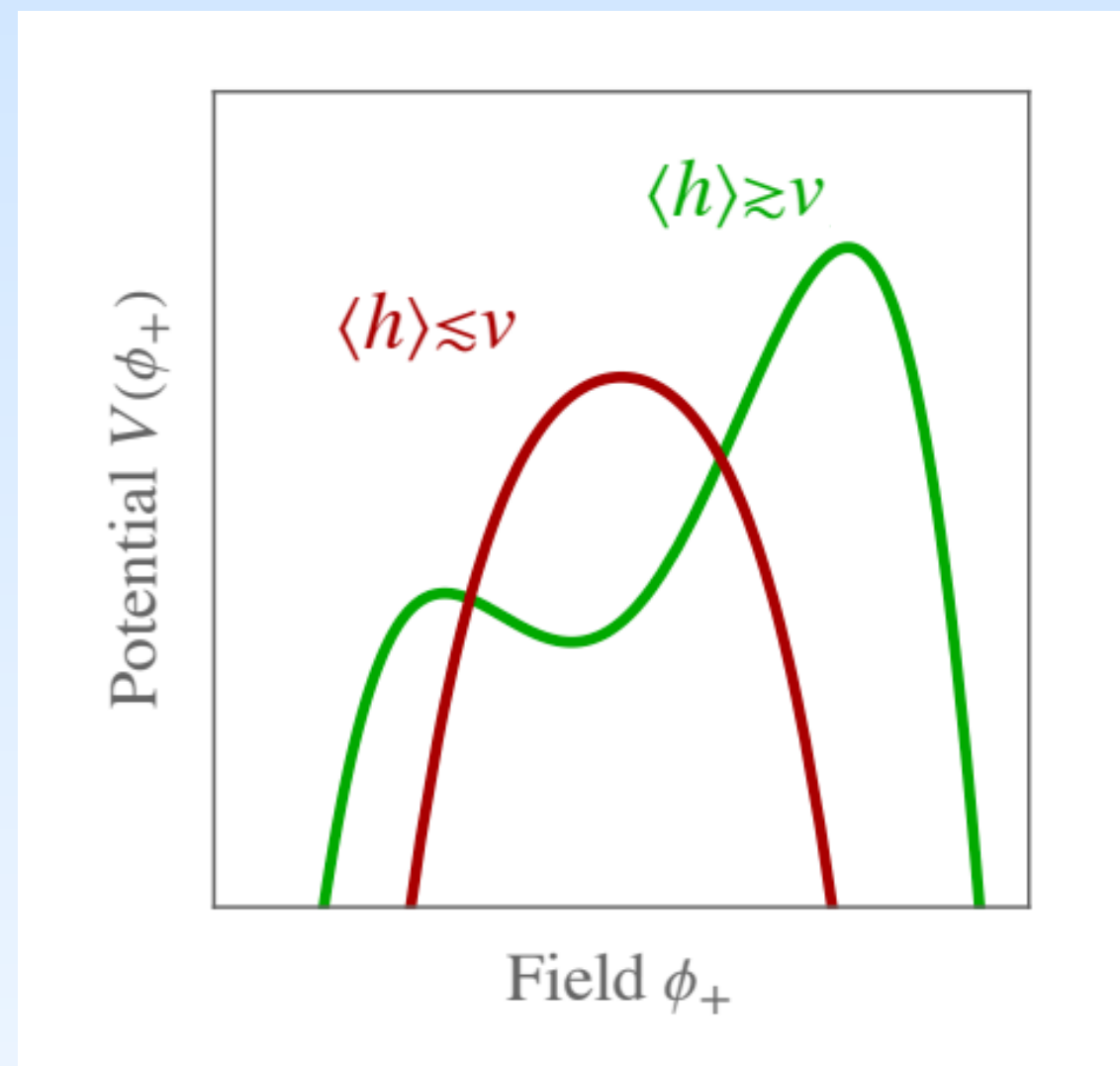
$$V_- = +\frac{m_-^2}{2}\phi_-^2 - \frac{m_-^2}{M_-^2}\phi_-^4$$



The Lagrangian

$$V_+ = -\frac{m_+^2}{2}\phi_+^2 - \frac{m_+^2}{M_+^2}\phi_+^4$$

$$V_- = +\frac{m_-^2}{2}\phi_-^2 - \frac{m_-^2}{M_-^2}\phi_-^4$$



$$V_{\phi H} = -\frac{\alpha_s}{8\pi} \left(\frac{\phi_+}{F_+} + \frac{\phi_-}{F_-} + \theta \right) G\tilde{G}$$

