



# Systematic studies of parton-Quark Gluon Plasma interactions with ATLAS

---

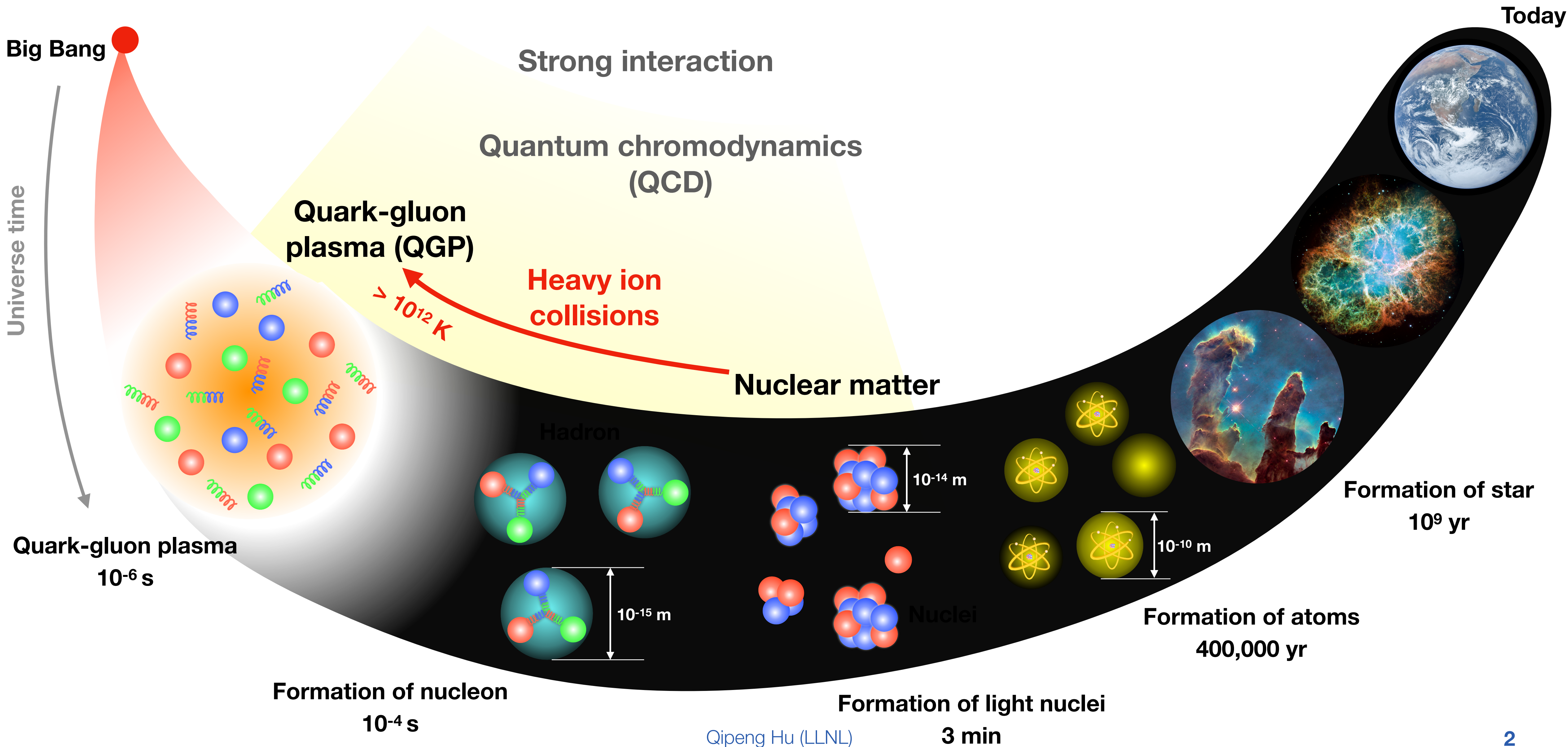
Qipeng Hu, LLNL

For the ATLAS Collaboration

LHC Seminar

May 10, 2022

# QCD matter — Quark-gluon plasma (QGP)





# QCD Parton model in vacuum

QCD factorization theorem (leading twist):

$$d\sigma(A + B \rightarrow C + X) = f_{a/A}(x_a, \mu_F^2) \otimes f_{b/B}(x_b, \mu_F^2) \otimes d\sigma_{ab \rightarrow cd}(Q^2, \mu_R) \otimes D_{C/c}(z_c, M^2)$$

Perturbative QCD (pQCD) is used to calculate

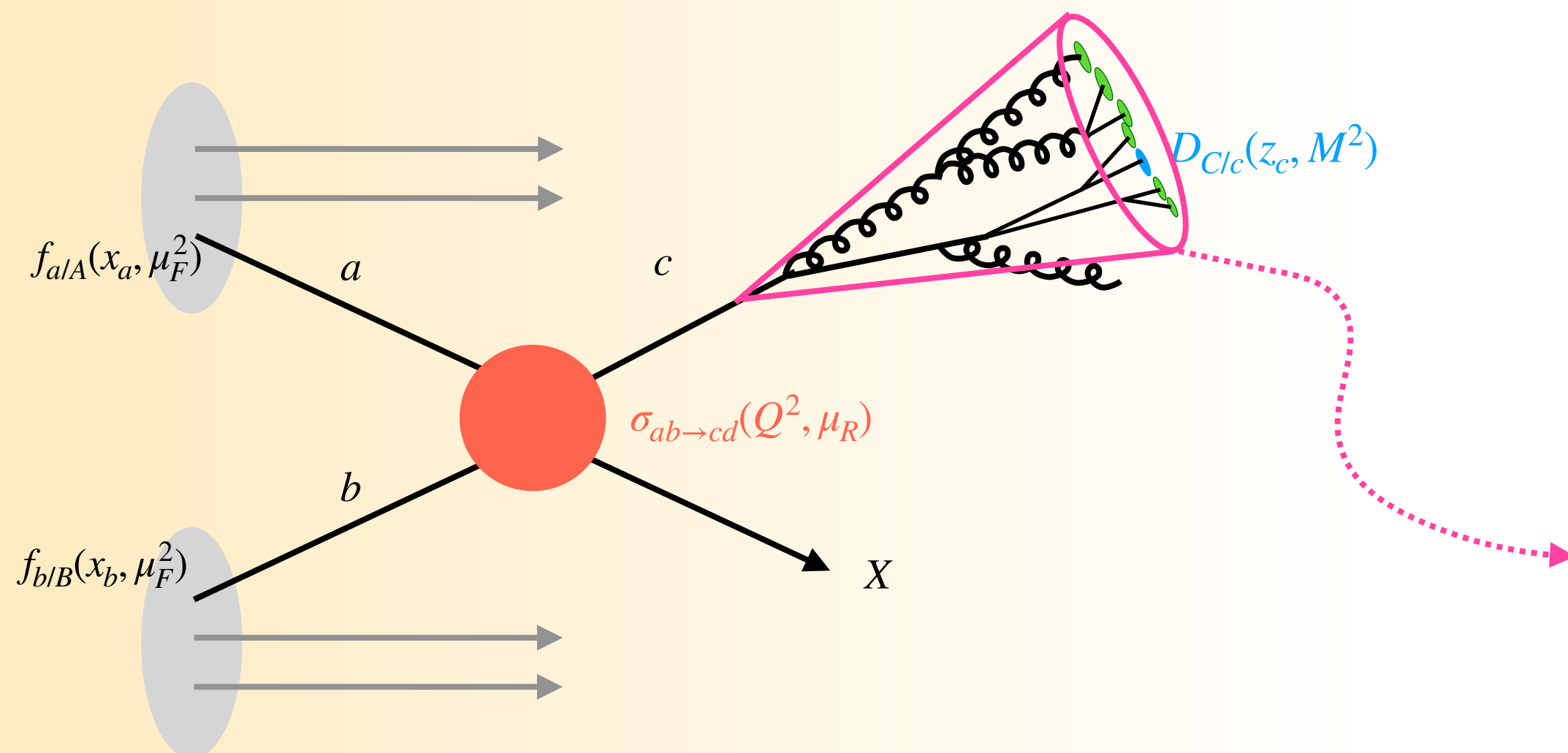
- Partonic cross section

Factorization scale  $\mu_F$

Renormalization scale  $\mu_R$

Non-perturbative, universal

- Parton distribution functions (PDF)
- Fragmentation function (FF)



**Jet:** access the whole shower of initial parton

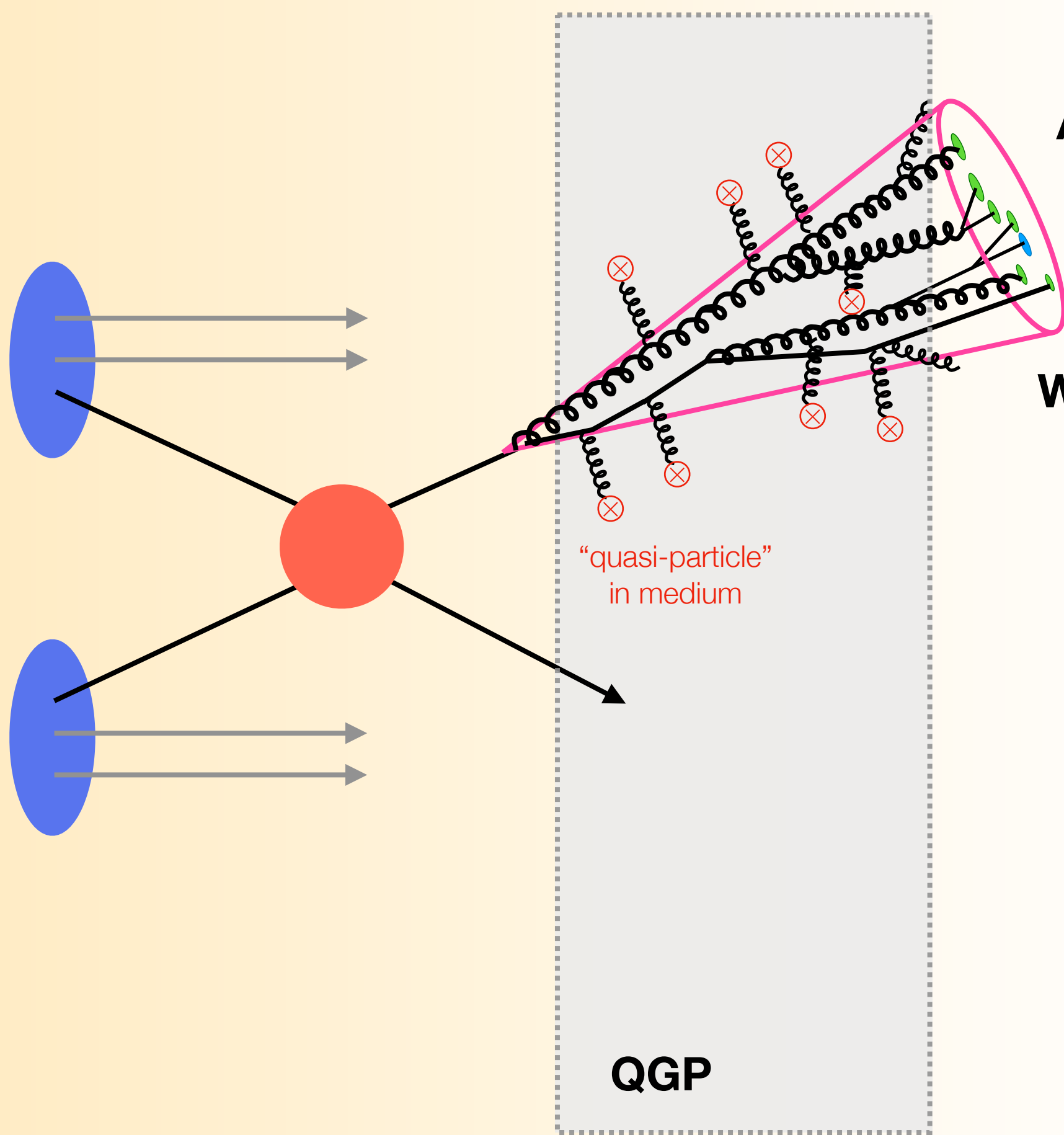
# QCD factorization in medium

QCD factorization in vacuum:

$$d\sigma(A + B \rightarrow C + X) = f_{a/A}(x_a, \mu_F^2) \otimes f_{b/B}(x_b, \mu_F^2) \otimes d\sigma_{ab \rightarrow cd}(Q^2, \mu_R) \otimes D_{C/c}(z_c, M^2)$$

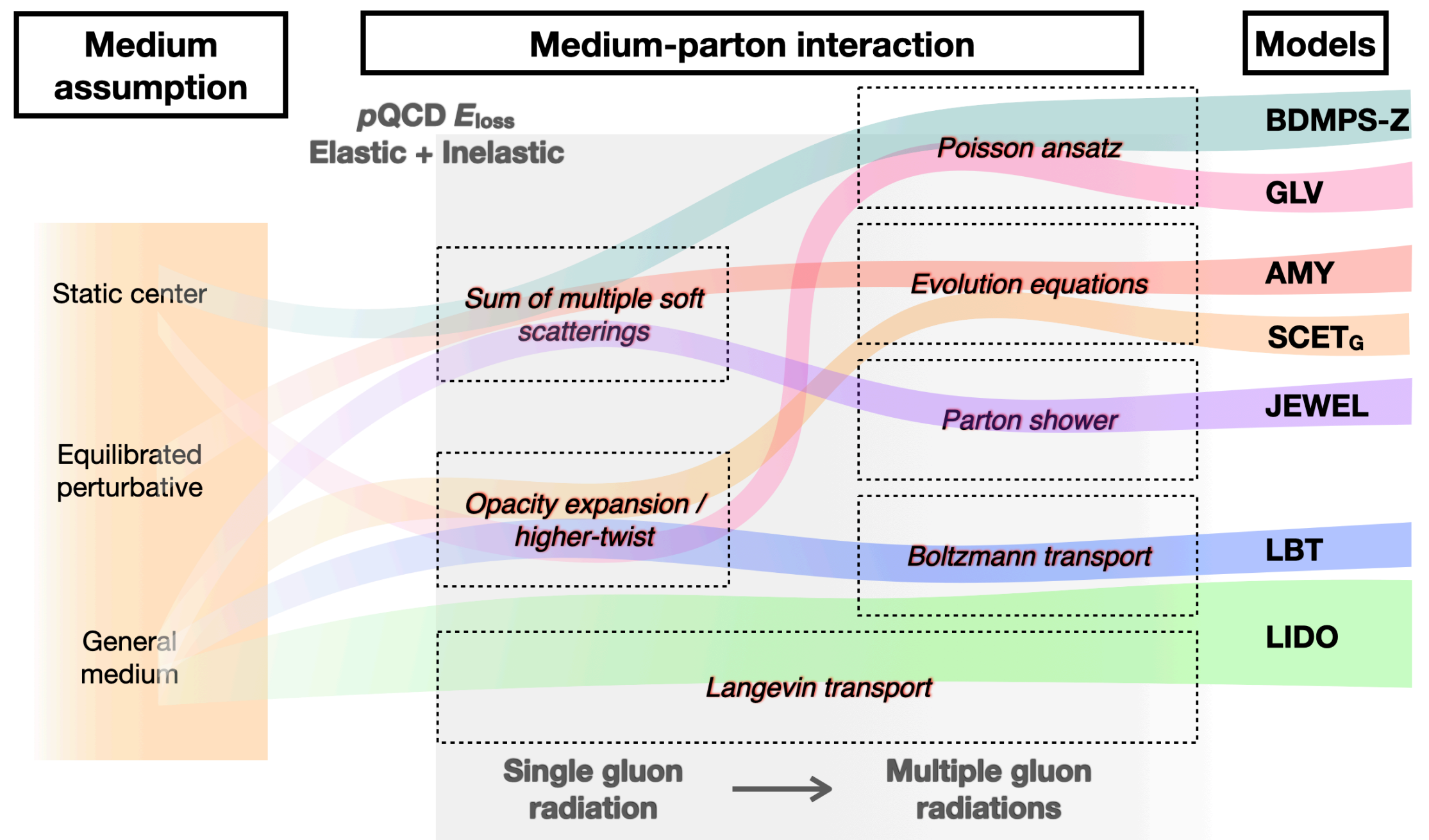
Assumed QCD factorization in QGP:

$$d\sigma(A + B \rightarrow C + X)|_{\text{geometry}} = f_{a/A}^{A,Z}(x_a, \mu_F^2) \otimes f_{b/B}^{A,Z}(x_b, \mu_F^2) \otimes d\sigma_{ab \rightarrow cd}(Q^2, \mu_R) \otimes P_{c \rightarrow c'}(T_{QGP}) \otimes D_{C/c'}(z_c, M^2)$$



Weakly coupled approach

Parton-medium interaction



\*\*\* Many models are more advanced than shown here



# QCD factorization in medium — Cont.

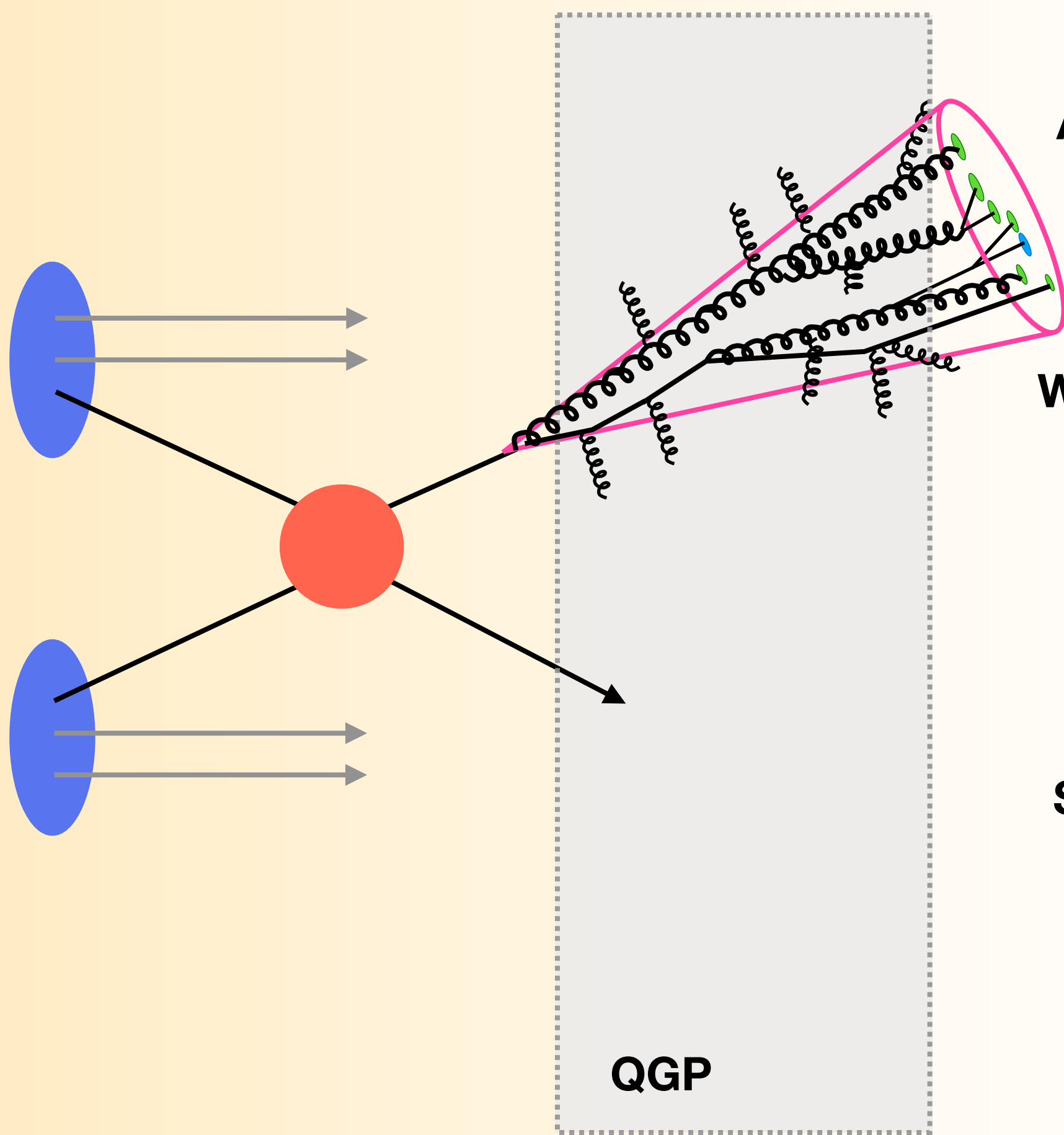
## QCD factorization in vacuum:

$$d\sigma(A + B \rightarrow C + X) = f_{a/A}(x_a, \mu_F^2) \otimes f_{b/B}(x_b, \mu_F^2) \otimes d\sigma_{ab \rightarrow cd}(Q^2, \mu_R) \otimes D_{C/c}(z_c, M^2)$$

## Assumed QCD factorization in QGP:

$$d\sigma(A + B \rightarrow C + X)|_{\text{geometry}} = f_{a/A}^{A,Z}(x_a, \mu_F^2) \otimes f_{b/B}^{A,Z}(x_b, \mu_F^2) \otimes d\sigma_{ab \rightarrow cd}(Q^2, \mu_R) \otimes P_{c \rightarrow c'}(T_{QGP}) \otimes D_{C/c'}(z_c, M^2)$$

Parton-medium interaction



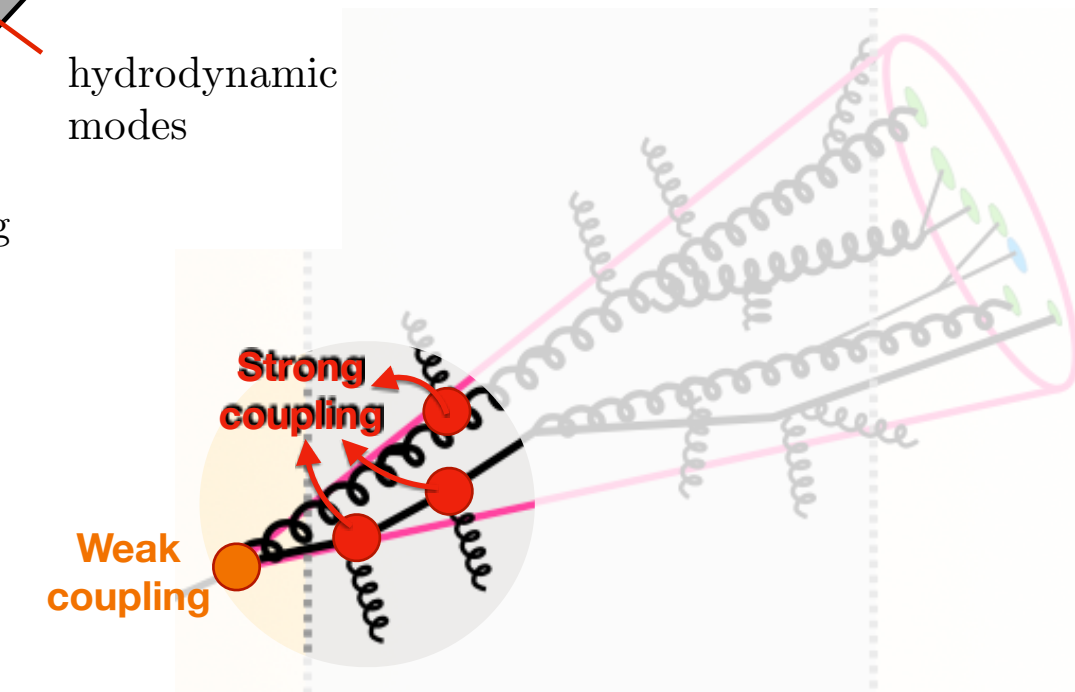
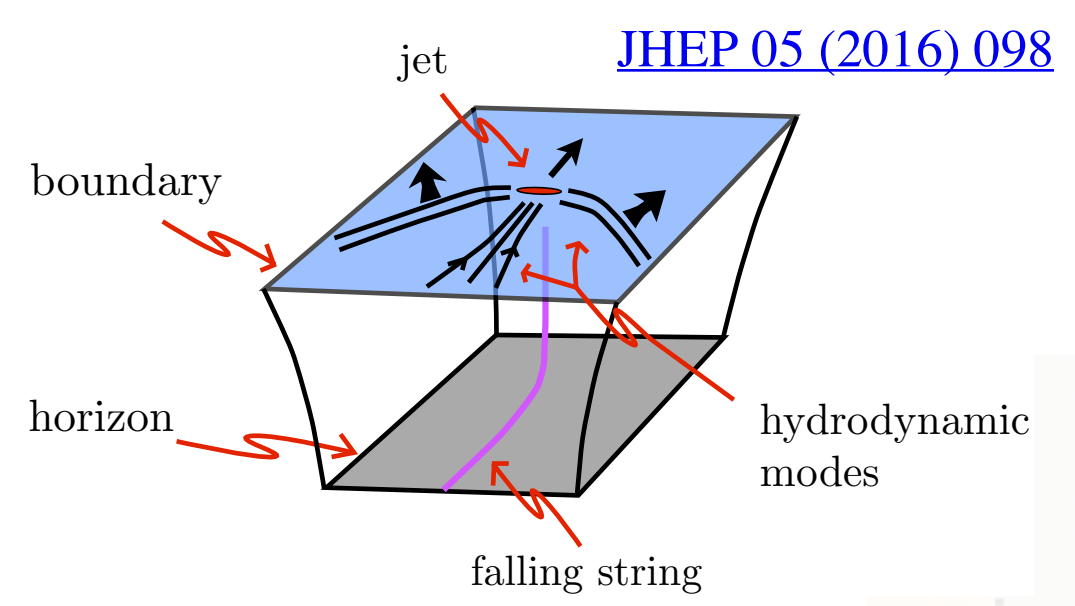
## Weakly coupled approach

- Quenching formalism: BDMPS-Z, GLV, AMY, SCET<sub>G</sub> etc.
- Transport: LBT, LIDO etc.

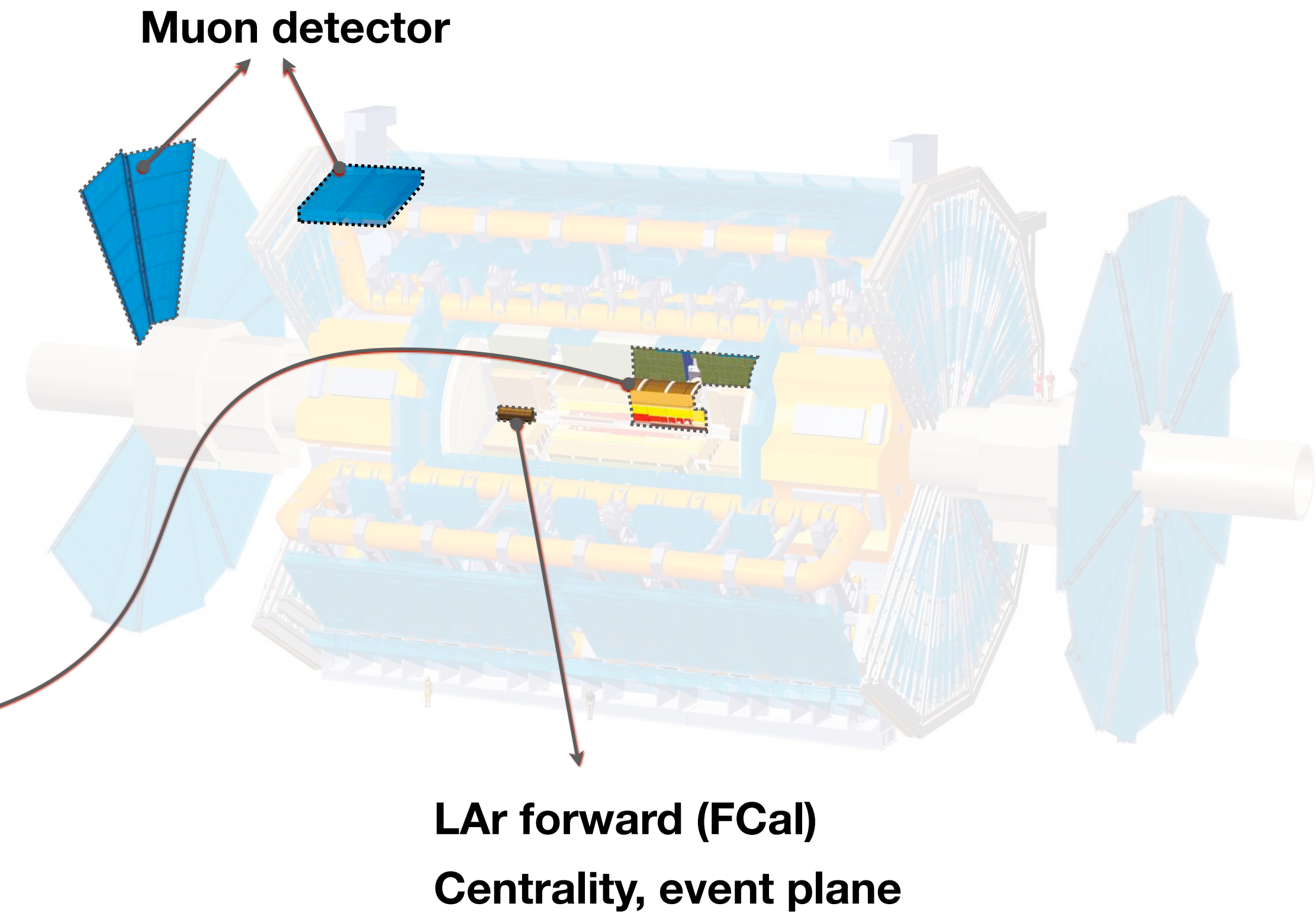
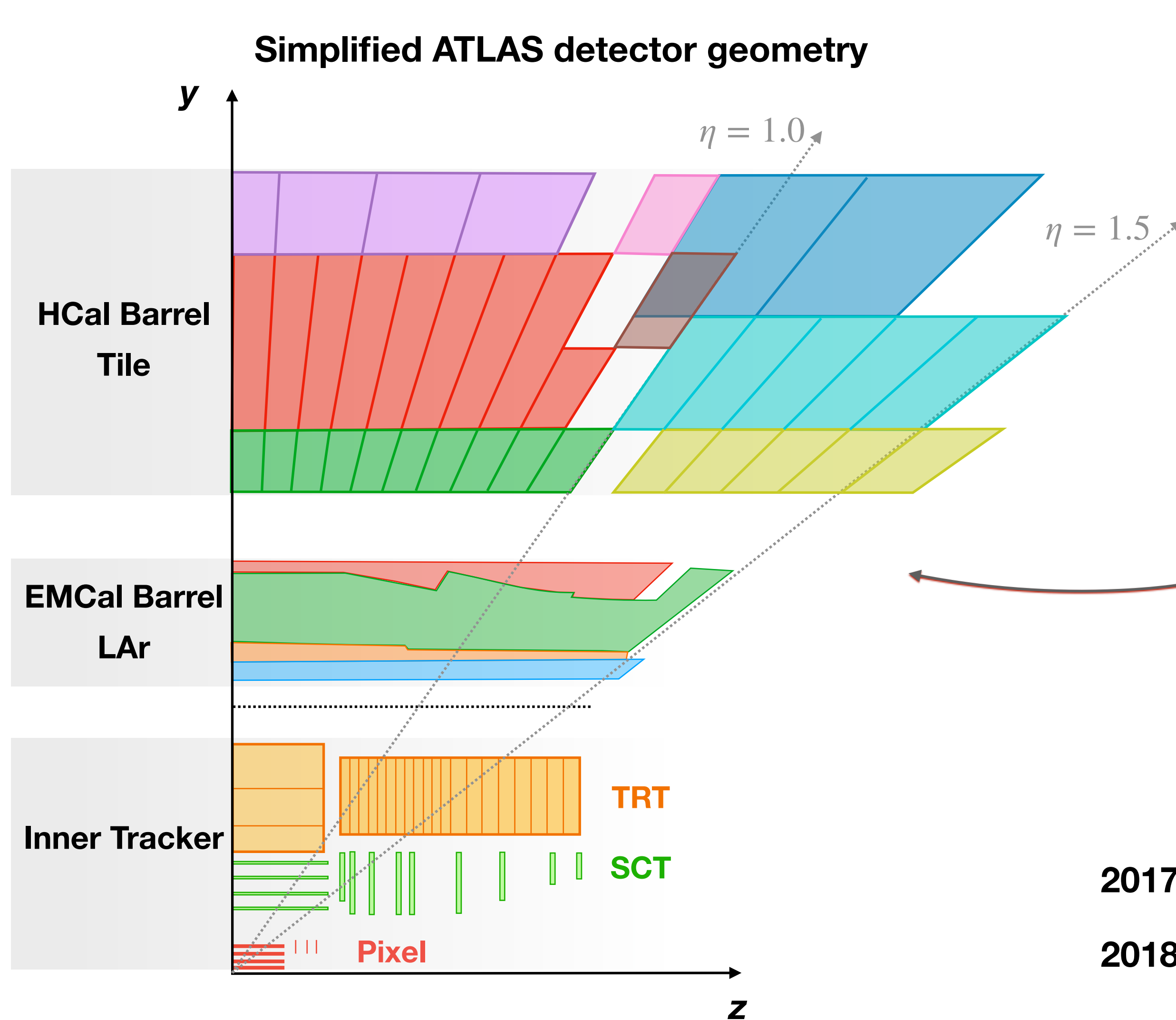
## Strongly coupled approach

- Holographic “jet”
- Strong/weak coupling hybrid

- Multiple scales involved:
- Vacuum QCD scales:  $\mu_F, \mu_R$
  - Medium scales:  $T, \Lambda_{\text{med}}, L_{\text{path}}$ , etc.
  - Leading parton:  $m^2, r_{\perp}$ , etc



# ATLAS detector



2017 *pp* collisions, 5.02 TeV,  $\int L dt = 272 \text{ pb}^{-1}$

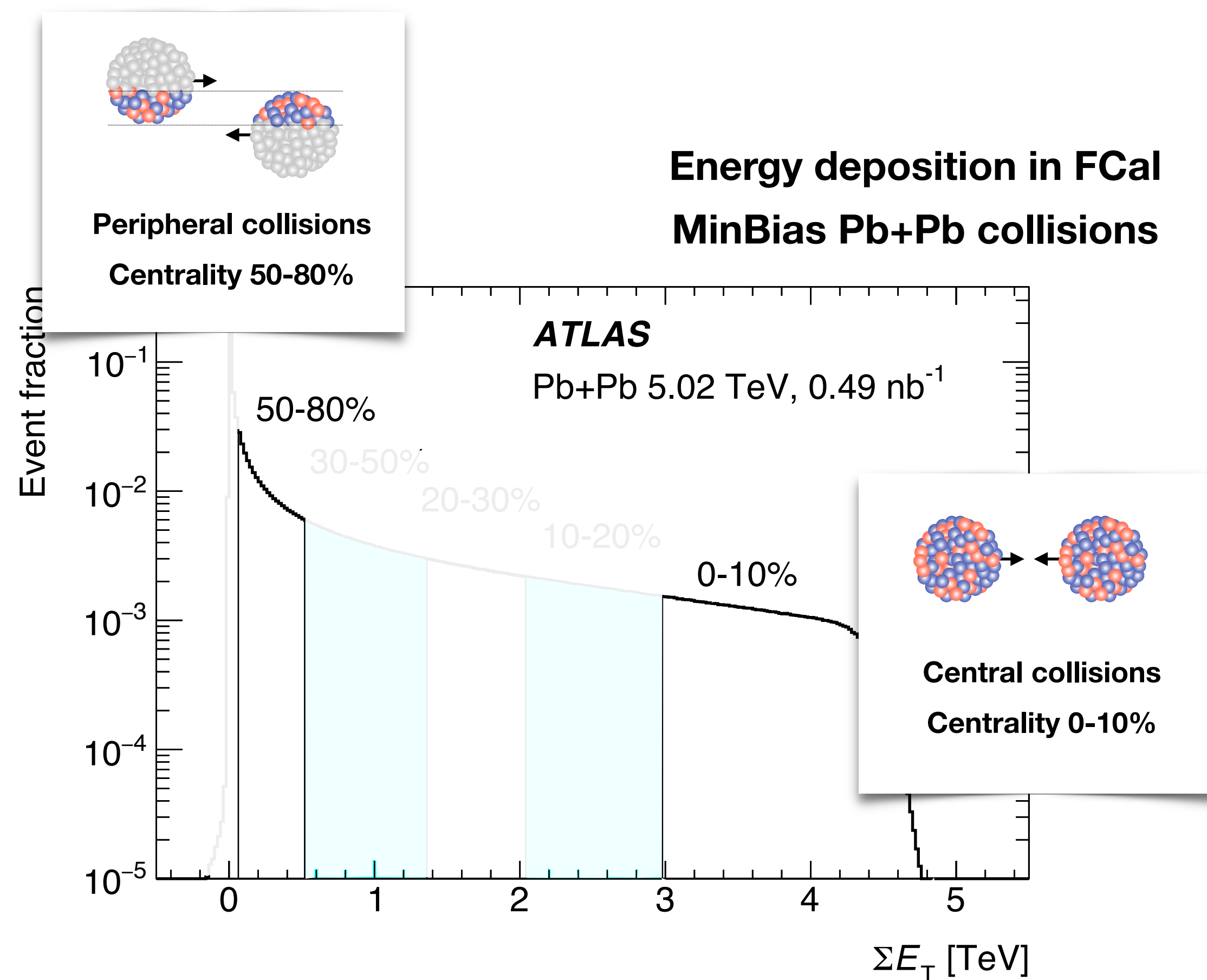
2018 Pb+Pb collisions, 5.02 TeV,  $\int L dt = 1.76 \text{ nb}^{-1}$



