

# Extracting the time of core bounce from core-collapse supernovae

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# Extracting $t_0$ using an unbinned maximum likelihood technique

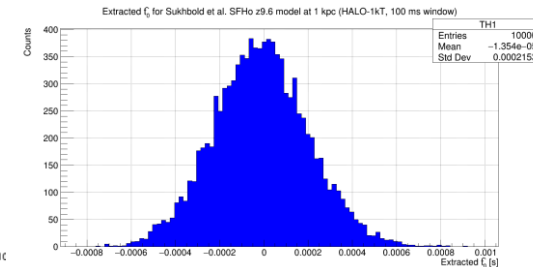
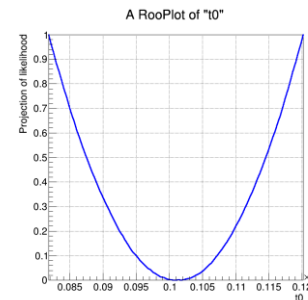
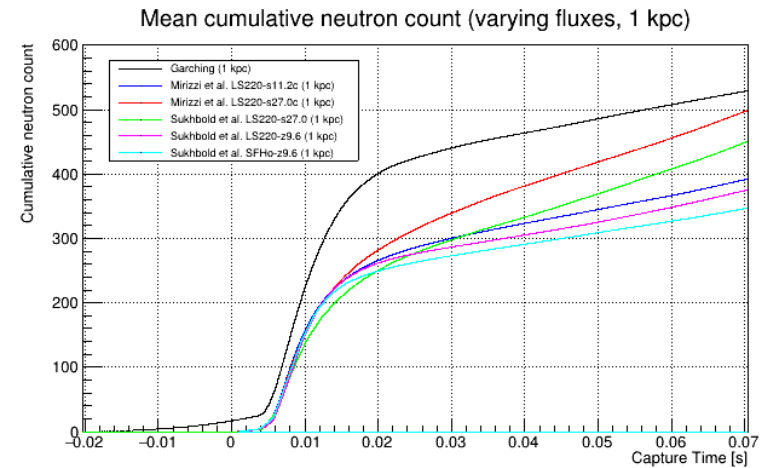
- Assume that the data derives from a Poisson PDF defined by:

$$P[n(t_i); \bar{n}(t_i - t_0)] = \frac{\bar{n}(t_i - t_0)^{n(t_i-t_0)} \cdot e^{-\bar{n}(t_i-t_0)}}{n(t_i)!}$$

- The negative ln likelihood (NLL) is therefore:

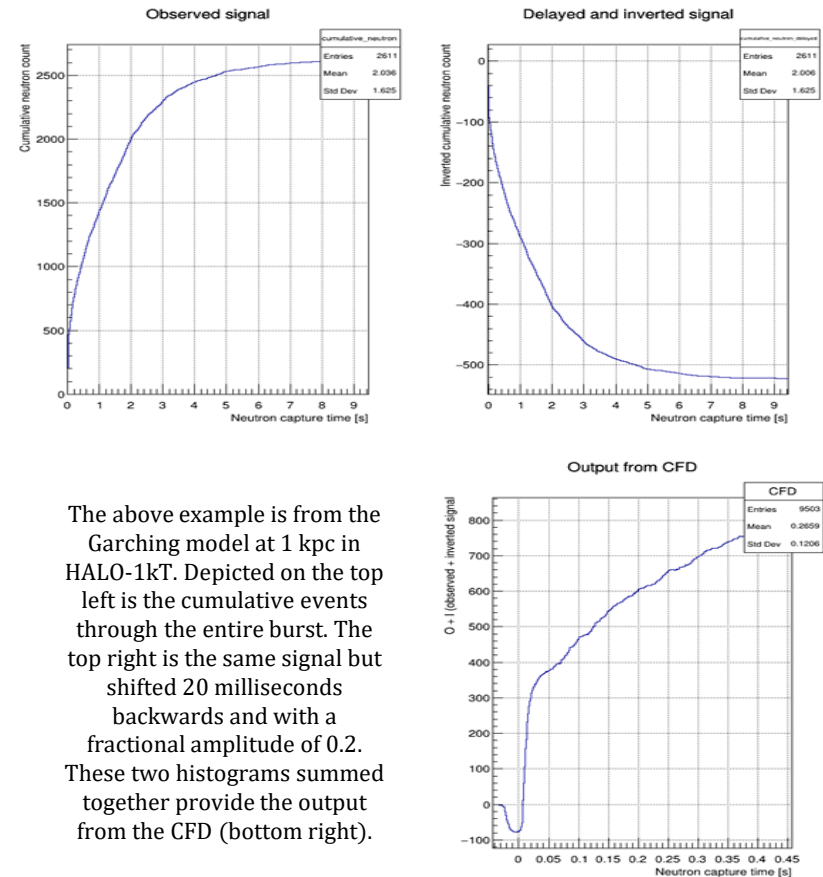
$$\ell(t_0) = 2 \cdot \sum_{i=1}^N \eta \cdot \bar{n}(t_i - t_0) + n(t_i) \cdot \ln(\eta \cdot \bar{n}(t_i - t_0))$$

