

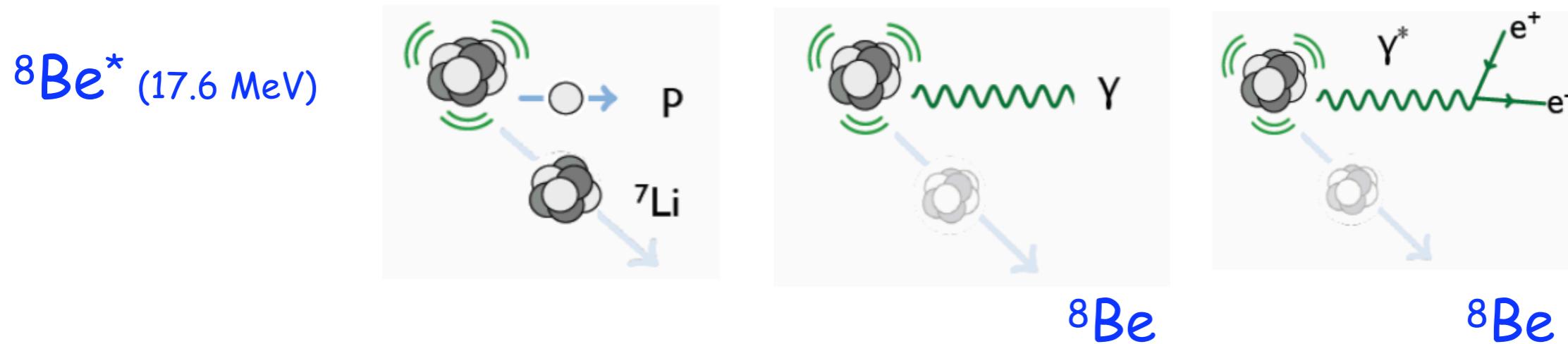
The X17 anomaly: status and prospect

Enrico Nardi



Silafae XIV - USFQ - Nov. 14-18, 2022

A nuclear physics anomaly: in ${}^8\text{Be}^*$ de-excitation process



Anomaly first observed in
 ${}^8\text{Be}$ Nuclear Transitions

=>

${}^8\text{Be}^* \rightarrow p + {}^7\text{Li}$ (mostly)

${}^8\text{Be}^* \rightarrow {}^8\text{Be} + \gamma$ ($B_\gamma = 1.4 \times 10^{-5}$)

${}^8\text{Be}^* \rightarrow {}^8\text{Be} + e^+e^-$ ($B_{e^\pm} = 4 \times 10^{-3} B_\gamma$)

$[{}^8\text{Be}^* \rightarrow {}^8\text{Be} + X_{17} \xrightarrow{\quad} e^+e^- \quad (B_X = 6 \times 10^{-6} B_\gamma)]$

Anomalies in nuclear transitions observed by the Atomki experiment

Summary:

Anomalies in nuclear transitions observed by the Atomki experiment

Summary:

- **2015:** First anomaly observed in the angular correlation of e^+e^- pairs emitted in nuclear transition of ${}^8Be^*(18.15\text{ MeV}) \rightarrow \text{ground state (g.s.)}$ [Phys. Rev. Lett. 116, 042501 (2016)]

Anomalies in nuclear transitions observed by the Atomki experiment

Summary:

- **2015:** First anomaly observed in the angular correlation of e^+e^- pairs emitted in nuclear transition of ${}^8\text{Be}^*(18.15 \text{ MeV}) \rightarrow$ ground state (g.s.) [Phys. Rev. Lett. 116, 042501 (2016)]
- **2017:** With improved exp. setup, similar anomaly in ${}^8\text{Be}^*(17.64 \text{ MeV}) \rightarrow$ to g.s. (previously not observed). Reported in Messina symposium (Oct 2016) and Bormio meeting (Jan 2017) [EPJ Web Conf. 142 (2017) 01019; PoS BORMIO 2017 (2017)]

Anomalies in nuclear transitions observed by the Atomki experiment

Summary:

- **2015:** First anomaly observed in the angular correlation of e^+e^- pairs emitted in nuclear transition of ${}^8\text{Be}^*(18.15 \text{ MeV}) \rightarrow$ ground state (g.s.) [Phys. Rev. Lett. 116, 042501 (2016)]
- **2017:** With improved exp. setup, similar anomaly in ${}^8\text{Be}^*(17.64 \text{ MeV}) \rightarrow$ to g.s. (previously not observed). Reported in Messina symposium (Oct 2016) and Bormio meeting (Jan 2017) [EPJ Web Conf. 142 (2017) 01019; PoS BORMIO 2017 (2017)]
- **2018:** Confirmation of ${}^8\text{Be}$ result (thinner target, 5+1 telescopes). First hint of similar anomaly in ${}^4\text{He}^*(21 \text{ MeV})$ transition [Zakopane Conf., Acta Phys.Polon.B 50 (2019) 3, 675]

Anomalies in nuclear transitions observed by the Atomki experiment

Summary:

- **2015:** First anomaly observed in the angular correlation of e^+e^- pairs emitted in nuclear transition of ${}^8Be^*(18.15\text{ MeV}) \rightarrow$ ground state (g.s.) [Phys. Rev. Lett. 116, 042501 (2016)]
- **2017:** With improved exp. setup, similar anomaly in ${}^8Be^*(17.64\text{ MeV}) \rightarrow$ to g.s. (previously not observed). Reported in Messina symposium (Oct 2016) and Bormio meeting (Jan 2017) [EPJ Web Conf. 142 (2017) 01019; PoS BORMIO 2017 (2017)]
- **2018:** Confirmation of 8Be result (thinner target, 5+1 telescopes). First hint of similar anomaly in ${}^4He^*(21\text{ MeV})$ transition [Zakopane Conf., Acta Phys.Polon.B 50 (2019) 3, 675]
- **2019:** Confirmation of 4He bump (7.2σ) consistent with $M_x \sim 17\text{ MeV}$ interpretation [Phys.Rev.C 104 (2021) 4, 044003 · (received 27 October 2019) - e-Print: [2104.10075](#) supersedes [1910.10459](#)]

Anomalies in nuclear transitions observed by the Atomki experiment

Summary:

- **2015:** First anomaly observed in the angular correlation of e^+e^- pairs emitted in nuclear transition of ${}^8Be^*(18.15\text{ MeV}) \rightarrow$ ground state (g.s.) [Phys. Rev. Lett. 116, 042501 (2016)]
- **2017:** With improved exp. setup, similar anomaly in ${}^8Be^*(17.64\text{ MeV}) \rightarrow$ to g.s. (previously not observed). Reported in Messina symposium (Oct 2016) and Bormio meeting (Jan 2017) [EPJ Web Conf. 142 (2017) 01019; PoS BORMIO 2017 (2017)]
- **2018:** Confirmation of 8Be result (thinner target, 5+1 telescopes). First hint of similar anomaly in ${}^4He^*(21\text{ MeV})$ transition [Zakopane Conf., Acta Phys.Polon.B 50 (2019) 3, 675]
- **2019:** Confirmation of 4He bump (7.2σ) consistent with $M_X \sim 17\text{ MeV}$ interpretation [Phys.Rev.C 104 (2021) 4, 044003 · (received 27 October 2019) - e-Print: [2104.10075](#) supersedes [1910.10459](#)]
- **2021:** Preliminary results for ${}^{12}C^*(17.2\text{ MeV})$ decaying to g.s.: excess of e^+e^- pairs at large angles ($\sim 160^\circ$). [A.J. Krasznahorkay, "Shedding light on X17" workshop, Centro Fermi, Rome, Sept. 2021]

Anomalies in nuclear transitions observed by the Atomki experiment

Summary:

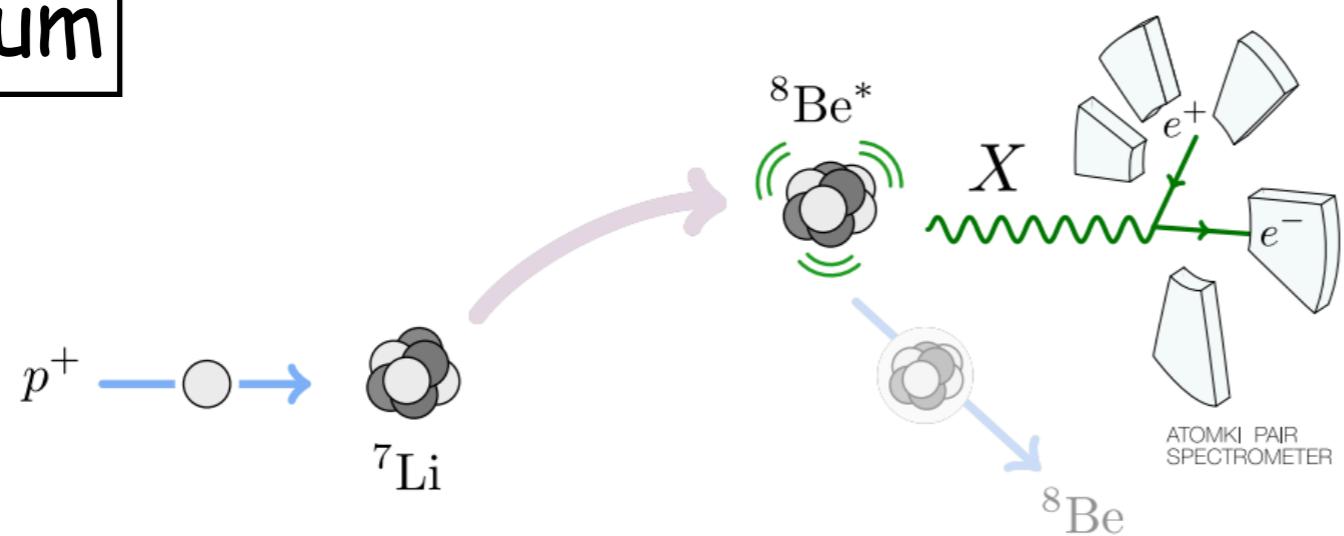
- **2015:** First anomaly observed in the angular correlation of e^+e^- pairs emitted in nuclear transition of ${}^8Be^*(18.15\text{ MeV}) \rightarrow$ ground state (g.s.) [Phys. Rev. Lett. 116, 042501 (2016)]
- **2017:** With improved exp. setup, similar anomaly in ${}^8Be^*(17.64\text{ MeV}) \rightarrow$ to g.s. (previously not observed). Reported in Messina symposium (Oct 2016) and Bormio meeting (Jan 2017) [EPJ Web Conf. 142 (2017) 01019; PoS BORMIO 2017 (2017)]
- **2018:** Confirmation of 8Be result (thinner target, 5+1 telescopes). First hint of similar anomaly in ${}^4He^*(21\text{ MeV})$ transition [Zakopane Conf., Acta Phys.Polon.B 50 (2019) 3, 675]
- **2019:** Confirmation of 4He bump (7.2σ) consistent with $M_X \sim 17\text{ MeV}$ interpretation [Phys.Rev.C 104 (2021) 4, 044003 · (received 27 October 2019) - e-Print: [2104.10075](#) supersedes [1910.10459](#)]
- **2021:** Preliminary results for ${}^{12}C^*(17.2\text{ MeV})$ decaying to g.s.: excess of e^+e^- pairs at large angles ($\sim 160^\circ$). [A.J. Krasznahorkay, "Shedding light on X17" workshop, Centro Fermi, Rome, Sept. 2021]
- **2022:** Confirmation of e^+e^- excess in ${}^{12}C^*(17.2\text{ MeV}) \rightarrow$ g.s. at large angles [$\sim 155^\circ - 160^\circ$] [A.J. Krasznahorkay, e-Print: [2209.10795](#) [nucl-ex], rev. v2 Nov. 2,2022]

The Atomki experimental apparatus

The Atomki experimental apparatus

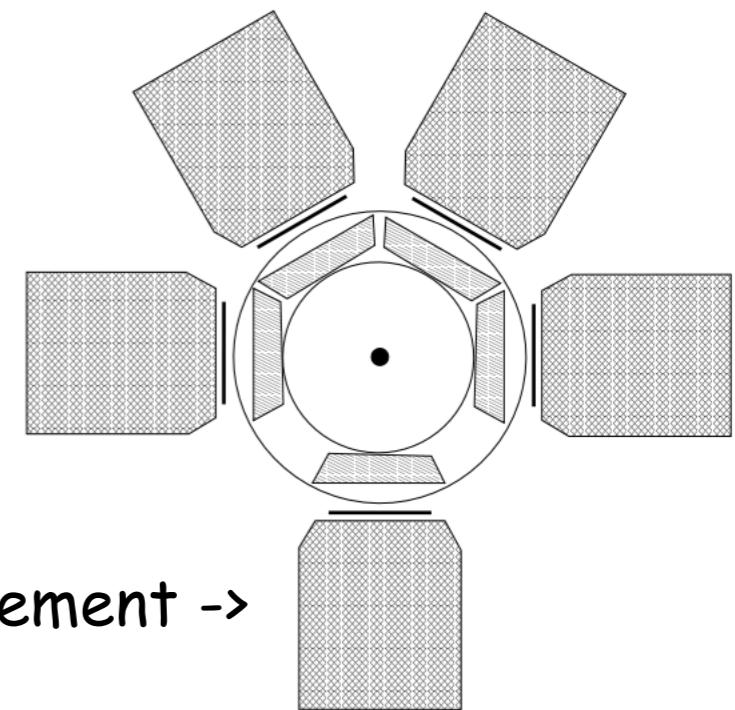
0 2 4 6 8 10 cm

Berillium

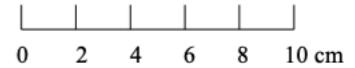


Feng+, 1608.0359

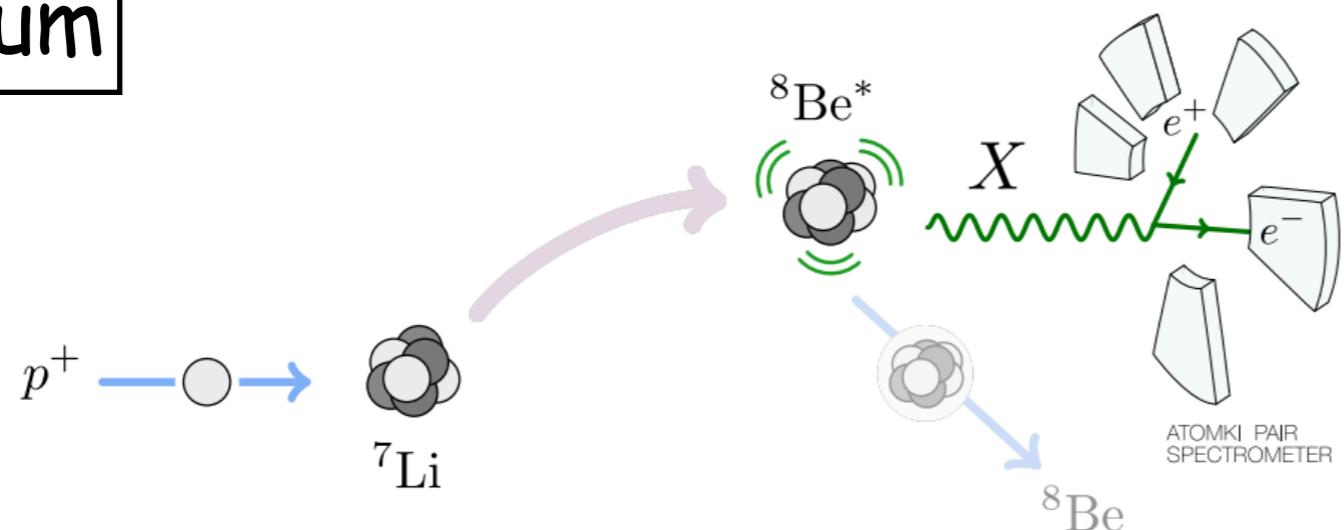
Five telescopes arrangement ->



The Atomki experimental apparatus

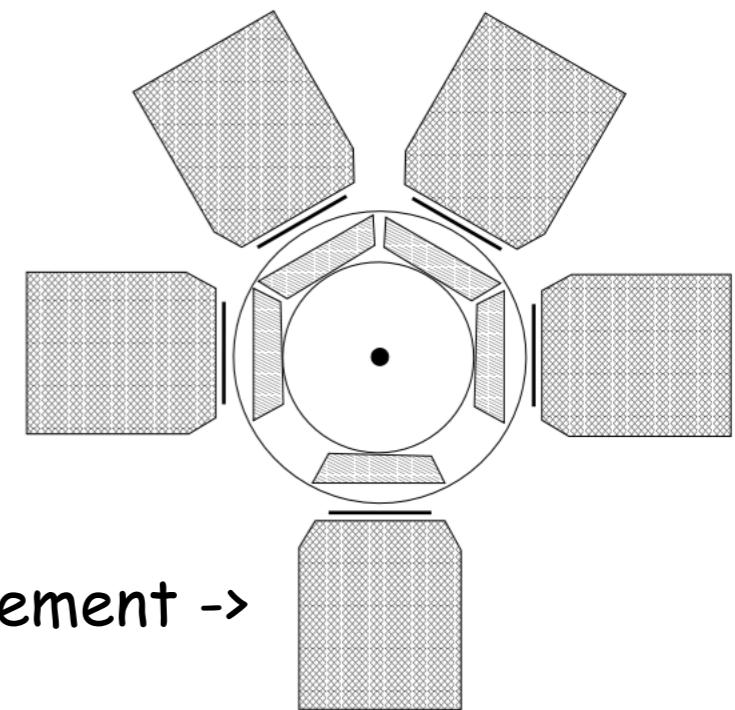


Berillium



Feng+, 1608.0359

Five telescopes arrangement ->

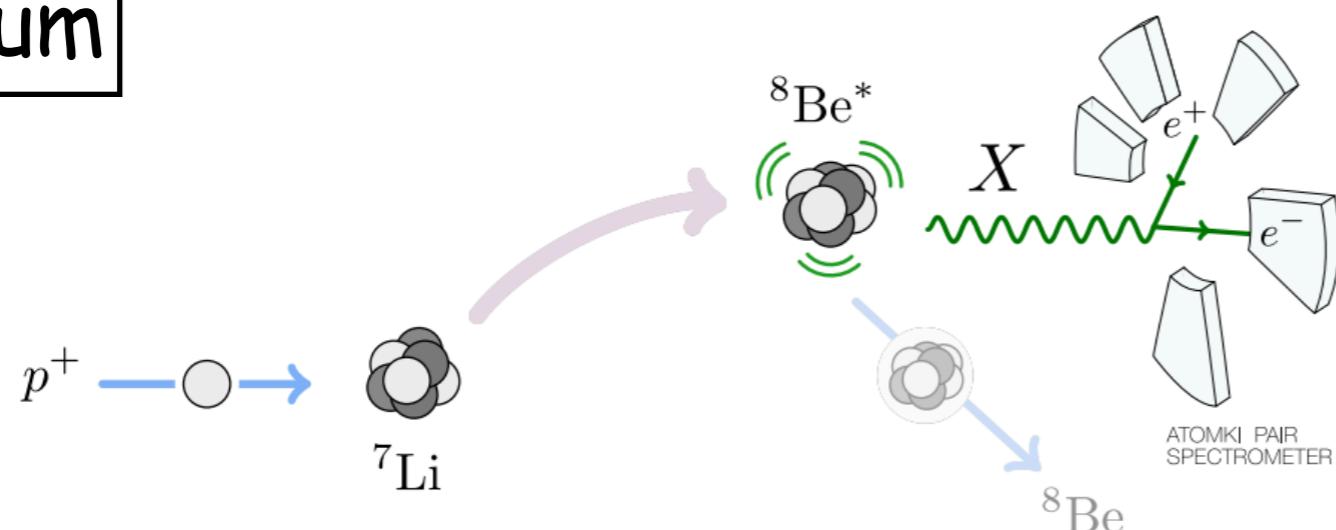


During the years, several improvements in the apparatus (accelerator, detectors, electronics)

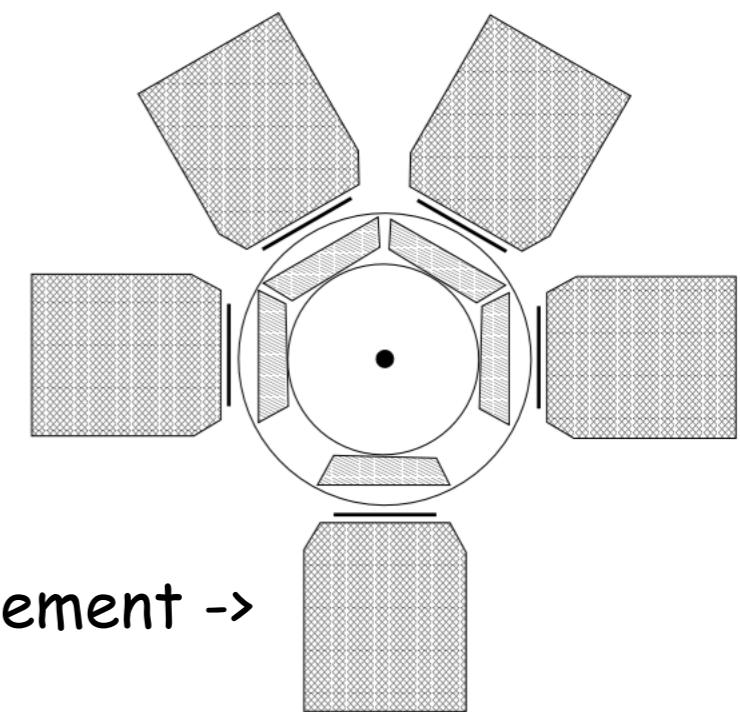
The Atomki experimental apparatus

0 2 4 6 8 10 cm

Berillium



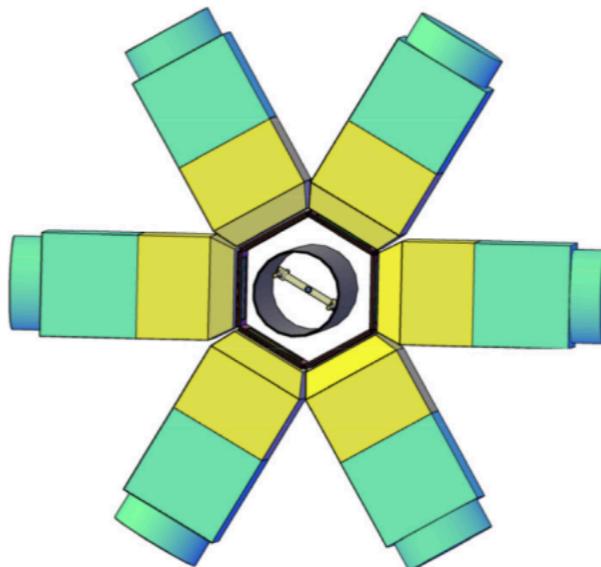
Feng+, 1608.0359



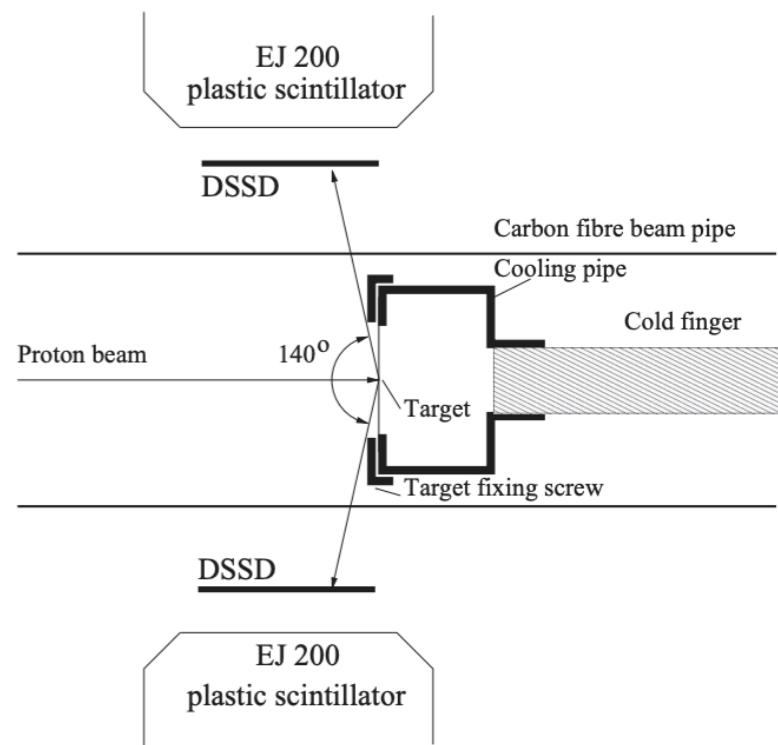
Five telescopes arrangement ->

During the years, several improvements in the apparatus (accelerator, detectors, electronics)