# Centrality

[Phys. Lett. B 789 \(2019\) 167](#)

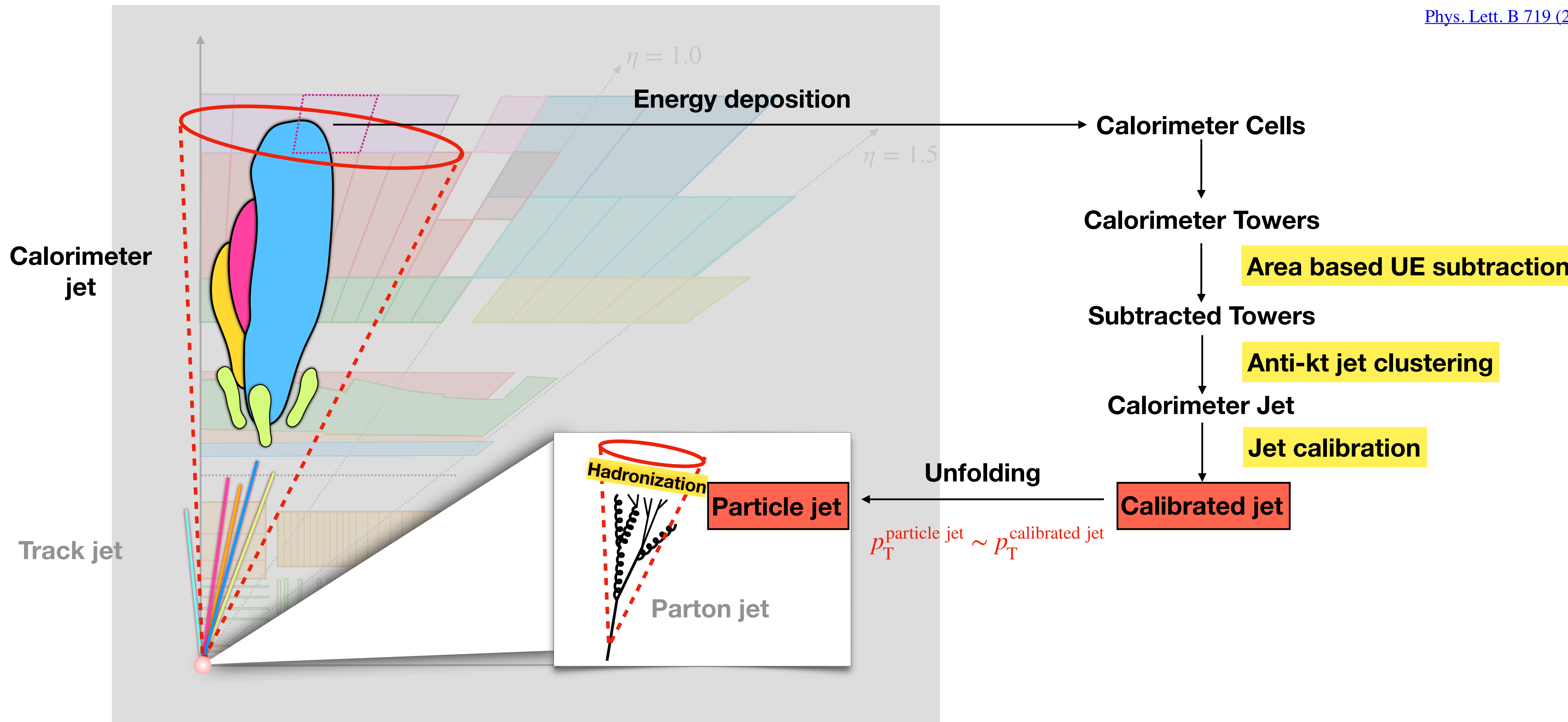
**Low activity**  
**Low Number of binary**  
**~ 50**



**10% events with highest activities**  
**High Number of binary collisions ~ 1600**

# Jet reconstruction — Heavy ion version

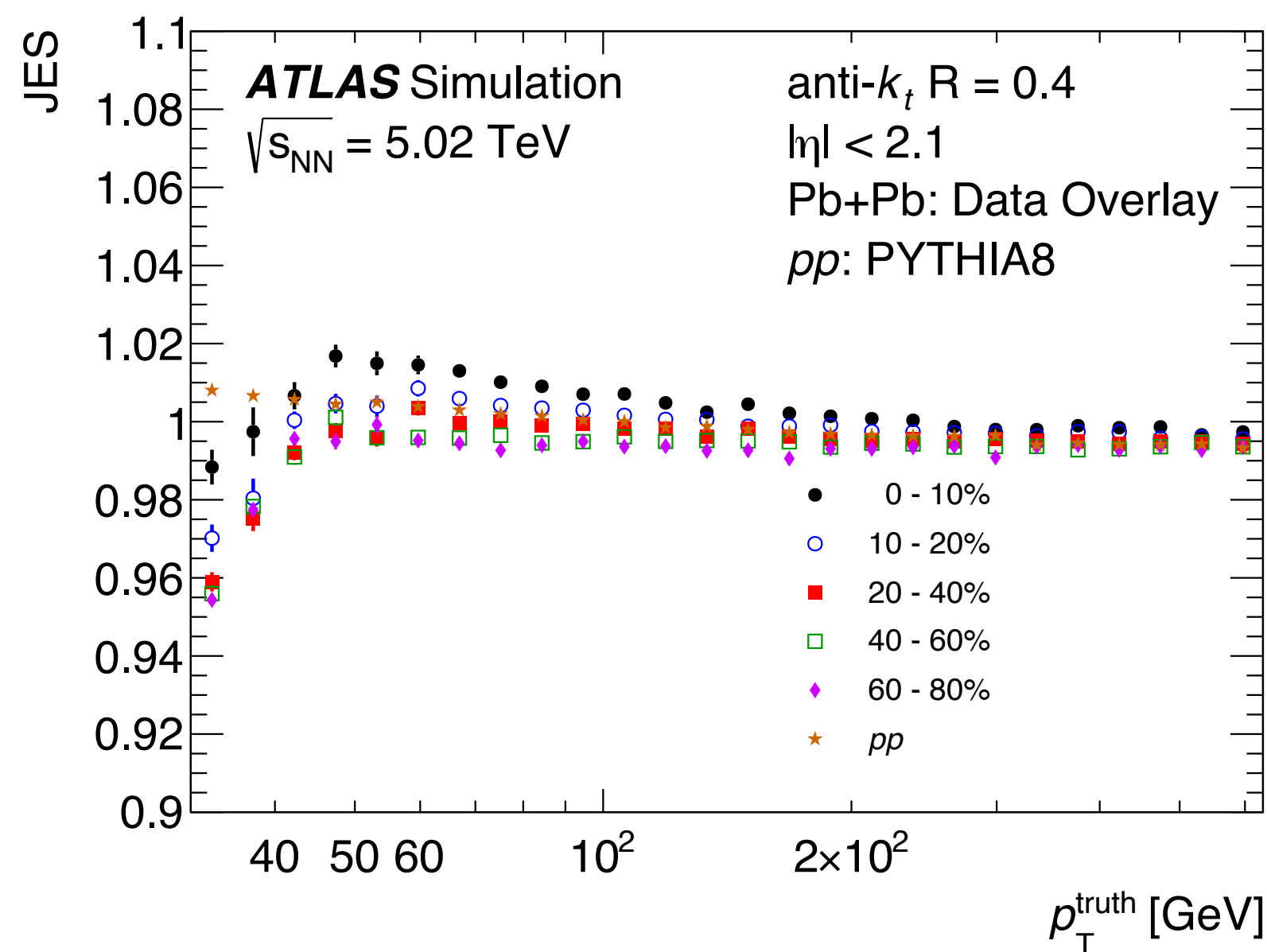
[Phys. Lett. B 719 \(2013\) 220-241](#)



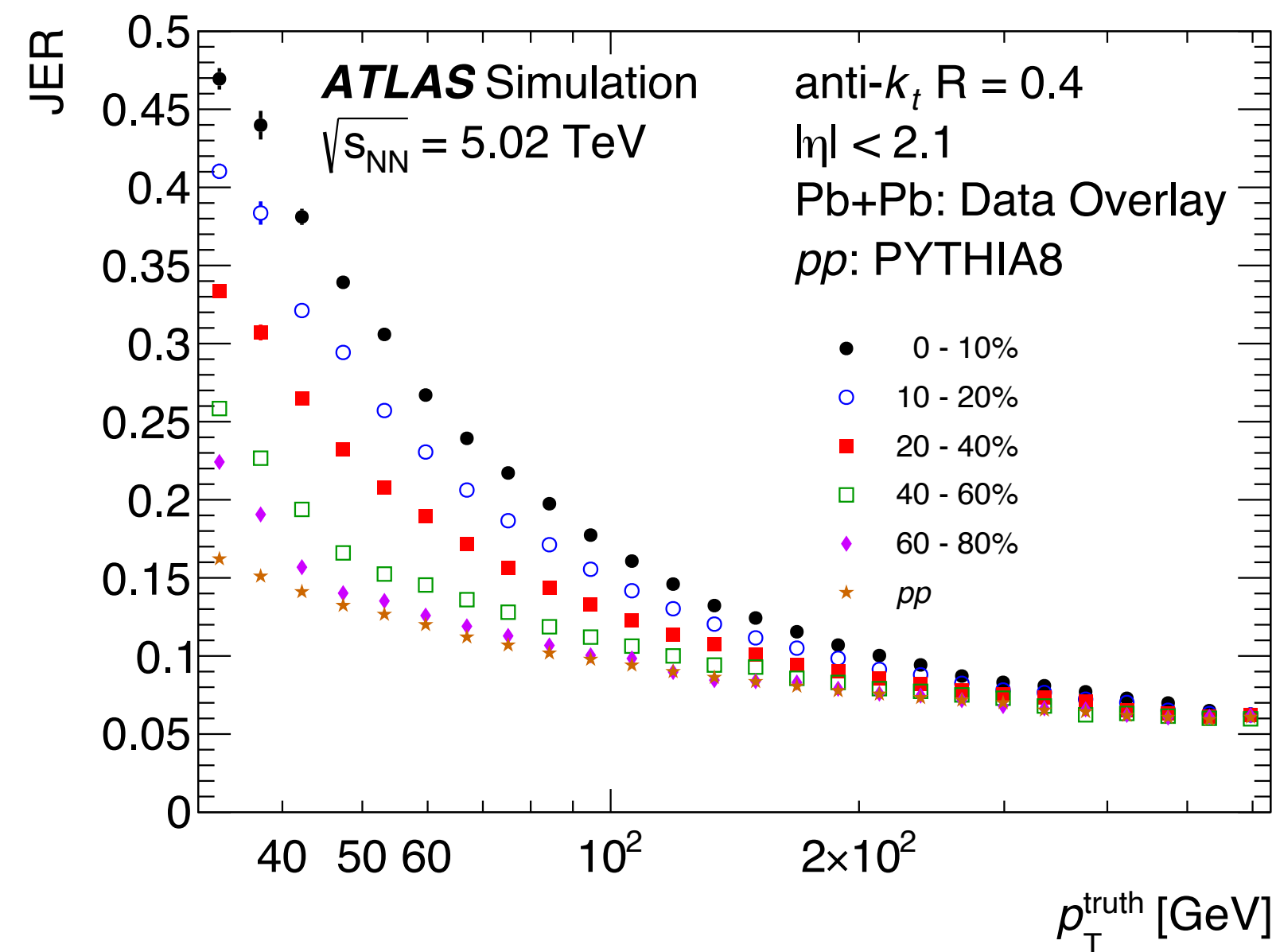


# Jet performance in heavy ion collisions

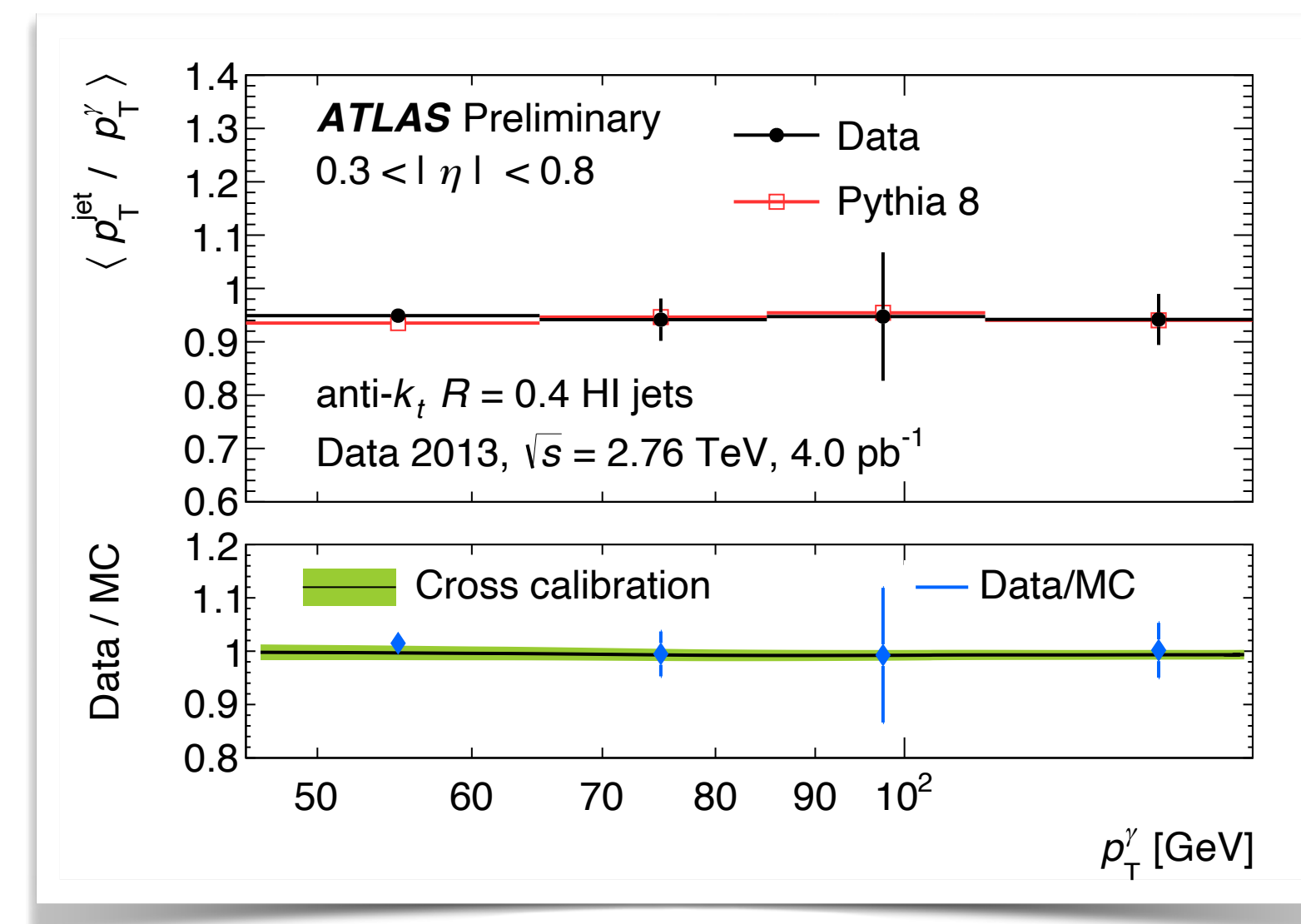
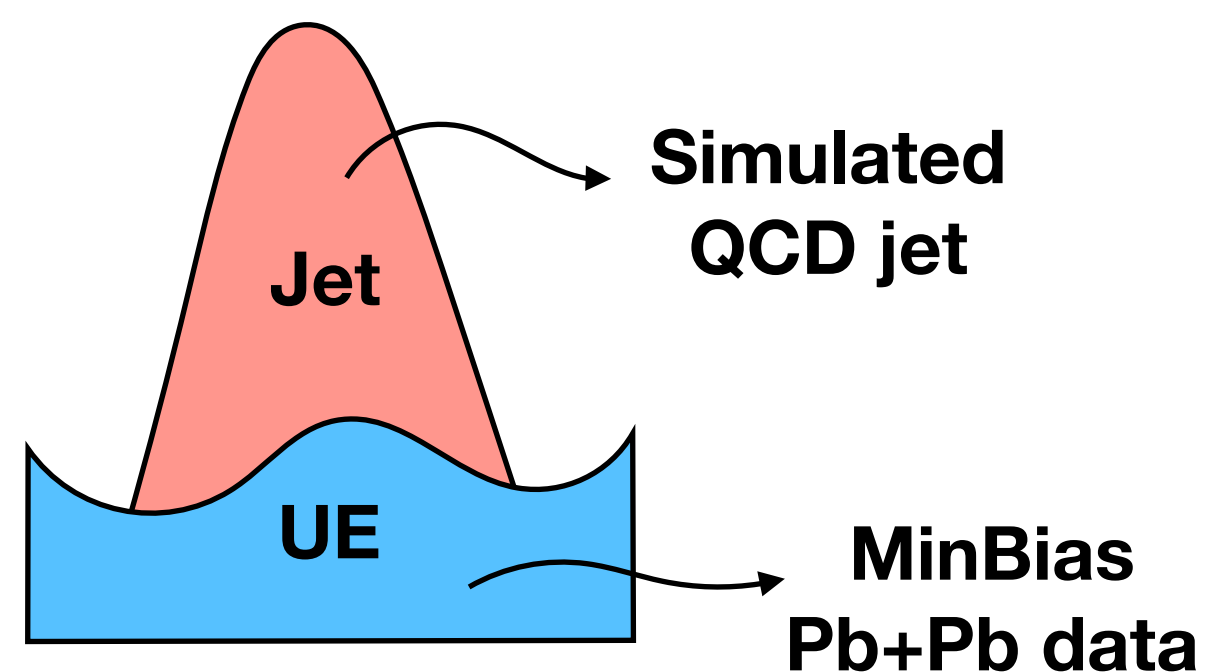
[arXiv:2205.00682](https://arxiv.org/abs/2205.00682)  
[ATLAS-CONF-2015-016/](https://atlas.conf-2015-016/)



**Jet Energy Scale (JES)**



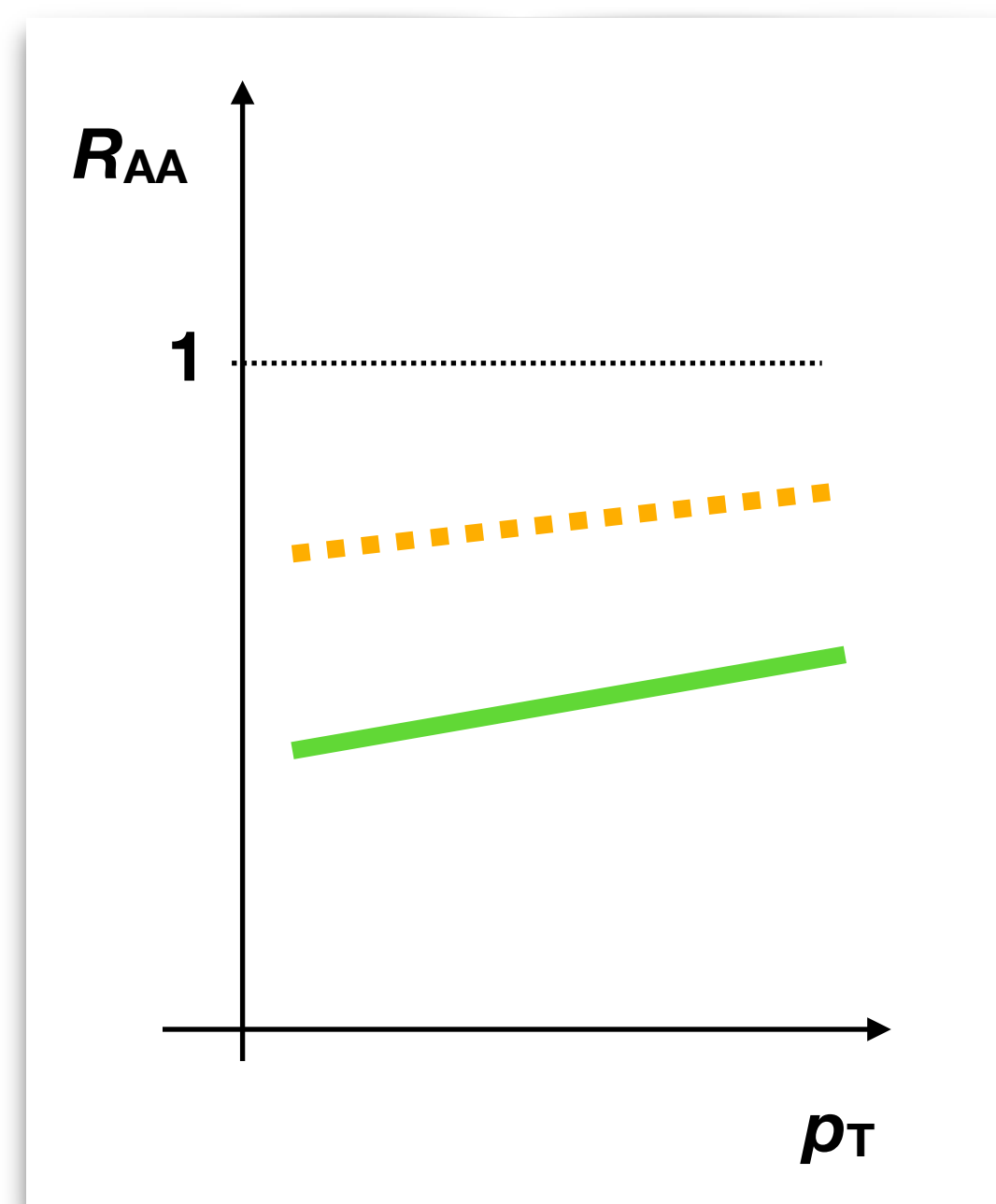
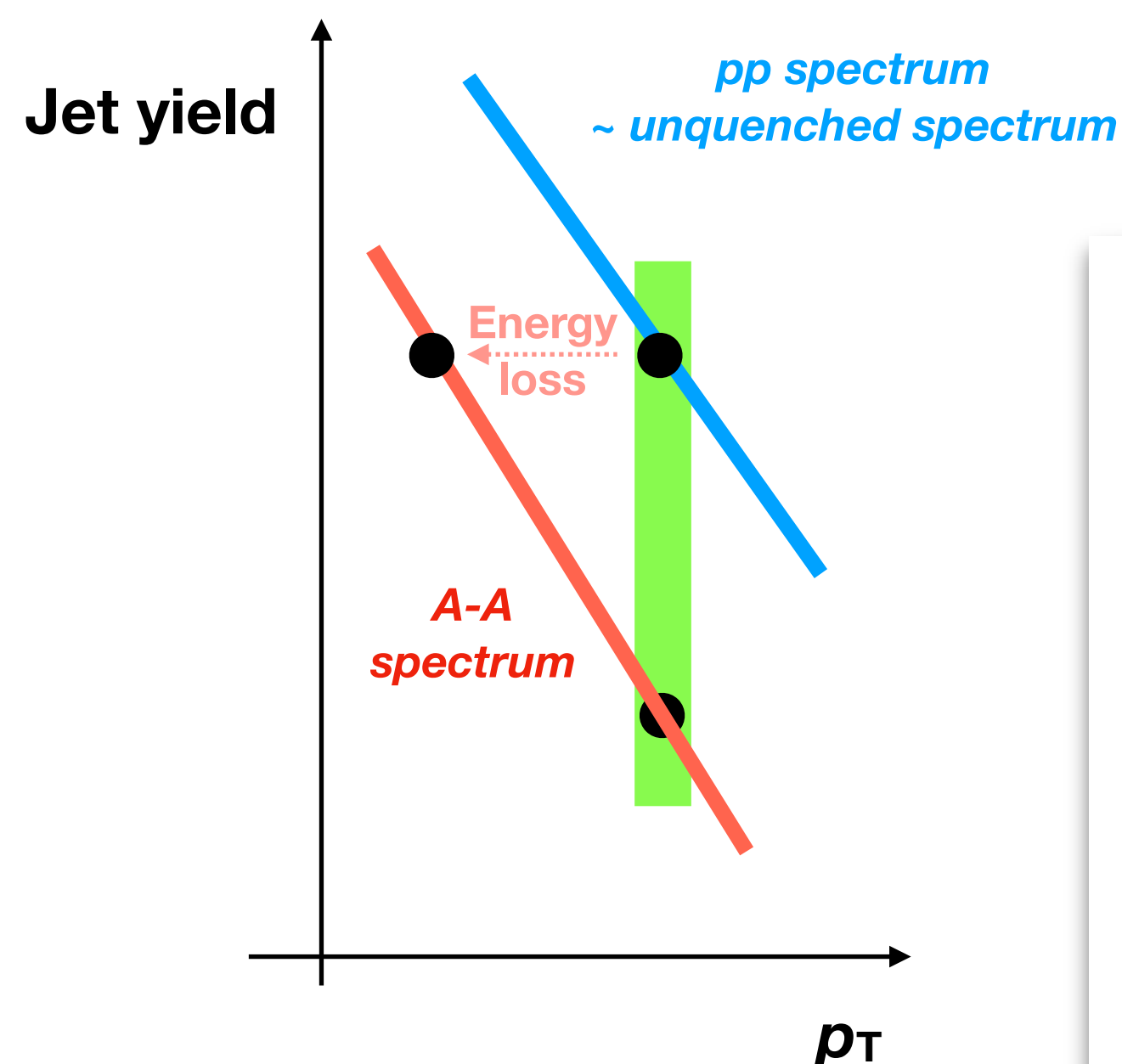
**Jet Energy Resolution (JER)**



**Good Data/MC agreement**

# Measuring jet quenching

PLB 790 (2019) 108



Jet production: heavy ion (**A-A**) collisions vs. *pp* collisions

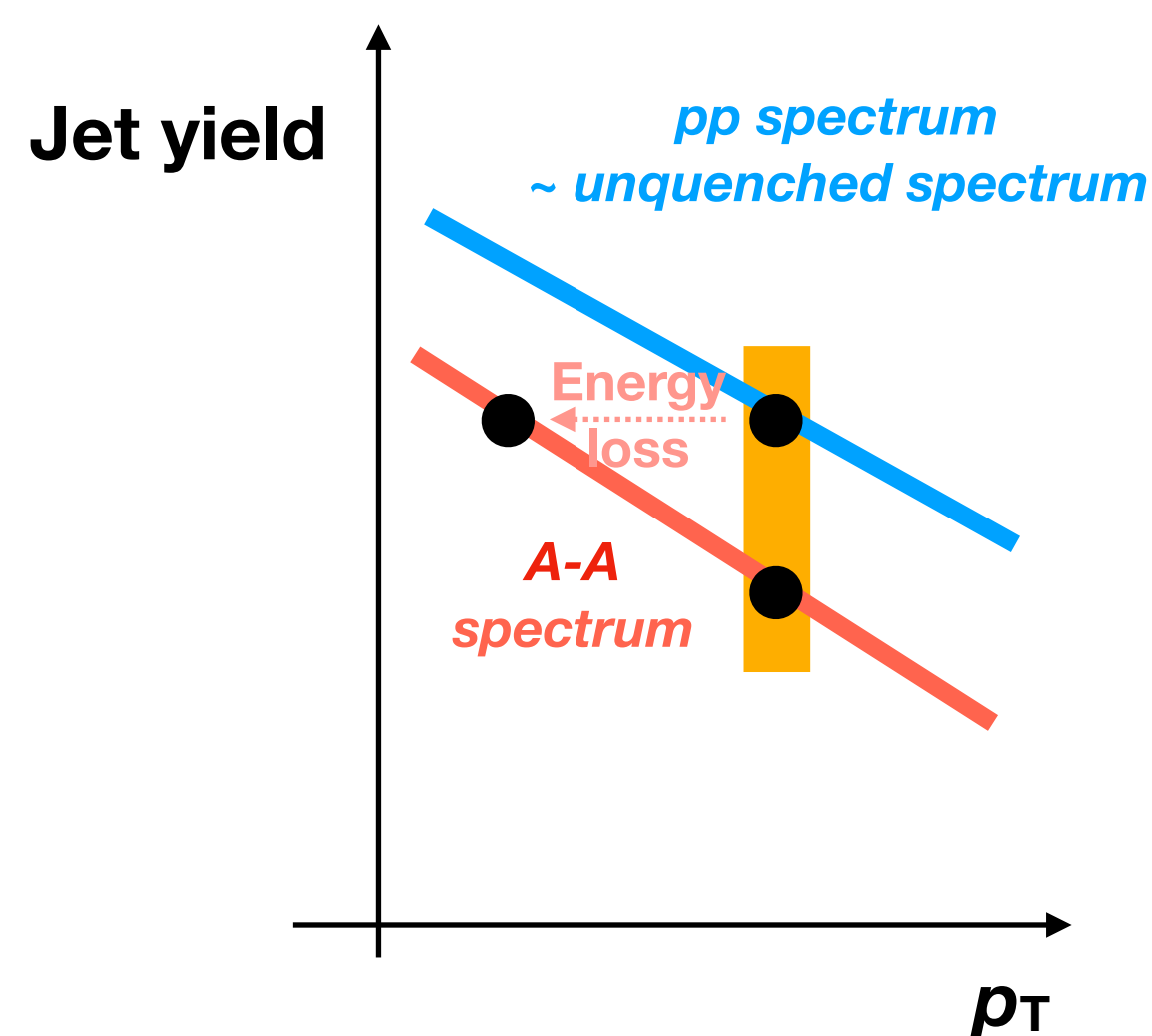
**Nuclear modification factor**

$$R_{AA} = \frac{\text{per-NN yields in A-A}}{\text{yields in } pp}$$

$R_{AA} < 1$  due to jet-QGP interaction: jets in A-A collisions are suppressed / **quenched**

$R_{AA}$  depends on:

- Energy loss
- Jet spectrum before quenching

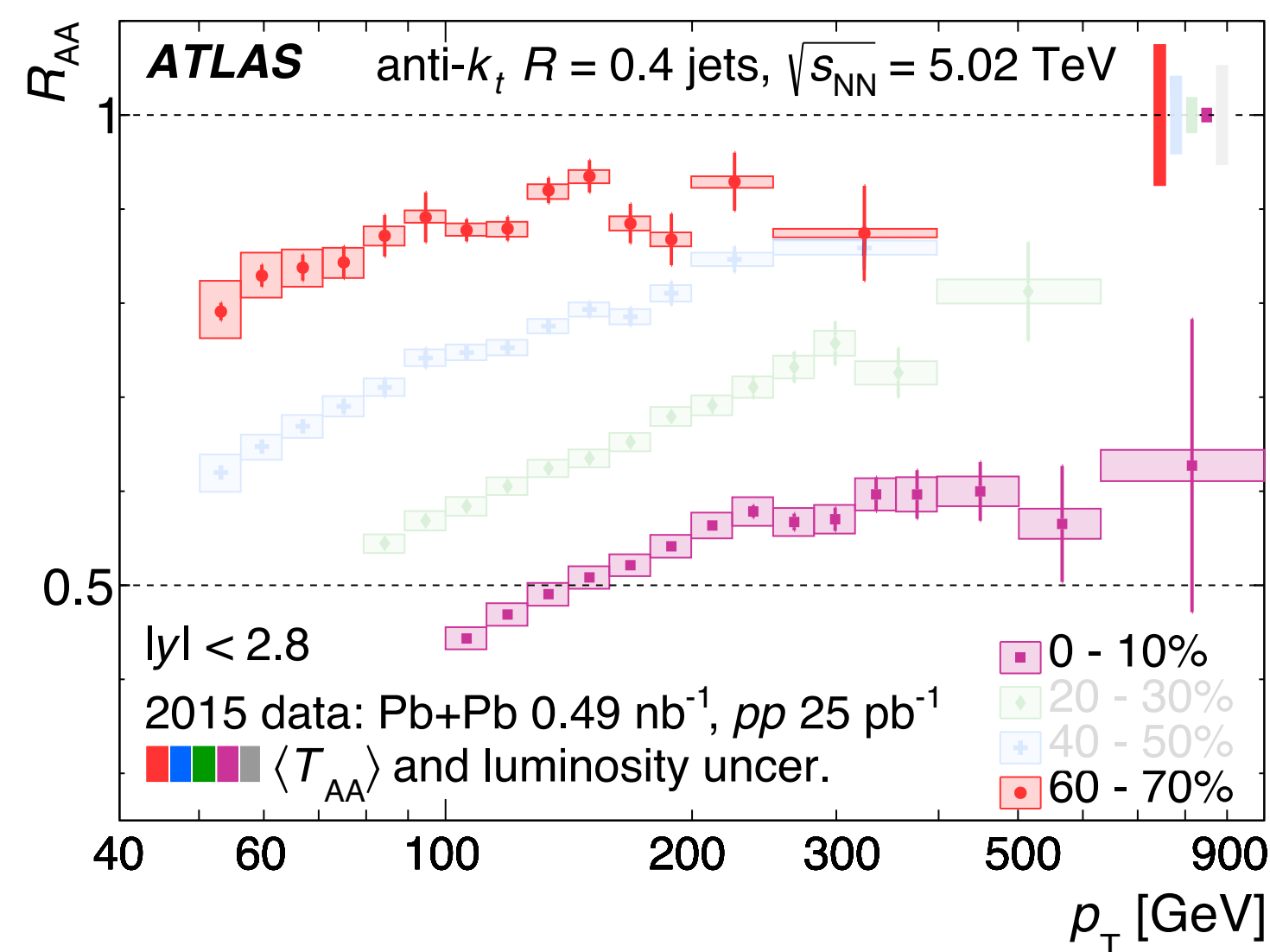




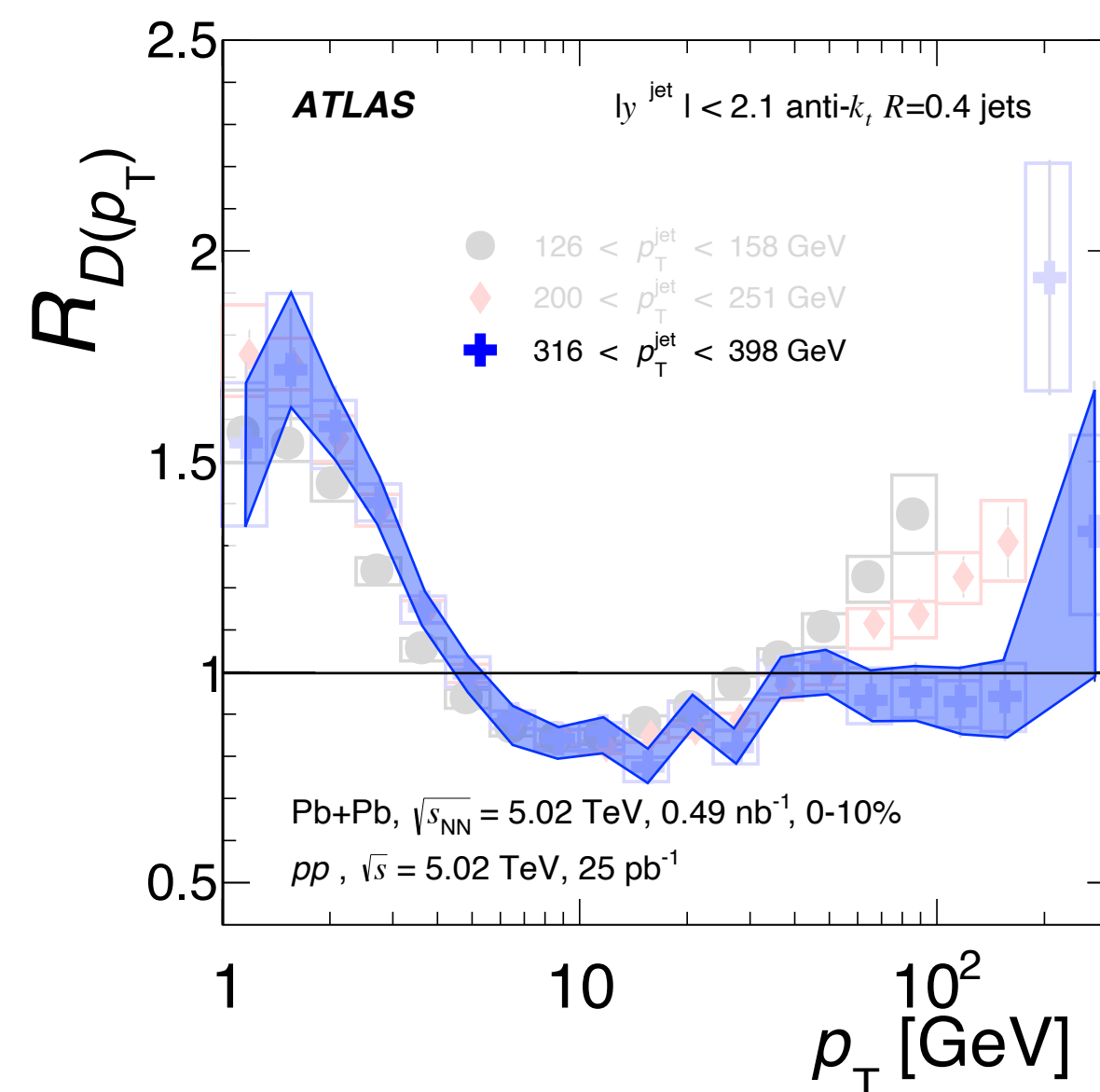
# What we learnt from inclusive jet quenching

