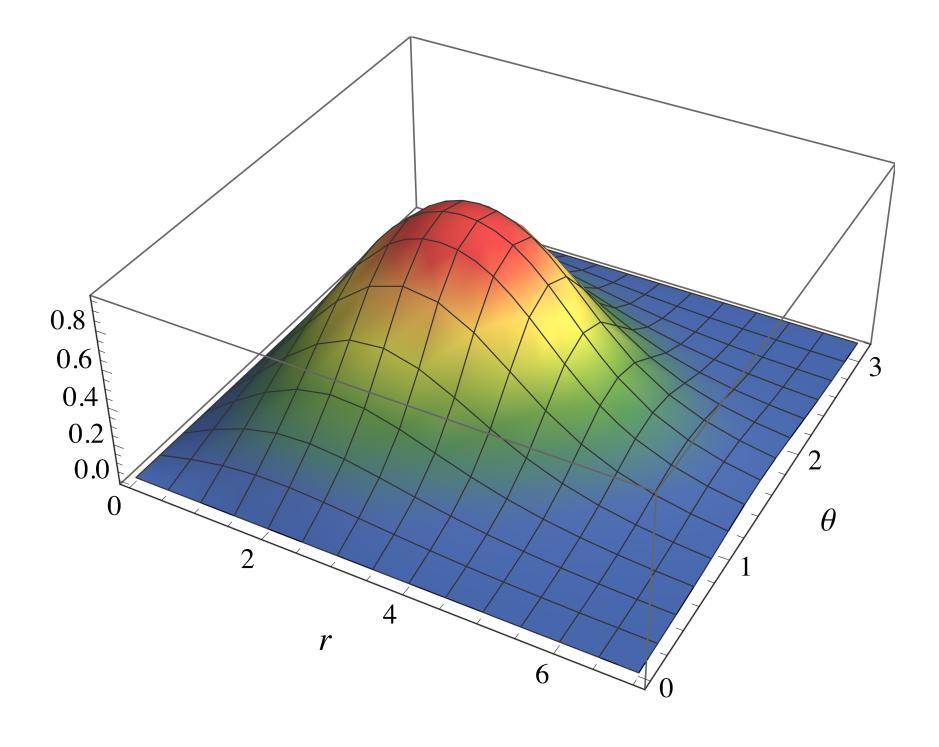
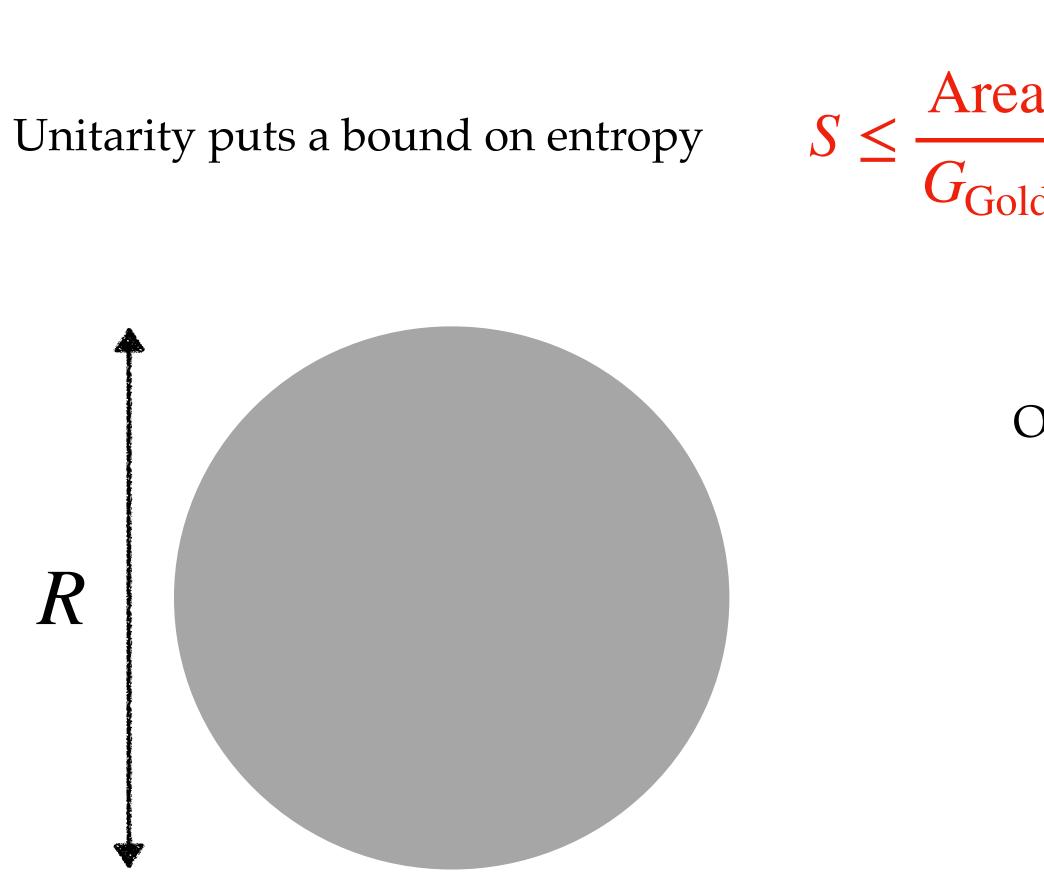
Vortexes and Black Holes

