

Exclusive J/ψ photoproduction in nucleus-nucleus UPCs at the LHC in NLO QCD



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ERC adG YoctoLHC

Based on K.J. Eskola, C.A. Flett, V. Guzey, T. Löytäinen, H. Paukkunen, PRC 106 (2022) 035202 and arXiv:2210.16048 [hep-ph]

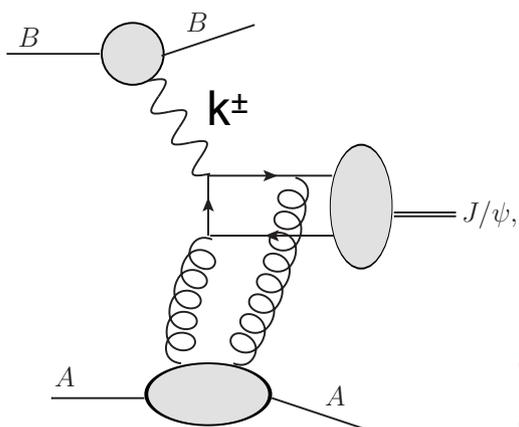
Outline:

- Ultraperipheral collisions (UPCs) at the LHC, J/ψ photoproduction and new constraints on small- x nuclear gluon density
- Exclusive J/ψ photoproduction in Pb-Pb and O-O UPCs@LHC in NLO pQCD: strong scale dependence, uncertainties due to nPDFs, and quark dominance
- Summary and outlook

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UPCs@LHC and J/ψ photoproduction

- **Ultraperipheral collisions (UPCs)**: ions pass each other at large impact parameters $b \gg R_A + R_B \rightarrow$ strong interactions suppressed \rightarrow interaction via quasi-real photons in Weizsäcker-Williams equivalent photon approximation, Budnev, Ginzburg, Meledin, Serbo, Phys. Rept. 15 (1975) 181
- UPCs@LHC allow one to study $\gamma\gamma$, γp and γA interactions at unprecedentedly high energies \rightarrow can be used to study open questions of **proton and nucleus structure in QCD** and search for new physics \rightarrow **new info on quark and gluon distributions in nuclei at small x** , Bertulani, Klein, Nystrand, Ann. Rev. Nucl. Part. Sci. 55 (2005) 271; Baltz et al, Phys. Rept. 480 (2008) 1; Contreras and Tapia-Takaki, Int. J. Mod. Phys. A 30 (2015) 1542012; Snowmass Lol, Klein et al, arXiv:2009.03838
- Most studied process is coherent J/ψ photoproduction in Pb-Pb/pp UPCs



$$\frac{d\sigma^{AB \rightarrow AJ/\psi B}}{dy} = \left[k \frac{dN_{\gamma/B}}{dk} \sigma^{\gamma A \rightarrow J/\psi A} \right]_{k=k^+} + \left[k \frac{dN_{\gamma/A}}{dk} \sigma^{\gamma B \rightarrow J/\psi B} \right]_{k=k^-}$$

Photon flux from QED:
 - high intensity $\sim Z^2$
 - high photon energy $\sim \gamma_L$

Photoproduction cross section

$$k^\pm = \frac{M_{J/\psi}}{2} e^{\pm y}$$

Photon momentum k^\pm from J/ψ rapidity y

Constraints on small-x gluon density in nuclei

- In leading logarithmic approximation (LLA) of pQCD, Ryskin, Z. Phys. C57 (1993) 89; Frankfurt, Koepf, Strikman, PRD 57 (1998) 512; Frankfurt, McDermott, Strikman, JHEP 03 (2001) 045

$$\frac{d\sigma^{\gamma p \rightarrow J/\psi p}(t=0)}{dt} = \frac{12\pi^3 \Gamma_V M_V^3}{\alpha_{\text{e.m.}} (4m_c^2)^4} [\alpha_s(Q_{\text{eff}}^2) x g(x, Q_{\text{eff}}^2)]^2 C(Q^2=0)$$

Γ_V is J/ψ leptonic decay width

gluon density at $x=(M_{J/\psi})^2/W^2$ and $Q_{\text{eff}}^2=2.5-3 \text{ GeV}^2$