PLB 790 (2019) 108  
 Phys. Rev. C 98 (2018) 024908  
 Phys. Rev. C 100 (2019) 064901

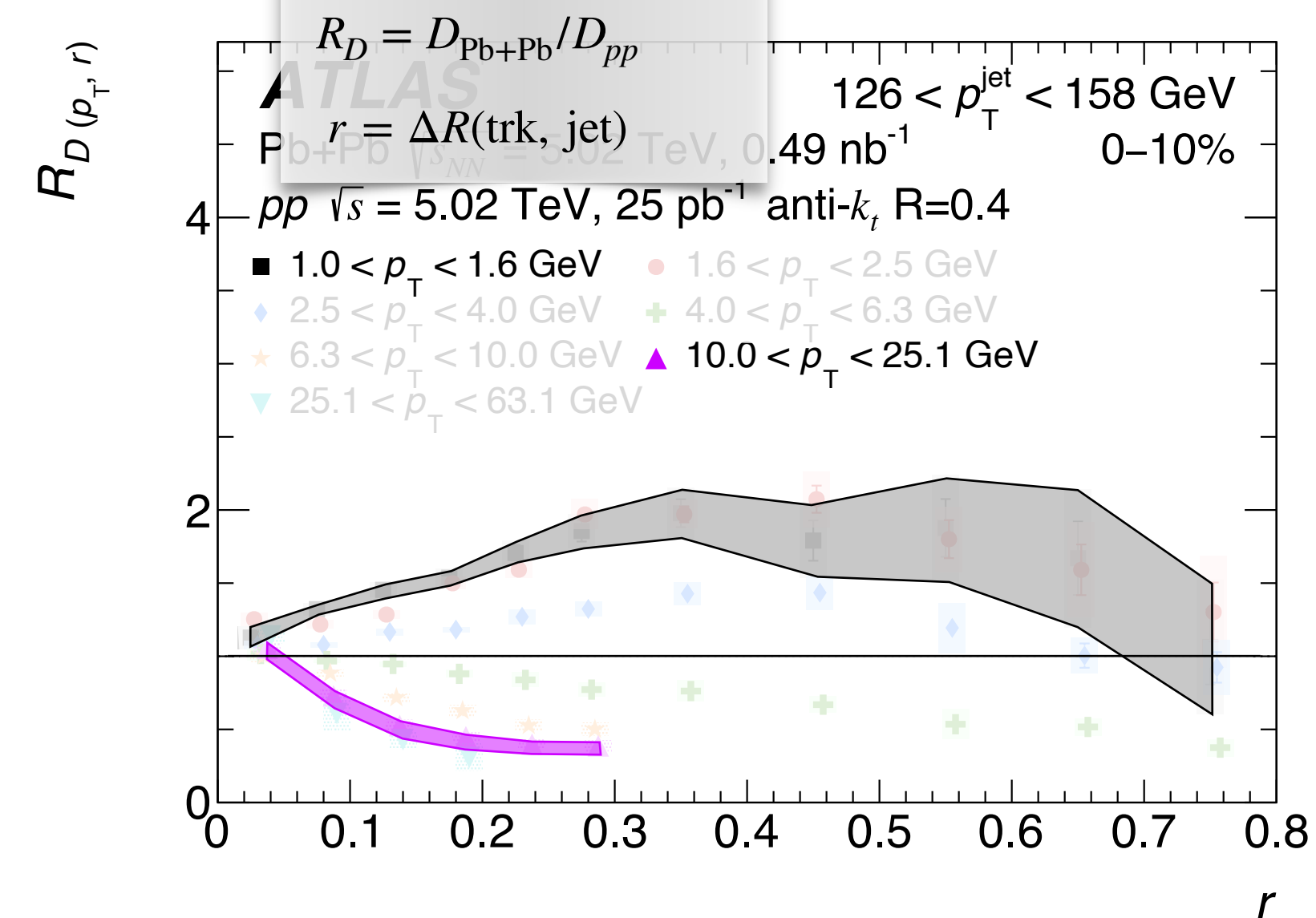
## Modifications in Spectrum



## Modifications in jet fragmentation



$$D(p_T, r) \propto \frac{1}{N_{\text{jet}}} \frac{d^2 n_{\text{trk}}}{dp_T^{\text{trk}} dr}$$

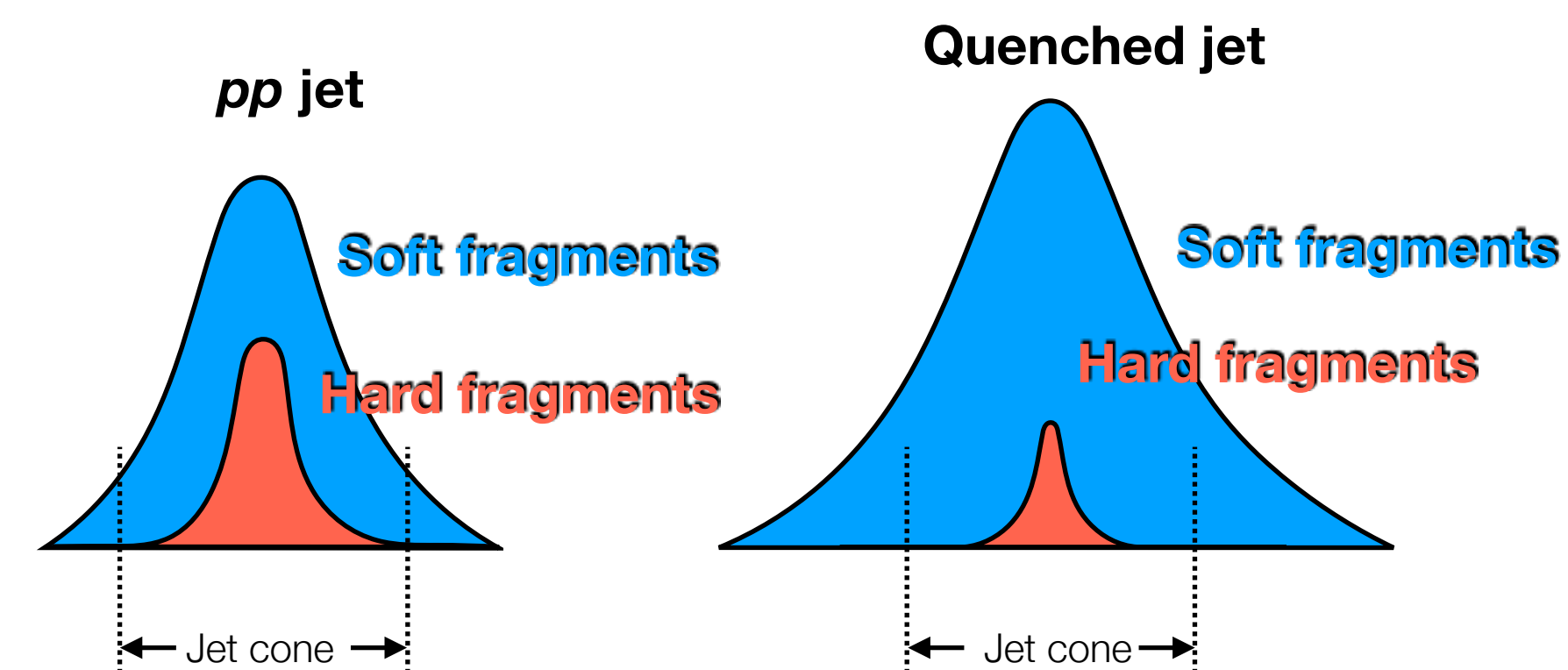


Jets are quenched in Pb+Pb, suppressed by a factor of 2 in central collisions

Quenched jet:

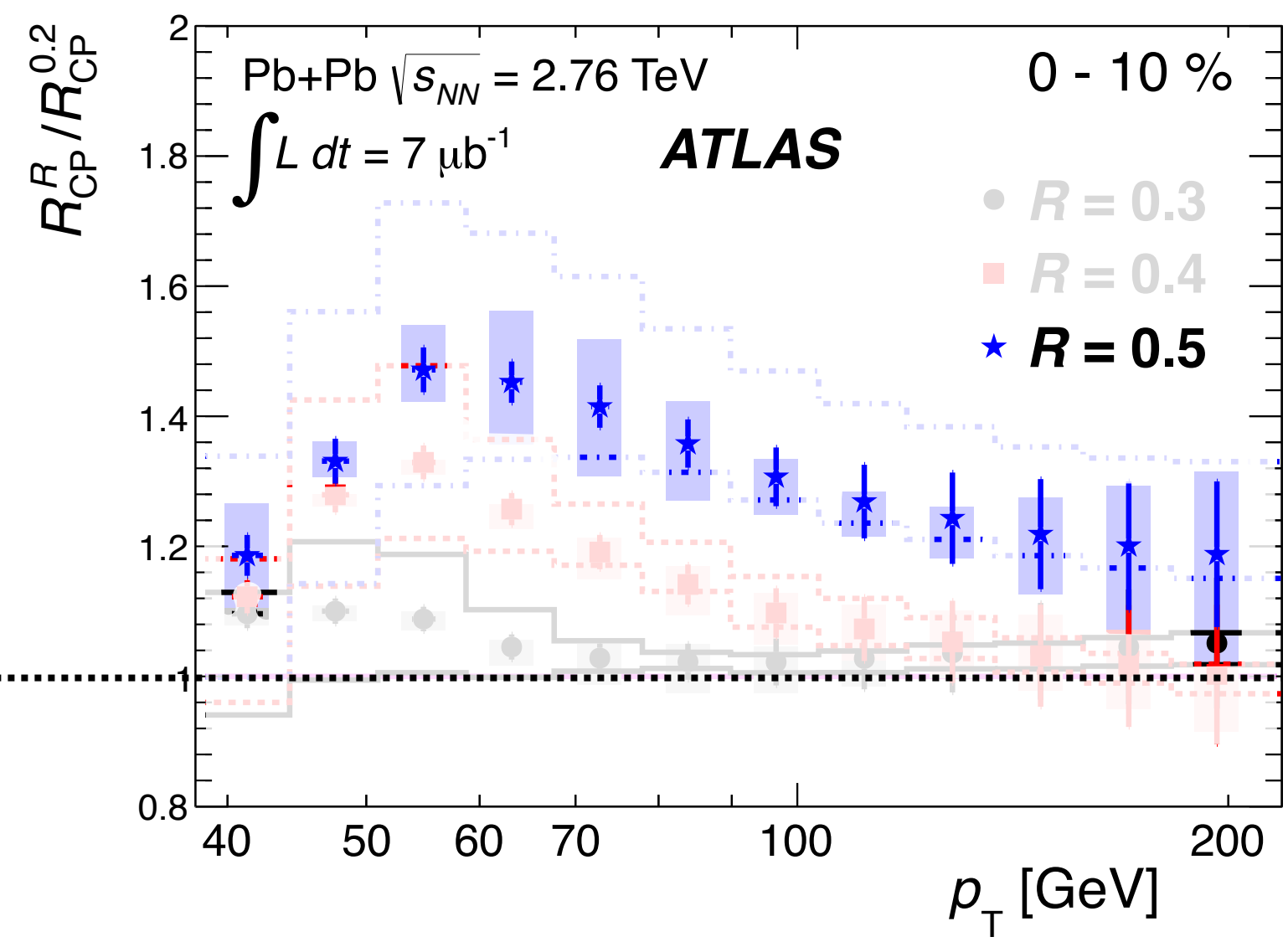
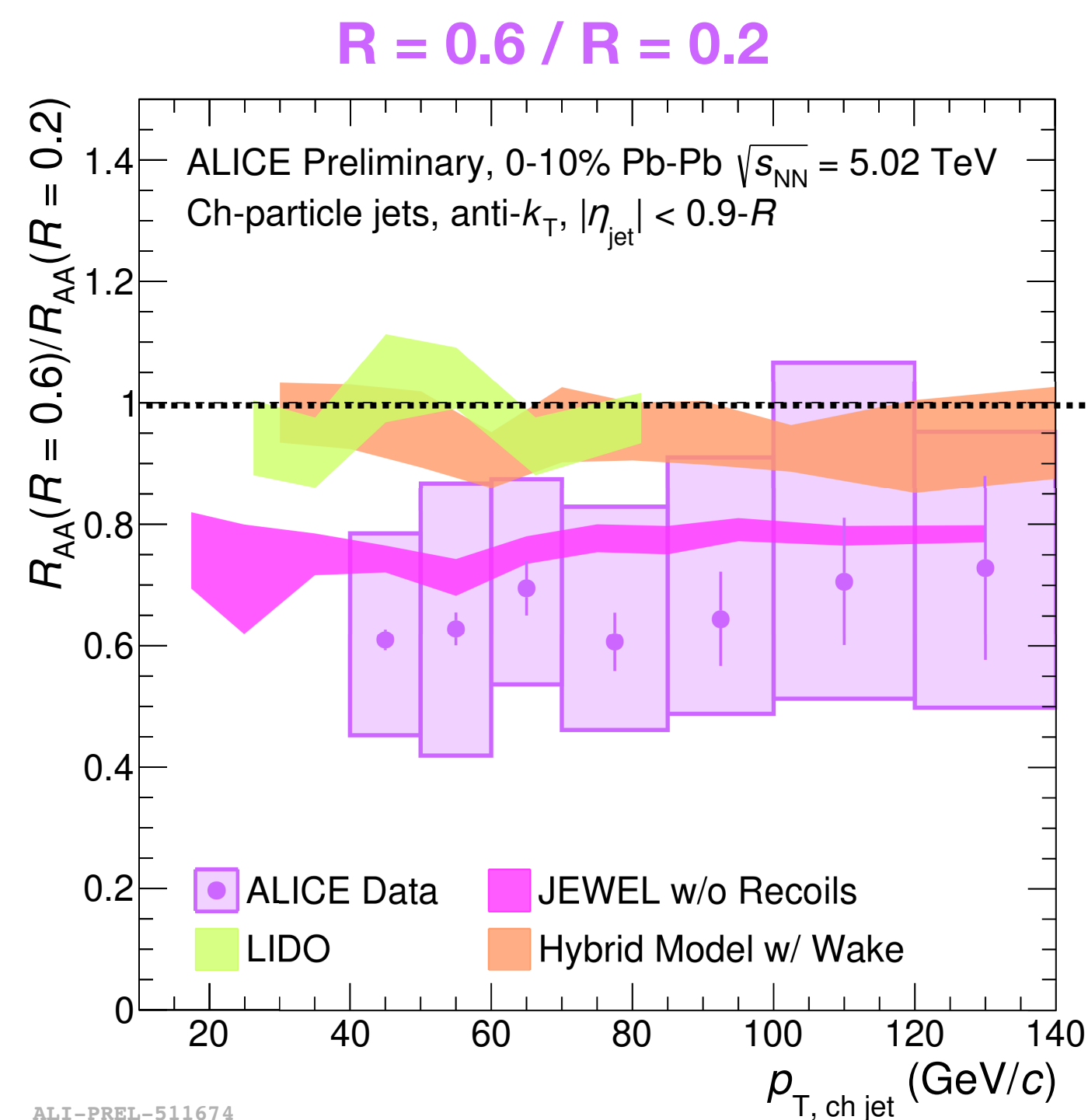
- Suppressed and narrowed **hard fragments**
- Enhanced and broadened **soft fragments**

Next generation jet quenching measurements to directly probe different scales:  $\Lambda_{\text{med}}$ ,  $m^2$ ,  $L_{\text{path}}$ , etc.



# R dependence of jet quenching

[Phys. Lett. B 719 \(2013\) 220-241](#)



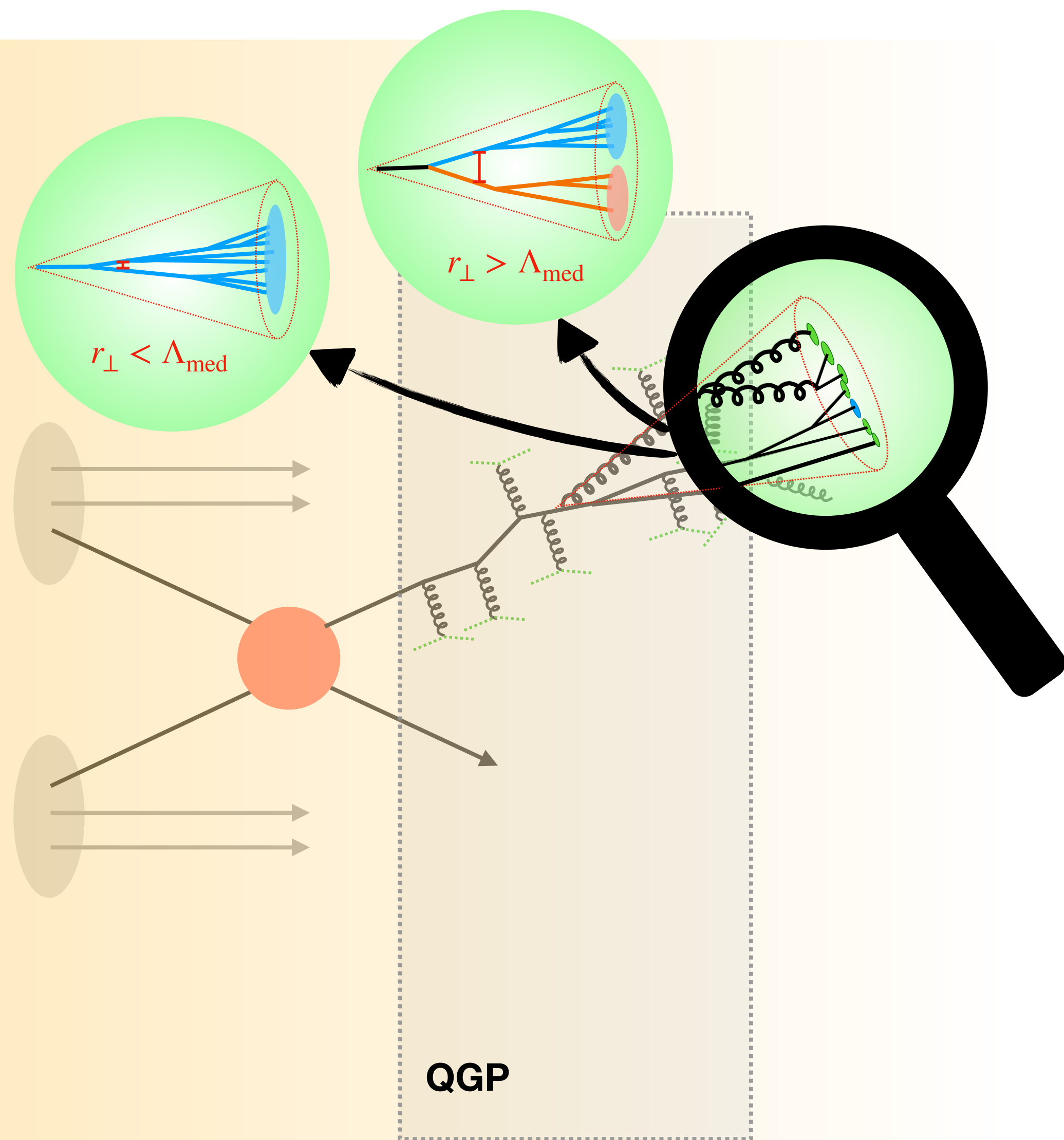
Different UE estimations: ML based (ALICE) vs. Area based (ATLAS)

Different Jets: track jet (ALICE) vs. calorimeter jet (ATLAS)

Larger jet radius:

- Wider area to recover lost energy
- Open phase space to jets with wide splittings

# Jet quenching — substructure dependence



## Jet quenching dictated by color coherence

Medium color coherence scale  $\Lambda_{med}$

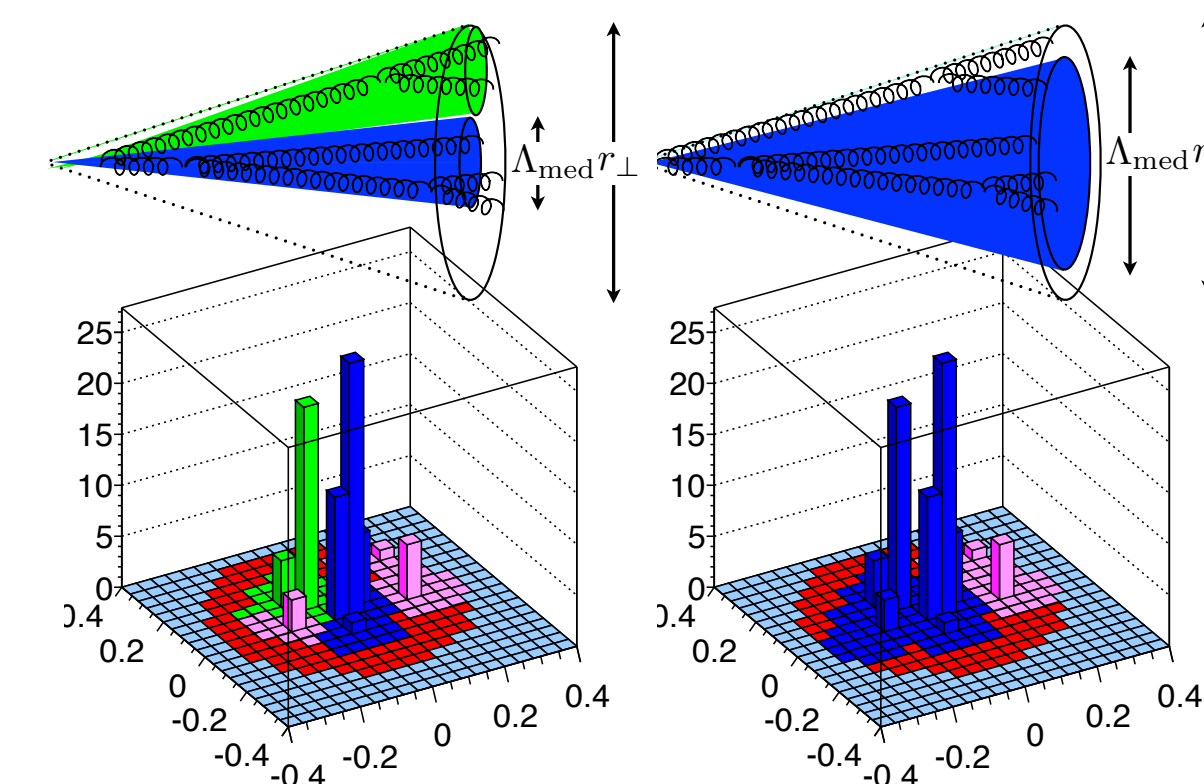
- Partons with separations below the scale  $\rightarrow$  radiate as a single parton

Medium color coherence  $\rightarrow$  Jet quenching depends on hard splitting of leading parton

Control jet substructure:

- Trimming of  $R=1.0$  jet
- Soft drop grooming of  $R=0.4$  jet

[Phys.Lett.B 725 \(2013\) 357-360](#)





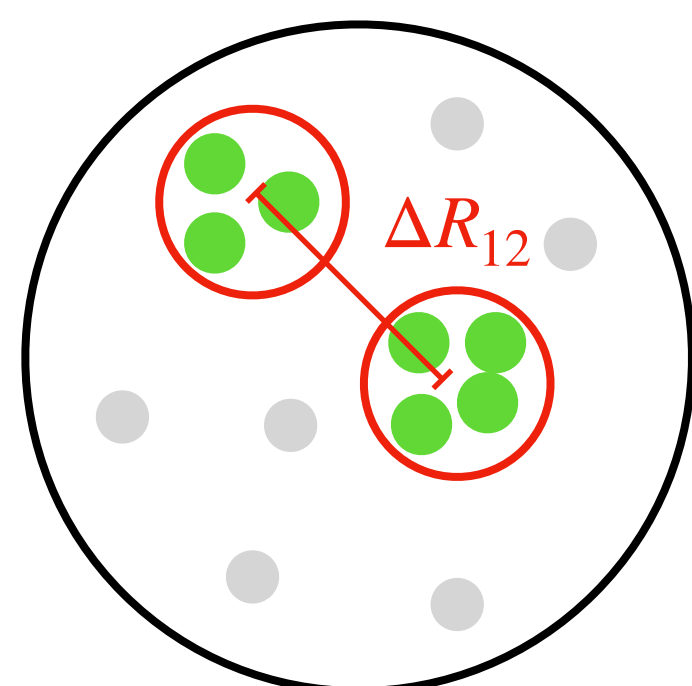
# $R_{AA}$ of trimmed large- $R$ jet

ATLAS-CONF-2019-056

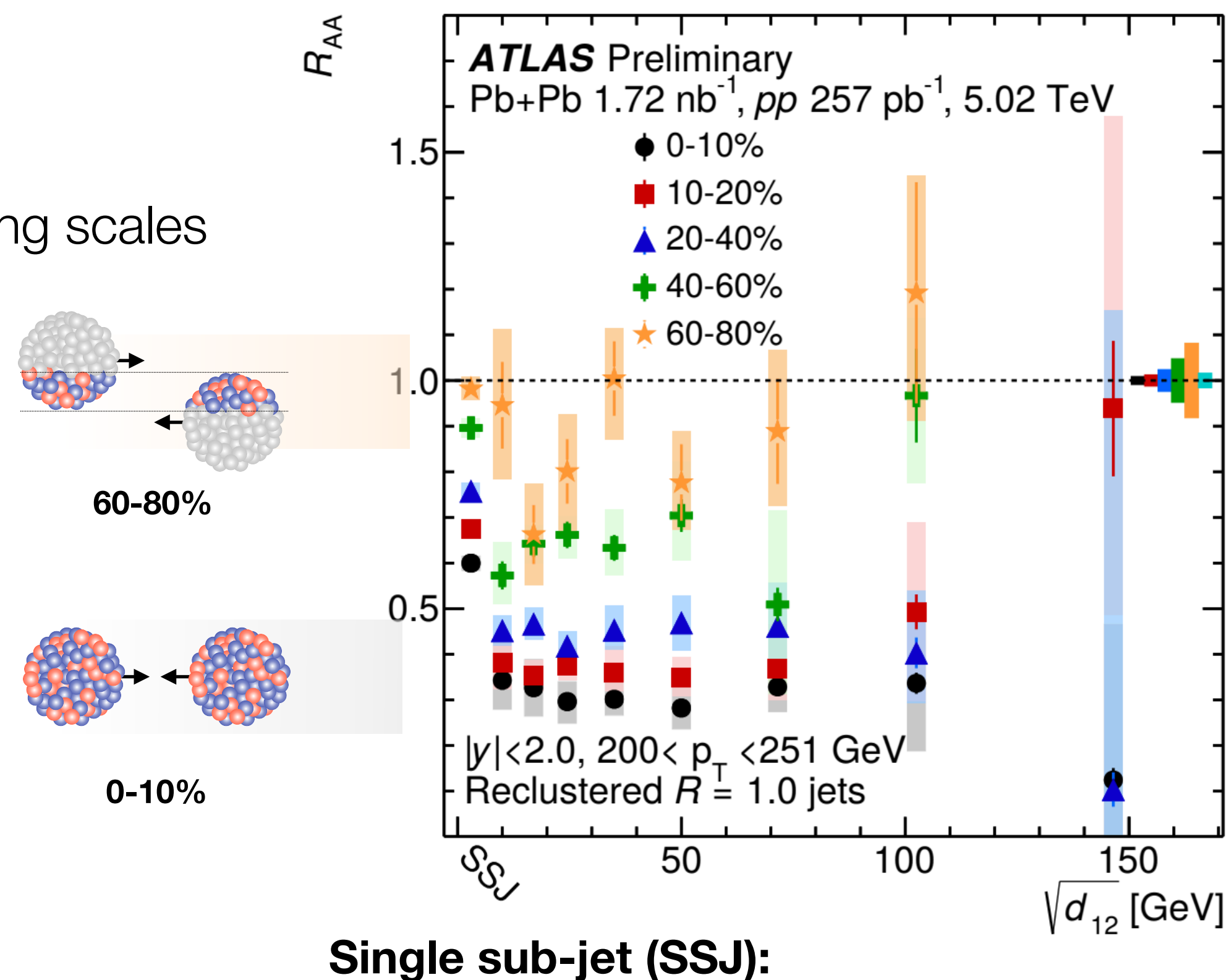
Trimming procedure

1. Cluster Anti- $k_T$   $R=0.2$  jets to form Anti- $k_T$   $R=1.0$  jet
2. Recluster  $R=0.2$  jets in  $R=1.0$  with algorithm to get  $k_T$  splitting scales

$$\sqrt{d_{12}} = \min(p_{T,1}, p_{T,2}) \times \Delta R_{12}$$



UE contribution suppressed by using  $R=0.2$  jet as constituents



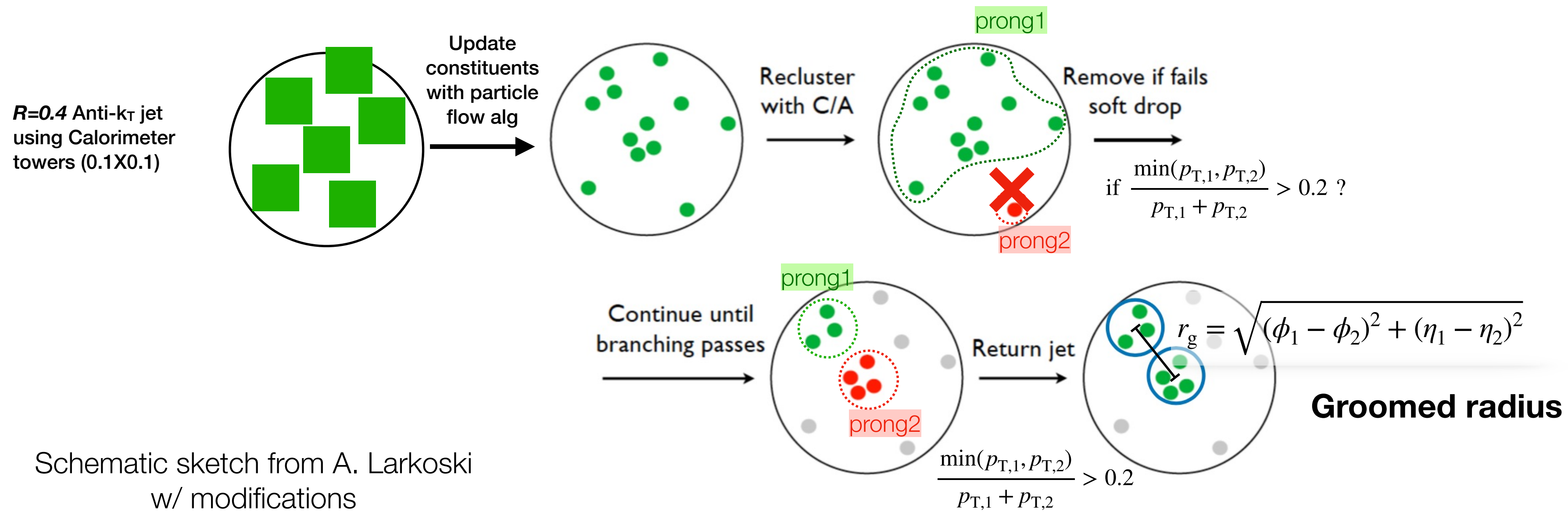
Two distinguished classes of jets:

- $\sqrt{d_{12}} > 0$ , more suppressed, weakly depend on  $\sqrt{d_{12}}$  value,  $r_{\perp} > \Lambda_{\text{med}}$
- $\sqrt{d_{12}} = 0$ , significant less suppressed,  $r_{\perp} > \Lambda_{\text{med}} + r_{\perp} < \Lambda_{\text{med}}$

# Jet grooming procedure

To further access  $\Delta R_{12} < 0.2$ , jet grooming is used:

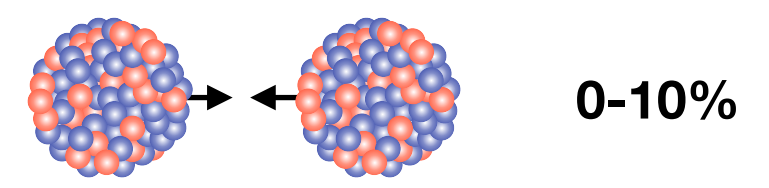
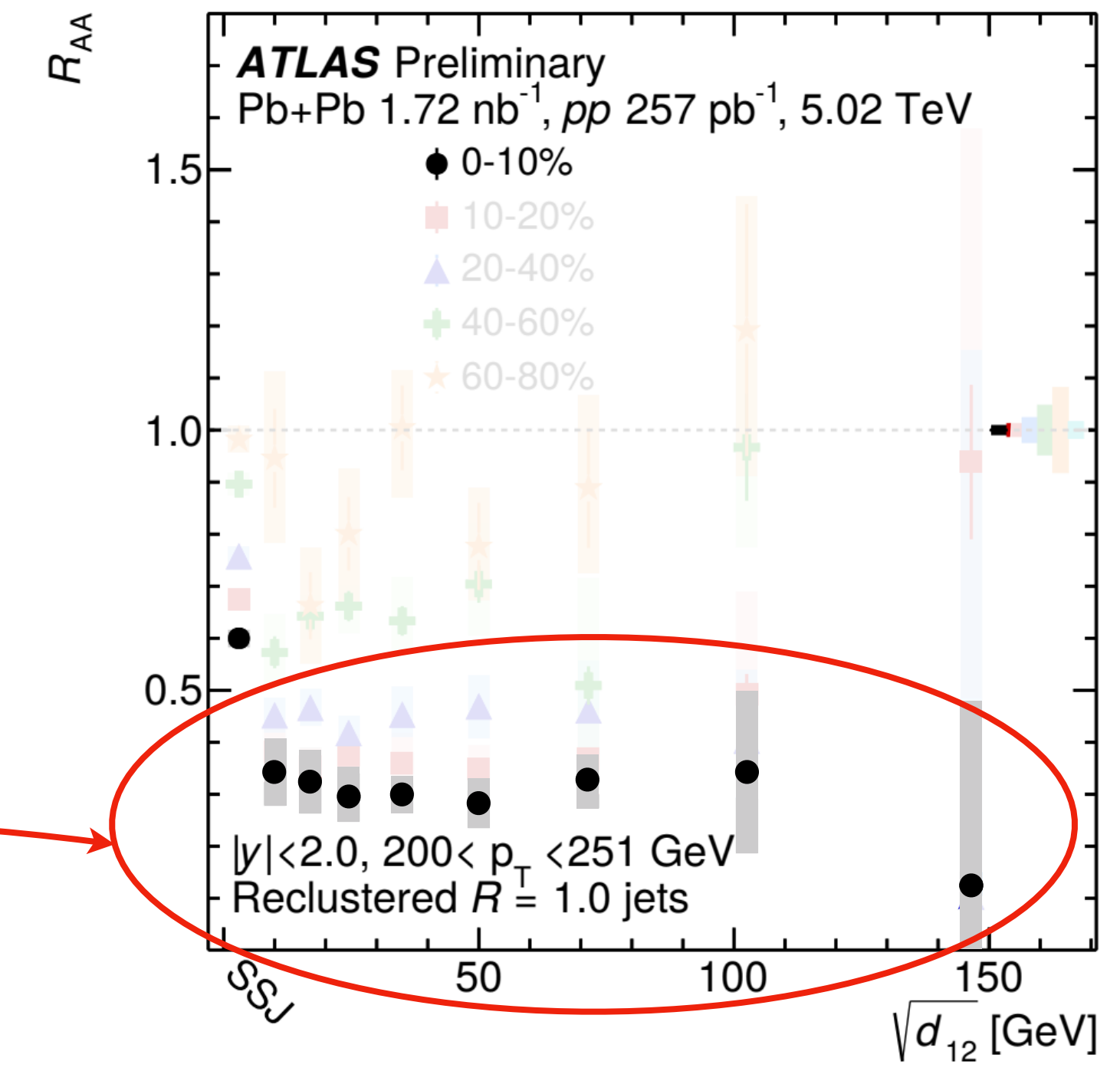
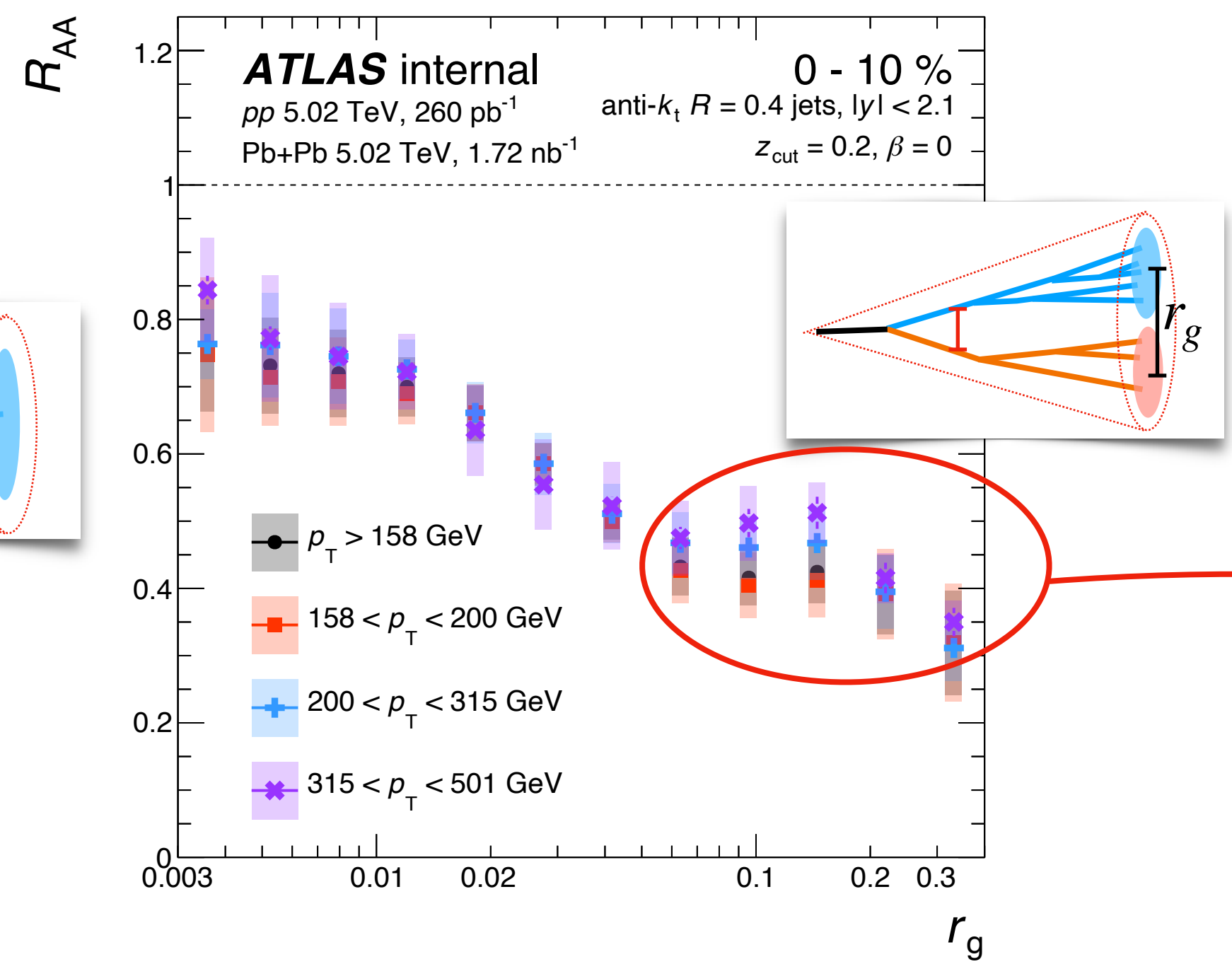
1. **Jet finding:**  $R=0.4$  Anti- $k_T$  jet
2. **Reclustering:** C/A cluster sequence
3. **Declustering:** Soft drop grooming with  $z_{\text{cut}} = 0.2$ ,  $\beta = 0$



Schematic sketch from A. Larkoski w/ modifications

# $R_{AA}$ of groomed jets

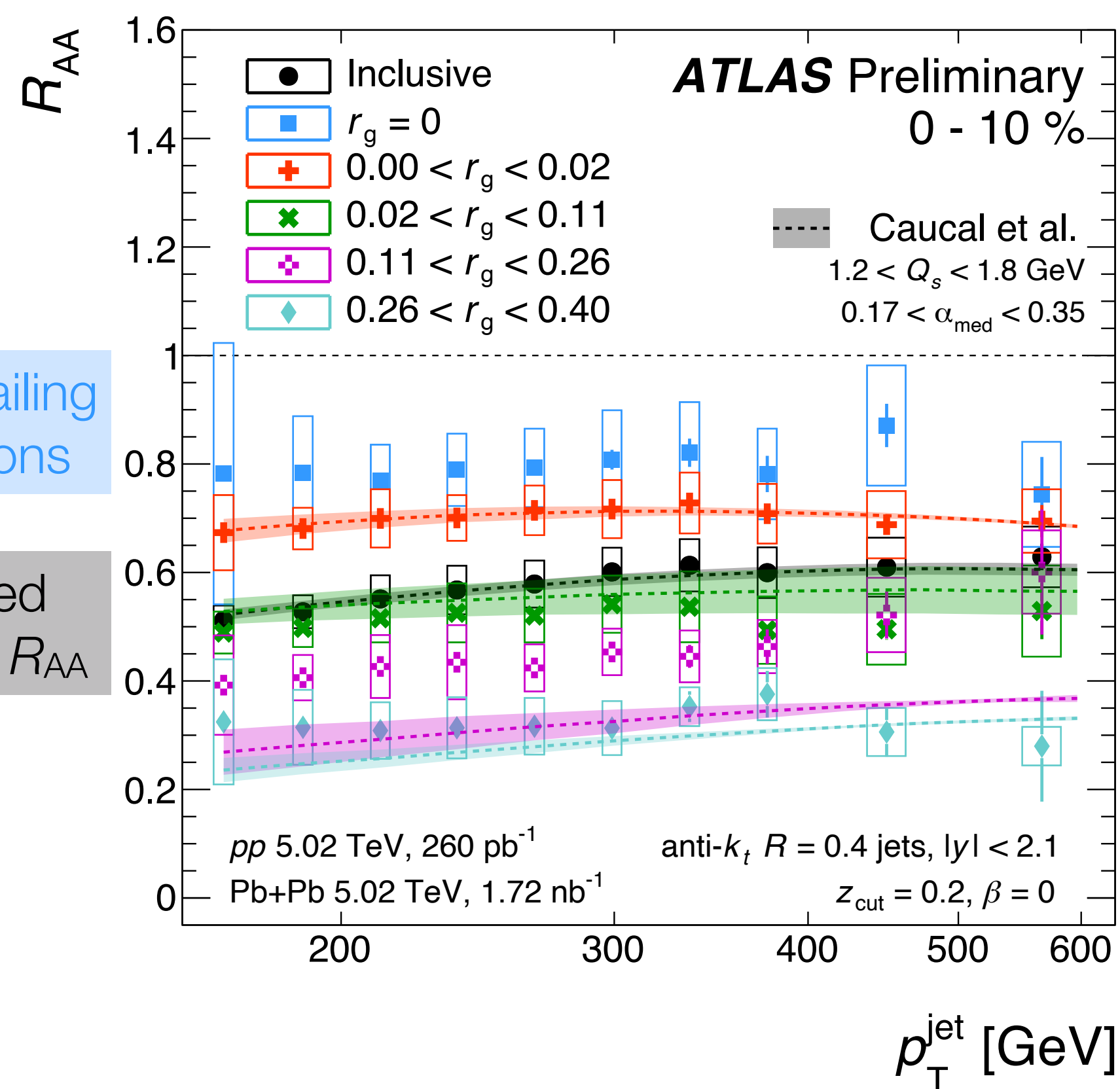
ATLAS-CONF-2022-026





# $R_{AA}$ of groomed jets — model comparison

ATLAS-CONF-2022-026



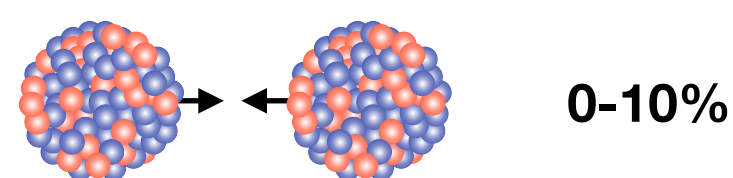
$r_g = 0$ : jets failing SD conditions

Ungroomed inclusive jet  $R_{AA}$

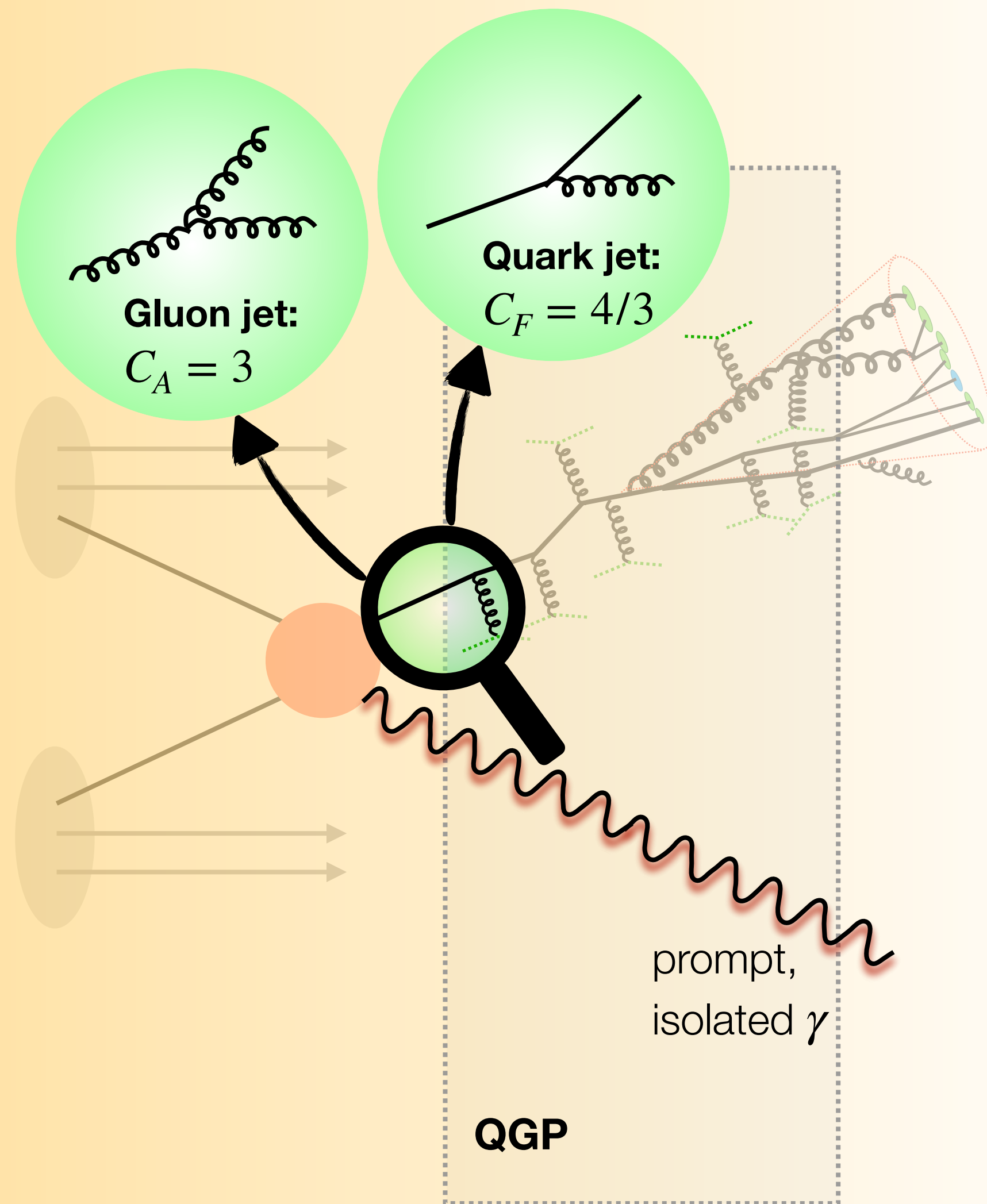
Groomed jet  $R_{AA}$   
 $0.00 < r_g < 0.02$

Groomed jet  $R_{AA}$   
 $0.26 < r_g < 0.4$

- Caucal et al.\*, MC generator with  $p$ QCD in-medium parton shower
- Very good agreement at small and large  $r_g$
- Discrepancies in  $0.11 < r_g < 0.26$

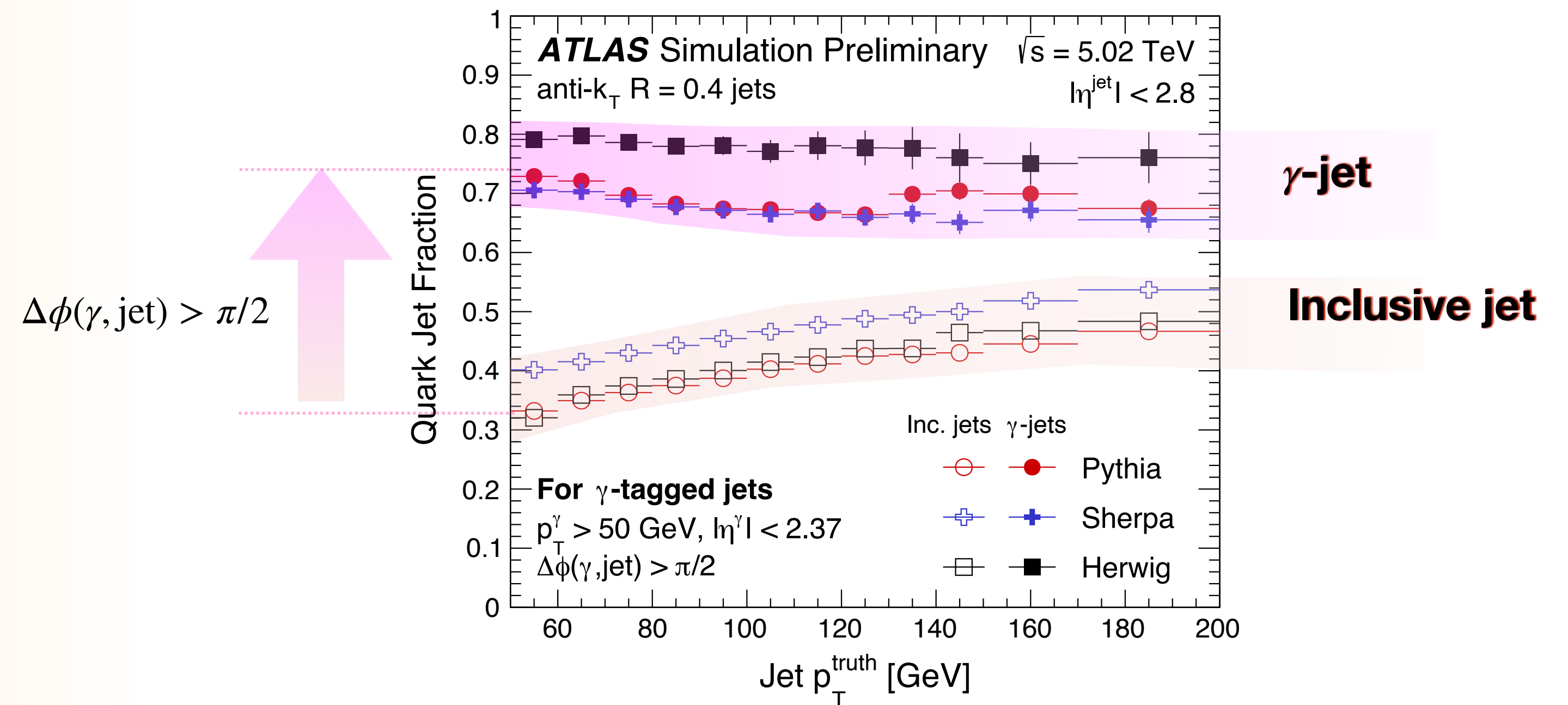


# Jet quenching — color charge dependence

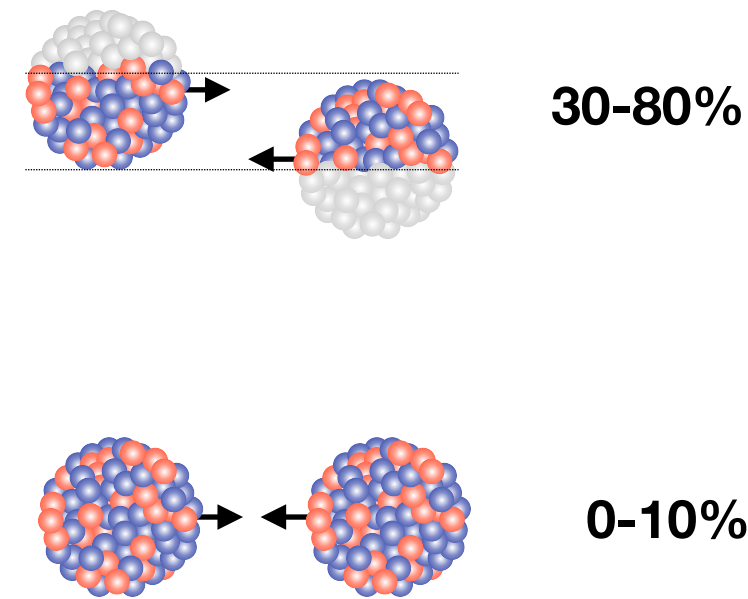
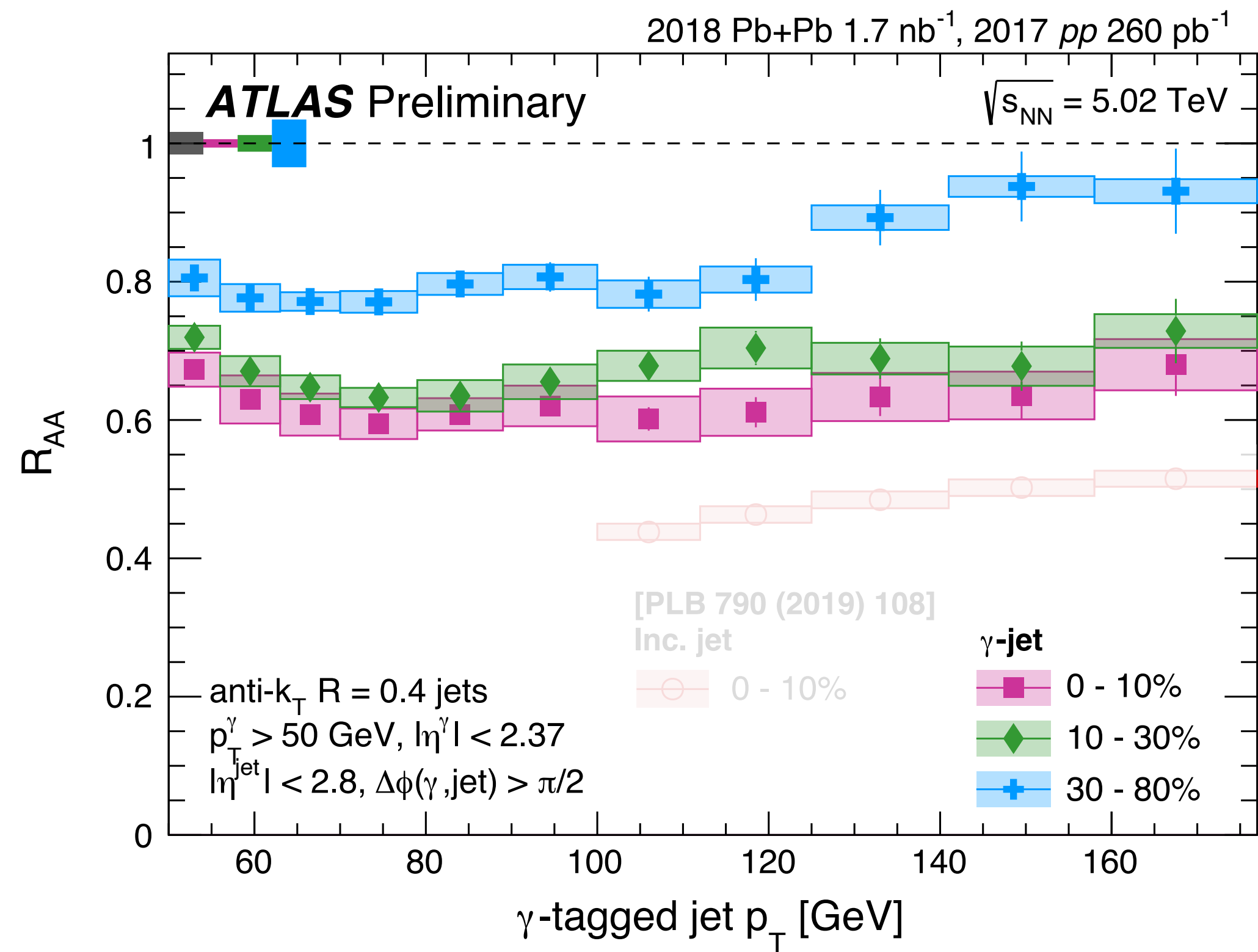


Gluon jets vs. quark jets:

- Gluon jets are more active: fatter and produce more particles
- Jet quenching depends on color charge



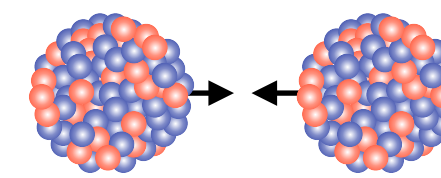
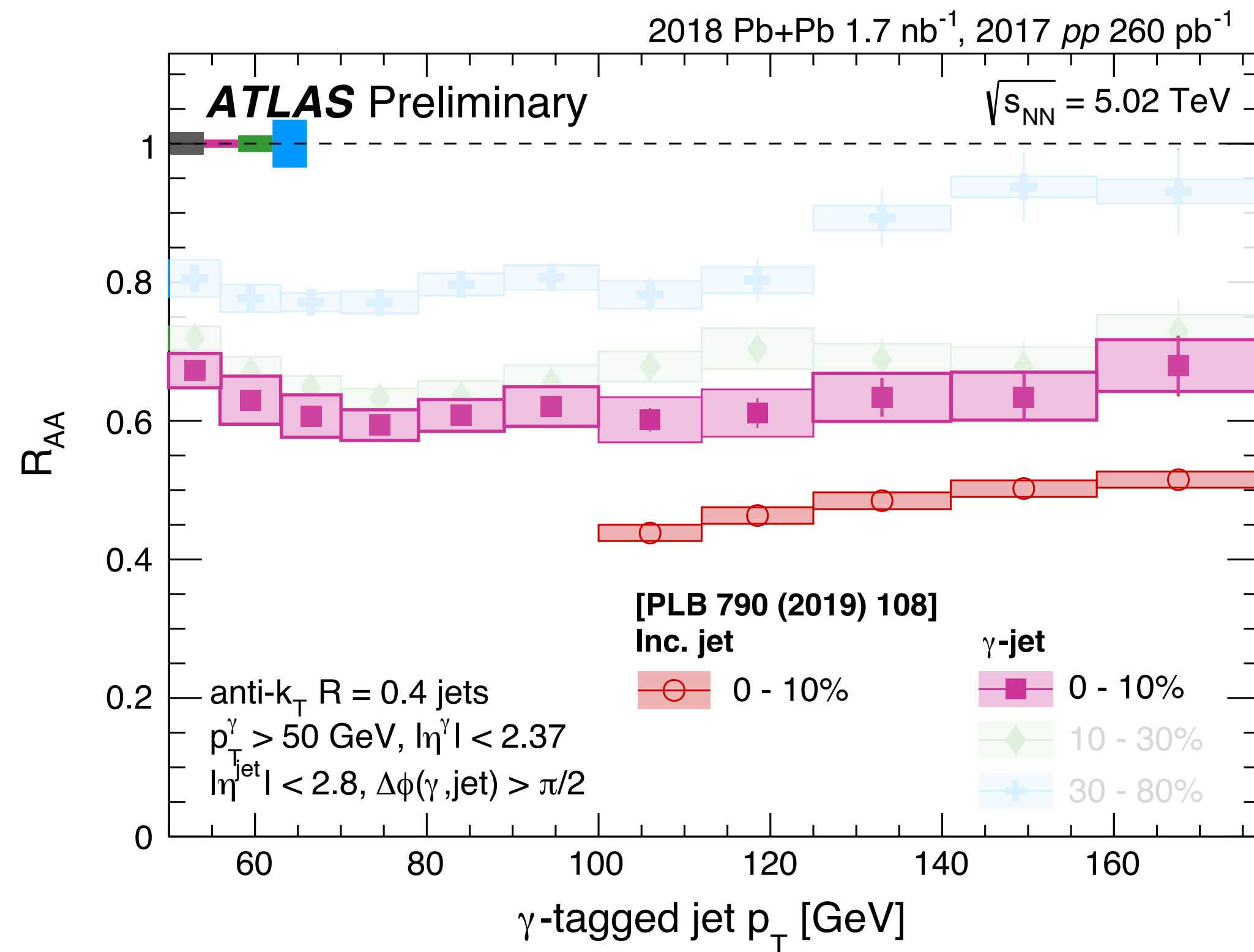
# $\gamma$ -tagged jet $R_{AA}$



- Strong centrality dependence
- Weak dependence on jet  $p_T$  in central collisions



# $\gamma$ -tagged jets vs. inclusive jets



0-10%

$\gamma$ -tagged jet

~80% quark jets

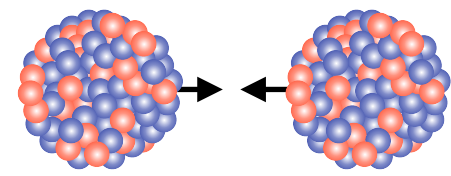
Inclusive jet

~40% quark jets

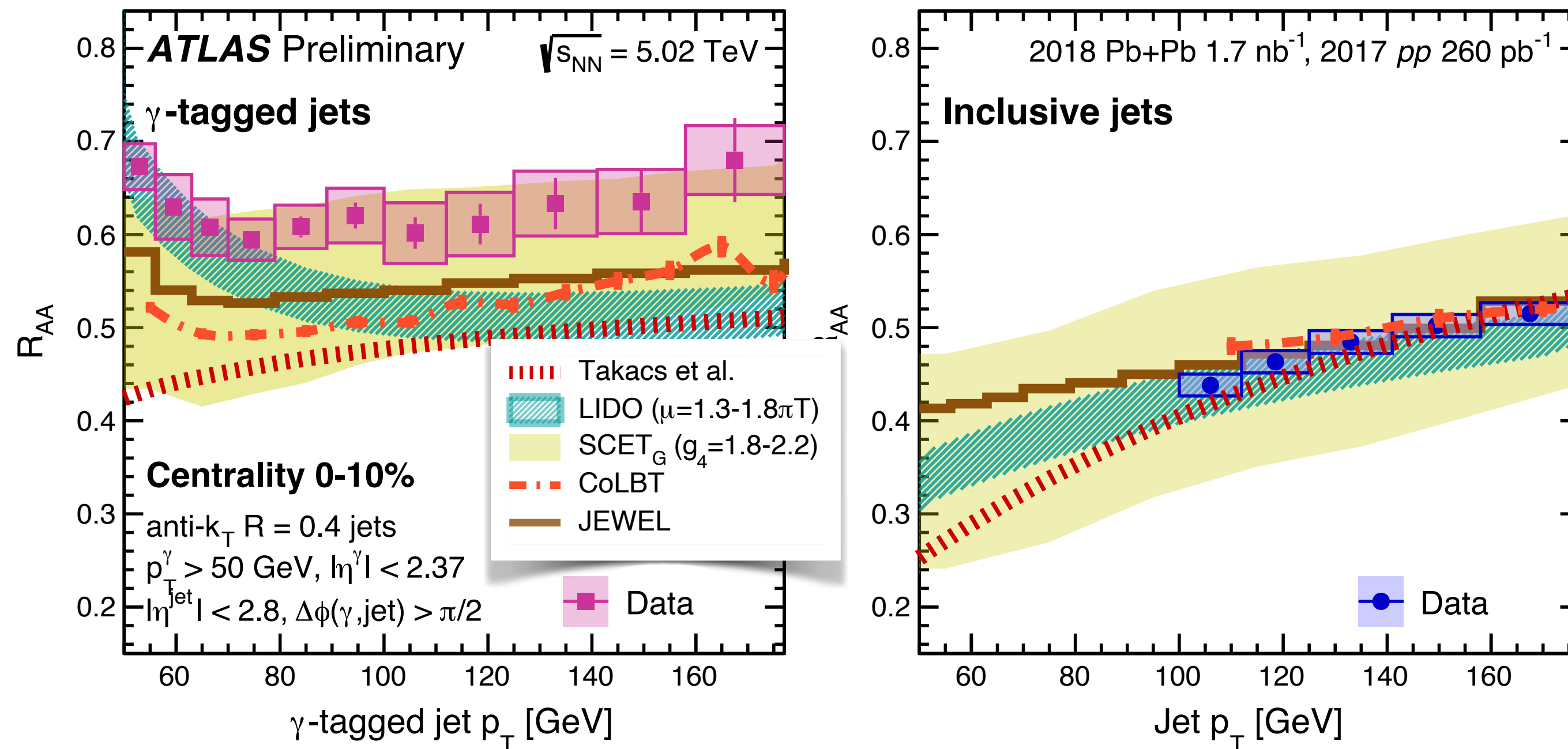
$\gamma$ -tagged jets are less suppressed than inclusive jets

- Different virtuality / substructures
- Different jet spectra

# $\gamma$ -tagged jet $R_{AA}$ — model comparisons



0-10%

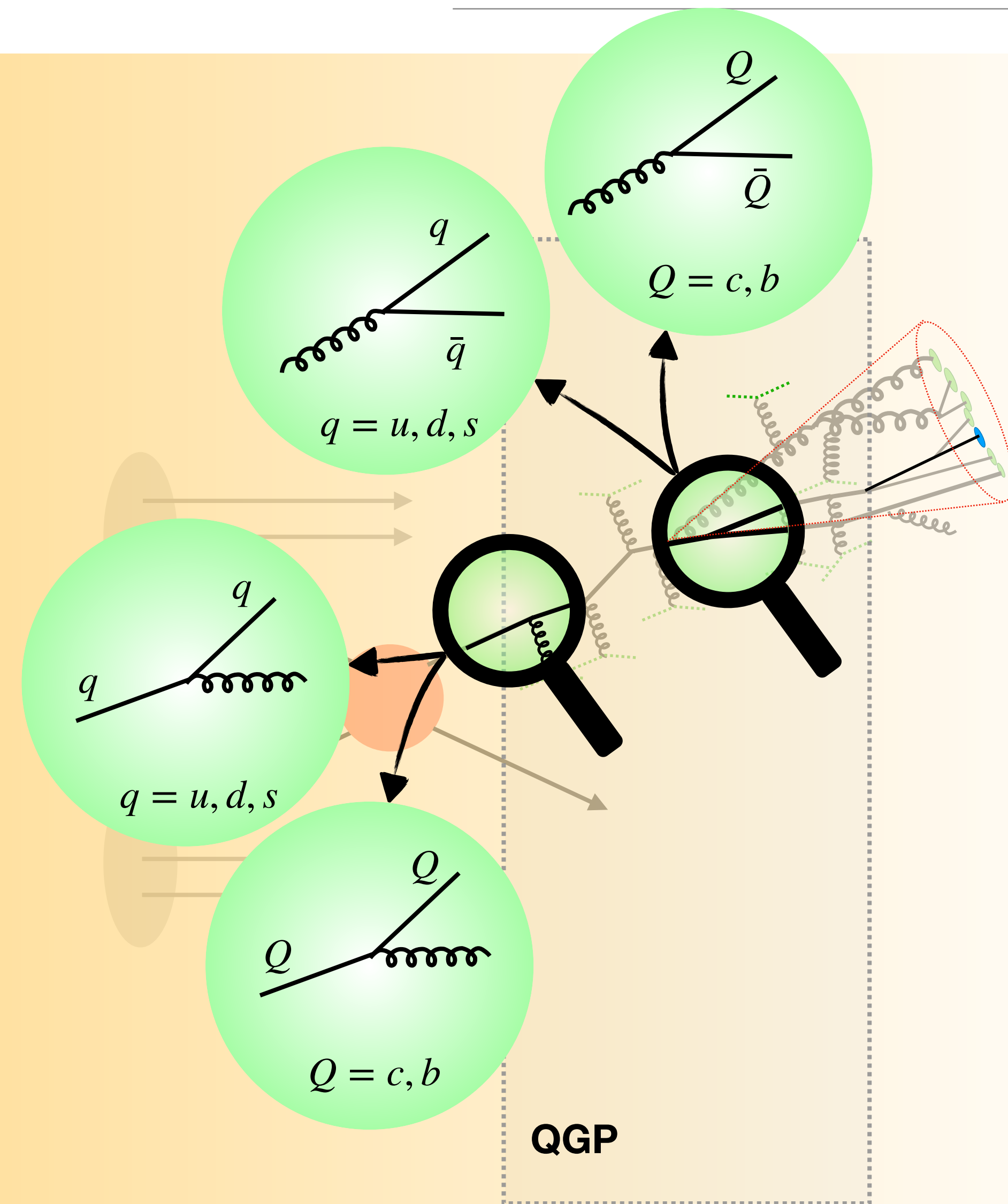


Models	Details	Uncertainty
LIDO*	Pythia8 + Boltzmann-Langevin transport	Medium-jet coupling
CoLBT*	Pythia8 + Higher-twist + Boltzmann transport	N/A
SCETG*	EFTs w/ medium modified splitting + Evolution equations	Medium-jet coupling
Takacs et al.*	Pythia8 + BDMPS-Z + GLV + Poisson ansatz	N/A
JEWEL*	Pythia6.4 + BDMPS + parton shower	N/A