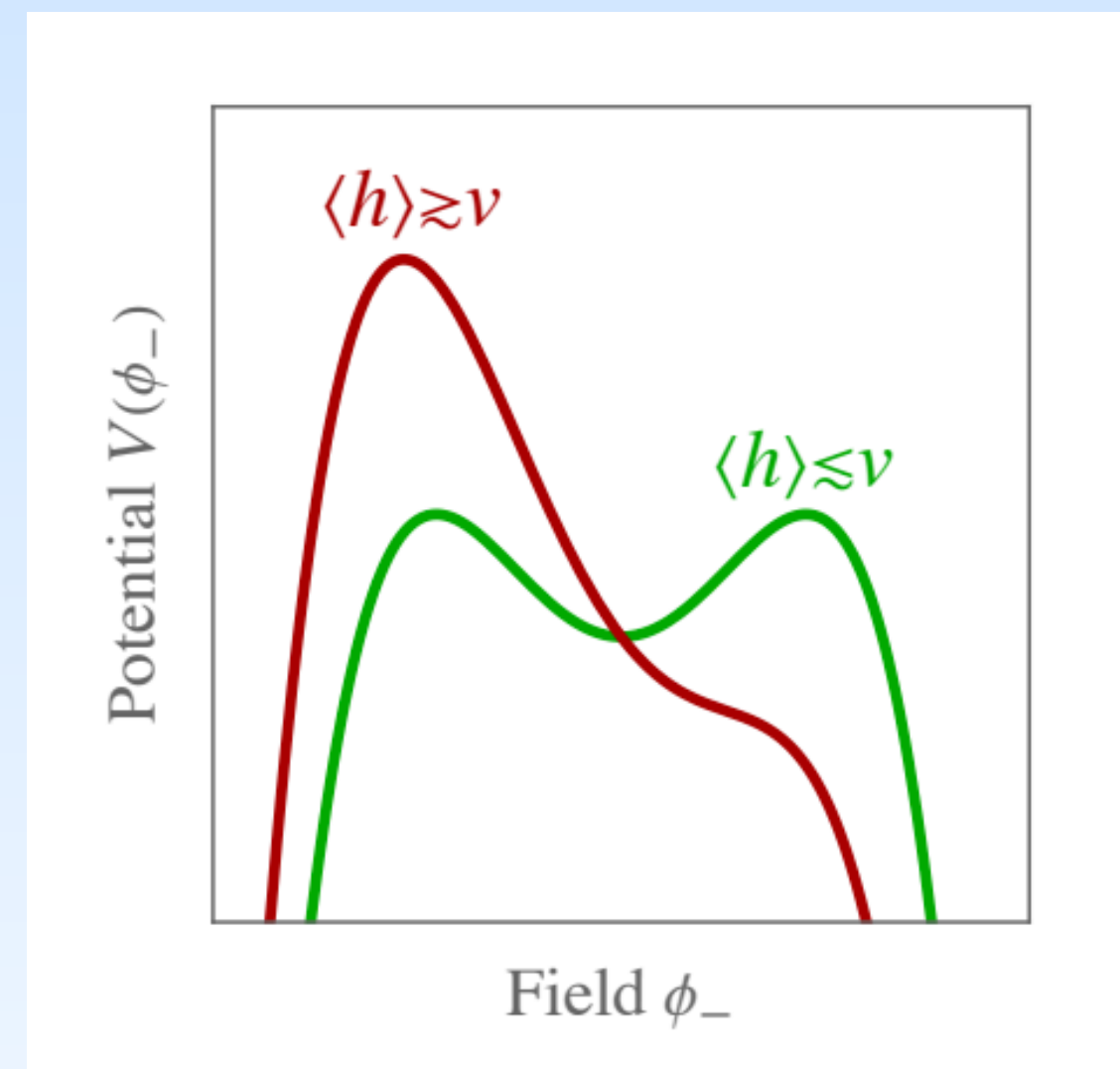
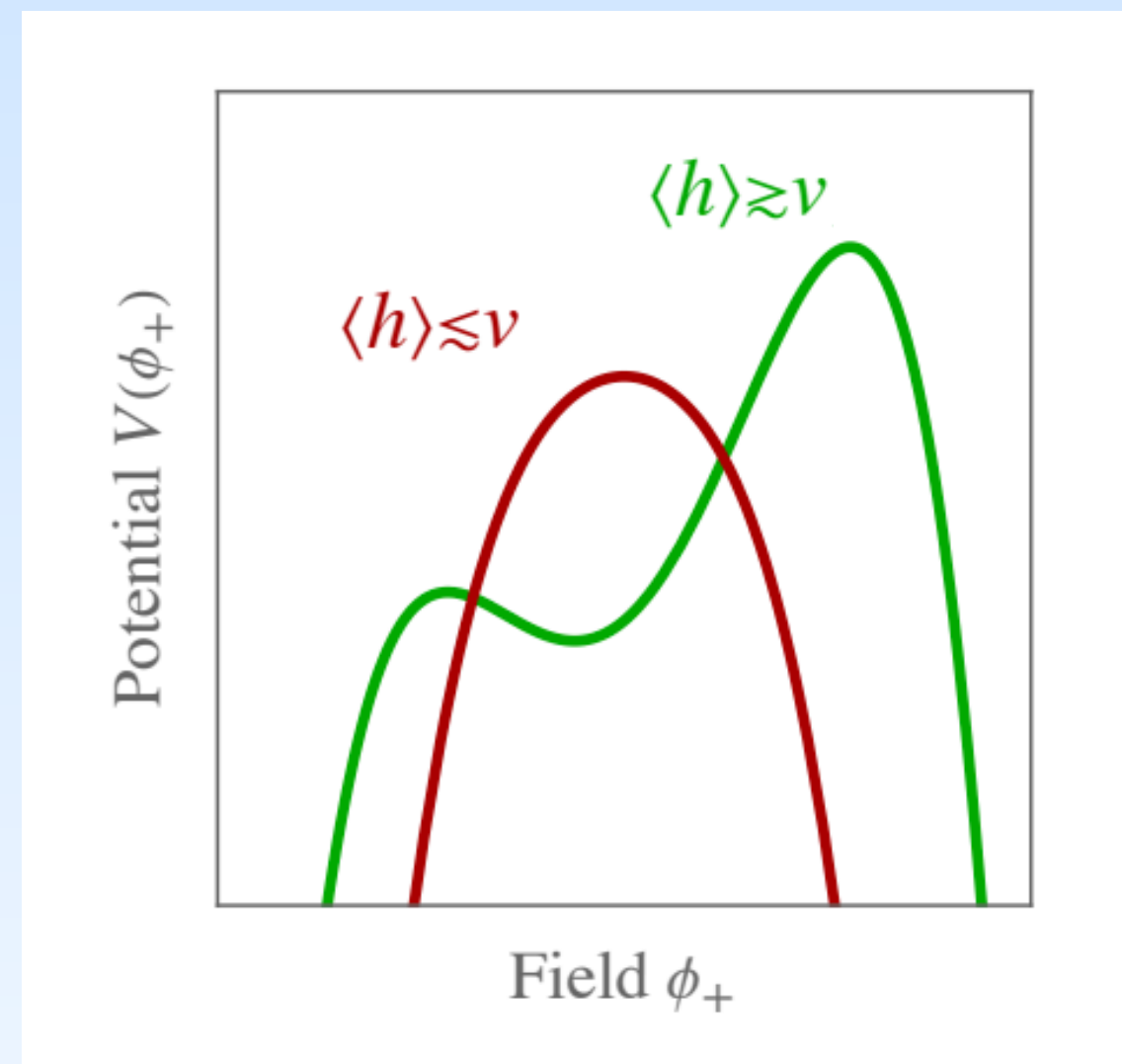
$$\longrightarrow -m_\pi^2 f_\pi^2 \cos(\dots)$$

$$\sim \frac{\Lambda(\langle h \rangle)^4}{2} \left(\frac{\phi_+}{F_+} + \frac{\phi_-}{F_-} + \theta \right)^2$$