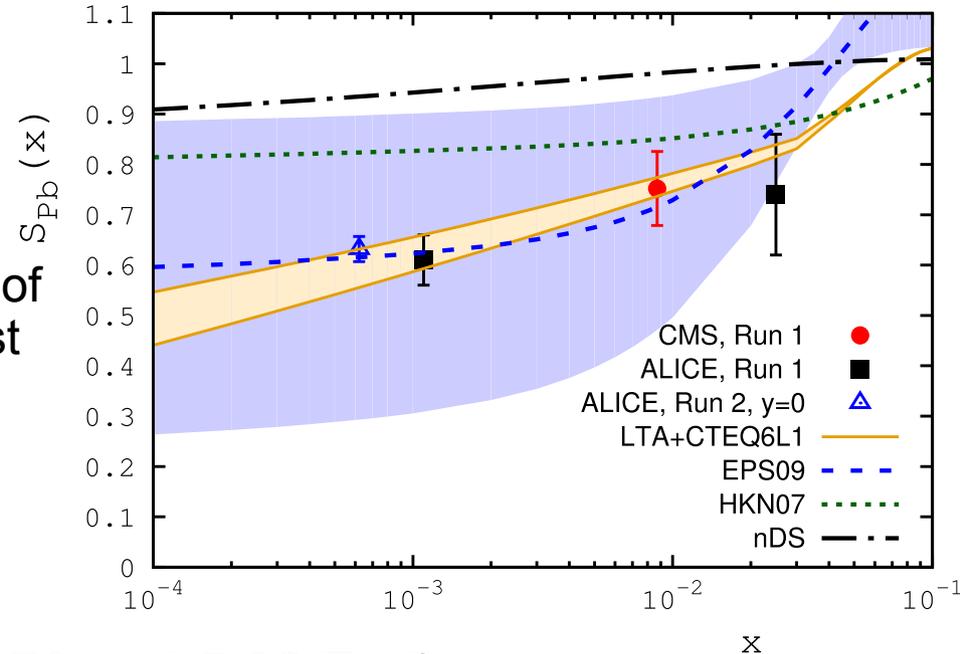
depends on details of J/ψ distribution amplitude

- Ratio of nucleus and IA cross sections \rightarrow nuclear suppression factor $S_{\text{Pb}}(x)$

$$S_{\text{Pb}}(x) = \left[\frac{\sigma_{\gamma \text{Pb} \rightarrow J/\psi \text{Pb}}}{\sigma_{\gamma \text{Pb} \rightarrow J/\psi \text{Pb}}^{\text{IA}}} \right]^{1/2} = \kappa_{A/N} \frac{g_A(x, \mu^2)}{A g_N(x, \mu^2)}$$

Model-independently using data on Pb-Pb UPCs@LHC, Abelev *et al.* [ALICE], PLB718 (2013) 1273; Abbas *et al.* [ALICE], EPJ C 73 (2013) 2617; [CMS] PLB 772 (2017) 489; Acharya *et al.* [ALICE], arXiv:2101:04577 [nucl-ex]

From global QCD fits of nPDFs or leading twist nuclear shadowing model, Guzey, Kryshen, Strikman, Zhalov, PLB 726 (2013) 290; Guzey, Zhalov, JHEP 1310 (2013) 207



- Good agreement with ALICE data at $y=0$ (2.76 and 5.02 TeV) \rightarrow direct evidence of large gluon shadowing, $R_g(x=6 \times 10^{-4} - 0.001) \approx 0.6$

Exclusive J/ψ photoproduction in NLO pQCD

- Collinear factorization for hard exclusive processes, [Collins, Frankfurt, Strikman, PRD 56 \(1997\) 2982](#): $\gamma A \rightarrow J/\psi A$ amplitude in terms of generalized parton distribution functions (GPDs), [Ji, PRD 55 \(1997\) 7114](#); [Radyushkin PRD 56 \(1997\) 5524](#); [Diehl, Phys. Rept. 388 \(2003\) 41](#)

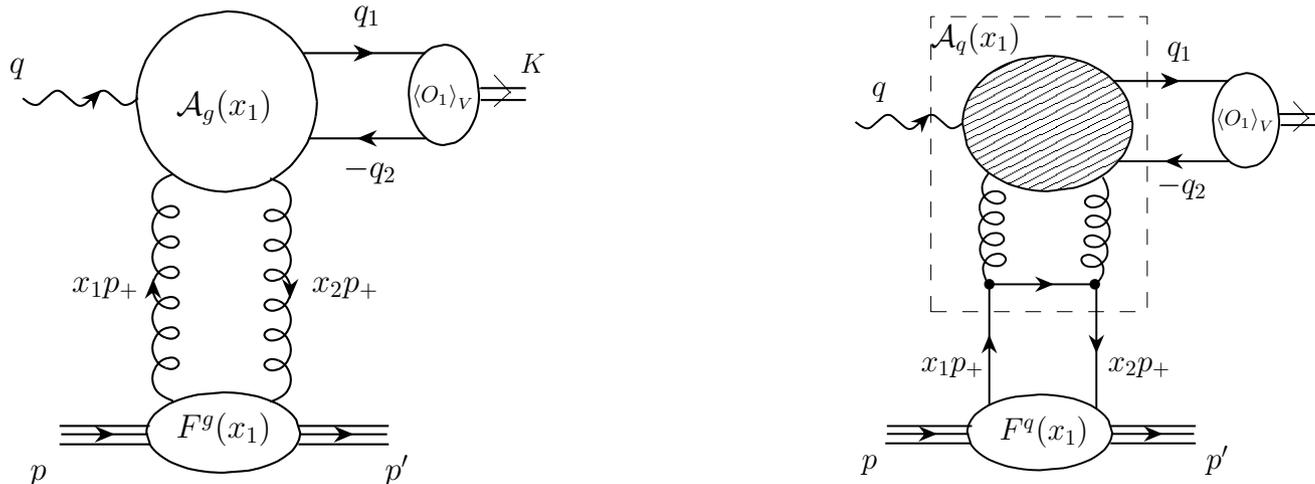
- To next-to-leading order (NLO) of perturbative QCD, [Ivanov, Schafer, Szymanowski, Krasnikov, EPJ C 34 \(2004\) 297, 75 \(2015\) 75 \(Erratum\)](#); [Jones, Martin, Ryskin, Teubner, J. Phys. G: Nucl. Part. Phys. 43 \(2016\) 035002](#)

