

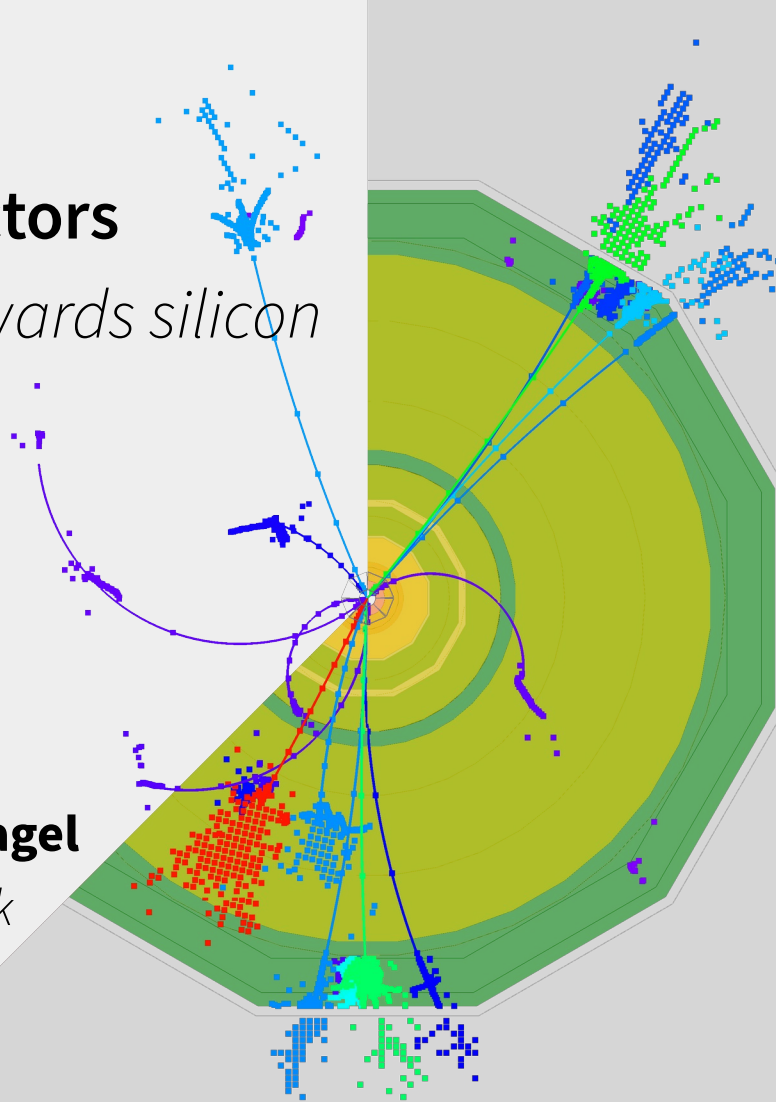


<https://www.desy.de/>

Modern Tracking Detectors

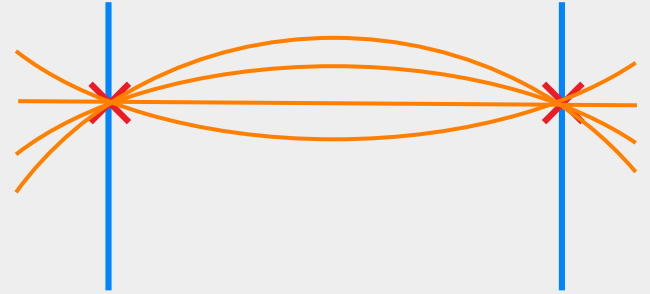
... with a strong bias towards silicon

Simon Spannagel
HighRR Lecture Week
29 June 2022



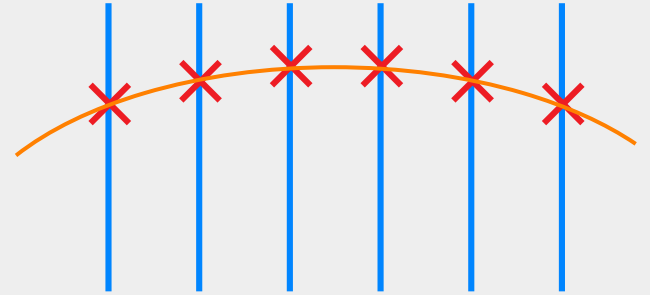
Motivation

- Efficient & precise tracking of relativistic particles always requires several measurement points



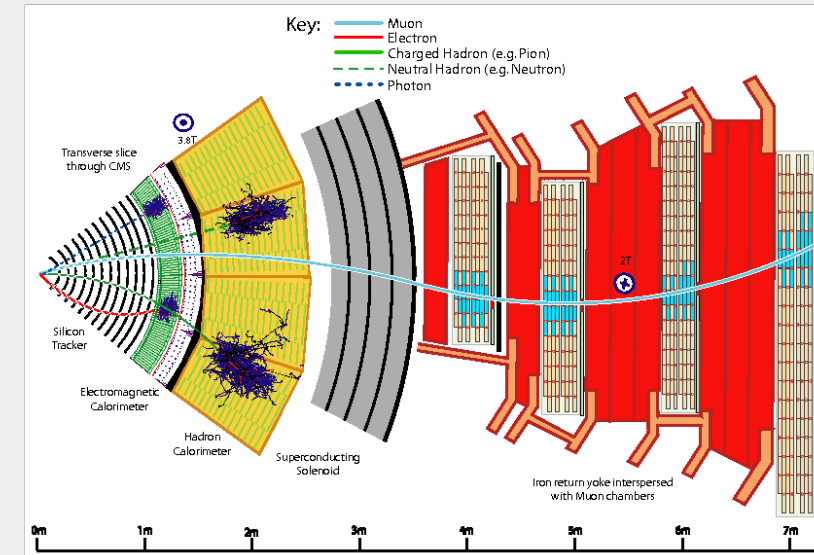
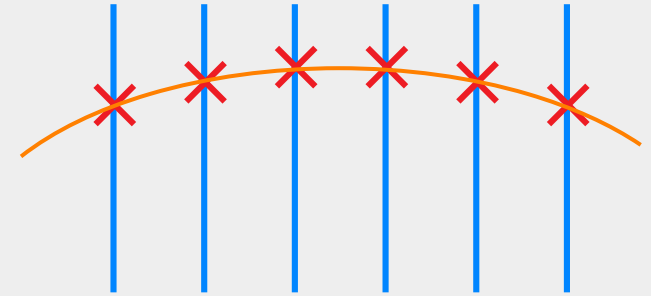
Motivation

- Efficient & precise tracking of relativistic particles always requires several measurement points



Motivation

- Efficient & precise tracking of relativistic particles always requires several measurement points
 - TPC: one data point per pad
 - Silicon tracker: one data point per layer
- Tracking detector systems:
 - Momentum determination
 - Vertex reconstruction
 - Pile-up mitigation, background suppression
 - Secondary vertices → B-tagging
- Supports identification of particle types



Momentum Reconstruction

Momentum Reconstruction

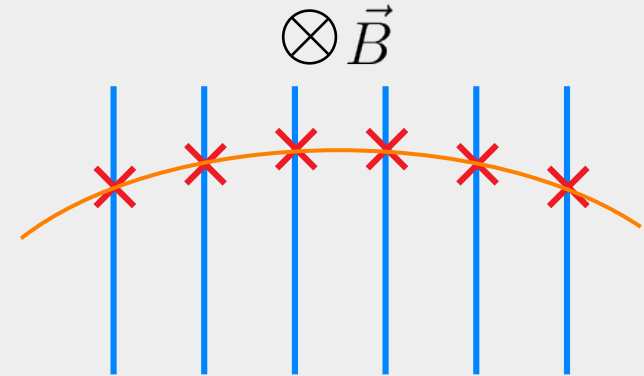
- Relativistic particles in magnetic field:

- $\vec{p} \perp \vec{B}$ circular track
- $\vec{p} \parallel \vec{B}$ straight track
- Else helical track

- Lorentz force and centripetal force in equilibrium:

$$\begin{aligned}
 |\vec{F}_L| &= |\vec{F}_Z| \\
 |q\vec{v} \times \vec{B}| &= \frac{m\vec{v}^2}{R} \\
 p_T &= qBR
 \end{aligned}
 \quad \longrightarrow$$

- Measurement of the radius of curved tracks yields *transverse momentum* p_T



Simple approximation

$$p_T = 0.3|z|BR$$

$$[p_T] = \text{GeV}/c$$

with $[R] = \text{m}$

$$[B] = \text{T}$$

$$z = q/e$$

Momentum Measurement

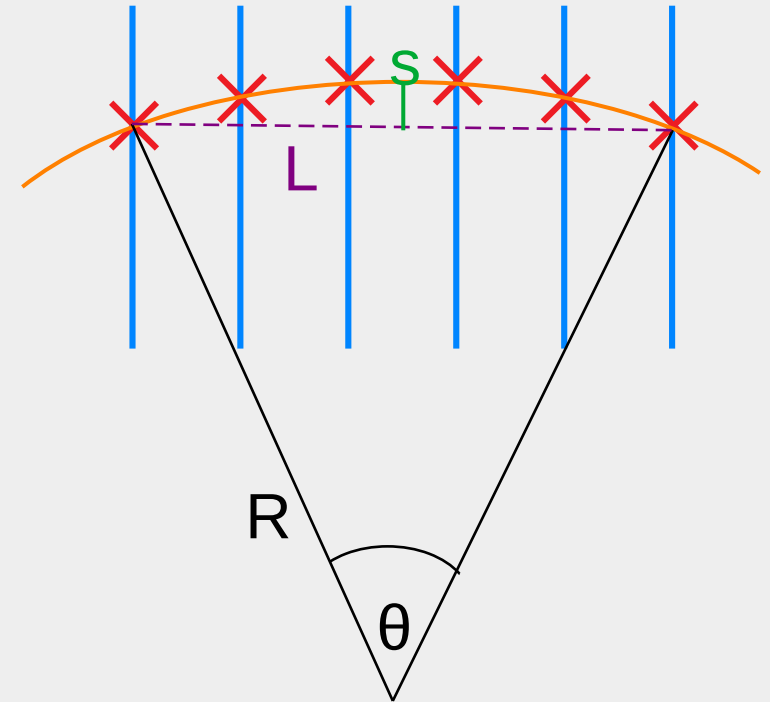
- Radius can be determined by measurement of *sagitta* length s at track length L

$$s \approx \frac{L^2}{8R}$$

- Uncertainty of momentum proportional to uncertainty of sagitta measurement:

$$p = qBR = \frac{8qB}{L^2}s \quad \Rightarrow \quad \sigma_p \propto \sigma_s$$

- Uncertainties of sagitta measurement:
 - Spatial detector resolution
 - Multiple Coulomb scattering



Momentum Measurement – Uncertainty

Contribution from detector resolution

- ✓ Spatial detector resolution σ_m
- ✓ Track length L
- ✓ Magnetic field B
- ✓ Number of measurements N
- ✗ Transverse momentum p_T

Gluckstern formula (for $N > 10$)

$$\frac{\sigma_{p_T}}{p_T} = \frac{p_T}{0.3|z|} \frac{1}{L^2 B} \sqrt{\frac{720}{N+4}} \sigma_m$$

➔ Large tracking detector inside a strong magnetic field essential!

Momentum Measurement – Uncertainty

Contribution from multiple Coulomb Scattering

- Scattering: $\theta_0 \propto \frac{1}{p} \sqrt{\frac{l}{X_0}}$
- ✓ Track length L
- ✓ Magnetic field B
- ✓ Number of measurements N
- ✗ Material budget $\epsilon=l/X_0$

Momentum Uncertainty (Scattering)

$$\frac{\sigma_{p_T}}{p_T} = \frac{p_T}{BL \sin \theta} \frac{\theta_{0,N}}{\sin \theta} \sqrt{C_N}$$

$$\theta = \angle(\vec{p}, \vec{B}) \quad \theta_{0,N} = \sum_N \theta_{0,i}$$

$$C_N = \frac{2N(2N - 1)}{3(N - 1)^2}$$

Momentum Measurement – Uncertainty

Contribution from multiple Coulomb Scattering

- Scattering: $\theta_0 \propto \frac{1}{p} \sqrt{\frac{l}{X_0}}$
- ✓ Track length L
- ✓ Magnetic field B
- ✓ Number of measurements N
- ✗ Material budget $\epsilon=l/X_0$

$$\frac{\sigma_{p_T}}{p_T}(p_T) \propto p_T \theta_{0,N} = \text{const.}$$

➔ Constant as a function of p_T

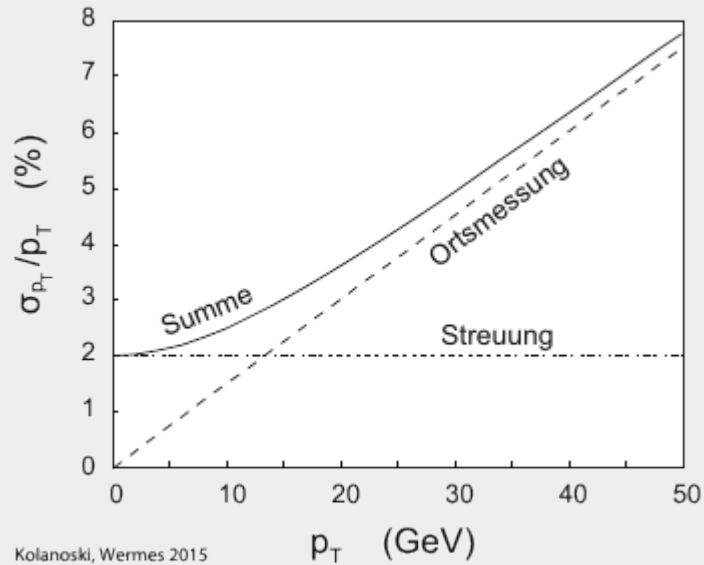
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$$C_N = \frac{2N(2N-1)}{3(N-1)^2}$$

Momentum Measurement – Uncertainty



Total Uncertainty

$$\frac{\sigma_{p_T}}{p_T} = \sqrt{\left(\frac{\sigma_{p_T}}{p_T}\right)_{\text{meas}}^2 + \left(\frac{\sigma_{p_T}}{p_T}\right)_{\text{scat}}^2}$$

- Resolution: linear dependency on p_T
- Scattering: constant as a function of p_T

Comparison

- Calorimeters: $\frac{\sigma_E}{E} \propto 1/\sqrt{E}$
- Tracking: $\frac{\sigma_{p_T}}{p_T} \propto p_T$
 \sim

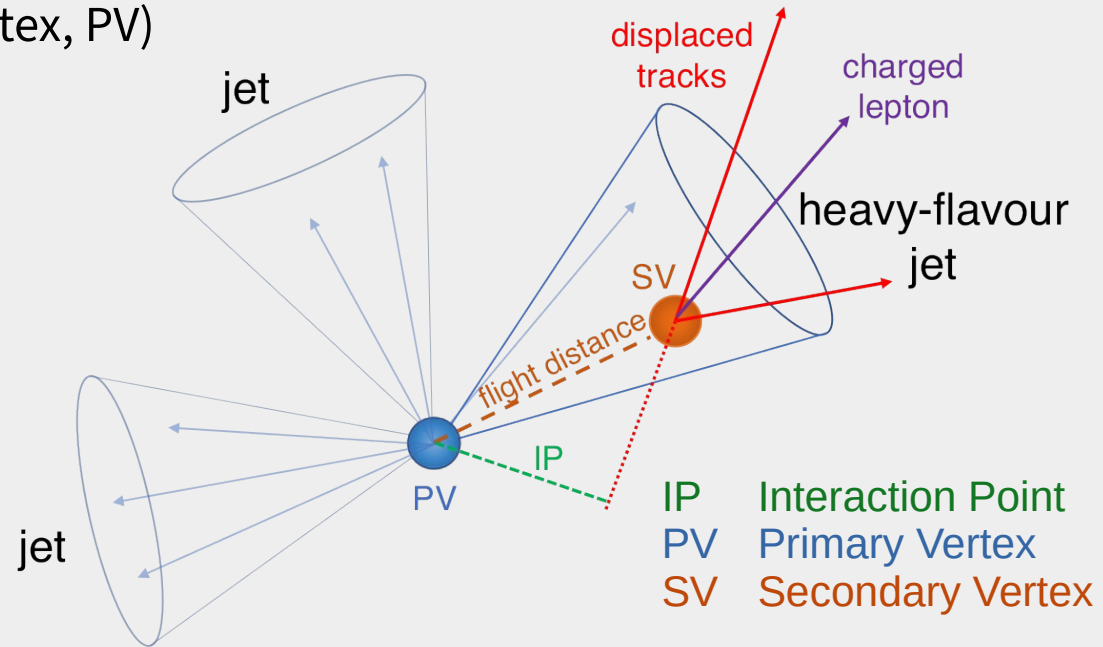
Vertex Reconstruction



Secondary Vertices – b-Tagging

- Not all particles decay immediately – example: **b-quarks**:

- Produced at collision (primary vertex, PV)
- Propagation during finite lifetime
- Decay → secondary vertex (SV)
- Lifetime of b-quarks: $\mathcal{O}(10^{-12} \text{ s})$
- Flight distance: $\mathcal{O}(100 \text{ } \mu\text{m})$
- Can be resolved with modern tracking detectors



- Indication of SV hints at b-quarks, used for e.g. b-Tagging

Secondary Vertices – b-Tagging

- Observable for tracking detectors:
Impact parameter d_0 of individual tracks w.r.t. main jet
- Resolution: $\sigma_{d_0} = \sigma_{d_0}(\text{meas}) \oplus \sigma_{d_0}(\text{scat})$

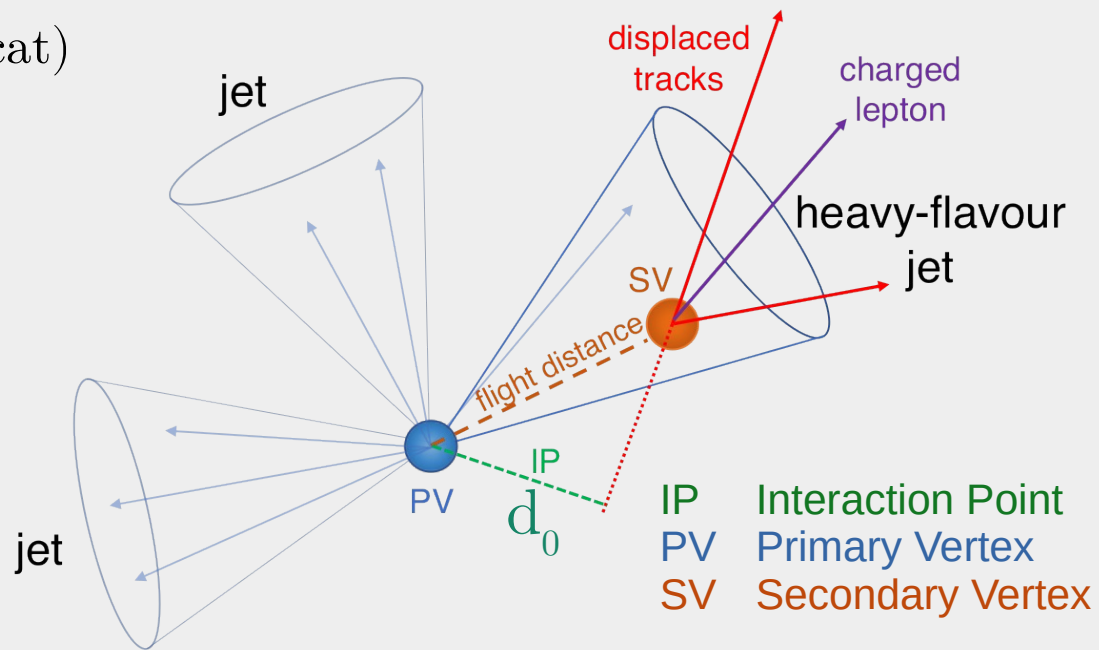
Example: ATLAS Pixel Detector

2008: 3 Layers, 10 μm resolution,
Radii: 5.05 cm, 8.85 cm, 12.25 cm.

$\rightarrow \sigma_{d_0} = 18.1 \mu\text{m} \oplus 142 \mu\text{m}/p$ (GeV/c)

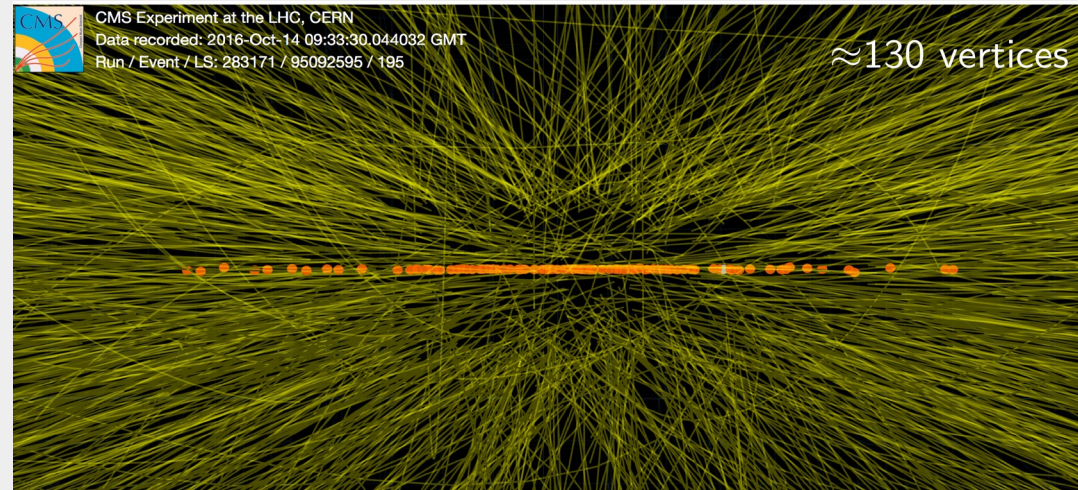
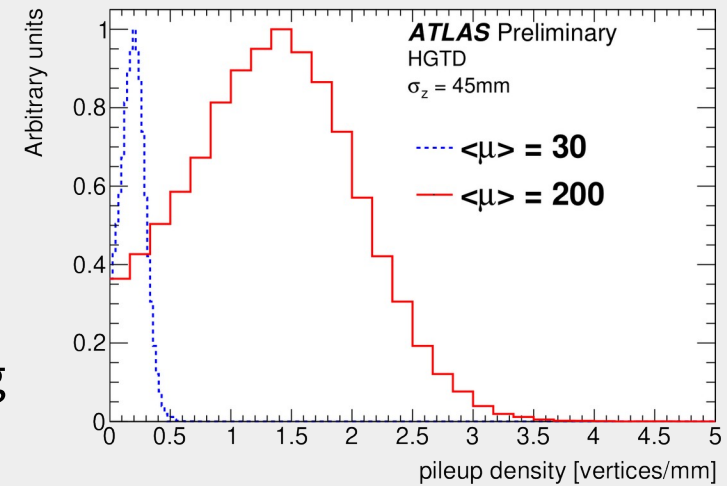
2014: Additional “B-Layer” / IBL @ 3.3 cm

$\rightarrow \sigma_{d_0} = 12.1 \mu\text{m} \oplus 93 \mu\text{m}/p$ (GeV/c)



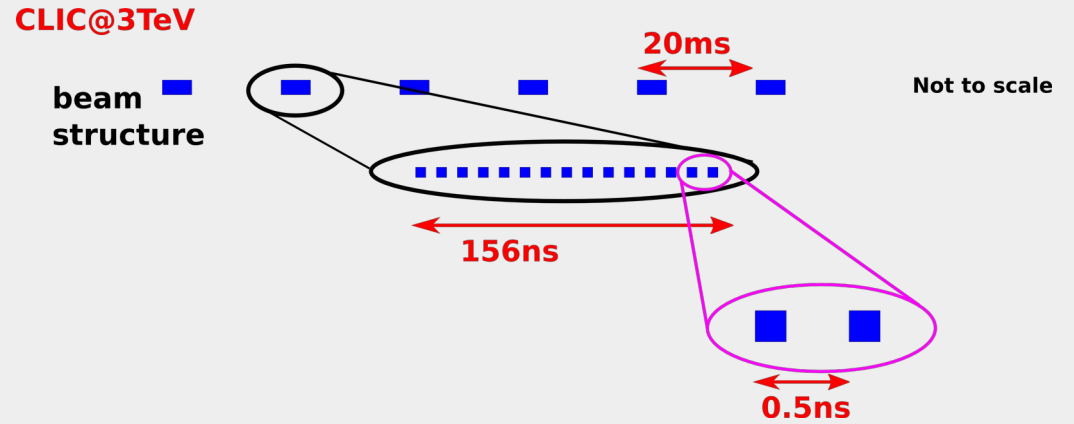
Timing – Background & Pile-Up

- Collisions become increasingly complex
 - Hadron colliders: Pile-Up
 - Multiple hard scattering collisions per bunch crossing
 - Lepton colliders: background
 - Often operate in bunch trains, need to separate out hard scattering
- Use time structure of collisions!
- Different approaches
 - Include timing in tracking detect.
 - Dedicated timing detectors



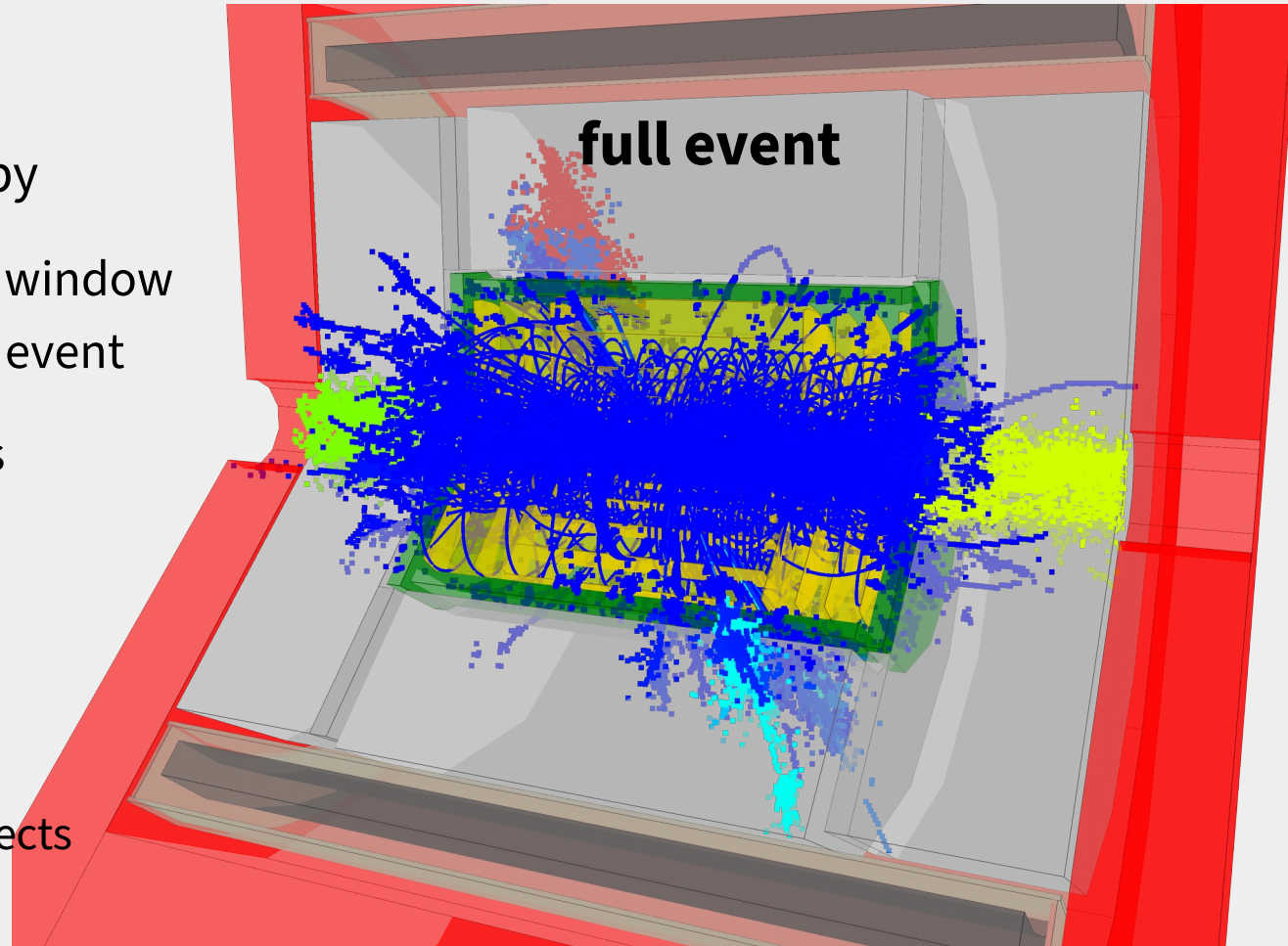
Example – CLIC @ 3 TeV: Experimental Conditions

