

The Higgs Boson

Daniel de Florian

7th ComHEP



Escuela de
Ciencia y Tecnología
ECyT_UNSAM

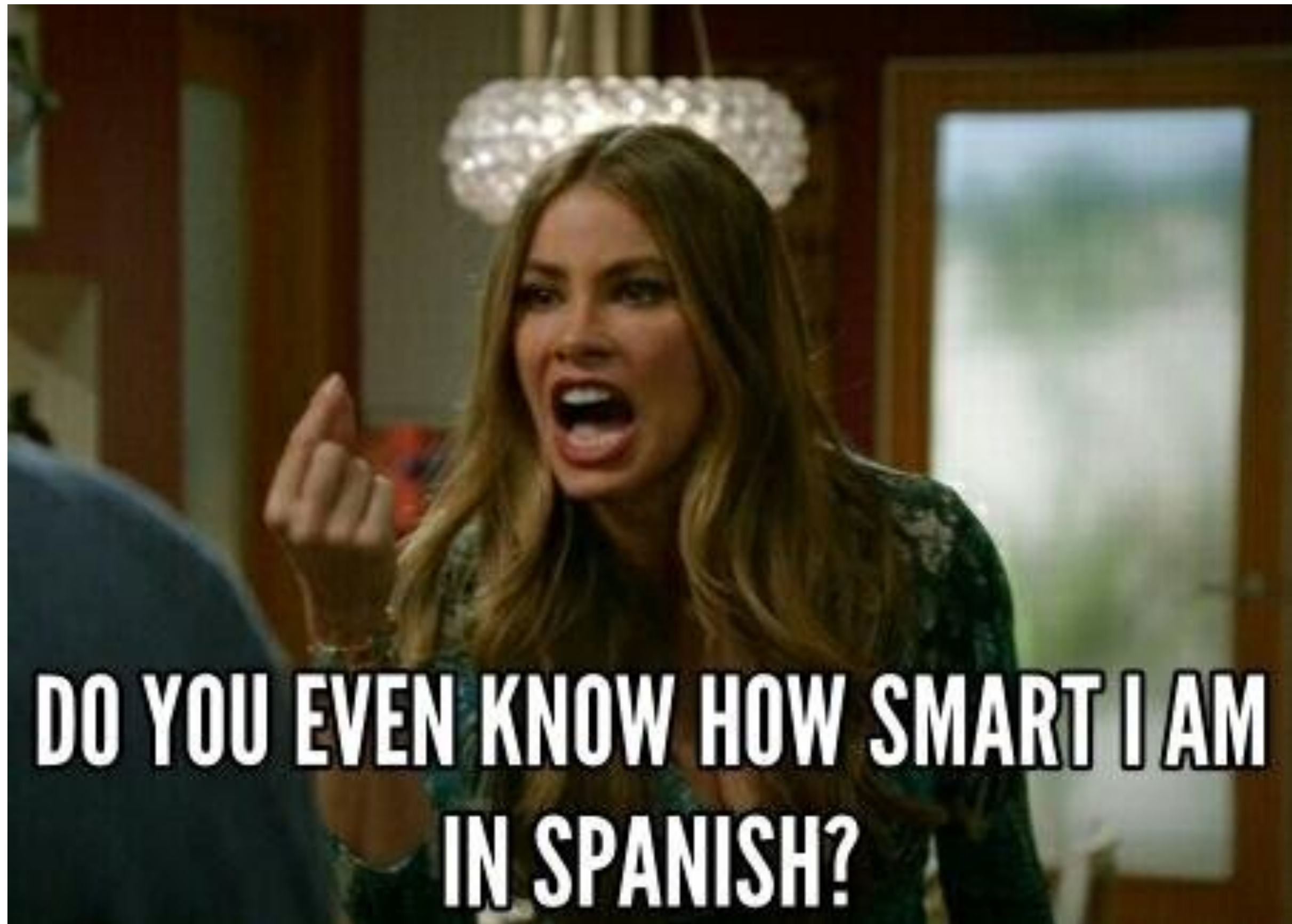
The Higgs Boson

Theoretical Introduction

Search and Discovery

Properties

Future(s)

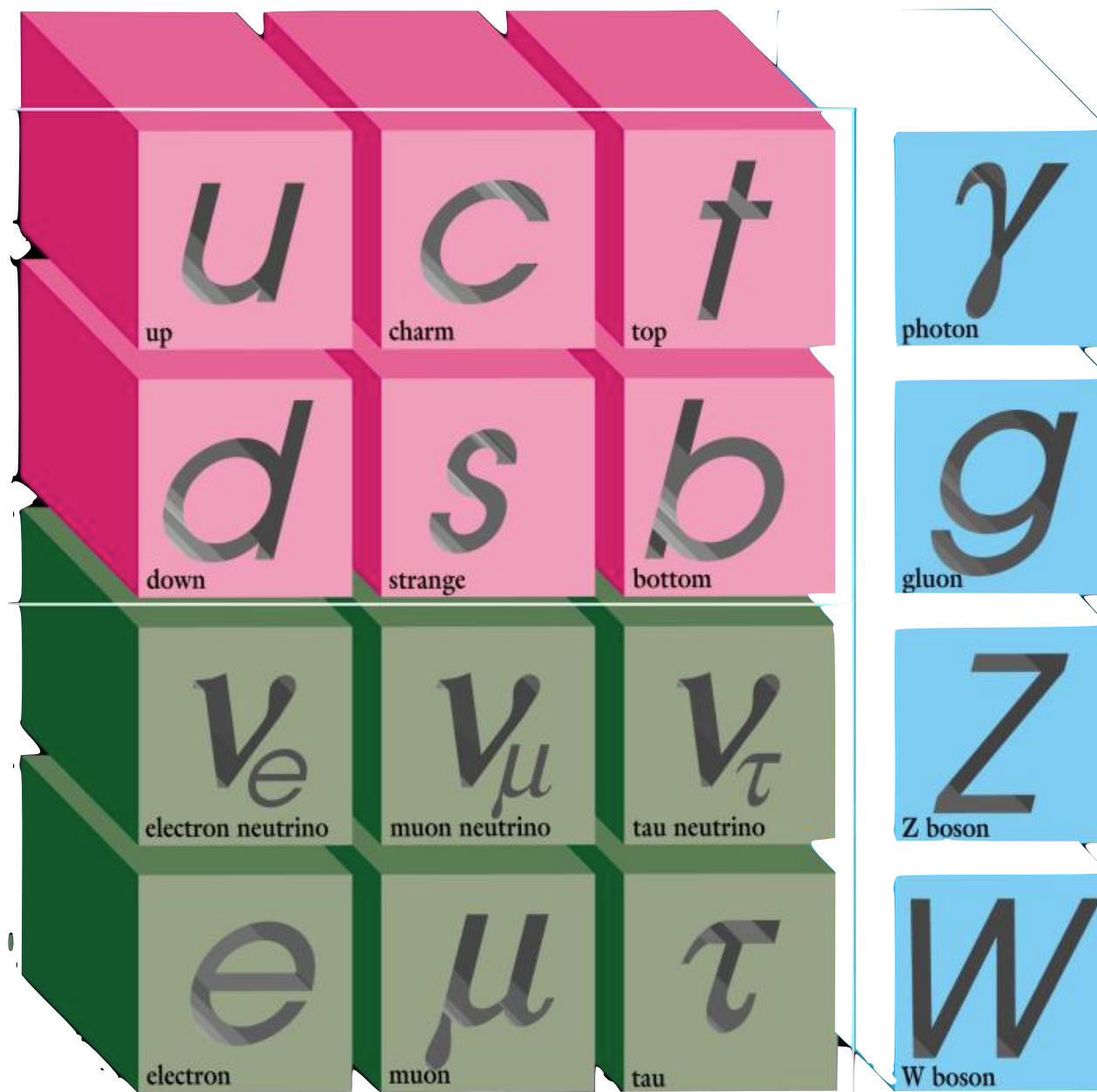


**DO YOU EVEN KNOW HOW SMART I AM
IN SPANISH?**

Theoretical Introduction

Standard Model

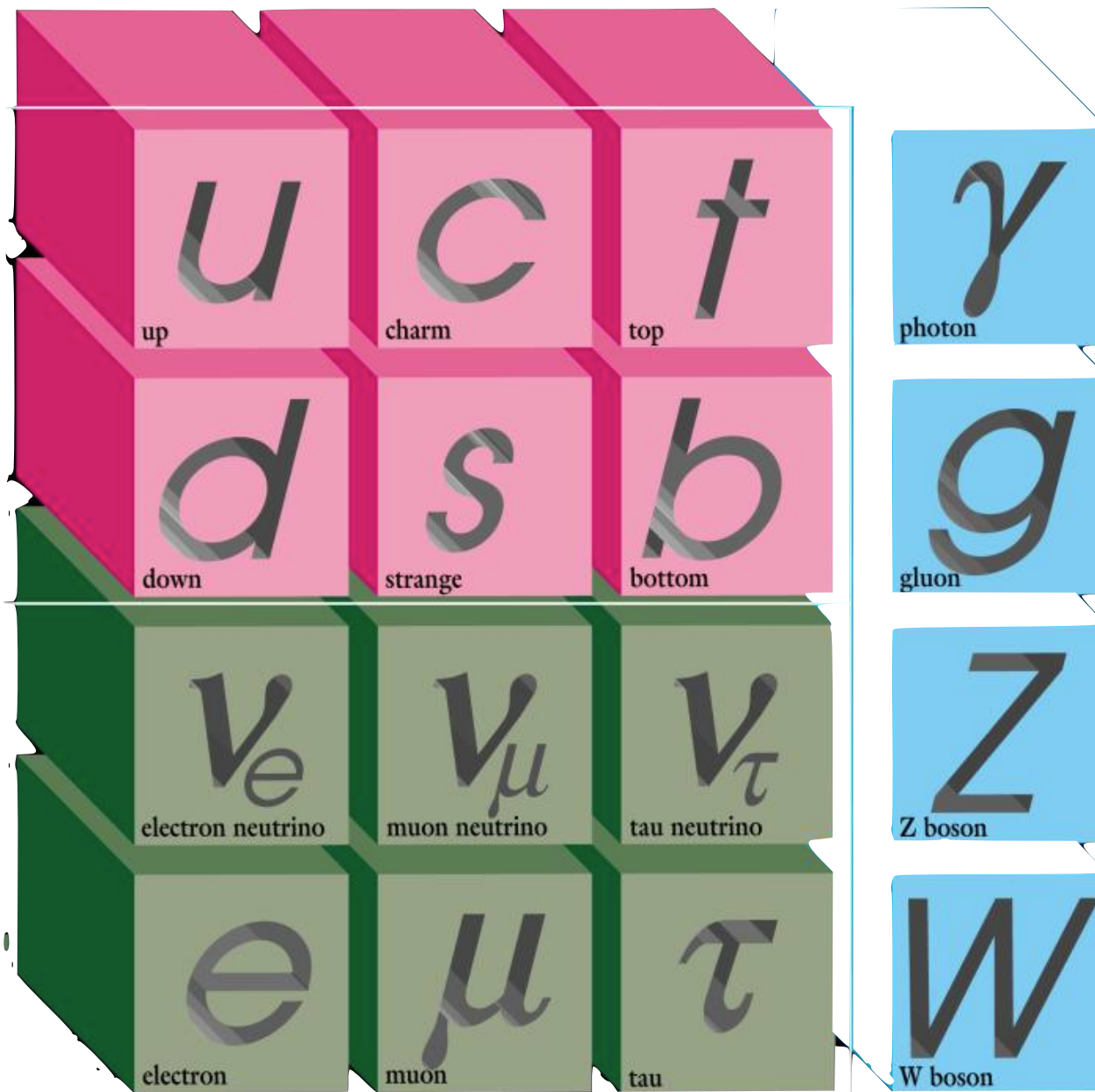
Strong + Weak + EM interactions



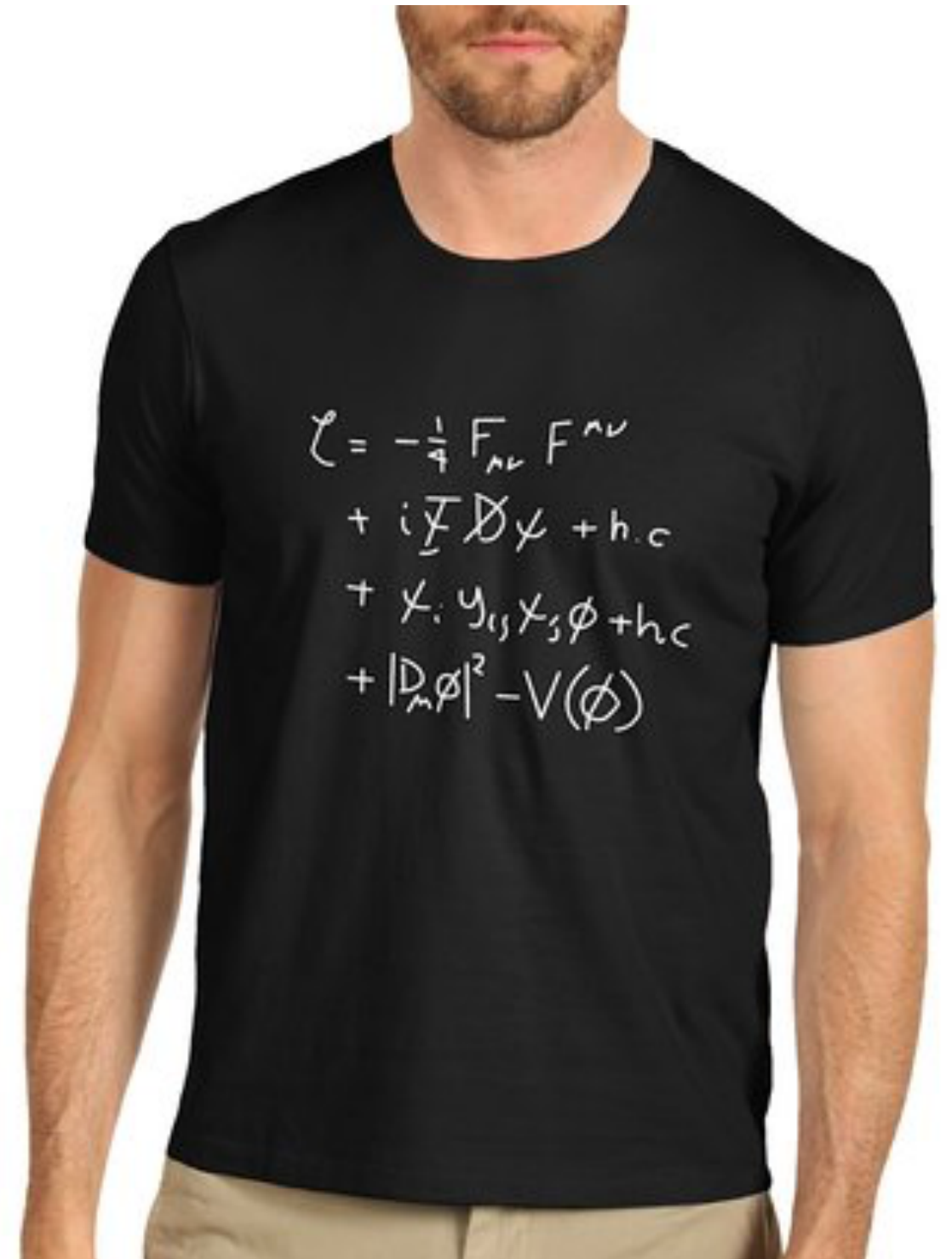
‘matter’ + gauge bosons

Standard Model

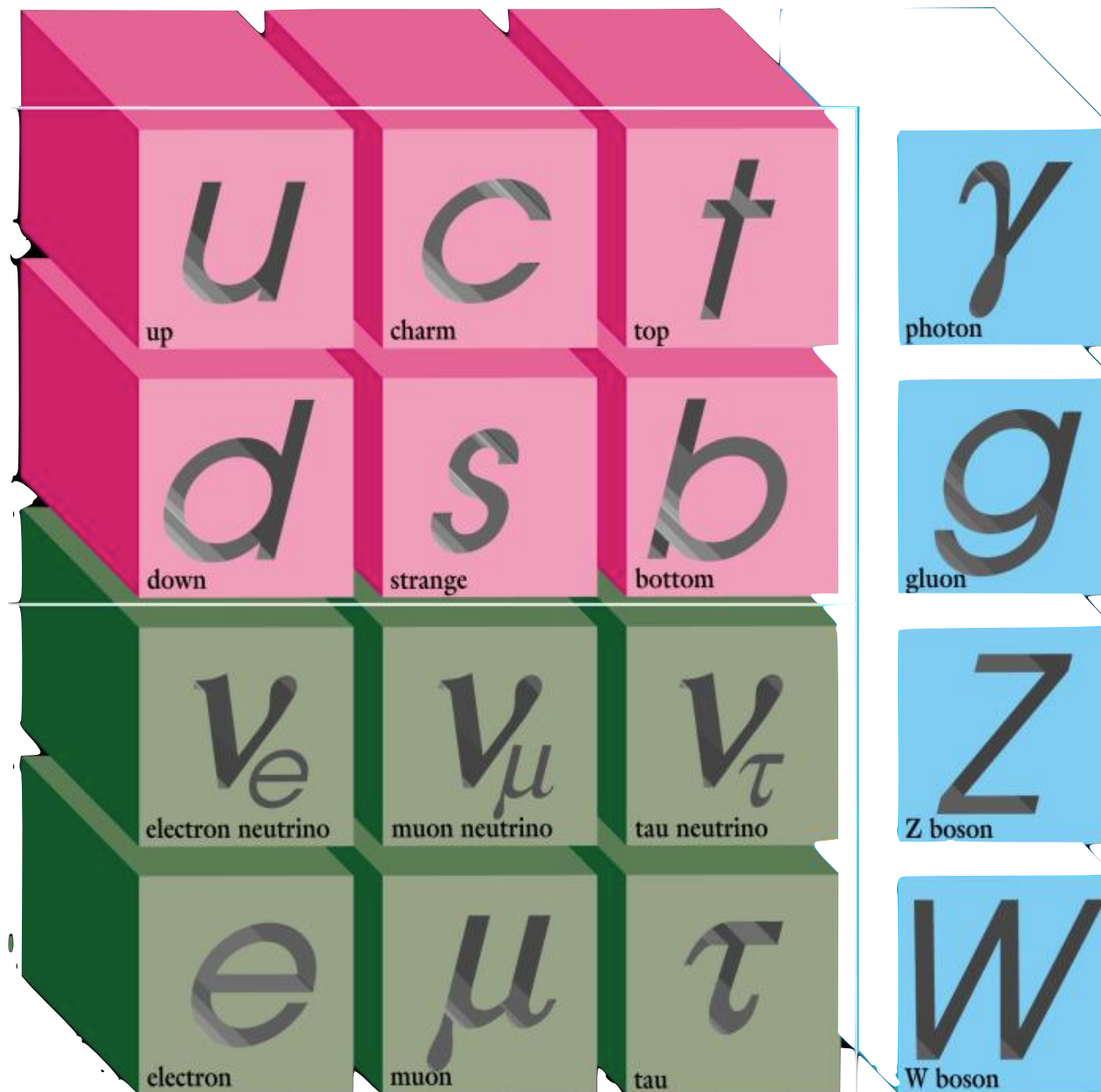
Strong + Weak + EM interactions



‘matter’ + gauge bosons



Standard Model Strong + Weak + EM interactions



‘matter’ + gauge bosons

Standard Model from Gauge Principle

✓ Each particle represented by a field $\langle \psi \rangle = \langle A^\mu \rangle = 0$

ψ Fermion (spin 1/2) A^μ Boson (spin 1)

✓ Recognize symmetry of Nature $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$

✓ Impose symmetry in local form (introduce gauge fields, just try U(1))

$$\psi(x) \rightarrow e^{i\alpha(x)} \psi(x) \quad A^\mu(x) \rightarrow A^\mu(x) + \partial^\mu \alpha(x)$$

✓ Minimal coupling between matter and gauge fields

$$\mathcal{L}_{matter}(\psi, \partial_\mu \psi) \rightarrow \mathcal{L}_{matter}(\psi, D_\mu \psi) \quad D_\mu = \partial_\mu - ig A_\mu^a(x) T^a$$

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$$\mathcal{L}_{mass}^{fermion} = -m_f \bar{\psi} \psi \equiv -m_f (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)$$

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Elementary Particles

Quarks	u up	c charm	t top	γ photon
	d down	s strange	b bottom	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	
Leptons	e electron	μ muon	τ tau	Z Z boson
				W W boson
	I	II	III	
Three Families of Matter				

With the particle content of SM
before July 4 2012 : **all massless**

Solution: **Spontaneous Symmetry Breaking**

Keep symmetry at the Lagrangian level
Break symmetry at the level of the ground state

Elementary Particles

Quarks	u up	c charm	t top	Force Carriers
	d down	s strange	b bottom	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	
Leptons	e electron	μ muon	τ tau	

	γ photon
	g gluon
	Z Z boson
	W W boson

I	II	III
---	----	-----

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Break symmetry at the level of the ground state