Helium



Six telescopes arrangement ->



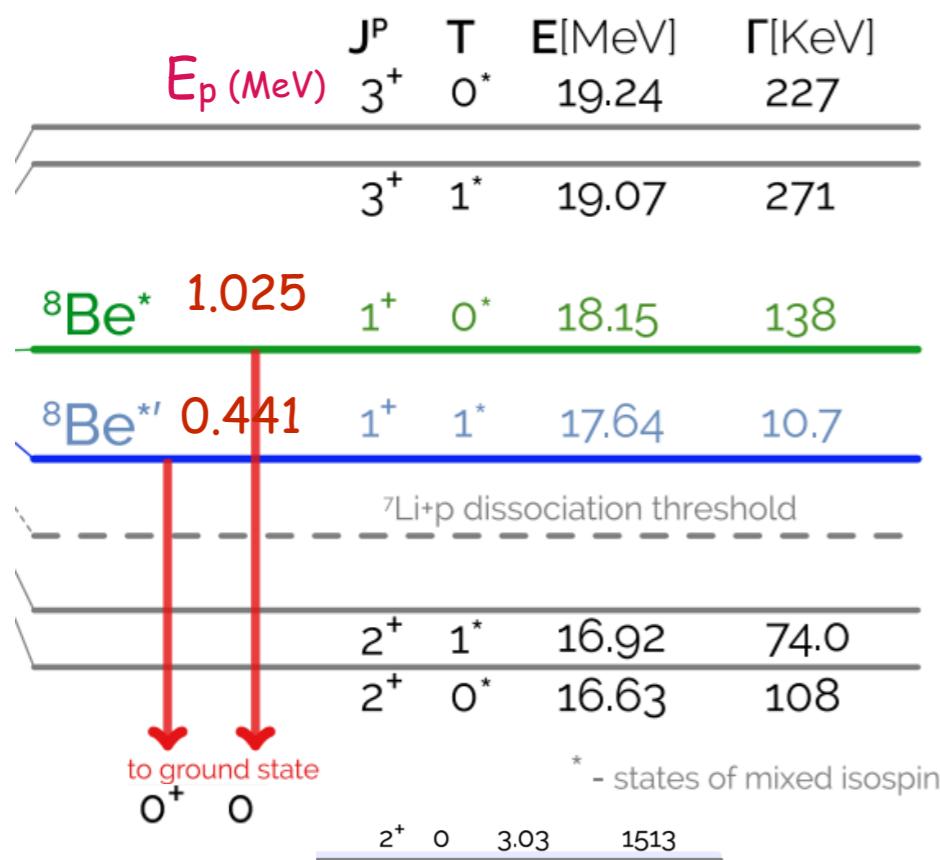
Atomki e^+e^- pairs measurements:

1. energy-sum spectrum $E_{\pm} = E_{e^+} + E_{e^-}$
2. e^+e^- angular correlations θ

Berillium nuclear transitions

Berillium nuclear transitions

Resonant transition $p + {}^7\text{Li} \rightarrow {}^8\text{Be}^* \rightarrow \dots$



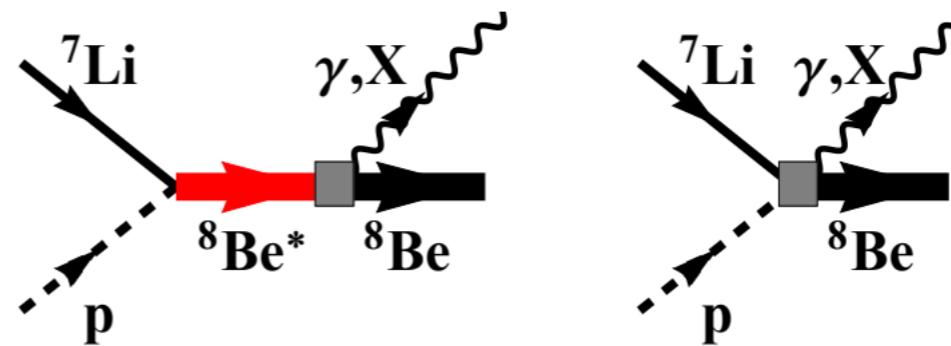
Berillium nuclear transitions

E_p (MeV)	J^P	T	E [MeV]	Γ [KeV]
	3^+	0^*	19.24	227
	3^+	1^*	19.07	271
${}^8\text{Be}^*$ 1.025	1^+	0^*	18.15	138
${}^8\text{Be}'^{*1}$ 0.441	1^+	1^*	17.64	10.7
${}^7\text{Li} + p$ dissociation threshold				
	2^+	1^*	16.92	74.0
	2^+	0^*	16.63	108
* - states of mixed isospin				
	0^+	0	3.03	1513

to ground state

Resonant transition $p + {}^7\text{Li} \rightarrow {}^8\text{Be}^* \rightarrow \dots$

Radiative $p + {}^7\text{Li} \rightarrow {}^8\text{Be} + \gamma$



M1 resonant transition - E1 direct p capture
(valid also for a Vector X_{17})

Berillium nuclear transitions

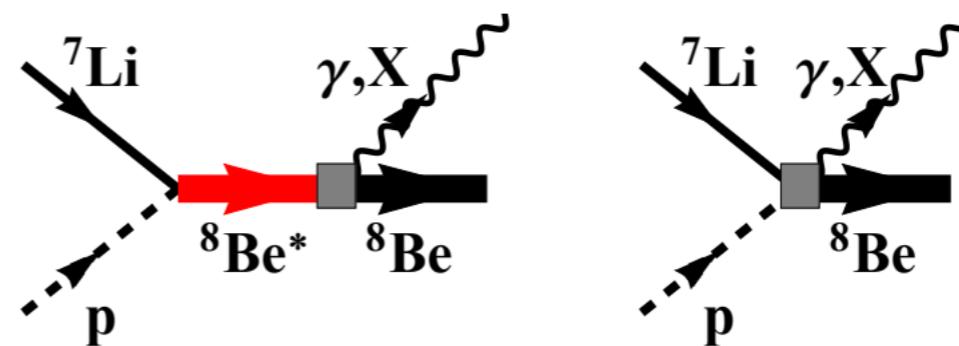
E_p (MeV)	J^P	T	E [MeV]	Γ [KeV]
	3^+	0^*	19.24	227
	3^+	1^*	19.07	271
$^{8}\text{Be}^*$ 1.025	1^+	0^*	18.15	138
$^{8}\text{Be}^{* \prime}$ 0.441	1^+	1^*	17.64	10.7
--- ---				
	2^+	1^*	16.92	74.0
	2^+	0^*	16.63	108

	0^+	0		
	2^+	0	3.03	1513

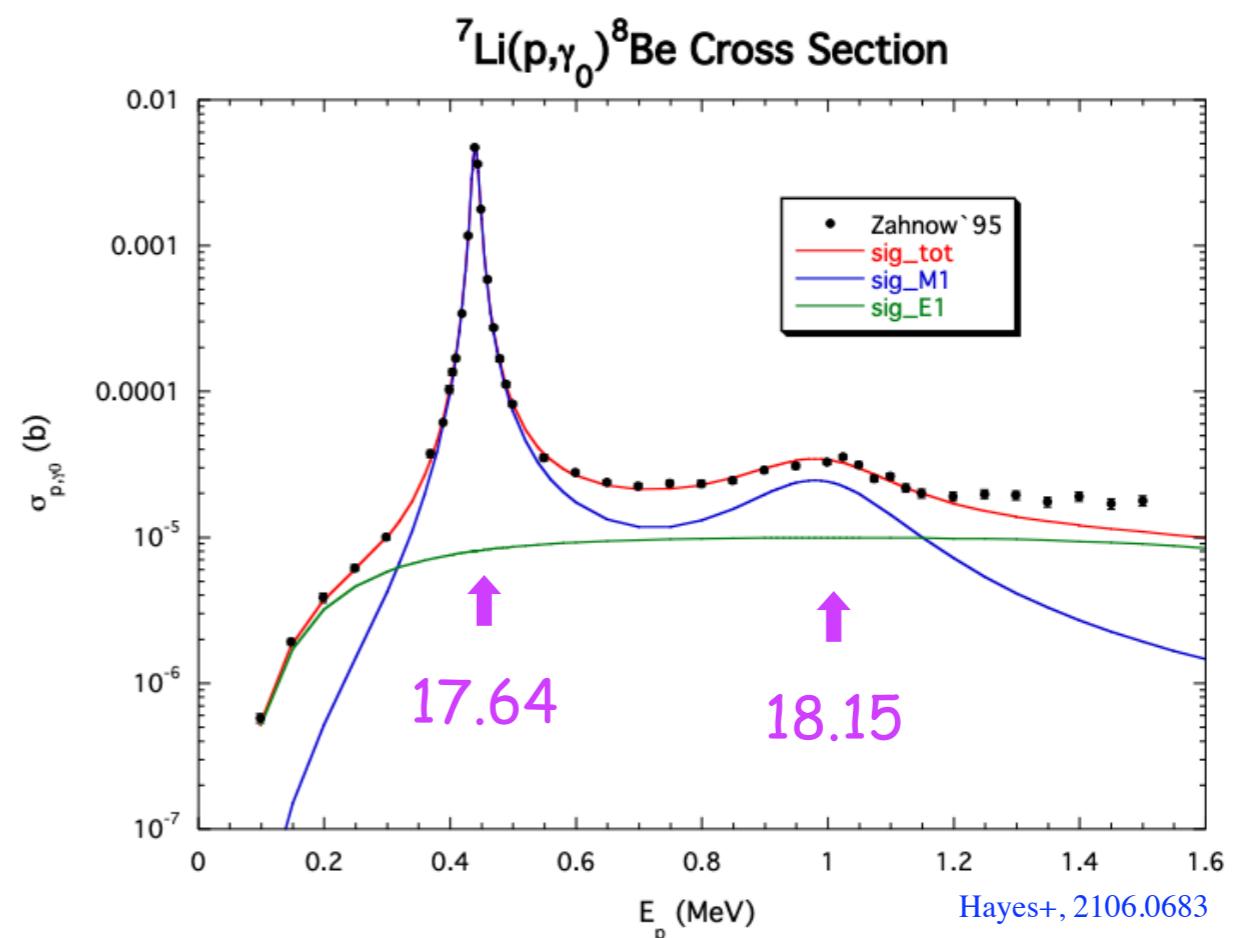
* - states of mixed isospin				

Resonant transition $p + {}^7\text{Li} \rightarrow {}^8\text{Be}^* \rightarrow \dots$

Radiative $p + {}^7\text{Li} \rightarrow {}^8\text{Be} + \gamma$



M1 resonant transition - E1 direct p capture
(valid also for a Vector X_{17})



Atomki results for ${}^8\text{Be}$

[PRL 116, 042501 (2016)]

Atomki results for ${}^8\text{Be}$

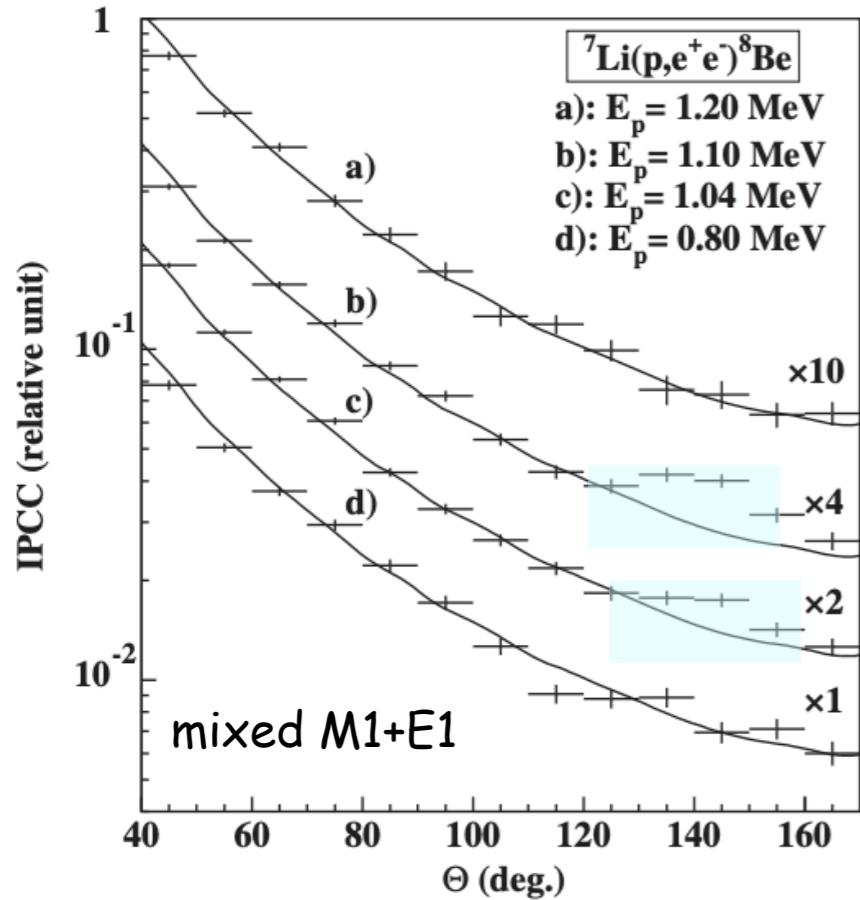
[PRL 116, 042501 (2016)]

${}^8\text{Be}^*(18.15\text{MeV}) \text{ IsoS}$

Atomki results for ${}^8\text{Be}$

[PRL 116, 042501 (2016)]

Angular correlation



${}^8\text{Be}^*(18.15\text{MeV}) \text{ IsoS}$

Energy gate: $E_{\pm} > 18 \text{ MeV}$

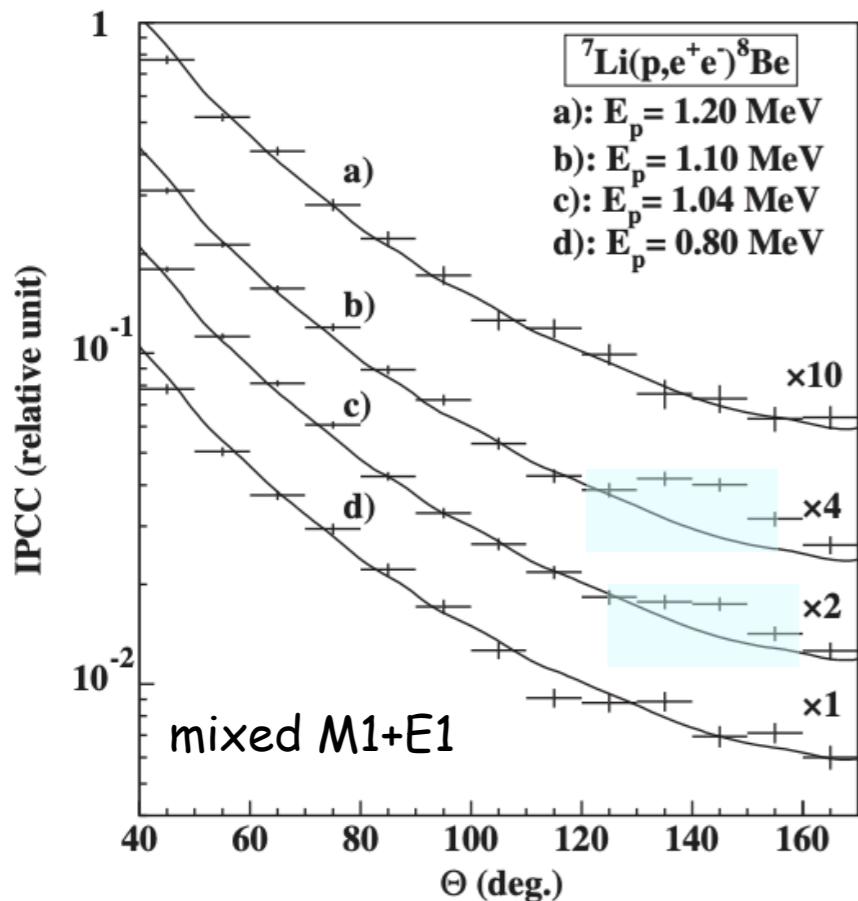
$$\gamma = \Delta E_{\pm} / E_{\pm} < 0.5$$



Atomki results for ${}^8\text{Be}$

[PRL 116, 042501 (2016)]

Angular correlation



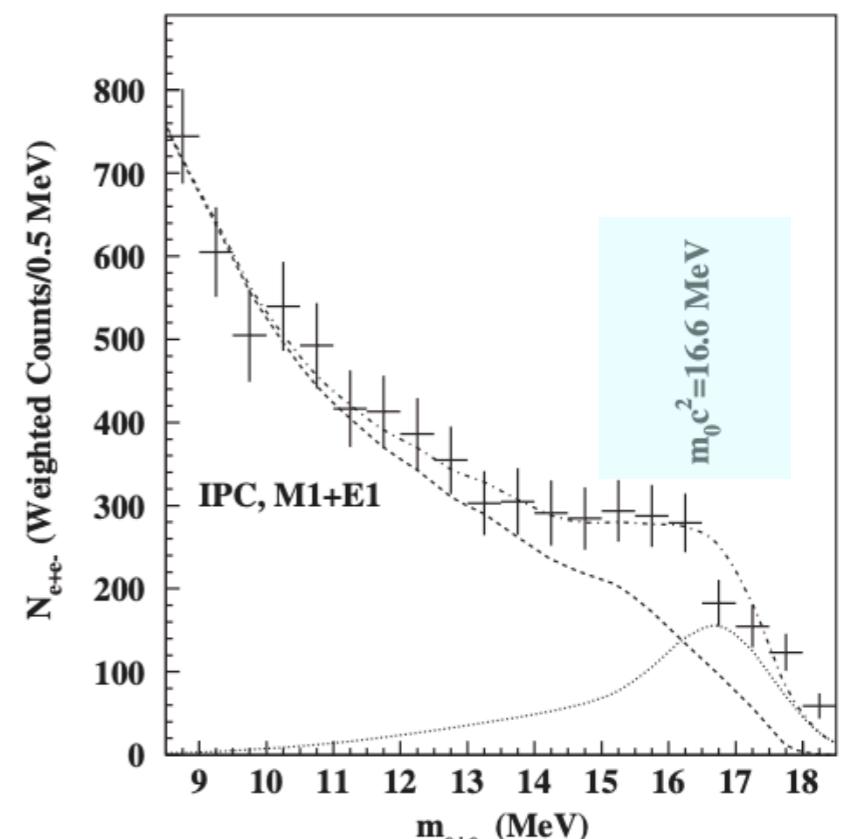
${}^8\text{Be}^*(18.15\text{MeV}) \text{ IsoS}$

Energy gate: $E_{\pm} > 18 \text{ MeV}$

$$\gamma = \Delta E_{\pm} / E_{\pm} < 0.5$$

$$m_{\pm}^2 \approx (1 - \gamma^2) E_{\pm}^2 \sin^2 \theta/2$$

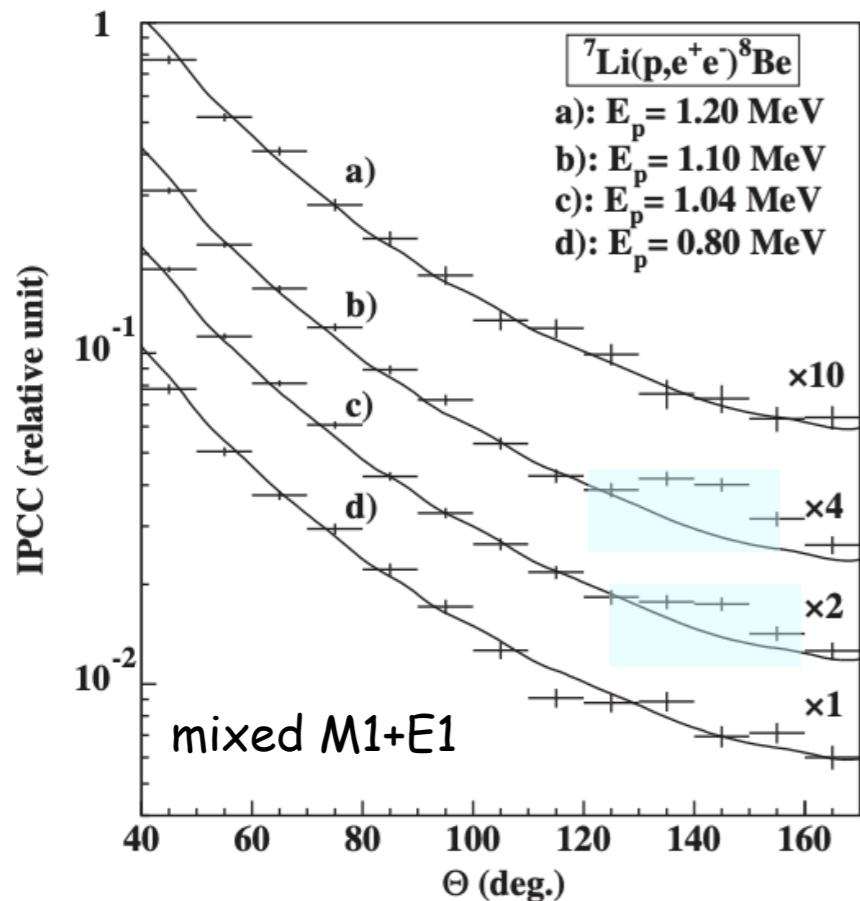
Invariant mass distribution



Atomki results for ${}^8\text{Be}$

[PRL 116, 042501 (2016)]