- CLIC operates in bunch trains, repetition rate of 50 Hz
 - 312 bunches within train
 - Bunch separation by 0.5 ns
- Bunch separation & cross-section of background events drive timing requirements for detector
 - 1 ns time resolution for calorimeters
 - 5 ns single-hit resolution for vertex/tracking detectors



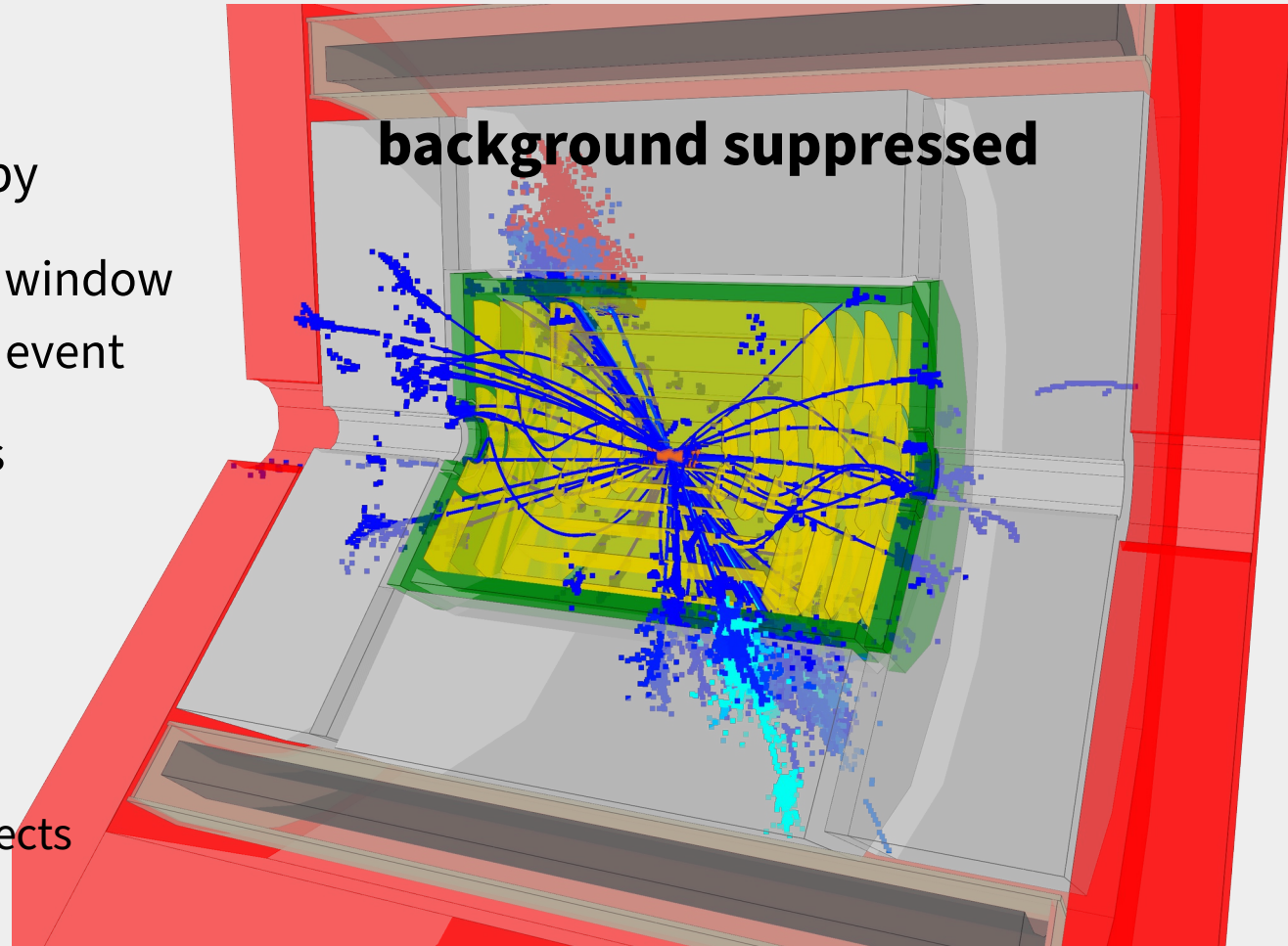
Example – CLIC @ **3 TeV**: Background suppression

- Fully-hadronic $t\bar{t}$ event
- Background suppression by
 - Defining reconstruction window
10 ns before, 30 ns after event
 - Building physics objects
 - Suppression via
 - Timing requirements
 - Particle type and p_T
 - Retaining high- p_T objects



Example – CLIC @ 3 TeV: Background suppression

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Very Short Recap

Gaseous Detectors vs. Semiconductor Detectors

Gaseous Detectors

Working principle

- Ions & electrons generated via ionization, ~ 30 eV ion. energy
- Electrons and ions drift in an electric field, current is induced by the drift of charge carriers
- Gas amplification for increase of signal

Types (selection)

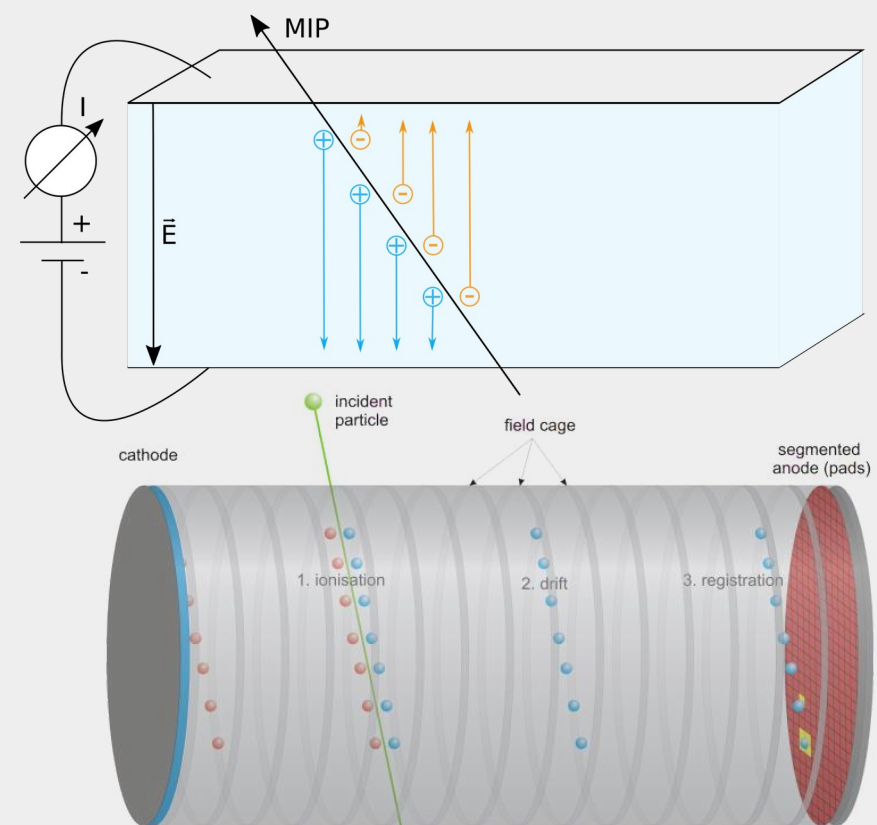
- Multi-wire proportional chamber
- Micro pattern gas detectors (MicroMegas, *GEM*)
- Drift Chamber, Time Projection Chamber



Very lightweight
Few issues with rad. damage
Large volume instrumentation

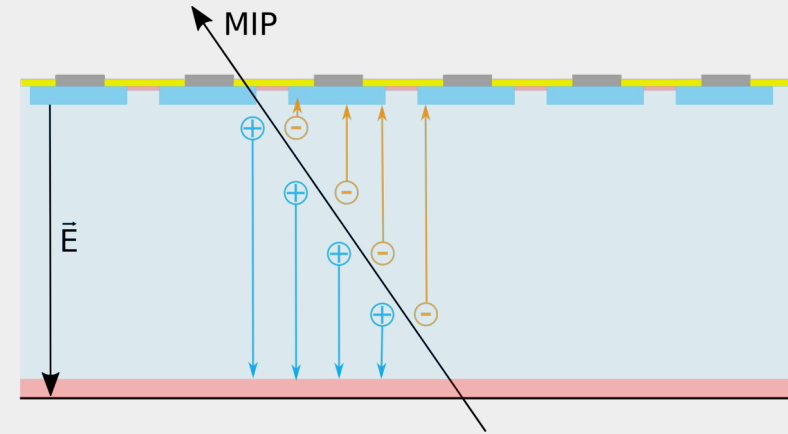


Relatively high ionization energy
Limited rate capabilities
Mechanical fragility (for some types)



Semiconductor Detectors

- Oversimplified: "Semiconductor ionization chamber"
 - Electron-hole pairs created from ionization energy, ~ few eV
 - Electrons/holes propagate in an electric field
 - Current is induced by drift of the charge carriers
- ! Technically very different from gaseous detectors!
→ Using established IC industry processes for production



Low signal generation energy
Nano-scale structures via lithography
Capable of running in high fluxes



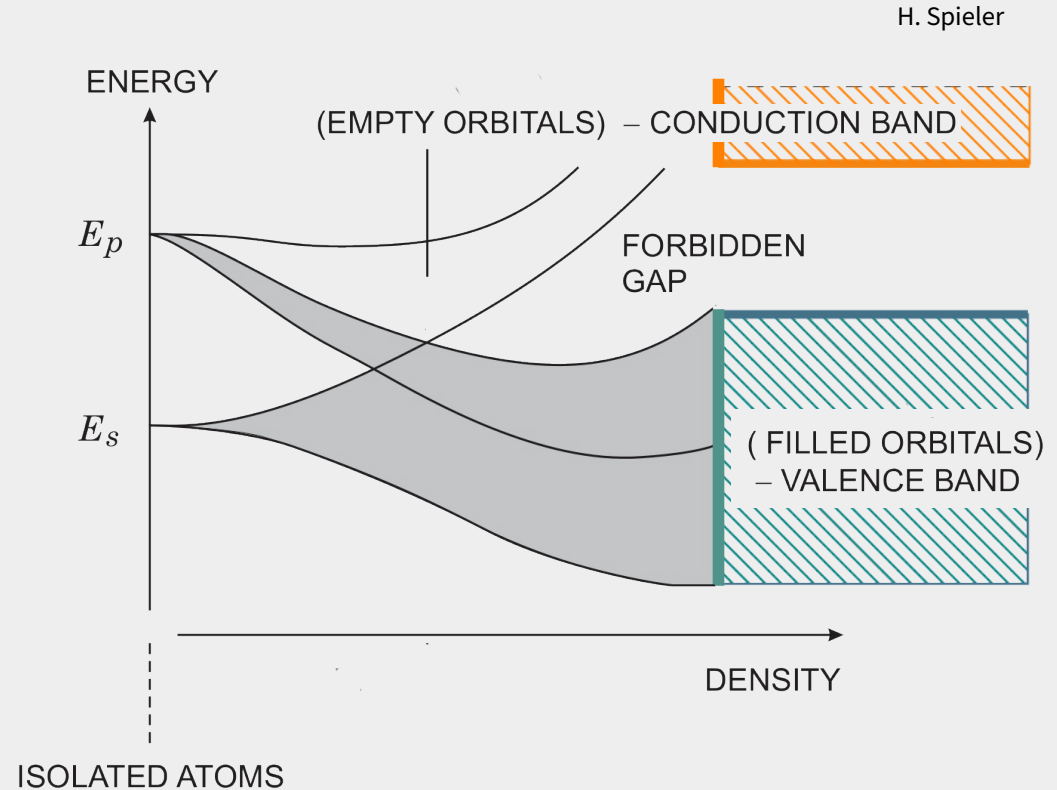
Expensive for large area instrumentation
Solid material placed in particle path
Often complex module / detector setup

Short Recap

Semiconductors, Bandgap, Doping

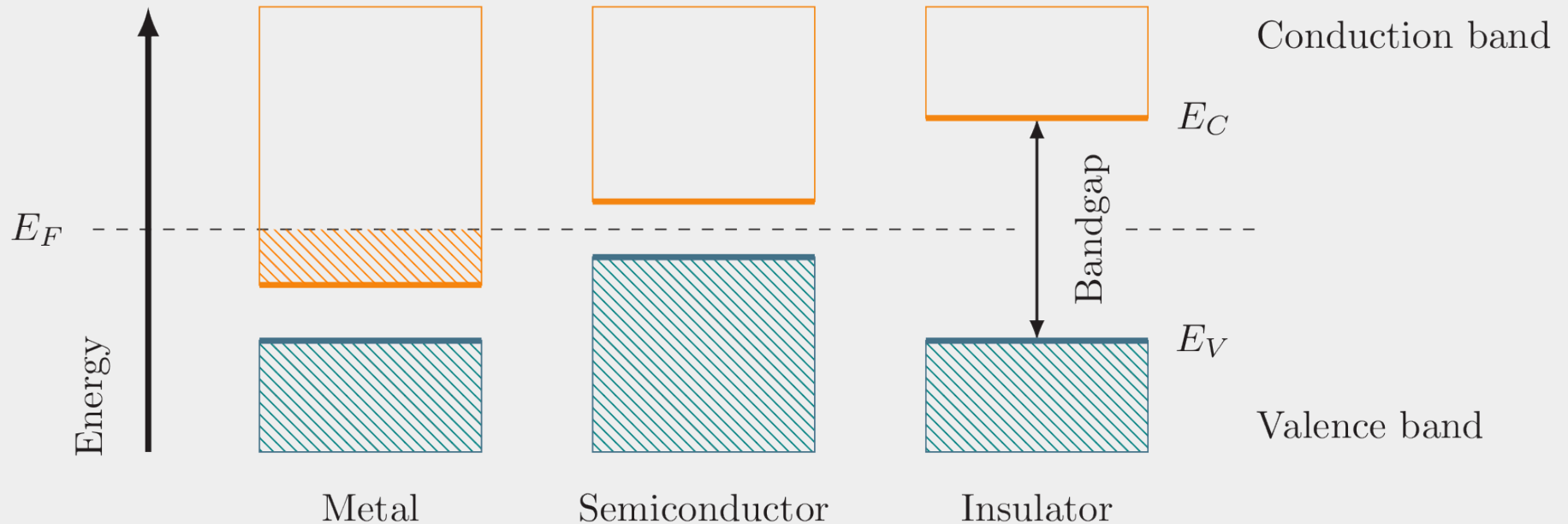
The Energy Band Model

- Atom: discrete energy levels, orbitals
- Crystal lattice: levels smear out
- Formation of energy bands
 - Valence band – (last) fully filled
 - Band gap
 - Conduction band



Metals, Insulators, Semiconductors

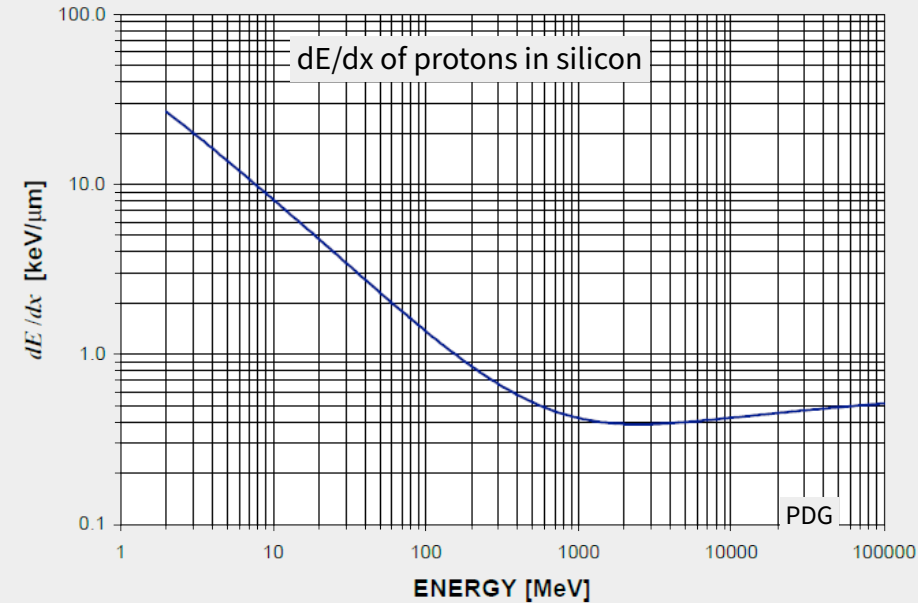
bandgap silicon:
 $E_G \approx 1.12 \text{ eV (300K)}$



Detecting a Particle with Intrinsic Silicon

- Silicon sensor: $A = 1 \text{ cm}^2$ and $d = 300 \text{ }\mu\text{m}$
- Signal of MIP:
 - Mean ionization: $E_0 = 3.6 \text{ eV}$ (silicon)
 - Mean energy loss: $dE/dx = 3.9 \text{ MeV/cm}$

$$\begin{aligned} \frac{dE}{dx} \cdot \frac{d}{E_0} &= 3.9 \cdot 10^6 \text{ eV/cm} \cdot 0.03 \text{ cm} / 3.6 \text{ eV} \\ &\approx 3 \cdot 10^4 \text{ e/h pairs} \end{aligned}$$



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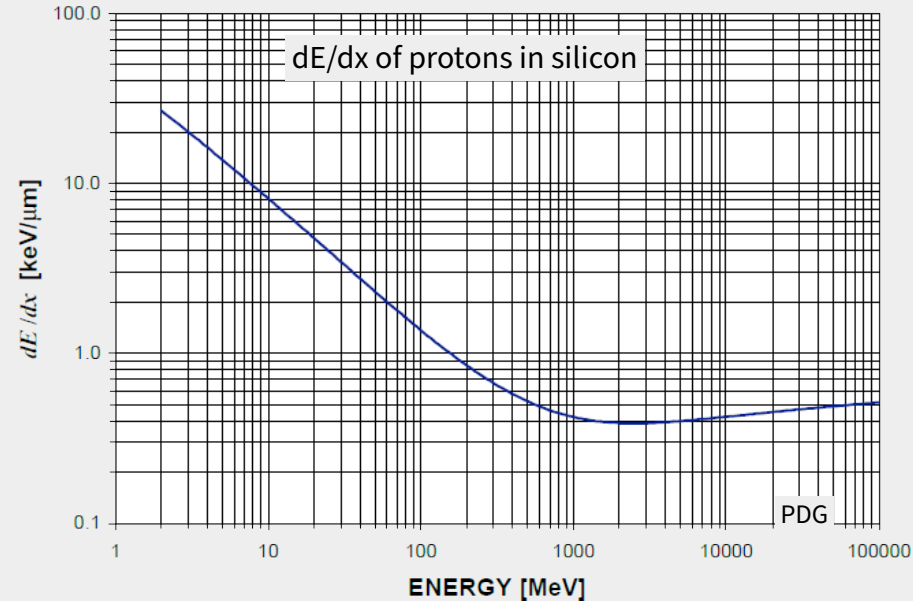
$$\approx 3 \cdot 10^4 \text{ e/h pairs}$$

- Thermally excited charge carriers in silicon: $n_i = 1.45 \times 10^{10} \text{ cm}^{-3}$ (at 300K)

$$n_i \cdot d \cdot A = 1.45 \cdot 10^{10} \text{ cm}^{-3} \cdot 0.03 \text{ cm} \cdot 1 \text{ cm}^2$$

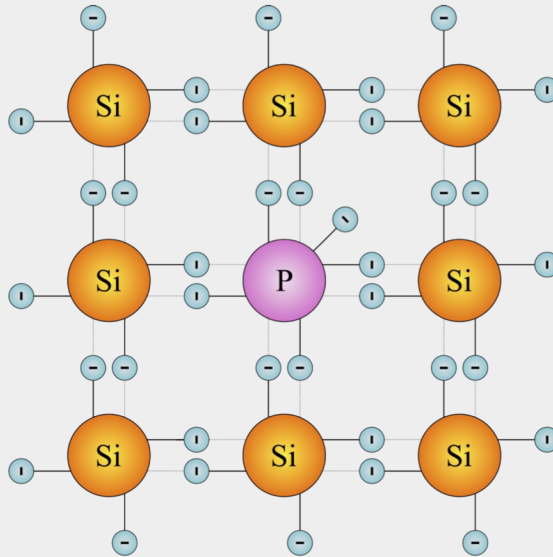
$$\approx 4 \cdot 10^8 \text{ e/h pairs}$$

density silicon:
 $N \approx 10^{22} \text{ atoms/cm}^3$

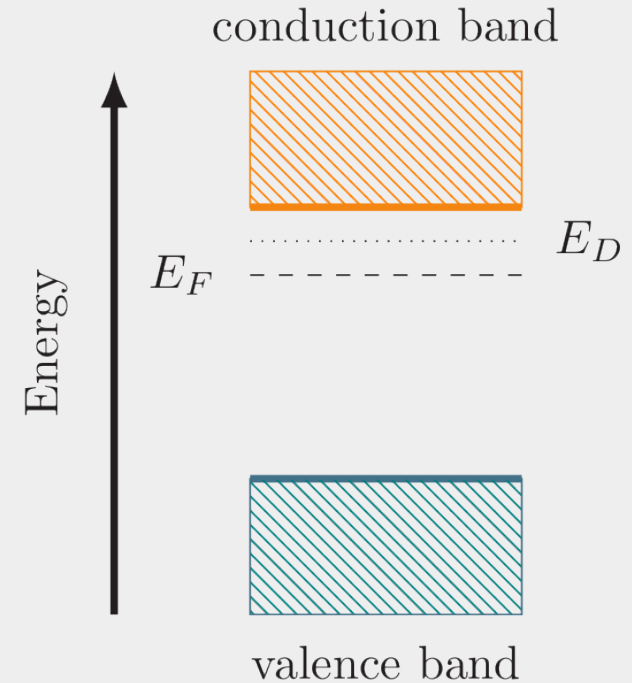


Doping Silicon – *n*-type

- Adding group-V element (phosphorus)
- Four covalent bonds, one “dangling” e
- Introduces “donor” state
- Negative majority charge carrier: “*n*”

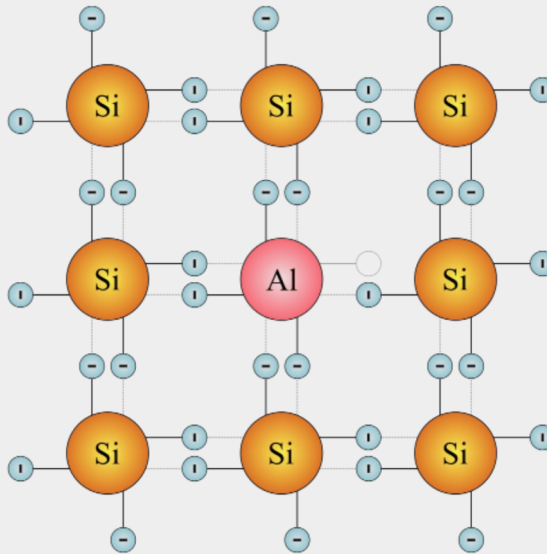


typical doping (*p-in-n* sensor):
 $N_D \approx 10^{12} \text{ cm}^{-3}$

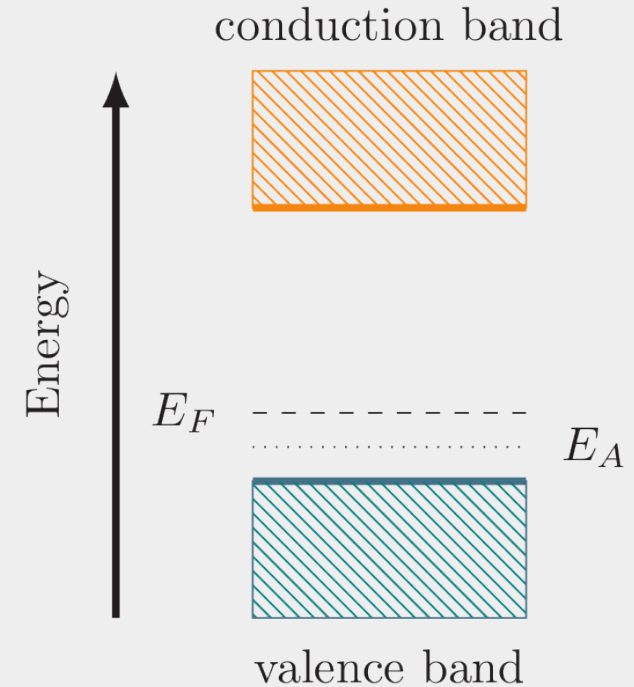


Doping Silicon – *p*-type

- Adding group-III element (boron, aluminum)
- Vacancy in covalent bonds – “hole”
- Introduces “acceptor” state
- Positive majority charge carrier: “*p*”



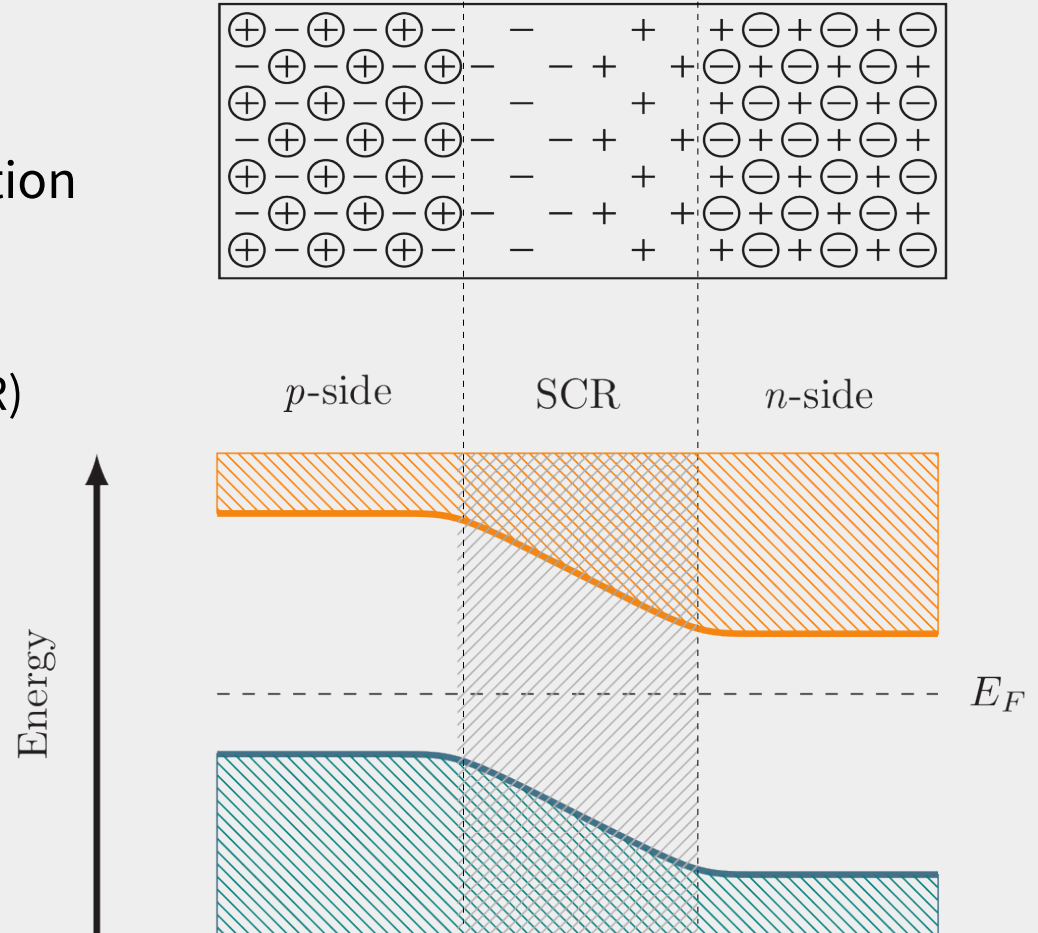
typical doping (*p-in-n* sensor):
 $N_A \approx 10^{15} \text{ cm}^{-3}$



Forming a pn -Junction

H. Spieler

- Electrons and holes diffuse over junction
- Donor/acceptor atoms remain
 - Depleted / space charge region (SCR)
 - Potential U_{bi} builds up
- Thermal equilibrium:
Built-in potential balances diffusion
- Constant Fermi level:
Deformation of energy bands