- Here,  $n(t)$  is the vector of data with  $N$  events (where  $t_i$  is the time stamp of the  $i$ -th event),  $\eta$  is the normalization condition and  $\bar{n}(t)$  is the mean cumulative neutron count.
  - The determination of the mean cumulative count is depicted on the top right (exact technique is left in the backup slides).
- $t_0$  can be obtained from a minimization of  $\ell(t_0)$  using MINOS/Minuit2.
  - An example of this can be found on the bottom left plot (where the errors on the extracted  $t_0$  come from the points at  $\ell(t_0) \pm 1$ ).
- The precision and resolution of the extracted  $t_0$  can be quantified by repeating the NLL fit for  $10^4$  SN bursts (see bottom right).
  - Fitting data to itself does not account for systematic uncertainties (no knowledge of 'model of best fit'). To do so, mix and match various models to quantify offset.



# Extracting $t_0$ using a digital CFD

- We noticed in our analysis that extracting  $t_0$  as a function of distance, introduced some 'walk' from the expected  $t_0$ .
  - As will be shown later.
  - Walk was gradual, less than  $\frac{1}{2}$  a millisecond from 1-5 kpc.
- Elected to implement a digital constant fraction discriminator (CFD) to attempt a correction for this walk.
  - Feed input signal (cumulative events from entire burst) into CFD.
  - Invert and delay the signal, then scale it down by some value (constant fraction).
  - Sum the two signals together (see right).
  - Locate  $t_0$  from the point where the signal crossed zero.
    - Performance of extracting  $t_0$  is comparable to that the likelihood fit.



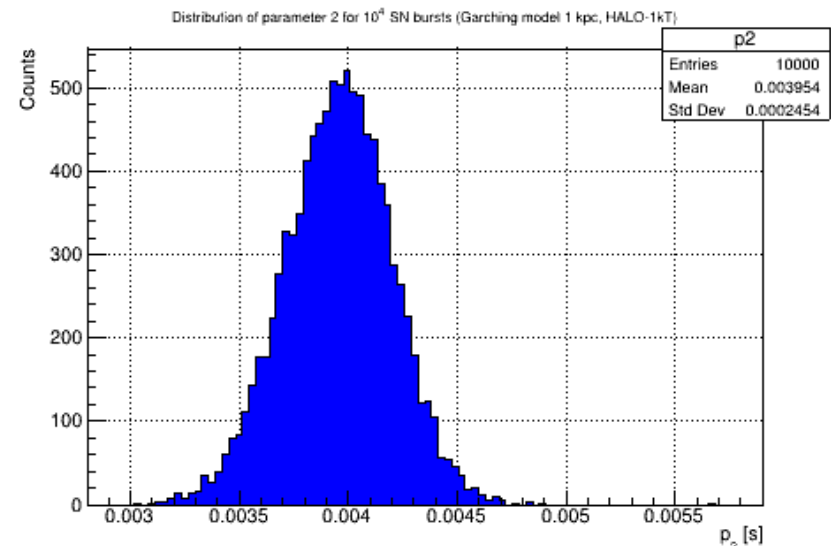
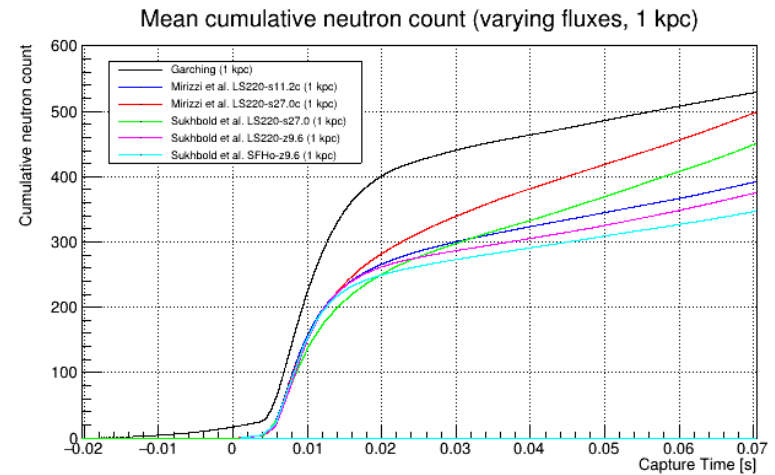
The above example is from the Garching model at 1 kpc in HALO-1kT. Depicted on the top left is the cumulative events through the entire burst. The top right is the same signal but shifted 20 milliseconds backwards and with a fractional amplitude of 0.2. These two histograms summed together provide the output from the CFD (bottom right).

