

Soft gluon resummation for the production of four top quarks at the LHC

Laura Moreno Valero

in collaboration with Melissa van Beekveld and Anna Kulesza

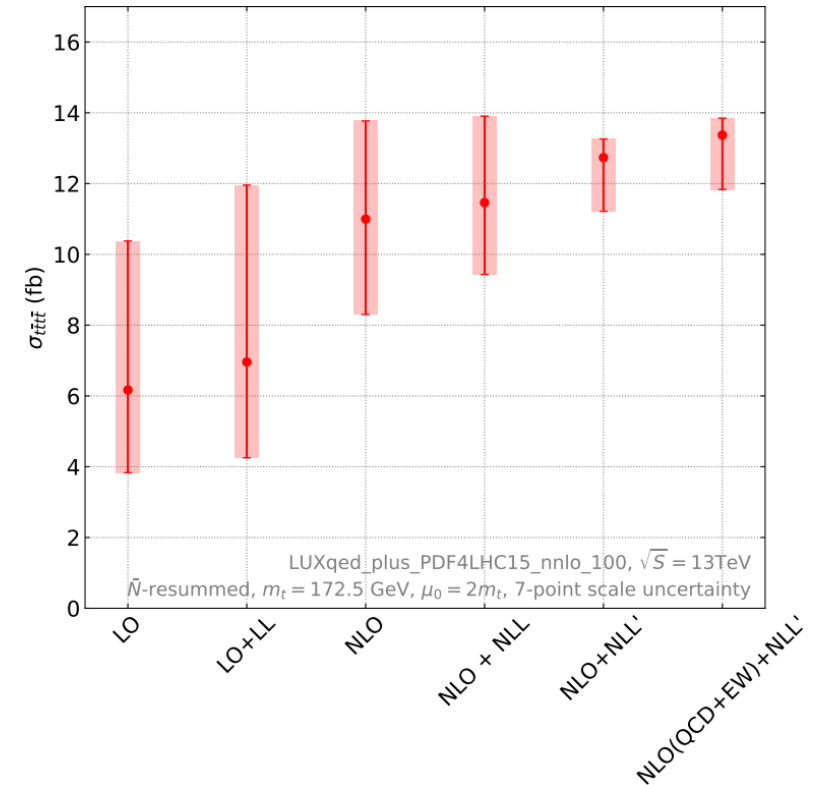
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QCD@LHC

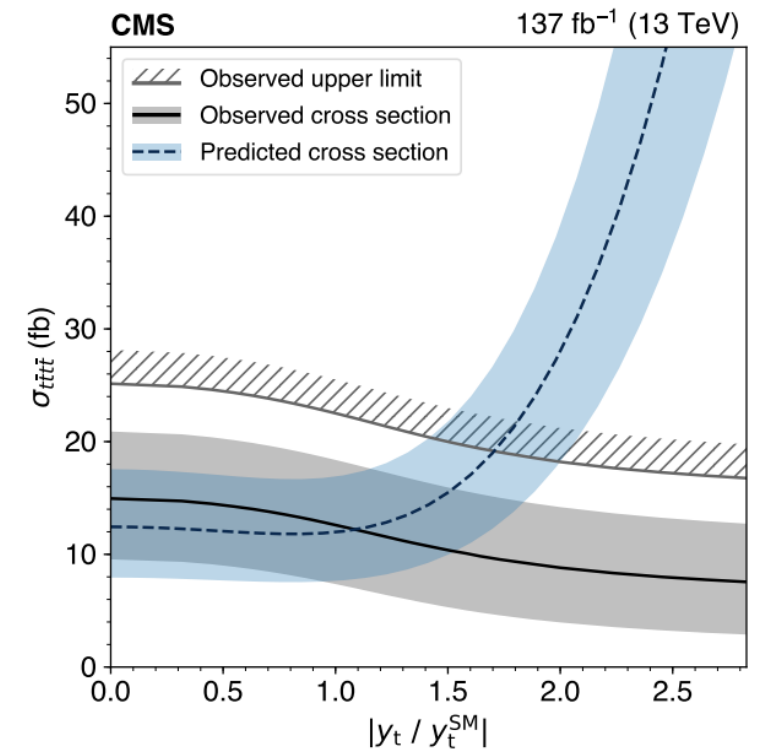
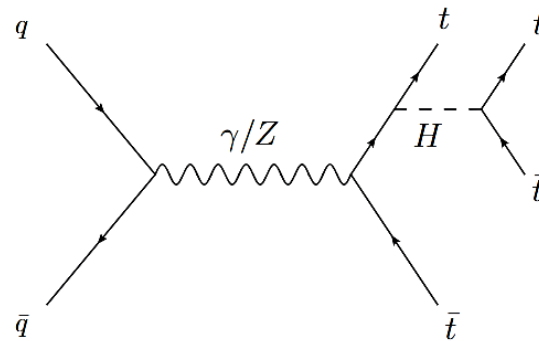
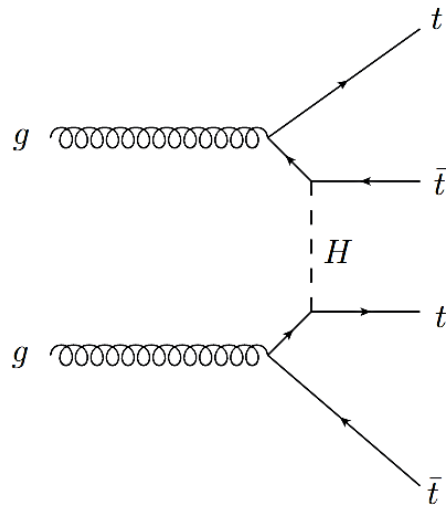
IJCLab Orsay

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- ▶ Sensitive to the Yukawa coupling



[Eur. Phys. J. C (2020) 80, 75]

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- ▶ Sensitive to new physics (gluinos, scalar gluons, heavy scalar bosons, ...)

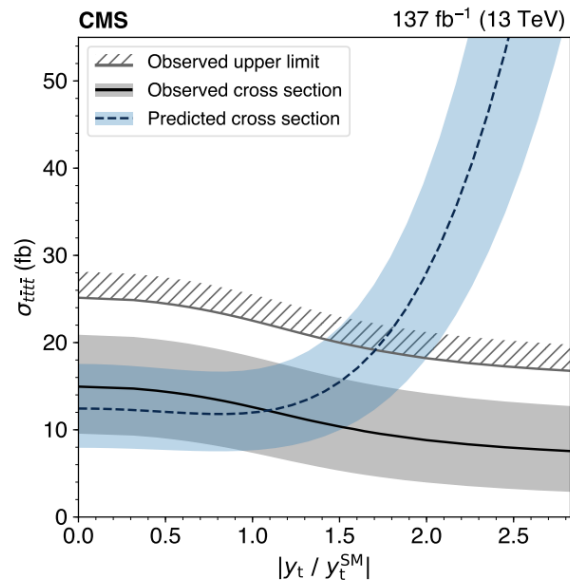
- ▶ Four coloured massive particles in the final state
- ▶ Sensitive to the Yukawa coupling
- ▶ Sensitive to new physics (gluinos, scalar gluons, heavy scalar bosons, ...)
- ▶ Constrains SMEFT coefficients: e.g. four-fermion operator

▶ Measured at the LHC

- ▶ ATLAS [Eur. Phys. J. C (2020) 80, 1085; JHEP 11 (2021), 118; Phys. Rev. D 99, 052009 (2019)]
- ▶ CMS [Eur. Phys. J. C (2020) 80, 75; JHEP 11 (2019), 082]

CMS

[Eur. Phys. J. C (2020) 80, 75]



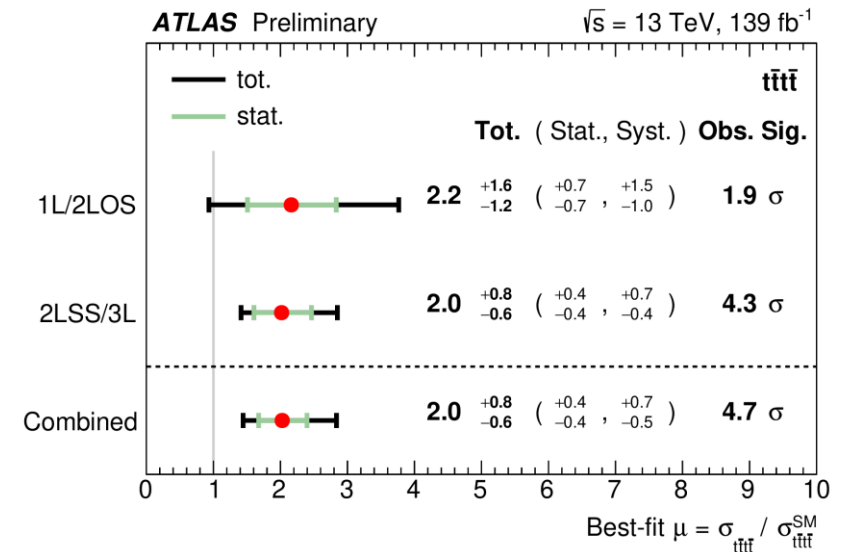
$$\sigma_{t\bar{t}t\bar{t}}^{\text{CMS}} = 12.6^{+5.8}_{-5.2} \text{ fb}$$

$$\sigma_{t\bar{t}t\bar{t}}^{\text{SM, NLO (QCD+EW)}} = 12.0 \pm 2.4 \text{ fb}$$

