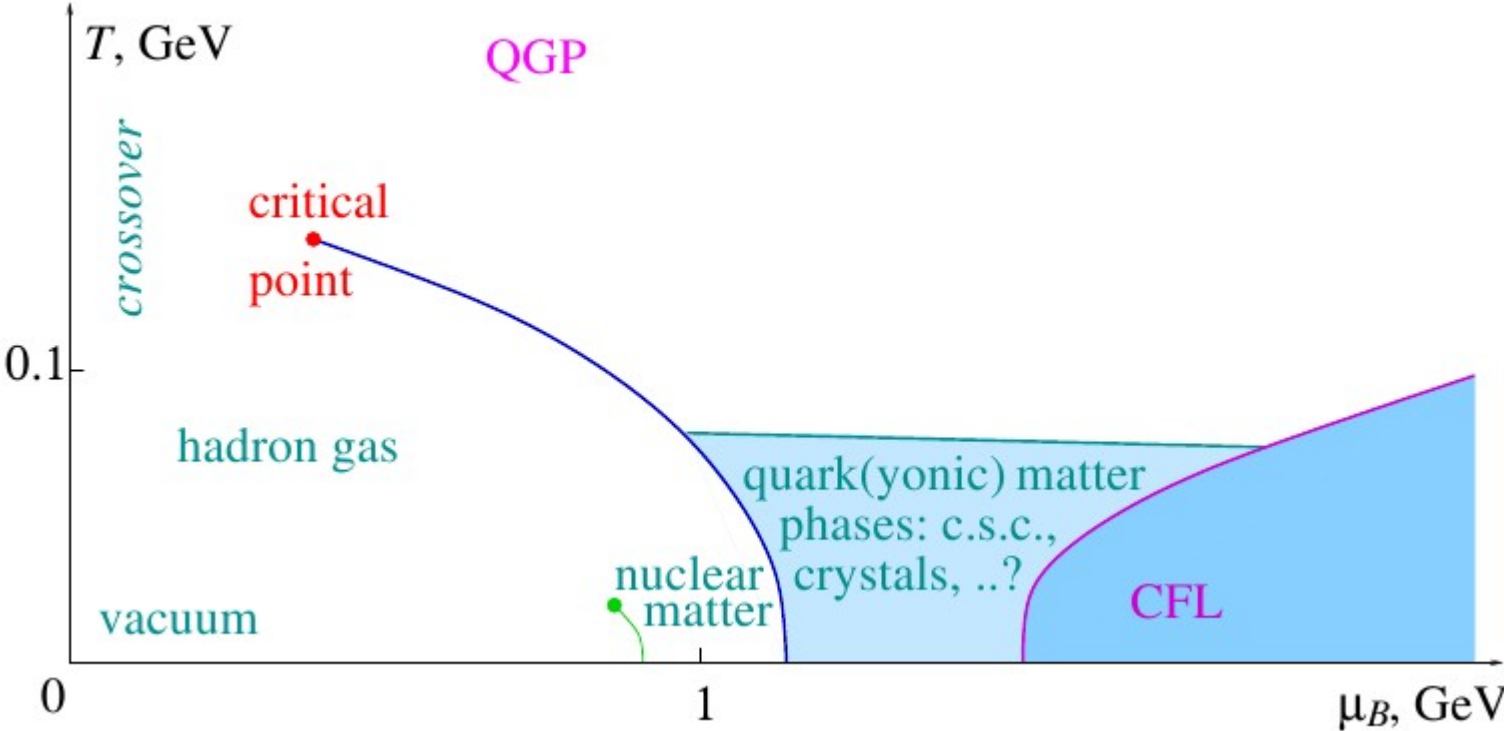


Higher order Fluctuations and Conservation Laws

- Introduction
- Details
- Results
- Summary

The quest



Higher moments (cumulants) and ξ

- Consider probability distribution for the order-parameter field:

$$P[\sigma] \sim \exp \{ -\Omega[\sigma]/T \},$$

$$\Omega = \int d^3x \left[\frac{1}{2} (\nabla \sigma)^2 + \frac{m_\sigma^2}{2} \sigma^2 + \frac{\lambda_3}{3} \sigma^3 + \frac{\lambda_4}{4} \sigma^4 + \dots \right]. \quad \Rightarrow \quad \xi = m_\sigma^{-1}$$

- Moments (connected) of $q = 0$ mode $\sigma_V \equiv \int d^3x \sigma(x)$:

$$\kappa_2 = \langle \sigma_V^2 \rangle = VT \xi^2; \quad \kappa_3 = \langle \sigma_V^3 \rangle = 2VT^2 \lambda_3 \xi^6;$$

$$\kappa_4 = \langle \sigma_V^4 \rangle_c \equiv \langle \sigma_V^4 \rangle - 3 \langle \sigma_V^2 \rangle^2 = 6VT^3 [2(\lambda_3 \xi)^2 - \lambda_4] \xi^8.$$

- Tree graphs. Each propagator gives ξ^2 .

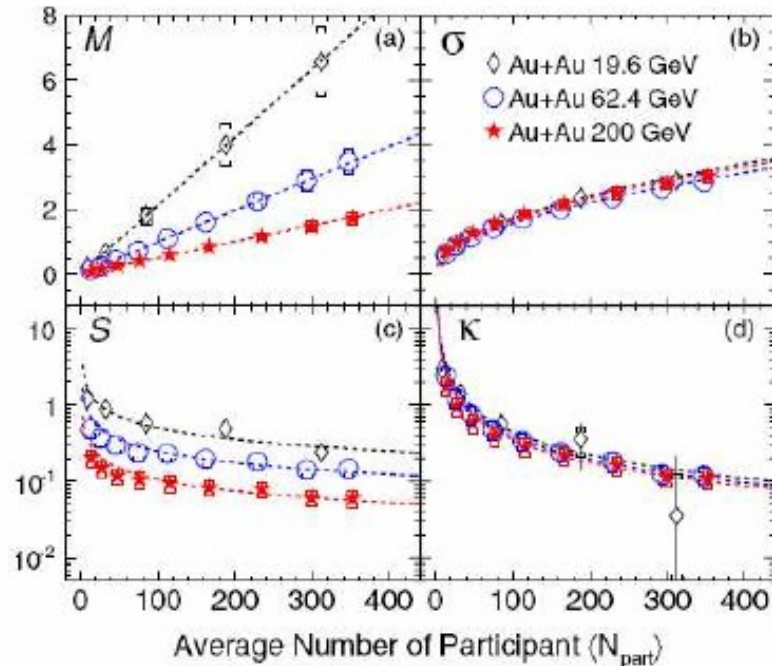
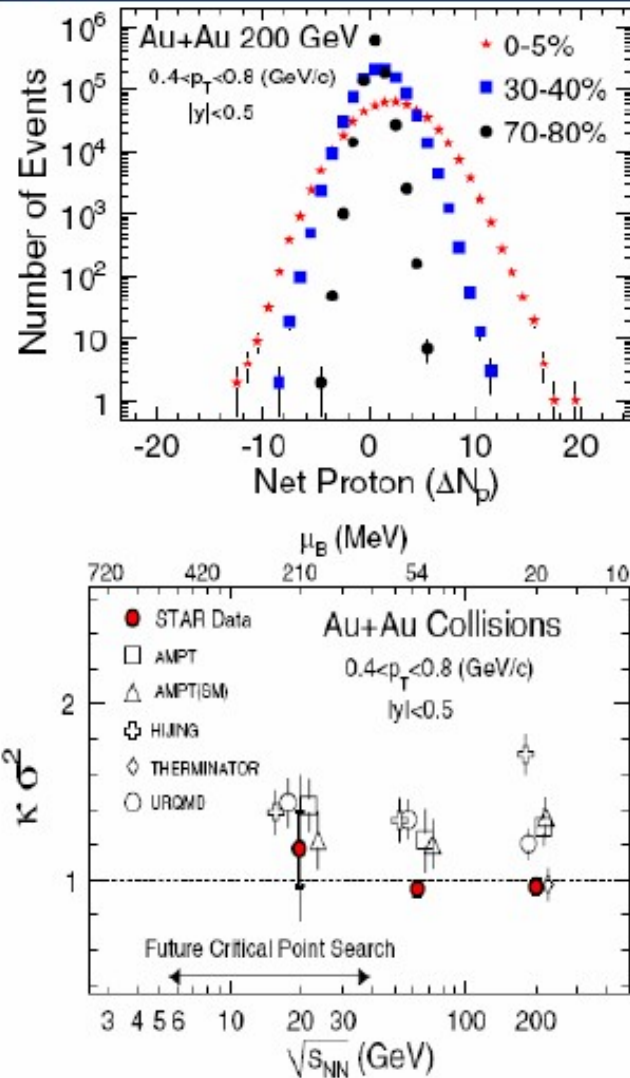


