

Optics measurements & correction in the old ESRF storage ring (2010-2018)

Andrea Franchi (ESRF, Grenoble)

mini-workshop on FCC-ee optics tuning and alignment, 11-12 May 2022

- Understanding & correcting linear coupling (2009-2010)
- Nonlinear optics & magnet calibration via turn-by-turn (TbT) BPM data (2010-2013)
- Error analysis of linear optics measurements via orbit & TbT analysis (2016-2017)
- Applications of AC Orbit data (2016-2019)
- Experience of AC dipole & TbT BPM data (2016-2018)

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- Misconception #1: In the presence of coupling we deal with vertical emittances **S** (plural is not a typo)
- Misconception #2: Coupling correction is a linear problem not needing CPU-intense random/genetic/non-linear optimizers
- Misconception #3: Guignard formulas don't apply to modern synchrotron light sources & damping rings (**integer part** of tunes $Q_x \gg Q_y$)
- Misconception #4: “indirect measurements” of vertical emittance should be avoided

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Vertical emittance in the **absence** of coupling

- Eigen-emittance \mathcal{E} : constant along the ring, with $\mathcal{E}_v \approx 0$
- Non measurable RMS emittance:

$$\epsilon_y = \sqrt{\sigma_y(s)\sigma_p(s) - \sigma_{yp}^2(s)}$$

- Measurable emittance from RMS beam size:

$$\mathbb{E}_y = \frac{\sigma_y^2(s)}{\beta_y(s)} = \frac{\langle y^2(s) \rangle - (\delta D_y(s))^2}{\beta_y(s)}$$

$\mathcal{E}_v = \mathcal{E}_y = E_y = \text{const.}$

With zero vertical dispersion, $\mathcal{E}_v = \mathcal{E}_y = E_y \approx 0$

Vertical emittances **s** in the **presence** of coupling

- Eigen-emittance \mathcal{E} : constant along the ring, but $\mathcal{E}_v \neq 0$
- Non measurable **projected s-dependent** RMS emittance:

$$\epsilon_y(s) = \sqrt{\sigma_y(s)\sigma_p(s) - \sigma_{yp}^2(s)}$$

- Measurable **apparent s-dependent** emittance from RMS beam size:

$$\mathbb{E}_y(s) = \frac{\sigma_y^2(s)}{\beta_y(s)} = \frac{\langle y^2(s) \rangle - (\delta D_y(s))^2}{\beta_y(s)}$$

$\mathcal{E}_v = \text{const} \neq \epsilon_y(s) \neq \mathbb{E}_y(s)$

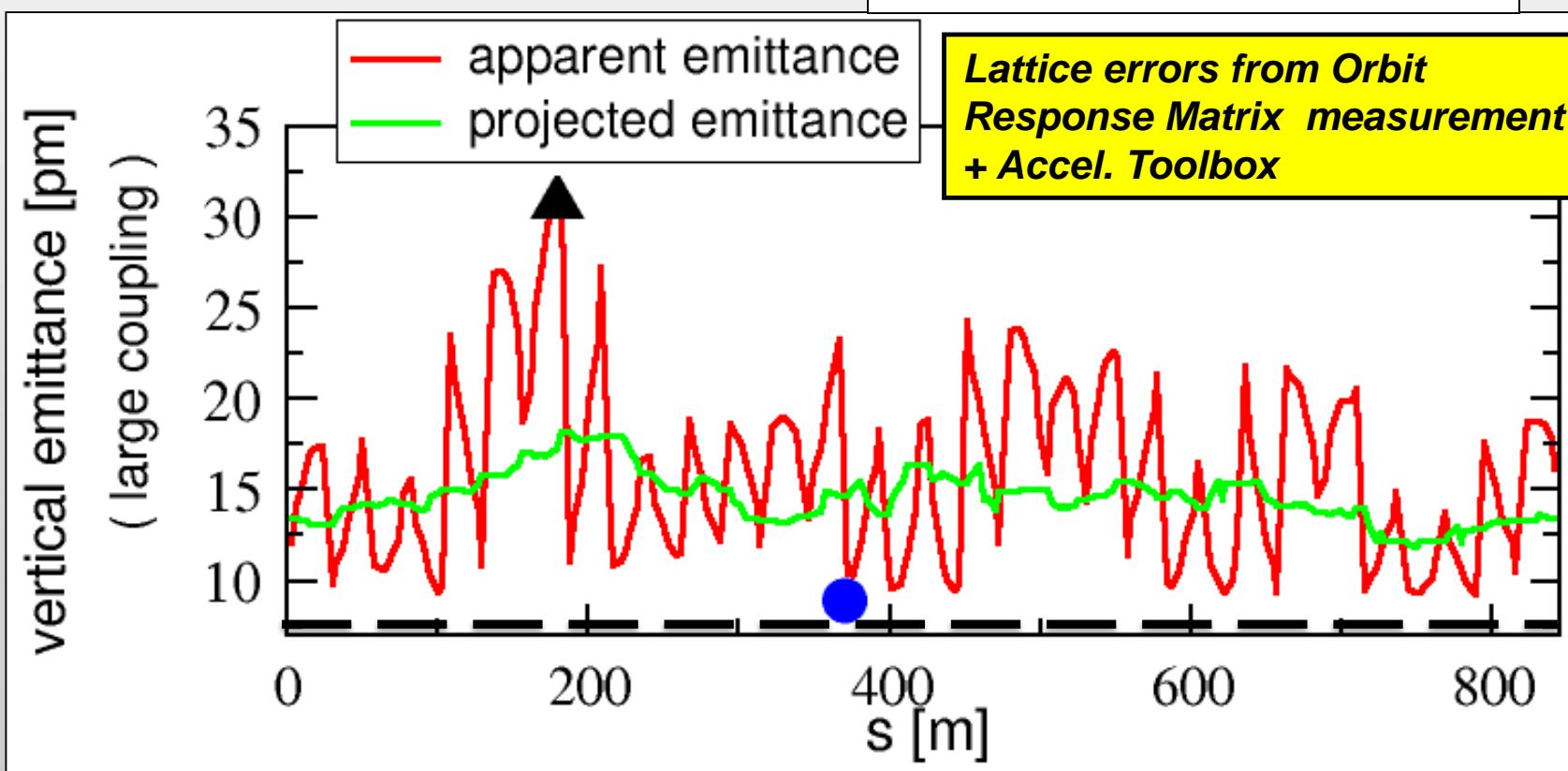
Vertical emittances in the presence of coupling

Measurable apparent emittance:

$$\mathbb{E}_y(s) = \frac{\sigma_y^2(s)}{\beta_y(s)} = \frac{\langle y^2(s) \rangle - (\delta D_y(s))^2}{\beta_y(s)}$$

Non measurable projected emittance:

$$\epsilon_y(s) = \sqrt{\sigma_y(s)\sigma_p(s) - \sigma_{yp}^2(s)}$$



Vertical emittances in the presence of coupling

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PHYSICAL REVIEW ACCELERATORS AND BEAMS 25, 044001 (2022)

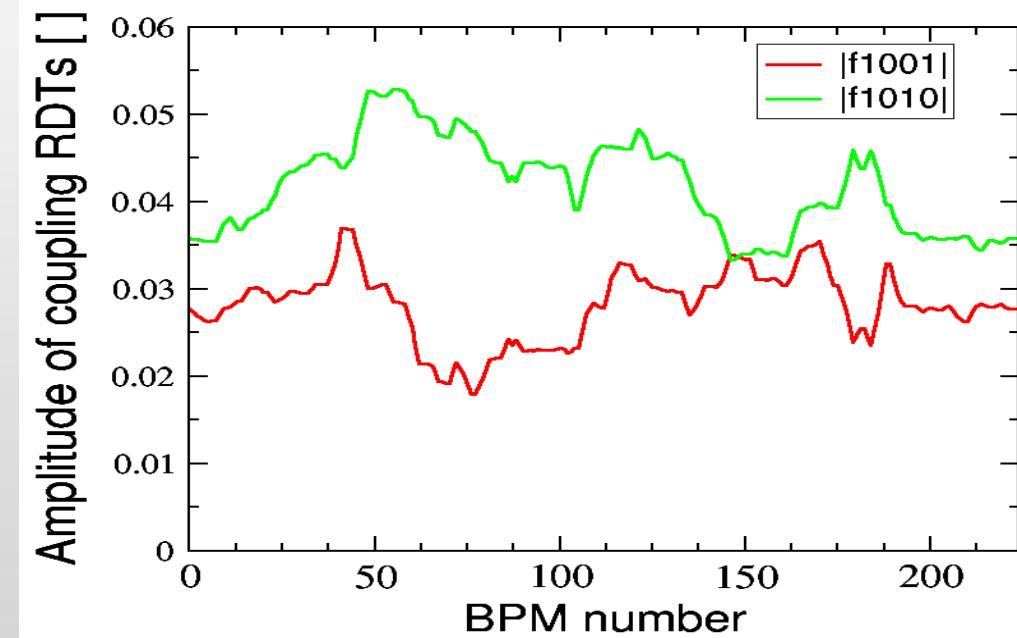
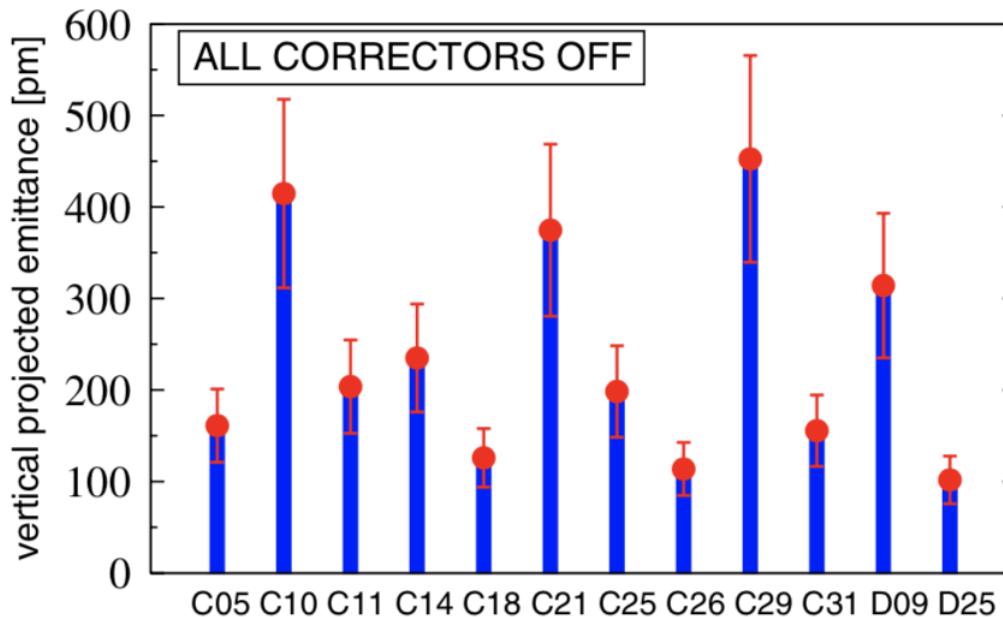
Demonstration of eigen-to-projected emittance mapping for an ellipsoidal electron bunch

consider the position-conjugate momentum pair (q_ℓ, p_ℓ) ,
with $\ell = x, y, z$, associated with d.o.f., the moment
invariant $\varepsilon_\ell \equiv [\langle q_\ell^2 \rangle \langle p_\ell^2 \rangle - \langle q_\ell p_\ell \rangle^2]^{1/2}$ is often introduced

In beam
physics, the eigenemittances are a generalization of the
projected emittances to the case of beams with coupled
d.o.f. [8].

Coupling correction is a linear problem not needing CPU-intense random/genetic/non-linear optimizers

all skew quad correctors OFF $\varepsilon_y/\varepsilon_x \sim 5\%$

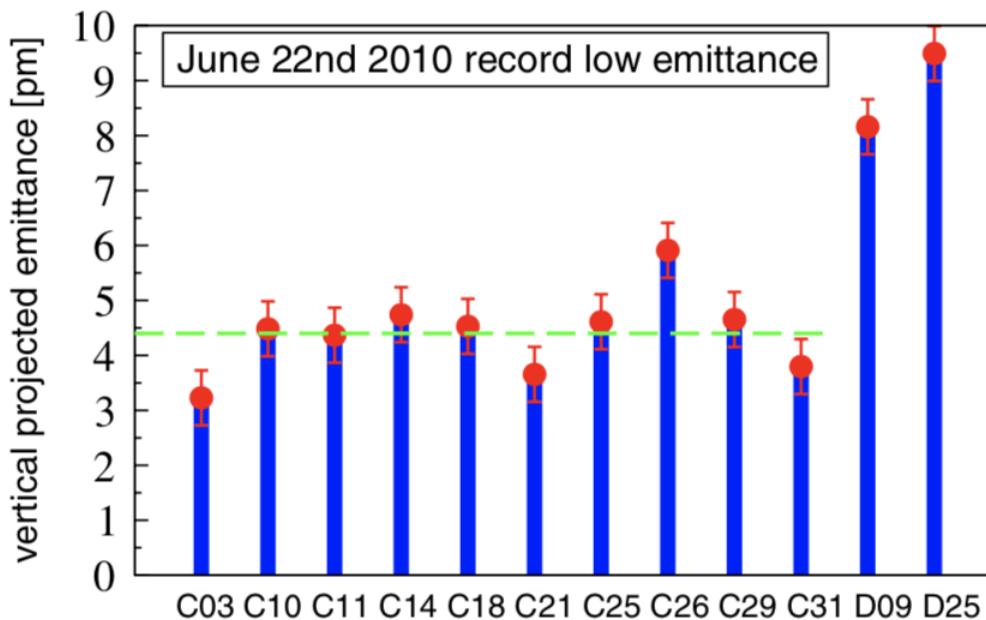


C^* , D^* : emittance monitors

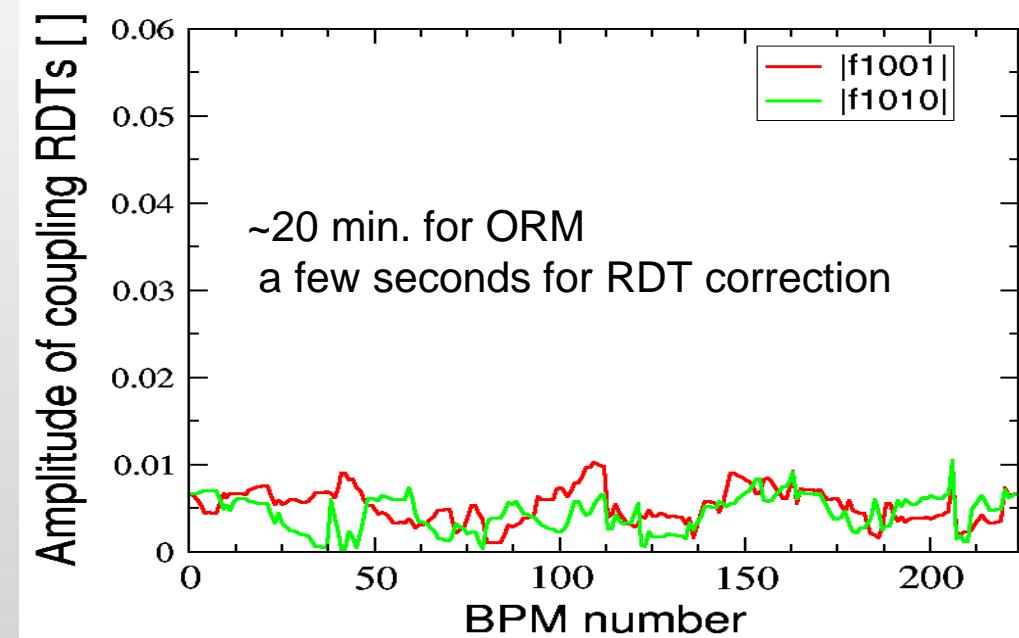
$$\begin{pmatrix} \vec{a}_1 \vec{f}_{1001} \\ \vec{a}_1 \vec{f}_{1010} \\ a_2 \vec{D}_y \end{pmatrix}_{\text{meas}} = -\mathbf{M} \vec{J}_c,$$

