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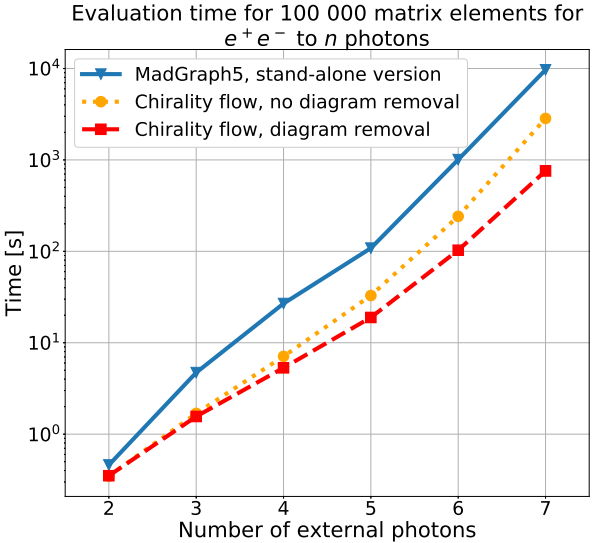
Speeding up SM Amplitude Calculations with Chirality Flow

QCD@LHC 2022 28 NOVEMBER 2022 - ANDREW LIFSON

BASED ON HEP-PH:2003.05877 (EPJC), HEP-PH:2011.10075 (EPJC), AND HEP-PH:2203.13618 (EPJC)

IN COLLABORATION WITH JOAKIM ALNEFJORD, CHRISTIAN REUSCHLE, MALIN SJÖDAHL, AND ZENNY WETTERSTEN

Our Main Numerical Result (so far) (hep-ph:2203.13618)



Spinor-Helicity: its Building Blocks

Lorentz algebra $so(3,1) \cong su(2) \oplus su(2)$
Consider massless particles: chirality $\sim$ helicity

Spinors (use chiral basis):

$$u^+(p) = v^-(p) = \begin{pmatrix} 0 \\ |p\rangle \end{pmatrix} \qquad u^-(p) = v^+(p) = \begin{pmatrix} |p] \\ 0 \end{pmatrix}$$

$$\bar{u}^+(p) = \bar{v}^-(p) = \begin{pmatrix} [p| & 0 \end{pmatrix} \qquad \bar{u}^-(p) = \bar{v}^+(p) = \begin{pmatrix} 0 & \langle p| \end{pmatrix}$$

- Amplitude written in terms of Lorentz-invariant spinor inner products

$$\langle ij\rangle = -\langle ji\rangle \equiv \langle i||j\rangle \text{ and } [ij] = -[ji] \equiv [i||j]$$

  - These are well known complex numbers, $\langle ij\rangle \sim [ij] \sim \sqrt{2p_i \cdot p_j}$

Spinor-Helicity: Vectors and Removing μ Indices

Lorentz algebra $so(3,1) \cong su(2) \oplus su(2)$

Dirac matrices in chiral basis

$$\gamma^\mu = \begin{pmatrix} 0 & \sqrt{2}\tau^\mu \\ \sqrt{2}\bar{\tau}^\mu & 0 \end{pmatrix} \qquad \sqrt{2}\tau^\mu = (1, \vec{\sigma}), \quad \sqrt{2}\bar{\tau}^\mu = (1, -\vec{\sigma}),$$

Remove vector indices with e.g.

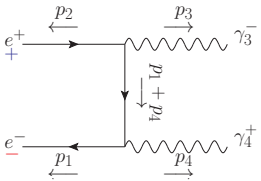
$$\underbrace{\langle i|\bar{\tau}^\mu|j][k|\tau_\mu|l\rangle = \langle il\rangle[kj]}_{\text{Fierz identity}}, \qquad \underbrace{\sqrt{2}p^\mu\tau_\mu \equiv \not{p} = |p]\langle p|}_{\text{Contraction with Pauli}}$$

Polarisation vectors ($r \equiv$ gauge choice, $r^2 = 0$, $r \cdot p \neq 0$):

$$\not{\epsilon}_+(p,r) = \frac{|p]\langle r|}{\langle rp\rangle}, \qquad \not{\epsilon}_-(p,r) = \frac{|r]\langle p|}{[pr]}$$

Every object can be written as left or right spinor(s)

An Illuminating Example: $e^+e^- \to \gamma\gamma$



Helas/MadGraph: explicit matrix multiplication

$$\sim \left[\bar{u}^-(p_1)\gamma^\mu \epsilon_\mu^+(p_4)\,(p_1^\nu + p_4^\nu)\,\gamma_\nu \gamma^\rho \epsilon_\rho^-(p_3) v^+(p_2)\right]$$

- Also cache and recycle various components
- No analytic simplification
- Most common numerical method

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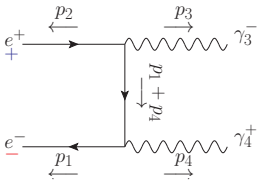

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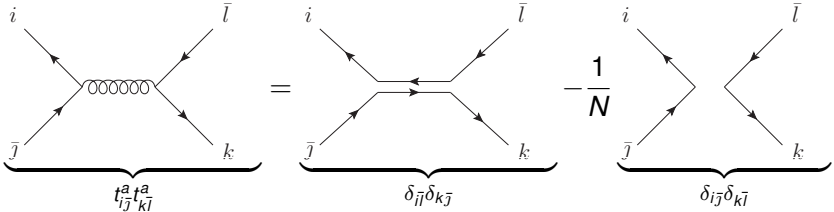
Can we systematically remove need for algebra or matrix multiplication?

