

Feynman integrals & special functions for $pp \rightarrow Hjj$ in NNLO QCD

Vasily Sotnikov

University of Zurich & Michigan State University

in preparation with

S. Abreu, D. Chicherin, H. Ita, B. Page, W. Tschernow, S. Zoia

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**Universität
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Introduction

Motivation

- LHC today is a precision machine
- Many measurements statistically limited \Rightarrow HL-LHC and future colliders
- **Theoretical understanding** of SM predictions is key to **interpret data**
- At least **NNLO QCD** and **NLO EW** corrections (\oplus parton shower, resummation, ...) must be included to achieve **percent level** theory uncertainties

$$d\sigma_{h_1 h_2 \rightarrow X}(p_1, p_2) = \sum_{i,j} \int dx_1 dx_2 f_i(x_1, \mu) f_j(x_2, \mu) \boxed{d\hat{\sigma}_{ij \rightarrow X}(x_1 p_1, x_2 p_2, \mu)} + \mathcal{O}(\Lambda_{\text{QCD}}/Q)$$

"Hard" partonic cross section

$$d\hat{\sigma}_0 \left(1 + \alpha_s \sigma^{(1,0)} + \alpha_s^2 \sigma^{(2,0)} + \alpha_s \sigma^{(0,1)} + \alpha_s^3 \sigma^{(3,0)} + \alpha_s \alpha_s \sigma^{(1,1)} + \dots \right)$$

Intrinsic uncertainty

$$\alpha_s(M_Z) \sim 0.1$$
$$\alpha(M_Z) \sim 0.01$$

- NLO QCD, EW conceptually solved, in practice $\lesssim 8$ partons
- NNLO QCD beyond $2 \rightarrow 2$ remarkably challenging

both **technical** and **conceptual**

What kind of functions loop integrals evaluate to?

process	known	desired
\vdots	\vdots	\vdots QCD precision
$pp \rightarrow V + 2j$	$NLO_{QCD} + NLO_{EW}$	$NNLO_{QCD}$
$pp \rightarrow V + b\bar{b}$	NLO_{QCD}	$NNLO_{QCD} + NLO_{EW}$ <small>[Rene Poncelet's talk]</small>
\vdots	\vdots	\vdots
$pp \rightarrow H + 2j$	$NLO_{HTL} \otimes LO_{QCD}$ $N^3LO_{QCD}^{(VBF^*)}$ (incl.) $NNLO_{QCD}^{(VBF^*)}$ $NLO_{EW}^{(VBF)}$	$NNLO_{HTL} \otimes NLO_{QCD} + NLO_{EW}$ $N^3LO_{QCD}^{(VBF^*)}$ $NNLO_{QCD}^{(VBF)}$
\vdots	\vdots	\vdots

VBF studies

$$\sigma_{NNLO}^{F+X} = \sigma_{NLO}^{F+X} + \int_{\Phi_{F+2}} d\sigma_{RR} + \int_{\Phi_{F+1}} d\sigma_{RV} + \int_{\Phi_F} d\sigma_{VV}$$

[Gloria Bertolotti's talk]

IR divergences

Two-loop amplitudes

Structure of analytic loop amplitudes

Rational/algebraic

Feynman rules, particle content
Integral & tensor reduction

Transcendental

Scattering kinematics
Feynman integrals

$$\mathcal{A} = \sum_{\mathbf{i}} r_{\mathbf{i}}(\mathbf{s}, \epsilon) g^{\mathbf{i}}(\mathbf{s}, \epsilon)$$

Compact analytic form

Function basis

~~Redundancy~~

~~Dim. reg. artifacts~~

✓ Analytic cancellation of IR divergences

✓ Enable modern finite-field based methods

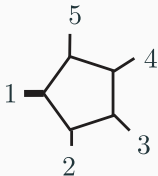
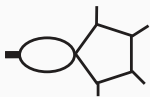
✓ Amplitudeology friendly

Goal: fully differential cross sections

fast and stable evaluation over whole physical phase space

Setup of the computation

Integral topologies & kinematics

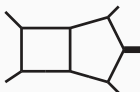


Variables

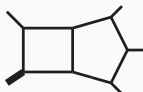
$$p_1^2, s_{12}, s_{23}, s_{34}, s_{45}, s_{15}$$