Coupling correction is a linear problem not needing CPU-intense random/genetic/non-linear optimizers

after correction $\varepsilon_y/\varepsilon_x \sim 1\%$



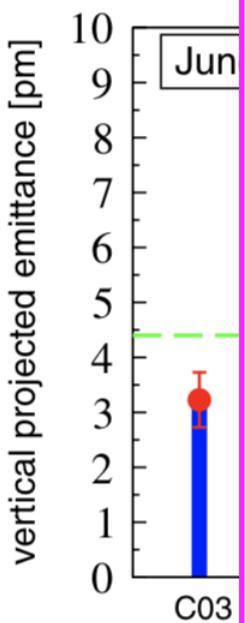
C*, D*: emittance monitors



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Coupling correction is a linear problem not needing CPU-intense random/genetic/non-linear optimizers

after c



Nuclear Instruments and Methods in Physics
Research Section A: Accelerators, Spectrometers,
Detectors and Associated Equipment



Volume 892, 1 June 2018, Pages 1-9

Coupling control and optimization at the Canadian Light Source

We

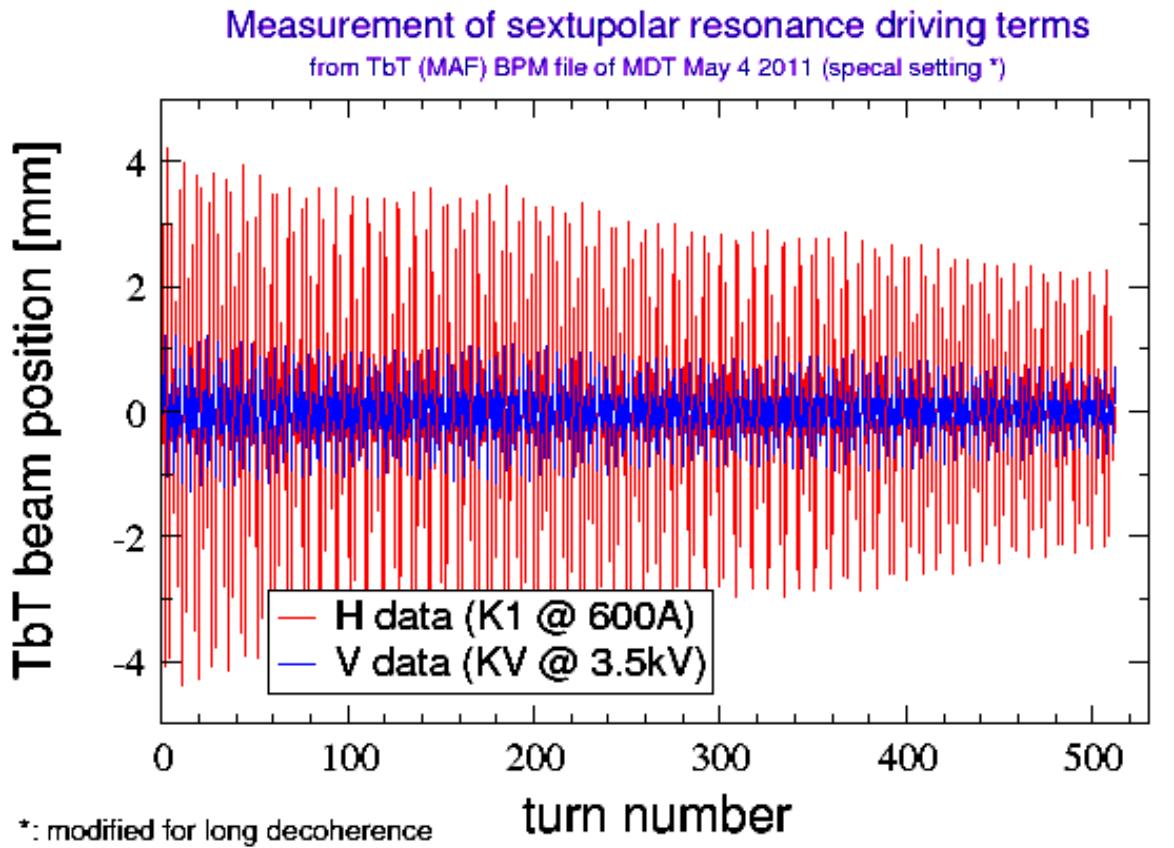
calculate the six-dimensional beam envelop matrix and use it to produce a variety of objective functions for optimization using the Multi-Objective Particle Swarm

C^* , D

Optimization (MOPSO) algorithm. MOPSO produces a number of skew quadrupole configurations that we apply to the storage ring. We use the X-ray synchrotron radiation diagnostic beamline to image the beam and we make measurements of the vertical dispersion and beam lifetime.



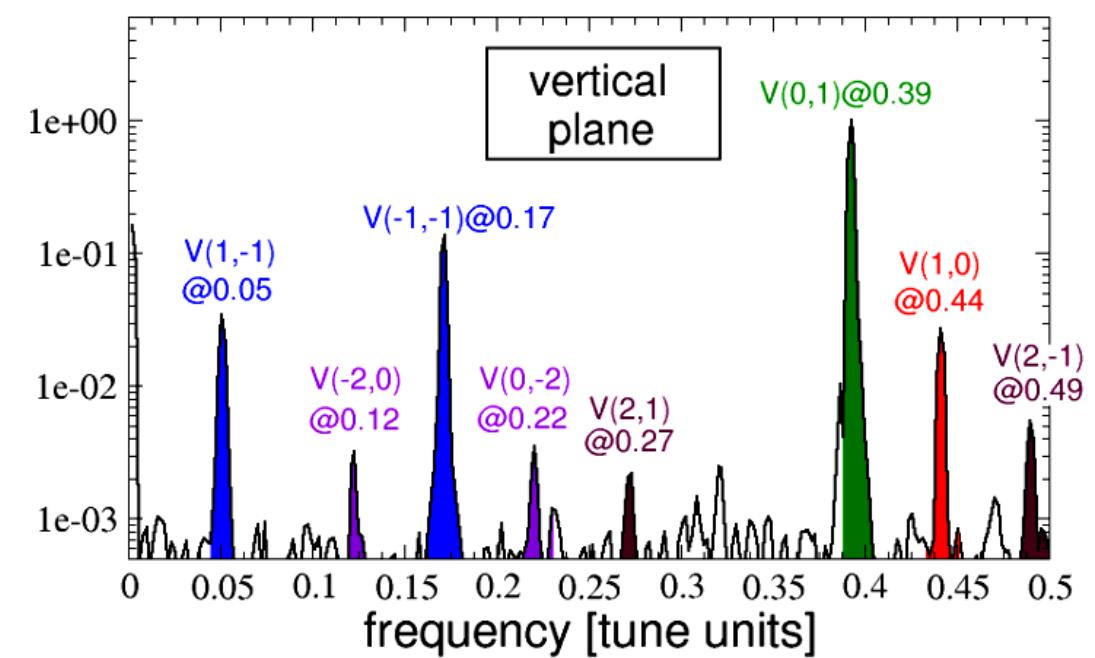
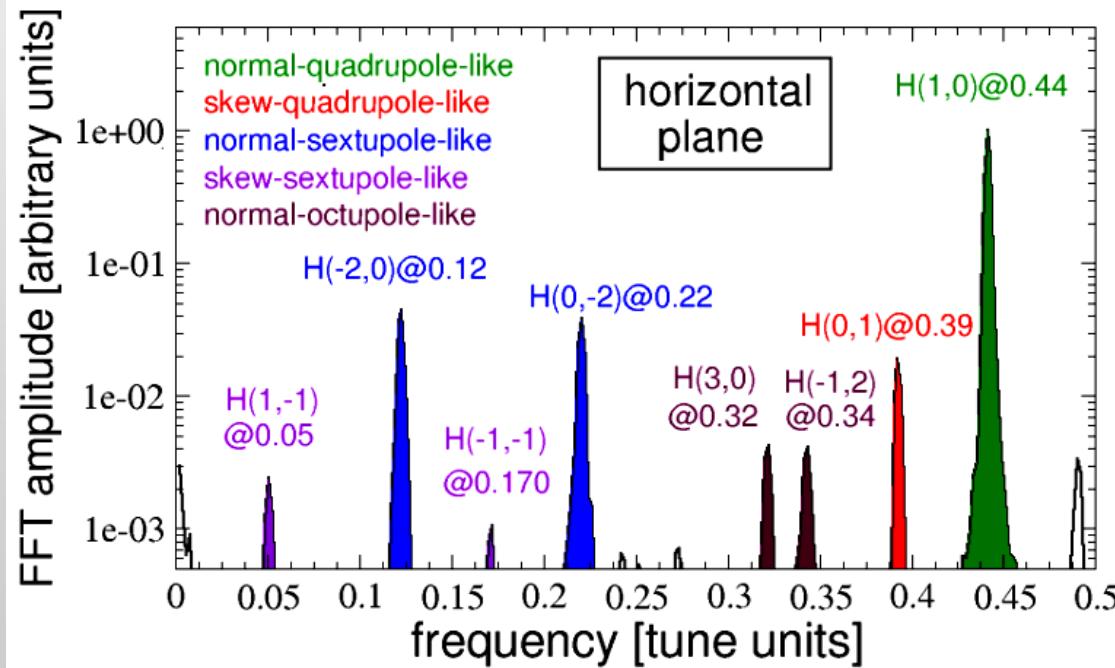
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1. Get synchronized and *decent*(*) TbT BPM data in both planes

(*) from nonlinear optics with zero chromaticity and zero (linear) amplitude-dependent detuning