Higgs; Brout-Englert; Guralnik–Hagen–Kibble (~1964) Nobel 2013

Spontaneous Symmetry Breaking

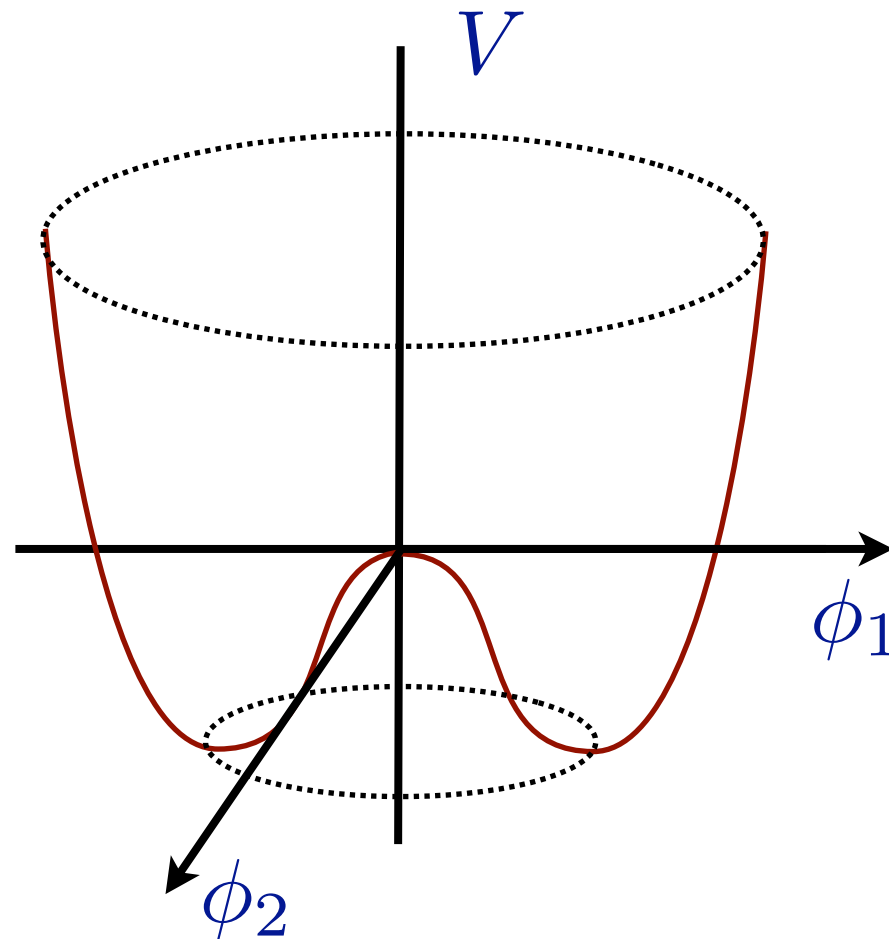
U(1) example: Introduce a complex scalar field $\phi = \frac{1}{\sqrt{2}} (\phi_1 + i \phi_2)$

gauge coupling

$$\mathcal{L}_\phi = |D_\mu \phi|^2 - V(\phi)$$

$$V(\phi) = \mu^2 |\phi|^2 + \lambda |\phi|^4 \quad \lambda > 0$$

$$\phi(x) \rightarrow \phi'(x) = e^{-i\alpha(x)} \phi(x)$$

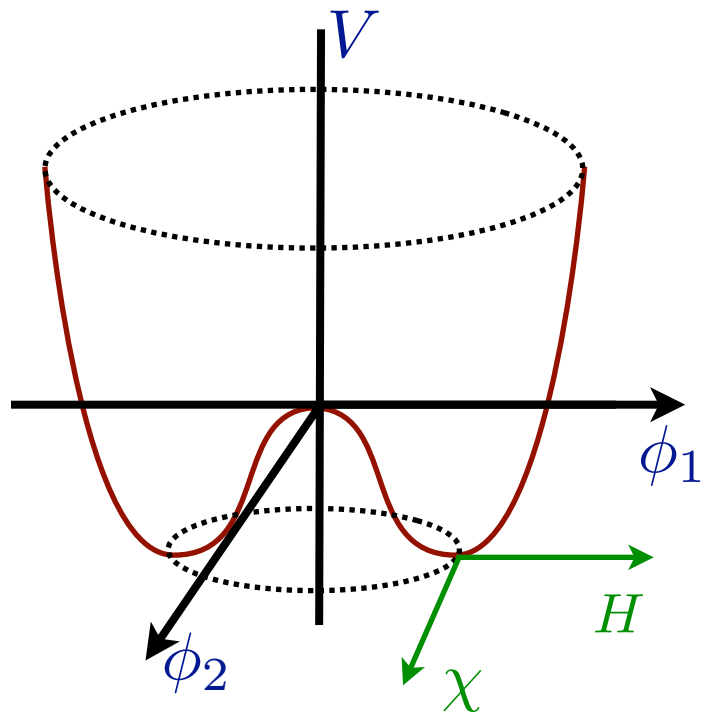


If $\mu^2 < 0$
infinite vacuums
not gauge invariant

$$|\langle \phi \rangle|^2 = \frac{-\mu^2}{2\lambda} \equiv v^2$$



Spontaneous Symmetry Breaking



$$\phi(x) \equiv \frac{1}{\sqrt{2}} e^{i \frac{\chi(x)}{v}} (v + H(x))$$

moves between vev's
VEV not gauge invariant

perturbation away from minimum

$\chi(x)$ field can be **rotated away** by gauge transformation (unitary gauge)

scalar boson mass term

$$m_H^2 = -2\mu^2$$

interactions

$$|D_\mu \phi|^2 = \frac{1}{2} \partial_\mu H \partial^\mu H - \frac{1}{2} (-2\mu^2 H^2) + \frac{q^2 v^2}{2} A_\mu A^\mu + \frac{q^2}{2} A_\mu A^\mu H^2 + v q^2 A_\mu A^\mu H$$

$m_A = qv$ Gauge boson mass term

Degree of freedom “eaten” to provide longitudinal polarization

Two parameters in the potential

$$\mu, \lambda \longrightarrow m_H, m_A$$

Spontaneous Symmetry Breaking $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$

The SM example: Introduce a doublet of complex scalar fields

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} \quad V(\Phi^\dagger \Phi) = \lambda \left(\Phi^\dagger \Phi - \frac{v^2}{2} \right)^2$$

Transforms as $(1, 2, 1/2)$ under $SU(3)_C \times SU(2)_L \times U(1)_Y$

$$\langle \Phi \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \quad \text{vacuum state neutral}$$

- Photon remains massless

- W and Z acquire masses

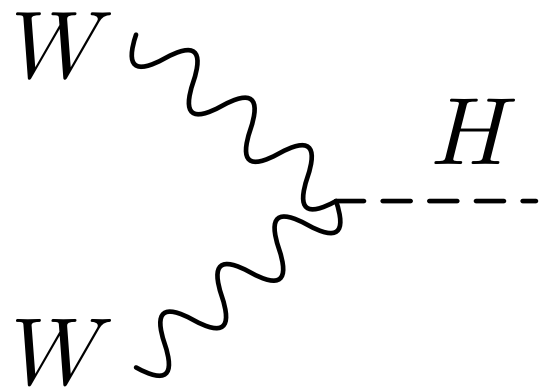
$$m_W^2 = \frac{g^2 v^2}{4} \quad m_Z^2 = \frac{(g^2 + g'^2) v^2}{4} = \frac{m_W^2}{\cos^2 \theta_W}$$

- **vev** known

$$v = \sqrt{\frac{1}{\sqrt{2} G_F}} \approx 246.22 \text{ GeV}$$

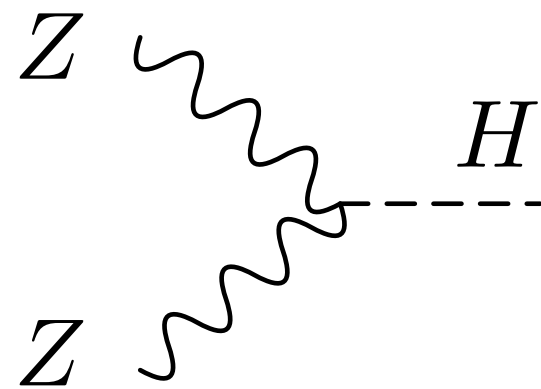
Interactions with Higgs Boson

WWH and ZZH couplings



A Feynman diagram showing two wavy lines labeled W meeting at a vertex with a dashed line labeled H extending from the vertex.

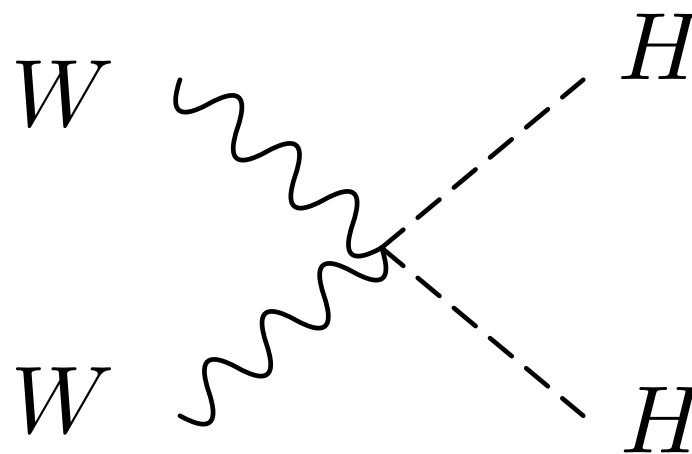
$$2i \frac{m_W^2}{v} g^{\mu\nu}$$



A Feynman diagram showing two wavy lines labeled Z meeting at a vertex with a dashed line labeled H extending from the vertex.

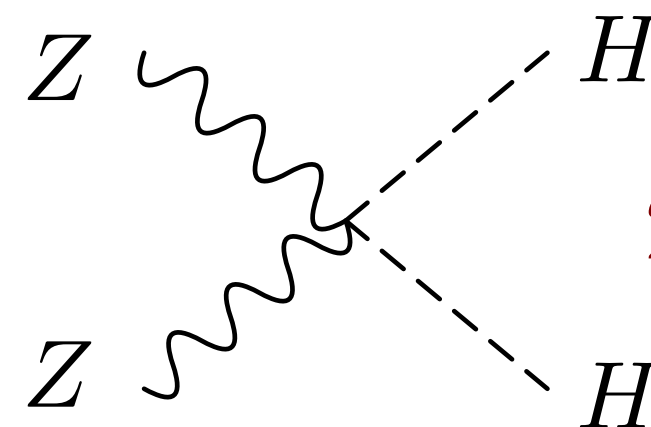
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$$2i \frac{m_Z^2}{v^2} g^{\mu\nu}$$

Fermion Masses

~~$$\mathcal{L}_{mass}^{fermion} = -m_f \bar{\psi}\psi \equiv -m_f (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)$$~~

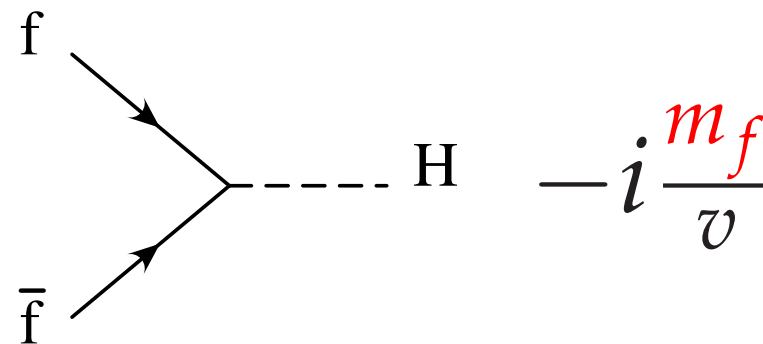
Generate Fermion masses through Yukawa interactions

$$\mathcal{L}_{Yukawa} = -\lambda_d \bar{Q}_L \Phi d_R - \lambda_u \bar{Q}_L (i\sigma_2 \Phi^*) u_R - \lambda_e \bar{L}_L \Phi e_R + \text{h.c.}$$

Fermion mass

$$m_d = \lambda_d \frac{v}{\sqrt{2}}$$

Fermion-Higgs interaction



All interactions with Higgs boson proportional to particle mass

Back to Higgs potential

$$V\left(\Phi^\dagger\Phi\right)=\lambda\left(\Phi^\dagger\Phi-\frac{v^2}{2}\right)^2$$

Expanding the potential around the vacuum

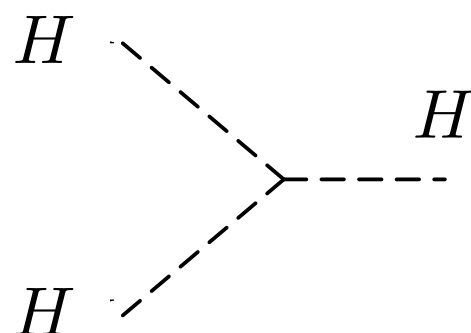
$$\Phi(x)=\frac{1}{\sqrt{2}}\begin{pmatrix}0\\v+H(x)\end{pmatrix}$$

results in

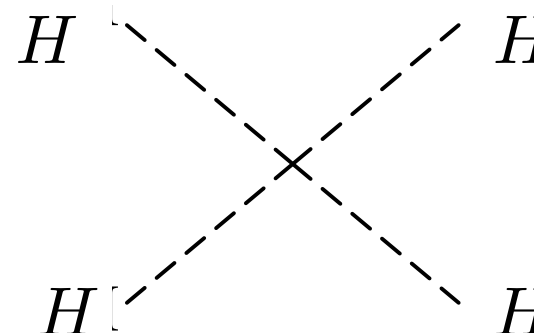
$$V=\frac{\lambda}{4}\left(2vH+H^2\right)^2=\frac{1}{2}(\textcolor{red}{2\lambda v^2})H^2+\lambda v\textcolor{blue}{H^3}+\frac{\lambda}{4}\textcolor{blue}{H^4}$$

Higgs boson mass $\textcolor{green}{m_H^2=2\lambda v^2}$ can not be predicted (*only unknown*)

Cubic and Quartic self-couplings



$$-3i\frac{m_H^2}{v}$$



$$-3i\frac{m_H^2}{v^2}$$

Search and Discovery

Higgs search was never easy....

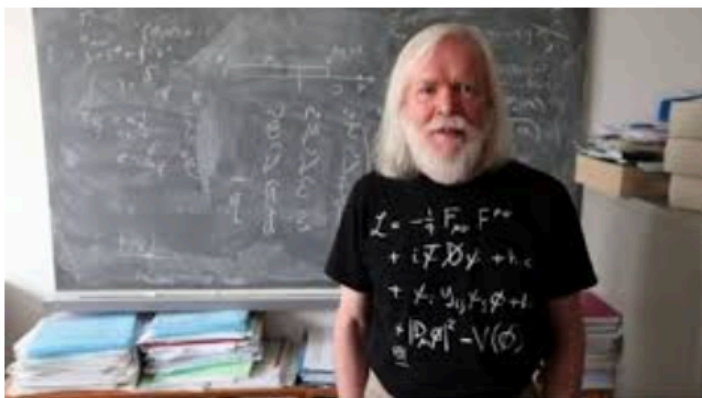
A PHENOMENOLOGICAL PROFILE OF THE HIGGS BOSON

John ELLIS, Mary K. GAILLARD * and D.V. NANOPOULOS **
CERN, Geneva

Received 7 November 1975

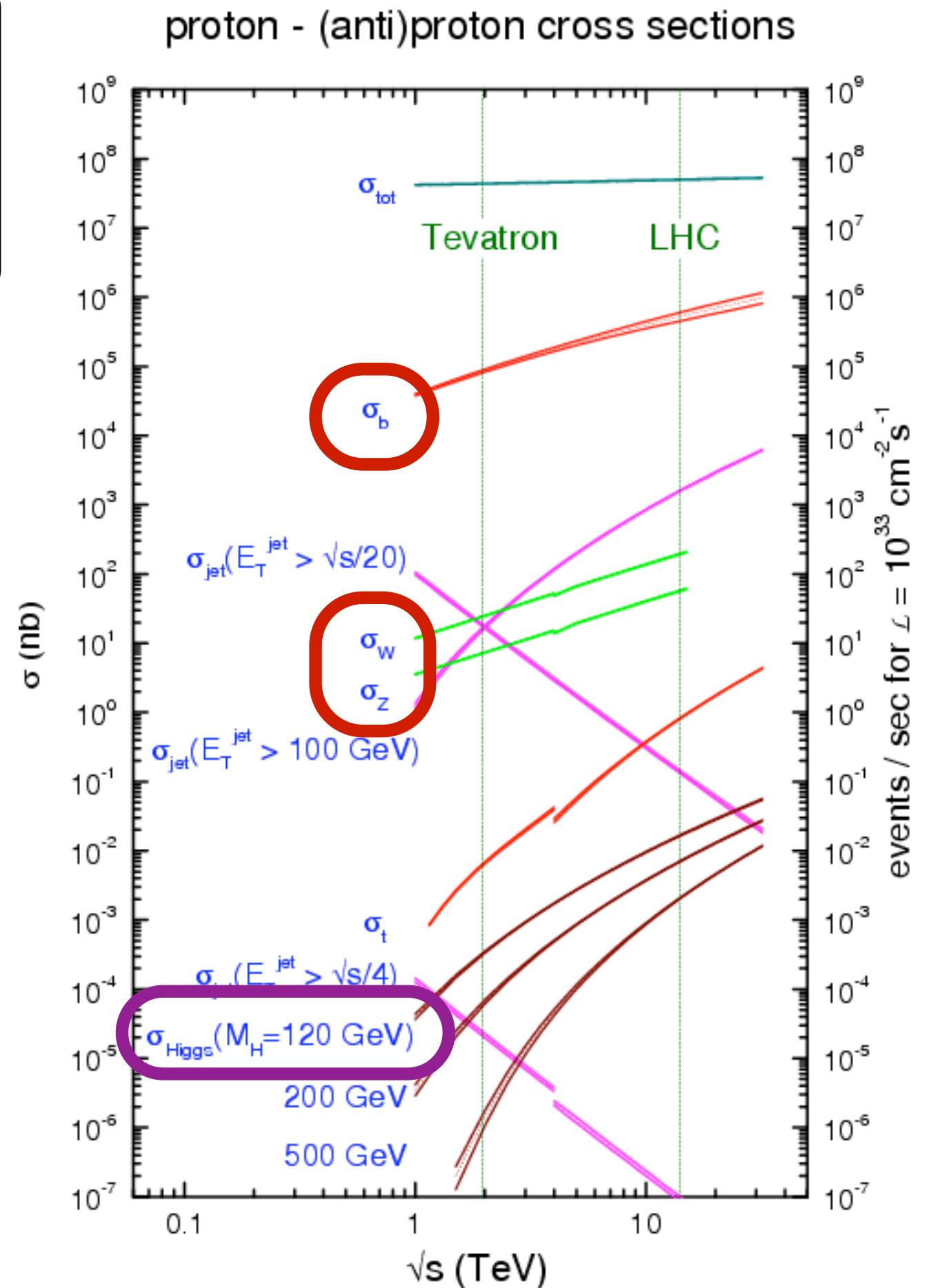
A discussion is given of the production, decay and observability of the scalar Higgs boson H expected in gauge theories of the weak and electromagnetic interactions such as the Weinberg-Salam model. After reviewing previous experimental limits on the mass of

We should perhaps finish with an apology and a caution. We apologize to experimentalists for having no idea what is the mass of the Higgs boson, unlike the case with charm [3,4] and for not being sure of its couplings to other particles, except that they are probably all very small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.



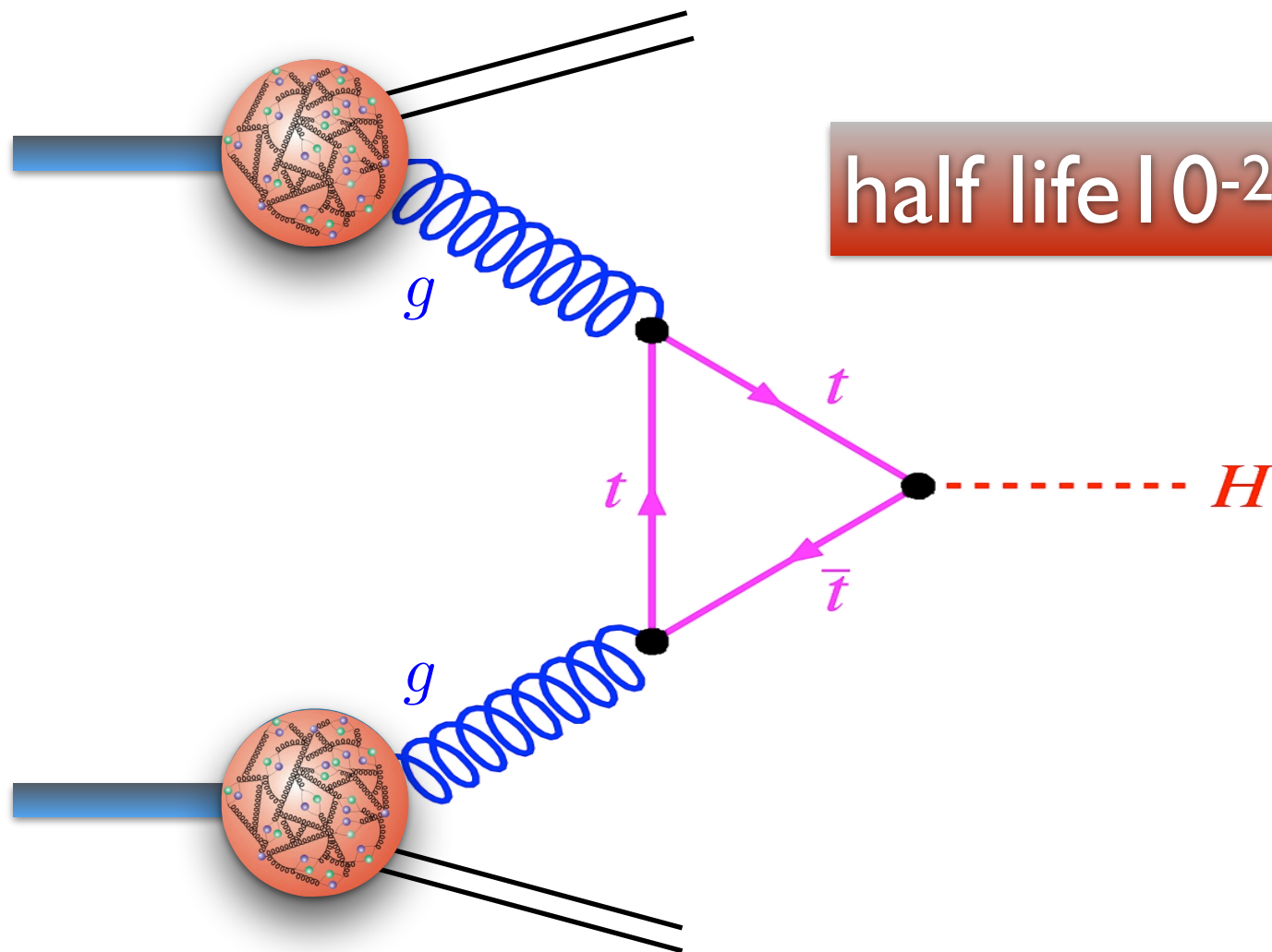
Particle production at hadronic colliders

- 1 Higgs in 10^{12} !