Colour Flow: a Quick Introduction

Standard method in $SU(N)$-colour calculations:

Write all objects in terms of $\delta_{i\bar{j}} \equiv$ flows of colour (for simplicity $\mathrm{T_R} = 1$)
Calculations done pictorially, not via indices



Chirality Flow Building Blocks

Key idea: su(2) = su(N) (hep-ph:2003.05877)

Draw & connect lines to directly obtain inner products $\langle ij \rangle \sim [ij] \sim \sqrt{2p_i \cdot p_j}$
Removes need to do algebra or matrix multiplication

- Spinor inner products defined as connected lines

$$\langle i|^{\alpha}|j\rangle_{\alpha} \equiv \langle ij \rangle = -\langle ji \rangle =$$

$$[i|_{\dot\beta}|j]^{\dot\beta} \equiv [ij] = -[ji] =$$

- Spinors definitions follow

$$\bar{u}_i^- = \bar{v}_i^+ = \langle i|^{\alpha} =$$

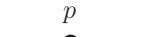
$$, \quad u_j^+ = v_j^- = |j\rangle_{\alpha} =$$

$$\bar{u}_i^+ = \bar{v}_i^- = [i|_{\dot\beta} =$$

$$, \quad u_j^- = v_j^+ = |j]^{\dot\beta} =$$

- Define slashed momentum as dot

$$\not{p} \equiv \sqrt{2}p^{\mu}\tau_{\mu}^{\dot\alpha\beta} =$$



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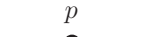
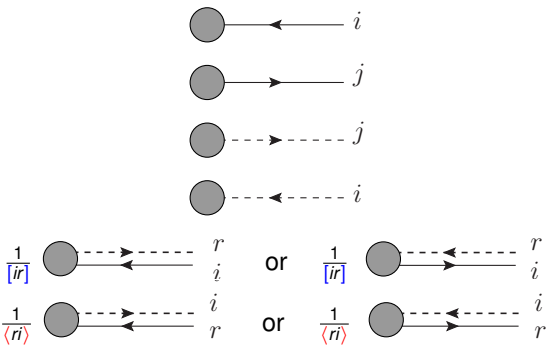
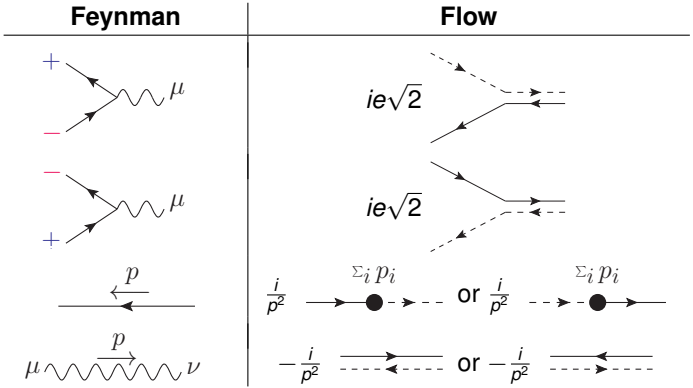


The Massless QED Flow Rules: External Particles

Species	Feynman	Flow
$\bar{u}^-(p_i)$		
$v^-(p_j)$		
$v^+(p_j)$		
$\bar{u}^+(p_i)$		
$\epsilon^\mu_-(p_i, r)$		
$\epsilon^\mu_+(p_i, r)$		

Left-chiral ≡ dotted lines right-chiral ≡ solid lines

The QED Flow Rules: Vertices and Propagators



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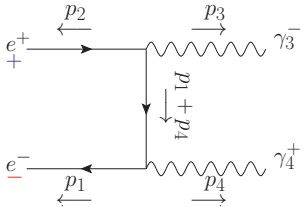

Spinor helicity analytically:

$$\sim \langle p_1|\bar{\tau}^\mu \underbrace{\left(|p_1\rangle\langle p_1| + |p_4\rangle\langle p_4|\right)}_{\not{p}_1+\not{p}_4} \bar{\tau}^\nu |p_2] \underbrace{\frac{\langle r_3|\bar{\tau}_\nu|p_3]}{\langle r_3 3\rangle}}_{\epsilon_3^-} \underbrace{\frac{[r_4|\tau_\mu|p_4\rangle}{[4r_4]}}_{\epsilon_4^+}$$

$$= \frac{\left(\langle p_1|\bar{\tau}^\mu|p_1] + \langle p_1|\bar{\tau}^\mu|p_4\rangle\right)[r_4|\tau_\mu|p_4\rangle \left(\langle p_1|\bar{\tau}^\nu|p_2] + \langle p_4|\bar{\tau}^\nu|p_2]\right)[p_3|\tau_\nu|r_3\rangle}{\langle r_3 3\rangle[4r_4]}$$

$$= \underbrace{\frac{\langle 1 r_4\rangle\left([41]\langle 13\rangle + [44]\langle 43\rangle\right)[r_3 2]}{\langle r_3 3\rangle[4r_4]}}_{\text{Fierz identities like } \langle i|\bar{\tau}^\mu|j][k|\tau_\mu|l\rangle = \langle il\rangle[kj]} = \underbrace{\frac{\langle 1 r_4\rangle[41]\langle 13\rangle[r_3 2]}{\langle r_3 3\rangle[4r_4]}}_{[ii]=0}$$