Built-in Voltage U_{bi}

- Potential across the junction: difference of Fermi energies

p-in-n sensor: $U_{bi} \approx 0.4 \text{ V}$

- Thickness of built-in SCR:

$$d(U_{bi}) = \sqrt{\frac{2 \epsilon_r \epsilon_0}{|N_D - N_A|} \cdot U_{bi}}$$

p-in-n sensor: $d \approx 20 \mu\text{m}$

$$\begin{aligned} U_{bi} &= E_{Fn} - E_{Fp} \\ &= k_B T \ln\left(\frac{N_A N_D}{n_i^2}\right) \end{aligned}$$

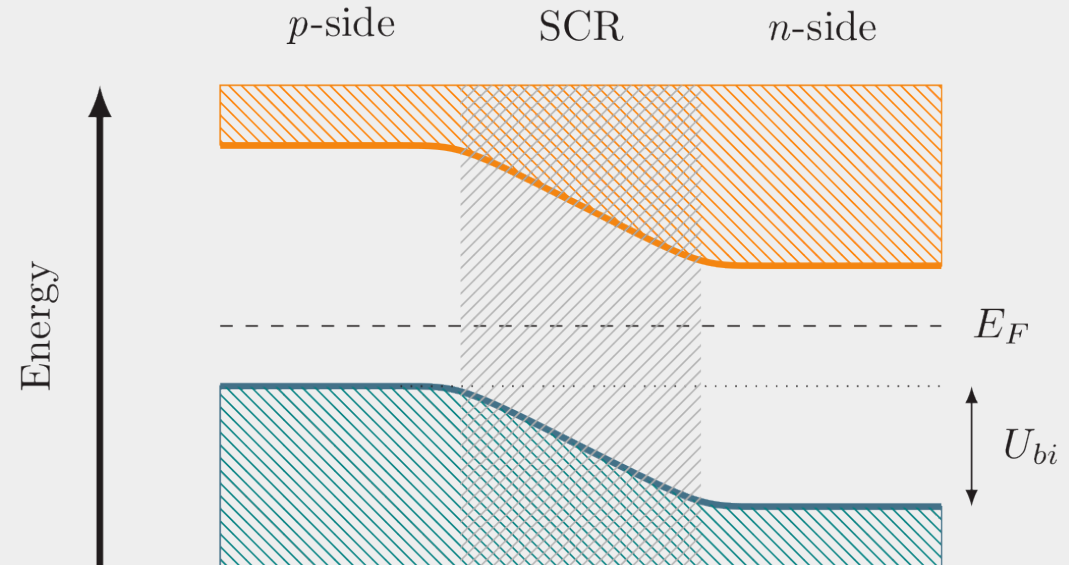
silicon p-in-n sensor:

$$N_A \approx 10^{15} \text{ cm}^{-3}$$

$$N_D \approx 10^{12} \text{ cm}^{-3}$$

$$k_B T \approx 0.026 \text{ V}$$

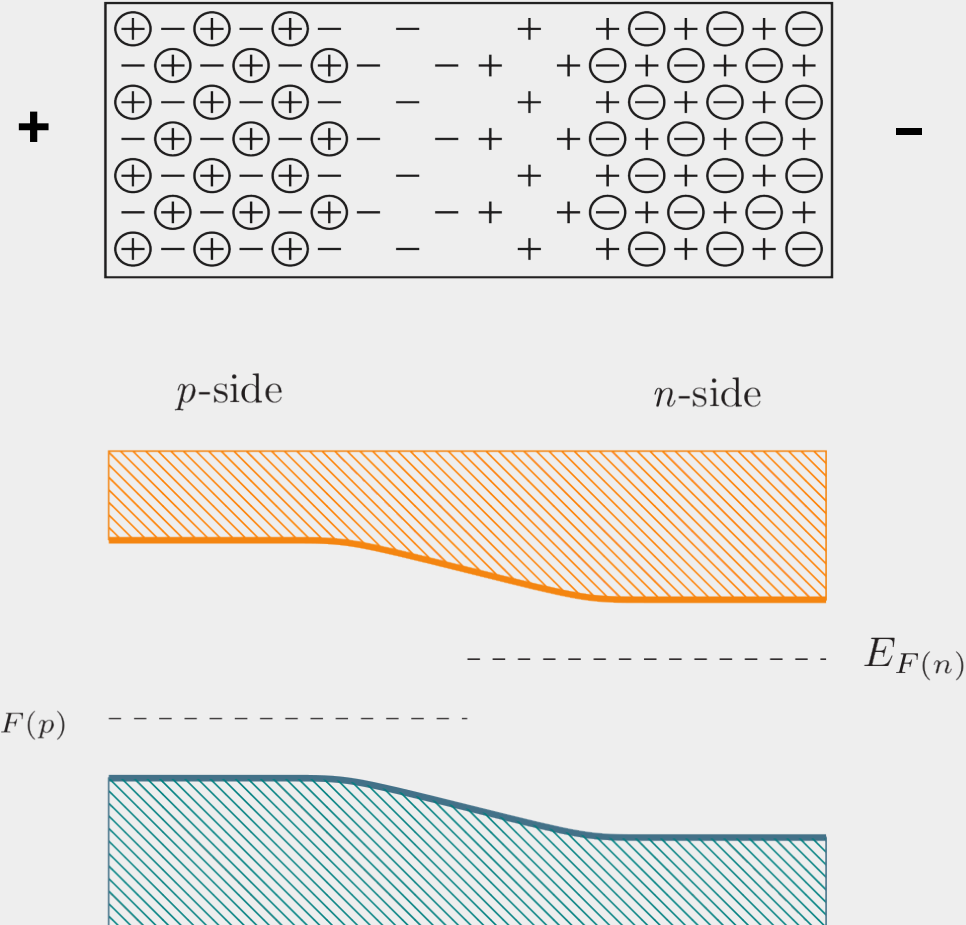
$$n_i = 1.45 \times 10^{10} \text{ cm}^{-3}$$



pn -Junction in Forward Bias

H. Spieler

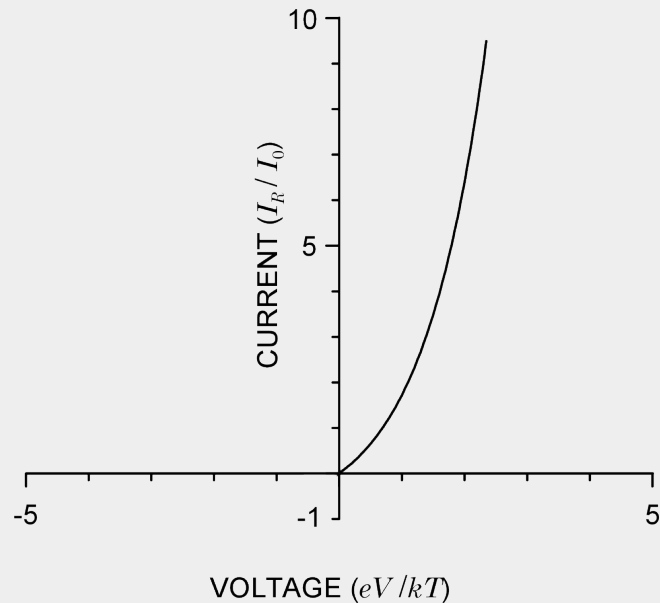
- Lowering potential difference
- Increases flow of electrons & holes



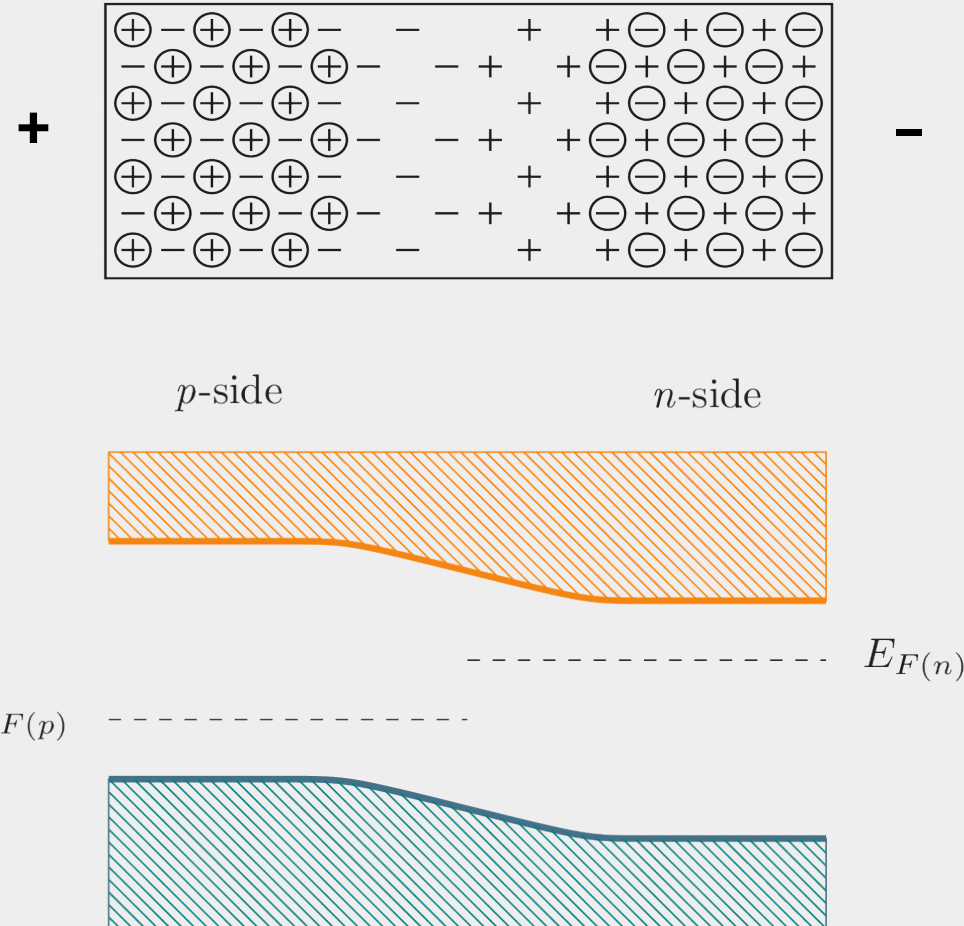
pn-Junction in Forward Bias

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- Lowering potential difference
- Increases flow of electrons & holes
- Shockley eq. $I = I_0 (e^{eU/k_B T} - 1)$



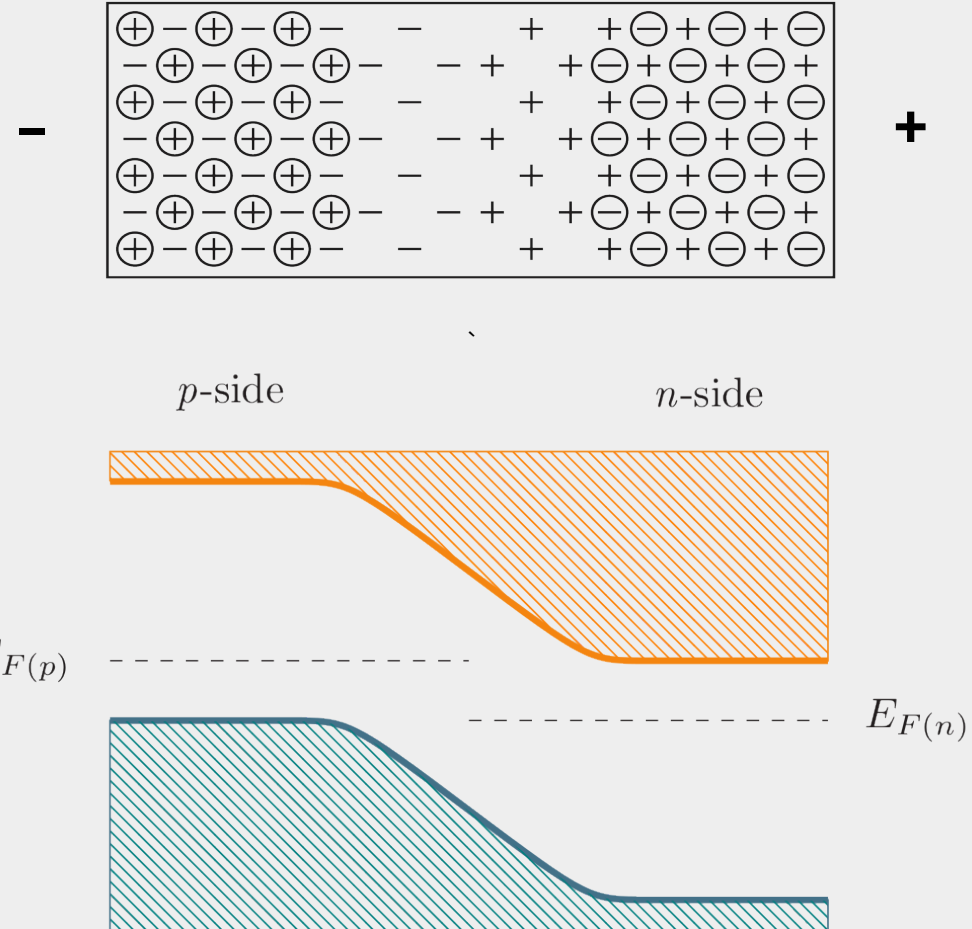
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pn-Junction in Reverse Bias

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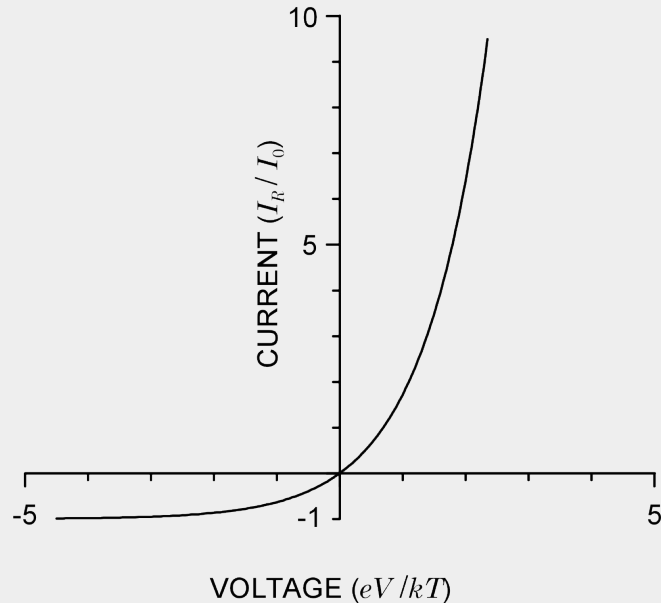
- Raising potential difference
- Widens depletion region



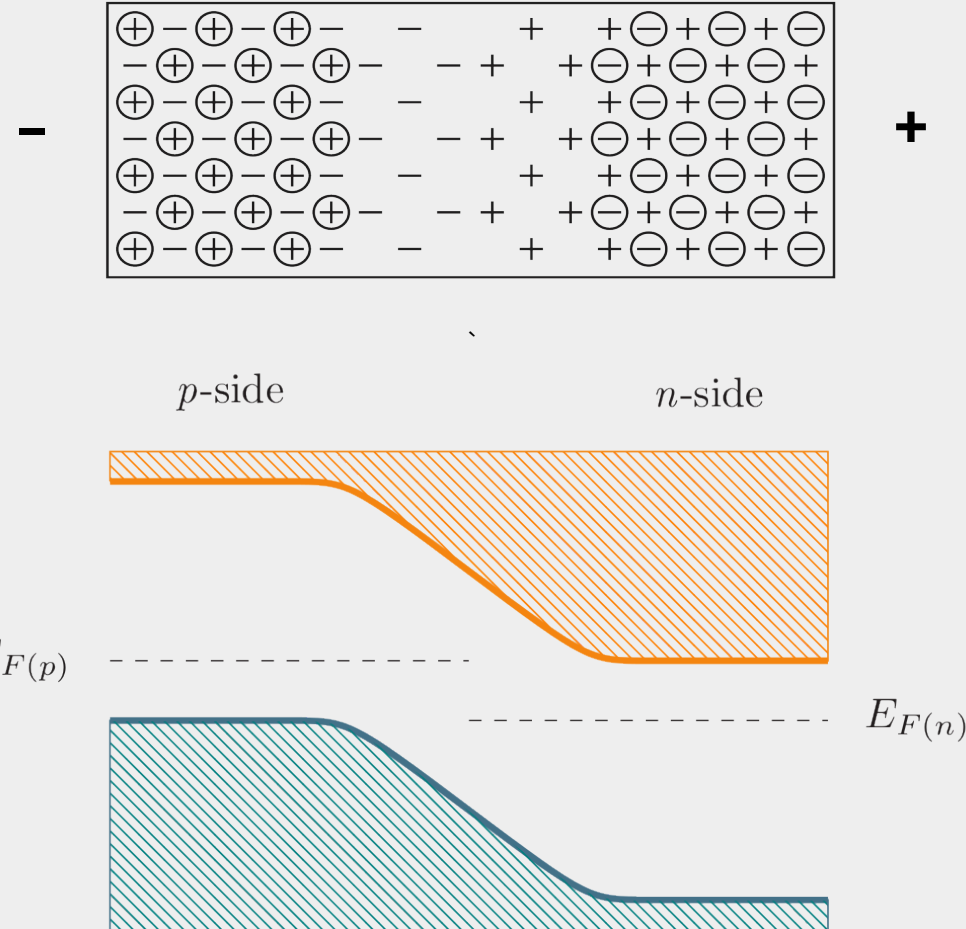
pn-Junction in Reverse Bias

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- Raising potential difference
- Widens depletion region
- Shockley eq. $I = I_0 (e^{eU/k_B T} - 1)$



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pn-Junction in a Sensor

- Asymmetric *pn*-junctions, here: *p*-in-*n*
- Lightly doped *n* bulk sensor material
- Thin, highly-doped *p* implant
- Depletion voltage:

$$V_d = \frac{Ned^2}{2\epsilon_r\epsilon_0}$$

$$e = 1.6 \times 10^{-19} \text{ As}$$

$$\epsilon_0 = 8.8 \times 10^{-12} \text{ As/V/m}$$

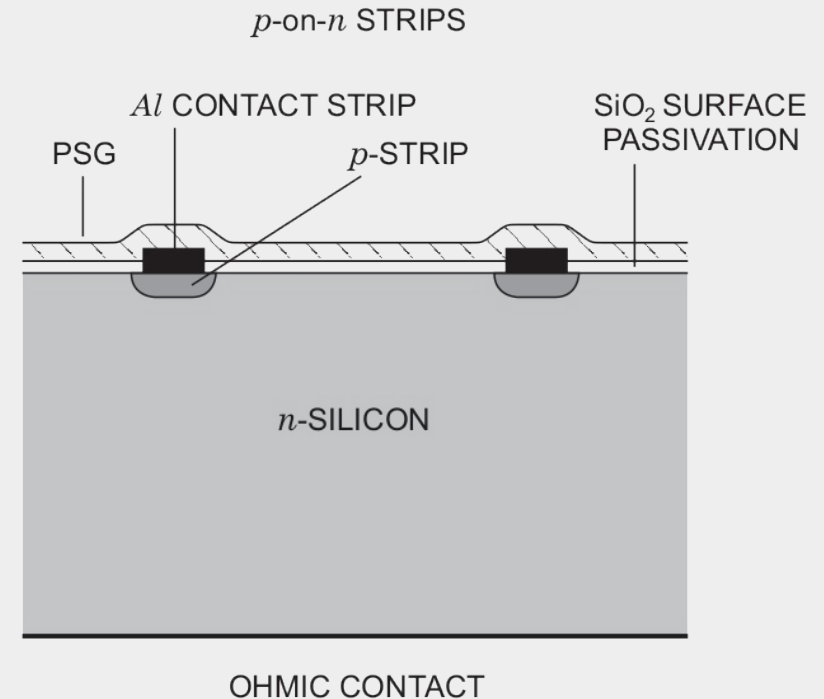
$$\epsilon_r(\text{Si}) = 11.7$$

- Segmentation of implant:
separate channels
- Backside: layer of highly doped *n*⁺
as ohmic contact

typical doping (p-in-n sensor):

$$N_A \approx 10^{15} \text{ cm}^{-3}$$

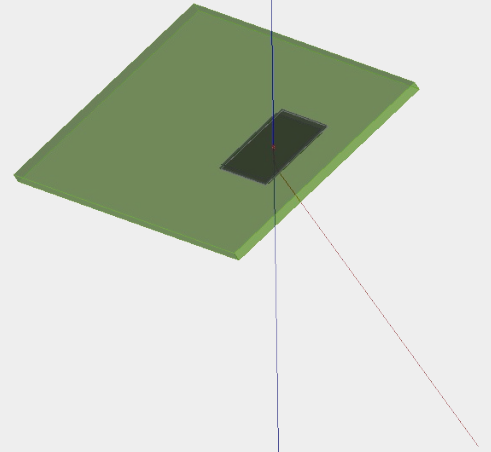
$$N_D \approx 10^{12} \text{ cm}^{-3}$$



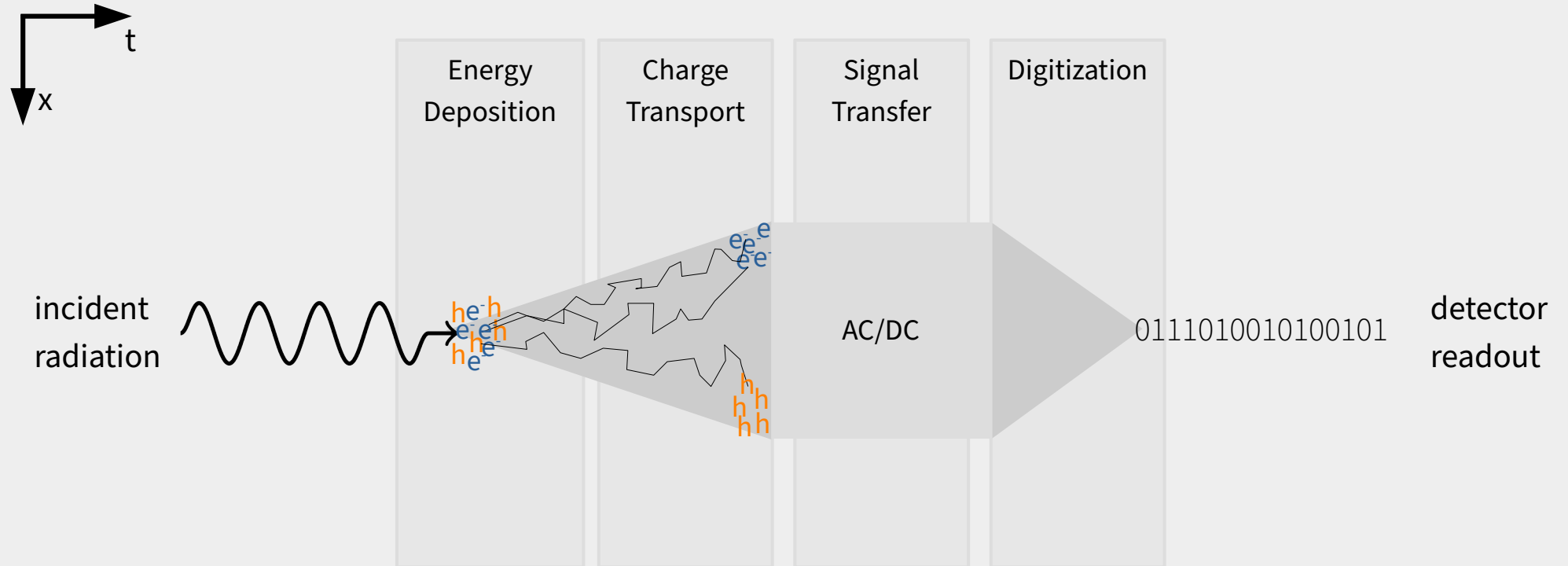
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Particle Detection

With Semiconductor Detectors



Particle Detection with Semiconductor Detectors



Energy Deposition – Energy Loss

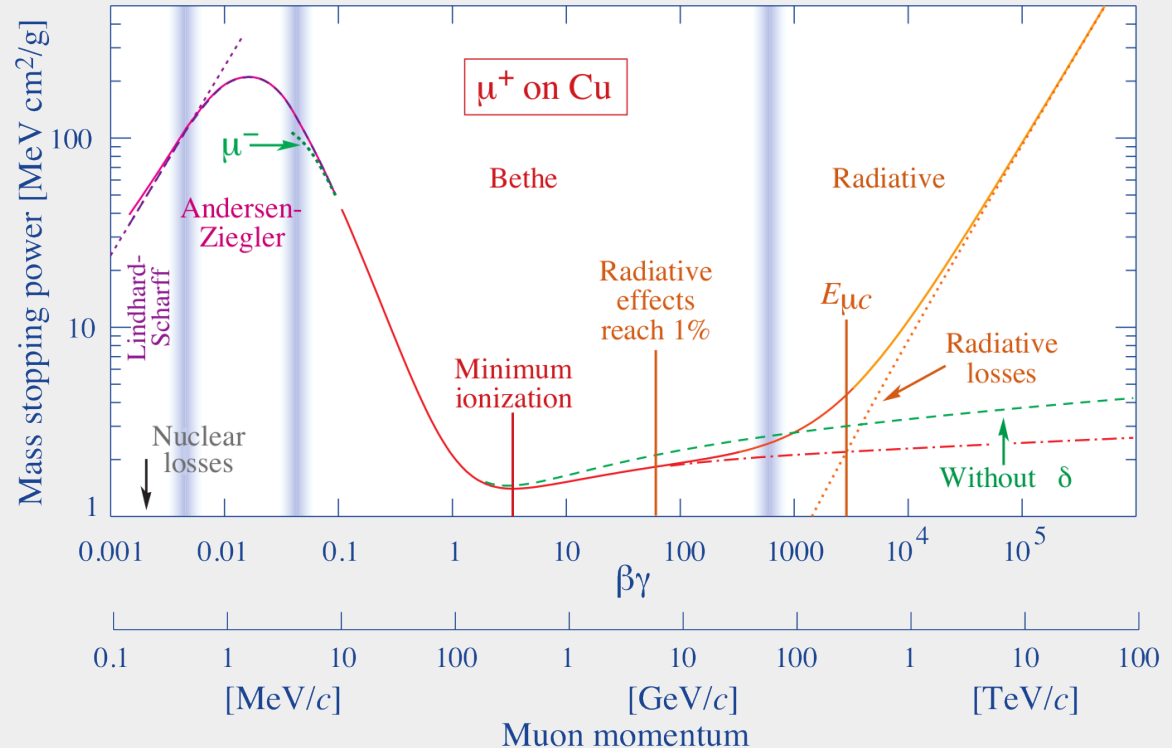
- (heavy) charged particles:
Mean energy loss described by **Bethe** formula

Phys. Rev. D 98, 030001
doi:10.1103/PhysRevD.98.030001

(sparing you the formula...)