Michael Zantedeschi - LMU and MPP PASCOS 2022

> G. Dvali, F. Kühnel and MZ arXiv:2112.08354



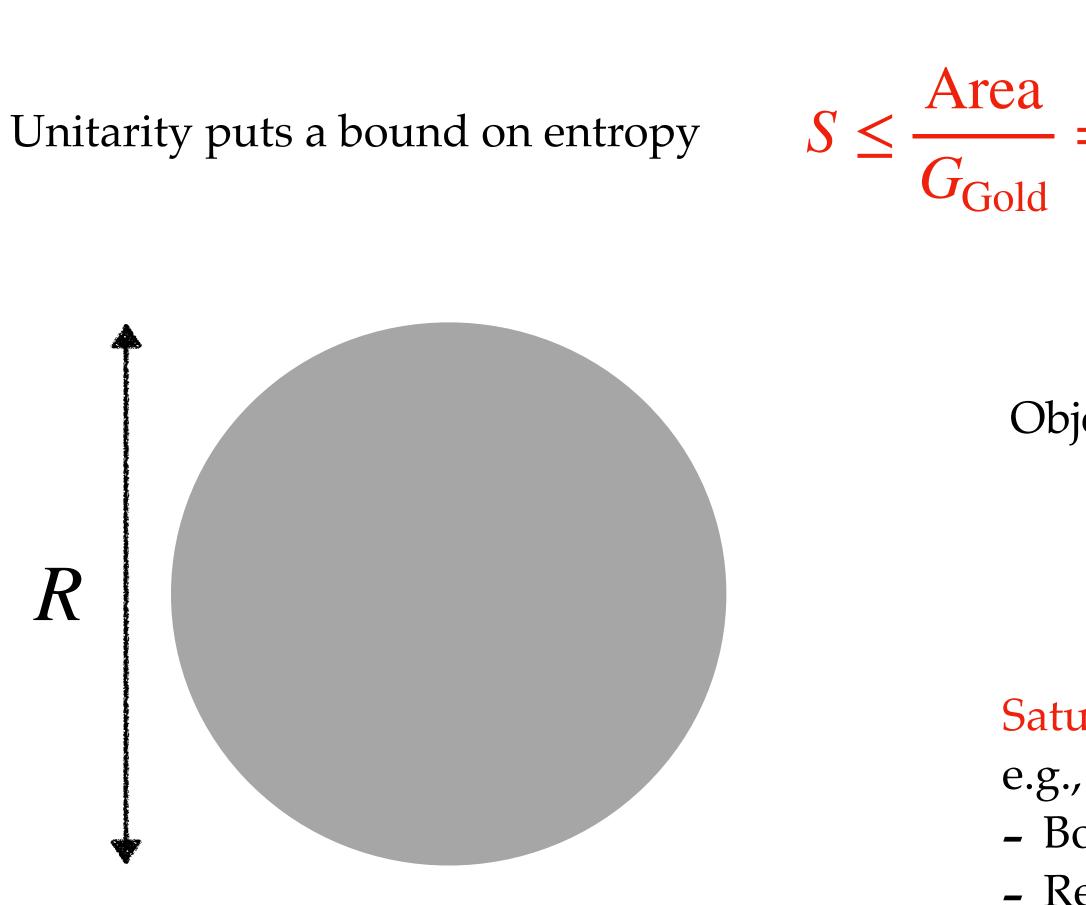




$S \leq \frac{\text{Area}}{G_{\text{Gold}}} = \text{Area}f^2$ $G_{\text{Gold}} = \text{Goldstone coupling} = f^{-2}$ Dvali arXiv:2003.05546arXiv:1907.07332

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Objects saturating the above bound are called saturons



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Objects saturating the above bound are called saturons

Black hole = Saturon $G_{\text{Gold}} = G_{\text{N}}$

Saturon configurations can be found also not in gravity,

- Bose-Einstein condensates

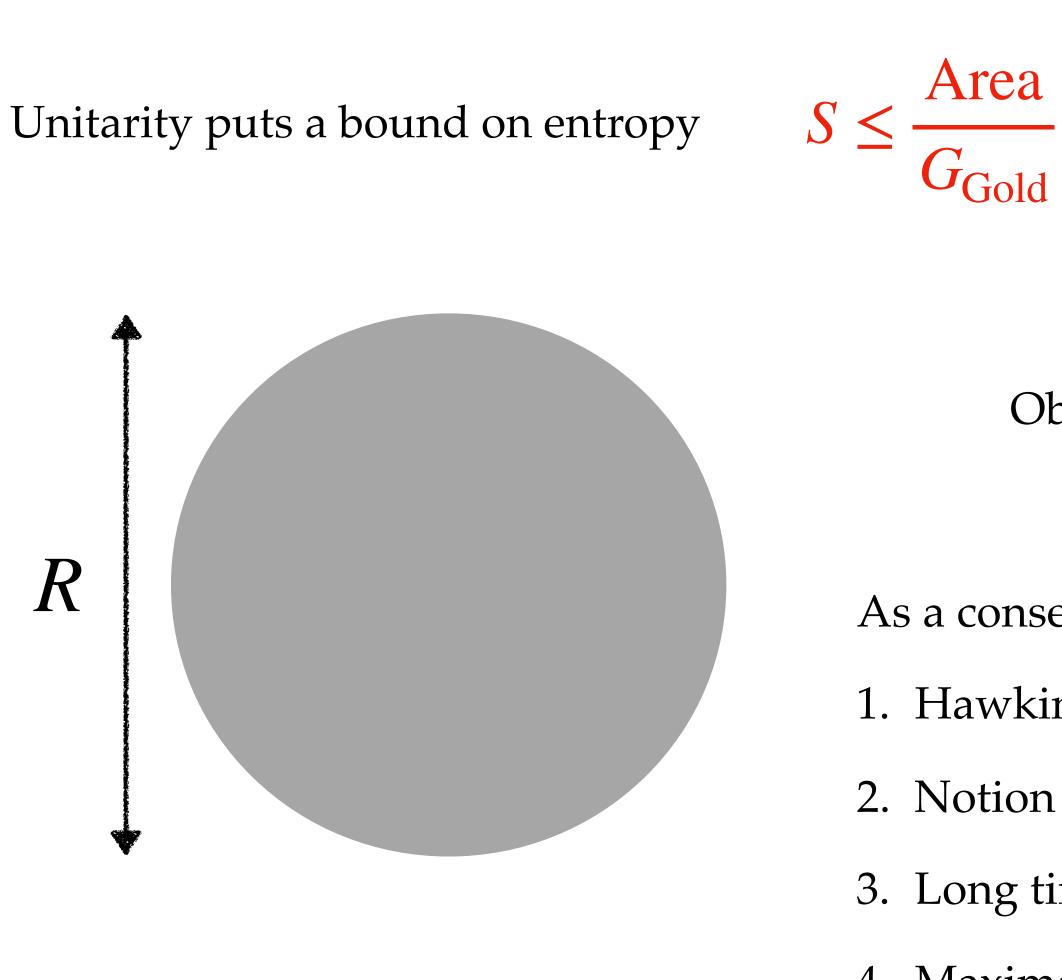
- Renormalizable QFTs (Valbeuena - previous talk)

- Gross-Neveu model (Shakelashvili)

- Color Glass condensate

- ...