[Frederix, Pagani, Zaro (2017)]

ATLAS

[JHEP 11 (2021), 118]



$$\sigma_{t\bar{t}t\bar{t}}^{\text{ATLAS}} = 24 \pm 5(\text{stat.})^{+5}_{-4}(\text{syst.}) \text{ fb}$$

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- ▶ Theoretical predictions for total rate: NLO (QCD + EW) available

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GOAL: extend the precision of theoretical predictions beyond NLO for $pp \rightarrow t\bar{t}t\bar{t}$ by means of **resummation** techniques

► Perturbative expansion cross-section

$$\sigma = \sum_n c_n \alpha_s^n = c_0 + c_1 \alpha_s + c_2 \alpha_s^2 + \dots \quad \text{with} \quad c_n = f_n + \sum_{k=0}^{2n-1} d_{nk} L^k, \quad L^k = \left[\frac{\ln^k(1 - M^2/s)}{1 - M^2/s} \right]_+$$

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- ▶ Large logs L can spoil predictive power, $L^k \rightarrow \infty$ for $M^2 \simeq s$

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▶ **Soft Gluon Resummation**

- ▶ Systematic treatment to all orders: resummation

- ▶ Relies on ME and Phase space factorization \rightarrow Mellin space $L := \left[\frac{\ln(1 - M^2/s)}{1 - M^2/s} \right]_+ \rightarrow \tilde{L} := \ln N$

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$$\hat{\sigma}^{\text{res}}(N) \sim \mathcal{F}(\alpha_s) \exp \left[\tilde{L} g_1(\alpha_s \tilde{L}) + g_2(\alpha_s \tilde{L}) + \dots \right]$$

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\ll
 $\alpha_s^n \ln^{2n} N$

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$$\hat{\sigma}^{\text{res}}(N) \sim \mathcal{F}(\alpha_s) \exp \left[\underbrace{\tilde{L} g_1(\alpha_s \tilde{L})}_{\substack{\text{LL} \\ \alpha_s^n \ln^{2n} N}} + \underbrace{g_2(\alpha_s \tilde{L})}_{\substack{\text{NLL} \\ \alpha_s^n \ln^{2n-1} N}} + \dots \right]$$

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$$\hat{\sigma}^{\text{res}}(N) \sim \underbrace{\mathcal{F}(\alpha_s)}_{\text{Finite contributions}} \exp \left[\underbrace{\tilde{L} g_1(\alpha_s \tilde{L})}_{\text{LL } \alpha_s^n \ln^{2n} N} + \underbrace{g_2(\alpha_s \tilde{L})}_{\text{NLL } \alpha_s^n \ln^{2n-1} N} + \dots \right]$$

- ▶ Resummed partonic cross section for $pp \rightarrow t\bar{t}t\bar{t}$ in Mellin space

$$\begin{aligned}\hat{\sigma}_{ij \rightarrow t\bar{t}t\bar{t}}^{\text{res}} &= \text{Tr} \left[\mathbf{H}_{ij \rightarrow t\bar{t}t\bar{t}} \mathbf{S}_{ij \rightarrow t\bar{t}t\bar{t}} \right] \Delta_i \Delta_j \\ &= \text{Tr} \left[\mathbf{H}_{ij \rightarrow t\bar{t}t\bar{t}} \bar{\mathbf{U}}_{ij \rightarrow t\bar{t}t\bar{t}} \tilde{\mathbf{S}}_{ij \rightarrow t\bar{t}t\bar{t}} \mathbf{U}_{ij \rightarrow t\bar{t}t\bar{t}} \right] \Delta_i \Delta_j\end{aligned}$$

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- Resummed partonic cross section for $pp \rightarrow t\bar{t}t\bar{t}$ in Mellin space

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(soft-)collinear enhancements

$$= \text{Tr} \left[\mathbf{H}_{ij \rightarrow t\bar{t}t\bar{t}} \bar{\mathbf{U}}_{ij \rightarrow t\bar{t}t\bar{t}} \tilde{\mathbf{S}}_{ij \rightarrow t\bar{t}t\bar{t}} \mathbf{U}_{ij \rightarrow t\bar{t}t\bar{t}} \right] \Delta_i \Delta_j$$

Hard function

constant contributions as $N \rightarrow \infty$

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Hard function **Soft function**

constant contributions as $N \rightarrow \infty$

soft wide-angle enhancements

$$\mathbf{U}_{ij \rightarrow t\bar{t}t\bar{t}}(N, M^2, \mu_F^2, \mu_R^2) = \text{Pexp} \left[\int_{\mu}^{M/\bar{N}} \frac{dq}{q} \mathbf{\Gamma}_{ij \rightarrow t\bar{t}t\bar{t}}(\alpha_s(q^2)) \right]$$

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- ▶ One-loop SAD matrix needed at NLL

$$\mathbf{U}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(N, M^2, \mu_F^2, \mu_R^2) = \text{Pexp} \left[\int_{\mu}^{M/\bar{N}} \frac{dq}{q} \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s(q^2)) \right]$$

$$\mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s) = \left(\frac{\alpha_s}{\pi} \right) \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(1)} + \left(\frac{\alpha_s}{\pi} \right)^2 \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(2)} + \dots$$

- ▶ $q\bar{q} \rightarrow t\bar{t}\bar{t}\bar{t} : \mathbf{3} \otimes \bar{\mathbf{3}} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{1} \oplus \mathbf{8} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

6-dimensional colour space

- ▶ $gg \rightarrow t\bar{t}\bar{t}\bar{t} : \mathbf{8} \otimes \mathbf{8} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{0} \oplus \mathbf{1} \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

14-dimensional colour space

► One-loop SAD matrix needed at NLL

$$\mathbf{U}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(N, M^2, \mu_F^2, \mu_R^2) = \text{Pexp} \left[\int_{\mu}^{M/\bar{N}} \frac{dq}{q} \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s(q^2)) \right]$$

Diagonal SAD matrix

$$\bar{\mathbf{U}}_R \tilde{\mathbf{S}}_R \mathbf{U}_R = \tilde{\mathbf{S}}_R \exp \left[\frac{2 \text{Re}(\mathbf{\Gamma}_R^{(1)})}{2\pi b_0} \ln(1 - 2\lambda) \right]$$

$$\mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s) = \left(\frac{\alpha_s}{\pi}\right) \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(1)} + \left(\frac{\alpha_s}{\pi}\right)^2 \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(2)} + \dots$$

with $\lambda = \alpha_s b_0 \ln(\bar{N})$

► $q\bar{q} \rightarrow t\bar{t}\bar{t}\bar{t}$: $\mathbf{3} \otimes \bar{\mathbf{3}} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{1} \oplus \mathbf{8} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

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$$C_1^{q\bar{q}} = \frac{1}{\sqrt{N_c^3}} \delta_{c_1 c_3} \delta_{c_2 c_4} \delta_{c_6 c_8}$$

$$C_2^{q\bar{q}} = \frac{1}{T_R \sqrt{N_c(N_c^2 - 1)}} \delta_{c_1 c_3} t_{c_2 c_4}^{a_1} t_{c_6 c_8}^{a_1}$$

$$C_3^{q\bar{q}} = \frac{1}{T_R \sqrt{N_c(N_c^2 - 1)}} t_{c_1 c_3}^{a_1} \delta_{c_2 c_4} t_{c_6 c_8}^{a_1}$$

$$C_4^{q\bar{q}} = \frac{1}{T_R \sqrt{N_c(N_c^2 - 1)}} t_{c_1 c_3}^{a_1} t_{c_2 c_4}^{a_1} \delta_{c_6 c_8}$$

$$C_5^{q\bar{q}} = \frac{\sqrt{N_c}}{T_R^2 \sqrt{2(N_c^4 - 5N_c^2 + 4)}} t_{c_1 c_3}^{a_1} d^{a_1 a_2 b_3} t_{c_2 c_4}^{a_2} t_{c_6 c_8}^{b_3}$$

$$C_6^{q\bar{q}} = \frac{1}{T_R^2 \sqrt{2N_c(N_c^2 - 1)}} t_{c_1 c_3}^{a_1} i f^{a_1 a_2 b_3} t_{c_2 c_4}^{a_2} t_{c_6 c_8}^{b_3},$$