- Scaling requires “running”: $\lambda_3 = \tilde{\lambda}_3 T (T\xi)^{-3/2}$ and $\lambda_4 = \tilde{\lambda}_4 (T\xi)^{-1}$, i.e.,

$$\kappa_3 = \langle \sigma_V^3 \rangle = 2VT^{3/2} \tilde{\lambda}_3 \xi^{4.5}; \quad \kappa_4 = 6VT^2 [2(\tilde{\lambda}_3)^2 - \tilde{\lambda}_4] \xi^7.$$



First result on higher moments of net-proton

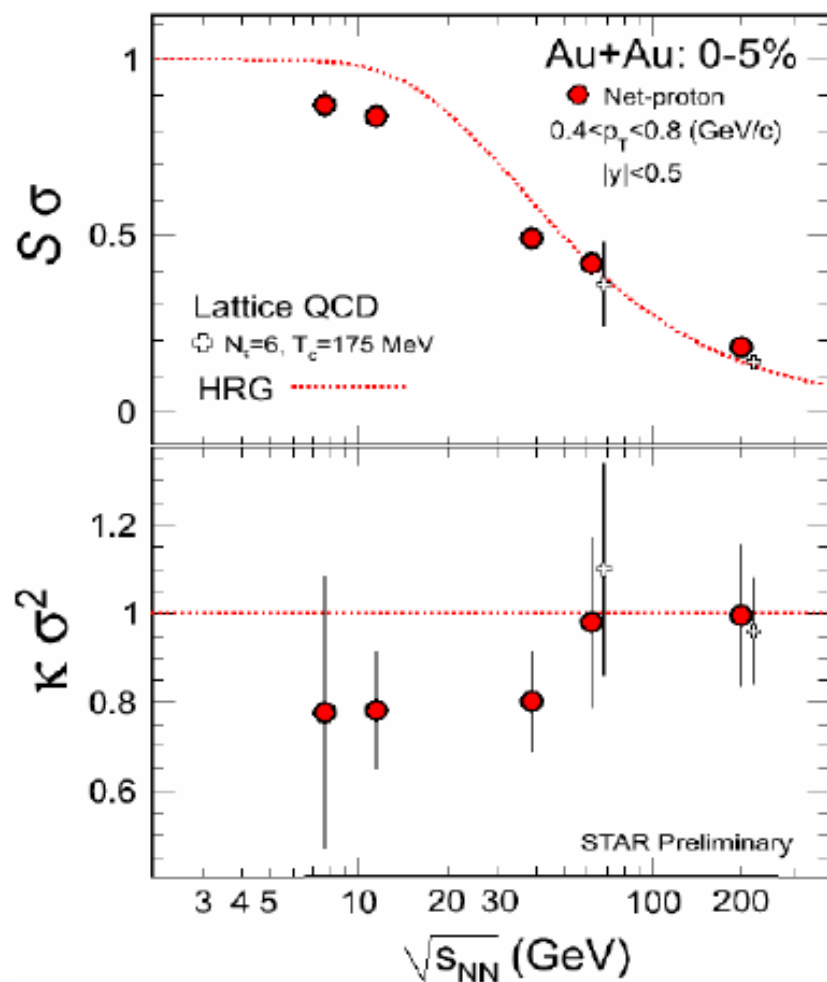


STAR Collaboration, PRL 105 (2010) 022302.

- STAR first results on higher moments analysis are up to fourth order.
- Using ratios used to establish base line measurements for the QCD critical point search.
- **This talk: C_6 / C_2 and C_4 / C_2 .**



Results (I): Energy dependence of product moments

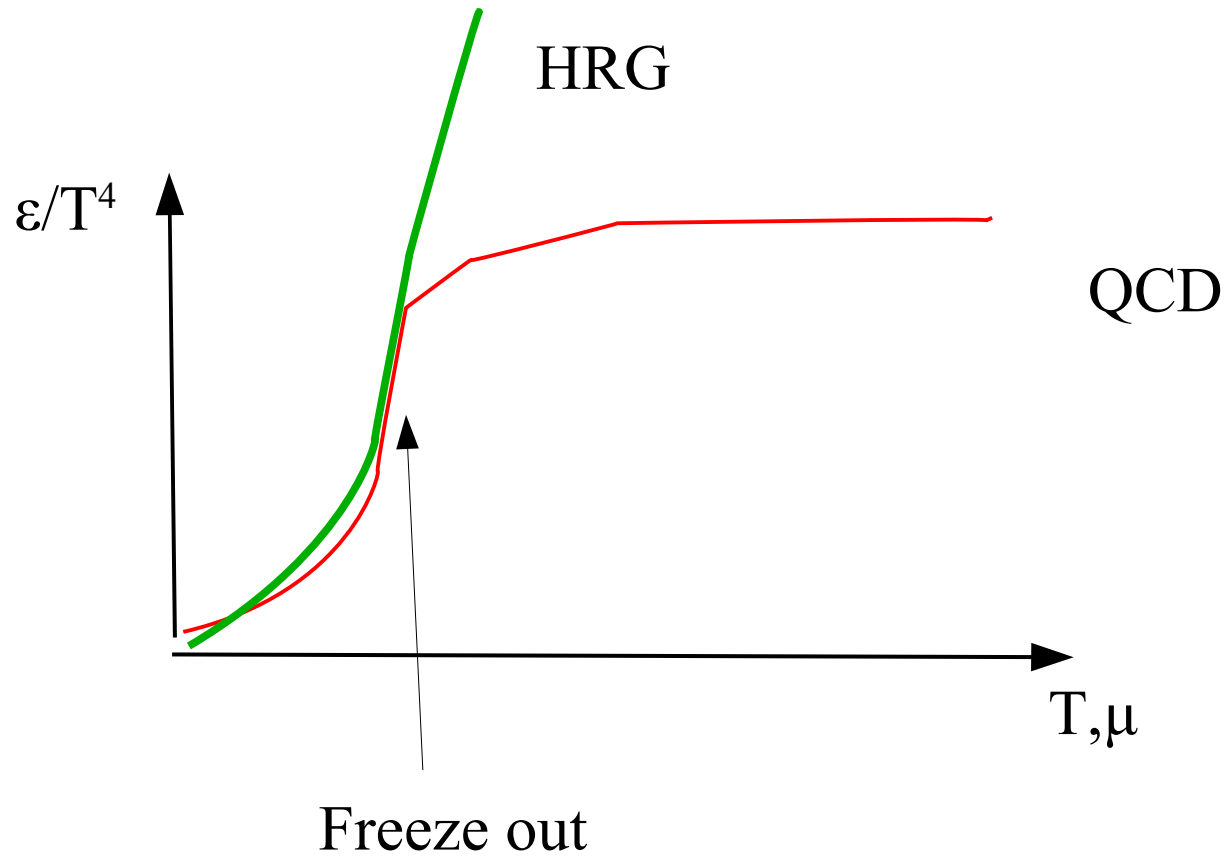


➤ Data are compared with HRG and Lattice QCD.

