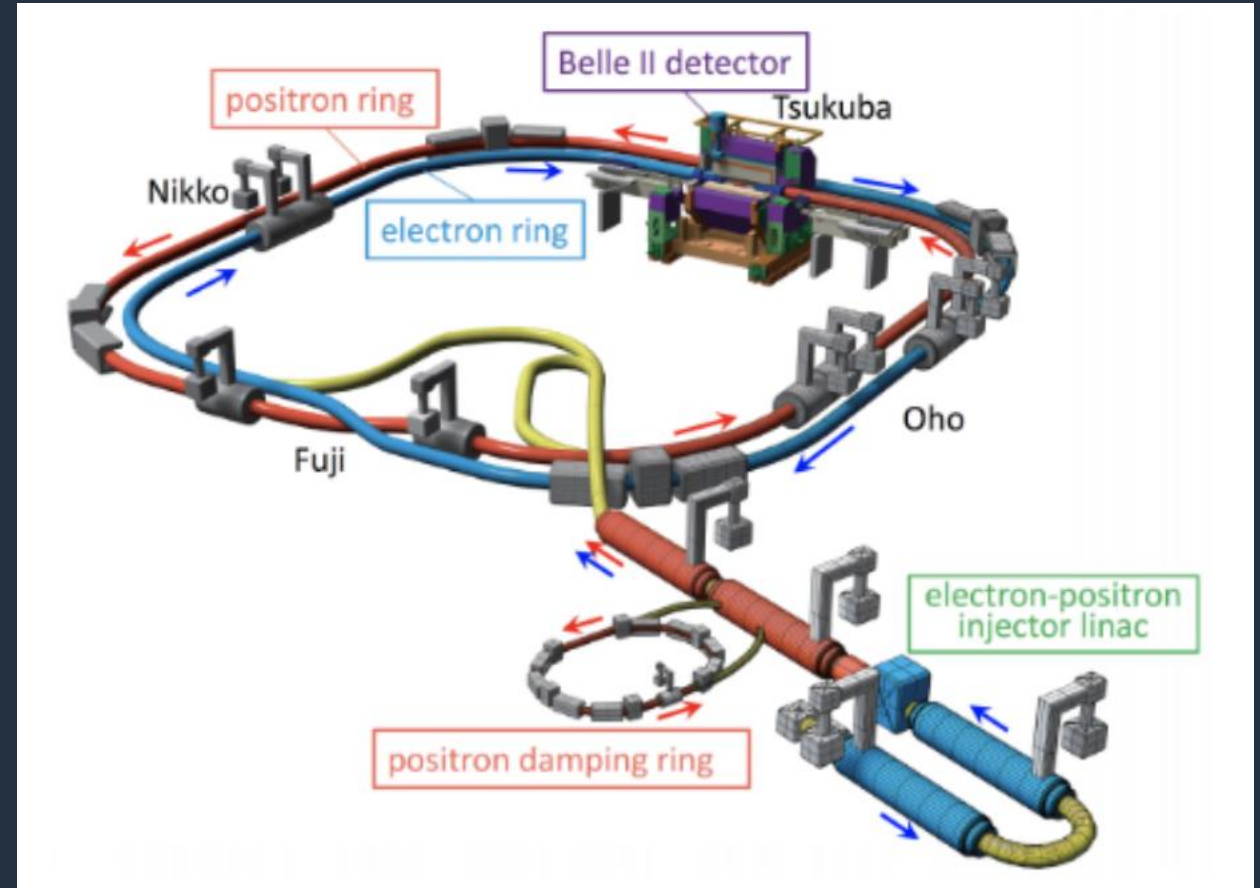
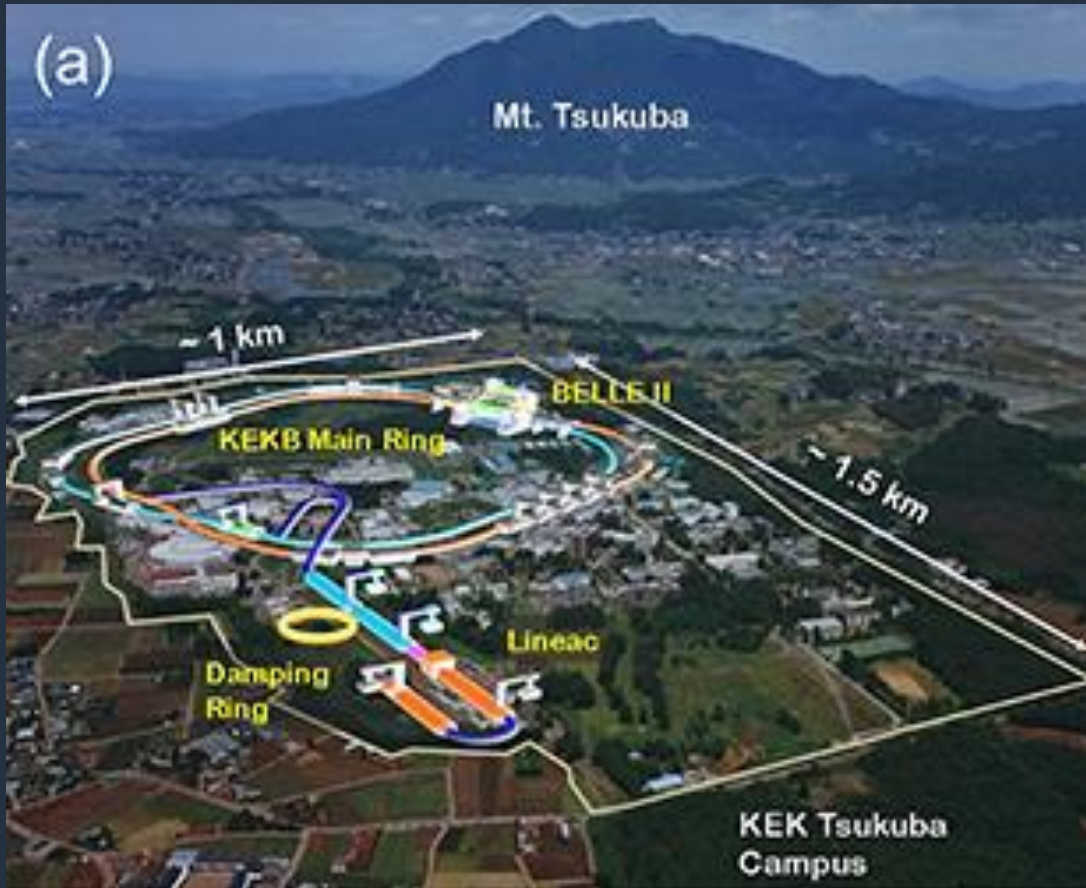


Experimental setup for high- precision laser polarisation determination

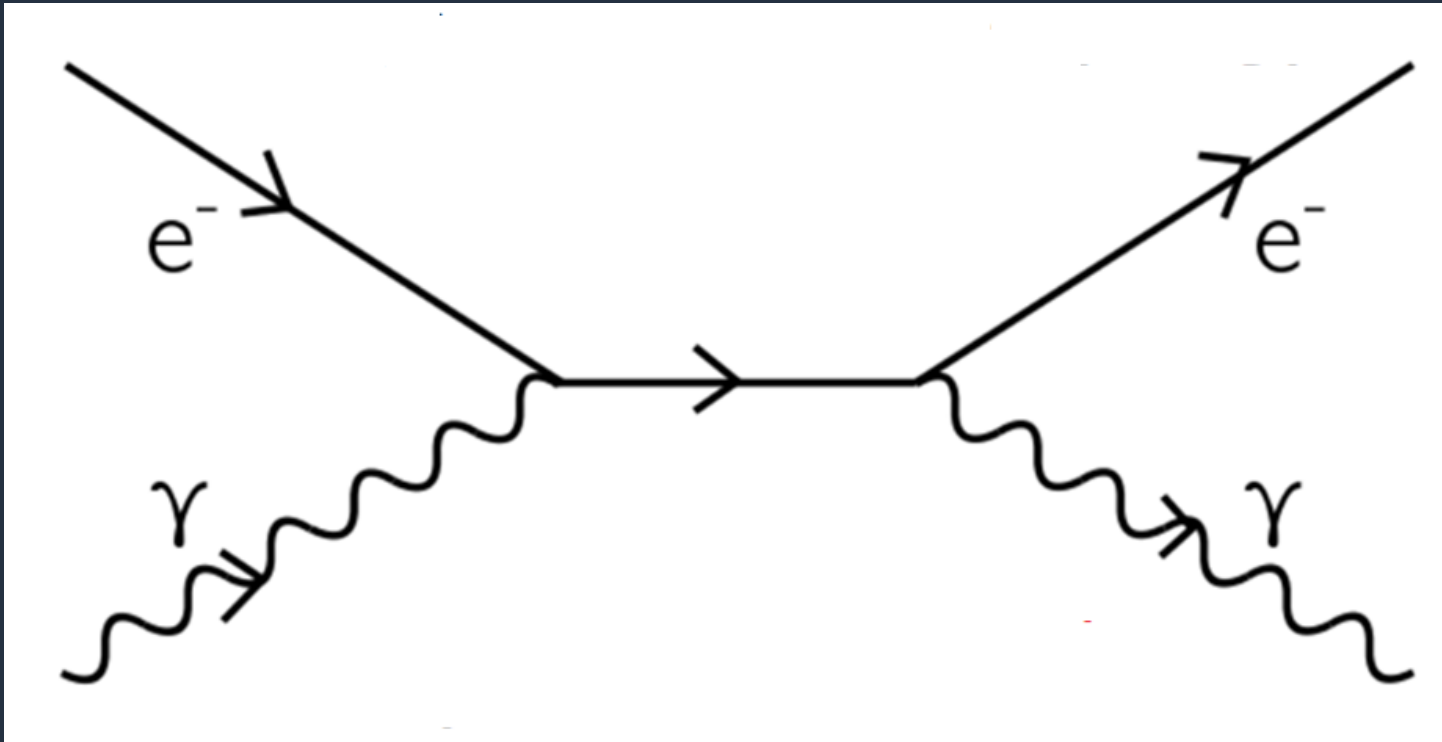
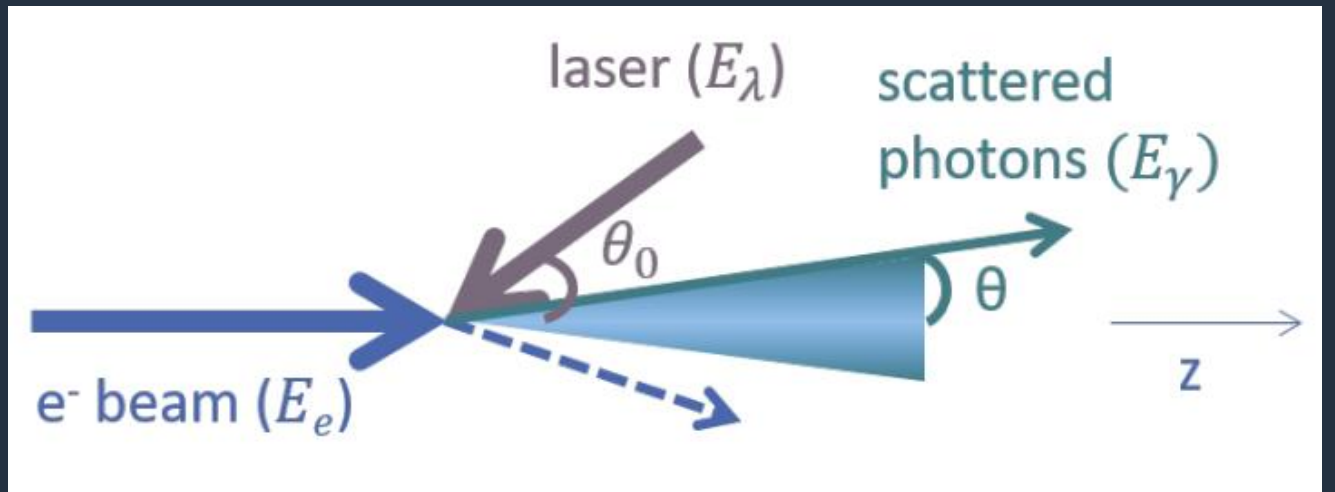
Santiago Rodríguez

