News from the strong interactions program of NA61/SHINE

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NA61/SHINE experiment

Strong interactions program:

- search for the critical point of strongly interacting matter,
- study of the properties of the onset of deconfinement.
- **•** Fixed-target experiment at CERN SPS.
- Large variety of beams and targets.
- Large acceptance: full forward hemisphere, down to $p_T=0$.

 $~13 m$

- **•** Particle identification: dE/dx in Time Projection Chambers, Time of Flight detector.
- Collision centrality measured by forward Projectile Spectator Detector (PSD).

NA61/SHINE, JINST 9, P06005, 2014

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Onset of deconfinement: step

- Qualitatively similar energy dependence is seen in p+p, Be+Be, Ar+Sc and Pb+Pb.
- Magnitude of T increases with the system size.

- NA61/SHINE, EPJC 81, 1, 73, 2021 and Ar+Sc preliminary results, APPB 30, 2705, 1999
- Sensitive to both the temperature and the radial flow.
- Kaons are only weakly affected by re-scattering and resonance decays during the post-hydro phase (at SPS and RHIC energies).
- Connected with the temperature of the freeze-out surface and not the early-stage fireball.

Onset of deconfinement: horn

- \bullet Be+Be close to $p+p$ in K^{+}/π^{+} .
- No horn-like structure in Ar+Sc.

 \bullet NA61/SHINE, EPJC 81, 1, 73, 2021 and Ar+Sc preliminary results

- **•** $p+p \approx Be+Be \neq Ar+Sc \ll Pb+Pb$
- Good measure of the strangeness to entropy ratio which is different in the confined phase (hadrons) and the QGP (quarks, anti-quarks and gluons) \rightarrow probe of the onset of deconfinement.

Results for $p+p$ interactions

- The sharp break in K^+/π^+ and inverse slope parameter T in p+p collisions at SPS energies.
- The break energy is \approx 7 GeV, close to the energy of the onset of deconfinement ≈ 8 GeV.
- The UrQMD model does not reproduce the sharpness of the break.

NA61/SHINE, PRC 102, 1, 011901, 2020

Results for Be+Be interactions

- The first world data for Be+Be collisions.
- No visible sharp break in K ⁺*/π*⁺ and inverse slope parameter T. Note the limited energy range of data.
- No models which describe all measured quantities.

NA61/SHINE, EPJC 81, 1, 73, 2021

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Onset of deconfinement: kink

- \bullet N+N interactions agree well with the world data.
- \bullet Be+Be collisions are mostly between measurements from $N+N$ and $Pb+Pb$ collisions.
- Ar+Sc collisions seem to be systematically higher than the results for N+N, Be+Be and Pb+Pb collisions at the lower energies.
- Ar+Sc close to the Pb+Pb results at the highest energies.

NA61/SHINE, EPJC 81, 5, 397, 2021

Width of the rapidity distribution

NA61/SHINE, EPJC 81, 5, 397, 2021
 \bullet Collision energy dependence of the width was derived by Shuryak (E. V. Shuryak. Yad.Fiz., 16, 395, 1972) from the Landau hydrodynamical model of high energy collisions:

$$
\sigma^2 = \frac{8}{3} \cdot \frac{c_s^2}{1 - c_s^4} \cdot \ln\left(\frac{\sqrt{s_{NN}}}{2m_p}\right),\tag{1}
$$

where c_s denotes the speed of sound.

- The dense matter produced in the collisions was predicted to show a minimum in the speed of sound energy dependence around the collision energy of the onset of deconfinement.
- \bullet Confirmed by Pb+Pb data in combination with results from central Au+Au collisions.
- The results of NA61/SHINE from central Ar+Sc, Be+Be collisions, and inelastic N+N reactions need to be extended to lower end energies for conclusion about a possible minimum.

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Onset of fireball – rapid change of observables when going from small $(p+p, Be+Be)$ to intermediate $(Ar+Sc)$ and large ones $(Pb+Pb) \rightarrow$ beginning of the creation of large clusters of strongly interacting matter?

None of the models reproduce K^+/π^+ ratio nor ${\cal T}$ for whole $\langle W\rangle$ range.