# Extracting $t_0$ with a liner fit

- Linear technique arose as a method of getting around the systematic uncertainties included in the likelihood technique.
  - Use a single function as the cumulative PDF for all models.
- A relatively simple approximation is to treat the cumulative event count in the neutronization region as linear.

$$f(x) = [0] + [1] * (x - [2]) * (x > [2]) * (x < [4] + [2]) + \dots$$

- Introduce a five-parameter fit (first three parameters constrain fit in neutronization region).
  - [0] - intercept of linear line.
  - [1] - average event rate of the tested models.
  - [2] - offset (wrt time) from the true  $t_0$ .
- Final two parameters constrain end point of fit and event rate immediately following neutronization.
- Results for the extraction of parameter [2] for  $10^4$  bursts from the Garching model at 1 kpc in HALO-1kT can be found on the bottom right plot.

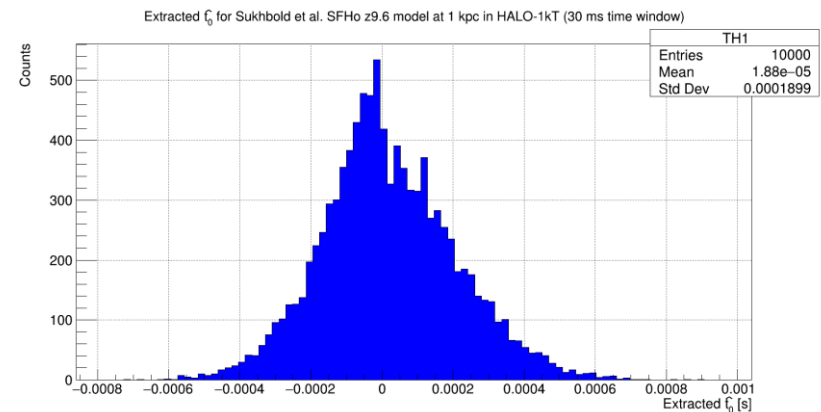
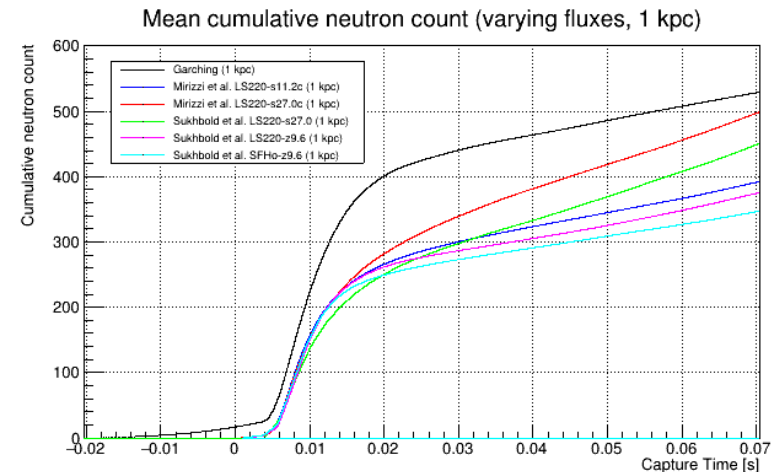


