

Dark Matter, Supersymmetry and the LHC

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Quick Review

Gaugino Sector Basics

- Gauginos part of vector supermultiplets: $A_a = \{\lambda_a, (A_\mu)_a, D_a\}$, $a = 1, 2, 3$

Names	spin 1/2	spin 1	$SU(3)_C, SU(2)_L, U(1)_Y$
gluino, gluon	\tilde{g}	g	$(\mathbf{8}, \mathbf{1}, 0)$
winos, W bosons	$\tilde{W}^\pm \tilde{W}^0$	$W^\pm W^0$	$(\mathbf{1}, \mathbf{3}, 0)$
bino, B boson	\tilde{B}^0	B^0	$(\mathbf{1}, \mathbf{1}, 0)$

- Supersymmetry breaking independent of EWSB
Thus in SUSY limit we have massless gauginos up to EWSB effects
- Soft SUSY-breaking gaugino masses: $\mathcal{L}_{\text{soft}} \ni -\frac{1}{2}M_a \lambda_a \lambda_a + \text{c.c.}$

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- Soft SUSY-breaking gaugino masses: $\mathcal{L}_{\text{soft}} \ni -\frac{1}{2}M_a \lambda_a \lambda_a + \text{c.c.}$

⇒ Neutralinos: Higgsinos $(\tilde{H}_u^0, \tilde{H}_d^0)$, W-ino \tilde{W}^0 and **B-ino** \tilde{B}

$$\psi^0 = \left(\tilde{B}, \tilde{W}^0, \tilde{H}_d^0, \tilde{H}_u^0 \right)$$

- Mass terms in 4×4 notation: $\mathcal{L} \ni -\frac{1}{2} (\psi^0)^T M_{\tilde{N}} (\psi^0) + \text{c.c.}$

$$M_{\tilde{N}} = \begin{pmatrix} M_1 e^{i\varphi_1} & 0 & -g'v_d/\sqrt{2} & g'v_u/\sqrt{2} \\ 0 & M_2 e^{i\varphi_2} & gv_d/\sqrt{2} & -gv_u/\sqrt{2} \\ -g'v_d/\sqrt{2} & gv_d/\sqrt{2} & 0 & -\mu e^{i\varphi_\mu} \\ g'v_u/\sqrt{2} & -gv_u/\sqrt{2} & -\mu e^{i\varphi_\mu} & 0 \end{pmatrix}$$

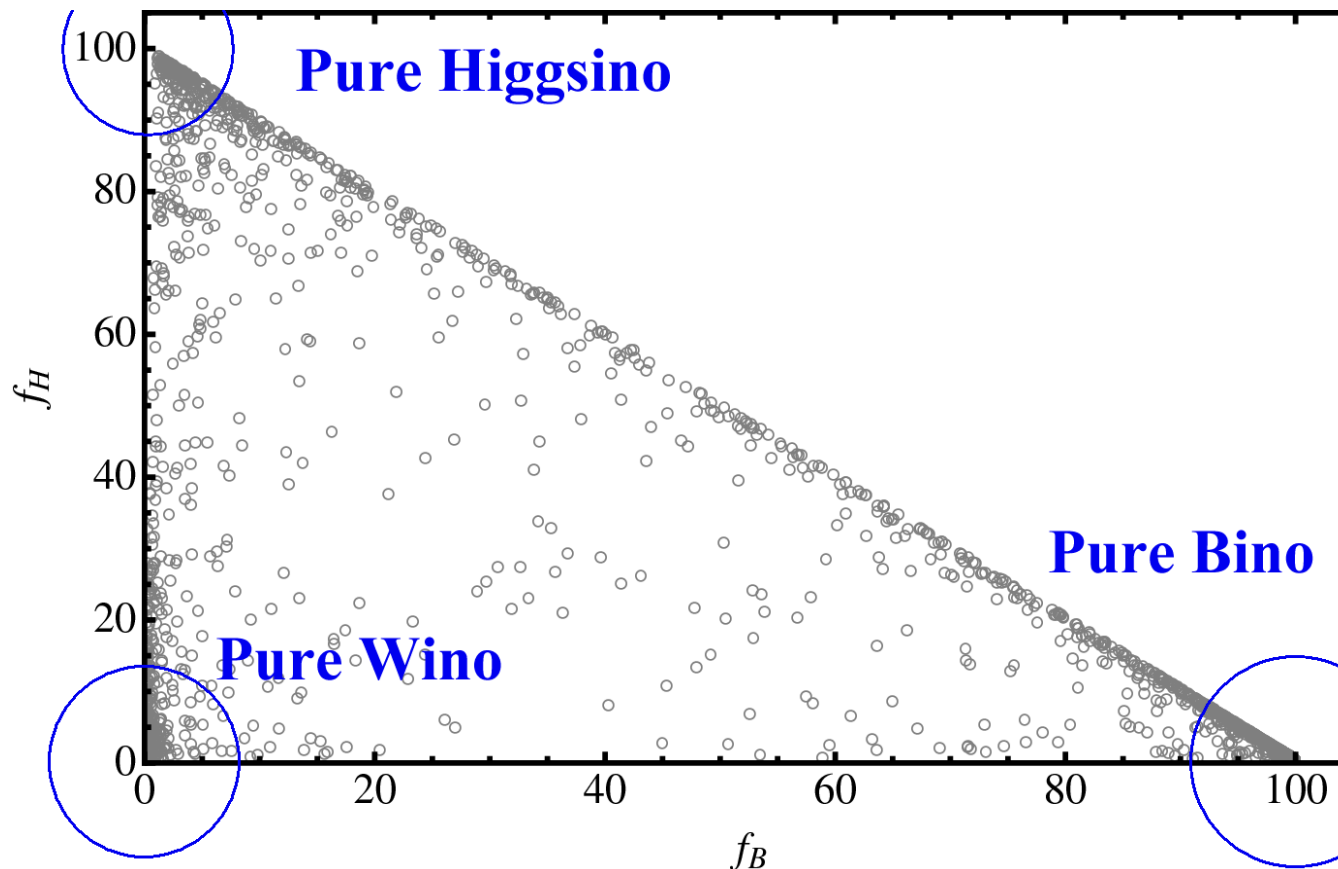
The Nature of the Lightest Supersymmetric Particle (LSP)

⇒ Can model possibilities via $M_a = m_{1/2} (1 + \delta_a)$

- $\delta_1 = \delta_2 = \delta_3$ produces bino-like LSP: $\tilde{N}_1 \sim \tilde{B}$; $\tilde{N}_2 \sim \tilde{W}^0$
- $\{\delta_1 = 0, \delta_2 < 0\}$ produces wino-like LSP
- $\{\delta_2 > 0, \delta_3 < 0\}$, $|\delta_3| < |\delta_2|$ produces Higgsino-like LSP via RGEs + EWSB

$$M_Z^2 = 5.9M_3^2 - 1.8\mu^2 + 0.4m_0^2 - 0.4M_2^2 + \dots$$

Kane, Lykken, BDN, Wang, PLB 551 (2003) 146



Relic Density: *mirabile dictu!*

Standard Thermal Cosmology

⇒ WIMP “miracle” never expected to be better than two orders of magnitude!

$$\Omega_\chi h^2 = \frac{m_\chi n_\chi}{\rho_c} \simeq \frac{3 \times 10^{-27} \text{ cm}^3/\text{s}}{\langle \sigma_{\text{ann}} v \rangle}$$

with $\langle \sigma_{\text{ann}} v \rangle |_{\text{weak}} \sim \alpha_{\text{weak}}^2 (100 \text{ GeV})^{-2} \sim 10^{-25} \text{ cm}^3/\text{s}$ for $\alpha \sim 1/100$.

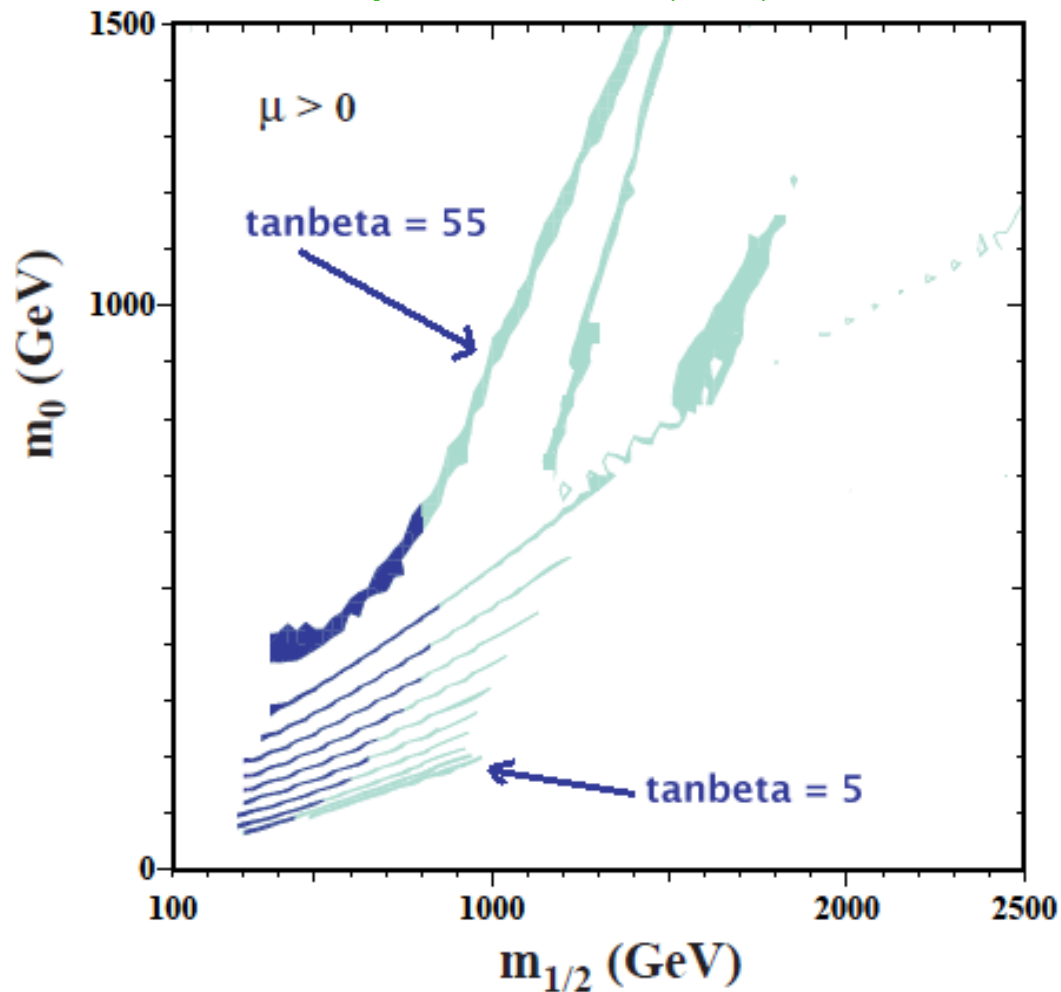
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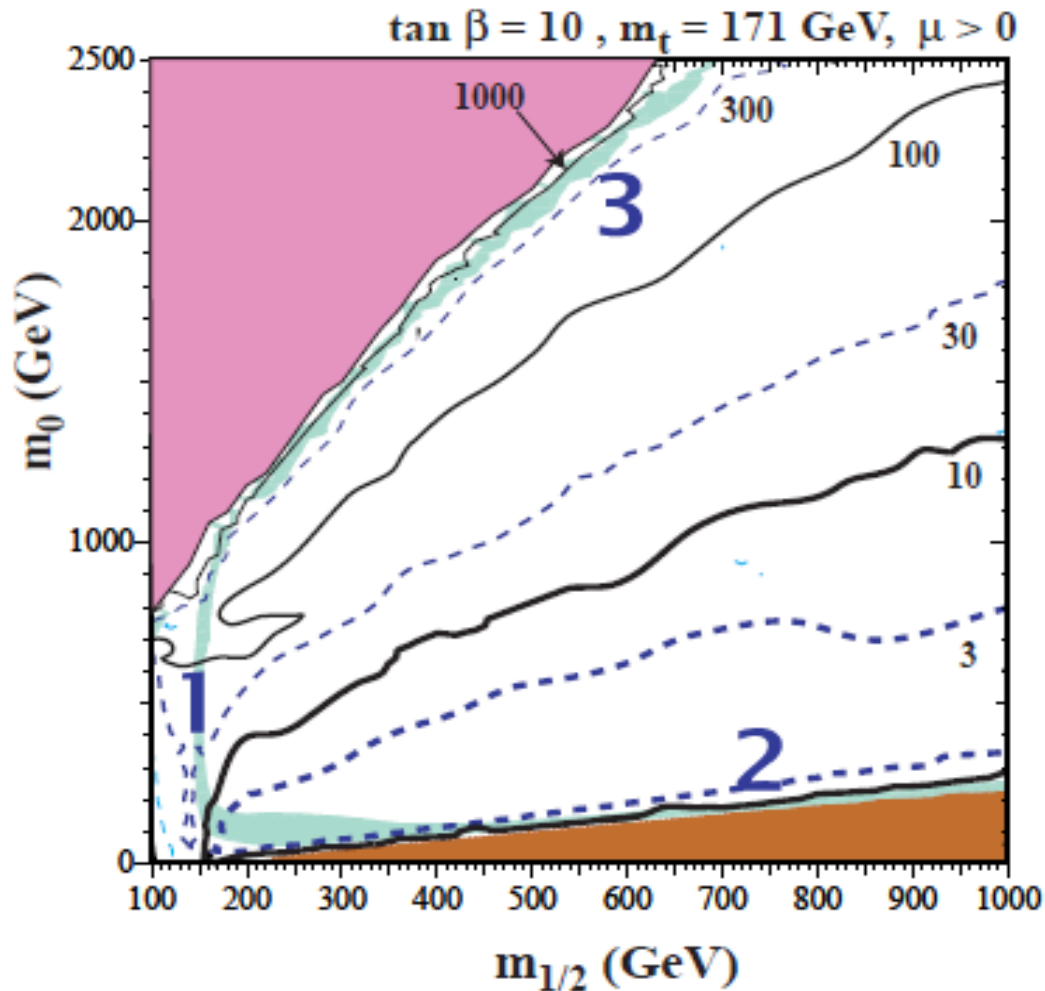
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Ellis, Olive, Santos, Spanos, PLB 565 (2003) 176



