

A novel method for computing correlators and spectral functions in perturbation theory

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Mikko Laine, Aleksi Vuorinen, YZ, 1108.1259

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Outline

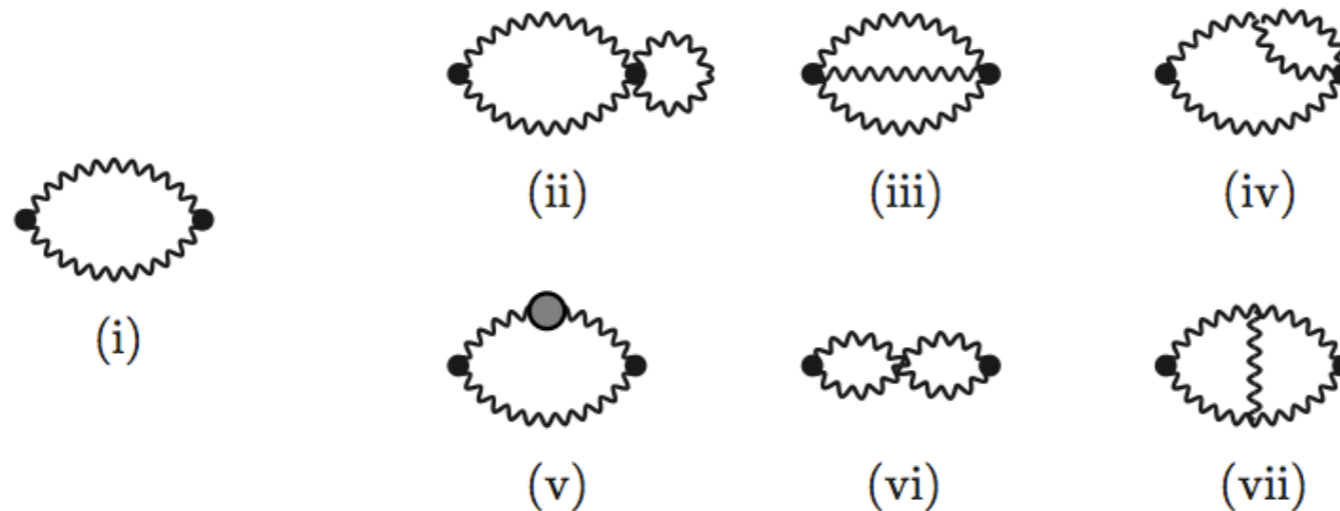
- Introduction & Motivation (Miao and Vuorinen's talk)
- Setup
- Correlators
- Spectral functions
- Summary and outlook

Setup

- Energy-Momentum Tensor $T_{\mu\nu} = \frac{1}{4}\delta_{\mu\nu}F_{\alpha\beta}^a F_{\alpha\beta}^a - F_{\mu\alpha}^a F_{\nu\alpha}^a$,
- Operators: $\theta \equiv c_\theta g_B^2 F_{\mu\nu}^a F_{\mu\nu}^a$, $\chi \equiv c_\chi \epsilon_{\mu\nu\rho\sigma} g_B^2 F_{\mu\nu}^a F_{\rho\sigma}^a$,
- Define
 - $G_\theta(x) \equiv \langle \theta(x)\theta(0) \rangle_c$,
 - $G_\chi(x) \equiv \langle \chi(x)\chi(0) \rangle$,
 - $G_\eta(x) \equiv 2c_\eta^2 X_{\mu\nu,\alpha\beta}(x) \langle T_{\mu\nu}(x) T_{\alpha\beta}(0) \rangle_c$,
 where $X_{\mu\nu,\alpha\beta} \equiv P_{\mu\nu}^T P_{\alpha\beta}^T - \frac{D-2}{2}(P_{\mu\alpha}^T P_{\nu\beta}^T + P_{\mu\beta}^T P_{\nu\alpha}^T)$,
 $G_\eta(x) = -16c_\eta^2 \langle T_{12}(x) T_{12}(0) \rangle_c$.
- We use $D = 4 - 2\epsilon$ to regularize our calculation.

Correlators

The LO and NLO Feynman graphs contributing to the correlators



- Wick contraction .
- Simplify the results to a minimal number of independent “master” sum-integrals.
- Carry out Matsubara Sums and expand them in terms of large P to get the results in UV limit (short distance and/or large frequency).

Results: Mikko Laine, Mikko Vepsäläinen, Aleksi Vuorinen, 1011.4439

York Schröder, Mikko Vepsäläinen, Aleksi Vuorinen, Yan Zhu, 1109.6548

Spectral Functions

$$\rho(\omega) = \text{Im} \left[\tilde{G}(P) \right]_{P \rightarrow (-i[\omega + i0^+], \mathbf{0})} .$$

- After Matsubara Sums, the imaginary part can be extracted with

$$\text{Im} \left[\frac{1}{\omega \pm i0^+} \right] = \mp \pi \delta(\omega) .$$

- Example:

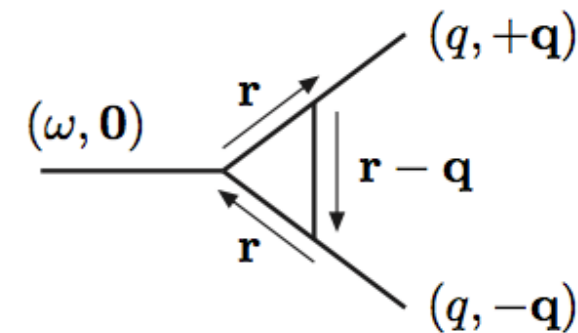
$$\mathcal{I}_j(P) \equiv \oint_{Q,R} \frac{P^6}{Q^2 R^2 [(Q - R)^2 + \lambda^2] (Q - P)^2 (R - P)^2} .$$

Denoting $E_q \equiv q$, $E_r \equiv r$, $E_{qr} \equiv \sqrt{(\mathbf{q} - \mathbf{r})^2 + \lambda^2}$,