# Extracting $t_0$ with the Anderson Darling test

- Anderson Darling (AD) test is intended to determine the probability a sample belongs to some parent CDF.
  - It is an extension of the Kolmogorov Smirnov (KS) test.
  - The KS test was investigated as a candidate technique, performance collapses past 1 kpc in HALO-1kT.
- We implement it here as:

$$S(t_0) = \sum_{i=1}^N \frac{2i-1}{N} [\ln \bar{n}(t_i - t_0) + \ln(1 - \bar{n}(t_{N+1-i} - t_0))] A^2(t_0) = -N - S(t_0)$$

- Where  $N$  is the number of events (each with timestamp  $t_i$ ), and  $\bar{n}$  is the mean light curve evaluated at some offset.
- Analogous to the NLL technique, must account for systematic uncertainties when mixing model and PDF.
- Distribution of extracted  $t_0$  for the Sukhbold et al. SFHo z9.6 model at 1 kpc can be found in the bottom right.
  - Of the five techniques implemented, the AD test has the most numerical issues.



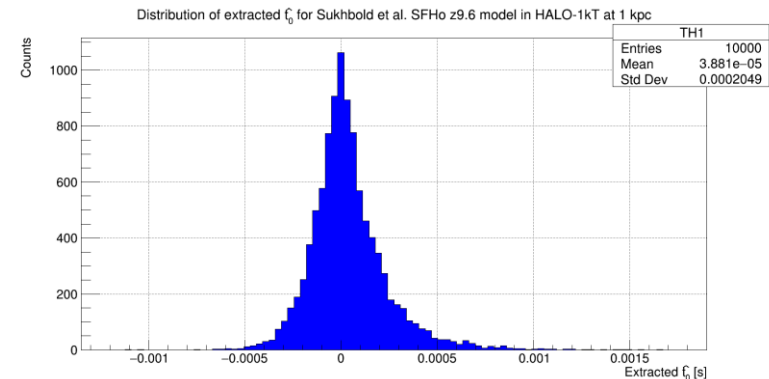
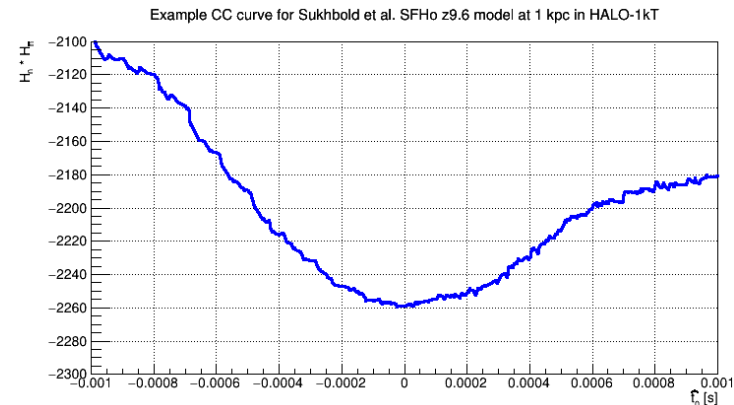
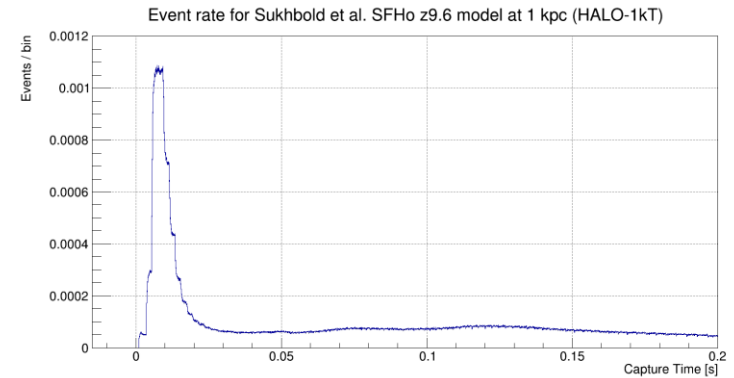
# Extracting $t_0$ via cross-correlation

- Originally intended to avoid binning our observed SN signal.
  - For HALO-1kT, relatively low statistics at 10 kpc, bins on the order 1 ms are sparsely populated.