$$\mathcal{M}^{\gamma A \rightarrow J/\psi A} \propto \sqrt{\langle O_1 \rangle_{J/\psi}} \int_{-1}^1 dx \left[T_g(x, \xi) F_A^g(x, \xi, t, \mu_F) + T_q(x, \xi) F_A^q(x, \xi, t, \mu_F) \right]$$

↓
↓
↓

NRQCD matrix element from J/ψ leptonic decay
pQCD coefficient function
Gluon GPD
Quark contribution

- To leading order (LO), only gluons; both quarks and gluons at NLO.



Exclusive J/ψ photoproduction in NLO pQCD (2)

- In the limit of **high W** corresponding to **small $\xi=(1/2)(M_{J/\psi})^2/W^2 \ll 1$**

$$\mathcal{M}^{\gamma A \rightarrow J/\psi A} \propto i \sqrt{\langle O_1 \rangle_{J/\psi}} \left[F_A^g(\xi, \xi, t, \mu_F) + \frac{\alpha_s N_c}{\pi} \ln \left(\frac{m_c^2}{\mu_F^2} \right) \int_{\xi}^1 \frac{dx}{x} F^g(x, \xi, t) \right. \\ \left. + \frac{\alpha_s C_F}{\pi} \ln \left(\frac{m_c^2}{\mu_F^2} \right) \int_{\xi}^1 dx (F^{q,S}(x, \xi, t) - F^{q,S}(-x, \xi, t)) \right]$$

→ helps to qualitatively understand the features of our numerical calculations.