6x6 SAD matrix

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},11}^{(1)} = -C_F(L_{\beta_{34}} + L_{\beta_{56}} + 2)$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},12}^{(1)} = \frac{\sqrt{N_c^2 - 1}}{2N_c} (L_{\beta_{35}} + L_{\beta_{46}} - L_{\beta_{36}} - L_{\beta_{45}})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},13}^{(1)} = \frac{\sqrt{N_c^2 - 1}}{2N_c} (T_{15} - T_{16} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},14}^{(1)} = \frac{\sqrt{N_c^2 - 1}}{2N_c} (T_{13} - T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},15}^{(1)} = 0$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},16}^{(1)} = 0$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},22}^{(1)} = -2C_F + \frac{1}{2N_c} (L_{\beta_{34}} + L_{\beta_{56}}) - \frac{1}{N_c} (L_{\beta_{35}} + L_{\beta_{46}}) - \frac{N_c^2 - 2}{2N_c} (L_{\beta_{36}} + L_{\beta_{45}})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},23}^{(1)} = \frac{1}{2N_c} (T_{13} - T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},24}^{(1)} = \frac{1}{2N_c} (T_{15} - T_{16} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},25}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{2\sqrt{2}N_c} (T_{13} - T_{14} + T_{15} - T_{16} - T_{23} + T_{24} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},26}^{(1)} = \frac{1}{2\sqrt{2}} (-T_{13} - T_{14} + T_{15} + T_{16} + T_{23} + T_{24} - T_{25} - T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},33}^{(1)} = -\frac{N_c^2 - 2}{2N_c} - C_F L_{\beta_{34}} + \frac{1}{2N_c} L_{\beta_{56}} + \frac{N_c^2 - 2}{2N_c} T_{15} + \frac{1}{N_c} T_{16} + \frac{N_c^2 - 2}{2N_c} T_{26} + \frac{1}{N_c} T_{25}$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},34}^{(1)} = \frac{1}{2N_c} (L_{\beta_{35}} + L_{\beta_{46}} - L_{\beta_{36}} - L_{\beta_{45}})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},35}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{2\sqrt{2}N_c} (L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}} + T_{13} - T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},36}^{(1)} = \frac{1}{2\sqrt{2}} (-L_{\beta_{35}} - L_{\beta_{36}} + L_{\beta_{45}} + L_{\beta_{46}} - T_{13} + T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},44}^{(1)} = -C_F + \frac{1}{2N_c} - C_F L_{\beta_{56}} + \frac{1}{2N_c} L_{\beta_{34}} + \frac{N_c^2 - 2}{2N_c} (T_{13} + T_{24}) + \frac{1}{N_c} (T_{14} + T_{23})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},45}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{2\sqrt{2}N_c} (L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}} + T_{15} - T_{16} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},46}^{(1)} = \frac{1}{2\sqrt{2}} (L_{\beta_{35}} - L_{\beta_{36}} + L_{\beta_{45}} - L_{\beta_{46}} + T_{15} - T_{16} + T_{25} - T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},55}^{(1)} = -C_F + \frac{1}{2N_c} + \frac{1}{2N_c} \left(L_{\beta_{34}} - 3L_{\beta_{35}} - 3L_{\beta_{46}} + L_{\beta_{56}} - \frac{N_c^2 - 6}{2} L_{\beta_{36}} - \frac{N_c^2 - 6}{2} L_{\beta_{45}} \right)$$

$$+ \frac{3}{2N_c} (T_{14} + T_{16} + T_{23} + T_{25}) + \frac{N_c^2 - 6}{4N_c} (T_{13} + T_{15} + T_{24} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},56}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{4} (-L_{\beta_{36}} + L_{\beta_{45}} - T_{13} + T_{15} + T_{24} - T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},66}^{(1)} = -C_F + \frac{1}{2N_c} - \frac{N_c^2 - 2}{4N_c} (L_{\beta_{36}} + L_{\beta_{45}}) + \frac{1}{2N_c} (L_{\beta_{34}} - L_{\beta_{35}} - L_{\beta_{46}} + L_{\beta_{56}})$$

$$+ \frac{N_c^2 - 2}{4N_c} (T_{13} + T_{15} + T_{24} + T_{26}) + \frac{1}{2N_c} (T_{14} + T_{16} + T_{23} + T_{25})$$

[Keppeler, Sjodahl (2012)]

► $q\bar{q} \rightarrow t\bar{t}t\bar{t}$: $\mathbf{3} \otimes \bar{\mathbf{3}} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{1} \oplus \mathbf{8} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

$$C_1^{q\bar{q}} = \frac{1}{\sqrt{N_c^3}} \delta_{c_1 c_3} \delta_{c_2 c_4} \delta_{c_6 c_8}$$

$$C_2^{q\bar{q}} = \frac{1}{T_R \sqrt{N_c(N_c^2 - 1)}} \delta_{c_1 c_3} t_{c_2 c_4}^{a_1} t_{c_6 c_8}^{a_1}$$

$$C_3^{q\bar{q}} = \frac{1}{T_R \sqrt{N_c(N_c^2 - 1)}} t_{c_1 c_3}^{a_1} \delta_{c_2 c_4} t_{c_6 c_8}^{a_1}$$

$$C_4^{q\bar{q}} = \frac{1}{T_R \sqrt{N_c(N_c^2 - 1)}} t_{c_1 c_3}^{a_1} t_{c_2 c_4}^{a_1} \delta_{c_6 c_8}$$

$$C_5^{q\bar{q}} = \frac{\sqrt{N_c}}{T_R^2 \sqrt{2(N_c^4 - 5N_c^2 + 4)}} t_{c_1 c_3}^{a_1} d^{a_1 a_2 b_3} t_{c_2 c_4}^{a_2} t_{c_6 c_8}^{b_3}$$

$$C_6^{q\bar{q}} = \frac{1}{T_R^2 \sqrt{2N_c(N_c^2 - 1)}} t_{c_1 c_3}^{a_1} i f^{a_1 a_2 b_3} t_{c_2 c_4}^{a_2} t_{c_6 c_8}^{b_3}$$



6x6 SAD matrix

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},11}^{(1)} = -C_F(L_{\beta_{34}} + L_{\beta_{56}} + 2)$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},12}^{(1)} = \frac{\sqrt{N_c^2 - 1}}{2N_c} (L_{\beta_{35}} + L_{\beta_{46}} - L_{\beta_{36}} - L_{\beta_{45}})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},13}^{(1)} = \frac{\sqrt{N_c^2 - 1}}{2N_c} (T_{15} - T_{16} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},14}^{(1)} = \frac{\sqrt{N_c^2 - 1}}{2N_c} (T_{13} - T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},15}^{(1)} = 0$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},16}^{(1)} = 0$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},22}^{(1)} = -2C_F + \frac{1}{2N_c} (L_{\beta_{34}} + L_{\beta_{56}}) - \frac{1}{2} (L_{\beta_{35}} + L_{\beta_{46}} - L_{\beta_{36}} - L_{\beta_{45}}) - \frac{N_c^2}{2N_c} (L_{\beta_{36}} + L_{\beta_{45}})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},23}^{(1)} = \frac{1}{2N_c} (T_{13} - T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},24}^{(1)} = \frac{1}{2N_c} (T_{15} - T_{16} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},25}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{2\sqrt{2}N_c} (T_{13} - T_{14} + T_{15} - T_{16} - T_{23} + T_{24} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},26}^{(1)} = \frac{1}{2\sqrt{2}} (-T_{13} - T_{14} + T_{15} + T_{16} + T_{23} + T_{24} - T_{25} - T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},33}^{(1)} = -\frac{N_c^2 - 2}{2N_c} - C_F L_{\beta_{34}} + \frac{1}{2N_c} L_{\beta_{56}} + \frac{N_c^2 - 2}{2N_c} T_{15} + \frac{1}{N_c} T_{16} + \frac{N_c^2 - 2}{2N_c} T_{26} + \frac{1}{N_c} T_{25}$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},34}^{(1)} = \frac{1}{2N_c} (L_{\beta_{35}} + L_{\beta_{46}} - L_{\beta_{36}} - L_{\beta_{45}})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},35}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{2\sqrt{2}N_c} (L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}} + T_{13} - T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},36}^{(1)} = \frac{1}{2\sqrt{2}} (-L_{\beta_{35}} - L_{\beta_{36}} + L_{\beta_{45}} + L_{\beta_{46}} - T_{13} + T_{14} - T_{23} + T_{24})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},44}^{(1)} = -C_F + \frac{1}{2N_c} - C_F L_{\beta_{56}} + \frac{1}{2N_c} L_{\beta_{34}} + \frac{N_c^2 - 2}{2N_c} (T_{13} + T_{24}) + \frac{1}{N_c} (T_{14} + T_{23})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},45}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{2\sqrt{2}N_c} (L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}} + T_{15} - T_{16} - T_{25} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},46}^{(1)} = \frac{1}{2\sqrt{2}} (L_{\beta_{35}} - L_{\beta_{36}} + L_{\beta_{45}} - L_{\beta_{46}} + T_{15} - T_{16} + T_{25} - T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},55}^{(1)} = -C_F + \frac{1}{2N_c} + \frac{1}{2N_c} (L_{\beta_{34}} - 3L_{\beta_{35}} - 3L_{\beta_{46}} + L_{\beta_{56}} - \frac{N_c^2 - 6}{2} L_{\beta_{36}} - \frac{N_c^2 - 6}{2} L_{\beta_{45}})$$

$$+ \frac{3}{2N_c} (T_{14} + T_{16} + T_{23} + T_{25}) + \frac{N_c^2 - 6}{4N_c} (T_{13} + T_{15} + T_{24} + T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},56}^{(1)} = \frac{\sqrt{N_c^2 - 4}}{4} (-L_{\beta_{36}} + L_{\beta_{45}} - T_{13} + T_{15} + T_{24} - T_{26})$$

$$\bar{\Gamma}_{q\bar{q} \rightarrow 4\text{top},66}^{(1)} = -C_F + \frac{1}{2N_c} - \frac{N_c^2 - 2}{4N_c} (L_{\beta_{36}} + L_{\beta_{45}}) + \frac{1}{2N_c} (L_{\beta_{34}} - L_{\beta_{35}} - L_{\beta_{46}} + L_{\beta_{56}})$$

$$+ \frac{N_c^2 - 2}{4N_c} (T_{13} + T_{15} + T_{24} + T_{26}) + \frac{1}{2N_c} (T_{14} + T_{16} + T_{23} + T_{25})$$