The Lagrangian

$$V_+ = -\frac{m_+^2}{2}\phi_+^2 - \frac{m_+^2}{M_+^2}\phi_+^4$$

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$$\sim \frac{\Lambda(\langle h \rangle)^4}{2} \left(\frac{\phi_+}{F_+} + \frac{\phi_-}{F_-} + \theta \right)^2$$

this selects a small and non-zero EW scale:

$$m_+^2 F_+^2 \lesssim \Lambda(\langle h \rangle)^4 \lesssim \frac{m_-^2 F_- M_-}{\theta}$$

A novel solution to the strong-CP problem

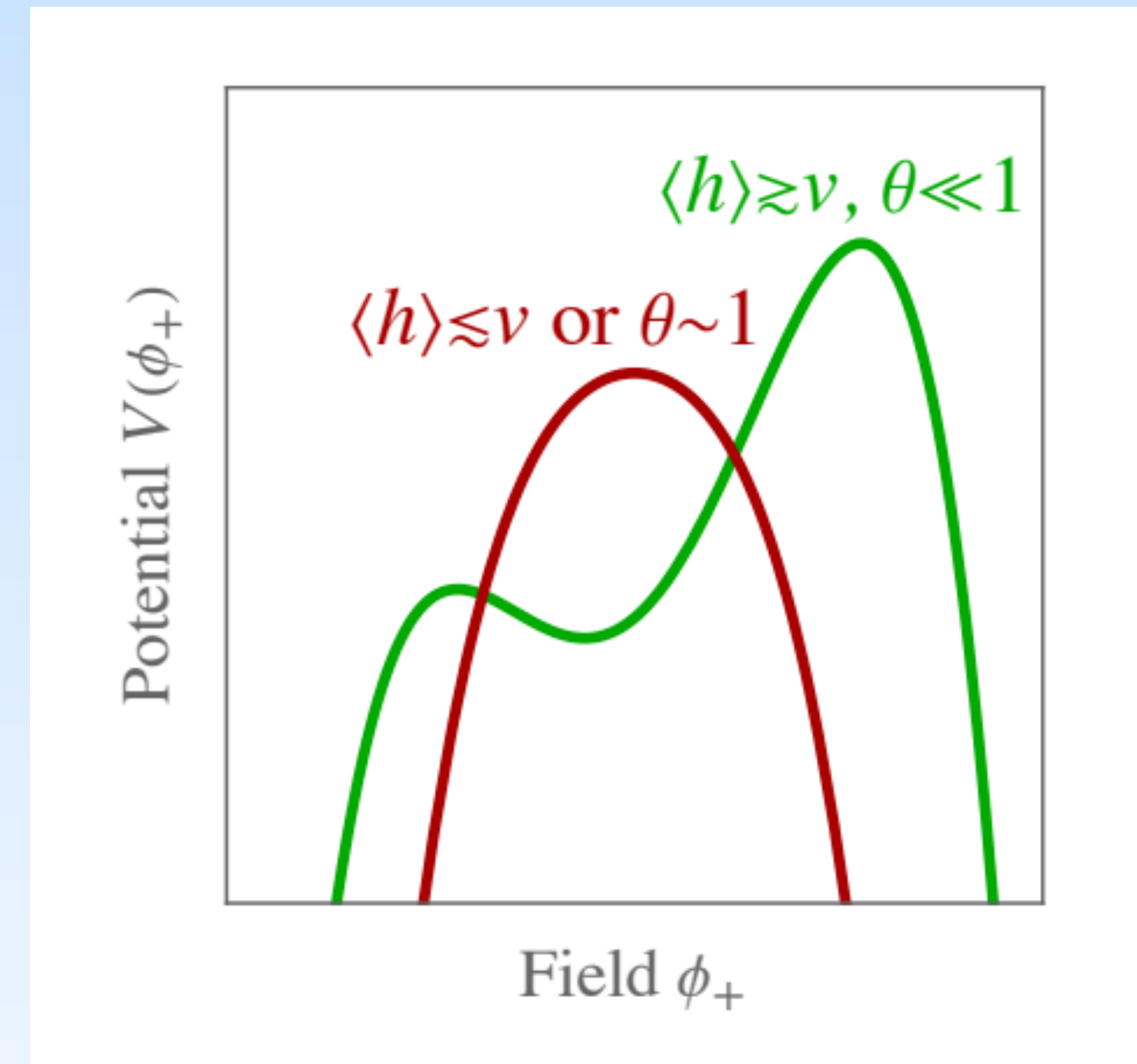
$$V_{\phi H} \sim \frac{\Lambda(\langle h \rangle)^4}{2} \left(\frac{\phi_+}{F_+} + \frac{\phi_-}{F_-} + \theta \right)^2$$

θ shifts the stabilizing effect for ϕ_+

metastable minimum possible only if $\theta \ll 1$
(otherwise negative quartic wins)

the same dynamics selects jointly:

$$v \lesssim \langle h \rangle \lesssim v \quad \text{and} \quad \theta \ll 1$$



Dark matter

- “wrong” universes crunch in $t \sim 1/m_{\pm}$
- $m_{+} \lesssim 1/t_{\text{QCD}} \sim 10^{-11}$ eV otherwise ϕ_{+} is doomed to crunch, before $\Lambda(\langle h \rangle) \neq 0$
- the two scalars are stable over cosmological scales, because very light
- nice scenario: ϕ_{+} or ϕ_{-} dark matter
- relic density from oscillations:

for instance:
$$\frac{\rho_{\phi_{+}}}{\rho_{\text{DM}}} \simeq \frac{\theta_0^2 \Lambda_{\text{QCD}}^4}{T_{\text{eq}} M_{\text{Pl}}^{3/2} m_{\phi_{+}}^{3/2}} \simeq \left(\frac{\theta_0}{10^{-10}} \right)^2 \left(\frac{10^{-19} \text{ eV}}{m_{\phi_{+}}} \right)^{3/2}$$