Models describe inclusive jet  $R_{AA}$  well while under-predict  $\gamma$ -tagged jet at high  $p_T$

# Jet quenching — parton mass dependence

[arXiv:2204.13530](https://arxiv.org/abs/2204.13530)



**Heavy flavor jet:** Particle jet contains **heavy flavor** hadrons

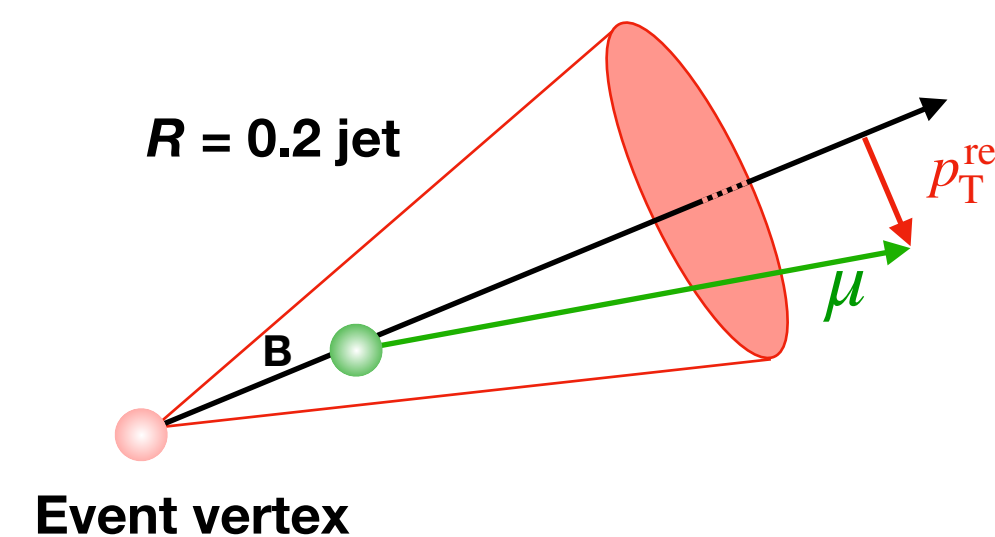
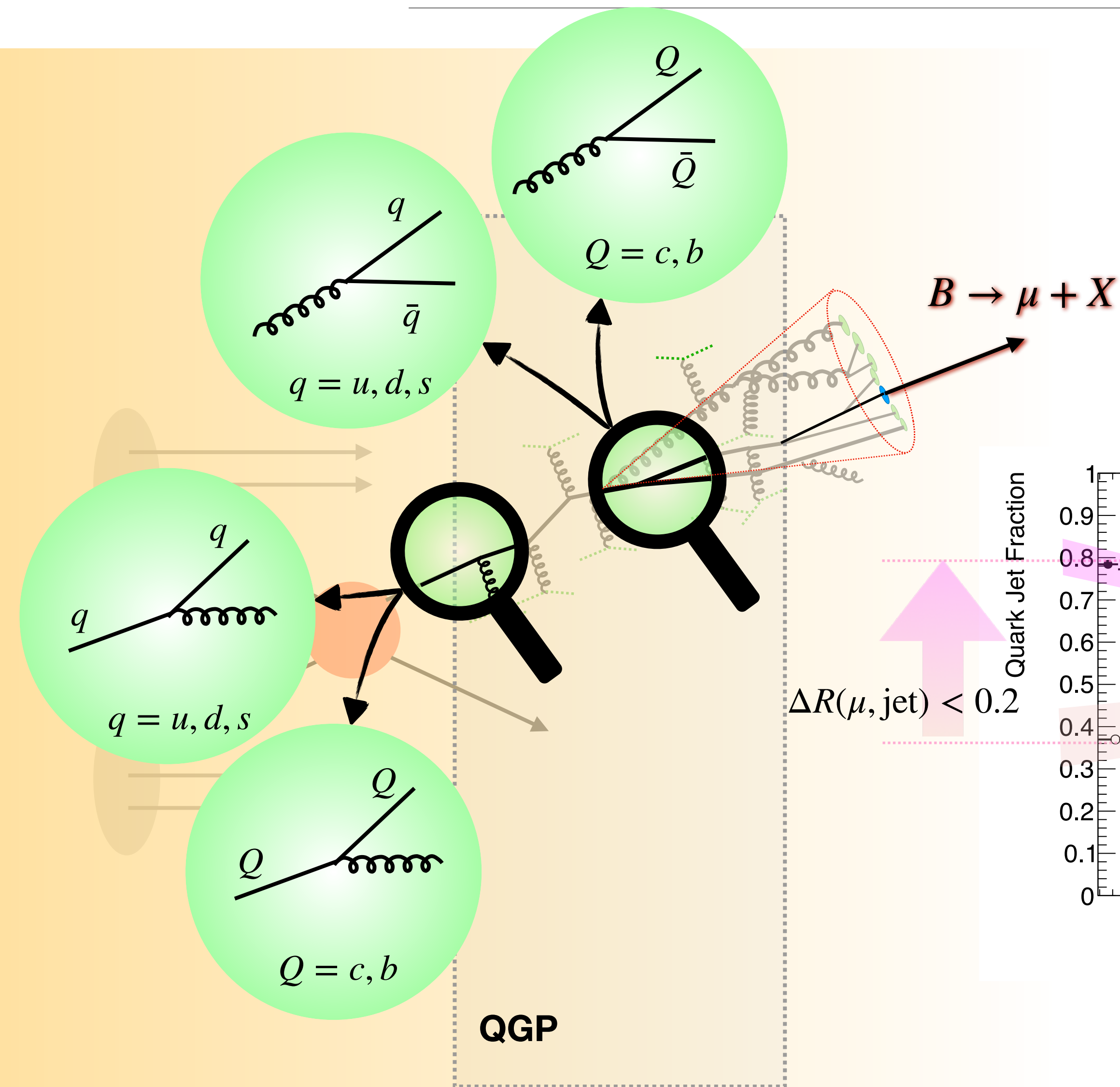
Heavy flavor quark created earlier in the heavy ion collisions, introduce variations on **parton mass** — dead cone effect

$$dP_{\text{Rad};Q}(\theta) \propto \left( 1 + \left( \frac{M_Q}{E_Q} \right)^2 \frac{1}{\theta^2} \right)^{-2}$$

Heavy flavor quark selection also introduces variations on initial parton **color charge**

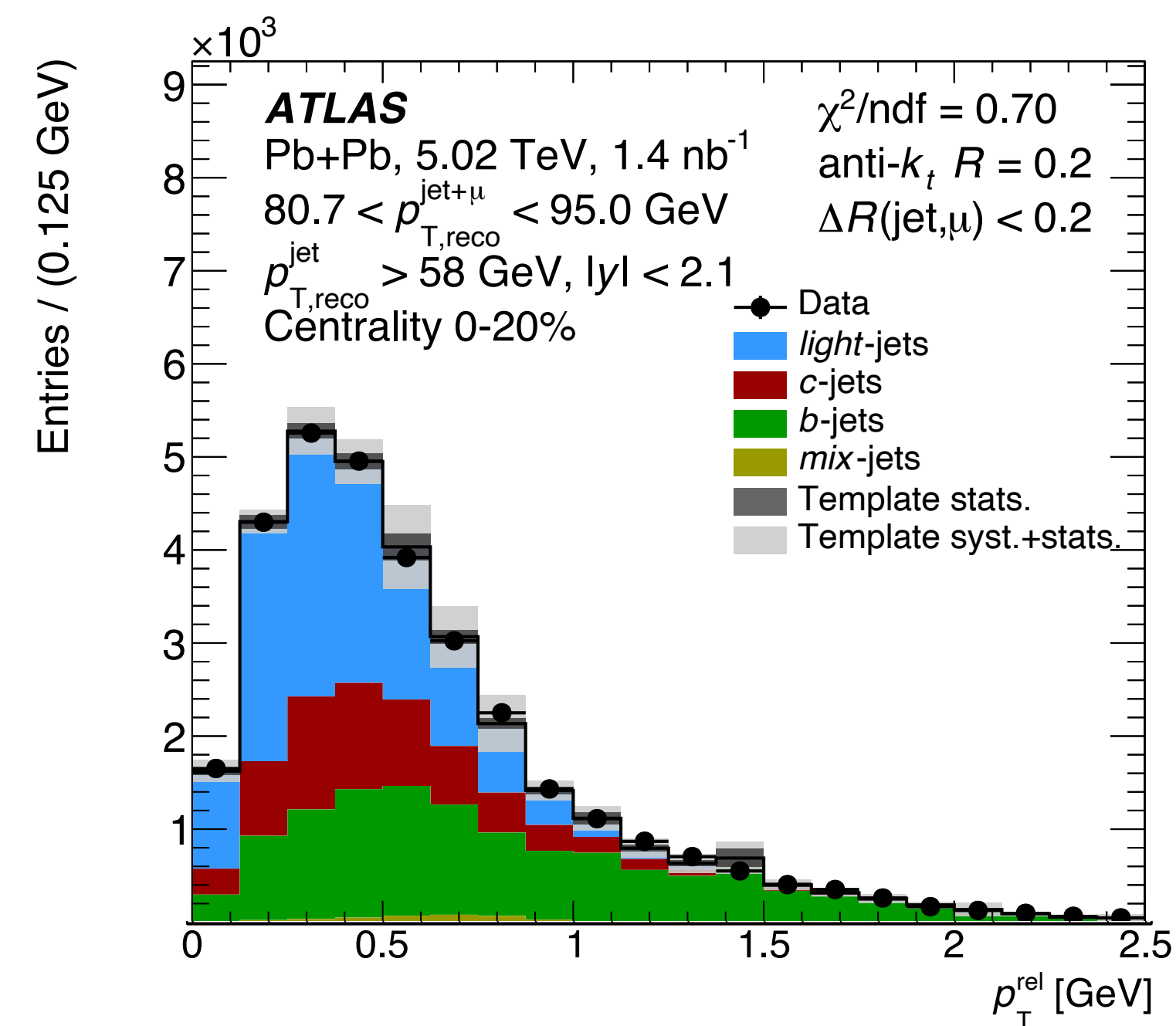
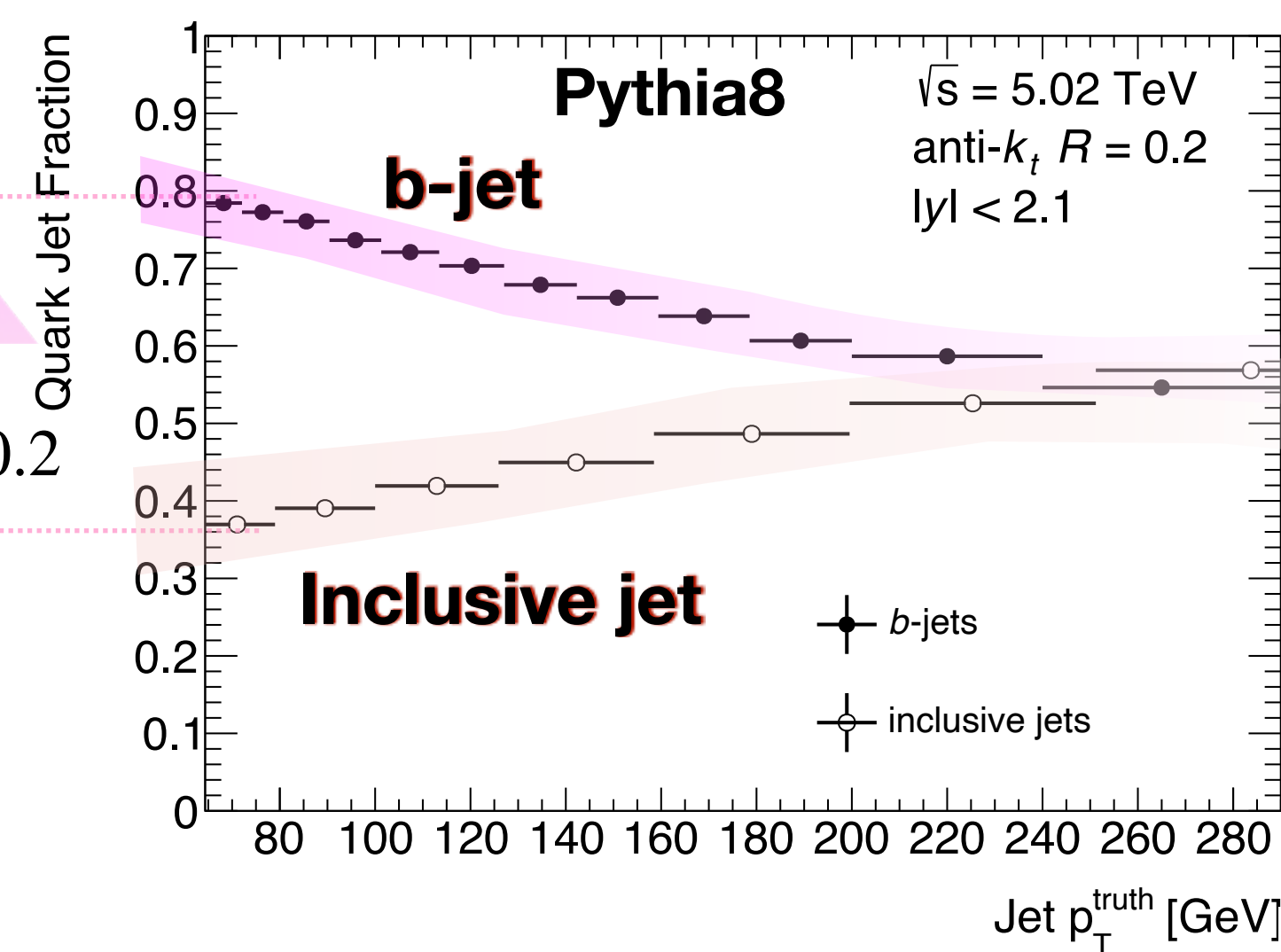


# Jet quenching — parton mass dependence



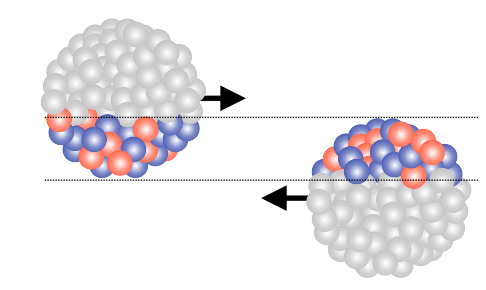
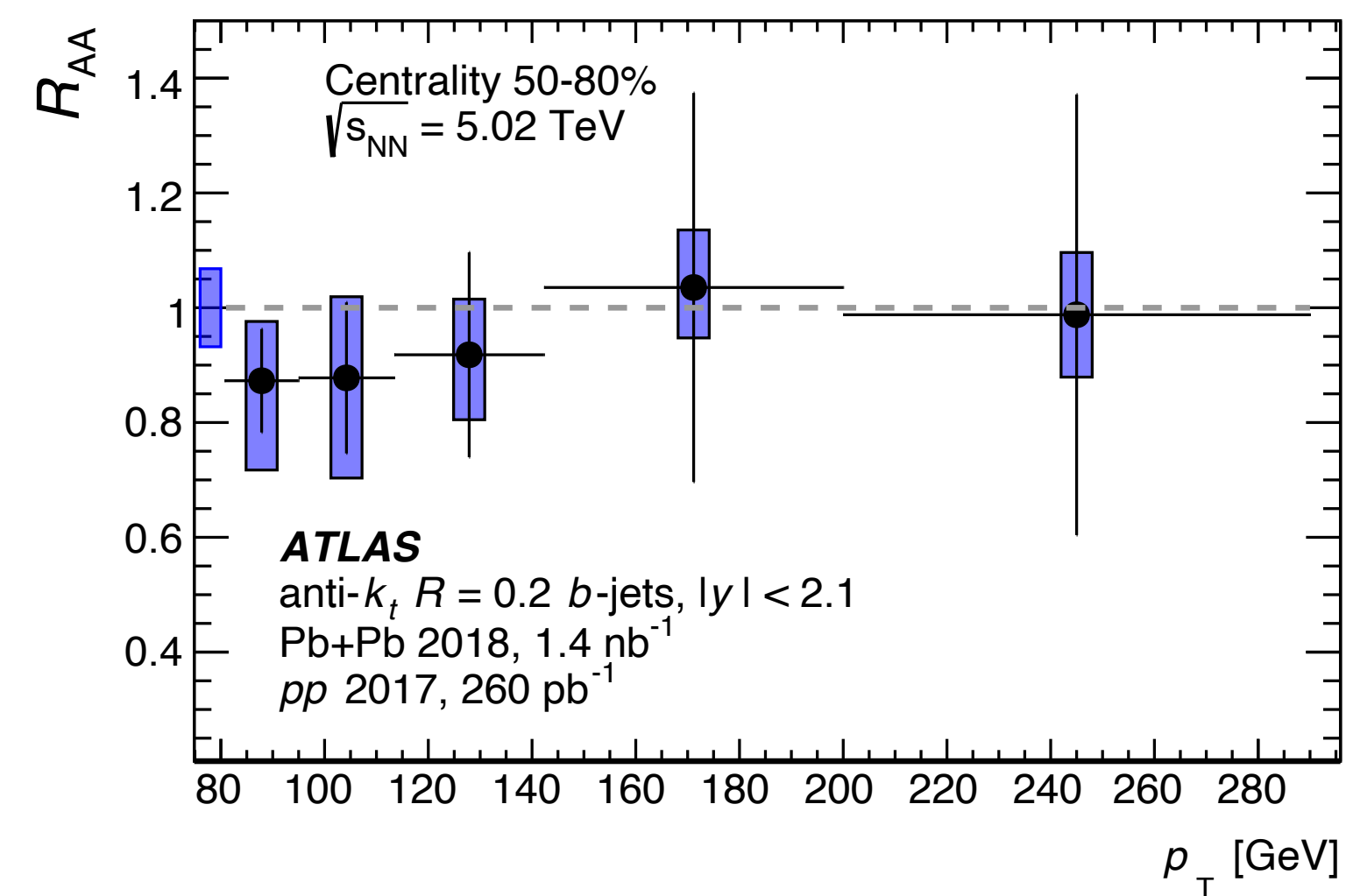
**Muon in jets:**  $\Delta R(\mu, \text{jet}) < 0.2$

**b-jet yields** are measured from fitting  $p_T^{\text{rel}}$



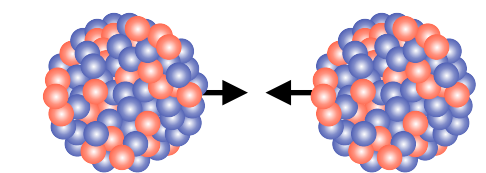
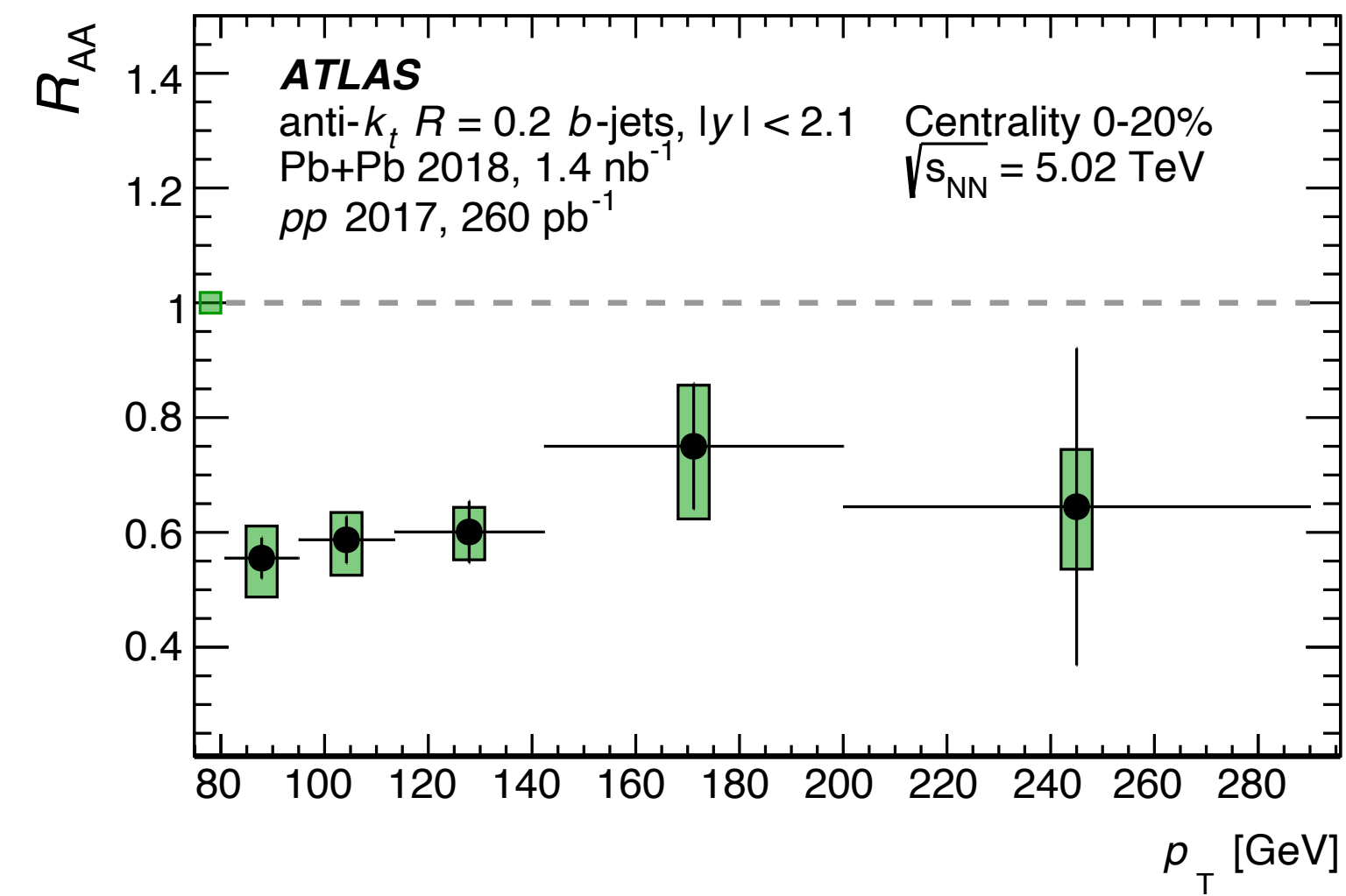
# $b$ -jet $R_{AA}$

[arXiv:2204.13530](https://arxiv.org/abs/2204.13530)



50-80%

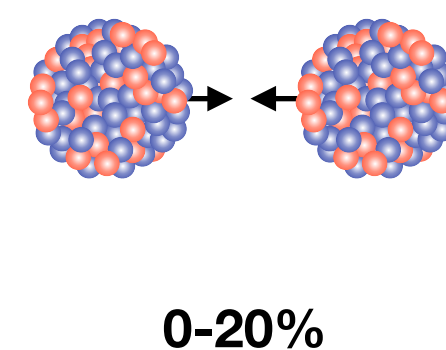
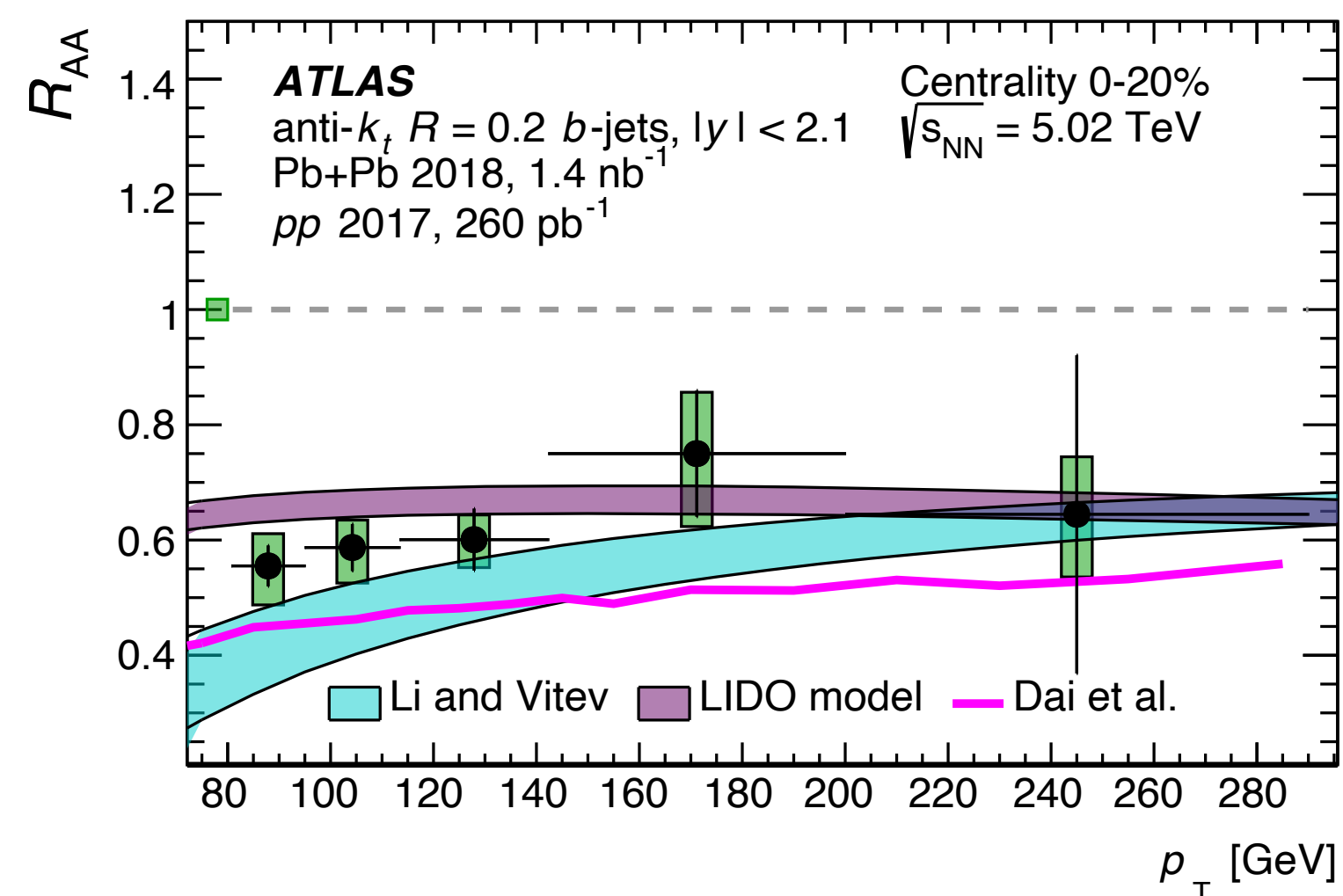
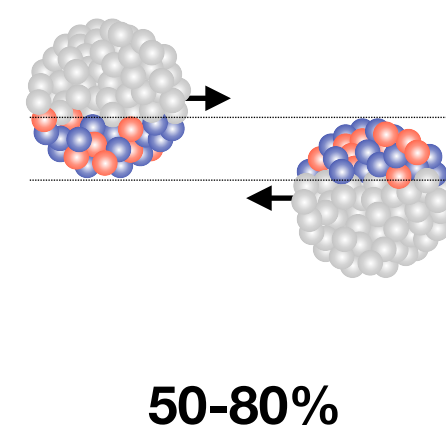
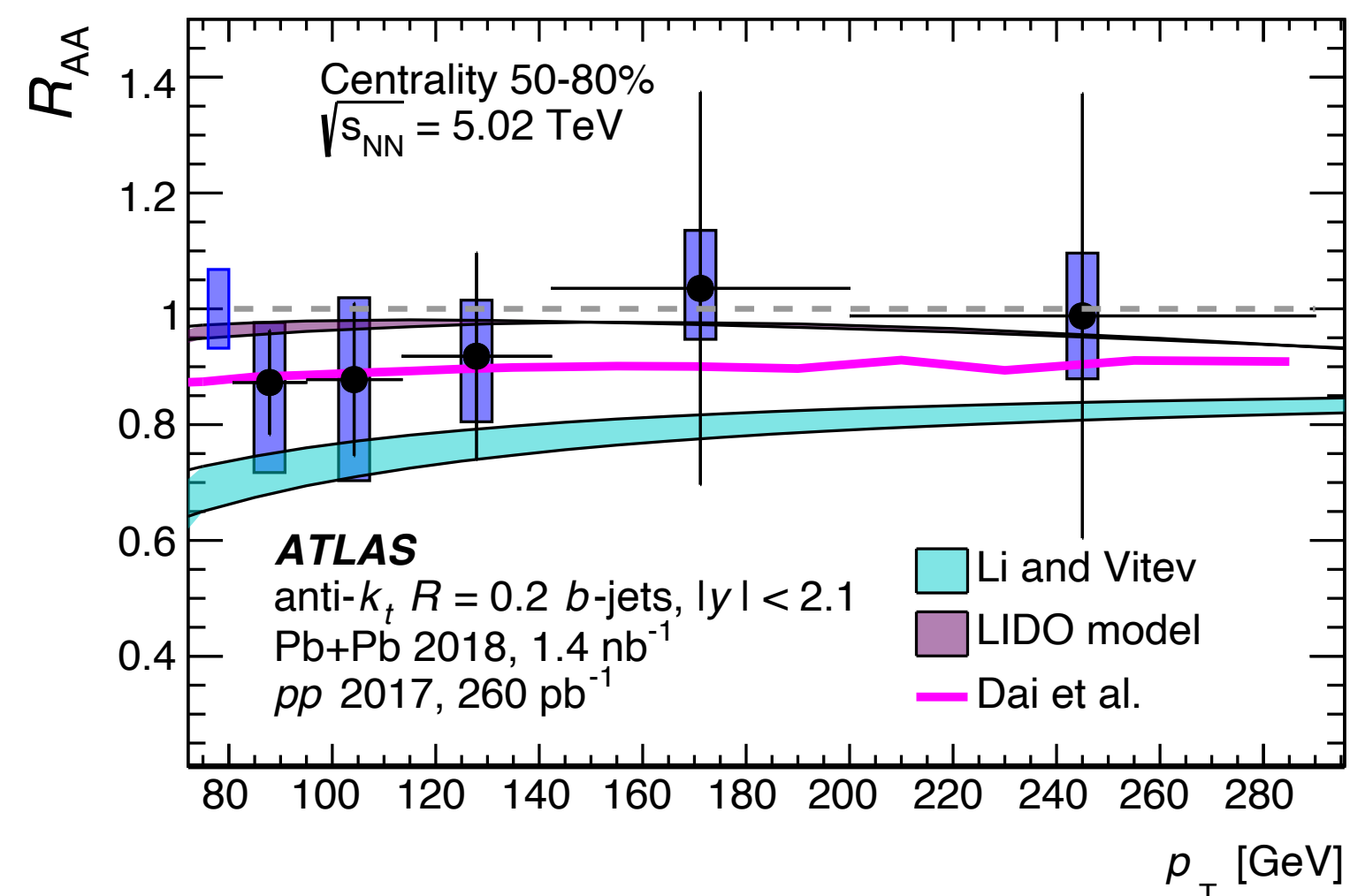
- Significantly more suppressed in 0-20% compare to 50-80%
- Weak dependence on  $b$ -jet  $p_T$



0-20%

# $b$ -jet $R_{AA}$ — model comparison

[arXiv:2204.13530](https://arxiv.org/abs/2204.13530)

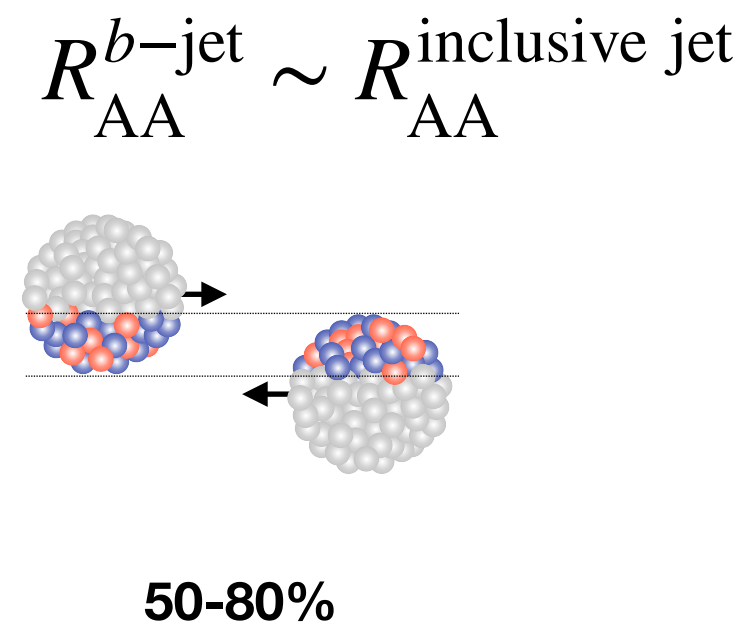
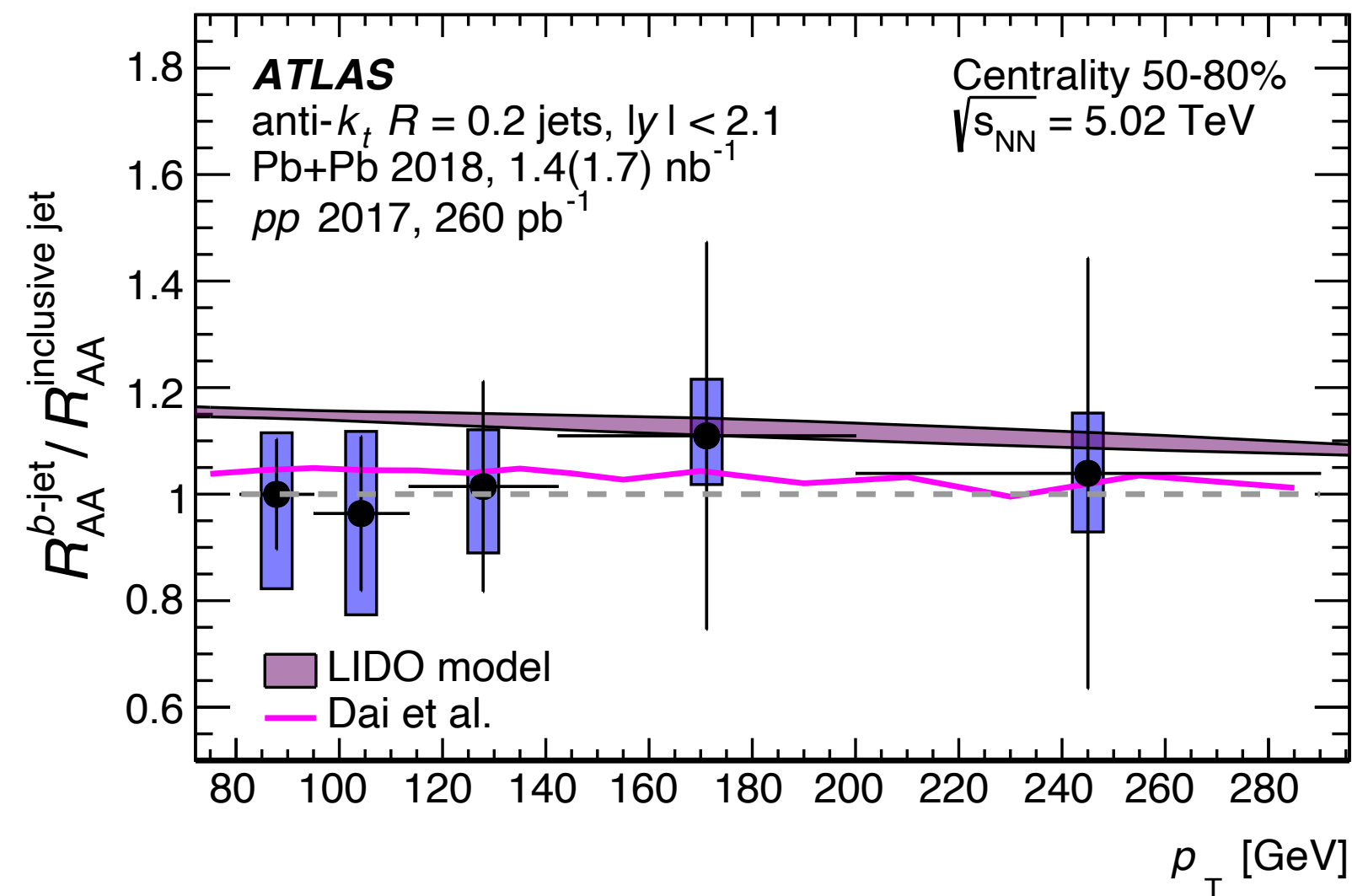