- WMAP Seven Year Data:
 $\Omega_\chi h^2 = 0.1109 \pm 0.0056$
- mSUGRA with Bino LSP:
 $\Omega_\chi h^2 \sim 10$
- Thin strip in $(m_0, m_{1/2})$ plane
which varies with $\tan\beta$
- At very large $\tan\beta$ reach
resonant annihilation
 $m_\chi \simeq 0.5m_A, 0.5m_H$

⇒ Three generic mSUGRA configurations consistent with WMAP data

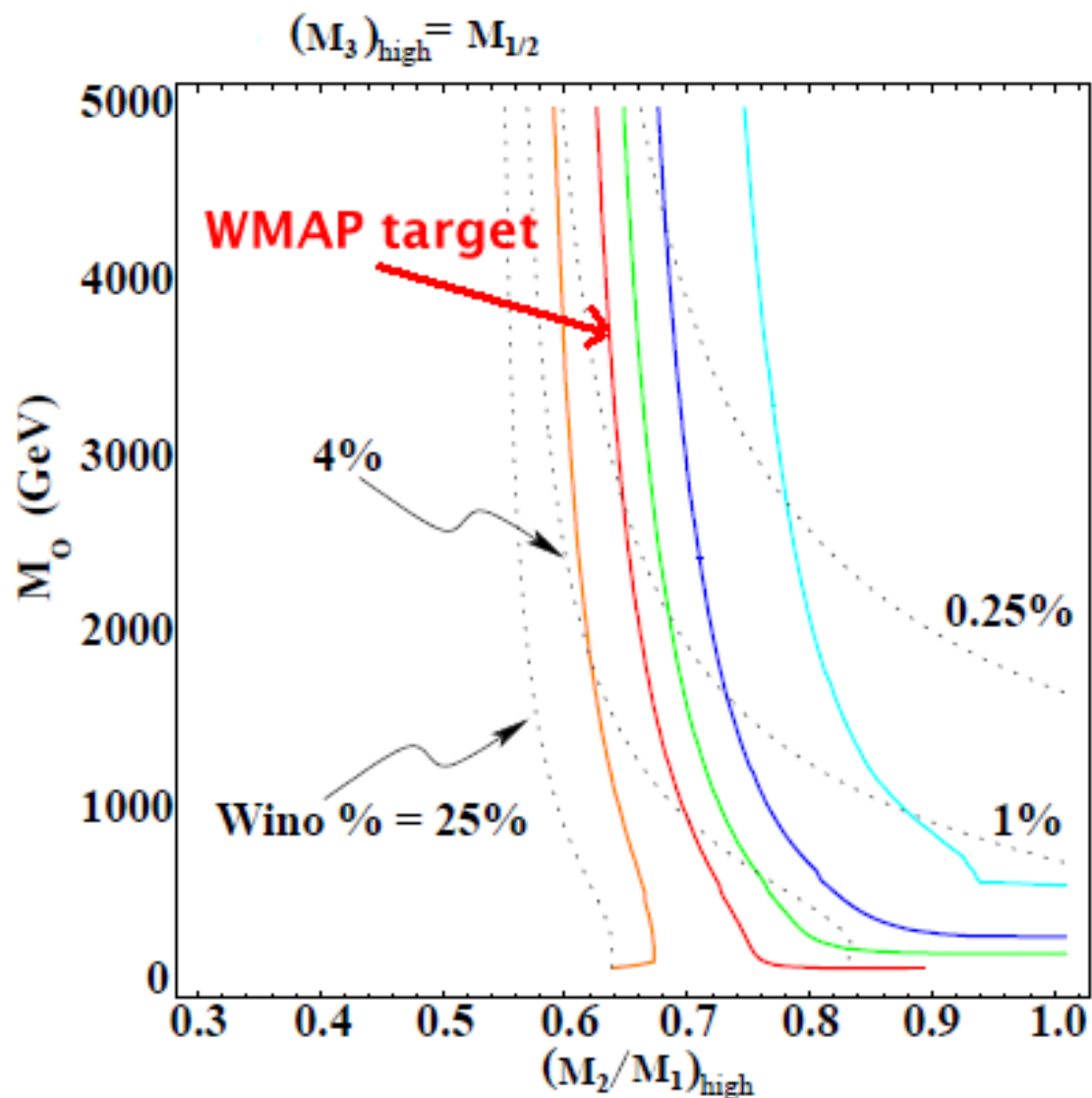


- **1** “Bulk region”: light sleptons facilitate annihilation in early universe; possible light CP-even Higgs resonance
- **2** “Stau Co-annihilation”: skirts region where $m_{\tilde{\tau}} < m_{\chi}$
- **3** “Gaugino Co-annihilation”: in focus point region where $\mu^2 \rightarrow 0$; degenerate chargino/neutralino states co-annihilate with Higgsino-like LSP

Impact of LSP Wavefunction

⇒ Previous conclusions heavily dependent on wavefunction of the LSP

Birkedal-Hansen & BDN, PRD 64 (2001) 015008



- Deviations from $\delta_1 = \delta_2 = \delta_3$ can have huge impacts on annihilation rates
- A small amount of wino content dramatically improves the thermal relic density result
- Anomaly-mediated extreme with 100% Wino LSP:
 $\Omega_\chi h^2 \sim 10^{-3}$
- Mildly dependent on value of M_3 – suggests a “sweet spot” in the space of gaugino masses (“well-tempered” neutralino)

- ⇒ If wino/Higgsino content of the LSP is large the model (probably) requires **non-thermal** production mechanisms
- Moduli (scalar fields with gravitational couplings to Standard Model) generally dominate energy density in early universe
 - Subsequent decays will (eventually) produce LSPs – and “reheat” the cosmos to some temperature T_R
 - Requiring $T_R \gtrsim 10 \text{ MeV}$ to preserve BBN predictions implies $m_\phi \sim m_{3/2} \gtrsim 10 \text{ TeV}$
- ⇒ What are the typical relic densities in this scenario?

$$\Omega_\chi h^2 \simeq 0.1 \times \left(\frac{m_\chi}{100 \text{ GeV}} \right) \left(\frac{3 \times 10^{-24} \text{ cm}^3/\text{s}}{\langle \sigma_{\text{ann}} v \rangle} \right) \left(\frac{100 \text{ TeV}}{m_\phi} \right)^{3/2}$$

Acharya, Kane, Kumar & Watson, PRD **80** (2009) 083529

- ⇒ Note: $\langle \sigma_{\text{ann}} v \rangle \simeq 2.4 \times 10^{-24} \text{ cm}^3/\text{s}$ for pure-wino LSP (AMSB-limit)

Moroi and Randall, NPB **570** (2000) 455

An Abundances of “Miracles”

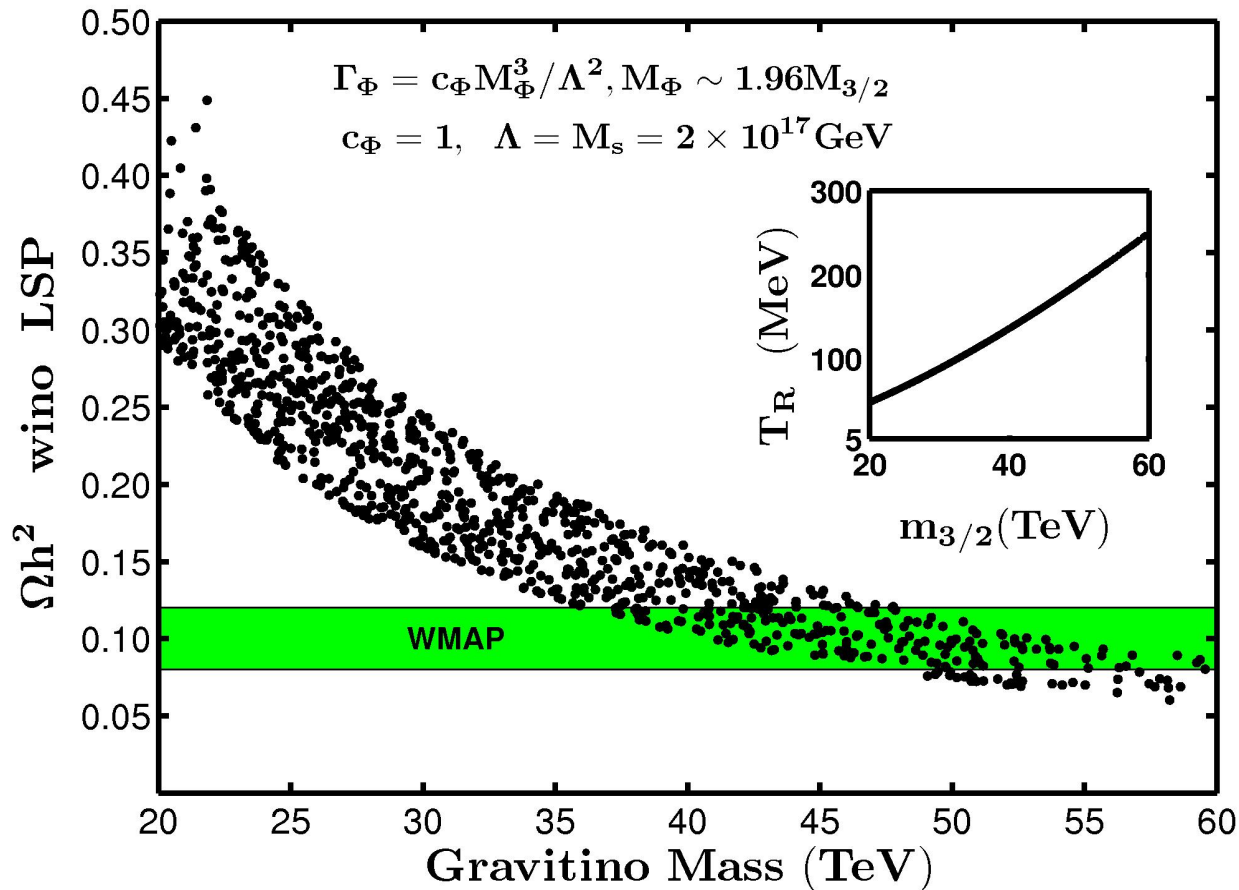
- Such basic ingredients are generically present in string-motivated effective field theories