- Definition of MIP:
Minimum Ionizing Particle

$$\beta\gamma \approx 3$$

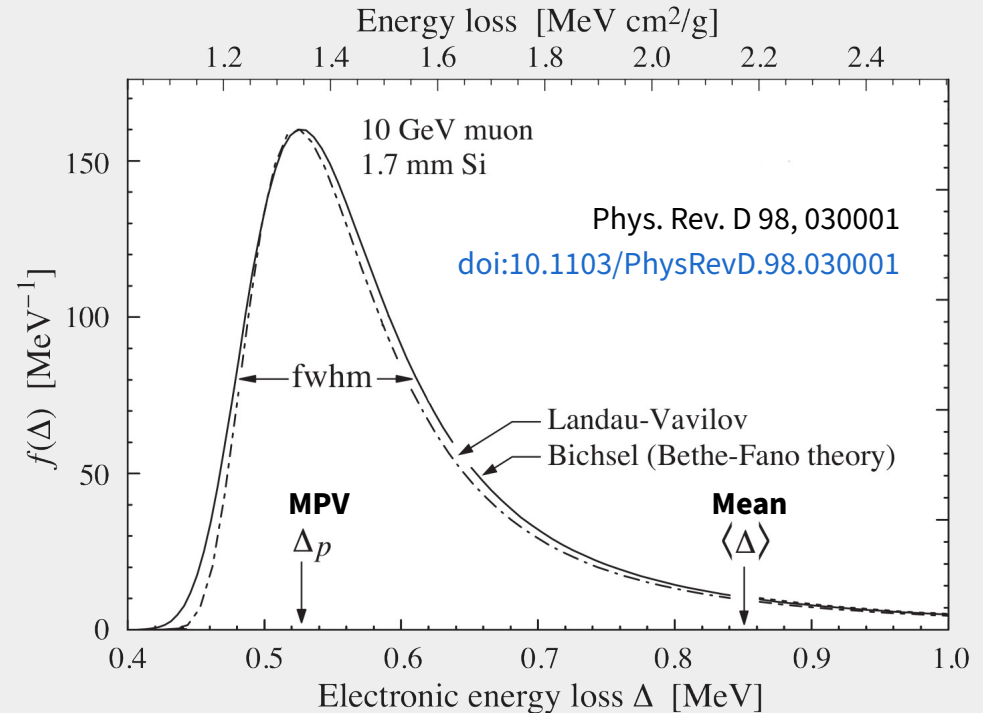


Energy Deposition – Fluctuations

- Strong fluctuations of energy loss: **Landau-Vavilov** distribution / **Bichsel** model
 - Varying number interactions, energy transfer
 - Secondary particles (e.g. delta rays)
 - Most probable value (MPV) < Mean
- Photons: Photo effect, Compton effect, pair production
- Creation of e/h pairs: 3.64 eV / pair

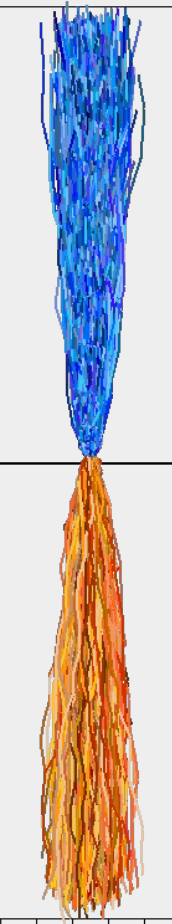
Fluctuations: **Fano** Factor

$$\sigma_{e/h} = \sqrt{N_{e/h}} \sqrt{F}$$



Signal Formation

- Sensor operated as diode in reverse bias → depleted volume
- Signal formed by motion of e/h pairs in electric field
- Contribution to motion:
 - **Diffusion** – Temperature-driven random motion, mean free path $\sim 0.1 \mu\text{m}$, mean 0
 - **Drift** – Directed motion, depending on electric field and charge carrier mobility, different parametrizations for mobility available, depending on temperature, silicon, ...
- Motion stops when...
 - Charge carriers reach readout electrode (conductor)
 - Charge carriers recombine/get trapped (depends on purity, doping, lattice defects, ...)



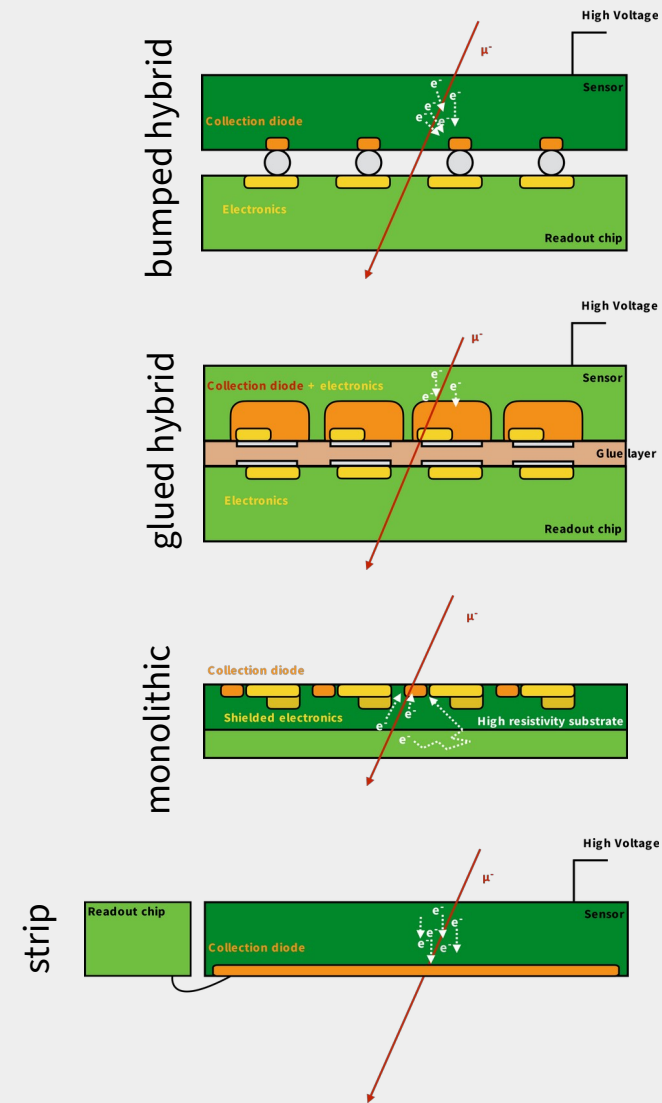
Signal Transfer

Coupling between sensor & front-end can be

- DC: bump bonds (hybrid pixel), direct (monolithic pixel), ...
- AC: glue layers (hybrid pixel), SiO_2 (strip detectors), ...

Digitization

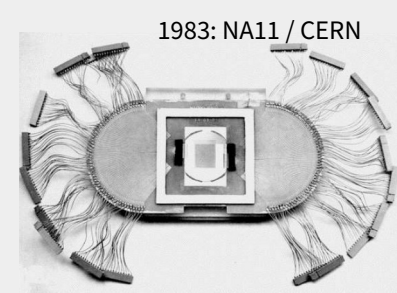
- Signal is amplified, shaped, zero-suppression (discriminator)
- Digitization of the signal via
 - Full ADC
 - Time-over -threshold
 - Threshold crossing (binary hit information)
- Buffering, encoding, data transmission...



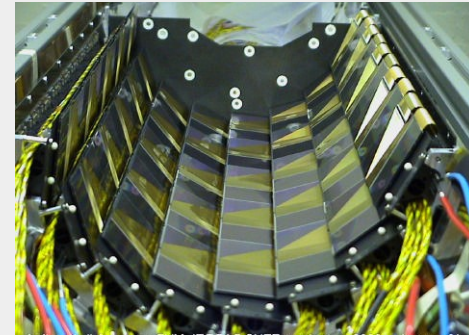
Requirements for Current & Future Tracking Detectors

Silicon Tracking Detectors in Particle Physics

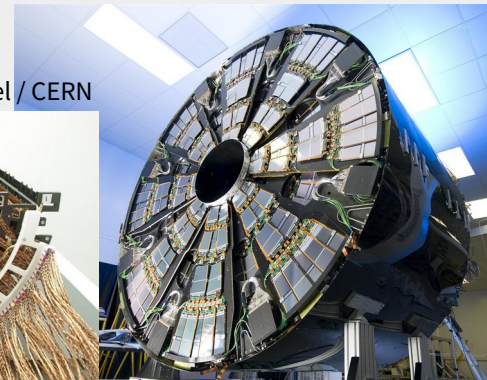
- Silicon tracking detectors have long history in particle physics
- Instrumental in discovery of Higgs boson at LHC
- Largest detectors installed in ATLAS & CMS
 - Tracking detectors: strips, 200 m² silicon, 70M channels
 - Vertex detectors: pixels, 1 m² silicon, 140M channels
- Detector upgrades for HL-LHC in preparation
 - More resilient against radiation-induced damage
 - Additional capabilities (e.g. triggering)



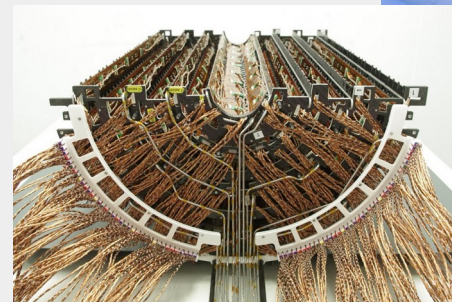
2000: ZEUS MVD / DESY



2007: CMS Tracker / CERN



2017: CMS Phase 1 Pixel / CERN



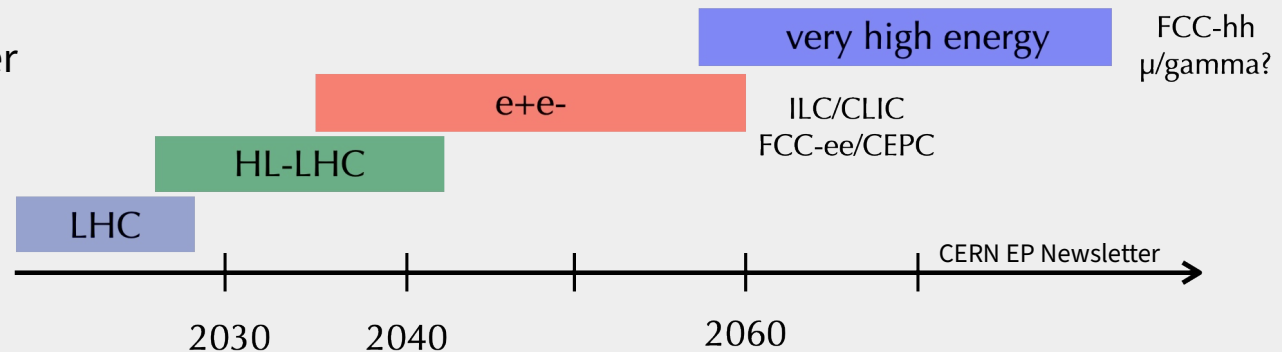
The Future of Particle Physics

- European Strategy Update: possible directions for particle physics
 - Importance of fundamental detector R&D specifically highlighted

- Higgs boson plays unique role in extending knowledge
 - Address questions within SM, provide sensitivity to new physics
 - Yukawa couplings, self-couplings, branching ratios
 - Precision measurements required



- Highest priority: future lepton collider
 - Different initial states
 - New opportunities & challenges



Silicon Detector Requirements at a Lepton Collider

- Precision measurements especially demanding on vertex & tracking detectors
 - Momentum resolution – large lever arm, minimum scattering
 - Impact parameter resolution – high resolution, minimum scattering
 - Time resolution – fast sensor response, large S/N
- Physics studies for lepton colliders provide guidelines:

	Lepton Colliders	(HL-) LHC (ATLAS/CMS)
Material budget	$< 1\% X_0$	10% X_0
Single-point resolution	$\leq 3 \mu\text{m}$	$\sim 15\mu\text{m}$
Time resolution	$\sim \text{ps} - \text{ns}$	25ns
Granularity	$\leq 25 \mu\text{m} \times 25 \mu\text{m}$	50 $\mu\text{m} \times 50\mu\text{m}$
Radiation tolerance	$< 10^{11} n_{\text{eq}} / \text{cm}^2$	O(10¹⁶ $n_{\text{eq}} / \text{cm}^2$)

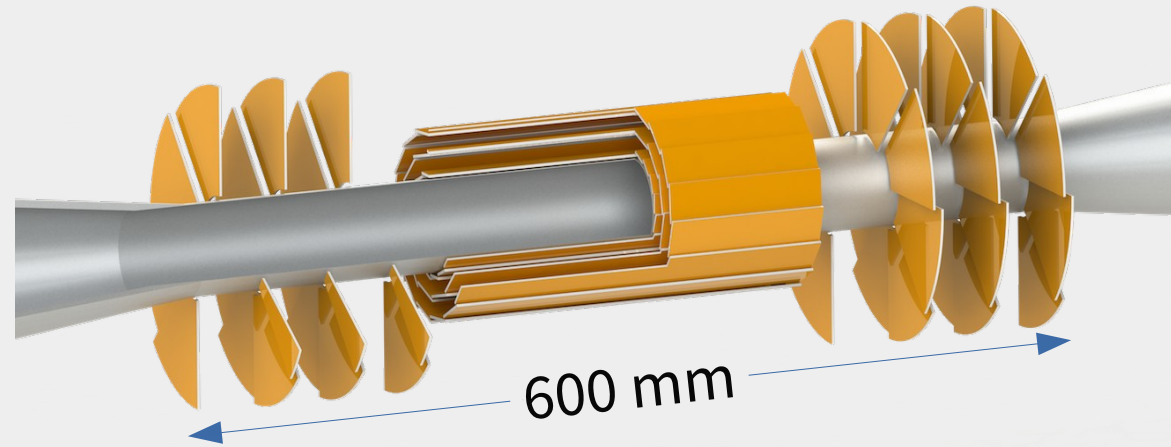
CLICdet Vertex Detector

Design driven by flavor tagging

- Minimal scattering
- High-resolution

Requirements

- **Low mass**
0.2% X_0 per layer
- **Low power consumption**
< 50 mW/cm⁻² for air-flow cooling
- **High single-point resolution**
 $\sigma_{SP} \sim 3 \mu\text{m}$
- **Precise time stamping** $\sim 5 \text{ ns}$



Current design:

- Double layer sensors
- 100 μm of silicon, 25 μm pixel pitch
- Surface area of $\sim 1 \text{ m}^2$
- Three barrel double-layers,
2x three spiral double-disks



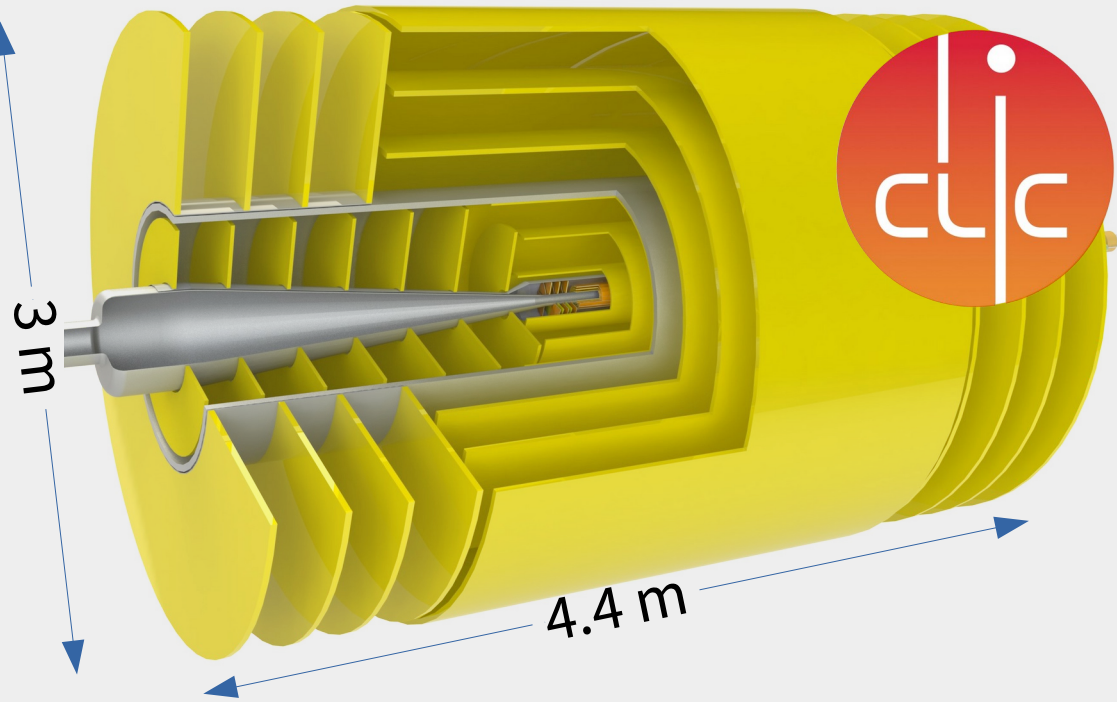
CLICdet Tracking Detector

Design driven by
efficiency & momentum resolution

- Many layers, large lever arm

Requirements

- **Low mass, high rigidity**
1 – 2% X_0 per layer
- **Good single-point resolution**
 $\sigma_{SP} \sim 7 \mu\text{m}$ (transverse plane)
- **High granularity**
few % occupancy from backgrounds
- **Precise time stamping** $\sim 5 \text{ ns}$



Current design:

- Detector with (elongated) pixels
- Max. 200 μm sensor, including electronics
- Surface area of approx. 140 m^2
- Leakless water cooling

Towards Next-Generation Tracking Detectors

Prospective R&D

- Define requirements from physics program, precision targets
- Explore ideas, new concepts
- Technology evaluation
- Simulations
- Proof-of-principle

Guided R&D

- Technology consolidation
- Demonstrators
- Design optimization
- Performance studies

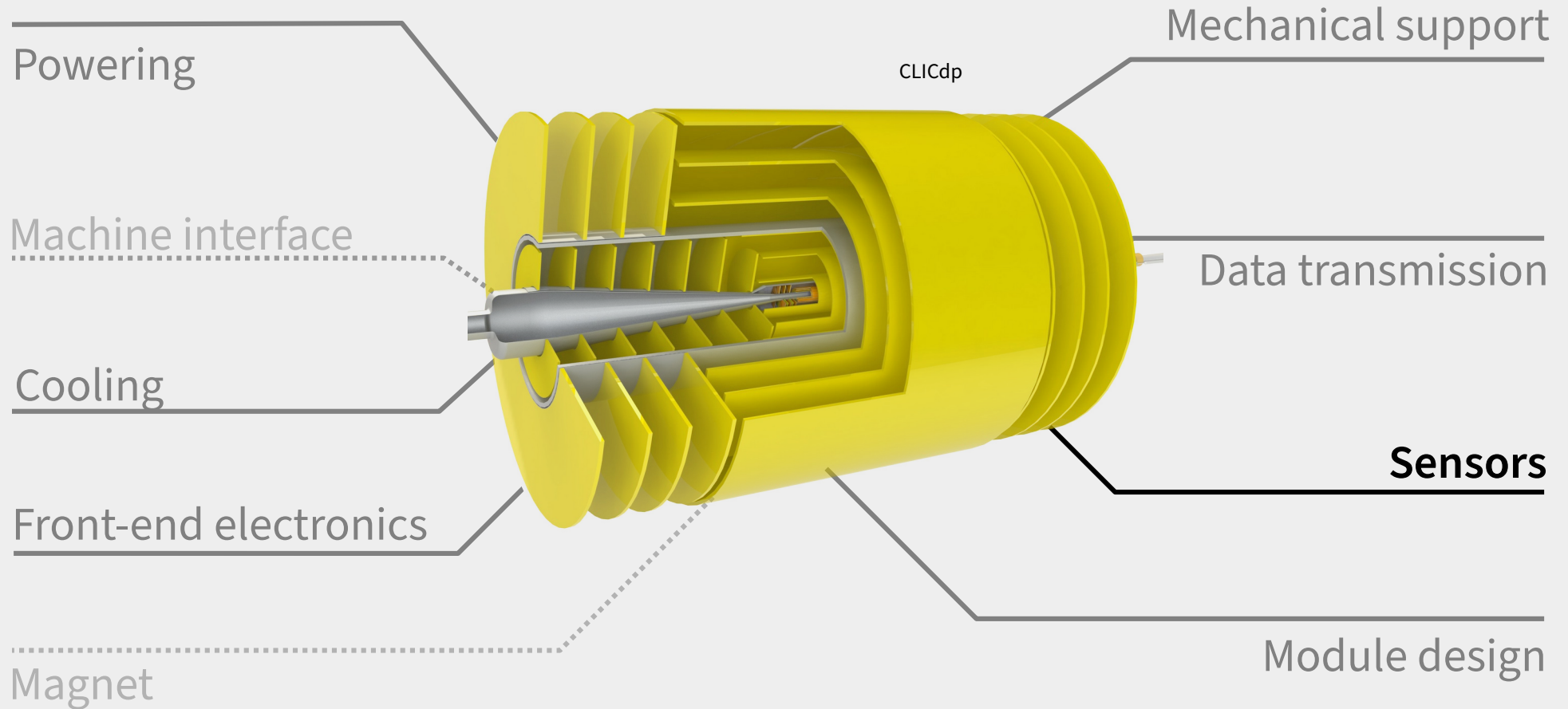
Targeted R&D

- Full-scale prototypes, engineering
- System integration

Construction

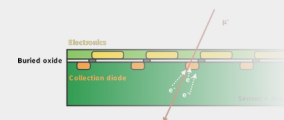
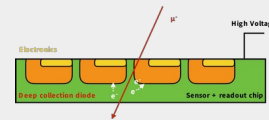
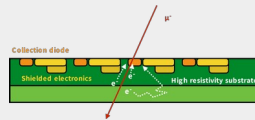
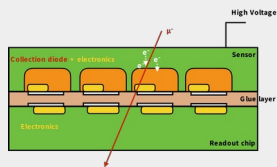
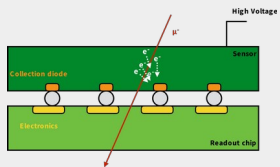
Collisions

Next-Generation Vertex & Tracking Detectors



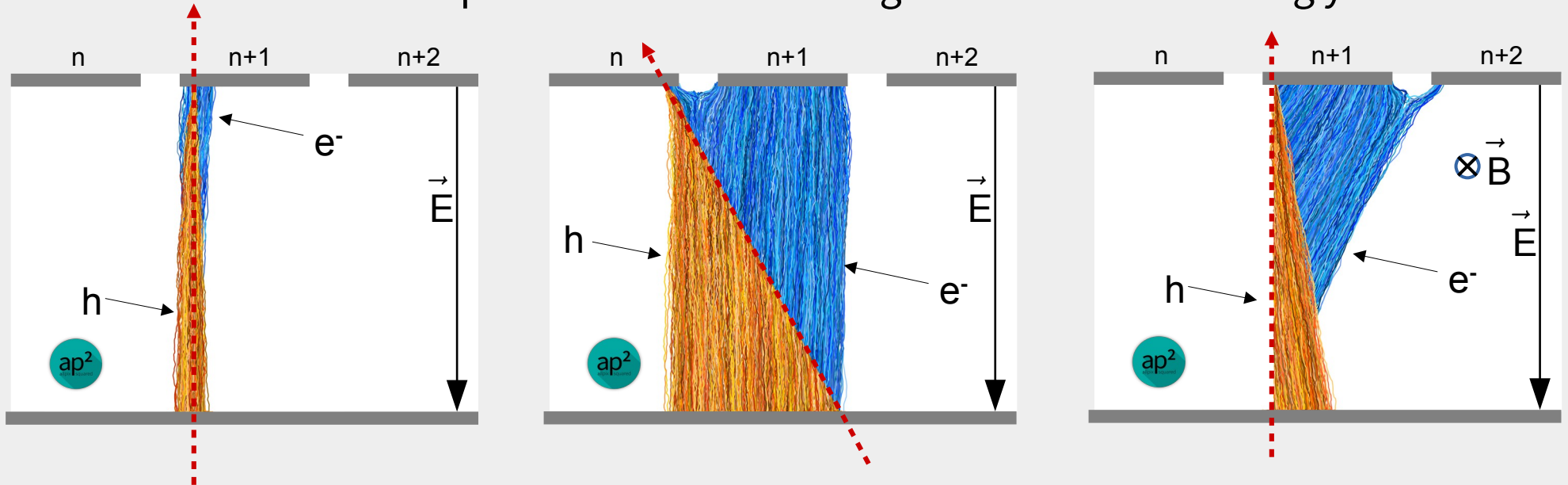
(Some) Sensor Technologies

for Future Tracking Detectors



Charge Sharing – Inclined Tracks & Lorentz Drift

- Charge sharing: distribution of charge carriers / signal over several strips (pixels)
- Can significantly improve the spatial resolution
- Often used: Inclined particle incidence along x & Lorentz drift along y



Position Resolution

- Precision of particle incidence prediction

Just a single pixel struck:

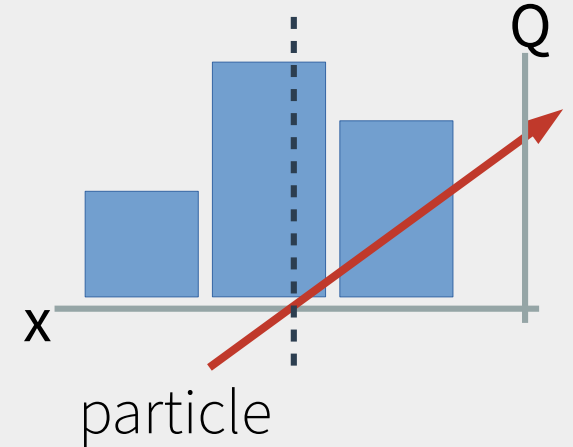
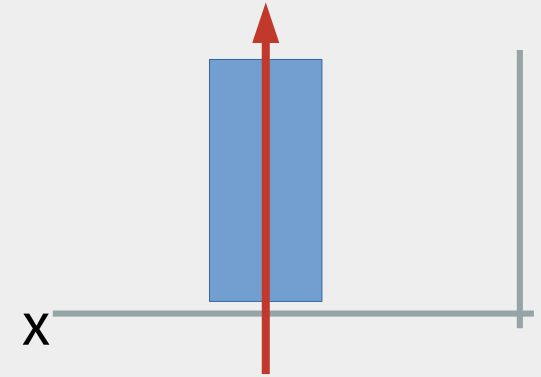
precision limited to variance of uniform distribution

$$\sigma = p/\sqrt{12}$$

- Multiple pixel struck (charge sharing):
interpolation using relative energy / charge distribution

$$\sigma = \frac{\sum_i Q_i x_i}{\sum_j Q_j}$$

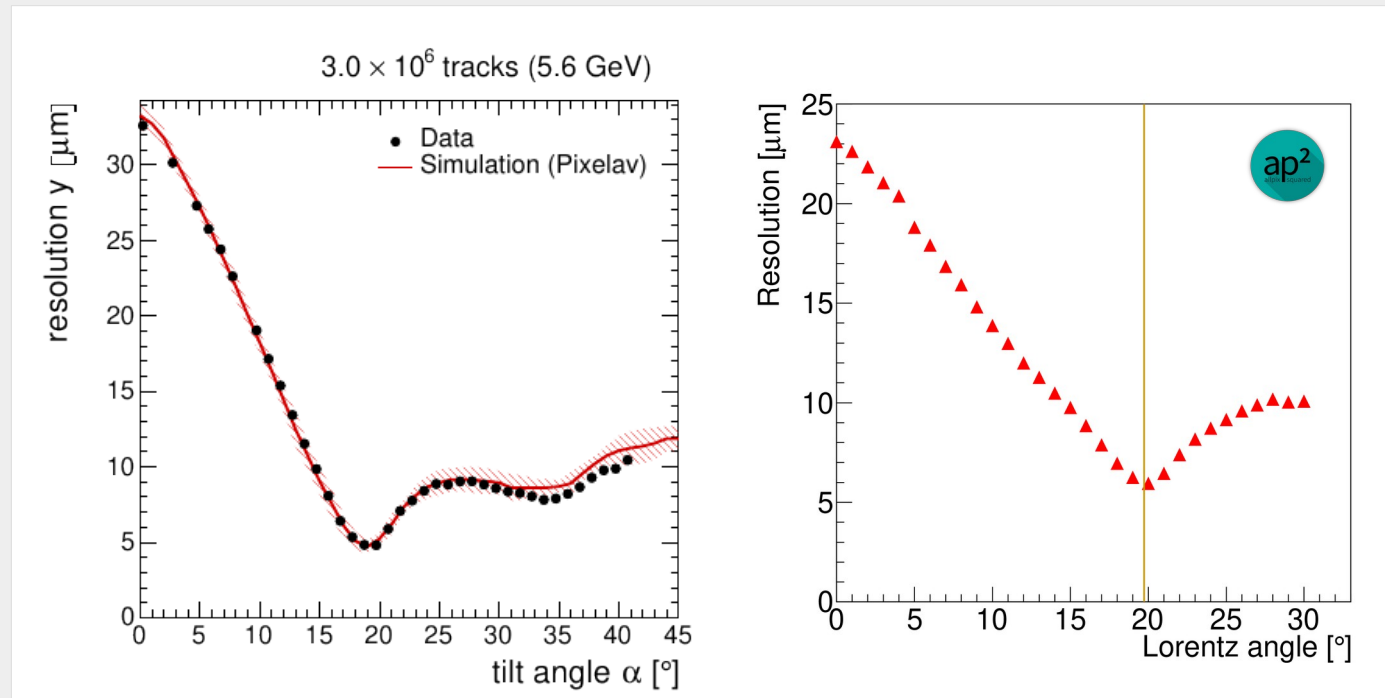
- Thinner sensors: less charge sharing...