Angular correlation



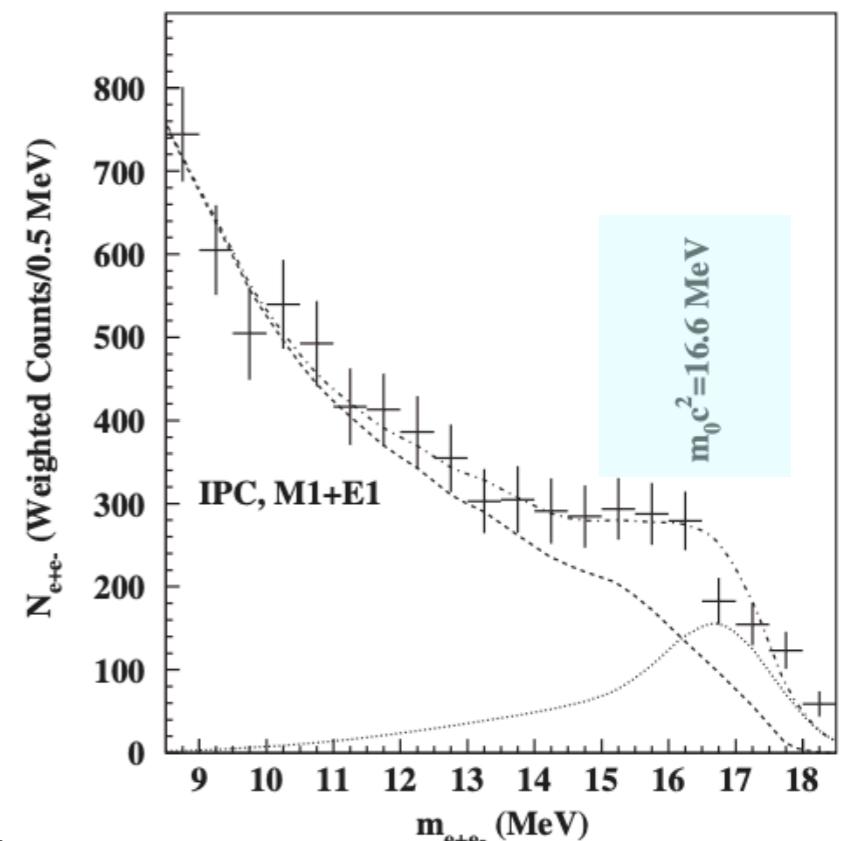
${}^8\text{Be}^*(18.15\text{MeV}) \text{ IsoS}$

Energy gate: $E_{\pm} > 18 \text{ MeV}$

$$\gamma = \Delta E_{\pm} / E_{\pm} < 0.5$$

$$m_{\pm}^2 \approx (1 - \gamma^2) E_{\pm}^2 \sin^2 \theta/2$$

Invariant mass distribution

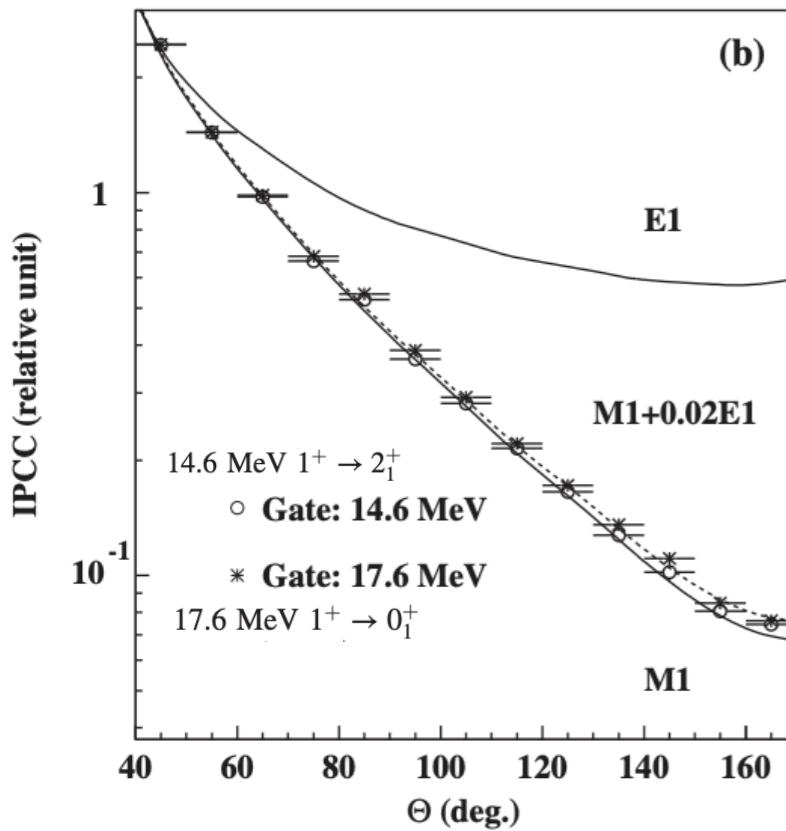
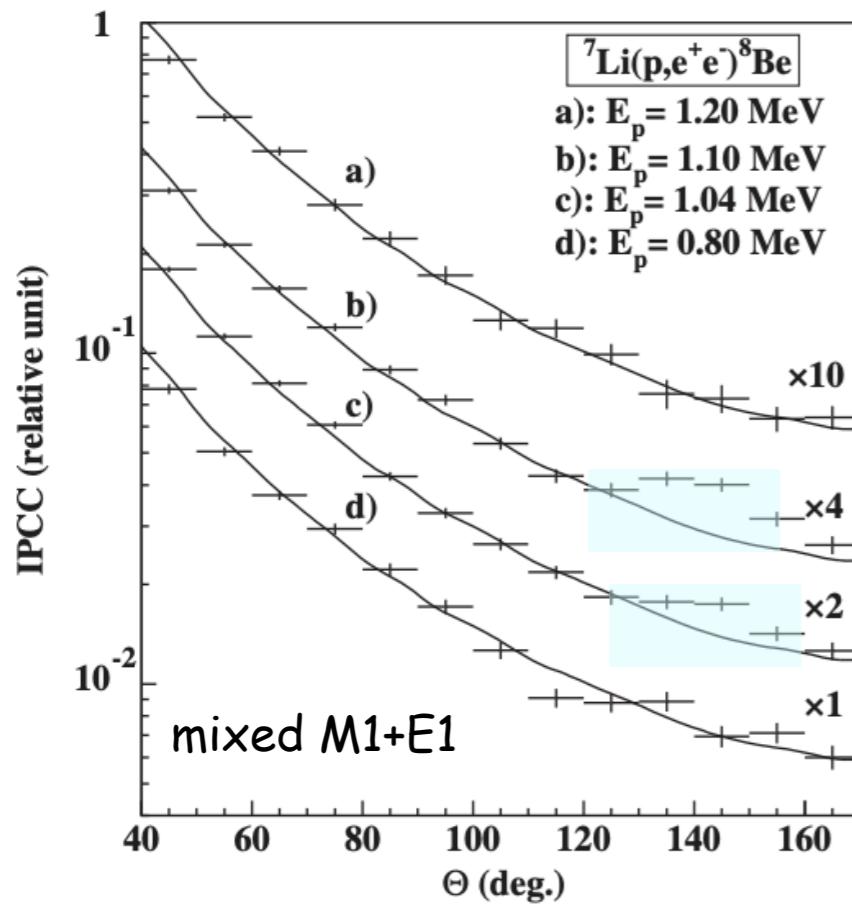


${}^8\text{Be}^*(17.64\text{MeV}) \text{ IsoV}$

Atomki results for ${}^8\text{Be}$

[PRL 116, 042501 (2016)]

Angular correlation



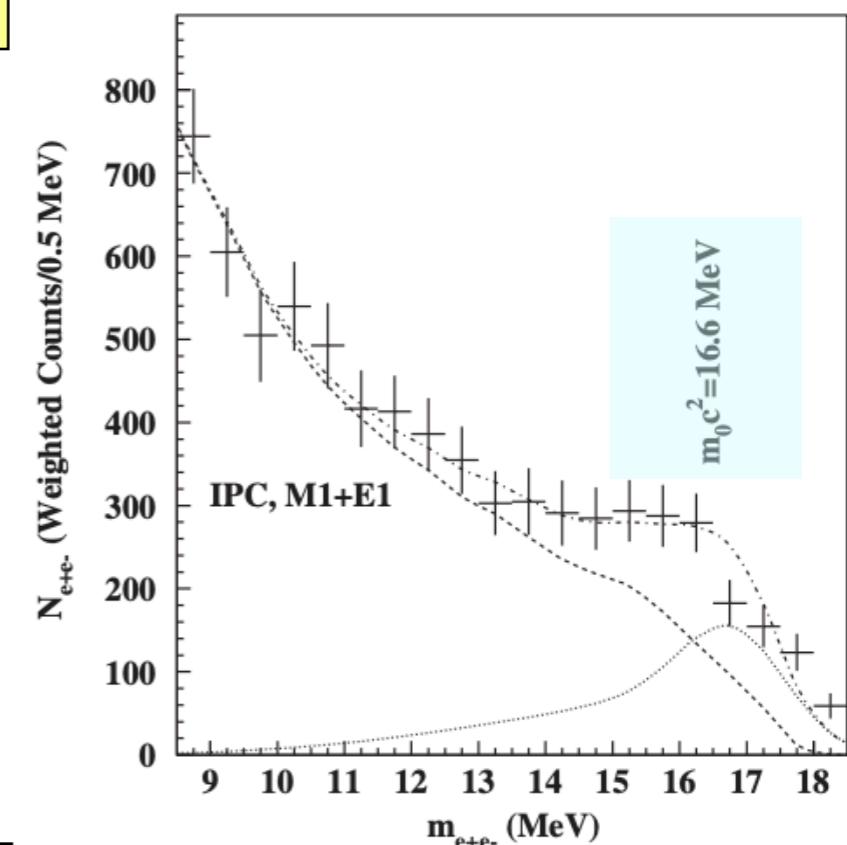
${}^8\text{Be}^*(18.15\text{MeV}) \text{ IsoS}$

Energy gate: $E_{\pm} > 18 \text{ MeV}$

$$\gamma = \Delta E_{\pm} / E_{\pm} < 0.5$$

$$m_{\pm}^2 \approx (1 - \gamma^2) E_{\pm}^2 \sin^2 \theta/2$$

Invariant mass distribution



${}^8\text{Be}^*(17.64\text{MeV}) \text{ IsoV}$

The contribution of the direct capture depends on the target thickness if the energy loss of the beam in the target is larger than the width of the resonance. The dashed simulated curve in Fig. 1(b) is obtained by fitting a small (2.0%) E1 contribution to the dominant M1 one, which describes the experimental data reasonably well.

One important theoretical input

[Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

As noted above, the decay ${}^8\text{Be}^* \rightarrow {}^8\text{Be} X$ is not seen. The protophobic gauge boson can mediate isovector transitions, so there is no dynamical suppression of this decay. However, its mass is near the 17.64 MeV threshold, so the decay is kinematically suppressed. For $m_X = 17.0$ (17.4) MeV, the $|\vec{p}_X|^3/|\vec{p}_\gamma|^3$ phase space suppression factor is 2.3 (5.2) times more severe for the ${}^8\text{Be}^*$ decay than for the ${}^8\text{Be}^*$ decay. In particular,

If the observed anomaly in ${}^8\text{Be}^*$ decays originates from a new particle, then the absence of new particle creation in the ${}^8\text{Be}^*$ decay combined with the isospin mixing discussed in Sec. IV strongly suggest that such decays are kinematically—not dynamically—suppressed and that the new particle mass is in the upper part of the range given in Eq. (1). It also suggests that with more data, a similar, but more phase space-suppressed, excess may appear in the IPC decays of the 17.64 state.

Feng+, PRD 1608.03591 [hep-ph];

Feng+, PRL 1604.07411 [hep-ph];

One important theoretical input

[Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

As noted above, the decay ${}^8\text{Be}^* \rightarrow {}^8\text{Be} X$ is not seen. The protophobic gauge boson can mediate isovector transitions, so there is no dynamical suppression of this decay. However, its mass is near the 17.64 MeV threshold, so the decay is kinematically suppressed. For $m_X = 17.0$ (17.4) MeV, the $|\vec{p}_X|^3/|\vec{p}_\gamma|^3$ phase space suppression factor is 2.3 (5.2) times more severe for the ${}^8\text{Be}^*$ decay than for the ${}^8\text{Be}^*$ decay. In particular,

If the observed anomaly in ${}^8\text{Be}^*$ decays originates from a new particle, then the absence of new particle creation in the ${}^8\text{Be}^*$ decay combined with the isospin mixing discussed in Sec. IV strongly suggest that such decays are kinematically—not dynamically—suppressed and that the new particle mass is in the upper part of the range given in Eq. (1). It also suggests that with more data, a similar, but more phase space-suppressed, excess may appear in the IPC decays of the 17.64 state.

Feng+, PRD 1608.03591 [hep-ph];

Feng+, PRL 1604.07411 [hep-ph];

New Atomki results for ${}^8\text{Be}^*(17.64)$

our experimental setup has been moved to a new accelerator laboratory and has also been improved.

we observed some smaller deviation also for the 17.6 MeV transition as was predicted by Feng et al.,
but which we did not see before

One important theoretical input

[Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

As noted above, the decay ${}^8\text{Be}^* \rightarrow {}^8\text{Be} X$ is not seen. The protophobic gauge boson can mediate isovector transitions, so there is no dynamical suppression of this decay. However, its mass is near the 17.64 MeV threshold, so the decay is kinematically suppressed. For $m_X = 17.0$ (17.4) MeV, the $|\vec{p}_X|^3/|\vec{p}_\gamma|^3$ phase space suppression factor is 2.3 (5.2) times more severe for the ${}^8\text{Be}^*$ decay than for the ${}^8\text{Be}^*$ decay. In particular,

If the observed anomaly in ${}^8\text{Be}^*$ decays originates from a new particle, then the absence of new particle creation in the ${}^8\text{Be}^*$ decay combined with the isospin mixing discussed in Sec. IV strongly suggest that such decays are kinematically—not dynamically—suppressed and that the new particle mass is in the upper part of the range given in Eq. (1). It also suggests that with more data, a similar, but more phase space-suppressed, excess may appear in the IPC decays of the 17.64 state.

Feng+, PRD 1608.03591 [hep-ph];

Feng+, PRL 1604.07411 [hep-ph];

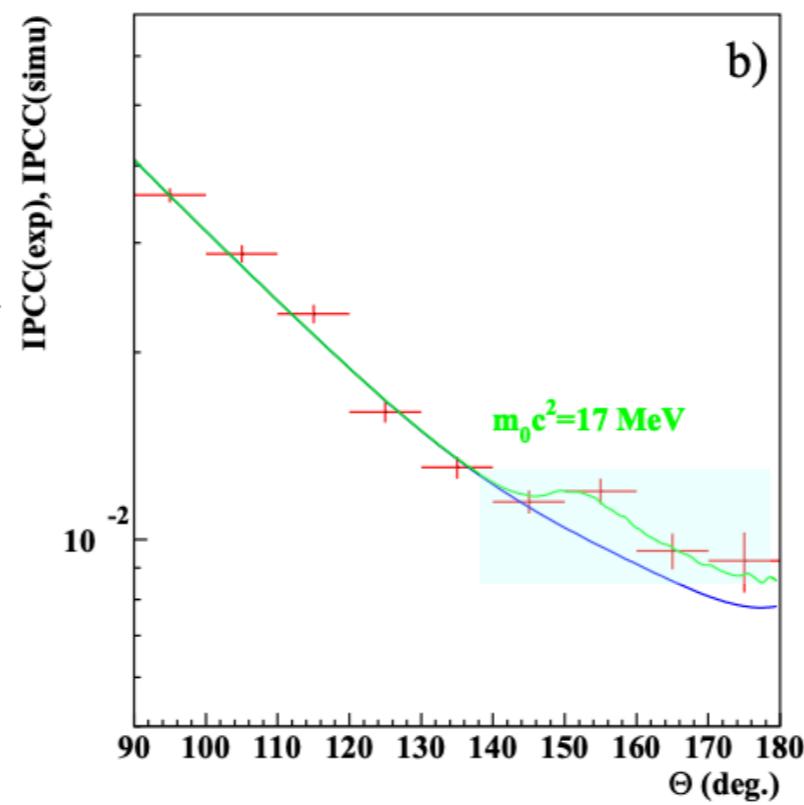
New Atomki results for ${}^8\text{Be}^*(17.64)$

our experimental setup has been moved to a new accelerator laboratory and has also been improved.

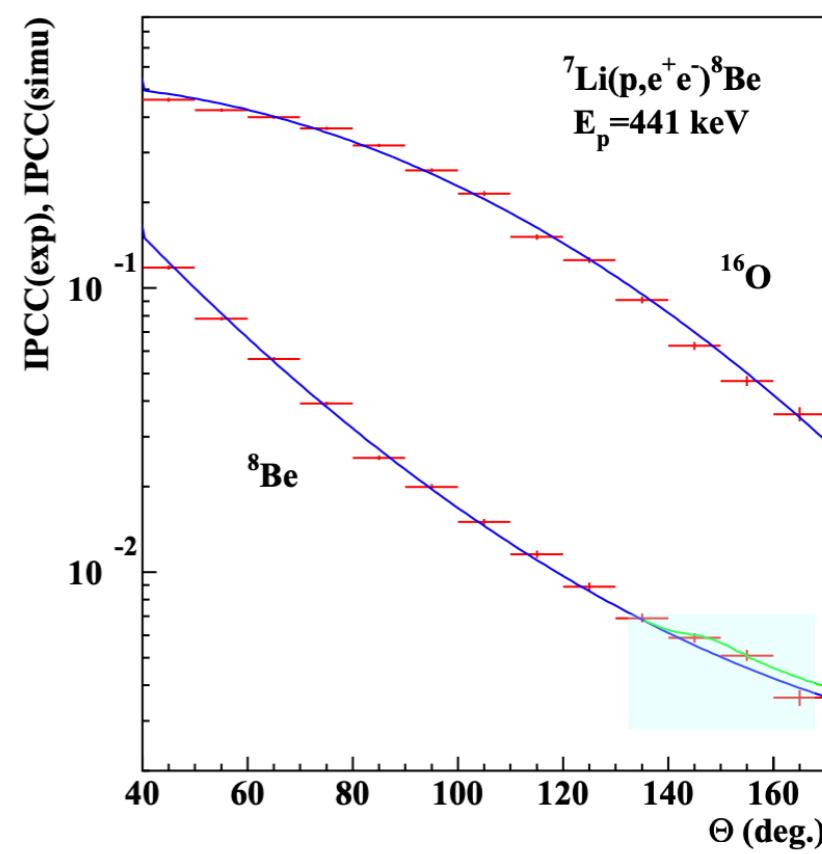
we observed some smaller deviation also for the 17.6 MeV transition as was predicted by Feng et al.,
but which we did not see before

Bump location:
 150° (17.64 MeV) vs.
 140° (18.15 MeV)

Messina symposium (Oct 2016)



Bormio meeting (Jan 2017)



One important theoretical input

[Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

As noted above, the decay ${}^8\text{Be}^* \rightarrow {}^8\text{Be}X$ is not seen. The protophobic gauge boson can mediate isovector transitions, so there is no dynamical suppression of this decay. However, its mass is near the 17.64 MeV threshold, so the decay is kinematically suppressed. For $m_X = 17.0$ (17.4) MeV, the $|\vec{p}_X|^3/|\vec{p}_\gamma|^3$ phase space suppression factor is 2.3 (5.2) times more severe for the ${}^8\text{Be}^*$ decay than for the ${}^8\text{Be}^*$ decay. In particular,

If the observed anomaly in ${}^8\text{Be}^*$ decays originates from a new particle, then the absence of new particle creation in the ${}^8\text{Be}^*$ decay combined with the isospin mixing discussed in Sec. IV strongly suggest that such decays are kinematically—not dynamically—suppressed and that the new particle mass is in the upper part of the range given in Eq. (1). It also suggests that with more data, a similar, but more phase space-suppressed, excess may appear in the IPC decays of the 17.64 state.