⇒ **Non-thermal WIMP miracle**

Acharya et al., JHEP 0806 (2008) 064

Acharya, Kane, Kumar & Watson, PRD 80 (2009) 083529

Acharya, Kane & Kuflik, arXiv:1006.3272



⇒ See also **WIMPlless miracle** in gauge-mediated models

Feng and Kumar, PRL 101 (2008) 231301

Dark Matter Detection

Dark Matter Signals – the Earliest SUSY Signature

- Assumption: lightest neutralino is stable LSP \Rightarrow dark matter
Goldberg, PRL 50 (1983) 1419
 - Prediction: annihilation into photons, positrons, anti-protons, neutrinos
Silk & Srednicki, PRL 53 (1984) 624
 - ★ Photons & neutrinos “point” back to source: high density areas such as galactic center or center of sun/earth
 - ★ Charged particles must be propagated from origin to earth numerically
 - ★ Both depend on the *halo profile* $\rho_\chi(r)$ assumed for the dark matter candidate, but to varying degrees
- \Rightarrow Differential recoil rate at direct detection experiments given by
Goodman & Witten, PRD 31 (1985) 3059

$$\frac{dR}{dE} = \sum_i c_i \frac{\rho_\chi \sigma_{\chi i}^{\text{SI}} |F_i(q_i)|^2}{2m_\chi \mu_{i\chi}^2} \int_{v_{\min}}^{\infty} \frac{f(\vec{v}, t)}{v} d^3v,$$

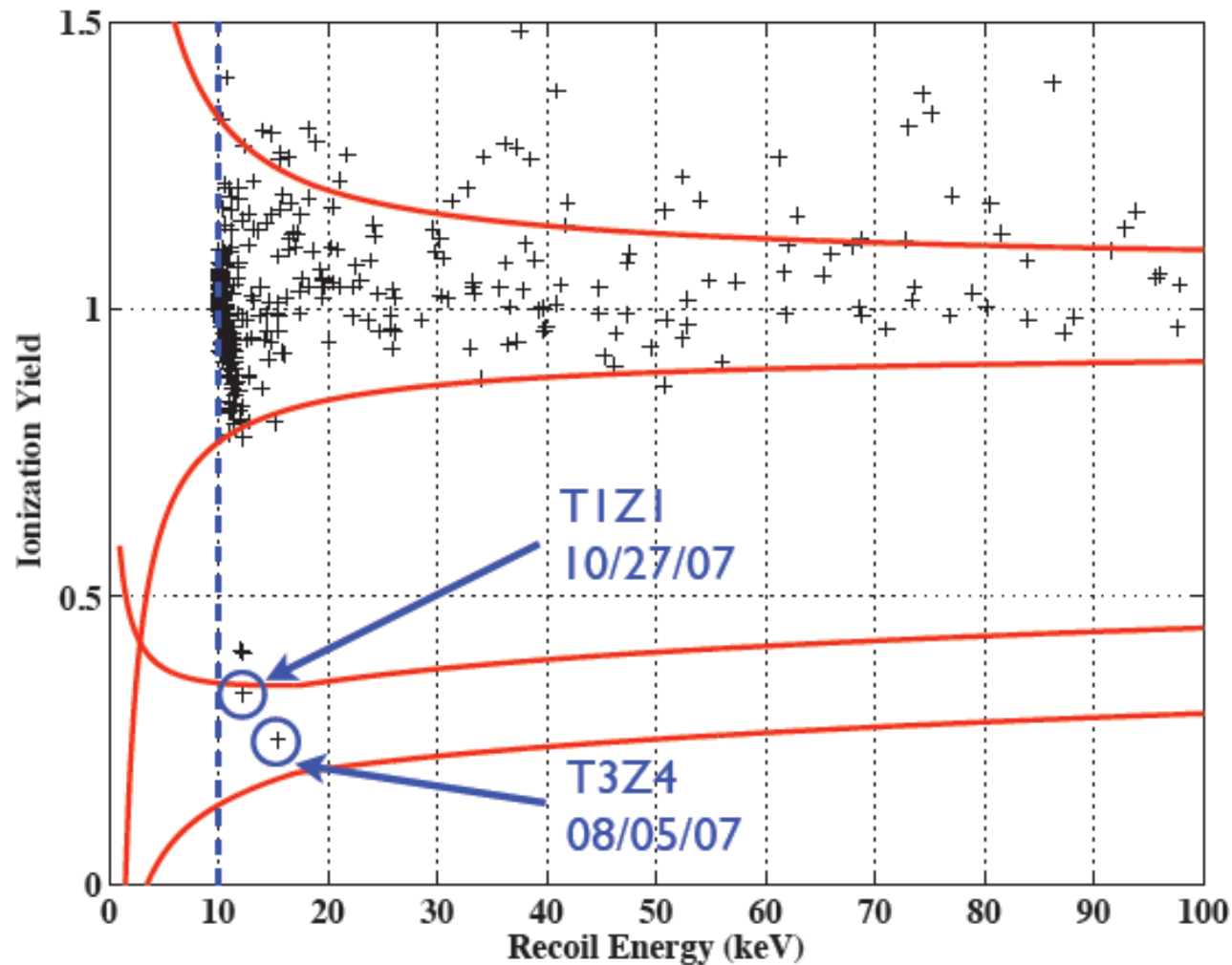
with $F_i(q_i)$ being a nuclear form factor for i -th target nucleus

- Calculation of integrated event rate depends on experimental configuration

$$R = \int_{E_{\min}}^{E_{\max}} \frac{dR}{dE} dE; \quad (\text{Germanium}) : 10 \text{ keV} \leq E_{\text{recoil}} \leq 100 \text{ keV}$$

Direct Detection Experiments: CDMS II

⇒ December 2009 data release for 14 Ge detectors by CDMS-II Collaboration

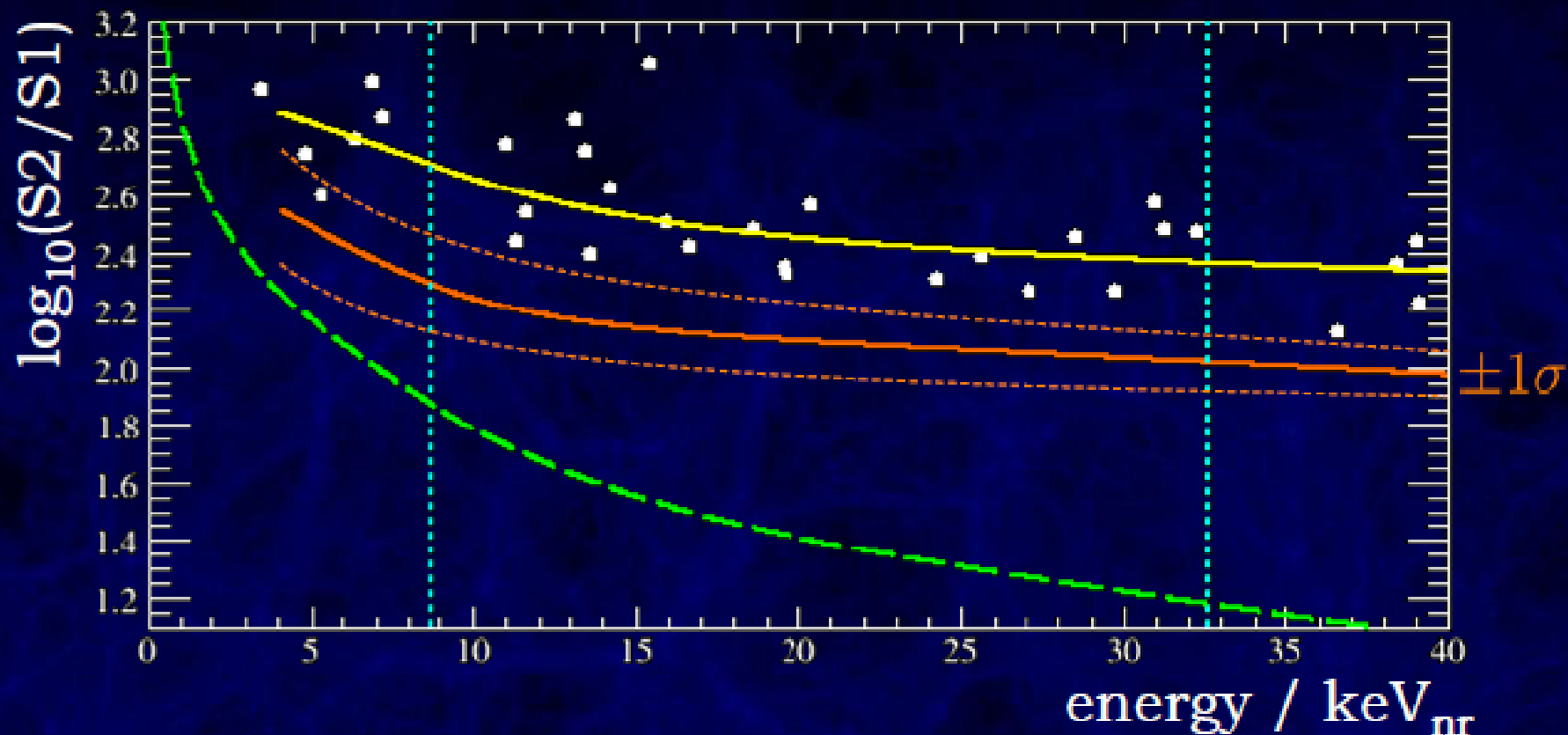


Jodi Cooley, data presentation, SLAC, 12/12/09

- Two events in signal region with (revised) background estimate of 0.8 ± 0.1 (stat) ± 0.2 (sys) events
- Implies an interaction cross-section $\sigma_{\chi p}^{\text{SI}} \sim 10^{-44} \text{ cm}^2 = 1 \times 10^{-8} \text{ pb}$

Confirmation in Xenon?

11.17 days, 40kg fiducial:



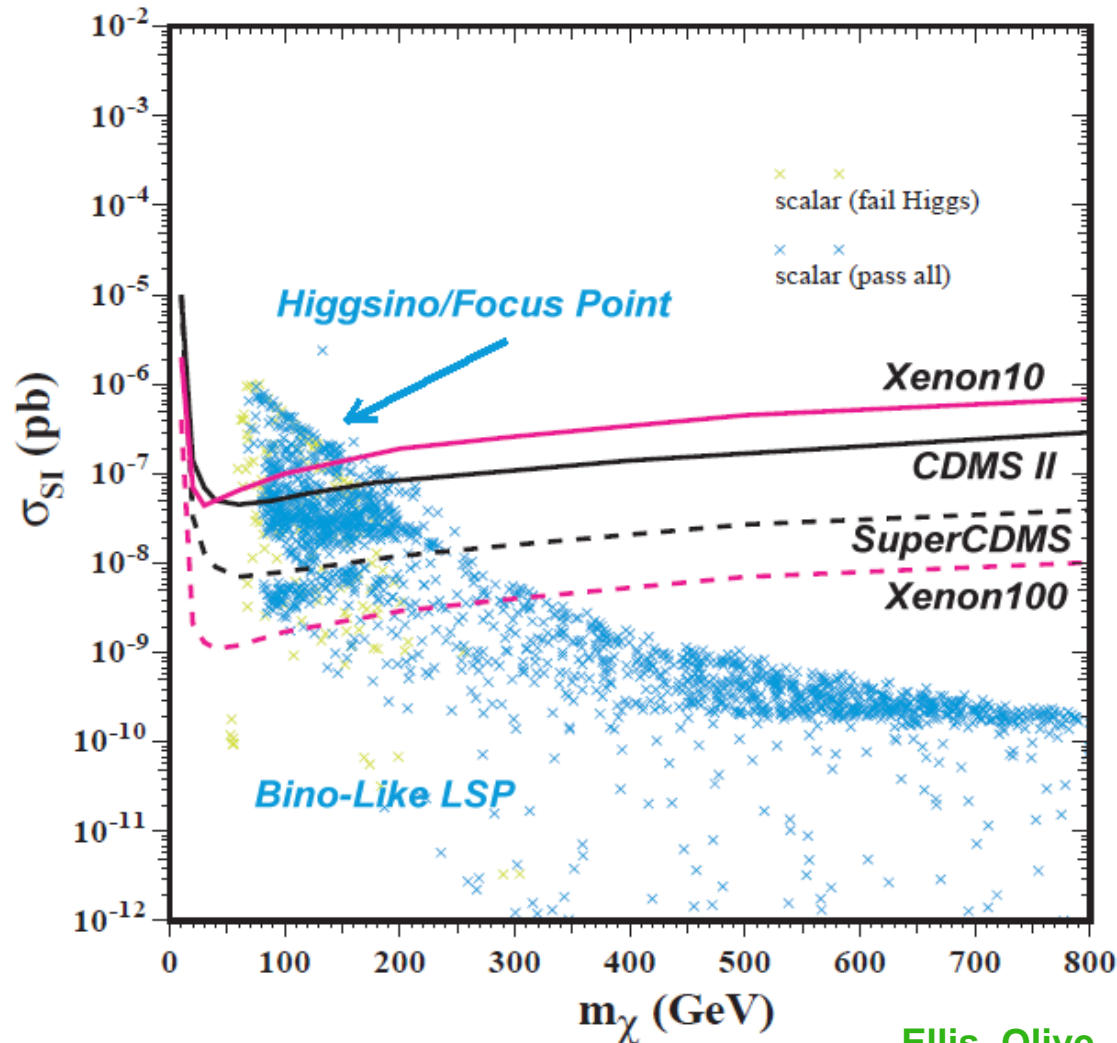
(no events even at 84% acceptance)

