

# Precise predictions for the trilinear Higgs coupling in arbitrary models

predicting  $\kappa_\lambda$  in *any* model at the one-loop order

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# Outline

Motivation: theories and tools

$\lambda_{hhh}^{\text{BSM}}$  in any BSM model at the one-loop level

Outlook: beyond  $\kappa_\lambda$

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- > SM:  $\lambda_{hhh} h^3$  experimentally no (very well) known.
- > BSM case: deformation of the scalar potential possible!

$$V_{\text{BSM}}(h, \dots) = \frac{m_h^2}{2} h^2 + 3 \frac{m_h^2}{v} \kappa_\lambda^{\text{BSM}} h^3 + 3 \frac{m_h^2}{v^2} \kappa_{2\lambda} h^4 + \dots$$

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- > Idea: use  $-1 < \kappa_\lambda^{\text{exp.}} < 6.6$  [ATLAS-CONF-2021-052] to constraint "BSM" parameter space.
- > Example: "BSM"=THDM (talk by Johannes Braathen)  
→ **higher-order corrections important**

# Higher-order corrections to $\lambda_{hhh}^{\text{BSM}}$

A few studies of the trilinear Higgs self-couplings (i.e.  $\kappa_\lambda$ ) already exists:

- > SM [Kanemura et al. '04]
- > SM + singlet [Kanemura et al. '16]
- > THDM [Kanemura et al. '04][Basler et al. '17][Braathen et al. '19]
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- > Triplet extensions [Aoki et al. '18] [Chiang et al. '18]
- > MSSM [Brucherseifer et al. '13]
- > NMSSM [Dao et al. '13] [Dao et al. '15]

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However, there many more models to explore:

- > THDM-variants
- > singlet-variants
- > extended gauge sectors
- > extended fermion sectors
- > non-minimal SUSY, e.g. seesaw extensions, dirac gauginos, SplitSUSY, ...
- > (+ combinations)
- > ...



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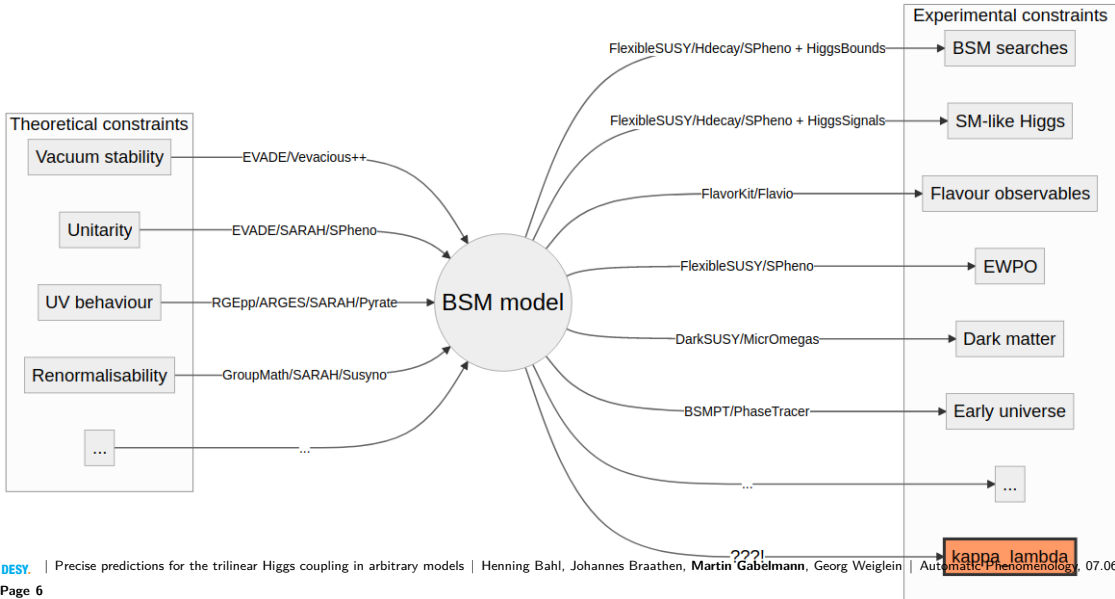
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→ **framework to calculate  $\lambda_{hhh}^{\text{BSM}}$  for large class of "BSM's"**