**DIAGONAL IN THE
ABSOLUTE MASS
THRESHOLD LIMIT**

[Keppeler, Sjodahl (2012)]

► $gg \rightarrow t\bar{t}t\bar{t} : 8 \otimes 8 = 3 \otimes \bar{3} \otimes 3 \otimes \bar{3} \rightarrow 0 \oplus 1 \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27 = \boxed{0 \oplus (2 \times 1) \oplus (2 \times 8) \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27}$

$$c_1^{gg} = \frac{1}{T_R} \frac{1}{N_c^2 - 1} t_{c_2 c_4}^{a_1} t_{c_6 c_8}^{a_2},$$

$$c_3^{gg} = \frac{1}{T_R} \frac{1}{\sqrt{2(N_c^4 - 5N_c^2 + 4)}} \delta_{c_2 c_4} d_{a_1 a_2 b_1} t_{c_6 c_8}^{b_1},$$

$$c_5^{gg} = \frac{1}{T_R} \frac{1}{\sqrt{2(N_c^4 - 5N_c^2 + 4)}} t_{c_2 c_4}^{b_1} d_{b_1 a_1 a_2} \delta_{c_6 c_8},$$

$$c_7^{gg} = \frac{1}{T_R^2} \frac{1}{2\sqrt{N_c^4 - 5N_c^2 + 4}} t_{c_2 c_4}^{b_1} d_{b_1 a_1 b_2} i f_{b_2 a_2 b_3} t_{c_6 c_8}^{b_3},$$

$$c_9^{gg} = \frac{1}{T_R^2} \frac{1}{2\sqrt{N_c^4 - 5N_c^2 + 4}} t_{c_2 c_4}^{b_1} i f_{b_1 a_1 b_2} d_{b_2 a_2 b_3} t_{c_6 c_8}^{b_3},$$

$$c_{11}^{gg} = \frac{1}{T_R} \frac{2}{\sqrt{N_c^4 - 5N_c^2 + 4}} t_{c_2 c_4}^{b_1} \mathbf{P}_{a_1 b_1 a_2 b_2}^{10} t_{c_6 c_8}^{b_2},$$

$$c_{13}^{gg} = \frac{1}{T_R} \frac{2}{N_c \sqrt{N_c^2 + 2N_c - 3}} t_{c_2 c_4}^{b_1} \mathbf{P}_{a_1 b_1 a_2 b_2}^{27} t_{c_6 c_8}^{b_2},$$

$$c_2^{gg} = \frac{1}{N_c \sqrt{N_c^2 - 1}} \delta_{a_1 a_2} \delta_{c_2 c_4} \delta_{c_6 c_8},$$

$$c_4^{gg} = \frac{1}{T_R} \frac{1}{N_c \sqrt{2(N_c^2 - 1)}} \delta_{c_2 c_4} i f_{a_1 a_2 b_1} t_{c_6 c_8}^{b_1},$$

$$c_6^{gg} = \frac{1}{T_R^2} \frac{N_c}{2(N_c^2 - 4)\sqrt{N_c^2 - 1}} t_{c_2 c_4}^{b_1} d_{b_1 a_1 b_2} d_{b_2 a_2 b_3} t_{c_6 c_8}^{b_3},$$

$$c_8^{gg} = \frac{1}{T_R} \frac{1}{N_c \sqrt{2(N_c^2 - 1)}} t_{c_2 c_4}^{b_1} i f_{b_1 a_1 a_2} \delta_{c_6 c_8},$$

$$c_{10}^{gg} = \frac{1}{T_R^2} \frac{1}{2N_c \sqrt{N_c^2 - 1}} t_{c_2 c_4}^{b_1} i f_{b_1 a_1 b_2} i f_{b_2 a_2 b_3} t_{c_6 c_8}^{b_3},$$

$$c_{12}^{gg} = \frac{1}{T_R} \frac{2}{\sqrt{N_c^4 - 5N_c^2 + 4}} t_{c_2 c_4}^{b_1} \mathbf{P}_{a_1 b_1 a_2 b_2}^{\bar{10}} t_{c_6 c_8}^{b_2},$$

$$c_{14}^{gg} = \frac{1}{T_R} \frac{2}{N_c \sqrt{N_c^2 - 2N_c - 3}} t_{c_2 c_4}^{b_1} \mathbf{P}_{a_1 b_1 a_2 b_2}^0 t_{c_6 c_8}^{b_2}.$$

[Keppeler, Sjo Dahl (2012)]

► $gg \rightarrow t\bar{t}t\bar{t}$: $8 \otimes 8 = 3 \otimes \bar{3} \otimes 3 \otimes \bar{3} \rightarrow 0 \oplus 1 \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27 = \boxed{0 \oplus (2 \times 1) \oplus (2 \times 8) \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27}$

14x14 SAD matrix

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},11}^{(1)} = \frac{1}{N_c} + \frac{1}{2N_c}(L_{\beta_{34}} + L_{\beta_{34}}) + \frac{N_c}{2}(T_{13} + T_{14} + T_{25} + T_{26})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},12}^{(1)} = \frac{1}{2N_c\sqrt{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},13}^{(1)} = \frac{1}{4N_c}\sqrt{\frac{2(N_c^2-4)}{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},14}^{(1)} = \frac{1}{2\sqrt{2(N_c^2-1)}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 2T_{23} - 2T_{24})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},15}^{(1)} = \frac{1}{4N_c}\sqrt{\frac{2(N_c^2-4)}{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},16}^{(1)} = \frac{N_c^2-4}{4N_c\sqrt{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},17}^{(1)} = \frac{1}{4}\sqrt{\frac{N_c^2-4}{N_c^2-1}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 2T_{23} - 2T_{24})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},18}^{(1)} = \frac{1}{2\sqrt{2(N_c^2-1)}}(-L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}} - 2T_{15} + 2T_{16})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},19}^{(1)} = -\frac{1}{4}\sqrt{\frac{N_c^2-4}{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} + L_{\beta_{45}} - L_{\beta_{46}} + 2T_{15} - 2T_{16})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},110}^{(1)} = -\frac{N_c}{4\sqrt{N_c^2-1}}(4 + L_{\beta_{35}} + L_{\beta_{36}} + L_{\beta_{45}} + L_{\beta_{46}} + 2T_{15} + 2T_{16} + 2T_{23} + 2T_{24})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},111}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},112}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},113}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},114}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},22}^{(1)} = -2C_F - C_F(L_{\beta_{34}} + L_{\beta_{56}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},23}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},24}^{(1)} = \frac{1}{\sqrt{2}}(T_{15} - T_{16} - T_{25} + T_{26})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},25}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},26}^{(1)} = \frac{1}{2N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},27}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},28}^{(1)} = \frac{1}{\sqrt{2}}(T_{13} - T_{14} - T_{23} + T_{24})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},29}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},210}^{(1)} = \frac{1}{2N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},211}^{(1)} = \frac{\sqrt{N_c^2-4}}{4N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},212}^{(1)} = \frac{\sqrt{N_c^2-4}}{4N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},213}^{(1)} = \frac{1}{4}\sqrt{\frac{N_c+3}{N_c+1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},214}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},33}^{(1)} = -\frac{N_c^2-2}{2N_c} - C_F L_{\beta_{34}} + \frac{1}{2N_c}L_{\beta_{56}} + \frac{N_c}{4}(T_{15} + T_{16} + T_{25} + T_{26})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},34}^{(1)} = \frac{\sqrt{N_c^2-4}}{4}(T_{15} - T_{16} - T_{25} + T_{26})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},35}^{(1)} = \frac{1}{2N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},36}^{(1)} = \frac{N_c^2-12}{4N_c\sqrt{2(N_c^2-4)}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},37}^{(1)} = \frac{1}{4\sqrt{2}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 2T_{23} - 2T_{24})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},38}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},39}^{(1)} = \frac{1}{4\sqrt{2}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 4T_{13} - 4T_{14} - 2T_{23} + 2T_{24})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},310}^{(1)} = \frac{\sqrt{2(N_c^2-4)}}{8N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},311}^{(1)} = \frac{N_c-2}{4N_c\sqrt{2}}(L_{\beta_{35}} - L_{\beta_{45}}) + \frac{N_c+2}{4N_c\sqrt{2}}(L_{\beta_{36}} - L_{\beta_{46}}) + \frac{1}{2\sqrt{2}}(T_{23} - T_{24})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},312}^{(1)} = \frac{N_c-2}{4N_c\sqrt{2}}(L_{\beta_{46}} - L_{\beta_{36}}) + \frac{N_c+2}{4N_c\sqrt{2}}(L_{\beta_{45}} - L_{\beta_{35}}) + \frac{1}{2\sqrt{2}}(T_{24} - T_{23})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},313}^{(1)} = \frac{1}{4\sqrt{2}}\sqrt{\frac{(N_c-2)(N_c+3)}{(N_c+1)(N_c+2)}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}})$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},314}^{(1)} = 0$$

$$\tilde{\Gamma}_{gg \rightarrow 4\text{top},44}^{(1)} = -\frac{N_c^2-2}{2N_c} - C_F L_{\beta_{34}} + \frac{1}{2N_c}L_{\beta_{56}} + \frac{N_c}{4}(T_{15} + T_{16} + T_{25} + T_{26})$$

+ 65 more components!