PHSD: EPJA 56, 9, 223, 2020, arXiv:1908.00451 and private communication; SMASH: JPG 47, 6, 065101, 2020 and private communication; UrQMD and HRG: PRC 99, 3, 034909, 2019; SMES: APPB 46, 10, 1991, 2015 p+p: NA61/SHINE, EPJC 77, 10, 671, 2017; Be+Be: NA61/SHINE, EPJC 81, 1, 73, 2021; Ar+Sc: NA61/SHINE preliminary; Pb+Pb: NA61/SHINE, PRC 66, 054902, 2002.

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Multiplicity and net-charge fluctuations in $p+p$, Be+Be and $Ar+Sc$ collisions

$$
\kappa_1 = \langle N \rangle
$$

\n
$$
\kappa_2 = \langle (\delta N)^2 \rangle = \sigma^2
$$

\n
$$
\kappa_3 = \langle (\delta N)^3 \rangle = S\sigma^3
$$

\n
$$
\kappa_4 = \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2 = K\sigma^4
$$

\nwhere:
\n
$$
N - \text{multiplicity: } \delta N = N - \langle N \rangle; \sigma -
$$

standard deviation; S – skewness; K – kurtosis.

- In case of h-, only the scaled variance show significant differences between heavier and lighter systems.
- In case of net-electric charge, the scaled skewness and scaled kurtosis indicate non-monotonic behaviour.
- **•** Currently, analysis is focused on reducing the considerable systematic uncertainties.

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Proton and charged hadron intermittency in Ar+Sc and Pb+Pb collisions

m_{th} bin	$p_{\text{m: number ofparticles in m_{\text{th}} bin$	p_{A}	p_{A}	p_{A}	p_{A}
$F_r(M) = \frac{\left(\frac{1}{M} \sum_{m=1}^{M} n_m (n_m-1)...(n_m-r+1)\right)}{\left(\frac{1}{M} \sum_{m=1}^{M} n_m\right)^r}$, • Statistics					
$F_r(M) = \frac{\left(\frac{1}{M} \sum_{m=1}^{M} n_m\right)^r}{\left(\frac{1}{M} \sum_{m=1}^{M} n_m\right)^r}$, • If the sys!					

where $\langle \ldots \rangle$ denotes averaging over events, M is the number of cells.

Ily independent points, cumulative

20

15

 $\widetilde{\boldsymbol{\epsilon}}$

control Ph+Ph at

protons

- If the system freezes-out in the vicinity of the critical point, $F_2(M)$ should reveal a power-law dependence \rightarrow not observed in these analyses.
- Work on more advanced methodology ongoing.

500

1000

Symmetric Lévy HBT correlations

 $C(a)$ KT = 0.22 GeV/c (0.20 - 0.25) GeV/c

- \bullet $A(q)$ pairs of pions from same event,
- $B(q)$ pairs of pions from mixed events,
- $C(q) = A(q)/B(q)$,
- $q = |p_1 p_2|$
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- **Bose-Einstein correlations are sensitive to** spatial extension of particle source.
- Usually correlation function assumes Gaussian source but it can be generalized by Lévy-shaped:

$$
C(q) = 1 + \lambda \cdot e^{-(qR)^{\alpha}} \qquad (2)
$$

where:

 $\alpha = 0.50 \pm 0.05$ – conjectured value at the critical point (CP), α < 2 – anomalous diffusion. $\alpha = 2$ – Gaussian (EPJC 36, 67, 2004).

Symmetric Lévy HBT correlations

- \bullet R Lévy-scale parameter:
	- **o** describes length of homogeneity,
	- $\frac{1}{2}$ from hydro: $R \sim 1/\sqrt{m_T}$ (For Gaussian source) PRC 54, 1390, 1996,
	- visible m_T dependence sign of transverse flow.
- **•** Lévy-stability index α :
	- shape of spatial correlation,
	- *α* does not indicate CP in Be+Be and Ar+Sc (far from 0.5),
	- *α* between Gaussian or Cauchy shape might be the sign of anomalous diffusion.