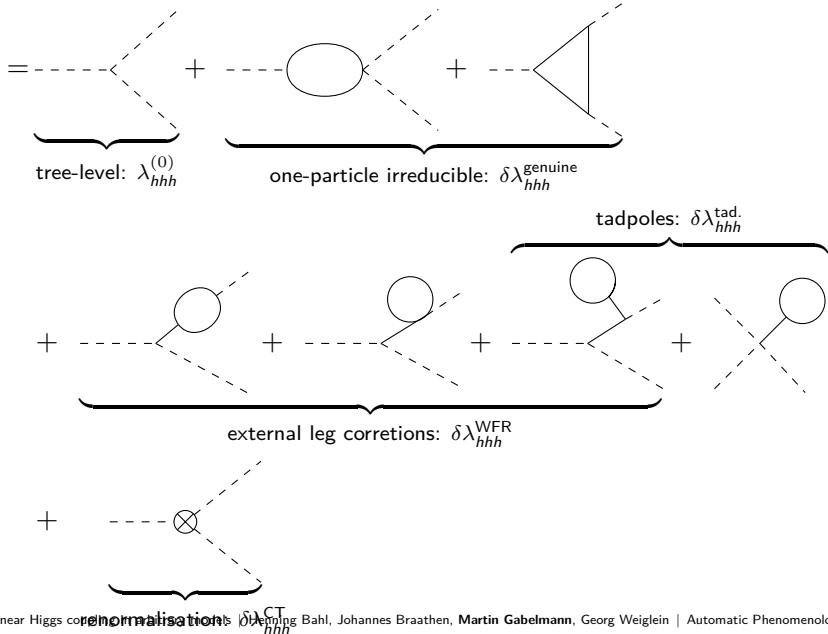
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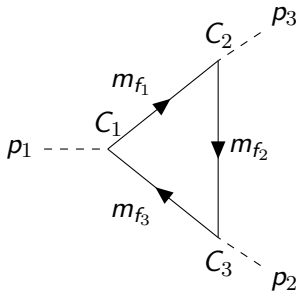
# Ingredients

$(\lambda_{hhh}^{\text{BSM}})^{\text{one-loop}}$



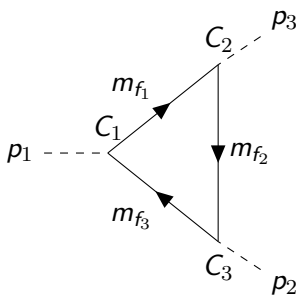
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Idea: compute *generic* diagrams i.e. assume most generic



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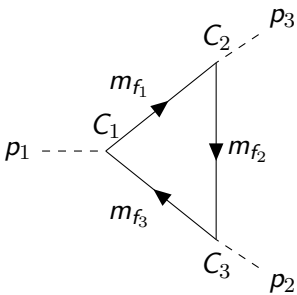
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- > couplings  $C_i = P_L C_i^L + P_R C_i^R$ ,  $P_{R/L} = \frac{1 \pm \gamma_5}{2}$
- > as well as loop-masses  $m_{f_i}$  and
- > external momenta  $p_i$ ,  $i = 1, 2, 3$ .

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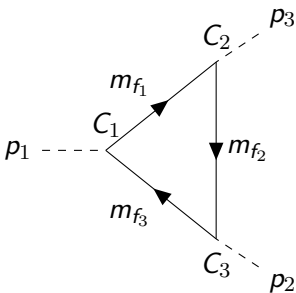
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$$\begin{aligned}
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 &C_2^L C_3^L m_{f_2} + C_2^L C_3^R m_{f_3})) + m_{f_1} \mathbf{C0}(p_2^2, p_3^2, p_1^2, m_1^2, m_3^2, m_2^2)((C_1^L C_2^L C_3^R + \\
 &C_1^R C_2^R C_3^L)(p_1^2 + p_2^2 - p_3^2) + 2(C_1^L C_2^L C_3^L + C_1^R C_2^R C_3^R)m_{f_2} m_{f_3} + \\
 &2m_{f_1}(C_1^L(C_2^L C_3^R m_{f_1} + C_2^R C_3^R m_{f_2} + C_2^R C_3^L m_{f_3}) + C_1^R(C_2^R C_3^L m_{f_1} + C_2^L C_3^L m_{f_2} + \\
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 \end{aligned}$$