2. Perform a *suitable* FFT of the TbT data & look at the lines popping up in the spectra



RDTs require $(x, p_x), (y, p_y)$

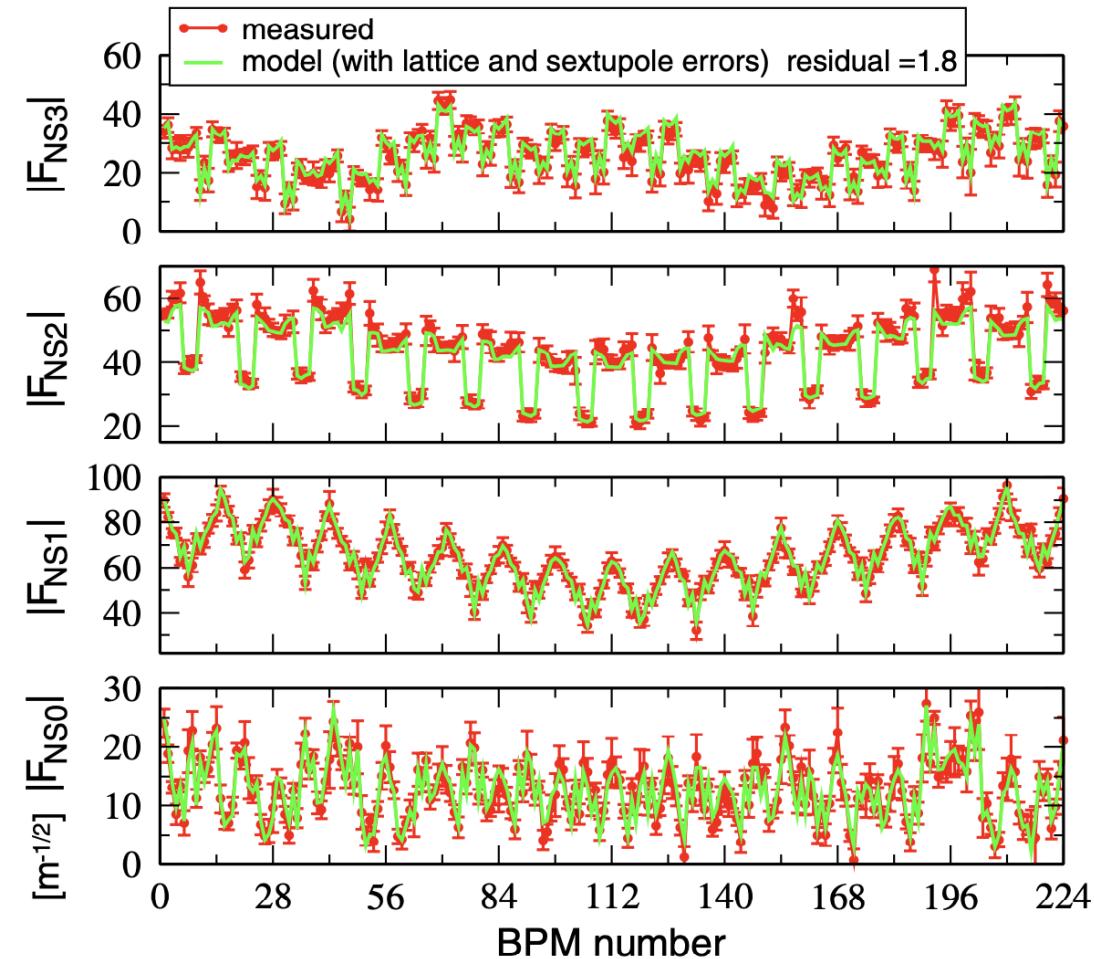
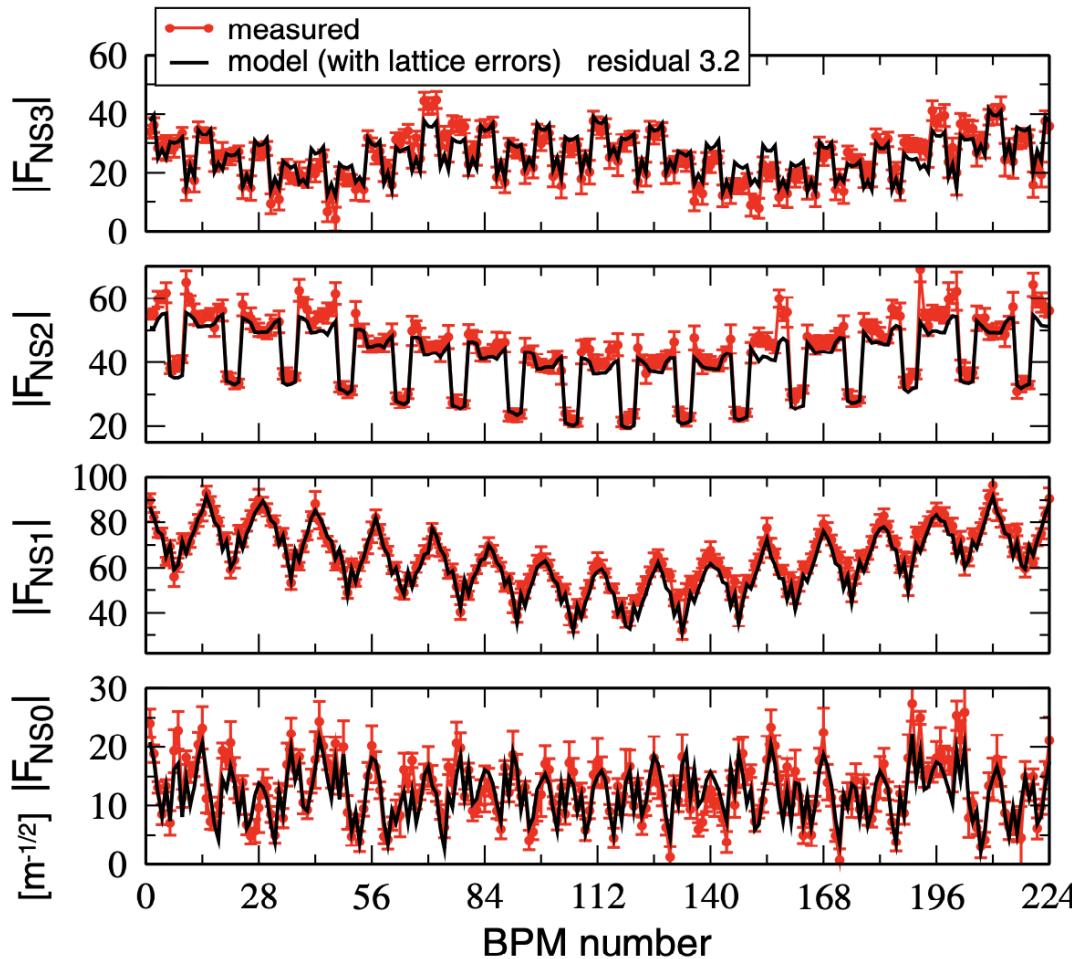
CRDTs require $(x, \cancel{p_x}), (y, \cancel{p_y})$

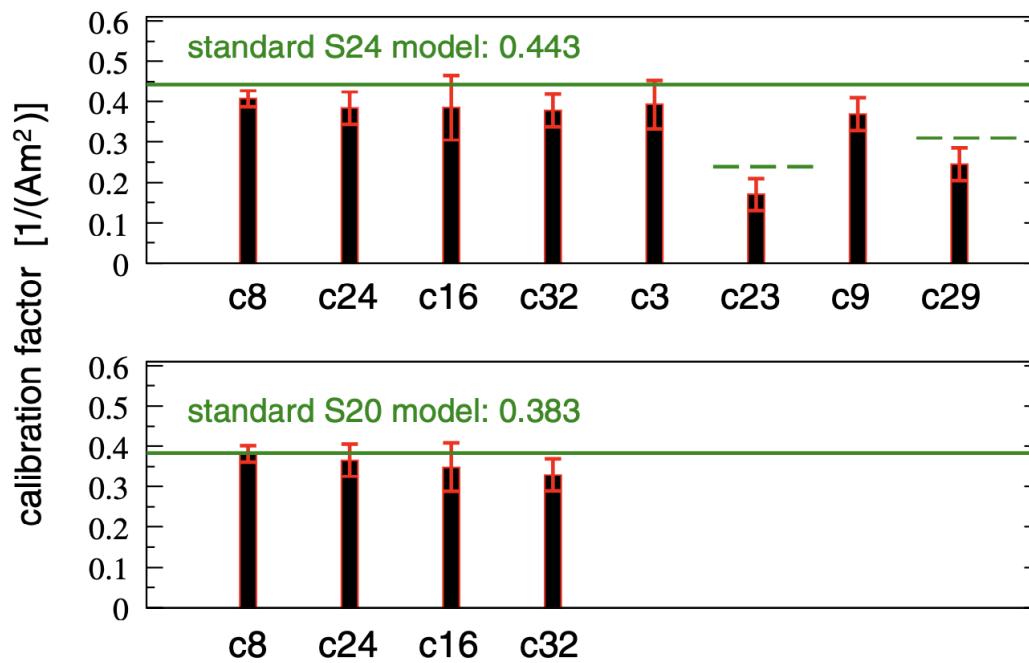
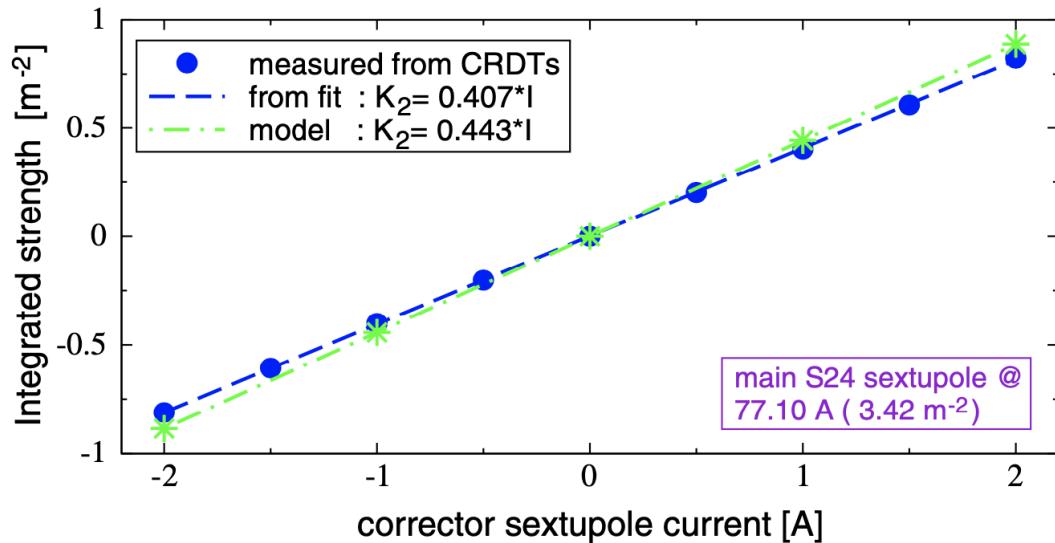
less information, but no
systematic error from $p_{x,y}$

3. from the spectral lines
infer the combined
resonance driving
terms (CRTDs) @ all
BPMs

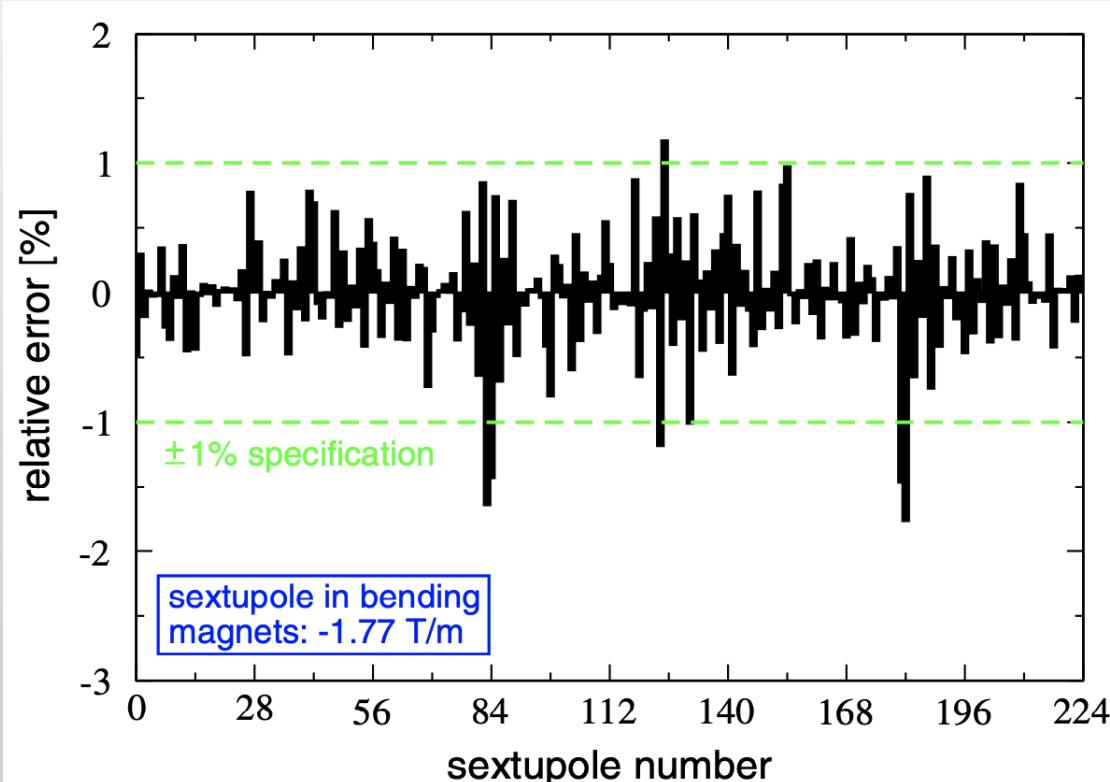
Combined RDT	Resonances	Magnetic term
$F_{xy} = f_{1001}^{(1)} - f_{1010}^{(1)*}$	$(1, 1), (1, -1)$	Skew quadrupole
$F_{yx} = f_{1001}^{(1)*} - f_{1010}^{(1)*}$	$(1, 1), (1, -1)$	Skew quadrupole
$F_{NS3} = 3f_{3000}^{(1)} - f_{1200}^{(1)*}$	$(1,0),(3,0)$	Normal sextupole
$F_{NS2} = f_{1020}^{(1)} - f_{0120}^{(1)}$	$(1, -2), (1, 2)$	Normal sextupole
...		Skew sextupole
...		Normal octupole

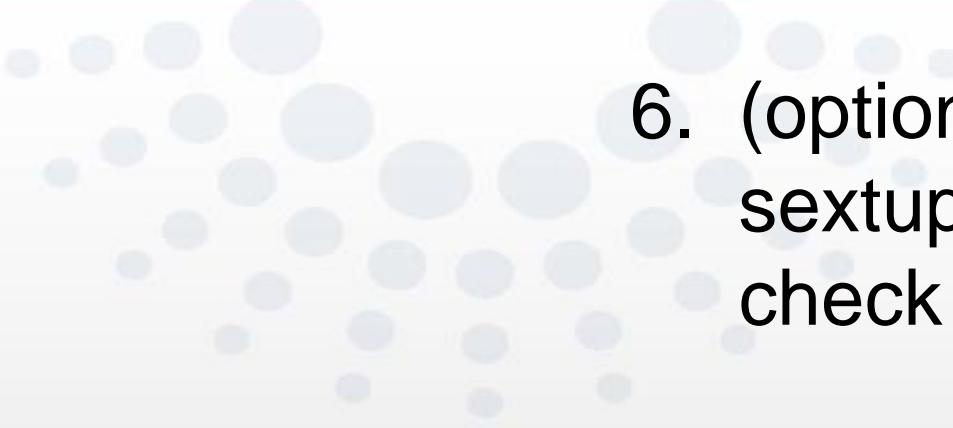
4. Compare & fit measured & model CRDTs





5. Compute your sextupole error model or calibrate individual magnets



- 
6. (optionally) correct your sextupole error model & check beam lifetime

PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS 11, 104002 (2008)

Correction of multiple nonlinear resonances in storage rings

@ Diamond light source

Applying these sextupole strength corrections to the real machine resulted in an increase in lifetime by 10%. This is a clear demonstration, in storage ring light sources, that a deterministic improvement of nonlinear beam dynamics leads to an improvement of the performance of the storage ring.