Feng+, PRD 1608.03591 [hep-ph];

Feng+, PRL 1604.07411 [hep-ph];

New Atomki results for ${}^8\text{Be}^*(17.64)$

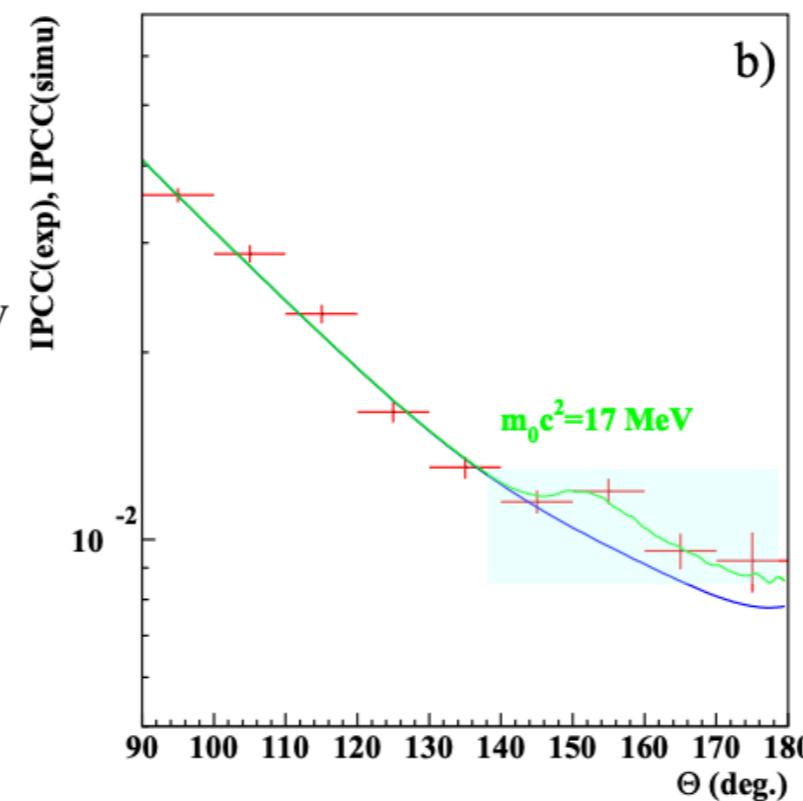
our experimental setup has been moved to a new accelerator laboratory and has also been improved.

we observed some smaller deviation also for the 17.6 MeV transition as was predicted by Feng et al.,
but which we did not see before

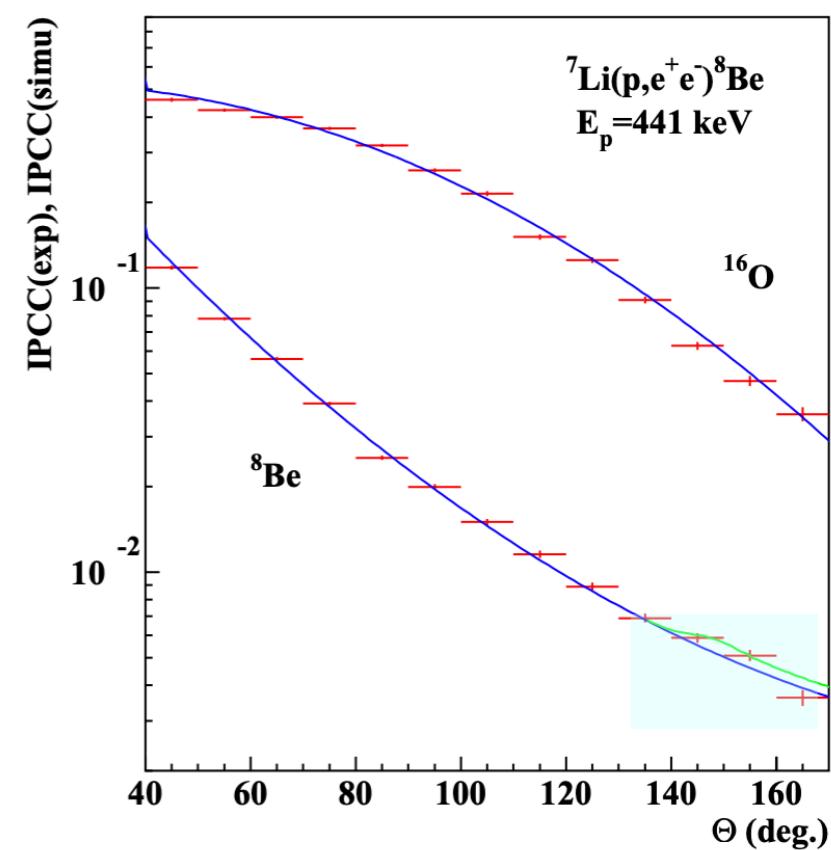
Bump location:
 150° (17.64 MeV) vs.
 140° (18.15 MeV)

Axial vector boson

Messina symposium (Oct 2016)



Bormio meeting (Jan 2017)



Calculation of relevant Nucl. Matrix Elements:

Kozaczuk+, PRD 1612.01525 [hep-ph]

the ${}^8\text{Be}^* \rightarrow {}^8\text{Be}+X$ transition rate can be suppressed

relative to that of the ${}^8\text{Be}^* \rightarrow {}^8\text{Be}+X$ mode for an axial vector. This effect is dynamical,

^8Be anomaly: Standard Model explanations ?

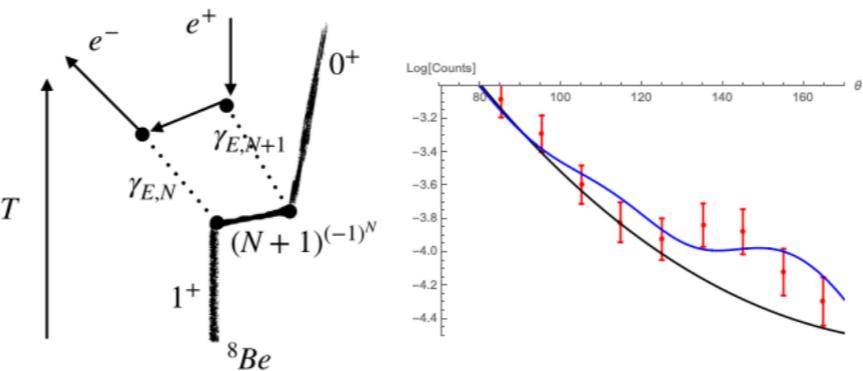
Zhang & Miller PLB, arXiv:1703.04588 [nucl-th]

Interferences between different multipoles. Possibility of using the nuclear transition form factor to explain the anomaly

We find that the model improvements are not able to explain the anomaly.

Koch, NPB, arXiv:2003.05722 [hep-ph]

Hypothesises nuclear chain reaction and conversion of two resulting highly energetic γ s into an electron-positron pair.



The kinematics fits perfectly the experimental result. No explanation for the isospin structure can be given. The process does not give a satisfying explanation of X17.

Kálmán & Keszthelyi EPJA, arXiv:2005.10643 [nucl-th]

Higher order processes, in which strong and electromagnetic interactions are coupled and govern jointly the system from the definite initial state to the definite final one
[Analyzed ^8Be and (qualitatively) also ^4He]

Enhancement can be generated by higher order processes. Lower energy nucl. transitions can cause peaked angle dependence in angular correlations.

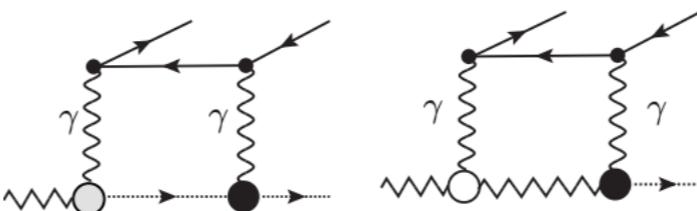
Zhang & Miller PLB, arXiv:2008.11288 [hep-ph]

Derived isospin relation between photon and (protophobic) X couplings to nucleons. X production dominated by direct transitions with a smooth energy dependence occurring for all proton beam energies above threshold

X bremsstrahlung occurs at all beam energies above threshold. The enhancement should have been seen at all four Atomki p-energies. The explanation of the anomaly in terms of protophobic vector boson cannot be correct.

Aleksejevs+, arXiv:2102.01127 [nucl-th]

Full second-order calculation of $^8\text{Be}^* \rightarrow ^8\text{Be} e^+e^-$ process:
interferences second-order corrections and the interference terms to the Born-level decay amplitudes



The observed ^8Be experimental structure can be reproduced within the Standard Model.

Hayes+, arXiv:2106.06834 [nucl-th]

Study of e^+e^- angular distributions for nuclear decay for several multipoles M1,E1 dominate, but the ratio of M1 to E1 strength strong function of energy (Atomki: M1/E1 assumed constant over the energy region $E_p = 0:8-1:2$ MeV)

The evidence of a new particle emitted from the 18.15 MeV resonance in ^8Be seems to be strongly dependent on the assumptions about the nuclear structure of this resonance. Atomki surplus events at large angles could be an artefact of the Atomki analysis nuclear structure assumptions.

About theoretical interpretation

[Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

About theoretical interpretation [Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

X₁₇ particle: Some simple possibilities are excluded:

Scalar: $J^P = 1^+({}^8Be^*) \rightarrow 0^+({}^8Be) 0^+(X) \Rightarrow L=1; P = +1 = (-1)^L$

Vector with no definite parity (Z'): APV constraints

$U(1)_{B-L}$ vector boson: ν -e scattering ($g_{B-L} \lesssim 10^{-5}$)

Kinetically mixed V': $g_f = \epsilon Q_f$ NA48/2 limit $\pi^0 \rightarrow X \gamma$

About theoretical interpretation [Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

X₁₇ particle: Some simple possibilities are excluded:

Scalar: $J^P = 1^+({}^8Be^*) \rightarrow 0^+({}^8Be) 0^+(X) \Rightarrow L=1; P = +1 = (-1)^L$

Vector with no definite parity (Z'): APV constraints

$U(1)_{B-L}$ vector boson: ν -e scattering ($g_{B-L} \lesssim 10^{-5}$)

Kinetically mixed V': $g_f = \epsilon Q_f$ NA48/2 limit $\pi^0 \rightarrow X \gamma$

Pionphobic/Protophobic vector particle interpretation:

About theoretical interpretation [Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

X₁₇ particle: Some simple possibilities are excluded:

Scalar: $J^P = 1^+({}^8Be^*) \rightarrow 0^+({}^8Be) 0^+(X) \Rightarrow L=1; P = +1 = (-1)^L$

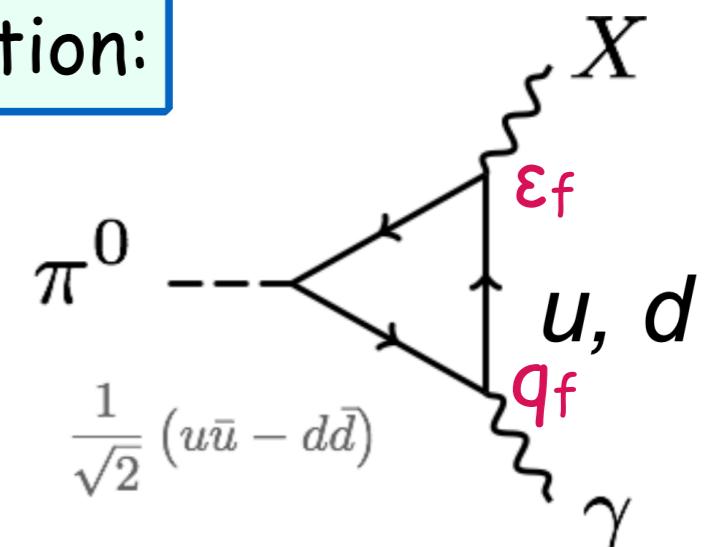
Vector with no definite parity (Z'): APV constraints

$U(1)_{B-L}$ vector boson: ν -e scattering ($g_{B-L} \lesssim 10^{-5}$)

Kinetically mixed V': $g_f = \epsilon Q_f$ NA48/2 limit $\pi^0 \rightarrow X \gamma$

Pionphobic/Protophobic vector particle interpretation:

$\pi^0 \rightarrow X \gamma : |2\epsilon_u + \epsilon_d| < 8 \times 10^{-4}$ (NA48/2)



About theoretical interpretation [Feng+, PRL 1604.07411 [hep-ph]; PRD 1608.03591 [hep-ph]]

X₁₇ particle: Some simple possibilities are excluded:

Scalar: $J^P = 1^+({}^8Be^*) \rightarrow 0^+({}^8Be) 0^+(X) \Rightarrow L=1; P = +1 = (-1)^L$

Vector with no definite parity (Z'): APV constraints

U(1)_{B-L} vector boson: ν-e scattering ($g_{B-L} \lesssim 10^{-5}$)

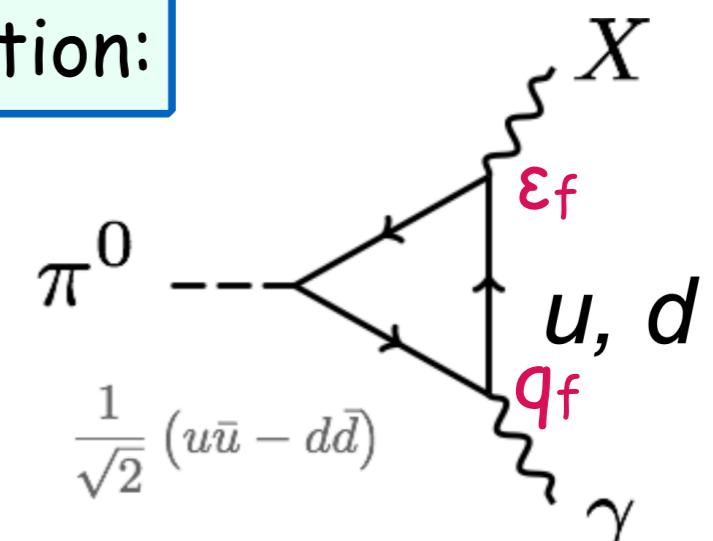
Kinetically mixed V': $g_f = \epsilon Q_f$ NA48/2 limit $\pi^0 \rightarrow X \gamma$

Pionphobic/Protophobic vector particle interpretation:

$$\pi^0 \rightarrow X \gamma : |2\epsilon_u + \epsilon_d| < \underline{8 \times 10^{-4}} \quad (\text{NA48/2})$$

$$B_X/B_\gamma \propto (\epsilon_p + \epsilon_n)^2 (p_X/p_\gamma)^3 \approx 6 \times 10^{-6} \quad (\text{Atomki})$$

$$\Rightarrow |\epsilon_u + \epsilon_d| \approx \underline{4 \times 10^{-3}}$$



$$\epsilon_d \approx -2 \epsilon_u \quad (\pm 10\%) \quad \Rightarrow \quad \epsilon_p = 2\epsilon_u + \epsilon_d \approx 0; \quad \epsilon_n = 2\epsilon_d + \epsilon_u \approx 1.2 \times 10^{-2}$$

[Feng+, 1608.0359 [hep-ph] (Aug. 2016)]

For protophobic vector, ${}^8\text{Be}$ data can be explained with:

$$\varepsilon_u = -\varepsilon_n/3 \approx \pm 3.7 \times 10^{-3}; \quad \varepsilon_d = 2\varepsilon_n/3 \approx \mp 7.4 \times 10^{-3}; \quad |\varepsilon_e| \in [2-14] \times 10^{-4}$$

[Feng+, 1608.0359 [hep-ph] (Aug. 2016)]

For protophobic vector, ${}^8\text{Be}$ data can be explained with:

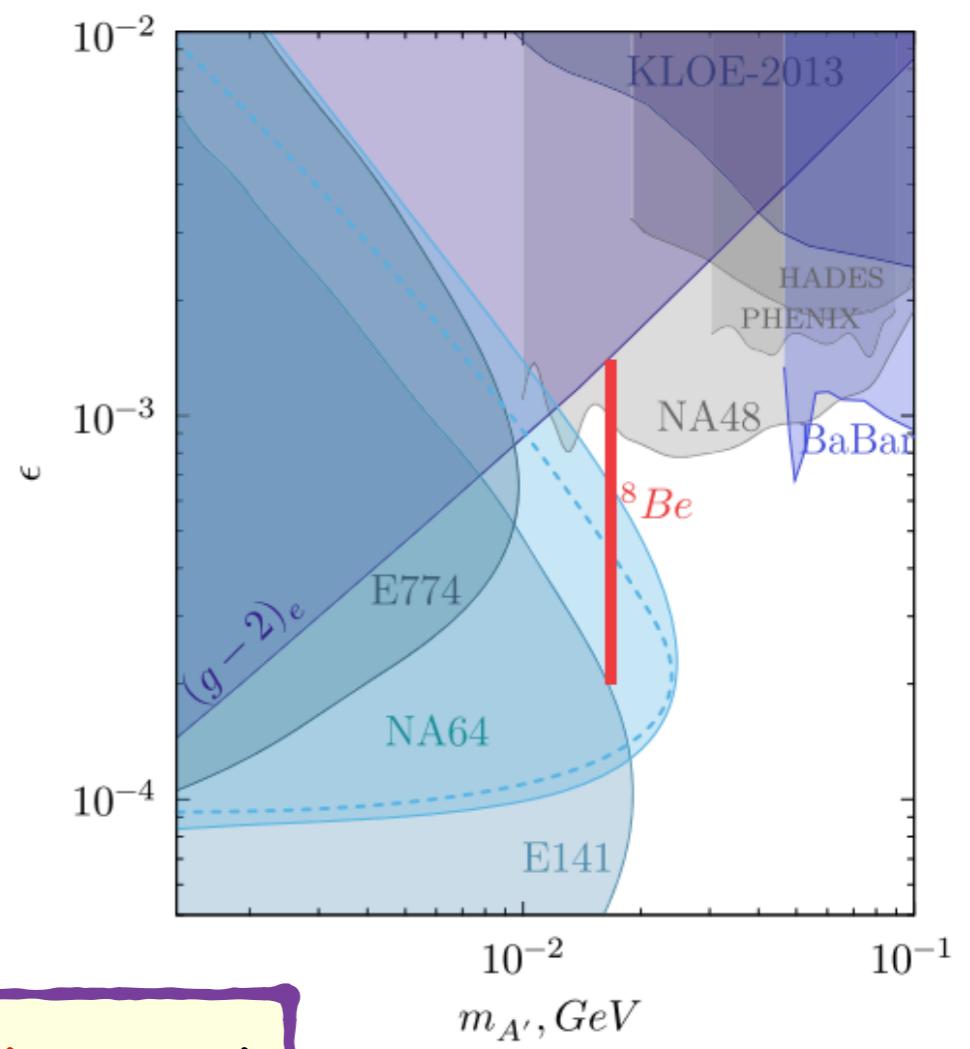
$$\varepsilon_u = -\varepsilon_n/3 \approx \pm 3.7 \times 10^{-3}; \quad \varepsilon_d = 2\varepsilon_n/3 \approx \mp 7.4 \times 10^{-3}; \quad |\varepsilon_e| \in [2-14] \times 10^{-4}$$

Current limits on X17

[NA64@ CERN, 1912.11389 [hep-ex] (Dec. 2019)]

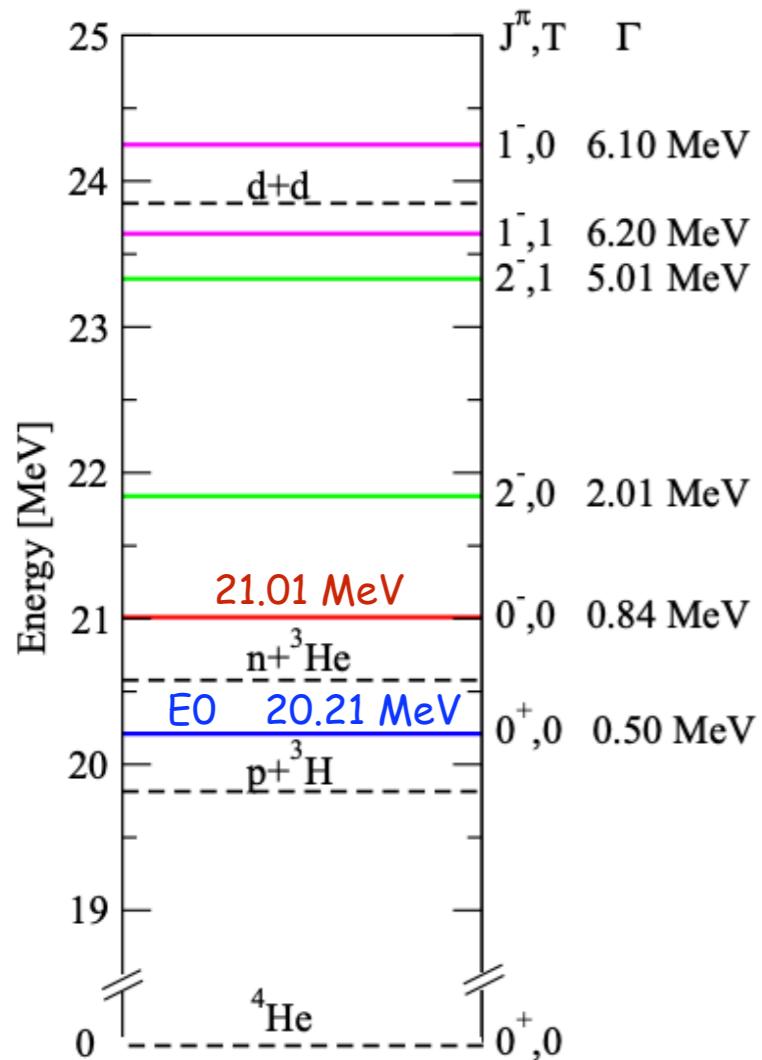
The X17 boson could be produced in the bremsstrahlung reaction $e^- Z \rightarrow e^- Z X$ by a high energy beam (150 GeV) of electrons incident on the active target in the NA64 experiment, and observed through its decay $X \rightarrow e^+e^-$

$$|\varepsilon_e| \in [2.0-6.8] \times 10^{-4} \text{ for } M_X = 16.7 \text{ MeV}$$

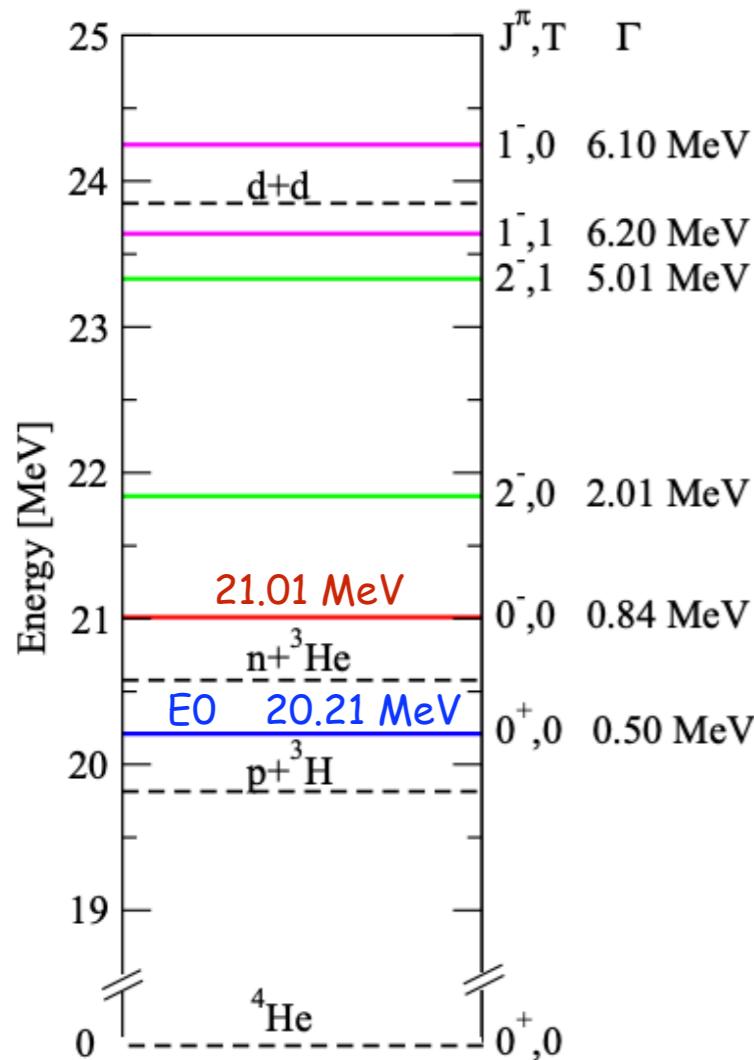


(In the meanwhile: $M_X({}^8\text{Be}) = (17.1 \pm 0.16) \text{ MeV}$)