Smoking-gun pattern

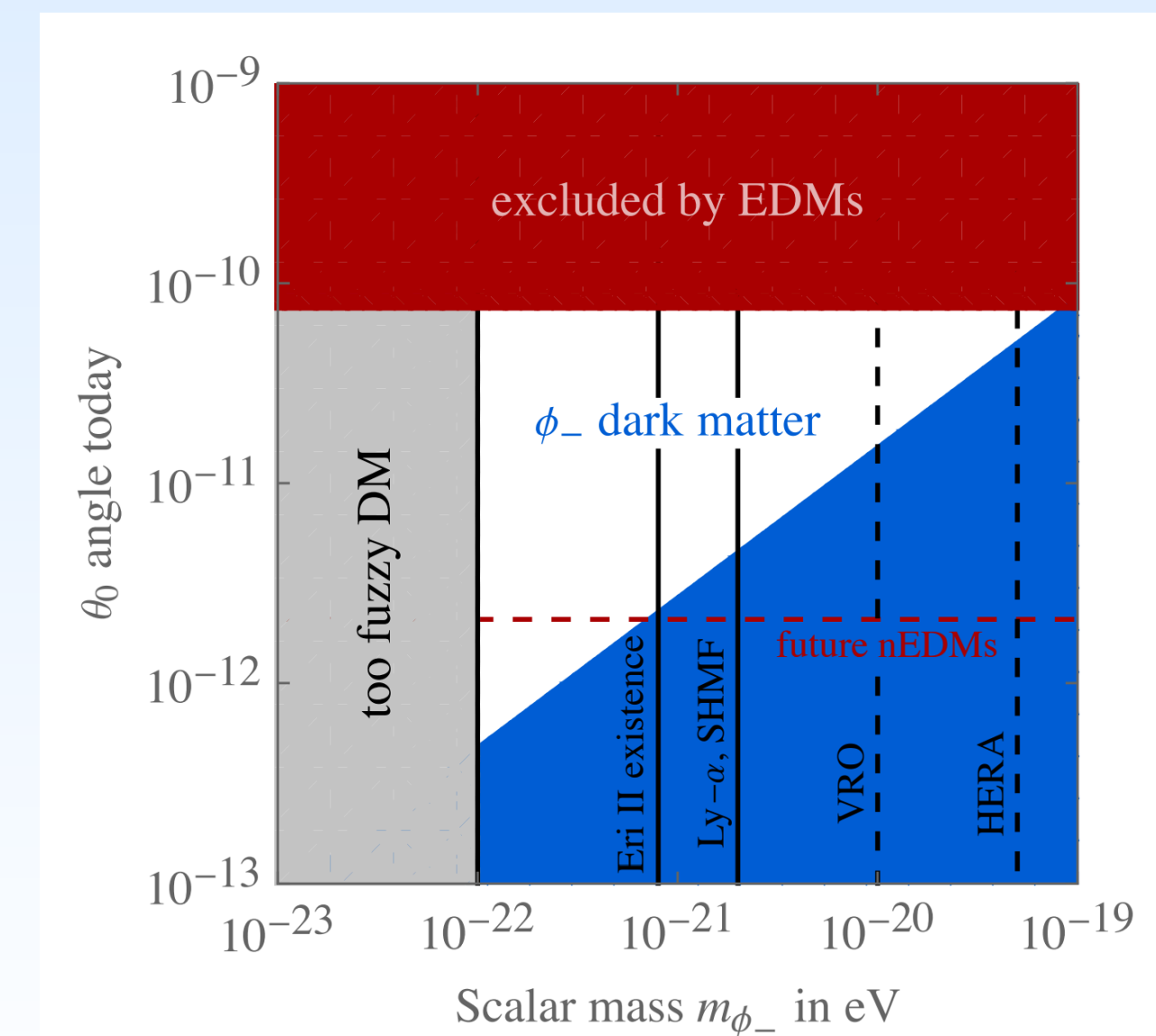
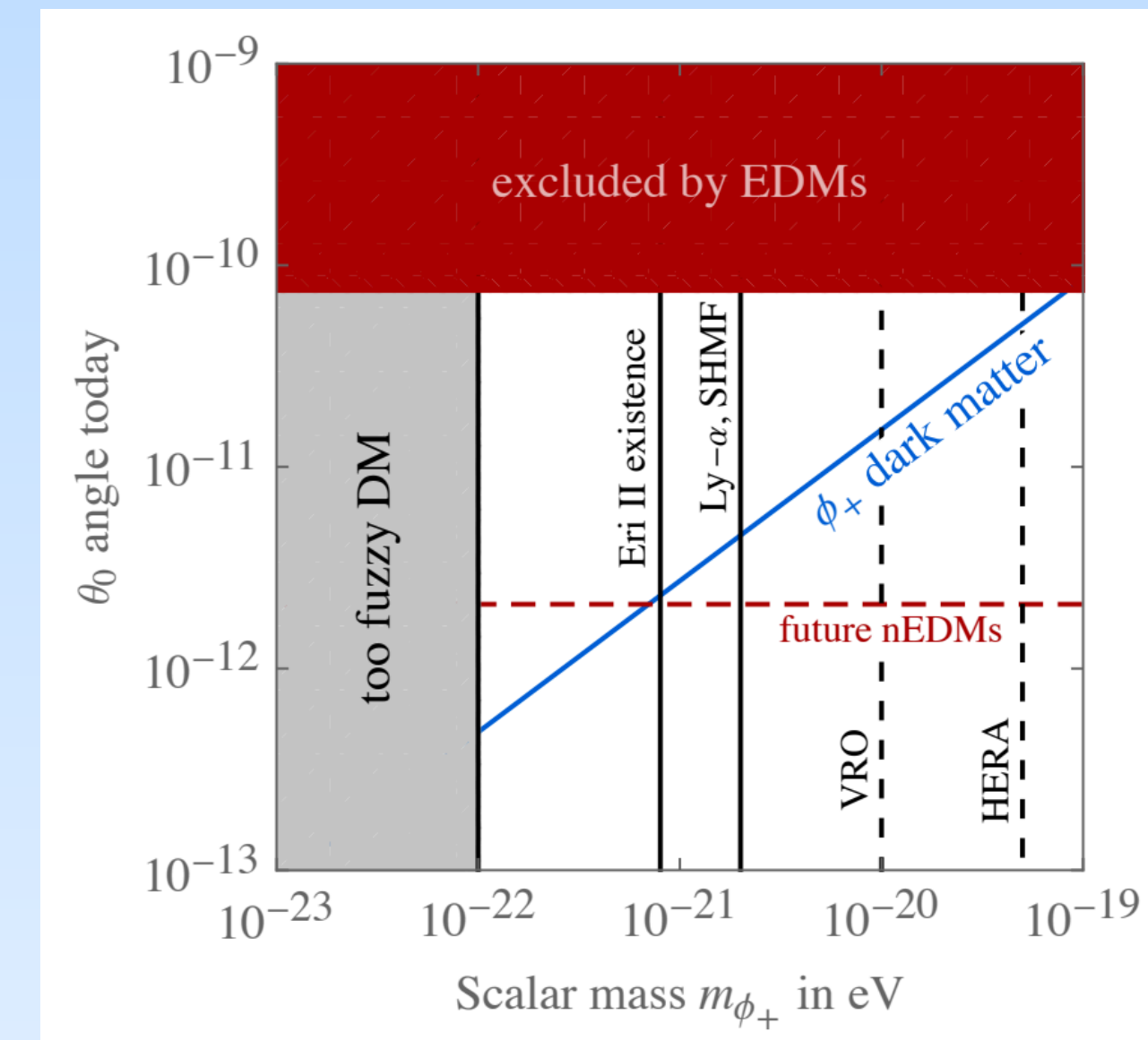
- ϕ_+ is a light scalar $m_{\phi_+} \lesssim 10^{-11}$ eV
with mass-couplings on the QCD line:

$$m_{\phi_+}^2 \sim \frac{\Lambda_{\text{QCD}}^4}{F_+^2}$$

- ϕ_- can be heavier, with mass larger than the QCD line:

$$m_{\phi_-}^2 \sim \theta \frac{\Lambda_{\text{QCD}}^4}{F_- M_-} \gtrsim \frac{\Lambda_{\text{QCD}}^4}{F_-^2} \quad \text{since } M_-/F_- \lesssim M_+/F_+ \sim \theta$$

- if DM, smoking-gun relation with EDMs



Virtues of the mechanism

Why, **in a totally unbiased way**, I would buy this story

- described by a simple potential and quite general (UV does not look painful)
- compatible with standard inflation (it does not need 10^{many} e-folds, or low $H\dots$)
- it can explain $v = 246 \text{ GeV}$ even if Higgs coupled at $\mathcal{O}(1)$ with particles at M_{Pl}
- it is not directly affected by the measure problem of eternal inflation
- compatible with swampland (dS and distance conjectures)
- **bonuses**: DM for free, strong-CP for free, smoking-gun pheno, ...

Conclusions

The hierarchy problem is still out there, more pressing than ever.

Traditional solutions have been failing, but new ideas are being developed. They involve cosmological dynamics.

The way to probe these ideas is different from traditional frameworks.

Sliding Naturalness seems to be a good option.

Backup slides

The ingredients of the game

general features of dynamical selection

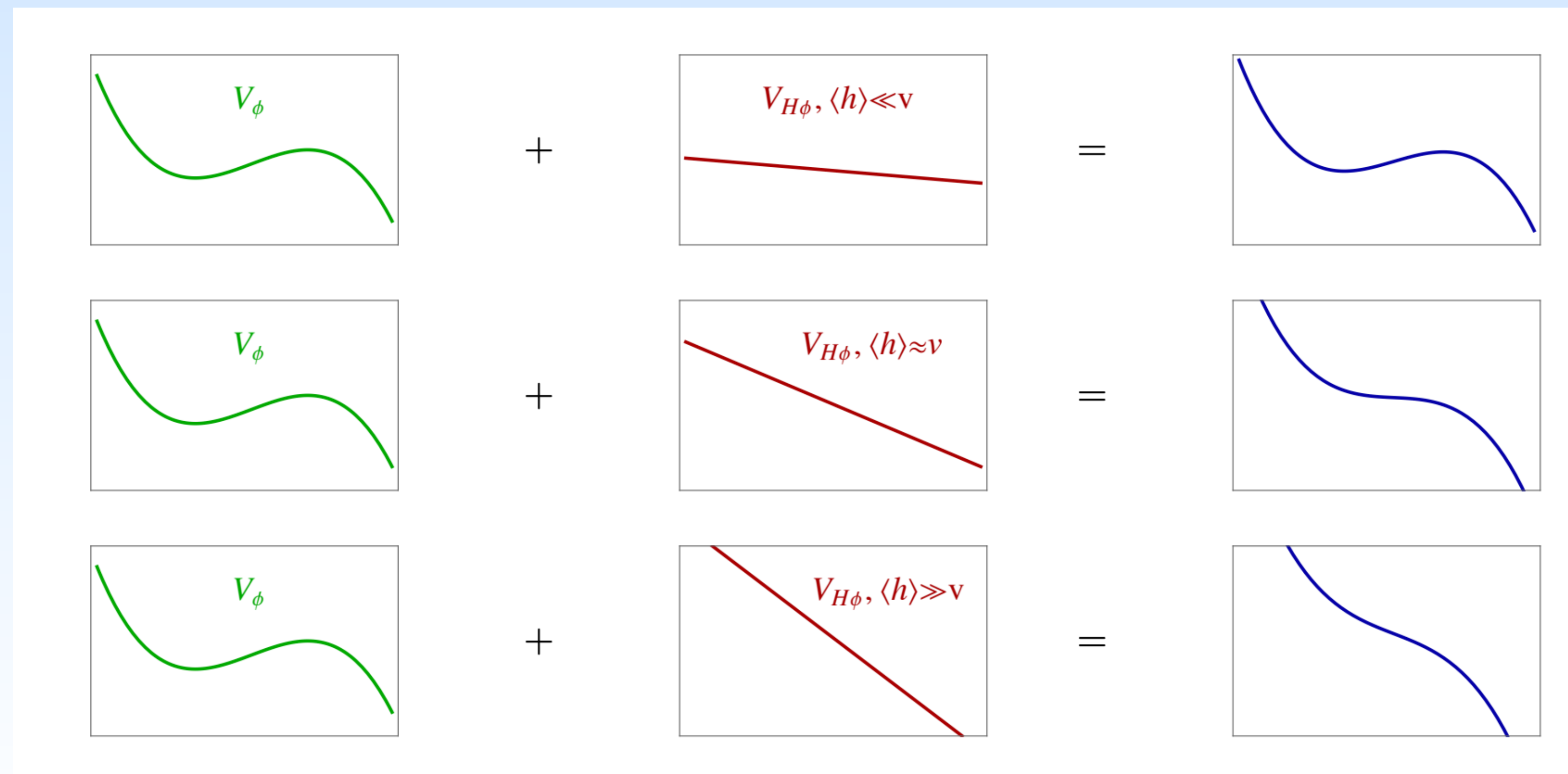
- A **landscape** for the Higgs mass (many vacua from string theory [Bousso, Polchinski] or $O(10-100)$ scalars [Arkani-Hamed, Dimopoulos, Kachru; Ghorbani, Strumia, DT, 1911.01441]). Difficult to observe.
- **light scalars** ϕ (more in next slide)
- **trigger** operator [Arkani-Hamed, D'Agnolo, Kim] $\mathcal{O}(\langle h \rangle) : \phi G\tilde{G}, \phi F\tilde{F}, \phi H_1 H_2 .$

Pheno signatures!