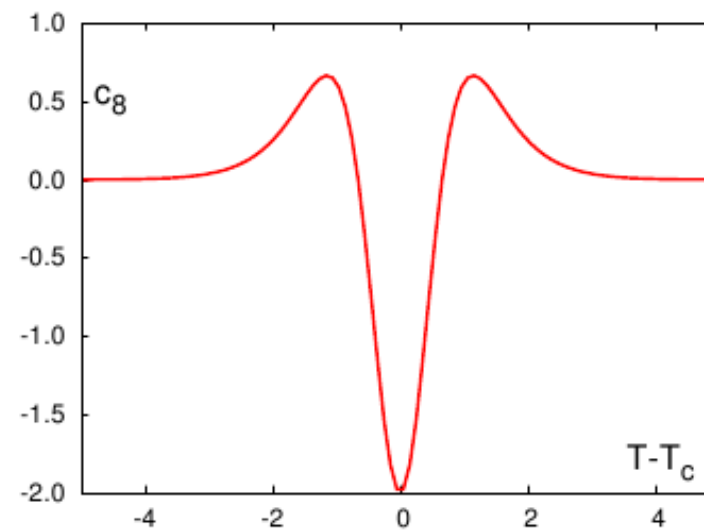
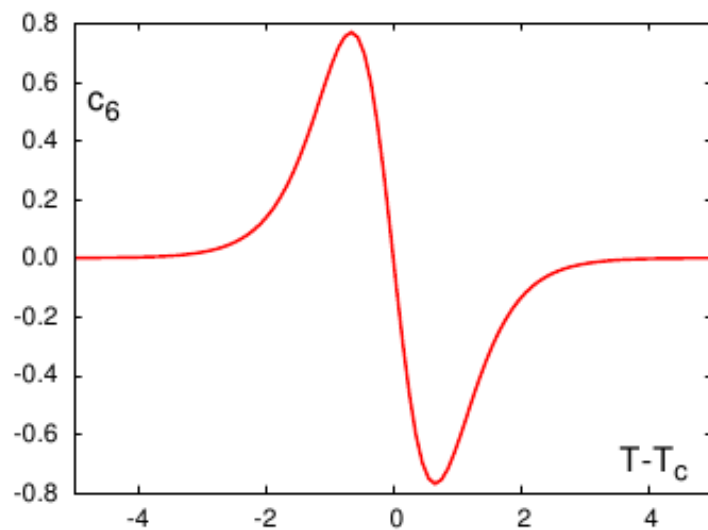
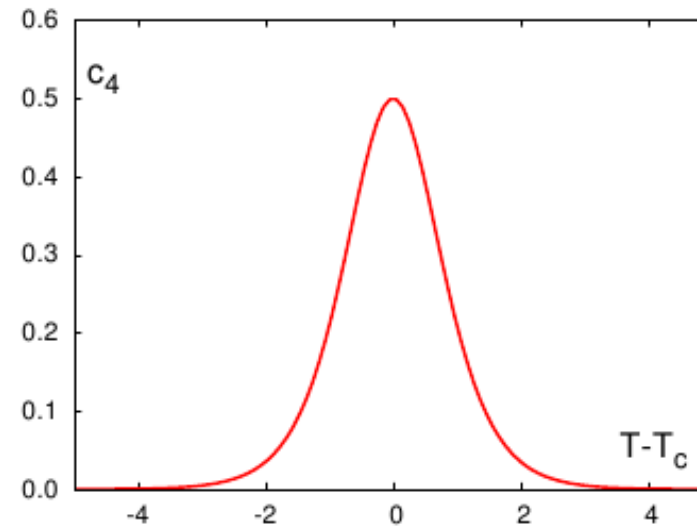
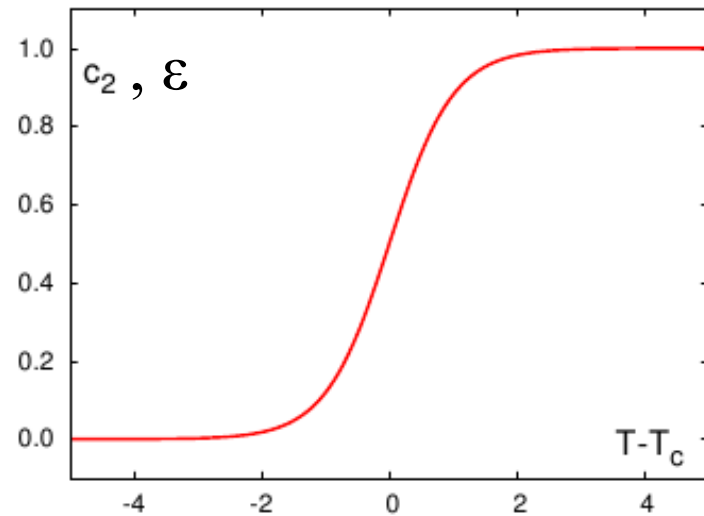
X. Luo, SQM 2011

Another use

- Study higher moments at large \sqrt{s}
- Establish that EOS bends over



Generic expansion coefficients



similar in PNJL model: S. Roessner et al, PR D75 (2007) 034007

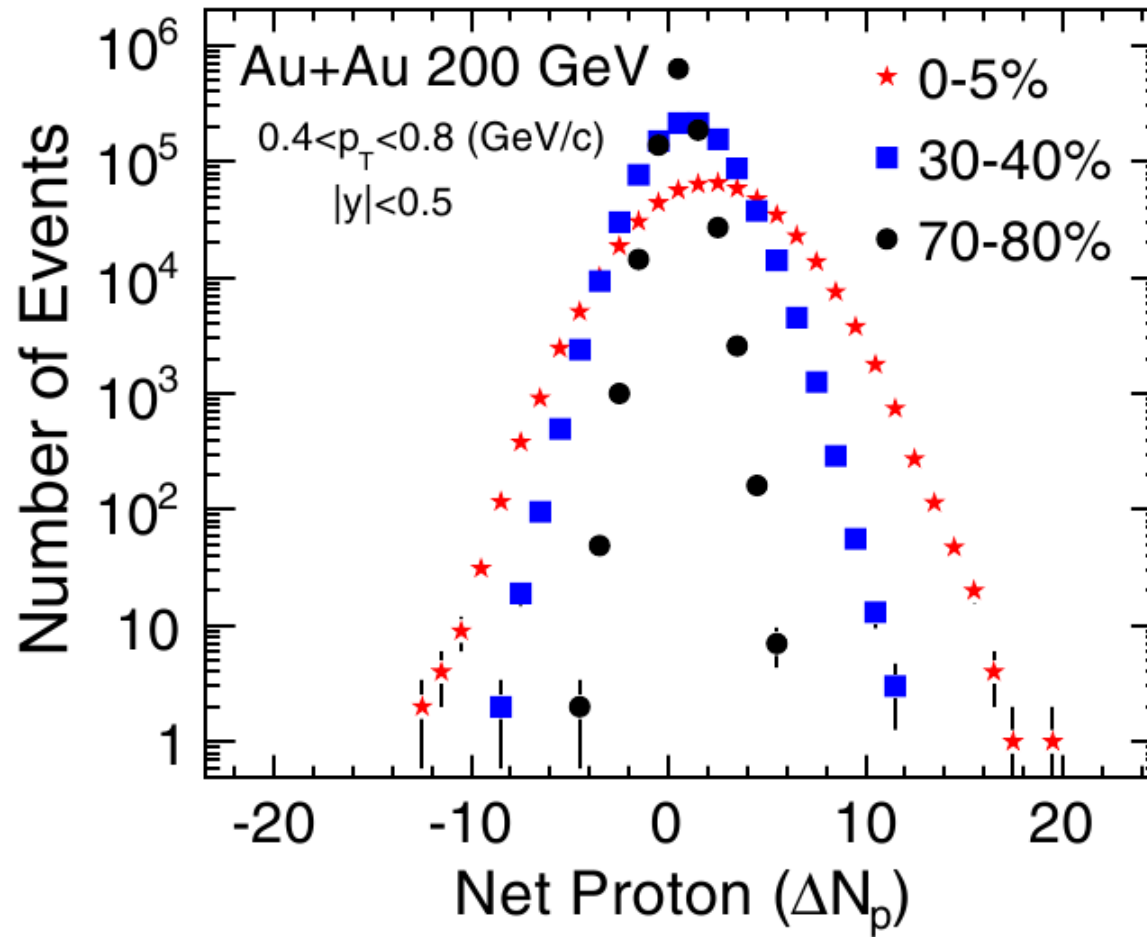
Higher cumulants can tell us a lot!