► $gg \rightarrow t\bar{t}t\bar{t}$: $8 \otimes 8 = 3 \otimes \bar{3} \otimes 3 \otimes \bar{3} \rightarrow 0 \oplus 1 \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27 = \boxed{0 \oplus (2 \times 1) \oplus (2 \times 8) \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27}$

14x14 SAD matrix

$$\begin{aligned} \tilde{\Gamma}_{gg \rightarrow 4\text{top},11}^{(1)} &= \frac{1}{N_c} + \frac{1}{2N_c}(L_{\beta_{34}} + L_{\beta_{34}}) + \frac{N_c}{2}(T_{13} + T_{14} + T_{25} + T_{26}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},12}^{(1)} &= \frac{1}{2N_c\sqrt{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},13}^{(1)} &= \frac{1}{4N_c}\sqrt{\frac{2(N_c^2-4)}{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},14}^{(1)} &= \frac{1}{2\sqrt{2(N_c^2-1)}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 2T_{23} - 2T_{24}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},15}^{(1)} &= \frac{1}{4N_c}\sqrt{\frac{2(N_c^2-4)}{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},16}^{(1)} &= \frac{N_c^2-4}{4N_c\sqrt{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},17}^{(1)} &= \frac{1}{4}\sqrt{\frac{N_c^2-4}{N_c^2-1}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 2T_{23} - 2T_{24}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},18}^{(1)} &= \frac{1}{2\sqrt{2(N_c^2-1)}}(-L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}} - 2T_{15} + 2T_{16}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},19}^{(1)} &= -\frac{1}{4}\sqrt{\frac{N_c^2-4}{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} + L_{\beta_{45}} - L_{\beta_{46}} + 2T_{15} - 2T_{16}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},110}^{(1)} &= -\frac{N_c}{4\sqrt{N_c^2-1}}(4 + L_{\beta_{35}} + L_{\beta_{36}} + L_{\beta_{45}} + L_{\beta_{46}} \\ &\quad + 2T_{15} + 2T_{16} + 2T_{23} + 2T_{24}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},111}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},112}^{(1)} &= 0 \end{aligned}$$

$$\begin{aligned} \tilde{\Gamma}_{gg \rightarrow 4\text{top},113}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},114}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},22}^{(1)} &= -2C_F - C_F(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},23}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},24}^{(1)} &= \frac{1}{\sqrt{2}}(T_{15} - T_{16} - T_{25} + T_{26}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},25}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},26}^{(1)} &= \frac{1}{2N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},27}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},28}^{(1)} &= \frac{1}{\sqrt{2}}(T_{15} - T_{16} - T_{25} + T_{26}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},29}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},210}^{(1)} &= \frac{1}{2\sqrt{2}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},211}^{(1)} &= \frac{1}{4N_c}\sqrt{\frac{N_c^2-4}{N_c^2-1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},212}^{(1)} &= \frac{\sqrt{N_c^2-4}}{4N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},213}^{(1)} &= \frac{1}{4}\sqrt{\frac{N_c+3}{N_c+1}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},214}^{(1)} &= 0 \end{aligned}$$

**NOT
DIAGONAL IN THE
ABSOLUTE MASS
THRESHOLD LIMIT**

$$\begin{aligned} \tilde{\Gamma}_{gg \rightarrow 4\text{top},33}^{(1)} &= -\frac{N_c^2-2}{2N_c} - C_F L_{\beta_{34}} + \frac{1}{2N_c}L_{\beta_{56}} + \frac{N_c}{4}(T_{15} + T_{16} + T_{25} + T_{26}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},34}^{(1)} &= \frac{\sqrt{N_c^2-4}}{4}(T_{15} - T_{16} - T_{25} + T_{26}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},35}^{(1)} &= \frac{1}{2N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},36}^{(1)} &= \frac{N_c^2-12}{4N_c\sqrt{2(N_c^2-4)}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},37}^{(1)} &= \frac{1}{4\sqrt{2}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 2T_{23} - 2T_{24}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},38}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},39}^{(1)} &= \frac{1}{4\sqrt{2}}(L_{\beta_{35}} + L_{\beta_{36}} - L_{\beta_{45}} - L_{\beta_{46}} + 4T_{13} - 4T_{14} - 2T_{23} + 2T_{24}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},310}^{(1)} &= \frac{\sqrt{2(N_c^2-4)}}{8N_c}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},311}^{(1)} &= \frac{N_c-2}{4N_c\sqrt{2}}(L_{\beta_{35}} - L_{\beta_{45}}) + \frac{N_c+2}{4N_c\sqrt{2}}(L_{\beta_{36}} - L_{\beta_{46}}) + \frac{1}{2\sqrt{2}}(T_{23} - T_{24}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},312}^{(1)} &= \frac{N_c-2}{4N_c\sqrt{2}}(L_{\beta_{46}} - L_{\beta_{36}}) + \frac{N_c+2}{4N_c\sqrt{2}}(L_{\beta_{45}} - L_{\beta_{35}}) + \frac{1}{2\sqrt{2}}(T_{24} - T_{23}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},313}^{(1)} &= \frac{1}{4\sqrt{2}}\sqrt{\frac{(N_c-2)(N_c+3)}{(N_c+1)(N_c+2)}}(L_{\beta_{35}} - L_{\beta_{36}} - L_{\beta_{45}} + L_{\beta_{46}}) \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},314}^{(1)} &= 0 \\ \tilde{\Gamma}_{gg \rightarrow 4\text{top},44}^{(1)} &= -\frac{N_c^2-2}{2N_c} - C_F L_{\beta_{34}} + \frac{1}{2N_c}L_{\beta_{56}} + \frac{N_c}{4}(T_{15} + T_{16} + T_{25} + T_{26}) \end{aligned}$$

**+ 65 more
components!**

► $gg \rightarrow t\bar{t}t\bar{t} : 8 \otimes 8 = 3 \otimes \bar{3} \otimes 3 \otimes \bar{3} \rightarrow 0 \oplus 1 \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27 = \boxed{0 \oplus (2 \times 1) \oplus (2 \times 8) \oplus 8_S \oplus 8_A \oplus 10 \oplus \bar{10} \oplus 27}$

$$\bar{c}_1^{gg} = \frac{3\sqrt{3}}{8}c_1^{gg} + \frac{3}{10}\sqrt{\frac{3}{2}}c_6^{gg} - \frac{1}{2}\sqrt{\frac{3}{2}}c_{10}^{gg} - \frac{1}{4}\sqrt{\frac{3}{10}}c_{11}^{gg} - \frac{1}{4}\sqrt{\frac{3}{10}}c_{12}^{gg} + \frac{7}{40}c_{13}^{gg},$$

$$\bar{c}_2^{gg} = -\frac{\sqrt{5}}{4}c_1^{gg} + \sqrt{\frac{2}{5}}c_6^{gg} - \frac{1}{2\sqrt{2}}c_{11}^{gg} - \frac{1}{2\sqrt{2}}c_{12}^{gg} + \frac{1}{4}\sqrt{\frac{3}{5}}c_{13}^{gg},$$

$$\bar{c}_3^{gg} = -\frac{1}{\sqrt{2}}c_7^{gg} + \frac{1}{\sqrt{2}}c_9^{gg},$$

Solution:

$$\bar{c}_5^{gg} = -\frac{1}{2\sqrt{2}}c_1^{gg} - \frac{1}{2}c_6^{gg} - \frac{1}{2}c_{10}^{gg} + \frac{1}{2}\sqrt{\frac{3}{2}}c_{13}^{gg},$$

$$\bar{c}_6^{gg} = -\frac{1}{2}\sqrt{\frac{5}{14}}c_1^{gg} + \frac{3}{2\sqrt{35}}c_6^{gg} - \frac{1}{2}\sqrt{\frac{5}{7}}c_{10}^{gg} + \frac{2}{\sqrt{7}}c_{12}^{gg} - \frac{3}{2}\sqrt{\frac{3}{70}}c_{13}^{gg},$$

$$\bar{c}_7^{gg} = -\frac{1}{2\sqrt{7}}c_1^{gg} + \frac{3}{5\sqrt{14}}c_6^{gg} - \frac{1}{\sqrt{14}}c_{10}^{gg} + \sqrt{\frac{7}{10}}c_{11}^{gg} - \frac{3}{\sqrt{70}}c_{12}^{gg} - \frac{3}{10}\sqrt{\frac{3}{7}}c_{13}^{gg},$$

$$\bar{c}_8^{gg} = \frac{1}{\sqrt{2}}c_7^{gg} + \frac{1}{\sqrt{2}}c_9^{gg},$$

$$\bar{c}_{13}^{gg} = \frac{1}{8}c_1^{gg} + \frac{1}{2\sqrt{2}}c_6^{gg} + \frac{1}{2\sqrt{2}}c_{10}^{gg} + \frac{1}{4}\sqrt{\frac{5}{2}}c_{11}^{gg} + \frac{1}{4}\sqrt{\frac{5}{2}}c_{12}^{gg} + \frac{3\sqrt{3}}{8}c_{13}^{gg},$$

$$\bar{c}_4^{gg} = c_{14}^{gg}, \quad \bar{c}_9^{gg} = c_8^{gg}, \quad \bar{c}_{10}^{gg} = c_5^{gg}, \quad \bar{c}_{11}^{gg} = c_4^{gg}, \quad \bar{c}_{12}^{gg} = c_3^{gg}, \quad \bar{c}_{14}^{gg} = c_2^{gg}.$$