Direct Detection in mSUGRA

⇒ Large direct detection rate implies Higgs exchange dominates

Hooper & Taylor, JCAP **0703** (2007) 017

$$\sigma_{\chi N} \propto |N_{11}|^2 |N_{13}|^2 \left(\frac{\tan \beta}{m_A} \right)^2$$



Ellis, Olive, Sandick, arXiv:0905.0107

- Decade's worth of data suggests an excess of positrons in cosmic ray data

Cirelli et al., NPB **800** (2008) 204; **813** (2009) 1

$$\Phi_{\bar{e}}(E) \simeq \frac{\tau_E B_{\bar{e}} c}{8\pi b(E)} \frac{\rho_\chi^2(r = R_0)}{m_{\tilde{N}_1}^2} F(E), \quad b(E) = 1 \text{ GeV} \left(\frac{E}{1 \text{ GeV}} \right)^2$$

$$F(E) = \int_E^{M_{\tilde{N}_1}} dE' \sum_k \langle \sigma v \rangle_{\text{halo}}^k \frac{dN_{\bar{e}}^k}{dE'} \cdot \mathcal{I}(E, E')$$

- ★ $B_{\bar{e}}$ = boost factor, $\tau_E = \tau \times 10^{16}$ sec is the diffusion time scale and $\mathcal{I}(E, E')$ is the halo function
- ★ For SUSY models, most important final state is usually $k = W^+W^-$
- Best fits require $\langle \sigma v \rangle_{WW} \simeq 2 \times 10^{-24}$ cm³/s and prefer NFW “min” profile
 - ★ Former easy to achieve if LSP significantly wino-like

Kane, L. Wang & Wells, PRD **65** (2002) 057701

Grajek, Kane, Phalen, Pierce, Watson, PRD **79** (2009) 043506

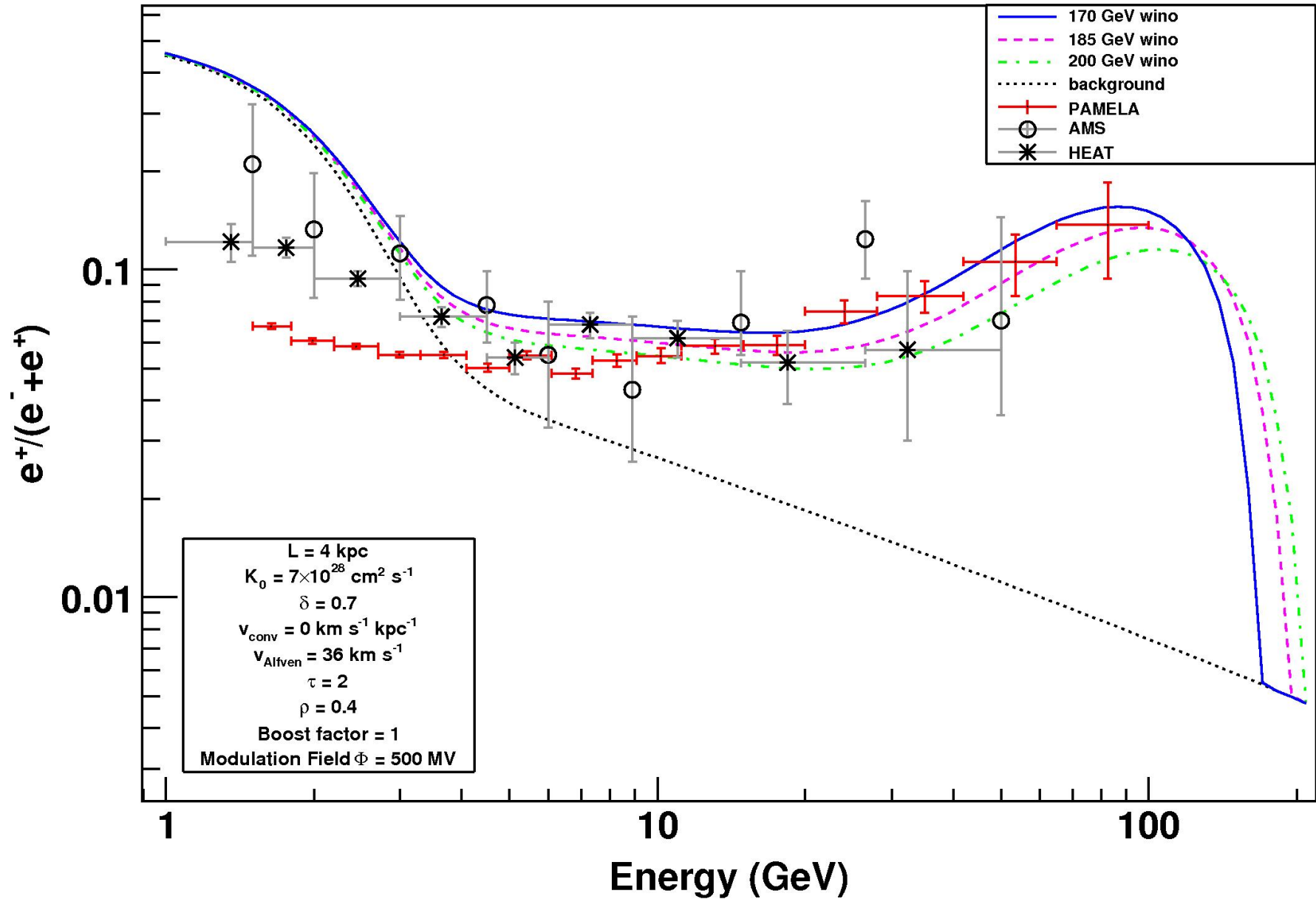
Kane, Lu, Watson, PLB **681** (2009) 151

SUSY Fits to Positron Flux Measurements

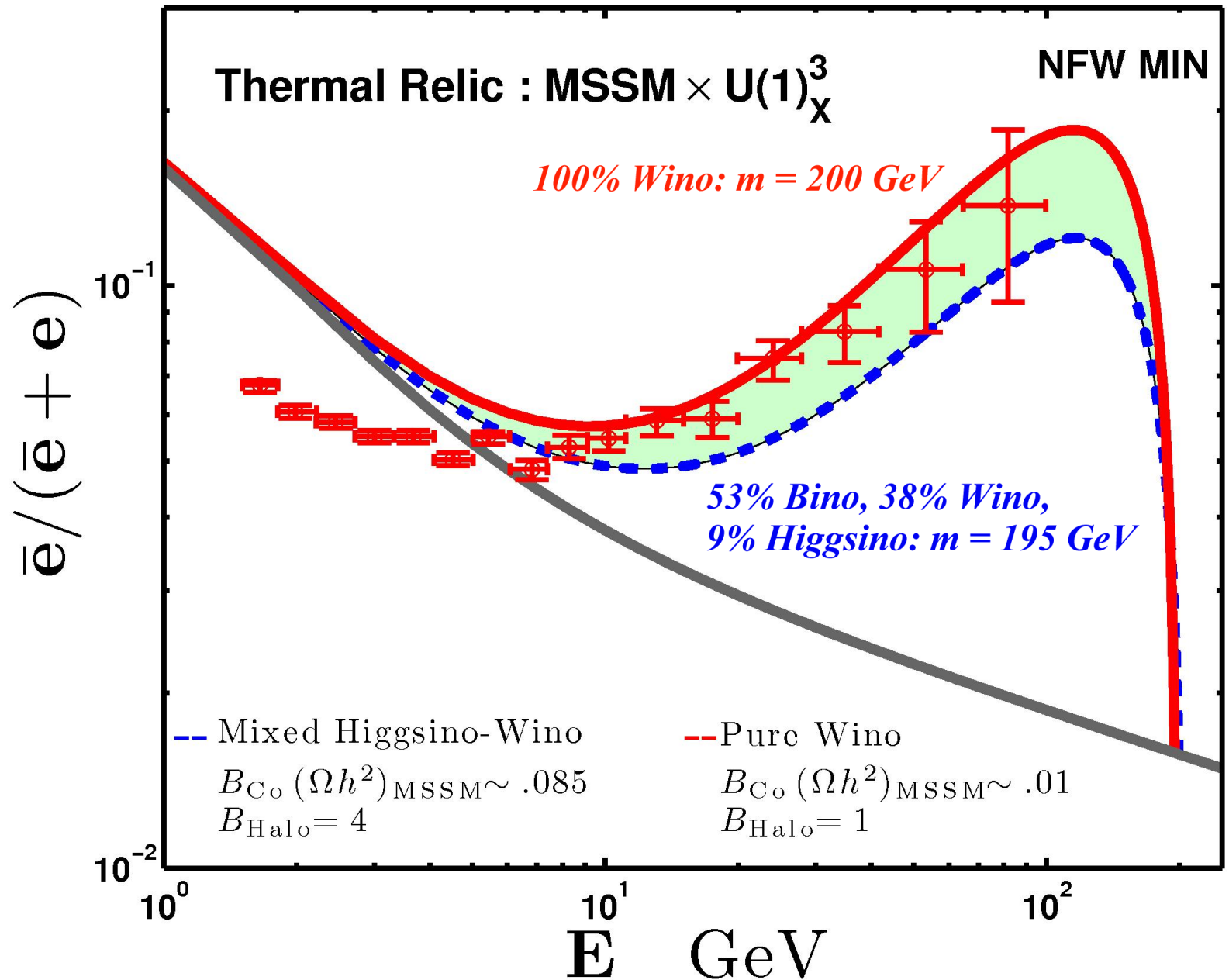
HEAT Collaboration, *Astrophys. J.* **559** (2001) 296; *PRL* **93** (2004) 241102

AMS Collaboration, *PLB* **646** (2007) 145

PAMELA Collaboration, *Nature* **458** (2009) 607; *PRL* **102** (2009) 051101



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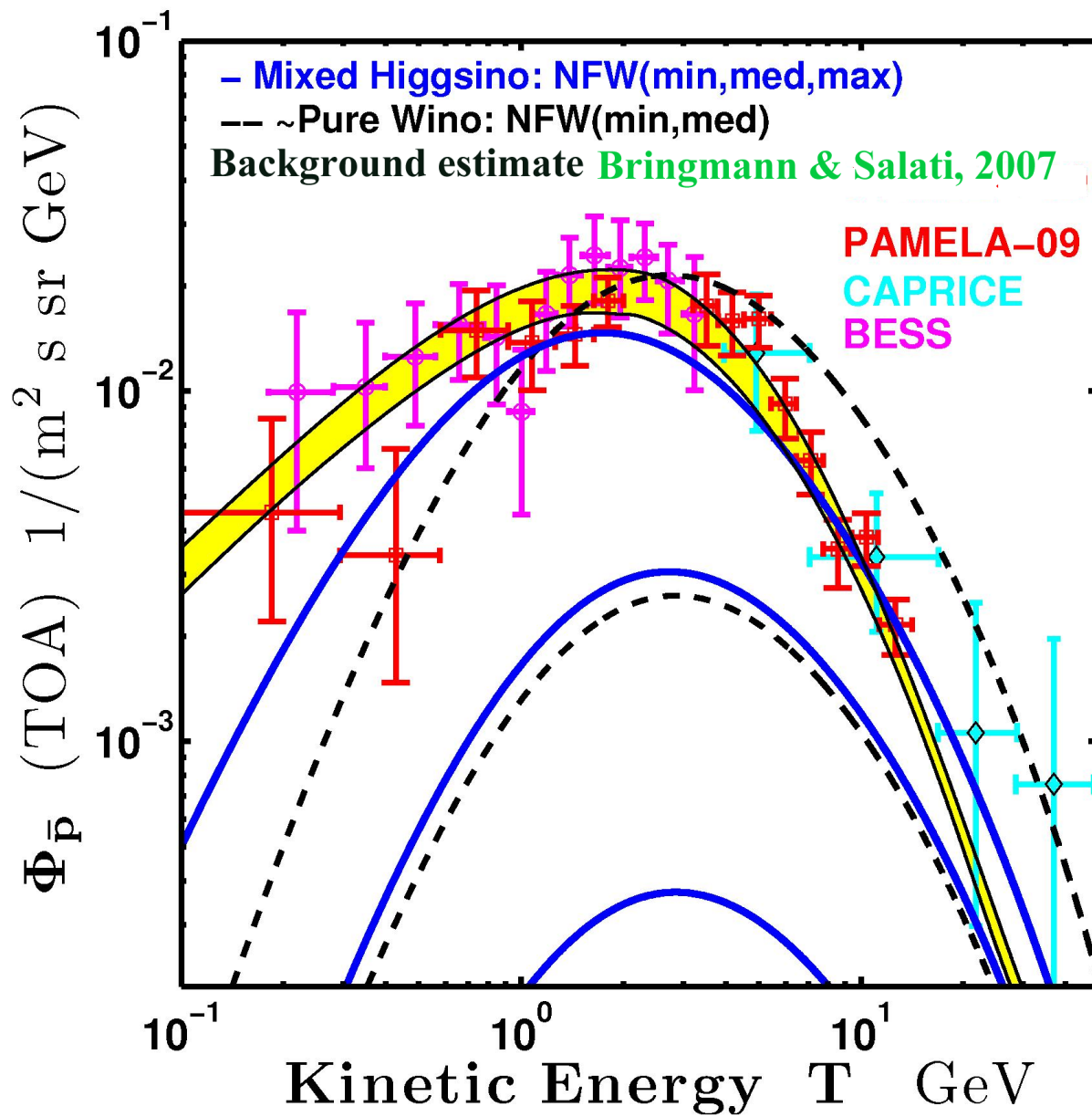