Details

D-



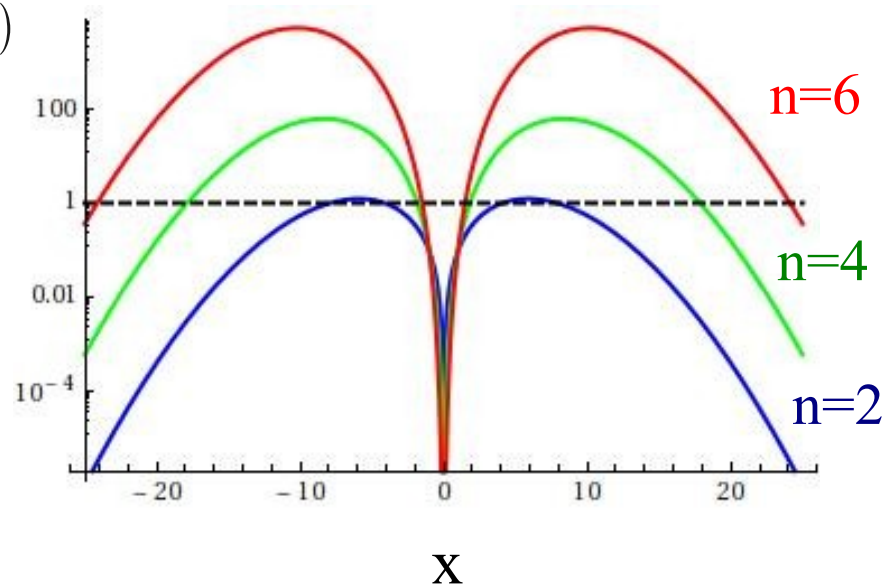
The tails matter



$$20^4 = 1.6 \cdot 10^5$$
$$20^6 = 6.4 \cdot 10^7$$

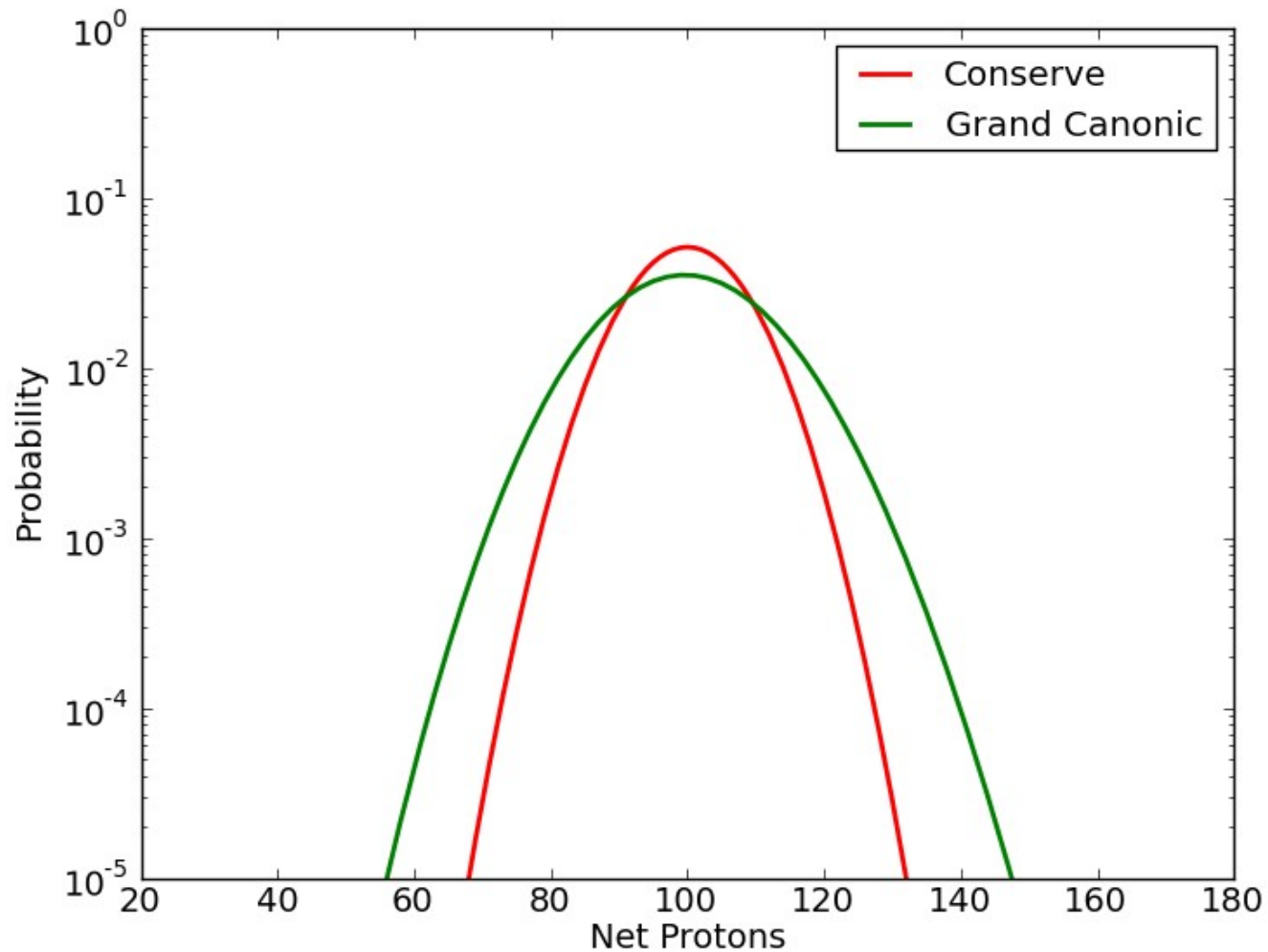
Higher Moments probe the tails

$$P_n(x) = x^n \text{Prob}(x)$$



Gaussian with width = 4.2,
corresponds roughly to STAR net proton distribution at 200 GeV

Conservation laws affect the tails



Baryon number conservation

- Affects all susceptibilities: Variance, Kurtosis....
- Proton Fluctuations are also affected
 - Distinguish from Isospin fluctuations
- Still LARGE baryon **DENSITY** fluctuations

Protons vs Baryons

All protons are baryons but not all baryons are protons

See also talks by Asakawa and Nahrgang

Partition Function

Consider only neutrons, protons, and charged “pions”

Conserved quantities: electric Charge Q and baryon number B

Still fluctuations: Protons \leftrightarrow neutrons and $\pi^+ \leftrightarrow \pi^-$

$$Z = \sum_{p, \bar{p}, n, \bar{n}, \pi^+, \pi^-} \frac{N_{nucl}^{p+\bar{p}+n+\bar{n}}}{p! \bar{p}! n! \bar{n}!} \frac{N_{pions}^{\pi^++\pi^-}}{\pi^+! \pi^-!} \delta(B - (p+n - \bar{p} - \bar{n})) \delta(Q - (p + \pi^+ - \bar{p} - \pi^-))$$

$$N_{nucl} = Z_1(Nucleon)$$

Single particle partition function

$$N_{pions} = Z_1(Pion)$$

Ignore strangeness, energy, and momentum conservation

Partition Function

$$Z = \sum_{p, \bar{p}, n, \bar{n}, \pi^+, \pi^-} \frac{N_{nucl}^{p+\bar{p}+n+\bar{n}}}{p! \bar{p}! n! \bar{n}!} \frac{N_{pions}^{\pi^++\pi^-}}{\pi^+! \pi^-!} \delta(B - (p+n-\bar{p}-\bar{n})) \delta(Q - (p+\pi^+-\bar{p}-\pi^-))$$