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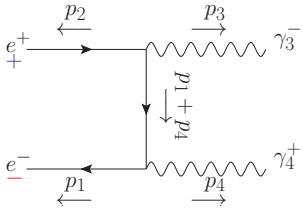
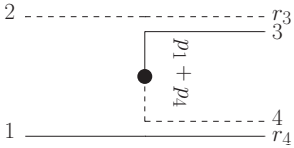


Chirality flow analytically:

$$\sim \frac{1}{\langle r_3 3\rangle [4 r_4]}$$



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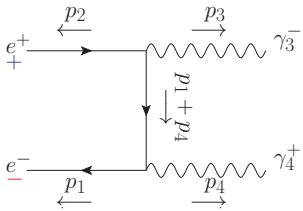

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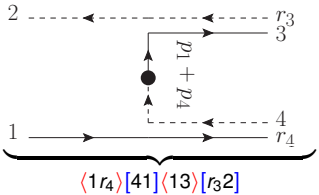
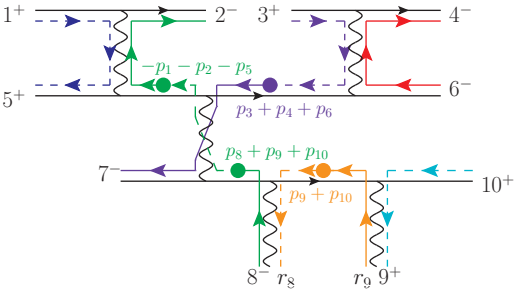


Recall:

$$\langle ij\rangle = -\langle ji\rangle = i \longrightarrow j$$

$$[ij] = -[ji] = i \dashrightarrow j$$

A complicated QED Example



Spinor-helicity analytic:

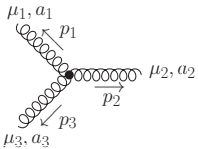
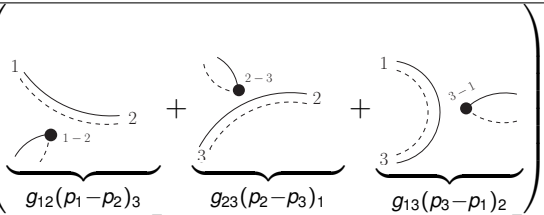
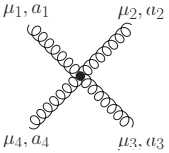
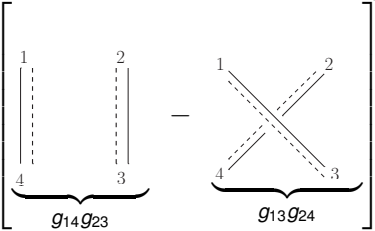
- 5 charge conjugation/Fierz + rearranging
- Not possible to fit on single slide!

$$= \underbrace{(\sqrt{2}ei)^8}_{\text{vertices}} \underbrace{\frac{(-i)^3}{s_{12}\, s_{34}\, s_{789\,10}}}_{\text{photon propagators}} \underbrace{\frac{(i)^4}{s_{125}\, s_{346}\, s_{89\,10}\, s_{9\,10}}}_{\text{fermion propagators}} \underbrace{\frac{1}{[8r_8]\langle r_9 9\rangle}}_{\text{polarization vectors}} [15]\langle 64\rangle[10\ 9]$$

$$\times \Big(\langle r_9 9\rangle[9r_8] + \langle r_9 10\rangle[10r_8]\Big)\Big(\underbrace{[33]}_{0}\langle 37\rangle + [34]\langle 47\rangle + [36]\langle 67\rangle\Big)$$

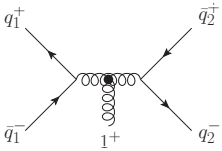
$$\times \Big(-\langle 89\rangle[91]\langle 12\rangle - \langle 89\rangle[95]\langle 52\rangle - \langle 8\,10\rangle[10\ 1]\langle 12\rangle - \langle 8\,10\rangle[10\ 5]\langle 52\rangle\Big)$$

The Non-abelian Massless QCD Flow Vertices

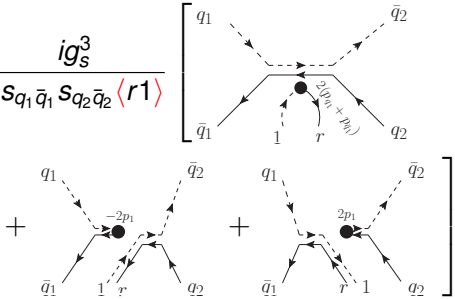
Feynman	Flow
	$-\frac{g_s f^{abc}}{2}$ 
	$ig_s^2 \sum_{Z(2,3,4)} f^{a_1 a_2 b} f^{b a_4 a_3}$ 

Arrow directions only consistently set within full diagram

QCD Example: $q_1\bar{q}_1 \to q_2\bar{q}_2 g$



$$= \frac{ig_s^3}{2s_{q_1\bar{q}_1}s_{q_2\bar{q}_2}\langle r1\rangle}$$



$$\left[\cdots\right] \equiv \Big\{2[q_1\bar{q}_2]\langle q_2\bar{q}_1\rangle\big([1q_1]\langle q_1 r\rangle + [1\bar{q}_1]\langle 1r\rangle\big)$$

$$- 2[q_1 1]\langle 1\bar{q}_1\rangle\langle q_2 r\rangle[1\bar{q}_2] + 2[q_1 1]\langle r\bar{q}_1\rangle\langle q_2 1\rangle[1q_2]\Big\}$$

MadGraph and the Automation of Chirality Flow

Summary

- So far: Numerical calculations use explicit multiplication rather than spin algebra analytically because quicker
- We have made the analytical spin algebra trivial
- Can we use this to make even faster numerics?