Not @ESRF: minor lifetime increase in low-intensity-per-bunch mode, detrimental for Touschek-dominated modes

6. (optionally) correct your sextupole error model & check beam lifetime

PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS 11, 104002 (2008)

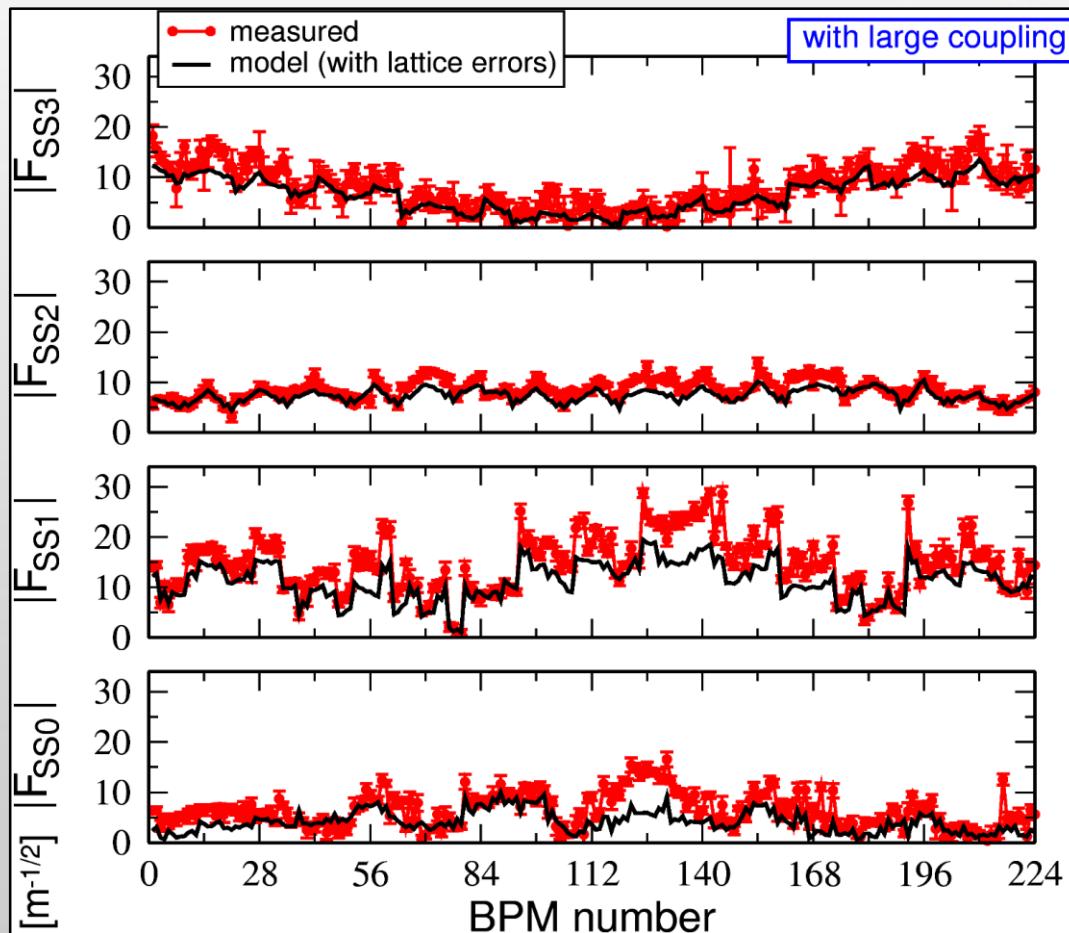
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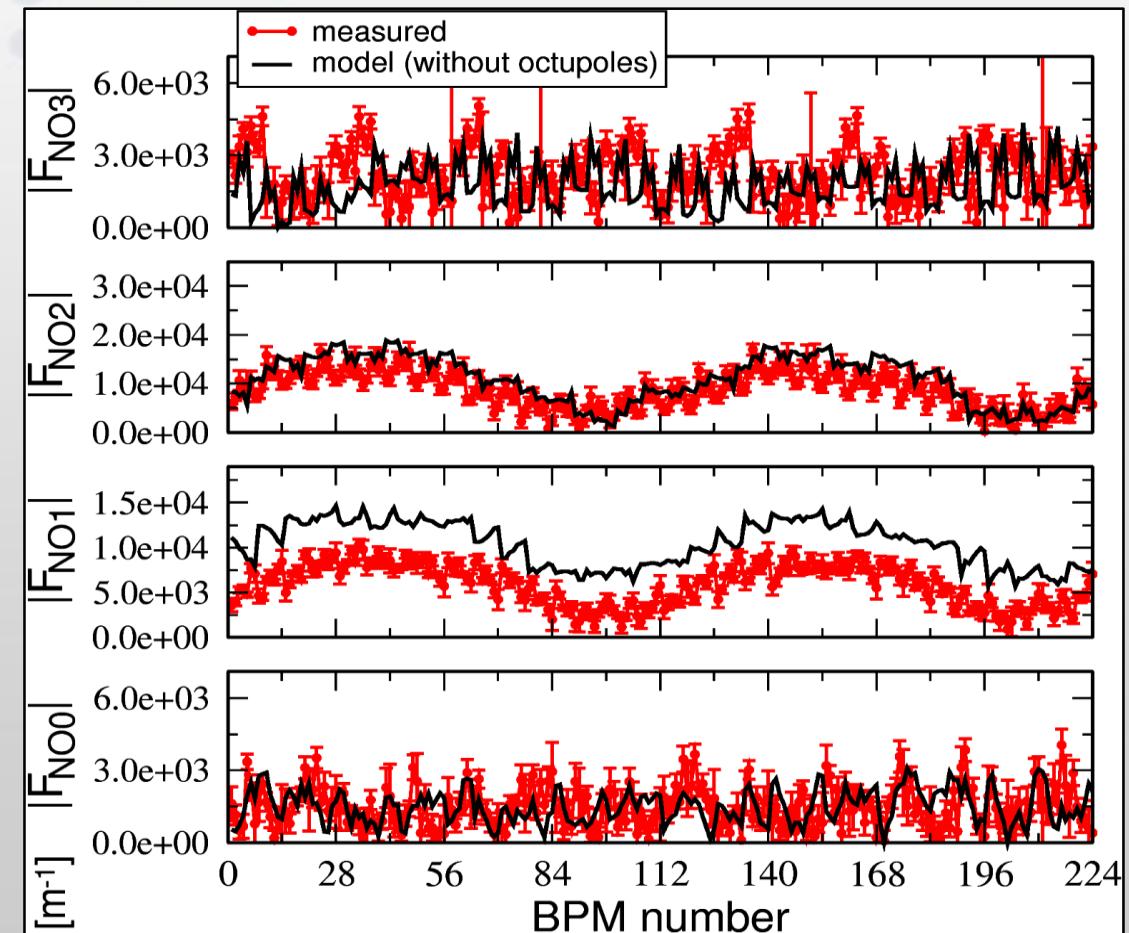
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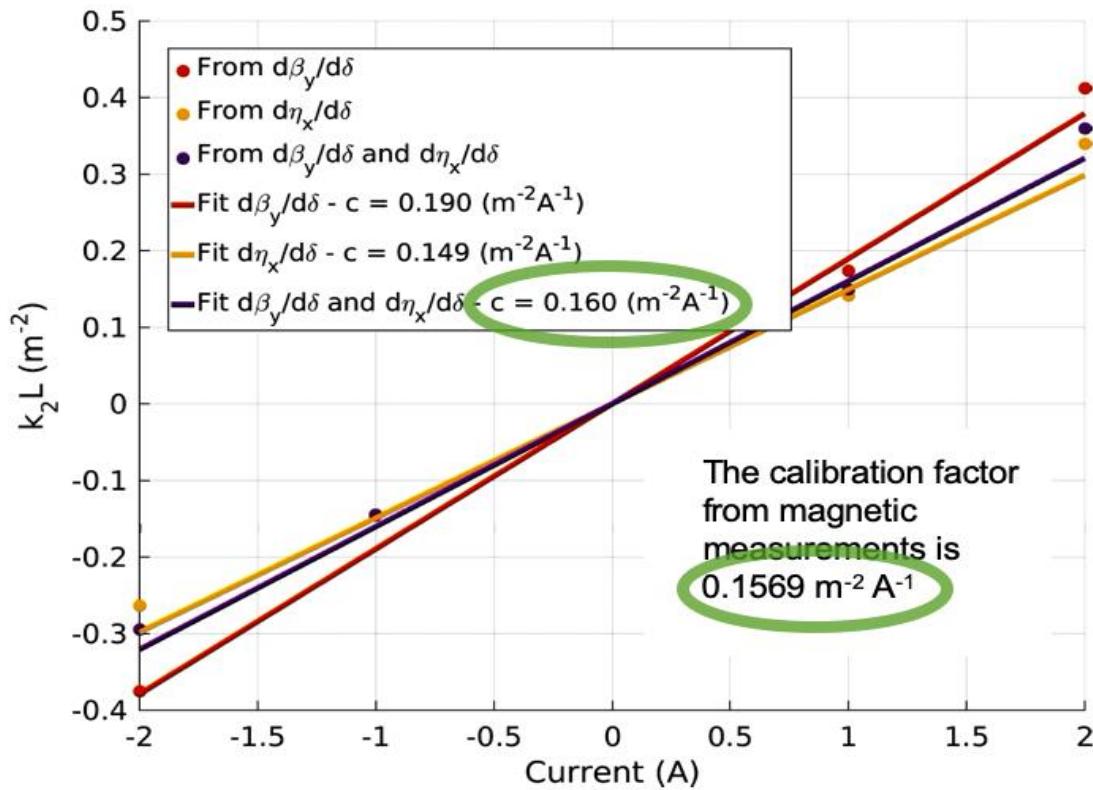
7. (for fun) measure second-order terms

skew sextupole CRDTs



normal octupole CRDTs

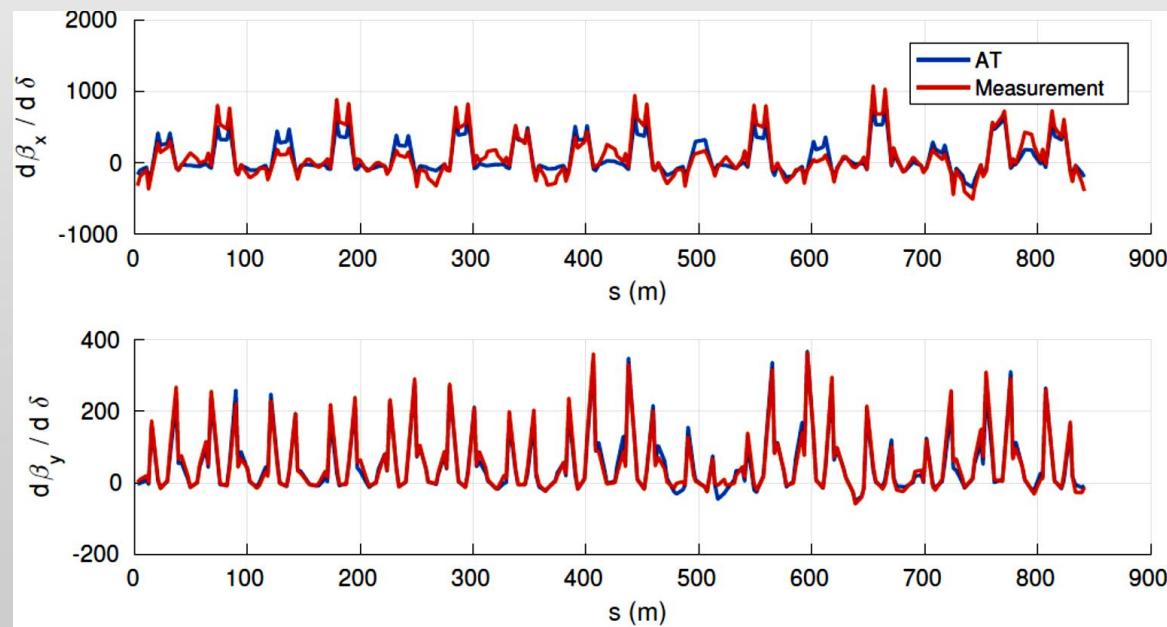




momentum compaction	α value (10^{-4})
Ideal model	1.7795
Model with errors	1.8316
ID 20	1.76 ± 0.14
ID 21	1.87 ± 0.11
hard x-ray camera	1.760 ± 0.003 (*)