PBmzz



PBzmm



PBzzz

Gram determinants:

$$\Delta_5 = 16 G(p_1, p_2, p_3, p_4)$$

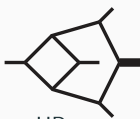
$$\Delta_3^{(1)} = -4 G(p_1, p_2 + p_3)$$

$$\Delta_3^{(2)} = -4 G(p_1, p_2 + p_4)$$

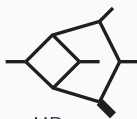
$$\Delta_3^{(3)} = -4 G(p_1, p_3 + p_4)$$



HBzzz



HBzmm



HBmzz

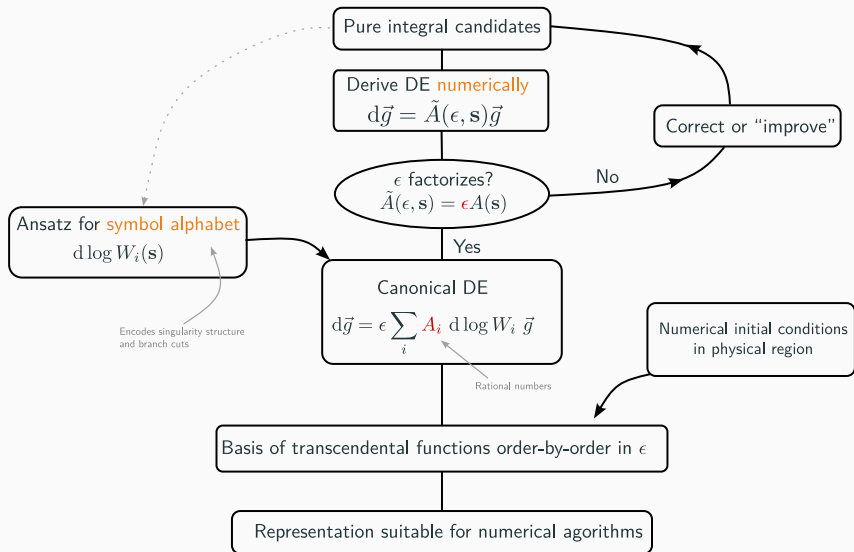


DPmz



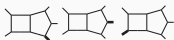
DPzz

Hope: the result is pure functions of uniform transcendentality [Henn '13]



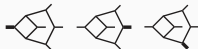
Canonical DE

Planar



[Abreu, Ita, Moriello, Page, Tschernow, Zeng '20]

Hexa-box



[Abreu, Ita, Page, Tschernow '21]

Results through generalized polylogarithms

[Papadopoulos, Tommasini, Wever '15]

[Canko, Papadopoulos, Syrrakos '20] [Syrrakos '20]

[Kardos, Papadopoulos, Smirnov, Syrrakos, Wever '22]

Function basis (planar)

[Badger, Hartanto, Zoia '21] color-ordered

[Chicherin, VS, Zoia '21]

Semi-numerical DE solution

DiffExp [Moriello '19] [Hidding '20]

AMFlow [Liu, Ma, Wang '17] [Liu, Ma '21]

SeaSyde [Armadiello, Bonciani, Devoto, Rana, Vicini '22]
[Hidding, Usovitsch '22]

- Initial (boundary) conditions
- Cross checks
- Numerical data for analytic work

Differential equations

Finding pure MIs

Great progress,
(semi-)automated in many cases

Canonica [Meyer '18] Epsilon [Prausa '17]
Fuchsia [Gituliar, Magerya '17]
DlogBasis [Henn, Mistlberger, V. Smirnov, Wasser '20]
Initial [Dlapa, Henn, Yan '20] [Dlapa, Henn, Wagner '22]
[Dlapa, X. Li, Y. Zhang '21]

Does not quite cut it for our case...