$$\Pi = p - \bar{p}$$

$$\nu = n - \bar{n}$$

$$q = \pi^+ - \pi^-$$

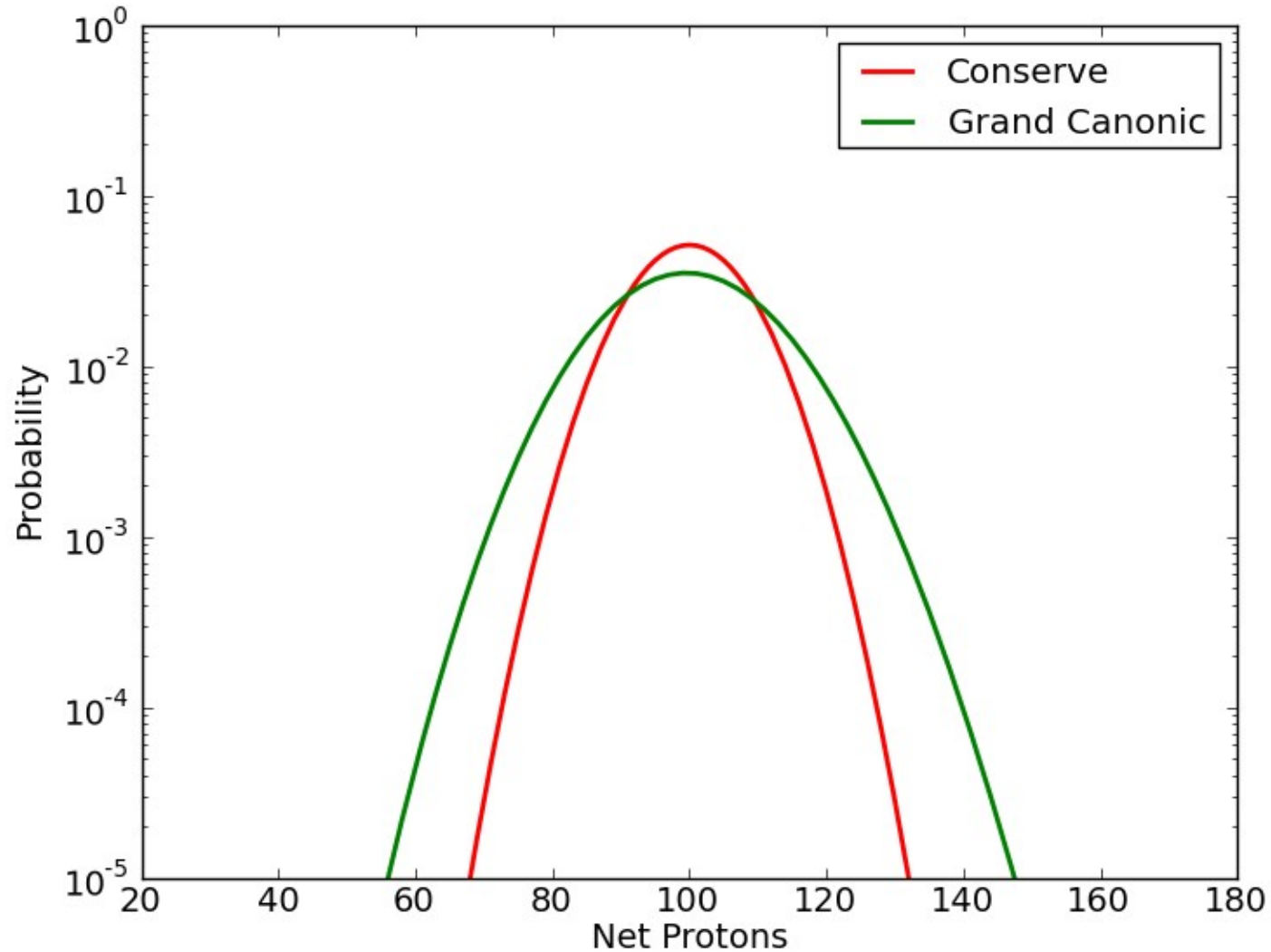
$$Z = \sum_{\Pi, \nu, q} I_{\Pi}(2N_{nucl}) I_{\nu}(2N_{nucl}) I_q(2N_{pions}) \delta(B - \Pi - \nu) \delta(Q - \Pi - q)$$

$$Z = \sum_{\Pi} I_{\Pi}(2N_{nucl}) I_{B-\Pi}(2N_{nucl}) I_{q-\Pi}(2N_{pions})$$

Probability distribution

$$P(\Pi) = \frac{I_{\Pi}(2N_{nucl}) I_{B-\Pi}(2N_{nucl}) I_{q-\Pi}(2N_{pions})}{Z}$$

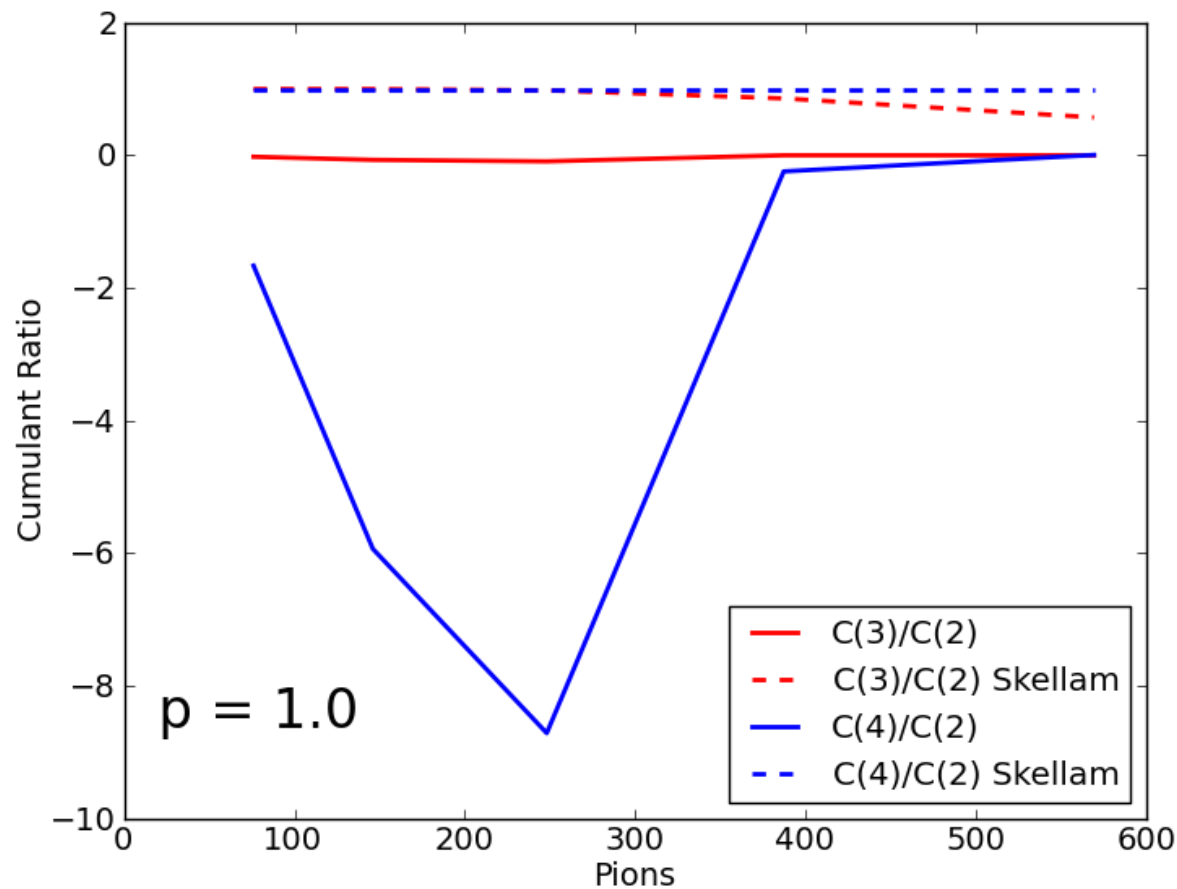
Probability distribution



$B=200$, $Q=100$, $Z_1(\text{nucleons})=40$, $Z_1(\text{pion})=500$

“Temperature Dependence”

Temperature (beam energy) dependence enters through the single particle Partition functions $Z_1(\text{pion})$ and $Z_1(\text{nucleon})$

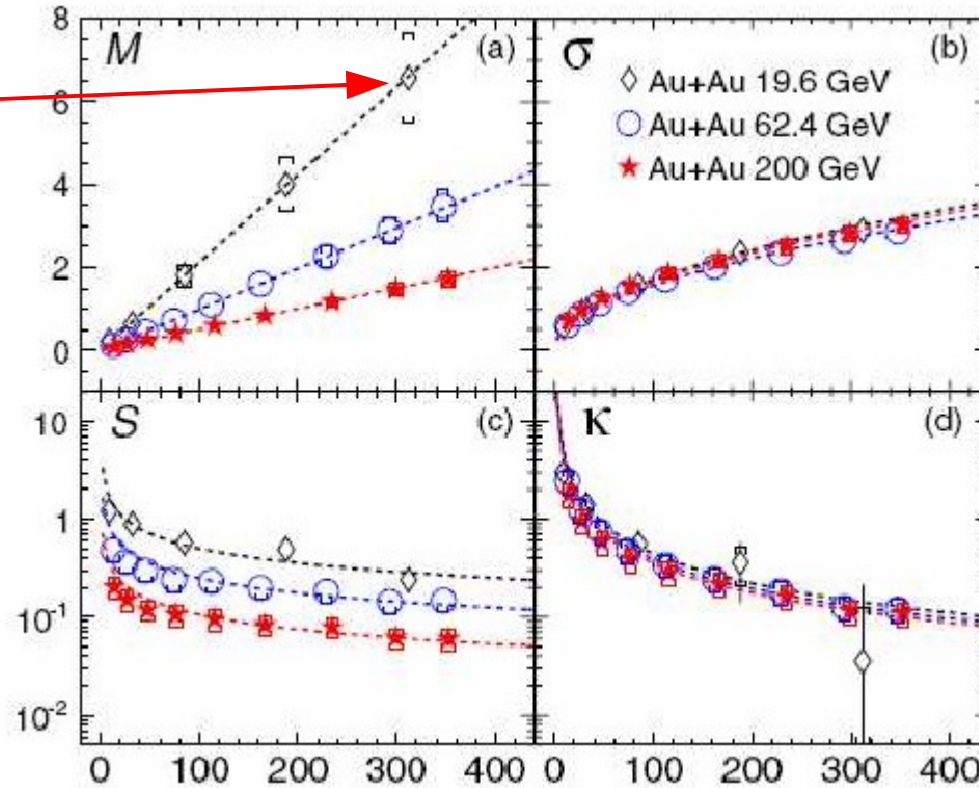


T=100 MeV

T=180 MeV

Finite acceptance

8 protons!!!