arXiv:1906.03530



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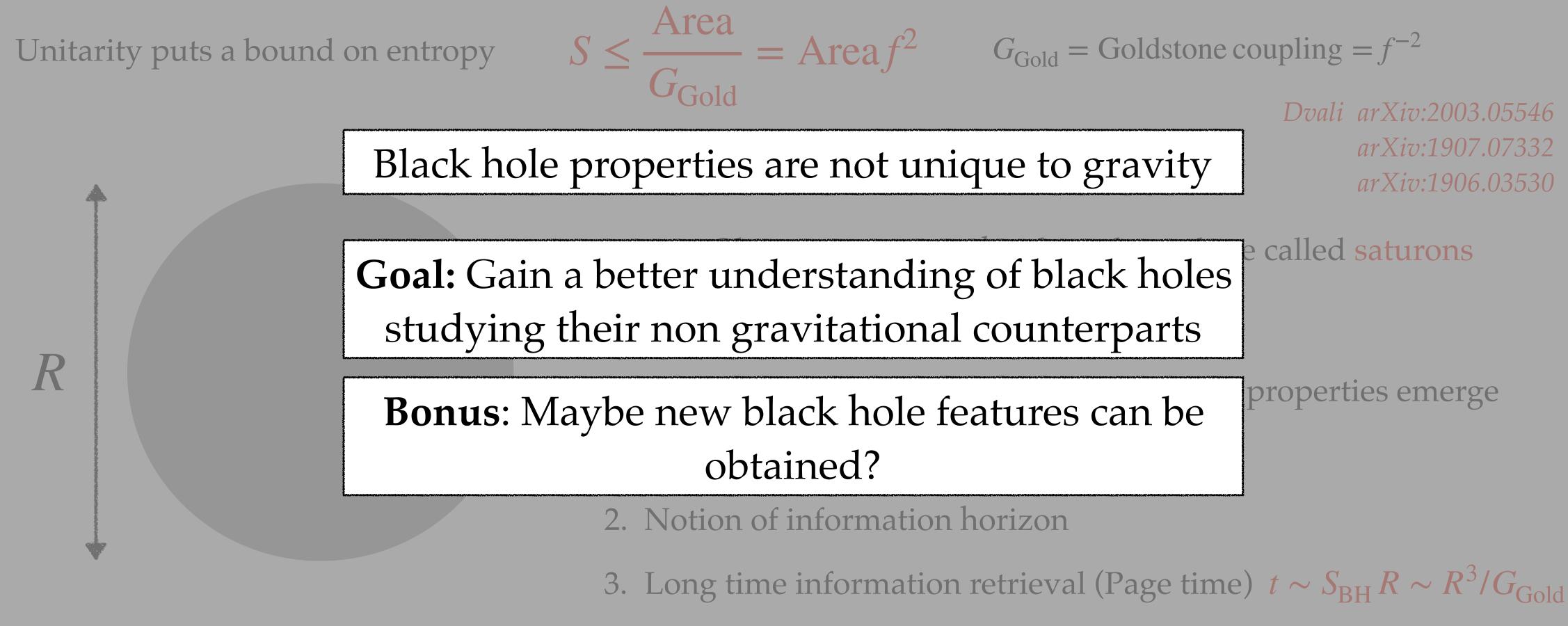
Objects saturating the above bound are called saturons

- As a consequence of saturation, the following properties emerge 1. Hawking thermal emission $T \sim 1/R$
- 2. Notion of information horizon
- 3. Long time information retrieval (Page time) $t \sim S_{BH} R \sim R^3 / G_{Gold}$
- 4. Maximal spin and halt of Hawking emission $J \leq S_{BH}$



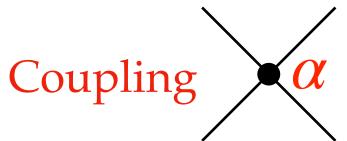
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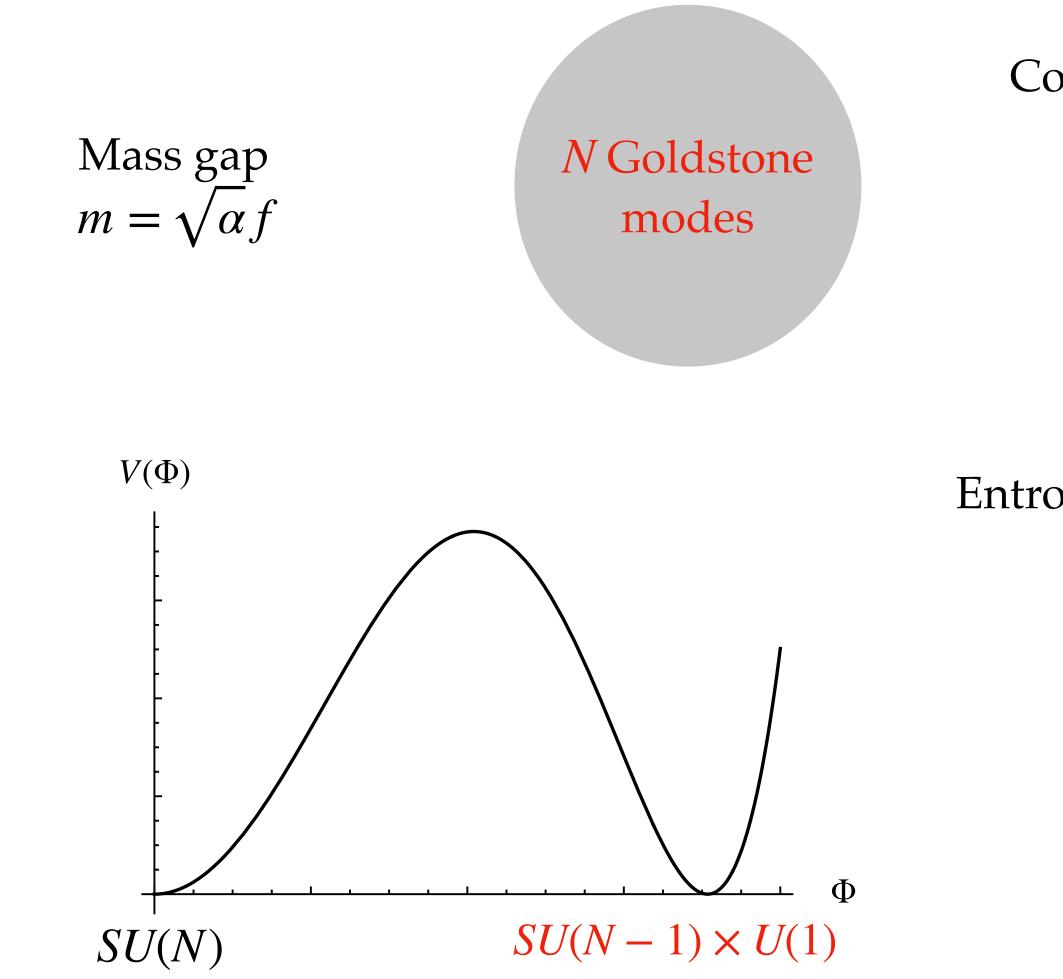