Why light scalars?

An NDA argument [D'Agnolo, DT, 2109.13249]

- $\langle \text{Higgs} \rangle$ -dependent “just-so” to compete with H -independent potential



Why light scalars?

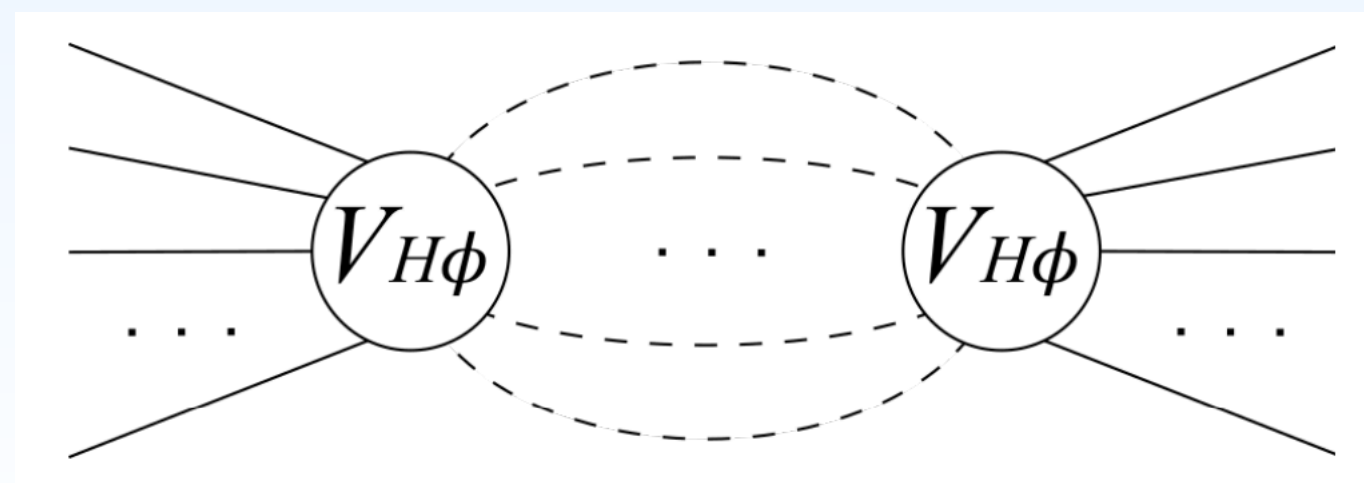
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- NDA: $V_\phi \sim m_\phi^2 M_*^2 \left(\frac{\phi}{M_*}\right)^n$ $V_{\langle H \rangle \phi} \sim g^2 M_*^4 \left(\frac{\phi}{M_*}\right)^m \left(\frac{\langle h \rangle}{\Lambda_H}\right)^{2q}$

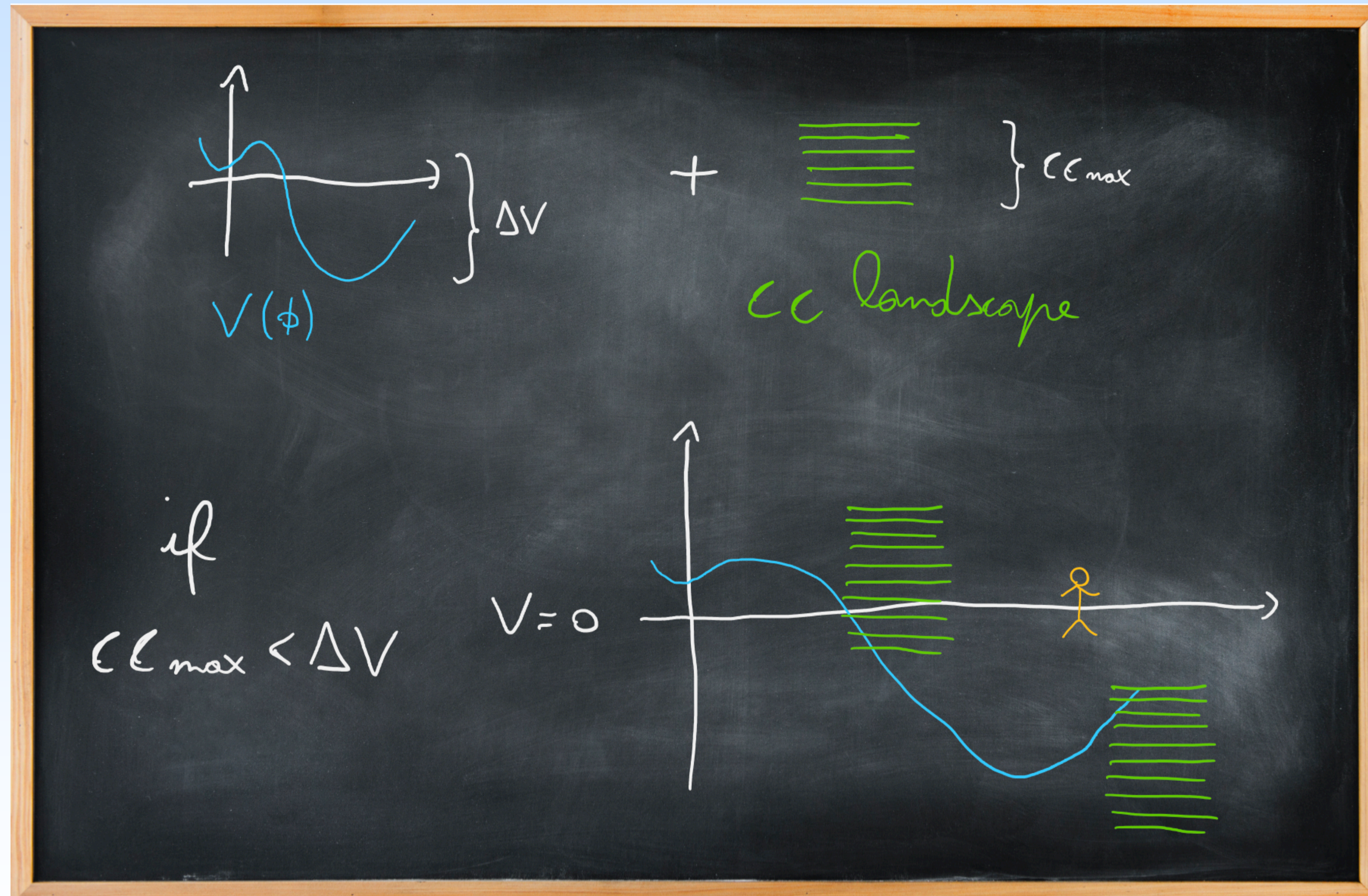
- light scalars \leftrightarrow cutoff $\gg v$: $m_\phi^2 \sim g^2 M_*^2 \left(\frac{v}{\Lambda_H}\right)^{2q}$