Models	Details	Uncertainty
LIDO*	FONLL + Boltzmann-Langevin transport	Medium-jet coupling
Dai et al.*	Sherpa + Langevin transport w/ radiation	N/A
Li and Vitev*	EFTs w/ medium modified splitting + Evolution equations	Medium-jet coupling

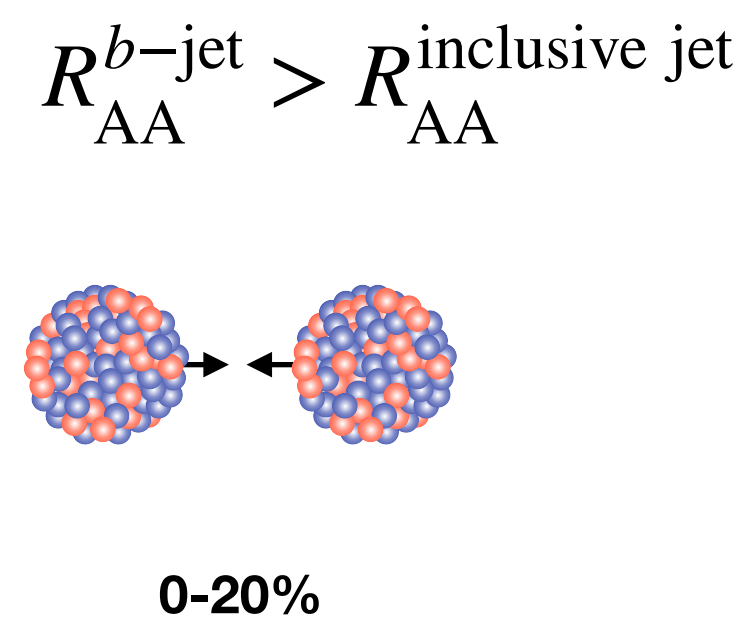
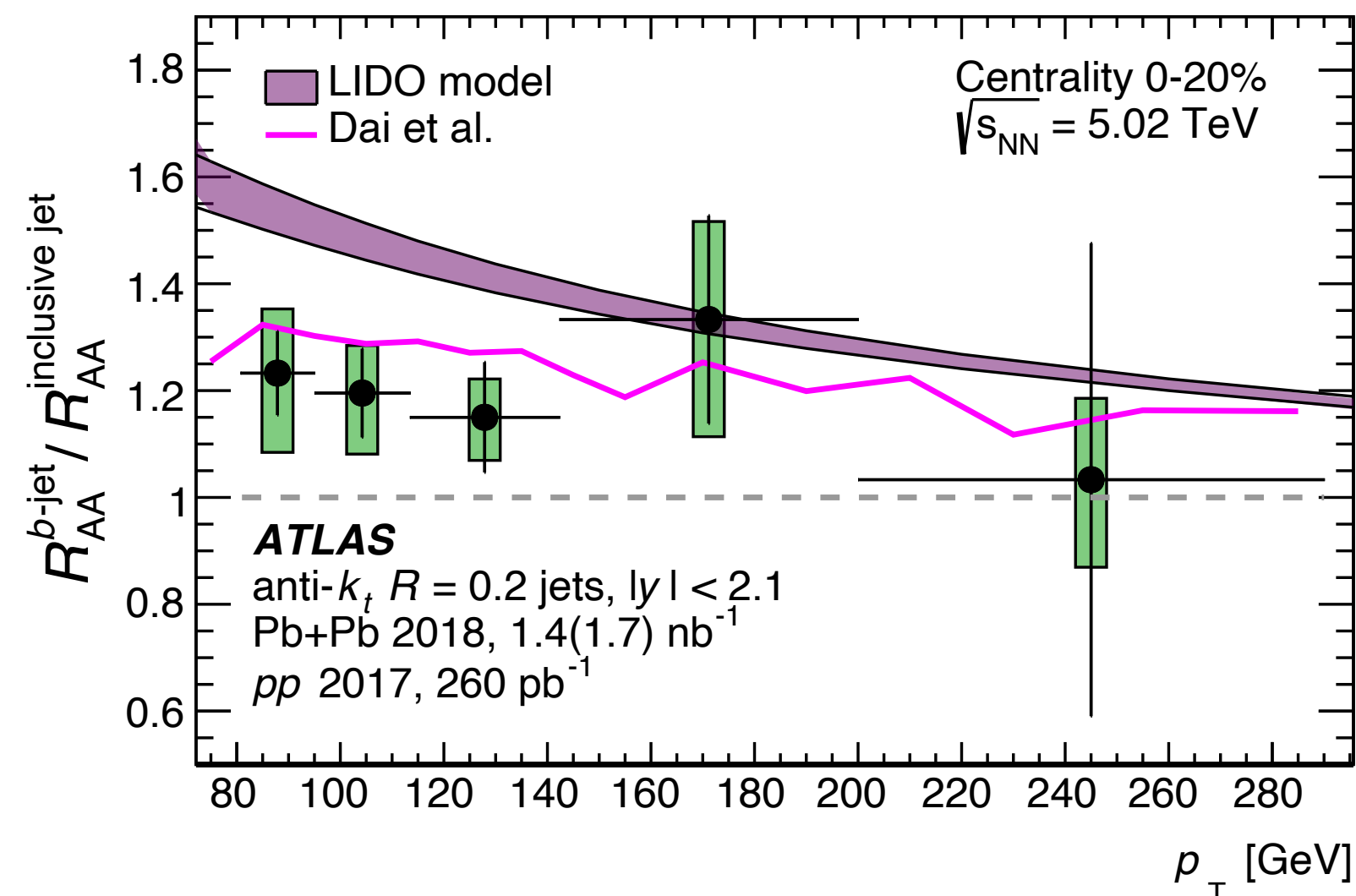
- LIDO (Boltzmann + Langevin) in good agreement with data
- Dai (modified Langevin) underpredicts  $b$ -jet  $R_{AA}$  in central collisions
- Li and Vitev (SCET) central values underpredict  $b$ -jet  $R_{AA}$  at low  $p_T$

# *b*-jet vs. inclusive jet

[arXiv:2204.13530](https://arxiv.org/abs/2204.13530)



Models	Details	Uncertainty
LIDO*	FONLL + Boltzmann-Langevin transport	Medium-jet coupling
Dai et al.*	Sherpa + Langevin transport w/ radiation	N/A



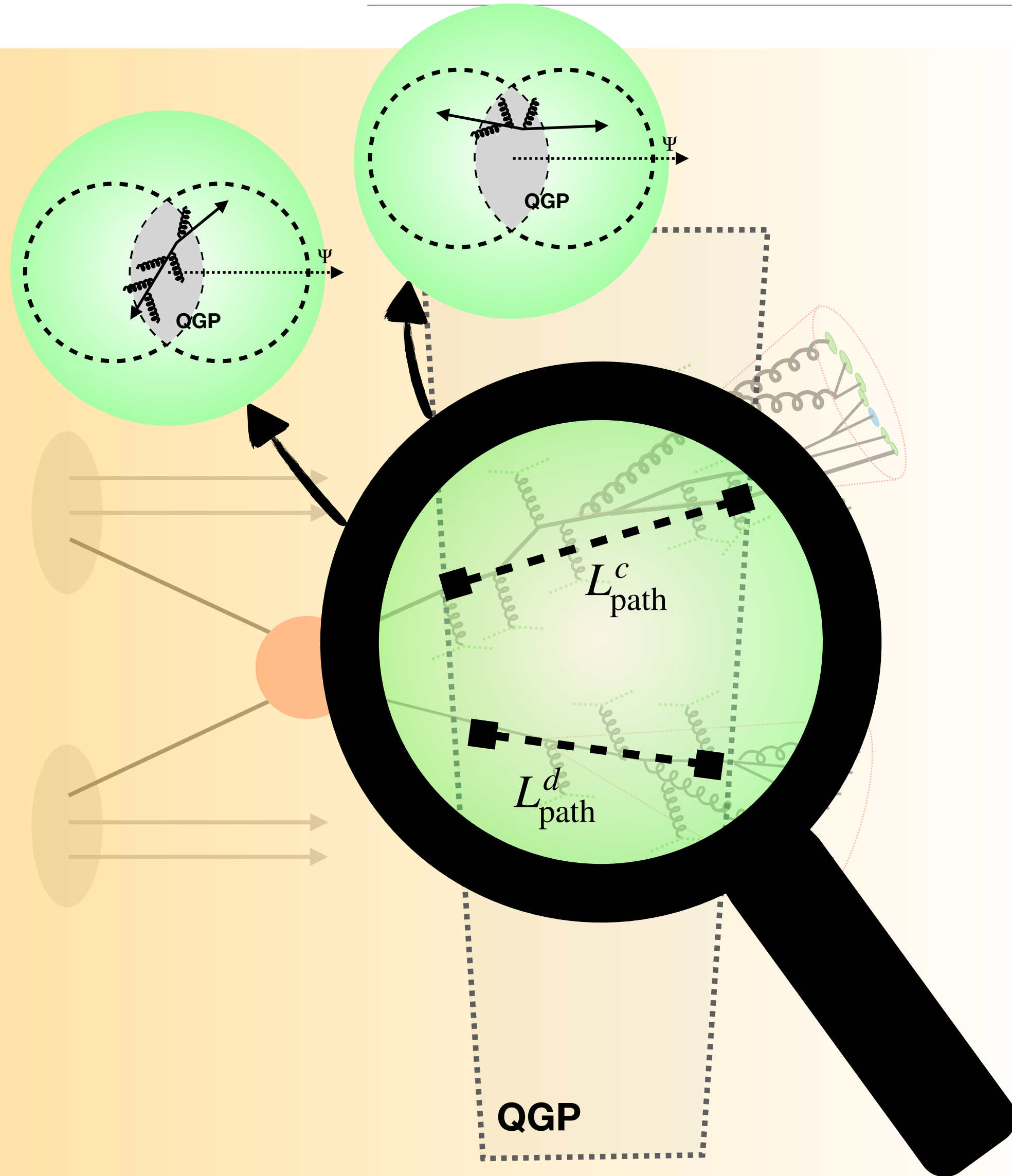
- LIDO (Boltzmann + Langevin) overpredicts relative difference
- Dai (modified Langevin) describes the relative difference better

Selected models cannot describe *b*-jet vs. inclusive jet simultaneously

Difference at different centralities: more likely a mass dependence driven effect



# Jet quenching — path-length dependence



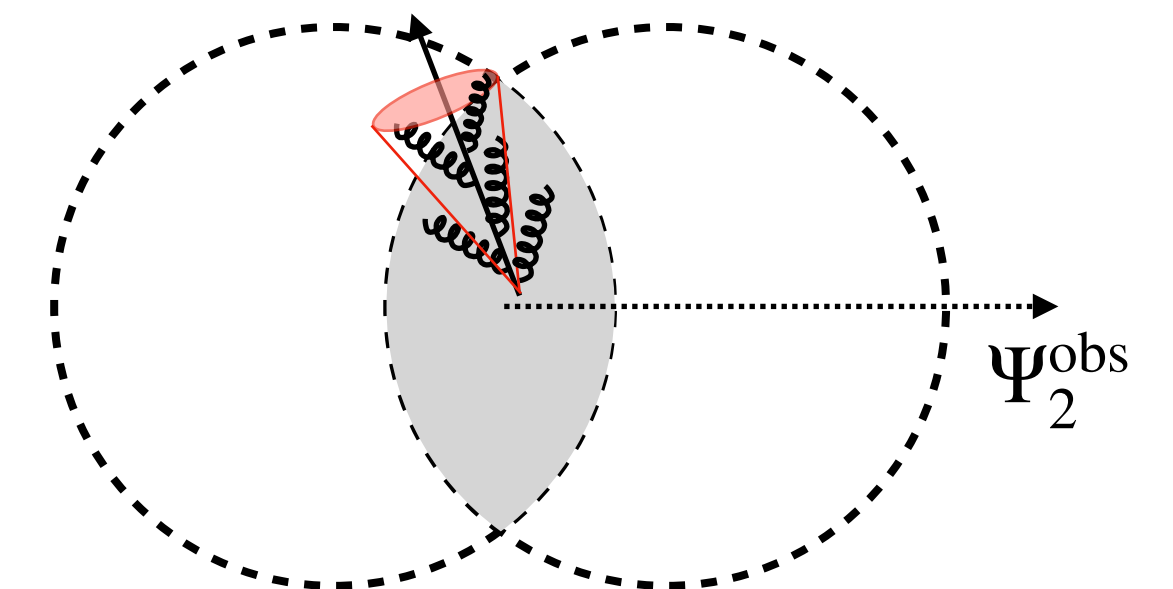
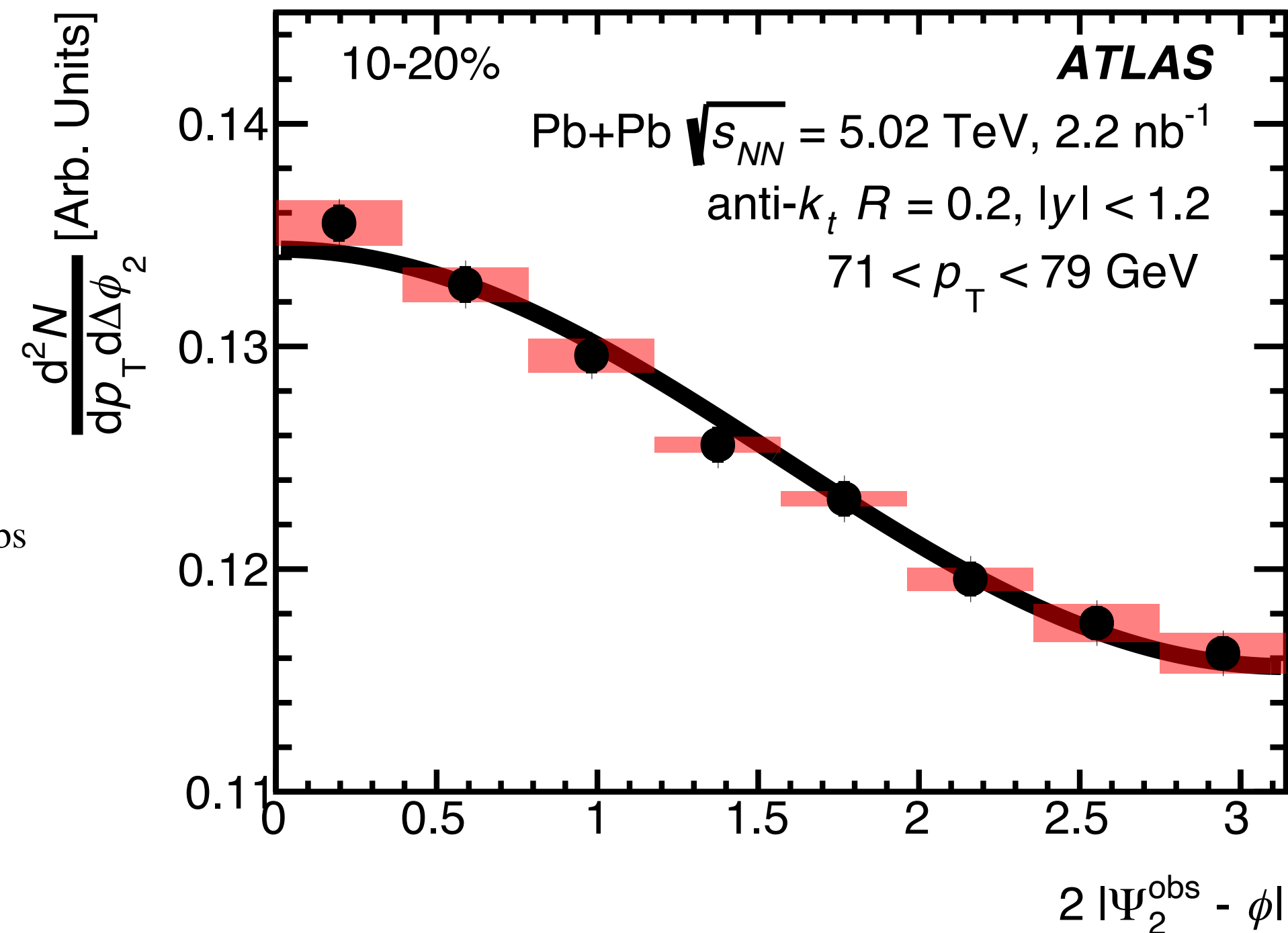
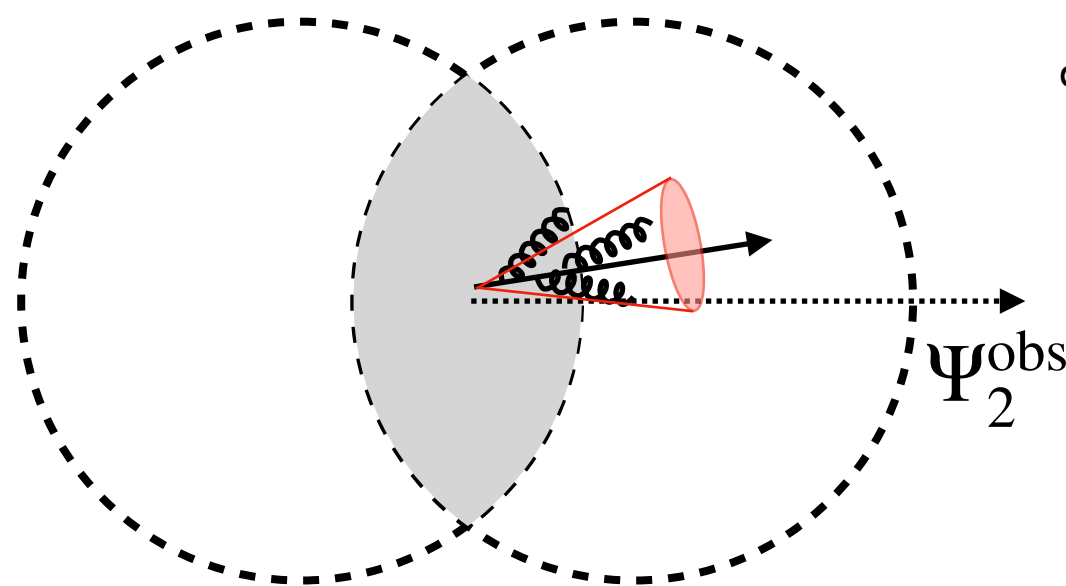
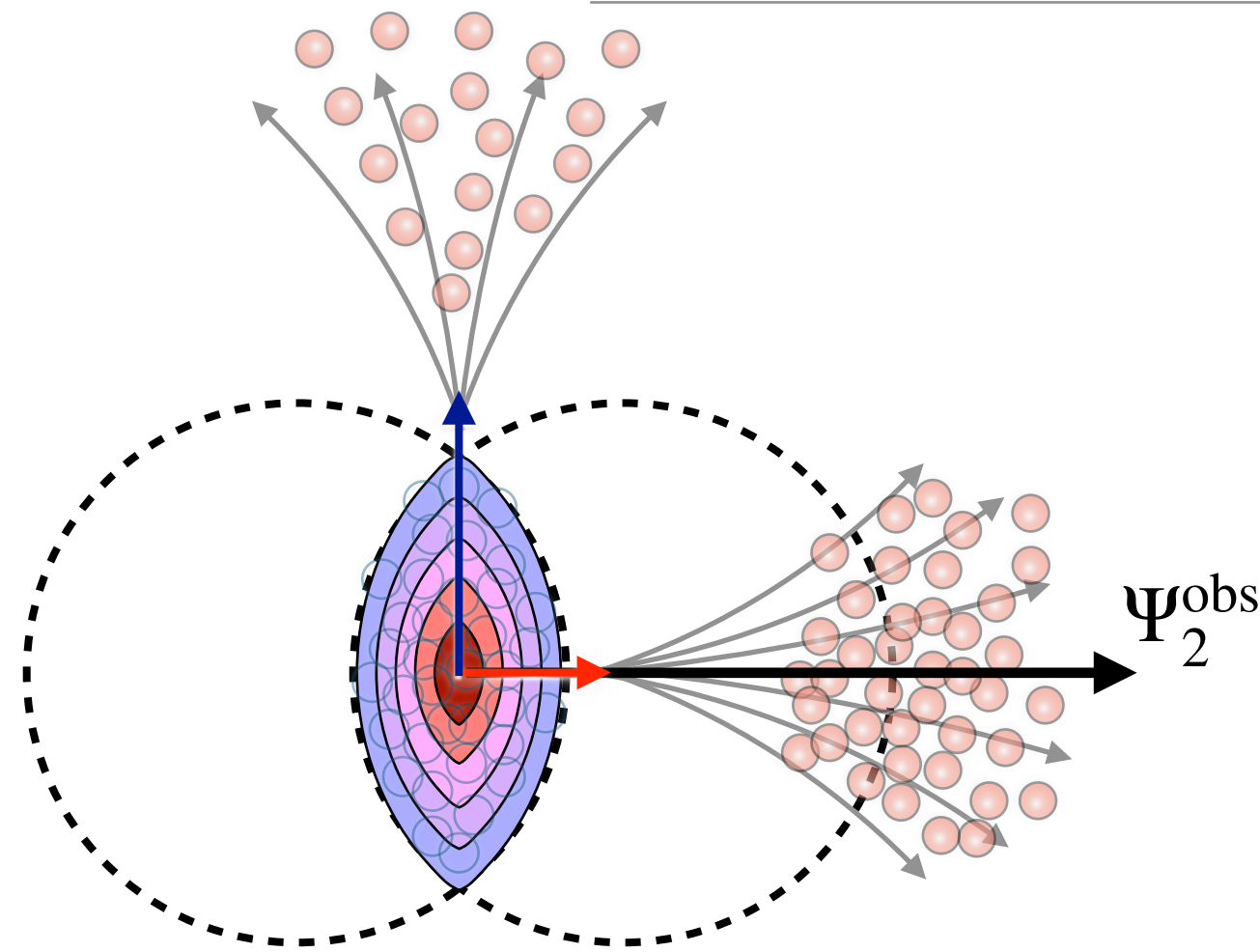
Control jet path-length through the QGP medium

- Correlate jet with event plane angle  $\rightarrow$
- Correlate back-to-back dijet events  $\rightarrow$

- Jet azimuthal anisotropy
- Dijet asymmetry

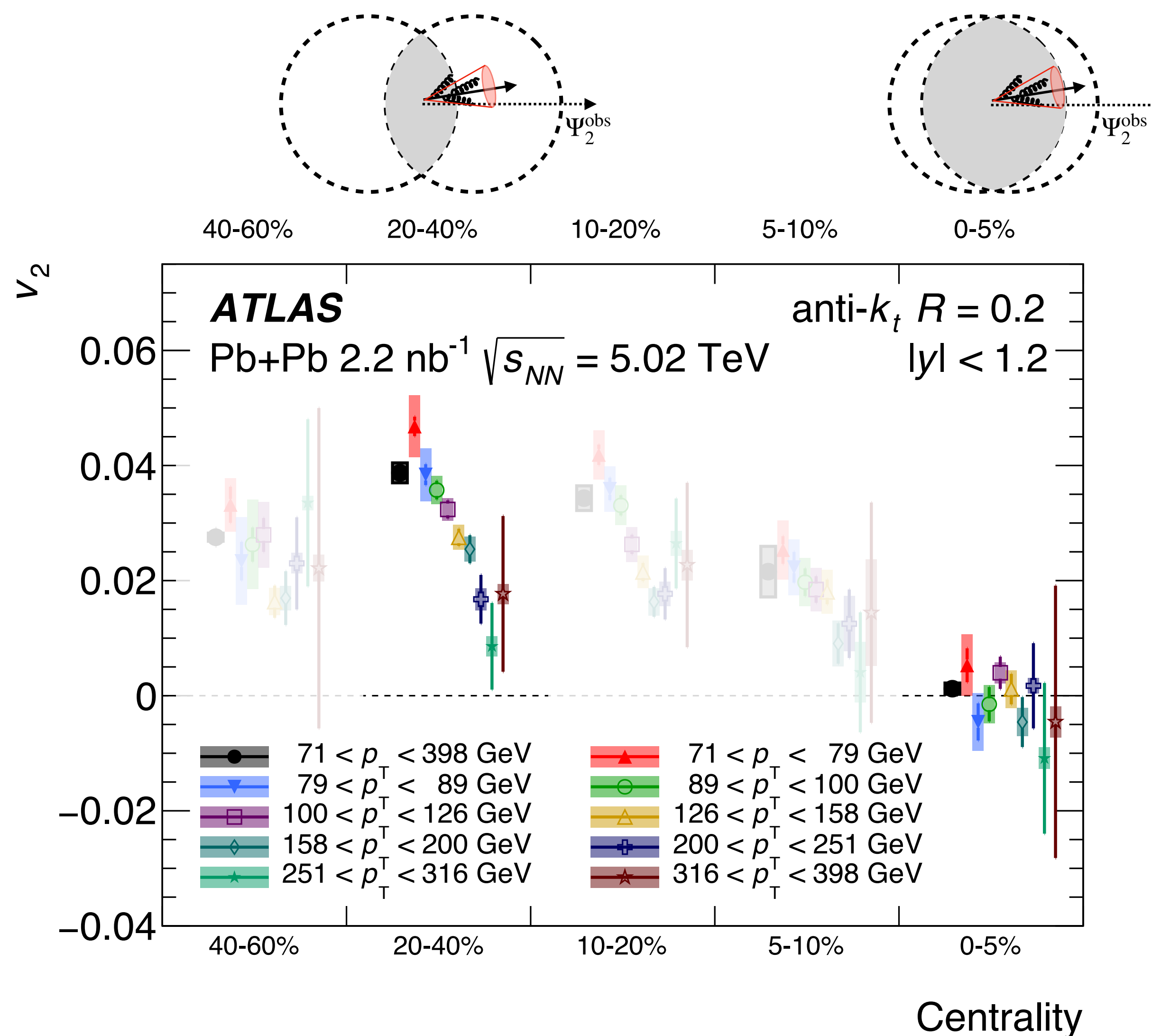
# Jet production wrt. event plane angle

$$\frac{dN_{\text{jet}}}{d\phi} \propto 1 + 2v_n \cos(n(\phi_{\text{jet}} - \Psi_2^{\text{obs}}))$$



# Jet $v_2$ results

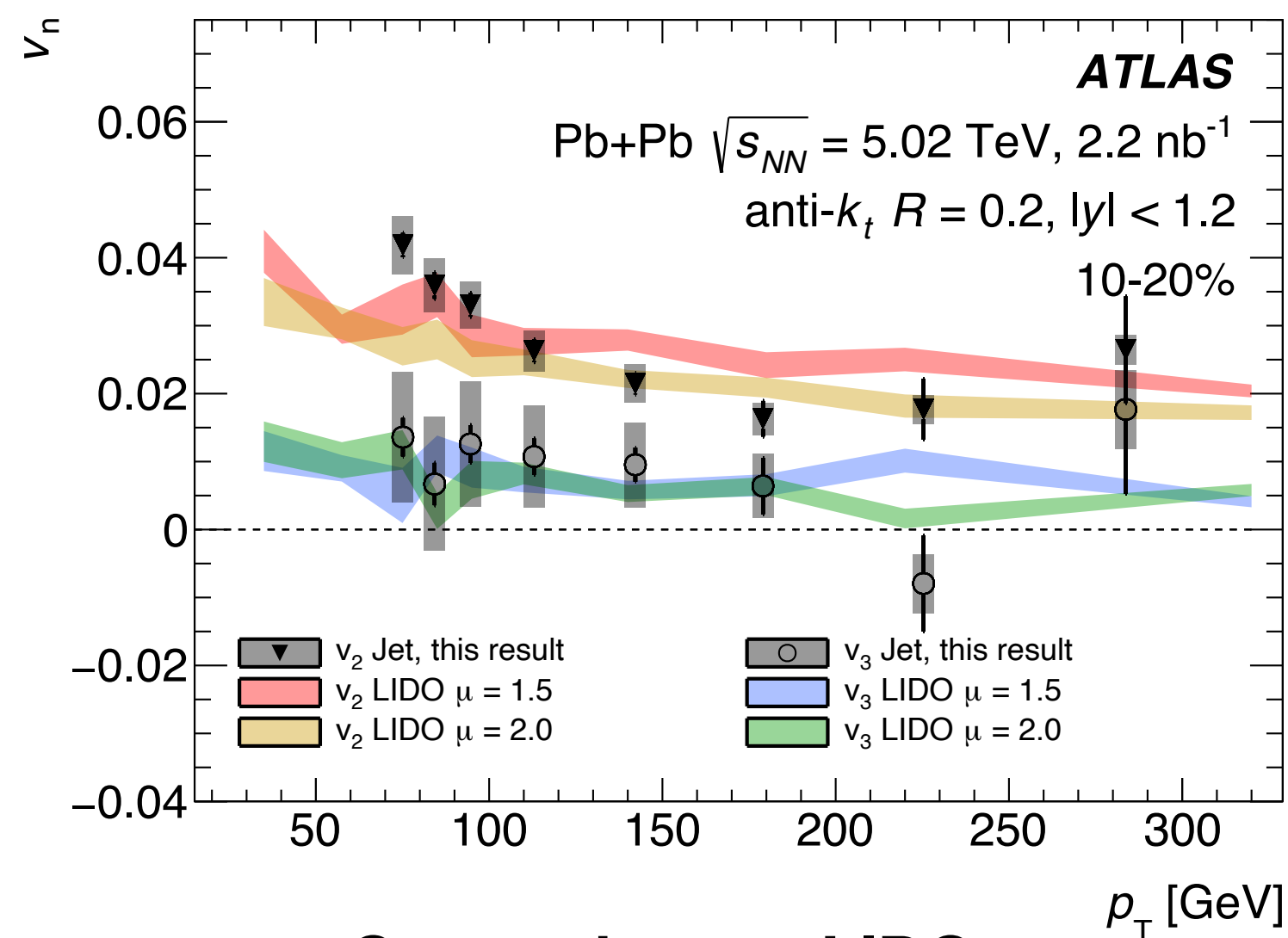
[arXiv:2111.06606](https://arxiv.org/abs/2111.06606)



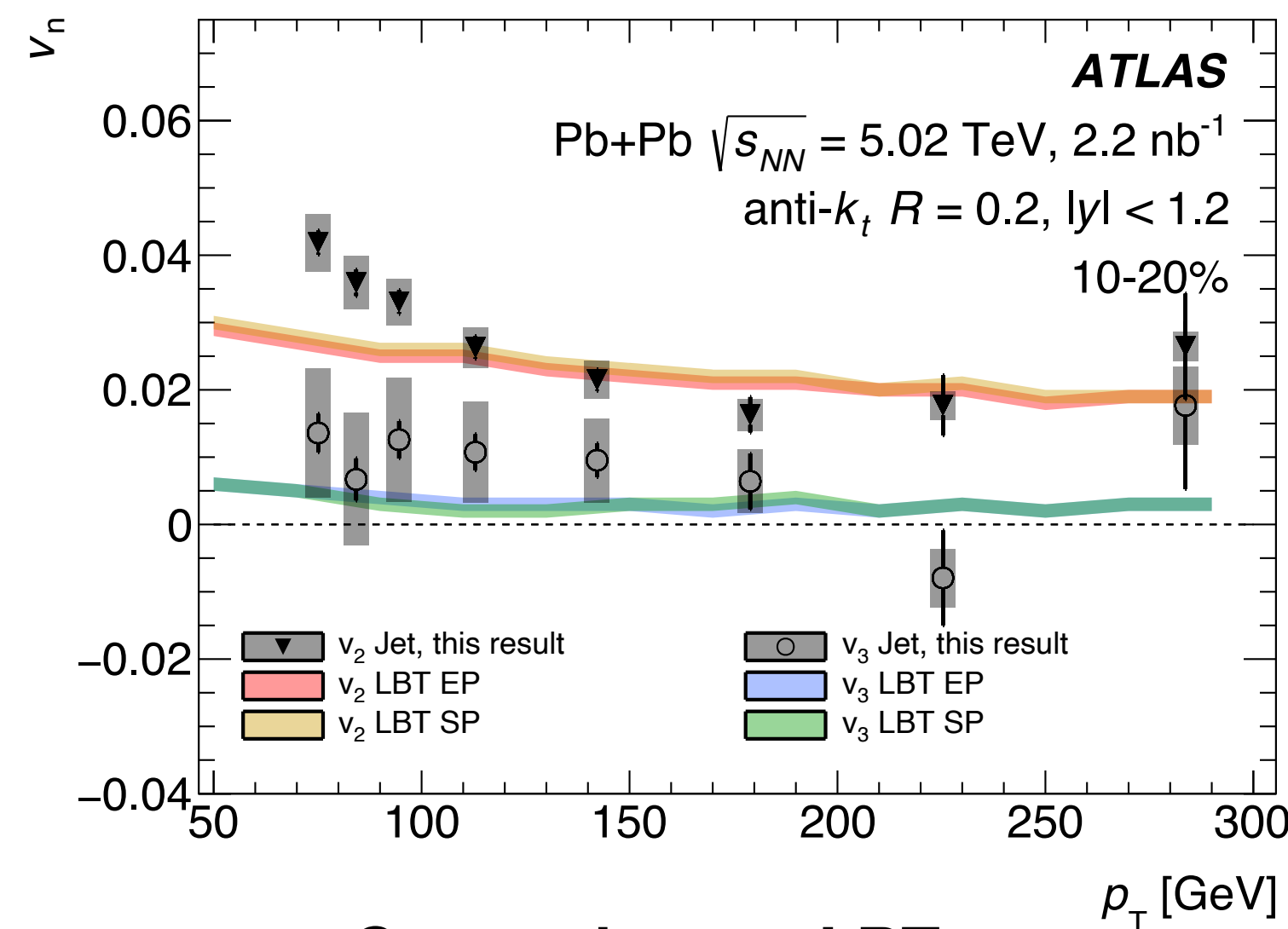
- Centrality dependence → Initial geometric anisotropy
- $p_T$  dependence → less suppression at higher  $p_T$ , surface bias