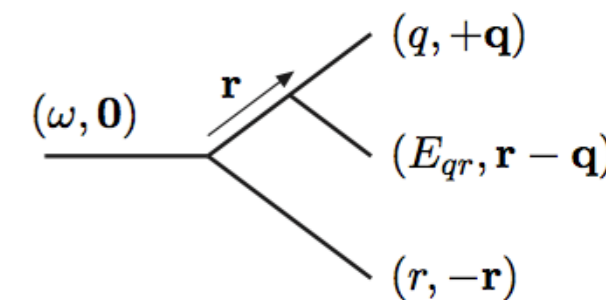
$\rho_{\mathcal{I}_j}(\omega)$

$$\rho_{\mathcal{I}_j}(\omega) = \int_{\mathbf{q}, \mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \left\{ \begin{aligned} & \frac{1}{8q^2} \left[\delta(\omega - 2q) - \delta(\omega + 2q) \right] \times \\ & \times \left[\left(\frac{1}{(q+r-E_{qr})(q+r)} - \frac{1}{(q-r+E_{qr})(q-r)} \right) (1+2n_q)(n_{qr}-n_r) \right. \\ & \left. + \left(\frac{1}{(q+r+E_{qr})(q+r)} - \frac{1}{(q-r-E_{qr})(q-r)} \right) (1+2n_q)(1+n_{qr}+n_r) \right] \\ & + \frac{1}{8r^2} \left[\delta(\omega - 2r) - \delta(\omega + 2r) \right] \times \\ & \times \left[\left(\frac{1}{(q+r-E_{qr})(q+r)} - \frac{1}{(q-r-E_{qr})(q-r)} \right) (1+2n_r)(n_{qr}-n_q) \right. \\ & \left. + \left(\frac{1}{(q+r+E_{qr})(q+r)} - \frac{1}{(q-r+E_{qr})(q-r)} \right) (1+2n_r)(1+n_{qr}+n_q) \right] \\ & + \left[\delta(\omega - q - r - E_{qr}) - \delta(\omega + q + r + E_{qr}) \right] \frac{(1+n_{qr})(1+n_q+n_r) + n_q n_r}{(q+r+E_{qr})^2 (q-r+E_{qr})(q-r-E_{qr})} \\ & + \left[\delta(\omega - q - r + E_{qr}) - \delta(\omega + q + r - E_{qr}) \right] \frac{n_{qr}(1+n_q+n_r) - n_q n_r}{(q+r-E_{qr})^2 (q-r+E_{qr})(q-r-E_{qr})} \\ & + \left[\delta(\omega - q + r - E_{qr}) - \delta(\omega + q - r + E_{qr}) \right] \frac{n_r(1+n_q+n_{qr}) - n_q n_{qr}}{(q-r+E_{qr})^2 (q+r+E_{qr})(q+r-E_{qr})} \\ & + \left[\delta(\omega + q - r - E_{qr}) - \delta(\omega - q + r + E_{qr}) \right] \frac{n_q(1+n_r+n_{qr}) - n_r n_{qr}}{(q-r-E_{qr})^2 (q+r+E_{qr})(q+r-E_{qr})} \end{aligned} \right\}$$

Factorized int./
Virtual correction



Phase space int./
Real correction



Virtual Correction

$$\begin{aligned}
 \rho_{\mathcal{I}_j}^{(\text{fz})}(\omega) &= \int_{\mathbf{q}, \mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \times \\
 &\times \frac{1}{8r^2} \left[\delta(\omega - 2r) - \delta(\omega + 2r) \right] \times \\
 &\times \left[\left(\frac{1}{(q+r-E_{qr})(q+r)} - \frac{1}{(q-r-E_{qr})(q-r)} \right) (1+2n_r)(n_{qr} - n_q) \right. \\
 &\left. + \left(\frac{1}{(q+r+E_{qr})(q+r)} - \frac{1}{(q-r+E_{qr})(q-r)} \right) (1+2n_r)(1+n_{qr} + n_q) \right]
 \end{aligned}$$

$$\star \int_{\mathbf{r}} \pi \delta(\omega - 2r) = \frac{\omega^2}{16\pi},$$

$$\star \int_{\mathbf{q}} \frac{1}{2qE_{qr}} = \frac{1}{4\pi^2\omega} \int_0^\infty dq \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr}, \quad E_{qr}^\pm \equiv \sqrt{\left(q \pm \frac{\omega}{2}\right)^2 + \lambda^2}.$$

$$\star \int_0^\infty dq \int_{\sqrt{(q-\frac{\omega}{2})^2 + \lambda^2}}^{\sqrt{(q+\frac{\omega}{2})^2 + \lambda^2}} dE_{qr} = \int_\lambda^\infty dE_{qr} \int_{|\frac{\omega}{2} - \sqrt{E_{qr}^2 - \lambda^2}|}^{\frac{\omega}{2} + \sqrt{E_{qr}^2 - \lambda^2}} dq.$$

Virtual Correction

$$\begin{aligned}
 \rho_{\mathcal{I}_j}^{(\text{fz})}(\omega) &= \int_{\mathbf{q}, \mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \times \\
 &\times \frac{1}{8r^2} \left[\delta(\omega - 2r) - \delta(\omega + 2r) \right] \times \quad \omega > 0 \\
 &\times \left[\left(\frac{1}{(q+r-E_{qr})(q+r)} - \frac{1}{(q-r-E_{qr})(q-r)} \right) (1+2n_r)(n_{qr} - n_q) \right. \\
 &\quad \left. + \left(\frac{1}{(q+r+E_{qr})(q+r)} - \frac{1}{(q-r+E_{qr})(q-r)} \right) (1+2n_r)(1+n_{qr} + n_q) \right]
 \end{aligned}$$

$$\star \int_{\mathbf{r}} \pi \delta(\omega - 2r) = \frac{\omega^2}{16\pi},$$

$$\star \int_{\mathbf{q}} \frac{1}{2qE_{qr}} = \frac{1}{4\pi^2\omega} \int_0^\infty dq \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr}, \quad E_{qr}^\pm \equiv \sqrt{\left(q \pm \frac{\omega}{2}\right)^2 + \lambda^2}.$$

$$\star \int_0^\infty dq \int_{\sqrt{(q-\frac{\omega}{2})^2 + \lambda^2}}^{\sqrt{(q+\frac{\omega}{2})^2 + \lambda^2}} dE_{qr} = \int_\lambda^\infty dE_{qr} \int_{|\frac{\omega}{2} - \sqrt{E_{qr}^2 - \lambda^2}|}^{\frac{\omega}{2} + \sqrt{E_{qr}^2 - \lambda^2}} dq.$$

Virtual Correction

$$\begin{aligned} \rho_{\mathcal{I}_j}^{(\text{fz})}(\omega) &= \int_{\mathbf{q}, \mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \times \\ &\times \frac{1}{8r^2} \left[\delta(\omega - 2r) - \delta(\omega + 2r) \right] \times \quad \omega > 0 \\ &\times \left[\left(\frac{1}{(q+r-E_{qr})(q+r)} - \frac{1}{(q-r-E_{qr})(q-r)} \right) (1+2n_r)(n_{qr} - n_q) \right. \\ &\left. + \left(\frac{1}{(q+r+E_{qr})(q+r)} - \frac{1}{(q-r+E_{qr})(q-r)} \right) (1+2n_r)(1+n_{qr} + n_q) \right] \end{aligned}$$

$$\star \int_r \pi \delta(\omega - 2r) = \frac{\omega^2}{16\pi},$$

$$\star \int_q \frac{1}{2qE_{qr}} = \frac{1}{4\pi^2\omega} \int_0^\infty dq \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr}, \quad E_{qr}^\pm \equiv \sqrt{\left(q \pm \frac{\omega}{2}\right)^2 + \lambda^2}.$$

$$\star \int_0^\infty dq \int_{\sqrt{(q-\frac{\omega}{2})^2 + \lambda^2}}^{\sqrt{(q+\frac{\omega}{2})^2 + \lambda^2}} dE_{qr} = \int_\lambda^\infty dE_{qr} \int_{|\frac{\omega}{2} - \sqrt{E_{qr}^2 - \lambda^2}|}^{\frac{\omega}{2} + \sqrt{E_{qr}^2 - \lambda^2}} dq.$$