4. Maximal spin and halt of Hawking emission $J \leq S_{BH}$

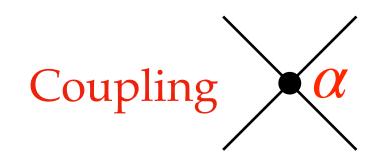
An example of Saturon Consider theory with large *N* symmetry (renormalizable!) e.g., SU(N) scalar adjoint Coupling α Dimensionful scale *f*



Build a localised configuration — bubble



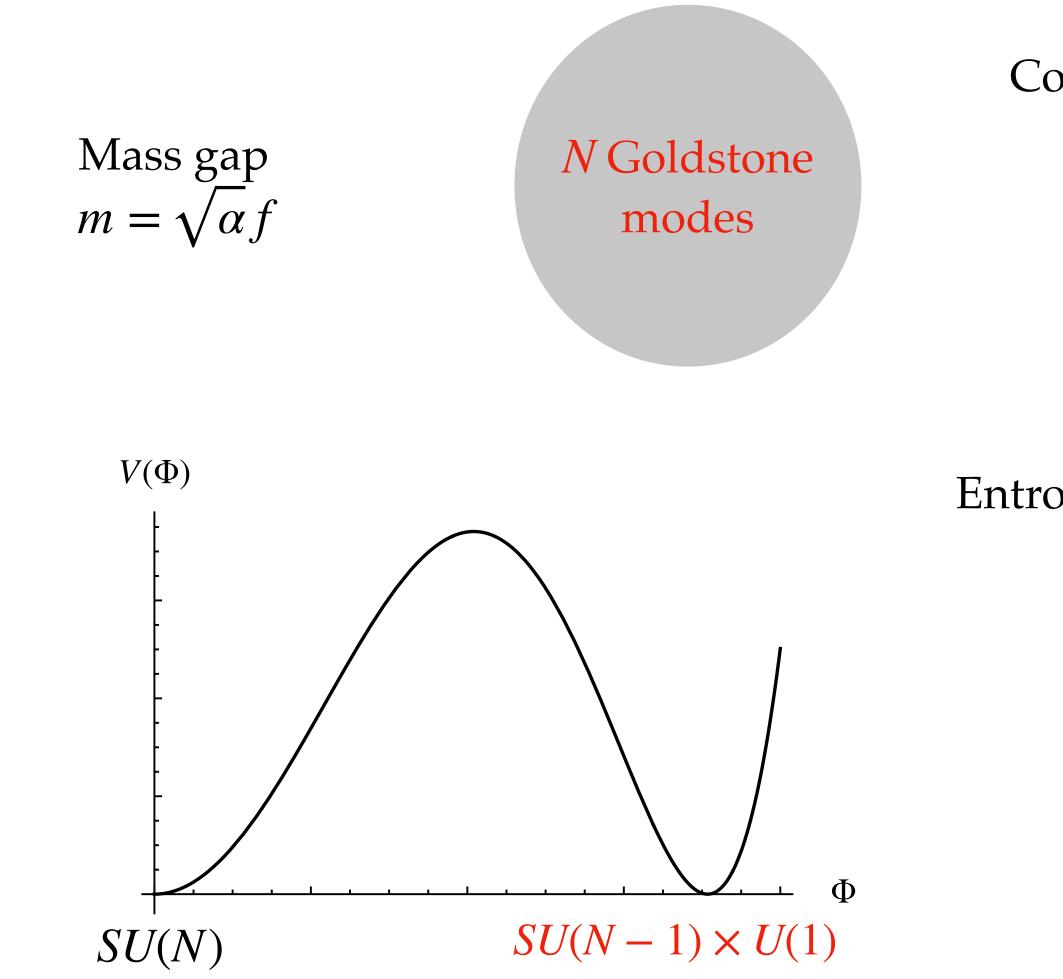
Consider theory with large *N* symmetry (renormalizable!) - e.g., *SU*(*N*) scalar adjoint



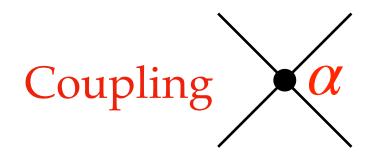
Dimensionful scale f

Build a localised configuration — bubble

Entropy = microstate degeneracy of Goldstone mode occupation



Consider theory with large *N* symmetry (renormalizable!) e.g., SU(N) scalar adjoint



Dimensionful scale f

Build a localised configuration — bubble Entropy = microstate degeneracy of Goldstone mode occupation