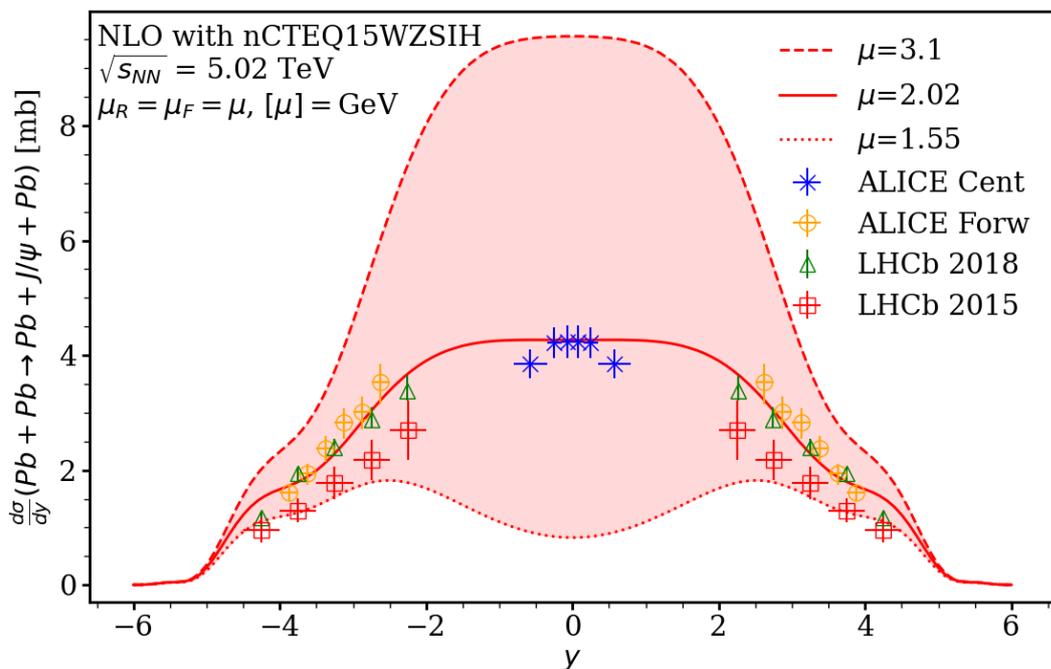
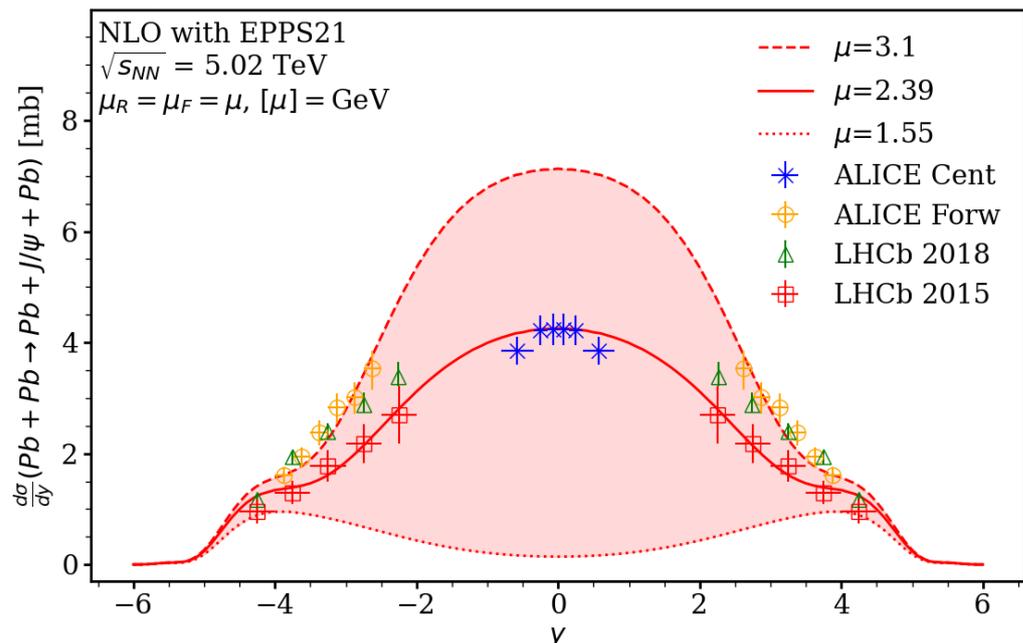
- GPDs are hybrid distributions interpolating between **usual PDFs** and **form factors** → depend on momentum fractions **x** and **ξ** and momentum transfer **t** .
- Connection between GPDs is necessarily model-dependent. In our analysis, we neglect dependence of GPDs on **ξ** and used the **forward model**, Freund, McDermott, Strikman, PRD 67 (2003) 036001. For gluons (quarks are similar):

$$F_A^g(x, \xi, t, \mu_F) = x g_A(x, \mu_F) F_A(t)$$

Nuclear PDFs: EPPS16, nCTEQ15,
nNNPDF2.0 + update with EPPS21,
nCTEQ15WZSIH, nNNPDF3.0