- g bound from pheno or naturalness:



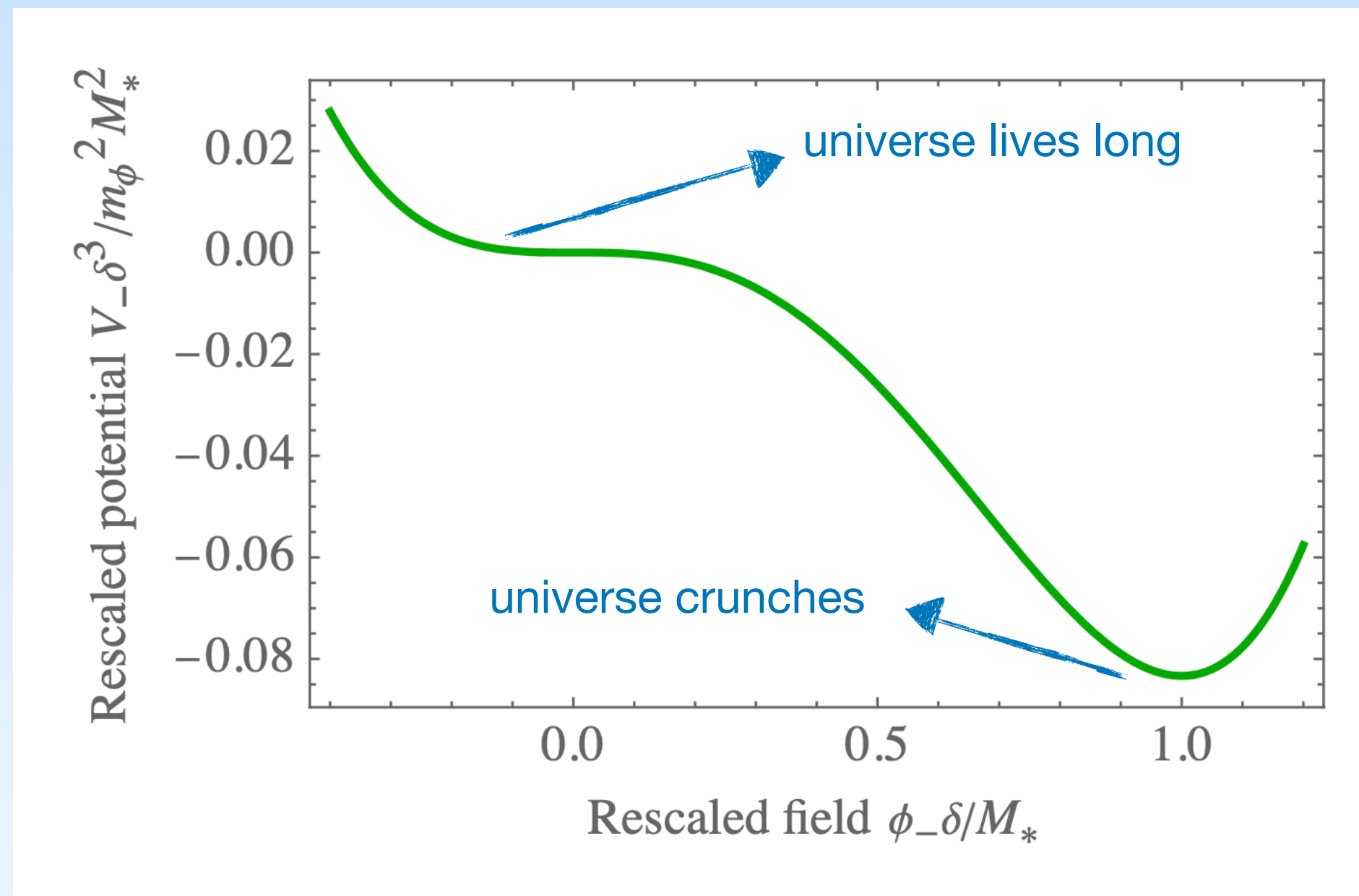
$$m_\phi \lesssim \frac{4\pi v^2}{M_*}$$