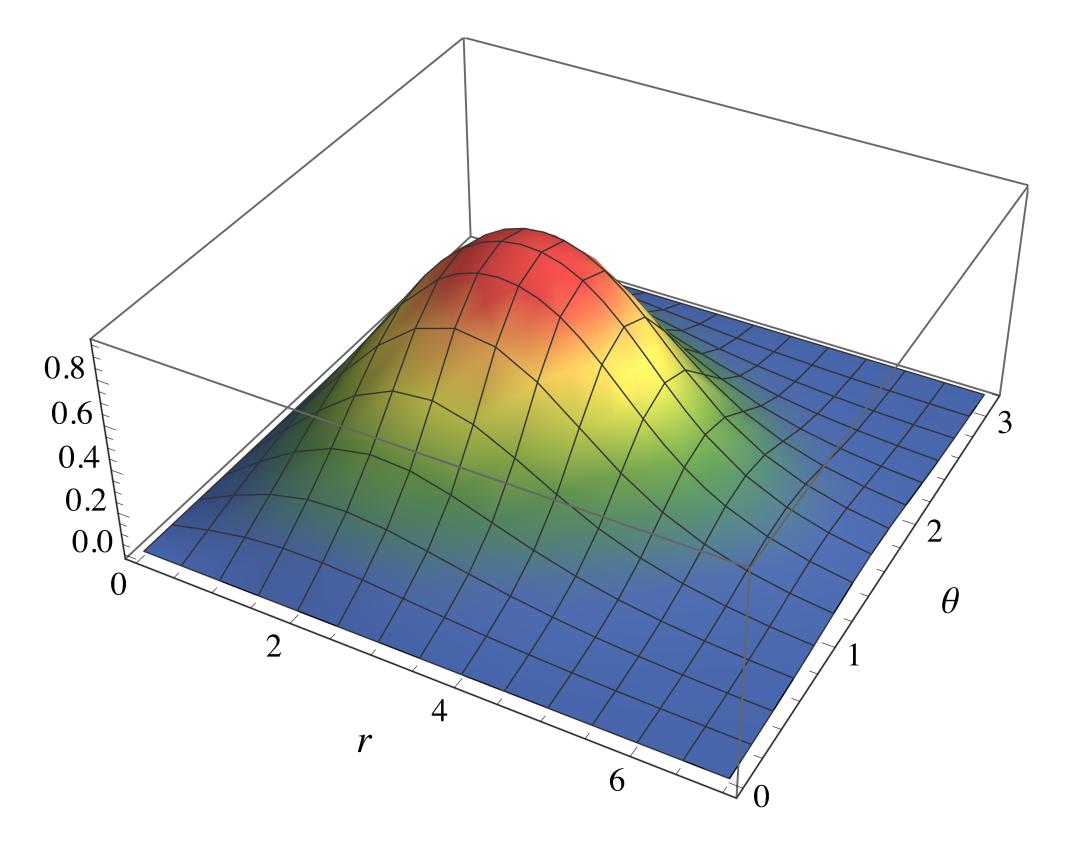
Saturation of unitarity — $\alpha N \sim 1$

$$S \simeq \operatorname{Area} f^2 \simeq N \simeq \frac{1}{\alpha}, \qquad M \simeq \frac{1}{R} N \simeq R f^2$$

Valbuena, previous t

Analogy with $\longrightarrow G_{\text{Gold}} = G_{\text{N}} \left(f = M_{\text{pl}} \right)$ black holes





Spin

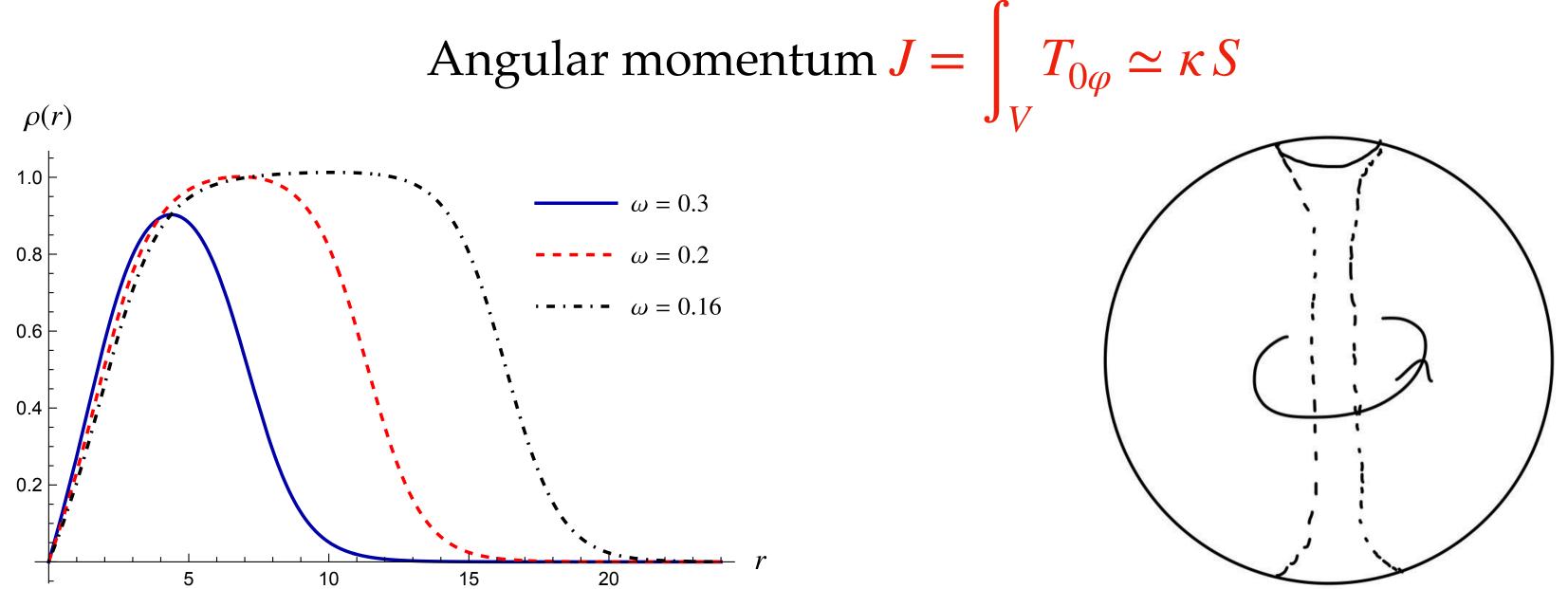
There is a way to spin a saturon bubble in an axial-symmetric way: **Vorticity**

 $\Phi = e^{i\kappa\varphi\hat{T}} \Phi_{Bubble} e^{i\kappa\varphi\hat{T}}$

winding number = $\kappa = 0, \pm 1, \pm 2,...$

 $\varphi = \text{polar angle}, \hat{T} = \text{broken generator inside bubble}$

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Similar study for Q-ball see Volkov, Wohnert '02



- Angular momentum $J \simeq \kappa S$
- Requiring vortex energy smaller than bubble energy

 $E_{\rm spin} \lesssim M_{\rm bubble} \to \kappa \sim \mathcal{O}(1)$

Requiring vortex energy smaller than bubble energy

Saturon bubble

 $\begin{aligned} J_{saturon} \lesssim M^2 \, G_{\text{Gold}} & J_{black \, hole} \lesssim M^2 \, G_{\text{N}} \\ \kappa \sim \mathcal{O}(1) \end{aligned}$

- Saturon and black hole obey the same bound on spin

- For maximally spin, topology prevents Hawking-like emission

Angular momentum $J \simeq \kappa S$

 $E_{\rm spin} \lesssim M_{\rm bubble} \to \kappa \sim \mathcal{O}(1)$

Black hole

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 $J_{saturon} \lesssim M^2 G_{Gold} \qquad J_{black \, hole} \lesssim M^2 G_{N}$ $\kappa \sim \mathcal{O}(1)$

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- Can novel properties be extrapolated to the substructure of black holes?