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Use MadGraph5_aMC@NLO (MG5aMC) for proof of concept automation

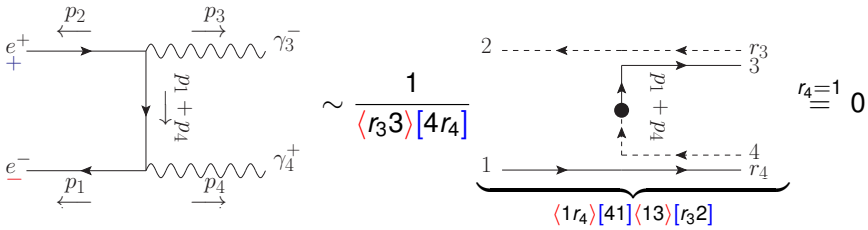
- Make minimal changes to massless QED in MG5aMC
- Pro: any difference in speed from our changes $\Rightarrow$ sound conclusions

Sources of Expect Speed Gains

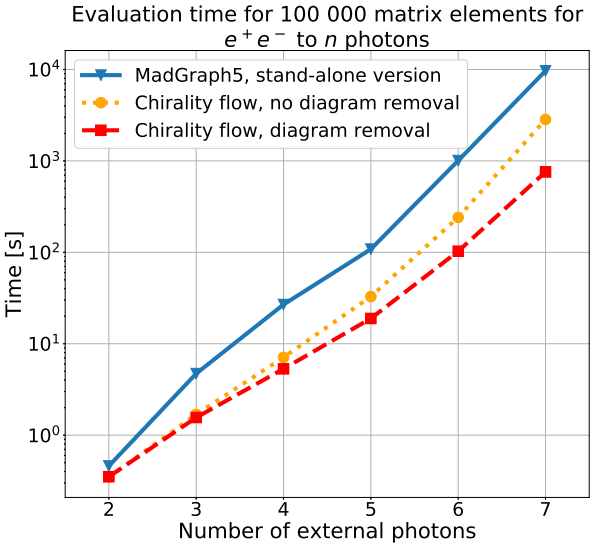
1. Simplified vertices and propagators
   - We minimise matrix multiplication
   - Each component of a calculation is simpler

Sources of Expect Speed Gains

1. Simplified vertices and propagators
   - We minimise matrix multiplication
   - Each component of a calculation is simpler
2. Gauge-based diagram removal
   - Polarisation vectors contain arbitrary gauge-reference spinor of momentum r
   - Spinor inner products antisymmetric $\Rightarrow \langle ii \rangle = [jj] = 0$
   - Chirality-flow makes optimal choice of r obvious $\Rightarrow$ remove diagrams!



Our Main Result (hep-ph:2203.13618)



Conclusions and Outlook

Conclusions:

- Chirality flow is the shortest route from Feynman diagram to complex number
- We have flow rules for full SM including masses at tree level
- We automised it for massless QED, found significant gains in MadGraph

Outlook and other work in this area:

- Use method analytically to calculate loop amplitudes
  - Ongoing work by AL, Simon Plätzer, and Malin Sjödahl,
- Automate for rest of (tree-level) Standard Model
  - Two current master students working to achieve this
  - First implementation of QCD working
- Simon Plätzer and Malin Sjödahl used chirality flow as basis for resummation (hep-ph:2204.03258)

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Spinor-Helicity: Gauge Bosons in Terms of Spinors

Lorentz algebra $so(3,1) \cong su(2) \oplus su(2)$
Consider massless particles: chirality $\sim$ helicity

Outgoing polarisation vectors:

$$\epsilon_+^\mu(p,r) = \frac{\langle r|\bar{\tau}^\mu|p]}{\langle rp\rangle}\,, \qquad \epsilon_-^\mu(p,r) = \frac{[r|\tau^\mu|p\rangle}{[pr]}$$

- r is a (massless) arbitrary reference momentum ($p \cdot r \neq 0$)
- Different r choices correspond to different gauges

$$\epsilon_+^\mu(p,r') - \epsilon_+^\mu(p,r) = -p^\mu \frac{\langle r'r\rangle}{\langle r'p\rangle\langle rp\rangle}$$

- Gauge invariant quantities must be r-invariant
  - Choose r as conveniently as possible (remember $\langle ij\rangle = -\langle ji\rangle$ s.t. $\langle ii\rangle = 0$)
    (4-gluon amplitude: can make 20/21 terms vanish)
  - Variance under $r \to r'$ good check of gauge invariance of (partial) amplitude

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Remove $\tau/\bar{\tau}$ matrices in amplitude with

$$\underbrace{\langle i|\bar{\tau}^\mu|j][k|\tau_\mu|l\rangle = \langle il\rangle[kj]}_{\text{Fierz identity}}, \qquad \underbrace{\langle i|\bar{\tau}^\mu|j] = [j|\tau^\mu|i\rangle}_{\text{Charge Conjugation}}$$