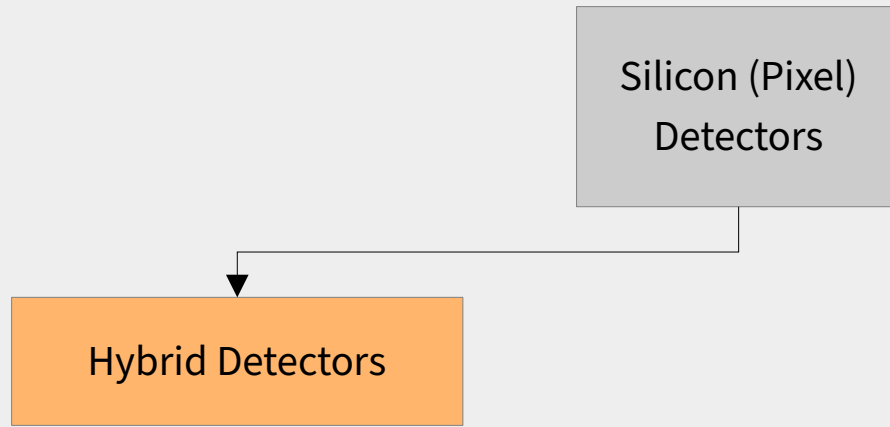


Position Resolution

- Charge sharing improves resolution in case of per-pixel charge information
- Similar effects from incidence angle and Lorentz drift

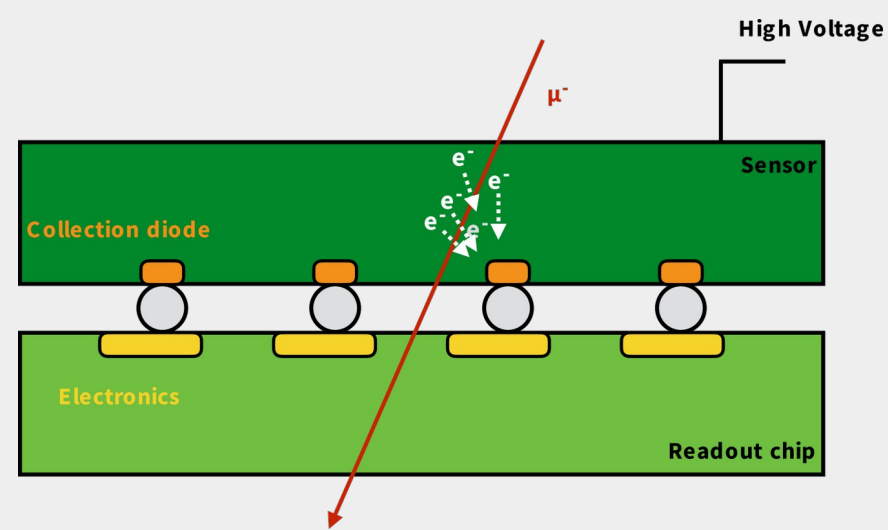
- Typical pixel pitches:
20 – 400 μm
- Typical resolutions:
5 – 15 μm





Hybrid Silicon Pixel Detectors

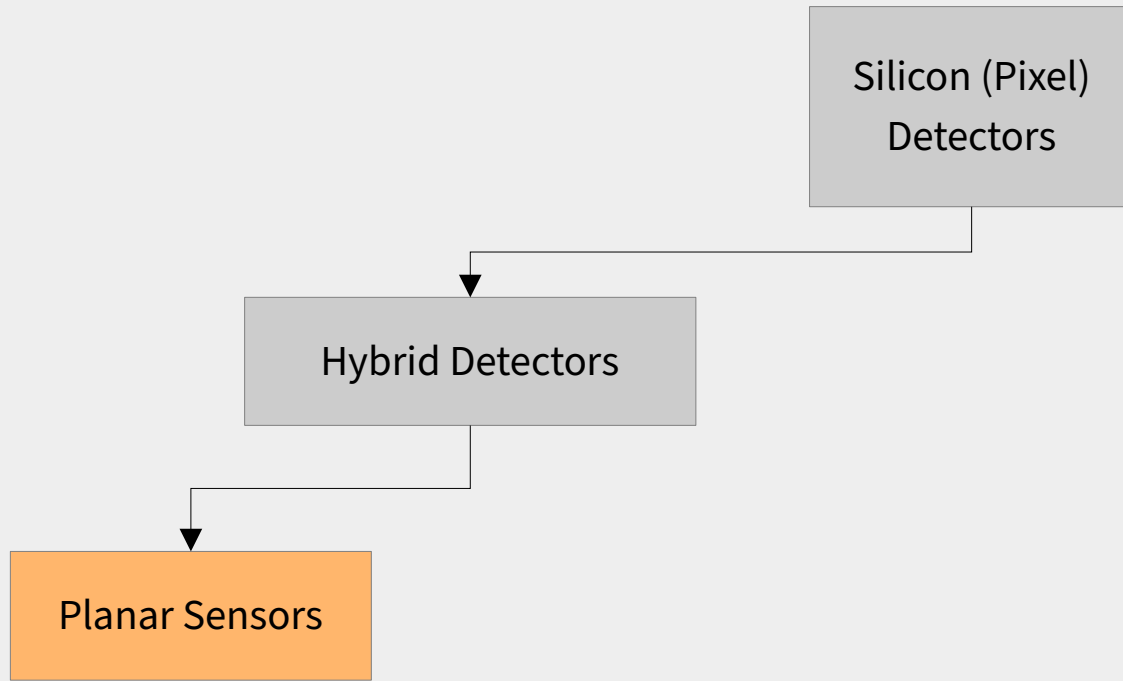
- Traditional design of HEP silicon pixel consist of sensor and separate readout chip
 - Sensor: pn-junction
 - Readout chip: front-end
 - Connection: small solder spheres – bump bonding
- Small pixel cell sizes achieved, $\sim 25 \mu\text{m}$ – limited by interconnects



Established mixed-mode CMOS
Complex circuits possible
Small technology nodes available

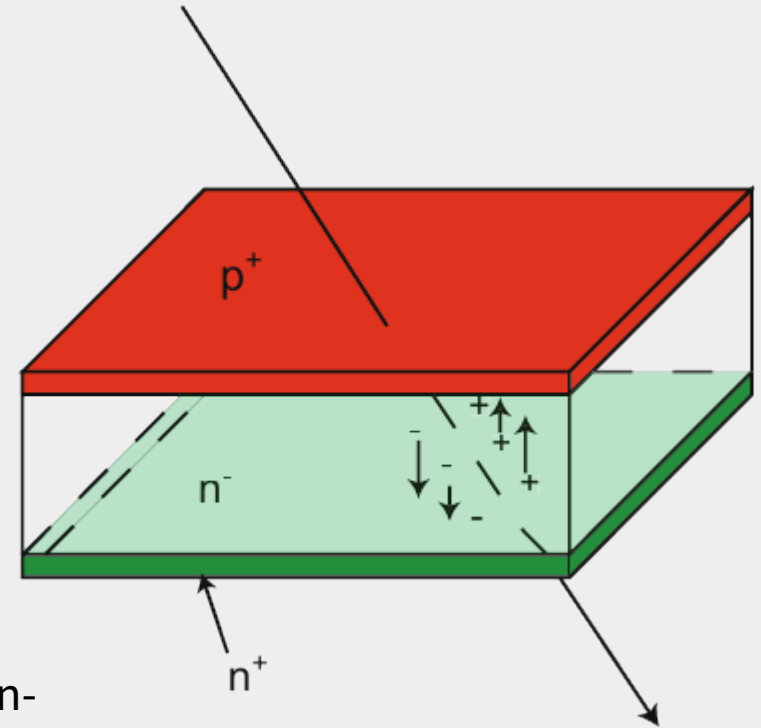


Relatively high material budget
Interconnects: cost-driver,
limits pixel pitch & thickness (stability)



Planar Silicon Sensors – The Diode

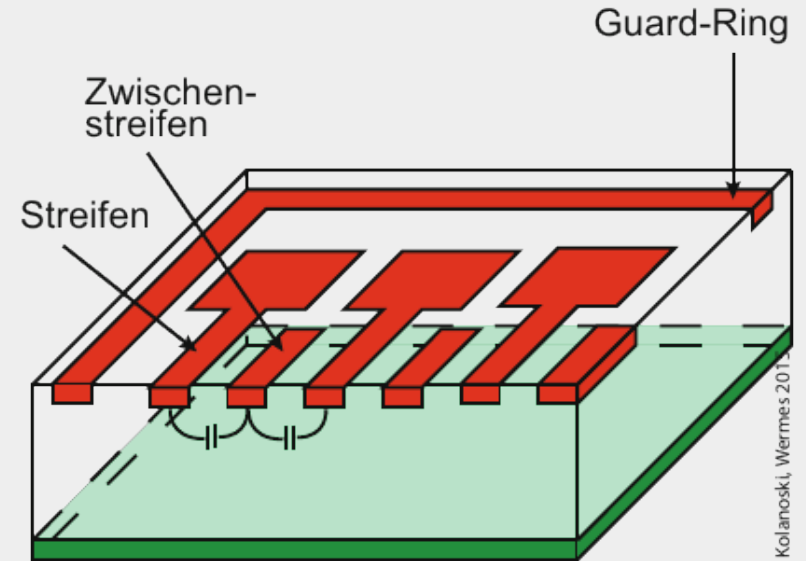
- Simplest semiconductor detector geometry
- Readout of a full area detector pad
- No spatial information
- Number of channels: 1
- Here:
 - Strong p^+ and weak n^- doping create asymmetric pn-junction at the sensor surface
 - Strong doping (n^+) at the backside for Ohmic contact to backside metalisation



Planar Silicon Sensors – Strip Detector

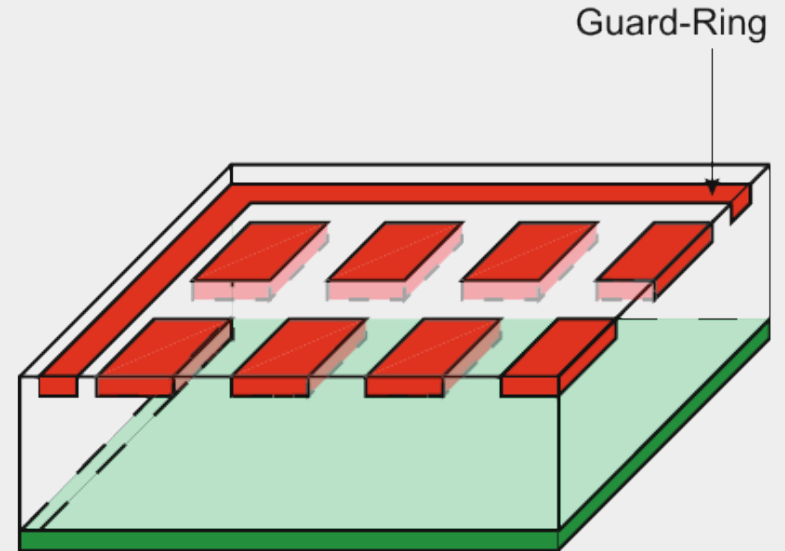
- Segmentation of sensor surface
- Implementation of strips
- Typical pitches: 50 – 100 μm

- Number of channels: N
- Charge carriers propagate towards one or few strips
 - ➔ 1D spatial information on particle traversal
 - ➔ Add second layer for 2D information



Planar Silicon Sensors – Pixel Detector

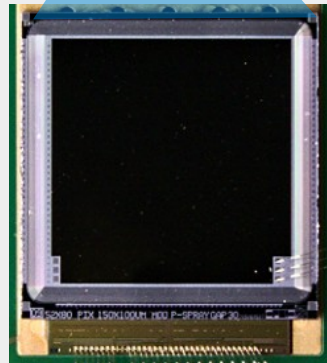
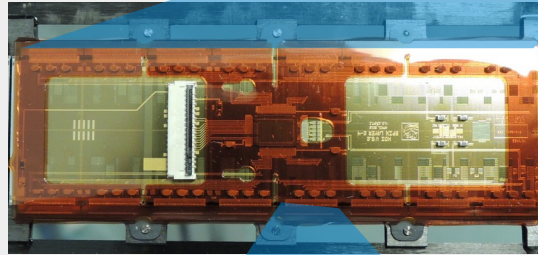
- Segmentation of sensor surface
- Implementation of pixels or pads
- Typical pitches: 50 – 400 μm



- Number of channels: N^2
- Charge carriers propagate towards one or few pixels
→ 2D spatial information on particle traversal

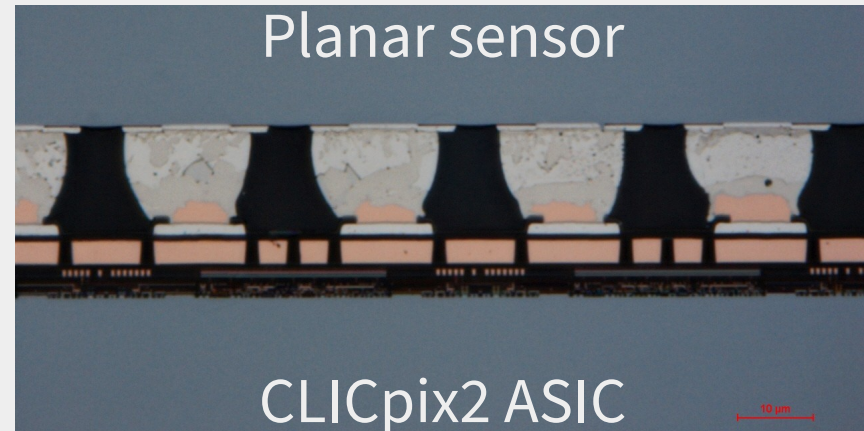
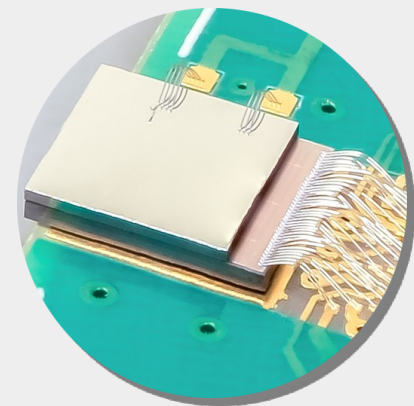
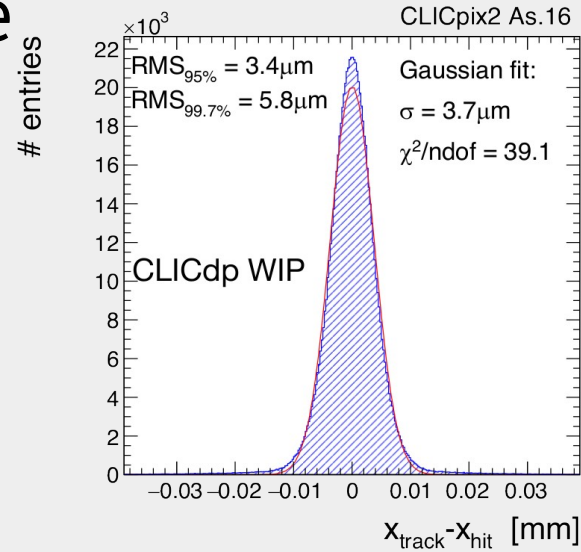
Example: CMS Phase I Pixel Detector

- Hybrid pixel detector with planar sensors
 - 4 barrel layers (30 – 160mm radius)
2x 3 end-cap disks
- ~1900 detector modules
~120M channels
- 150 μm x 100 μm pixels
8 μm (z) / 5 μm ($r\phi$) resolution
- Material budget:
~ 0.3% X_0 per module

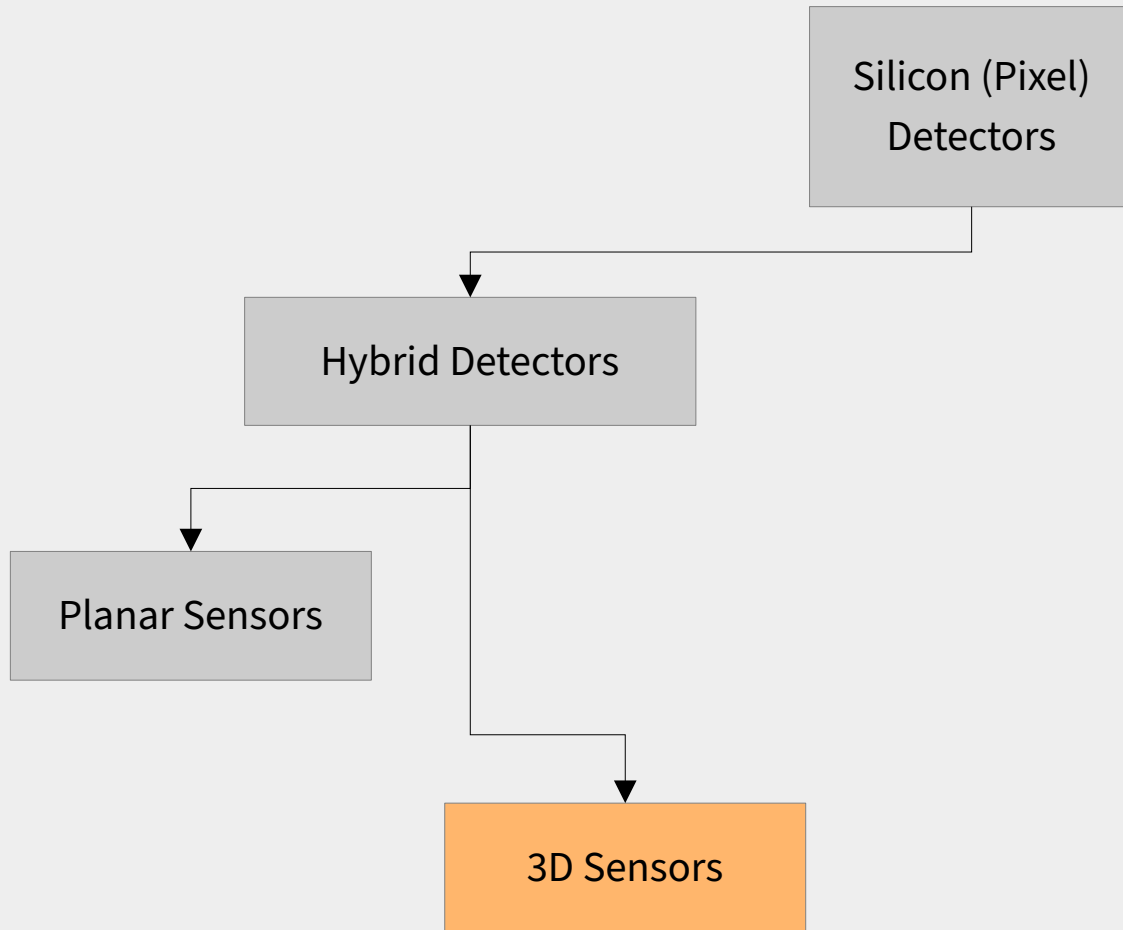


The CLICpix2 Hybrid Prototype

- Readout ASIC designed for CLIC vertex detector
 - Designed in 65nm CMOS process (same as **RD53** for **ATLAS/CMS**)
 - Matrix of 128 x 128 pixels of 25 x 25 μm
- Goal: development of high-resolution detector for detector at linear accelerator
- Chip-level bump bonding difficult, limitation on pixel pitch

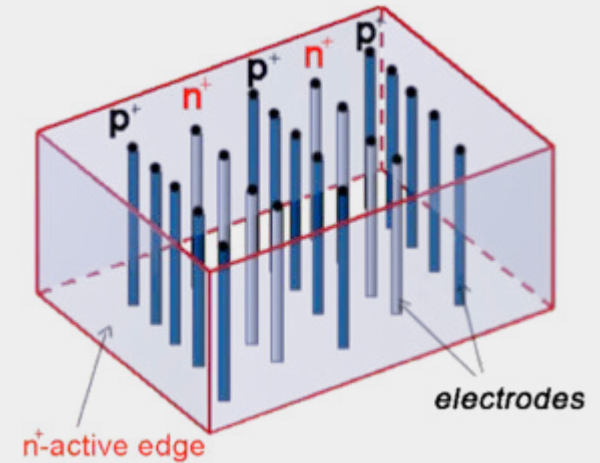


<https://doi.org/10.1088/1748-0221/14/06/C06003>

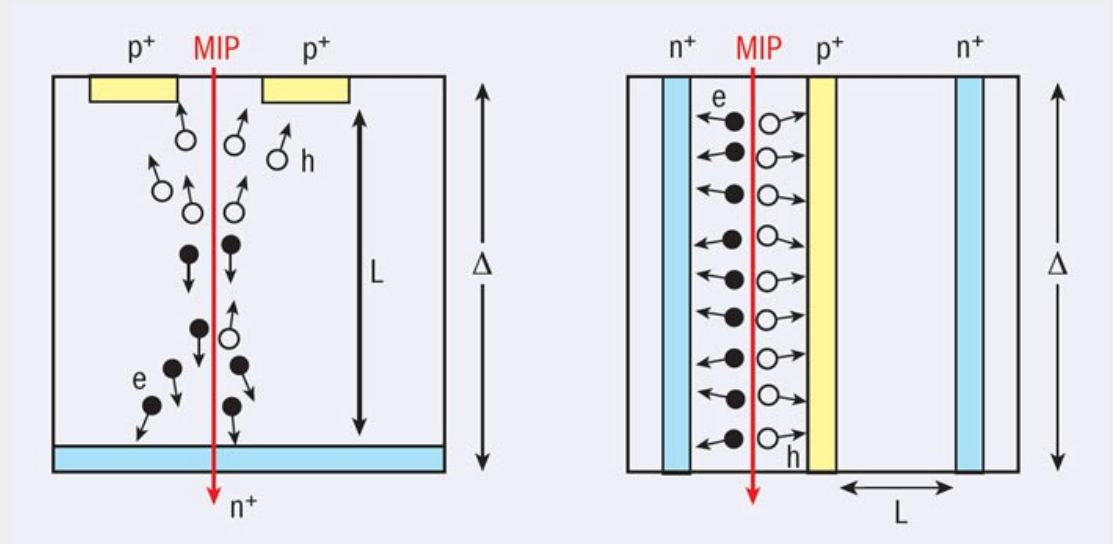


3D Silicon Sensors

- p- and n-implants implemented as columns through the sensor volume
→ Generation of a horizontal pn-junction

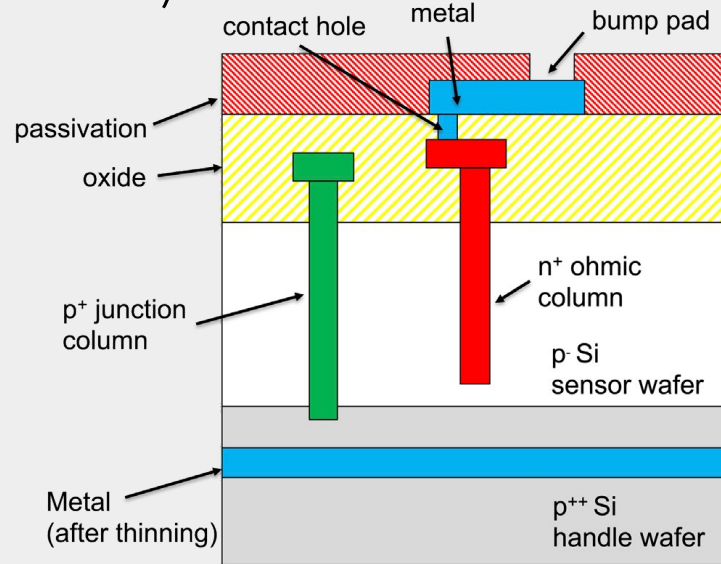


- ✓ Short drift time → fast!
- ✓ High (not reduced) signal
- ✓ High radiation tolerance
- ✗ High production costs & time

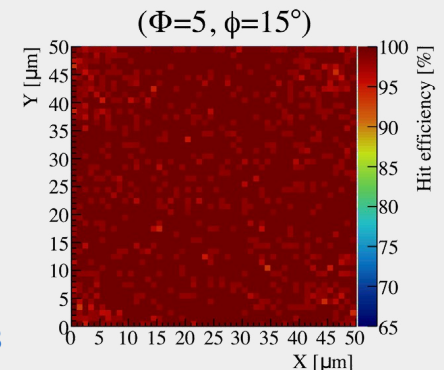
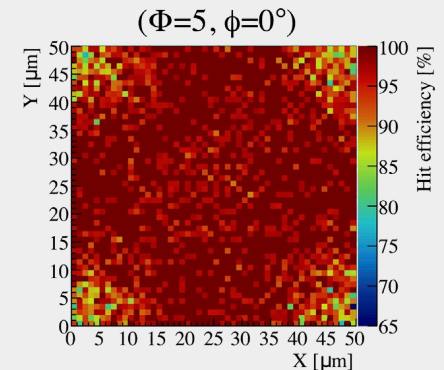
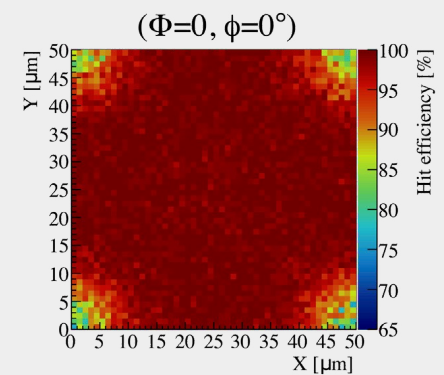


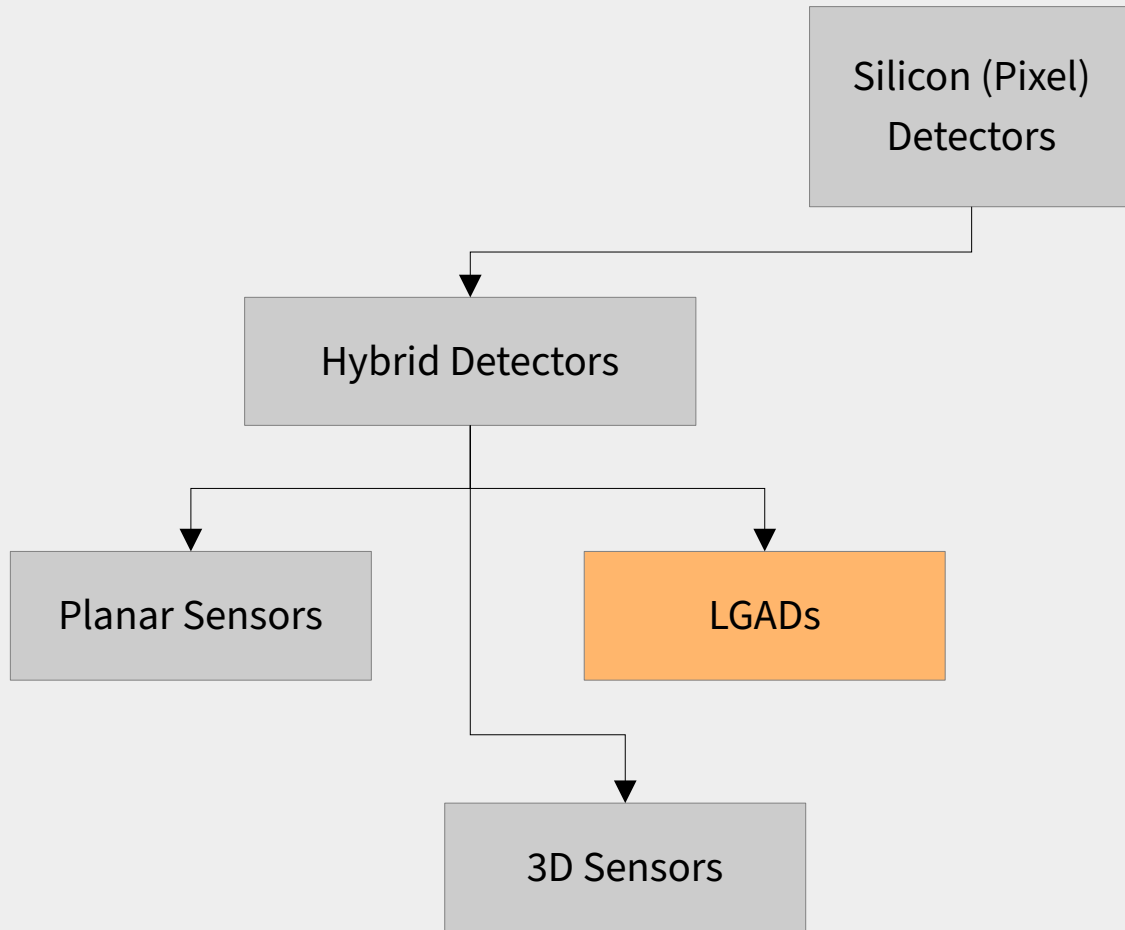
ATLAS ITk 3D Pixel Sensors

- Sensor for new inner tracker of ATLAS experiment
- 3D sensors for innermost layer of pixel detector
 - Very radiation hard (short drift times)
 - Different sensor layouts:
 - 50 x 50 μm
 - 25 x 100 μm
- At vertical incidence: inefficiencies at backside columns