- We implement cross-correlation as:

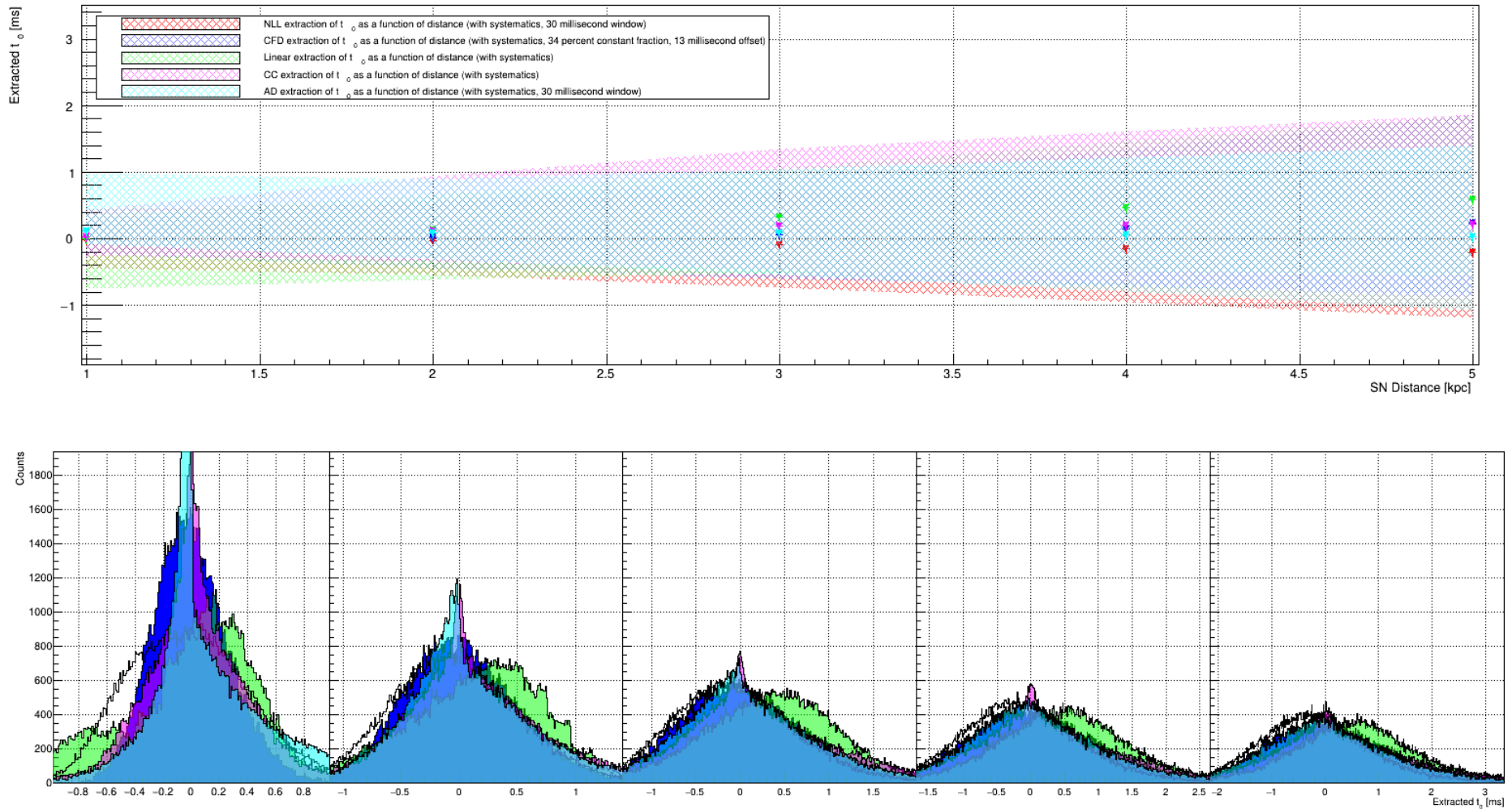
$$(H_n * H_{\bar{n}})(t - t_0) = - \sum_{-\infty}^{\infty} H_n(t) \cdot H_{\bar{n}}(t - t_0)$$