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New data on hadron spectra in $p+p$ reactions

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NA61/SHINE, EPJC 82, 4, 322, 2022.

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NA61/SHINE, EPJC 80, 5, 460, 2020

NA61/SHINE, EPJC 82, 4, 322, 2022

System size dependence of $\mathit{K}^{*}(892)^{0}$ to charged kaon ratio

 K^*/K^- or $K^*/K^+ \rightarrow$ time between chemical and kinetic freeze-outs (STAR, PRC 71, 064902, 2005; C. Blume, APPB 43, 577, 2012):

$$
\left.\frac{K^*}{K}\right|_{kinetic} = \left.\frac{K^*}{K}\right|_{chemical} e^{-\frac{\Delta t}{\tau}}
$$
\n(3)

Assumption: no regeneration processes; ratio for kinetic freeze-out from Pb+Pb interactions; ratio for chemical freeze-out from p+p interactions.

- Lorentz boosted time interval between chemical and kinetic freeze-outs for Pb+Pb at 158A GeV/ c :
	- 5.3 fm/*c* for $\mathsf{K}^*(892)^0/\mathsf{K}^+$, 4.6 fm/*c* for $\mathsf{K}^*(892)^0/\mathsf{K}^-$
- ∆t **at SPS** *>* ∆t **at RHIC** (at corresponding centrality) NA61/SHINE, EPJC 80, 5, $460, 2020 \rightarrow$ regeneration effects may be significant at higher energies
- Regeneration effects may exist also at SPS → obtained ∆t **is a lower limit of time between freeze-outs**
- Reference ion data are needed to estimate Δt at lower energies $(K^*/K^{\pm}$ for p+p data already exist – left plot).

\equiv production in inelastic p+p collisions at 158 GeV/c

- The only results on Ξ^- and $\overline{\Xi}^+$ production in p+p
- at the SPS energy (NA61/SHINE, EPJC 81, 10, 911, 2021).
- Suppression of $\overline{\Xi}^+$ production at mid-rapidity.

Strangeness enhancement factors - Ξ production

Erratum: EPJC 82, 174, 2022.

- **The enhancement recalculated** based on the NA61/SHINE data.
- The strangeness enhancement factor (NPB 111, 461, 1976)):

$$
E = \frac{2}{\langle N_W \rangle} \frac{dn/dy (A+A)}{dn/dy (p+p)}, \quad (4)
$$

• The NA61/SHINE $p+p$ data is new baseline for Ξ production at 158A GeV/c .

 $\Xi(1530)^0$ production in inelastic p+p collisions at 158 GeV/c

Mean multip. $\langle \overline{\Xi}(1530)^{0}\rangle/\langle \Xi(1530)^{0}\rangle$ $= 0.40 \pm 0.03 \pm 0.05.$

- The only results on $\Xi(1530)^0$ production in p+p at the SPS energy (NA61/SHINE, EPJC 81, 10, 911, 2021).
- The second result on $\Xi(1530)^0$ production in p+p (ALICE at 7 TeV – EPJC 75, 1, 2015).

NA61/SHINE upgrade

- Main goal: first ever open charm measurements at SPS. Open questions:
	- . What is the mechanism of open charm production?
	- How does the onset of deconfinement impact open charm production?
	- **How does the formation of** quark-gluon plasma impact J*/*Ψ production?
- To answer these questions mean number of charm quark pairs $\langle c\bar{c}\rangle$ produced in the full phase space in $A+A$ collisions has to be known.

Summary

- \bullet 2D scan in system size and collision energy was completed in 2017 with $Xe+La$.
- NA61/SHINE delivers reach information related to the onset of deconfinement in the light and medium-size system.
- The onset of fireball unexpected system size dependence.
- So far no convincing indication of the critical point.
- Detector upgrade almost done, open charm measurements starting this year.

Thank you!

This work was supported by the Polish Ministry of Science and Higher Education (grant WUT ID-UB), the Norwegian Financial Mechanism 2014–2021 (grant 2019/34/H/ST2/00585), the Polish Minister of Education and Science (contract No. 2021/WK/10).