- Example: beyond semiclassicality in the N-portrait

Angular momentum $J \simeq \kappa S$

 $E_{\rm spin} \lesssim M_{\rm bubble} \to \kappa \sim \mathcal{O}(1)$

Black hole

Requiring vortex energy smaller than bubble energy

Conjecture: *highly rotating black holes correspond* to graviton condensates endowed with vorticity.

Saturon bubble

 $J_{saturon} \lesssim M^2 G_{Gold} \qquad J_{black hole} \lesssim M^2 G_{N}$ $\kappa \sim O(1)$

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- Example: beyond semiclassicality in the N-portrait

An example of Saturon

Angular momentum $J \simeq \kappa n_{Gold} \sim \kappa S$ @ saturation

Black hole

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Saturon bubble

 $\kappa \sim \mathcal{O}(1)$

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Black hole

Caveat: black hole might just be special

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- Sat

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Black hole

Caveat: black hole might just be special

Are there observable consequences?

- Can novel properties be extrapolated to the substructure of black holes?

Global vortex, when interacting with charged matter under arbitrary $U(1)_{gauge}$, traps the associated magnetic field

Dvali, Senjanovic '93

 $\chi_+ = \rho_+ e^{\gamma}$

 $\psi \chi_+ \chi_- + h.c. = 2$

Asymptotically for the gauge field

$$A_{\mu} = \frac{1}{eq} \frac{\langle \rho_{+} \rangle^{2} \partial_{\mu} \theta_{+} - \langle \rho_{-} \rangle^{2} \partial_{\mu} \theta_{-}}{\langle \rho_{-} \rangle^{2} + \langle \rho_{+} \rangle^{2}}$$

Global vortex, when interacting with charged matter under arbitrary $U(1)_{gauge}$, traps the associated magnetic field Dvali, Senjanovic '93

Consider two oppositely $U(1)_{gauge}$ charged field

$$i\theta_+, \quad \chi_- = \rho_- e^{i\theta_-}$$

Effectively coupled to the vortex order parameter $\psi = \rho e^{i\theta}$ as

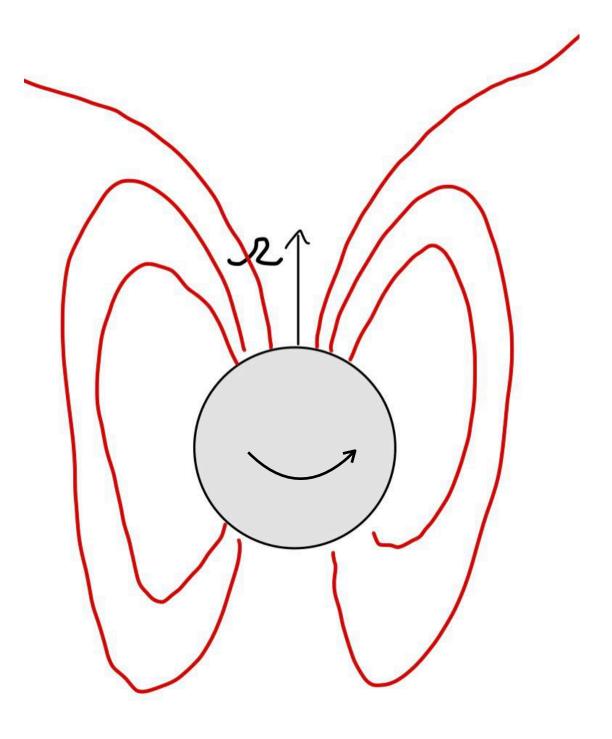
$$2\rho\rho_+\rho_-\cos\left(\theta+\theta_++\theta_-\right)$$

Flux =
$$\int dx^{\mu} A_{\mu} = \frac{2\pi}{eq} \left[\kappa_{+} + \kappa \frac{\langle \rho_{-} \rangle^{2}}{\langle \rho_{-} \rangle^{2} + \langle \rho_{+} \rangle^{2}} \right]$$
$$\kappa_{\pm} = \frac{1}{2\pi} \int dx^{\mu} \partial_{\mu} \theta_{\pm} \qquad \kappa_{+} + \kappa_{-} = -\kappa$$

E.g., charged components of neutral plasma or charged dark matter

Fractional flux has further stabilising properties

If $U(1)_{gauge} = U(1)_{em}$ natural support for magnetic field Black hole pierced by magnetic field lines



Highly rotating BHs endowed with a magnetosphere produce extremely powerful jets (BZ)

The mechanism providing magnetic field remains to present day a mystery (although BZ-like emissions are observed)

Intermezzo: Blandford Znajek emission

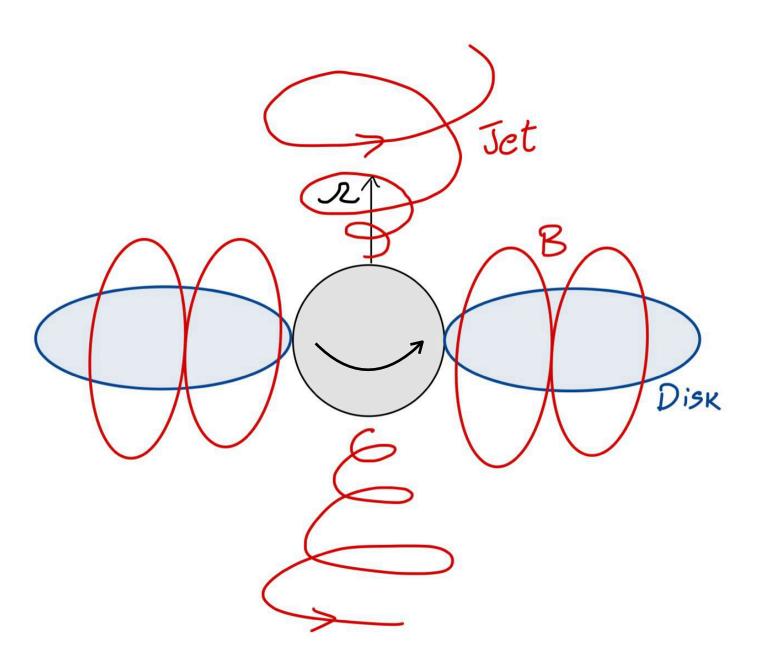
$P_{\rm BZ} \sim {\rm Flux}^2 \Omega^2$

Classically, a BH cannot have magnetic hairs.