Virtual Correction

$$\begin{aligned}
 \rho_{\mathcal{I}_j}^{(\text{fz})}(\omega) &= \int_{\mathbf{q}, \mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \times \\
 &\times \frac{1}{8r^2} \left[\delta(\omega - 2r) - \delta(\omega + 2r) \right] \times \quad \omega > 0 \\
 &\times \left[\left(\frac{1}{(q+r-E_{qr})(q+r)} - \frac{1}{(q-r-E_{qr})(q-r)} \right) (1+2n_r)(n_{qr} - n_q) \right. \\
 &\quad \left. + \left(\frac{1}{(q+r+E_{qr})(q+r)} - \frac{1}{(q-r+E_{qr})(q-r)} \right) (1+2n_r)(1+n_{qr} + n_q) \right]
 \end{aligned}$$

$$\star \rho_{\mathcal{I}_j}^{(\text{fz,p})}(\omega) \approx \frac{\omega^4}{(4\pi)^3} (1+2n_{\frac{\omega}{2}}) \int_{\lambda}^{\frac{\omega}{2}} \frac{dq}{q} \ln \left| \frac{q + \sqrt{q^2 - \lambda^2}}{q - \sqrt{q^2 - \lambda^2}} \right|.$$

$$\star \int_{\mathbf{q}} \frac{1}{2qE_{qr}} = \frac{1}{4\pi^2\omega} \int_0^\infty dq \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr}, \quad E_{qr}^\pm \equiv \sqrt{\left(q \pm \frac{\omega}{2}\right)^2 + \lambda^2}.$$

$$\star \int_0^\infty dq \int_{\sqrt{(q-\frac{\omega}{2})^2 + \lambda^2}}^{\sqrt{(q+\frac{\omega}{2})^2 + \lambda^2}} dE_{qr} = \int_{\lambda}^\infty dE_{qr} \int_{|\frac{\omega}{2} - \sqrt{E_{qr}^2 - \lambda^2}|}^{\frac{\omega}{2} + \sqrt{E_{qr}^2 - \lambda^2}} dq.$$

Virtual Correction

$$\begin{aligned}
 \rho_{\mathcal{I}_j}^{(\text{fz})}(\omega) &= \int_{\mathbf{q}, \mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \times \\
 &\times \frac{1}{8r^2} \left[\delta(\omega - 2r) - \delta(\omega + 2r) \right] \times \quad \omega > 0 \\
 &\times \left[\left(\frac{1}{(q+r-E_{qr})(q+r)} - \frac{1}{(q-r-E_{qr})(q-r)} \right) (1+2n_r)(n_{qr} - n_q) \right. \\
 &\left. + \left(\frac{1}{(q+r+E_{qr})(q+r)} - \frac{1}{(q-r+E_{qr})(q-r)} \right) (1+2n_r)(1+n_{qr} + n_q) \right]
 \end{aligned}$$

$$\begin{aligned}
 \star \rho_{\mathcal{I}_j}^{(\text{fz,p})}(\omega) &\approx \frac{\omega^4}{(4\pi)^3} (1+2n_{\frac{\omega}{2}}) \int_{\lambda}^{\frac{\omega}{2}} \frac{dq}{q} \ln \left| \frac{q + \sqrt{q^2 - \lambda^2}}{q - \sqrt{q^2 - \lambda^2}} \right|. \\
 \star \rho_{\mathcal{I}_j}^{(\text{fz,e})}(\omega) &= \frac{\omega^4}{(4\pi)^3} (1+2n_{\frac{\omega}{2}}) \left\{ \right. \\
 \star &\int_0^{\infty} dq n_q \mathbb{P} \left[\frac{1}{q + \frac{\omega}{2}} \ln \left| \frac{\lambda^2}{2q\omega - \lambda^2} \right| + \frac{1}{q - \frac{\omega}{2}} \ln \left| \frac{\lambda^2}{2q\omega + \lambda^2} \right| \right] \\
 &\left. + \int_{\lambda}^{\infty} dq n_q \left[\frac{1}{q} \ln \left| \frac{q + \frac{\lambda^2}{\omega} + \sqrt{q^2 - \lambda^2}}{q + \frac{\lambda^2}{\omega} - \sqrt{q^2 - \lambda^2}} \right| + \frac{1}{q} \ln \left| \frac{q - \frac{\lambda^2}{\omega} + \sqrt{q^2 - \lambda^2}}{q - \frac{\lambda^2}{\omega} - \sqrt{q^2 - \lambda^2}} \right| \right] \right\}.
 \end{aligned}$$

Real Correction

$$\begin{aligned}
 \rho_{\mathcal{I}_j}^{(\text{ps})}(\omega) &\equiv \int_{\mathbf{q},\mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \left\{ \right. \\
 &+ \left[\delta(\omega - q - r - E_{qr}) - \delta(\omega + q + r + E_{qr}) \right] \frac{(1 + n_{qr})(1 + n_q + n_r) + n_q n_r}{(q + r + E_{qr})^2 (q - r + E_{qr})(q - r - E_{qr})} \\
 &+ \left[\delta(\omega - q - r + E_{qr}) - \delta(\omega + q + r - E_{qr}) \right] \frac{n_{qr}(1 + n_q + n_r) - n_q n_r}{(q + r - E_{qr})^2 (q - r + E_{qr})(q - r - E_{qr})} \\
 &+ \left[\delta(\omega - q + r - E_{qr}) - \delta(\omega + q - r + E_{qr}) \right] \frac{n_r(1 + n_q + n_{qr}) - n_q n_{qr}}{(q - r + E_{qr})^2 (q + r + E_{qr})(q + r - E_{qr})} \\
 &+ \left. \left[\delta(\omega + q - r - E_{qr}) - \delta(\omega - q + r + E_{qr}) \right] \frac{n_q(1 + n_r + n_{qr}) - n_r n_{qr}}{(q - r - E_{qr})^2 (q + r + E_{qr})(q + r - E_{qr})} \right\}.
 \end{aligned}$$

$$\star \int_{\mathbf{q},\mathbf{r}} \frac{\pi}{4qr E_{qr}} = \frac{2}{(4\pi)^3} \int_0^\infty dq \int_0^\infty dr \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr}, \quad E_{qr}^\pm \equiv \sqrt{(q \pm r)^2 + \lambda^2}.$$

$$\begin{aligned}
 \star (1 + n_{qr})(1 + n_q + n_r) + n_q n_r &= n_q n_r n_{qr} (e^{q+r+E_{qr}} - 1), \\
 n_{qr}(1 + n_q + n_r) - n_q n_r &= n_q n_r n_{qr} (e^{q+r} - e^{E_{qr}}), \\
 n_r(1 + n_q + n_{qr}) - n_q n_{qr} &= n_q n_r n_{qr} (e^{q+E_{qr}} - e^r), \\
 n_q(1 + n_r + n_{qr}) - n_r n_{qr} &= n_q n_r n_{qr} (e^{r+E_{qr}} - e^q),
 \end{aligned}$$