⇒ Pure wino not necessary – but must compensate with $B_{\bar{e}}$ (here B_{HALO})

Feldman, Liu, Nath, BDN, PRD D80 (2009) 075001

What About Anti-protons?

BESS Collaboration, PRL 84 (2000) 1078; CAPRICE Collaboration, Astrophys. J. 561 (2001) 787

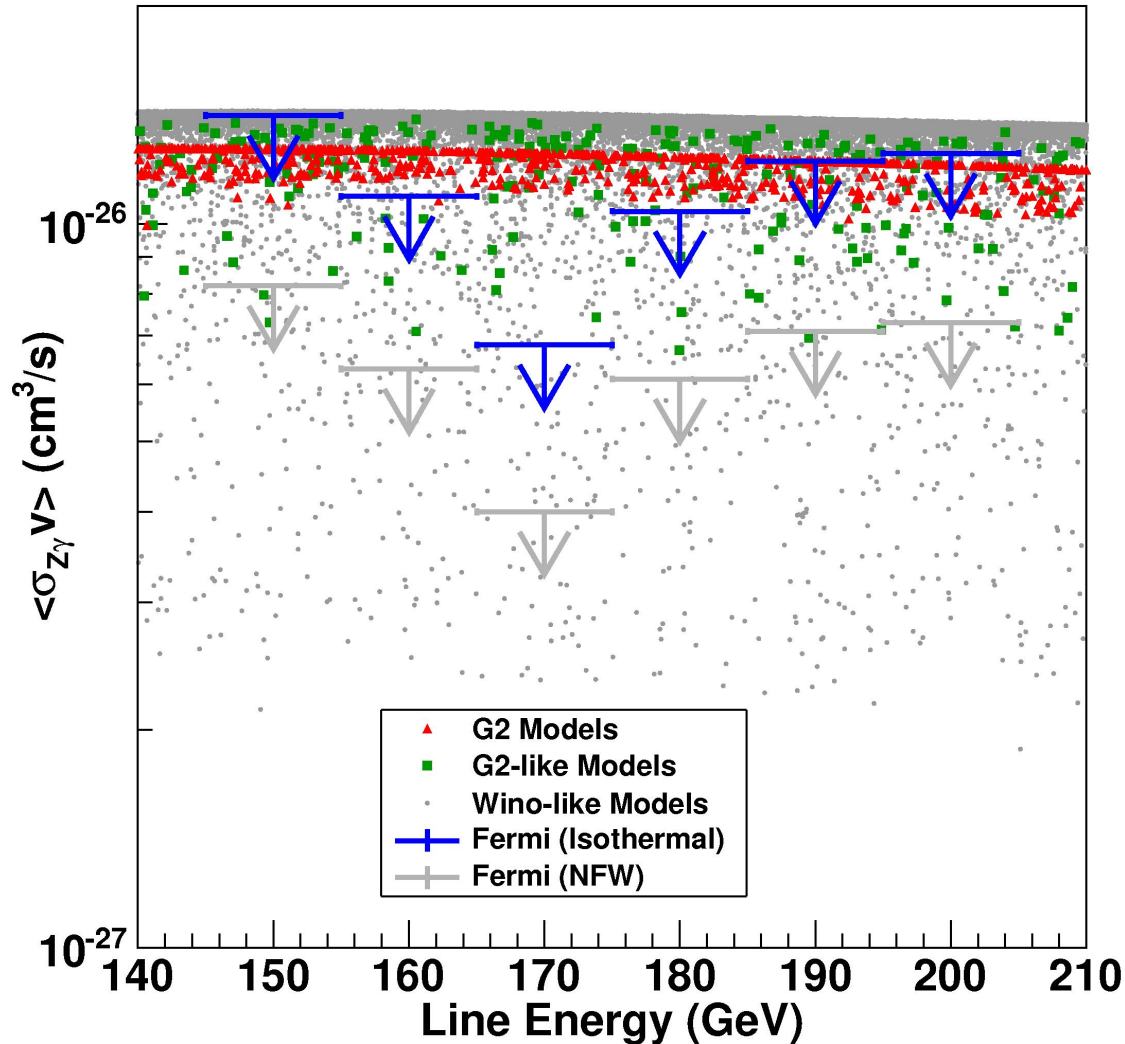


⇒ Higgsino has no problems; Wino OK for NFW “min” and “med” halo profiles

Monochromatic Signals

⇒ Monochromatic gamma ray signals a “smoking gun” for dark matter

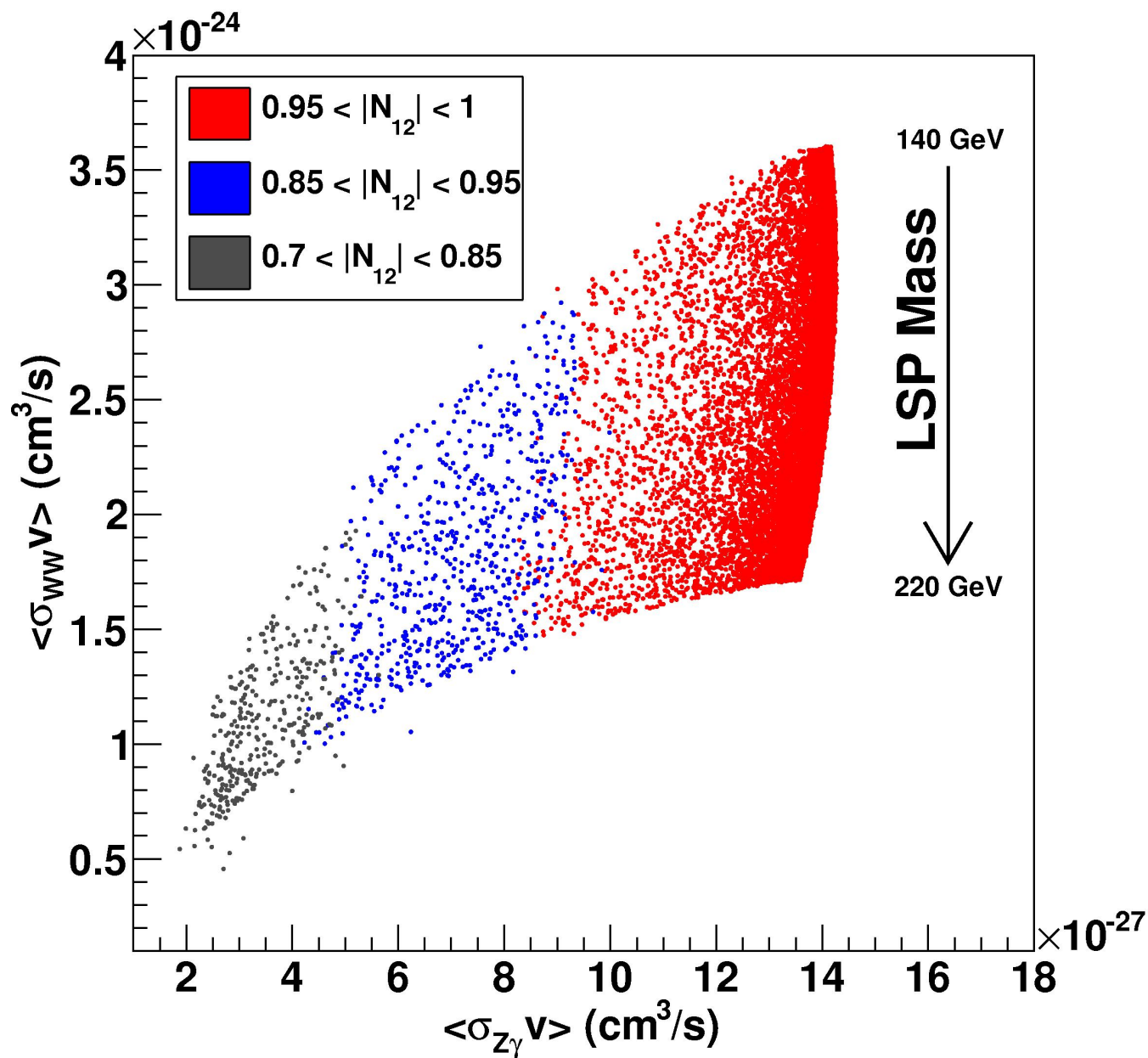
Fermi-LAT Collaboration, arXiv: 1001.4531



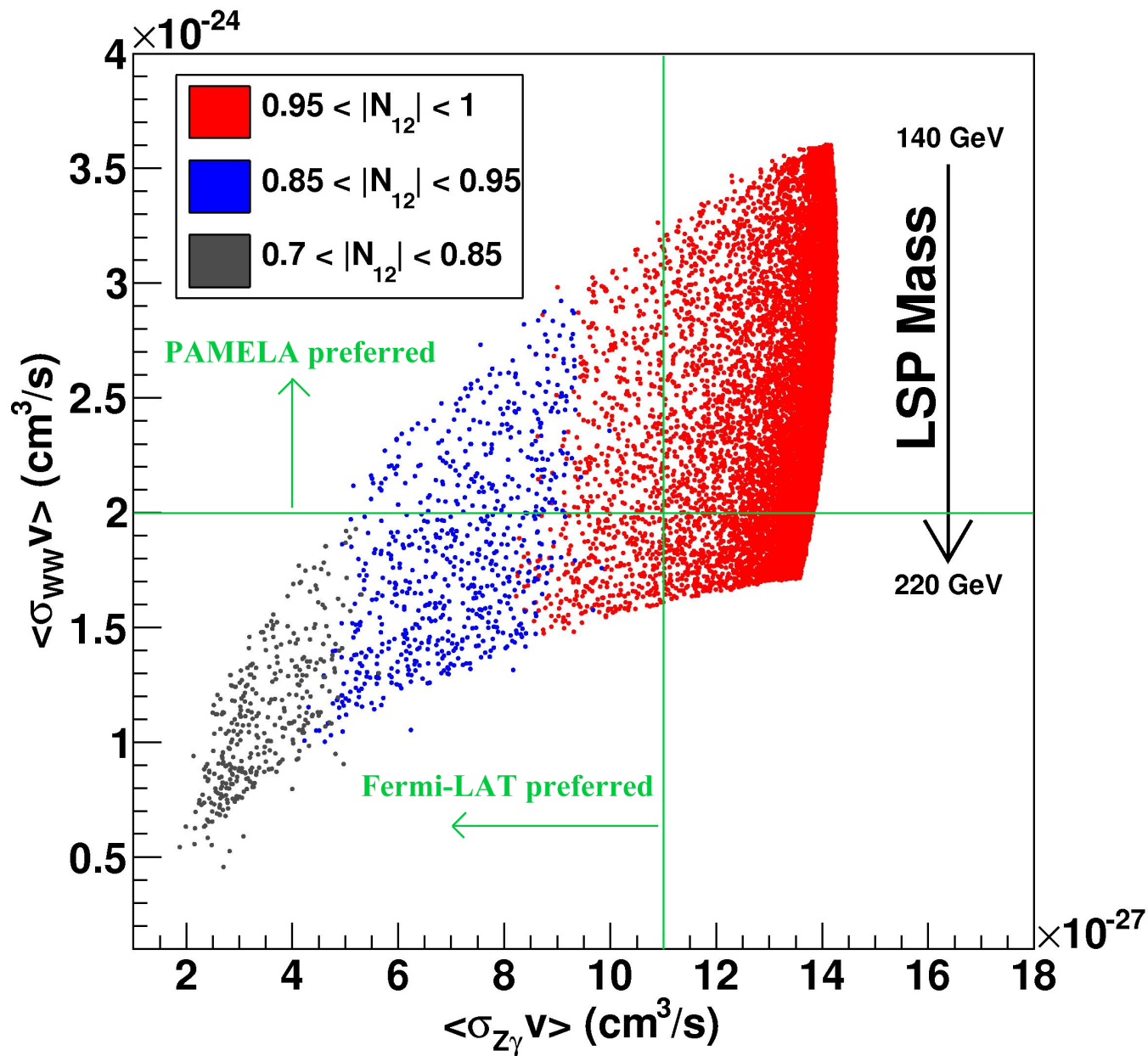
- Loop-induced diagrams provide annihilation into $\gamma\gamma$ and γZ final states
- Monoenergetic signals with $E_{\gamma\gamma} = m_\chi$ and $E_{\gamma Z} = m_\chi - M_Z^2/4m_\chi$
- Easy to pick out over background, but branching fractions reduce rate by factors of $10^3 - 10^4$
- Pure-wino models capable of getting PAMELA correct in trouble!