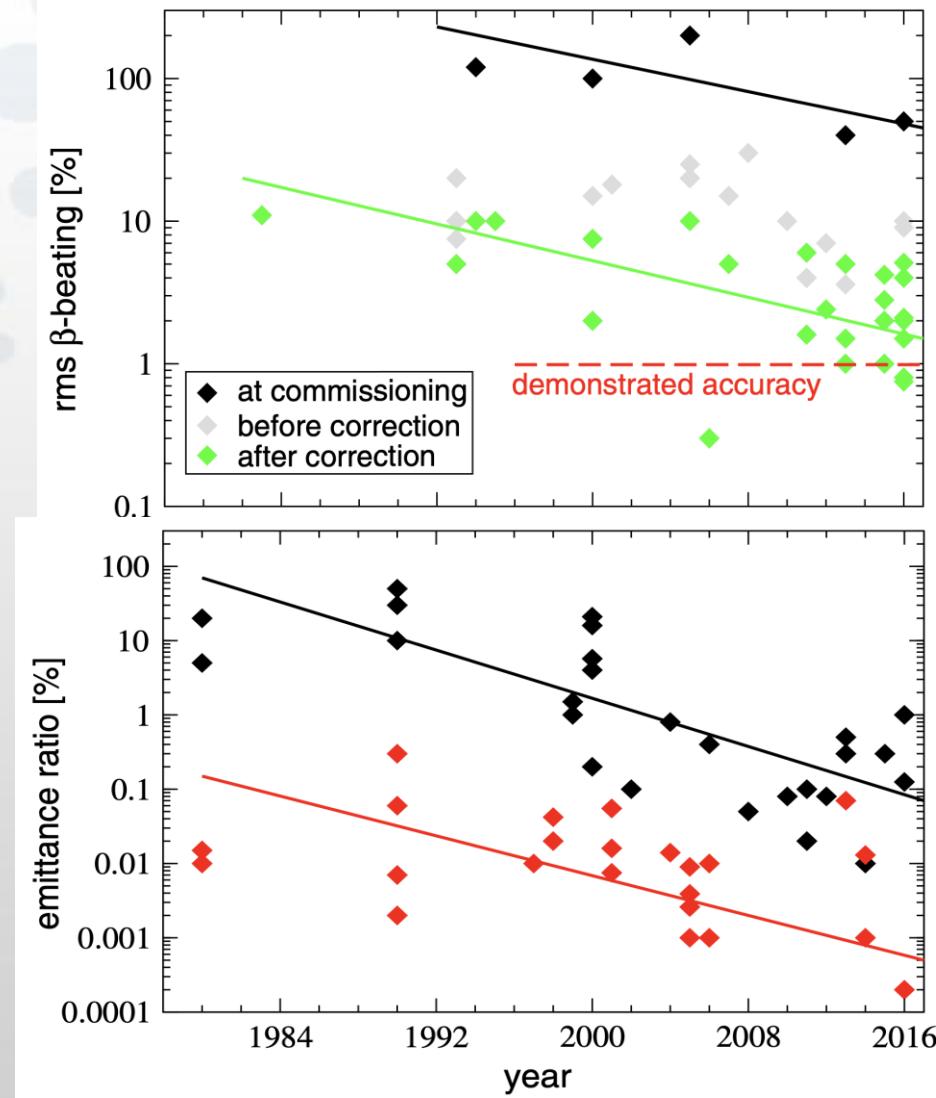
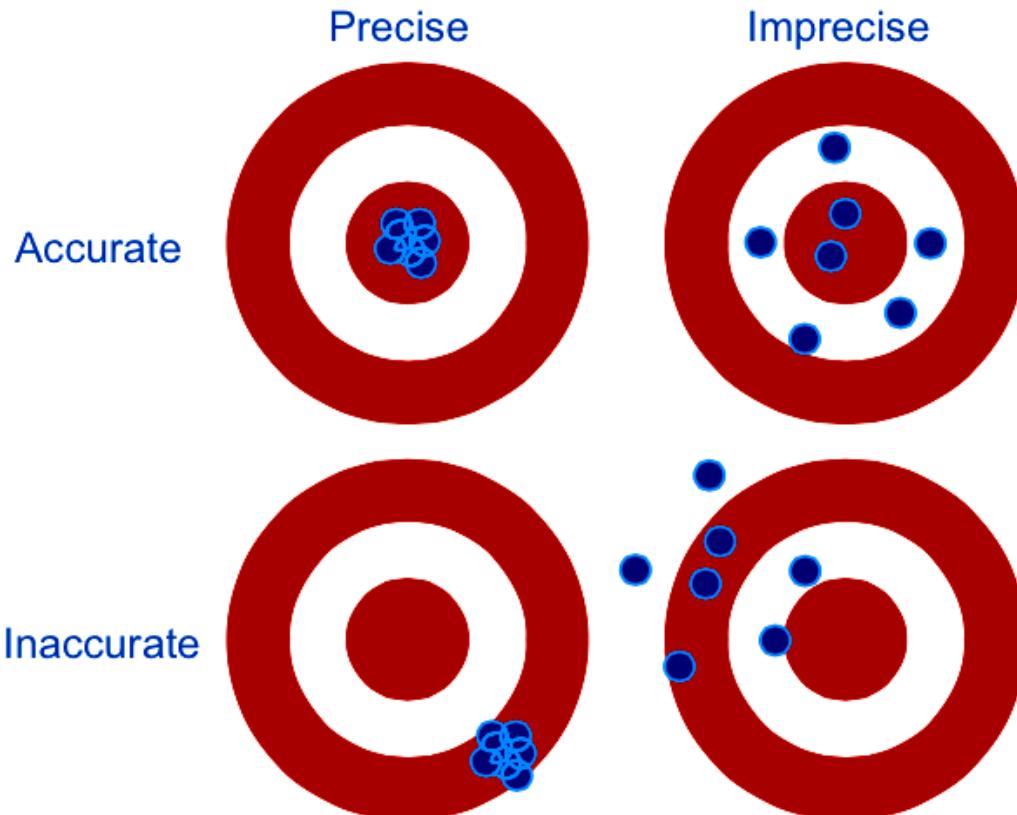
collaboration with
Laura Torino &
Nicola Carmignani

Digression: a similar beam-based sextupole calibration performed via measurement & fit of off-energy ORM (chromatic functions $d\beta/d\delta$, D' , Q')



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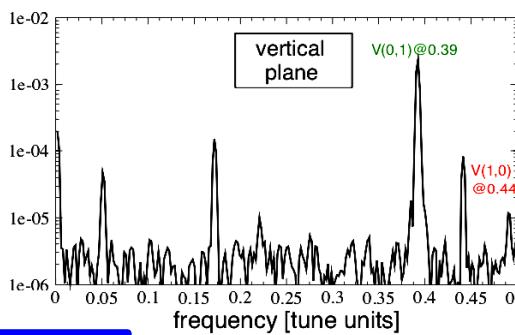
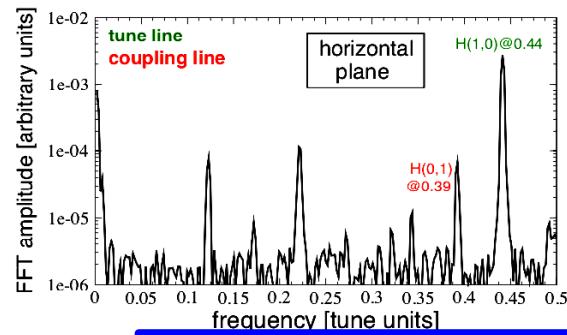
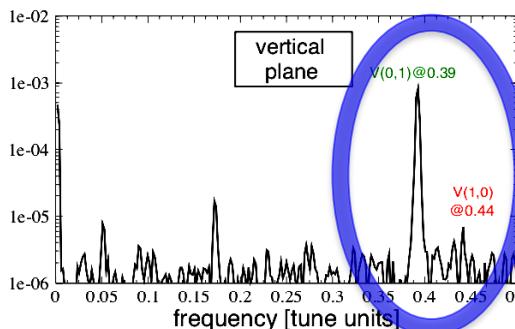
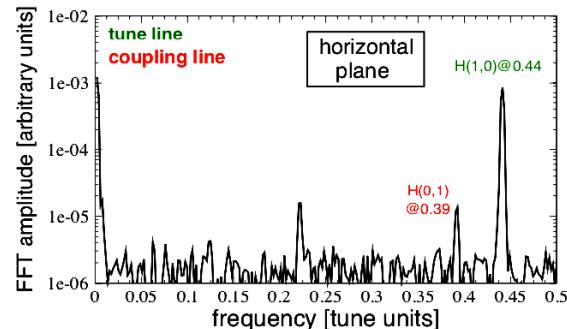
Precision / Accuracy



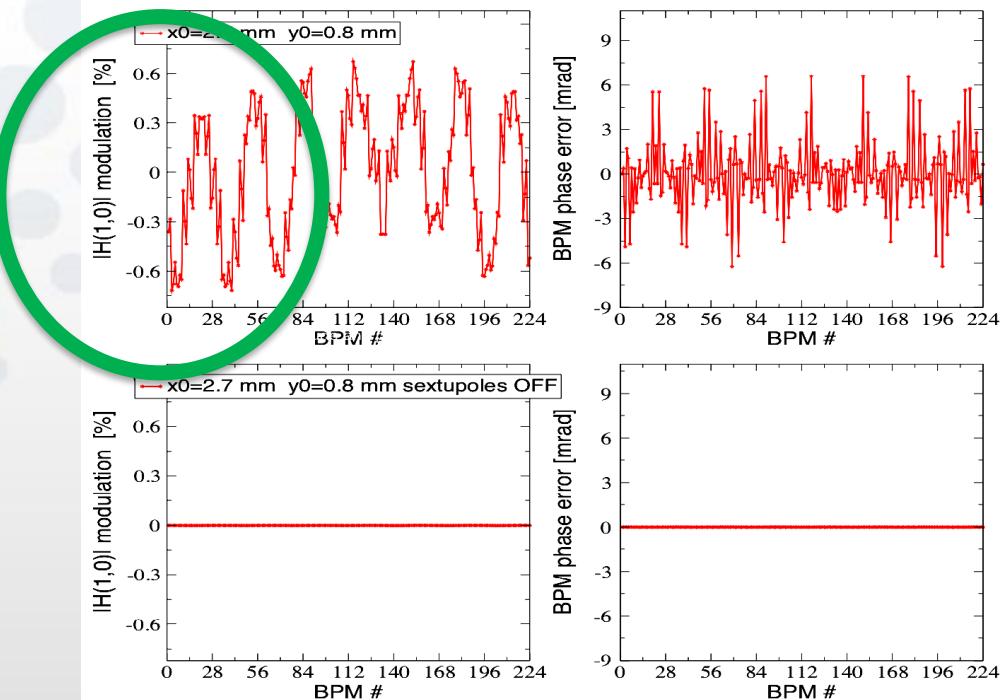
L. Malina et. al. PRAB 20, 082802 (2017)

R. Tomás et. al. PRAB 20, 054801 (2017)

Error analysis of linear optics measurements



low beam excitation
incompatible with
ultra-low coupling



large beam excitation
incompatible with strong
chromatic sextupoles

<https://arxiv.org/abs/1603.00281>

Error analysis of linear optics measurements

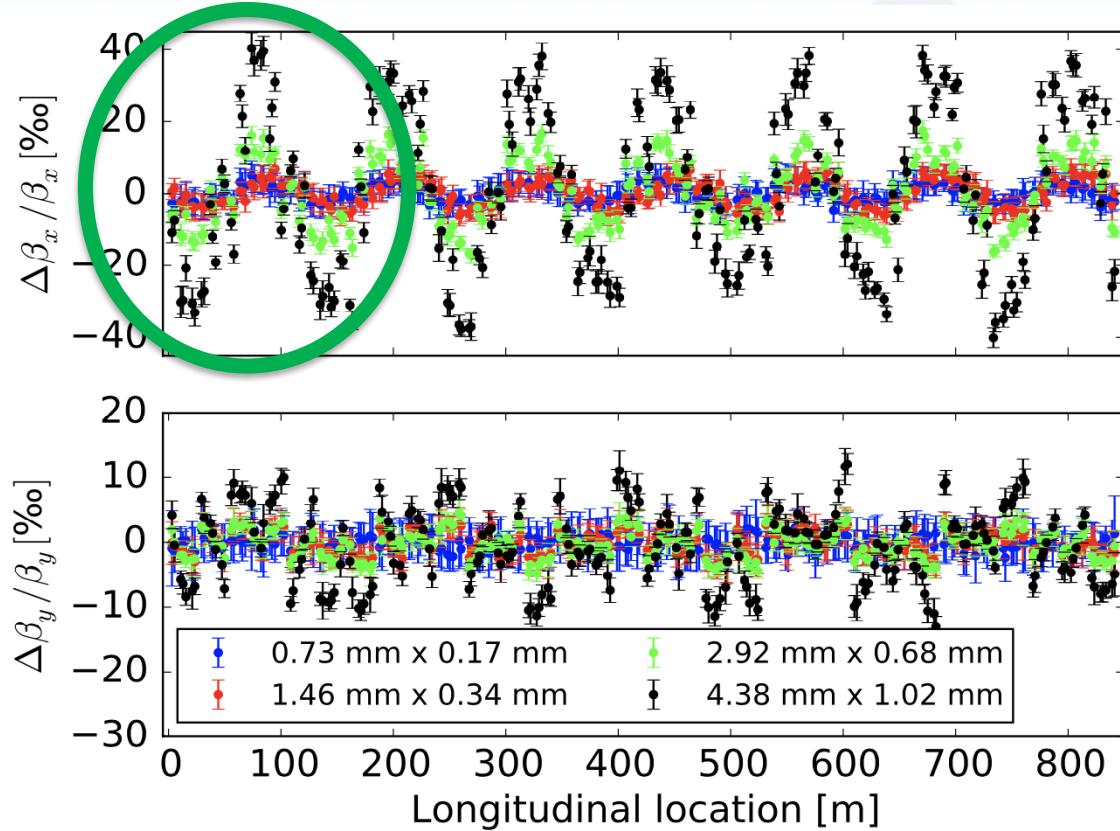
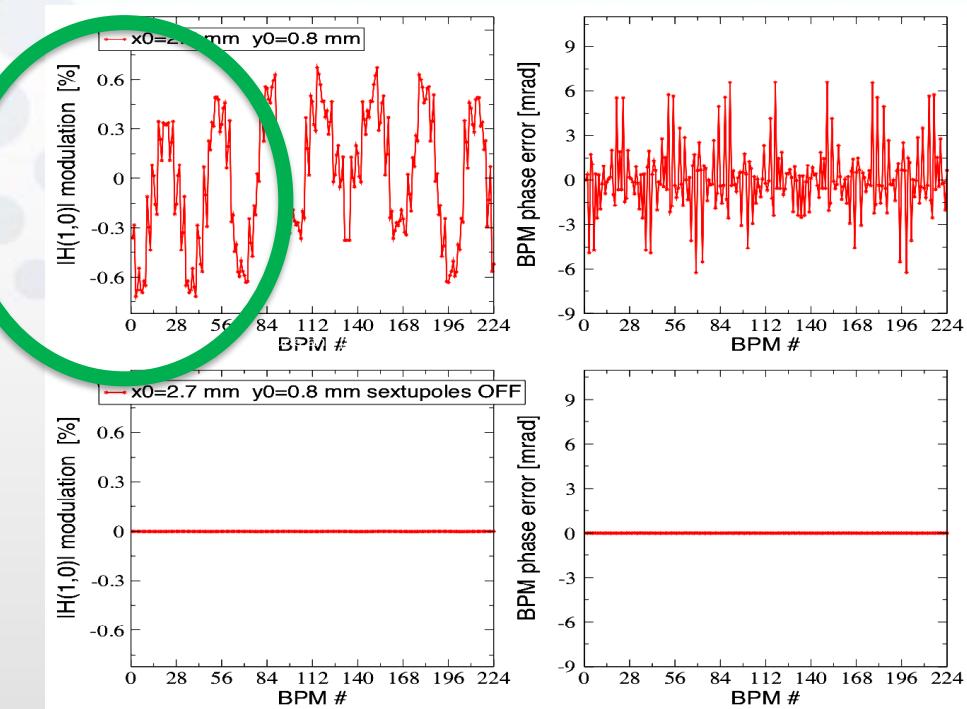
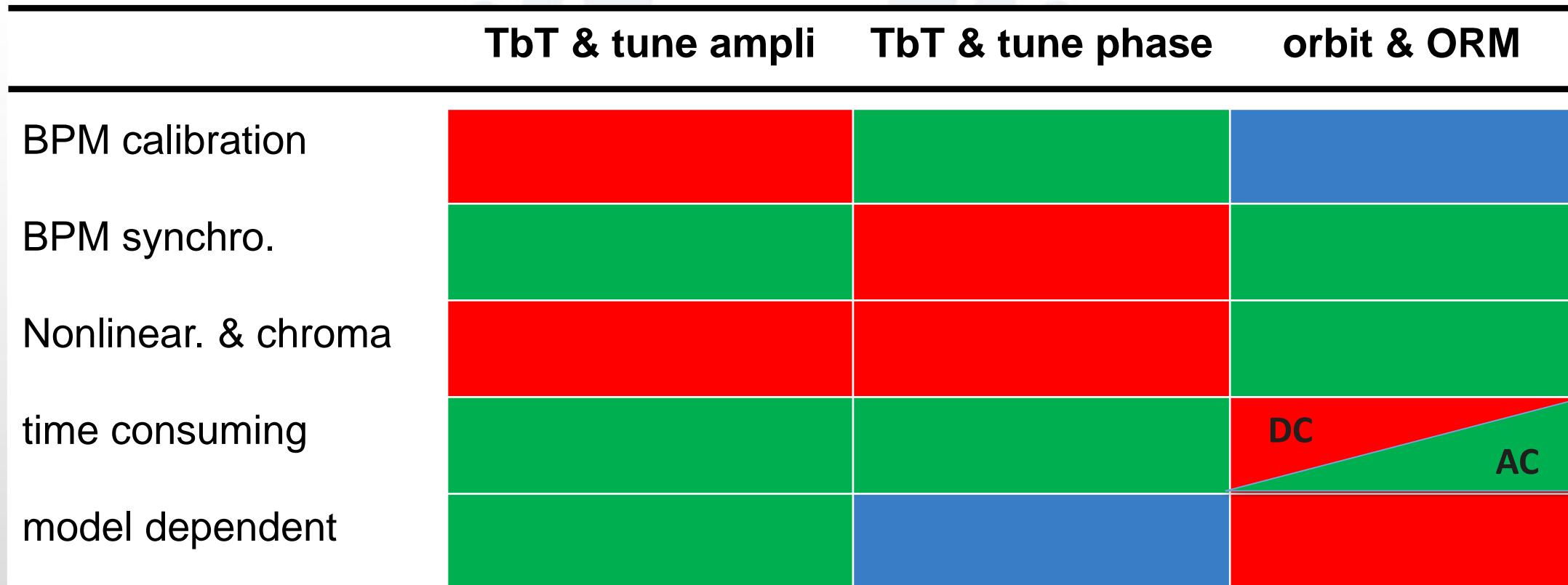


