

The QCD phase diagram at vanishing baryon density

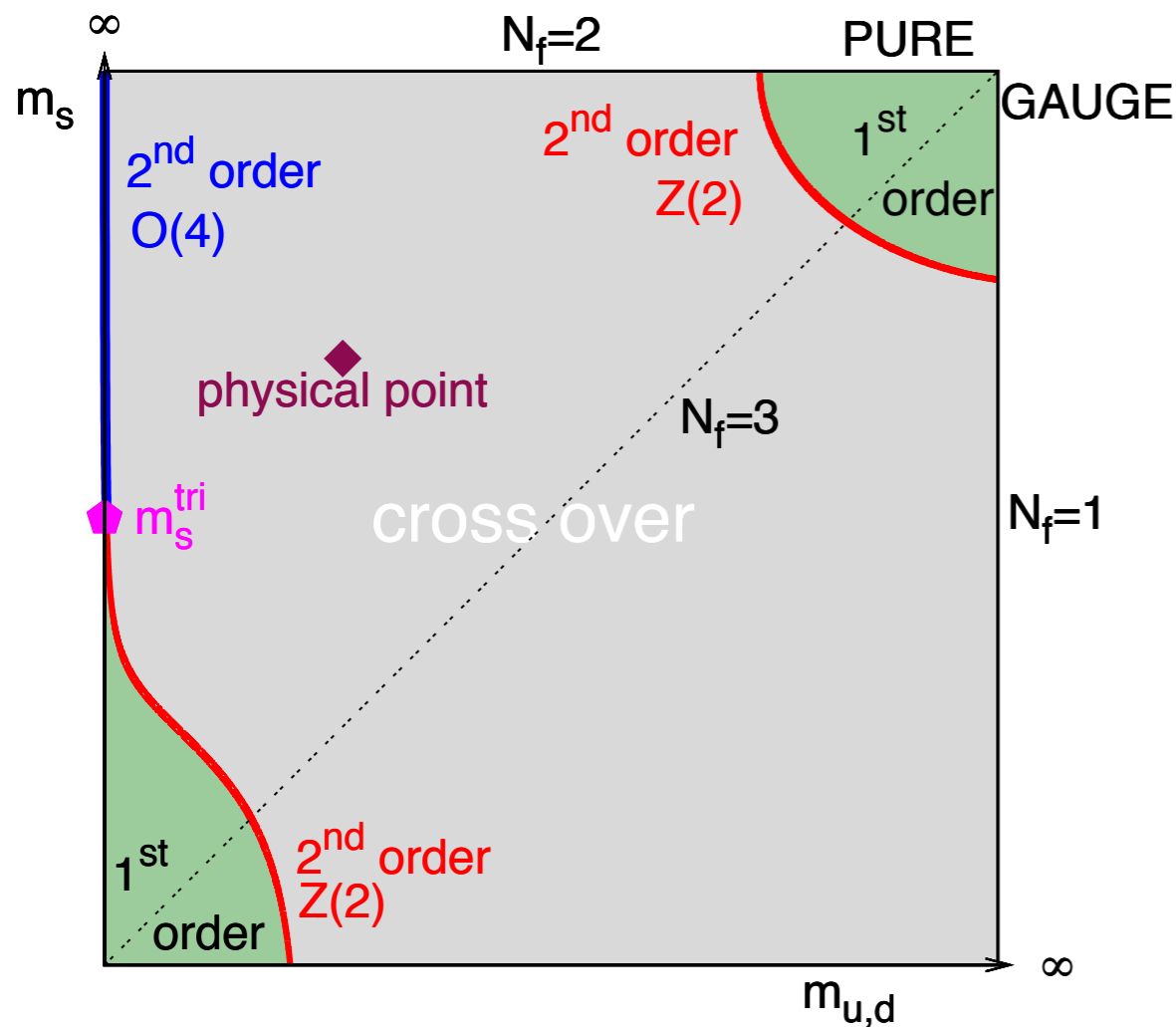
Heng-Tong Ding

Brookhaven National Laboratory

done in collaboration with Alexei Bazavov, Prasad Hegde, Frithjof Karsch,
Swagato Mukherjee, Peter Petreczky

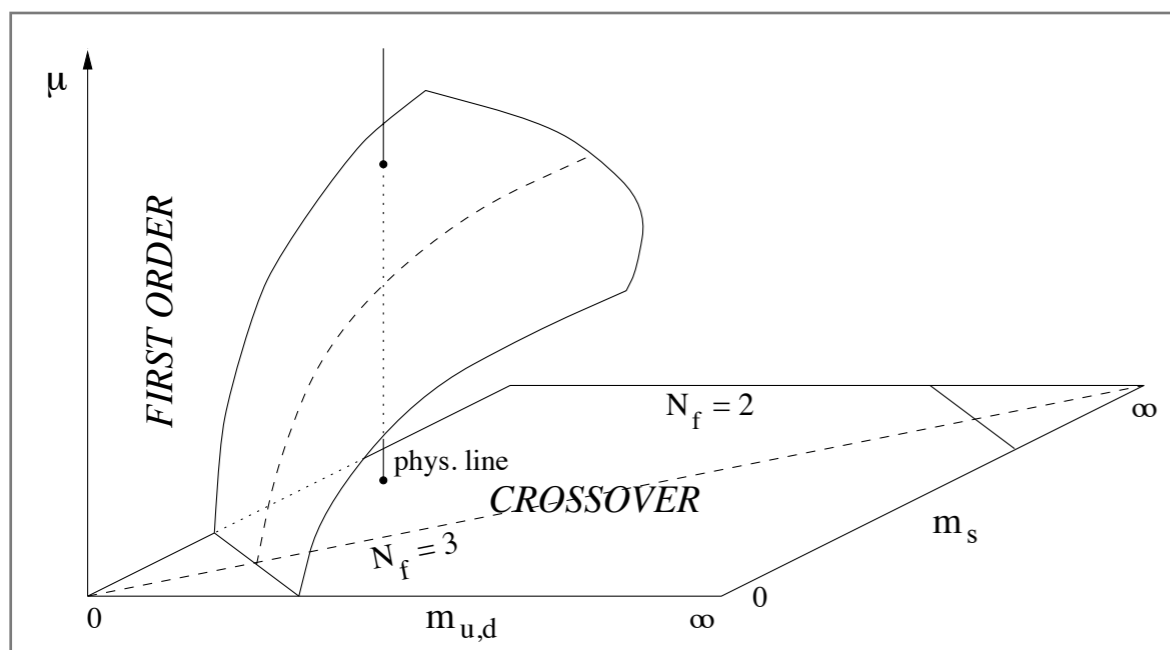
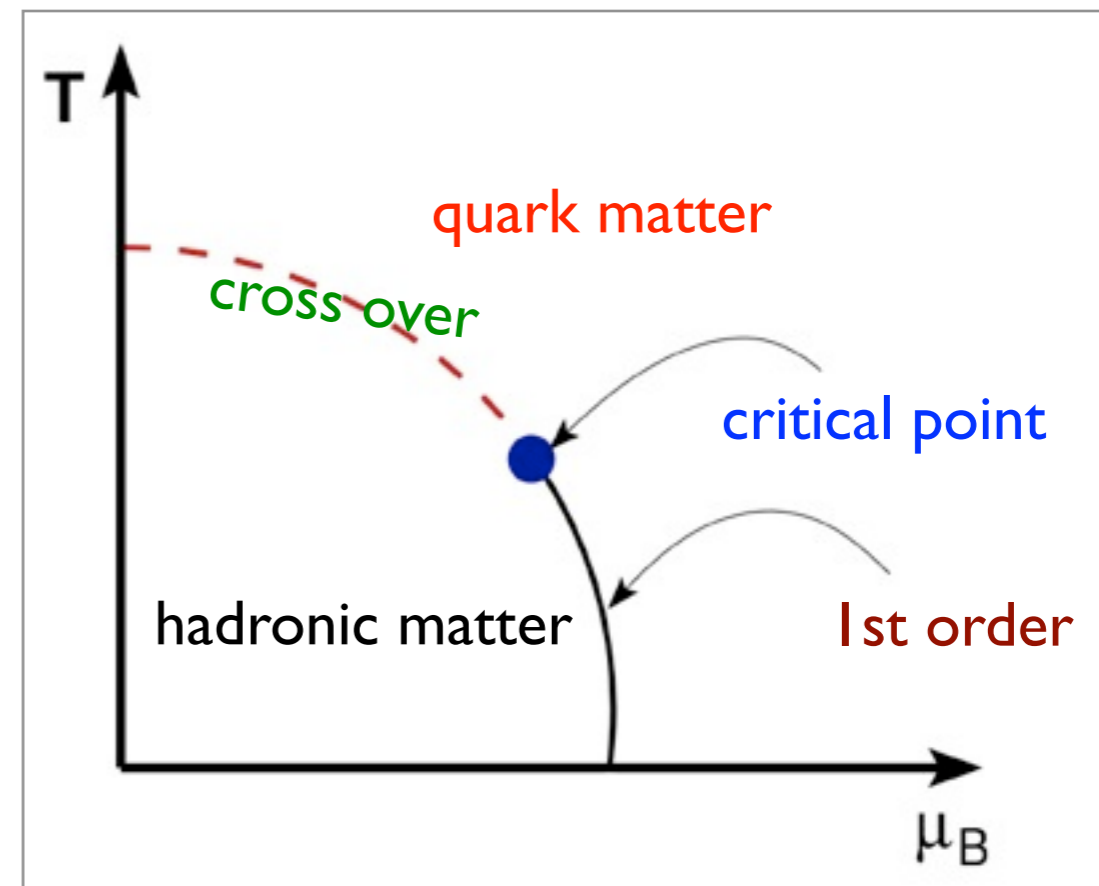
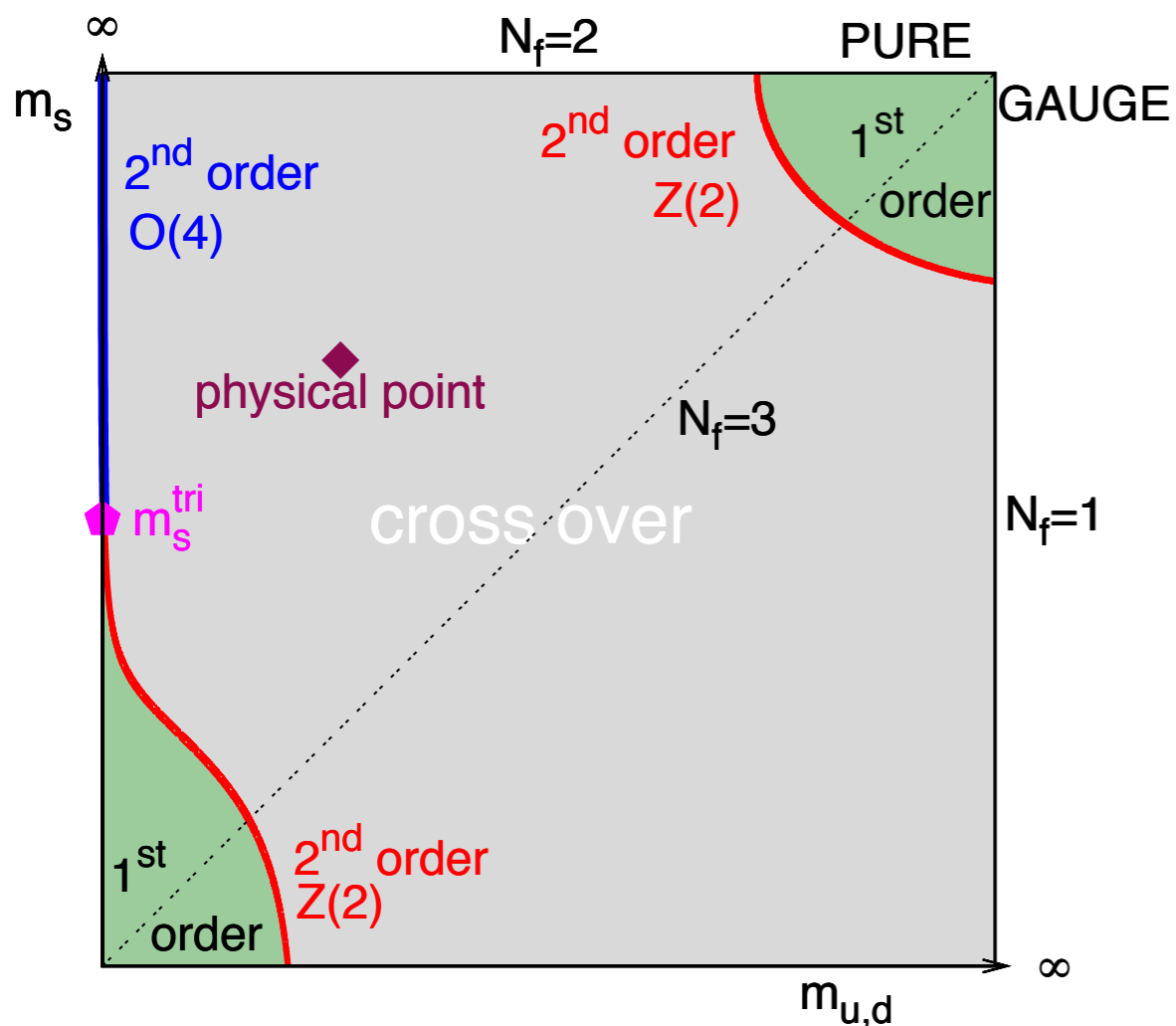
arXiv:1111.0185