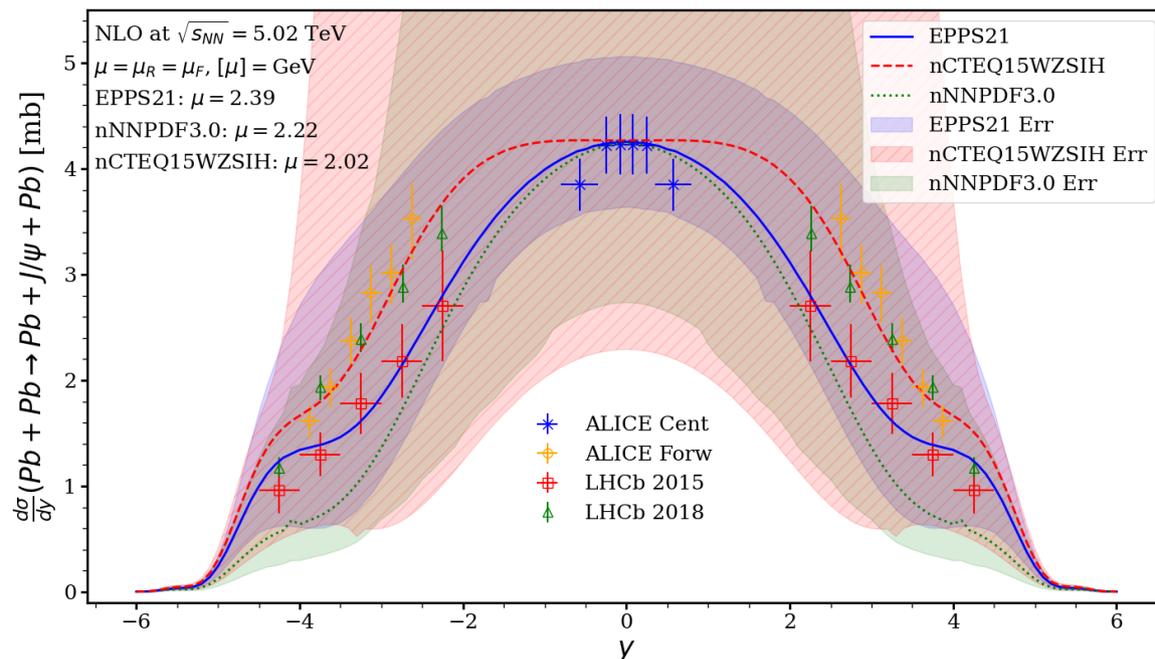
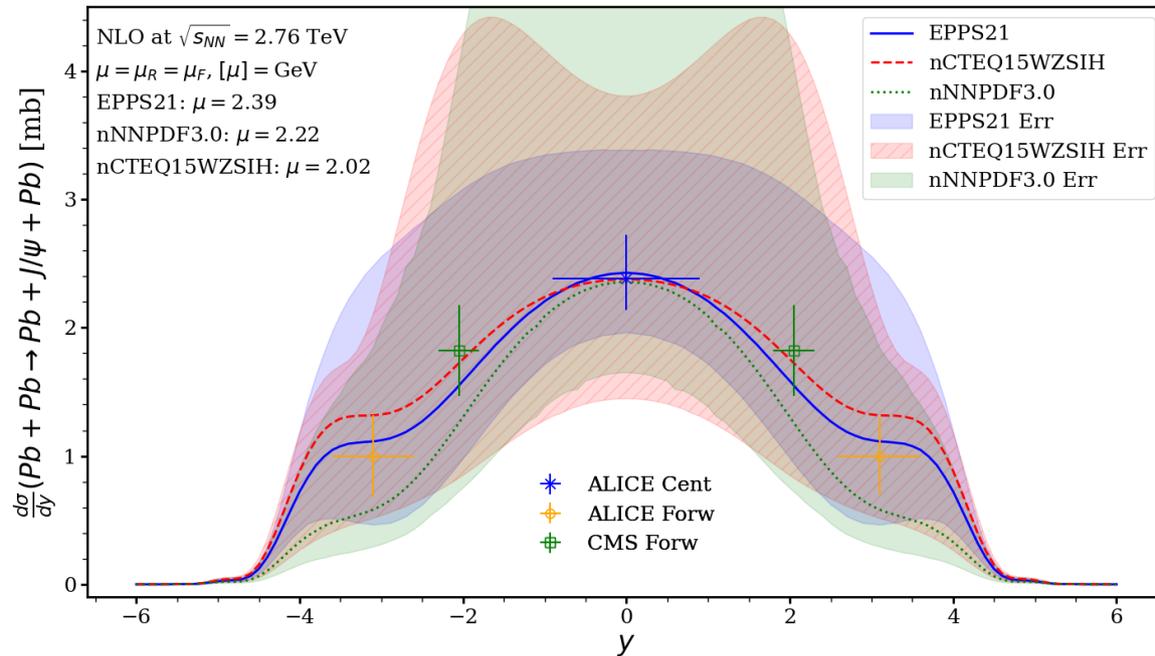
Nucleus form factor
(Woods-Saxon form)

Scale dependence and comparison to data on J/ψ photoproduction in Pb-Pb UPCs (Runs 1&2)



- Scale dependence of our NLO pQCD results for $m_c \leq \mu_F \leq M_{J/\psi}$ is **very strong**.
- One can find an “**optimal scale**” $\mu_F=2.39$ GeV (EPPS21) giving simultaneous good description of Run 1&2 UPC data → **note that $\gamma+p \rightarrow J/\psi+p$ proton data is somewhat overestimated**.
- Note that updated LHCb data have moved up worsening the agreement with **EPPS21**.
- The agreement is restored by using **nCTEQ15WZSIH** nPDFs characterized by large strange quark density → **sensitivity to strange quarks in nuclei?**