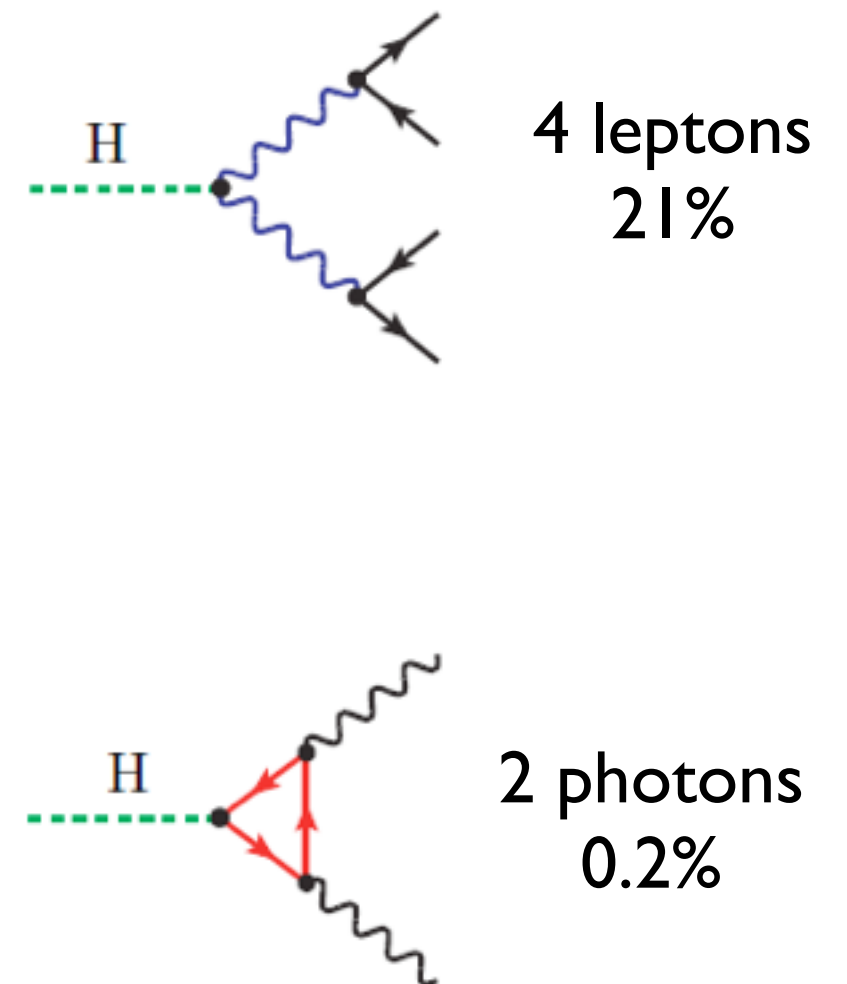


Higgs in hadronic colliders

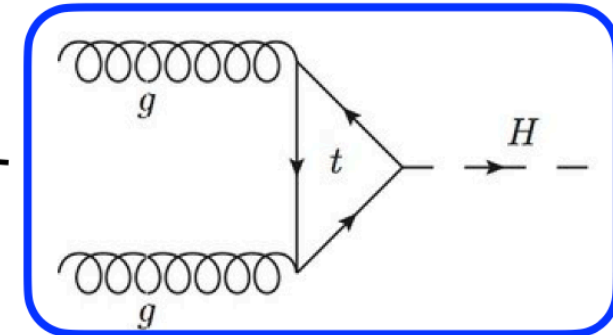
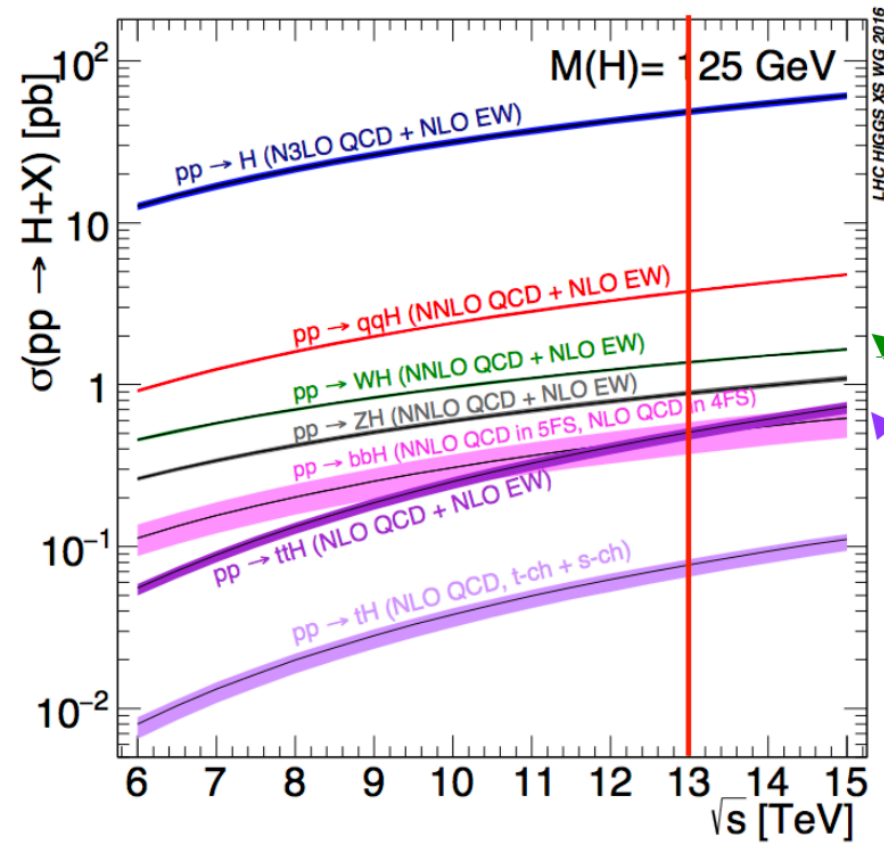
Production



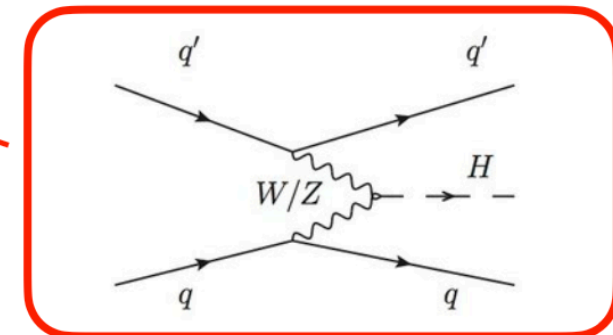
Decay



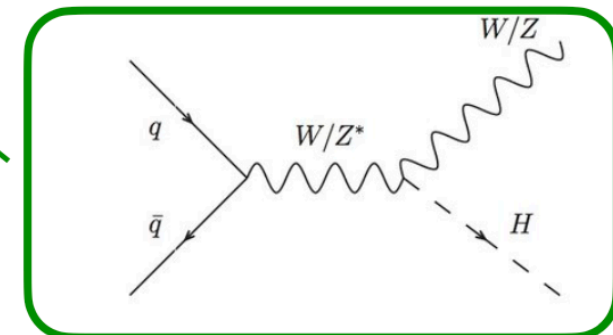
Higgs in hadronic colliders (all channels)



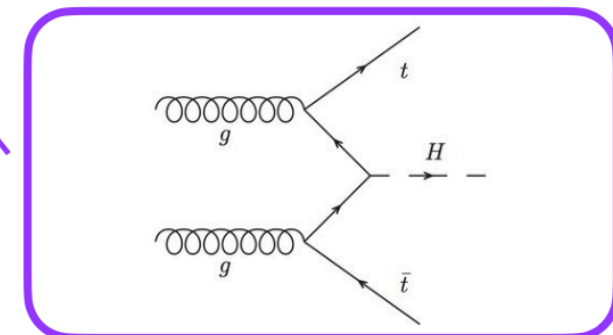
N³LO



N³LO



NNLO

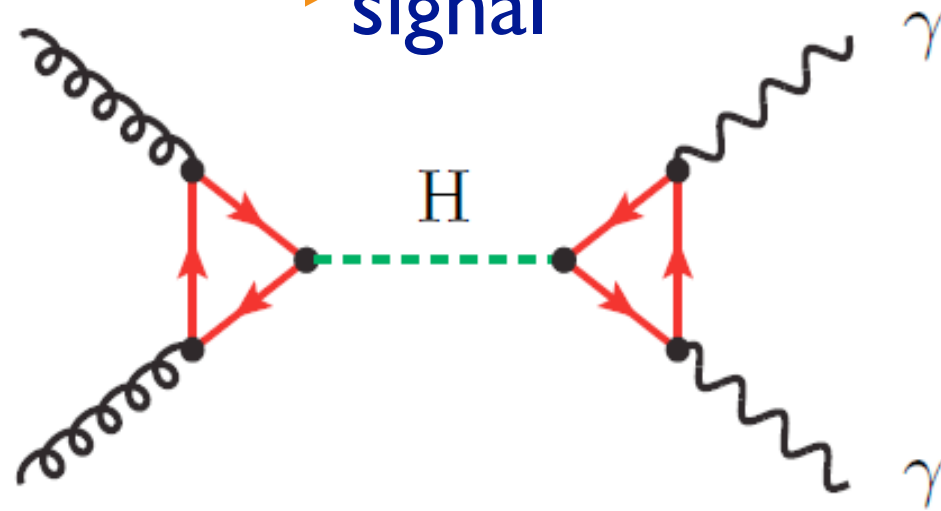


**NLO +
approx
NNLO**

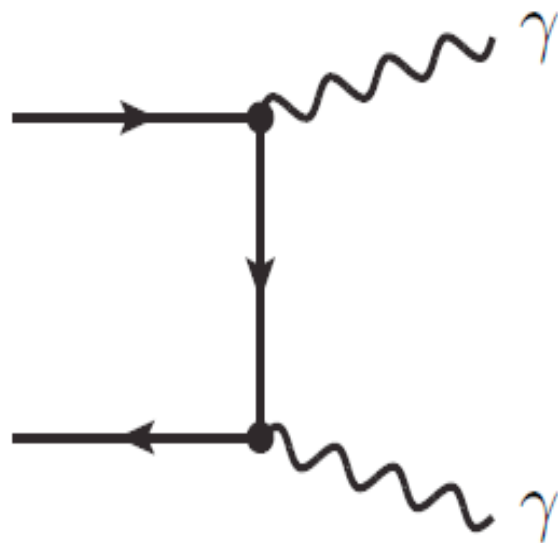
\sqrt{s} (TeV)	Production cross section (in pb) for $m_H = 125$ GeV					total
	ggF	VBF	WH	ZH	$t\bar{t}H$	
1.96	$0.95^{+17\%}_{-17\%}$	$0.065^{+8\%}_{-7\%}$	$0.13^{+8\%}_{-8\%}$	$0.079^{+8\%}_{-8\%}$	$0.004^{+10\%}_{-10\%}$	1.23
7	$16.9^{+5\%}_{-5\%}$	$1.24^{+2\%}_{-2\%}$	$0.58^{+3\%}_{-3\%}$	$0.34^{+4\%}_{-4\%}$	$0.09^{+8\%}_{-14\%}$	19.1
8	$21.4^{+5\%}_{-5\%}$	$1.60^{+2\%}_{-2\%}$	$0.70^{+3\%}_{-3\%}$	$0.42^{+5\%}_{-5\%}$	$0.13^{+8\%}_{-13\%}$	24.2
13	$48.6^{+5\%}_{-5\%}$	$3.78^{+2\%}_{-2\%}$	$1.37^{+2\%}_{-2\%}$	$0.88^{+5\%}_{-5\%}$	$0.50^{+9\%}_{-13\%}$	55.1
14	$54.7^{+5\%}_{-5\%}$	$4.28^{+2\%}_{-2\%}$	$1.51^{+2\%}_{-2\%}$	$0.99^{+5\%}_{-5\%}$	$0.60^{+9\%}_{-13\%}$	62.1

HIGGS boson discovery at the LHC

► signal

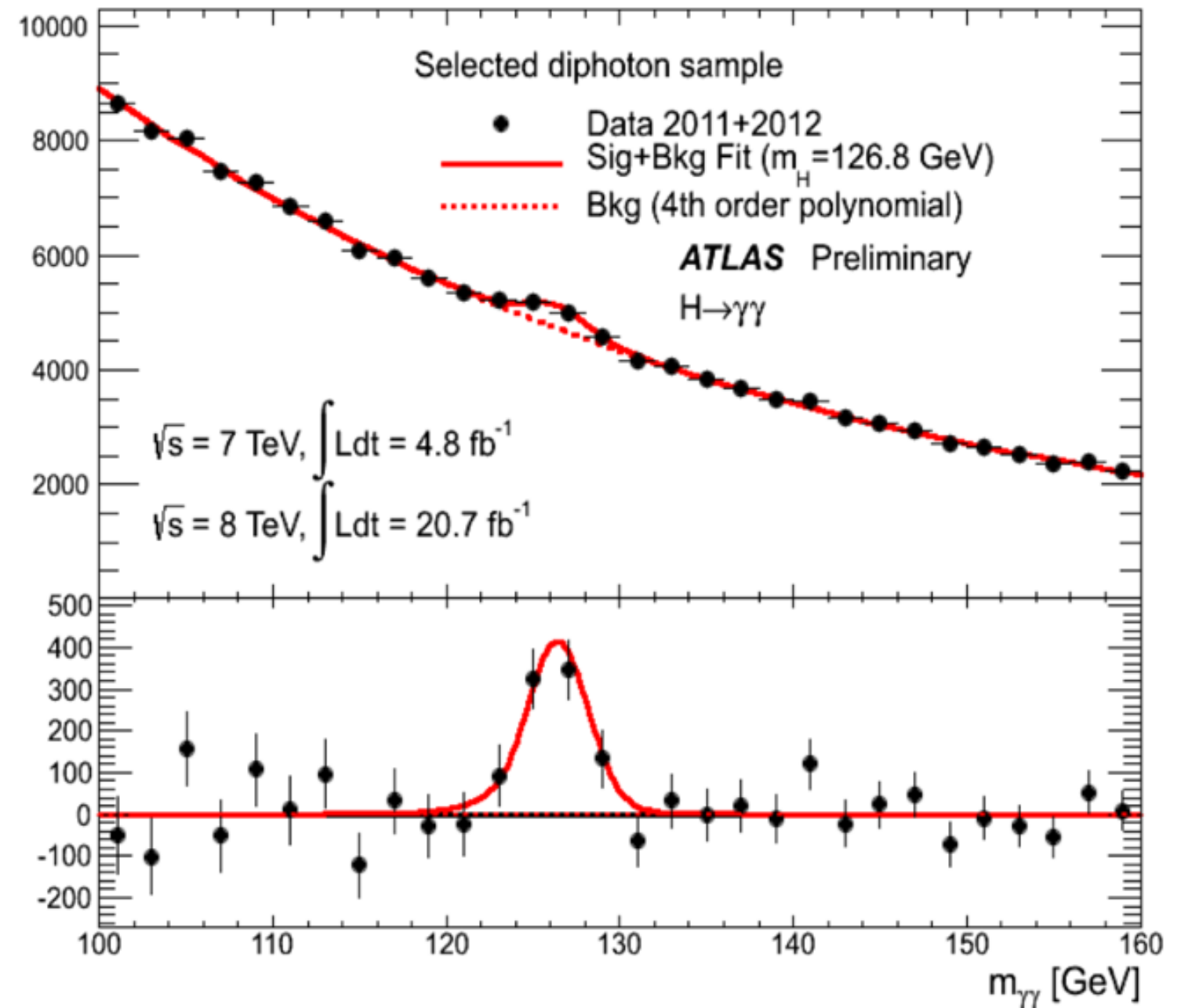


► background



Events / 2 GeV

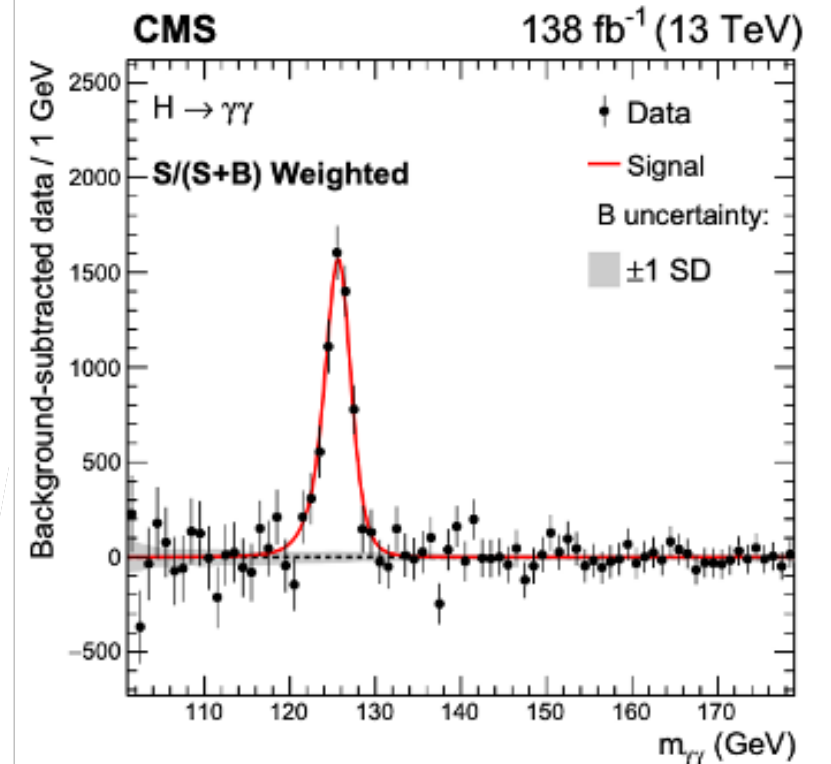
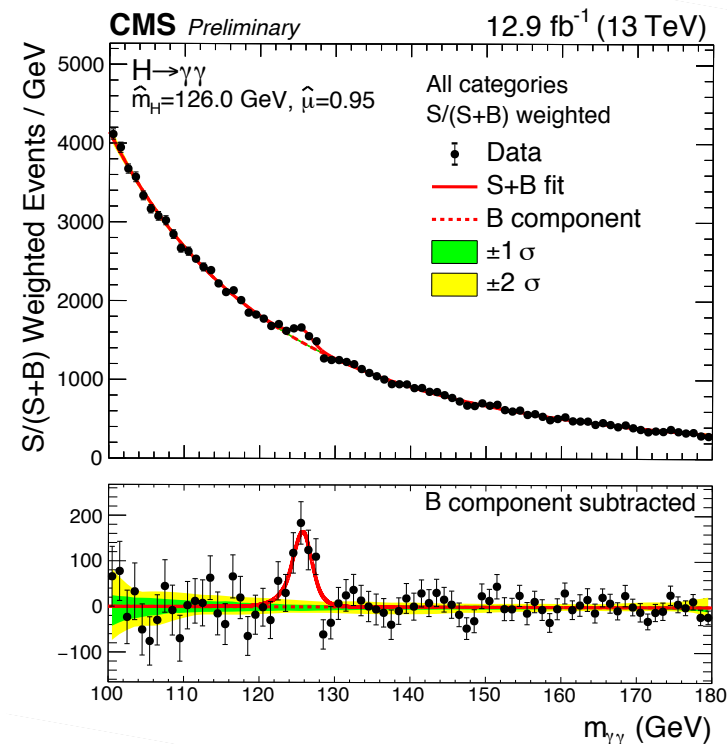
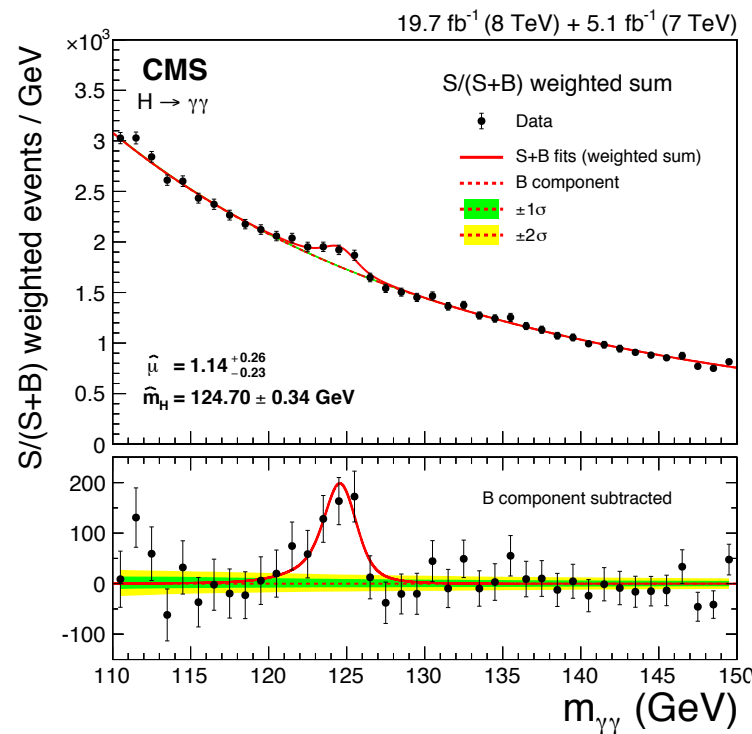
Events - Fitted bkg



$$(p_{\gamma_1} + p_{\gamma_2})^2 = m_{\gamma\gamma}^2$$

Properties

Higgs discovered, re-discovered and re-re-discovered



today: 30 x Higgs

- ▶ Mass uncertainty (125 GeV) ~ 0.2% !