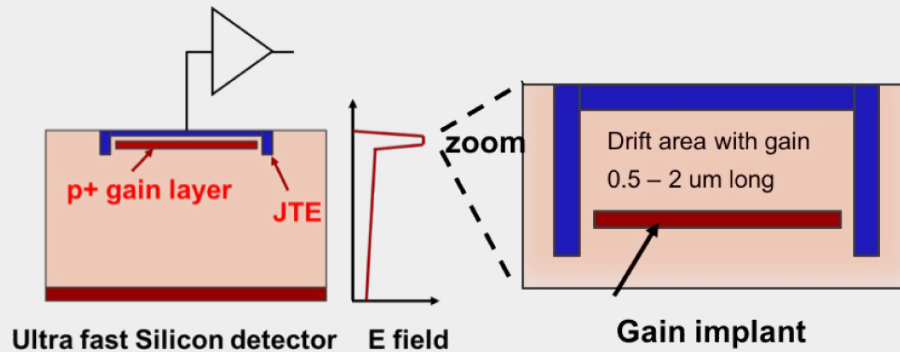
<https://doi.org/10.3389/fphy.2021.624668>



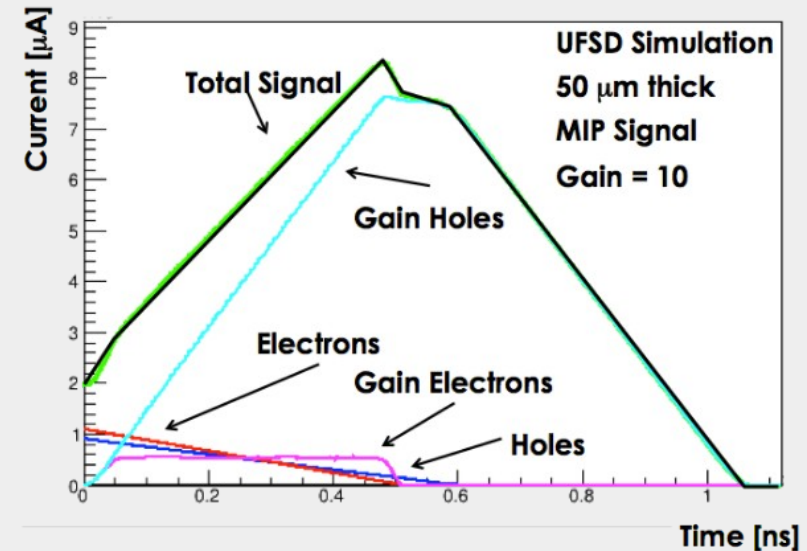


Low Gain Avalanche Diodes (LGADs)

- High electric fields: secondary ionization by charge carriers becomes possible → **Impact Ionization**
 - Similar to charge multiplication in gaseous detectors
- High electric fields in small sensor volume fraction generated via thin doping layer

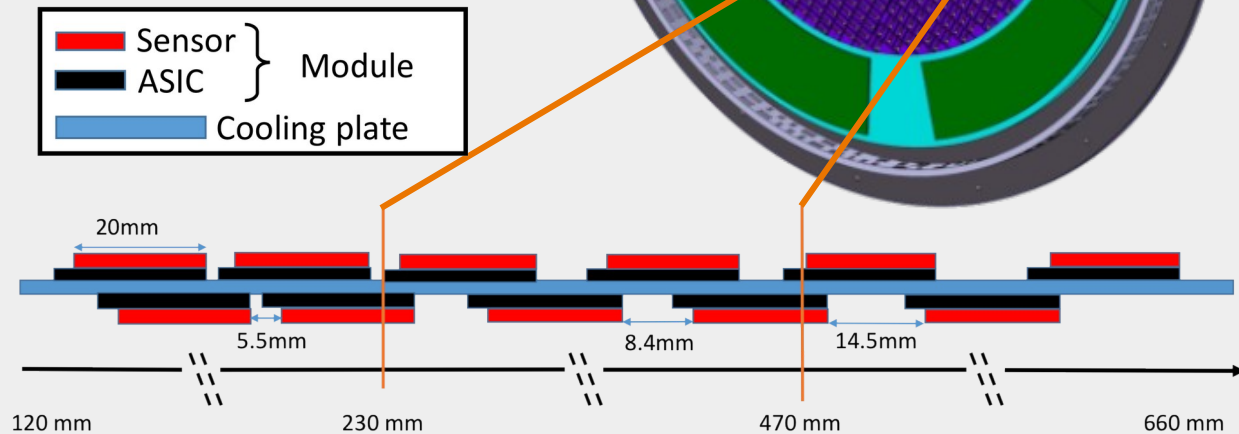
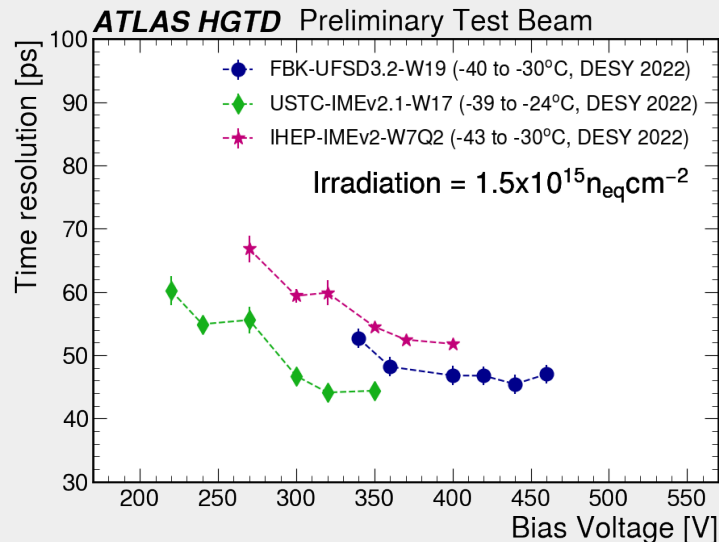
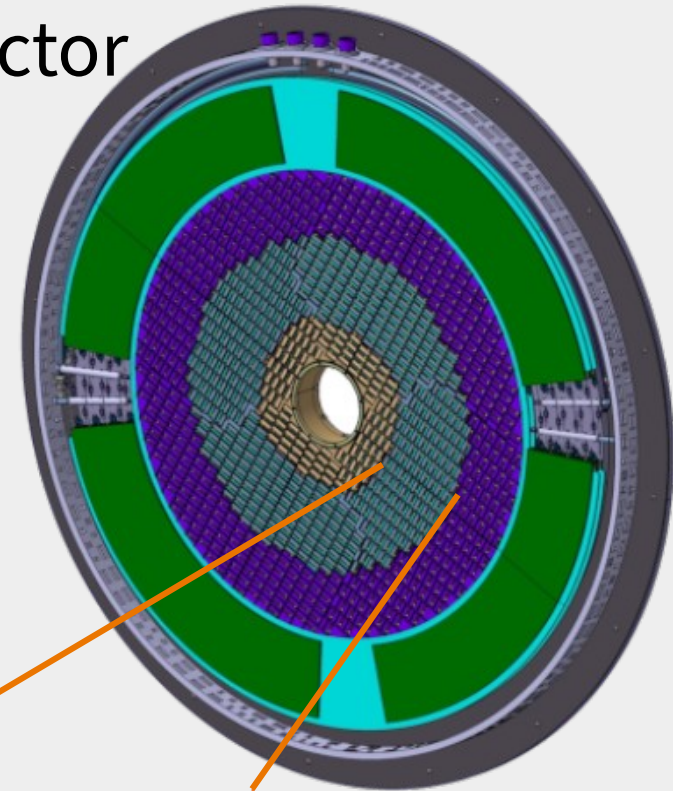


- Different types of LGADs: pads, iLGADs, TI-LGADs, ...
- “Typical” silicon detectors: $\sigma > 1 \text{ ns}$
- LGADs (Low Gain Avalanche Diodes) / UFSDs: $\sigma > 30 \text{ ps}$



The ATLAS High Granularity Timing Detector

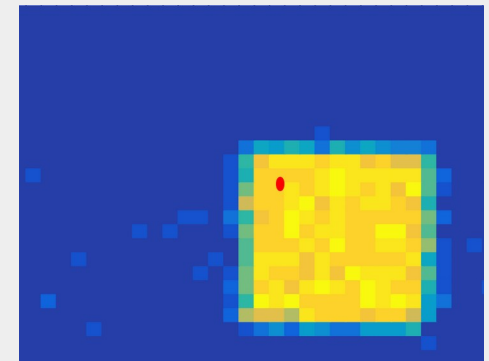
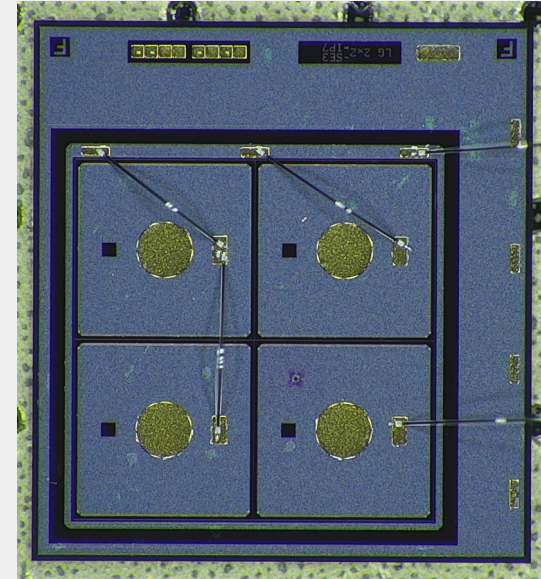
- Mitigate tracking issues from high pile-up at HL-LHC
 - Required timing resolution better than 50 ps/track
 - $\sim 3.7 \times 10^6$ channels with 6.4 m² area
 - Radiation hardness 2.5×10^{15} N_{eq}/cm² and 2.0 MGy
- LGADs with dedicated readout ASIC (ALTIROC)



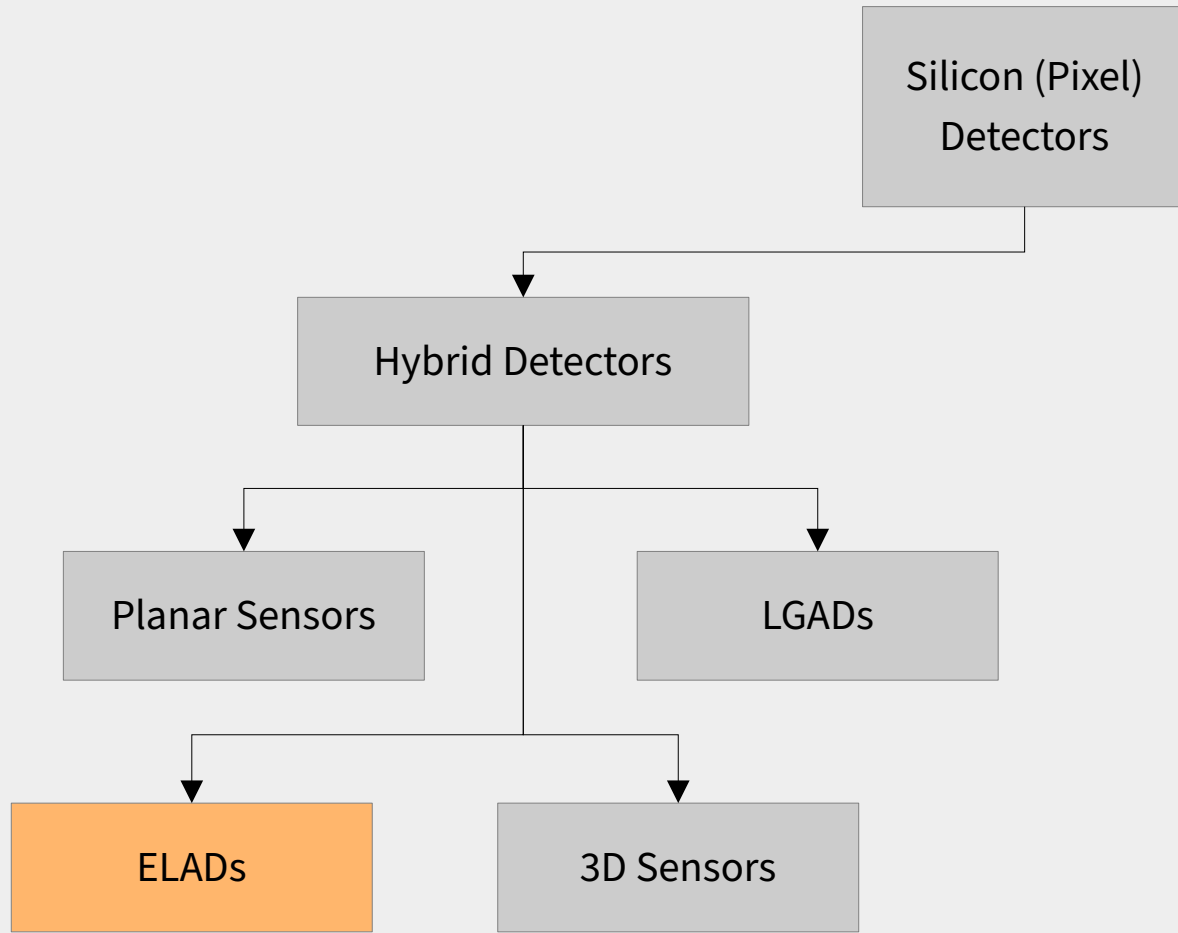
<https://doi.org/10.1016/j.nima.2022.166628>

LGAD Mortality – Burnout Events

- Observation: LGADs sometimes suddenly “die”, indication of single destructive event
- Proton testbeam allowed to pin-point impact point of track
- Clear correlation with “crater” location
- Preliminary explanation:
 - With high bias voltages, strong electric fields form locally
 - Combined with high-ionization events (Landau tail)
 - Enough to create conductive path across diode, burnout due to high current density
- Establishing of safe-operation conditions

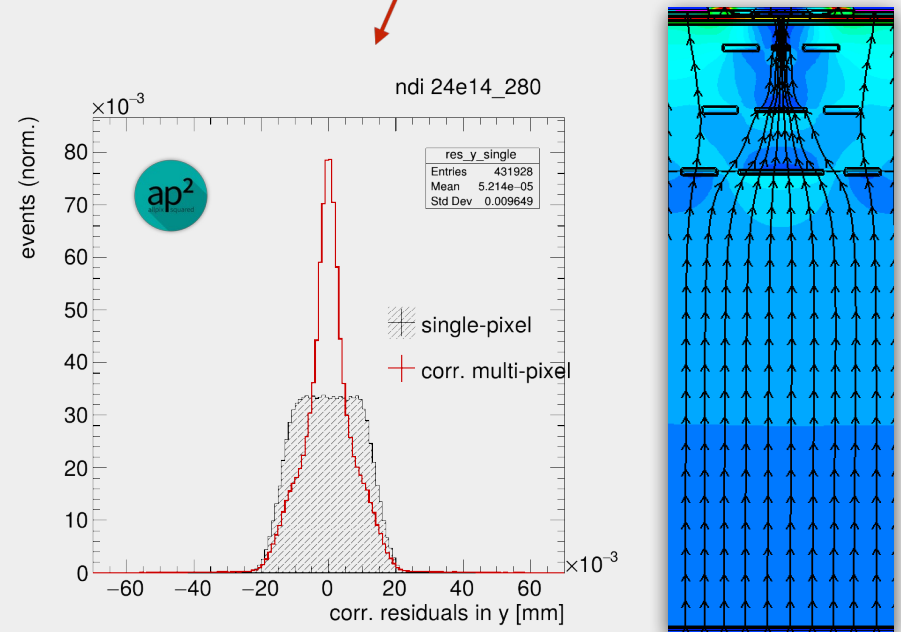
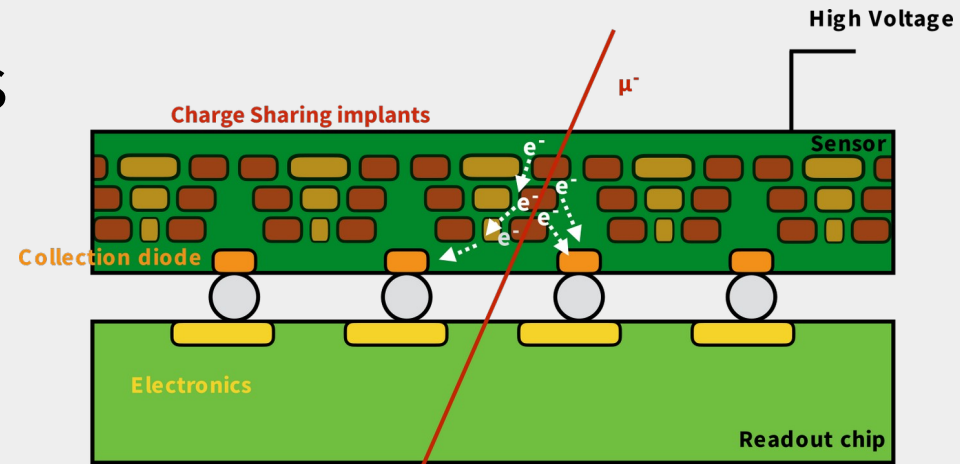
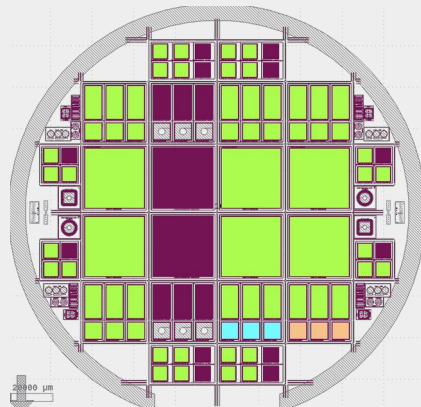


<https://indico.cern.ch/event/1058977/contributions/4865300/>



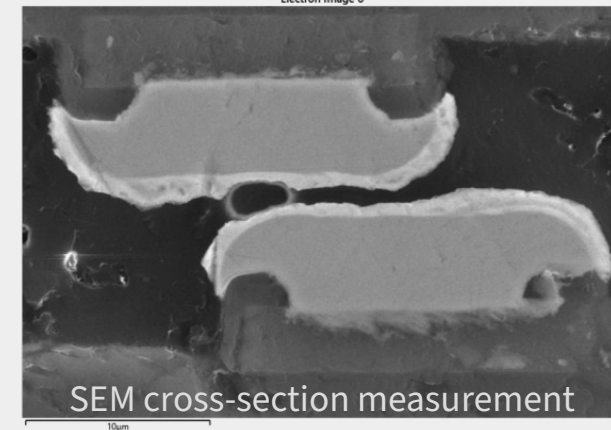
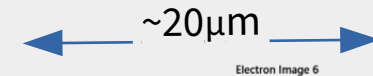
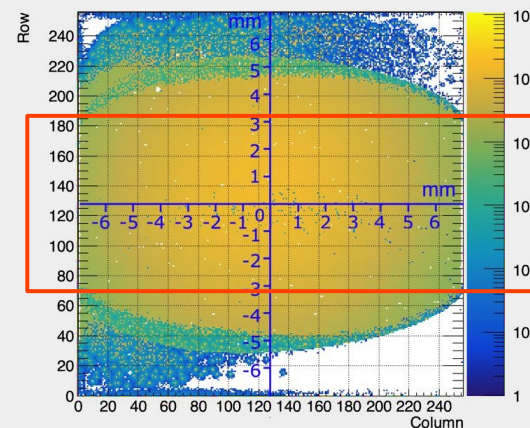
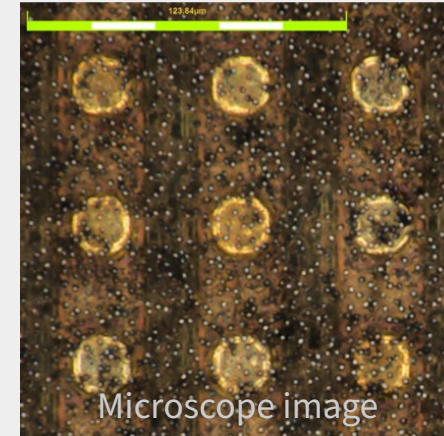
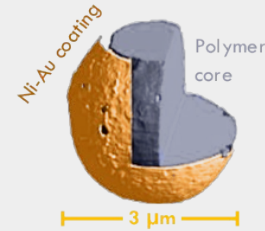
Enhanced Lateral Drift Detectors

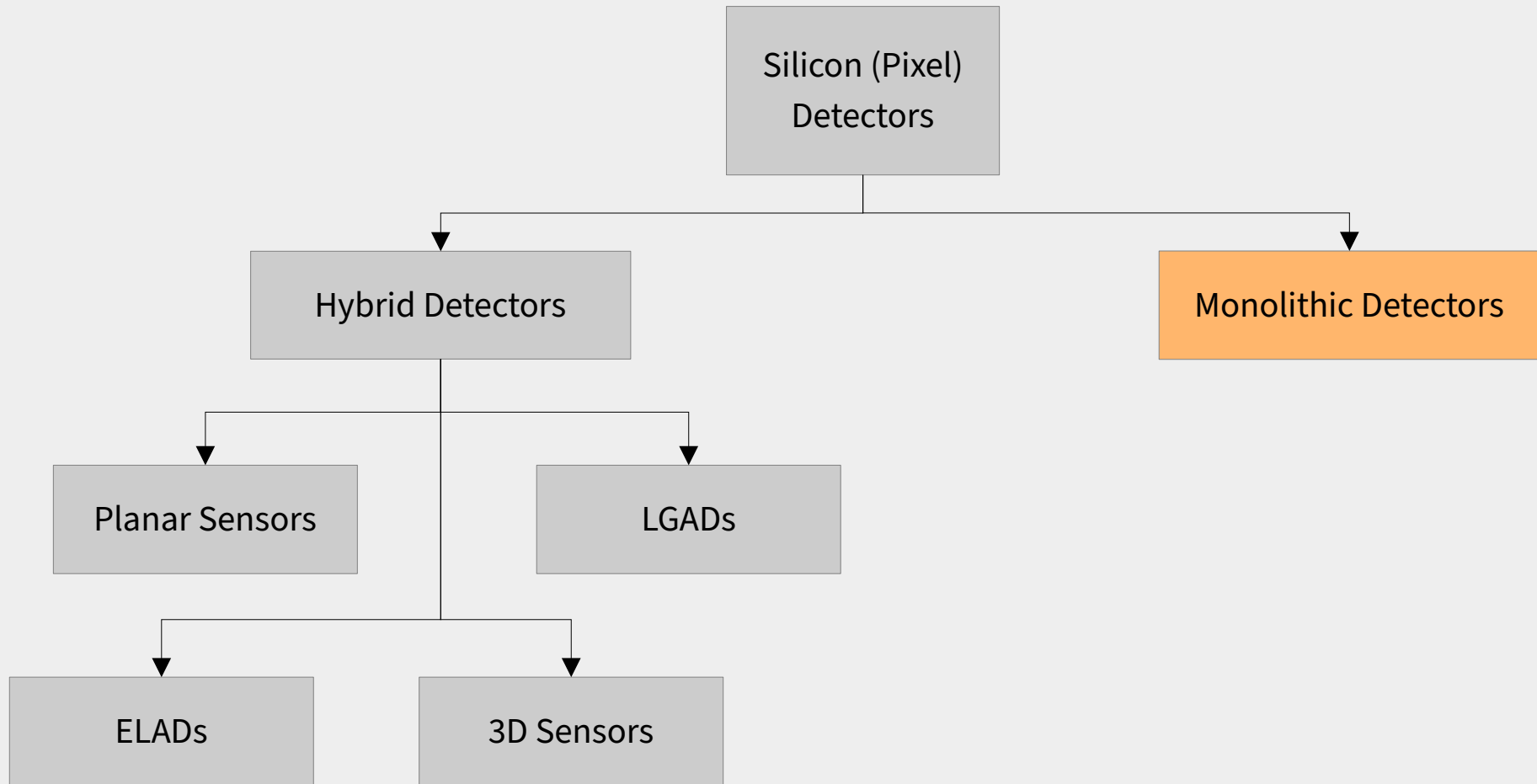
- Position resolution in **thin sensors** limited to $\text{pitch} / \sqrt{12}$ (almost no charge sharing)
- New concept: **enhance charge sharing** in Enhanced Lateral Drift sensors (ELAD)
 - Close to theoretical optimum: linear charge sharing
 - Deep implants to alter field, improve resolution
 - Lateral spread of charges during drift, cluster size ~ 2
- Challenges:
 - Complex production process
 - Low-field regions (recombination)



Hybridization with Anisotropic Conductive Film

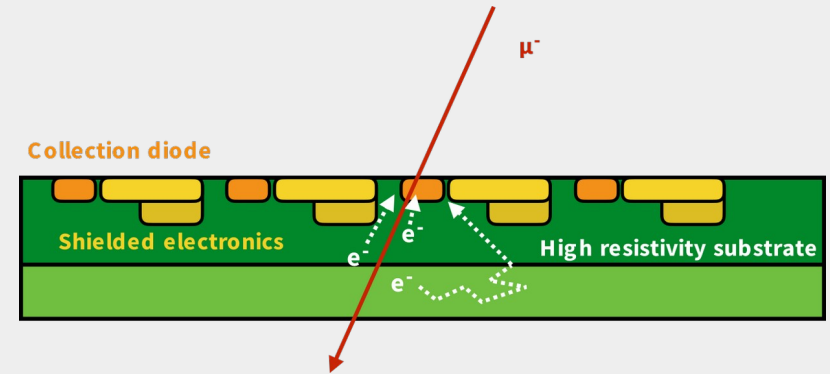
- Alternative to traditional solder-bump bonding
- Adhesive film with conductive micro-particles
 - Stochastically distributed in film
 - Some spheres end up under bond pads, get deformed, establish contact
- Widely used in display industry in one dimension, challenge: 2D distribution
- Requires careful optimization of
 - Film thickness
 - # spheres/area
 - Bonding force...
- Currently early R&D phase





Monolithic Silicon Detectors

- Also called Monolithic Active Pixel Sensors (MAPS)
 - Electronics & sensor: same wafer
 - Fully integrated: amplification, discrimination & readout
- Shield electronics via additional implants
- Different approaches
 - Deep collection diode surrounding electronics
 - Separate shielding & collection diode

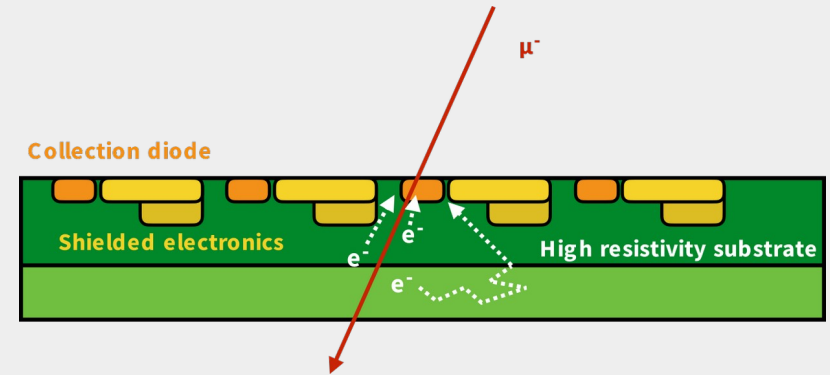
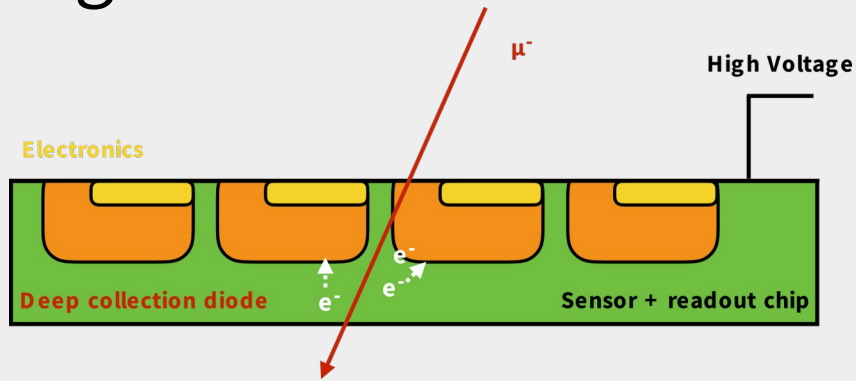