Aurélien Martens

MOTIVATION: We need to measure the polarization of high energy e- beams

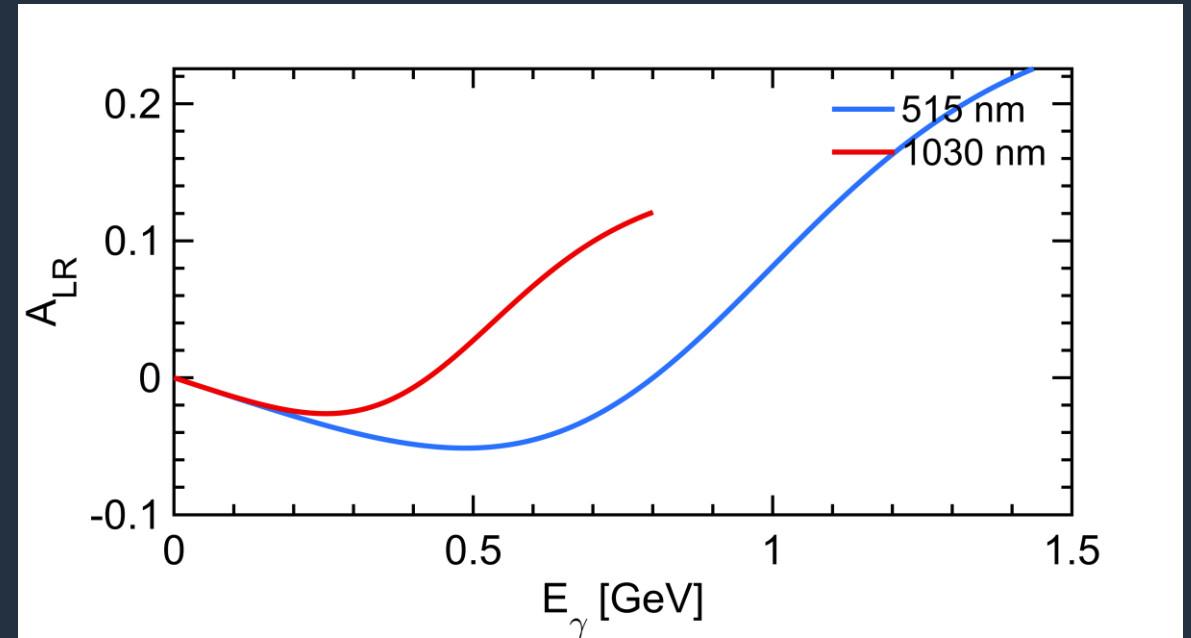
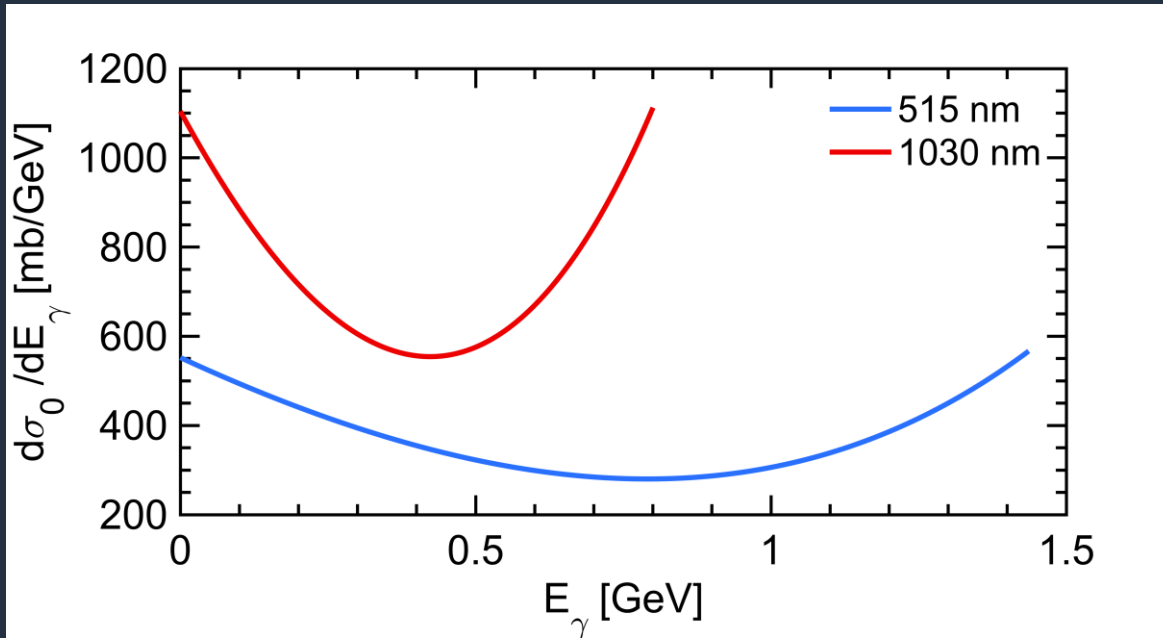


Inverse Compton Scattering



Differential cross section

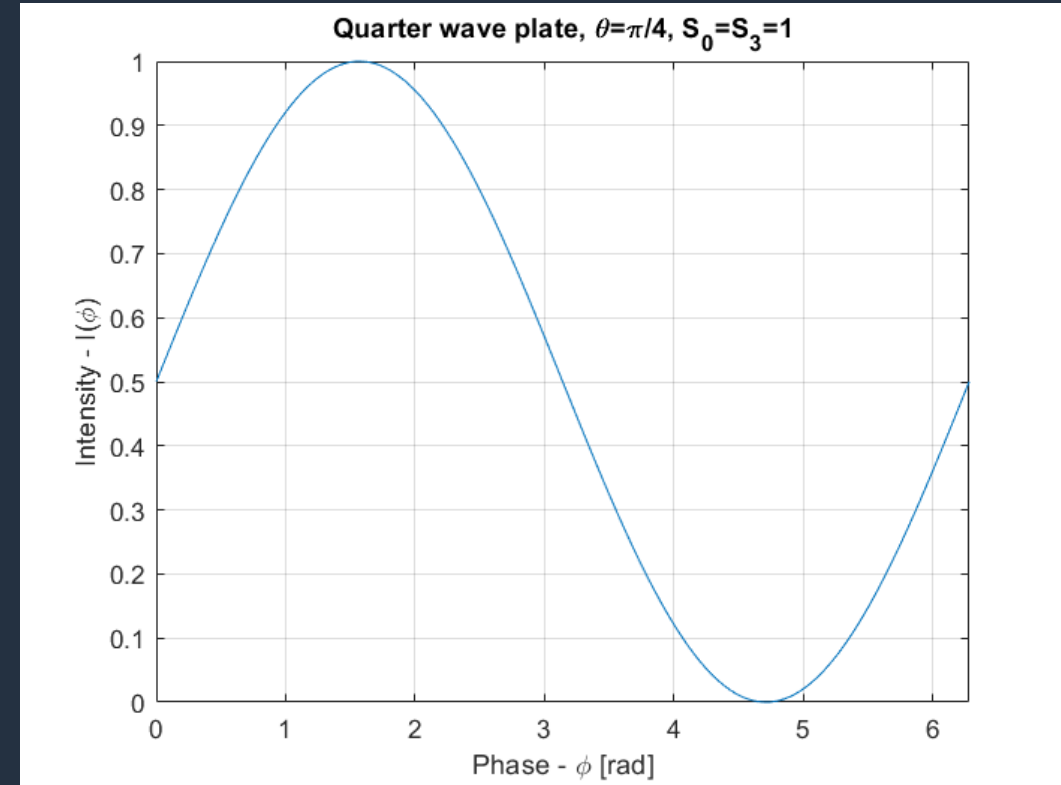
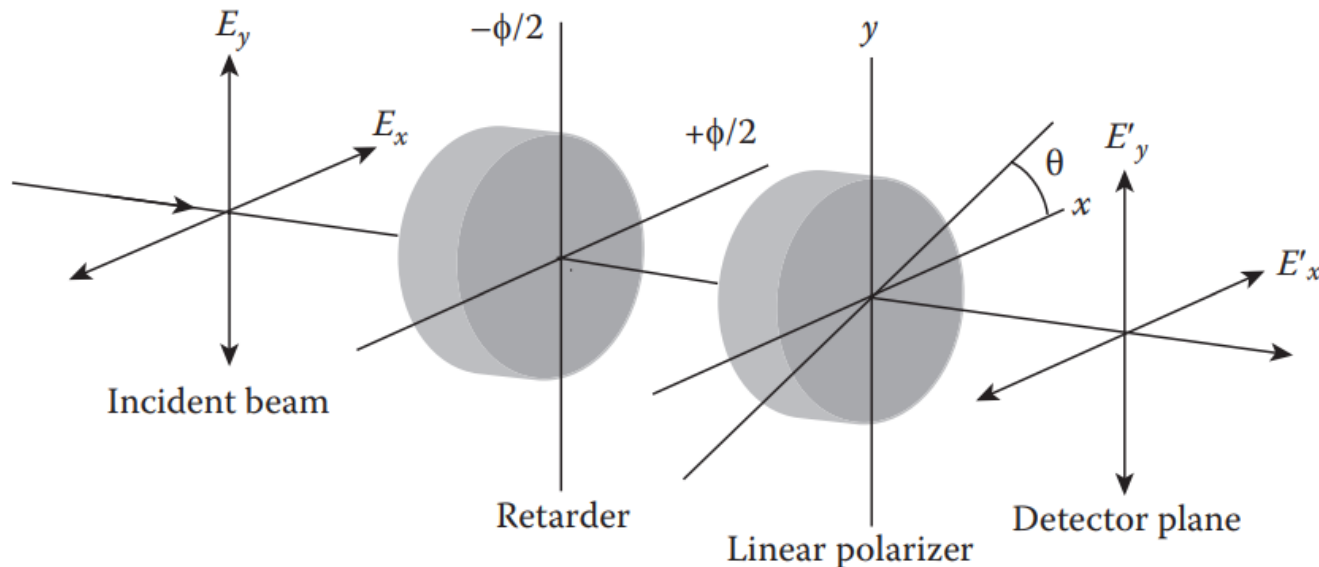
$$\frac{d\sigma}{dE_\gamma}(E_\gamma) \approx \frac{d\sigma_0}{dE_\gamma}(1 + P_Z P_C A_{LR})$$



- **Asymmetric term is obtained through QED**
- **If we want 0.1% accuracy and precision of P_z**
=> we need 0.1% precision and accuracy of P_c

Classical measurement - Slow

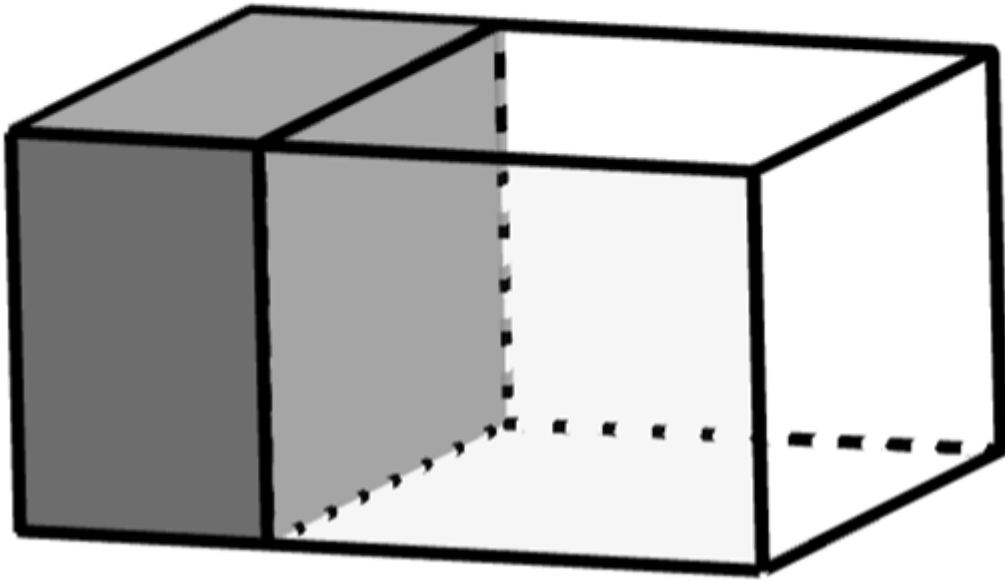
$$I(\theta, \phi) = \frac{1}{2} [S_0 + S_1 \cos 2\theta + S_2 \sin 2\theta \cos \phi + S_3 \sin 2\theta \sin \phi].$$



Dennis H Goldstein, Polarized Light - CRC Press Taylor & Francis Group

Photo-elastic modulator

Piezo-excited glass



$$n_x = n_0 \left(1 - \frac{n_0^2 q_{11}}{2} P \right) \quad n_y = n_0 \left(1 - \frac{n_0^2 q_{12}}{2} P \right)$$