- ▶ Signal strength

$$\text{ATLAS } \mu = 1.05 \pm 0.04 \text{ (th)} \pm 0.03 \text{ (exp)} \pm 0.03 \text{ (stat)}$$

$$\text{CMS } \mu = 1.002 \pm 0.036 \text{ (th)} \pm 0.033 \text{ (exp)} \pm 0.029 \text{ (stat)}$$

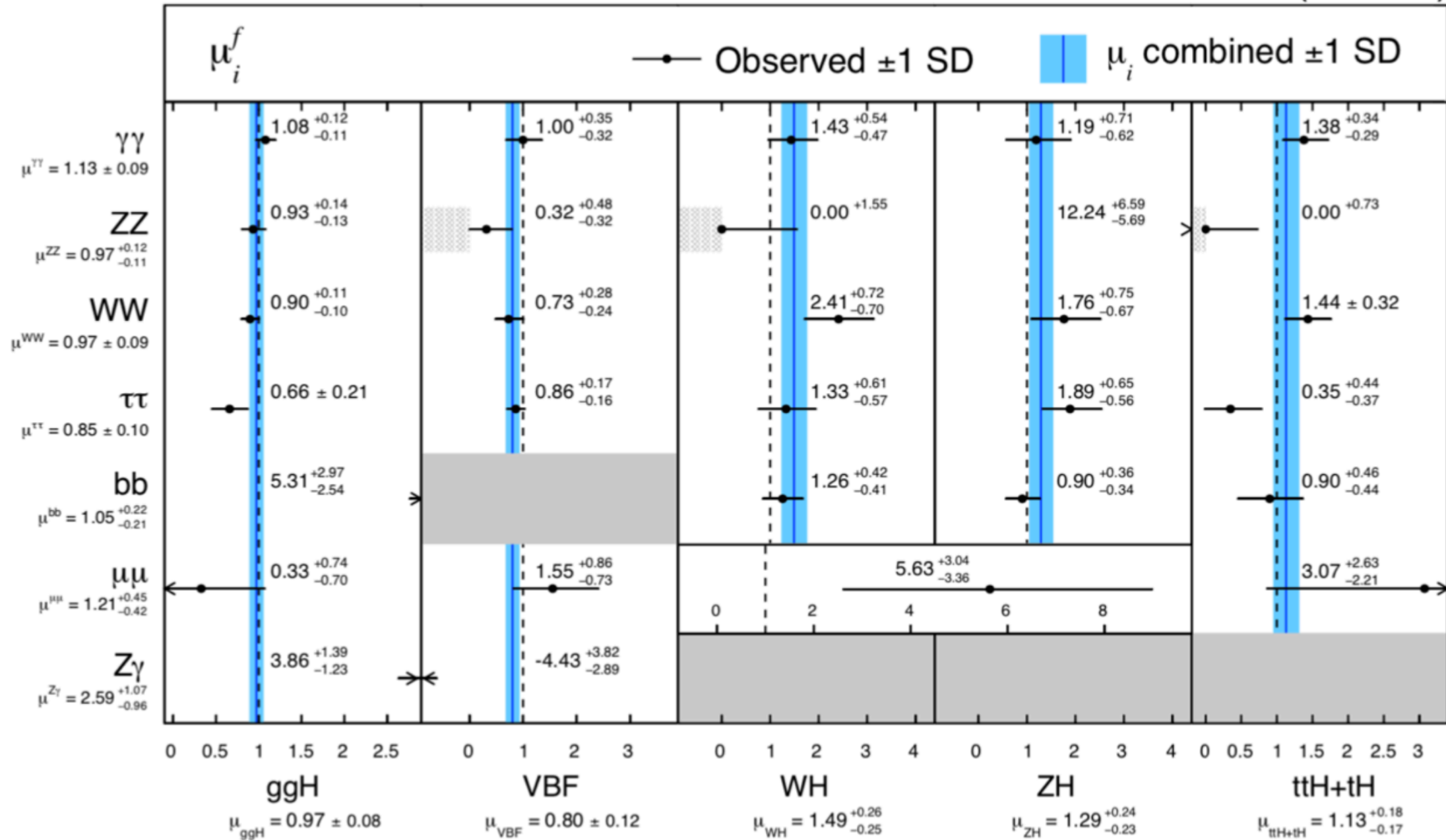
- ▶ Cross sections in agreement with SM

- ▶ Electric charge = 0

- ▶ Spin = 0

CMS

138 fb⁻¹ (13 TeV)



Overall signal strength normalised to SM expectation:

$$\mu = 1.002 \pm 0.057$$

Higgs couplings

- ▶ Observation in all production channels, including $t\bar{t}H$!
- ▶ Observation of decay in $\gamma\gamma$, WW , ZZ , $b\bar{b}$, $\tau\bar{\tau}$
- ▶ Second generation more difficult but on the way to $c\bar{c}$, $\mu^+\mu^-$
- ▶ No measurement of HHH and HHHH couplings (see later)

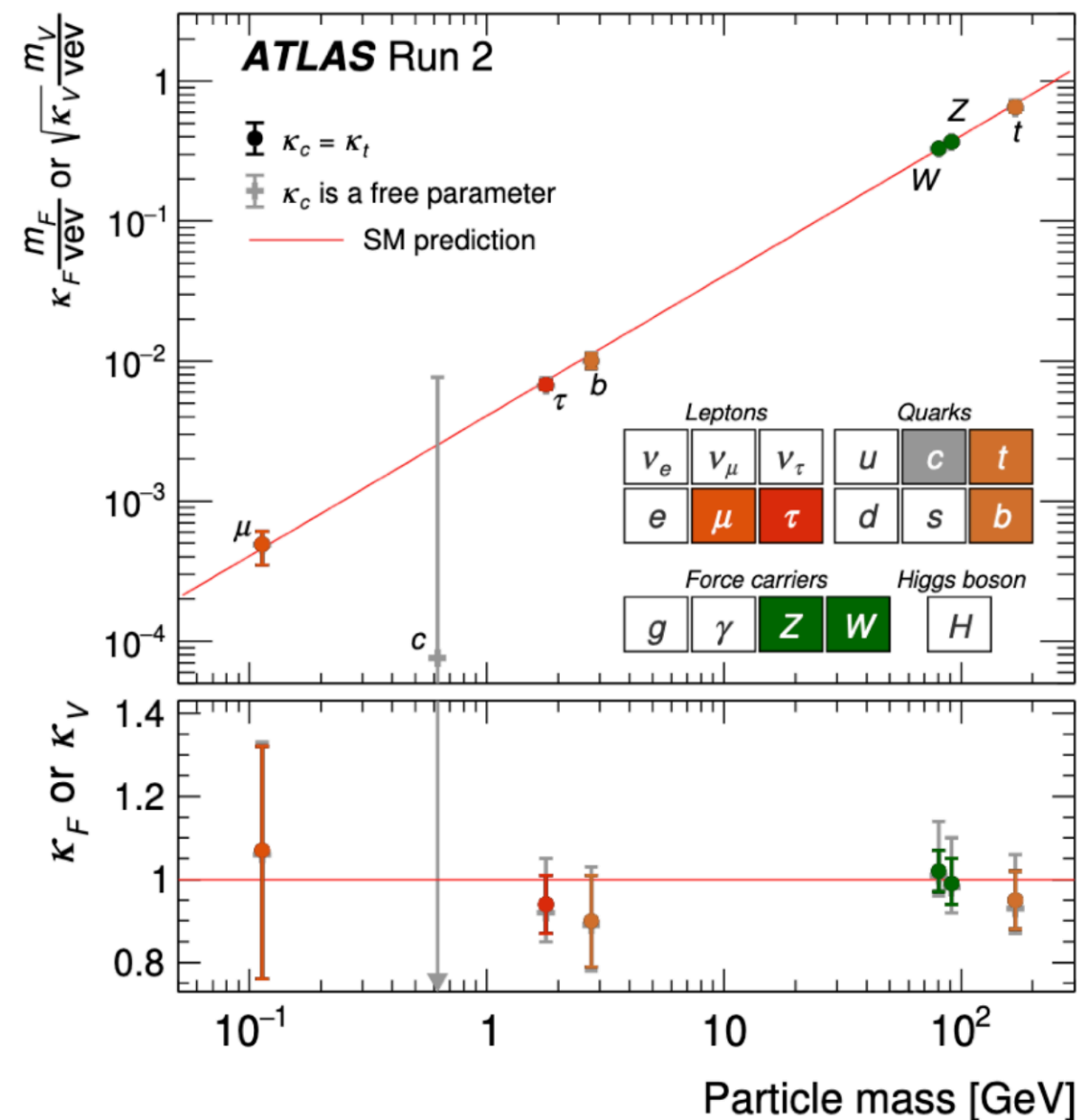
Three Generations of Matter (Fermions)

	I	II	III	
mass	2.4 MeV/c ²	1.27 GeV/c ²	171.2 GeV/c ²	0
charge	2/3	2/3	2/3	0
spin	1/2	1/2	1/2	1
name	u up	c charm	t top	γ photon
			✓	✓
	4.8 MeV/c ²	104 MeV/c ²	4.2 GeV/c ²	0
	-1/3	-1/3	-1/3	0
	1/2	1/2	1/2	1
	d down	s strange	b bottom	g gluon
			✓	✓
	<2.2 eV/c ²	<0.17 MeV/c ²	<15.5 MeV/c ²	91.2 GeV/c ²
	0	0	0	0
	1/2	1/2	1/2	1
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z⁰ Z boson
				✓
	0.511 MeV/c ²	105.7 MeV/c ²	1.777 GeV/c ²	80.4 GeV/c ²
	-1	-1	-1	±1
	1/2	1/2	1/2	1
	e electron	μ muon	τ tau	W[±] W boson
			✓	✓

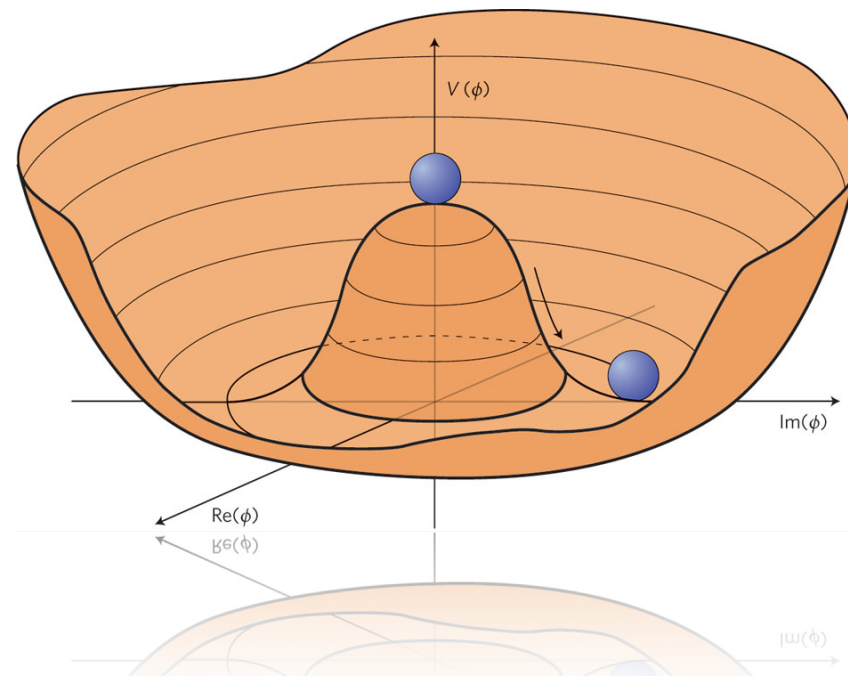
Quarks

Leptons

Gauge Bosons



- **NO information about Higgs self-couplings**
- Need to study **multiple Higgs production** to explore mechanism for Spontaneous Symmetry Breaking



$$V = \frac{\lambda}{4} (2vH + H^2)^2 = \frac{1}{2} (2\lambda v^2) H^2 + \lambda v H^3 + \frac{\lambda}{4} H^4$$

$$H \text{---} H \text{---} H \quad -3i \frac{m_H^2}{v}$$

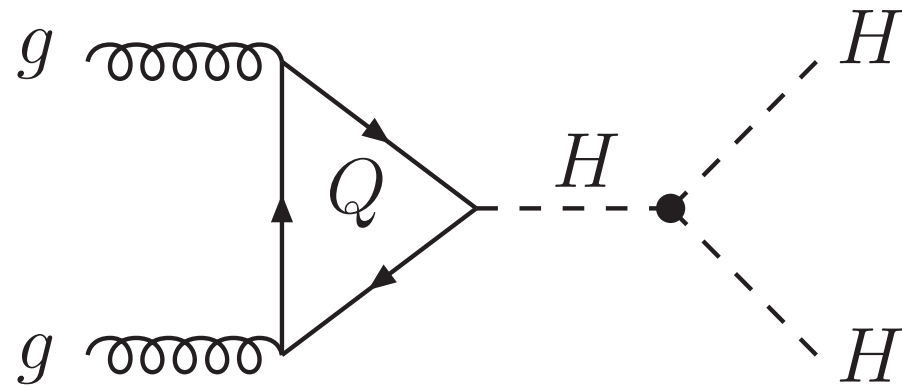
$$H \text{---} H \text{---} H \quad -3i \frac{m_H^2}{v^2}$$

Higgs pair production

- To measure $3H$ coupling one needs to look at $2H$ production

much smaller **cross sections**

@ 14 TeV



$\sim 40 \text{ fb}$

difficult but
next discovery?

HL-LHC $4000 \text{ fb}^{-1} \sim 160.000 \text{ HH}$

- Production rate drops 3 orders of magnitude for each extra H

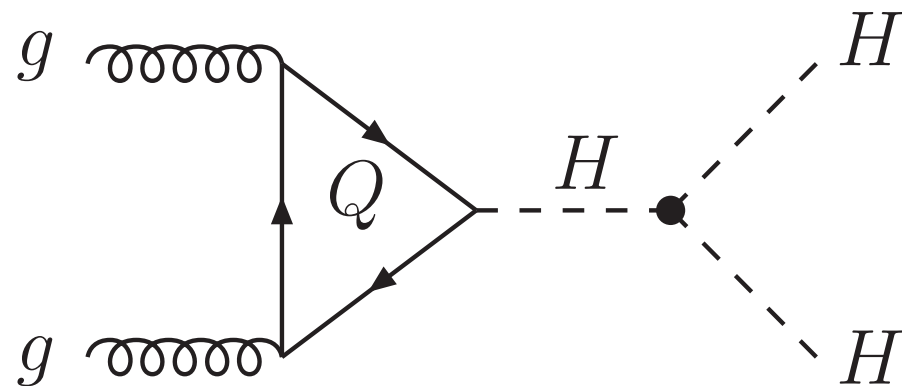
Experimental limit in self-coupling today $\sim 5 \times \text{SM}$

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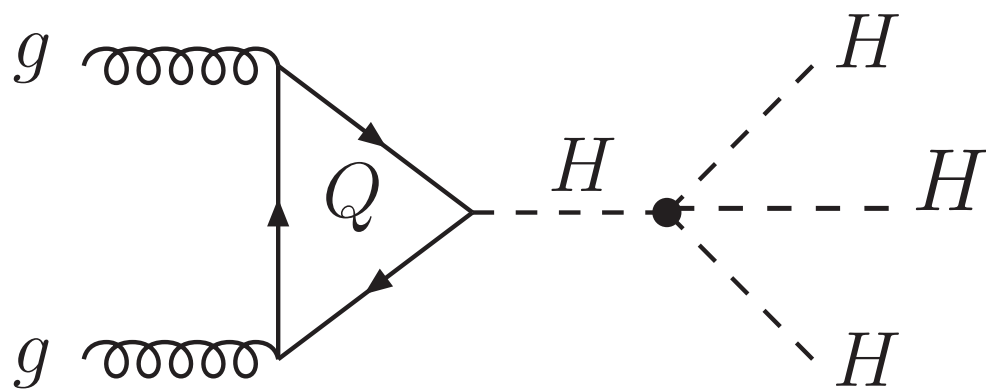
@ 14 TeV



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HL-LHC $4000 \text{ fb}^{-1} \sim 160.000 \text{ HH}$



$\sim 0.05 \text{ fb}$

impossible at LHC

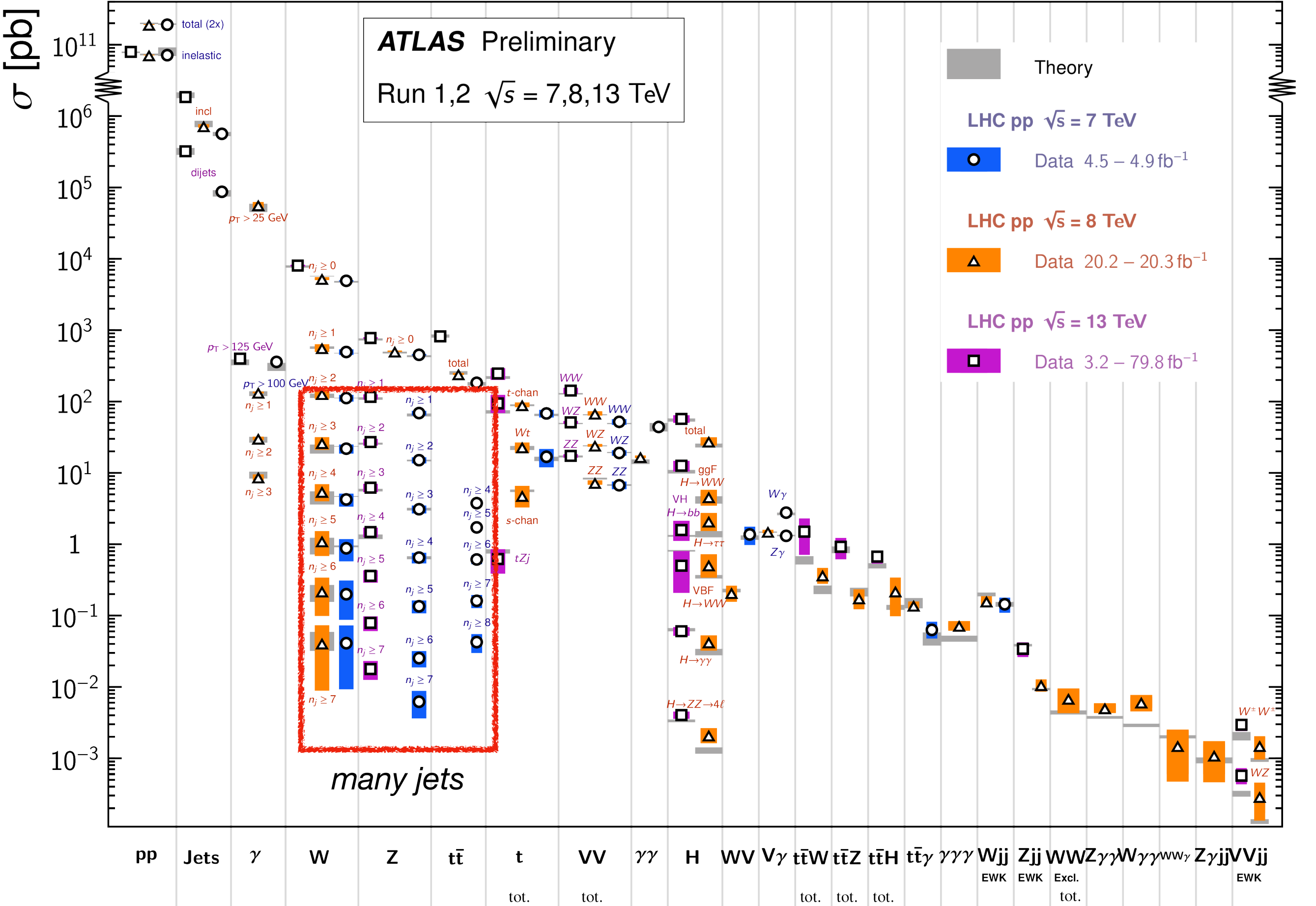
- Production rate drops 3 orders of magnitude for each extra H