► One-loop SAD matrix needed at NLL

Diagonal SAD matrix

$$\mathbf{U}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(N, M^2, \mu_F^2, \mu_R^2) = \text{Pexp} \left[\int_{\mu}^{M/\bar{N}} \frac{dq}{q} \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s(q^2)) \right] \longrightarrow \bar{\mathbf{U}}_R \tilde{\mathbf{S}}_R \mathbf{U}_R = \tilde{\mathbf{S}}_R \exp \left[\frac{2 \text{Re}(\mathbf{\Gamma}_R^{(1)})}{2\pi b_0} \ln(1 - 2\lambda) \right]$$

$$\mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s) = \left(\frac{\alpha_s}{\pi}\right) \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(1)} + \left(\frac{\alpha_s}{\pi}\right)^2 \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(2)} + \dots$$

with $\lambda = \alpha_s b_0 \ln(\bar{N})$

► $q\bar{q} \rightarrow t\bar{t}\bar{t}\bar{t} : \mathbf{3} \otimes \bar{\mathbf{3}} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{1} \oplus \mathbf{8} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

► $gg \rightarrow t\bar{t}\bar{t}\bar{t} : \mathbf{8} \otimes \mathbf{8} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{0} \oplus \mathbf{1} \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

► One-loop SAD matrix needed at NLL

Diagonal SAD matrix

$$\mathbf{U}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(N, M^2, \mu_F^2, \mu_R^2) = \text{Pexp} \left[\int_{\mu}^{M/\bar{N}} \frac{dq}{q} \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s(q^2)) \right] \longrightarrow \bar{\mathbf{U}}_R \tilde{\mathbf{S}}_R \mathbf{U}_R = \tilde{\mathbf{S}}_R \exp \left[\frac{2 \text{Re}(\mathbf{\Gamma}_R^{(1)})}{2\pi b_0} \ln(1 - 2\lambda) \right]$$

$$\mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}(\alpha_s) = \left(\frac{\alpha_s}{\pi} \right) \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(1)} + \left(\frac{\alpha_s}{\pi} \right)^2 \mathbf{\Gamma}_{ij \rightarrow t\bar{t}\bar{t}\bar{t}}^{(2)} + \dots$$

with $\lambda = \alpha_s b_0 \ln(\bar{N})$

► $q\bar{q} \rightarrow t\bar{t}\bar{t}\bar{t} : \mathbf{3} \otimes \bar{\mathbf{3}} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{1} \oplus \mathbf{8} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

$$2\text{Re} \left[\mathbf{\Gamma}_{R, q\bar{q} \rightarrow t\bar{t}\bar{t}\bar{t}}^{(1)} \right] = \text{diag} \left(0, 0, -3, -3, -3, -3 \right)$$

► $gg \rightarrow t\bar{t}\bar{t}\bar{t} : \mathbf{8} \otimes \mathbf{8} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{0} \oplus \mathbf{1} \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

$$2\text{Re} \left[\mathbf{\Gamma}_{R, gg \rightarrow t\bar{t}\bar{t}\bar{t}}^{(1)} \right] = \text{diag} \left(-8, -6, -6, -4, -3, -3, -3, -3, -3, -3, -3, -3, 0, 0 \right)$$

► Quadratic Casimir Invariants $N_C = 3$

$$C_2(\mathbf{1}) = 0$$

$$C_2(\mathbf{8}_{(S/A)}) = 3$$

$$C_2(\mathbf{10}, \bar{\mathbf{10}}) = 6$$

$$C_2(\mathbf{27}) = 8$$

$$C_2(\mathbf{0}) = 4$$

► $q\bar{q} \rightarrow t\bar{t}t\bar{t}$: $\mathbf{3} \otimes \bar{\mathbf{3}} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{1} \oplus \mathbf{8} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

$$2\text{Re} \left[\Gamma_{R, q\bar{q} \rightarrow t\bar{t}t\bar{t}}^{(1)} \right] = \text{diag} \left(0, 0, -3, -3, -3, -3 \right)$$

► $gg \rightarrow t\bar{t}t\bar{t}$: $\mathbf{8} \otimes \mathbf{8} = \mathbf{3} \otimes \bar{\mathbf{3}} \otimes \mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow \mathbf{0} \oplus \mathbf{1} \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27} = \mathbf{0} \oplus (2 \times \mathbf{1}) \oplus (2 \times \mathbf{8}) \oplus \mathbf{8}_S \oplus \mathbf{8}_A \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}$

$$2\text{Re} \left[\Gamma_{R, gg \rightarrow t\bar{t}t\bar{t}}^{(1)} \right] = \text{diag} \left(-8, -6, -6, -4, -3, -3, -3, -3, -3, -3, -3, -3, 0, 0 \right)$$

$$\mathbf{H} = \mathbf{H}^{(0)} + \frac{\alpha_s(\mu_R)}{\pi} \mathbf{H}^{(1)} + \dots$$

$$\tilde{\mathbf{S}} = \tilde{\mathbf{S}}^{(0)} + \frac{\alpha_s(\mu_R)}{\pi} \tilde{\mathbf{S}}^{(1)} + \dots$$

- ▶ **NLL** accuracy: exponential functions at NLL together with

$$\text{Tr} [\mathbf{H}\tilde{\mathbf{S}}] = \text{Tr} [\mathbf{H}^{(0)}\tilde{\mathbf{S}}^{(0)}]$$

- ▶ **NLL'** accuracy: exponential functions at NLL together with

$$\text{Tr} [\mathbf{H}\tilde{\mathbf{S}}] = \text{Tr} \left[\mathbf{H}^{(0)}\tilde{\mathbf{S}}^{(0)} + \frac{\alpha_s(\mu_R)}{\pi} \mathbf{H}^{(1)}\tilde{\mathbf{S}}^{(0)} + \frac{\alpha_s(\mu_R)}{\pi} \mathbf{H}^{(0)}\tilde{\mathbf{S}}^{(1)} \right]$$

$$\sigma_{t\bar{t}t\bar{t}}^{\text{NLL}(\prime)}(\tau) = \int_{\mathcal{C}} \frac{dN}{2\pi i} \tau^{-N} f_i(N+1, \mu_F^2) f_j(N+1, \mu_F^2) \hat{\sigma}_{ij \rightarrow t\bar{t}t\bar{t}}^{\text{res}}(N)$$

- Combination fixed-order + resummation → **Matching**

$$\sigma_{t\bar{t}t\bar{t}}^{\text{NLO+NLL}(\prime)}(\tau) = \boxed{\sigma_{t\bar{t}t\bar{t}}^{\text{NLO}}(\tau)} + \int_{\mathcal{C}} \frac{dN}{2\pi i} \tau^{-N} f_i(N+1, \mu_F^2) f_j(N+1, \mu_F^2) \times \left[\frac{\hat{\sigma}_{ij \rightarrow t\bar{t}t\bar{t}}^{\text{res}}(N) - \hat{\sigma}_{ij \rightarrow t\bar{t}t\bar{t}}^{\text{res}}(N)|_{\text{NLO}}}{\text{avoid double counting with NLO!}} \right]$$