FIG. 4. Simulated artificial β -beating computed by N-BPM from single-particle simulations of the ESRF storage ring lattice



large beam excitation
incompatible with strong
chromatic sextupoles



L. Malina et. al. PRAB 20, 082802 (2017)

measured β -functions
best precision TbT: 0.4%
best precision ORM: 0.5%
best accuracy ORM-TbT: ~1%

	TbT & tune ampli	TbT & tune phase	orbit & ORM
BPM calibration	Red	Green	Blue
BPM synchro.	Green	Red	Green
Nonlinear. & chroma	Red	Red	Green
time consuming	Green	Green	DC AC
model dependent	Green	Blue	Red

	TbT & tune ampli	TbT & tune phase	orbit & ORM
$\beta_1^{(meas)} = \beta_1^{(mod)} \frac{\cot \Delta\phi_{12}^{(meas)} - \cot \Delta\phi_{13}^{(meas)}}{\cot \Delta\phi_{12}^{(mod)} - \cot \Delta\phi_{13}^{(mod)}} + O(\delta K_1)$			
Castro's formula (no error δk_1 between BPMs, model needed)			
time consuming			DC
model dependent		AC	

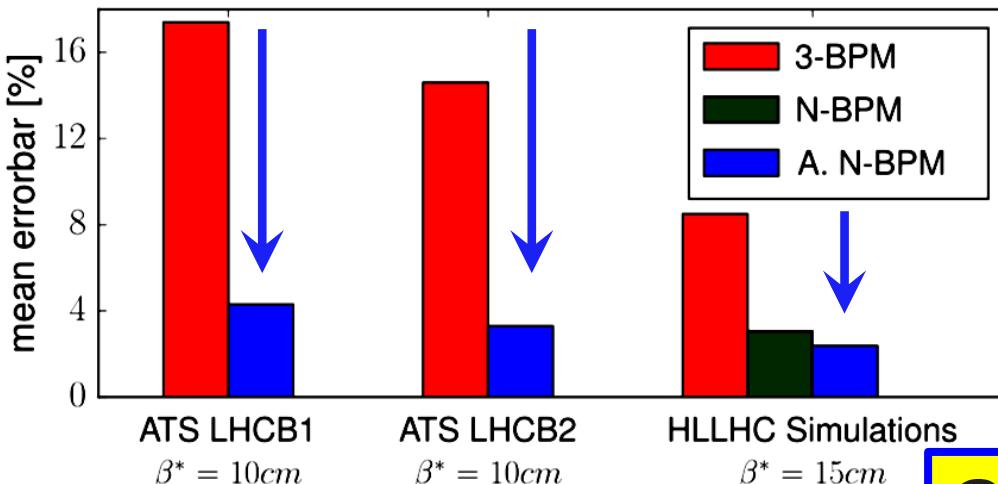
TbT & tune ampli

TbT & tune phase

orbit & ORM

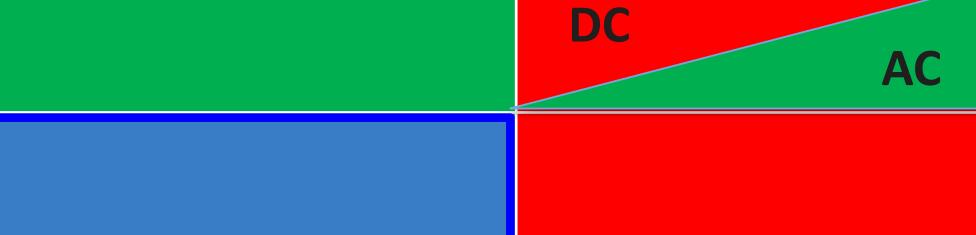
$$\beta_1^{(meas)} = \beta_1^{(mod)} \frac{\cot \Delta\phi_{12}^{(meas)} - \cot \Delta\phi_{13}^{(meas)}}{\cot \Delta\phi_{12}^{(mod)} - \cot \Delta\phi_{13}^{(mod)} + (\bar{h}_{12} - \bar{h}_{13})} + O(\delta K_1^2)$$

New formula (with error δk_1 between BPMs, model needed)

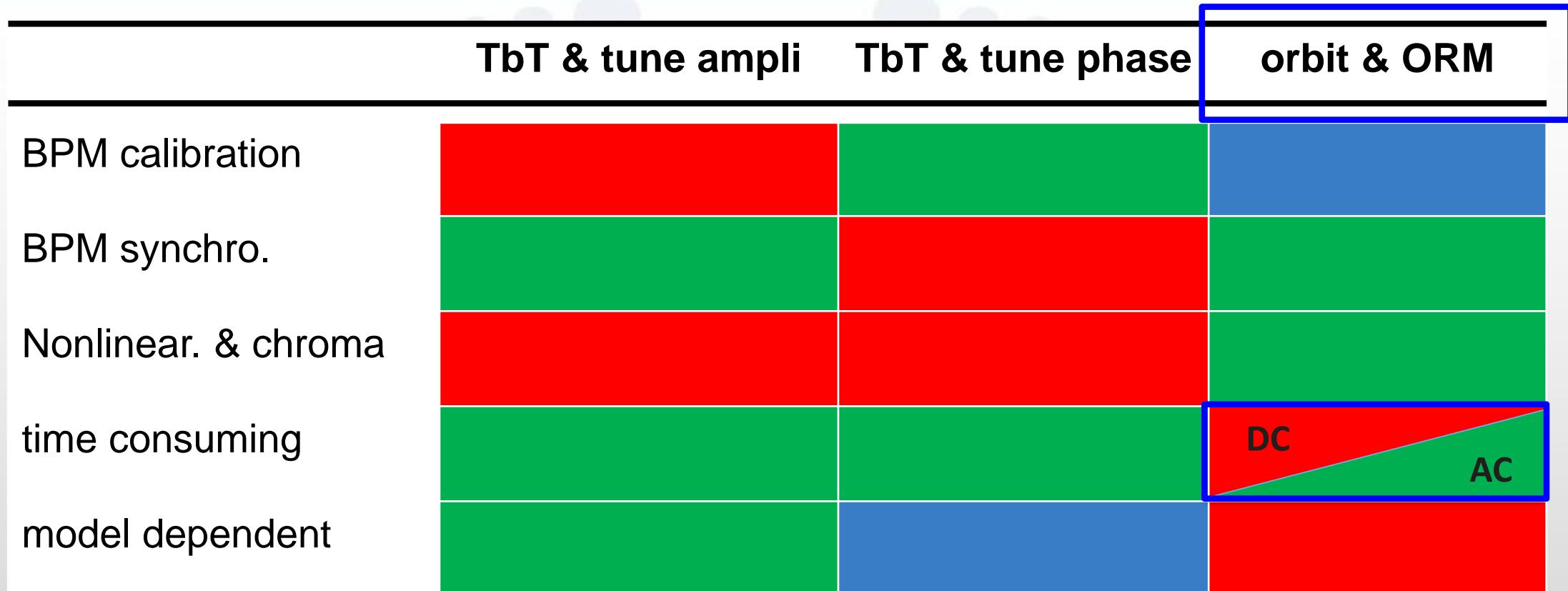


@LHC

A. Wegscheider et. al. PRAB 20, 111002 (2017)

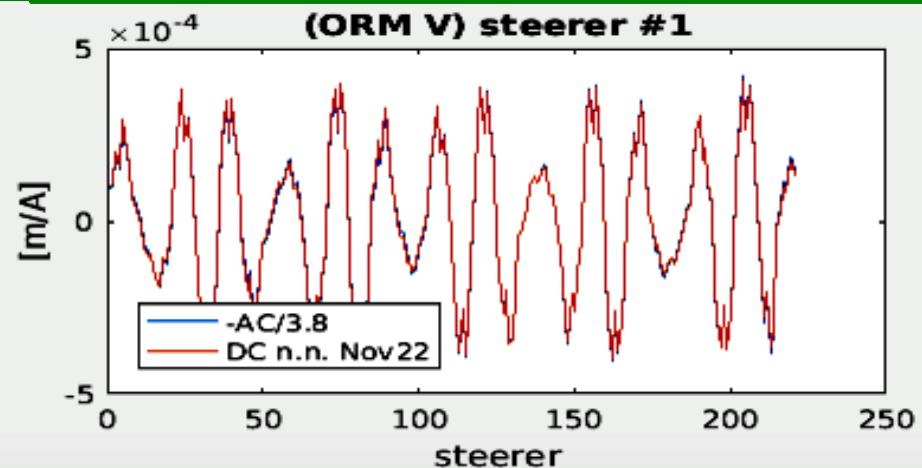


$$\bar{h}_{ij} = \mp \frac{\sum_{i < w < j} \beta_w^{(mod)} \delta K_{w,1} \sin^2 \Delta\phi_{wj}^{(mod)}}{\sin^2 \Delta\phi_{ij}^{(mod)}}$$



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TbT & tune ampli TbT & tune phase orbit & ORM

ESRF ORM column: AC Vs DC

measurement
analysis

$$\begin{pmatrix} \delta\vec{O}^{(xx)} \\ \delta\vec{O}^{(yy)} \\ \delta\vec{D}_x \end{pmatrix} = \mathbf{N} \begin{pmatrix} \delta\vec{K}_1 \\ \delta\vec{K}_0 \end{pmatrix}$$

$$\begin{pmatrix} \delta\vec{O}^{(xy)} \\ \delta\vec{O}^{(yx)} \\ \delta\vec{D}_y \end{pmatrix} = \mathbf{S} \begin{pmatrix} \vec{J}_1 \\ \vec{J}_0 \end{pmatrix}.$$

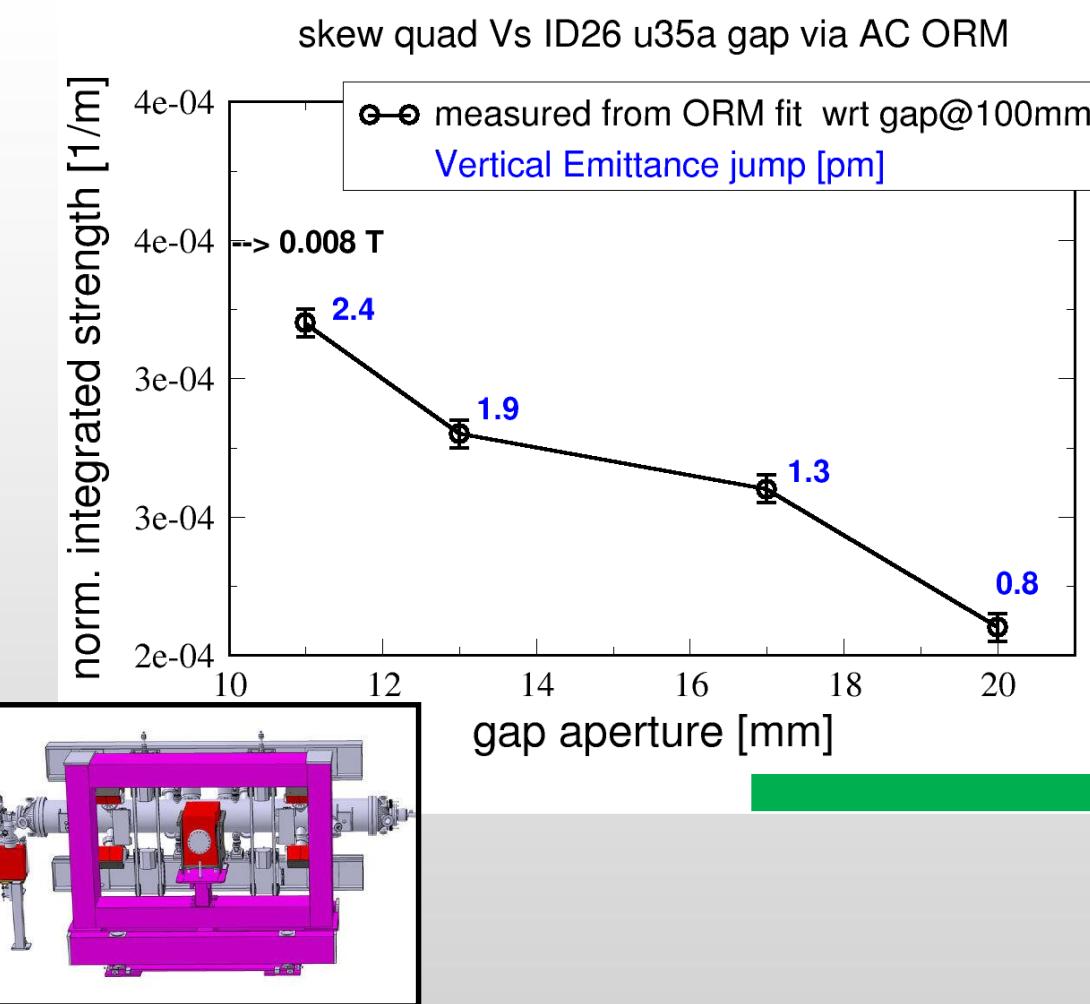
DC
AC

Orbit response matrix (ORM) @ESRF:
 ~1 h (DC steerer excitation + numerical response*)
 ~3' (AC steerer excitation + analytical response*)