The QCD phase diagram

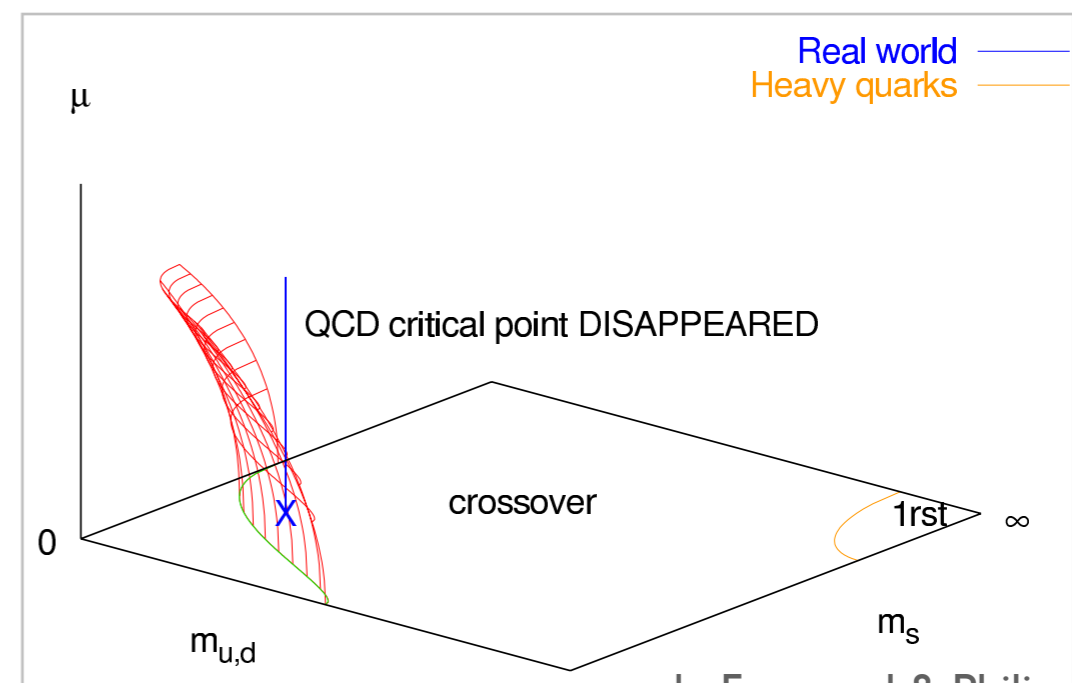


- $N_f=2+1$ theory: at $m=0$ or ∞ has a first order phase transition Pisarski, Wilczek PRD '84, Alexandrou et al., PRD'99...
- Intermediate quark mass region an analytic cross over is expected
- At physical quark masses, a cross over is confirmed Bernard et al., PRD '05, Cheng et al., PRD '06, Aoki et al., Nature '06...
- Critical lines of second order transition
 - $N_f=2$: O(4) universality class Ejiri et al., PRD '09, ...
 - $N_f=3$: Ising universality class Karsch, Laermann, Schmidt PLB '04, ...