Uncertainties due to nuclear PDFs

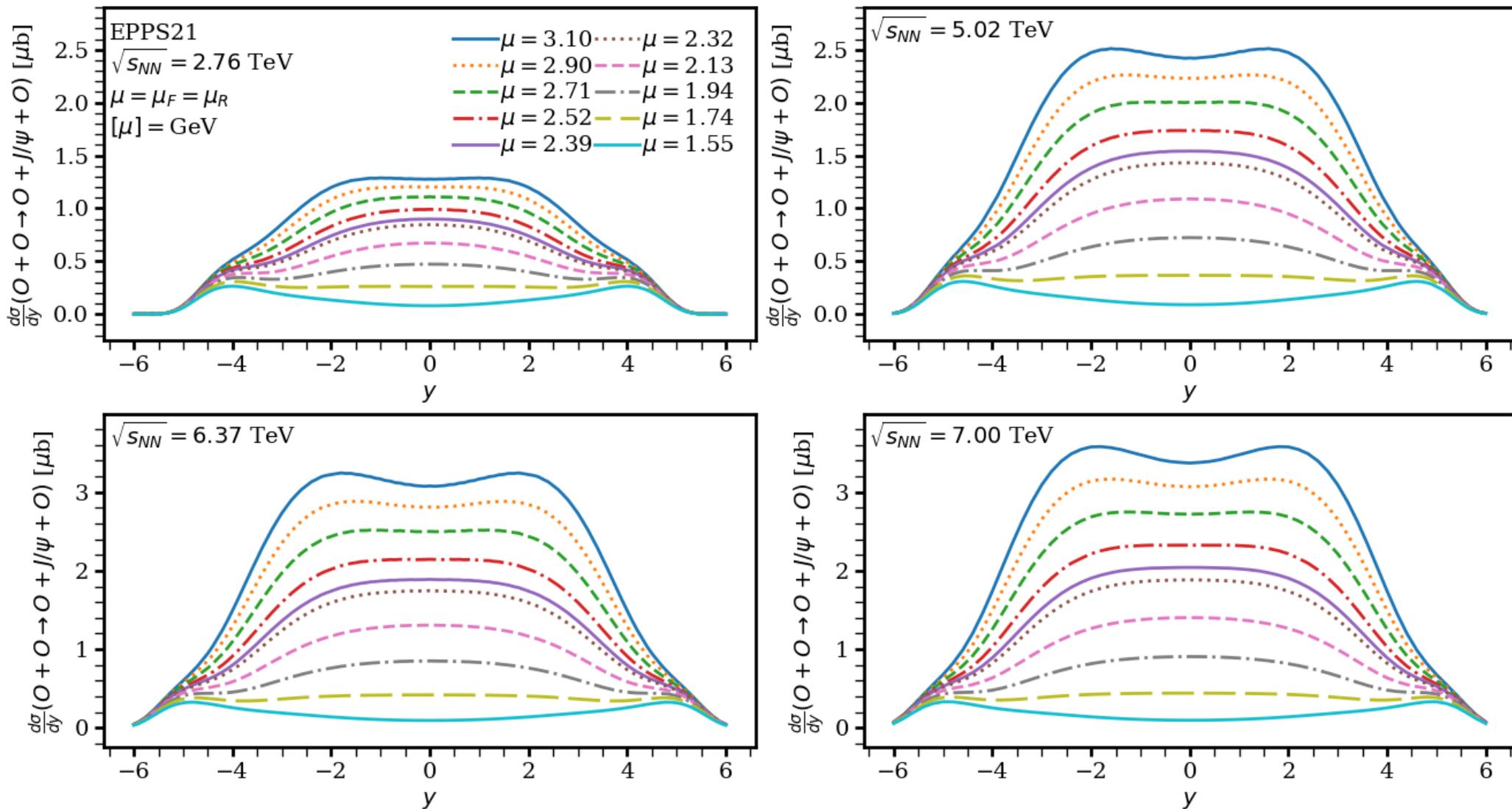


- Uncertainties due to nPDFs are quite significant → **opportunity to reduce** them using the data on J/ψ photoproduction in AA UPCs.

- Compared to our original calculations, abnormally large uncertainty associated with **EPPS16** disappears when using more recent **EPPS21**.

- The **nNNPDF3.0** nPDFs correspond to much less constrained fit → large uncertainties.