<http://arxiv.org/abs/1711.06589>

* response : **N** & **S** (large matrices)
(AC excitation pioneered @Diamond)

Applications of AC Orbit data



beam-based measurement of skew-quadrupole fields induced by undulators & skew quad correctors

Orbit response matrix (ORM) @ESRF:
~1 h (DC steerer excitation)
~3' (AC steerer excitation)

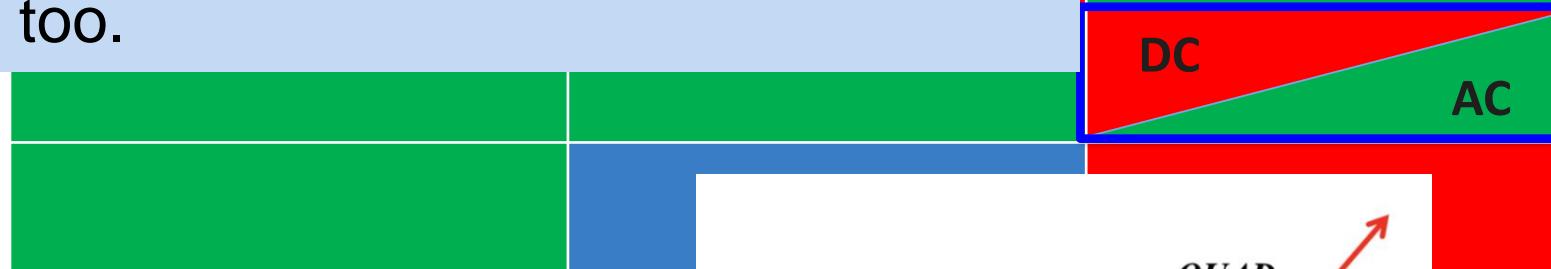
Digression: fast AC beam-based quadrupole alignment derived & applied to ALBA light source (120 quadrupoles):

- serial DC steerer excitation: 5h, precision $\sim 50 \mu\text{m}$
- parallel AC steerer excitation: 10', precision $\sim 15 \mu\text{m}$

applied to sextupoles too.

model dependent

orbit & ORM



Z. Marti e. al. PRAB 23, 012802 (2019)

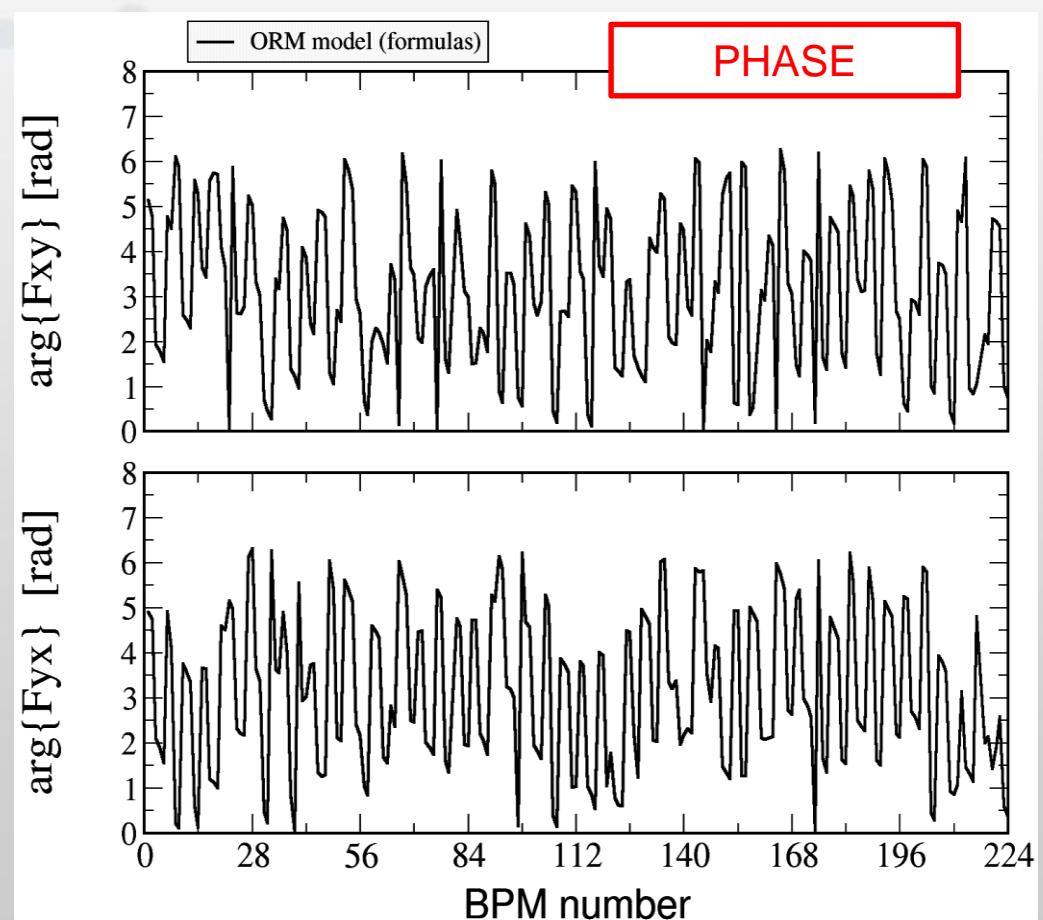
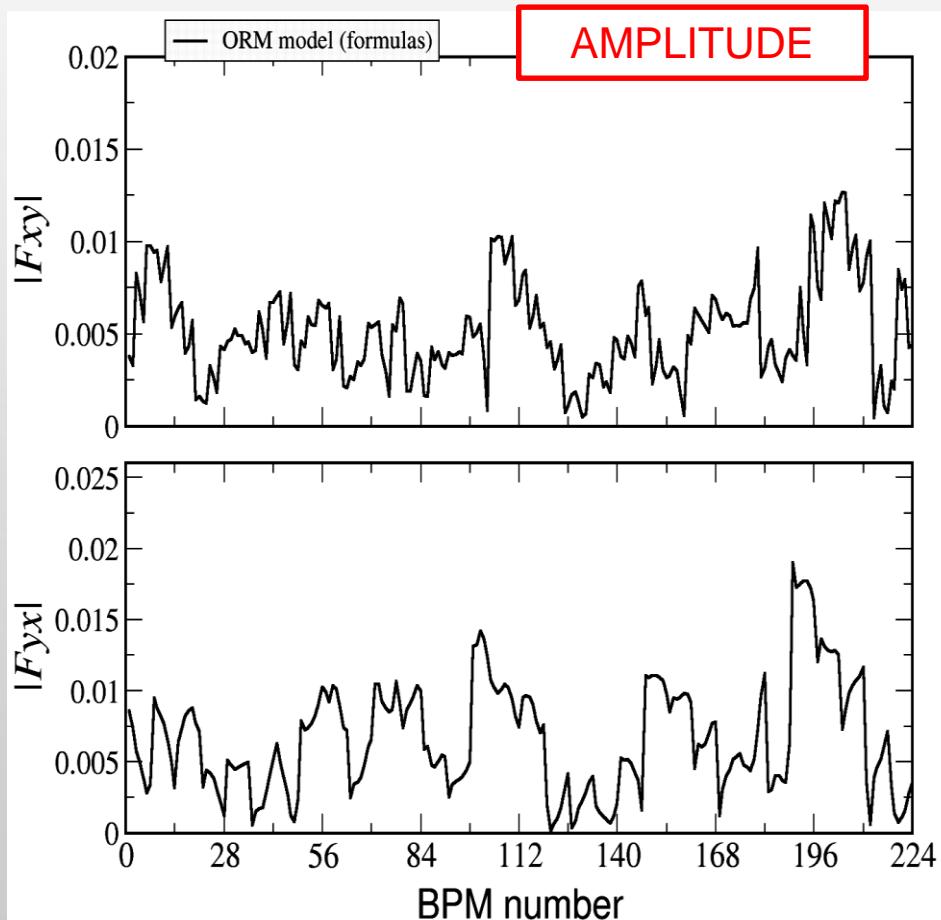
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Betatron coupling described by two CRDTs

$$F_{xy} = (f_{1001} - f_{1010}^*) \quad \& \quad F_{yx} = (f_{1001}^* - f_{1010}^*)$$

Measurement with **low chroma (0,0) & detuning sext. optics**

compare ($\epsilon_y/\epsilon_x \sim 1\%$) ORM model ...

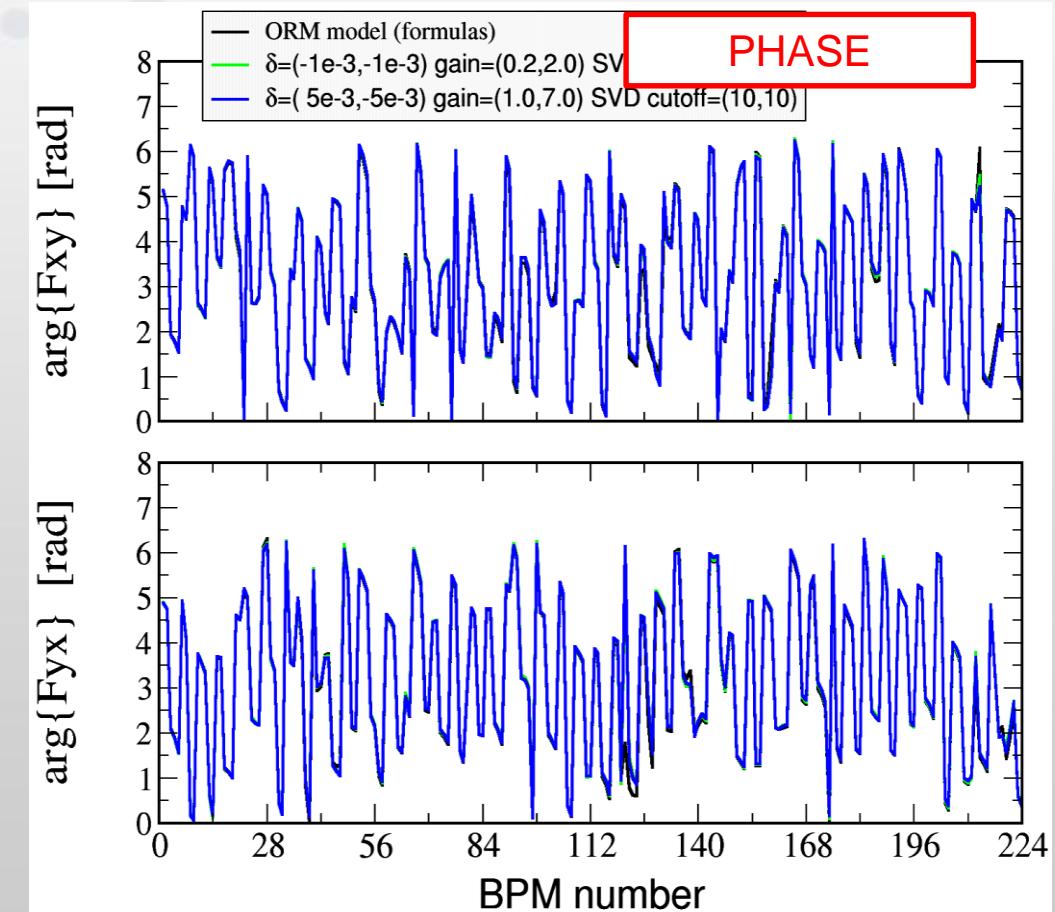
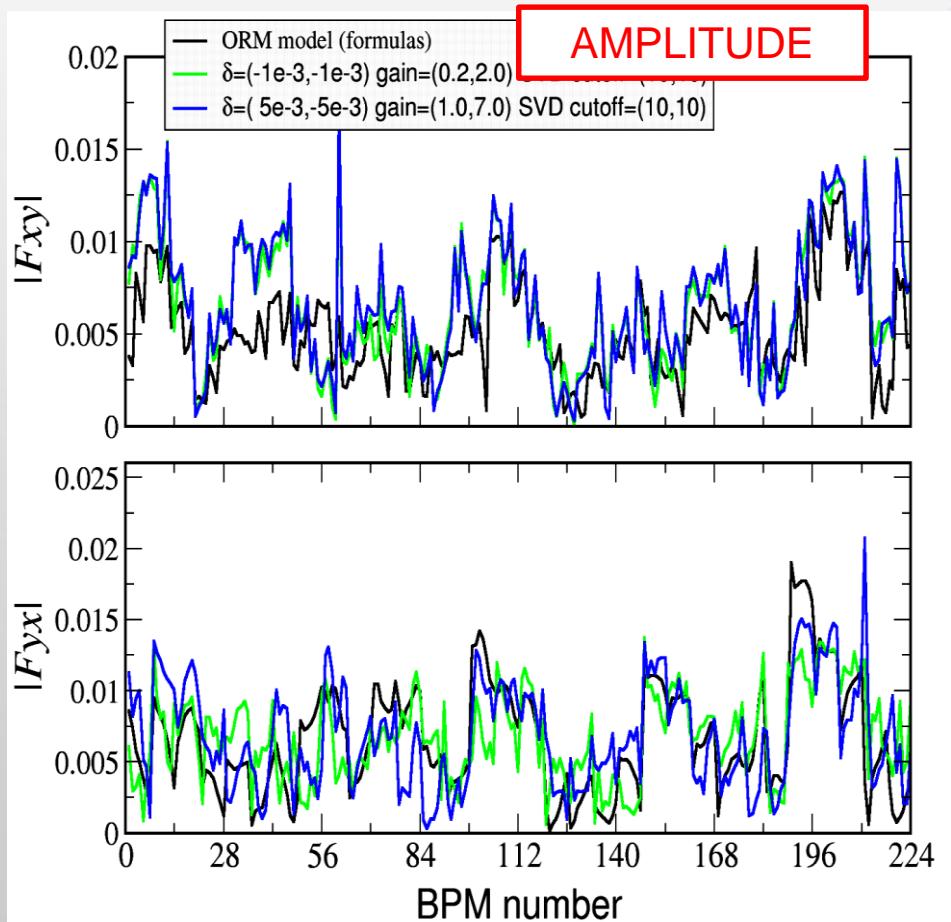