Express (massless) p^μ in terms of spinors

$$p^\mu = \frac{[p|\tau^\mu|p\rangle}{\sqrt{2}} = \frac{\langle p|\bar{\tau}^\mu|p]}{\sqrt{2}}, \quad \sqrt{2}p^\mu\tau_\mu \equiv \not{p} = |p]\langle p|, \quad \sqrt{2}p^\mu\bar{\tau}_\mu \equiv \bar{\not{p}} = |p\rangle[p|$$

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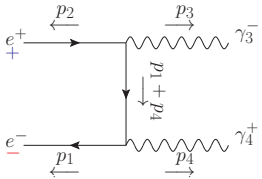

- $|p\rangle \equiv$ right-chiral spinor
- $|p] \equiv$ left-chiral spinor
- $\tau^\mu, \bar{\tau}^\mu \equiv$ Pauli matrices
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Spinor helicity: analytic

$$\sim \langle p_1|\bar{\tau}^\mu \underbrace{(|p_1]\langle p_1| + |p_4]\langle p_4|)}_{\not{p}_1+\not{p}_4} \bar{\tau}^\nu |p_2] \underbrace{\frac{\langle r_3|\bar{\tau}_\nu|p_3]}{\langle r_3 3\rangle}}_{\epsilon_3^-} \underbrace{\frac{[r_4|\tau_\mu|p_4\rangle}{[4r_4]}}_{\epsilon_4^+}$$

$$= \frac{(\langle p_1|\bar{\tau}^\mu|p_1] + \langle p_1|\bar{\tau}^\mu|p_4])[r_4|\tau_\mu|p_4\rangle\ (\langle p_1|\bar{\tau}^\nu|p_2] + \langle p_4|\bar{\tau}^\nu|p_2])[p_3|\tau_\nu|r_3\rangle}{\langle r_3 3\rangle[4r_4]}$$

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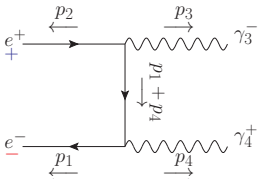

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$$\sim \left[\bar{u}^-(p_1)\gamma^\mu \epsilon^+_\mu(p_4)\left(p_1^\nu + p_4^\nu\right)\gamma_\nu\gamma^\rho\epsilon^-_\rho(p_3)v^+(p_2)\right]$$

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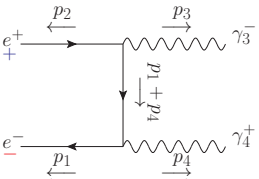

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Outgoing polarisation vectors ($r \equiv$ gauge choice, $r^2 = 0$, $r \cdot p \neq 0$):

$$\epsilon_+^\mu(p,r) = \frac{\langle r|\bar{\tau}^\mu|p]}{\langle rp\rangle}, \qquad \epsilon_-^\mu(p,r) = \frac{[r|\tau^\mu|p\rangle}{[pr]}$$

$$p\cdot\epsilon_+(p,r) = \underbrace{\frac{\langle r|p^\mu\bar{\tau}_\mu|p]}{\langle rp\rangle} = 0}_{\text{Weyl eq. } p^\mu\bar{\tau}_\mu|p]=0} \qquad p\cdot\epsilon_-^\mu(p,r) = \underbrace{\frac{[r|p^\mu\tau_\mu|p\rangle}{[pr]} = 0}_{\text{Weyl eq. } p^\mu\tau_\mu|p\rangle=0}$$

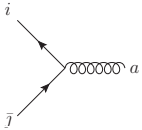
$$\epsilon_+(p,r)\cdot(\epsilon_-)^*(p,r) = \underbrace{\frac{\langle r|\bar{\tau}^\mu|p]}{\langle rp\rangle}\frac{[r|\tau_\mu|p\rangle}{[pr]}}_{\epsilon_\pm=(\epsilon_\mp)^*} = \frac{\langle rp\rangle[rp]}{\langle rp\rangle[pr]} = \underbrace{-1}_{[pr]=-[rp]}$$

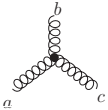
Colour Flow: a Quick Introduction

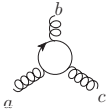
Standard method in $SU(N)$-colour calculations:

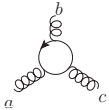
Write all objects in terms of $\delta_{i\bar{j}} \equiv$ flows of colour (for simplicity $\mathrm{T_R} = 1$)
Calculations done pictorially, not via indices

$\delta_{i\bar{j}} =$ 

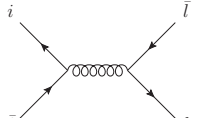
, $\sum_i \delta_{ii} = N =$ 
, $t^a_{i\bar{j}} =$ 