The QCD phase diagram

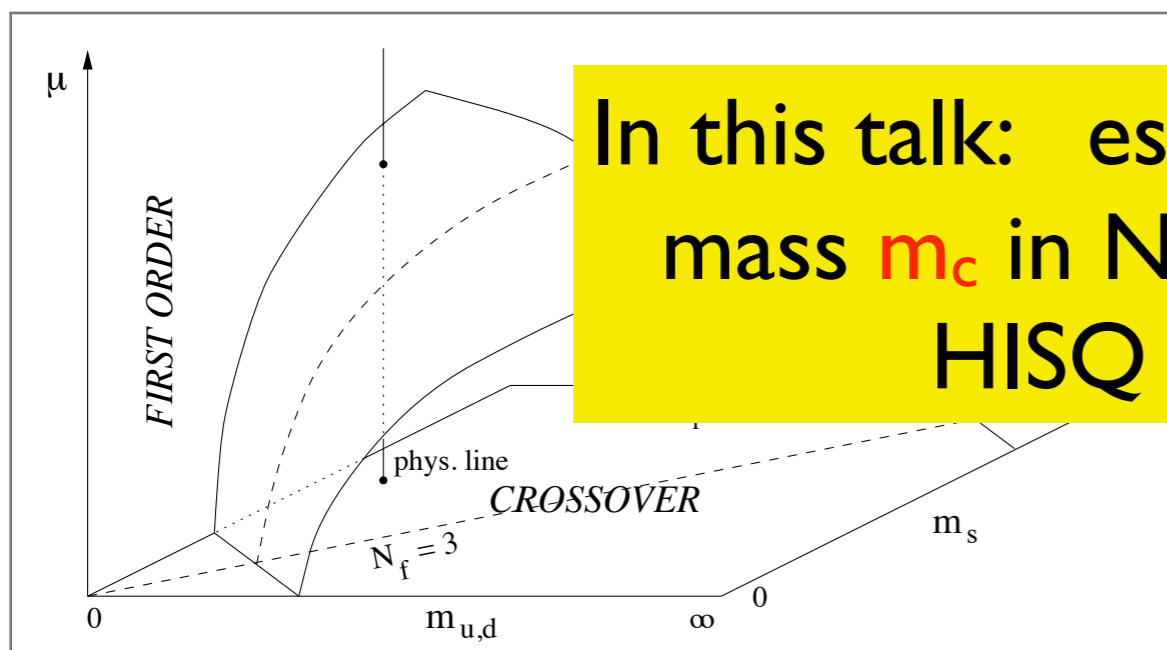
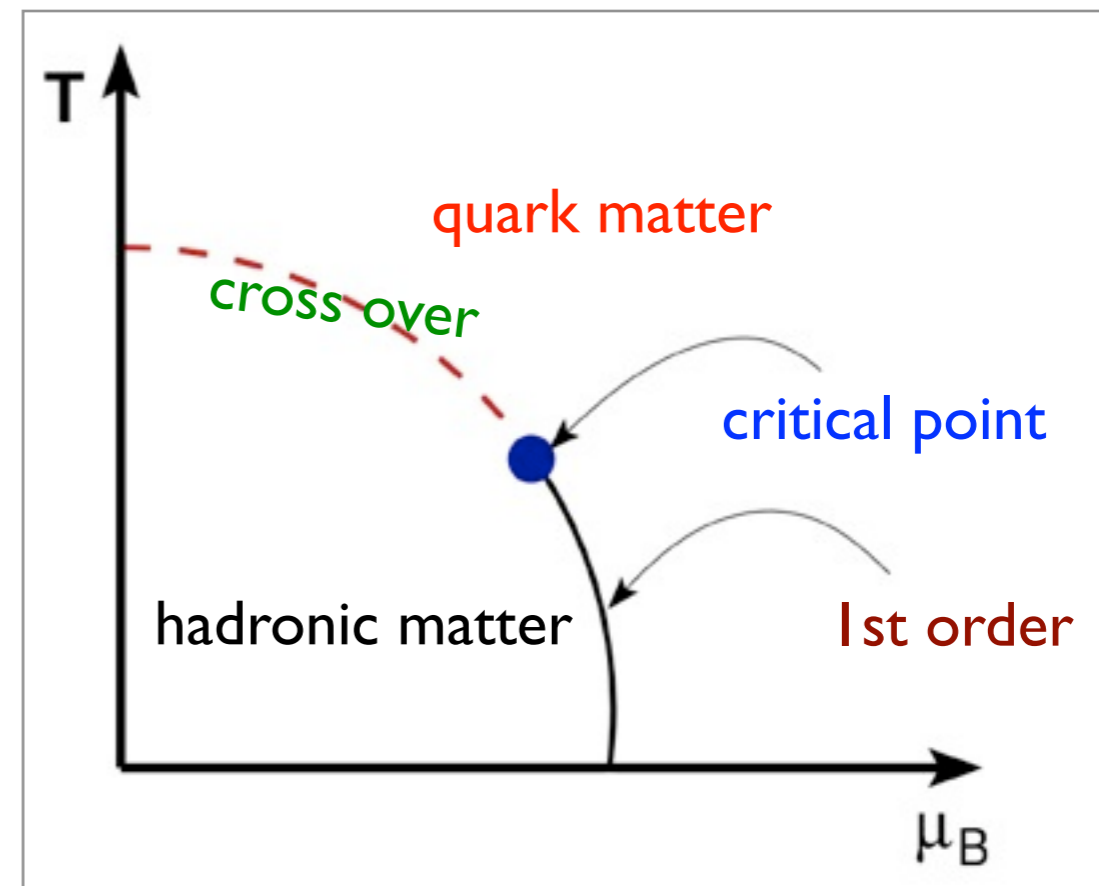
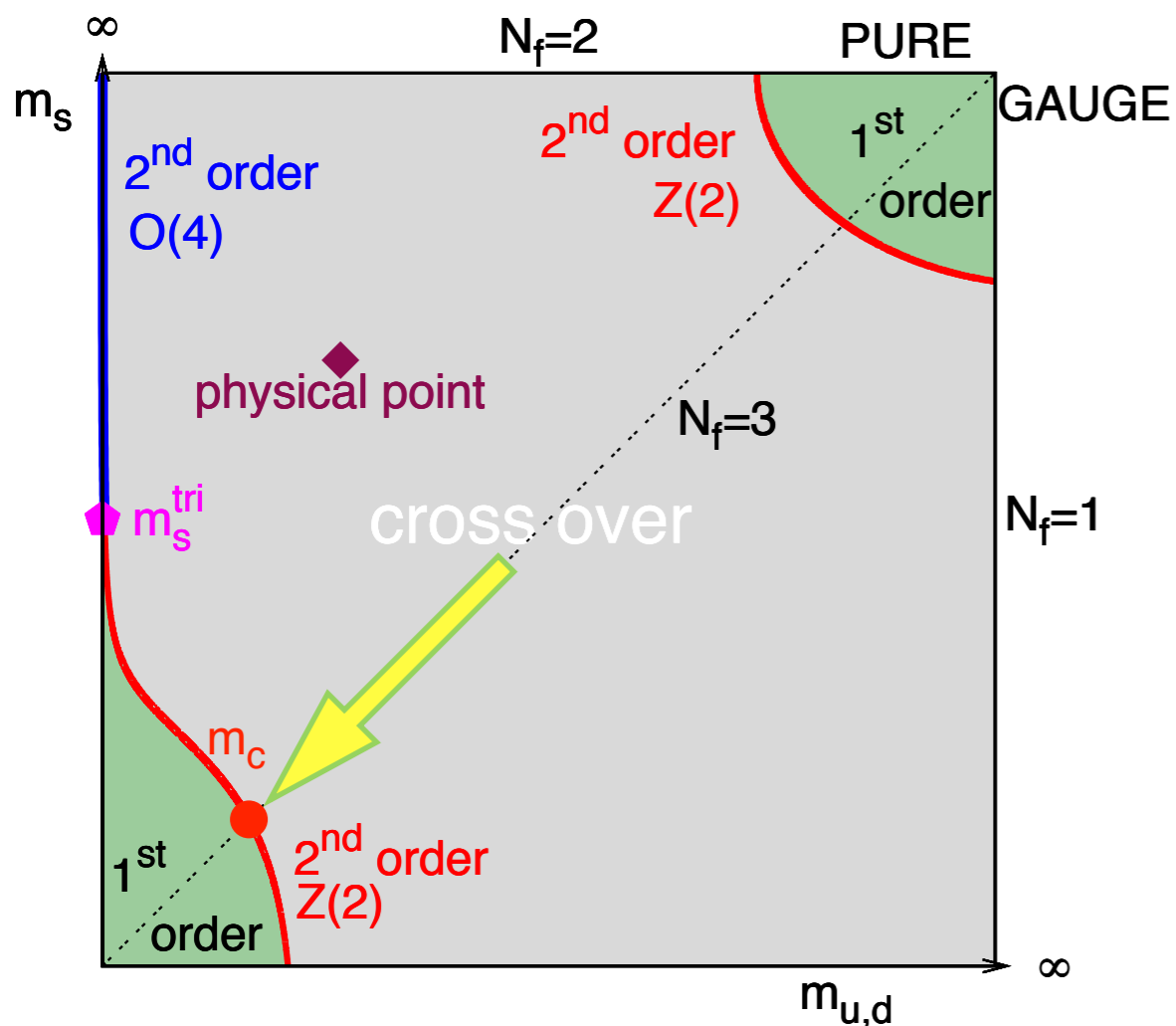


Karsch et al., '03



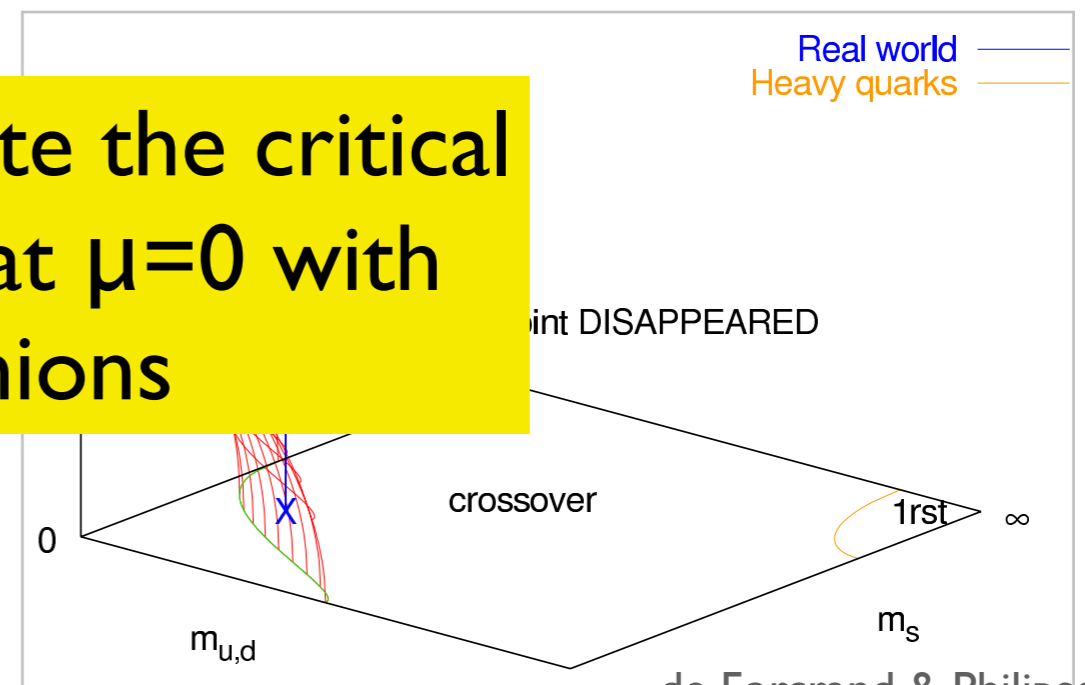
de Forcrand & Philipsen, '07

The QCD phase diagram



Karsch et al., '03

In this talk: estimate the critical mass m_c in $N_f=3$ at $\mu=0$ with HISQ fermions



de Forcrand & Philipsen, '07

Universality class near critical lines

Behavior of the free energy close to critical lines

$$f(m, T) = h^{1+1/\delta} f_s(z) + f_{\text{reg}}(m, T)$$

$$h \sim m - m_c$$

$z = t/h^{1/\beta\delta}$, h : external field, t : reduced scaling variable

$$t \sim T - T_c$$

$f_s(z)$: universal scaling function, $Z(2)$ etc.

β, δ : universal critical exponents

Magnetic Equation of State (MEoS):

$$M = -\partial f_s(t, h) / \partial h = h^{1/\delta} f_G(z)$$

M : order parameter

- $N_f=2, m_c=0$

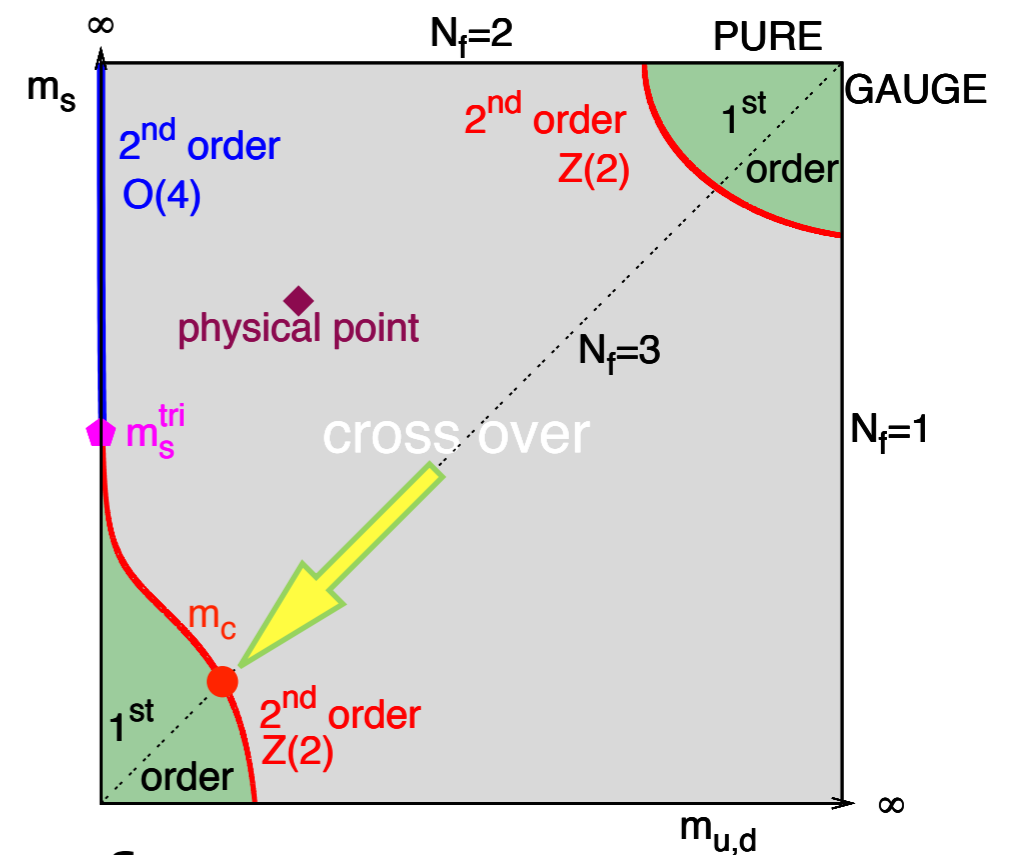
$$M = \langle \bar{\Psi} \Psi \rangle / T^3 \Big|_{\text{fixed } z} \sim m^{1/\delta}$$