NLO pQCD predictions for O-O UPCs



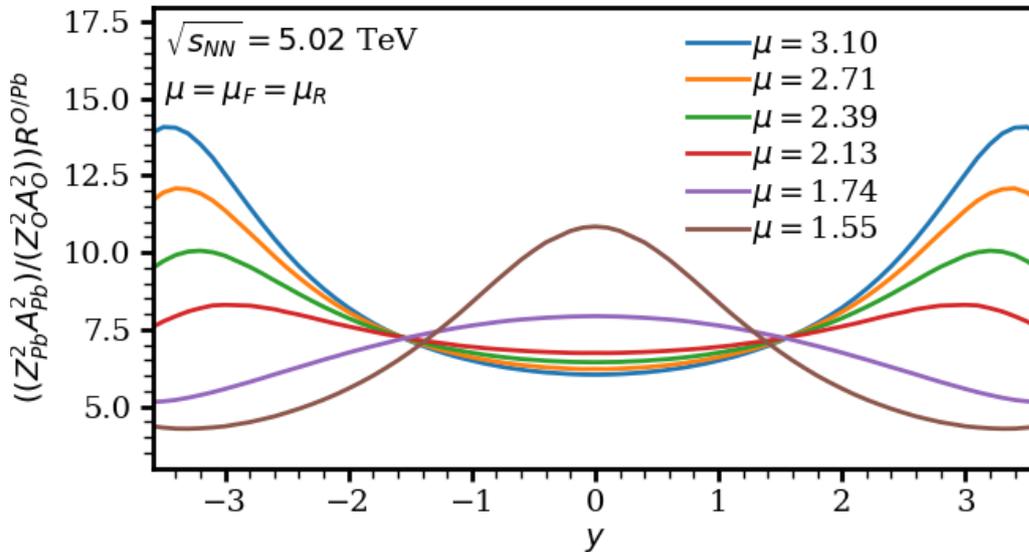
- NLO pQCD predictions for anticipated O-O run with 4 options for $\sqrt{s_{NN}} \rightarrow$ similar trends as for Pb-Pb UPCs \rightarrow large scale and nuclear PDF uncertainties.

Reduction of uncertainties using O/Pb ratio

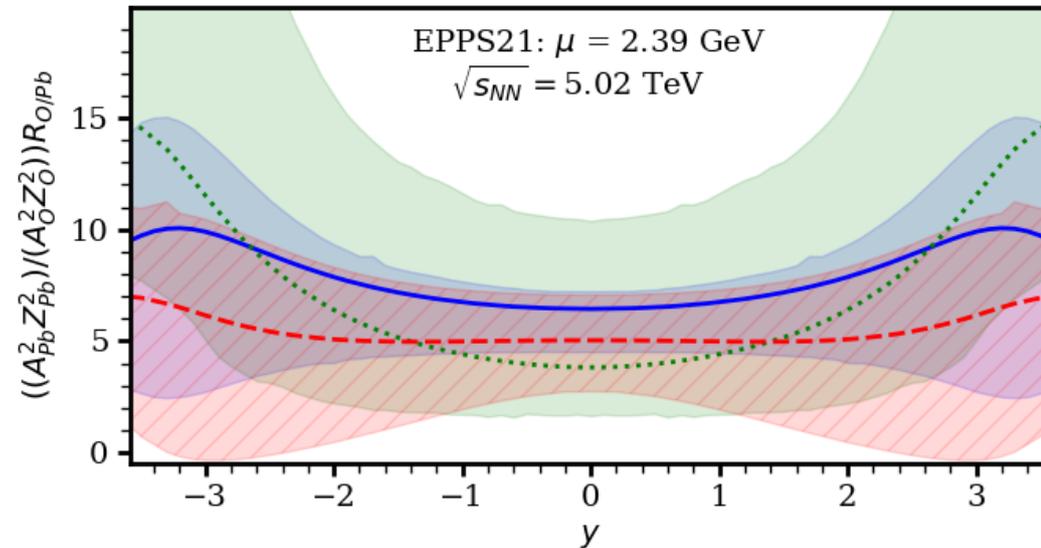
- One can reduce the significant scale μ_F and nPDF uncertainties by considering the ratio of oxygen to lead UPC cross sections:

$$R^{O/Pb} = \left(\frac{208Z_{Pb}}{16Z_O} \right)^2 \frac{d\sigma(O + O \rightarrow O + J/\psi + O)/dy}{d\sigma(Pb + Pb \rightarrow Pb + J/\psi + Pb)/dy}$$