$$P = P(t) = P_m \sin(2\pi f t)$$

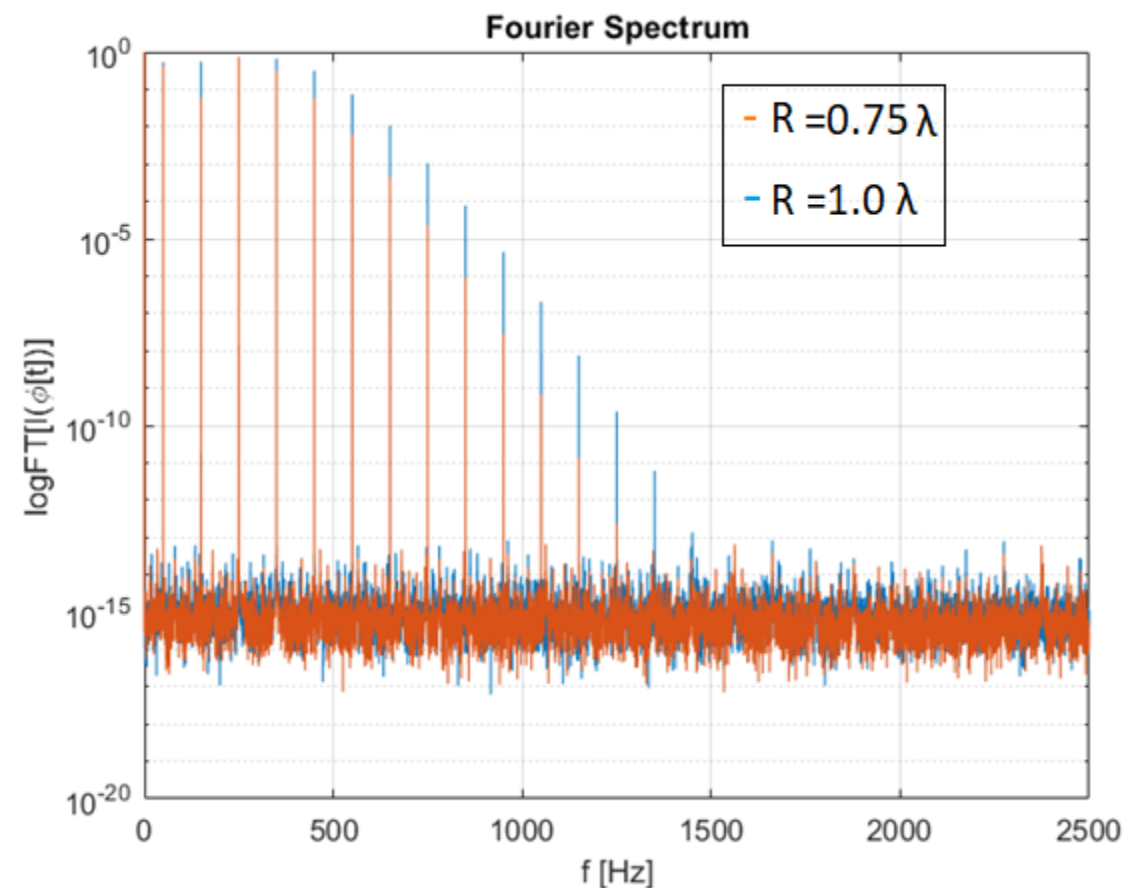
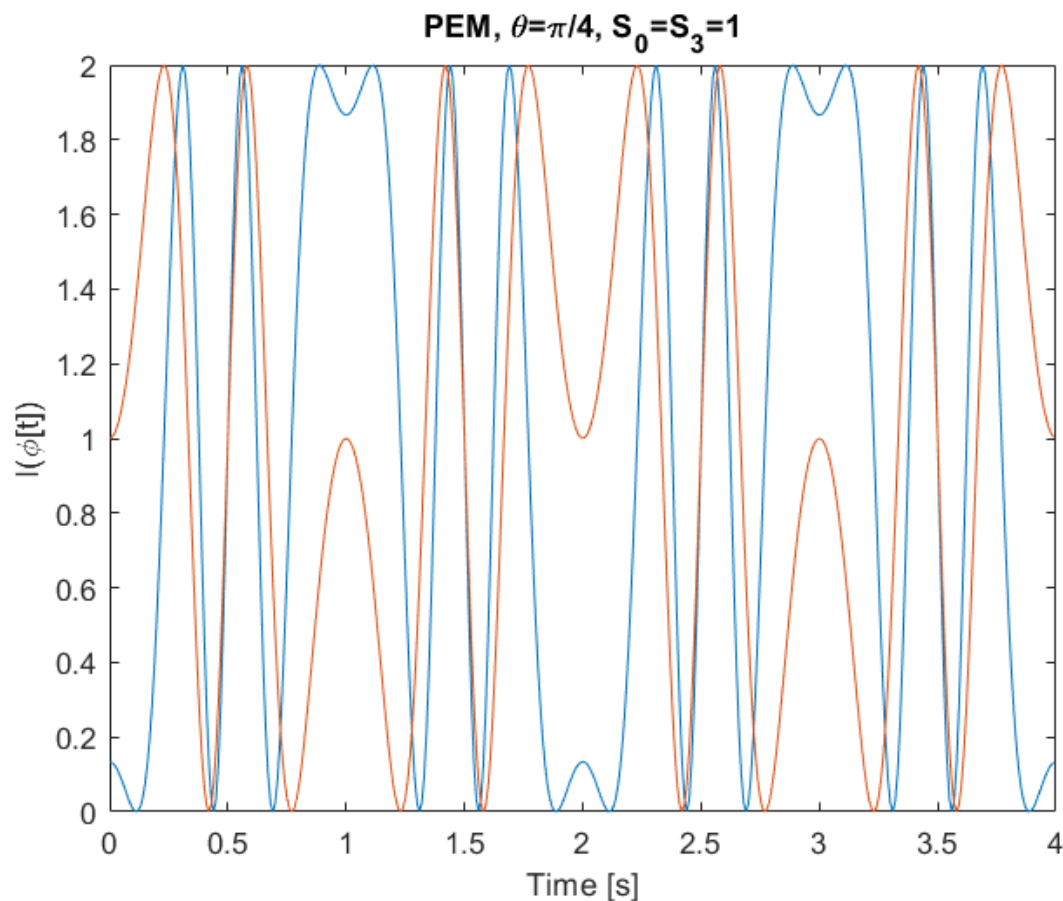
$$\phi(t) = \frac{2\pi e}{\lambda} (n_x - n_y) = \phi_0 \cos(2\pi f t)$$

Modulated difference of refraction indices

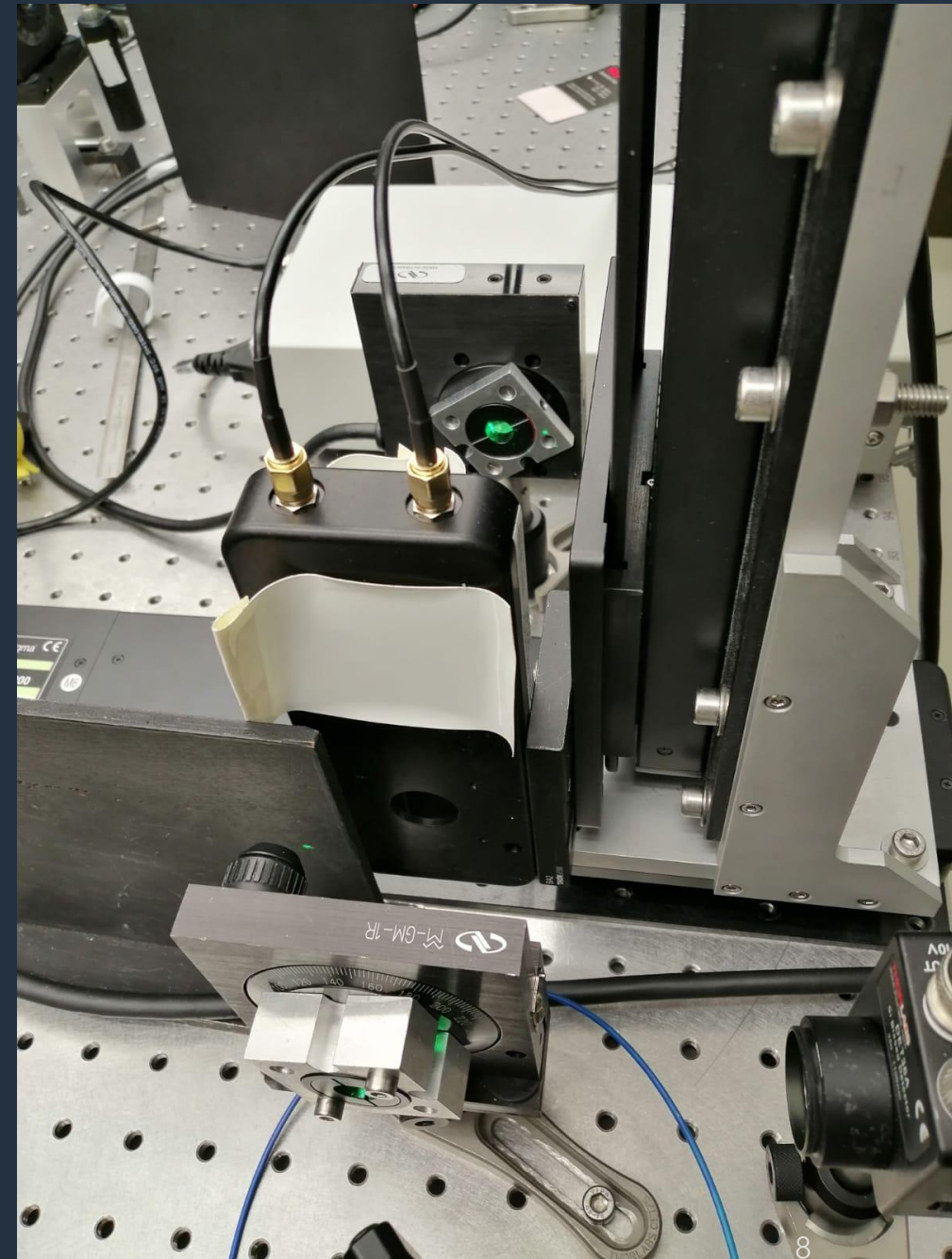
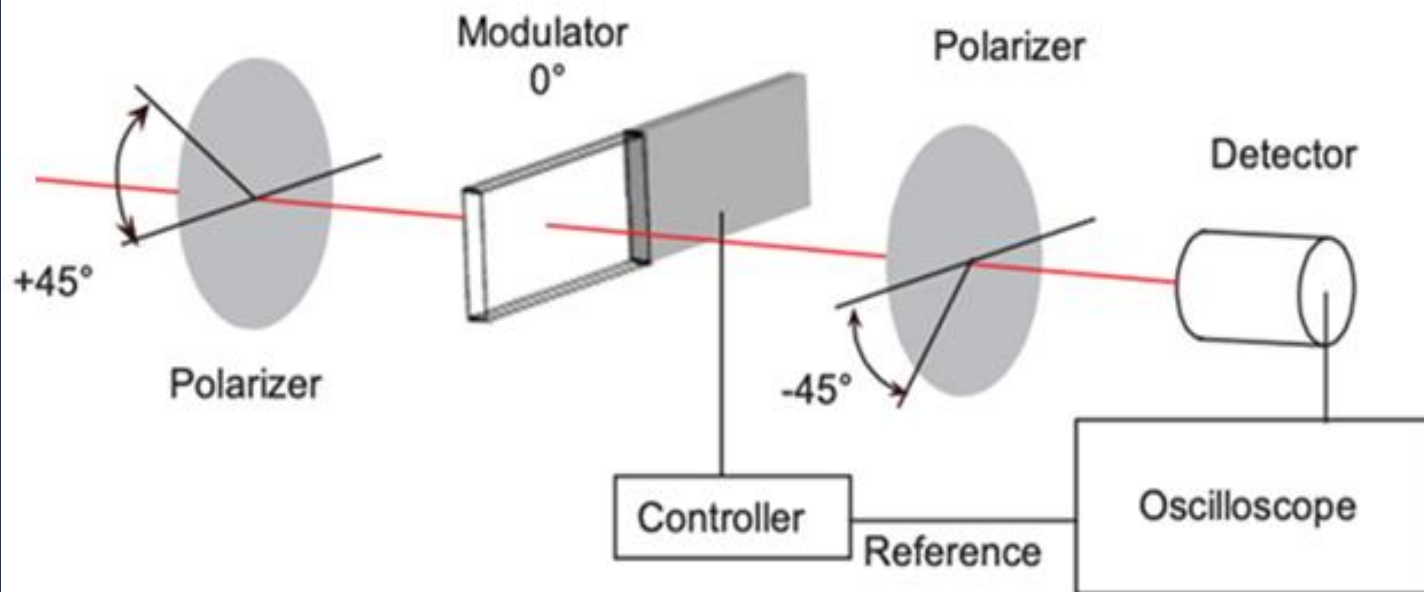
D. Yang, J. C. Canit, E. Gagnebet ; Photoelastic modulator - J. Optics (Paris)

Intensity measured with PEM

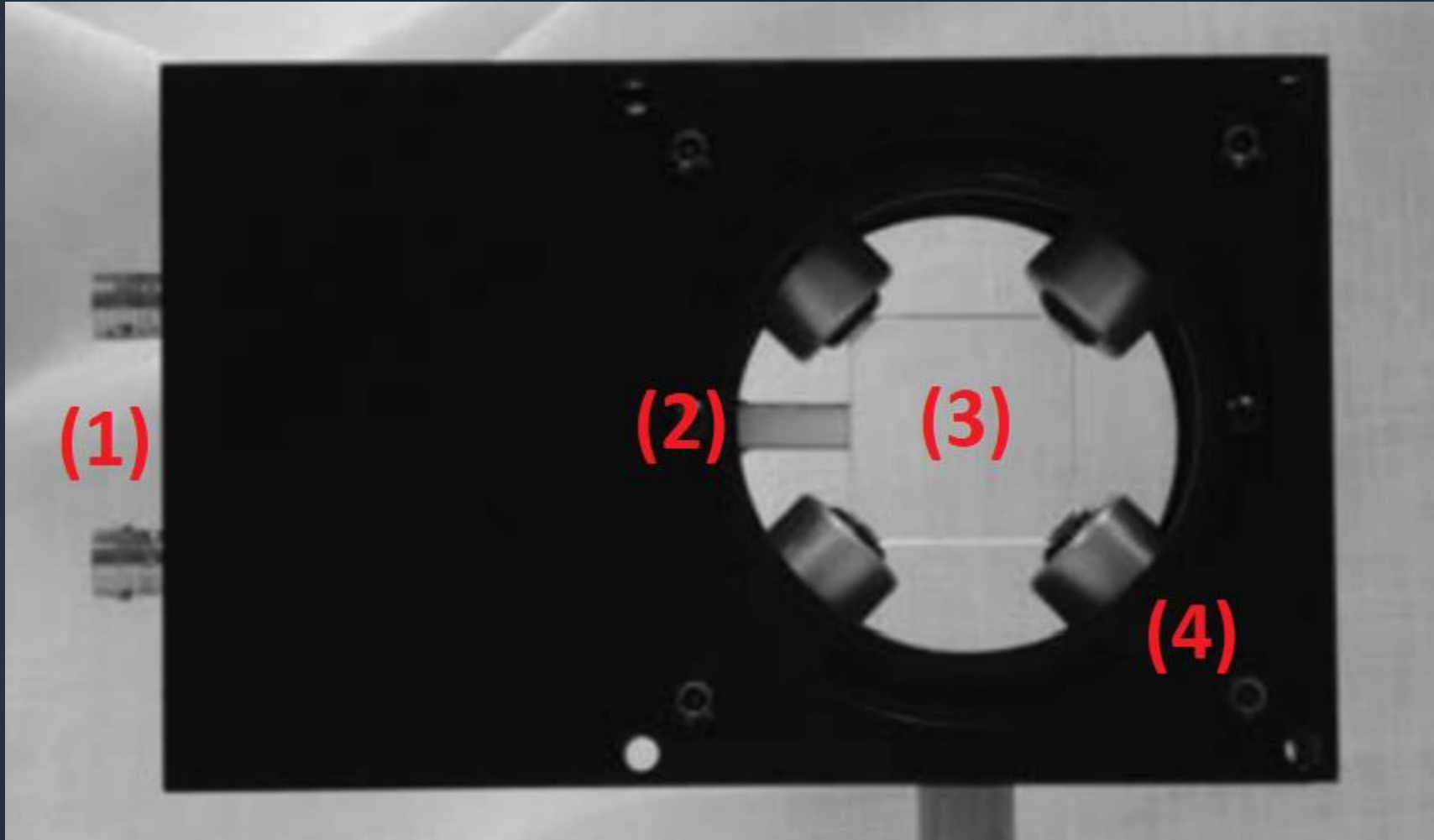
$$I(\theta, \phi[t]) = \frac{1}{2} [S_0 + S_1 \cos(2\theta) + \sin(2\theta)(S_2 \cos(\phi[t]) - S_3 \sin(\phi[t]))] \quad \left| \quad \phi(t) = \phi_0 \cos(2\pi f t) \quad \right| \quad \phi_0 = \frac{2\pi R}{\lambda}$$