$if^{abc} =$ 

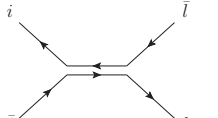
$=$ 

$-$ 

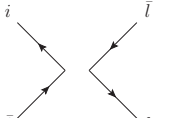
$= \mathrm{Tr}(t^a[t^b, t^c])$



$t^a_{i\bar{j}} t^a_{k\bar{l}}$

$=$ 

$\delta_{i\bar{l}}\delta_{k\bar{j}}$

$-\frac{1}{N}$ 

$\delta_{i\bar{j}}\delta_{k\bar{l}}$

Colour Flow: a Quick Introduction

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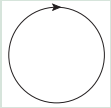
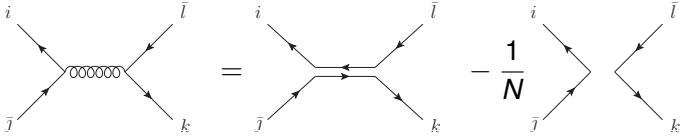
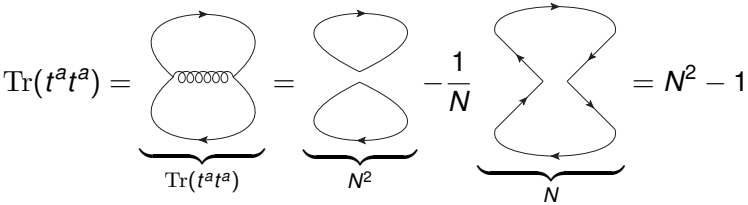
Calculations done pictorially, not via indices $\sum_i \delta_{ii} = N =$ 





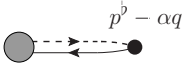
Massive Chirality Flow (hep-ph:2011.10075)

Decompose massive momentum into massless ones

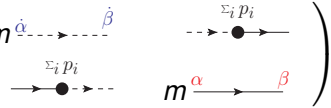
$$p^\mu = p^{\flat,\mu} + \alpha q^\mu \,, \quad (p^\flat)^2 = q^2 = 0 \,, \quad \alpha = \frac{p^2}{2p^\flat \cdot q}$$

- Spinors contain both chiralities, e.g.

$\bar{v}^-(p) =$ 

- Add new polarisation vector $\not{\epsilon}_0 = \frac{1}{m\sqrt{2}}$ 

- Need matrix structure in fermion propagators and vertices, e.g.

$p^\mu \gamma_\mu - m \sim$ 

Massive Chirality Flow (hep-ph:2011.10075)

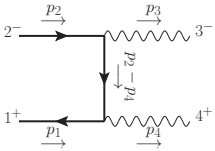
Main conclusion

Matrix structure unavoidable with massive fermions
Proceed as before to calculate without algebra

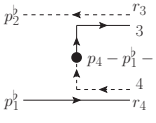
A Massive *Illuminating* Example

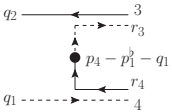
Consider the same diagram of $f_1^+ \bar{f}_2^- \to \gamma_3^+ \gamma_4^-$ as before but include mass m_f

- Obtain 3 new terms
- Simplify with choices of q_1, q_2, r_3, r_4
- $e^{i\varphi_i}\sqrt{\alpha_i} = \frac{m_i}{\langle p_i^\flat q_i \rangle}$, $\quad e^{-i\varphi_i}\sqrt{\alpha_i} = \frac{m_i}{[q_i p_i^\flat]}$

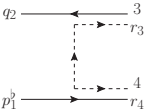


$$= \frac{-2ie^2}{(s_{23} - m_f^2)\langle r_3 3\rangle[4 r_4]} \Bigg\{ \sqrt{\alpha_1\alpha_2}\, e^{i(\varphi_2-\varphi_1)}$$

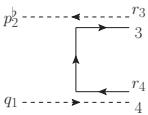




$$+ m_f \Bigg(\sqrt{\alpha_2}\, e^{i\varphi_2}$$



$$- \sqrt{\alpha_1}\, e^{-i\varphi_2}$$

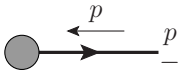
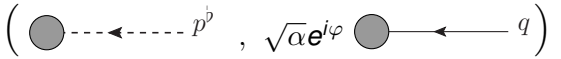
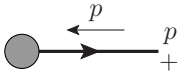
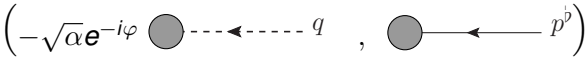
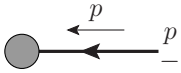
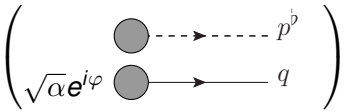
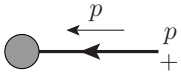
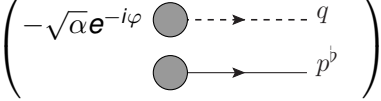


$$\Bigg) \Bigg\}$$

Incoming Massive Spinors in Chirality Flow

$$p^\mu = p^{\flat,\mu} + \alpha q^\mu\,, \quad (p^\flat)^2 = q^2 = 0\,, \quad e^{i\varphi}\sqrt{\alpha} = \frac{m}{\langle p^\flat q\rangle}\,, \qquad e^{-i\varphi}\sqrt{\alpha} = \frac{m}{[q p^\flat]}$$

Spin operator $-\frac{\Sigma^\mu s_\mu}{2} = \frac{\gamma^5 s^\mu \gamma_\mu}{2}, \quad s^\mu = \frac{1}{m}(p^{\flat,\mu} - \alpha q^\mu)$

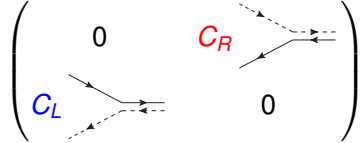
Spinor	Feynman	Flow
$\bar{v}^-(p)$		
$\bar{v}^+(p)$		
$u^-(p)$		
$u^+(p)$		