$$\chi_q / T^2 \Big|_{\text{fixed } z} \sim m^{1/\delta - 1}$$

- $N_f=3, m_c \neq 0$

$$M = (\langle \bar{\Psi} \Psi \rangle + r \square) \Big|_{T=T_c, m_c} \sim (m - m_c)^{1/\delta}$$

$$\chi_q / T^2 \Big|_{T=T_c, m_c} \sim (m - m_c)^{1/\delta - 1}$$

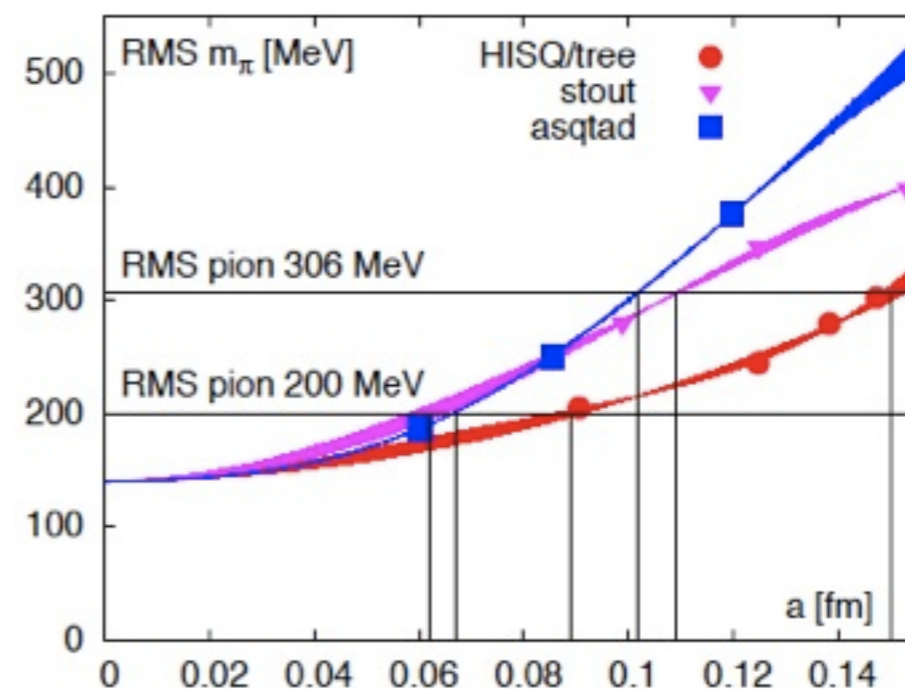
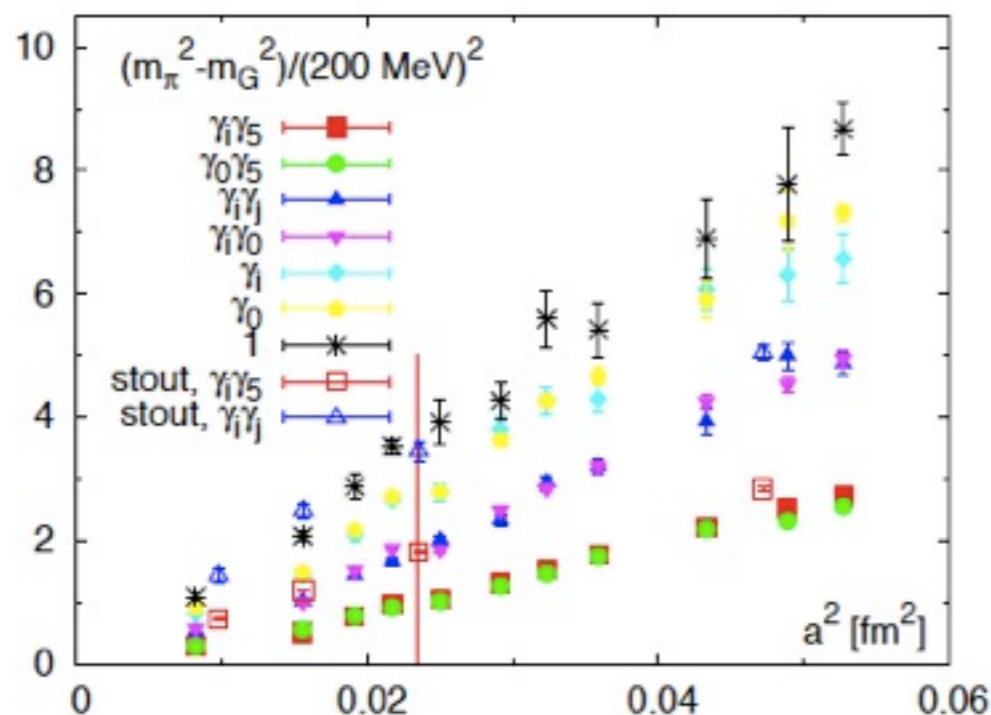


the value of m_π^c investigated so far...

Dependence on the choice of action

- Naive action: $N_\tau = 4 \Rightarrow m_\pi^c \approx 290 \text{ MeV}$ F. Karsch et al., Nucl.Phys.Proc.Suppl. 129 (2004) 614
- Naive action: $N_\tau = 6 \Rightarrow m_\pi^c \approx 140 \text{ MeV}$ P. de Forcrand et al, PoS LATTICE2007 (2007) 178
- p4fat3 action: $N_\tau = 4 \Rightarrow m_\pi^c \approx 67 \text{ MeV}$ F. Karsch et al., Nucl.Phys.Proc.Suppl. 129 (2004) 614
- stout action: $N_\tau = 6 \Rightarrow m_\pi^c \approx 50 \text{ MeV}$ G. Endrodi et al., PoS LAT2007 (2007) 228

Highly Improved staggered quarks (HISQ) used in this work



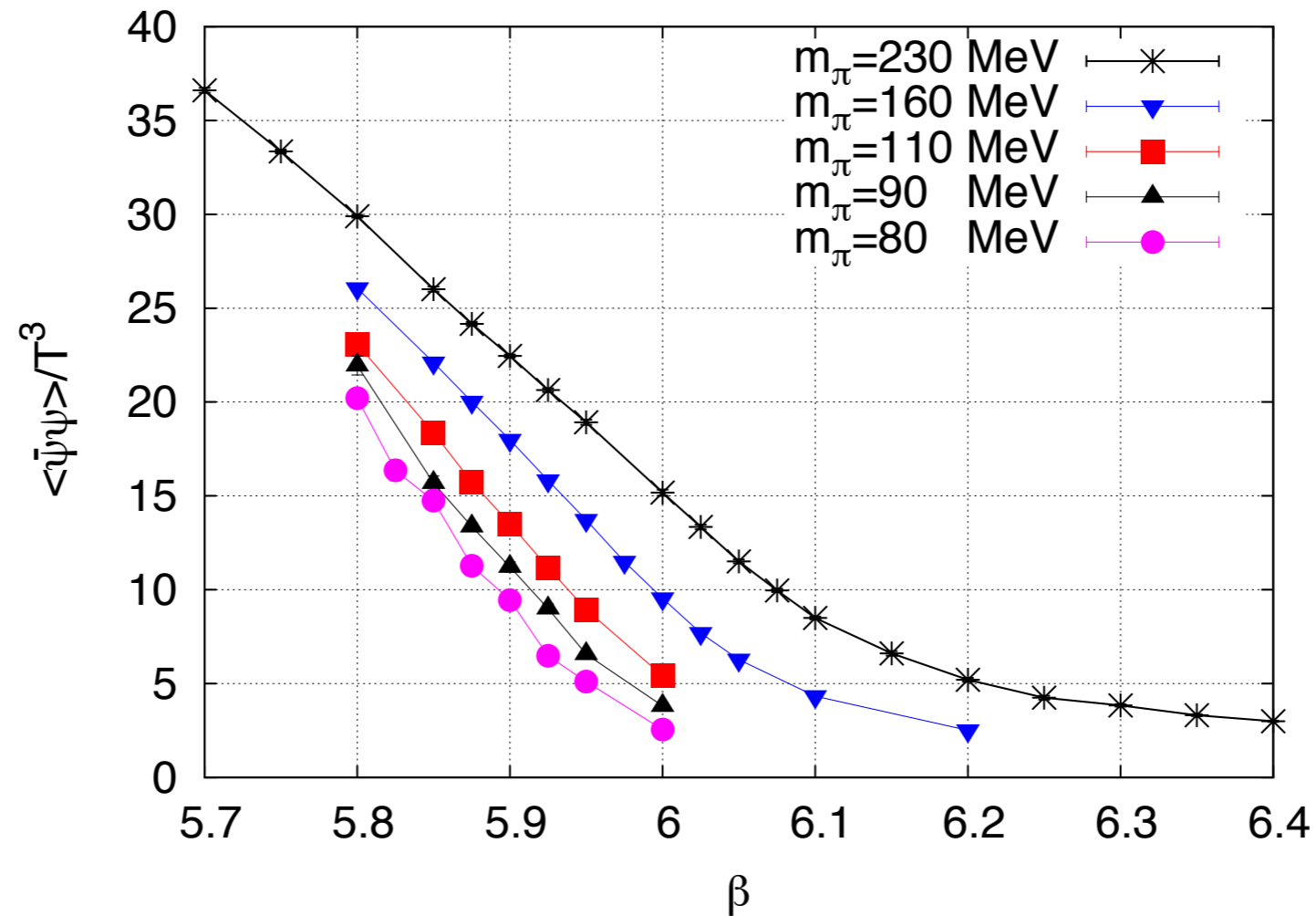
Bazavov and Petreczky, PoS LAT2010,169(2010)

Lattice data for $N_f=3$ and $N_T=6$

- ★ Highly Improved staggered fermions/tree action used
- ★ 3 degenerate quarks
- ★ $N_T=6$ lattices