PEM calibration setup



Example of a real PEM



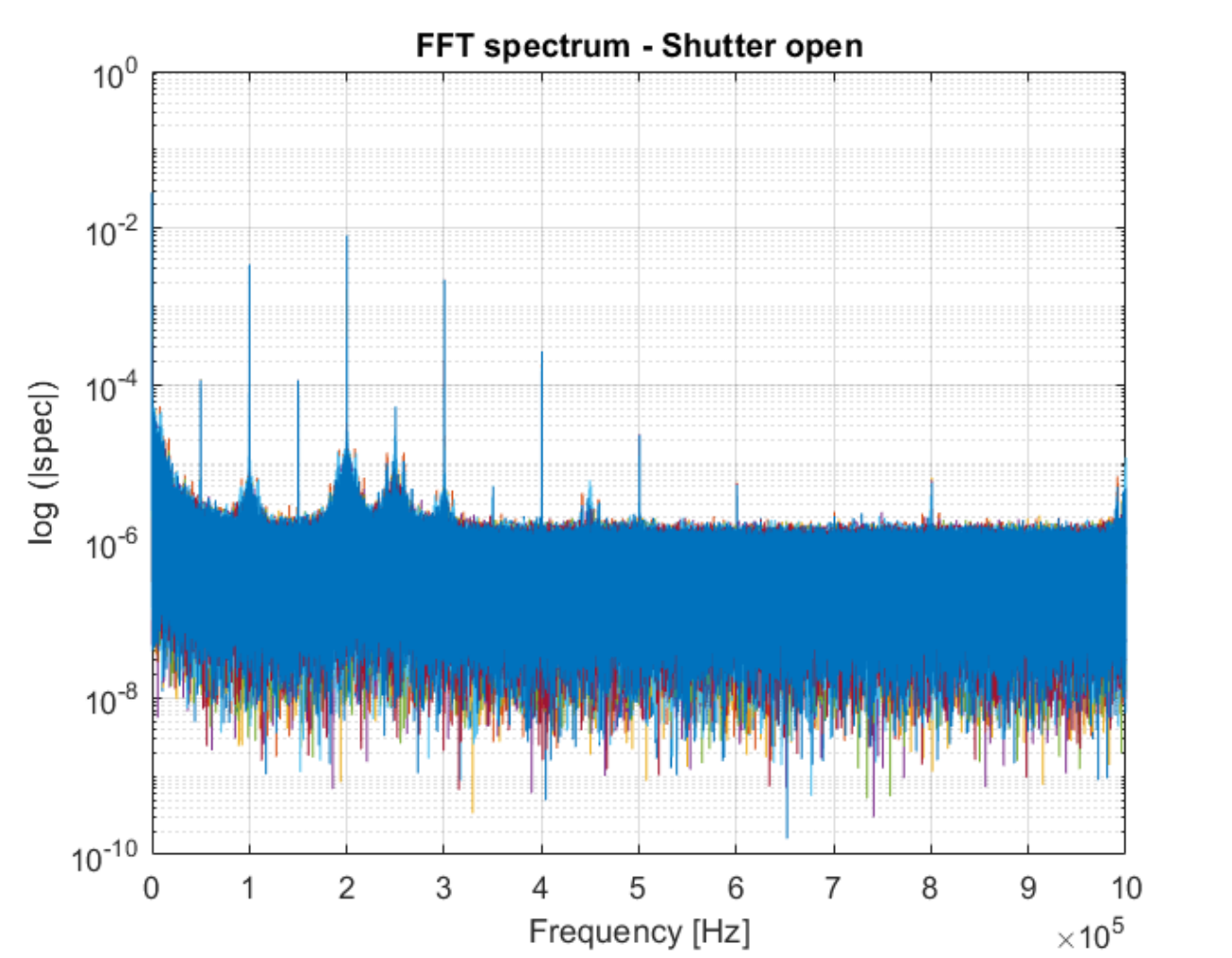
Dennis H Goldstein, Polarized Light - CRC Press Taylor & Francis Group

$$\Phi = \delta_0 + \delta_1 \cos(2\pi ft + \phi_1) + \delta_2 \cos(4\pi ft + \phi_2)$$

$$\begin{aligned} I(\delta_0, \delta_1, \delta_2, t) = I_o \Big\{ & 1 - J_0(\delta_2)J_0(\delta_1) + 2\delta_0 J_0(\delta_2)J_1(\delta_1)\cos(A) + 2\delta_0 J_1(\delta_2)J_0(\delta_1)\cos(B) \\ & + 2J_1(\delta_2)J_1(\delta_1)\left(\cos(A+B) + \cos(A-B)\right) + 2\sum_{n=1}^{\infty}(-1)^n \left[\delta_0 J_{2n+1}(\delta_2)J_0(\delta_1)\cos((2n+1)B) \right. \\ & + J_{2n+1}(\delta_2)J_1(\delta_1)\left(\cos(A+(2n+1)B) + \cos(A-(2n+1)B)\right) - J_{2n}(\delta_2)J_0(\delta_1)\cos(2nB) \\ & + \delta_0 J_{2n}(\delta_2)J_1(\delta_1)\left(\cos(A+2nB) + \cos(A-2nB)\right) - J_0(\delta_2)J_{2n}(\delta_1)\cos(2nA) \\ & + \delta_0 J_1(\delta_2)J_{2n}(\delta_1)\left(\cos(2nA+B) + \cos(2nA-B)\right) + \delta_0 J_0(\delta_2)J_{2n+1}(\delta_1)\cos((2n+1)A) \\ & + J_1(\delta_2)J_{2n+1}(\delta_1)\left(\cos((2n+1)A+B) + \cos((2n+1)A-B)\right) \Big] + 2\sum_{k=1}^{\infty}\sum_{n=1}^{\infty}(-1)^k(-1)^n \\ & \cdot \left[J_{2k+1}(\delta_2)J_{2n+1}(\delta_1)\left(\cos((2n+1)A+(2k+1)B) + \cos((2n+1)A-(2k+1)B)\right) \right. \\ & + \delta_0 J_{2k+1}(\delta_2)J_{2n}(\delta_1)\left(\cos(2nA+(2k+1)B) + \cos(2nA-(2k+1)B)\right) \\ & - J_{2k}(\delta_2)J_{2n}(\delta_1)\left(\cos(2nA+2kB) + \cos(2nA-2kB)\right) \\ & \left. + \delta_0 J_{2k}(\delta_2)J_{2n+1}(\delta_1)\left(\cos((2n+1)A+2kB) + \cos((2n+1)A-2kB)\right) \right] \Big\} \end{aligned}$$

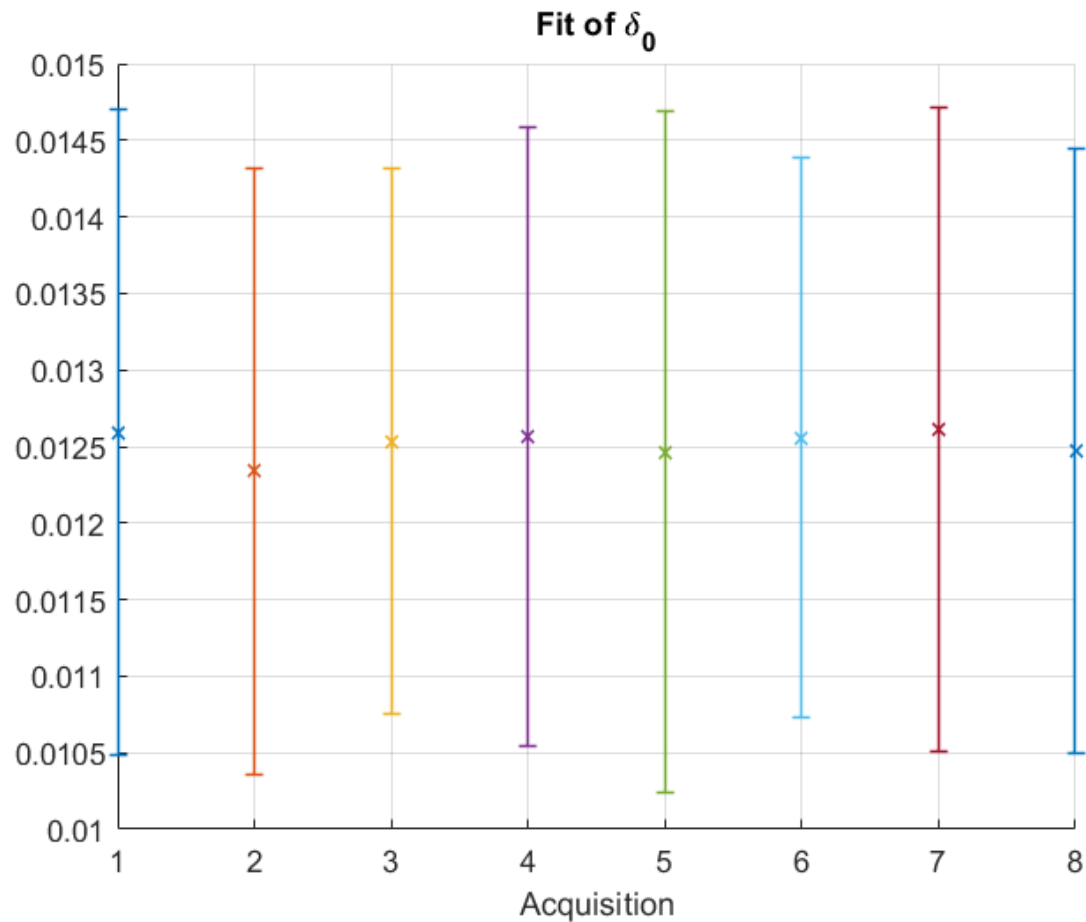
Frequency	Amplitude
n=0	$1 - J_0(\delta_2)J_0(\delta_1) + 2\left[J_2(\delta_2)J_4(\delta_1) - J_4(\delta_2)J_8(\delta_1) + J_6(\delta_2)J_{12}(\delta_1) - J_8(\delta_2)J_{16}(\delta_1) + \delta_0\left(J_3(\delta_2)J_6(\delta_1) - J_5(\delta_2)J_{10}(\delta_1) + J_7(\delta_2)J_{14}(\delta_1) - J_1(\delta_2)J_2(\delta_1)\right)\right]$
n=1	$2\left[J_1(\delta_2)J_1(\delta_1) - J_1(\delta_2)J_3(\delta_1) - J_3(\delta_2)J_5(\delta_1) + J_3(\delta_2)J_7(\delta_1) + J_5(\delta_2)J_9(\delta_1) - J_5(\delta_2)J_{11}(\delta_1) - J_7(\delta_2)J_{13}(\delta_1) + J_7(\delta_2)J_{15}(\delta_1) + J_9(\delta_2)J_{17}(\delta_1) + \delta_0\left(J_0(\delta_2)J_1(\delta_1) + J_2(\delta_2)J_3(\delta_1) - J_2(\delta_2)J_5(\delta_1) - J_4(\delta_2)J_7(\delta_1) + J_4(\delta_2)J_9(\delta_1) + J_6(\delta_2)J_{11}(\delta_1) - J_6(\delta_2)J_{13}(\delta_1) - J_8(\delta_2)J_{15}(\delta_1) + J_8(\delta_2)J_{17}(\delta_1)\right)\right]$
n=2	$2\left[J_0(\delta_2)J_2(\delta_1) - J_2(\delta_2)J_2(\delta_1) - J_2(\delta_2)J_6(\delta_1) + J_4(\delta_2)J_6(\delta_1) + J_4(\delta_2)J_{10}(\delta_1) - J_6(\delta_2)J_{10}(\delta_1) - J_6(\delta_2)J_{14}(\delta_1) + J_8(\delta_2)J_{14}(\delta_1) + \delta_0\left(J_1(\delta_2)J_0(\delta_1) + J_1(\delta_2)J_4(\delta_1) - J_3(\delta_2)J_4(\delta_1) - J_3(\delta_2)J_8(\delta_1) + J_5(\delta_2)J_8(\delta_1) + J_5(\delta_2)J_{12}(\delta_1) - J_7(\delta_2)J_{12}(\delta_1) - J_7(\delta_2)J_{16}(\delta_1) + J_9(\delta_2)J_{16}(\delta_1)\right)\right]$