Lower mass than hybrids
No bump-bonding
Cheaper (large scale) manufacturing

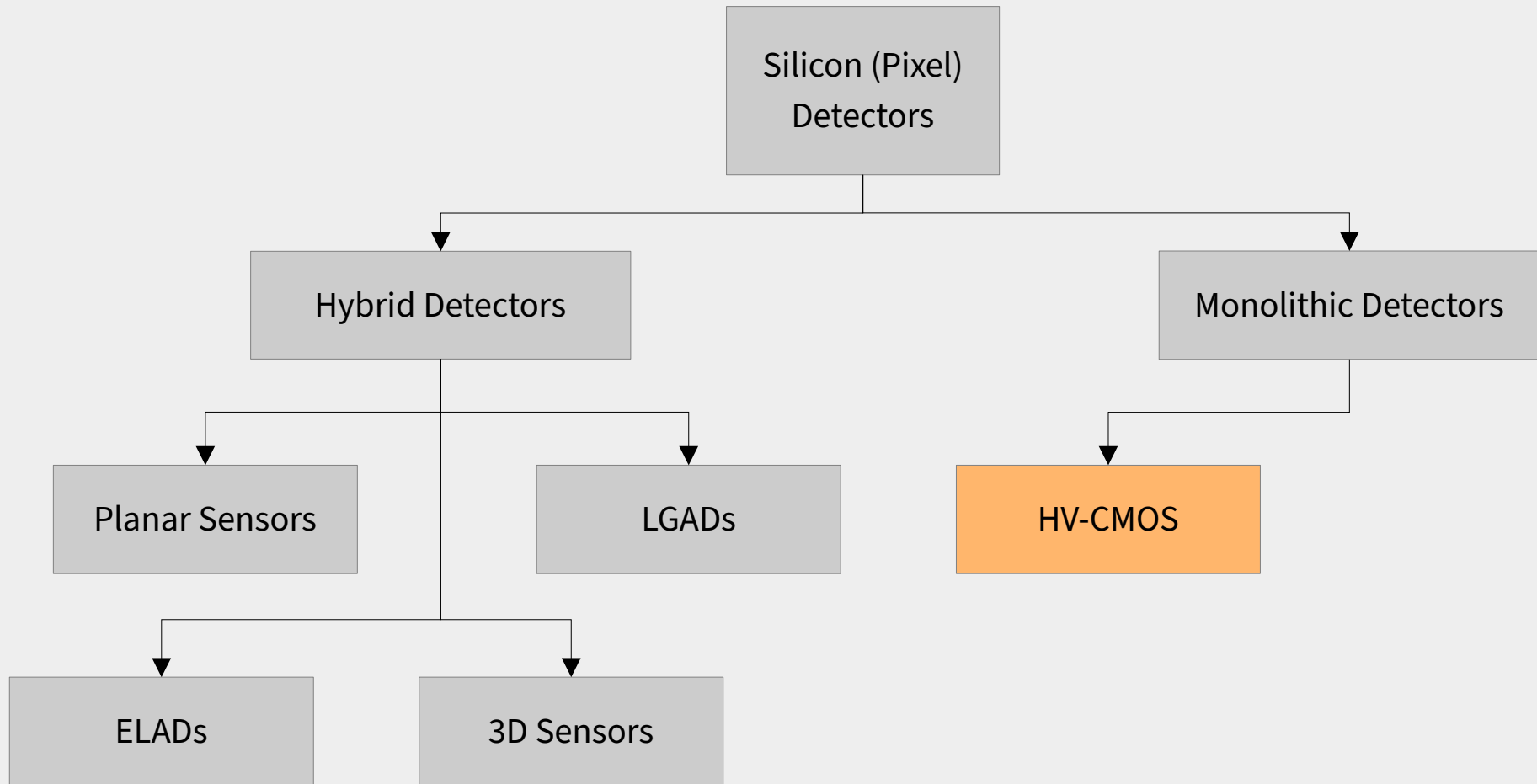


Smaller depletion volume & signal
Intricate sensor design
Limited in-pixel functionality

Large & Small Collection Electrode

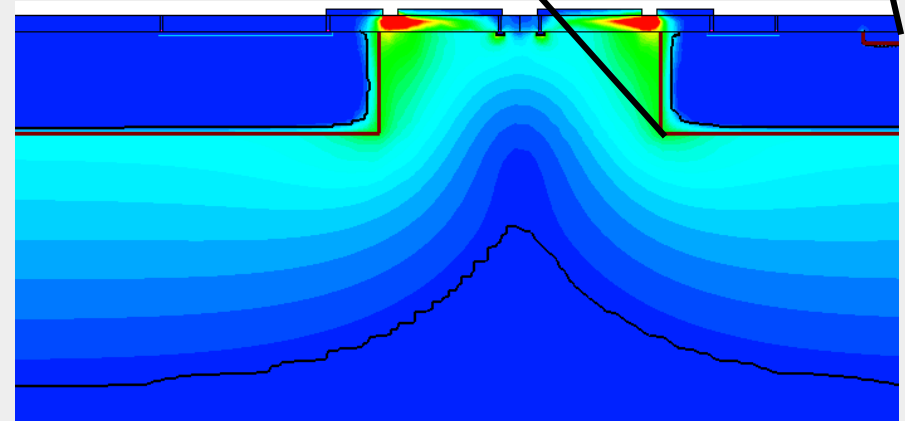
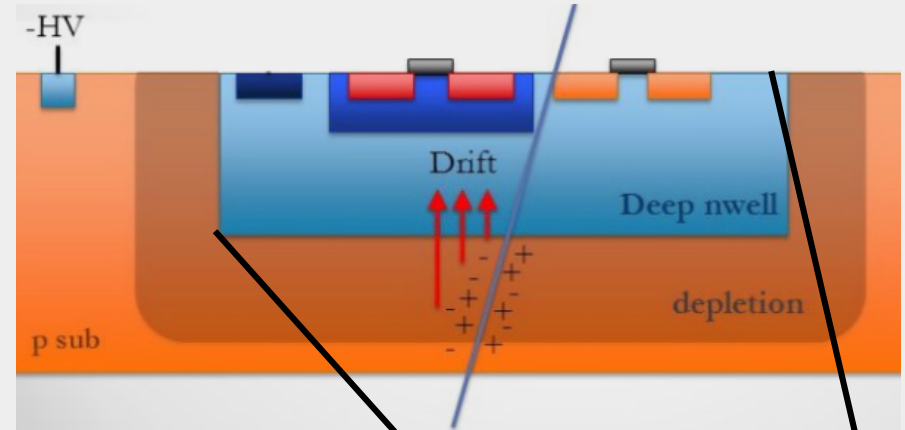


- Shield electronics via deep collection diode around electronics
 - Allows high bias voltage to be applied
 - Fast & large signal, large depletion volume
- Challenge: large collection diode leads to
 - large input capacitance
 - increased power consumption
- Electronics outside charge-collection well
 - Requires high-resistivity material (e.g. epitaxial layer) to allow depletion
 - Small collection diode leads to small capacitance
- Challenge: effect of p-well potential on electric field / charge collection



HV-CMOS Sensors

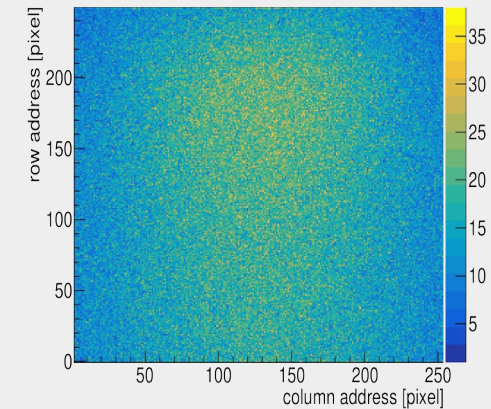
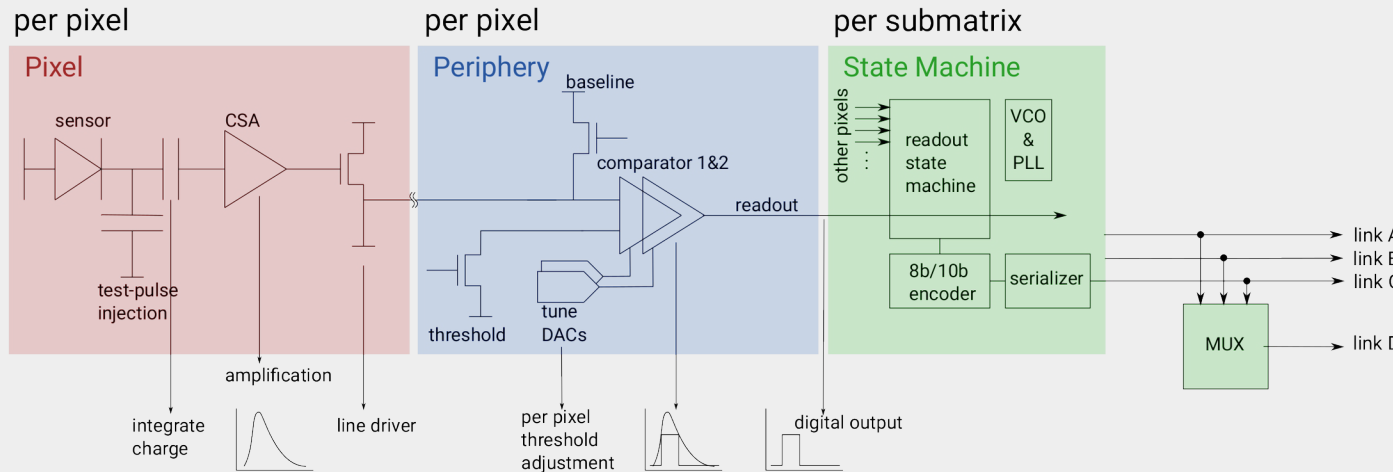
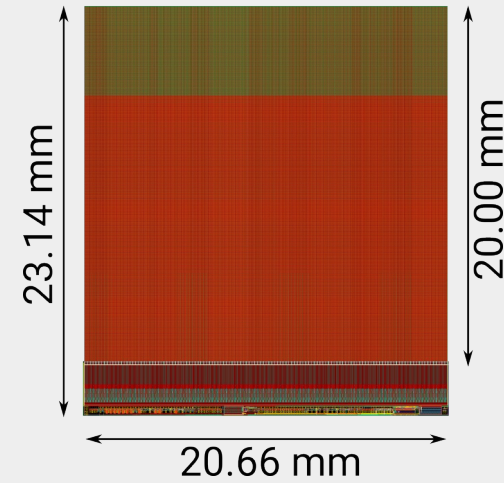
- Place all transistors inside deep well
- Deep well acts as electrode
 - pn -junction forms between deep well and substrate
- Possible to apply high voltage $O(10-100V)$ over this diode
 - High electric field allows for fast charge collection via drift
 - Need to be careful about strong fields near surface



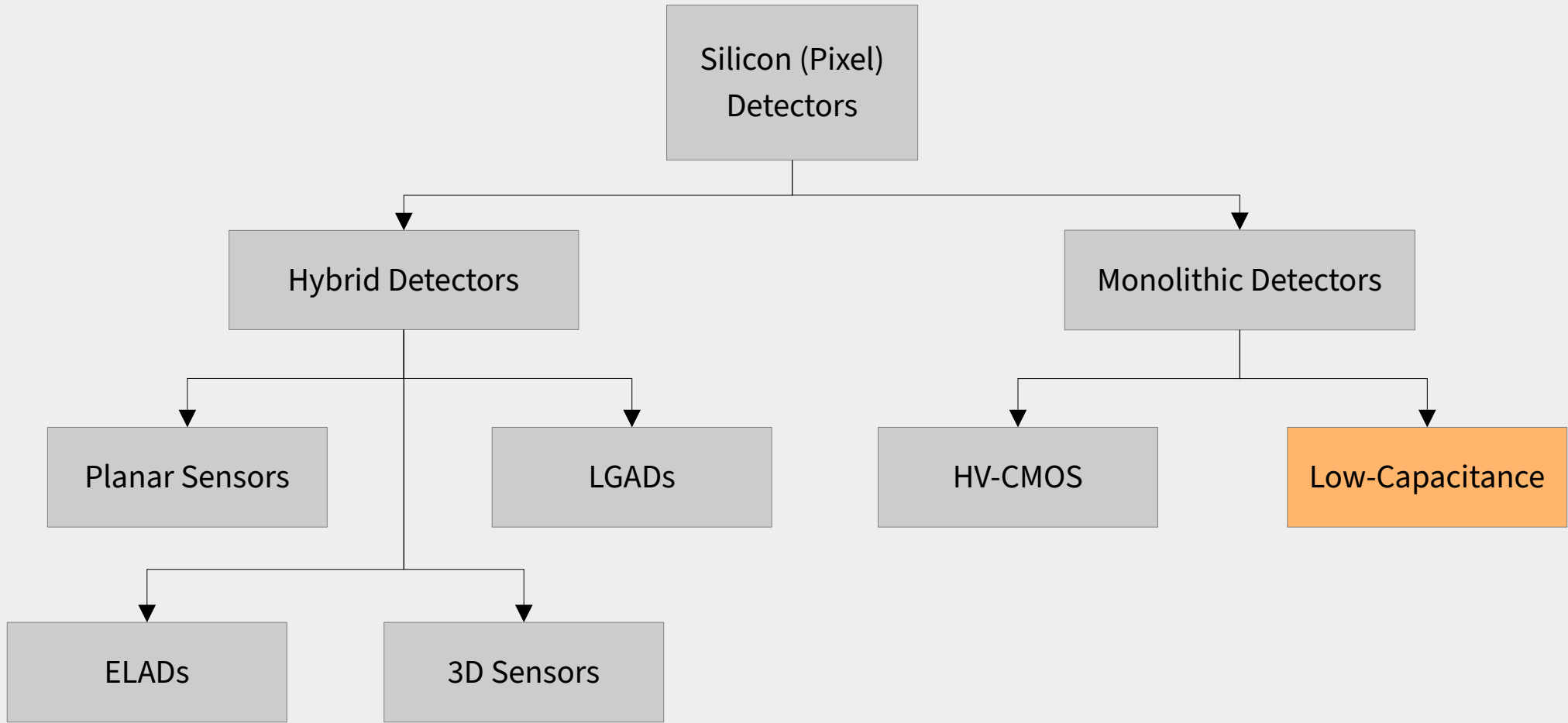
Annie Meneses

The MuPix10 HV-MAPS

- Designed for the mu3e experiment at PSI
 - 180 nm HV-CMOS technology
 - Matrix of 256 x 250 pixels with pitch of 80 x 80 μm^2
 - Energy + time measurement per pixel

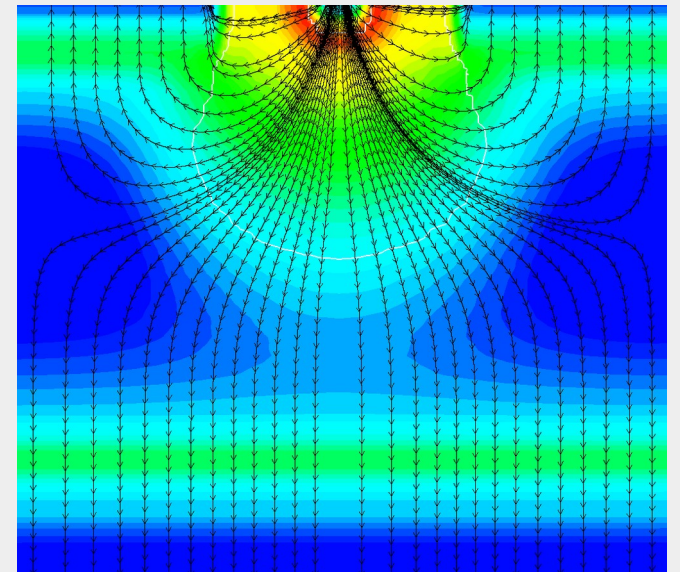
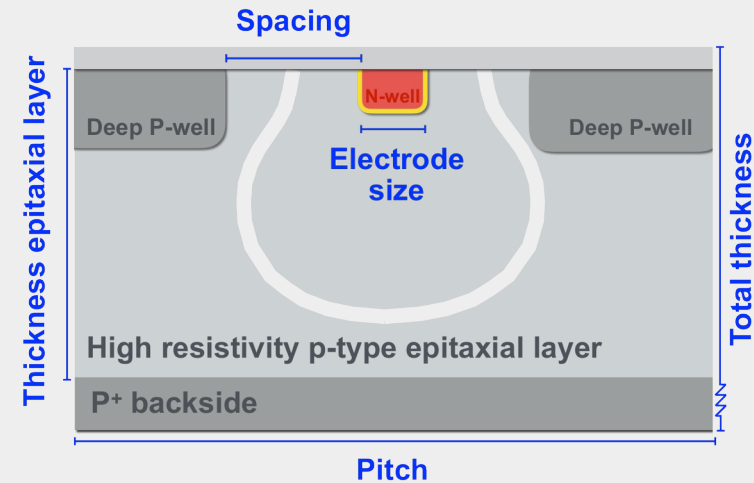


<https://arxiv.org/abs/2012.05868>



Low-Capacitance CMOS Sensors

- Instead of using well as electrode:
place next to deep well
 - Small sensor capacitance, low noise $O(\sim 10e)$
 - Requires high-resistivity material
 - No requirement for HV-compatible process
- Can only apply low bias voltage $O(< 6V)$
 - Only small volume depleted,
partially collecting charge via diffusion
 - Requires detailed field optimization

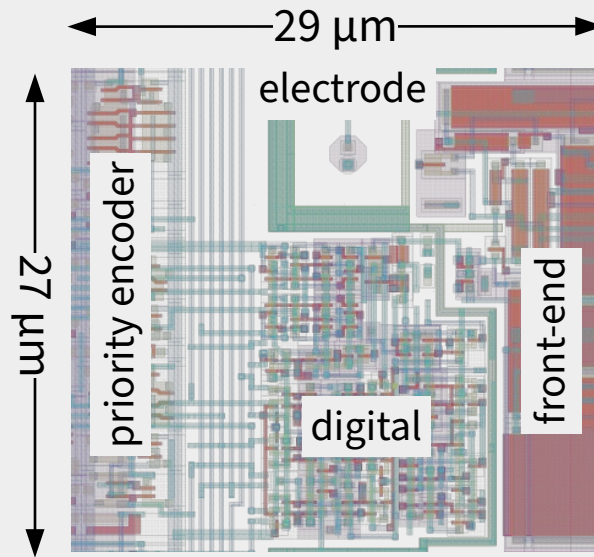
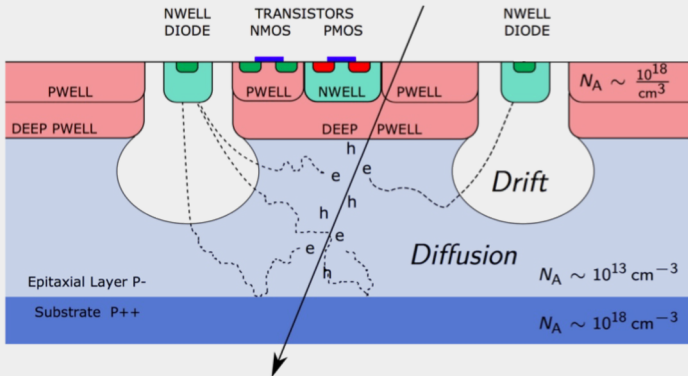
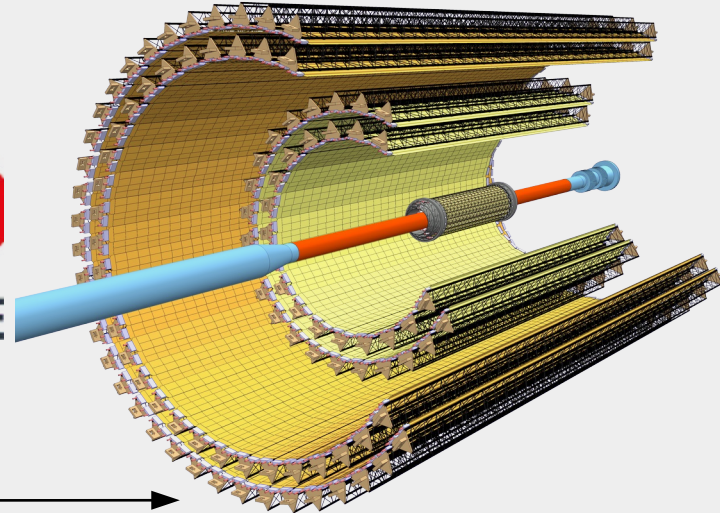


The ALPIDE Sensor of the ALICE ITS2

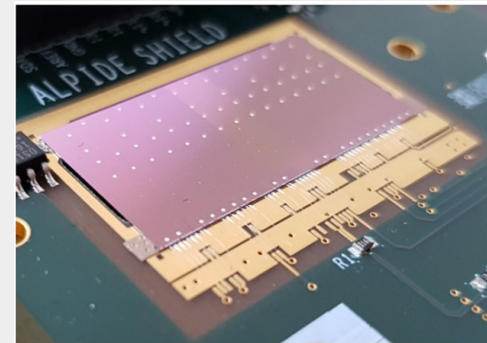
- Full Inner Tracking System: 24'000 ALPIDE chips, one of the first large-scale detectors with MAPS
- ALPIDE – MAPS in 180 nm CMOS imaging technology
 - 512 × 1024 pixels of 29 μm x 27 μm pitch
 - Binary detection & readout (hit/no hit)
 - Optimized for low power consumption
 - Produced on epitaxial layers of 18 – 30 μm



ALICE



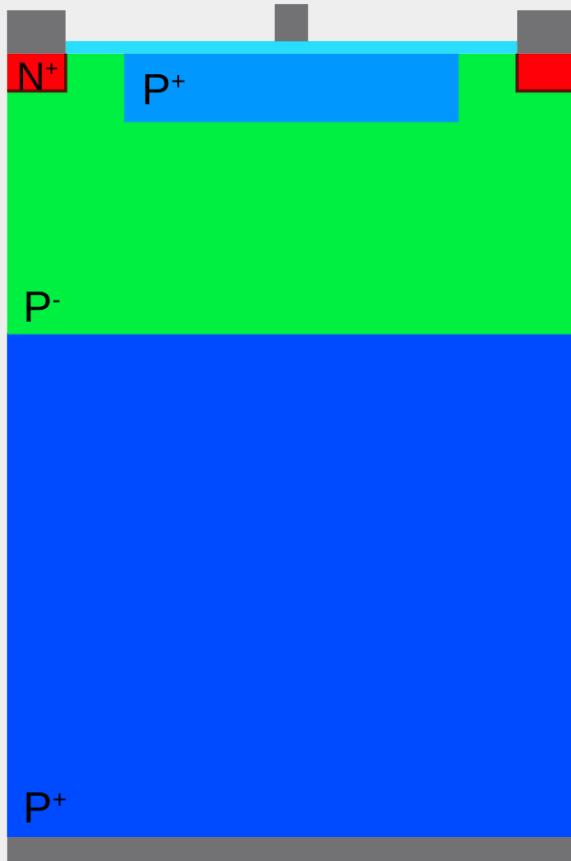
ALPIDE pixel cell



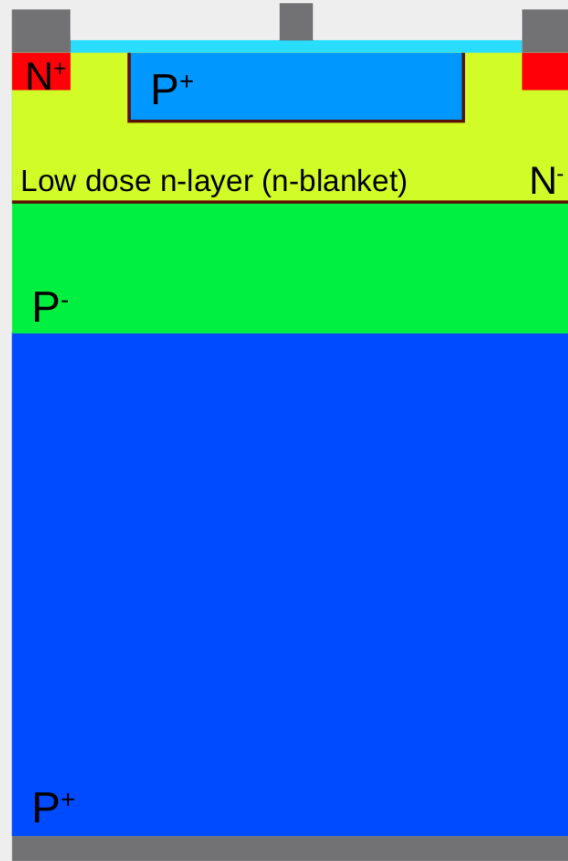
<http://dx.doi.org/10.1016/j.nima.2016.05.016>

Sensor Layouts – Improving Signal Formation

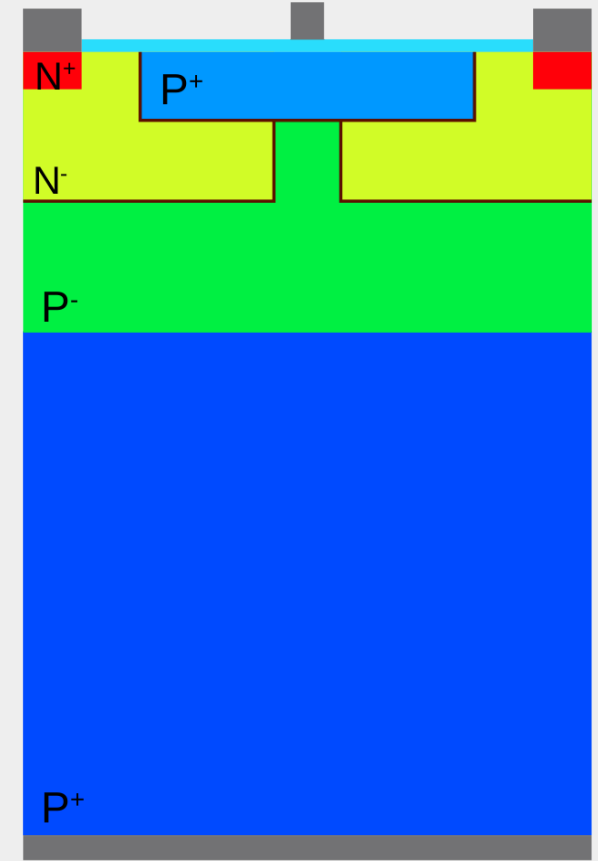
Standard Layout



N-blanket Layout

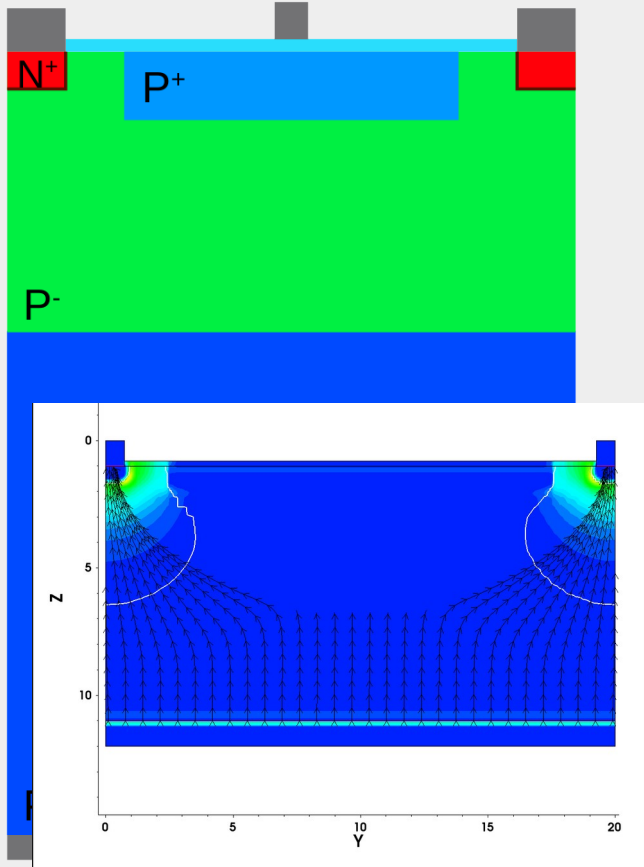


N-gap Layout

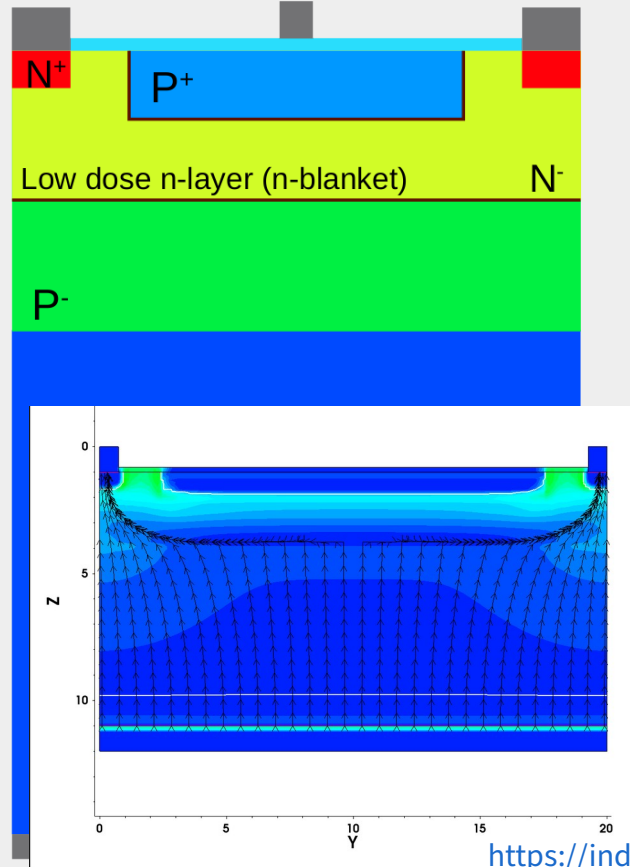


Sensor Layouts – Improving Signal Formation

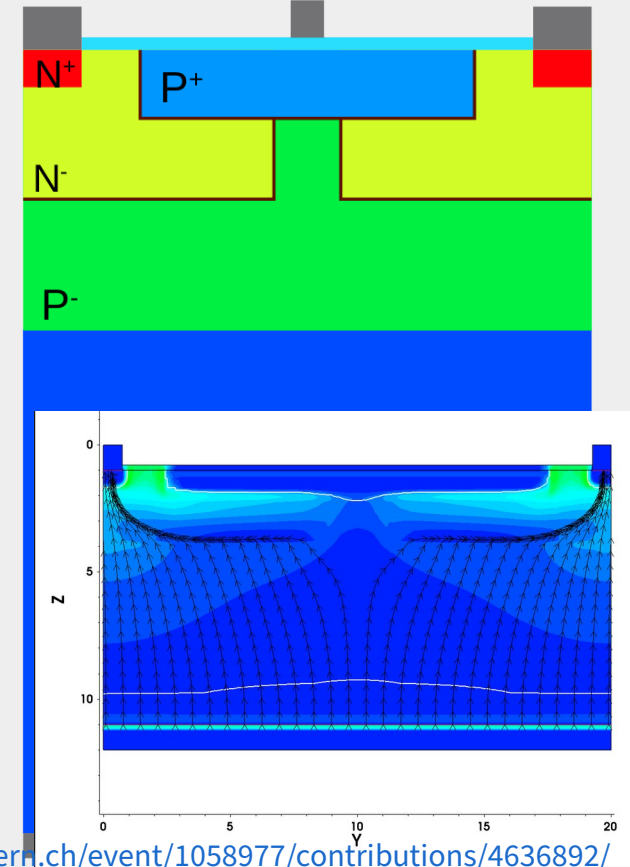
Standard Layout



N-blanket Layout



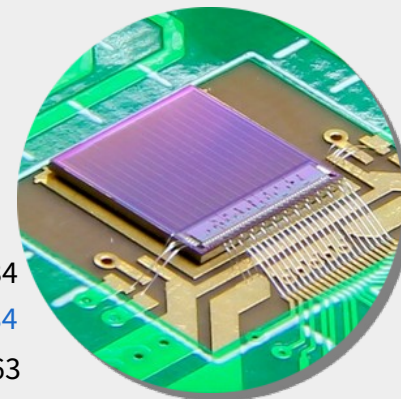
N-gap Layout



<https://indico.cern.ch/event/1058977/contributions/4636892/>

The CLICTD MAPS Prototype

- Silicon detector prototype for CLIC tracking detector
 - 180 nm CMOS imaging process, small collection electrode
 - Pixel pitch: $37.5 \mu\text{m} \times 30 \mu\text{m}$, $30 \mu\text{m}$ epitaxial layer
 - Fully-integrated sensors, ToA/ToT measurement