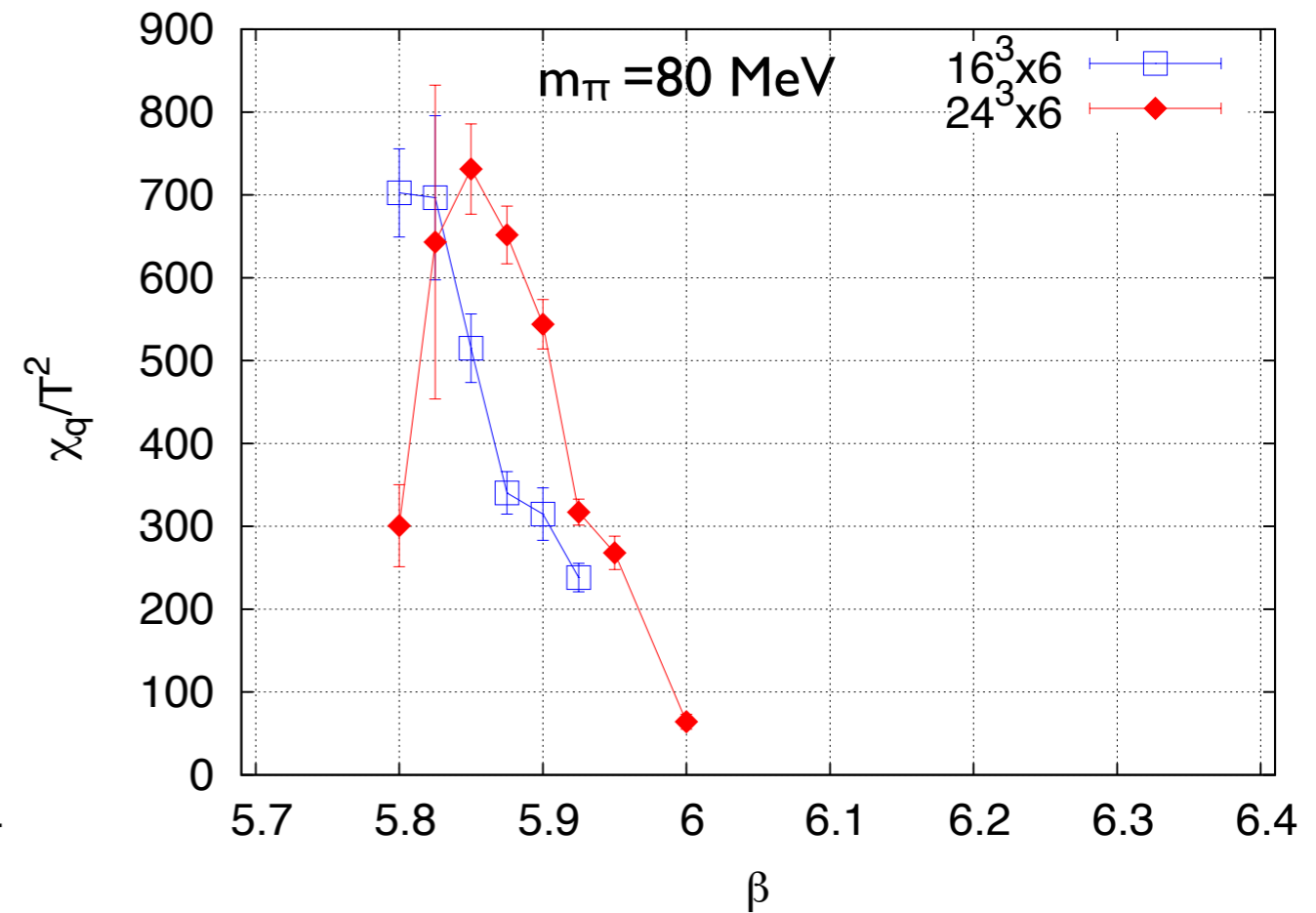
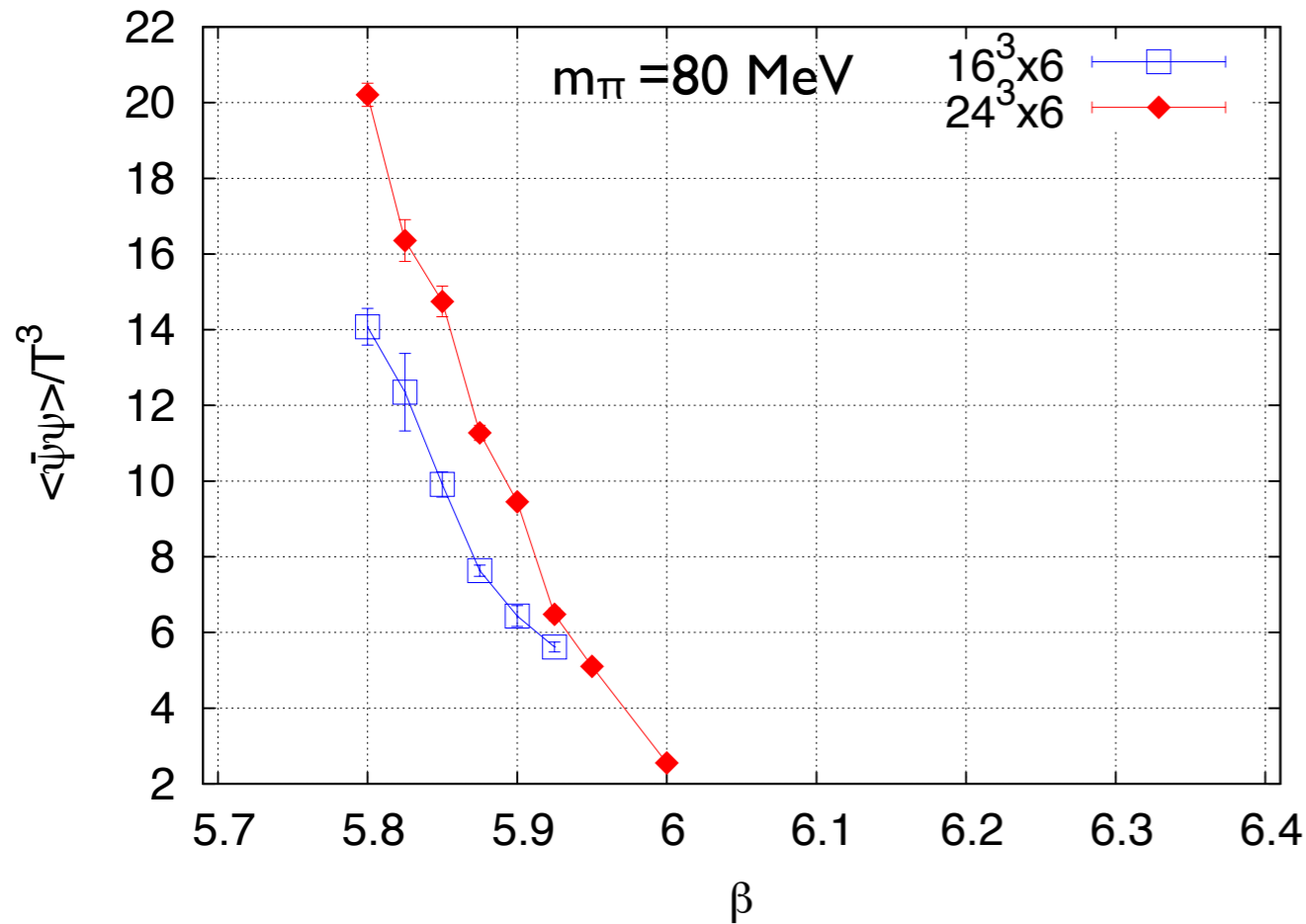
lattice dim	quark mass	m_π	# T	statistics
$16^3 \times 6$	$ma=0.0075$	230 MeV	17	~6000
$24^3 \times 6$	$ma=0.00375$	160 MeV	12	~8000
$24^3 \times 6$	$ma=0.001875$	110 MeV	7	~8000
$24^3 \times 6$	$ma=0.00125$	90 MeV	7	~5000
$24^3 \times 6$	$ma=0.0009375$	80 MeV	8	~6000
$16^3 \times 6$	$ma=0.0009375$	80 MeV	6	~7000