Experimental limit in self-coupling today $\sim 5 \times \text{SM}$

Future 1

Standard Model Production Cross Section Measurements

Status: July 2018

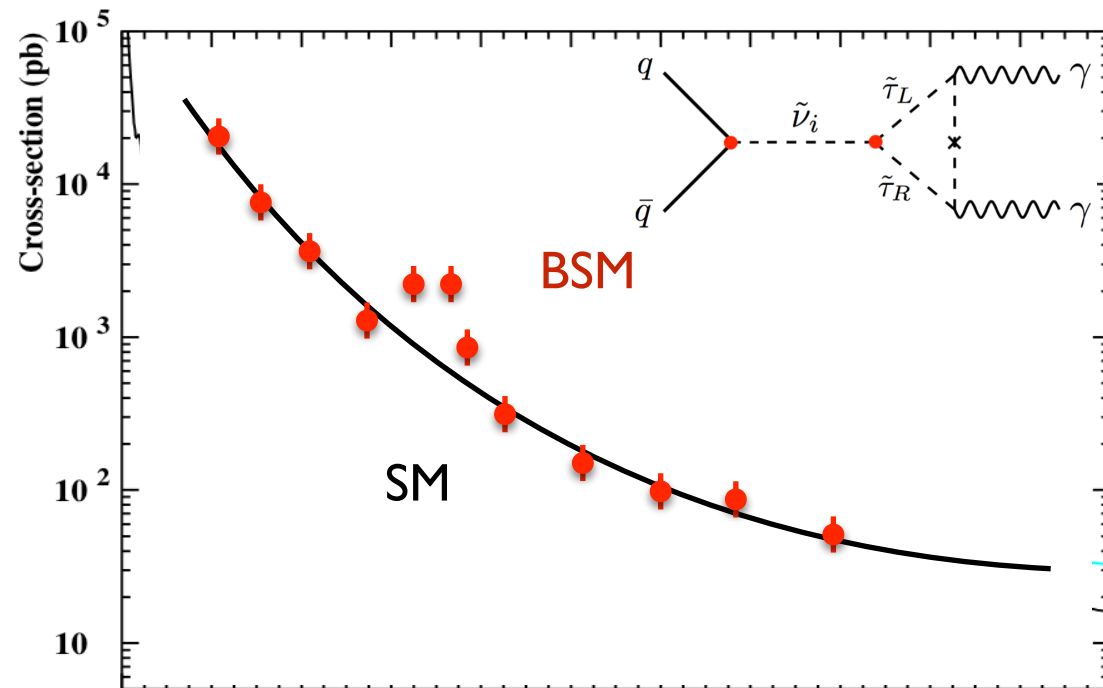


Everything looks SM-like within (large) uncertainties

- ▶ There is plenty of room for discoveries yet

Everything looks SM-like within (large) uncertainties

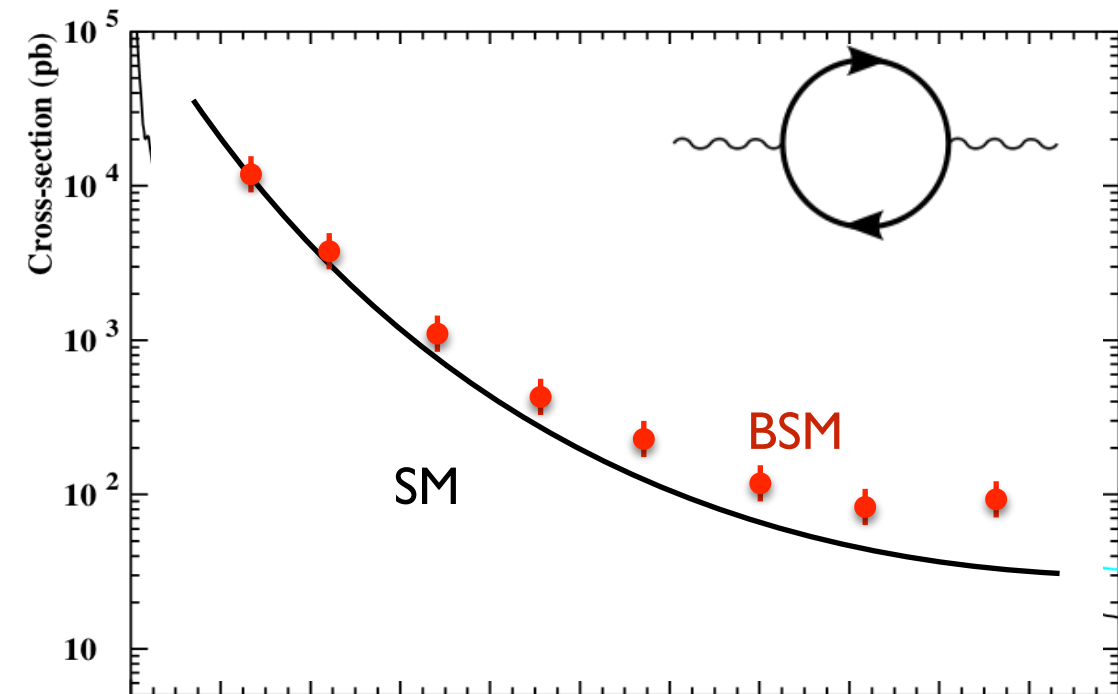
► There is plenty of room for discoveries yet



Search for new *states*
Resonances: “Descriptive TH”

• Explore Higgs sector with precision

- less known (room for surprises!)
- more sensitive (Portal) to new physics
- *Potential* : look at multiple Higgs production



Search for new *interactions*
Deviations: “Precision TH”
+ EFT description

► **EXP and TH : (for Higgs) Precision is the name of the game**

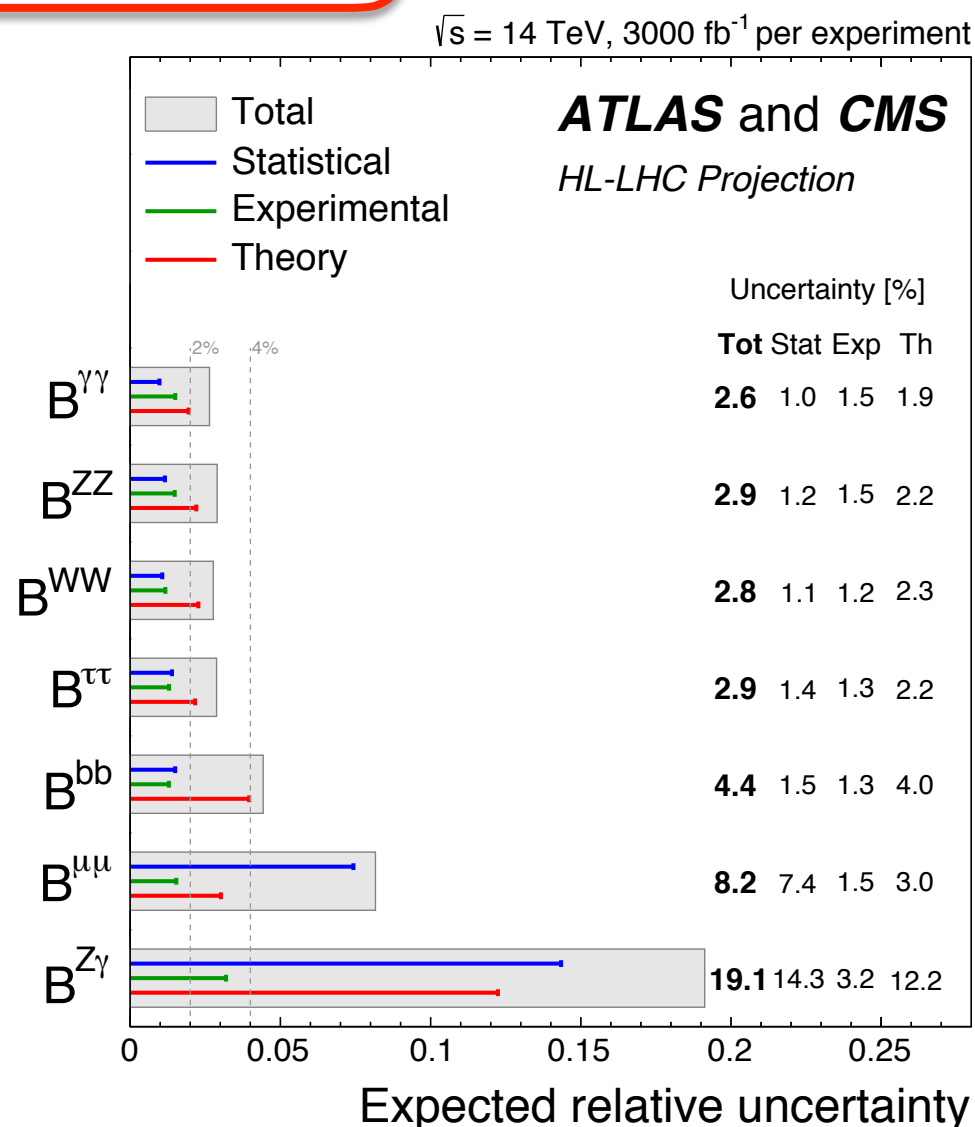
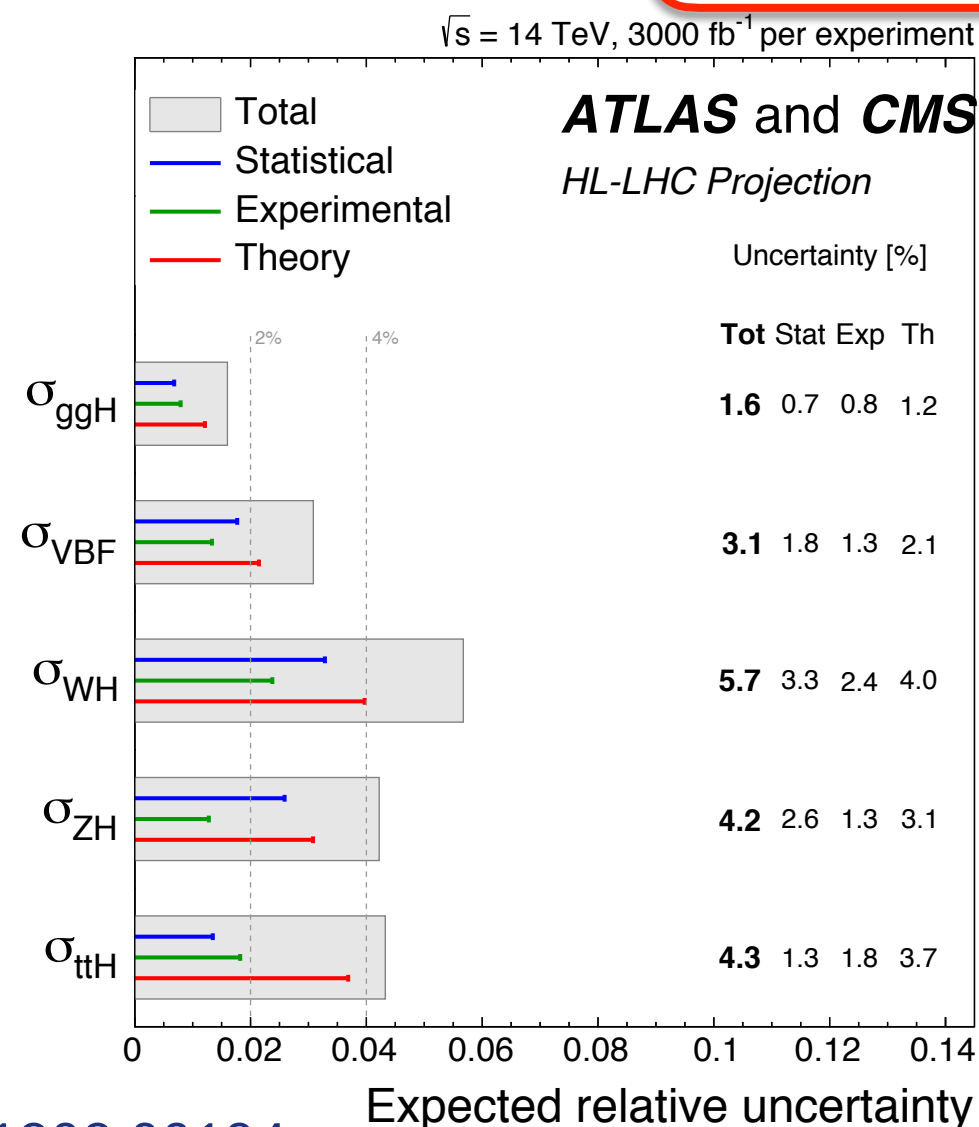
electron	(1897) Thompson
positron	(1932) Anderson
muon	(1937) Cosmic radiation-Cloud chamber
neutrino electron	(1956) Savannah River Plant
neutrino muon	(1962) BNL
u,d,s	(1969) SLAC
charm	(1974) SLAC-BNL
tau	(1975) SLAC-SPEAR-LBL
bottom	(1977) E288
gluon	(1979) DORIS/PETRA
W/Z	(1983) UA1
top	(1995) Tevatron
neutrino tau	(2000) DONUT
Higgs	(2012) LHC

One big discovery per experiment...

► HL-LHC projections

~20 years from now!

(S2) TH uncertainties scaled down by factor 2, EXP scaled according to $\sqrt{\mathcal{L}}$



1902.00134

- Theoretical uncertainties on SM predictions generally largest component
- Precision becomes critical
- TH: can we improve calculations? Where? How?

Very nice Lagrangian, but can not be solved analytically
perturbative expansion



$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi + h.c. \\ & + \chi_i y_{ij} \chi_j \phi + h.c. \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

Actual
Taylor expansion

$$f(x) = \overbrace{f(0) + f'(0)x}^{\text{Actual Taylor expansion}} + \frac{f''(0)}{2!}x^2 + \frac{f'''(0)}{3!}x^3 \dots$$

Lies invented by mathematicians
to feel superior to physicists

$$\frac{f''(0)}{2!}x^2 + \frac{f'''(0)}{3!}x^3 \dots$$

Very nice Lagrangian, but can not be solved analytically



perturbative expansion

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► Perturbative parameter

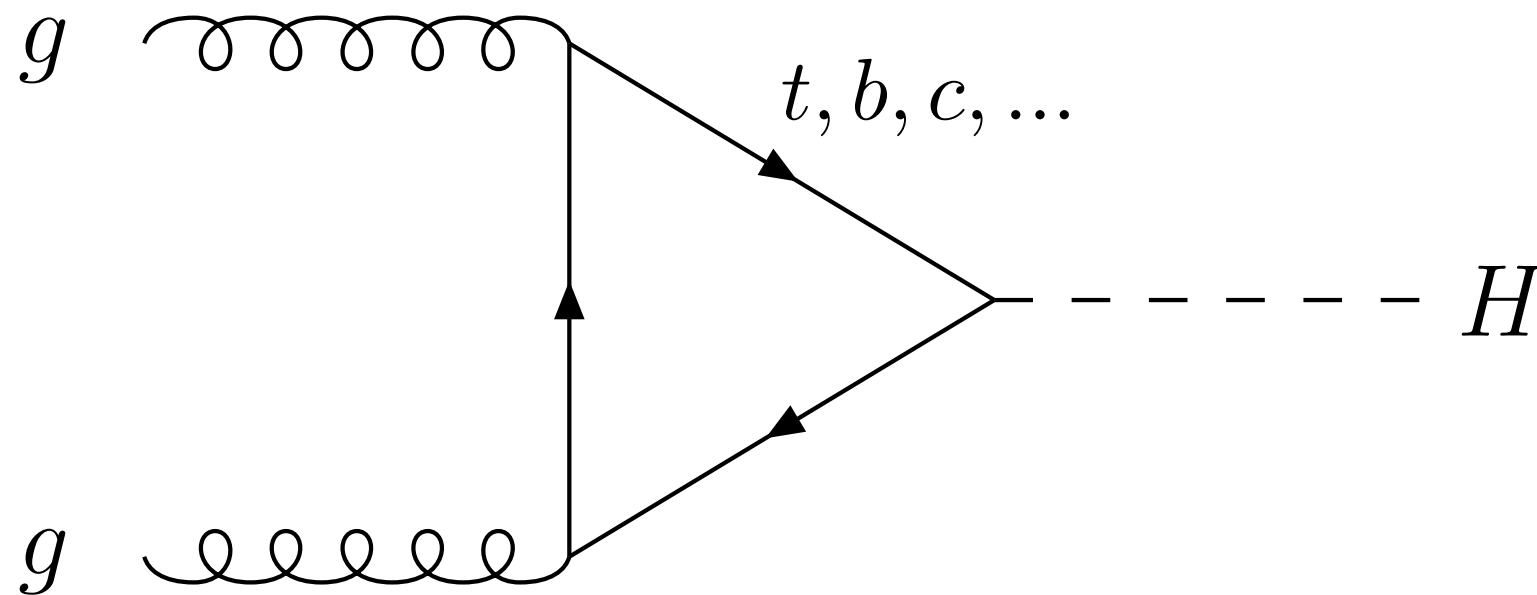
$$\alpha_s \sim 0.11$$

► Strong coupling

QCD requires high orders...

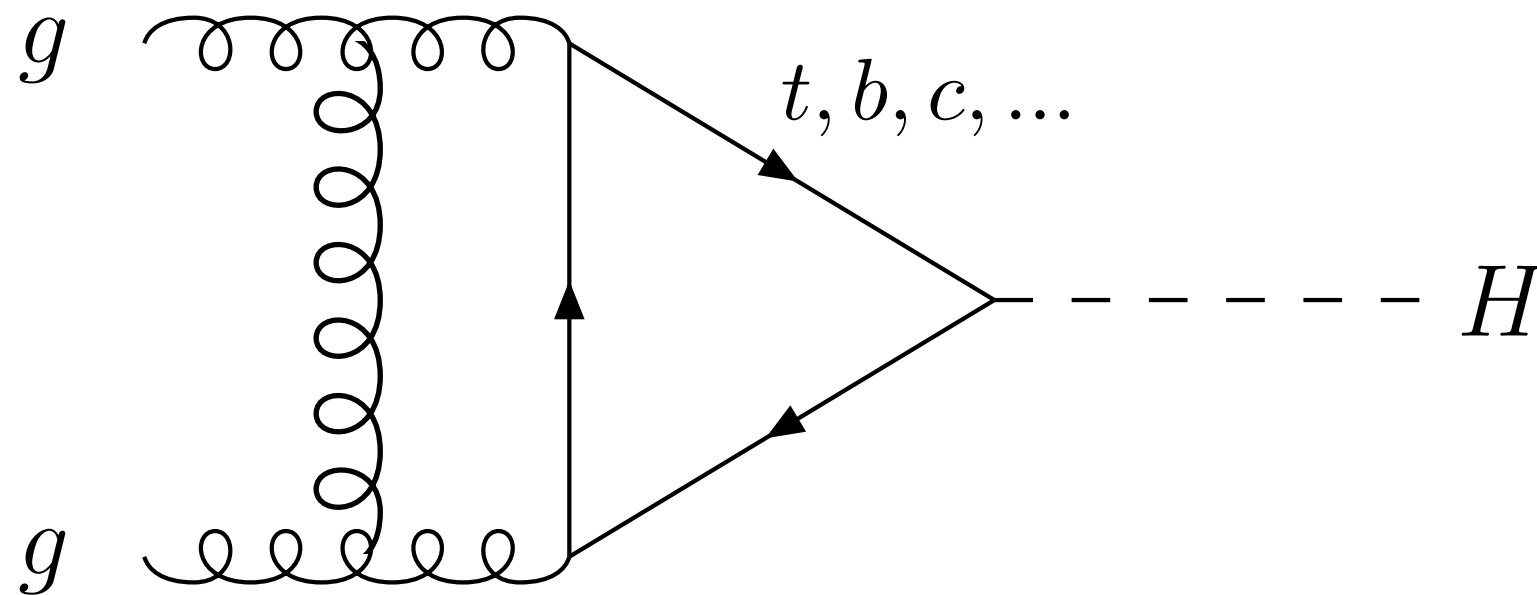
- Perturbative expansion in QCD coupling $\alpha_s \sim 0.11$