Non-linear approach

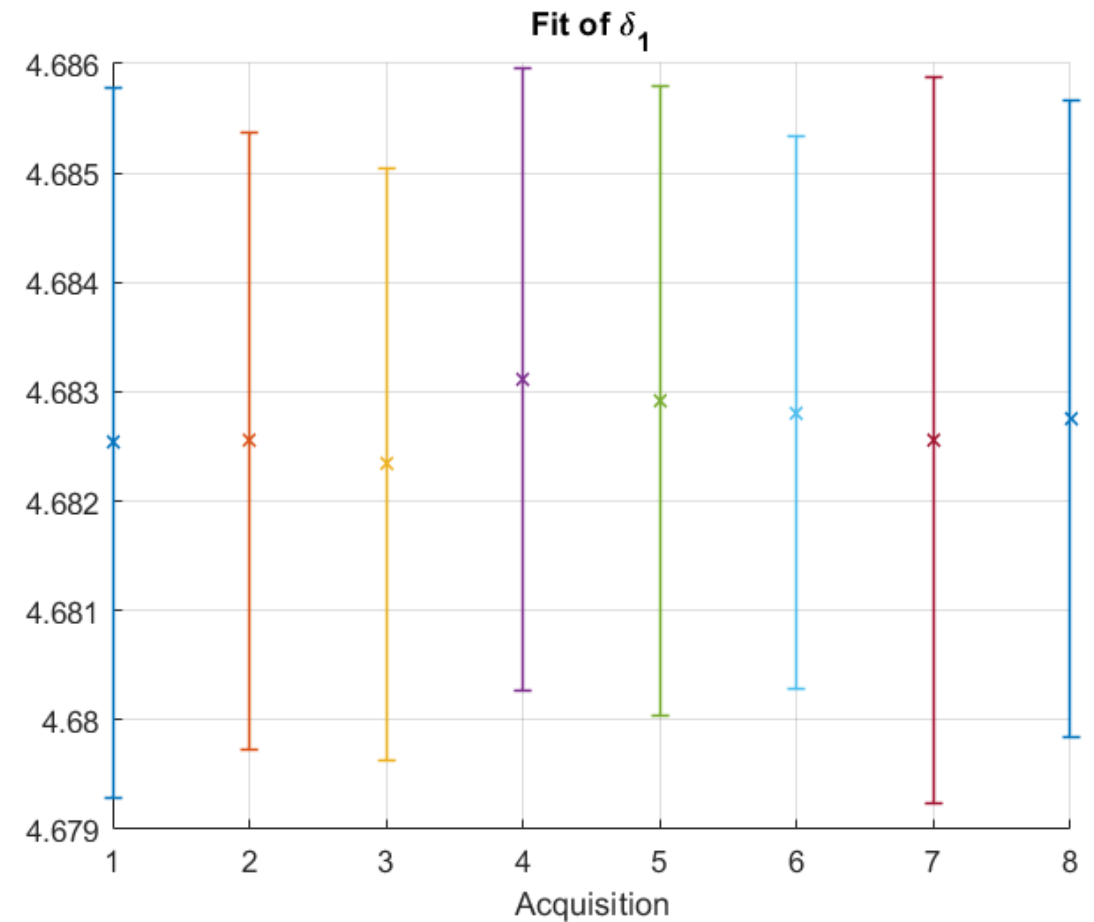


Experimental results

$$\delta_0 = \frac{2\pi}{\lambda} \Delta n \cdot e ; e(\text{width}) \rightarrow 10\text{mm}$$



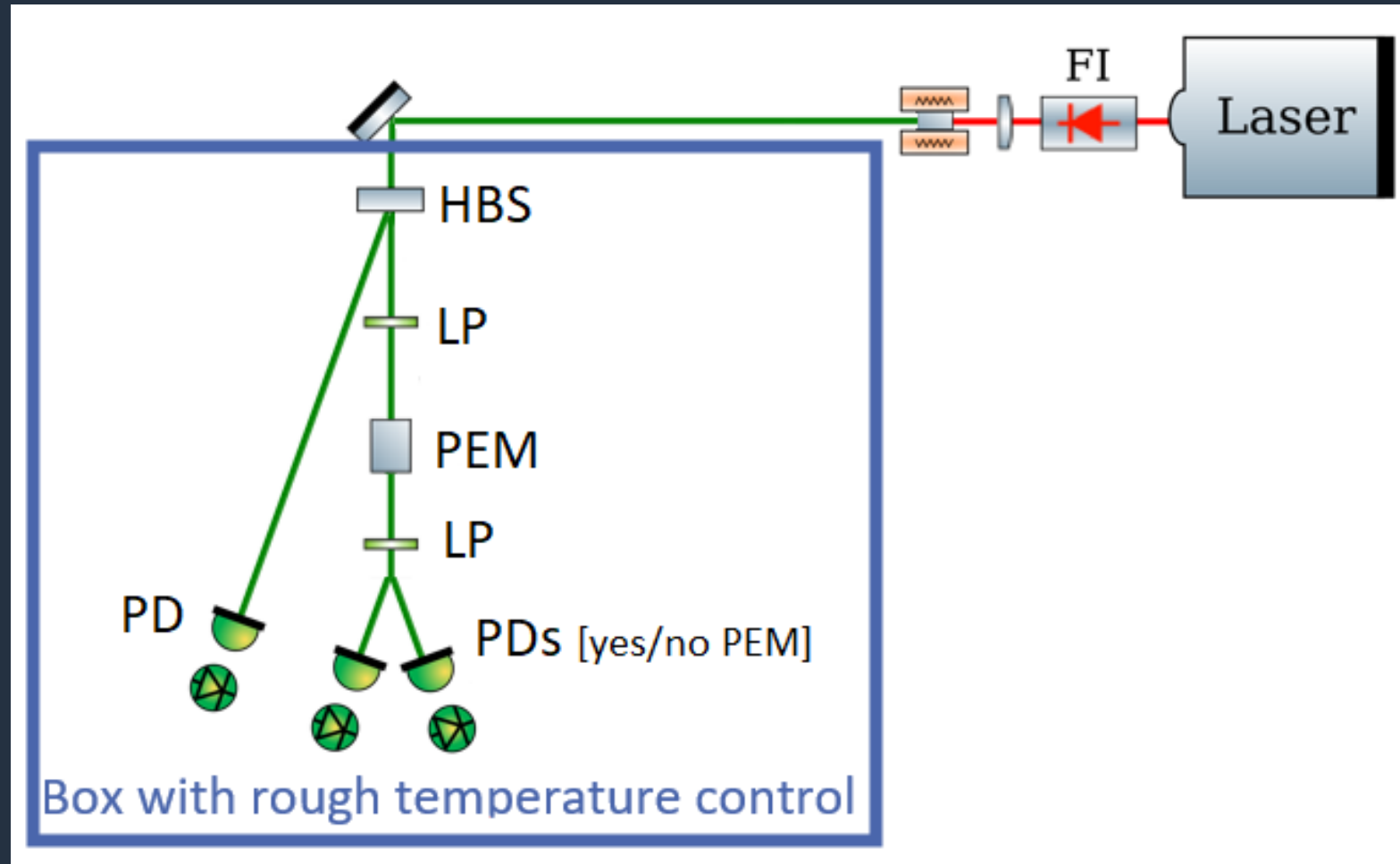
$$\delta_1 = \frac{2\pi R}{\lambda} \approx 4.712$$

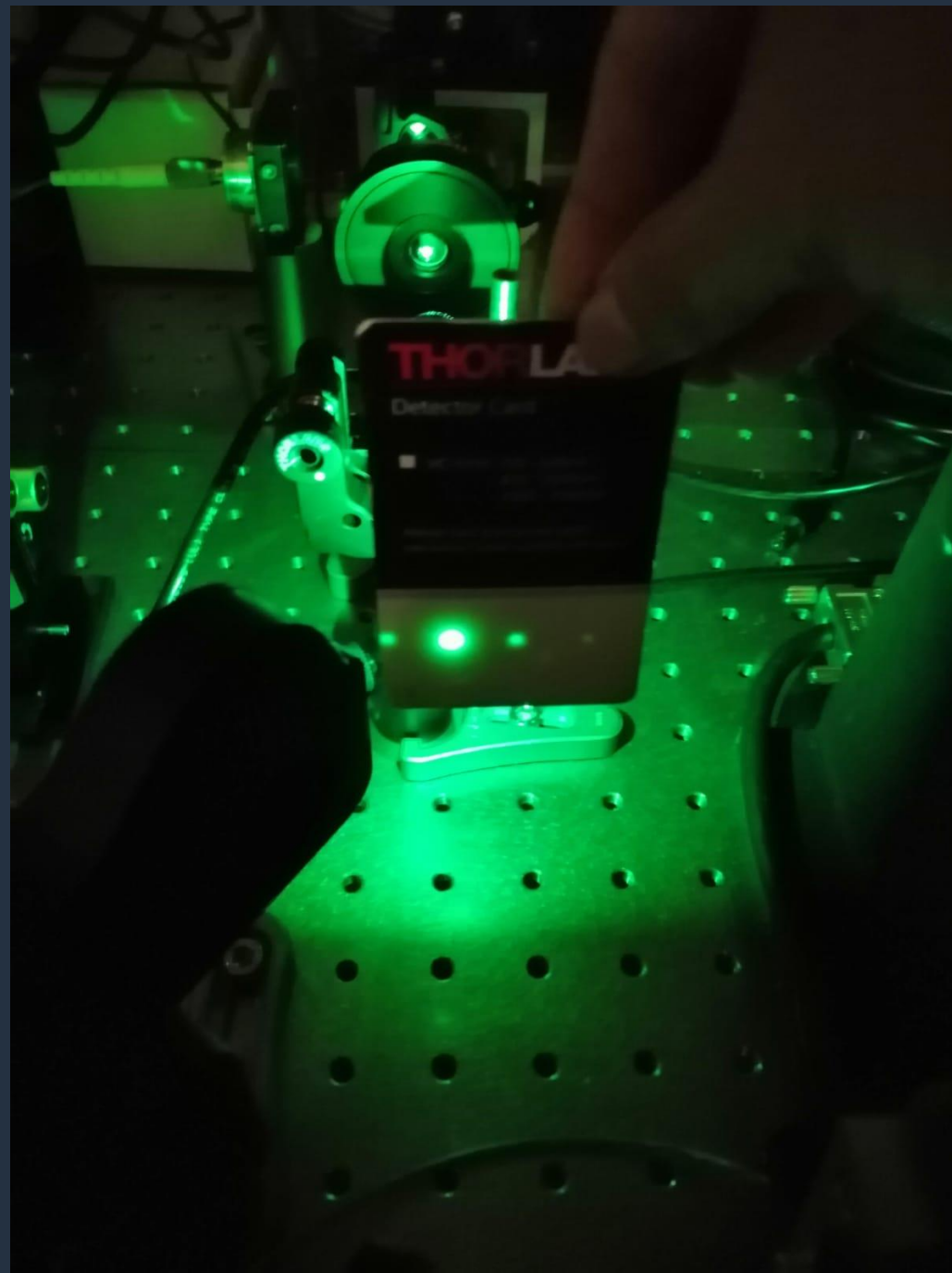
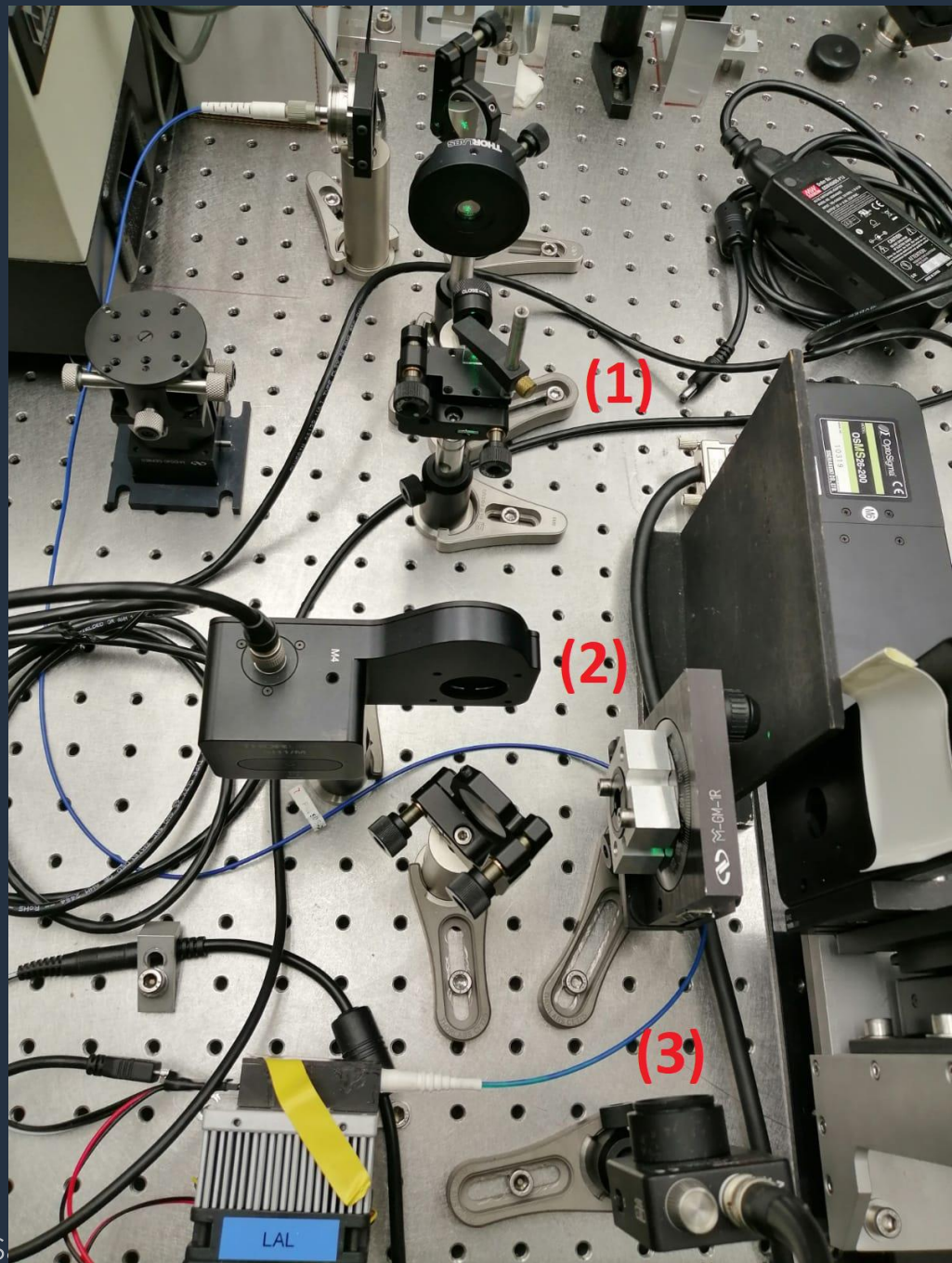