chiral condensate as an approx. order parameter



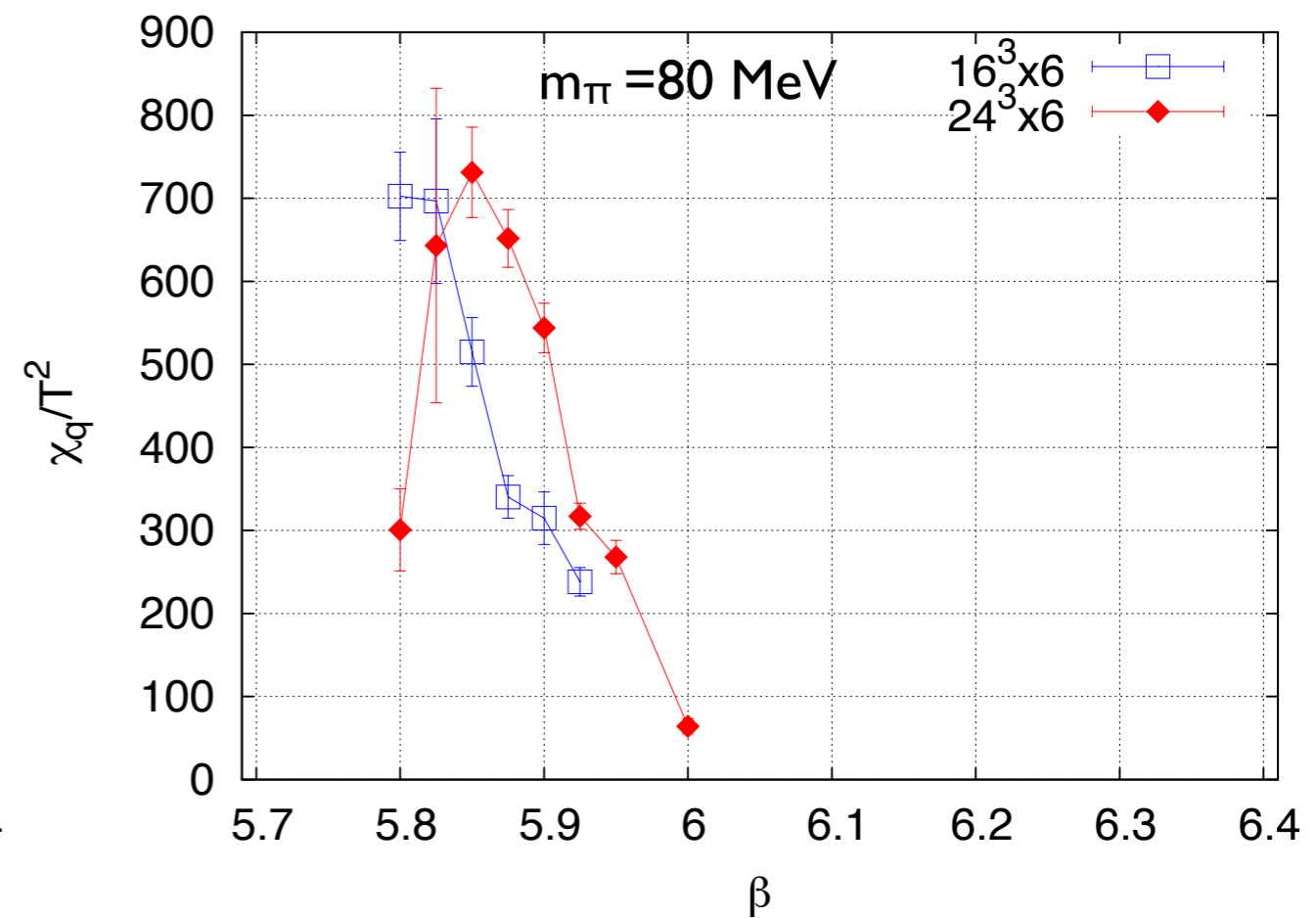
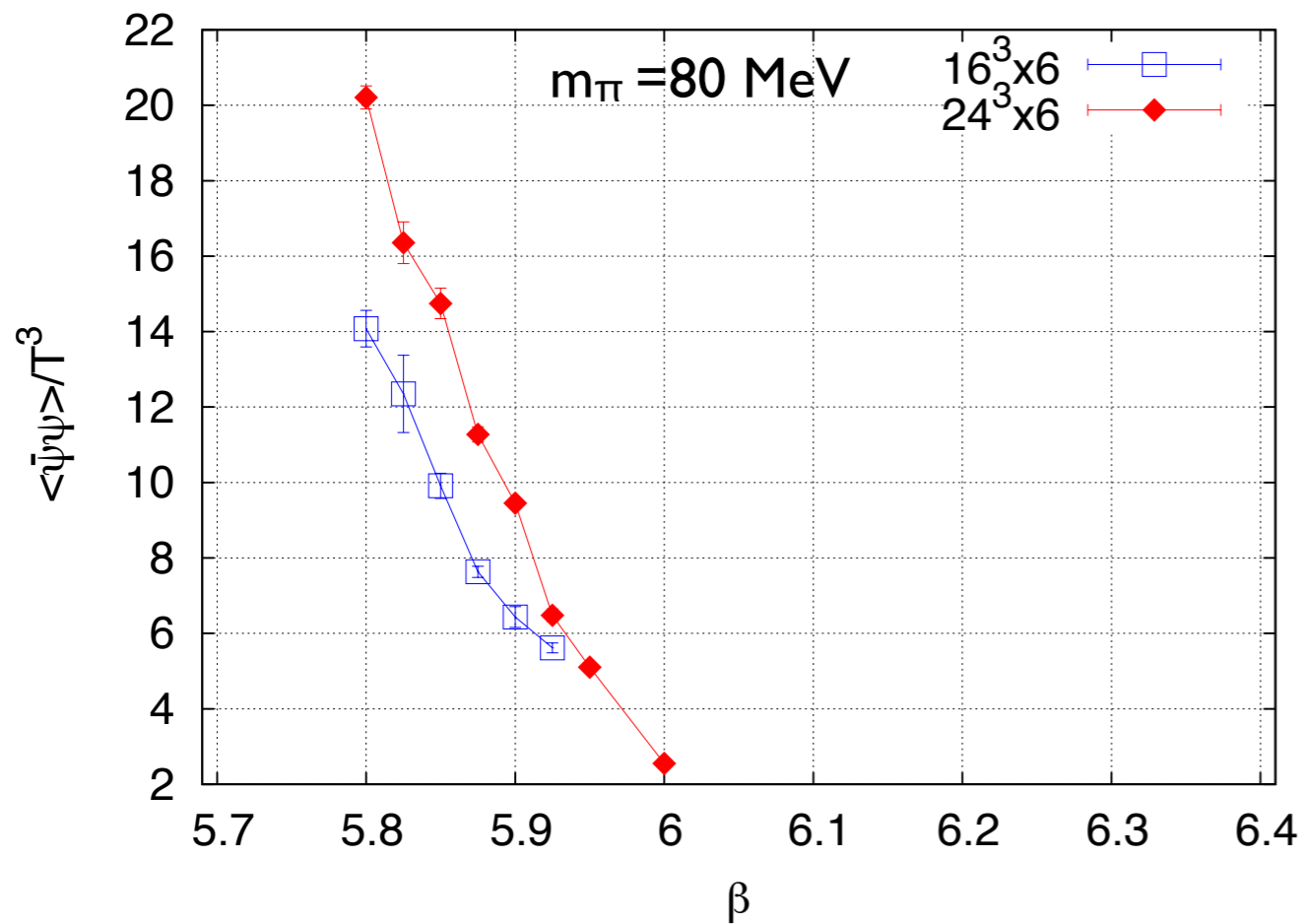
- No evidence for discontinuity is found

chiral condensate as an approx. order parameter



- finite volume effects bring chiral condensates down
- No evidence of volume scaling from chiral susceptibility

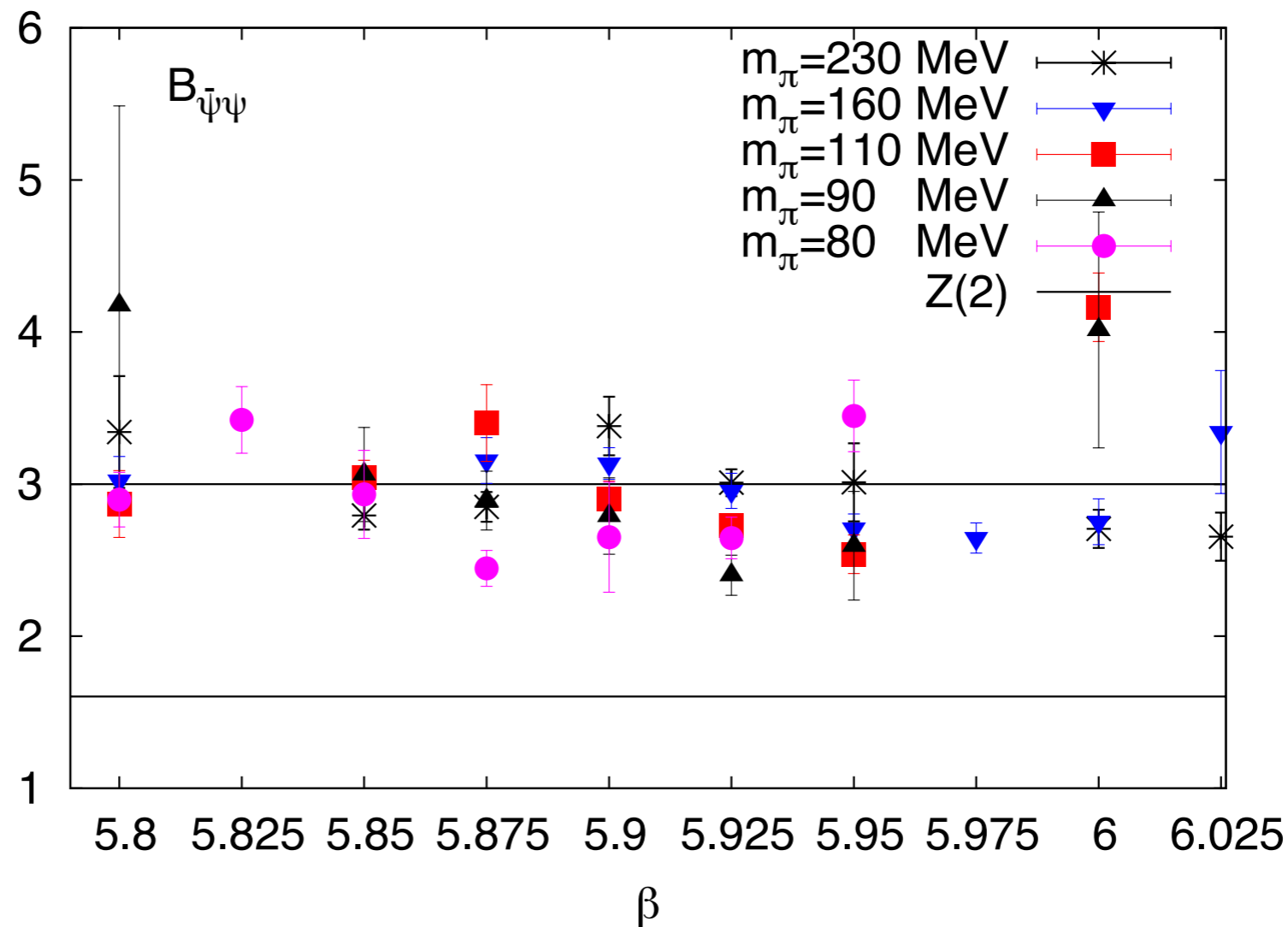
chiral condensate as an approx. order parameter



- finite volume effects bring chiral condensates down
- No evidence of volume scaling from chiral susceptibility
- With $230 \geq m_\pi \geq 80$ MeV, no direct signal of first order phase transition is found

Binder cumulant of chiral condensate

$$B_{\bar{\psi}\psi} \equiv \frac{\langle (\delta\bar{\psi}\psi)^4 \rangle}{\langle (\delta\bar{\psi}\psi)^2 \rangle^2}$$