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 \end{aligned}$$

> insert concrete BSM model

> evaluate with the help of  
COLLIER [Denner et al. '16]



## Another ingredient: Renormalisation of $\lambda_{hhh}$

> one-loop  $\rightarrow$  renormalisation of all parameters entering  $\lambda_{hhh}^{(0),\text{BSM}}$

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- > user's choice:
  - SM sector: fully OS **or**  $\overline{\text{MS}}/\overline{\text{DR}}$
  - BSM masses: OS **or**  $\overline{\text{MS}}/\overline{\text{DR}}$
  - additional couplings/vevs: most likely  $\overline{\text{MS}}$  **but also custom ren. conditions possible!**

couplings etc. (that can't be expressed in terms of masses)

$$\delta\lambda_{hhh}^{\text{CT}} = \sum_x \frac{\partial \lambda_{hhh}^{(0),\text{BSM}}}{\partial x} \delta x, \quad x = (m_h^{\text{SM}})^{\text{OS}/\overline{\text{MS}}}, (v^{\text{SM}})^{\text{OS}/\overline{\text{MS}}}, (m_{H_i}^{\text{BSM}})^{\text{OS}/\overline{\text{MS}}}, (\dots)^{\overline{\text{MS}}/\text{custom}}$$

## (Default) Renormalization choice of $(v^{\text{SM}})^{\text{OS}}$ and $(m_i^2)^{\text{OS}}$

>  $v^{\text{OS}} \equiv \frac{2M_W^{\text{OS}}}{e} \sqrt{1 - \frac{M_W^{2\text{OS}}}{M_Z^{2\text{OS}}}}$  with

•  $\delta^{(1)} M_V^{2\text{OS}} = \frac{\Pi_V^{(1),T}}{M_V^{2\text{OS}}}(p^2 = M_V^{2\text{OS}}), V = W, Z$

•  $\delta^{(1)} e^{\text{OS}} = \frac{1}{2} \dot{\Pi}_\gamma(p^2 = 0) + \text{sign}(\sin \theta_W) \frac{\sin \theta_W}{M_Z^2 \cos \theta_W} \Pi_{\gamma Z}(p^2 = 0)$

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**All bosonic one- & two-point functions and their derivatives for general QFTs are required for flexible OS renormalisation.**



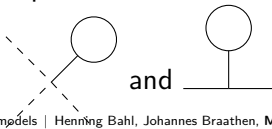
## Tadpole contributions to $\lambda_{hhh}$

- > In the SM: once  $\lambda_{hhh}$  is expressed in terms of *physical* input parameters, its result is independent of the treatment of the tadpoles:

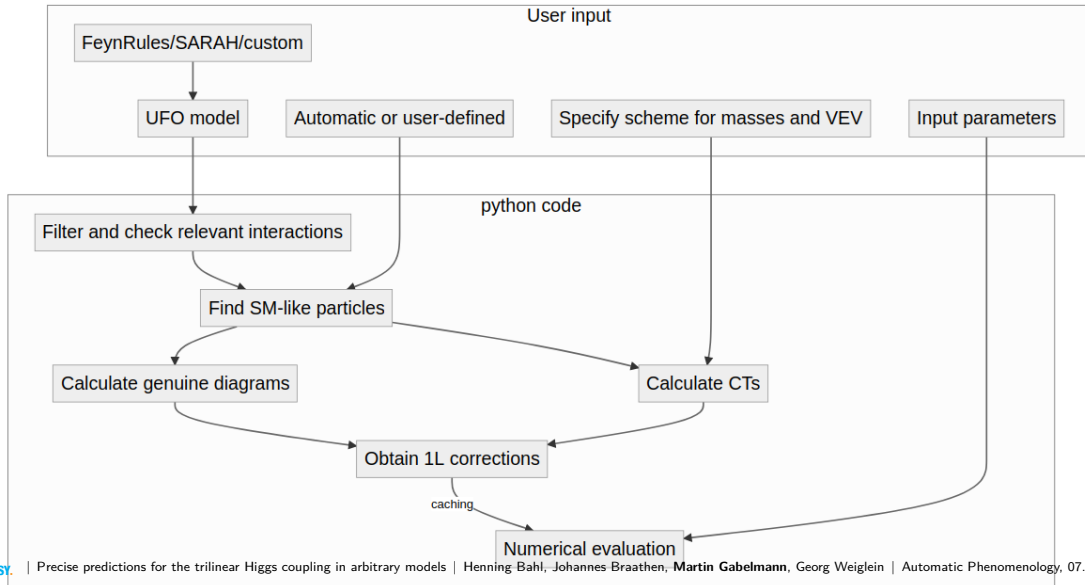
$$\delta^{(1)}\lambda_{hhh} \supset -\frac{3}{v^2}\delta^{(1)}t_{\text{finite}}$$

- > However: UFO models do (often) **not** contain the explicit dependence on the tree-level/one-loop/one-loop-CT tadpoles.
- > Thus: we choose the Fleischer-Jegerlehner treatment  $t^{\text{tree-level}} = 0$  and renormalize  $\delta^{(1)}t^{\text{CT}}|_{\text{finite}} = 0$  in the  $\overline{\text{MS}}$  scheme per default (can also turn-off automatic tadpoles and implement own scheme).

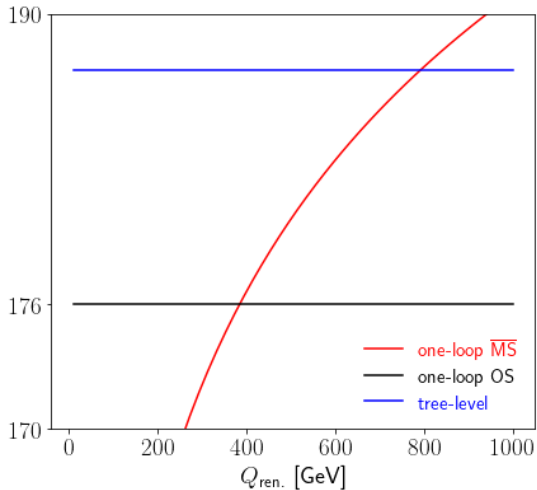
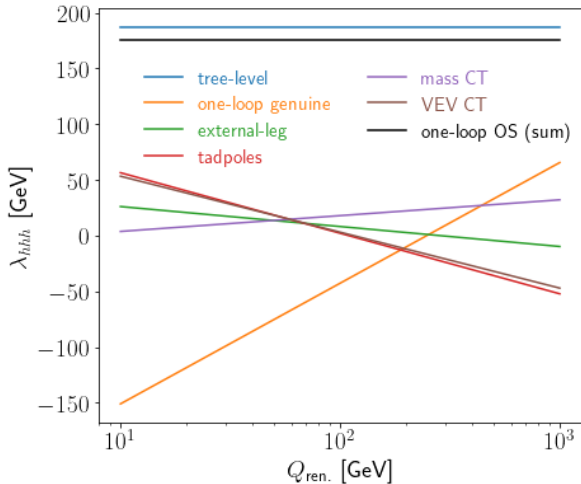
→ only need to take into account tadpole contributions to all two- and three-point functions:



# Workflow



# Simplest model to consider for cross-check: SM



Leading two-loop  $\mathcal{O}(\alpha_t(\alpha_t + \alpha_s))^{\text{OS}}$ :  $\mathcal{O}(+1.4\%)$  [Senaha '18] [Braathen et al. '19]

# Code demo

# Feature list (so far)

- > import/convert arbitrary UFO models
- > definition of renormalisation schemes

```
# schemes.yml
renormalization_schemes:
  OS:
    mass_counterterms:
      h1: OS
      h2: OS
    VEV_counterterm: OS
  MS:
    mass_counterterms:
      h1: MS
      h2: MS
    VEV_counterterm: MS
```

- > command line interface or python-library
- ```
from anyH3 import anyH3
myfancymodel = anyH3(
    'path/to/UFO/model',
    evaluation = 'numerical',
    scheme = 'OS')
result = myfancymodel.lambdahhh()
```
- > optional: full  $p^2$  dependence
  - > numerical / analytical /  $\LaTeX$  outputs
  - > restrict to certain topologies
  - > restrict to certain particles in the loop
  - > ...