$$C_0\alpha_s^0 + C_1\alpha_s^1 + C_2\alpha_s^2 + C_3\alpha_s^3 + \dots$$



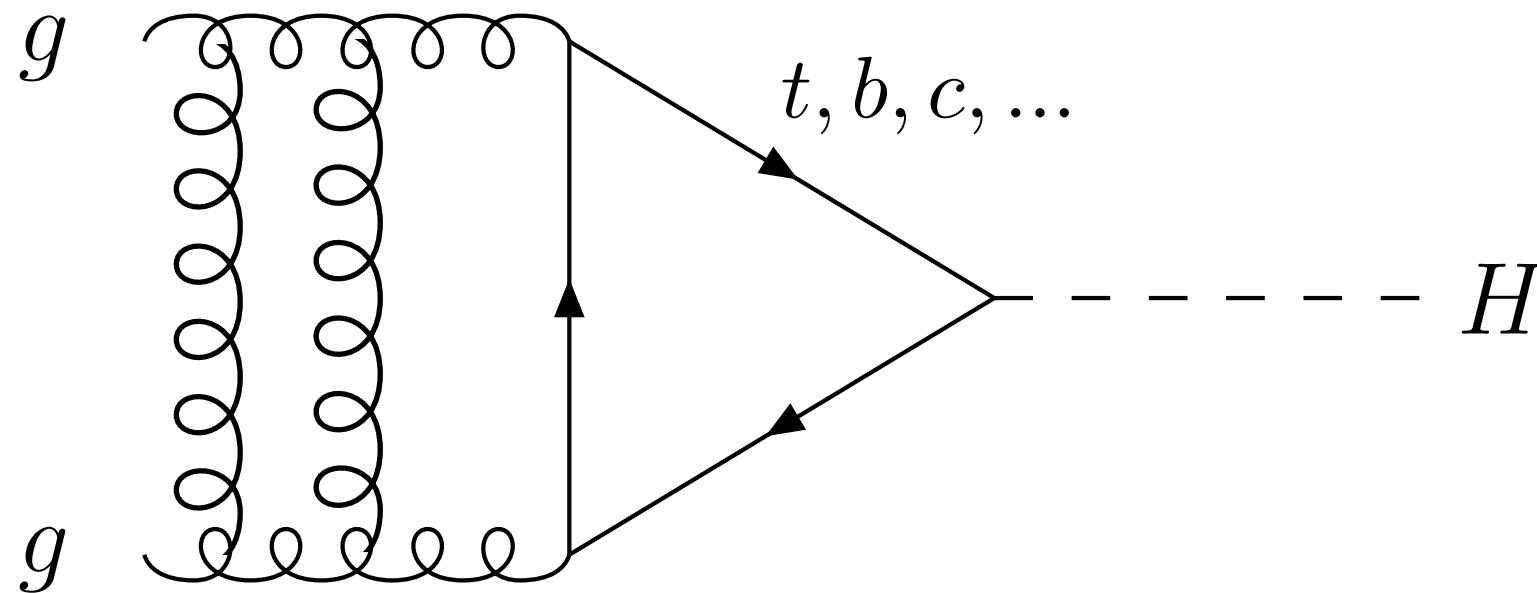
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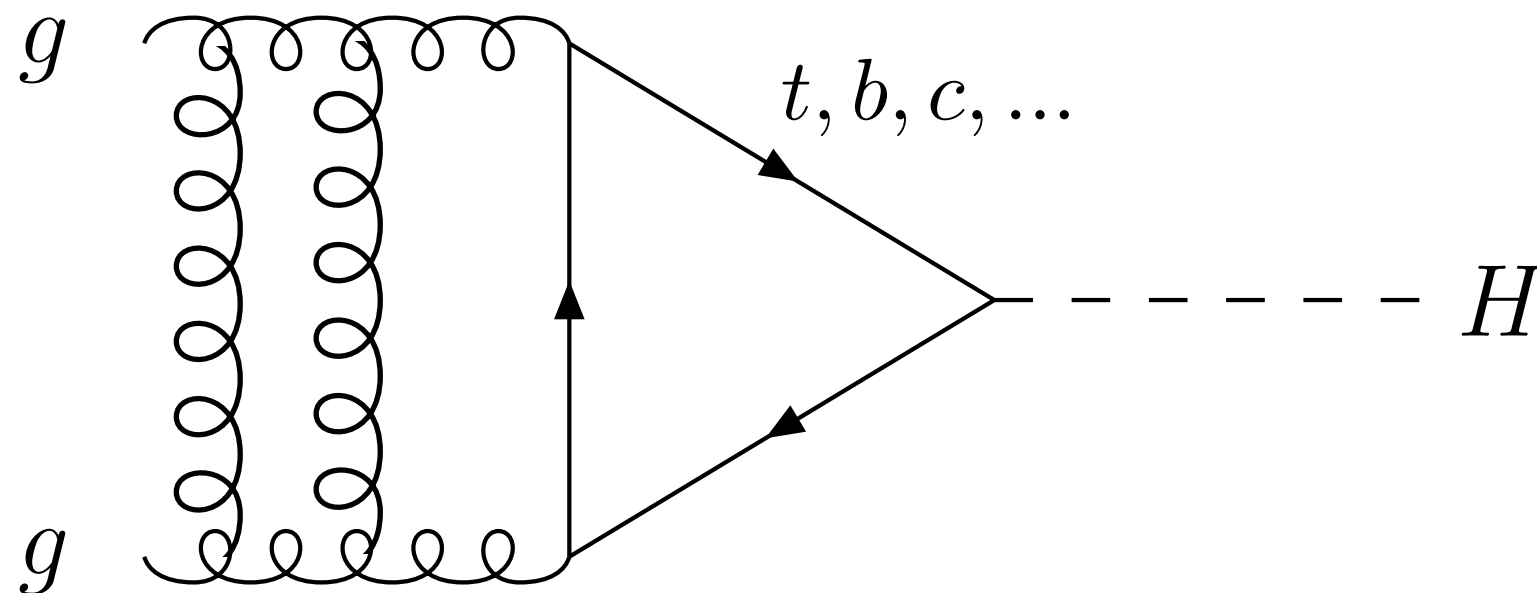
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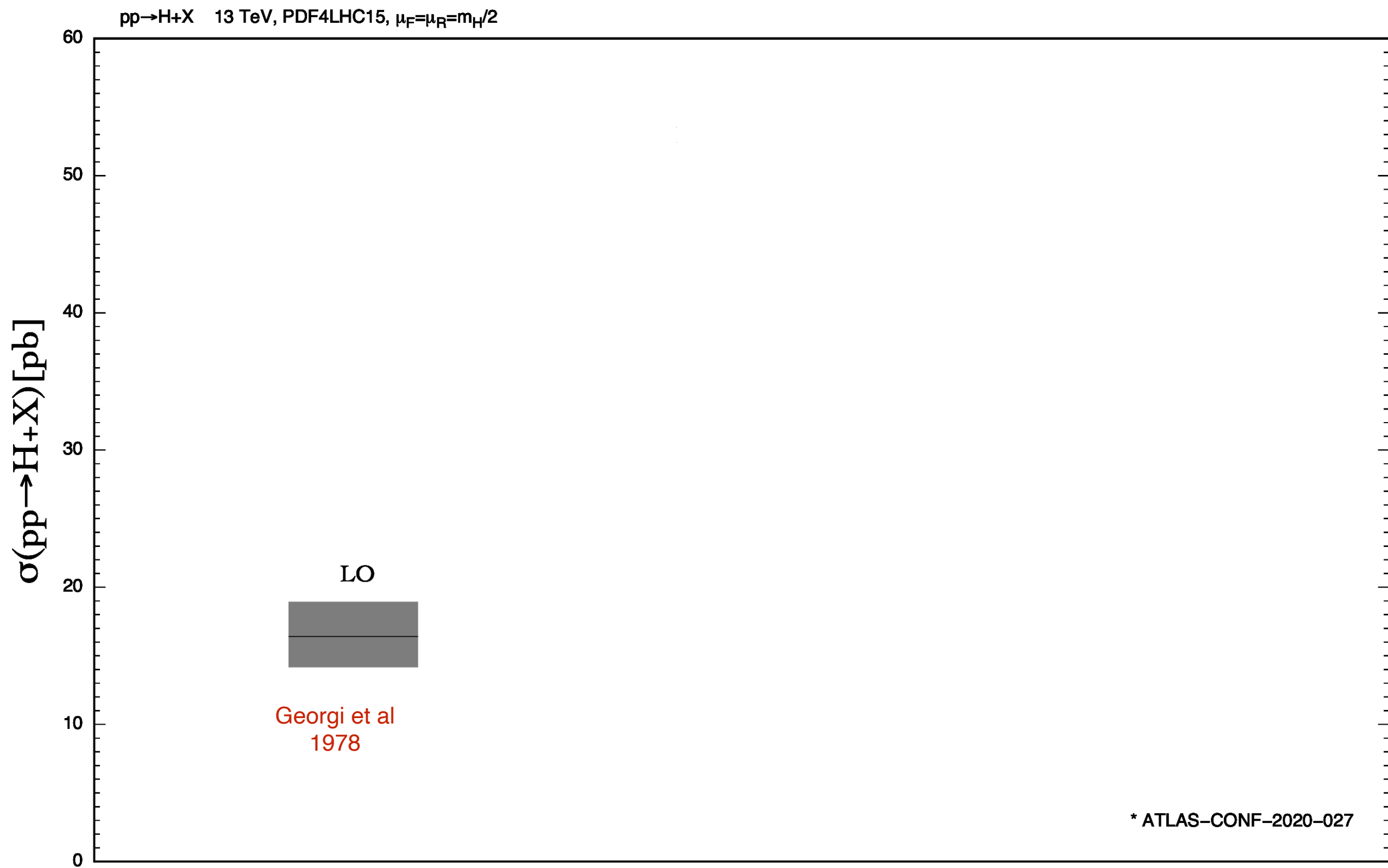
$$C_0 \alpha_s^0 + C_1 \alpha_s^1 + C_2 \alpha_s^2 + C_3 \alpha_s^3 + \dots$$



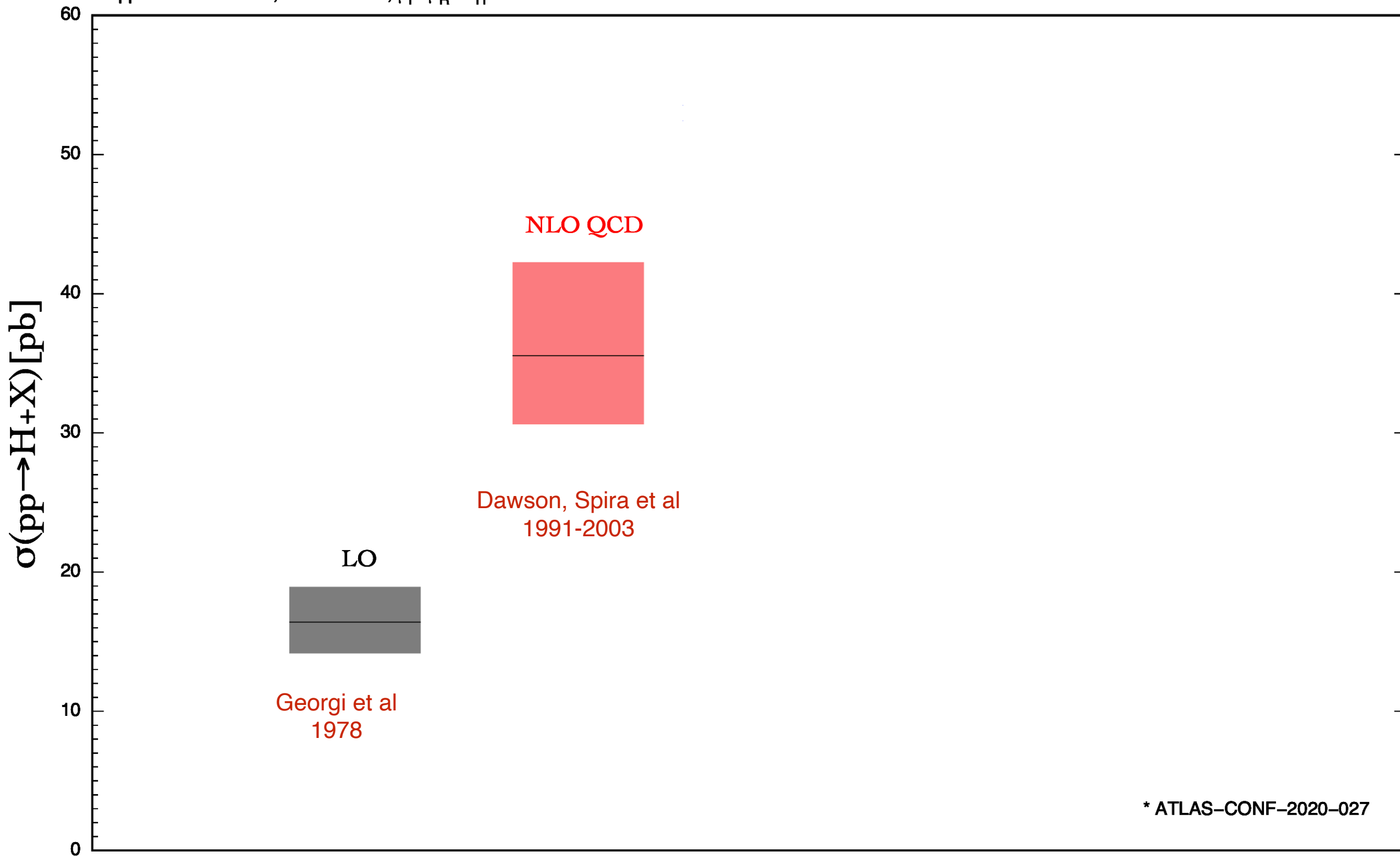
$$\sigma = \sigma^{(0)} (1 + 0.89 + 0.55 + 0.3 + \dots)$$

$$\alpha_s^0 + \alpha_s^1 + \alpha_s^2 + \alpha_s^3$$

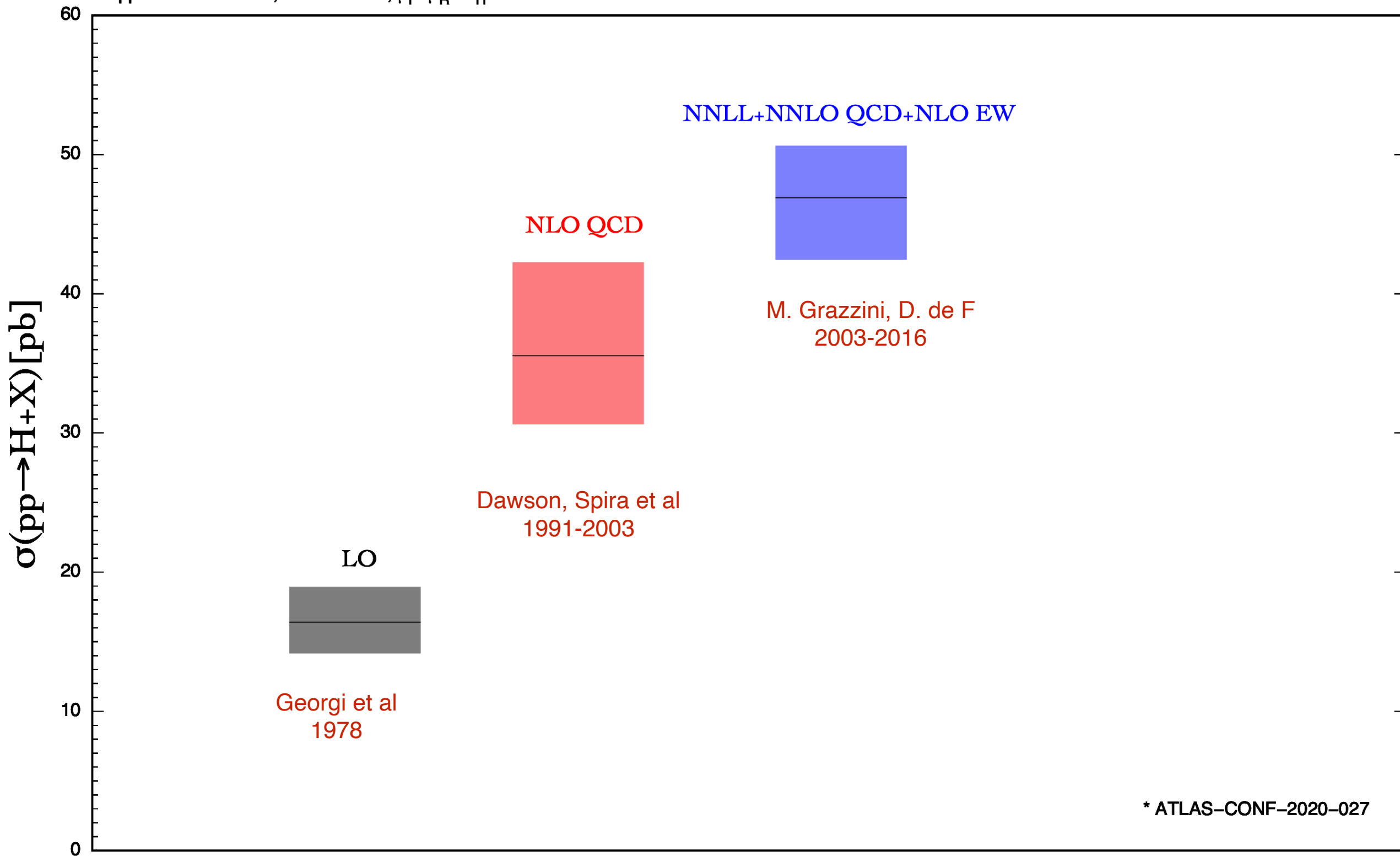
Very slow convergence (and millions of integrals needed...)



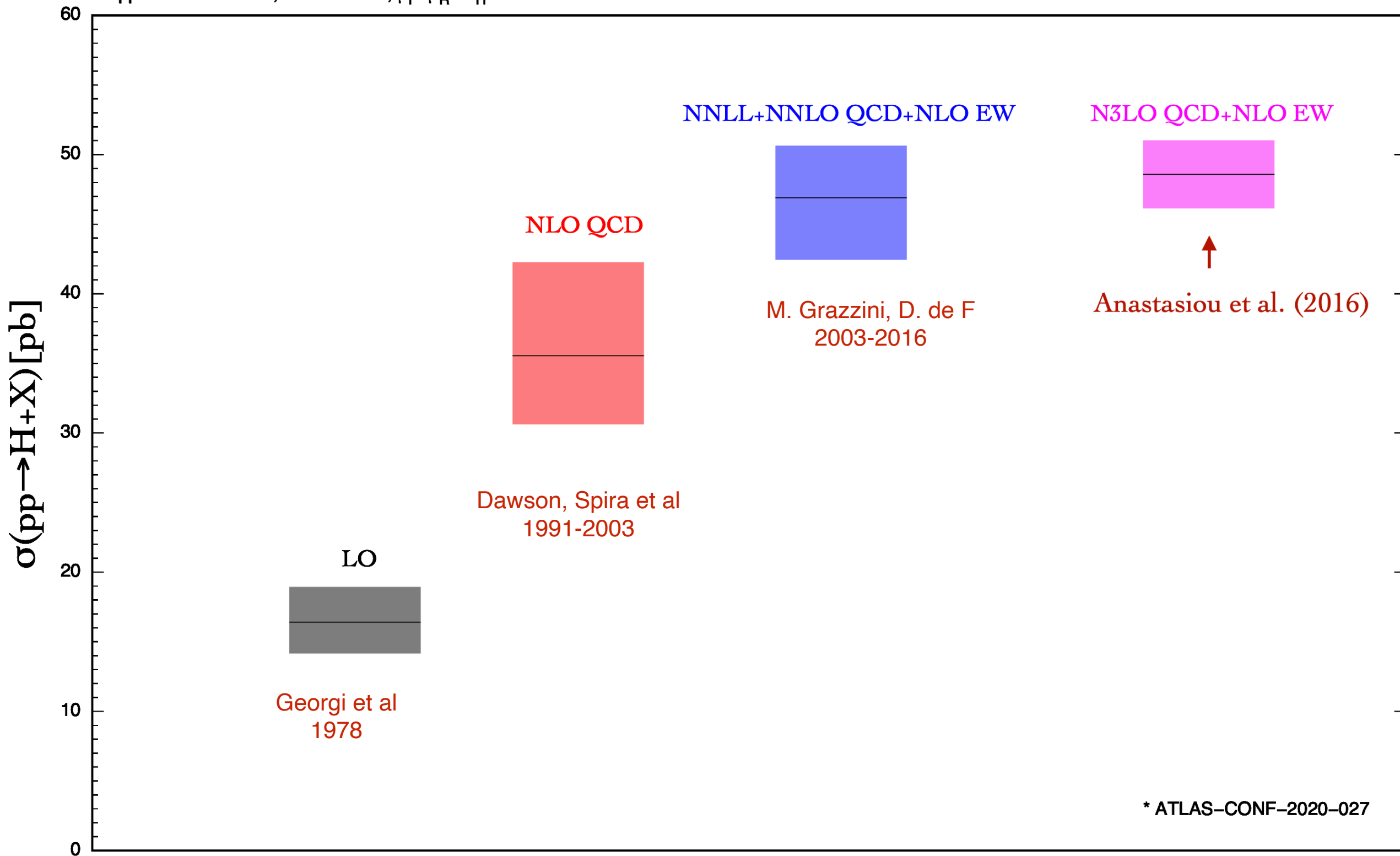
pp→H+X 13 TeV, PDF4LHC15, $\mu_F=\mu_R=m_H/2$



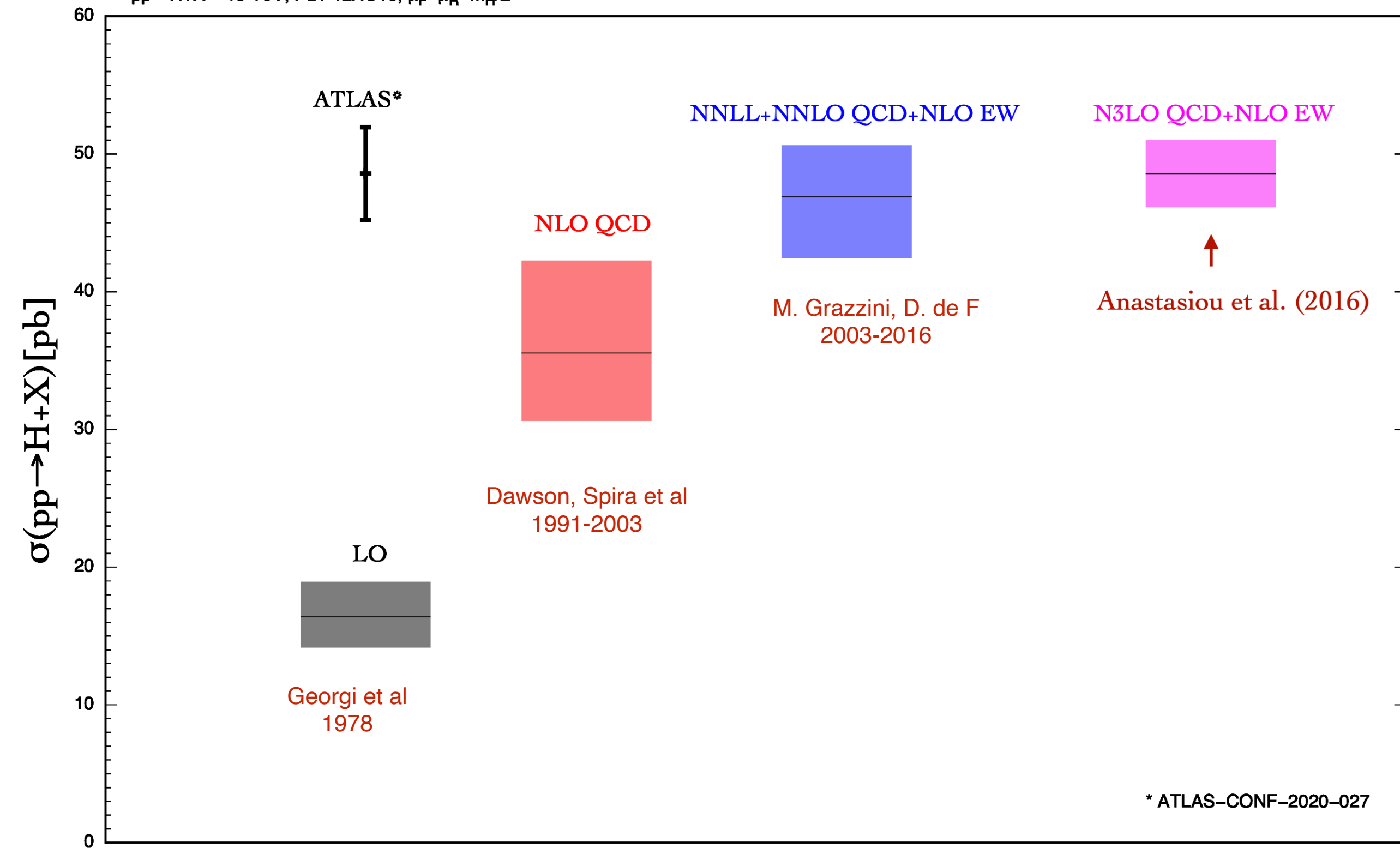
pp→H+X 13 TeV, PDF4LHC15, $\mu_F=\mu_R=m_H/2$