# Jet $v_n$ results — model comparisons

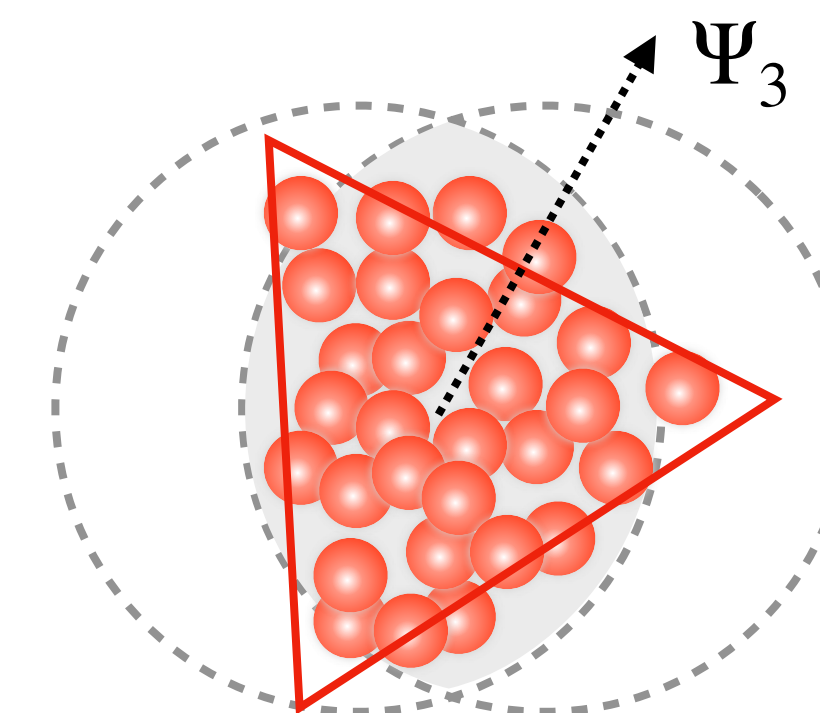
[arXiv:2111.06606](https://arxiv.org/abs/2111.06606)



**Comparison to LIDO**



**Comparison to LBT**

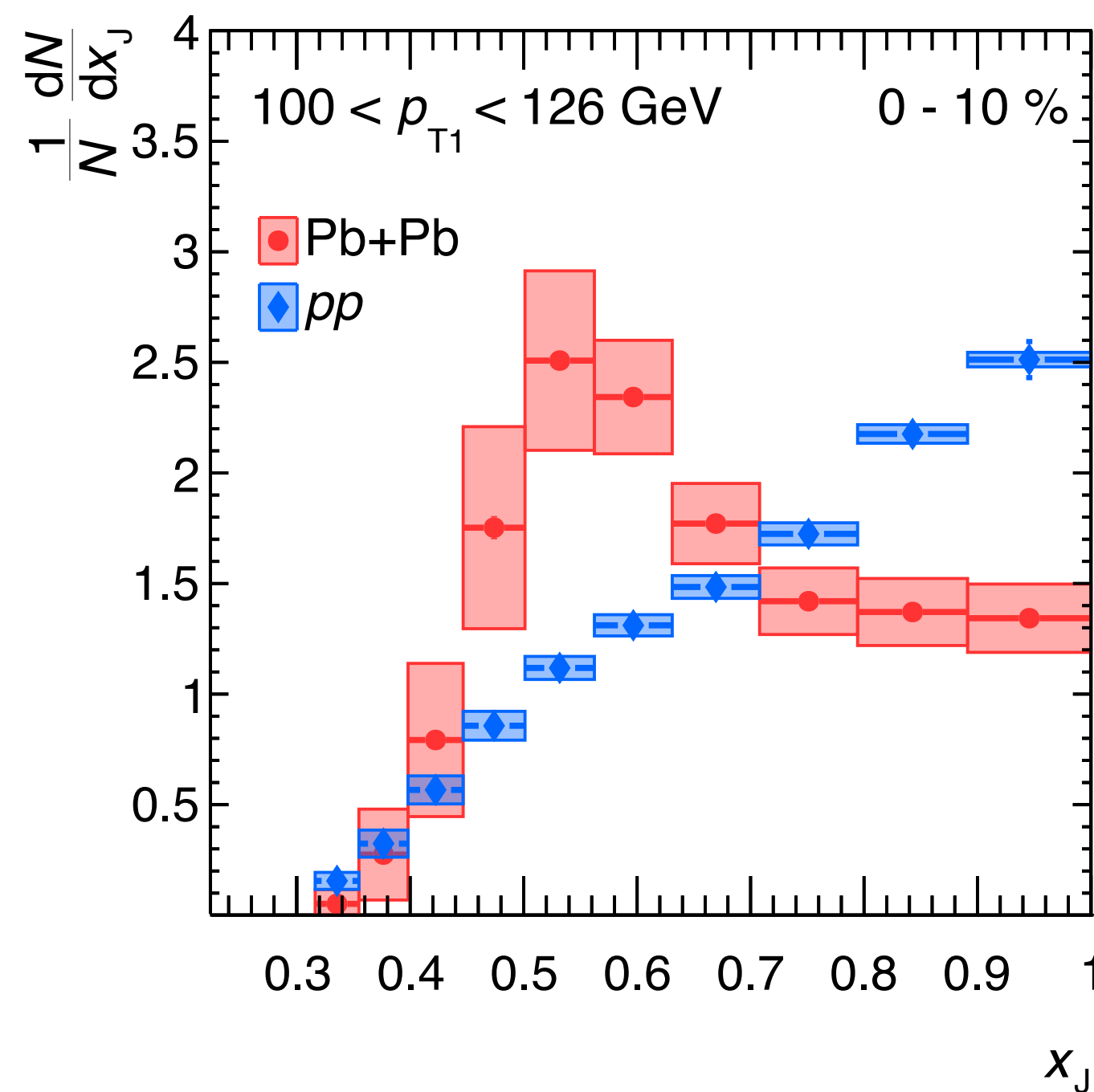


- LIDO<sup>\*</sup>: Pythia8 + Boltzmann + Langevin transport
- LBT<sup>\*</sup>: Pythia8 + Boltzmann transport
- Both transport models get the size of the  $v_2$  &  $v_3$  well, except for  $v_2$  below  $\sim 100$  GeV

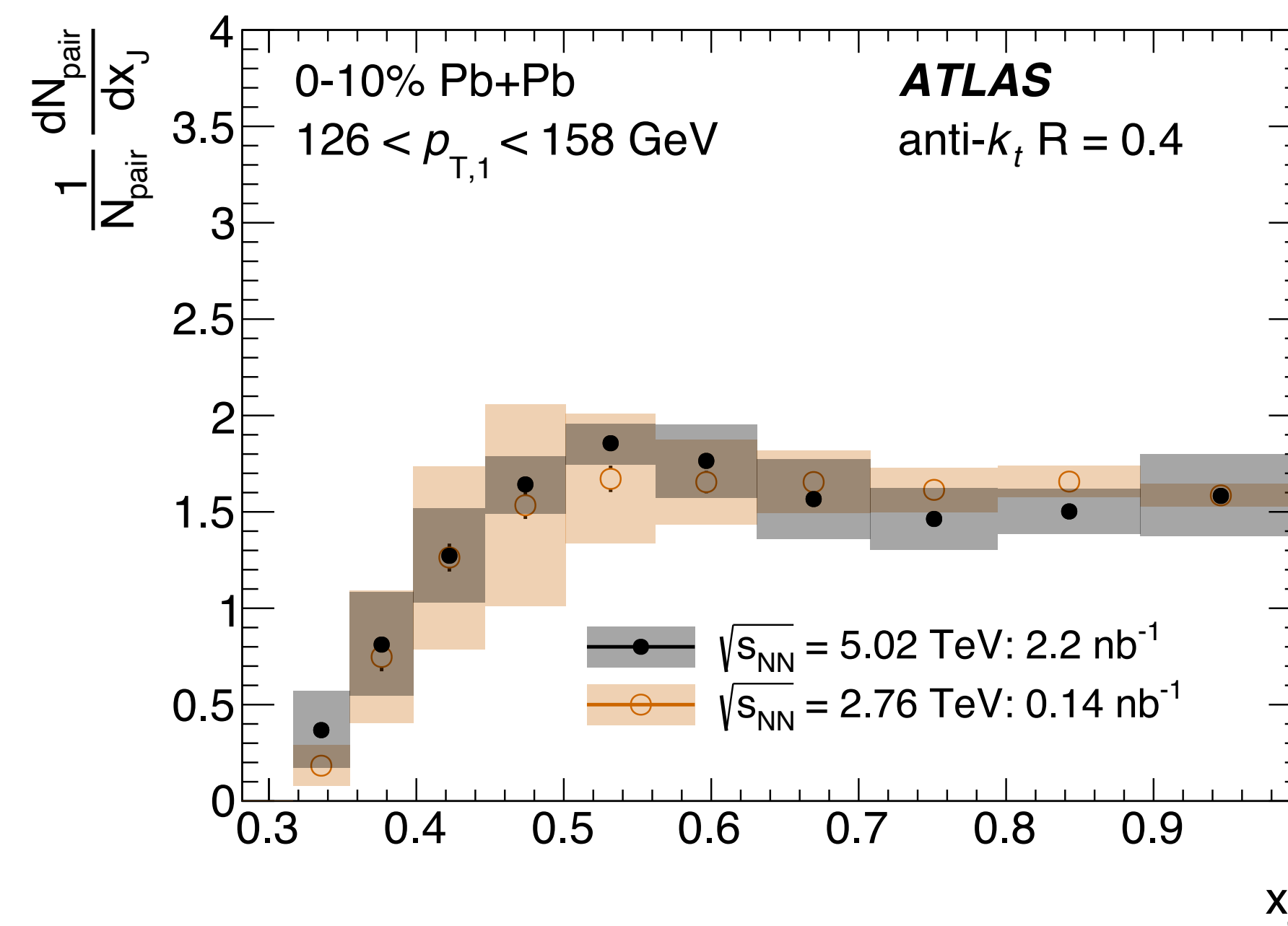


# Normalized dijet $x_J$ distribution

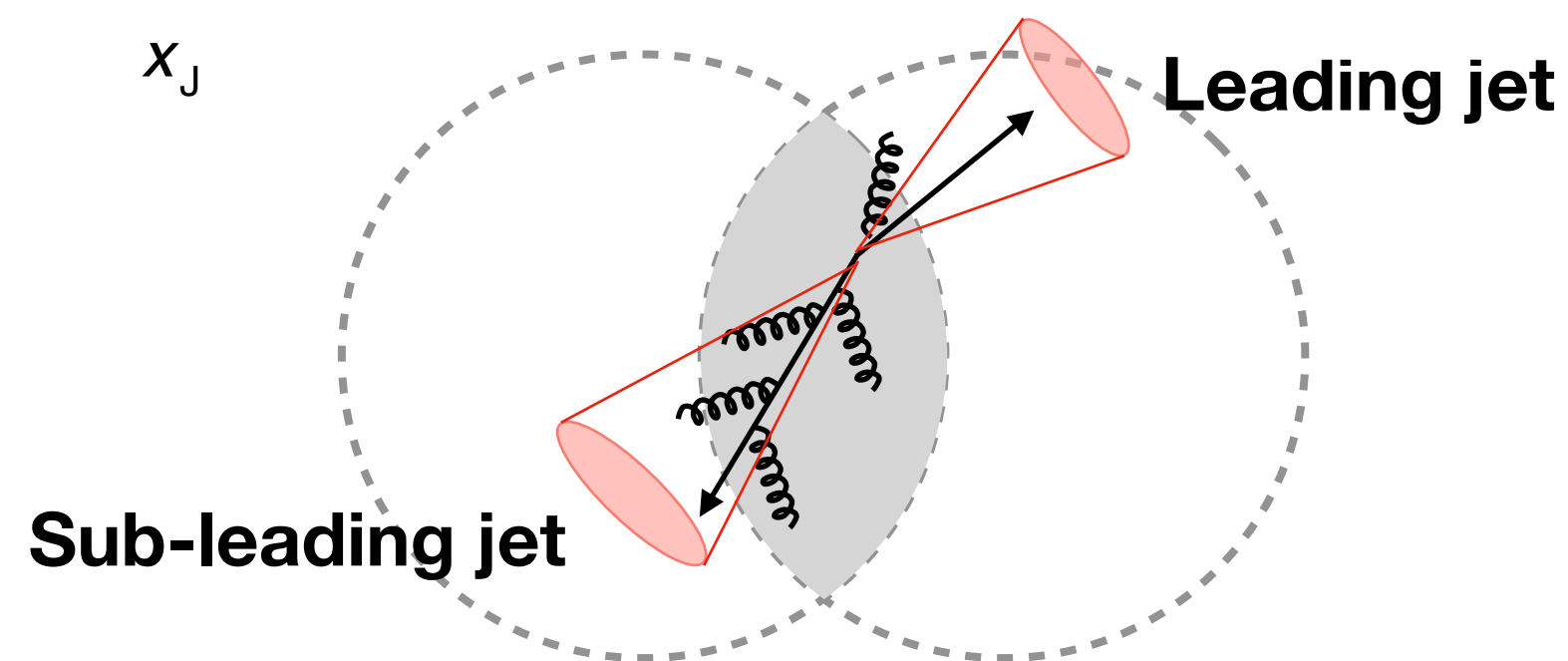
Physics Letters B 774 (2017) 379  
[arXiv:2205.00682](https://arxiv.org/abs/2205.00682)



**ATLAS**  
 anti- $k_t$   $R = 0.4$  jets  
 $\sqrt{s_{NN}} = 2.76 \text{ TeV}$   
 2011 Pb+Pb data,  $0.14 \text{ nb}^{-1}$   
 2013 pp data,  $4.0 \text{ pb}^{-1}$

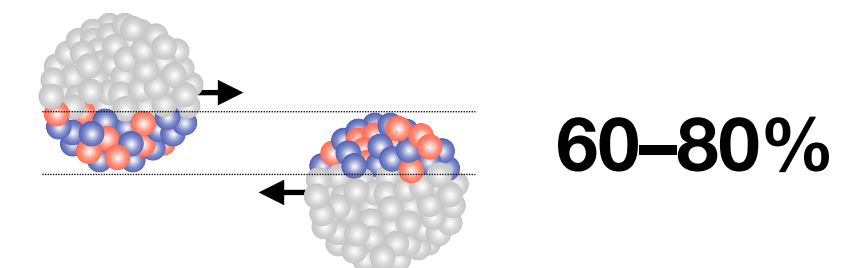
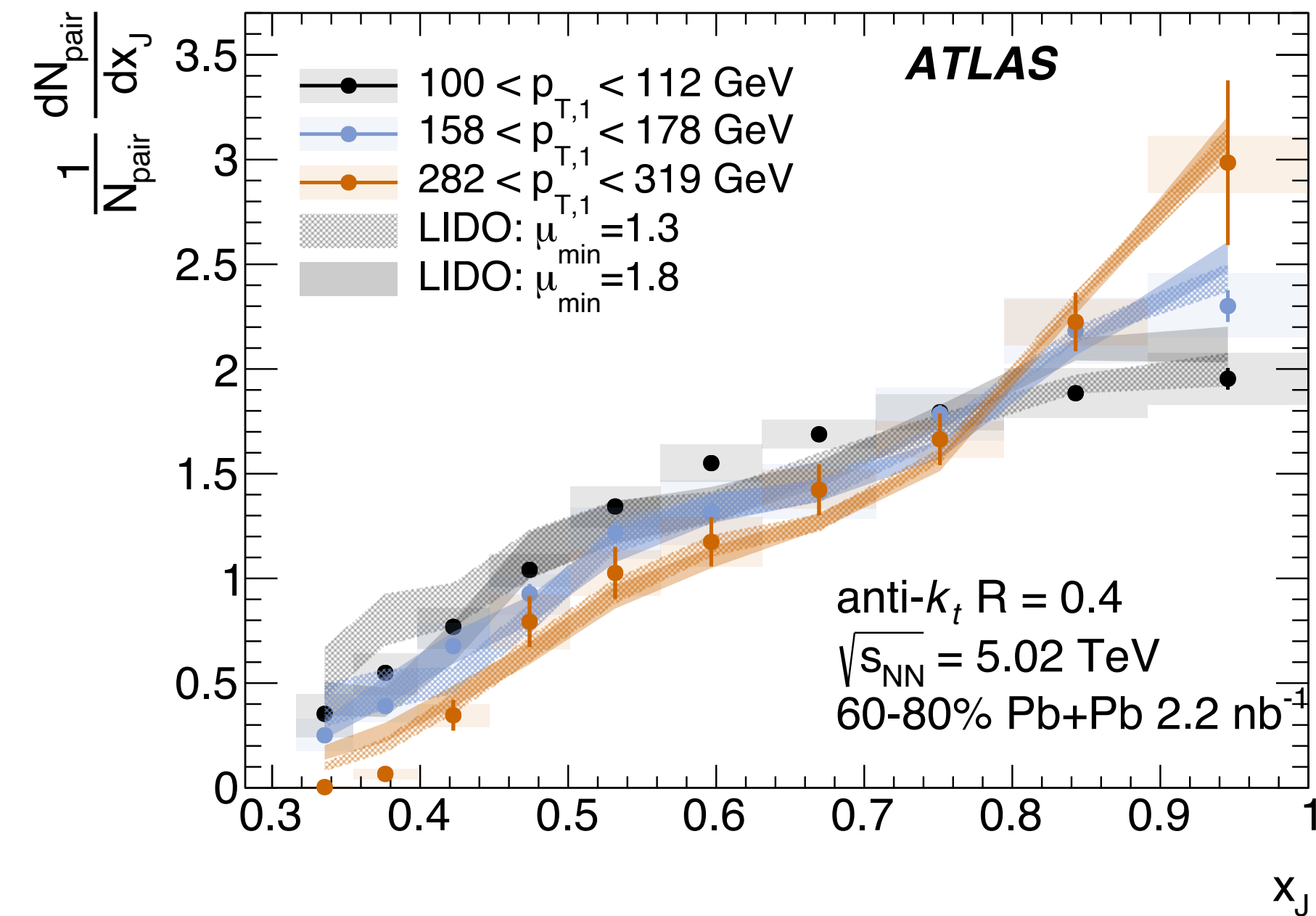
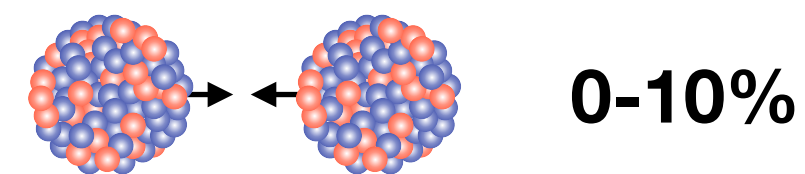
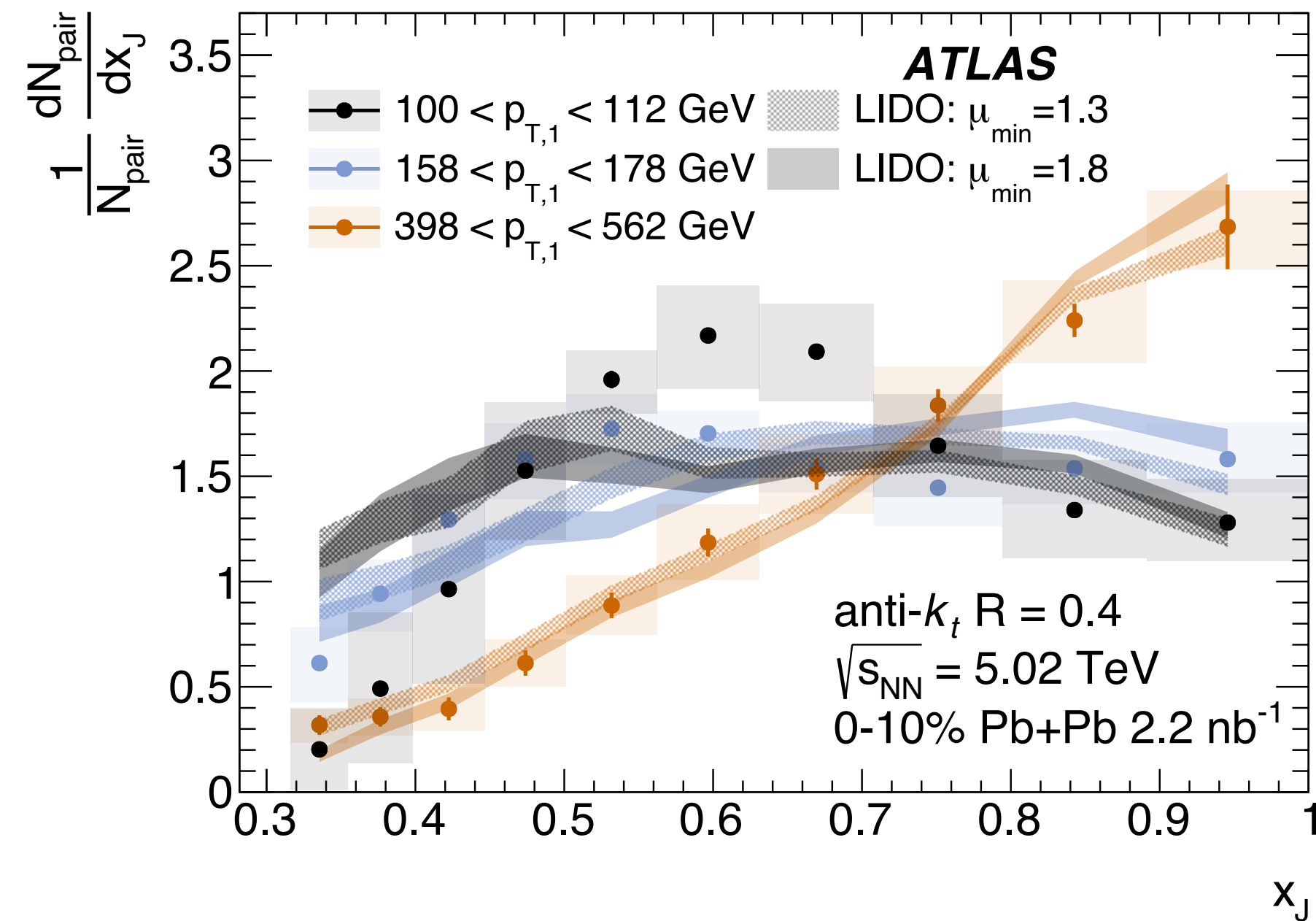


$$x_J = \frac{p_T^{\text{sub-leading}}}{p_T^{\text{leading}}}$$



# Normalized dijet $x_J$ model comparison

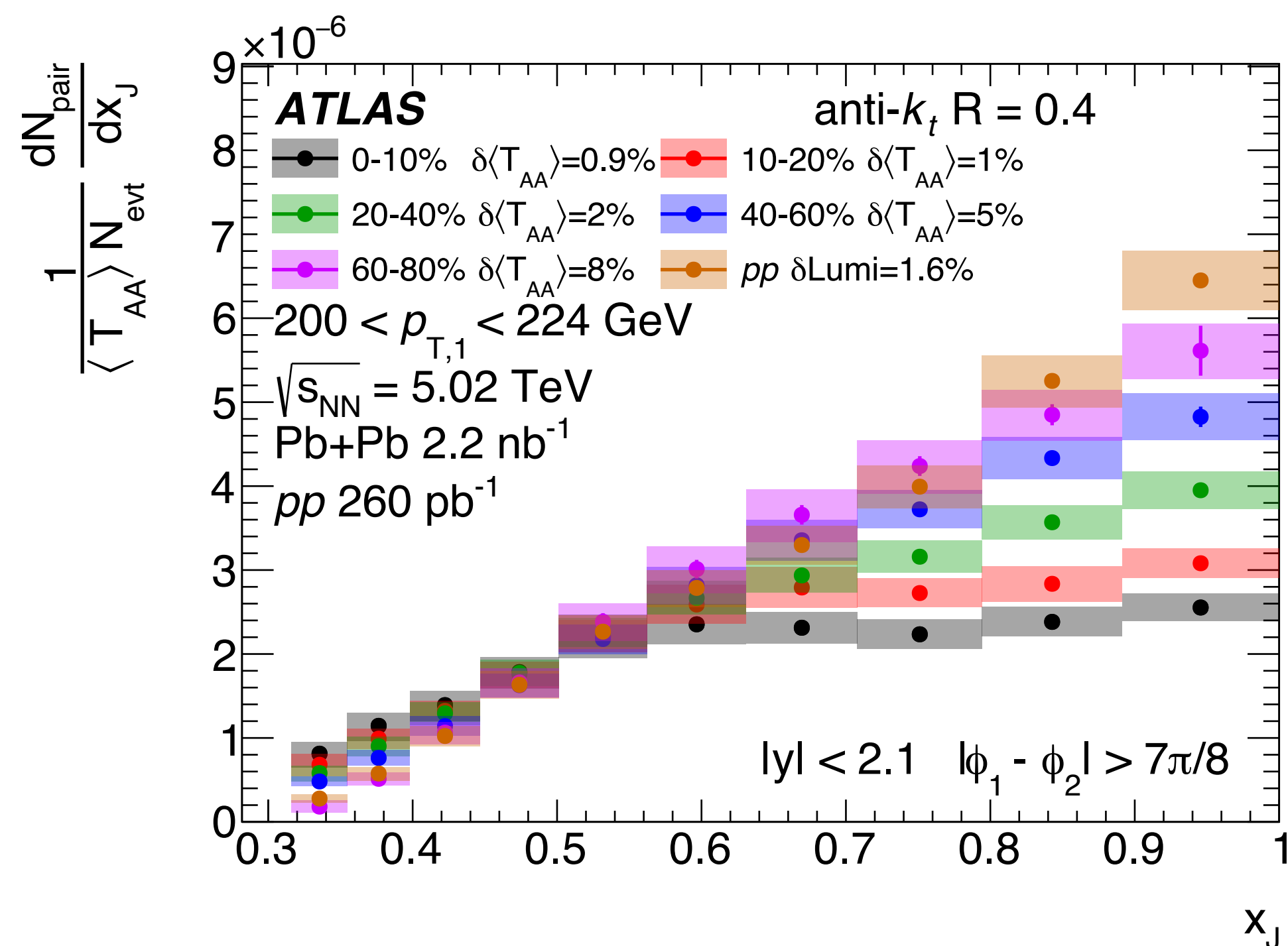
[arXiv:2205.00682](https://arxiv.org/abs/2205.00682)



LIDO transport model reproduces high  $p_T$  and peripheral results, not low  $p_T$  and in central collision

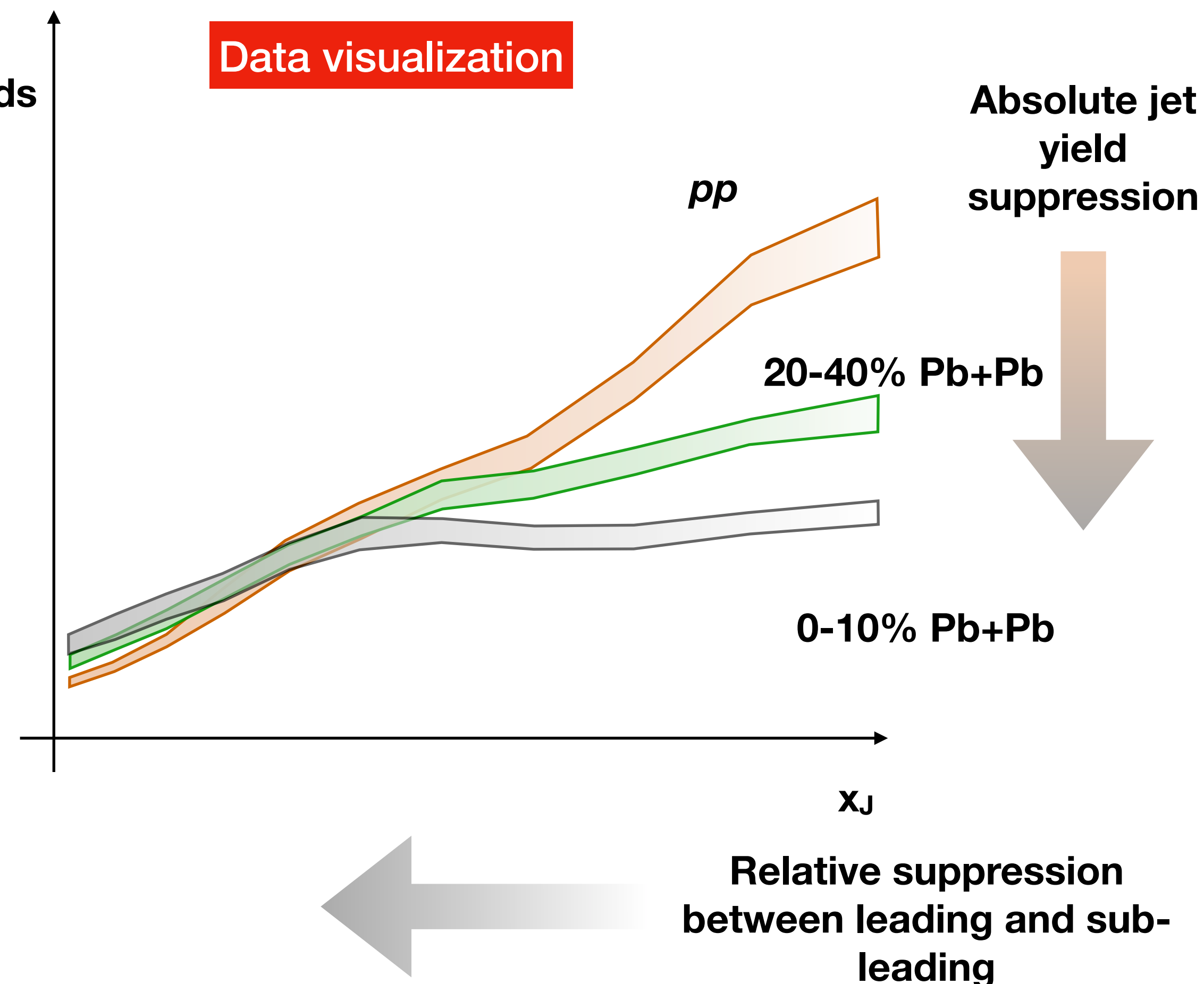
# Absolutely normalized dijet $x_J$ distribution

[arXiv:2205.00682](https://arxiv.org/abs/2205.00682)



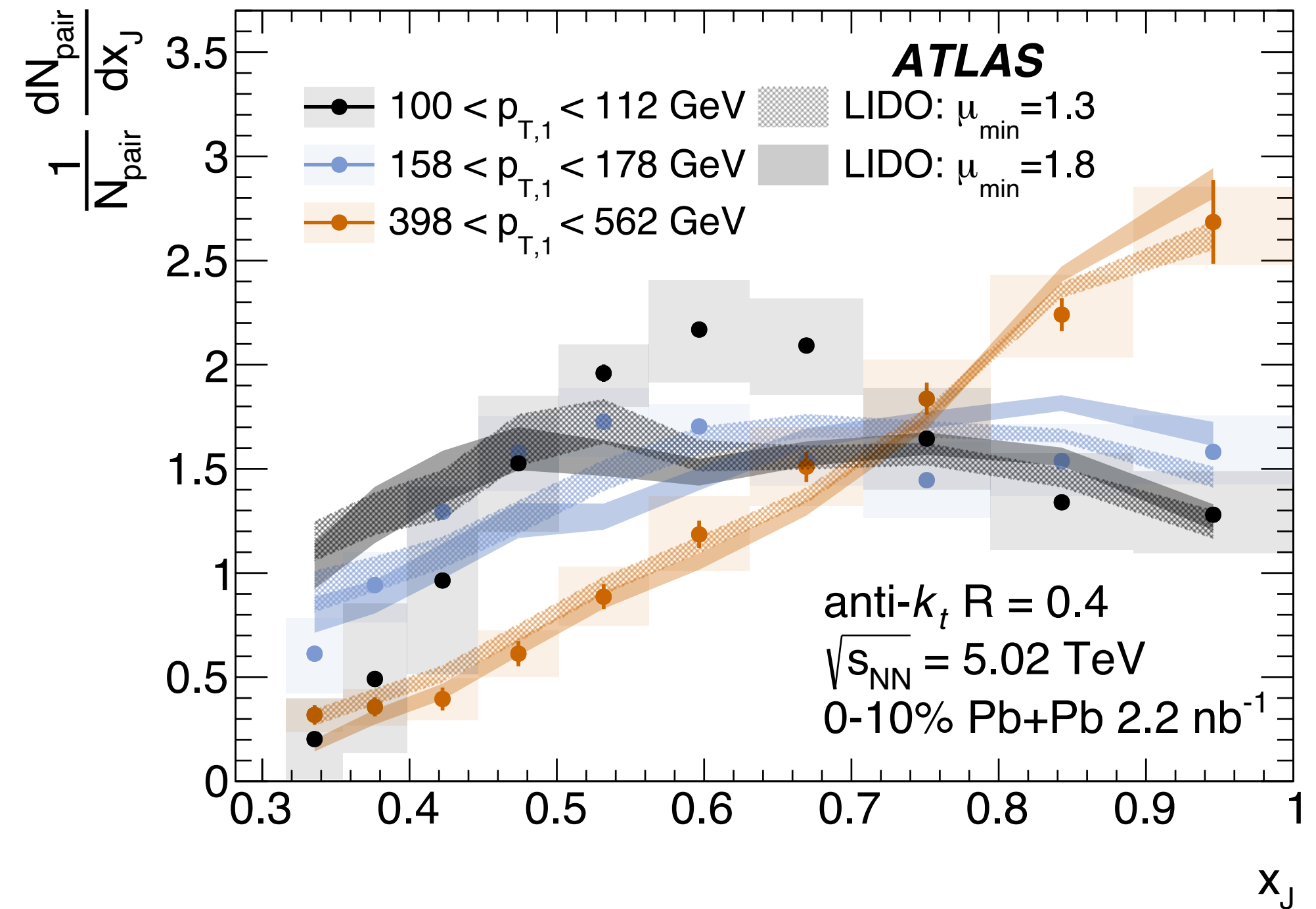
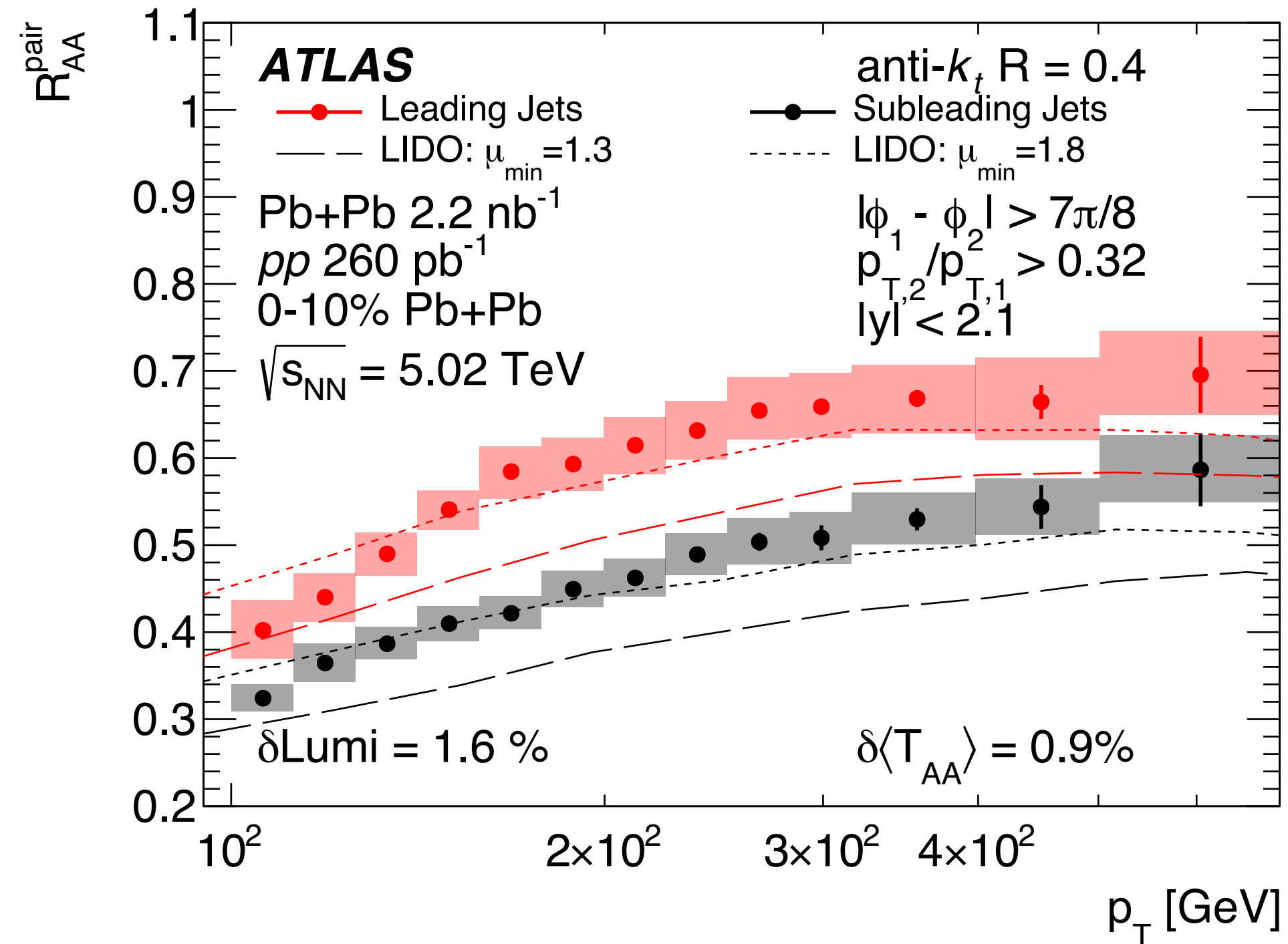
Di-jet yields