Some Fermion Flow Rules

$$p^\mu = p^{\flat,\mu} + \alpha q^\mu \,, \quad (p^\flat)^2 = q^2 = 0 \,, \quad \alpha = \frac{p^2}{2p\cdot q} \neq 0$$

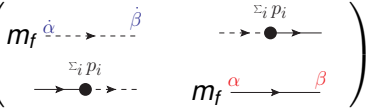
Fermion-vector vertex

$$= ie(P_L C_L + P_R C_R)\gamma^\mu = ie\sqrt{2}$$



Fermion propagator

$$\frac{i}{p^2 - m_f^2}\begin{pmatrix} m_f \delta^{\dot\alpha}{}_{\dot\beta} & \sqrt{2}p^{\dot\alpha\beta} \\ \sqrt{2}\bar{p}_{\alpha\dot\beta} & m_f \delta_\alpha{}^\beta \end{pmatrix} = \frac{i}{p^2 - m_f^2}$$

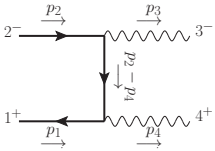


Left and right chiral couplings may differ

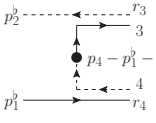
A Massive *Illuminating* Example

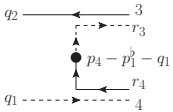
Consider the same diagram of $f_1^+ \bar{f}_2^- \to \gamma_3^+ \gamma_4^-$ as before but include mass m_f

- Obtain 3 new terms
- Simplify with choices of q_1, q_2, r_3, r_4
- $e^{i\varphi_i}\sqrt{\alpha_i} = \frac{m_i}{\langle p_i^\flat q_i \rangle}$, $\quad e^{-i\varphi_i}\sqrt{\alpha_i} = \frac{m_i}{[q_i p_i^\flat]}$

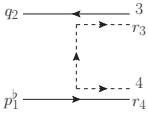


$$= \frac{-2ie^2}{(s_{23} - m_f^2)\langle r_3 3\rangle[4 r_4]} \Bigg\{ \sqrt{\alpha_1\alpha_2}\, e^{i(\varphi_2-\varphi_1)}$$

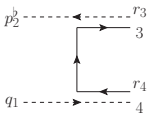




$$+ m_f \Bigg(\sqrt{\alpha_2}\, e^{i\varphi_2}$$



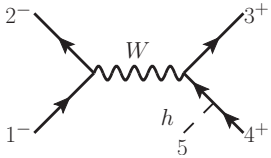
$$- \sqrt{\alpha_1}\, e^{-i\varphi_2}$$



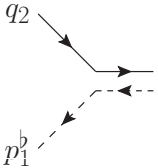
$$\Bigg) \Bigg\}$$

A Second Massive Example: $f_1\bar{f}_2 \to W \to f_3\bar{f}_4 h_5$

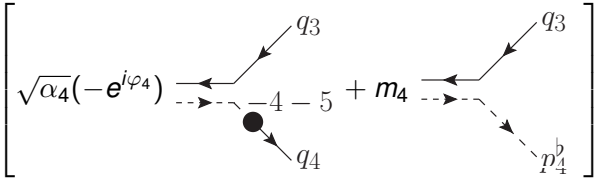
- W bosons simplifies ($C_R = 0$)
- Simplify with choices of $q_1, \cdots q_5$
- $e^{i\varphi_i}\sqrt{\alpha_i} = \frac{m_i}{\langle p_i^\flat q_i\rangle}$, $\quad e^{-i\varphi_i}\sqrt{\alpha_i} = \frac{m_i}{[q_i p_i^\flat]}$
- Scalar has no flow line



Step 1: Draw fermion lines: $\sim C_{L,12}\sqrt{\alpha_2}e^{i\varphi_2}$



$\times\, C_{L,34}\sqrt{\alpha_3}(-e^{i\varphi_3})$



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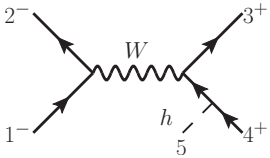
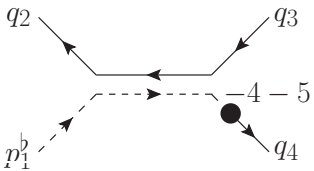
- Scalar has no flow line

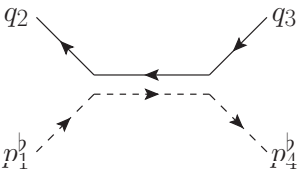


Step 2: Flip arrows and connect: $C_{L,12}C_{L,34}\sqrt{\alpha_2\alpha_3}e^{i(\varphi_2+\varphi_3)}$

$\times \Bigg[\sqrt{\alpha_4}e^{i\varphi_4}$



$- \; m_4$



$\Bigg]$

Lorentz Group Representations

Lorentz group elements: $e^{i(\theta_i J_i + \eta_i K_i)}$ $\quad J_i \equiv$ rotations, $\quad K_i \equiv$ boosts

- Lorentz group generators $\simeq$ 2 copies of su(2) generators
  - $so(3,1)_{\mathbb{C}} \cong su(2) \oplus su(2)$