Real Correction

$$\begin{aligned}
 \rho_{\mathcal{I}_j}^{(\text{ps})}(\omega) &\equiv \int_{\mathbf{q},\mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \left\{ \begin{array}{l} 0 < \lambda < \omega \\ + \left[\delta(\omega - q - r - E_{qr}) - \delta(\omega + q + r + E_{qr}) \right] \frac{(1 + n_{qr})(1 + n_q + n_r) + n_q n_r}{(q + r + E_{qr})^2 (q - r + E_{qr})(q - r - E_{qr})} \\ + \left[\delta(\omega - q - r + E_{qr}) - \delta(\omega + q + r - E_{qr}) \right] \frac{n_{qr}(1 + n_q + n_r) - n_q n_r}{(q + r - E_{qr})^2 (q - r + E_{qr})(q - r - E_{qr})} \\ + \left[\delta(\omega - q + r - E_{qr}) - \delta(\omega + q - r + E_{qr}) \right] \frac{n_r(1 + n_q + n_{qr}) - n_q n_{qr}}{(q - r + E_{qr})^2 (q + r + E_{qr})(q + r - E_{qr})} \\ + \left[\delta(\omega + q - r - E_{qr}) - \delta(\omega - q + r + E_{qr}) \right] \frac{n_q(1 + n_r + n_{qr}) - n_r n_{qr}}{(q - r - E_{qr})^2 (q + r + E_{qr})(q + r - E_{qr})} \end{array} \right\}.
 \end{aligned}$$

$$\star \int_{\mathbf{q},\mathbf{r}} \frac{\pi}{4qr E_{qr}} = \frac{2}{(4\pi)^3} \int_0^\infty dq \int_0^\infty dr \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr}, \quad E_{qr}^\pm \equiv \sqrt{(q \pm r)^2 + \lambda^2}.$$

$$\begin{aligned}
 \star (1 + n_{qr})(1 + n_q + n_r) + n_q n_r &= n_q n_r n_{qr} (e^{q+r+E_{qr}} - 1), \\
 n_{qr}(1 + n_q + n_r) - n_q n_r &= n_q n_r n_{qr} (e^{q+r} - e^{E_{qr}}), \\
 n_r(1 + n_q + n_{qr}) - n_q n_{qr} &= n_q n_r n_{qr} (e^{q+E_{qr}} - e^r), \\
 n_q(1 + n_r + n_{qr}) - n_r n_{qr} &= n_q n_r n_{qr} (e^{r+E_{qr}} - e^q),
 \end{aligned}$$

Real Correction

$$\begin{aligned}
 \rho_{\mathcal{I}_j}^{(\text{ps})}(\omega) &\equiv \int_{\mathbf{q},\mathbf{r}} \frac{\omega^6 \pi}{4qr E_{qr}} \left\{ \begin{array}{l} 0 < \lambda < \omega \\ + \left[\delta(\omega - q - r - E_{qr}) - \delta(\omega + q + r + E_{qr}) \right] \frac{(1 + n_{qr})(1 + n_q + n_r) + n_q n_r}{(q + r + E_{qr})^2 (q - r + E_{qr})(q - r - E_{qr})} \\ + \left[\delta(\omega - q - r + E_{qr}) - \delta(\omega + q + r - E_{qr}) \right] \frac{n_{qr}(1 + n_q + n_r) - n_q n_r}{(q + r - E_{qr})^2 (q - r + E_{qr})(q - r - E_{qr})} \\ + \left[\delta(\omega - q + r - E_{qr}) - \delta(\omega + q - r + E_{qr}) \right] \frac{n_r(1 + n_q + n_{qr}) - n_q n_{qr}}{(q - r + E_{qr})^2 (q + r + E_{qr})(q + r - E_{qr})} \\ + \left[\delta(\omega + q - r - E_{qr}) - \delta(\omega - q + r + E_{qr}) \right] \frac{n_q(1 + n_r + n_{qr}) - n_r n_{qr}}{(q - r - E_{qr})^2 (q + r + E_{qr})(q + r - E_{qr})} \end{array} \right\}.
 \end{aligned}$$

$$\star \int_{\mathbf{q},\mathbf{r}} \frac{\pi}{4qr E_{qr}} = \frac{2}{(4\pi)^3} \int_0^\infty dq \int_0^\infty dr \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr}, \quad E_{qr}^\pm \equiv \sqrt{(q \pm r)^2 + \lambda^2}.$$

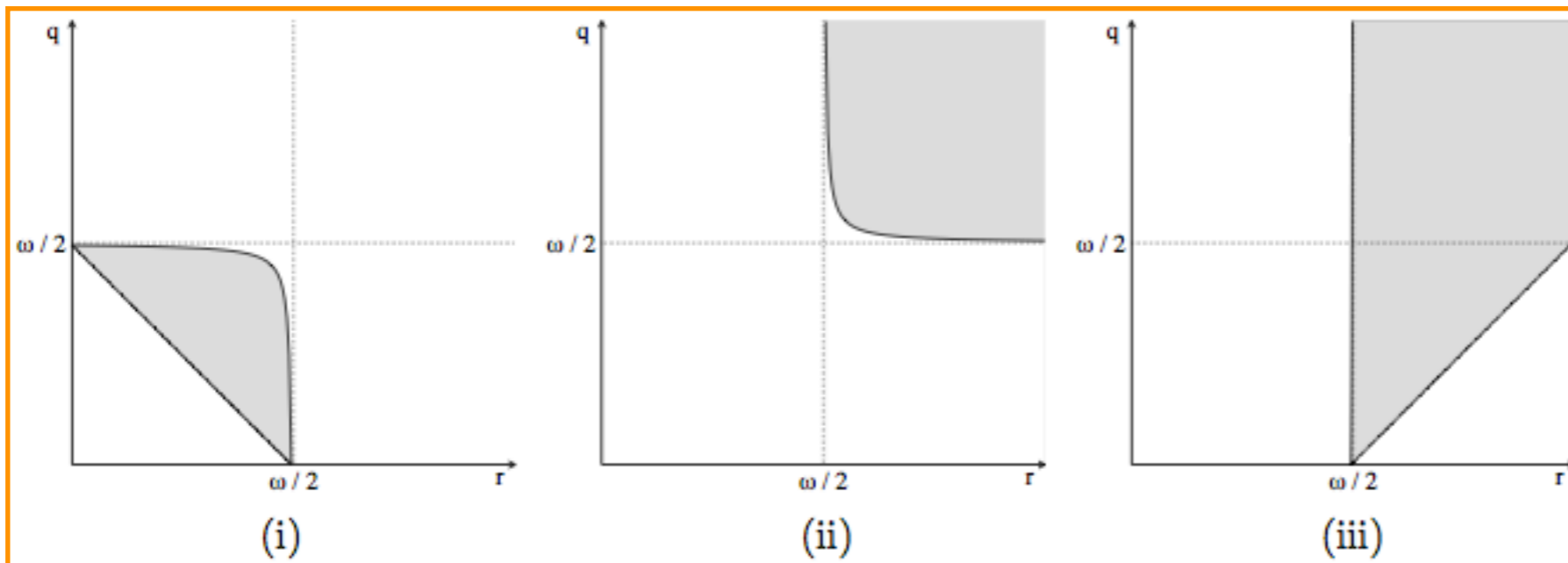
$$\begin{aligned}
 \star (1 + n_{qr})(1 + n_q + n_r) + n_q n_r &= n_q n_r n_{qr} (e^{q+r+E_{qr}} - 1), \\
 n_{qr}(1 + n_q + n_r) - n_q n_r &= n_q n_r n_{qr} (e^{q+r} - e^{E_{qr}}), \\
 n_r(1 + n_q + n_{qr}) - n_q n_{qr} &= n_q n_r n_{qr} (e^{q+E_{qr}} - e^r), \\
 n_q(1 + n_r + n_{qr}) - n_r n_{qr} &= n_q n_r n_{qr} (e^{r+E_{qr}} - e^q),
 \end{aligned}$$