Rely on **educated guessing**, main guides:

- Integrands with **constant leading singularities**
[Arkani-Hamed, Bourjaily, Cachazo, Trnka '10]
- Generalized (unitarity) cuts
- Recycle known simpler results
- Get inspired by symmetries ($\mathcal{N} = 4$ sYM, conformal)
- Guesses are easy to check numerically

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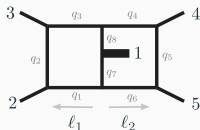
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Example:

9 pure integrals (179 in whole family)



$$\rho_i := q_i^2$$

$$\mu_{ij} := \ell_i^{(D-4)} \ell_j^{(D-4)}$$

$$\sqrt{\Delta_5} \sqrt{\Delta_3^{(1)}} \mu_{12} \quad , \quad \frac{\sqrt{\Delta_3^{(1)}}}{\sqrt{\Delta_5}} \frac{\partial \mathcal{B}}{\partial \rho_8}$$

$$\sqrt{\Delta_5} (2s_{23} \mu_{22} + (s_{23} + s_{45} - p_1^2) \mu_{12}), \quad \begin{matrix} 2 \leftrightarrow 5 \\ 3 \leftrightarrow 4 \end{matrix}$$

$$\frac{1}{\sqrt{\Delta_5}} (2s_{23} \frac{\partial \mathcal{B}}{\partial \rho_3} - (p_1^2 + s_{23} - s_{45}) \left(\frac{\partial \mathcal{B}}{\partial \rho_7} + \Delta_5 \mu_{12} \right)), \quad \begin{matrix} 2 \leftrightarrow 5 \\ 3 \leftrightarrow 4 \end{matrix}$$

$$\text{tr} \left(\not{p}_2 \not{q}_2 \not{q}_3 \not{q}_8 \not{q}_7 \not{q}_6 \not{q}_5 \not{p}_4 \right), \quad \begin{matrix} 2 \leftrightarrow 5 \\ 3 \leftrightarrow 4 \end{matrix}$$

$$\sqrt{\Delta_3^{(1)}} \text{tr} \left(\not{q}_5 \not{p}_5 \not{p}_2 \not{q}_2 \not{q}_3 \not{q}_4 \right)$$

The alphabet

Hexa-box alphabet [Abreu, Ita, Page, Tschernow '21], **no new letters from DP**

$$A = \sum_{i=1}^{204} d \log W_i A_i$$

c.f. 31 letters,
1 square root
for massless!

Algebraic letters, odd under root sign flip

$$\text{e.g. } W_{118} = \frac{p_1^2 - s_{23} + s_{45} + \sqrt{\Delta_3^{(1)}}}{p_1^2 - s_{23} + s_{45} - \sqrt{\Delta_3^{(1)}}}$$

Degree	Letters	
linear	27	} 117 rational
quadratic	66	
cubic	24	

Roots	Letters	
$\sqrt{\Delta_5}$	32	} 68 one square root
$\sqrt{\Delta_3^{(i)}}$	12	
$\sqrt{\Sigma_5^{(i)}}$	24	
$\sqrt{\Delta_5}, \sqrt{\Delta_3^{(i)}}$	3	} 9 two square roots
$\sqrt{\Delta_5}, \sqrt{\Sigma_5^{(i)}}$	6	

Root	Degree	Perms./ Letters	
$\sqrt{\Delta_5}$	4	1	} 10 square roots
$\sqrt{\Delta_3^{(i)}}$	2	3	
$\sqrt{\Sigma_5^{(i)}}$	4	6	

$$\Sigma_5^{(1)} = (s_{12}s_{15} - s_{12}s_{23} - s_{15}s_{45} + s_{34}s_{45} + s_{23}s_{34})^2 - 4s_{23}s_{34}s_{45}(s_{34} - s_{12} - s_{15})$$

Function basis

Basis construction

[Chicherin, VS, Zoia '21] (see also [Chicherin, VS '20] [Badger, Hartanto, Zoia '21])

Pure Feynman integrals


Canonical DE

$$d\vec{g} = \epsilon A \vec{g}$$
$$A = \sum_i d \log W_i(s) A_i$$

Initial conditions
physics, limits, PSLQ

weight = length = ϵ order

Chen iterated integrals [Chen '77]

$$[W_1, \dots, W_n]_\gamma = \int_0^1 d \log W_n(t_n) \dots \int_0^{t_2} d \log W_n(t_1)$$