pp→H+X 13 TeV, PDF4LHC15, $\mu_F=\mu_R=m_H/2$



pp→H+X 13 TeV, PDF4LHC15, $\mu_F=\mu_R=m_H/2$



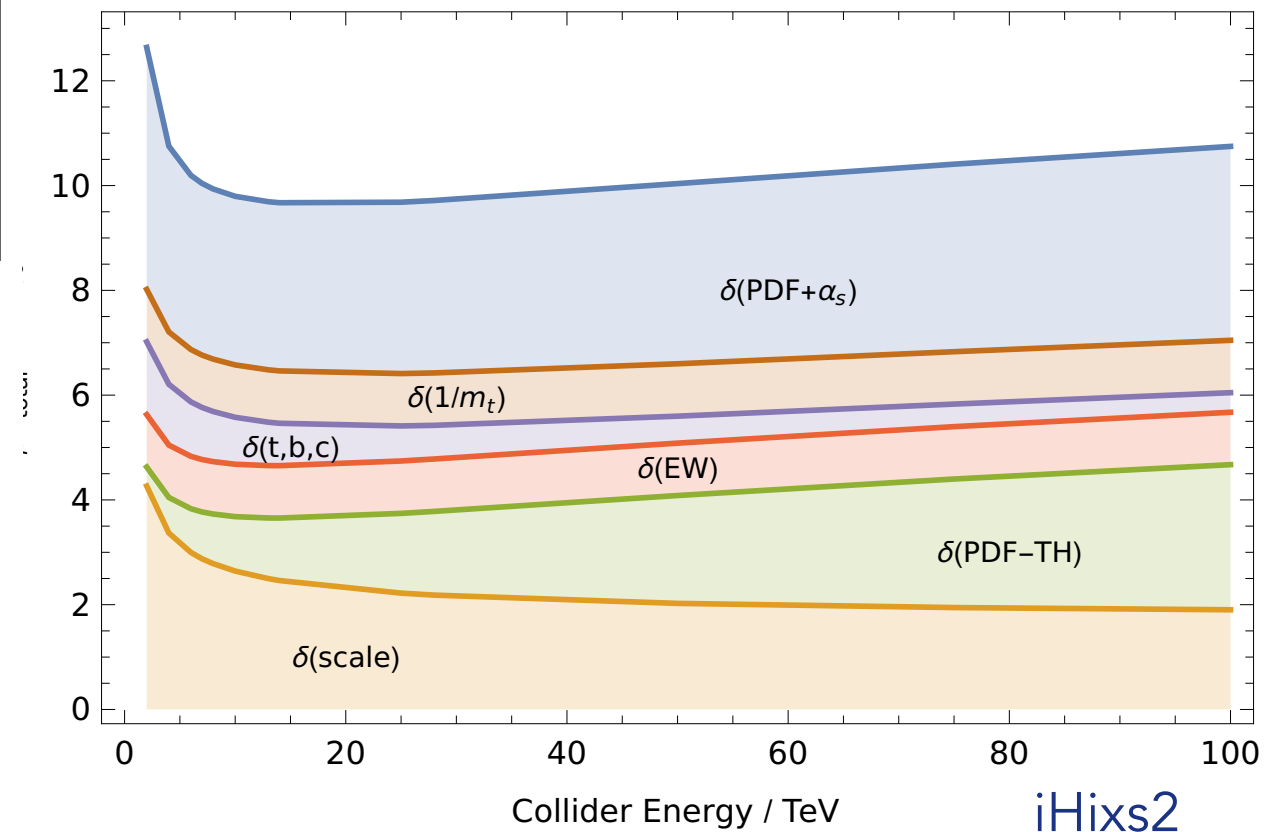
Improved Higgs Cross-section @ LHC

✓ at 13 TeV Higgs Cross-Section WG

$\sigma_{PP \rightarrow H+X}$	=	16.00 pb	(+32.87%)	LO, rEFT
	+	20.84 pb	(+42.82%)	NLO, rEFT
	+	9.56 pb	(+19.64%)	NNLO, rEFT
	+	1.62 pb	(+3.32%)	N ³ LO, rEFT
	−	2.07 pb	(−4.25%)	(t,b,c) corr. to exact NLO
	+	0.34 pb	(+0.70%)	1/ m_t corr. to NNLO
	+	2.37 pb	(+4.87%)	EWK corr.
	=	48.67 pb		

$\delta(\text{theory})$	=	+0.13pb	(+0.28%)	$\delta(\text{scale})$
		−1.20pb	(−2.50%)	
	+	$\pm 0.56pb$	($\pm 1.16\%$)	$\delta(\text{PDF-TH})$
	+	$\pm 0.49pb$	($\pm 1.00\%$)	$\delta(\text{EWK})$
	+	$\pm 0.41pb$	($\pm 0.85\%$)	$\delta(\text{t,b,c})$
	+	$\pm 0.49pb$	($\pm 1.00\%$)	$\delta(1/m_t)$
	=	+2.08pb	(+4.28%)	
		−3.16pb	(−6.5%)	
$\delta(\text{PDF})$	=	$\pm 0.89pb$	($\pm 1.85\%$)	
$\delta(\alpha_S)$	=	+1.25pb	(+2.59%)	
		−1.26pb	(−2.62%)	

● uncertainties



► Great improvement over the last years

► Without QCD corrections : fail by more than a factor of 2

► Need to attack in many fronts to further improve: pdf, top mass, EW

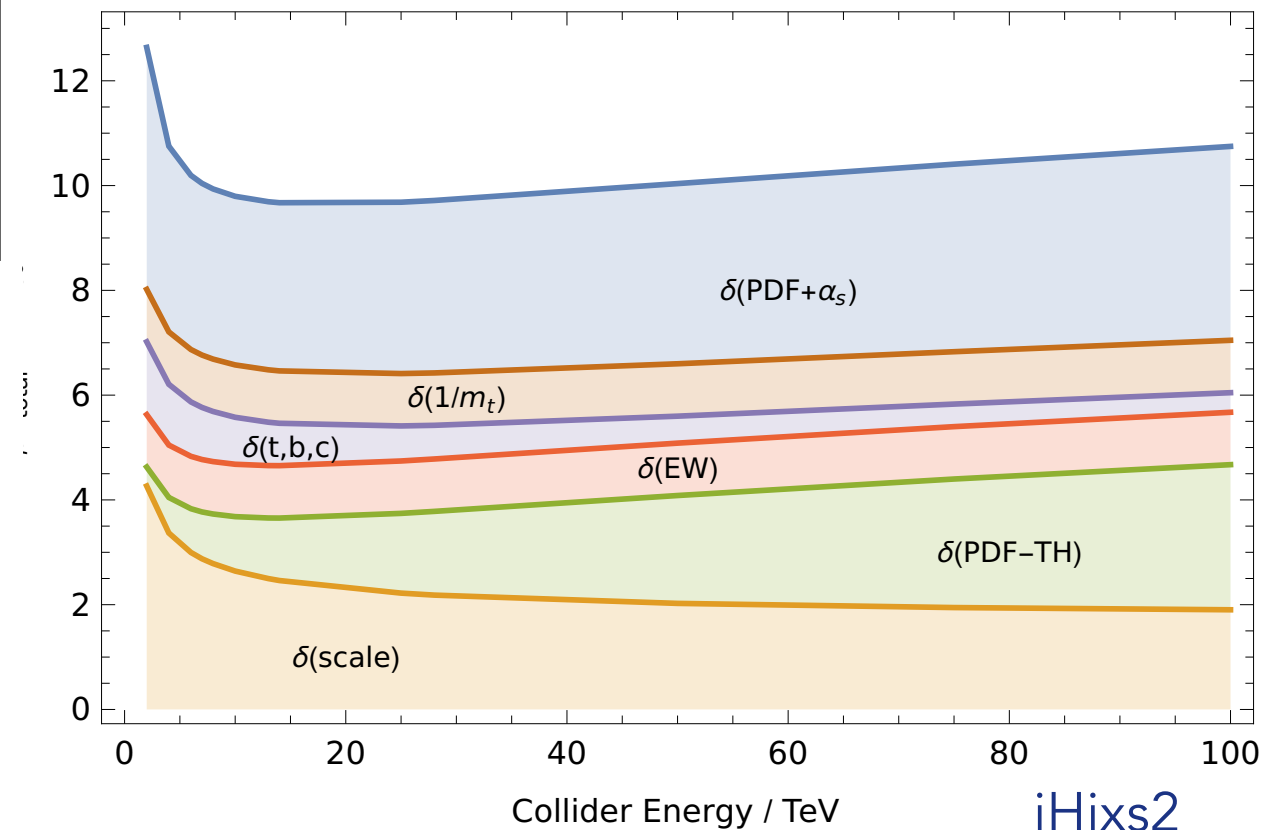
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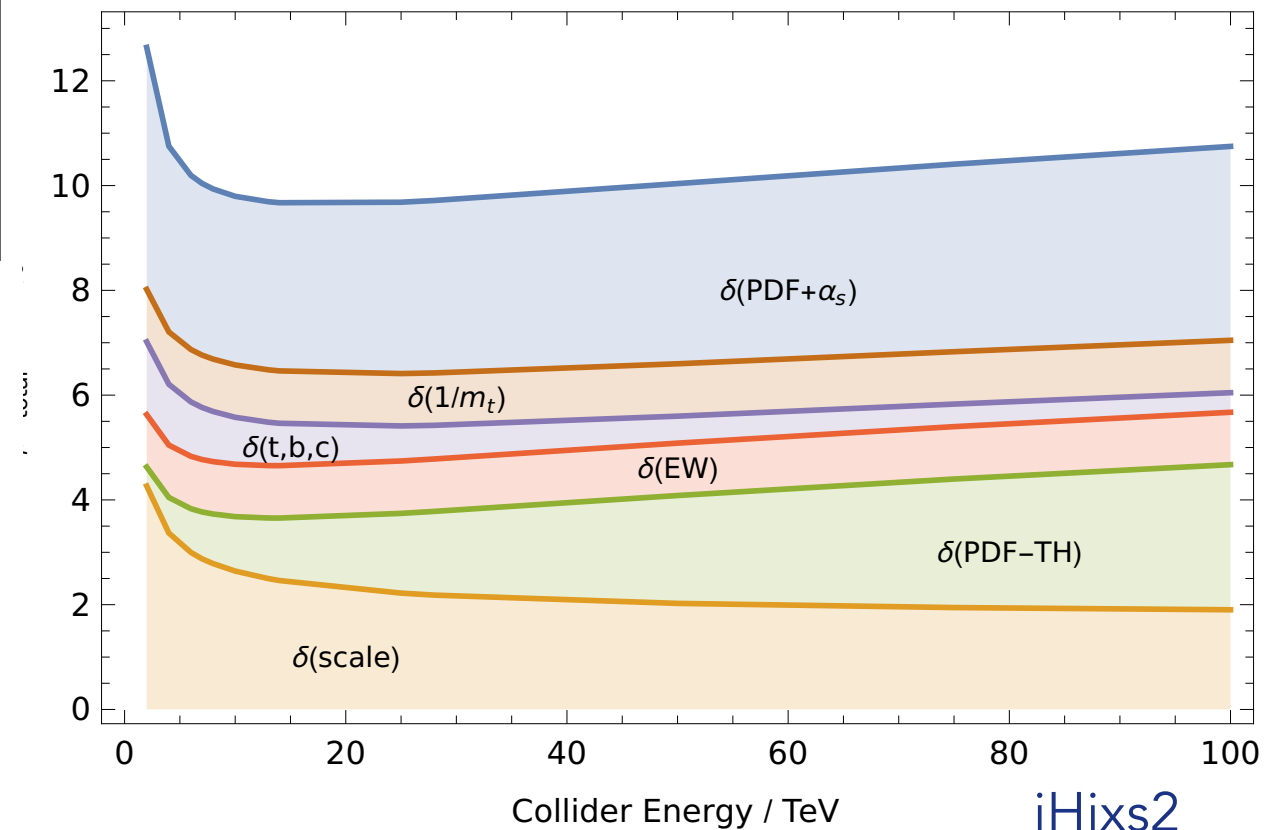
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► Without QCD corrections : fail by more than a factor of 2

► Need to attack in many fronts to further improve: pdf, top mass, EW

▶ TH : precision calculations

- ▶ QCD at high orders NNLO or more
- ▶ EW corrections
- ▶ Resummation
- ▶ MonteCarlos matched to fixed order calculations

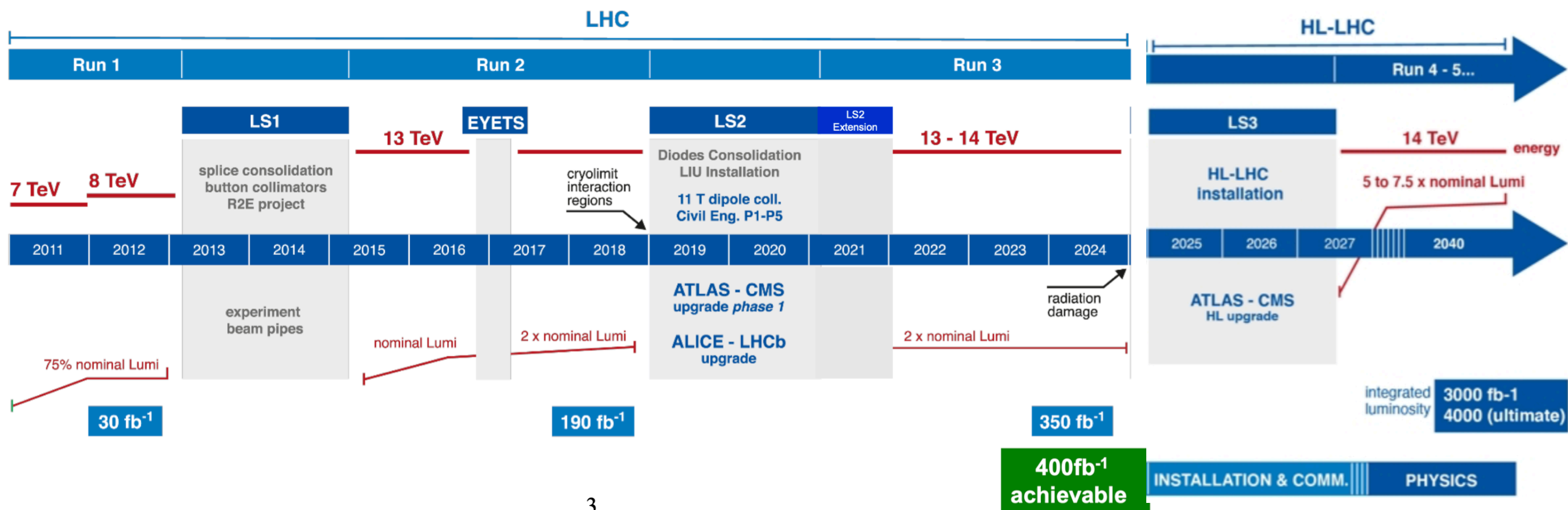
Future 2

Next years (>20x number of events) Discovery and **precision**

Run 3 outlook

Slide from F. Bordry

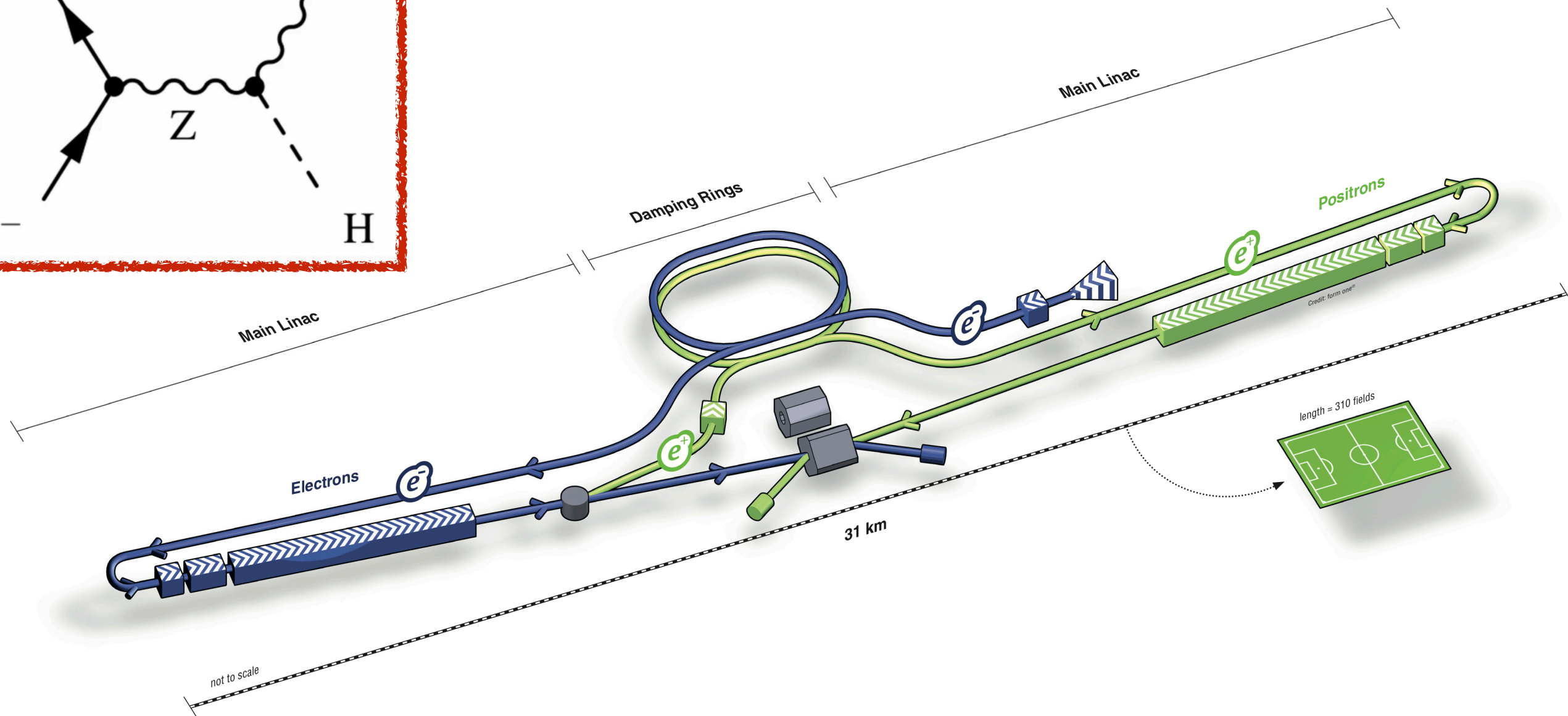
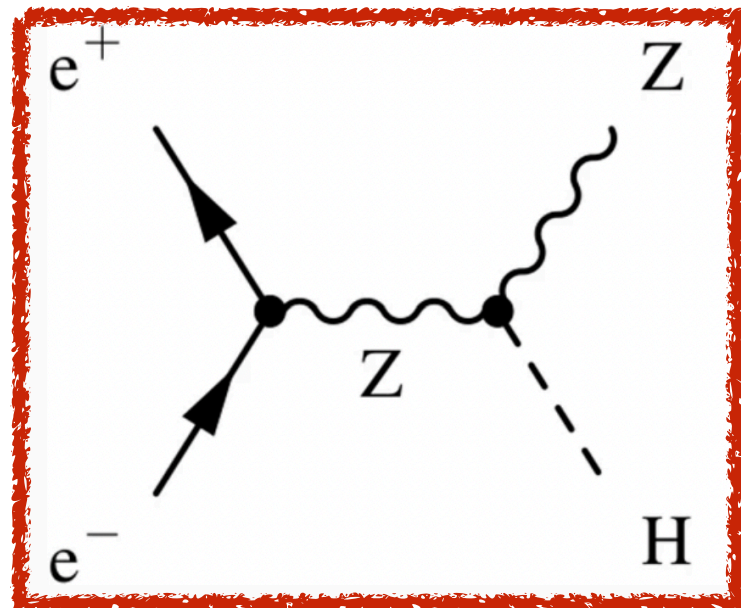
Goal: $\Sigma(\text{Run1} + \text{Run2} + \text{Run 3}) > 350 \text{ fb}^{-1}$



► Approved >2040?