Feldman, Kane, Lu, BDN, PLB 687 (2010) 363

Photons versus Positrons



Photons versus Positrons



⇒ PAMELA preferred models with potentially large direct detection rates prefer **non-universal** gaugino masses

The **mirage pattern** (competition between tree and anomaly-mediated contributions to soft masses) appears in a number of phenomenologically successful string constructions:

- Kähler stabilized heterotic string models Binetruy, Gaillard, Wu, **NPB 481** (1996) 109
Gaillard and BDN, **IJMP A22** (2007) 1451
 - Type-IIB flux compactifications with anti- D_3 branes Kachru, Kallosh, Linde, Trivedi, **PRD 68** (2003) 046005
Choi, Falkowski, Nilles, Olechowski, **NPB 718** (2005) 113
 - M -theory compactified on fluxless G_2 manifolds Acharya, Kane, et al.,
PRL 97(2006) 191601
PRD 76 (2007) 126010
PRD 78 (2008) 065038
- ⇒ Common features:

- Single modulus stabilized by gaugino condensation
- Kähler potential for this modulus substantially altered from tree-level value
- Tuning of cosmological constant ($\langle V \rangle$) to zero by adjusting parameters

LHC Implications

mSUGRA Stau Coannihilation Region

⇒ Most promising mSUGRA scenario: stau co-annihilation region

Arnowitz, Dutta, Kamon, et al., PLB **639** (2006) 46; PLB **649** (2007) 73; PRL **100** (2008) 231802

- Physical hierarchies very stable in this region

Konar et al., PRL **105** (2010) 221801

- Characterized by $m_{N_2} > m_{\tilde{\tau}_1} > m_{N_1}$ and $\Delta M = m_{\tilde{\tau}_1} - m_{N_1} = 5 - 15 \text{ GeV}$

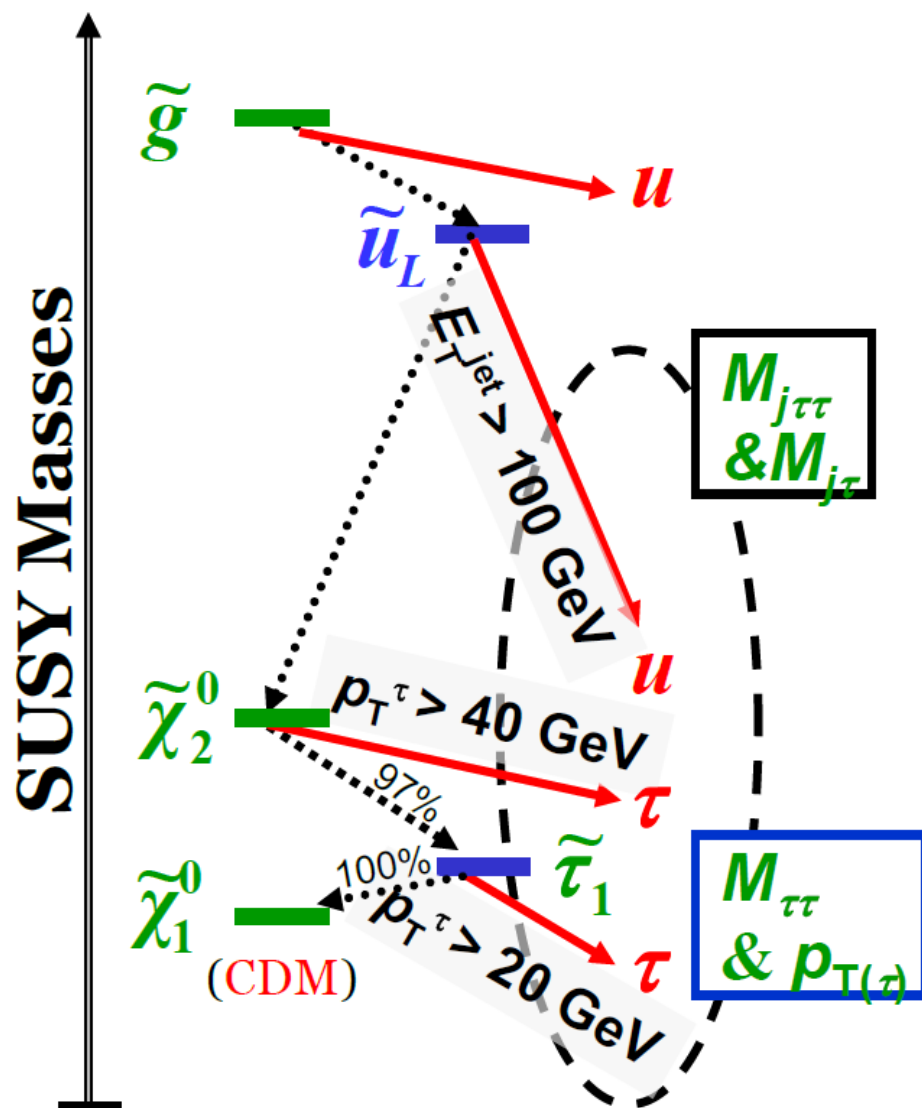
- Main production mode: $\tilde{g}\tilde{q}$ ($\sigma_{\text{SUSY}} \sim 10 \text{ fb}$)

- Discovery channels:

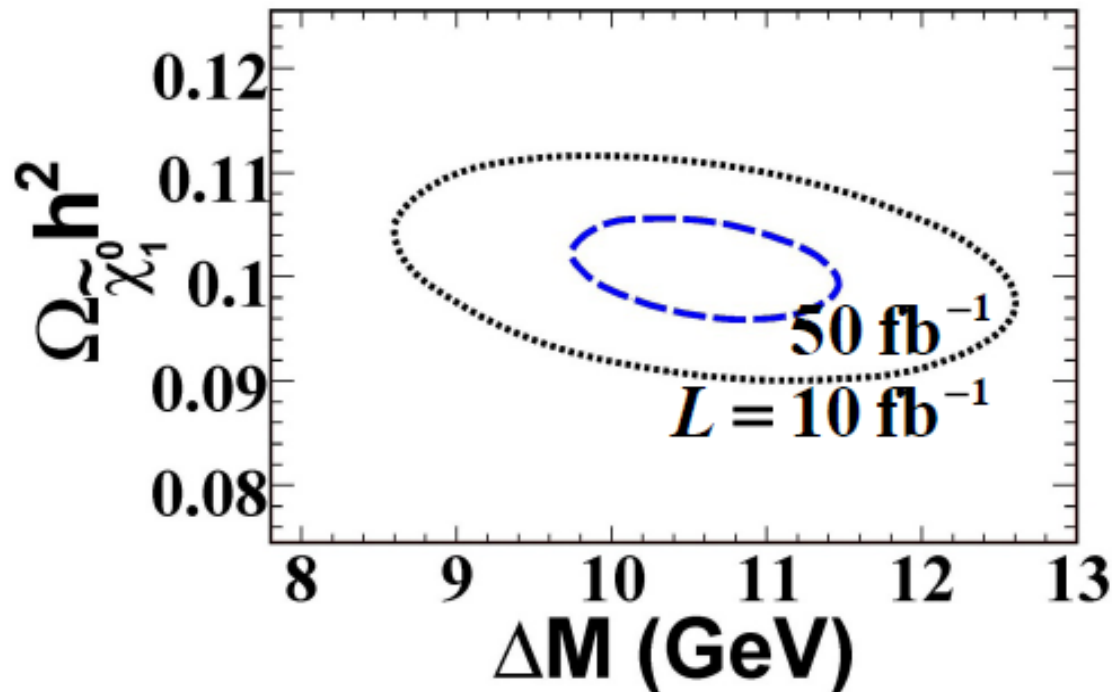
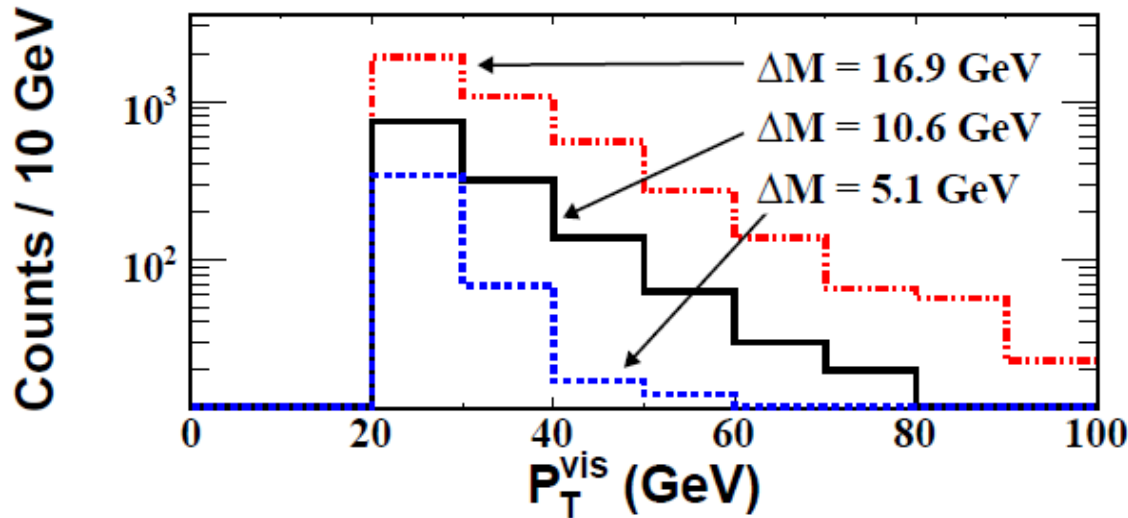
- ★ 4^+ jets + \cancel{E}_T
- ★ 2^+ jets + 2τ + \cancel{E}_T

- $\tilde{\chi}_2^0 \rightarrow \tau\tilde{\tau}$ 97% of time!