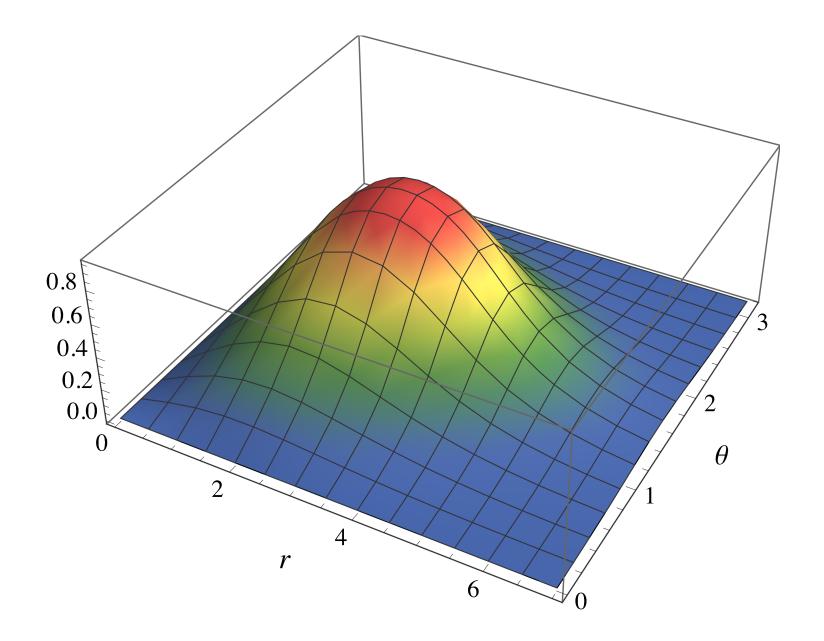
Intermezzo: Blandford Znajek emission

Highly rotating BHs endowed with a magnetosphere produce extremely powerful jets (BZ)

Traditionally it is assumed the presence of an accreting disk with very specific coherent magnetization which can temporary source a magnetic field on the BH



A possible source could also be vorticity



Flux =
$$\int dx^{\mu} A_{\mu} = \frac{2\pi}{eq} \left[\kappa_{+} + \kappa \frac{\langle \rho_{-} \rangle^{2}}{\langle \rho_{-} \rangle^{2} + \langle \rho_{+} \rangle^{2}} \right]$$

$$P_{\rm BZ} \sim {\rm Flux}^2 \Omega^2 \sim P_{\rm M_{87}} \sim 10^2$$

- Jet emission (e.g., à la Blandford Znajek) can take place without the need of an accreting magnetized disk providing a smoking gun for the scenario
 - Example: milli-charged dark matter
 - Consider the axial-symmetric solution found before and $eq \sim 10^{-39}$
 - 44 erg s⁻¹ for maximally rotating BH

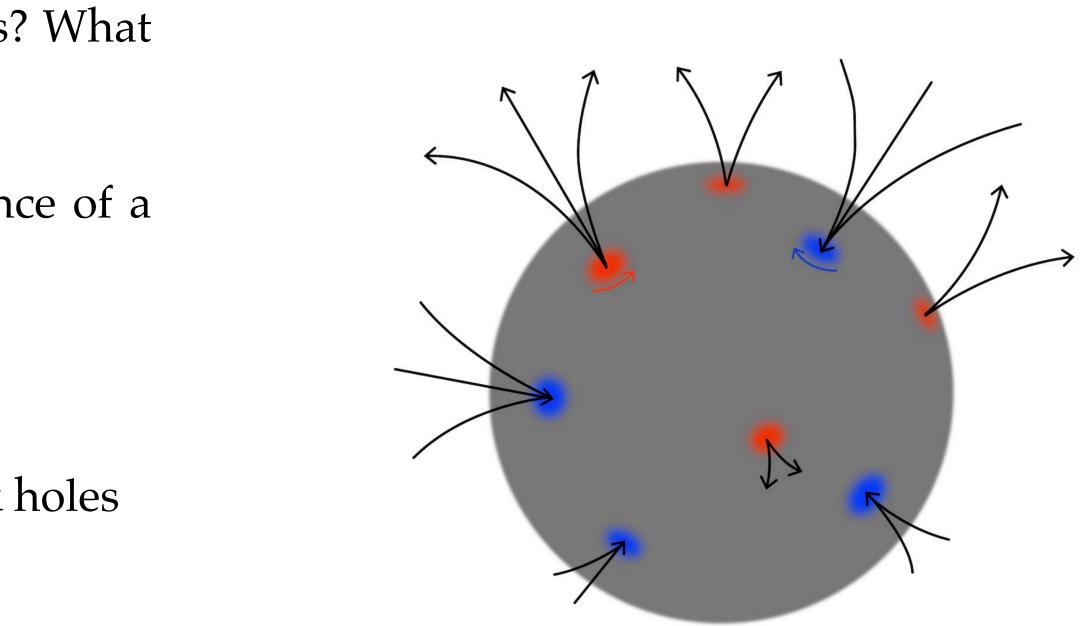
Reminder: $\kappa_{+} + \kappa_{-} = -\kappa \sim \mathcal{O}(1)$

Outlook

- Can we have saturons with multiple vortexes? What about their dynamics?
- How does the emission change in the presence of a magnetized accretion disk?
- Electromagnetic counter part in mergers
- Early Universe relevance for primordial black holes
- •