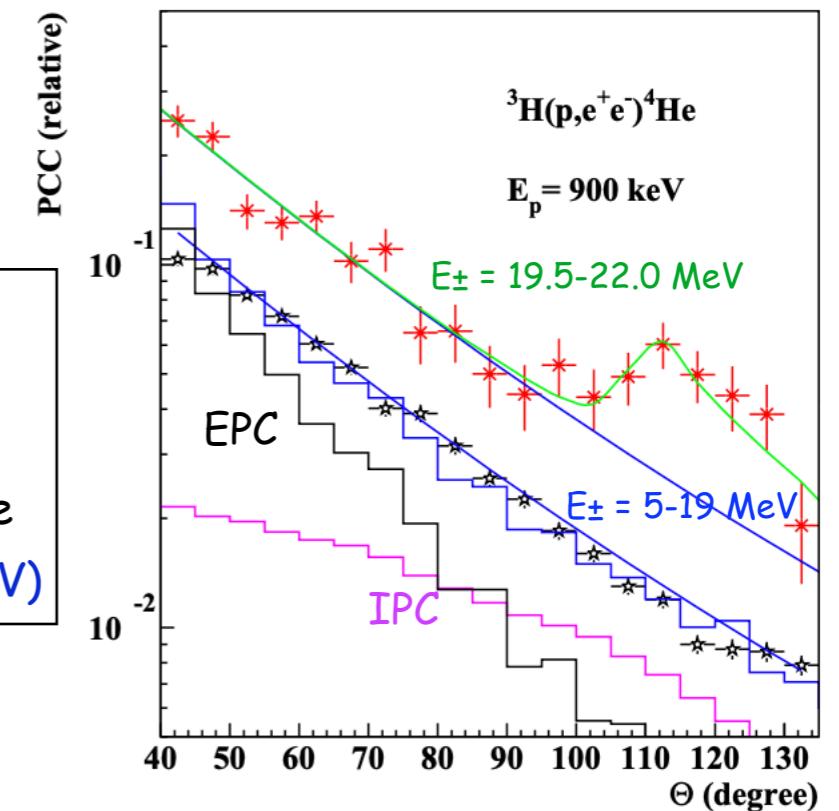
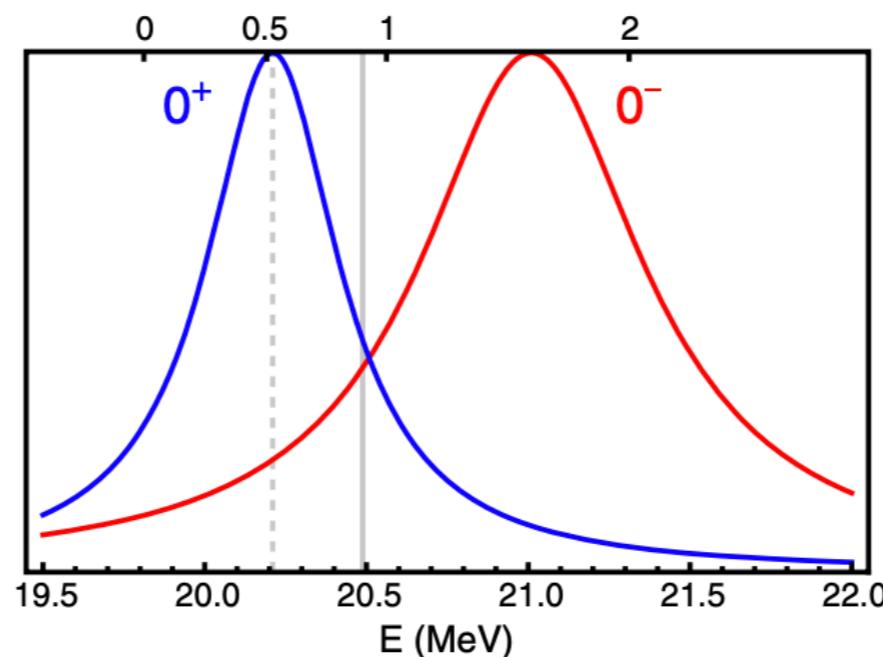
Helium 4 nuclear transitions



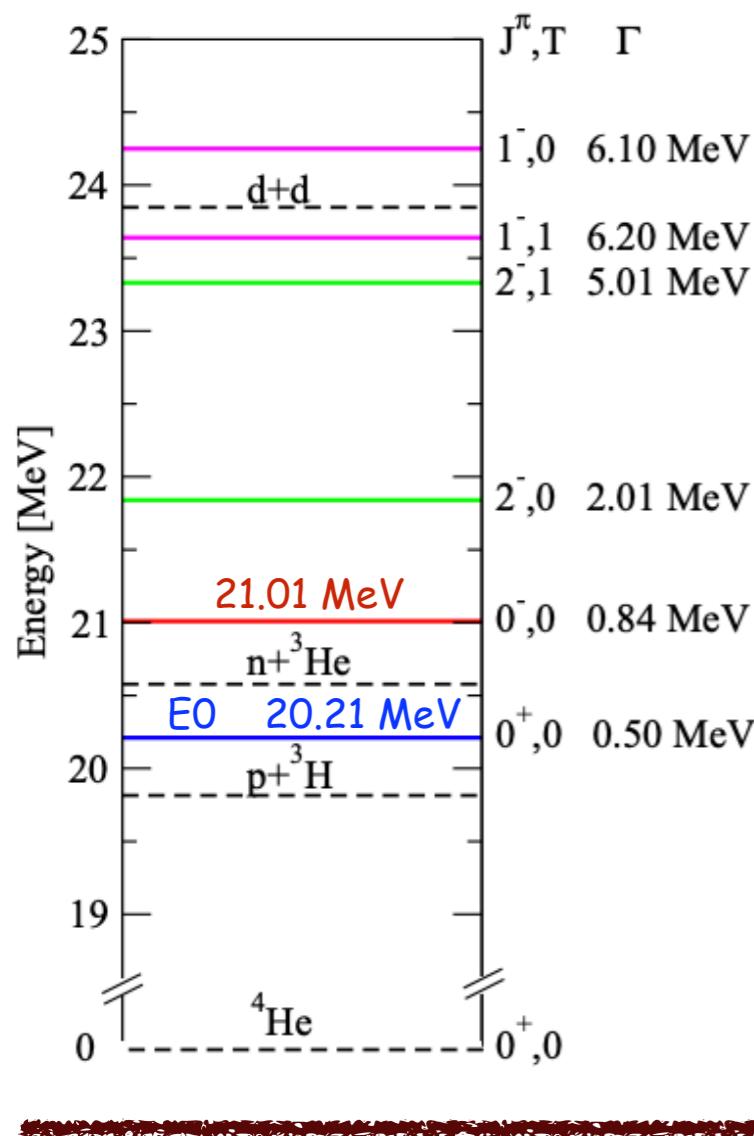
Helium 4 nuclear transitions



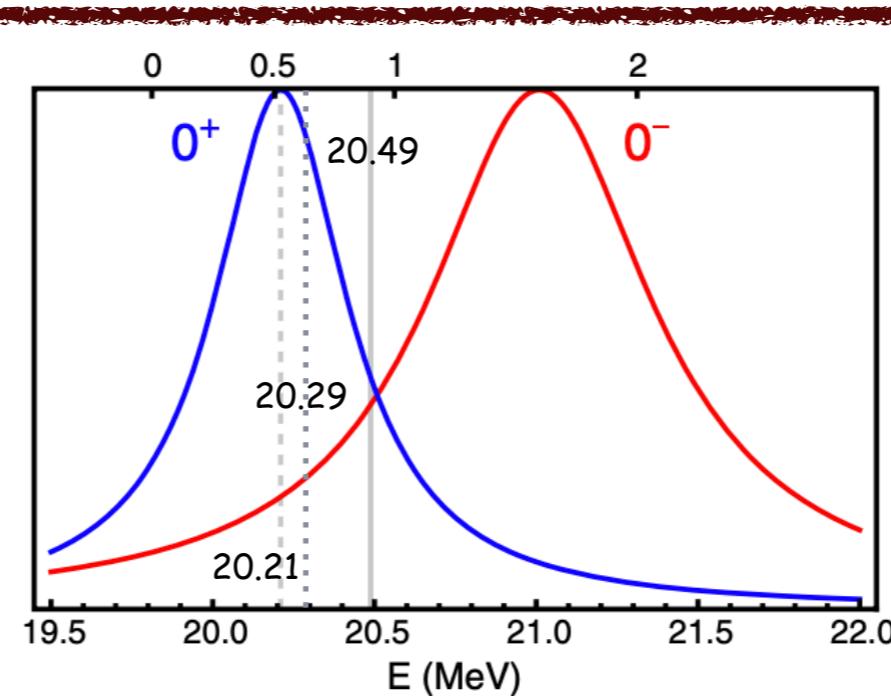
[arXiv:1910.10459](https://arxiv.org/abs/1910.10459) [nucl-ex] $E_p=900 \text{ keV}$
 (below $E(p,n) = 1.018 \text{ keV}$ threshold)
 excites the ${}^4\text{He}$ nucleus to $E^*=20.49 \text{ MeV}$
 and populates the second ${}^4\text{He}$ excited state
 $0^- (21.01 \text{ MeV})$ overlapped with $0^+ (20.21 \text{ MeV})$



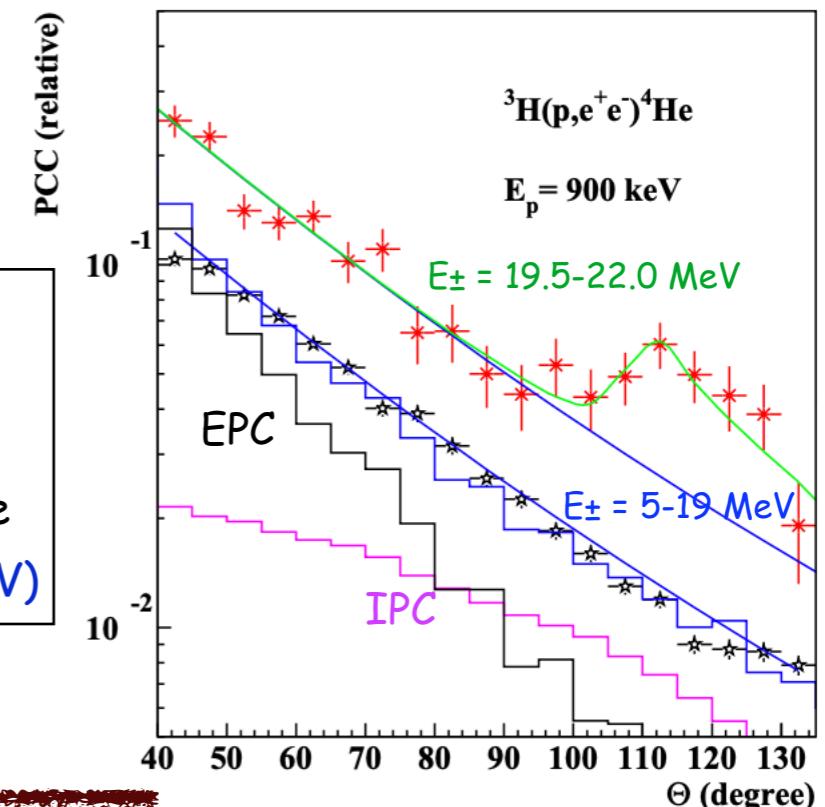
Helium 4 nuclear transitions



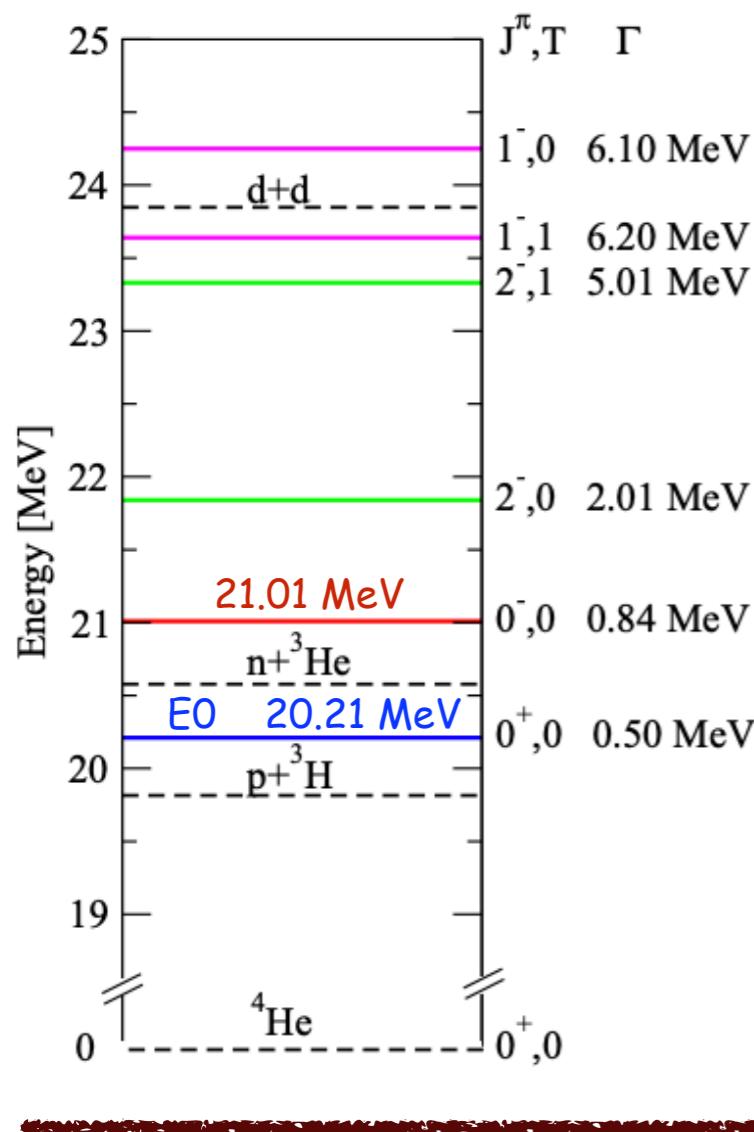
[arXiv:1910.10459](https://arxiv.org/abs/1910.10459) [nucl-ex] $E_p=900 \text{ keV}$
 (below $E(p,n) = 1.018 \text{ keV}$ threshold)
 excites the ${}^4\text{He}$ nucleus to $E^*=20.49 \text{ MeV}$
 and populates the second ${}^4\text{He}$ excited state
 $0^- (21.01 \text{ MeV})$ overlapped with $0^+ (20.21 \text{ MeV})$



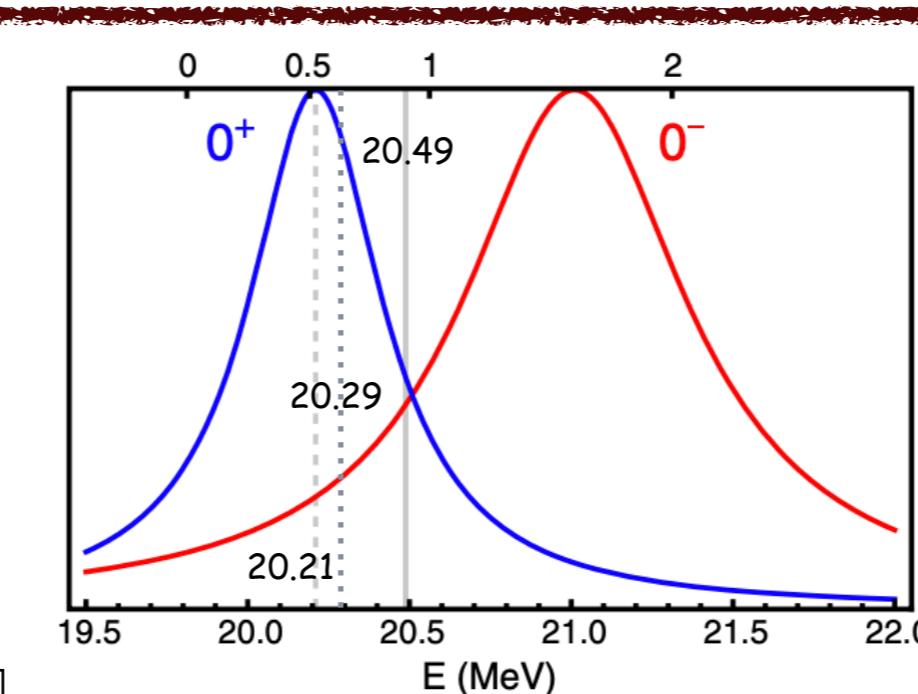
PRC (2021) [[arXiv:2104.1075](https://arxiv.org/abs/2104.1075) [nucl-ex]]
 $E_p = 510, 610, 900 \text{ keV}$ to induce direct & resonant radiative capture ${}^3\text{H} (\text{p},\gamma) {}^4\text{He}$ and populate the overlapping 1st 0^+ and 2nd 0^- ${}^4\text{He}$ excited states



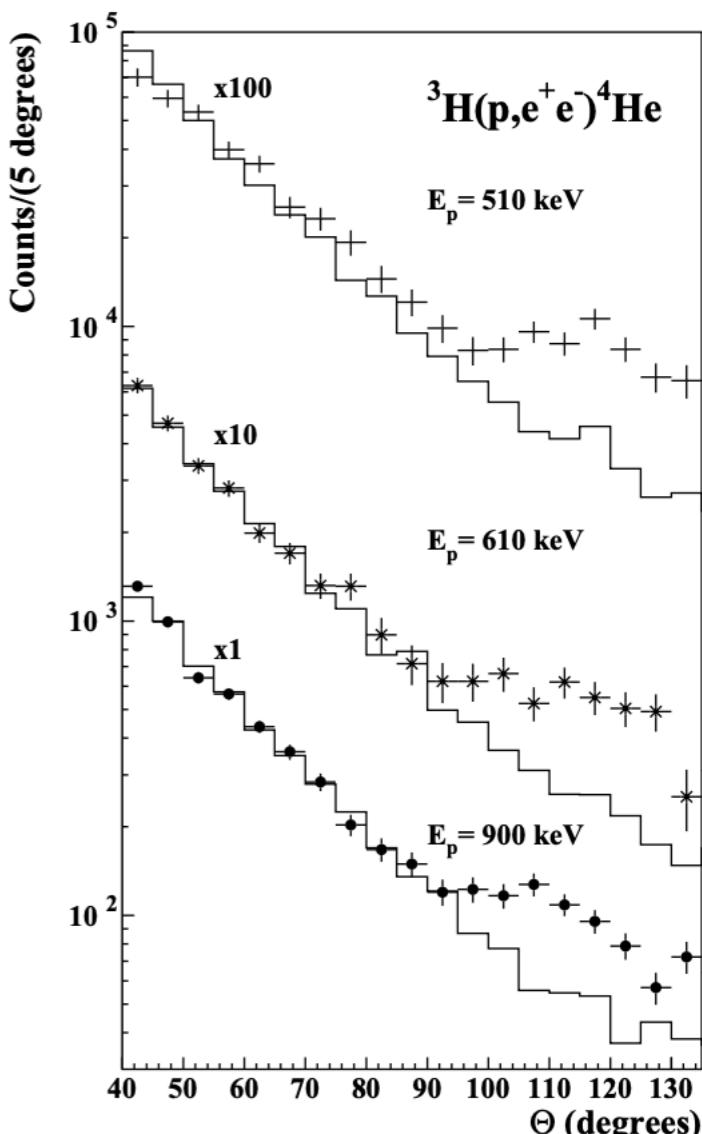
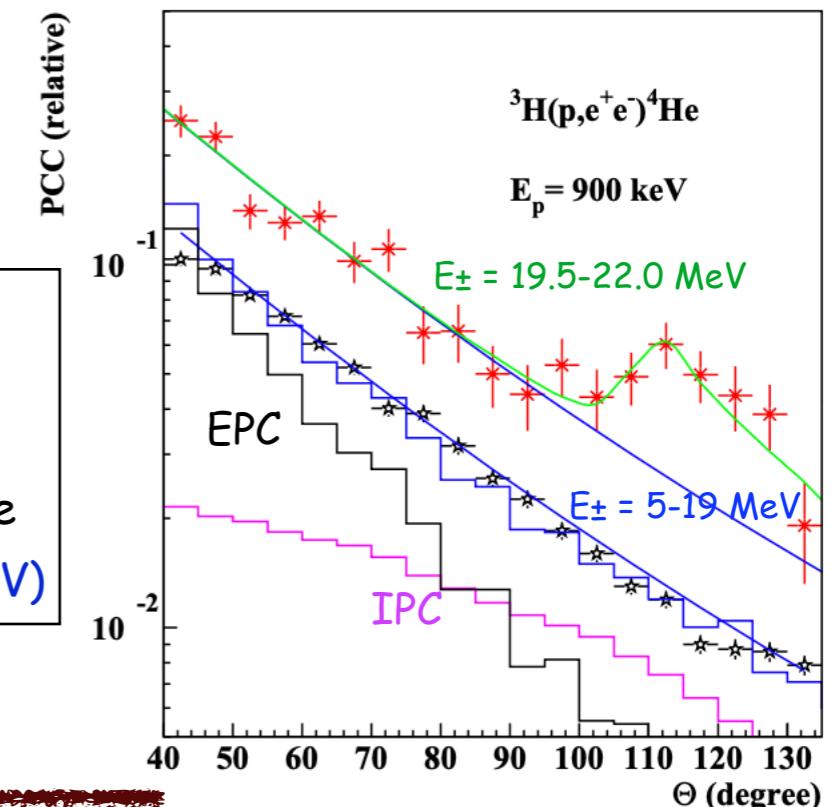
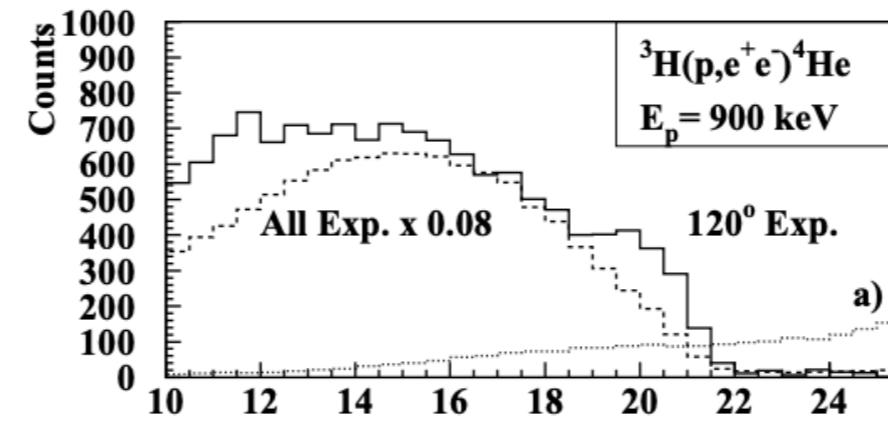
Helium 4 nuclear transitions



arXiv:1910.10459 [nucl-ex] $E_p=900$ keV
 (below $E(p,n) = 1.018$ keV threshold)
 excites the ${}^4\text{He}$ nucleus to $E^*=20.49$ MeV
 and populates the second ${}^4\text{He}$ excited state
 0^- (21.01 MeV) overlapped with 0^+ (20.21 MeV)



PRC (2021) [arXiv:2104.1075 [nucl-ex]]
 $E_p = 510, 610, 900$ keV to induce direct & resonant radiative capture ${}^3\text{H}(\text{p},\gamma){}^4\text{He}$ and populate the overlapping 1st 0^+ and 2nd 0^- ${}^4\text{He}$ excited states



^4He anomaly: Standard Model explanations ?