Real Correction

$$\rho_{I_j}^{(ps)}(\omega) \equiv \frac{2\omega^4}{(4\pi)^3} \int_0^\infty dq \int_0^\infty dr \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr} n_q n_r n_{qr} \left\{ \begin{array}{l} \text{(i)} \quad \frac{\delta(\omega - q - r - E_{qr})}{(2r - \omega)(2q - \omega)} (1 - e^{q+r+E_{qr}}) \\ \text{(ii)} \quad + \frac{\delta(\omega - q - r + E_{qr})}{(2r - \omega)(2q - \omega)} (e^{E_{qr}} - e^{q+r}) \\ \text{(iii)} \quad + \frac{\delta(\omega + q - r - E_{qr})}{(2r - \omega)(2q + \omega)} (e^{r+E_{qr}} - e^q) \\ \text{(iv)} \quad + \frac{\delta(\omega - q + r - E_{qr})}{(2r + \omega)(2q - \omega)} (e^{q+E_{qr}} - e^r) \end{array} \right\}.$$

$$\begin{array}{l} \text{(i)} \quad \int_0^\infty dq \int_0^\infty dr \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr} \delta(\omega - q - r - E_{qr}) \phi(q, r, E_{qr}) \\ \quad = \int_0^{\frac{\omega^2 - \lambda^2}{2\omega}} dq \int_{\frac{\omega(\omega - 2q) - \lambda^2}{2(\omega - 2q)}}^{\frac{\omega(\omega - 2q) - \lambda^2}{2\omega}} dr \phi(q, r, \omega - q - r), \\ \text{(ii)} \quad \int_0^\infty dq \int_0^\infty dr \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr} \delta(\omega - q - r + E_{qr}) \phi(q, r, E_{qr}) \\ \quad = \int_{\frac{\omega}{2}}^\infty dq \int_{\frac{\omega(2q - \omega) + \lambda^2}{2(2q - \omega)}}^\infty dr \phi(q, r, -\omega + q + r), \\ \text{(iii)} \quad \int_0^\infty dq \int_0^\infty dr \int_{E_{qr}^-}^{E_{qr}^+} dE_{qr} \delta(\omega + q - r - E_{qr}) \phi(q, r, E_{qr}) \\ \quad = \int_\omega^\infty dq \int_{\frac{\omega(\omega + 2q) - \lambda^2}{2(\omega + 2q)}}^{\frac{\omega(\omega + 2q) - \lambda^2}{2\omega}} dr \phi(q, r, \omega + q - r), \end{array}$$

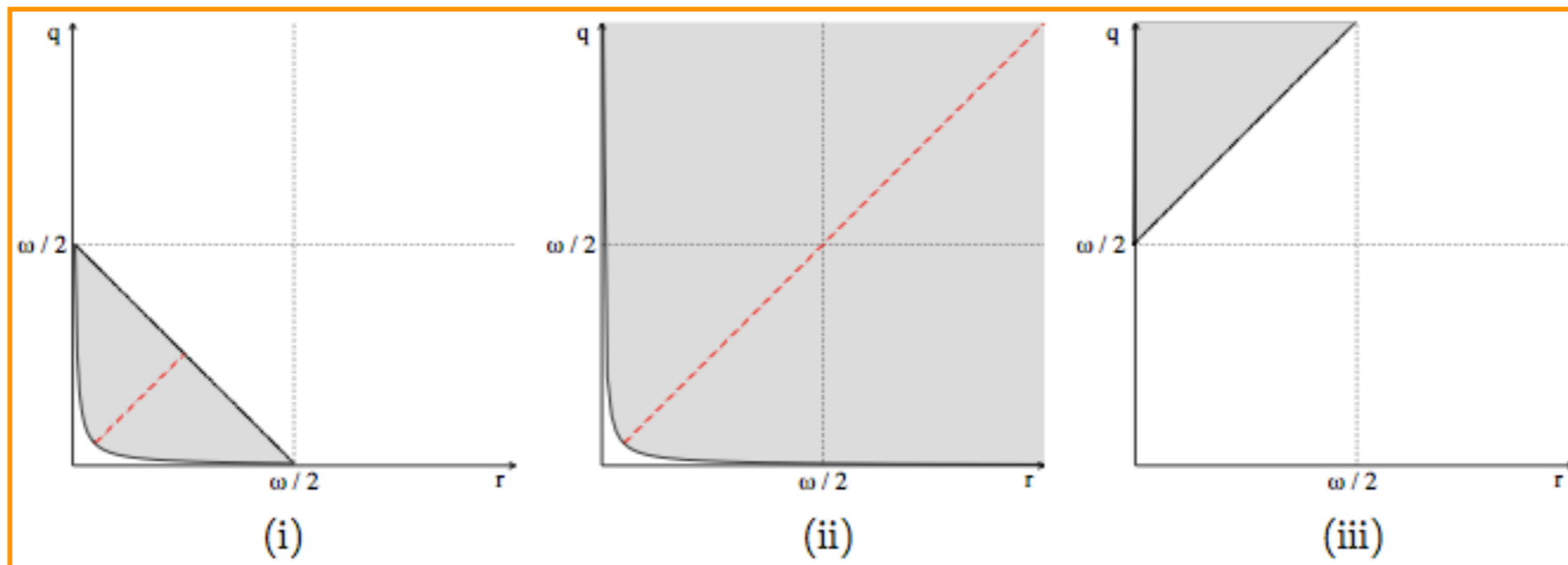
Original integration ranges for $\lambda = \omega/10$.



Real Correction

For turning all denominators into $1/(4qr)$:

$$\begin{aligned} \text{(i)} : \quad & q \rightarrow \frac{\omega}{2} - q, \quad r \rightarrow \frac{\omega}{2} - r, \\ \text{(ii)} : \quad & q \rightarrow \frac{\omega}{2} + q, \quad r \rightarrow \frac{\omega}{2} + r, \\ \text{(iii)} : \quad & q \rightarrow -\frac{\omega}{2} + q, \quad r \rightarrow \frac{\omega}{2} + r, \end{aligned}$$

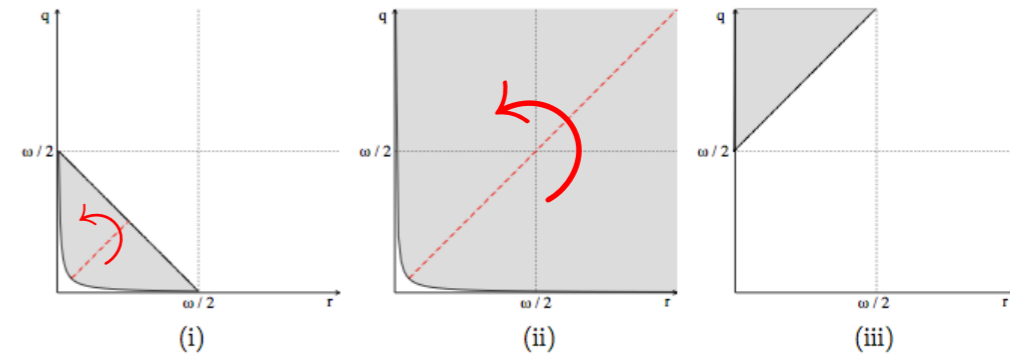


Integration ranges after the shifts, for $\lambda = \omega/10$.