QCD-only NLO and
QCD + EW NLO

▶ $\sqrt{S} = 13 \text{ TeV}$

▶ $\mu_R = \mu_F$

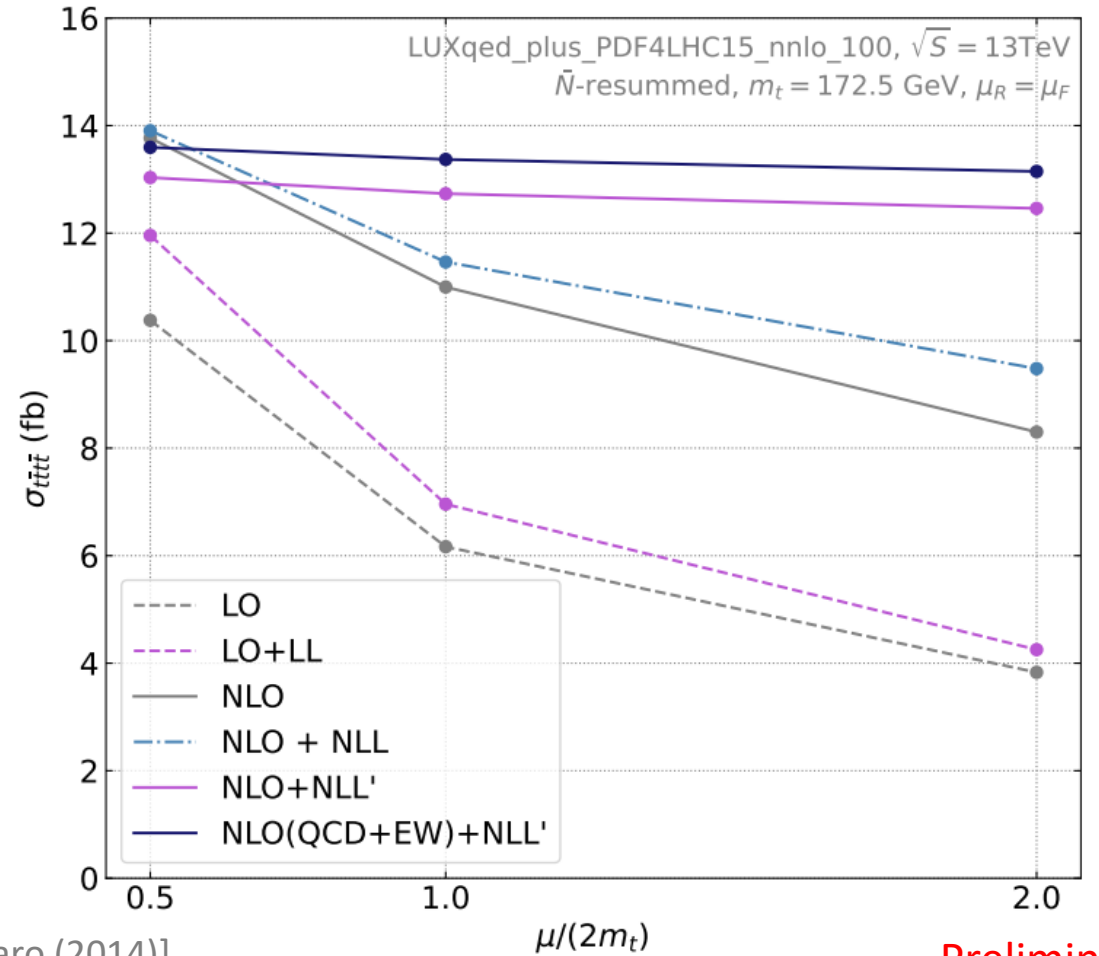
▶ Big impact of resummation

NLO (QCD+EW): EW corrections up to $\mathcal{O}(\alpha^2)$

Fixed-order results obtained with

[Frederix, Frixione, Hirschi, Pagani, Shao, Zaro (2018)]

[Alwall, Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Shao, Stelzer, Torrielli, Zaro (2014)]



Preliminary

► 7-point scale variation

	$\sigma_{t\bar{t}t\bar{t}}$ [fb]	K -factor
NLO	$11.00(2)^{+25.2\%}_{-24.5\%}$	1.04
NLO+NLL	$11.46(2)^{+21.3\%}_{-17.7\%}$	
NLO+NLL'	$12.73(2)^{+4.1\%}_{-11.8\%}$	1.16
NLO (QCD+EW)	$11.64(2)^{+23.2\%}_{-22.8\%}$	1.15
NLO (QCD+EW)+NLL'	$13.37(2)^{+3.6\%}_{-11.4\%}$	

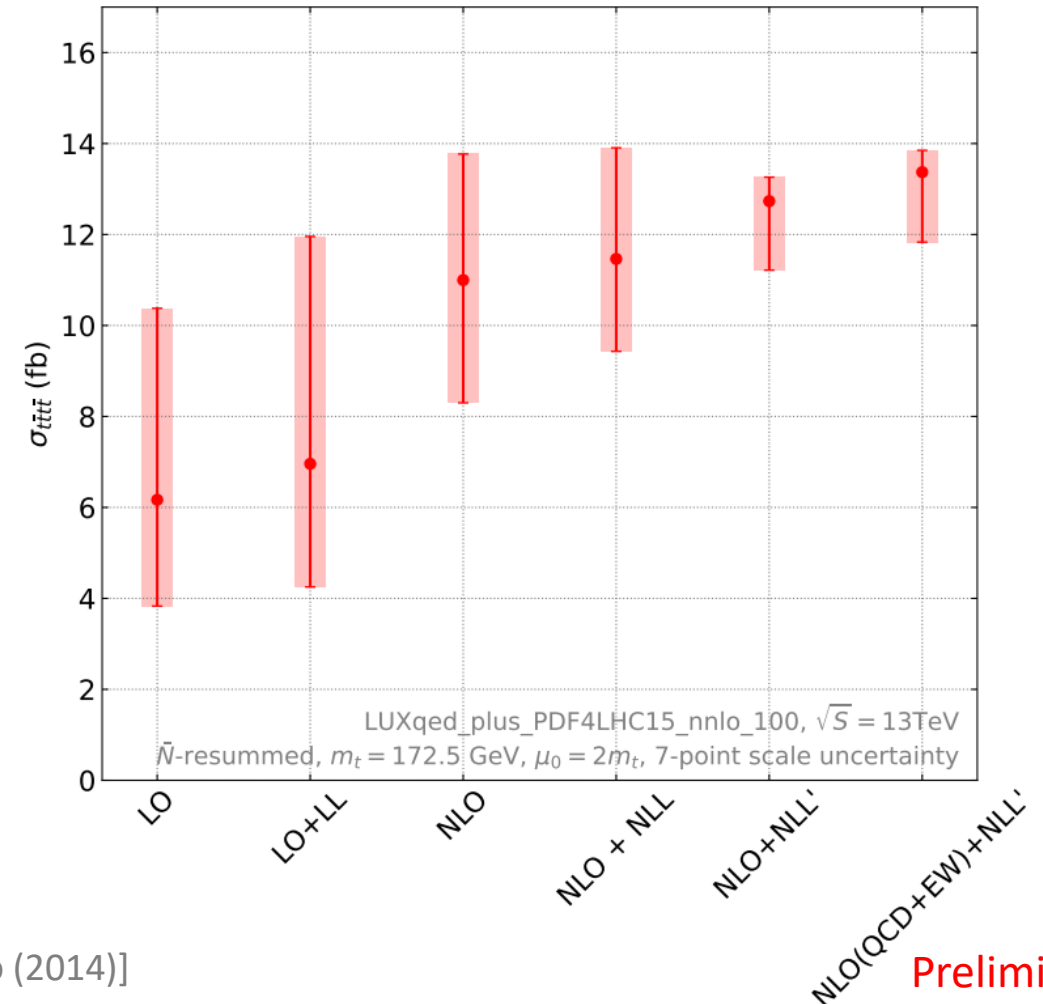
NLO (QCD+EW): EW corrections up to $\mathcal{O}(\alpha^2)$

PDF error: $\pm 6.9\%$

Fixed-order results obtained with

[Frederix, Frixione, Hirschi, Pagani, Shao, Zaro (2018)]

[Alwall, Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Shao, Stelzer, Torrielli, Zaro (2014)]



Preliminary

► 7-point scale variation

	$\sigma_{t\bar{t}t\bar{t}}$ [fb]	K -factor
NLO	$13.14(2)^{+25.1\%}_{-24.4\%}$	
NLO+NLL	$13.81(2)^{+20.7\%}_{-20.1\%}$	1.05
NLO+NLL'	$15.16(2)^{+4.3\%}_{-11.9\%}$	1.15
NLO (QCD+EW)	$13.80(2)^{+22.9\%}_{-22.6\%}$	
NLO (QCD+EW)+NLL'	$15.81(2)^{+3.6\%}_{-11.6\%}$	1.14