- ★ Need 50% efficiency or better for hadronically-decaying taus
- ★ Second tau likely to be very soft

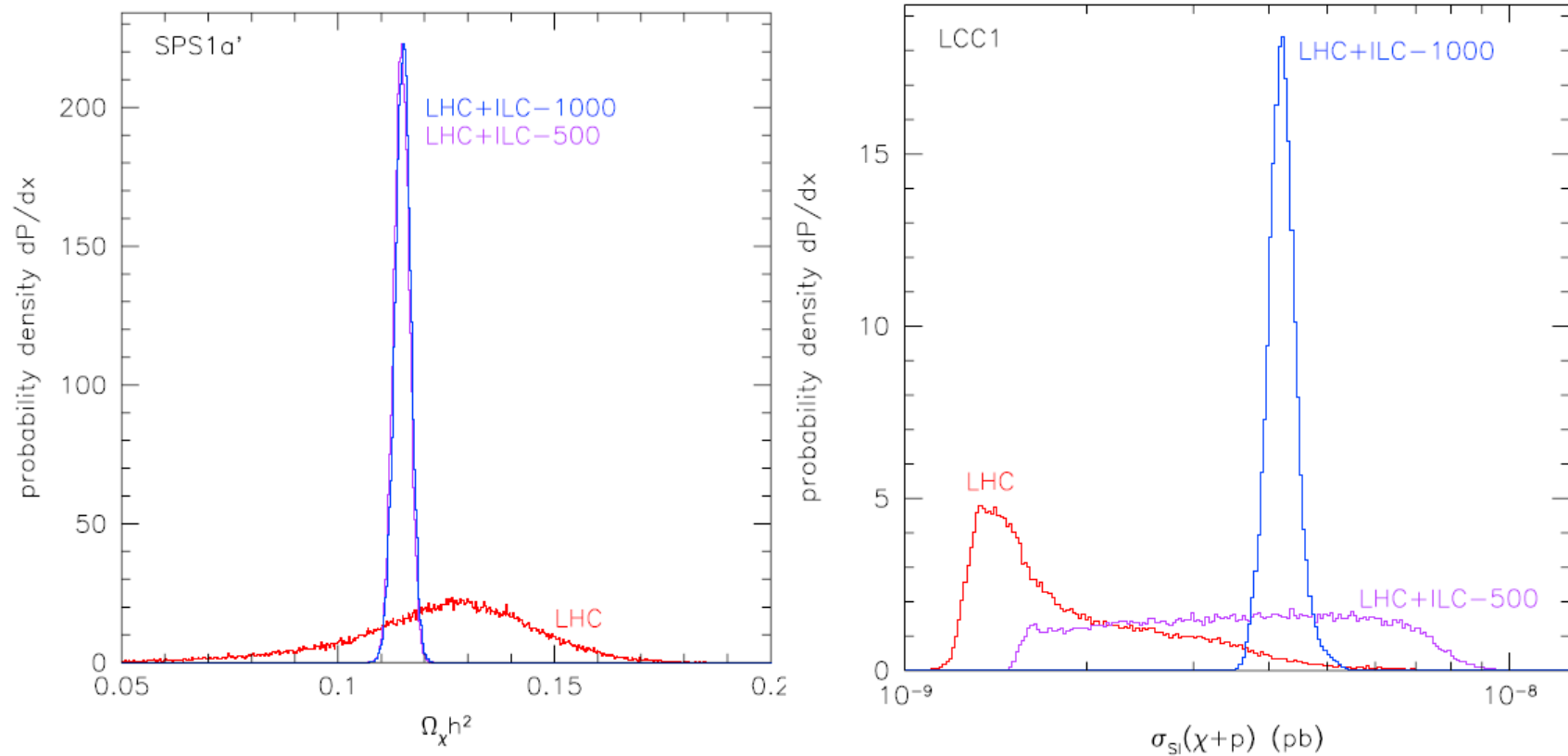


mSUGRA Stau Coannihilation Region



- Global fit to extract four mSUGRA parameters from the kinematics of dominant processes
 - ★ Slope of p_T for softer tau
 - ★ Di-tau invariant mass
 - ★ Jet + ditau invariant mass
 - ★ Jet + single tau invariant mass (two of these)
 - ★ M_{eff} for 4^+ jet events (b-jet veto)
 - ★ M_{eff}^b for 4^+ jet events (leading jet b-tagged)
- Invert to obtain $m_0, m_{1/2}, A_0$ and $\tan \beta$
- Use these as inputs (with uncertainties) to compute thermal relic abundance

Using LHC to Predict (Cross-Check) Dark Matter



⇒ Numerous studies of LHC/ILC influence on making dark matter predictions in mSUGRA – much fewer in generic SUSY models

Baltz, Battaglia, Peskin, Wizansky, PRD **74** (2006) 103521

Goodman, Ibe, Rajaraman, Shepherd, Tait, Yu, arXiv:1008.1783

Bertone, Cerdeno, Fornasa, Ruiz de Austri, Trota, PRD **82** (2010) 055008

Konar, Kong, Matchev, Perelstein, New J. Physics **11** (2009) 105004

Chung, Everett, Kong, Matchev, JHEP **0710** (2007) 016

Nojiri, Polesello, Tovey, JHEP **0603** (2006) 063

Allanach, Bellanger, Boudjema, Pukhov, JHEP **0412** (2004) 020

Birkedal, Matchev, Perelstein, PRD **70** (2004) 077701

Benchmark Models I: PAMELA Examples

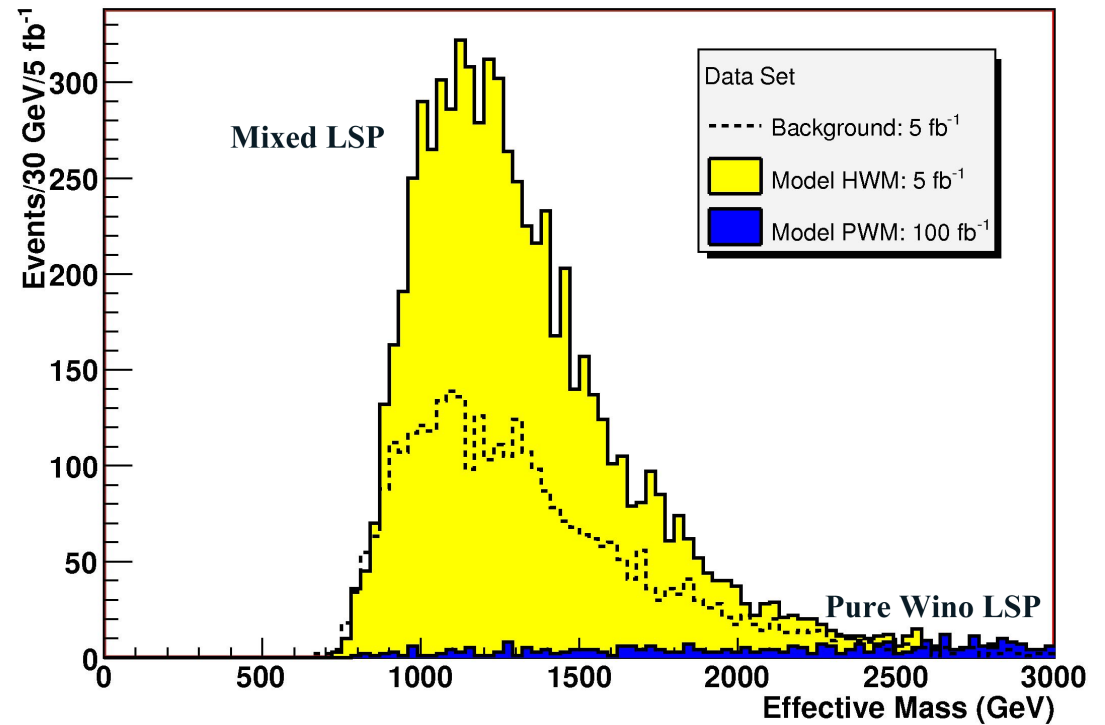
Feldman, Liu, Nath, BDN, PRD 80 (2009) 075001

Mass	Mixed LSP	Pure Wino LSP
$m_{\tilde{N}_1}$	198.9	195.2
$m_{\tilde{N}_2}$	217.0	357.0
$m_{\tilde{N}_3}$	429.9	1025
$m_{\tilde{N}_4}$	451.3	1029
$m_{\tilde{C}_1}$	208.8	195.5
$m_{\tilde{C}_2}$	448.6	1036
$m_{\tilde{t}_1}$	648.5	1516
$m_{\tilde{t}_2}$	866.8	1749
$m_{\tilde{b}_1}$	841.4	1729
$m_{\tilde{b}_2}$	970.2	1902
$m_{\tilde{\tau}_1}$	817.7	1011
$m_{\tilde{\tau}_2}$	822.8	1041
$m_{\tilde{g}}$	707.1	1929

Signature	Mixed LSP		Pure Wino LSP	
	Events	S/\sqrt{B}	Events	S/\sqrt{B}
Multijets	8766	183.74	50	1.05
Lepton + jets	2450	32.25	26	0.34
OS dileptons + jets	110	6.39	4	0.23
SS dileptons + jets	60	11.77	0	NA
Trileptons + jets	14	2.47	0	NA

- Big impact of gluino mass in number of multijet events
- Small mass gaps significantly reduce number of leptonic events

M_{eff} Distribution



10 fb⁻¹ at $\sqrt{s} = 14$ TeV

Benchmark Models II: CDMS-II Examples

Holmes and BDN, PRD 81 (2010) 055002

Point	C	D	E
δ_2	-0.6	0.82	-0.47
δ_3	-0.3	-0.35	-0.3
B%	0.3%	5.4%	40.9%
W%	95.8%	0.5%	53.0%
H%	3.9%	94.1%	6.1%
$m_{\tilde{N}_1}$	175	112	230
$m_{\tilde{N}_2}$	235	130	239
$m_{\tilde{N}_3}$	505	252	504
$m_{\tilde{N}_4}$	513	846	515
$m_{\tilde{C}_1}$	175	123	234
$m_{\tilde{C}_2}$	514	846	515
$m_{\tilde{g}}$	952	890	951
$m_{\tilde{t}_1}$	719	544	709
$m_{\tilde{t}_2}$	862	964	865
$m_{\tilde{b}_1}$	809	766	812
$m_{\tilde{b}_2}$	874	943	871
$m_{\tilde{\tau}_1}$	344	338	352
$m_{\tilde{\tau}_2}$	414	752	424
m_h	113	114	113
$\sigma_{\text{SUSY}}^{7 \text{ TeV}}$ (pb)	1.2	2.7	0.4
$\sigma_{\text{SUSY}}^{10 \text{ TeV}}$ (pb)	2.5	5.1	1.3
$\sigma_{\text{SUSY}}^{14 \text{ TeV}}$ (pb)	5.7	10.0	3.7

- All models can produce signals at CDMS II – C & E can fit PAMELA data as well
- Signal simulated: 1 fb^{-1} at $\sqrt{s} = 14 \text{ TeV}$
- Again, healthy multijets but disappearance of leptonic events

Numbers of Events

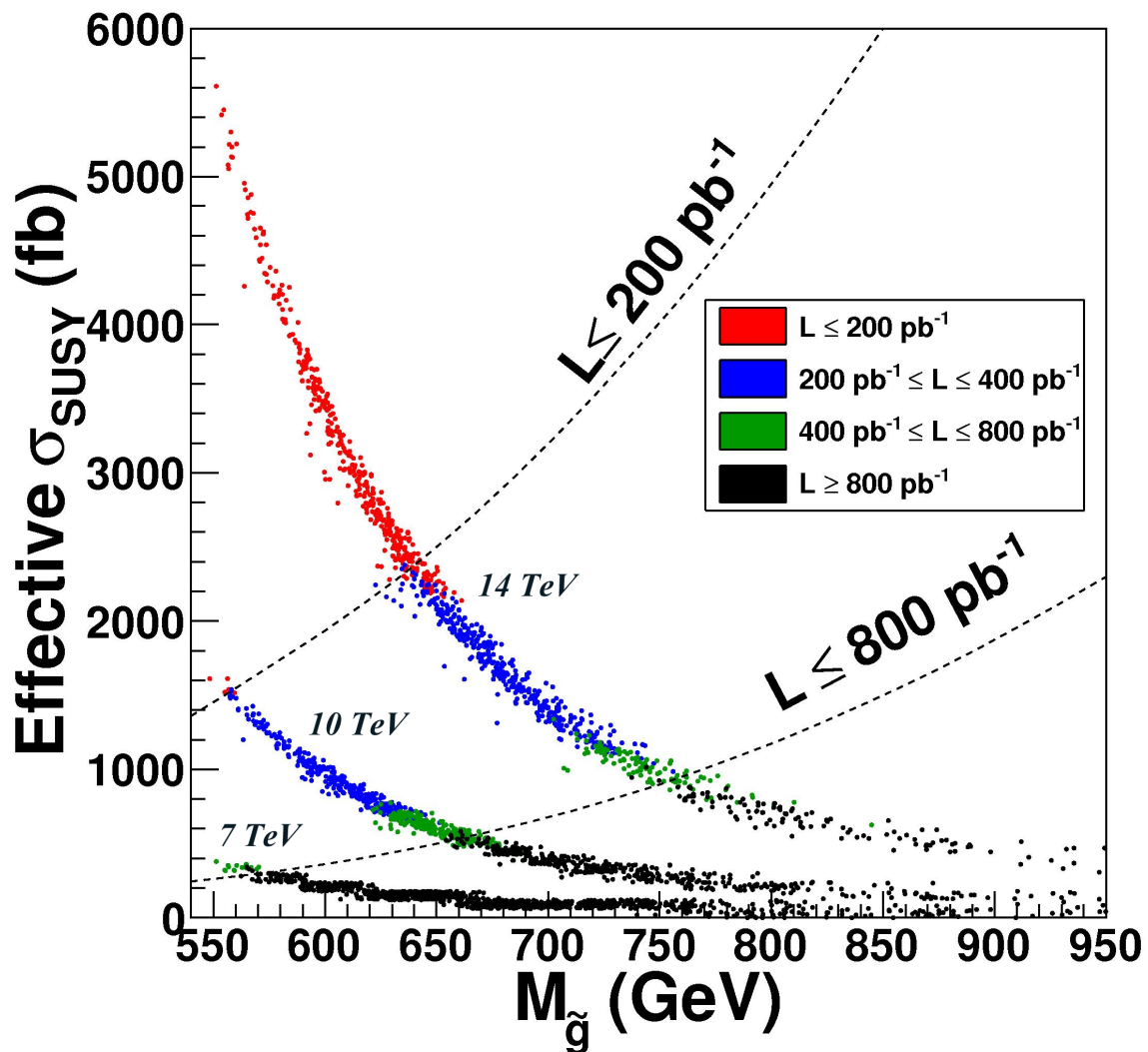
Point	C	D	E
Multijets	402	436	298
$1l + \text{jets}$	202	310	111
OS $2l + \text{jets}$	12	45	7
SS $2l + \text{jets}$	6	16	3
$3l + \text{jets}$	4	6	1