$$\sigma^2 \approx N_p + N_{\bar{p}}$$
$$\Rightarrow N_p + N_{\bar{p}} \approx \text{const}$$

Finite Acceptance

Effect of conservation laws get reduced with finite acceptance:
“Equilibration via ignorance”

Model with binomial distribution: p = probability to see particle

$$P(n; N) = \frac{N!}{n!(N-n)!} p^n (1-p)^{N-n}$$

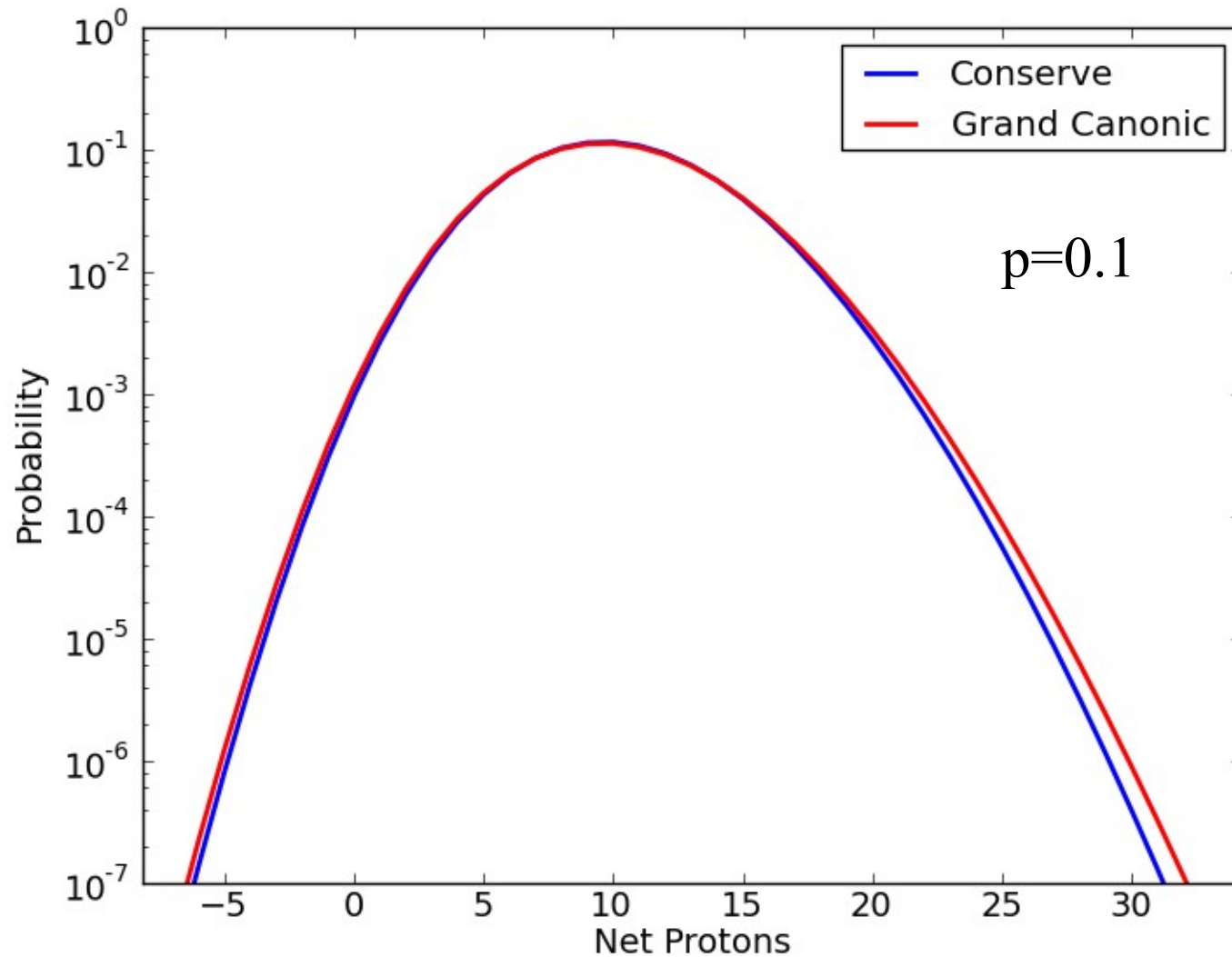
$$C_{2,bin} = p\mu + p^2 (-\mu + k_2)$$

$$C_{3,bin} = p\mu + p^2 (-3\mu + 3k_2) + \underline{p^3} (2\mu - 3k_2 + \underline{k_3})$$

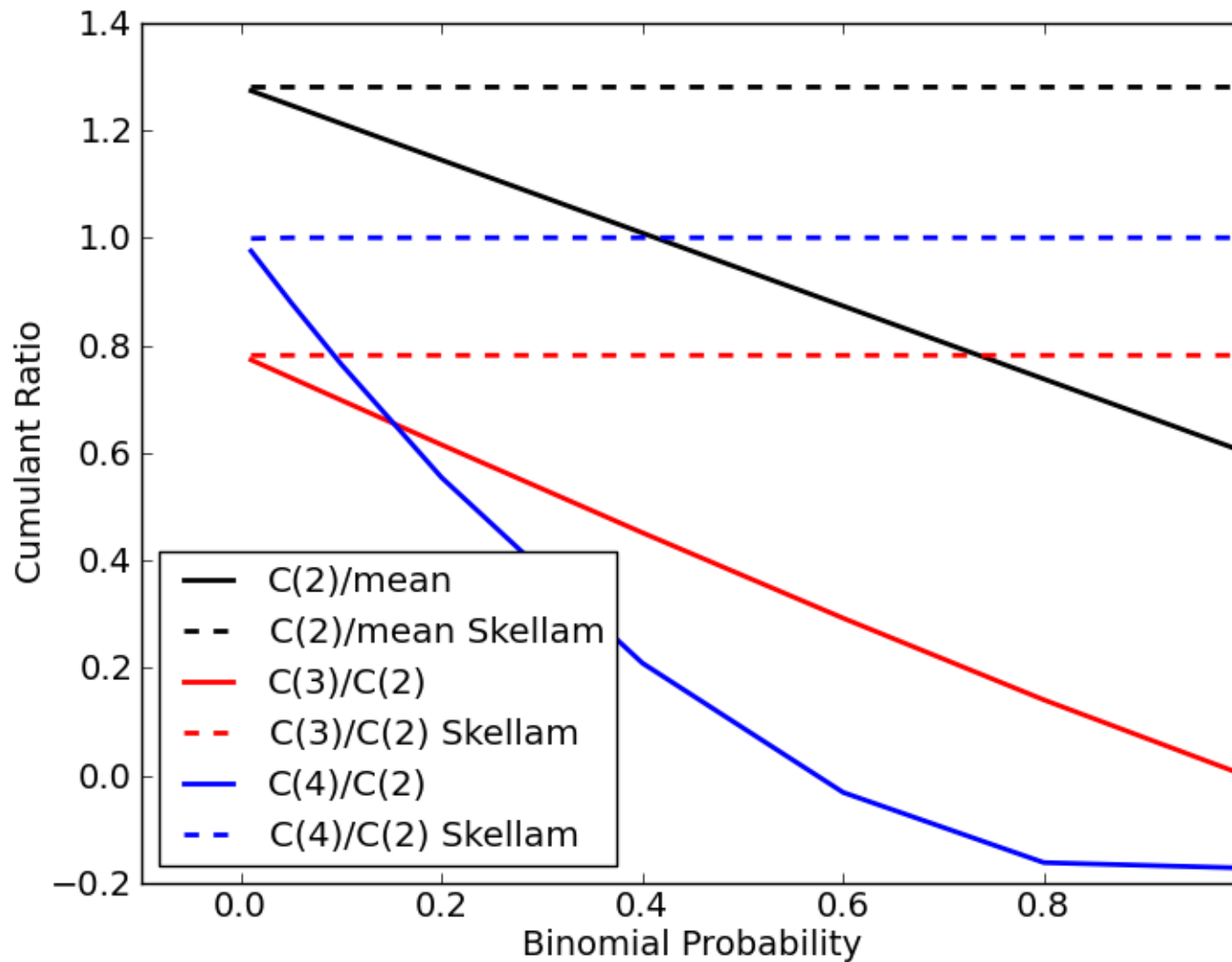
$$C_{4,bin} = p\mu + p^2 (-7\mu + 7k_2) + p^3 (12\mu - 18k_2 + 6k_3) \\ + \underline{p^4} (-6\mu + 11k_2 - 6k_3 + \underline{k_4})$$