Real Correction

$$\rho_{I_j}^{(ps)}(\omega) = \frac{\omega^4}{2(4\pi)^3} (e^\omega - 1) \left\{ \begin{array}{l} \text{(i)} \quad - \int_{\frac{\lambda^2}{2\omega}}^{\frac{\omega}{2}} dq \int_{\frac{\lambda^2}{4q}}^{\frac{\omega(\omega-2q)+\lambda^2}{2\omega}} dr \mathbb{P}\left(\frac{1}{qr}\right) n_{\frac{\omega}{2}-q} n_{\frac{\omega}{2}-r} n_{q+r} \\ \text{(ii)} \quad - \int_0^\infty dq \int_{\frac{\lambda^2}{4q}}^\infty dr \mathbb{P}\left(\frac{1}{qr}\right) n_{\frac{\omega}{2}+q} n_{\frac{\omega}{2}+r} n_{q+r} e^{q+r} \\ \text{(iii)} \quad + \int_{\frac{\omega}{2}}^\infty dq \int_{-\frac{\lambda^2}{4q}}^{\frac{\omega(2q-\omega)-\lambda^2}{2\omega}} dr \mathbb{P}\left(\frac{2}{qr}\right) n_{q-\frac{\omega}{2}} n_{\frac{\omega}{2}+r} n_{q-r} e^{q-\frac{\omega}{2}} \end{array} \right\}.$$

Divergent!



(i)	$n_{\frac{\omega}{2}-q} n_{\frac{\omega}{2}-r} n_{q+r} (1 - e^\omega) = -(1 + 2n_{\frac{\omega}{2}}) \left[1 + n_{q+r} + n_{\frac{\omega}{2}-q} + (1 + n_{\frac{\omega}{2}-r}) \frac{n_{q+r} n_{\frac{\omega}{2}-q}}{n_r^2} \right].$
(ii)	$n_{\frac{\omega}{2}+q} n_{\frac{\omega}{2}+r} n_{q+r} e^{q+r} (1 - e^\omega) = (1 + 2n_{\frac{\omega}{2}}) \left[-n_{q+r} + n_{q+\frac{\omega}{2}} - (1 + n_{q+\frac{\omega}{2}}) \frac{n_{q+r} n_{r+\frac{\omega}{2}}}{n_r^2} \right].$
(iii)	$n_{q-\frac{\omega}{2}} n_{\frac{\omega}{2}+r} n_{q-r} e^{q-\frac{\omega}{2}} (e^\omega - 1) = (1 + 2n_{\frac{\omega}{2}}) \left[n_{q-\frac{\omega}{2}} - n_q - n_{q-\frac{\omega}{2}} \frac{(1 + n_{q-r})(n_q - n_{r+\frac{\omega}{2}})}{n_r n_{-\frac{\omega}{2}}} \right].$

For integrals:

- Depend on $q+r$ only, we can integrate them by integrating $(x,y)=(q+r,q-r)$,
- Independent on r , one can integrate r first.

$$\rho \mathcal{I}_j(\omega)$$

- Collect every part together and simplify them with $\lambda \ll \omega$,
- All the divergent terms cancel each other, we can set $\lambda \rightarrow 0$ in the end.

$$\begin{aligned} & \frac{(4\pi)^3 \rho \mathcal{I}_j(\omega)}{\omega^4(1+2n_{\frac{\omega}{2}})} = \\ & \int_0^{\frac{\omega}{4}} dq n_q \left[\left(\frac{1}{q - \frac{\omega}{2}} - \frac{1}{q} \right) \ln\left(1 - \frac{2q}{\omega}\right) - \frac{\frac{\omega}{2}}{q(q + \frac{\omega}{2})} \ln\left(1 + \frac{2q}{\omega}\right) \right] \\ & + \int_{\frac{\omega}{4}}^{\frac{\omega}{2}} dq n_q \left[\left(\frac{2}{q - \frac{\omega}{2}} - \frac{1}{q} \right) \ln\left(1 - \frac{2q}{\omega}\right) - \frac{\frac{\omega}{2}}{q(q + \frac{\omega}{2})} \ln\left(1 + \frac{2q}{\omega}\right) - \frac{1}{q - \frac{\omega}{2}} \ln\left(\frac{2q}{\omega}\right) \right] \\ & + \int_{\frac{\omega}{2}}^{\infty} dq n_q \left[\left(\frac{2}{q - \frac{\omega}{2}} - \frac{2}{q} \right) \ln\left(\frac{2q}{\omega} - 1\right) - \frac{\frac{\omega}{2}}{q(q + \frac{\omega}{2})} \ln\left(1 + \frac{2q}{\omega}\right) + \left(\frac{1}{q} - \frac{1}{q - \frac{\omega}{2}} \right) \ln\left(\frac{2q}{\omega}\right) \right] \\ & + \int_0^{\frac{\omega}{2}} dq \int_0^{\frac{\omega}{4} - |q - \frac{\omega}{4}|} dr \left(-\frac{1}{qr} \right) \frac{n_{\frac{\omega}{2} - q} n_{q+r} (1 + n_{\frac{\omega}{2} - r})}{n_r^2} \\ & + \int_{\frac{\omega}{2}}^{\infty} dq \int_0^{q - \frac{\omega}{2}} dr \left(-\frac{1}{qr} \right) \frac{n_{q - \frac{\omega}{2}} (1 + n_{q-r}) (n_q - n_{r + \frac{\omega}{2}})}{n_r n_{-\frac{\omega}{2}}} \\ & + \int_0^{\infty} dq \int_0^q dr \left(-\frac{1}{qr} \right) \frac{(1 + n_{q + \frac{\omega}{2}}) n_{q+r} n_{r + \frac{\omega}{2}}}{n_r^2} + \mathcal{O}(\lambda \ln \lambda). \end{aligned}$$