^4He anomaly: Standard Model explanations ?

The X17 boson and the $^3\text{H}(\text{p},\text{e}^+\text{e}^-)^4\text{He}$ and $^3\text{He}(\text{n},\text{e}^+\text{e}^-)^4\text{He}$ processes: a theoretical analysis [Viviani+, PRD 2104.04808 [nucl-th]]

- Analysis of the process in the standard theory (ab initio nuc. phys. calculations)
- Study of how the exchange of $X_{17}(V,A,S,P)$ would impact such a process
- Beyond the resonance-saturation approach (justified for ^8Be but not for ^4He)
- Detailed study of the behaviour of the (V,A,S,P) induced angular correlations

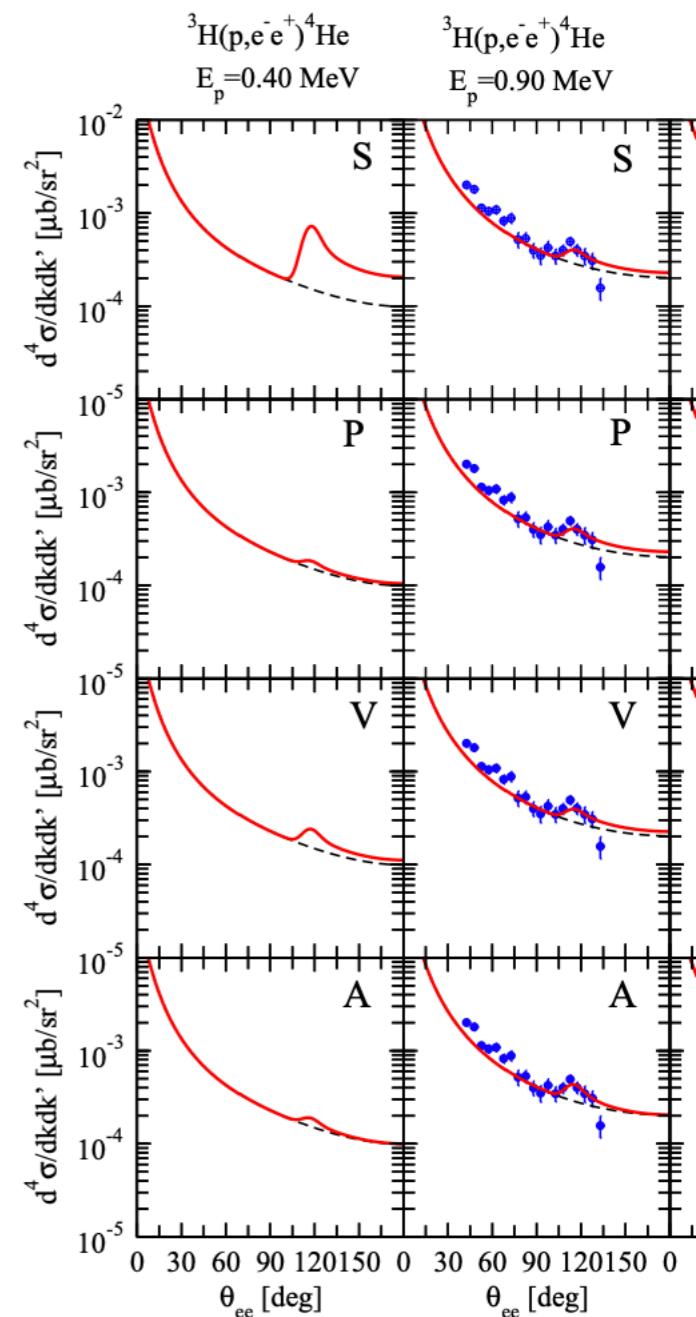
^4He anomaly: Standard Model explanations ?

The X17 boson and the $^3\text{H}(\text{p},\text{e}^+\text{e}^-)^4\text{He}$ and $^3\text{He}(\text{n},\text{e}^+\text{e}^-)^4\text{He}$ processes: a theoretical analysis [Viviani+, PRD 2104.04808 [nucl-th]]

- Analysis of the process in the standard theory (ab initio nuccl. phys. calculations)
- Study of how the exchange of $X_{17}(V,A,S,P)$ would impact such a process
- Beyond the resonance-saturation approach (justified for ^8Be but not for ^4He)
- Detailed study of the behaviour of the (V,A,S,P) induced angular correlations

Main results:

- The predicted cross sections are monotonically decreasing as function of the e^+e^- opening angle.
- Absence of any resonance-like structure
- Measurements at $\theta_{\text{vp}} \neq 90^\circ$ can discriminate $X=V,A,S,P$



$$M_x = 17 \text{ MeV}$$

$$\theta_{\text{vp}} = 90^\circ$$

^4He anomaly: Standard Model explanations ?

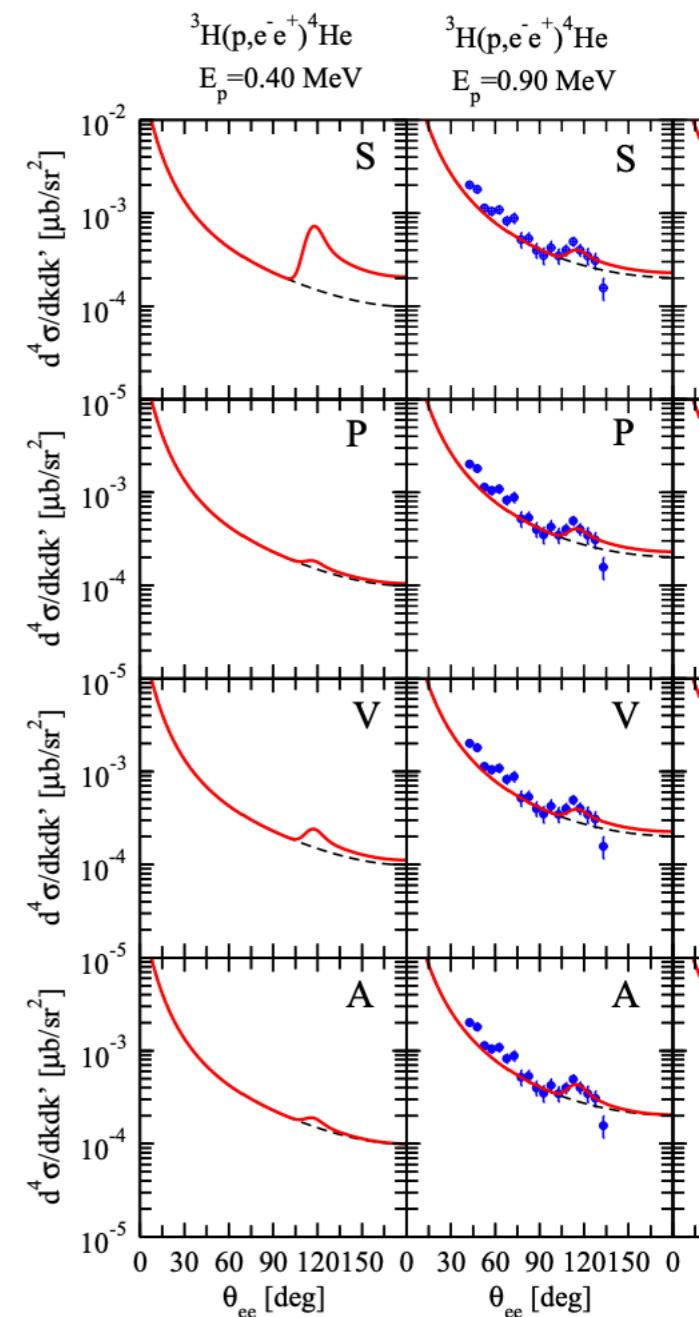
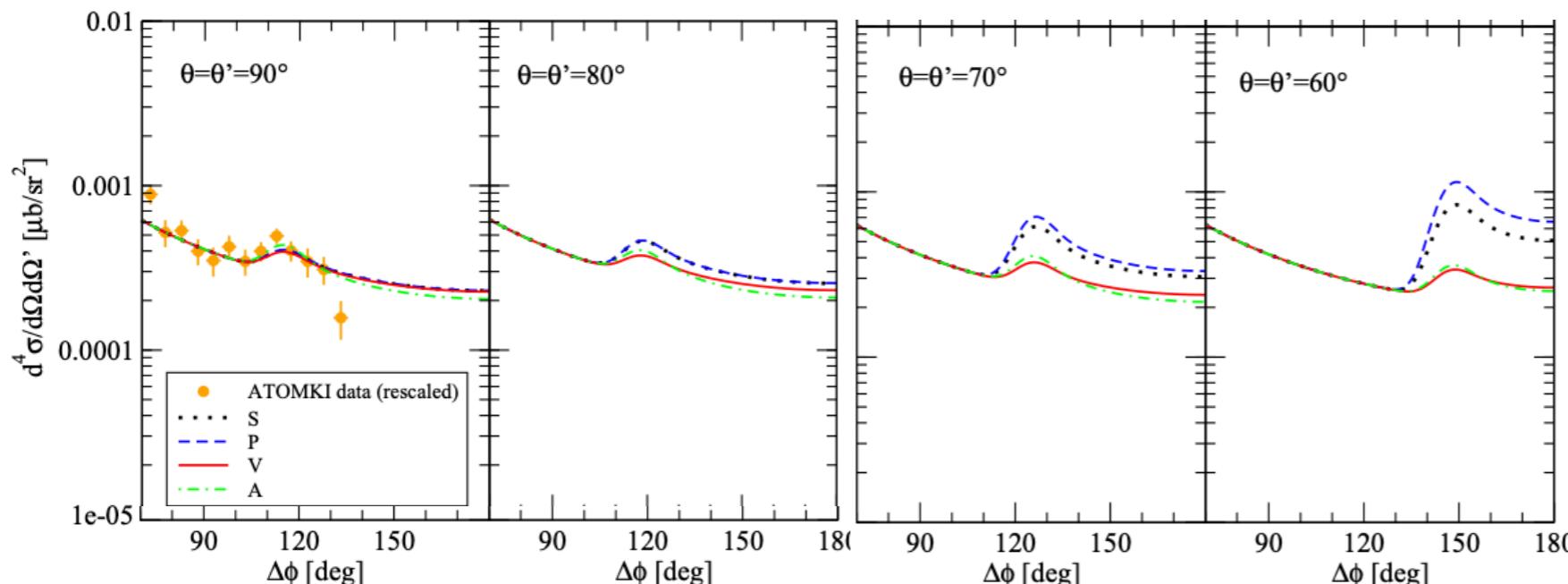
The X17 boson and the $^3\text{H}(\text{p},\text{e}^+\text{e}^-)^4\text{He}$ and $^3\text{He}(\text{n},\text{e}^+\text{e}^-)^4\text{He}$

processes: a theoretical analysis [Viviani+, PRD 2104.04808 [nucl-th]]

- Analysis of the process in the standard theory (ab initio nuccl. phys. calculations)
- Study of how the exchange of $X_{17}(V,A,S,P)$ would impact such a process
- Beyond the resonance-saturation approach (justified for ^8Be but not for ^4He)
- Detailed study of the behaviour of the (V,A,S,P) induced angular correlations

Main results:

- The predicted cross sections are monotonically decreasing as function of the e^+e^- opening angle.
- Absence of any resonance-like structure
- Measurements at $\theta_{\text{vp}} \neq 90^\circ$ can discriminate $X=V,A,S,P$



$$M_x = 17 \text{ MeV}$$

$$\theta_{\text{vp}} = 90^\circ$$

${}^8\text{Be}$ vs. ${}^4\text{He}$: kinematic consistency [Feng+, PRD 2006.01151 [hep-ph]]

${}^8\text{Be}$ vs. ${}^4\text{He}$: kinematic consistency [Feng+, PRD 2006.01151 [hep-ph]]

For $M_x = 17 \text{ MeV}$ and uniform distrib. in $\cos \varphi(e^\pm \text{ c.o.m. axis vs. } v_x)$ the Lab. opening angle distrib. will be strongly peaked near their minimal values (when $e^\pm \text{ axis} \perp v_x$)
The theor. values are: $\Theta_{\pm}^{\min} = 112^\circ$ [${}^4\text{He}(20.49)$]; 139° [${}^8\text{Be}(18.15)$]; 161° [${}^{12}\text{C}(17.23)$].
[Exact for spin 0, approximate for spin 1]

${}^8\text{Be}$ vs. ${}^4\text{He}$: kinematic consistency [Feng+, PRD 2006.01151 [hep-ph]]

For $M_x = 17 \text{ MeV}$ and uniform distrib. in $\cos \varphi(e^\pm \text{ c.o.m. axis vs. } v_x)$ the Lab. opening angle distrib. will be strongly peaked near their minimal values (when $e^\pm \text{ axis} \perp v_x$)
 The theor. values are: $\Theta_{\pm}^{\min} = 112^\circ$ [${}^4\text{He}(20.49)$]; 139° [${}^8\text{Be}(18.15)$]; 161° [${}^{12}\text{C}(17.23)$].
 [Exact for spin 0, approximate for spin 1]

${}^4\text{He}$: $M_x = 16.94 \pm 0.24$, $\theta \sim 115^\circ$

${}^8\text{Be}$: $M_x = 17.01 \pm 0.16$, $\theta \sim 140^\circ$ [$\theta(17.64 \text{ MeV}) \sim 150^\circ$]

${}^{12}\text{C}$: M_x broadly consistent, $\theta \sim 160^\circ$ [prediction]

N_*	J_*^P	T_*	Γ_{N_*} [keV]	$B(N_* \rightarrow N_0 \gamma)$
${}^8\text{Be}(18.15)$	1^+	0	138	1.4×10^{-5}
${}^8\text{Be}(17.64)$	1^+	1	10.7	1.4×10^{-3}
${}^{12}\text{C}(17.23)$	1^-	1	1150	3.8×10^{-5}
${}^4\text{He}(21.01)$	0^-	0	840	0
${}^4\text{He}(20.21)$	0^+	0	500	6.6×10^{-10} (E0)

$({}^8\text{Be}, {}^{12}\text{C}, {}^4\text{He})_{\text{gs}} 0^+ \quad 0$

${}^8\text{Be}$ vs. ${}^4\text{He}$: kinematic consistency

[Feng+, PRD 2006.01151 [hep-ph]]

For $M_x = 17 \text{ MeV}$ and uniform distrib. in $\cos \varphi(e^\pm \text{ c.o.m. axis vs. } v_x)$ the Lab. opening angle distrib. will be strongly peaked near their minimal values (when $e^\pm \text{ axis} \perp v_x$)
 The theor. values are: $\Theta_{\pm}^{\min} = 112^\circ$ [${}^4\text{He}(20.49)$]; 139° [${}^8\text{Be}(18.15)$]; 161° [${}^{12}\text{C}(17.23)$].
 [Exact for spin 0, approximate for spin 1]

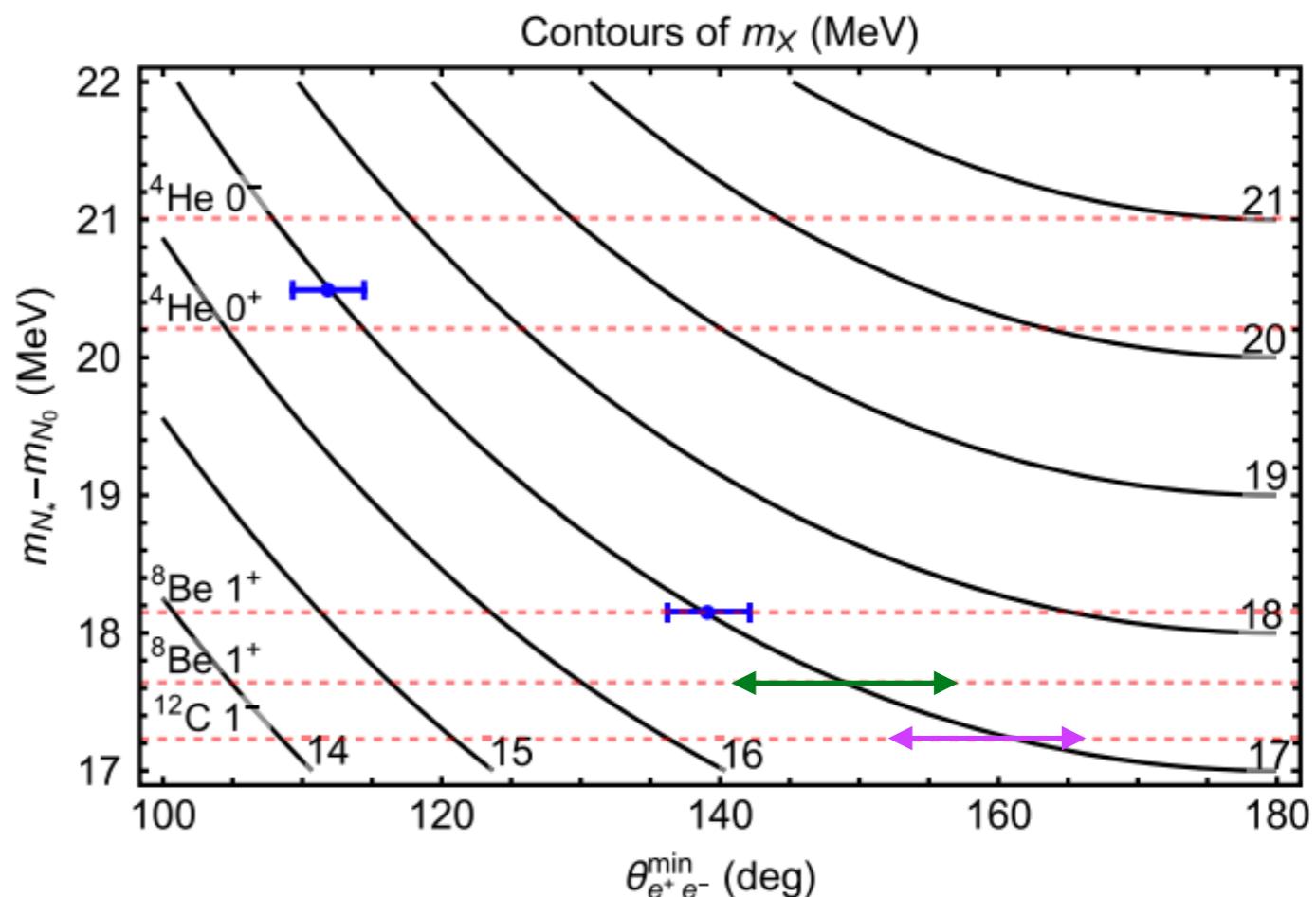
${}^4\text{He}$: $M_x = 16.94 \pm 0.24$, $\theta \sim 115^\circ$

${}^8\text{Be}$: $M_x = 17.01 \pm 0.16$, $\theta \sim 140^\circ$ [$\theta(17.64 \text{ MeV}) \sim 150^\circ$]

${}^{12}\text{C}$: M_x broadly consistent, $\theta \sim 160^\circ$ [prediction]

N_*	J_*^P	T_*	Γ_{N_*} [keV]	$B(N_* \rightarrow N_0 \gamma)$
${}^8\text{Be}(18.15)$	1^+	0	138	1.4×10^{-5}
${}^8\text{Be}(17.64)$	1^+	1	10.7	1.4×10^{-3}
${}^{12}\text{C}(17.23)$	1^-	1	1150	3.8×10^{-5}
${}^4\text{He}(21.01)$	0^-	0	840	0
${}^4\text{He}(20.21)$	0^+	0	500	6.6×10^{-10} (E0)