μ, k_i = original cumulants

Probability Distribution



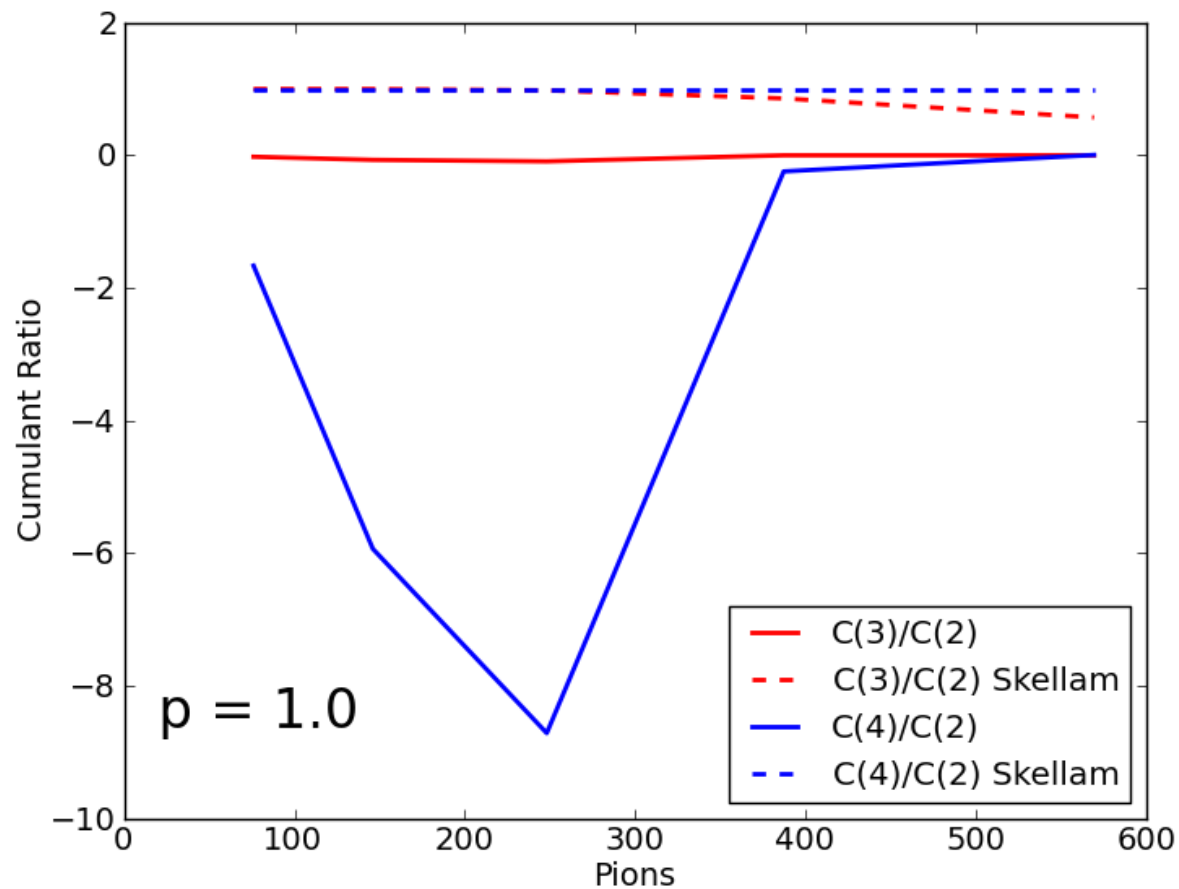
Cumulants



$B=200, Q=100, Z_1(\text{nucleons})=40, Z_1(\text{pion})=500$

“Temperature Dependence”

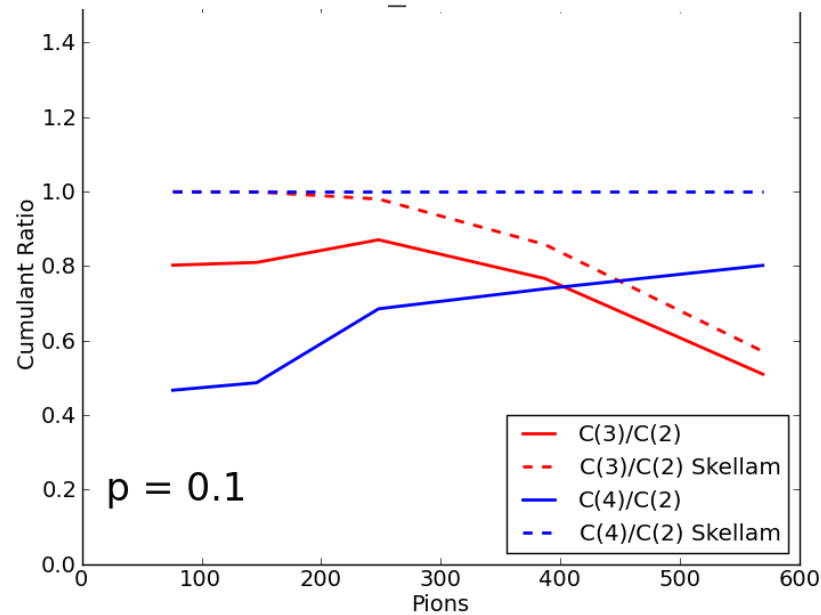
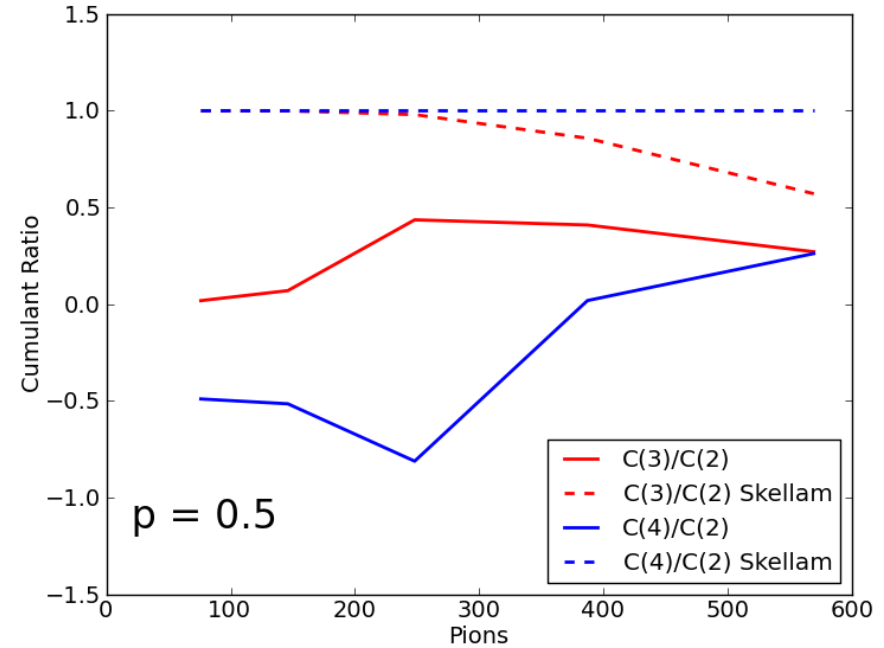
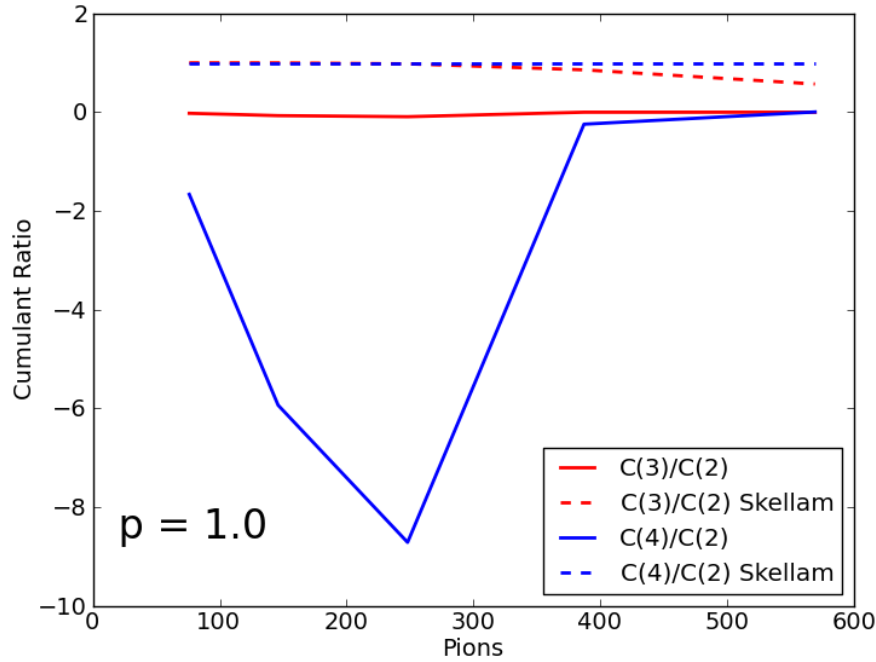
Temperature (beam energy) dependence enters through the single particle Partition functions $Z_1(\text{pion})$ and $Z_1(\text{nucleon})$