Vector subspace, weight-graded

$$\mathbf{G} = \bigoplus_w \mathbf{G}^{(w)}$$

+ shuffle product

$$\mathbf{G}^{(w_1)} \times \mathbf{G}^{(w_2)} \mapsto \mathbf{G}^{(w_1+w_2)}$$

$$[W_1, \dots, W_r]_\gamma [W_{r+1}, \dots, W_n]_\gamma$$
$$= \sum_{\mathbf{i} \in \text{shuffles}} [W_{i_1}, \dots, W_{i_n}]_\gamma$$

Basis in $\mathbf{G}^{(w)}$ mod products

- ✓ complete
- ✓ non-redundant

Initial values

Function basis construction requires **algebraic relations** between **initial values** $\vec{g}(X_0)$

Previously

Serious bottleneck [Chicherin, VS, Zoia '21]

- Rely on MPL expressions [Canko, Papadopoulos, Syrrakos '20] to calculate $\vec{g}(X_0)$ to $\mathcal{O}(3000)$ digits
- Pushing the limits of most advanced PSLQ algorithms [Bailey, Broadhurst '01] [Bailey, Borwein, Kimberley, Ladd '17] (~ 400 constants with ~ 2000 digits)

New approach

1. Construct **symbol-level** function basis, i.e. setting $\vec{g}^{(w)}(X_0)$ with $w > 0$
2. Use their definitions through MI components to **upgrade to iterated integrals**
3. Presume that MI expressions are polynomials in basis functions and $\langle i\pi \rangle \oplus \langle \pi^2 \rangle \oplus \langle i\pi^3, \zeta_3 \rangle \oplus \langle \pi^4, i\pi\zeta_3 \rangle \implies$ **derive constraints**
4. Match to numerical evaluation from AMFlow to validate and fix remaining rational numbers

- ✓ Precision for $\vec{g}(X_0)$ need not exceed final target precision
- ✓ PSLQ fit just one rational number at a time

Non-analyticity within physical region

Consider iterated integral along $\gamma : t \in [0, 1] \rightarrow \mathcal{P}_{\text{phys}}$, and $W_i(t^*) = 0$,

$$\int_{\gamma} d \log W_i h = \int_{\gamma} \frac{dW_i}{W_i} h \xrightarrow{t \rightarrow t^*} \frac{W'(t)}{t - t^*} \left(h^{(0)} + h^{(1)}(t - t^*) + \mathcal{O}\left((t - t^*)^2\right) \right)$$

Planar scattering

Only **linear or quadratic** letters vanish in $\mathcal{P}_{\text{phys}}$, poles always canceled, i.e. $h^{(0)} = 0$

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Planar scattering

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New feature of nonplanar scattering


Square roots of quartic polynomials $\sqrt{\Sigma_5}$ can vanish in $\mathcal{P}_{\text{phys}} \implies$ new types of **divergences**

1. Integrable square root: $d \log \frac{a + \sqrt{\Sigma_5}}{a - \sqrt{\Sigma_5}} \xrightarrow{\Sigma_5 \rightarrow 0} \frac{d\Sigma_5}{a\sqrt{\Sigma_5}} \xrightarrow{t \rightarrow t^*} \frac{C}{\sqrt{t - t^*}} + \dots$
2. Uncompensated poles: $d \log \sqrt{\Sigma_5} \xrightarrow{\Sigma_5 \rightarrow 0} \frac{d\Sigma_5}{2\Sigma_5} \xrightarrow{t \rightarrow t^*} \frac{C}{t - t^*} + \dots \implies$ **log divergence!**

- Choose basis functions to localize non-analytic behavior
- Expectation: functions with type 2 divergence cancel out in physical results
- Numerical evaluation more challenging

Basis structure

Weight	P ∪ PB	+HB	+DP	Total
1	11	0	0	11
2	25	10	0	35
3	145	72	0	217
4	675	305	48	1028
#MIs	1361	542	345	2248