The global minimum crunches



Sliding Naturalness

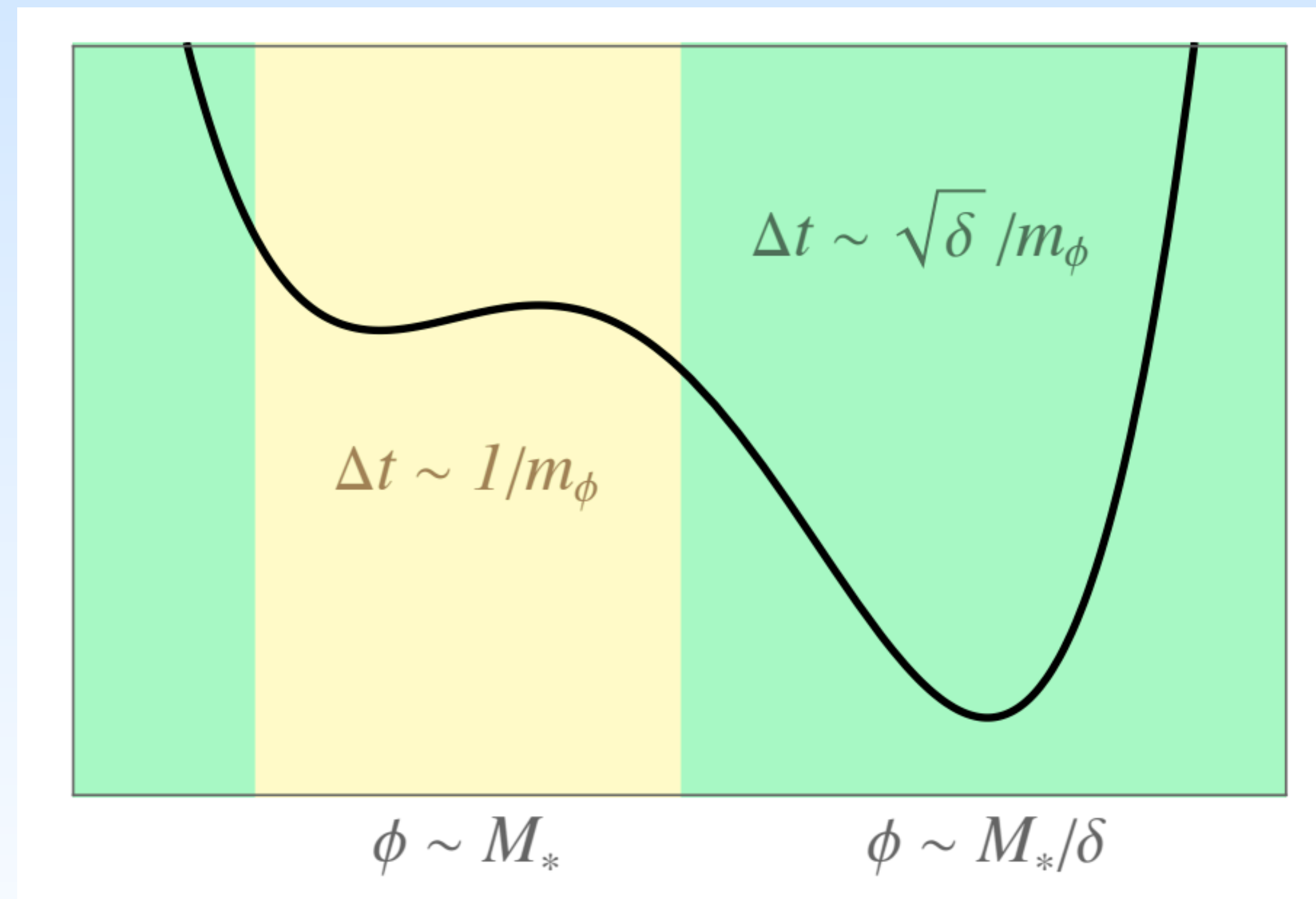
Two Scalars to Rule Them All



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Dark matter

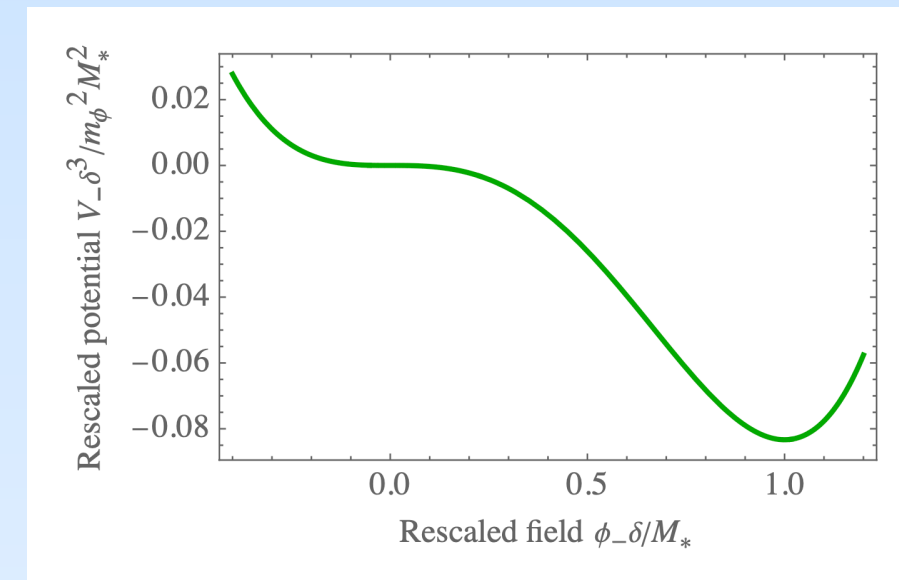
- “wrong” universes crunch in $t \sim 1/m_{\pm}$



More in general

- Many different potentials will do the job (two widely split minima),

for instance
$$V_{\phi_-} = m_{\phi_-}^2 M_*^2 \left(\frac{\phi_-}{M_*} + \frac{\phi_-^2}{2M_*^2} - \frac{\phi_-^3}{3M_*^3} + \frac{\delta \phi_-^4}{4M_*^4} \right) + \dots \quad \text{with } \delta \ll 1$$



- Hierarchy between minima needs to be stabilized by a symmetry,

for instance
$$W_{\phi_-} = L\Phi_- + \mu\Phi_-^2 + \lambda\Phi_-^3 \quad \text{and} \quad V_B = \epsilon\mu\phi_-^3$$

- Hierarchy problem \longrightarrow stabilizing symmetry for $V_{\phi_{\pm}}$ (secluded from SM!)

- ϕ_{\pm} coupled to any “trigger operator” $V \supset \phi_{\pm} \mathcal{O}(\langle h \rangle)$


The $H_1 H_2$ trigger

just a few words

[Arkani-Hamed, D'Agnolo, Kim, '20]

marginally alive, fully tested at HL-LHC

2HDM becomes a trigger if a \mathbb{Z}_2 forbids " $H_1 H_2$ " operators $\longrightarrow V \supset \kappa \phi_{\pm} H_1 H_2$ but $V \not\supset \kappa \phi_{\pm} \times \text{cutoff}^2$

$\phi |H|^2$ doesn't work:  $\Rightarrow V \supset \phi M_*^2$ dominates w.r.t. $\phi \langle h \rangle^2$

Same for standard 2HDM:  $V \supset B_{\mu} H_1 H_2$

If \mathbb{Z}_2 forbids $*$: V sensitive to $\langle h \rangle$

The H_1H_2 trigger just a few words

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