T=100 MeV

T=180 MeV

“Temperature Dependence”



Furthermore

- Impact parameter (volume) fluctuations
 - $C(4)/C(2)$ is independent of volume but NOT of volume fluctuations
- “Glauber” fluctuations
 - At high energy large fraction of baryons not part of the system
 - Can be addressed with transport (in principle)
- Can we unfold the effects of charge conservation?
 - Maybe.... at least at high energies.

Summary

- Higher cumulants are very sensitive to “charge” conservation
- Finite acceptance suppresses higher cumulants more strongly
- Not all baryons are protons
- Need to study dependence of e.g. size of rapidity window
- Sixth order cumulants for central Au+Au needs LOTS of statistics

Volume Fluctuations

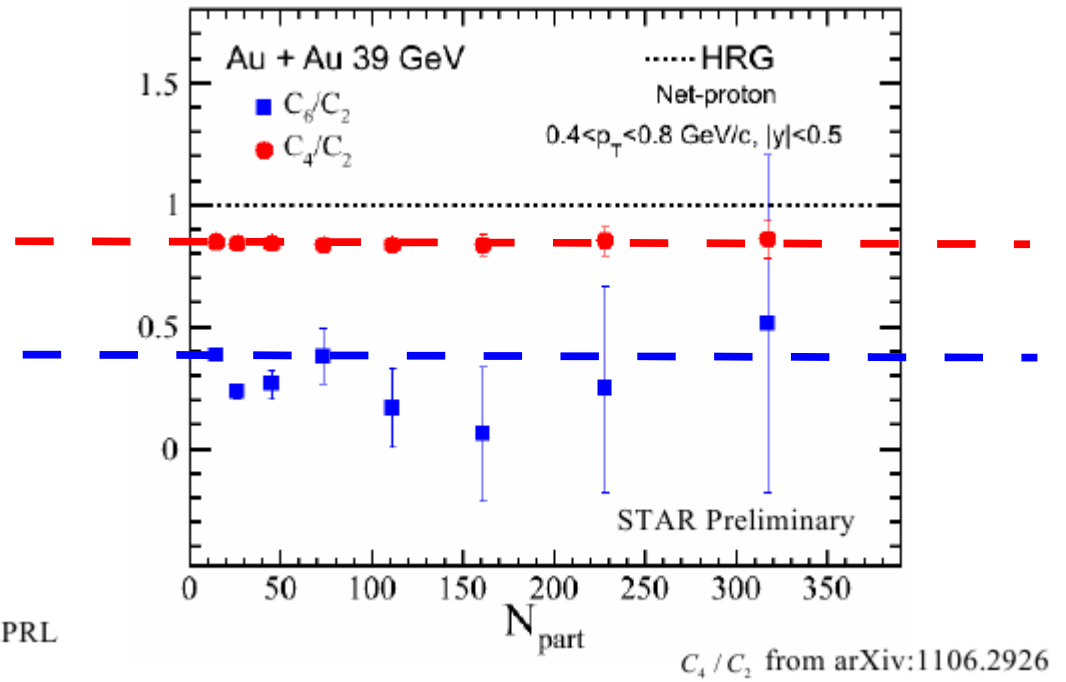
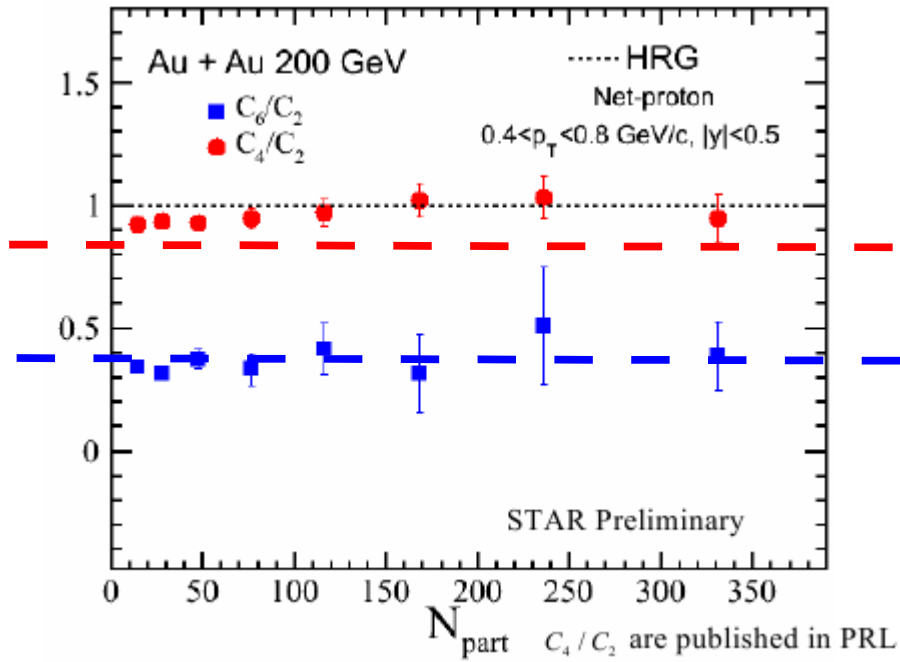
$$N = \rho V$$
$$\delta N = V \delta \rho + \rho \delta V$$

$$\langle (\delta N)^2 \rangle = \langle V \rangle^2 \langle (\delta \rho)^2 \rangle + \langle \rho \rangle^2 \langle (\delta V)^2 \rangle$$

$$\langle (\delta N)^4 \rangle = \langle V \rangle^4 \langle (\delta \rho)^4 \rangle + 6 \langle V \rangle^2 \langle \rho \rangle^2 \langle (\delta V)^2 \rangle \langle (\delta \rho)^2 \rangle + \langle \rho \rangle^4 \langle (\delta V)^4 \rangle$$

$$K \sigma^2 = \frac{\langle (\delta N)^4 \rangle}{\langle (\delta N)^2 \rangle} - 3 \langle (\delta N)^2 \rangle \text{ is independent of the volume but not independent of volume fluctuations}$$

The End



Centrality dependence ???