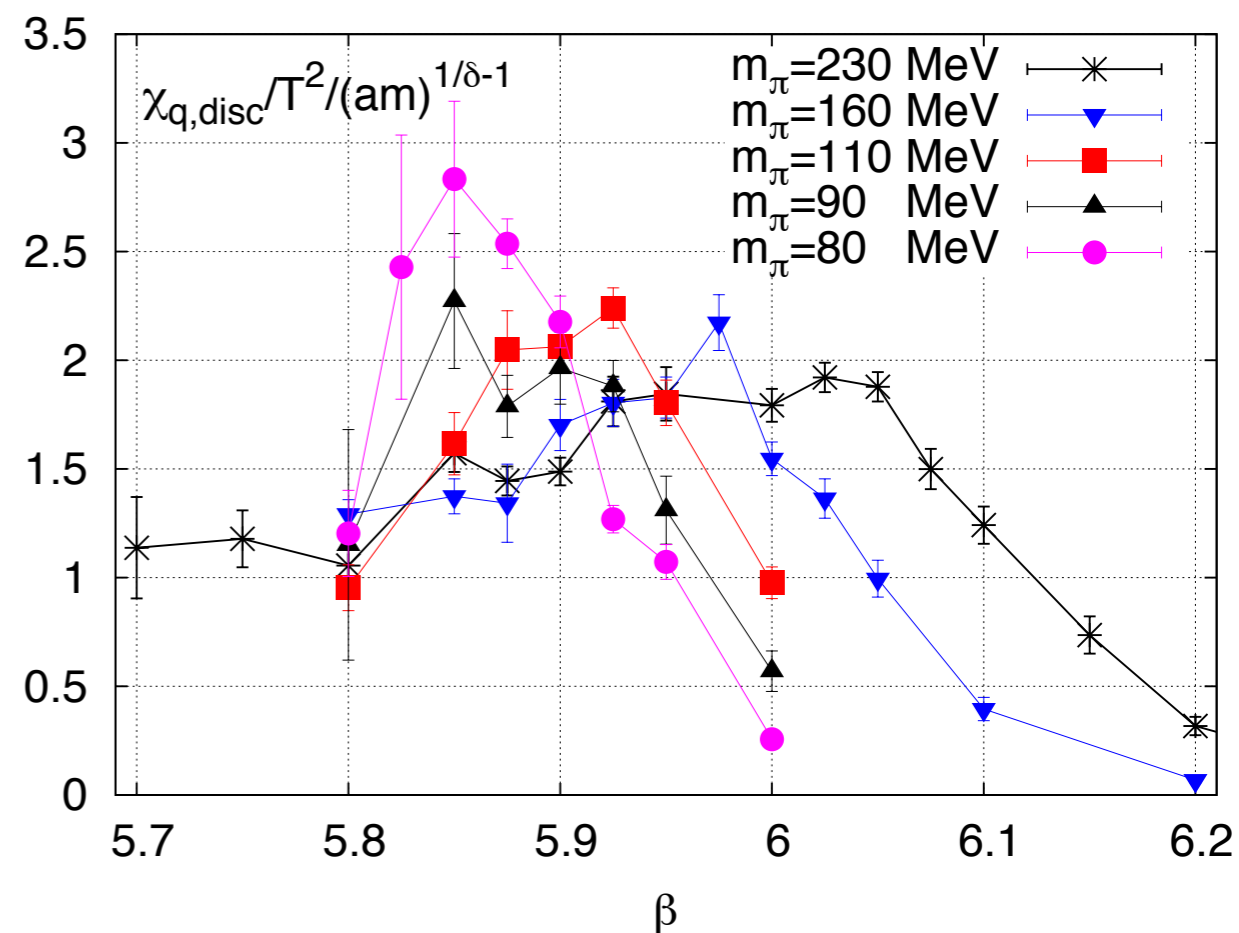
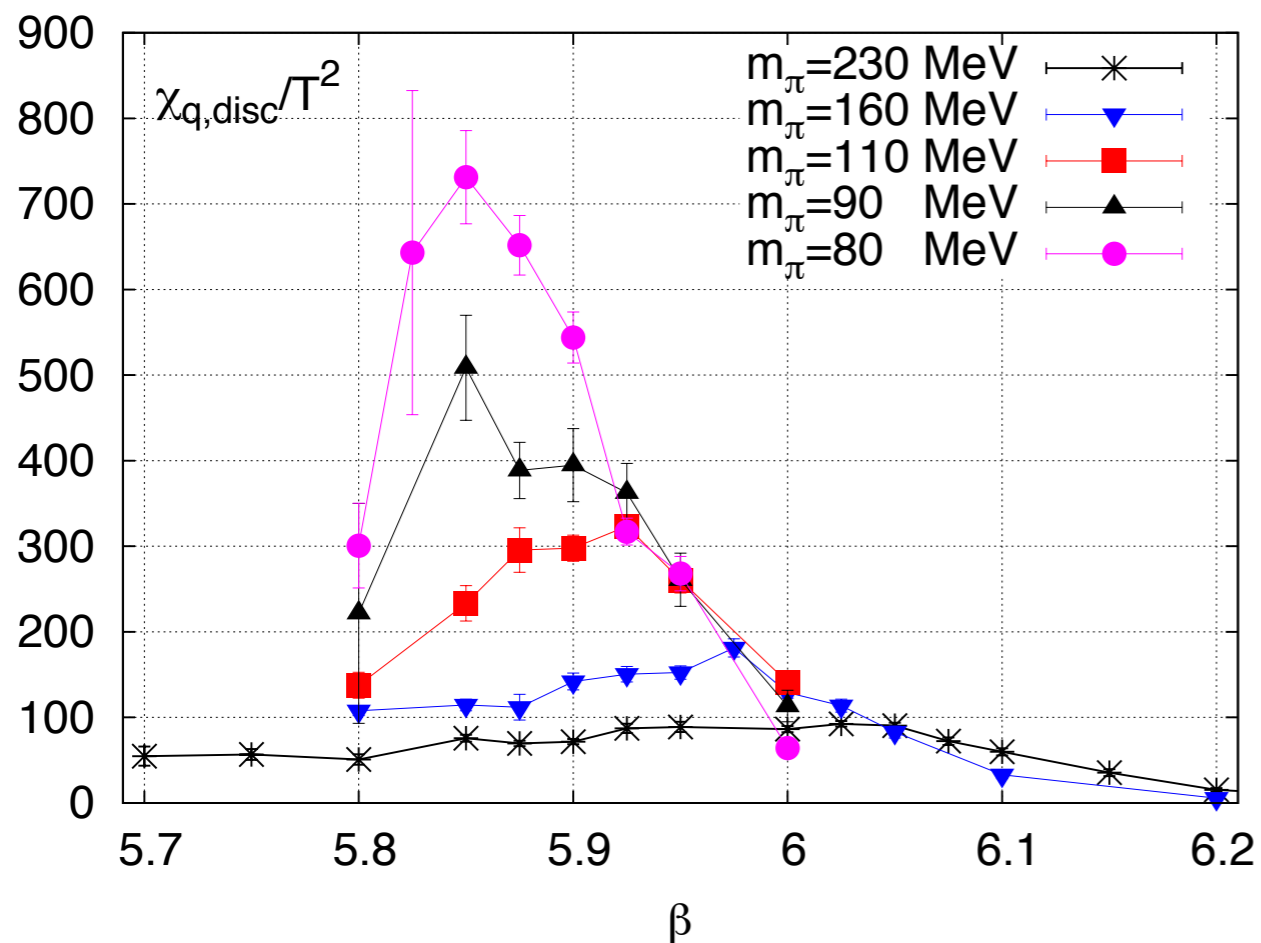


In the crossover region $B=3$

2nd order transition in the Ising universal class $B=1.604$

1st order transition $B=1$

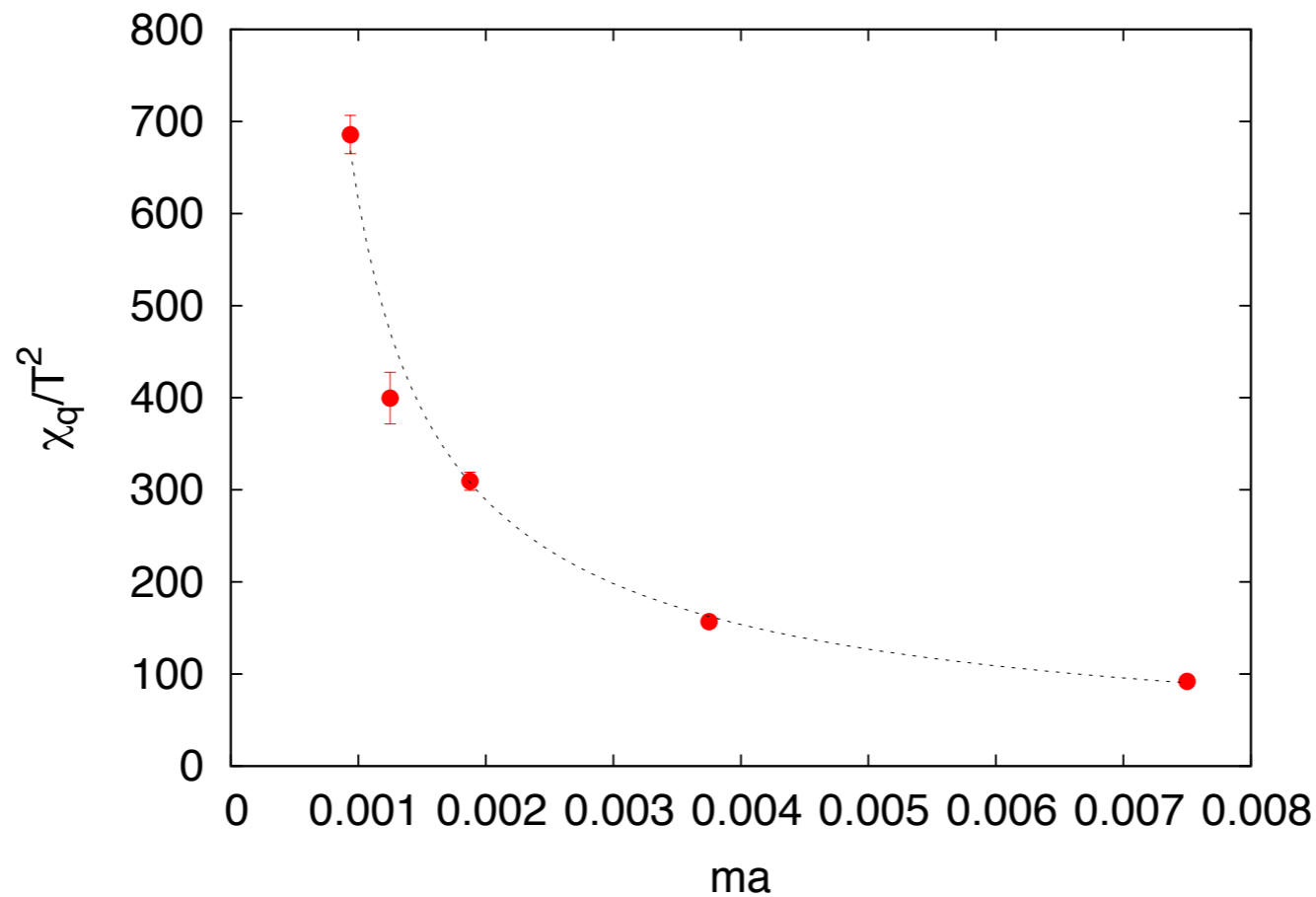
chiral susceptibilities



- peak locations in chiral susceptibilities shift to lower temperature with decreasing quark mass
- peak heights in chiral susceptibilities grow faster than $(am)^{1/\delta-1}$:
indication of a non-zero critical mass m_c