Data visualization

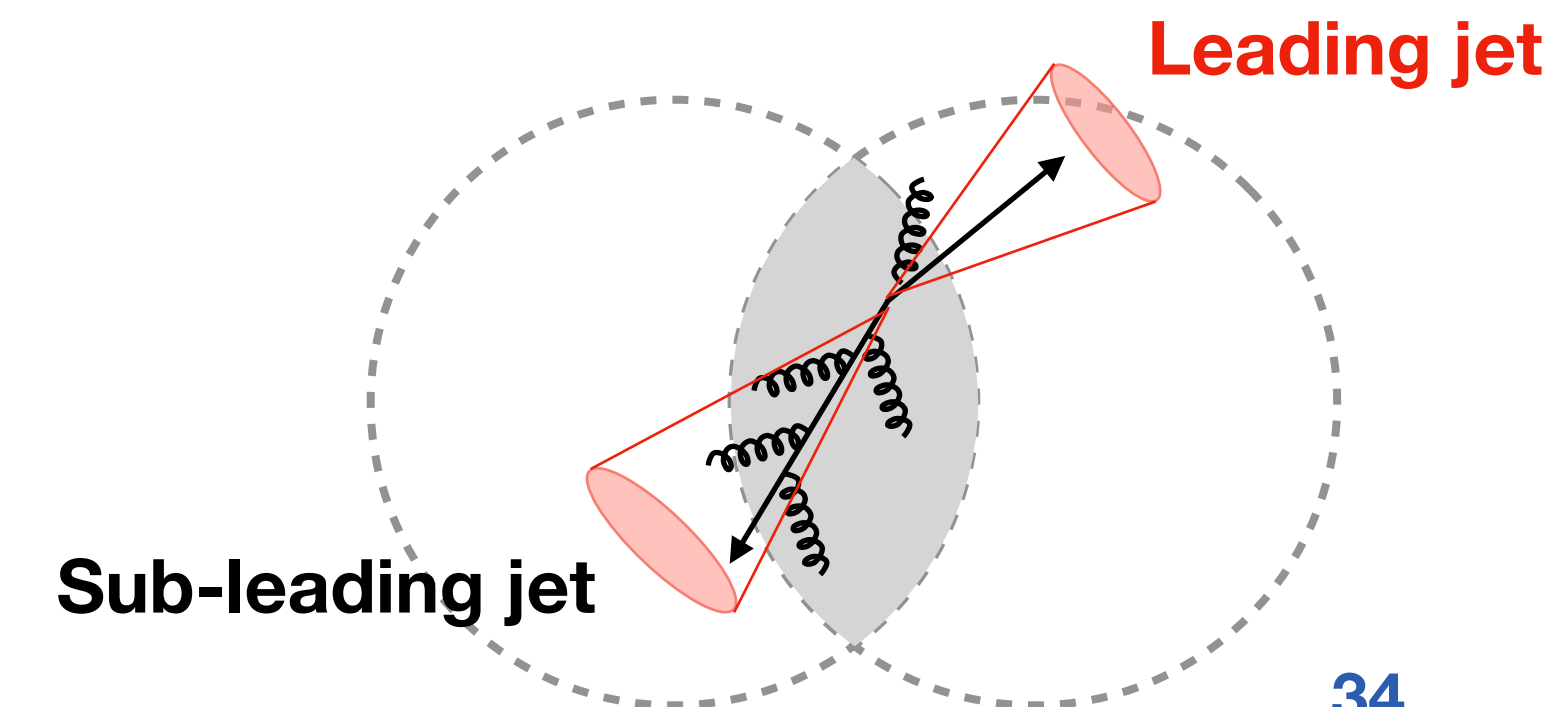


# Dijet $R_{AA}$

arXiv:2205.00682



$$R_{AA}^{\text{pair}}(p_T^{\text{leading}}) \propto \frac{\int_{\text{Pb+Pb}} N(p_T^{\text{leading}}, p_T^{\text{sub}}) dp_T^{\text{sub}}}{\int_{pp} N(p_T^{\text{leading}}, p_T^{\text{sub}}) dp_T^{\text{sub}}}$$



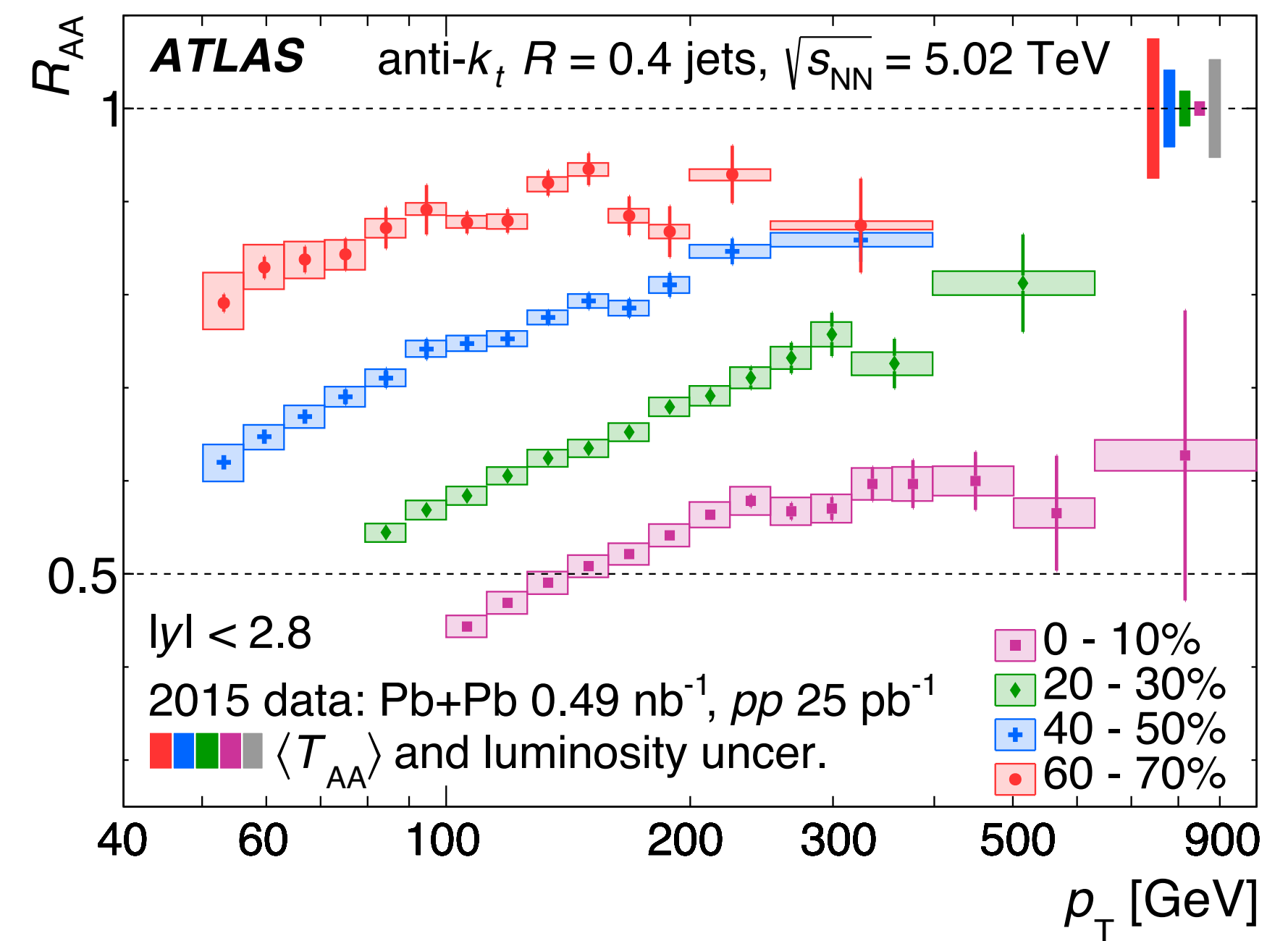
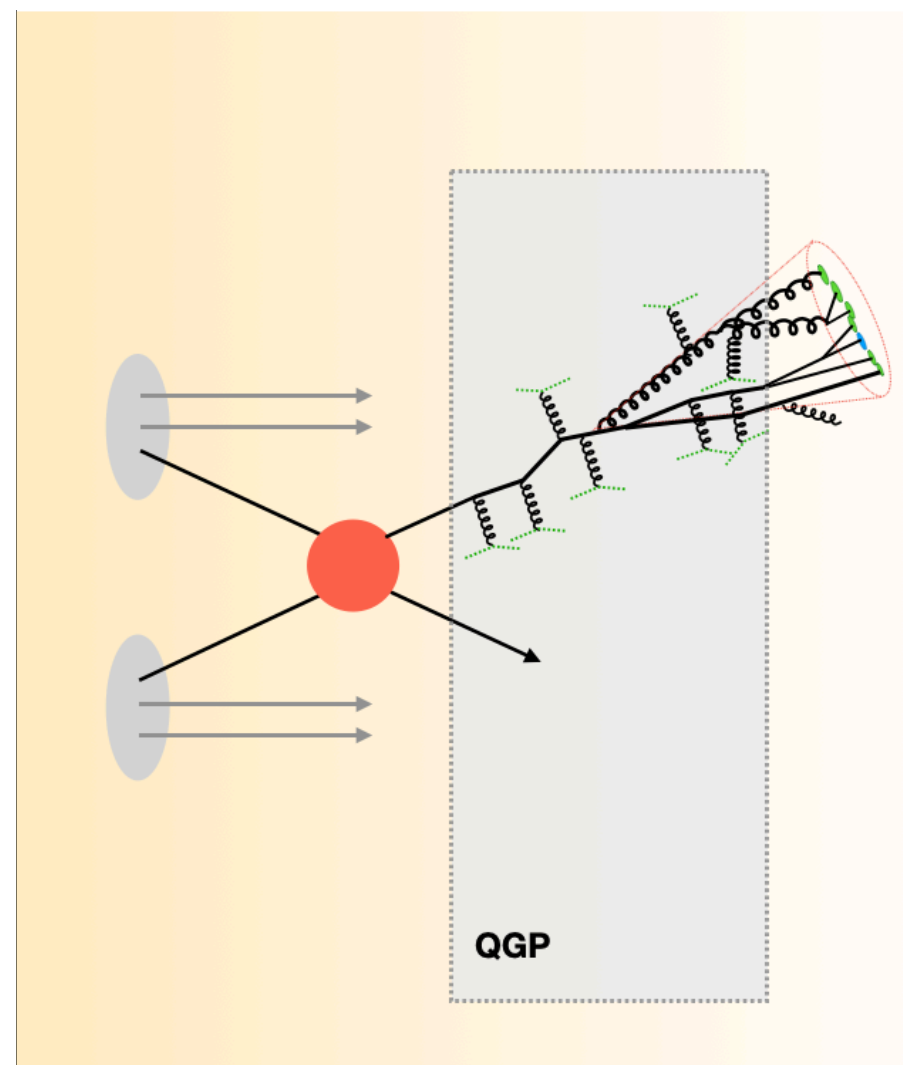


# Summary

Parton-QGP interaction  $\xrightarrow{\text{observation}}$  Jet quenching

QGP size/temperature

Centrality dependence

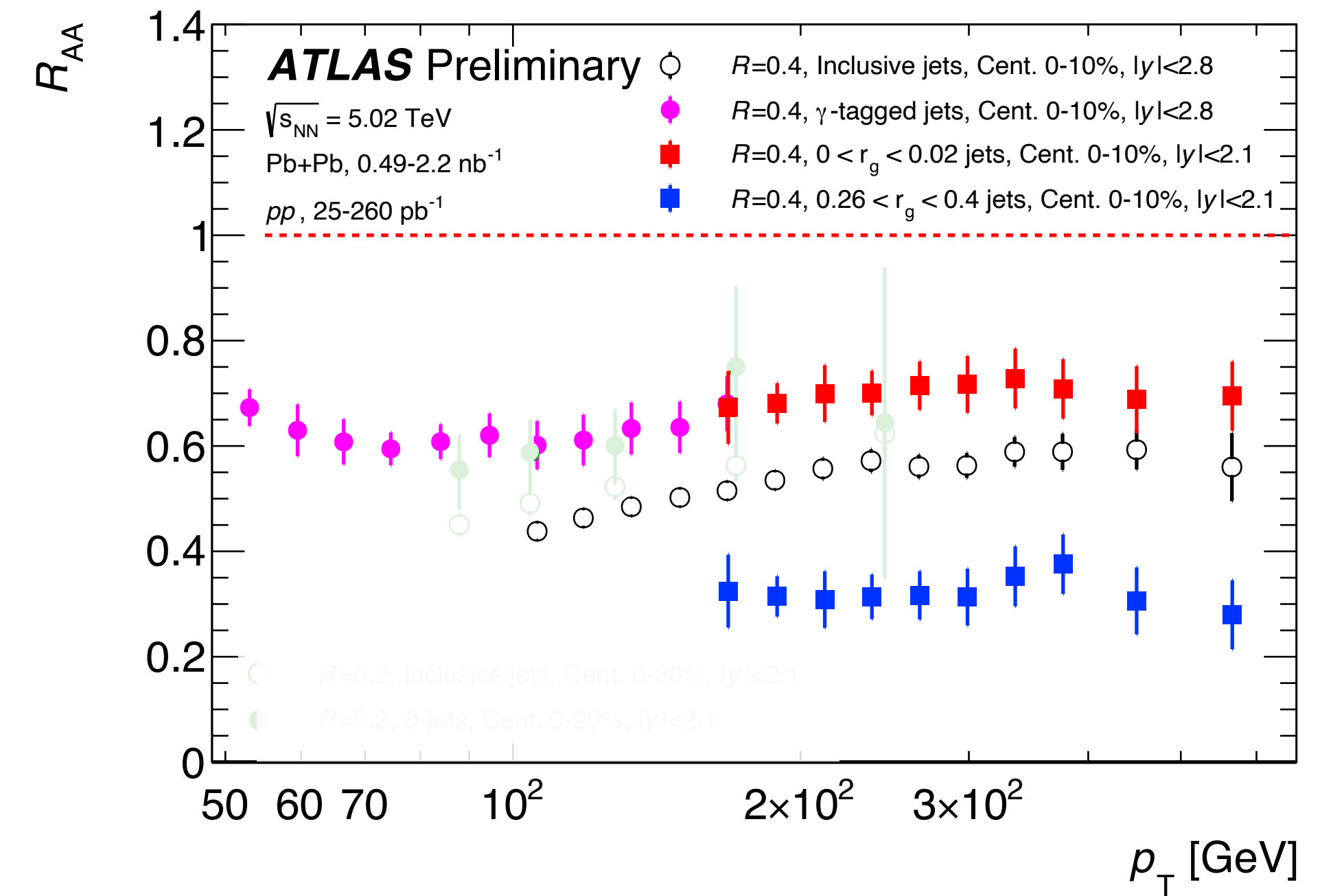
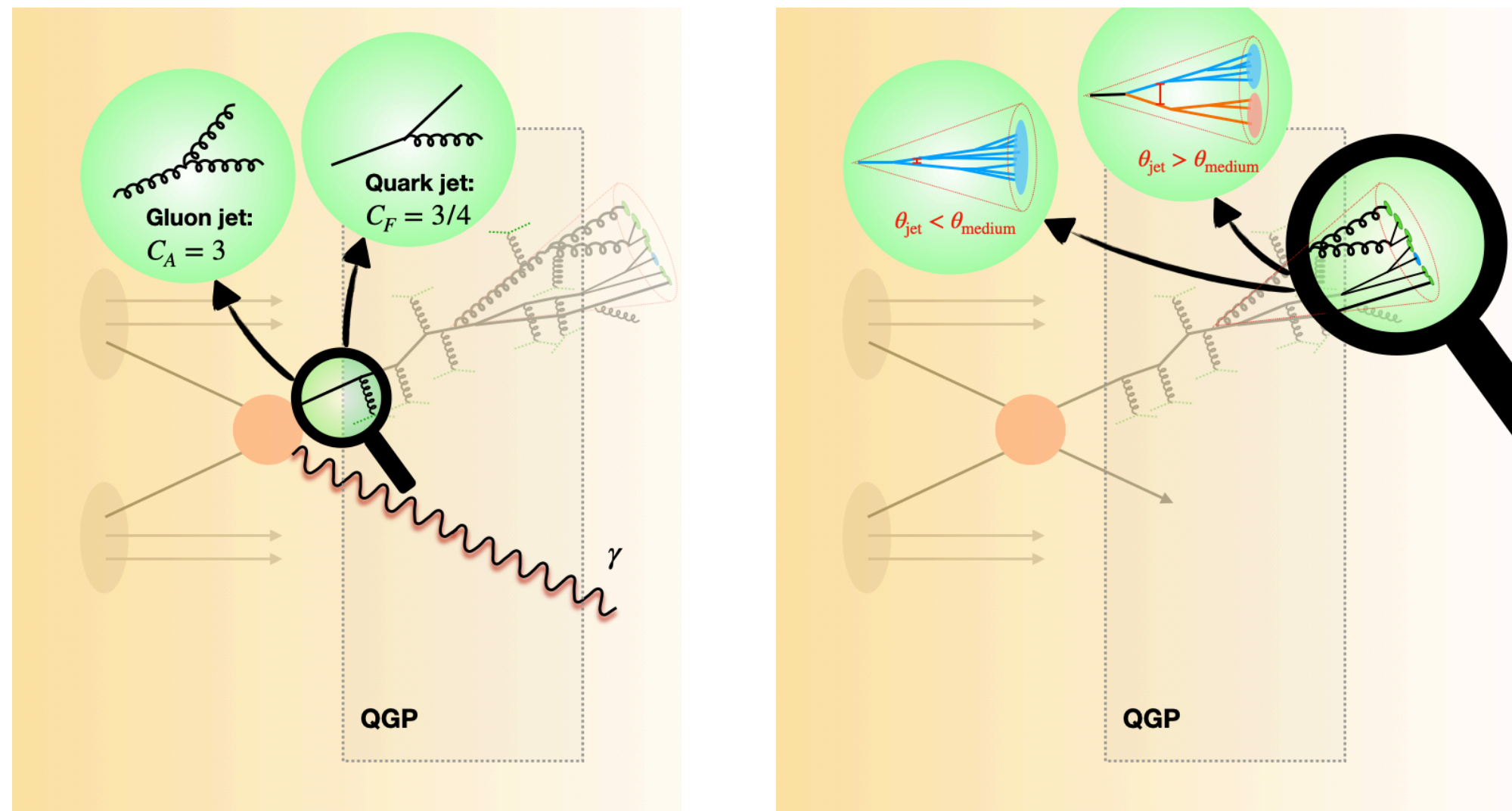


# Summary

Parton-QGP interaction  $\xrightarrow{\text{observation}}$  Jet quenching

QGP size/temperature  
QGP color coherence

Centrality dependence  
Color charge / substructure dependence



# Summary

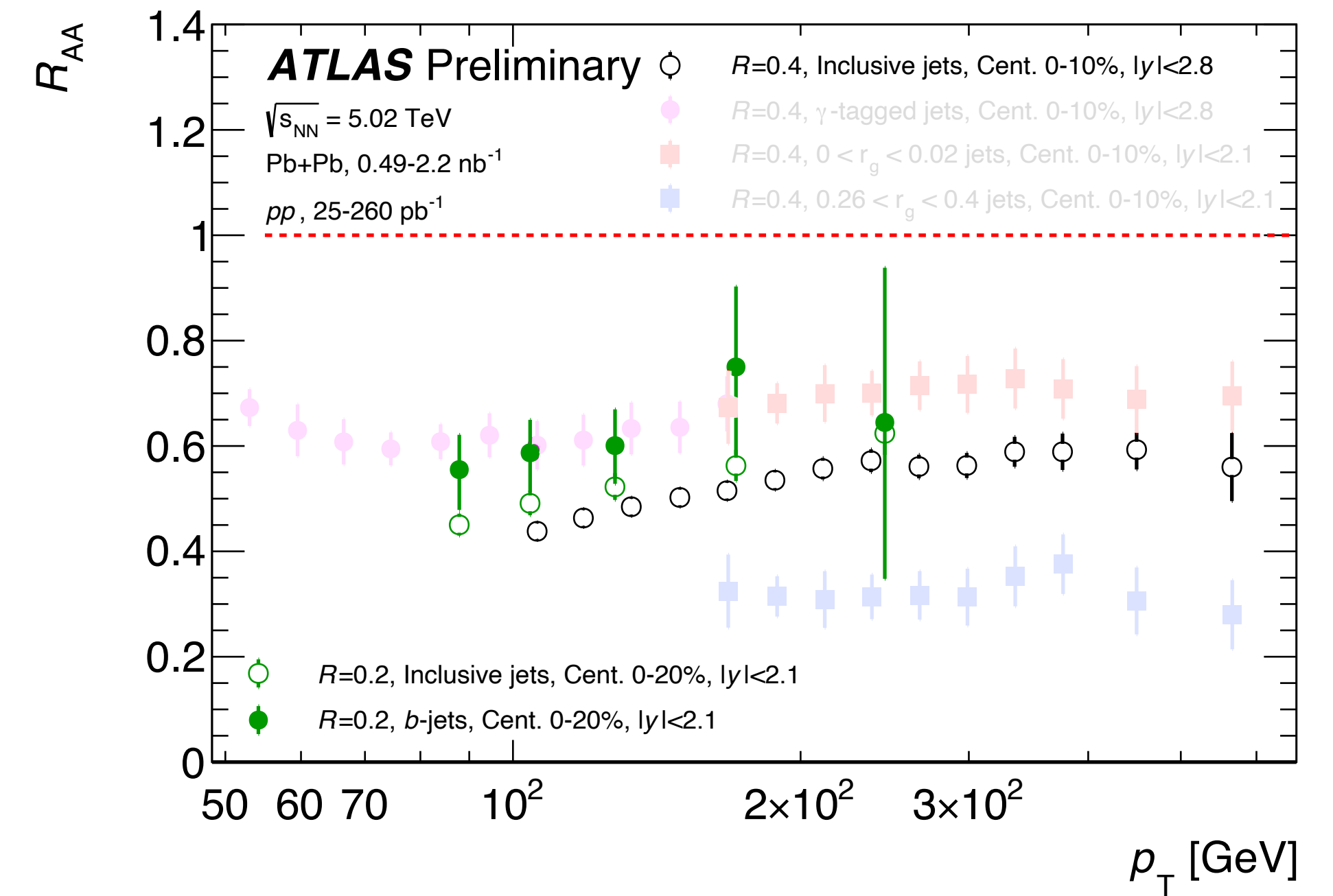
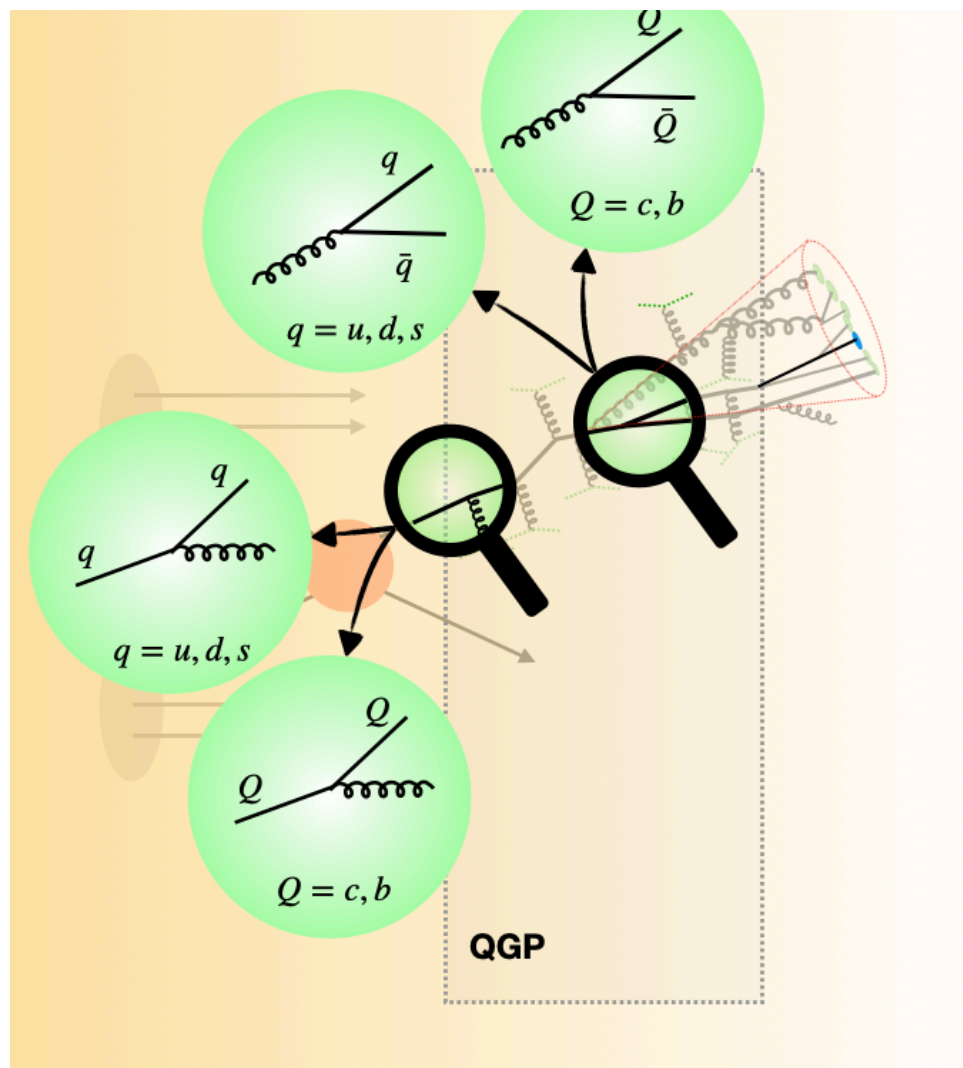
## Parton-QGP interaction

observation →

## Jet quenching

- QGP size/temperature
- QGP color coherence
- Parton mass

- Centrality dependence
- Color charge / substructure dependence
- Flavor dependence





# Summary

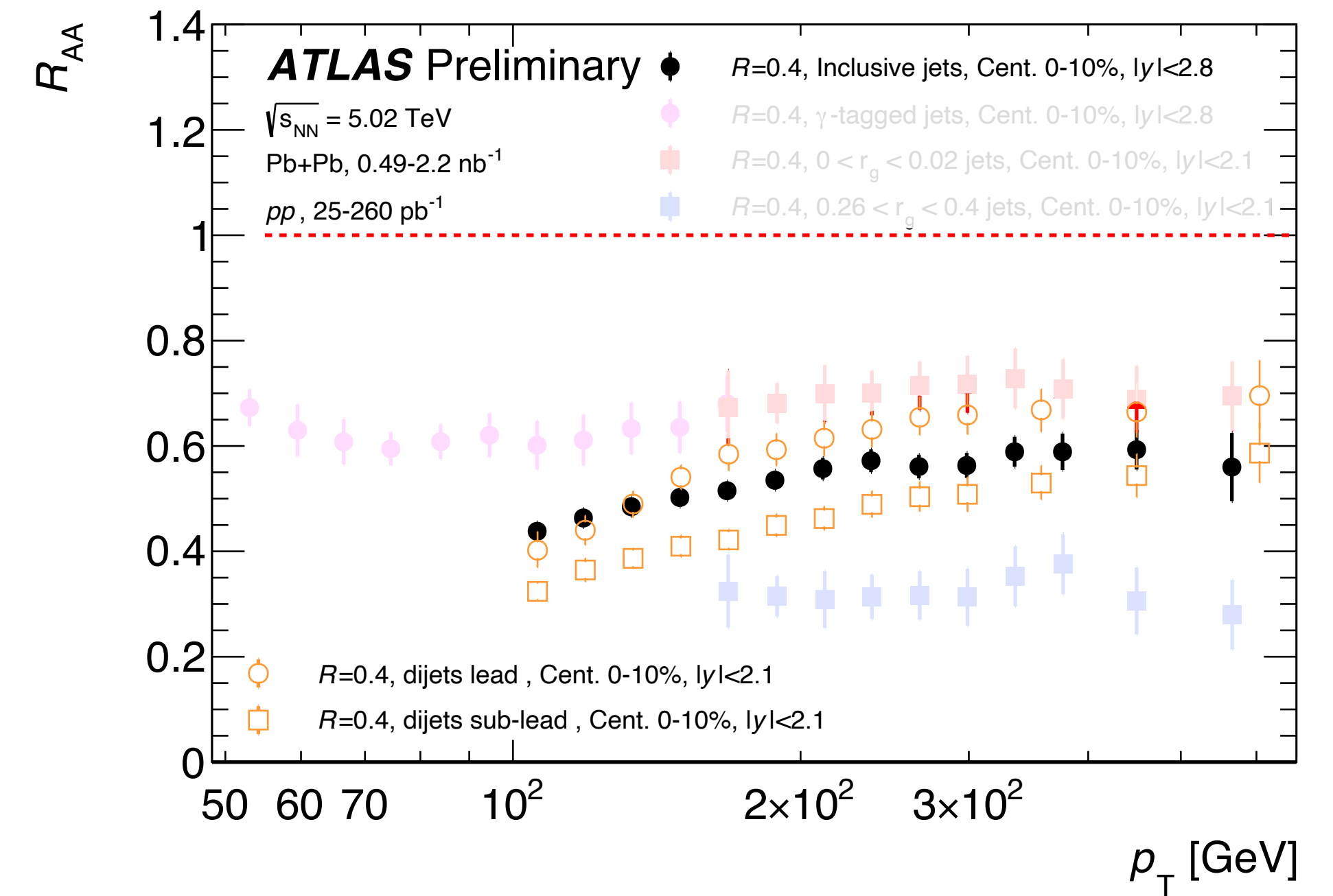
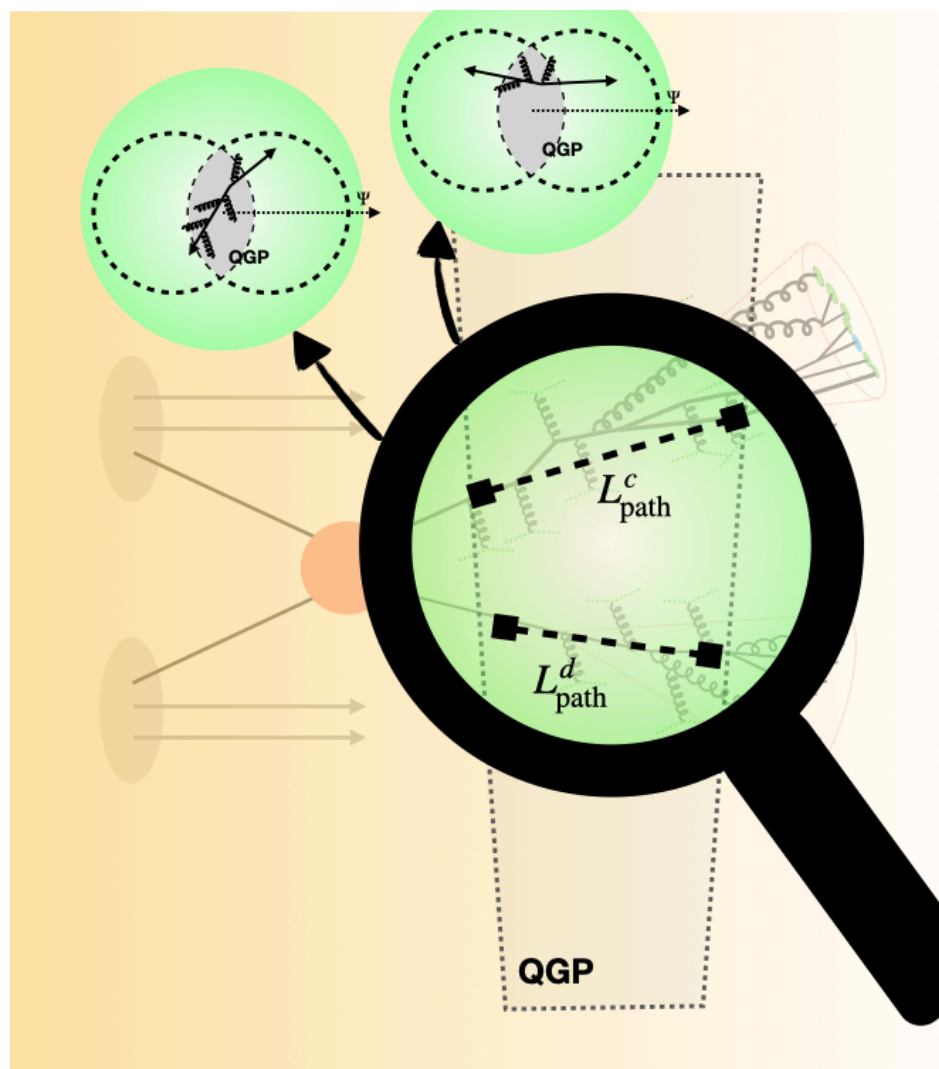
## Parton-QGP interaction

observation →

## Jet quenching

- QGP size/temperature
- QGP color coherence
- Parton mass
- Parton path-length in QGP

- Centrality dependence
- Color charge / substructure dependence
- Flavor dependence
- Inclusive jet  $v_n$  / dijet asymmetry





# Recent ATLAS HI jet measurements and beyond

- Jet  $v_n$ , [arXiv:2111.06606](https://arxiv.org/abs/2111.06606)
- Di-jet asymmetry, [arXiv:2205.00682](https://arxiv.org/abs/2205.00682) **Final**
- $b$ -jet suppression, [arXiv:2204.13530](https://arxiv.org/abs/2204.13530)
- $\gamma$ -tagged jet suppression, [ATLAS-CONF-2022-019](https://atlas.cern/ATLAS-CONF-2022-019) **Preliminary**
- Substructure dependence of jet suppression, [ATLAS-CONF-2022-026](https://atlas.cern/ATLAS-CONF-2022-026)

## Run3 outlook:

- More data → better statistical precision
- Improved jet reconstruction and calibration → better systematic precision
- New collision system (O-O) → explore jet quenching in smaller systems

see all ATLAS Heavy Ion results here: <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/HeavyIonsPublicResults>

**Backup Slides**