Group algebra defined by commutator relations

$$[J_i, J_j] = i\epsilon_{ijk} J_k, \quad [J_i, K_j] = i\epsilon_{ijk} K_k, \quad [K_i, K_j] = -i\epsilon_{ijk} J_k$$

$$N_i^{\pm} = \frac{1}{2}(J_i \pm iK_i)\ , \quad [N_i^-, N_j^+] = 0\ ,$$

$$[N_i^-, N_j^-] = i\epsilon_{ijk} N_k^-\ , \qquad [N_i^+, N_j^+] = i\epsilon_{ijk} N_k^+$$

- Representations (i.e. realisations of $N_i^{\pm}$)
  - $(0,0)$ scalar particles
  - $(\frac{1}{2},0)$ left-chiral and $(0,\frac{1}{2})$ right-chiral Weyl (2-component) spinors.
  - $(\frac{1}{2},0) \oplus (0,\frac{1}{2})$, Dirac (4-component) spinors.
  - $(\frac{1}{2},\frac{1}{2})$ vectors, e.g. gauge bosons

How to Calculate? Spinor-Helicity

Give each particle a defined helicity $\Rightarrow$ amplitude now a number!

Spinors (in chiral basis):

$$u^+(p) = v^-(p) = \begin{pmatrix} 0 \\ |p\rangle \end{pmatrix} \qquad u^-(p) = v^+(p) = \begin{pmatrix} |p] \\ 0 \end{pmatrix}$$

$$\bar{u}^+(p) = \bar{v}^-(p) = \begin{pmatrix} [p| & 0 \end{pmatrix} \qquad \bar{u}^-(p) = \bar{v}^+(p) = \begin{pmatrix} 0 & \langle p| \end{pmatrix}$$

$$\gamma^\mu = \begin{pmatrix} 0 & \sqrt{2}\tau^\mu \\ \sqrt{2}\bar{\tau}^\mu & 0 \end{pmatrix} \qquad \sqrt{2}\tau^\mu = (1, \vec{\sigma}), \quad \sqrt{2}\bar{\tau}^\mu = (1, -\vec{\sigma}),$$

- Amplitude written in terms of Lorentz-invariant spinor inner products

$$\langle ij\rangle = -\langle ji\rangle \equiv \langle i||j\rangle \text{ and } [ij] = -[ji] \equiv [i||j]$$

  - These are well known complex numbers, $\langle ij\rangle \sim [ij] \sim \sqrt{2p_i \cdot p_j}$
- Remove $\tau/\bar{\tau}$ matrices in amplitude with

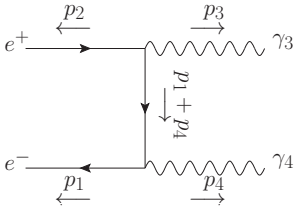
$$\langle i|\bar{\tau}^\mu|j][k|\tau_\mu|l\rangle = \langle il\rangle[kj] \,, \qquad \langle i|\bar{\tau}^\mu|j] = [j|\tau^\mu|i\rangle$$

How to Calculate a Process

Sum all Feynman diagrams, square, and integrate

Often spin structure is non-trivial

e.g.



$$\sim \underbrace{\left[\bar{u}(p_1)\gamma^\mu \left(p_1^\nu + p_4^\nu\right) \gamma_\nu \gamma^\rho v(p_2)\right] \epsilon_\rho(p_3)\epsilon_\mu(p_4)}_{\text{A mathematical expression we have simplify and square}}$$

Most common method: use helicity basis

Each diagram is a complex number, easy to square
Can use algebra to simplify first, or brute force matrix multiplication

Define Problem

Kinematic part of amplitude slowed by spin and vector structures

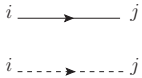
- Can we still improve on this?
  - Deriving spinor inner products $\langle ij \rangle$, $[kl]$ requires at least 2 steps
    - Re-write every object as spinors
    - Use Fierz identity $\bar{\tau}^{\mu}_{\alpha\dot{\beta}} \tau_{\mu}^{\dot{\alpha}\beta} = \delta_{\alpha}^{\beta} \delta^{\dot{\alpha}}_{\dot{\beta}}$
  - Not intuitive which inner products we obtain
- In SU(N) use graphical reps for calculations
  - E.g. using the colour-flow method
  - (Also birdtracks etc.)
- Spinor-helicity $\equiv su(2) \oplus su(2)$
  - Can we use graphical reps?

Creating Chirality Flow: Building Blocks

A flow is a directed line from one object to another
$su(2)$ objects have dotted indices and $su(2)$ objects undotted indices

- First step: Ansatz for spinor inner products (only possible Lorentz invariant)

$$\langle i|^{\alpha}|j\rangle_{\alpha} \equiv \langle ij\rangle = -\langle ji\rangle =$$



$$[i|_{\dot\beta}|j]^{\dot\beta} \equiv [ij] = -[ji] =$$

- Spinors and Kronecker deltas follow

$$\langle i|^{\alpha} =$$

,

$$|j\rangle_{\alpha} =$$

$$[i|_{\dot\beta} =$$

,

$$|j]^{\dot\beta} =$$

$$\delta_{\alpha}{}^{\beta} \equiv \mathbb{1}_{\alpha}{}^{\beta} =$$

,

$$\delta^{\dot\beta}{}_{\dot\alpha} \equiv \mathbb{1}^{\dot\beta}{}_{\dot\alpha} =$$