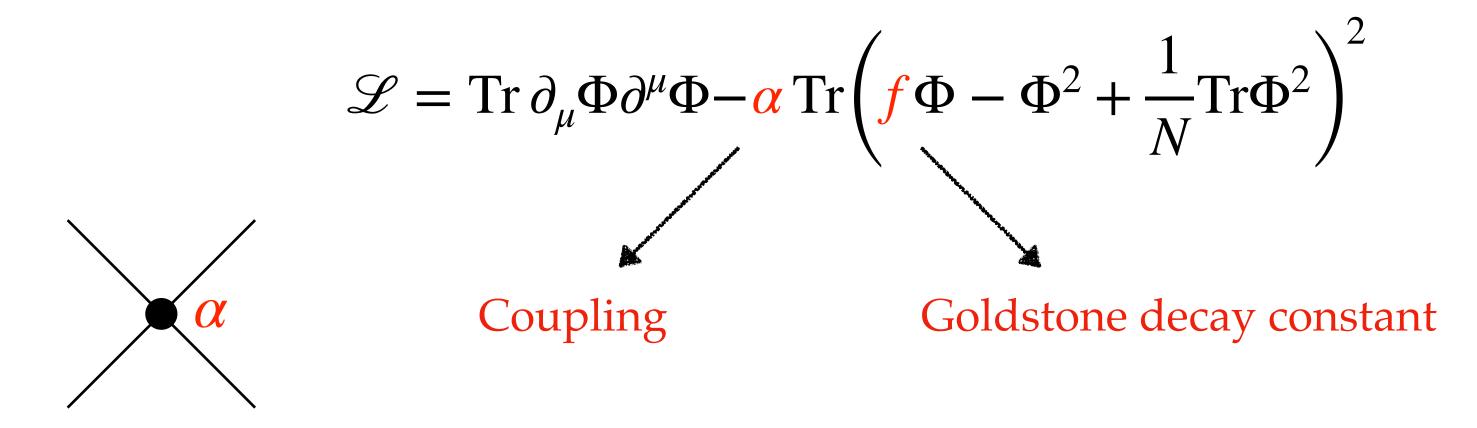
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Thank you

 $i, j = 1, ..., N \rightarrow$ "flavour" indices, $\operatorname{Tr} \Phi = 0$, $\Phi = \Phi^{\dagger}$



Collective coupling is controlled by unitarity

Explicit model

Consider scalar field Φ_i^j in the adjoint representation of SU(N) theory

 $\alpha N \leq 1$

NB The model is renormalizable! NNB Double scaling limit $N \to \infty, \alpha \to 0$

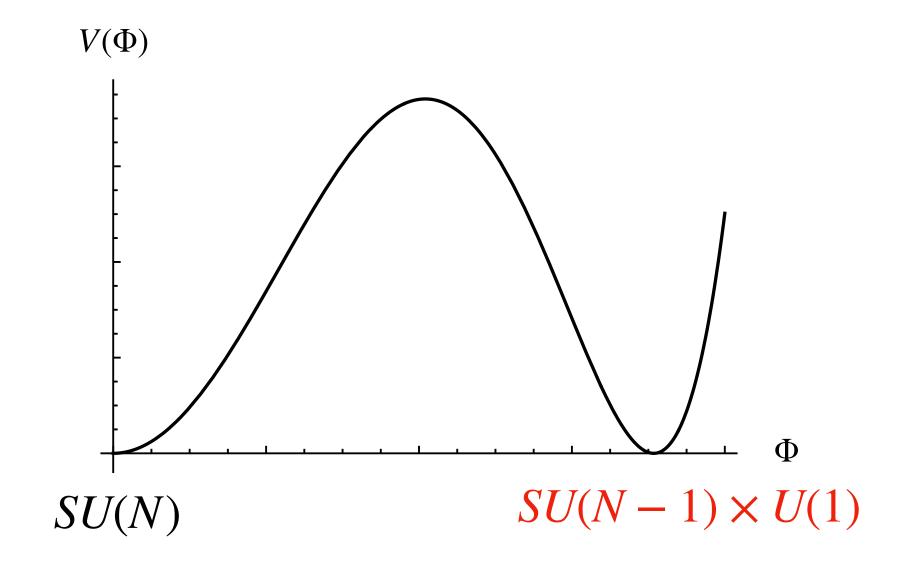
Vacuum structure - many degenerate vacua

$$V(\Phi) = 0$$

$$\int \Phi - \Phi^2 + \frac{1}{N} \operatorname{Tr} \Phi^2 = 0$$

1) SU(N) symmetric vacuum:

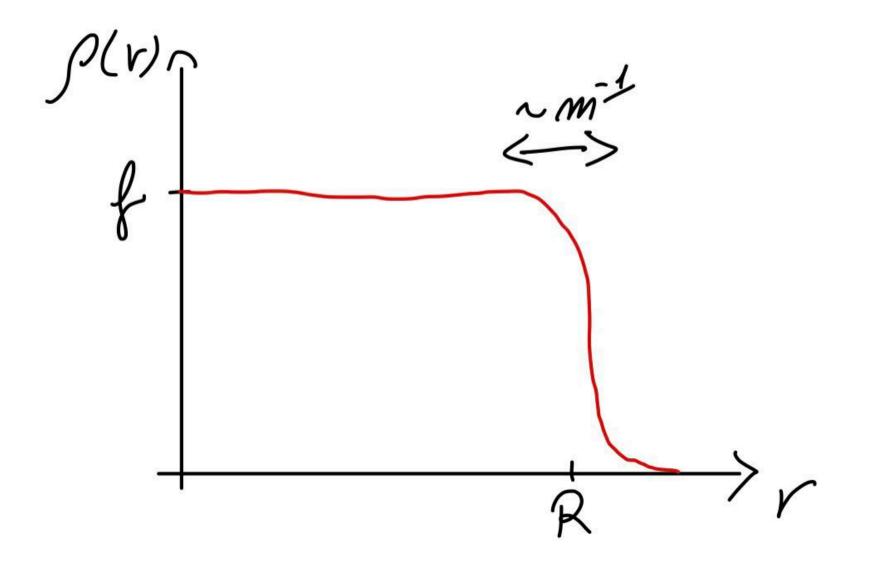
$$\langle \Phi \rangle = 0$$
, mass gap $m = \sqrt{\alpha} f$



2) SU(N-1)xU(1) symmetric vacuum $\langle \Phi \rangle = \frac{f}{N} \text{diag} (N - 1, -1, ..., -1)$ ~ 2N gapless modes

SU(N)
mass gap $m = \sqrt{\alpha} f$

$SU(N-1) \times U(1)$ 2N masslessmodes



An example of Saturon

$$\Phi = \frac{\rho(r)}{f} \langle \Phi \rangle$$

Bubble is highly degenerate and can store information in Goldstone modes

Exciting Goldstones stationarizes the bubble (adding information)

$$\Phi \to U^{\dagger} \Phi U = \frac{\rho(r)}{f} e^{i\omega t \hat{T}} \langle \Phi \rangle e^{-i\omega t \hat{T}}$$

 \hat{T} being one of the broken generators Similar construction in non-topological solitons