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compare ($\epsilon_y/\epsilon_x \sim 1\%$) ORM model with TbT harmonic analysis

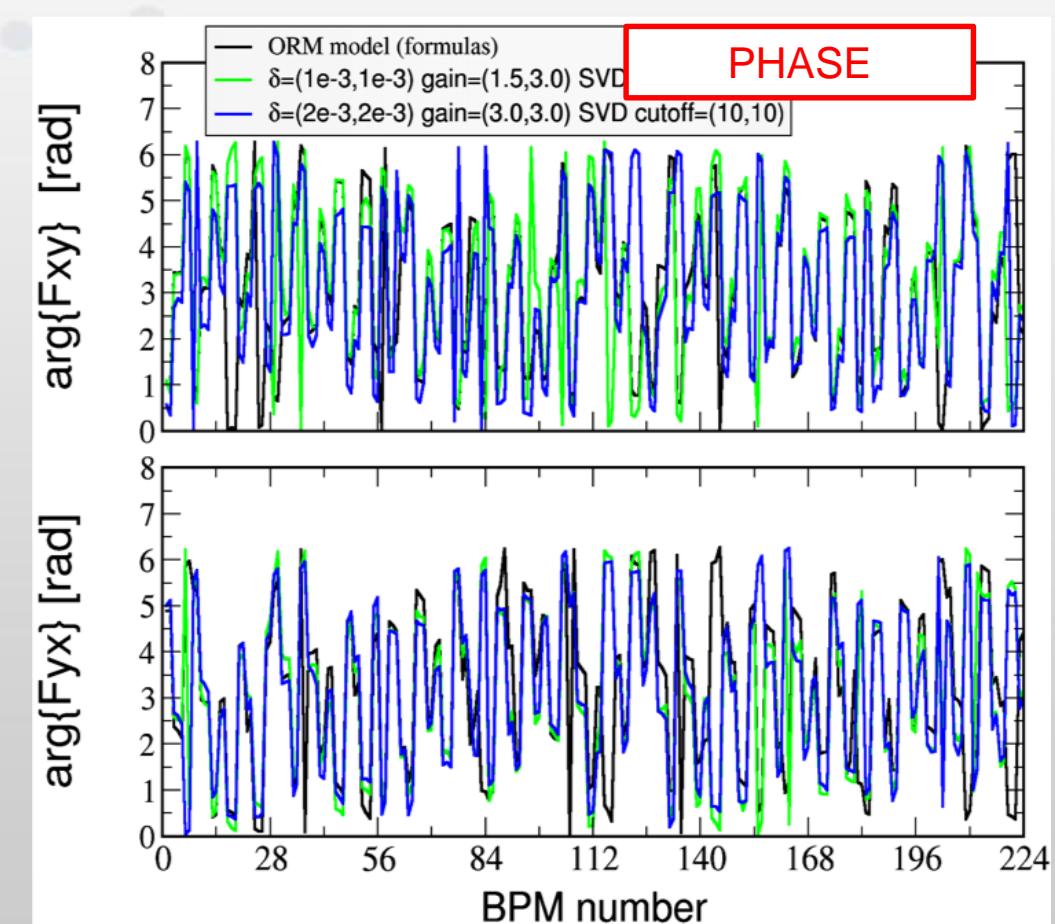
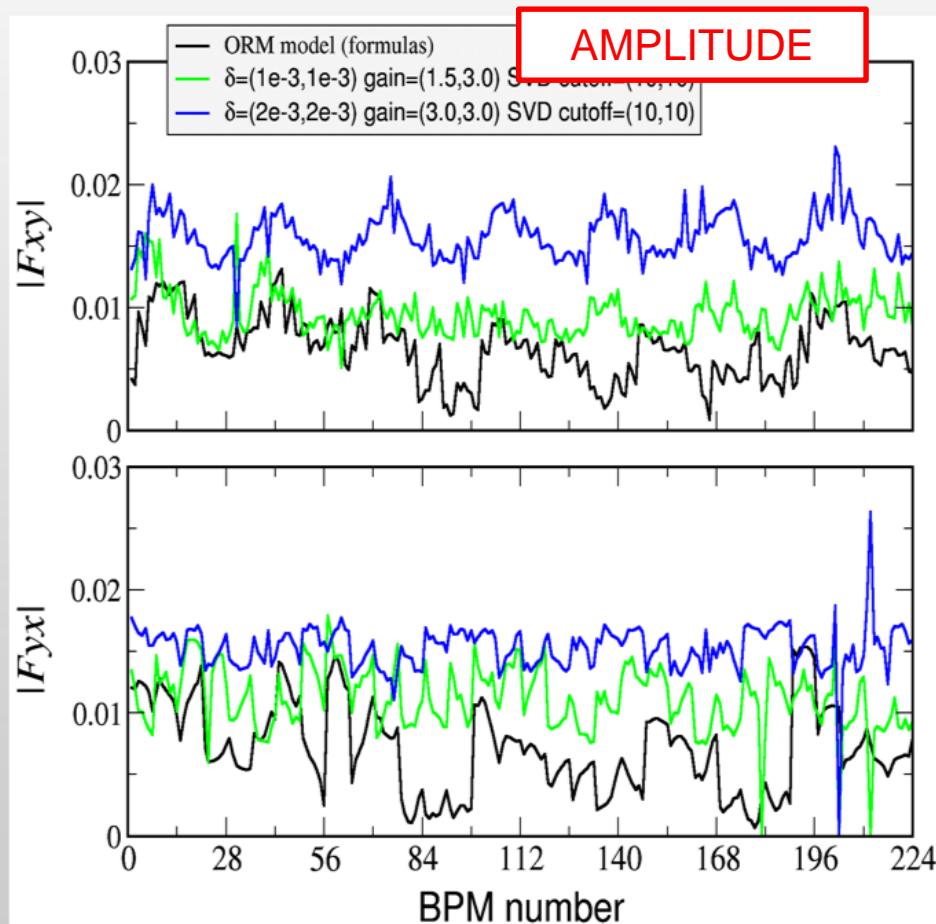


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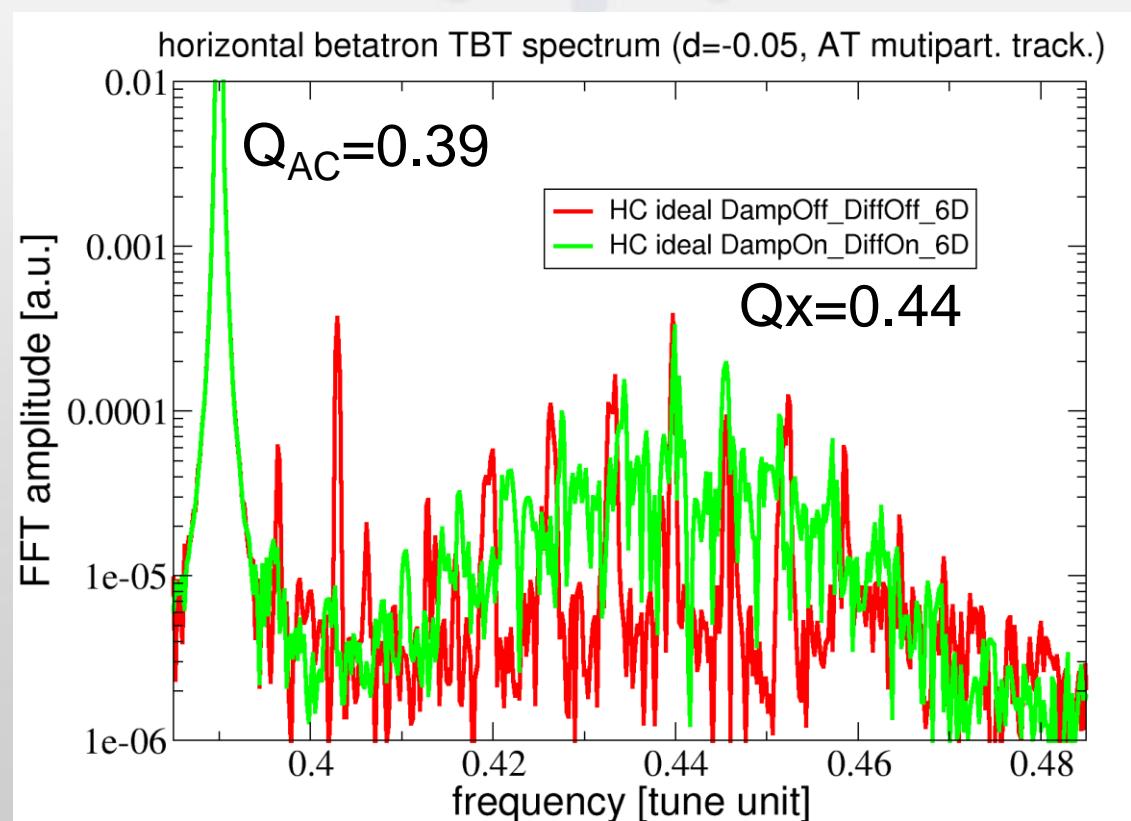
Measurement with **large chroma (8,13) operational optics**

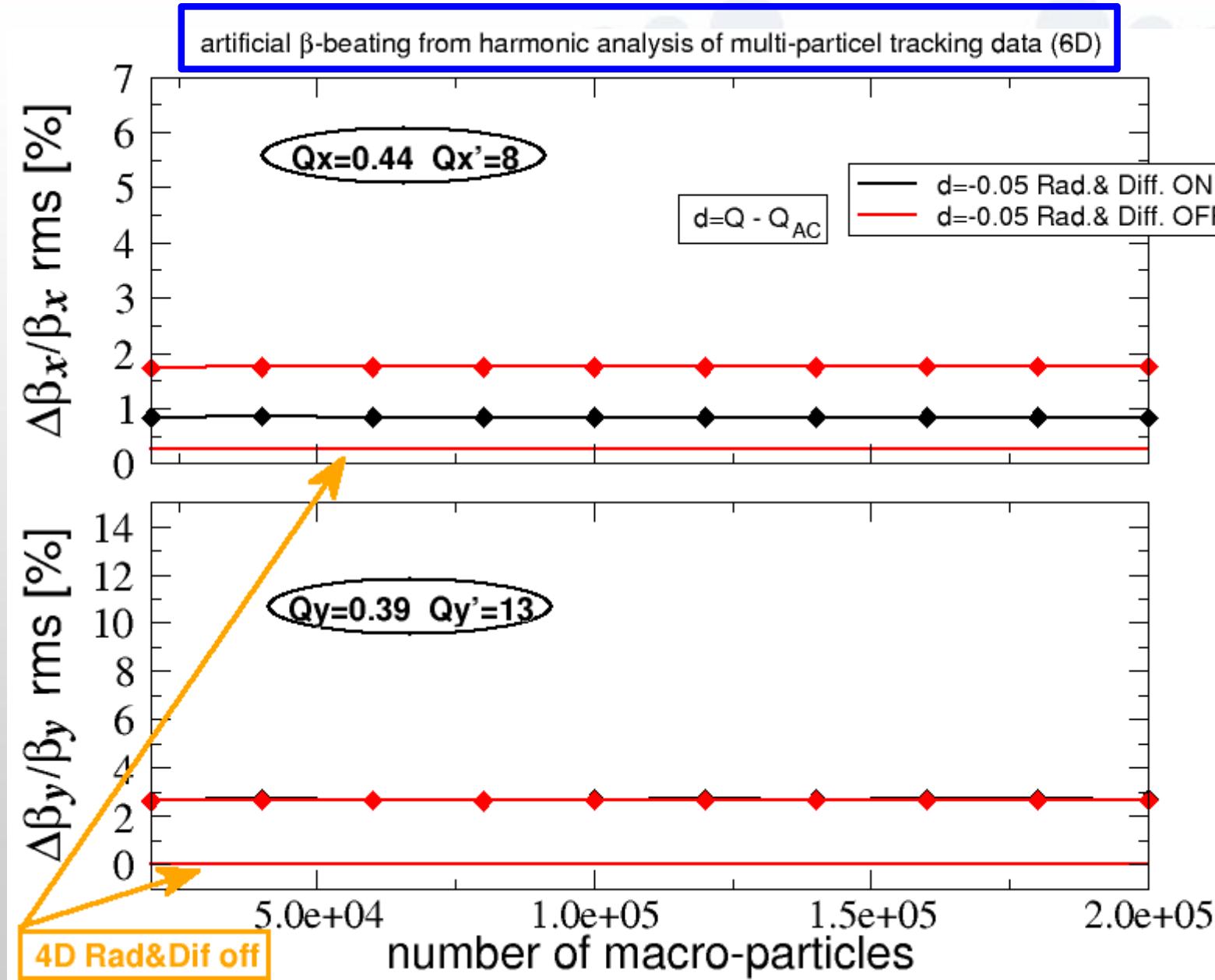
compare ($\epsilon_y/\epsilon_x \sim 1\%$) ORM model with TbT harmonic analysis



1st experimental observation: AC dipole tuning tedious with high chromaticity: beam lost with setting used for low-chroma optics.

High (8,13) chroma sextupole optics: RAD OFF & ON (multi-particle simulations)





2nd experimental & numerical observation:
synchrotron radiation + diffusion & high chroma
=> limited accuracy.

AC dipole & data cleaning OK for low-chroma lepton rings.

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- Nonlinear optics & magnet calibration via turn-by-turn (TbT) BPM data (2010-2013)
- Error analysis of linear optics measurements via orbit & TbT analysis (2016-2017)
- Applications of AC Orbit data (2016-2019)
- Experience of AC dipole & TbT BPM data (2016-2018)

Thank you for your attention