Permutation closed

$$\sigma \left(f_i^{(w)} \right) \rightarrow \sum_j c_{ij} f_j^{(w)} + \dots$$

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Permutation closed

$$\sigma \left(f_i^{(w)} \right) \rightarrow \sum_j c_{ij} f_j^{(w)} + \dots$$

$d \log \sqrt{\Delta_5}$ likely cancel in finite remainders
[Chicherin, Henn, Papathanasiou '20]

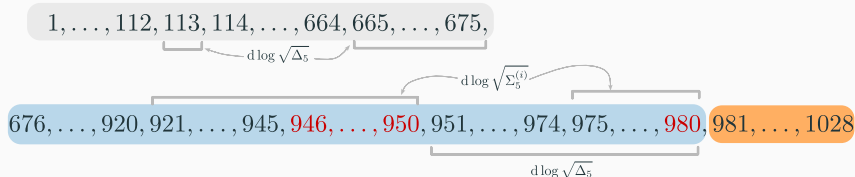
$d \log \sqrt{\Sigma_5^{(i)}}$ also cancel?

7 functions diverge at $\Sigma_5^{(3)} = 0$

Weight 3



Weight 4



Numerical evaluation

Weights 1 and 2

Well-defined combinations of
log, Li₂ functions

$$f_{13}^{(2)} = \text{Li}_2 \left(1 - \frac{s_{15} - s_{23} - s_{34}}{s_{15}} \right)$$

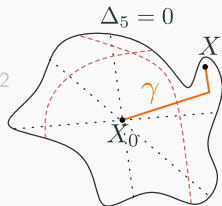
> 0

Weights 3 and 4

$$f_i^{(3)} = \int_0^1 \sum_j \frac{\partial \log W_j(t)}{\partial t} h_{i,j}^{(2)} dt$$

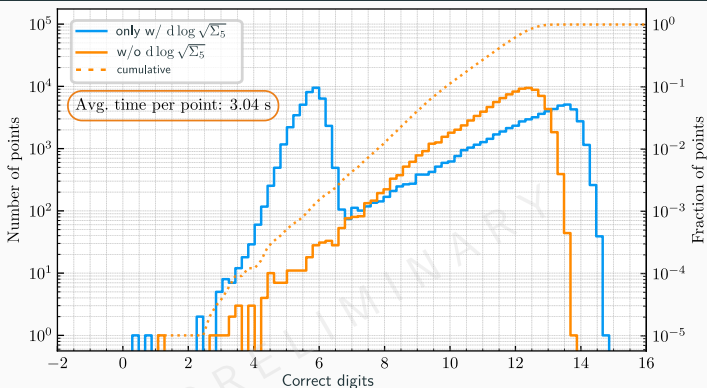
$$f_i^{(4)} = \int_0^1 \sum_{j,k} \frac{\partial \log W_j(t)}{\partial t} \log \frac{W_k(1)}{W_k(t)} h_{i,jk}^{(2)} dt$$

weight 2



- Numerical **one-fold** integration [Caron-Huot, Henn '14] of **analytic** integrands \implies **exponential** convergence [Takahasi, Mori '73]
- No crossing of physical thresholds \implies **no analytic continuation** needed
- Dedicated **series expansions** around (spurious) singularities

Numerical performance



- Sample over physical phase space for NLO Wjj production at the LHC
- Evaluate **all** functions on each point, plot the worst accuracy per point
- Timing for all functions on one CPU
- Worse stability for functions with $d \log \sqrt{\Sigma_5}$ (hopefully irrelevant!)

Very promising performance!

Soon available in `PentagonFunctions++`

<https://gitlab.com/pentagon-functions/PentagonFunctions-cpp>

Conclusions

Conclusions

- A complete set of special functions describing double virtual corrections for two-loop five-point one-mass processes is available.
- Enables application of modern techniques for analytic calculation of two-loop amplitudes.
- Ready for cross section calculations.

Outlook

- Paper to appear soon.
- Possibility of calculating NNLO QCD corrections for a large class of $2 \rightarrow 3$ processes is open. Hopefully exciting phenomenology in near future!
- Important contribution towards N³LO QCD corrections for Vj , $V\gamma$ production.

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Backup

Numerical performance: conservative precision rescue