$({}^8\text{Be}, {}^{12}\text{C}, {}^4\text{He})_{\text{gs}} 0^+$ 0



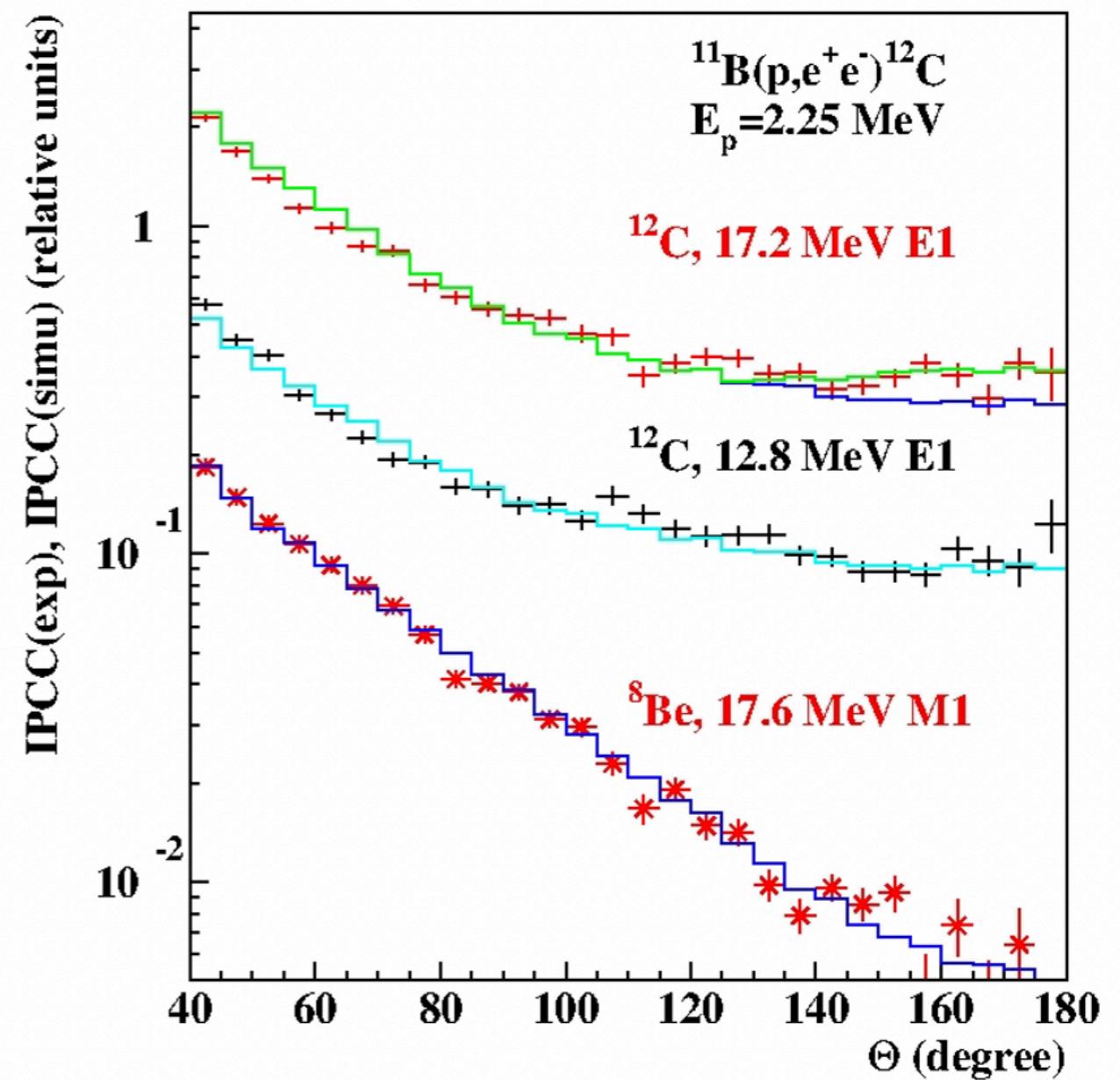
Preliminary results for ^{12}C

Nuclear reaction: $p + ^{11}B \rightarrow ^{12}C^*(17.23 \text{ MeV}) \rightarrow ^{12}C + e^+e^-$
 $E_p = 2.25 \text{ MeV}$ $J^P(^{12}C^*) = 1^-$

Preliminary results for ^{12}C

Nuclear reaction: $p + ^{11}B \rightarrow ^{12}C^*(17.23 \text{ MeV}) \rightarrow ^{12}C + e^+e^-$
 $E_p = 2.25 \text{ MeV}$ $J^P(^{12}C^*) = 1^-$

A. Krasznahorkay
"Shedding light on X17 Workshop"
Rome, September 6-8, 2021

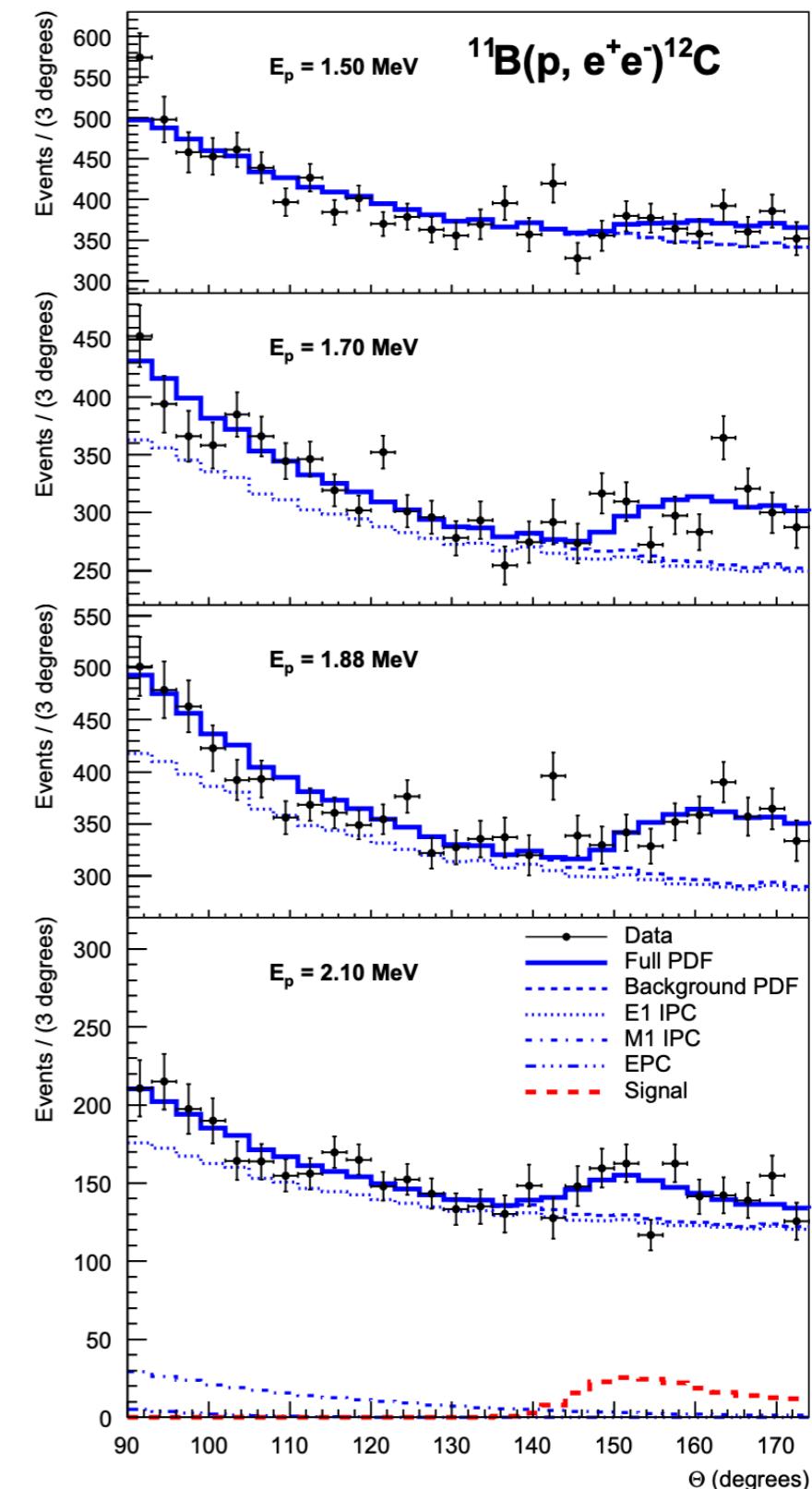


Nuclear reaction: $p + ^{11}B \rightarrow ^{12}C^*(17.23 \text{ MeV}) \rightarrow ^{12}C + e^+e^-$

$E_p = 2.25 \text{ MeV}$ $J^P(^{12}C^*) = 1^-$

E_p (MeV)	B_x $\times 10^{-6}$	Mass (MeV/c ²)	Confidence
1.50	1.1(6)	16.81(15)	3σ
1.70	3.3(7)	16.93(8)	7σ
1.88	3.9(7)	17.13(10)	8σ
2.10	4.9(21)	17.06(10)	3σ
Averages	3.6(3)	17.03(11)	
Previous [14]	5.8	16.70(30)	
Previous [28]	5.1	16.94(12)	

$M_X = 17.3 \pm 0.11 \pm 0.20 \text{ MeV}$ and B_X are consistent with the same X_{17} particle suggested by the 8Be and 4He anomalies



${}^8\text{Be}$ vs. ${}^4\text{He}$: dynamical consistency [Feng+, PRD 2006.01151 [hep-ph]]

${}^8\text{Be}$ vs. ${}^4\text{He}$: dynamical consistency [Feng+, PRD 2006.01151 [hep-ph]]

Allowed nuclear transitions and X_{17} mediators

N_*	J_*^P	Scalar X	Pseudoscalar X	Vector X	Axial Vector X
${}^8\text{Be}(18.15)$	1^+	✗	✓	✓	✓
${}^{12}\text{C}(17.23)$	1^-	✓	✗	✓	✓
${}^4\text{He}(21.01)$	0^-	✗	✓	✗	✓
${}^4\text{He}(20.21)$	0^+	✓	✗	✓	✗

Selection rules:
 $J^* = L \oplus J_X$
 $P^* = (-1)^L P_X$

${}^8\text{Be}$ vs. ${}^4\text{He}$: dynamical consistency [Feng+, PRD 2006.01151 [hep-ph]]

Allowed nuclear transitions and X_{17} mediators

N_*	J_*^P	Scalar X	Pseudoscalar X	Vector X	Axial Vector X
${}^8\text{Be}(18.15)$	1^+	✗	✓	✓	✓
${}^{12}\text{C}(17.23)$	1^-	✓	✗	✓	✓
${}^4\text{He}(21.01)$	0^-	✗	✓	✗	✓
${}^4\text{He}(20.21)$	0^+	✓	✗	✓	✗

Selection rules:
 $J^* = L \oplus J_X$
 $P^* = (-1)^L P_X$

Measured X_{17} production rates

$$\frac{\Gamma_X^{\text{Be}}}{\Gamma_\gamma^{\text{Be}}} = \frac{\Gamma({}^8\text{Be}^* \rightarrow {}^8\text{Be} + X)}{\Gamma({}^8\text{Be}^* \rightarrow {}^8\text{Be} + \gamma)} \simeq 6 \times 10^{-6} \quad {}^8\text{Be}^*(18.15)$$

$$\frac{\Gamma_X^{\text{He}}}{\Gamma_\pm^{\text{He}}} \equiv \frac{\Gamma({}^4\text{He}' \rightarrow {}^4\text{He} + X)}{\Gamma({}^4\text{He}^* \rightarrow {}^4\text{He} e^+e^-)} \simeq 4 \times 10^{-5} \quad {}^4\text{He}'(20.49), {}^4\text{He}^*(20.21)$$

Are these branchings consistent with a single set of X_{17} couplings?

${}^8\text{Be}$ vs. ${}^4\text{He}$: dynamical consistency

[Feng+, PRD 2006.01151 [hep-ph]]

Allowed nuclear transitions and X_{17} mediators

N_*	J_*^P	Scalar X	Pseudoscalar X	Vector X	Axial Vector X
${}^8\text{Be}(18.15)$	1^+	✗	✓	✓	✓
${}^{12}\text{C}(17.23)$	1^-	✓	✗	✓	✓
${}^4\text{He}(21.01)$	0^-	✗	✓	✗	✓
${}^4\text{He}(20.21)$	0^+	✓	✗	✓	✗

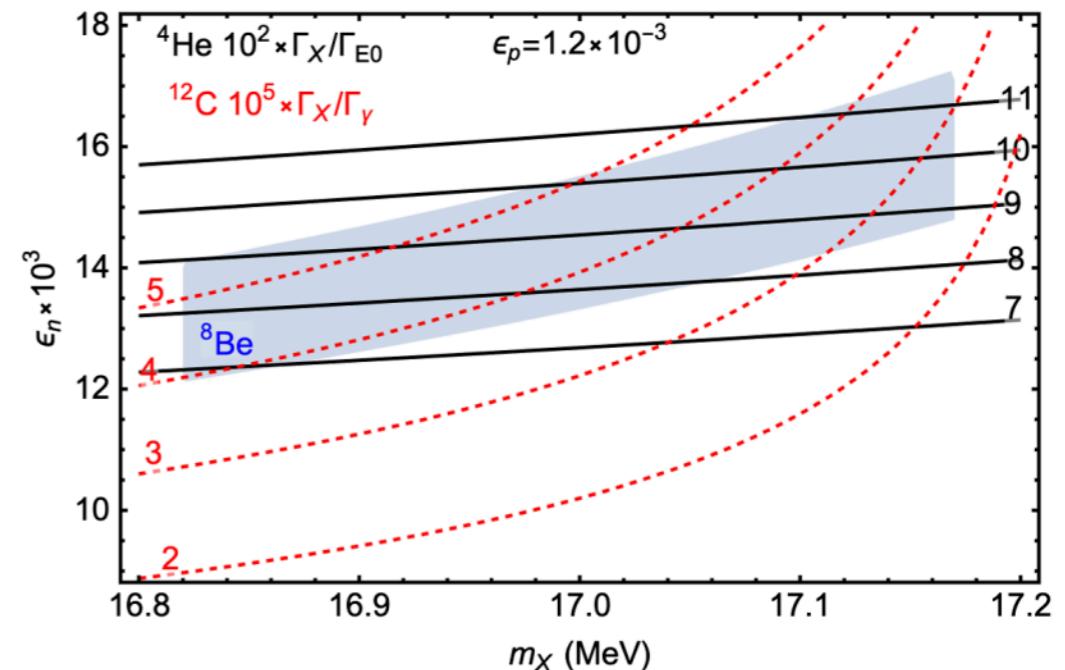
Selection rules:
 $J^* = L \oplus J_X$
 $P^* = (-1)^L P_X$

Measured X_{17} production rates

$$\frac{\Gamma_X^{\text{Be}}}{\Gamma_\gamma^{\text{Be}}} = \frac{\Gamma({}^8\text{Be}^* \rightarrow {}^8\text{Be} + X)}{\Gamma({}^8\text{Be}^* \rightarrow {}^8\text{Be} + \gamma)} \simeq 6 \times 10^{-6} \quad {}^8\text{Be}^*(18.15)$$

$$\frac{\Gamma_X^{\text{He}}}{\Gamma_\pm^{\text{He}}} \equiv \frac{\Gamma({}^4\text{He}' \rightarrow {}^4\text{He} + X)}{\Gamma({}^4\text{He}^* \rightarrow {}^4\text{He} e^+e^-)} \simeq 4 \times 10^{-5} \quad {}^4\text{He}'(20.49), {}^4\text{He}^*(20.21)$$

Are these branchings consistent with a single set of X_{17} couplings?



${}^8\text{Be}$ vs. ${}^4\text{He}$: dynamical consistency

[Feng+, PRD 2006.01151 [hep-ph]]

Allowed nuclear transitions and X_{17} mediators

N_*	J_*^P	Scalar X	Pseudoscalar X	Vector X	Axial Vector X
${}^8\text{Be}(18.15)$	1^+	✗	✓	✓	✓
${}^{12}\text{C}(17.23)$	1^-	✓	✗	✓	✓
${}^4\text{He}(21.01)$	0^-	✗	✓	✗	✓
${}^4\text{He}(20.21)$	0^+	✓	✗	✓	✗

Selection rules:
 $J^* = L \oplus J_X$
 $P^* = (-1)^L P_X$

Measured X_{17} production rates

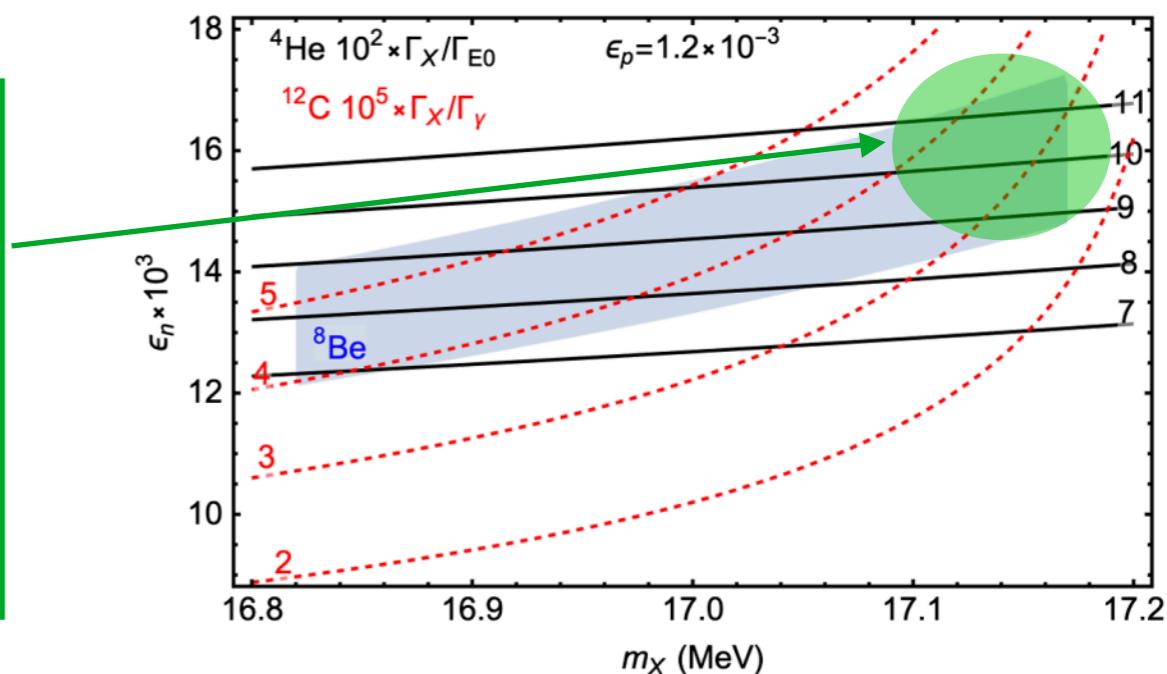
$$\frac{\Gamma_X^{\text{Be}}}{\Gamma_\gamma^{\text{Be}}} = \frac{\Gamma({}^8\text{Be}^* \rightarrow {}^8\text{Be} + X)}{\Gamma({}^8\text{Be}^* \rightarrow {}^8\text{Be} + \gamma)} \simeq 6 \times 10^{-6} \quad {}^8\text{Be}^*(18.15)$$

$$\frac{\Gamma_X^{\text{He}}}{\Gamma_\pm^{\text{He}}} \equiv \frac{\Gamma({}^4\text{He}' \rightarrow {}^4\text{He} + X)}{\Gamma({}^4\text{He}^* \rightarrow {}^4\text{He} e^+e^-)} \simeq 4 \times 10^{-5} \quad {}^4\text{He}'(20.49), {}^4\text{He}^*(20.21)$$

Are these branchings consistent with a single set of X_{17} couplings?

Protophobic Vector: ${}^8\text{Be} - {}^4\text{He} - {}^{12}\text{C}$
dynamical consistency region

Axial vector: might also explain
 ${}^8\text{Be} - {}^4\text{He}$ (with more difficulties)



Summarising:

- All the three anomalies $\gtrsim 7\sigma$, not a statistical fluctuation
- Bumps, not general excesses. Not a last bin effect
- By Introducing a new particle, remarkable improvement of the fits
- SM explanation seems disfavoured ${}^8\text{Be}$ [Zhang+, (2017)]; ${}^4\text{He}$ [Viviani+, (2021)]
- ${}^8\text{Be}$ - ${}^4\text{He}$ - ${}^{12}\text{C}$ anomalies kinematically & dynamically consistent for V (and A)
- For ${}^{12}\text{C}$ the effect was predicted, and confirmed by experimental data

Experimental perspective: Nuclear Physics

MEGII @ PSI: (search for CLFV $\mu^+ \rightarrow e^+ \gamma$)

^8Be : CW accelerator $E_p = 1.1 \text{ MeV}$, MEGII spectrometer, Li_2O target

Measurement during main HIPA 2022 shutdown (5σ , 50h DAQ)

Performed in Jan/Feb 2022 (possibly problems with ^7Li target ?)

LUNA-MV @ LNGS: high intensity proton beam and very low background

^4He via $^3\text{H}(\text{p},e^+e^-)^4\text{He}$ reaction. (RICH detector under study)

Measurements: 2023-5 (LoI in preparation)

n_ToF @ CERN: pulsed neutron beam in a wide energy range.

^4He via $^3\text{He}(\text{n},e^+e^-)^4\text{He}$. Measurements: 2022-24 (CERN LoI approved)

AN2000 @ LNL (INFN): Focus on ^8Be and, possibly, ^{12}C cases (timescale ?)

IPN@ORSAY (?)

Validation/confutation from a particle physics experiment

Validation/confutation from a particle physics experiment

PHYSICAL REVIEW D 97, 095004 (2018)

Resonant production of dark photons in positron beam dump experiments

Enrico Nardi,^{1,*} Cristian D. R. Carvajal,² Anish Ghoshal,^{1,3} Davide Meloni,^{3,4} and Mauro Raggi⁵

Since $X_{17} \rightarrow e^+ e^-$,
then $e^+ e^- \rightarrow X_{17}$

via positron-electron resonant
annihilation (CERN-EPFL-Korea
Theory Institute, early 2017)

Validation/confutation from a particle physics experiment

PHYSICAL REVIEW D 97, 095004 (2018)

Resonant production of dark photons in positron beam dump experiments

Enrico Nardi,^{1,*} Cristian D. R. Carvajal,² Anish Ghoshal,^{1,3} Davide Meloni,^{3,4} and Mauro Raggi⁵

BTF@LNF: $E_+ \sim 250 - 500 \text{ MeV}$

$\sqrt{s} \sim 15.8 - 22.4 \text{ MeV}$

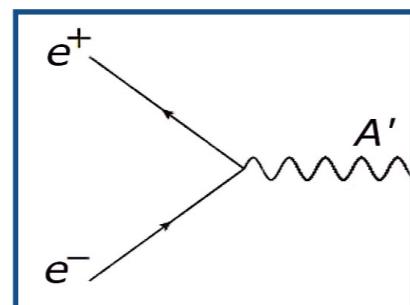
$M_X = 17 \text{ MeV}$ $E_+ = 289 \text{ MeV}$

Since $X_{17} \rightarrow e^+ e^-$,
then $e^+ e^- \rightarrow X_{17}$

via positron-electron resonant
annihilation (CERN-EPFL-Korea
Theory Institute, early 2017)

$$\sigma_{\text{res}} = \sigma_{\text{peak}} \frac{\Gamma_X}{2m_X} \delta\left(1 - \frac{\sqrt{s}}{M_X}\right) \quad \Gamma_X = 0.05 \left(\frac{\epsilon}{10^{-3}}\right)^2 \text{ eV}$$
$$\sigma_{\text{peak}} \sim 50 \text{ b}$$

"Huge" cross section !