- Where  $H_n$  is the 1D histogram of the observed signal, and  $H_{\bar{n}}$  is the event rate (top) evaluated at some offset.
  - Observed time series is placed in a histogram with 1 ms bins.
  - $(H_n * H_{\bar{n}})$  should reach a minimum around the Monte Carlo truth  $t_0 = 0$ .
  - An example can be found in the middle figure.
- The extracted  $t_0$  distribution for the Sukhbold et al. SFHo z9.6 model at 1 kpc in HALO-1kT can be found in the bottom right.

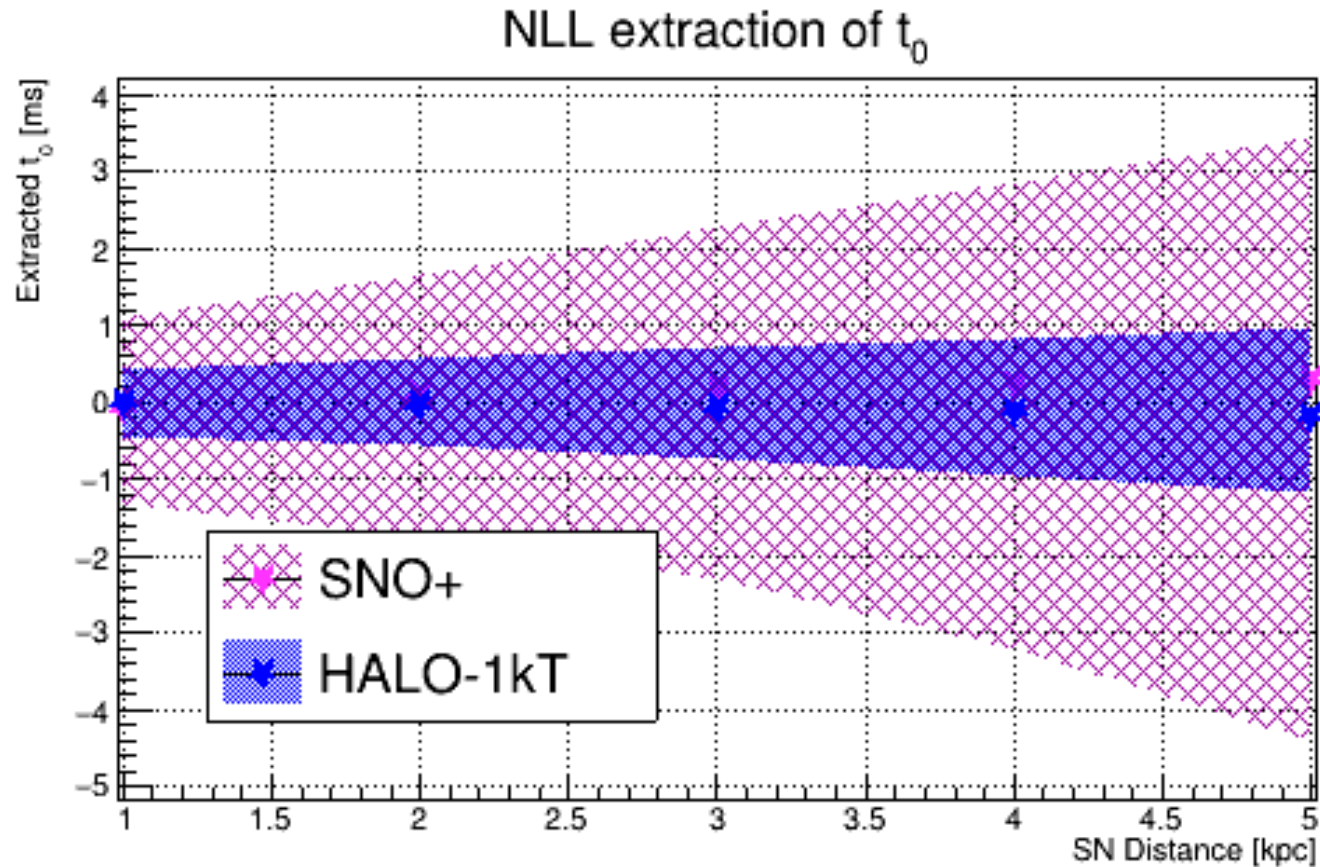




# Projection of techniques



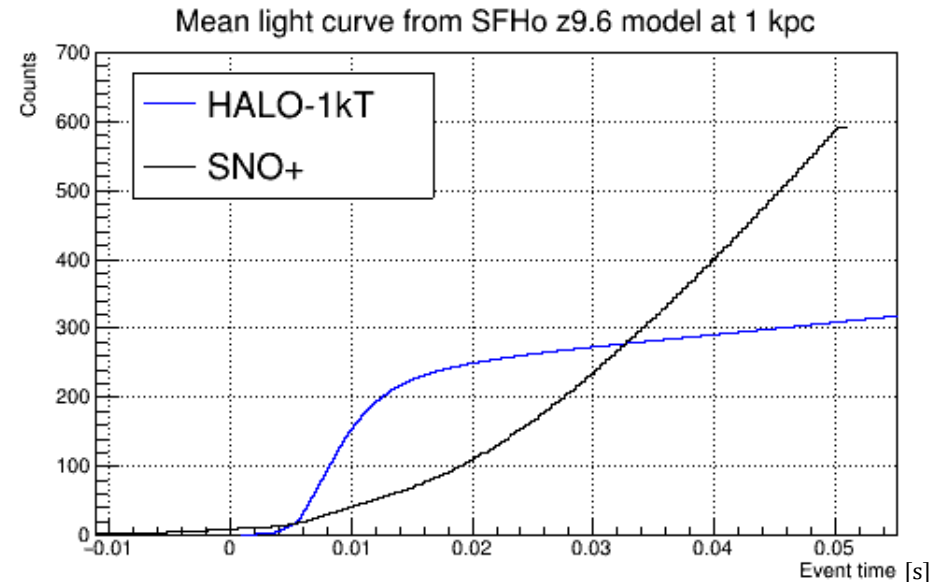
# SNO+ vs HALO-1kT





# Final remarks on SNO+

- As it stands, SNO+ fails to extract  $t_0$  with the same precision as HALO-1kT.
  - It is not a question of statistics, as SNO+ expects more events in the time window used in this analysis.
- My initial thoughts are that the neutronization burst in HALO-1kT (primarily  $\nu_e$  sensitive) is much easier to fit.
- A time window is necessary to reduce systematic uncertainties introduced when mixing models.
- However, with the current time window, there isn't a pronounced feature that the likelihood can easily locate.
- I have begun running some preliminary simulations of the entire burst to see if this improves our results.
  - For context, in both HALO-1kT and HALO, performance degrades as more of the burst is used. Maybe things will be different for SNO+.



Depicted above is the mean light curves as observed in HALO-1kT and SNO+ for the SFHo model at 1 kpc. The abrupt cut-off at 50 ms in SNO+ is a result of the upper time in sntools being set to 50 ms. Additionally, for context, in HALO-1kT the time of the events used comes from the capture time on  $^3\text{He}$  in the proportional counters. For SNO+, it comes from the event times as registered on the GPS clock in UTC. This is then shifted back to the Monte Carlo truth values.