Significance S/\sqrt{B}

Point	C	D	E
Multijets	26.9	29.1	19.9
$1l + \text{jets}$	8.2	12.5	4.5
OS $2l + \text{jets}$	2.0	7.4	1.2
SS $2l + \text{jets}$	2.3	6.0	1.1
$3l + \text{jets}$	1.6	2.5	0.4

⇒ General rule: Discovery of DM-motivated models needs a light gluino

Feldman, Kane, Lu, BDN, PLB 687 (2010) 363

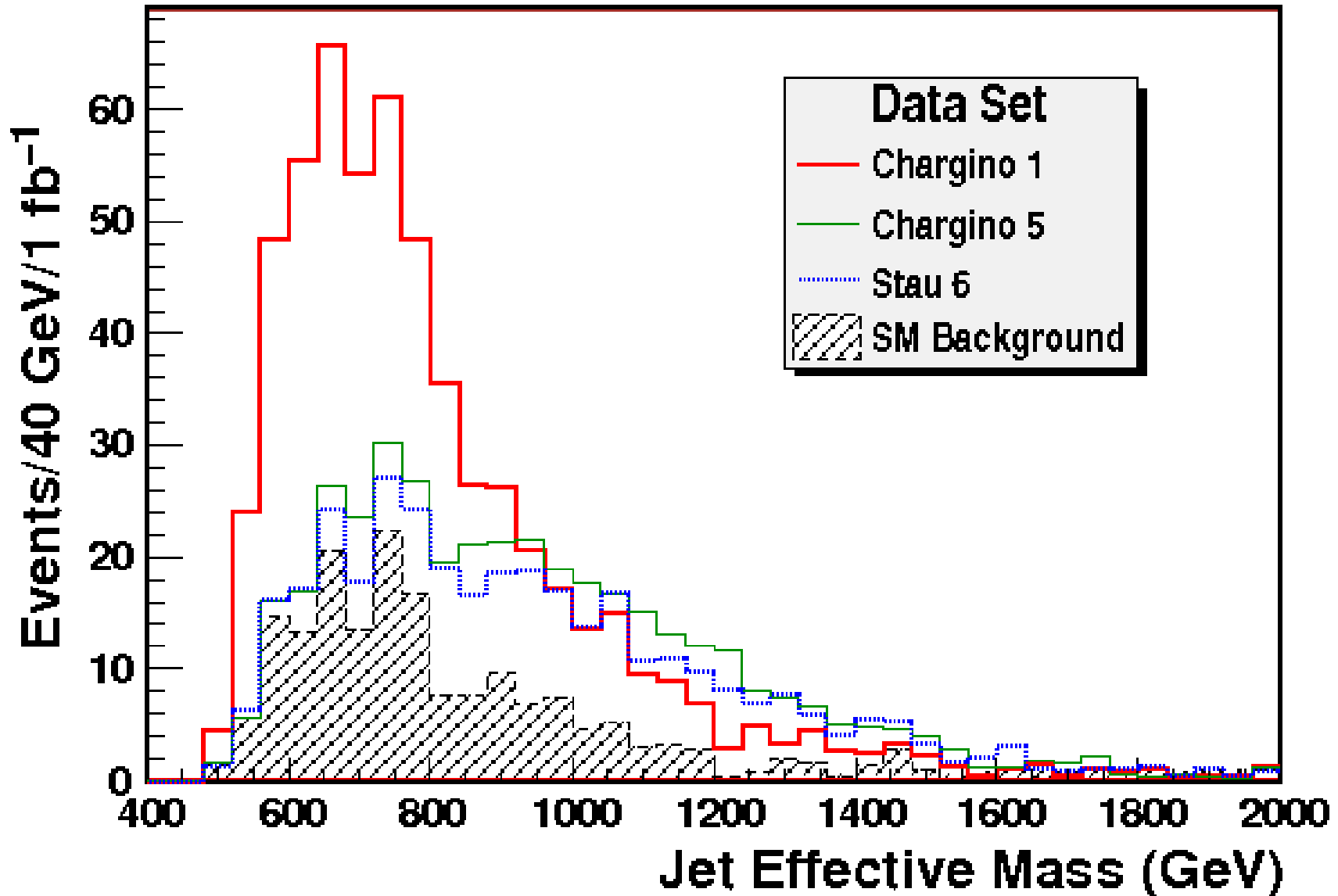


- High wino-content (for PAMELA) implies small mass gap between \tilde{C}_1/\tilde{N}_2 and LSP
- Result: major reduction in expected leptonic SUSY signatures
- Increasing Higgsino content to match CDMS (and photon data) requires a light gluino
- Result: multijet signals may be our only handle

⇒ We will need to learn how to do more with less!

Must look for new signatures targeted to non-universalities in gaugino sector

Altunkaynak, Holmes, Nath, BDN & Peim, arXiv:1008.3423



⇒ All three are mirage-type models; “Stau 6” is in the PAMELA-preferred regime

Epilogue: Dark Matter and the LHC Inverse Problem

The “LHC Inverse Problem” – in One Slide

Arkani-Hamed, Kane, Thaler & Wang, “Supersymmetry and the LHC inverse problem,” JHEP **0608**, 070 (2006)

⇒ Basic premise: multiple SUSY parameter sets likely to fit the LHC data

- After many years of data-taking at LHC we have a well-defined “signature space” bounded only by experimental errors
- A global fit to a “minimal” multi-parameter MSSM model is performed

$$\left\{ \begin{array}{l} \tan \beta, \mu, M_1, M_2, M_3 \\ m_{Q_{1,2}}, m_{U_{1,2}}, m_{D_{1,2}}, m_{L_{1,2}}, m_{E_{1,2}} \\ m_{Q_3}, m_{U_3}, m_{D_3}, m_{L_3}, m_{E_3} \\ m_A, A_t, A_b, A_\tau \end{array} \right\}$$

- Several (isolated) points in “parameter space” will be good fits to the observed data

⇒ This conjecture was explicitly verified by brute force

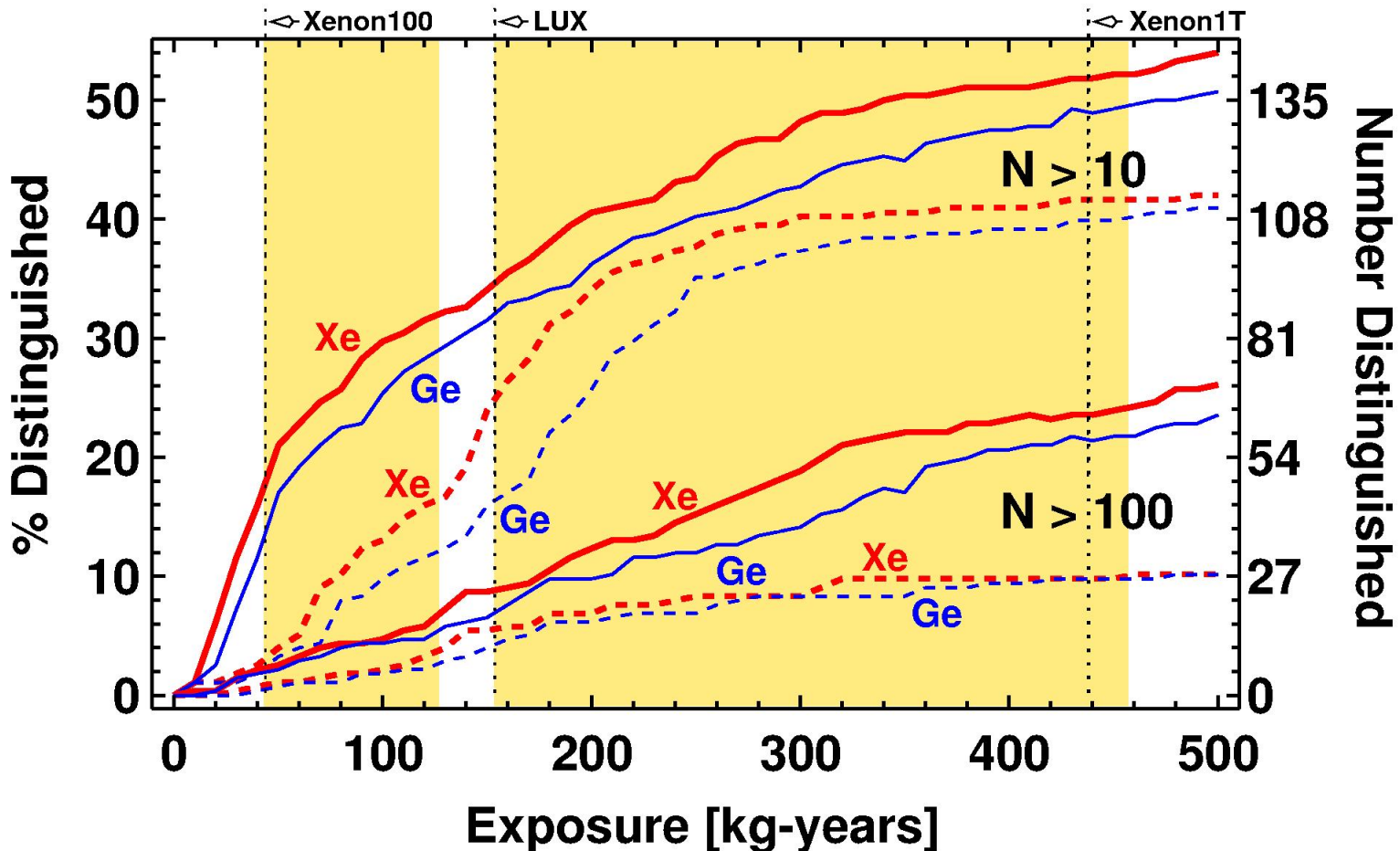
- 45,000 SUSY **models** sampled and 10 fb^{-1} of simulated data generated
- Out of 45,000 test models they found 276 degenerate **pairs**

Direct Detection and SUSY Degeneracies

⇒ Requirements to consider a degenerate pair resolved:

Altunkaynak, Holmes, BDN, JHEP 0810 (2008) 013

1. Counts N_A and N_B ($N_i = \text{rate}_i \times \text{exposure}$) must *both* exceed N_{\min} events
2. The two quantities N_A and N_B must differ by at least $5\sigma^{AB}$
3. For the moment assume statistical errors only: $\sigma^{AB} = \sqrt{N_A + N_B}$



- Dark matter hints may disfavor pure Bino LSP (i.e. mSUGRA)
- PAMELA needs wino predominance; CDMS/photons want strong Higgsino admixture
- Such models find a natural home in many (all?) semi-realistic string constructions
- Likely that mass gaps between \tilde{C}_1/\tilde{N}_2 and LSP small, so leptonic signatures a bust
- Will need to learn to do more with jet-based signatures and hope the gluino is lighter than in mSUGRA models

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- ⇒ *Gaugino sector is truly a window on the high-energy world: we may be on the verge of revolutionary discoveries!*