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Outlook: beyond  $\kappa_\lambda$

# Electroweak precision observables

- >  $M_W$  (via  $\Delta r \approx \delta\rho$ )
- >  $\sin\theta_W^{\text{eff.}}$
- >  $\Gamma_Z$
- > ...

**Status:** reproduced full one-loop result for  $M_W/\delta\rho/\Delta r$  in the SM.

**To Do:** vertex + box diagrams for arbitrary QFTs (small in many extended Higgs sectors (+ alignment)); incorporation of higher-order SM results.

# W Mass

> start with HO corrections to muon decay:  $M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi\alpha_{em}}{\sqrt{2}G_F} [1 + \Delta r]$

> and solve for:  $M_W^2 = M_Z^2 \left[ \frac{1}{2} + \sqrt{\frac{1}{4} - \frac{\pi\alpha_{em}}{\sqrt{2}G_F M_Z^2} (1 + \Delta r)} \right]$

> with:  $\Delta r^{(1)} = 2\delta^{(1)} e + \frac{\Pi_W^{(1),T}(0) - \delta^{(1)} M_W^2}{M_W^2} - \frac{\delta^{(1)} \sin^2 \theta_W}{\sin^2 \theta_W} + \delta_{\text{vertex+box}}$

> and:  $\frac{\delta^{(1)} \sin^2 \theta_W}{\sin^2 \theta_W} = \frac{\cos^2 \theta_W}{\sin^2 \theta_W} \left( \frac{\Pi_W^{(1),T}(M_W^2)}{M_W^2} - \frac{\Pi_W^{(1),T}(M_Z^2)}{M_Z^2} \right)$

It's all there but:

>  $\delta_{\text{vertex+box}}^{\text{SM}} = -\frac{2 \text{sign}(\sin \theta_W)}{\cos \theta_W \sin \theta_W M_Z^2} \Pi_{Z\gamma}(p^2 = 0) + \frac{\alpha_{QED}}{4\pi \sin^2 \theta_W} \left( 6 + \frac{7-4 \sin^2 \theta_W}{2 \sin^2 \theta_W} \right) \log(\cos^2 \theta_W)$

>  $\delta_{\text{vertex+box}}^{\text{BSM}} = \text{needs to be implemented}$

However:

> in many models  $\Delta r \supset \frac{\delta \sin^2 \theta_W}{\sin \theta_W} \approx \delta\rho$  is the dominant effect!



## Other coupling modifiers

Question: how do you know that only  $\kappa_\lambda \neq 1$  while e.g.  $\kappa_t = 1$ ?

$$\mathcal{L}_{HH} = -\frac{3m_H^2}{v}(1 + \kappa_\lambda)h^3 - \left[ \frac{m_t}{v}(1 + \kappa_t)\bar{t}_L t_R h + \text{h.c.} \right] + \dots$$

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  - at the full one-loop order
  - with arbitrary choice of renormalization schemes
- > reproduced SM results
- > uses UFO input (generate with SARAH, FeynRules or use a custom one)
- > analytical results; fast numerical results; SM:  $\mathcal{O}(0.2\text{ s})$ , MSSM:  $\mathcal{O}(0.5\text{ s})$
- > many models already implemented:  
SM, SM+**singlet**, **THDM**, **NTHDM**, various **triplet extensions**, **MSSM** and **NMSSM**

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Future ideas / todos:

- > more models / cross-checks
- > go beyond one-loop
- > non-SM self-couplings (e.g.  $\kappa\lambda_{hhH}$ )
- >  $\kappa_t$  and  $\kappa_{tt}$