Conclusions and prospects

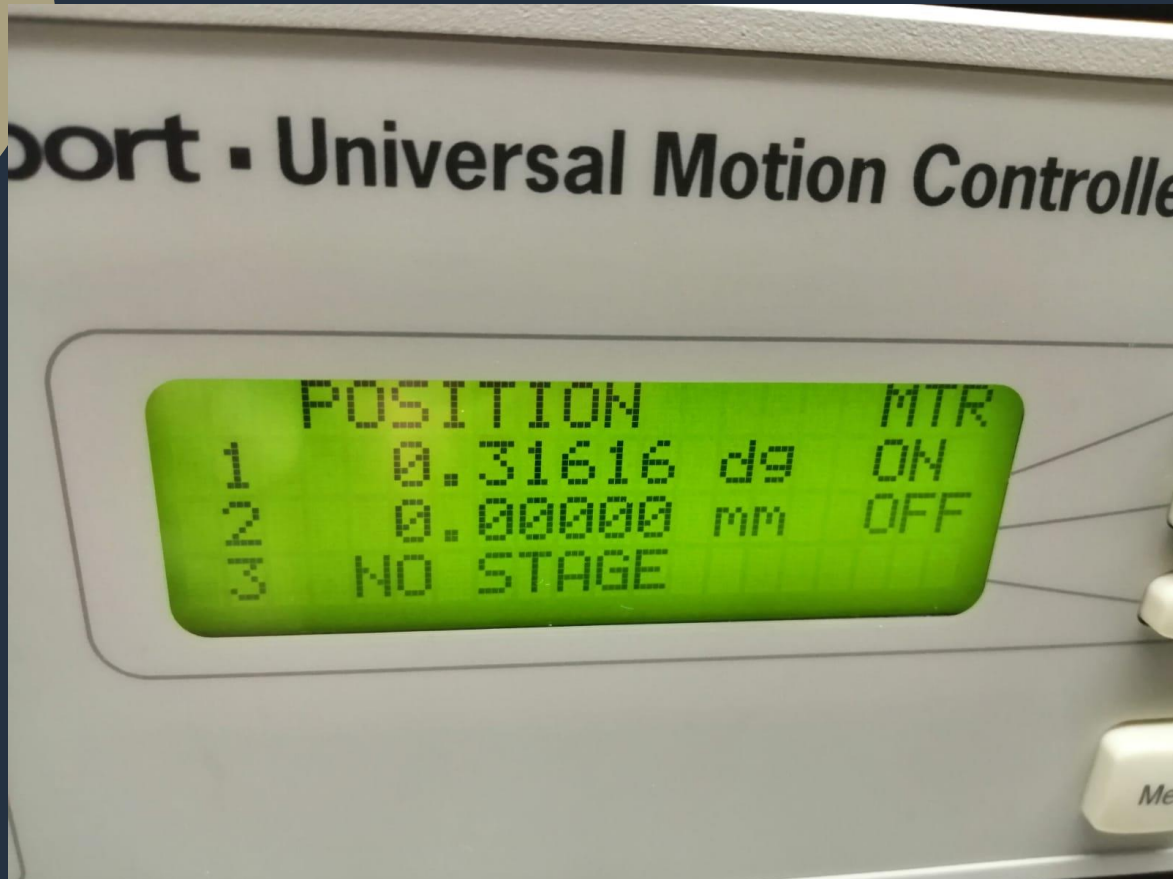
- The first order treatment gives stable and consistent results, especially for even harmonics.
- The use of the Fourier transform to analyze the signals represents a great advantage.
- For future work, higher order terms should be considered, keeping in mind the relative phase that is generated.

BACKUP – Optical setup

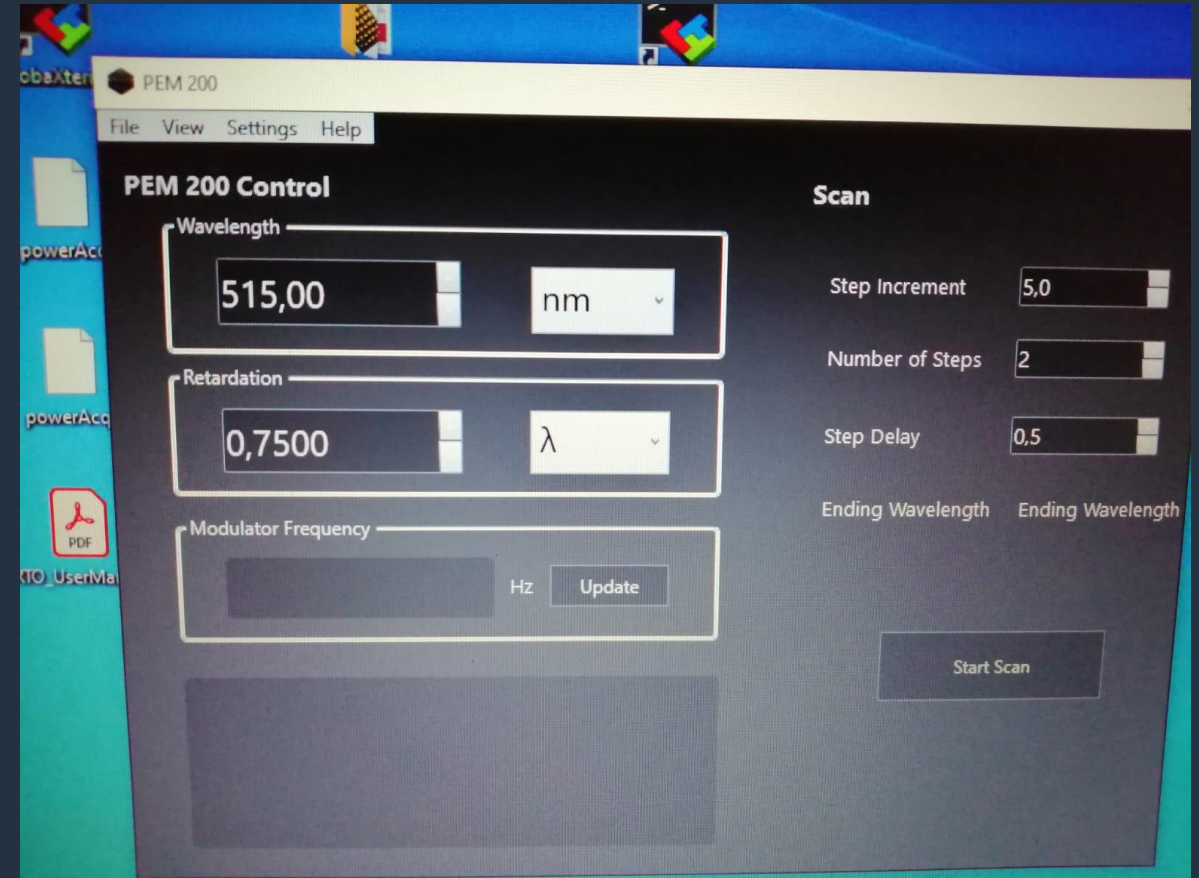




Drivers used in the optical room



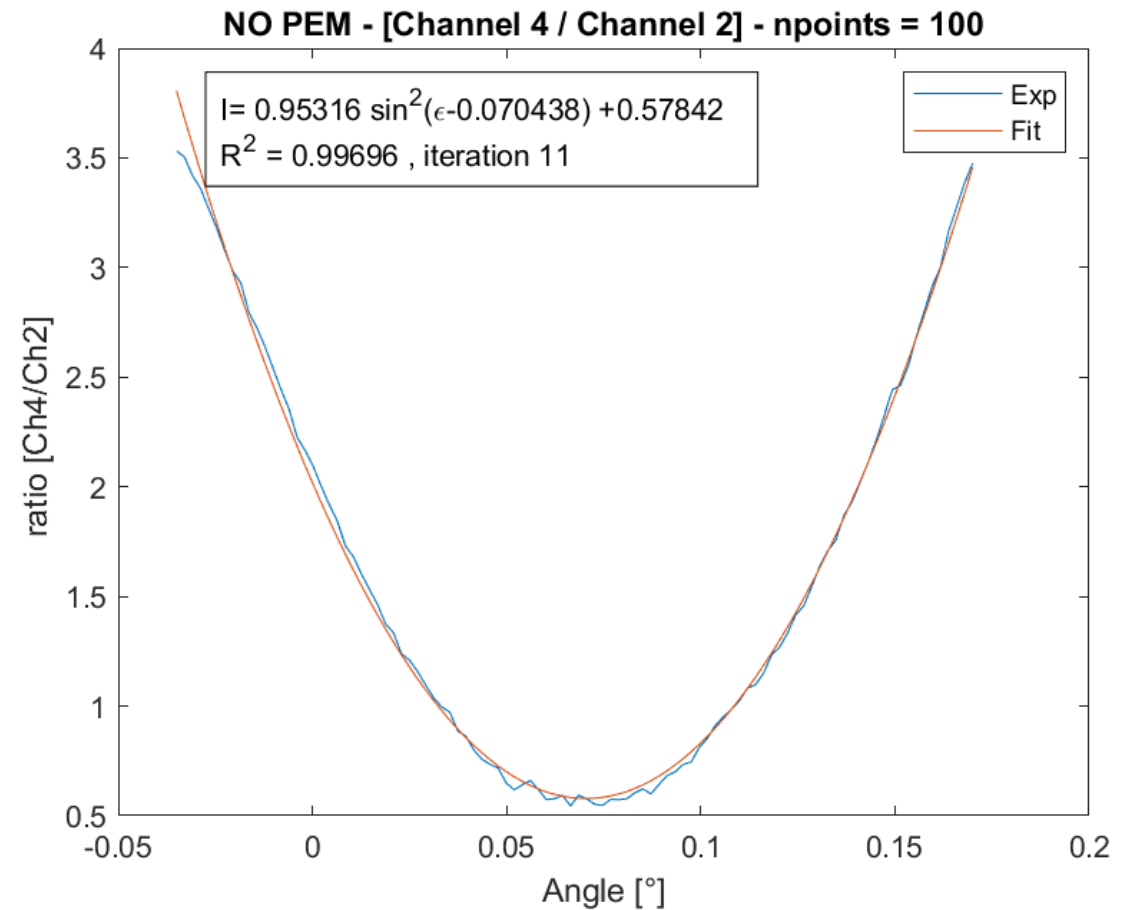
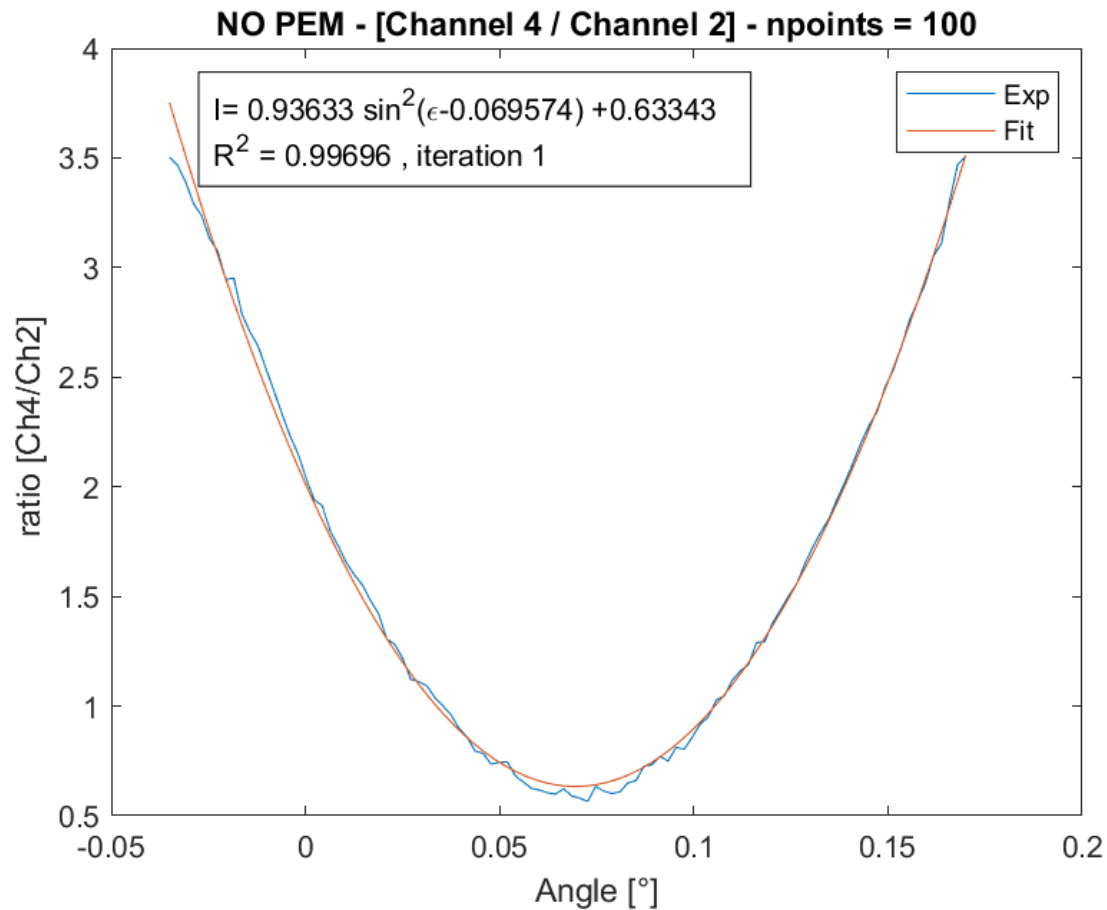
Second polarizer rotator