Numerical Calculation

Complex integral function

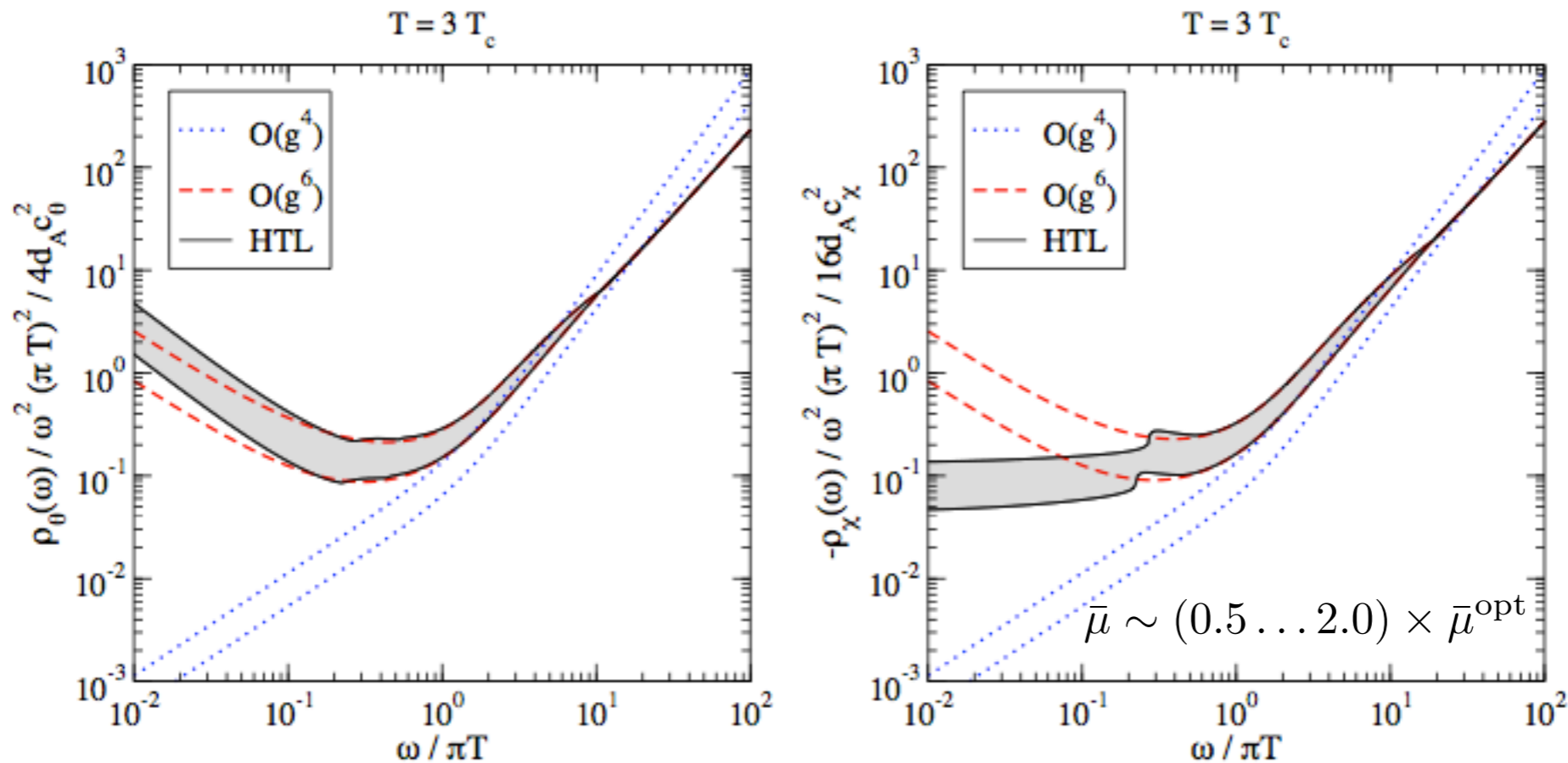
$$\frac{\rho_\theta(\omega)}{4d_A c_\theta^2} = \frac{\pi\omega^4}{(4\pi)^2} (1 + 2n_{\frac{\omega}{2}}) \left\{ g^4 + \frac{g^6 N_c}{(4\pi)^2} \left[\frac{22}{3} \ln \frac{\bar{\mu}^2}{\omega^2} + \frac{73}{3} + 8\phi_T(\omega) \right] \right\} + \mathcal{O}(g^8)$$

$$\frac{-\rho_\chi(\omega)}{16d_A c_\chi^2} = \frac{\pi\omega^4}{(4\pi)^2} (1 + 2n_{\frac{\omega}{2}}) \left\{ g^4 + \frac{g^6 N_c}{(4\pi)^2} \left[\frac{22}{3} \ln \frac{\bar{\mu}^2}{\omega^2} + \frac{97}{3} + 8\phi_T(\omega) \right] \right\} + \mathcal{O}(g^8)$$

In this results, the terms with the structure $\omega^n \delta(\omega)$ have been omitted!

- In the regime $\omega \gg \pi T$, “Fast apparent convergence”: $\ln(\bar{\mu}_\theta^{\text{opt}(\omega)}) \equiv \ln(\omega) - \frac{73}{44}$, $\ln(\bar{\mu}_\chi^{\text{opt}(\omega)}) \equiv \ln(\omega) - \frac{97}{44}$.
- When $\omega \ll \pi T$, EQCD: $\ln(\bar{\mu}_{\theta,\chi}^{\text{opt}(T)}) \equiv \ln(4\pi T) - \gamma_E - \frac{1}{22}$.
- The scale parameter $\Lambda_{\overline{\text{MS}}} \equiv \lim_{\bar{\mu} \rightarrow \infty} \bar{\mu} [b_0 g^2]^{-b_1/2b_0^2} \exp\left[-\frac{1}{2b_0 g^2}\right]$, $T_c = 1.25\Lambda_{\overline{\text{MS}}}$,

Spectral Functions



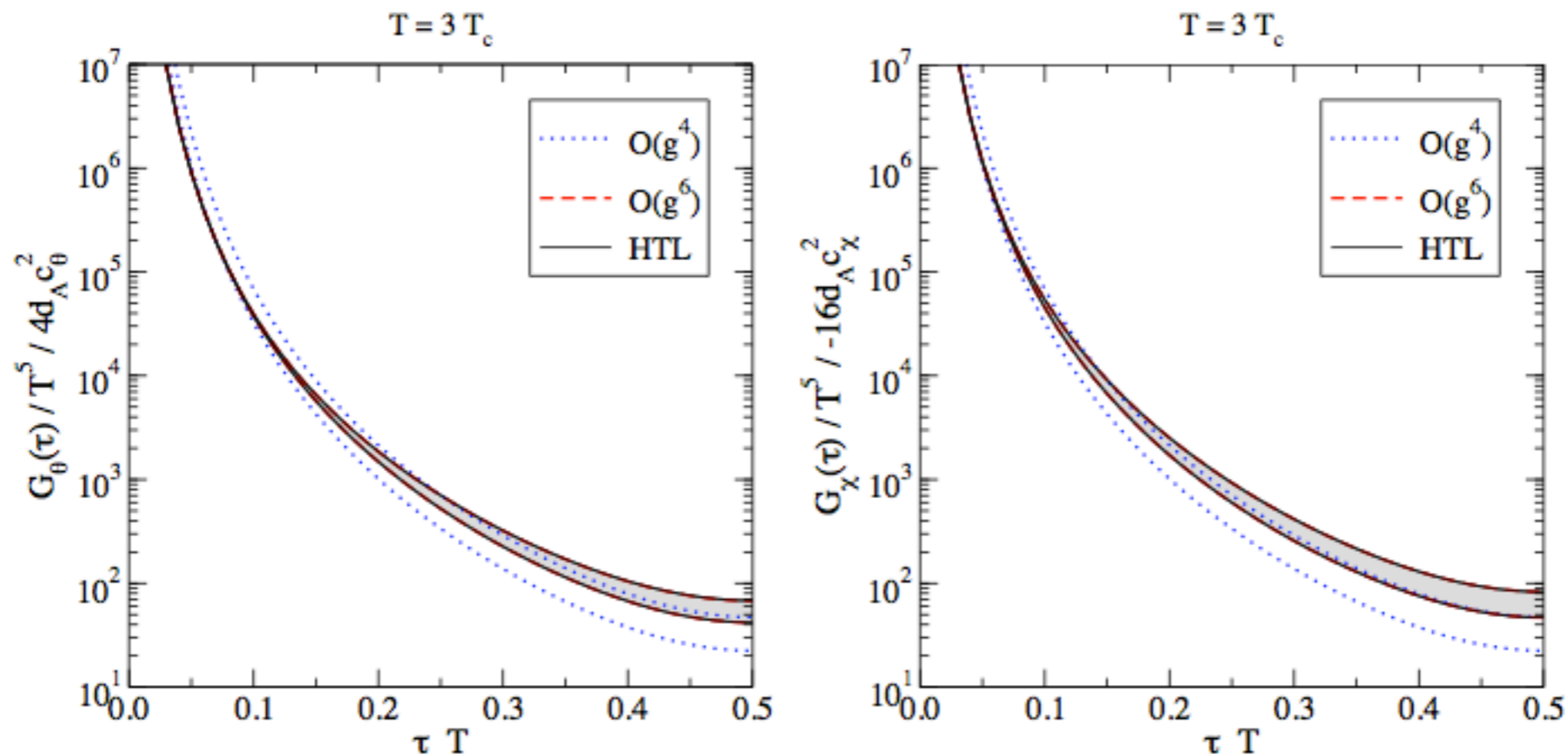
$$\frac{\rho_\theta(\omega)}{4d_A c_\theta^2} = \frac{\pi\omega^4}{(4\pi)^2} (1 + 2n_{\frac{\omega}{2}}) \left\{ g^4 + \frac{g^6 N_c}{(4\pi)^2} \left[\frac{22}{3} \ln \frac{\bar{\mu}^2}{\omega^2} + \frac{73}{3} + 8\phi_T(\omega) \right] \right\} + \mathcal{O}(g^8)$$