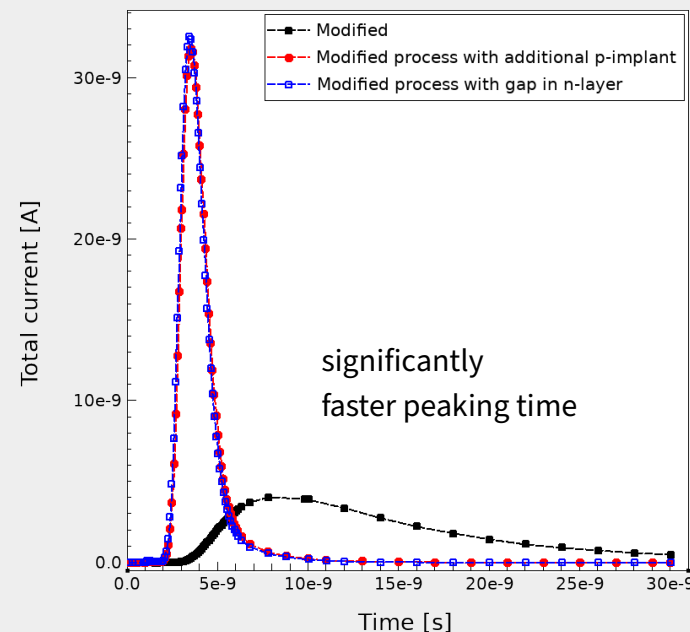
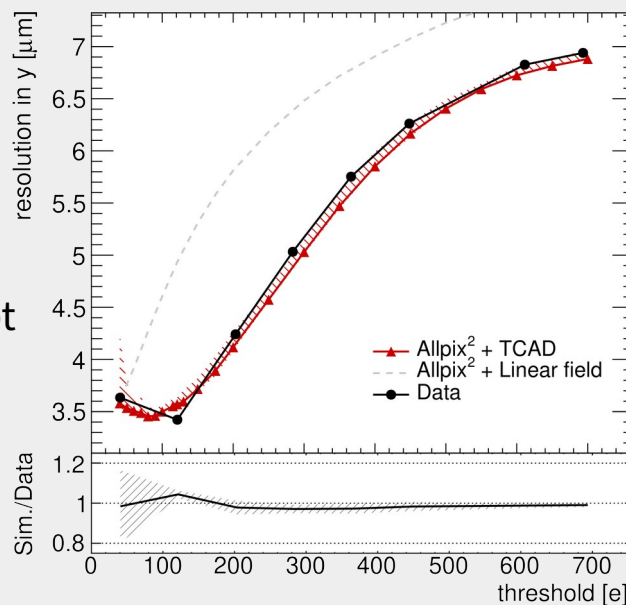
NIMA 964 (2020) 163784

[doi:10.1016/j.nima.2020.163784](https://doi.org/10.1016/j.nima.2020.163784)

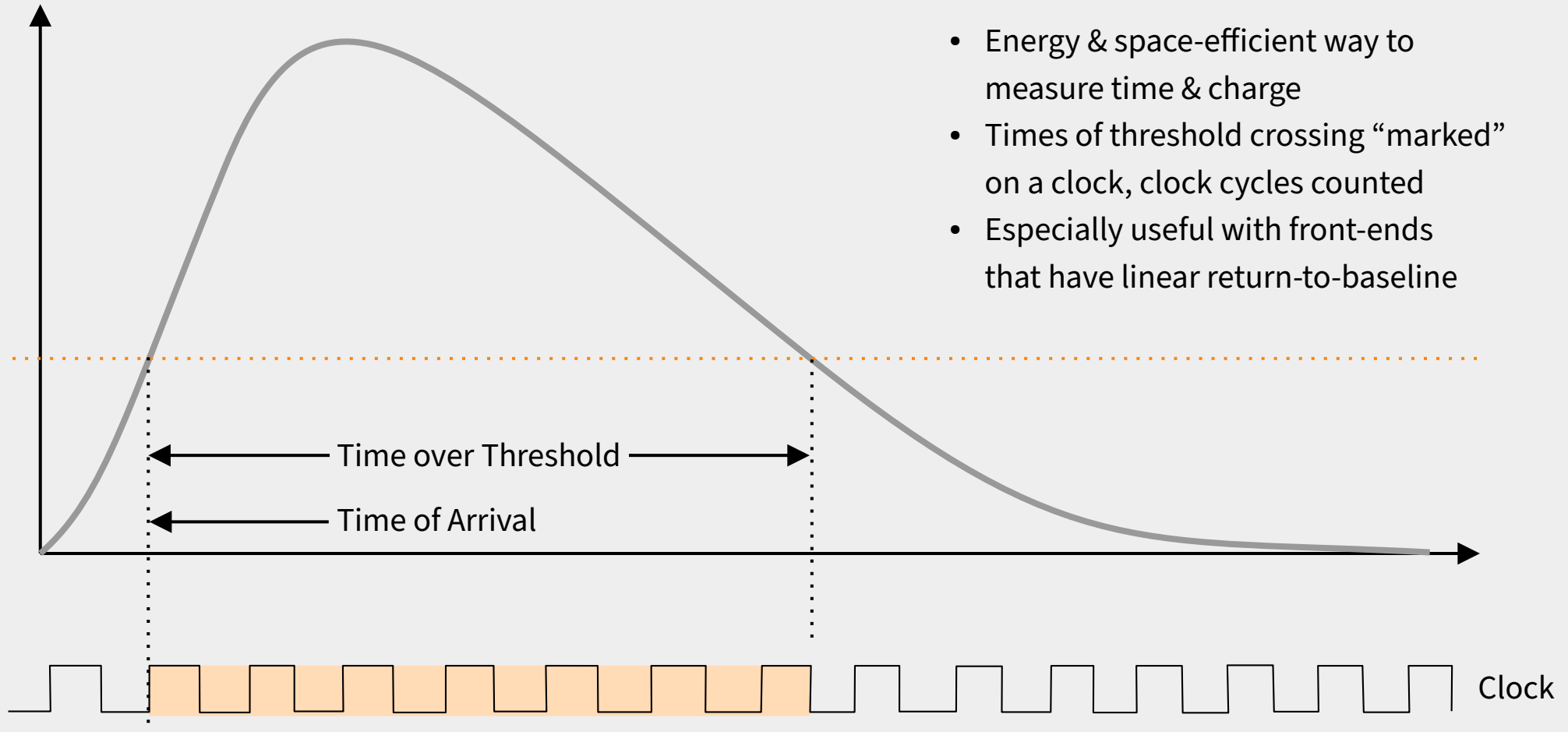
IEEE TNS, vol. 67, no. 10 (2020), 2263

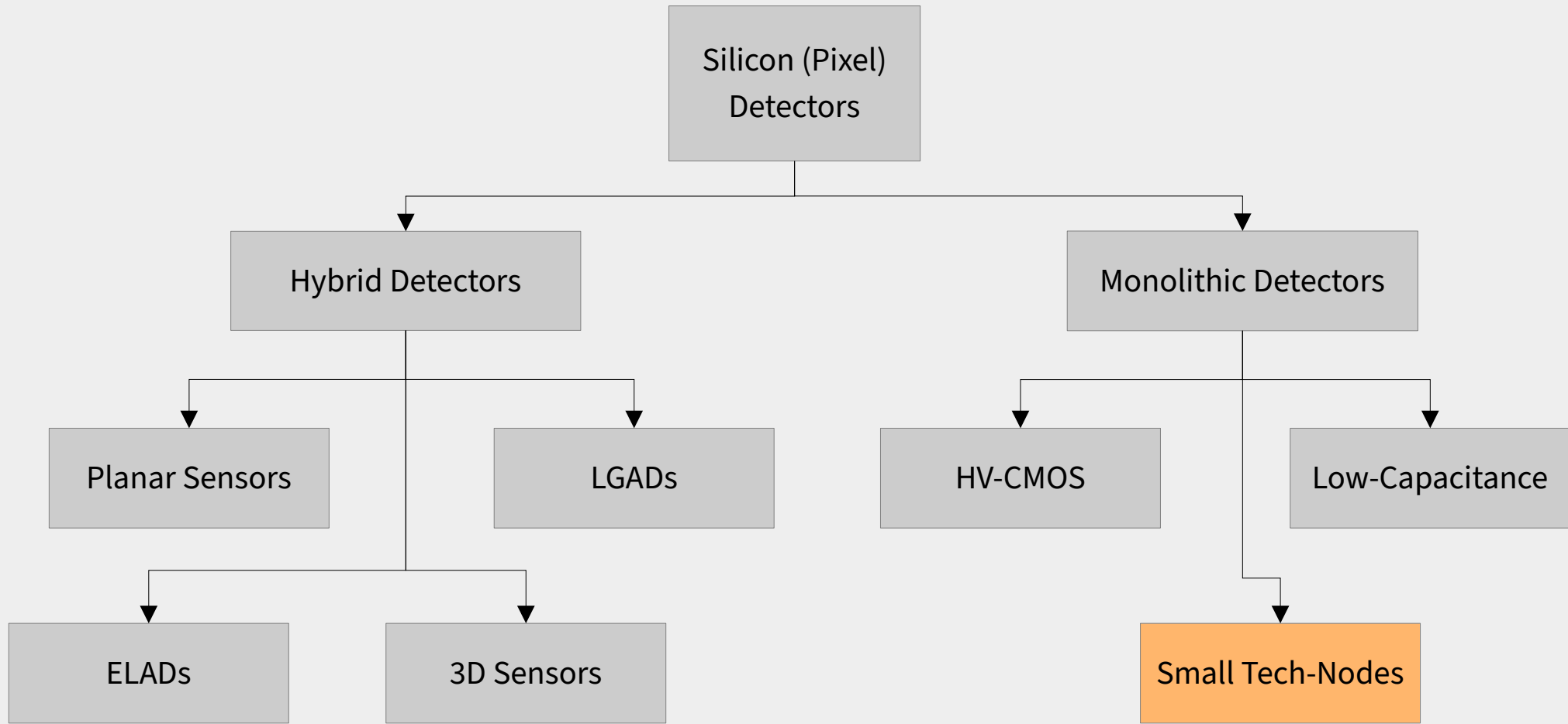
[doi:10.1109/TNS.2020.3019887](https://doi.org/10.1109/TNS.2020.3019887)

- Test bench for sensor designs
 - Optimizations e.g. for prompt signal formation
 - Spatial resolution studies



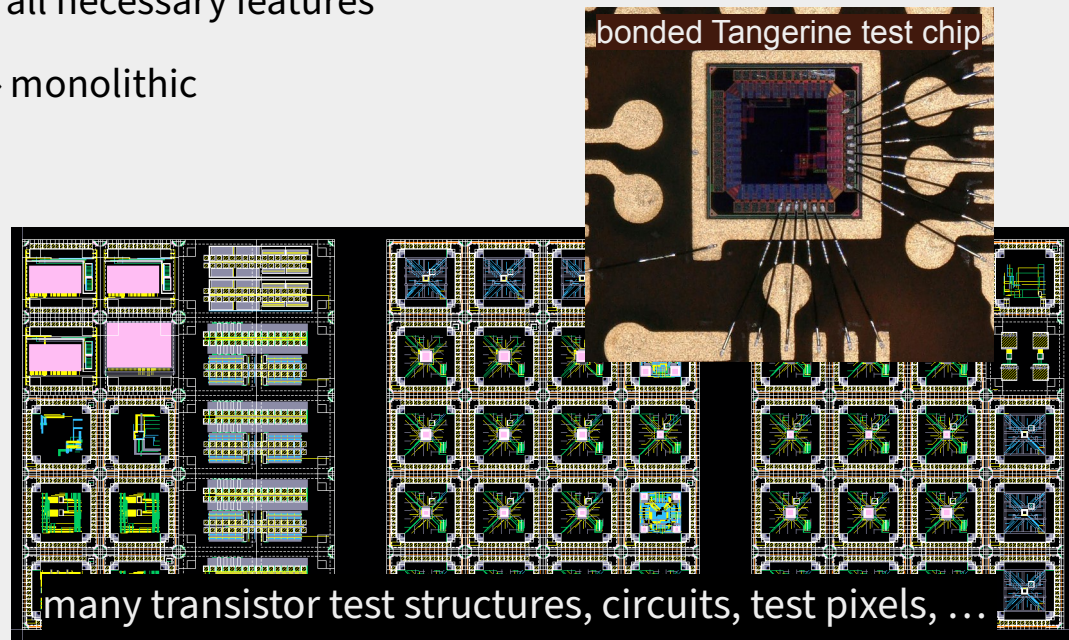
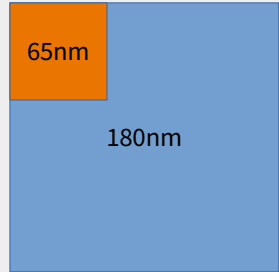
Time of Arrival & Time over Threshold





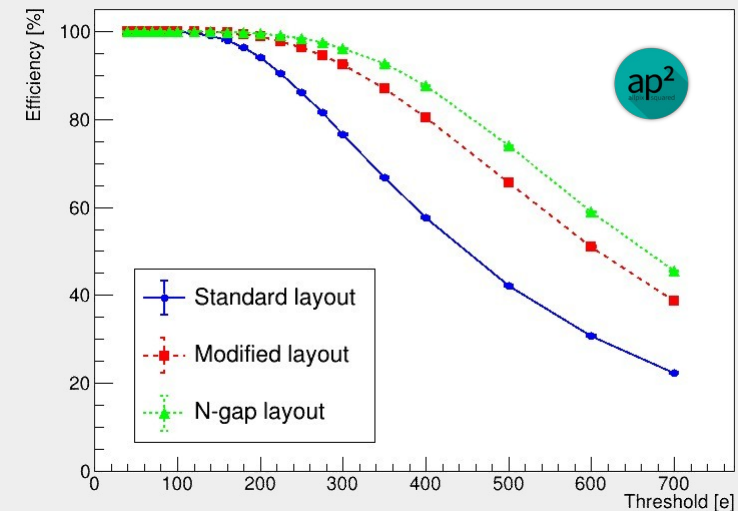
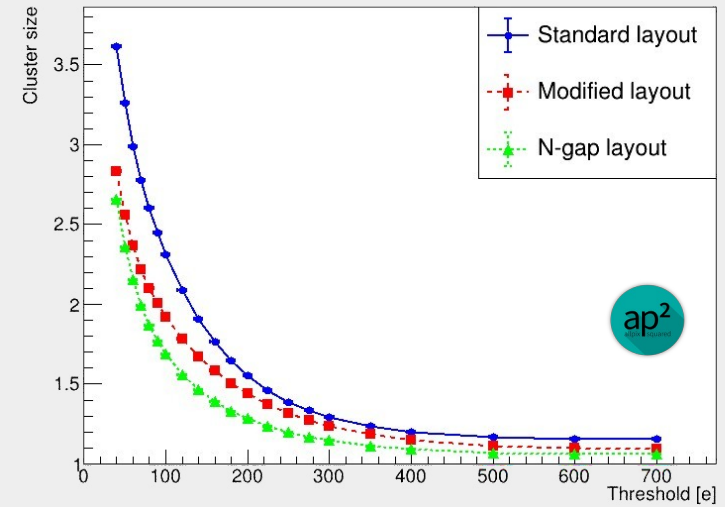
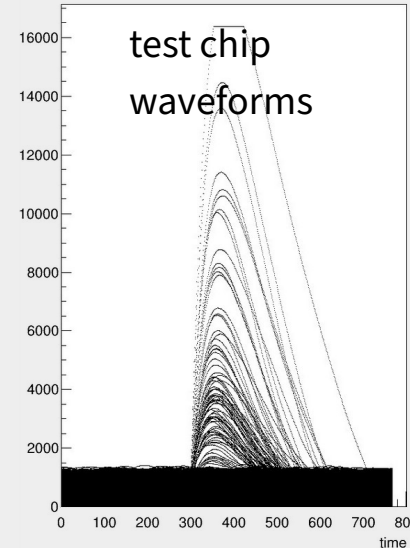
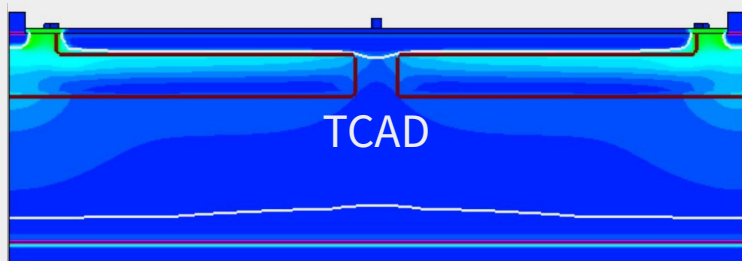
Smaller CMOS Nodes for Increased Logic Density

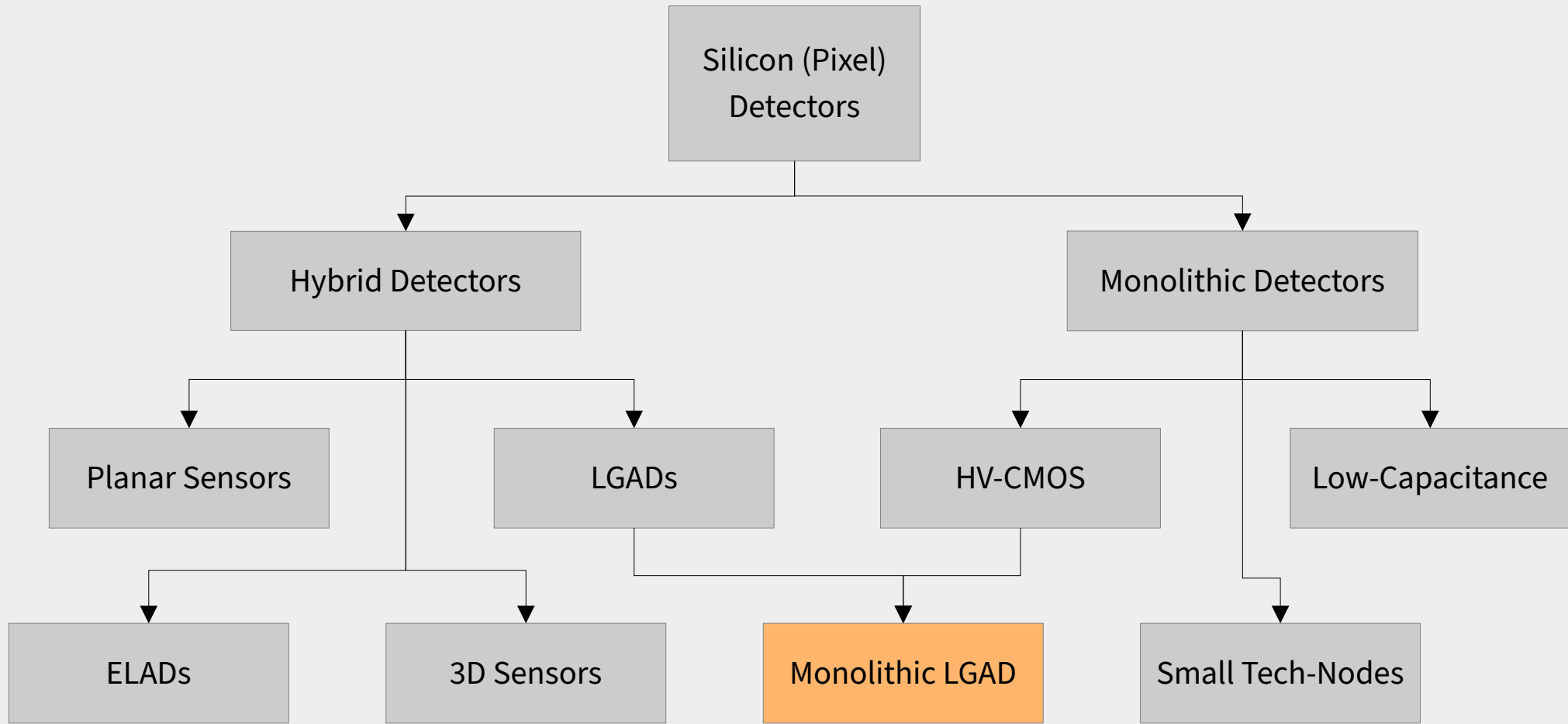
- **New CMOS imaging technologies** become available from foundries
- E.g. 65nm CMOS imaging process – under investigation for HEP detectors
 - Lower analog/digital power consumption
 - Higher logic density, allow pixels $\leq 25 \mu\text{m}$ with all necessary features
 - Bridging gap in front-end capabilities hybrid \leftrightarrow monolithic
- Issues to overcome:
 - Accessibility of technologies from vendors
 - Analog circuitry stays almost same size, only digital parts scale well
 - Needs to be **tested as particle detectors!**



Understanding a new Technology

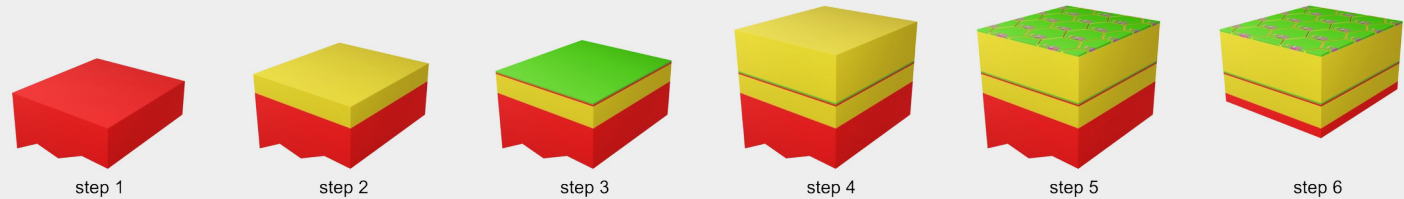
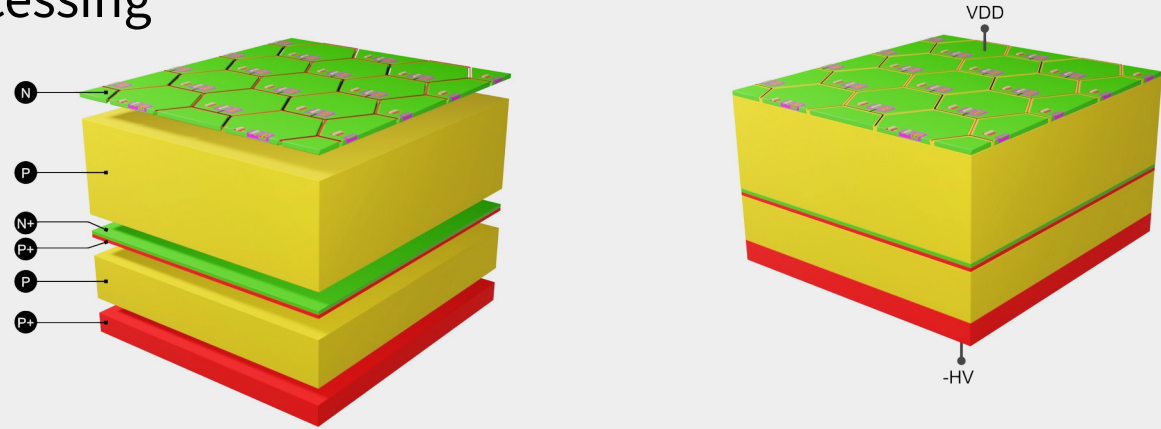
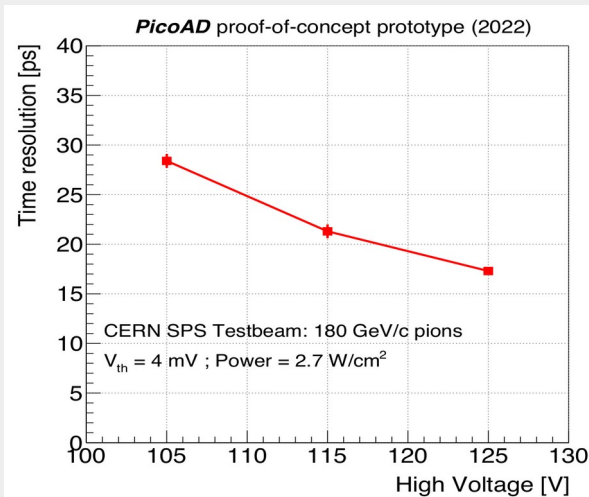
- Extensive **TCAD & Monte Carlo simulations** to optimize sensor response
 - Signal formation time
 - Efficiency
 - Cluster size, resolution
- **Compare to data from prototypes** to gain understanding of the technology





Timing in MAPS Sensors

- Investigation of embedding LGAD-like gain layer in MAPS
- Requires multi-step processing
- Promising results from PicoAD prototype



<https://indico.cern.ch/event/1058977/contributions/4631550/>

New Sensors Open Up New Opportunities