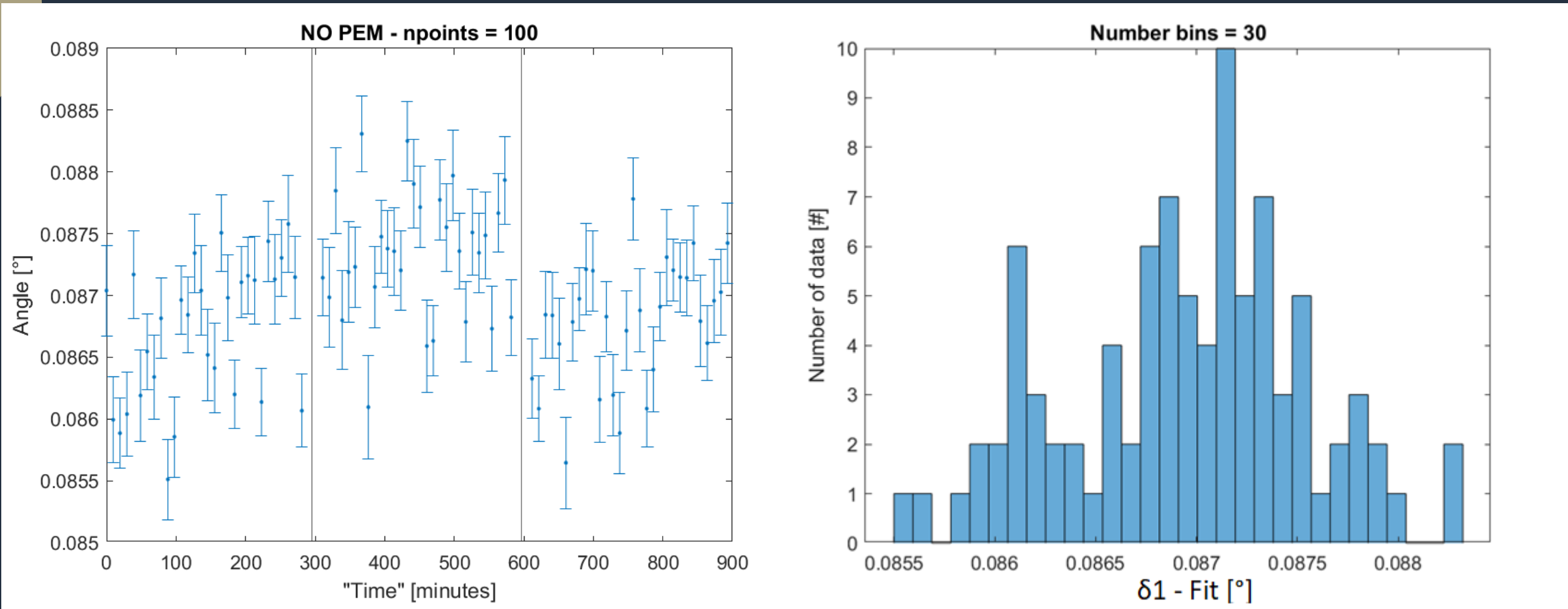
Photoelastic modulator controller



Looking for the cross angle

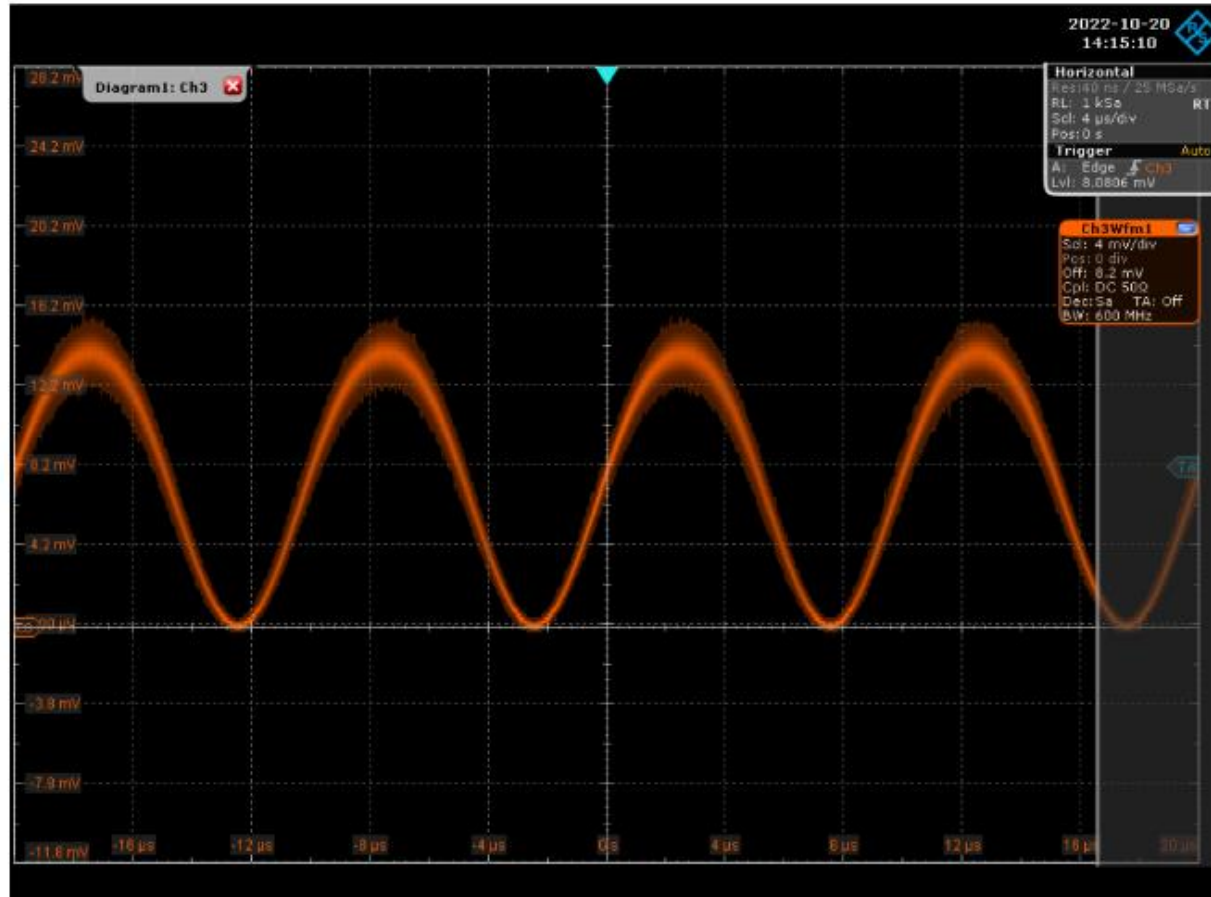


Collecting the data of three days

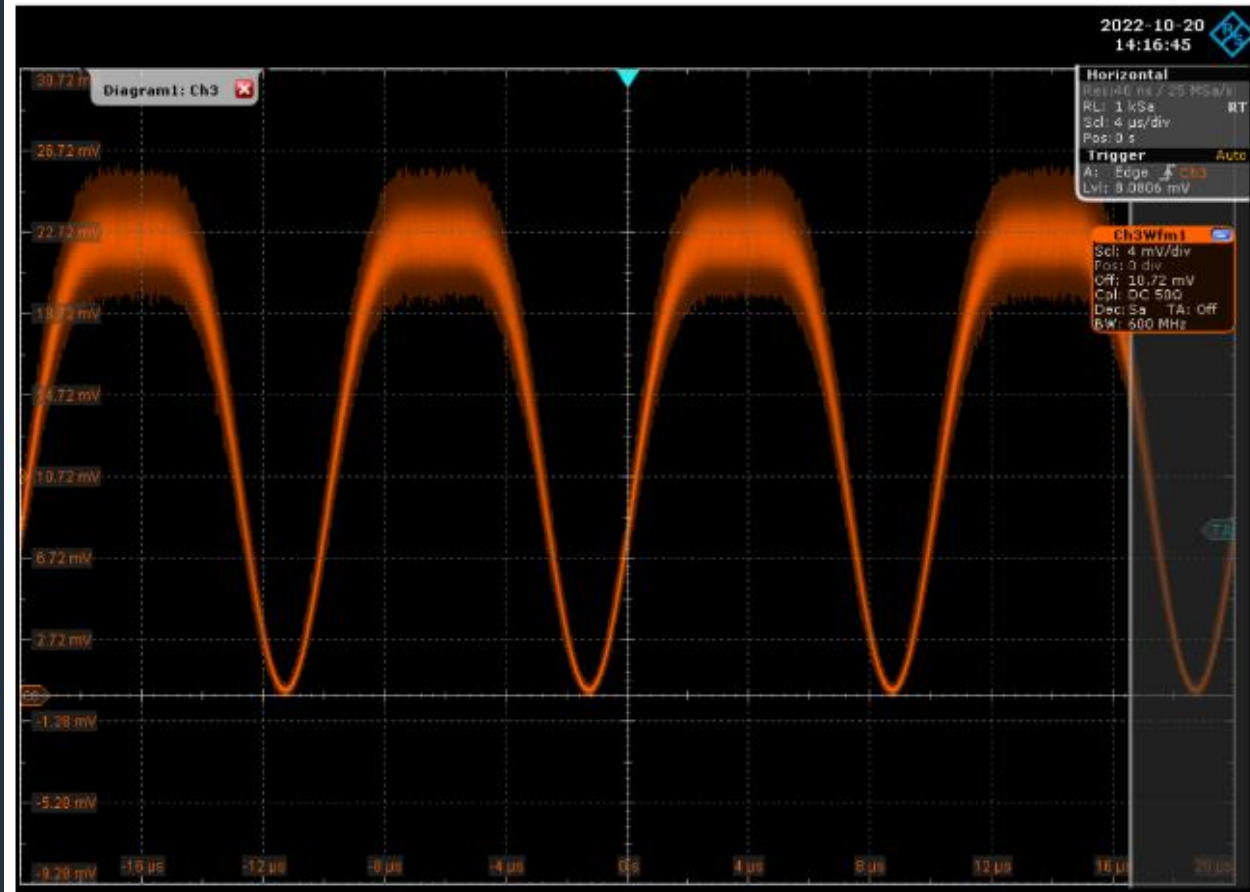


Using the PEM driver

$\lambda = 0.25$:

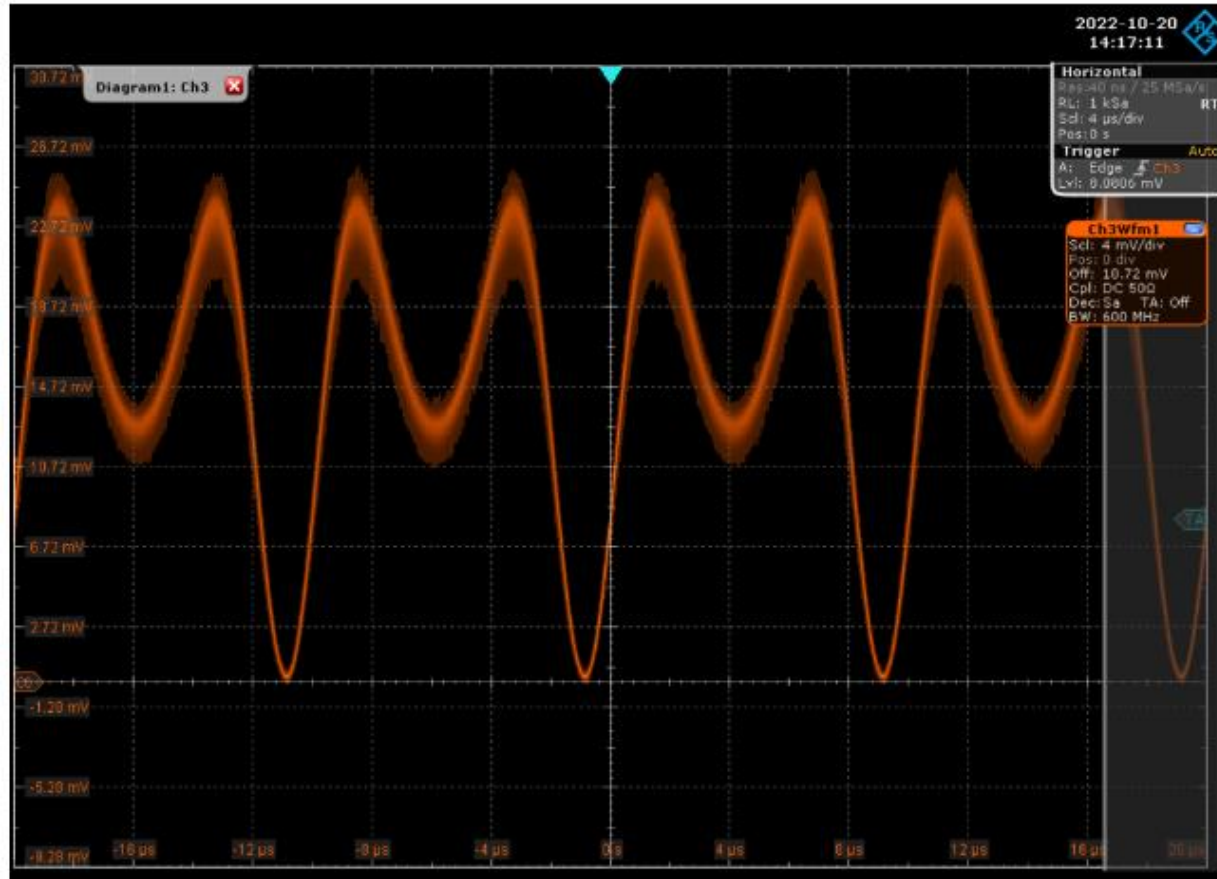


$\lambda = 0.5$:



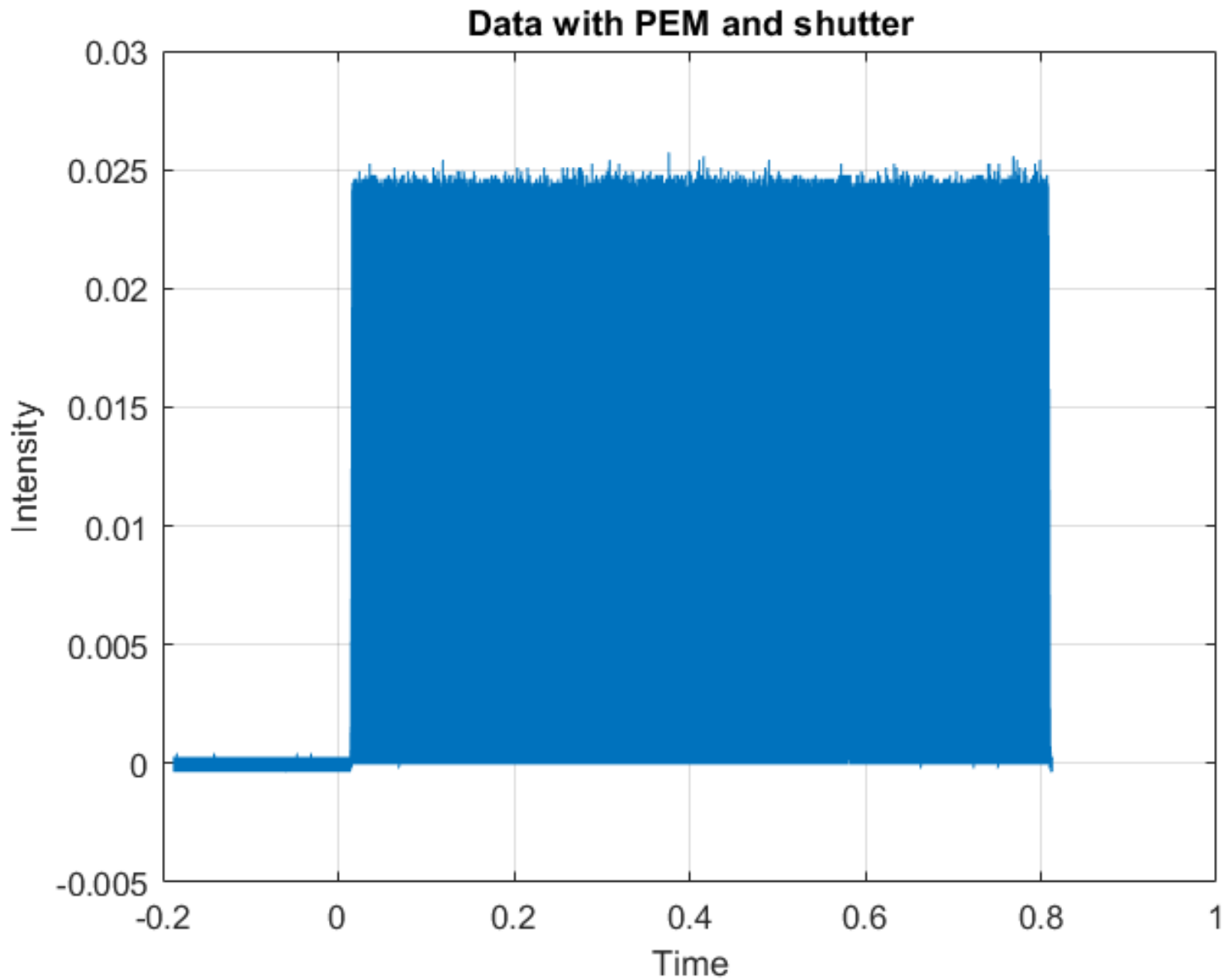
Relevance of trigger position !!!

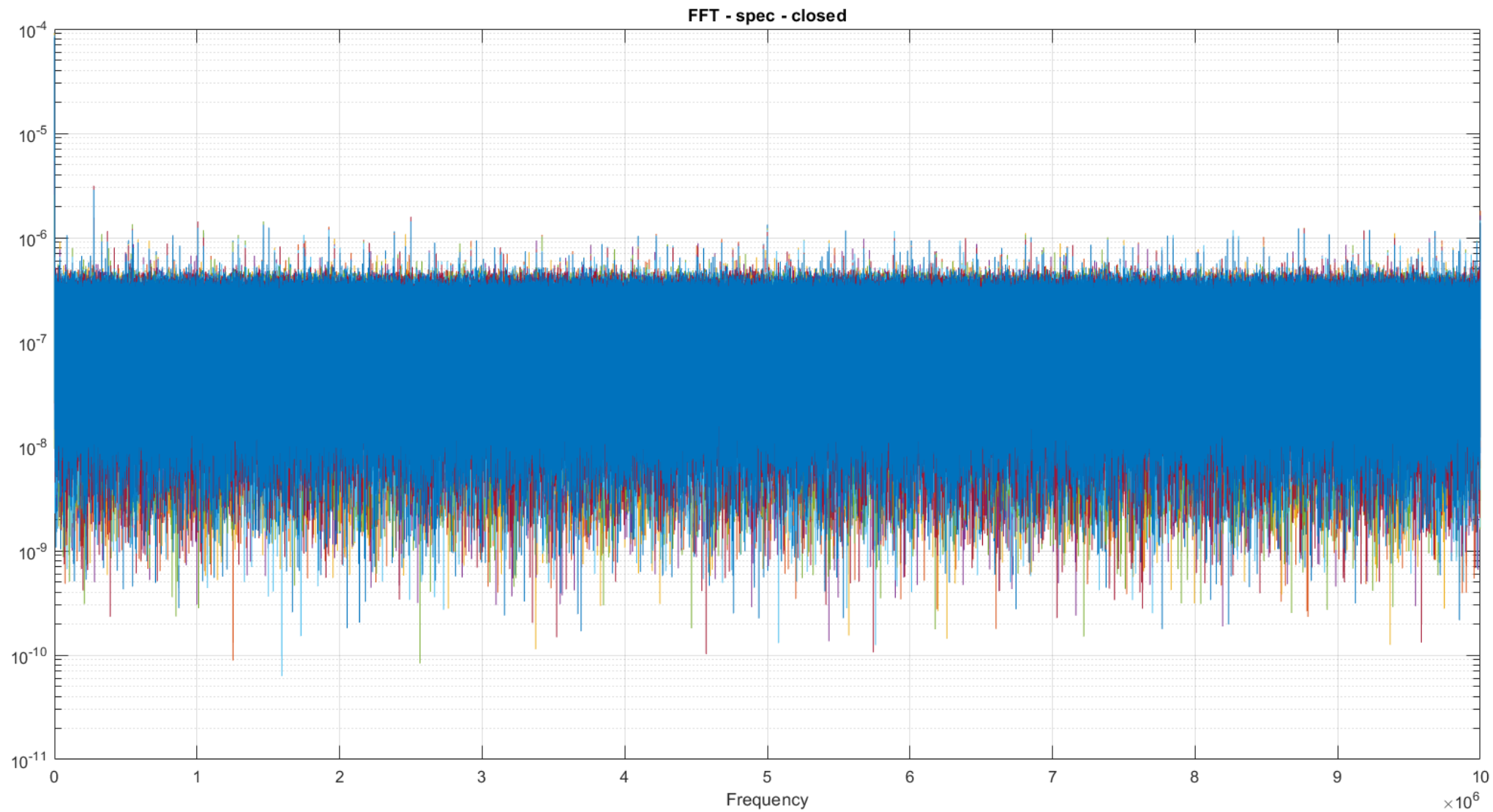
$\lambda = 0.75$:

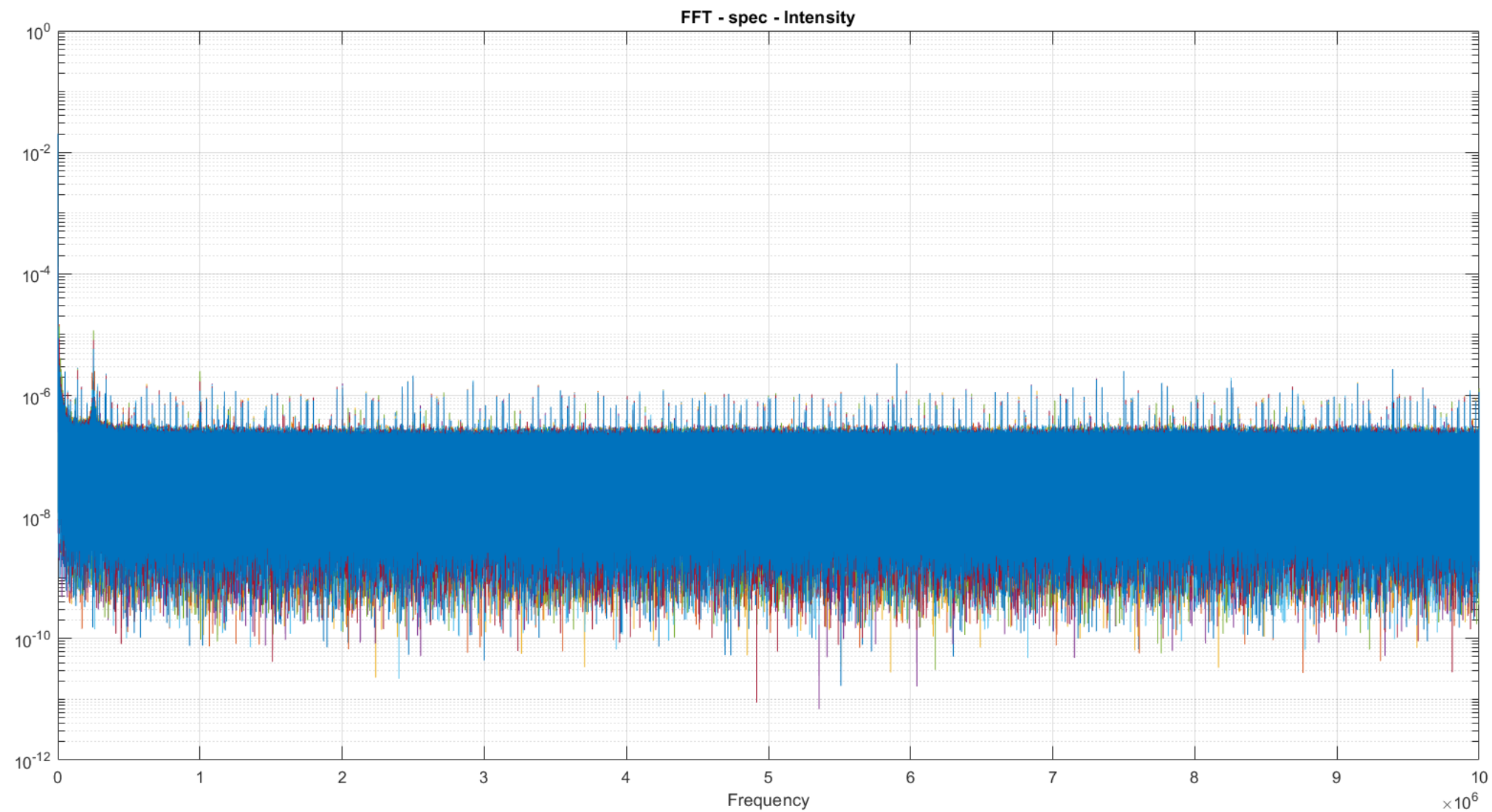


$\lambda = 1$:









Harmonic peaks and background zones