# Integration into SNEWPDAG

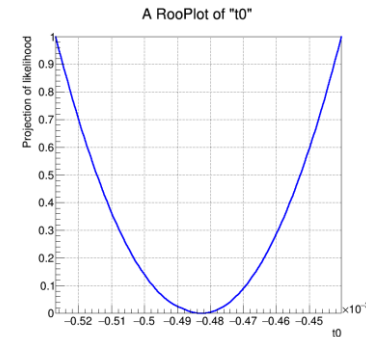
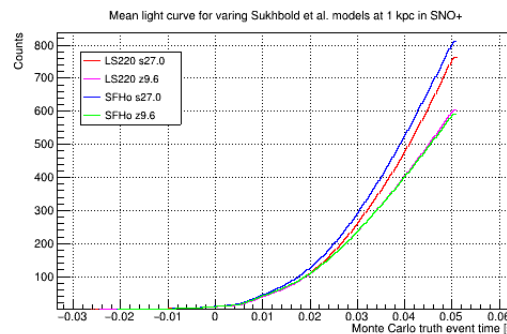
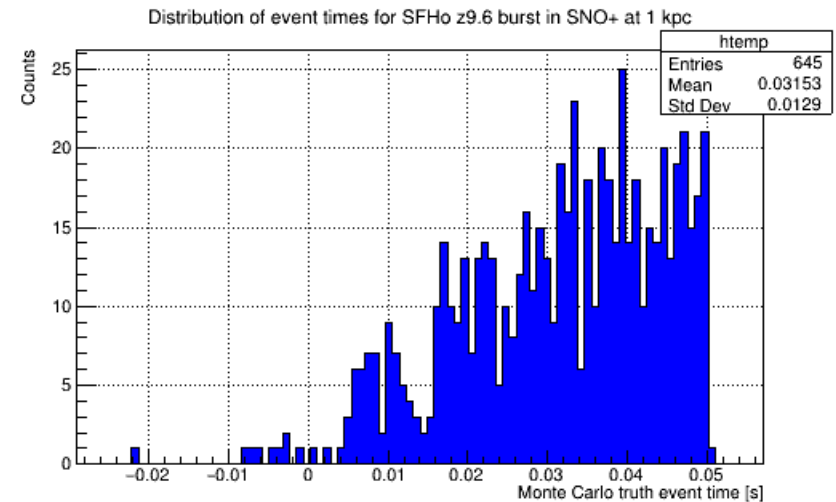
- Throughout the hackathon I have worked on implementing some of these techniques into SNEWPDAG.
  - Constant fraction discriminator
  - Negative log likelihood
  - Linear fit
  - Anderson Darling test
  - Cross-correlation
- Will move away from the full Monte Carlo treatment done in HALO-1KT, HALO and SNO+.
  - Will still have to account for detector response, background, etc.
- Run other  $\nu_e$  sensitive detectors through this pipeline to determine if my initial thoughts regarding the significance of the neutronization burst are true.
  - Will also be running some simulations with  $\nu$  oscillations.
- Each detector will be unique and formal recommendations on the technique required to extract  $t_0$  can be made once each detector is tested.
  - For HALO-1kT, the ideal candidate appears to be the CFD technique.

# Thank you for listening! Questions?



# Negative Log Likelihood in SNO+

- We apply the same technique used in HALO-1kT (slide 2) to the output from the SNO+ Monte Carlo.
  - The background is added in, Poisson fluctuated and sampled from a physics run.
  - However, the thresholds of this analysis are raised to a cleaned  $n_{hit} > 800$ , effectively removing its contribution.
- On top of the level 2 and level 3 triggers applied to the data, we introduce an additional cut.
  - Require 5 events with greater than 500  $n_{hit}$  each in a rolling 20 ms window.
  - This is intended to remove the sparse neutrino signal in the leading edge.
  - For both the likelihood and CFD technique these events caused severe issues in our extraction techniques.
  - Also cuts out many of elastic scattering on protons as a result.
- Applying these cuts will yield the event distribution on the top right.
  - The resulting likelihood function formed with the mean light curve can be found on the bottom right.



# Comparison to HALO-1kT

- Depicted on the top right is the extracted  $t_0$  from the SFHo z9.6 model in both HALO-1kT (left) and SNO+ (right) at 1 kpc.
  - Precision of extracted  $t_0$  is reduced by a factor of 3 in SNO+.
  - Difference in statistics is not attributed to failures to fit but rather the amount of simulations done.
- On the bottom right is the extracted  $t_0$  with the systematics of fitting all models to one another included.
  - With the systematics included the difference in precision is reduced to a factor of  $\sim 2.5$ .
  - With the current techniques, SNO+ does not obtain sub 1 ms precision at 1 kpc.
  - The comparison between the two experiments is not much improved at further distances.