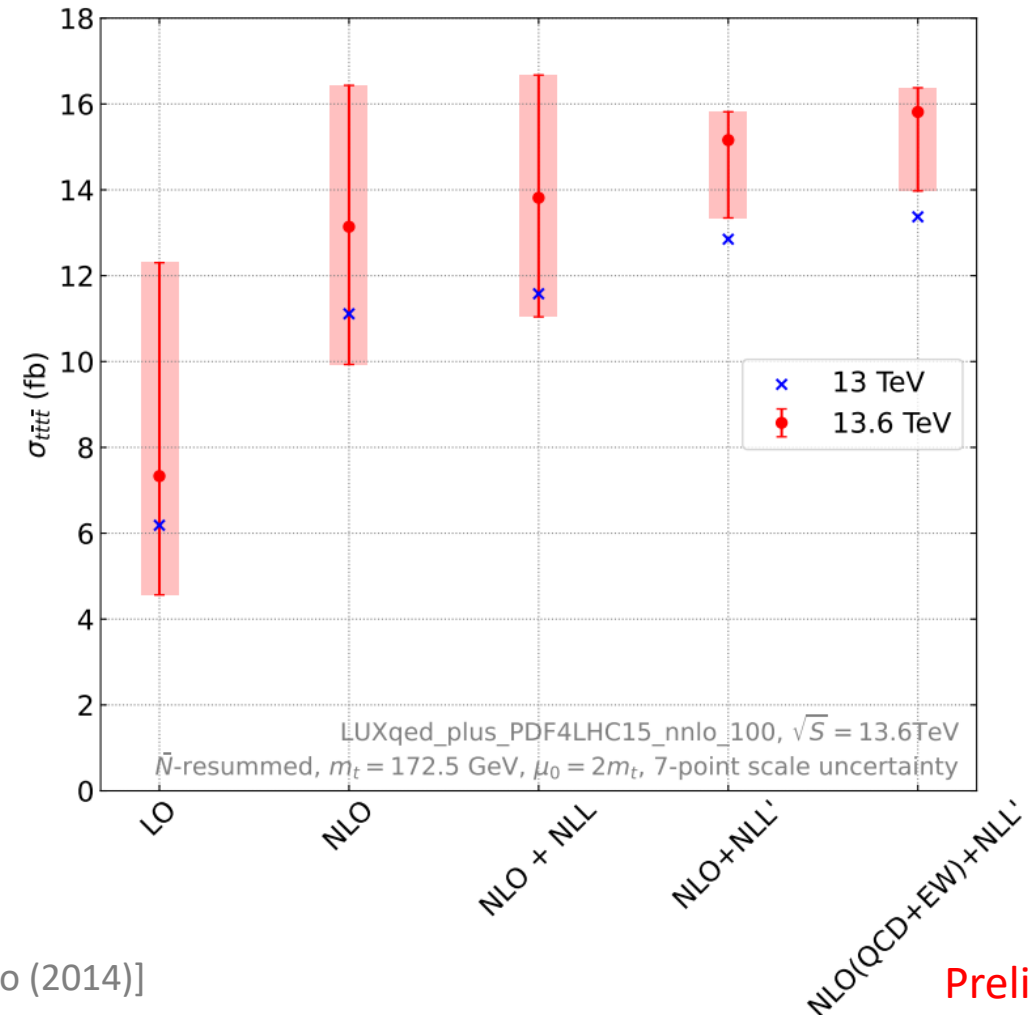
NLO (QCD+EW): EW corrections up to $\mathcal{O}(\alpha^2)$

PDF error: $\pm 6.7\%$

Fixed-order results obtained with

[Frederix, Frixione, Hirschi, Pagani, Shao, Zaro (2018)]

[Alwall, Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Shao, Stelzer, Torrielli, Zaro (2014)]

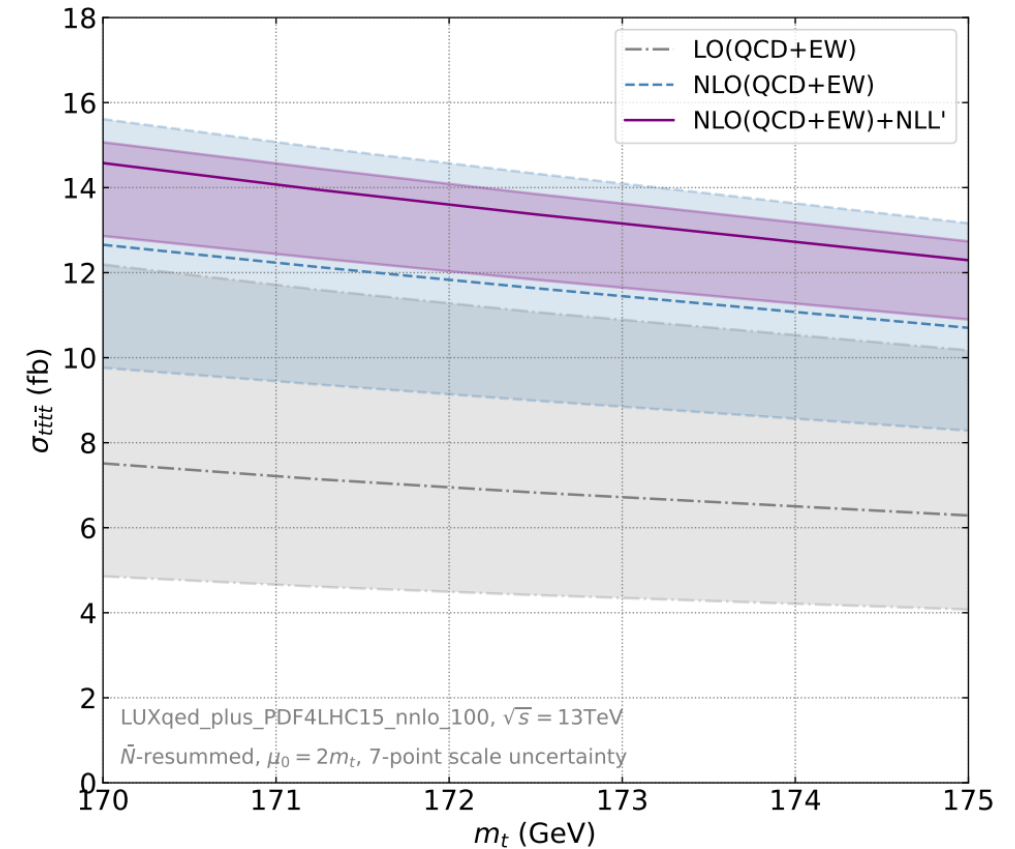


Preliminary

► Top mass $\in [170,175]$ GeV

Error bands: 7-point scale uncertainty

NLO (QCD+EW): EW corrections up to $\mathcal{O}(\alpha^2)$

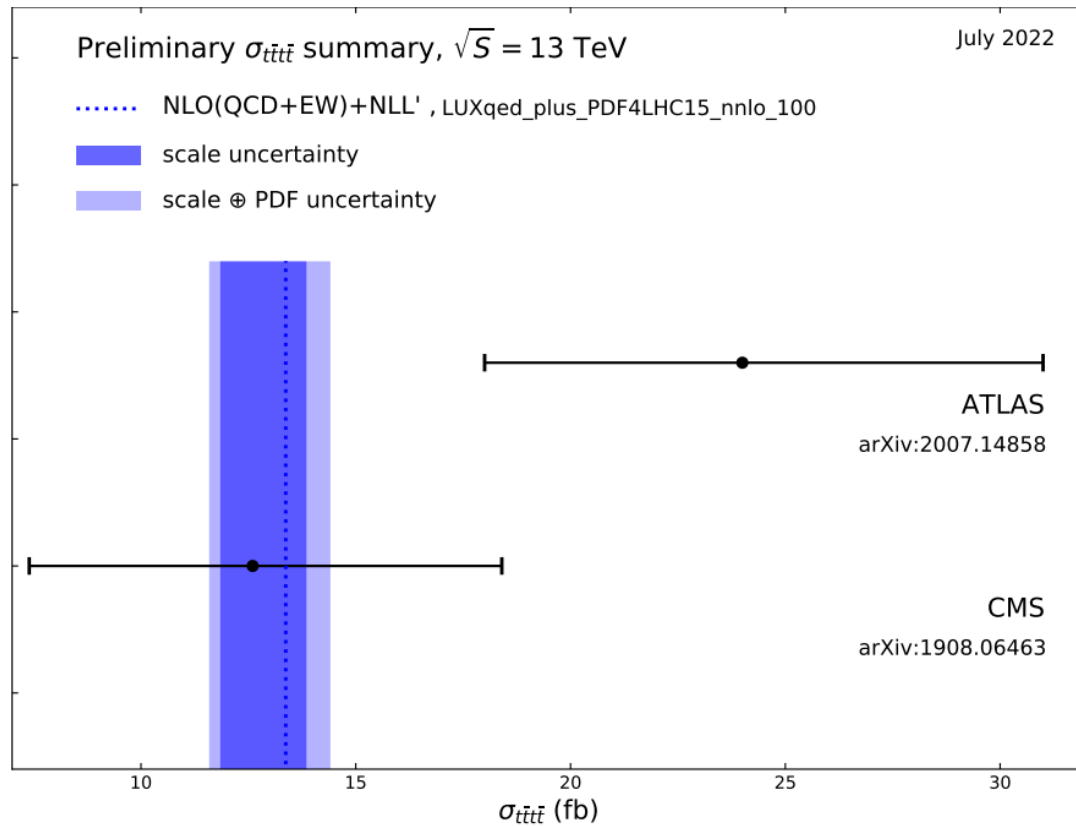


Fixed-order results obtained with

[Frederix, Frixione, Hirschi, Pagani, Shao, Zaro (2018)]

[Alwall, Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Shao, Stelzer, Torrielli, Zaro (2014)]

Preliminary



- ▶ Soft gluon resummation at NLL' accuracy
- ▶ Significant reduction of the total scale uncertainty

NLO (QCD + EW) + NLL'

$13.37 (2) \begin{matrix} +3.6\% & +6.9\% \\ -11.4\% & -6.9\% \end{matrix} \text{ fb}$

Preliminary

Soft gluon resummation for the production of four top quarks at the LHC

Laura Moreno Valero

in collaboration with Melissa van Beekveld and Anna Kulesza

Institute for Theoretical Physics, University of Münster



QCD@LHC

IJCLab Orsay

28/11/2022 - 2/12/2022