Validation/confutation from a particle physics experiment

PHYSICAL REVIEW D 97, 095004 (2018)

Resonant production of dark photons in positron beam dump experiments

Enrico Nardi,^{1,*} Cristian D. R. Carvajal,² Anish Ghoshal,^{1,3} Davide Meloni,^{3,4} and Mauro Raggi⁵

BTF@LNF: $E_+ \sim 250 - 500 \text{ MeV}$

$\sqrt{s} \sim 15.8 - 22.4 \text{ MeV}$

$M_X = 17 \text{ MeV}$ $E_+ = 289 \text{ MeV}$

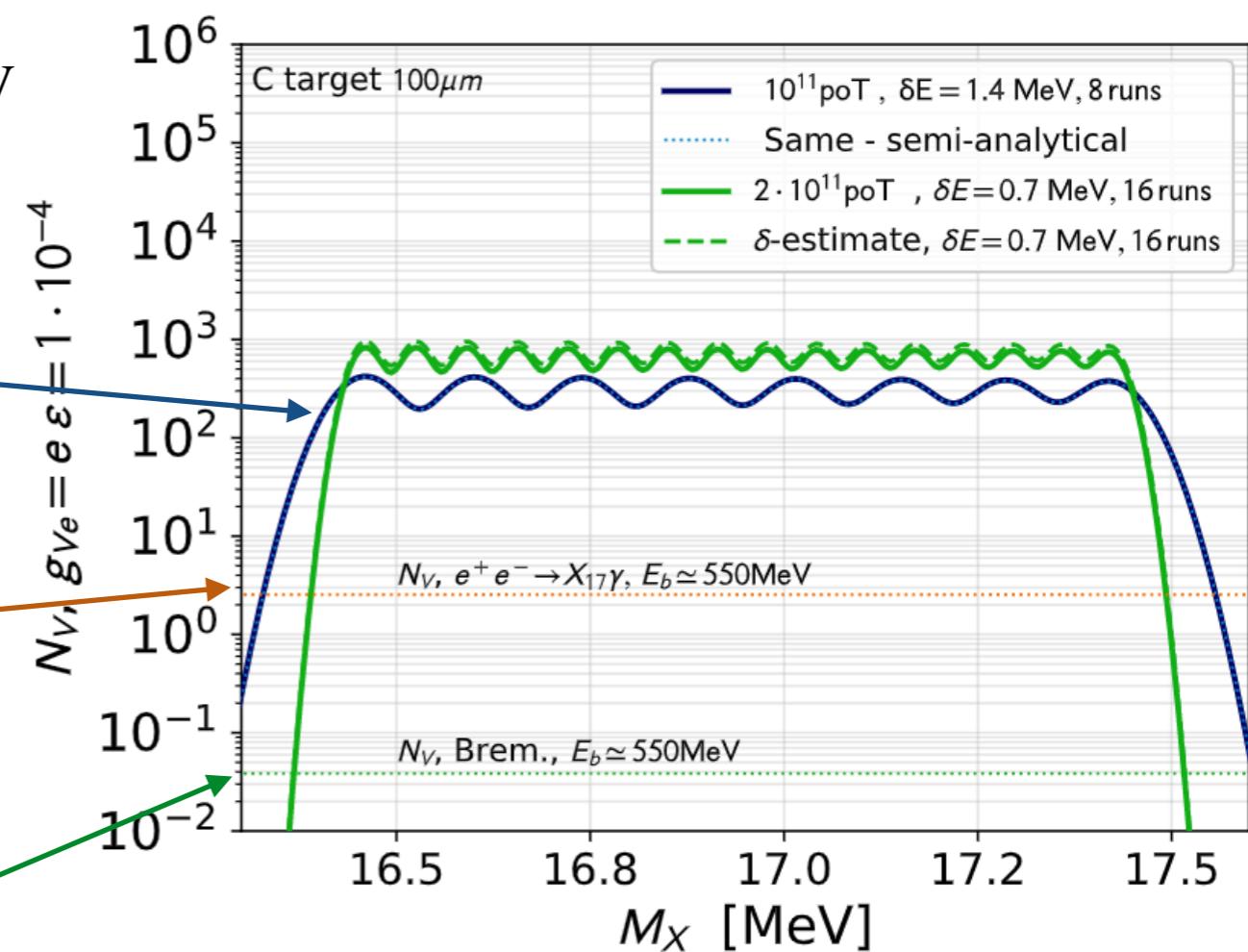
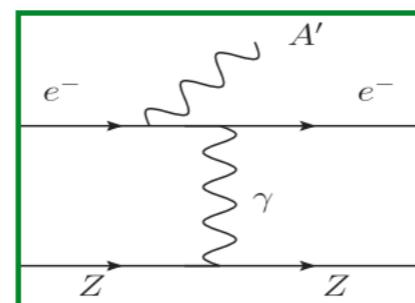
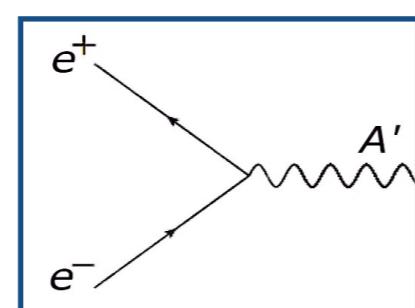
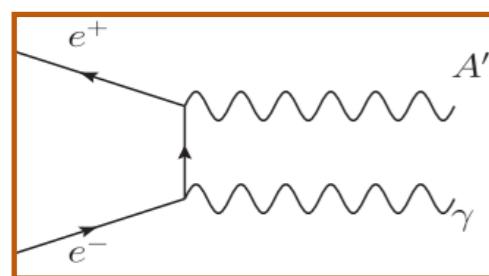
Since $X_{17} \rightarrow e^+ e^-$,
then $e^+ e^- \rightarrow X_{17}$
via positron-electron resonant
annihilation (CERN-EPFL-Korea
Theory Institute, early 2017)

$$\sigma_{\text{res}} = \sigma_{\text{peak}} \frac{\Gamma_X}{2m_X} \delta \left(1 - \frac{\sqrt{s}}{M_X} \right)$$

$$\Gamma_X = 0.05 \left(\frac{\epsilon}{10^{-3}} \right)^2 \text{ eV}$$

$$\sigma_{\text{peak}} \sim 50 \text{ b}$$

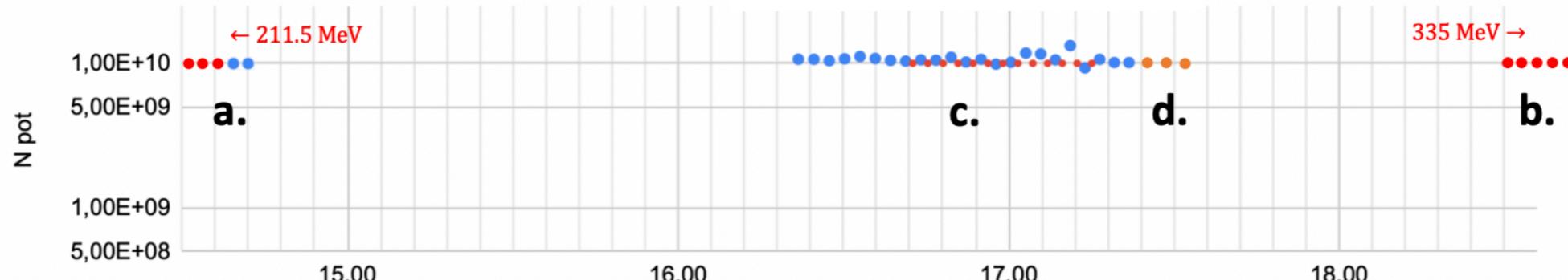
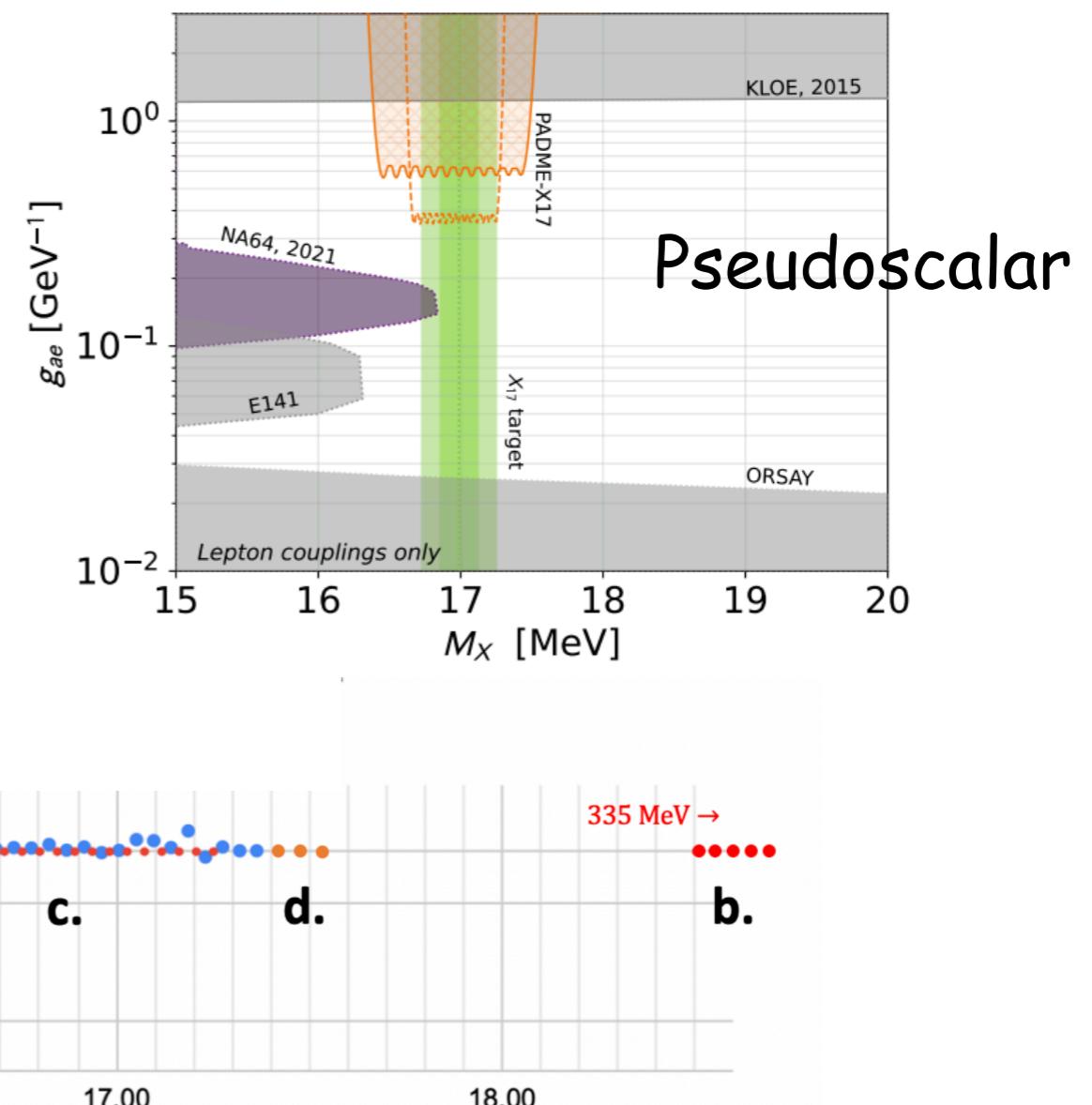
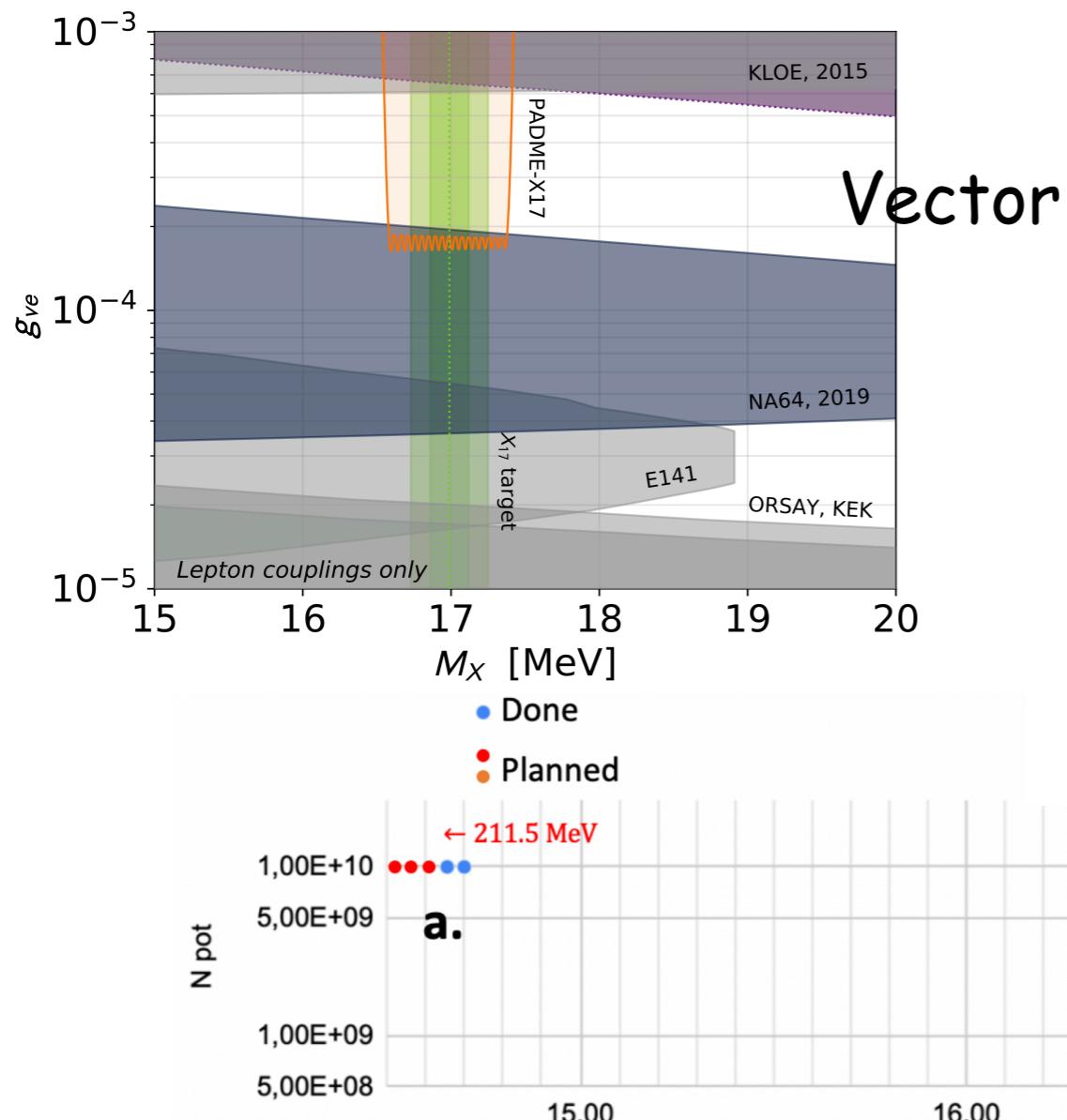
"Huge" cross section!



Resonant search for the X₁₇ boson at PADME

- E. Nardi, C. Carvajal, A. Ghoshal, D. Meloni, M. Raggi PRD97 095004 (2018)
- L. Darme, M. Mancini, E. Nardi, M. Raggi arXiv:2209.09261 [hep-ph]

- Our exp. colleagues are presently collecting data $E_{\text{beam}} \sim 290 \text{ MeV}$
- Control of beam parameters is excellent, background understood
- Our projections indicate that the spin-1 X_{17} can be fully tested
- Spin-0 pseudoscalar only partially (but a 0- particle is ^{12}C disfavoured)

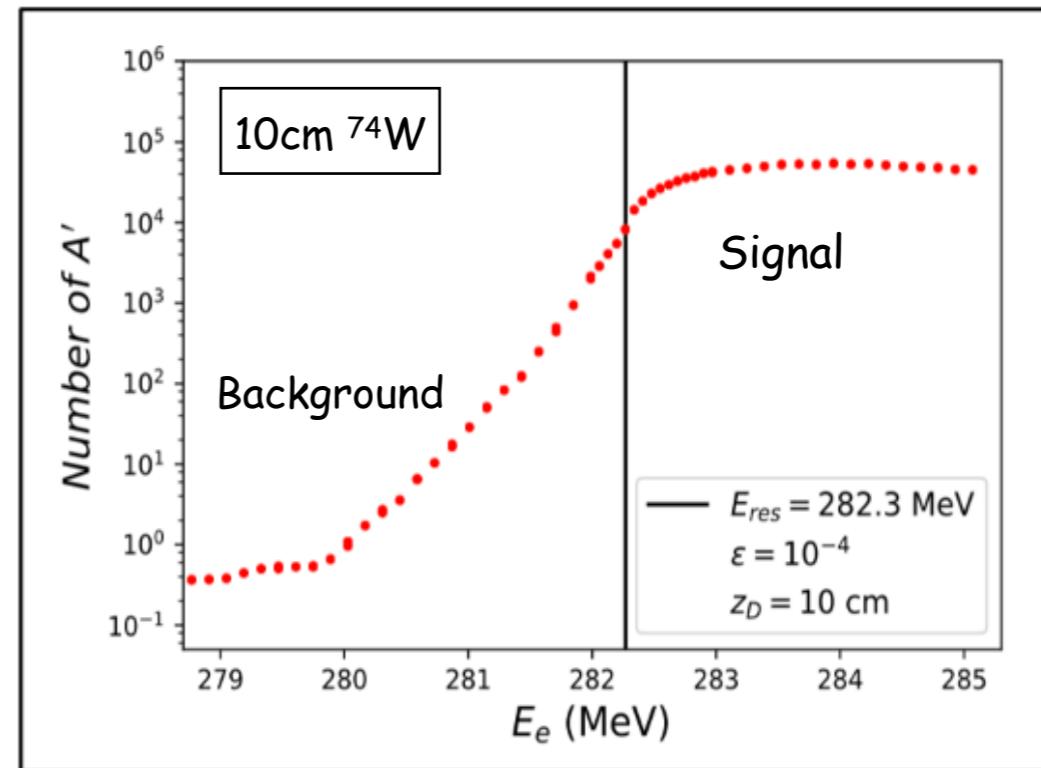


Conclusions

- Three anomalies observed in nuclear transitions appear to be consistent with a particle physics interpretation (X_{17})
- Statistical evidence is very strong ($> 7\sigma$)
- Explanations via higher order nuclear physics effects, interferences, higher multipoles contributions, are theoretically (strongly) disfavoured...
- Present data are from a single experiment. Independent verifications are needed.
- Intense effort for new Nucl. Phys. experiments is ongoing
First results expected probably in late 2023.
- Being of a completely different nature, a particle physics experiment can be decisive to validate the X_{17} hypothesis.

Several other advantages, as e.g. measurement of background

- E_{beam} below/above resonance
- Shoot with an e⁻ beam



- Although not optimal for $X \rightarrow e^+e^-$ detection/reconstruction (conceived for $e^+e^- \rightarrow \gamma X_{\text{invis.}}$) the existing PADME detector can be used (with minor upgrades)
- Beam tests at 280-290 MeV will be performed soon (weeks)
- Physics run most probably only after the summer 😠

LKB 2020 result from ^{87}Rb recoil velocity

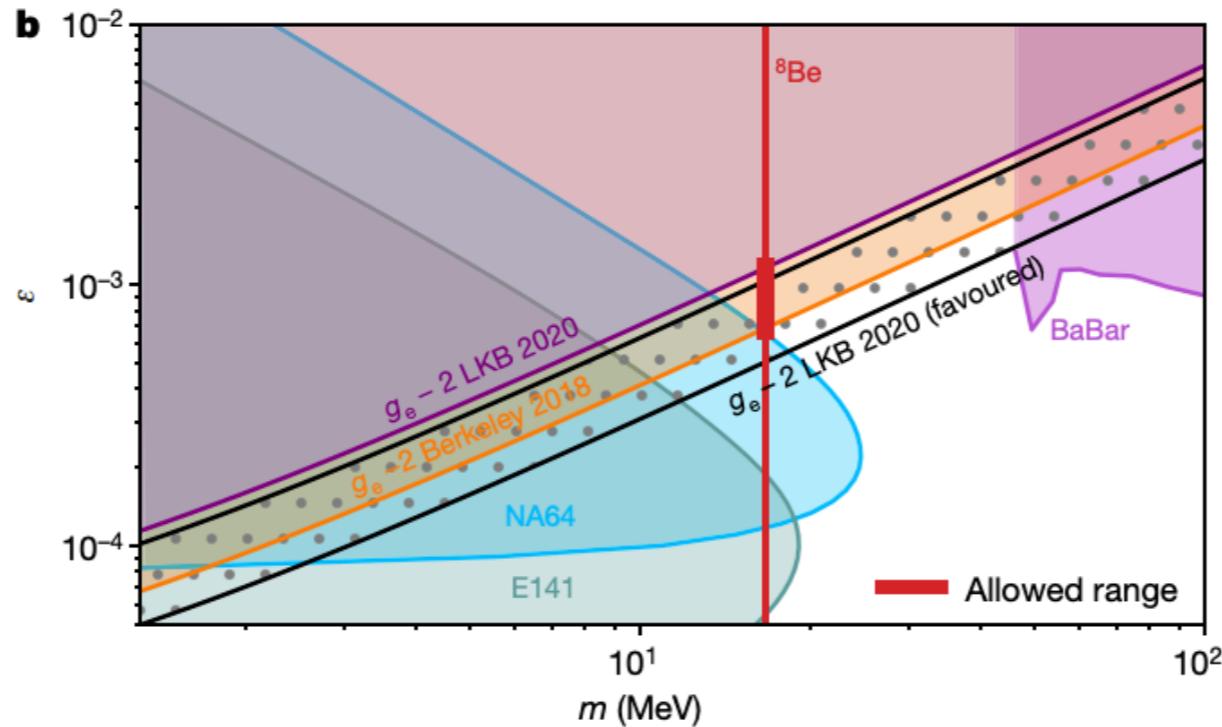


Fig. 4 | Impact on the test of the standard-model prediction of α_e and limits on hypothetical X boson. **a**, Summary of contributions to the relative uncertainty on $\delta\alpha_e$. The horizontal green line corresponds to the $\delta\alpha_e$ value obtained by taking into account the muon magnetic moment discrepancy and using a naive scaling model. Previous data from ref.⁹ (Harvard 2008), ref.¹⁸ (LKB 2011), ref.³ (Berkeley 2018), ref.¹³ (Atomic Mass Evaluation, AME 2016), ref.¹⁴ (Max-Planck-Institut für Kernphysik, MPIK 2014) and ref.² (RIKEN 2019). Also shown are the 10th-order and hadronic contributions in the calculation of the electron moment anomaly. **b**, Exclusion area in (ϵ, m_X) space for the X boson. The grey, blue and light purple regions are ruled out by the E141³¹, NA64³² and BaBar³⁵ experiments, respectively. A test based on the magnetic moment of the electron rules out the orange region when using the Berkeley measurement³ and the purple region when using the present result. Disregarding the Berkeley measurement, the remaining allowed range at 16.7 MeV is depicted by the thick red line. The zone favoured by $\delta\alpha_e > 0$, as deduced from this work, is shown by grey dots.