$$-\frac{\rho_\chi(\omega)}{16d_A c_\chi^2} = \frac{\pi\omega^4}{(4\pi)^2} (1 + 2n_{\frac{\omega}{2}}) \left\{ g^4 + \frac{g^6 N_c}{(4\pi)^2} \left[\frac{22}{3} \ln \frac{\bar{\mu}^2}{\omega^2} + \frac{97}{3} + 8\phi_T(\omega) \right] \right\} + \mathcal{O}(g^8)$$

$$\rho_{\text{resummed}}^{\text{QCD}} = \rho_{\text{resummed}}^{\text{QCD}} - \rho_{\text{resummed}}^{\text{HTL}} + \rho_{\text{resummed}}^{\text{HTL}} \approx \rho_{\text{naive}}^{\text{QCD}} - \rho_{\text{naive}}^{\text{HTL}} + \rho_{\text{resummed}}^{\text{HTL}}.$$

Imaginary-time Correlators

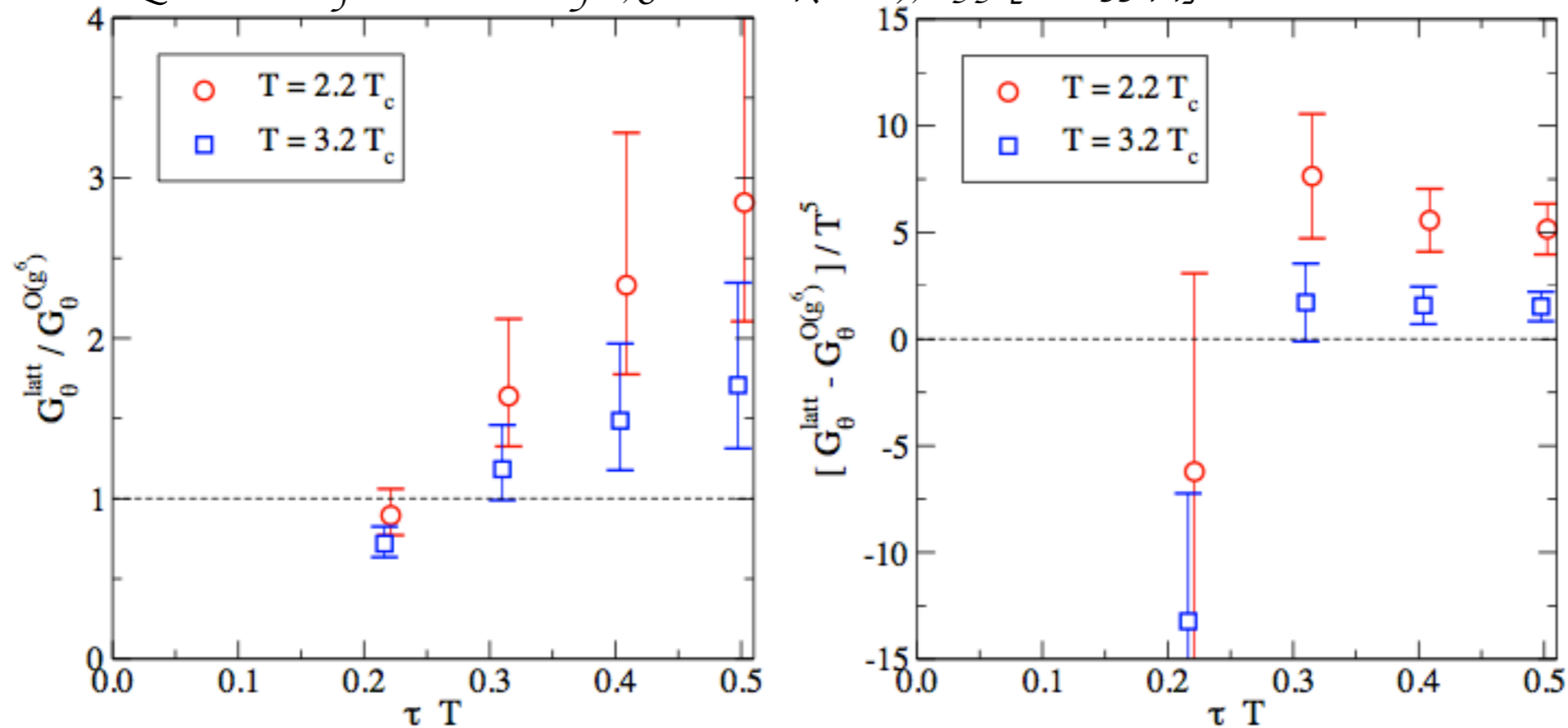
$$G(\tau) = \int_0^\infty \frac{d\omega}{\pi} \rho(\omega, \mathbf{0}) \frac{\cosh\left(\frac{\beta}{2} - \tau\right) \omega}{\sinh \frac{\beta\omega}{2}} .$$



Considerable difference between LO and NLO in spectral function leads to a small correction to the imaginary-time correlators.

Comparing with LQCD

Lattice data from H.B. Meyer, JHEP 04(2010), 099 [10023344]



- The ratio shows good agreement at short distance.
- The difference no longer shows the short distance divergence.
A model independent analytic continuation could be attempted.

Summary and Outlook

- A novel method has been evaluated to calculate the correlators and spectral functions to NLO in pure Yang-Mills theory.
- Spectral functions in bulk channel has been computed with this method, which is very helpful to constrain the corresponding correlator and determine the thermal coefficient.
- ★ The application to shear channel is underway.
- ★ Production rate of axion and dilation in cosmology, as well as the spectral function of electromagnetic current in QCD, ..., can be evaluated with this method.



Thanks for your attention!