Scale uncertainty of $R^{O/Pb}$



nPDF uncertainty of $R^{O/Pb}$

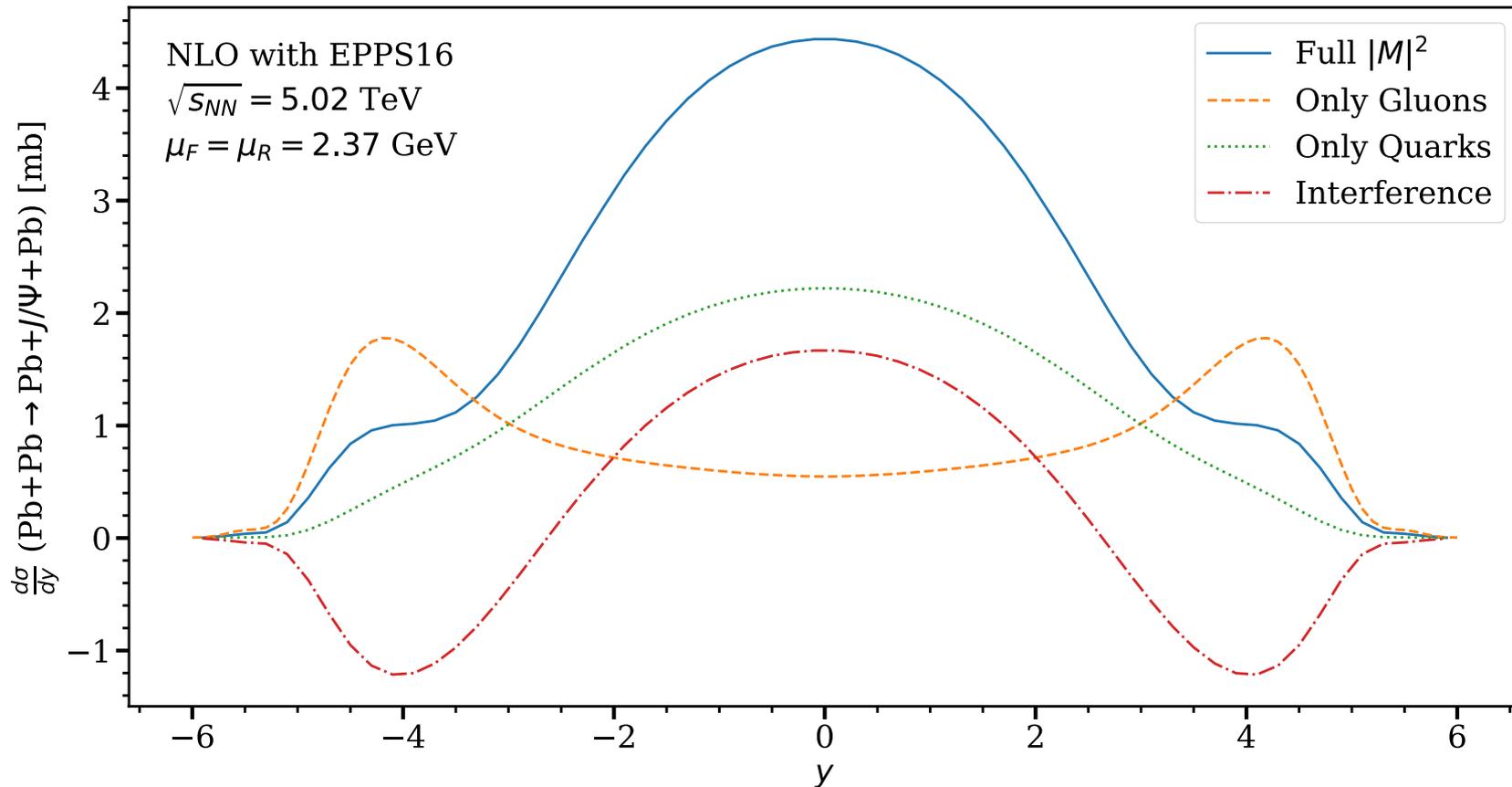


- Hard scattering coefficient functions for O and Pb are the same \rightarrow differences come from O and Pb nPDFs \rightarrow scale dependence reduced by factor of 10 compared to individual UPC cross sections.

- Reduction of nPDF uncertainties is also large due to additional partial cancellation of uncertainties associated with proton PDFs.

Dominance of quark contribution

- The most striking result is strong cancellations between LO and NLO gluons → **dominance of quark contribution** at central rapidities.



- At the face value, **this totally changes** the interpretation of data on coherent J/ψ photoproduction in heavy-ion UPCs as a probe of small- x nuclear gluons.

Summary and Outlook

- First NLO pQCD calculation of exclusive J/ψ photoproduction in Pb-Pb and O-O UPCs@LHC in the framework of collinear factorization.
- Our analysis confirmed strong scale dependence noticed earlier, quantified uncertainty due to nuclear PDFs, observed the dominance of the quark contribution, and provided simultaneous description of Run 1&2 LHC data.
- From phenomenology point of view, the ultimate goal is to use these UPC data in global QCD fits.
- In the present form, this is challenging. Possible solutions:
 - ❖ Consider ratio of AA to OO/pp UPC cross sections, where most of complications (scale dependence, uncertainties of nPDFs, details of GPD modeling and its Q^2 evolution, relativistic corrections to the charmonium wave function) partially cancel.
 - ❖ Even in the case of the UPC cross section ratios, non-relativistic corrections to charmonium wave function do not cancel exactly and should be taken into account, Eskobedo, Lappi, PRD 101 (2020) 3, 034030; Lappi, Mantysaari, Penttala, PRD 102 (2020) 5, 054020