estimate of the critical mass

Fitting ansatz: $\chi_q/T^2 = c (m-m_c)^{1/\delta-1}$



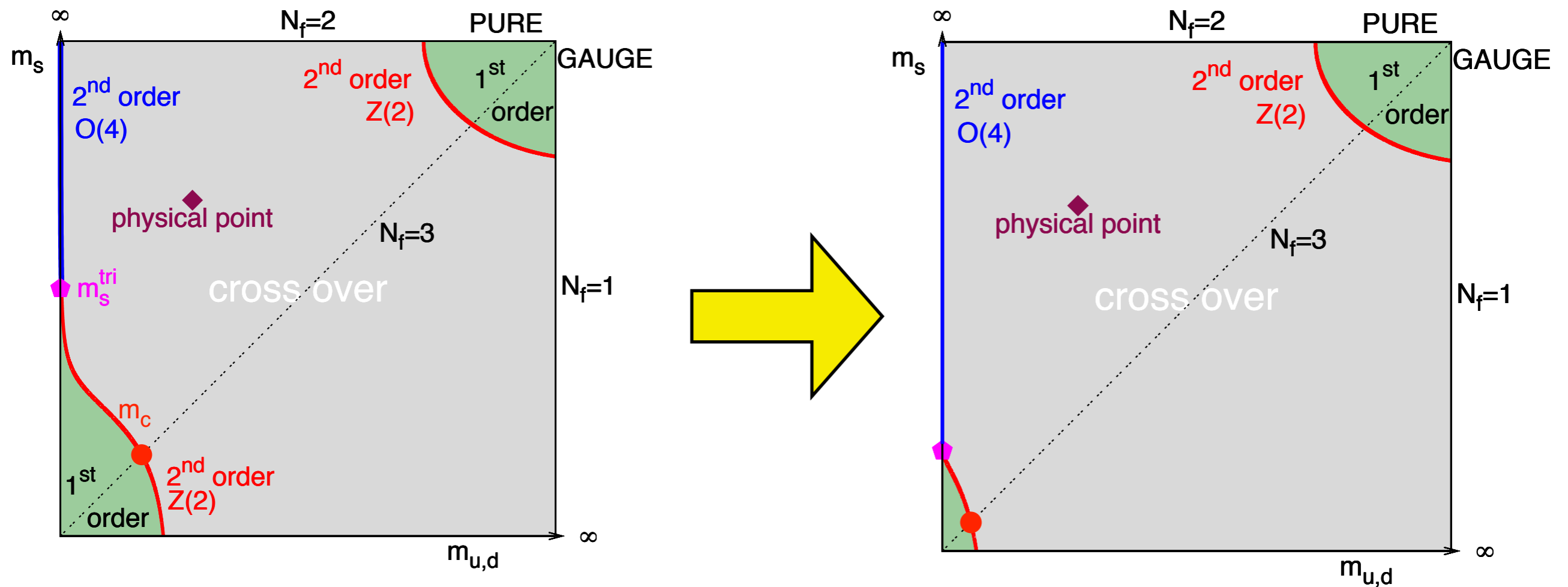
Navigation: $am=0.0009375 \Leftrightarrow m_\pi=80 \text{ MeV}$

$am_c \approx 0.00037$ $m_\pi^c \approx 45 \text{ MeV}$

finite V effects

$m_\pi^c \approx 45 \text{ MeV}$

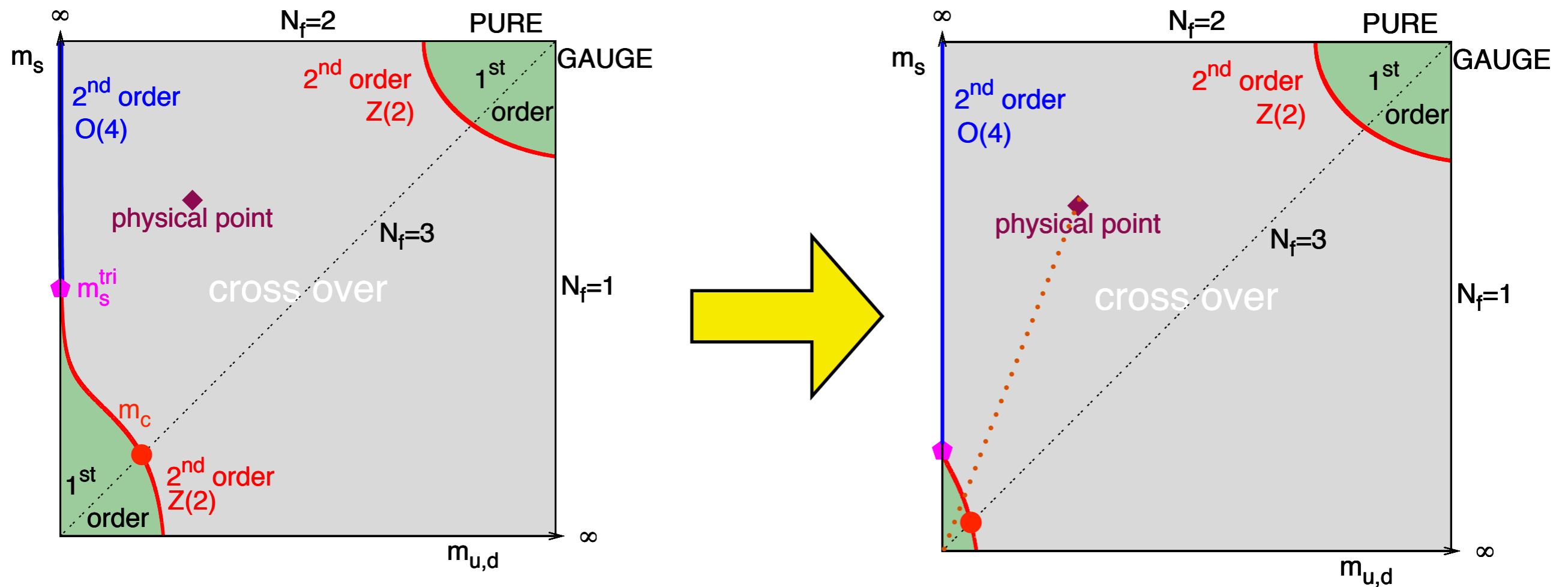
the first order phase transition region revisited



coordinates of the physical point: $(\bar{m}_s/27, \bar{m}_s)$

3 degenerate quarks: coordinates of $m_{max}^c \approx (\bar{m}_s/270, \bar{m}_s/270)$

the first order phase transition region revisited

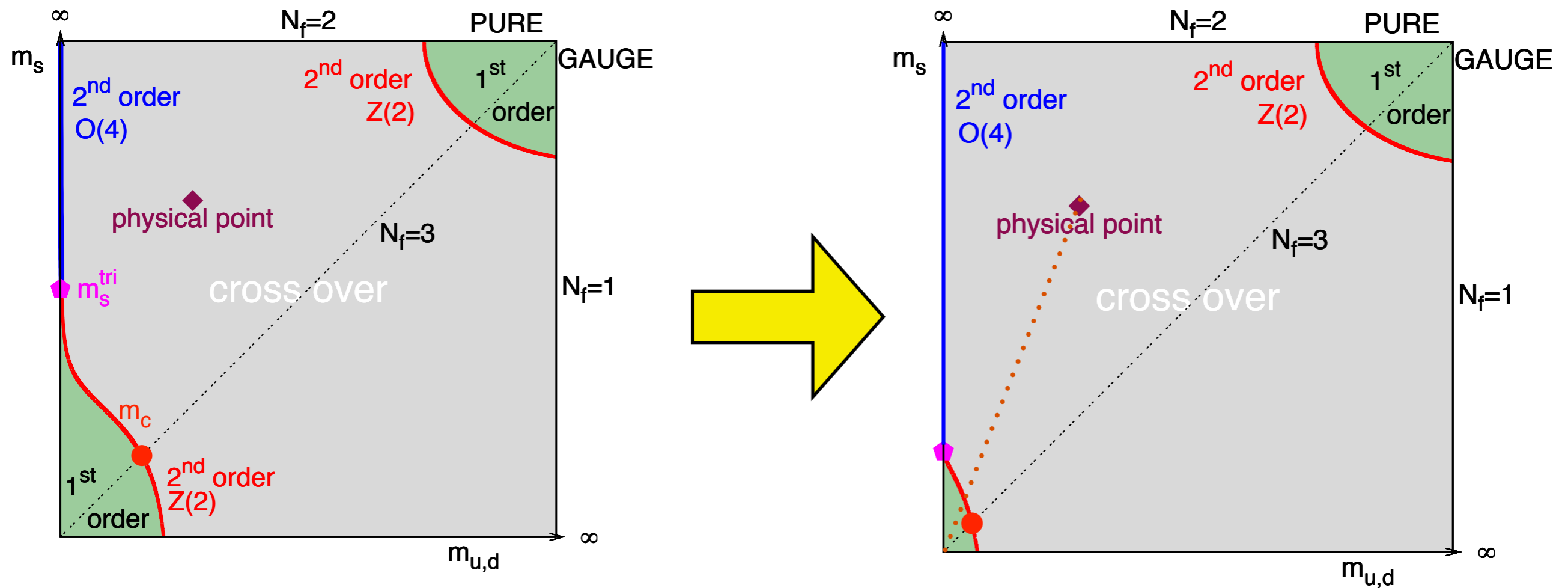


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non-degenerate quarks: coordinates of $m_{\text{max}}^c \approx (\bar{m}_s/225, \bar{m}_s/8)$ Endrodi et al., '07

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Consequences: to have a critical point at small μ , the critical surface has to bend towards to the physical point with a very large curvature

Summary & Outlook

- We study the direct signal for first order phase transition with m_π from 230 MeV down to 80 MeV for $N_f=3$ on $N_\tau=6$ lattices using the HISQ action
- No evidence for first order phase transition is found with $230 \text{ MeV} \geq m_\pi \geq 80 \text{ MeV}$
- From scaling analysis on chiral susceptibility, current estimation gives $m_\pi^c \approx 45 \text{ MeV}$, which indicates the first order chiral phase transition region is far away from the physical point