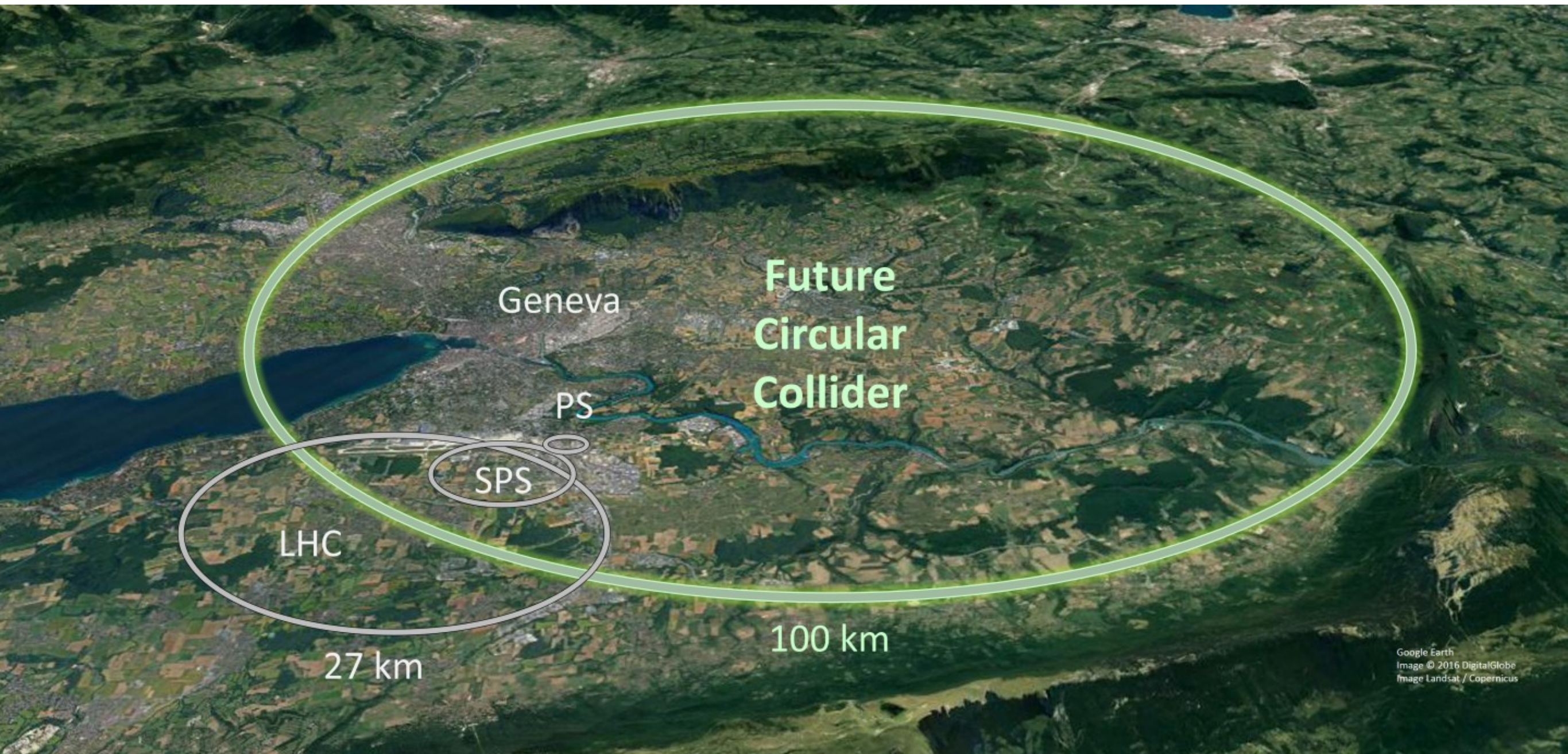
International Linear Collider (e^+e^-)

~250 GeV : Higgs Factory (precision)

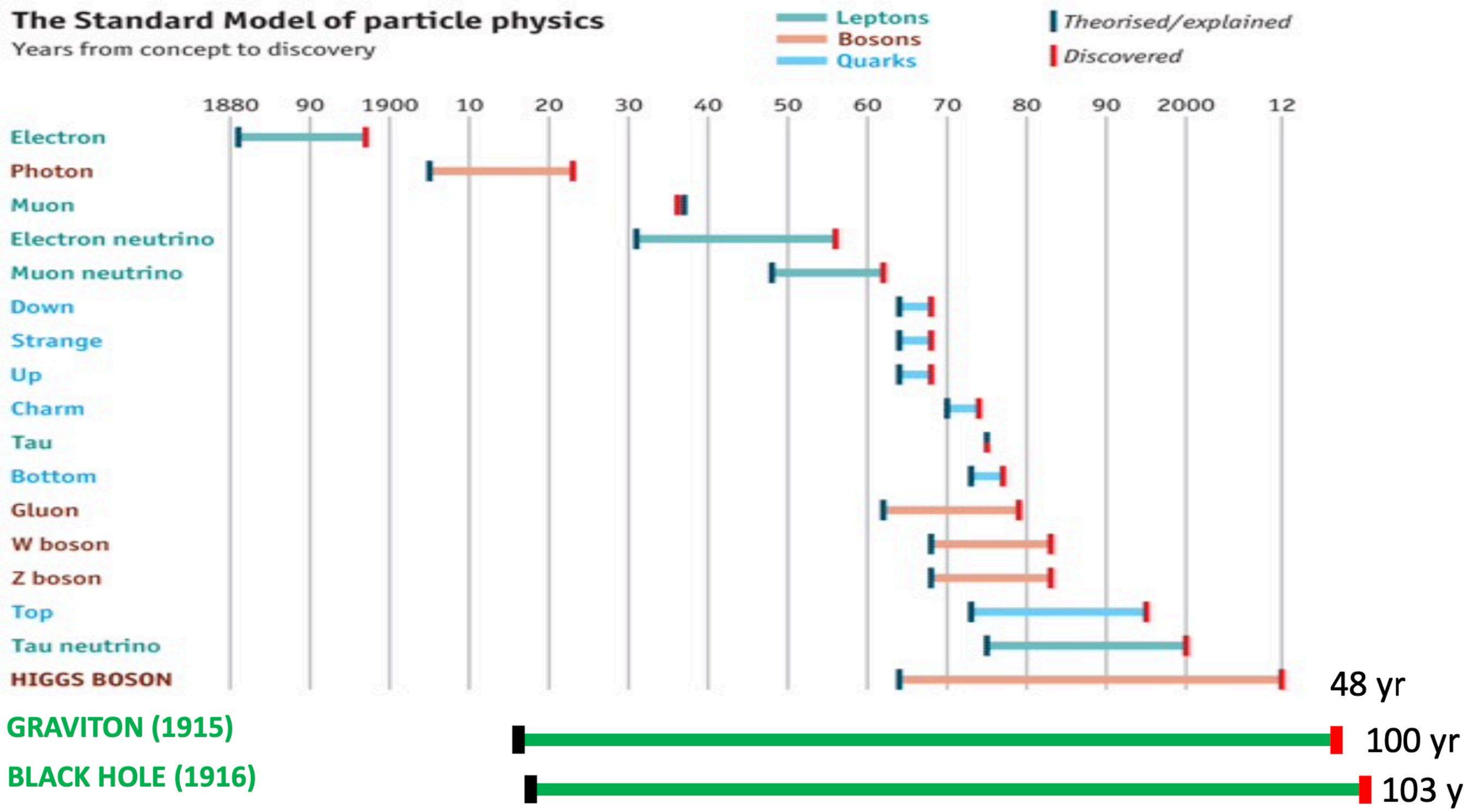


- Technology available, ready to build (lack of funds and site...)

FCC pp collider (y e^+e^-) ~ 100 TeV (descubrimiento)

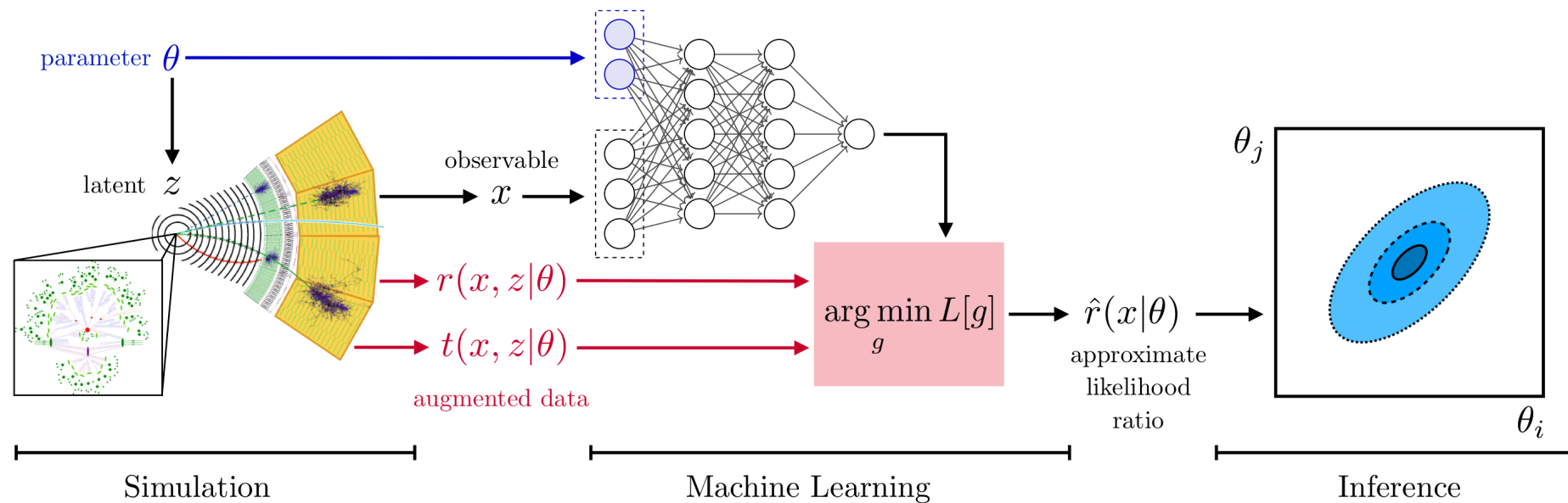


► For a long time search was guided by Theoretical expectations



Now is experimentally-driven

- new methods for discovering new physics based on data rather than theory: machine learning



- Search for variables that maximize signal vs background
- Search for "anomalies" and new physics without theoretical bias

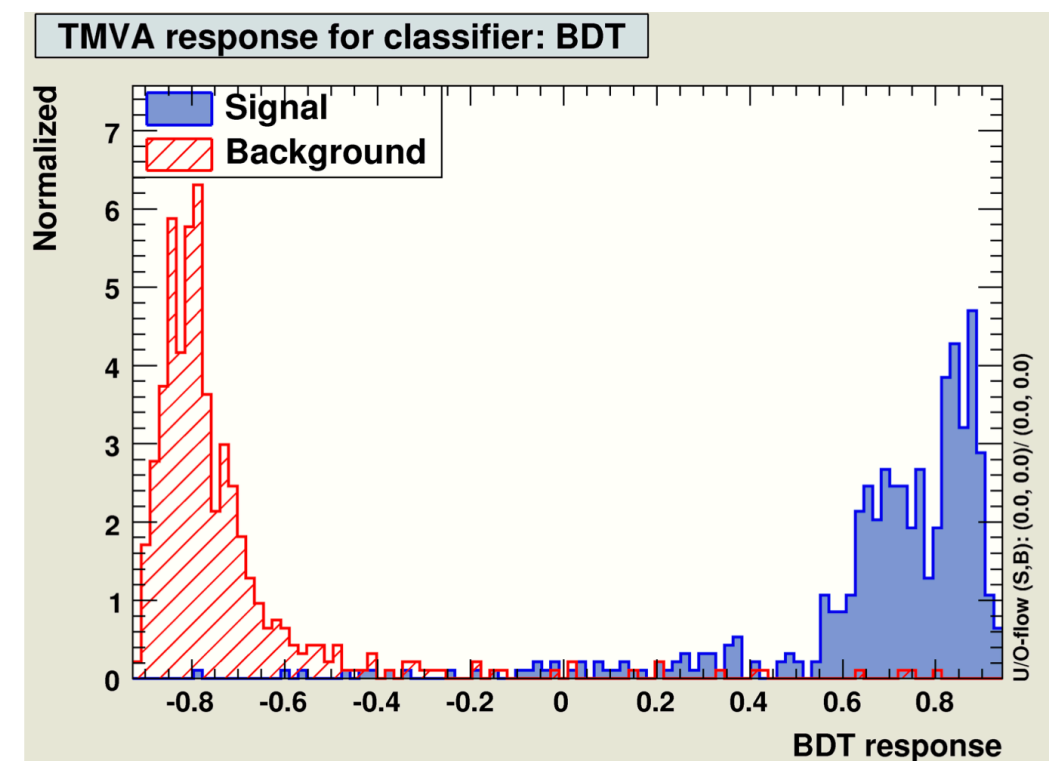


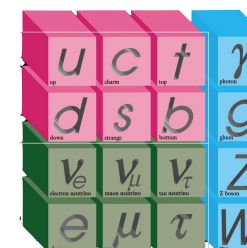
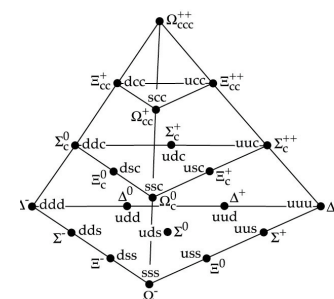
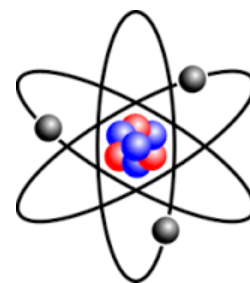


Tabla Periódica de los Elementos

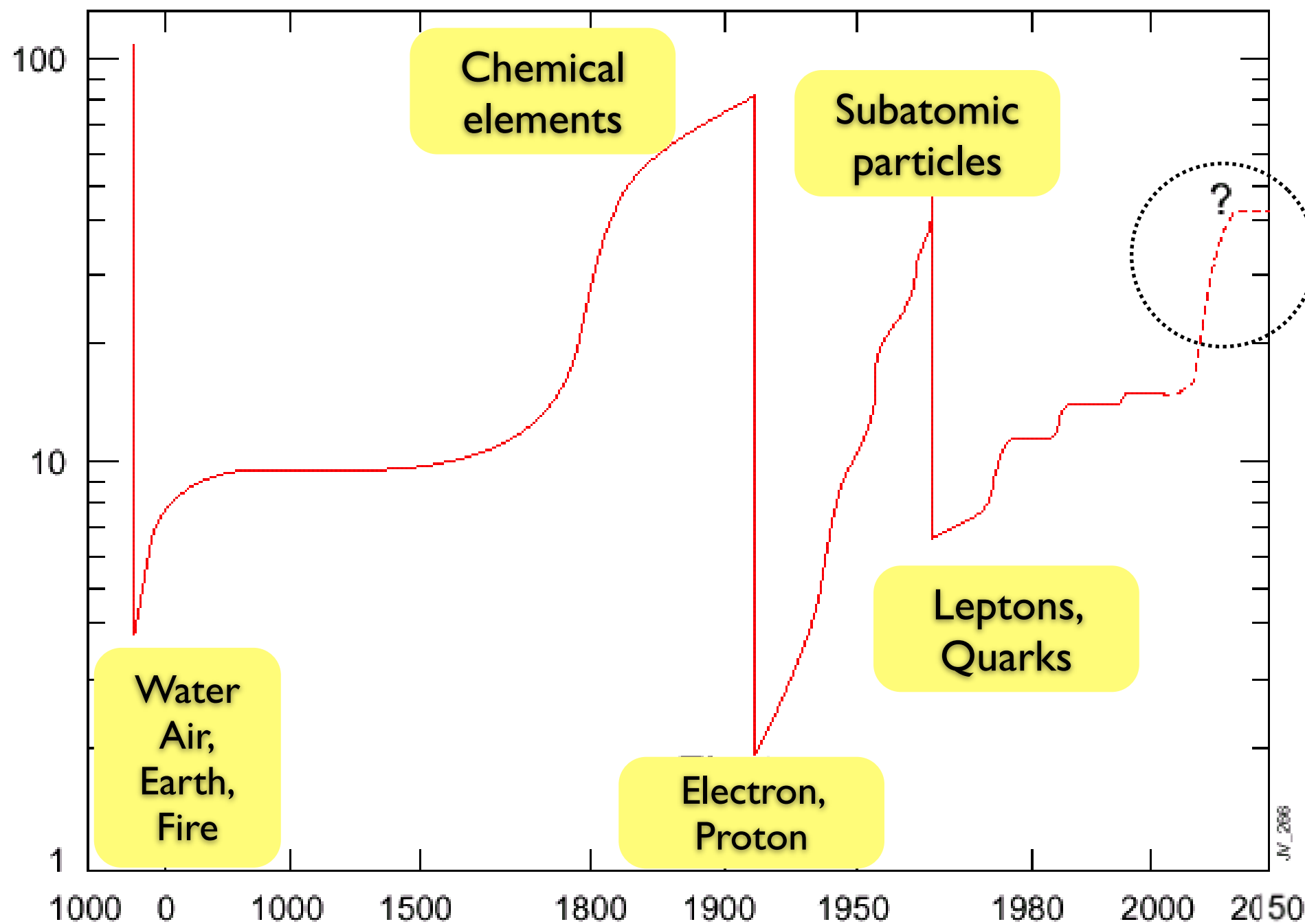
This image shows a periodic table of elements, color-coded by groups. The legend on the left identifies the following categories:

- Metales** (Metals): Yellow background.
- Alcalinos** (Alkali metals): Light blue background.
- Alcalinotérreos** (Alkaline earth metals): Light green background.
- Metales de transición** (Transition metals): Pink background.
- Aluminos** (Aluminum): Light orange background.
- Calcógenos** (Chalcogens): Light purple background.
- Halogenuros** (Halogens): Light red background.
- Gases nobles** (Noble gases): Light yellow background.

The table includes atomic numbers, element symbols, and names. The noble gases (Group 18) are highlighted in green. The lanthanide and actinide series are shown at the bottom, separated from the main body of the table.



Fundamental Particles



LHC
+
...

¡Gracias!

Physics

History

Aristotle said a bunch of stuff that was wrong. Galileo and Newton fixed things up. Then Einstein broke everything again. Now, we've basically got it all worked out, except for small stuff, big stuff, hot stuff, cold stuff, fast stuff, heavy stuff, dark stuff, turbulence and the concept of time.

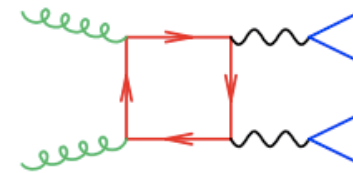
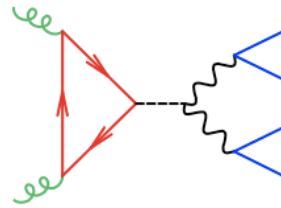
Summary

- ▶ A Higgs boson was found at the LHC
- ▶ Consistent with SM Higgs (large uncertainties yet)
- ▶ Great TH work to match EXP precision
 - High order corrections
- ▶ Can still expect surprises at the LHC
- ▶ Need to analyze Higgs sector with detail
- ▶ **PRECISION** for the next decades
- ▶ Studies underway for new colliders

Off-shell effects and interference

$$\mathcal{A}_{ij \rightarrow X} = \mathcal{A}_{ij \rightarrow H} \Delta_H \mathcal{A}_{H \rightarrow X} + \mathcal{A}_{\text{continuum}}$$

Propagator

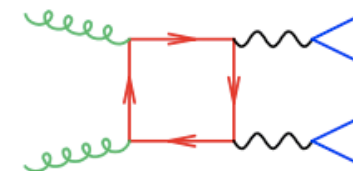
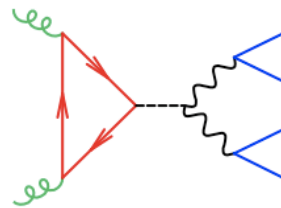


$$\Delta_H^2(q^2) \sim \frac{1}{(q^2 - M_H^2)^2 + \Gamma_H^2 M_H^2} \sim \frac{\pi}{M_H \Gamma_H} \delta(q^2 - M_H^2) + \mathcal{O}\left(\frac{\Gamma_H}{M_H}\right) \text{ ZWA}$$

Off-shell effects and interference

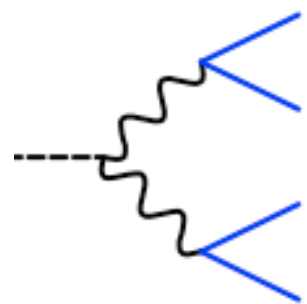
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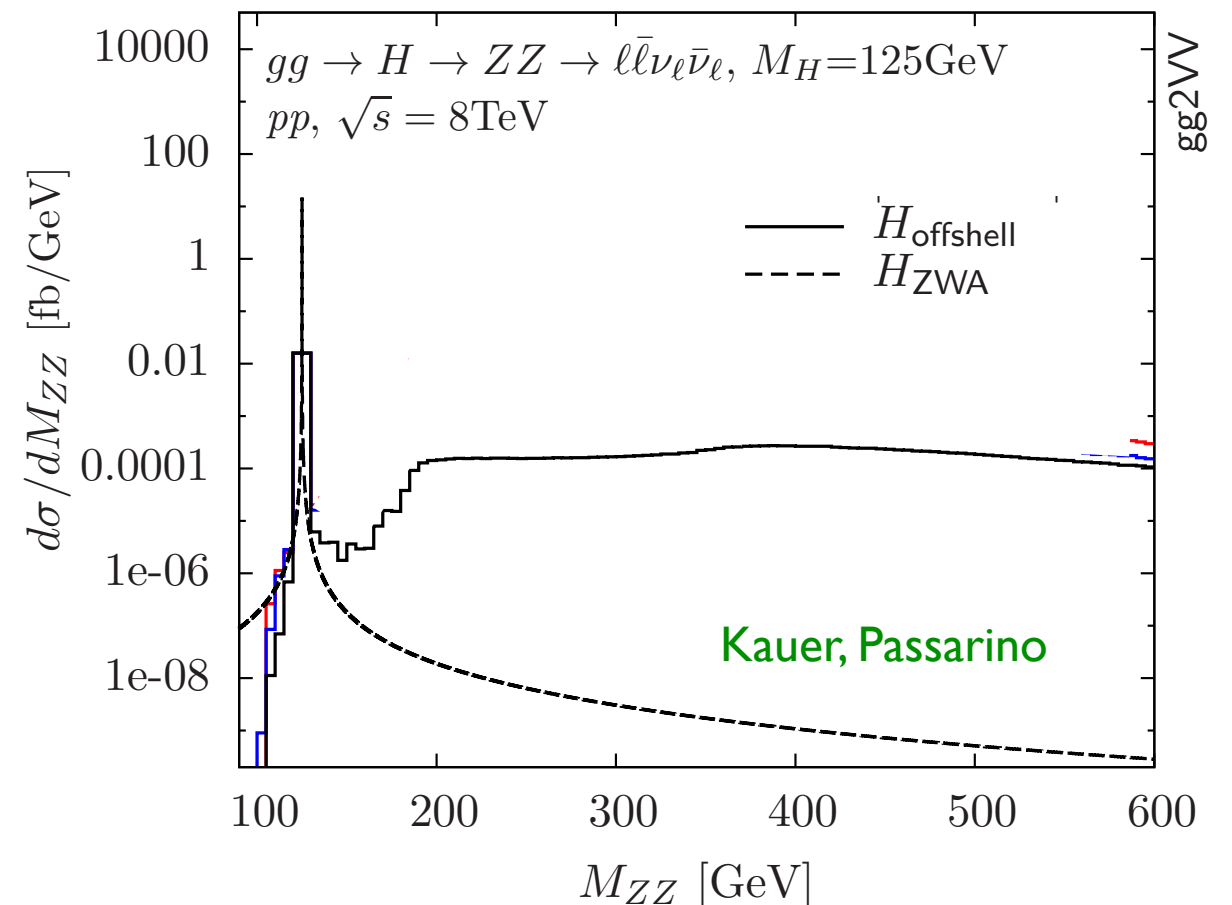
$$\Delta_H^2(q^2) \sim \frac{1}{(q^2 - M_H^2)^2 + \Gamma_H^2 M_H^2} \sim \frac{\pi}{M_H \Gamma_H} \delta(q^2 - M_H^2) + \mathcal{O}\left(\frac{\Gamma_H}{M_H}\right) \text{ ZWA}$$

But above threshold decay amplitude compensates $1/(q^2)^2$



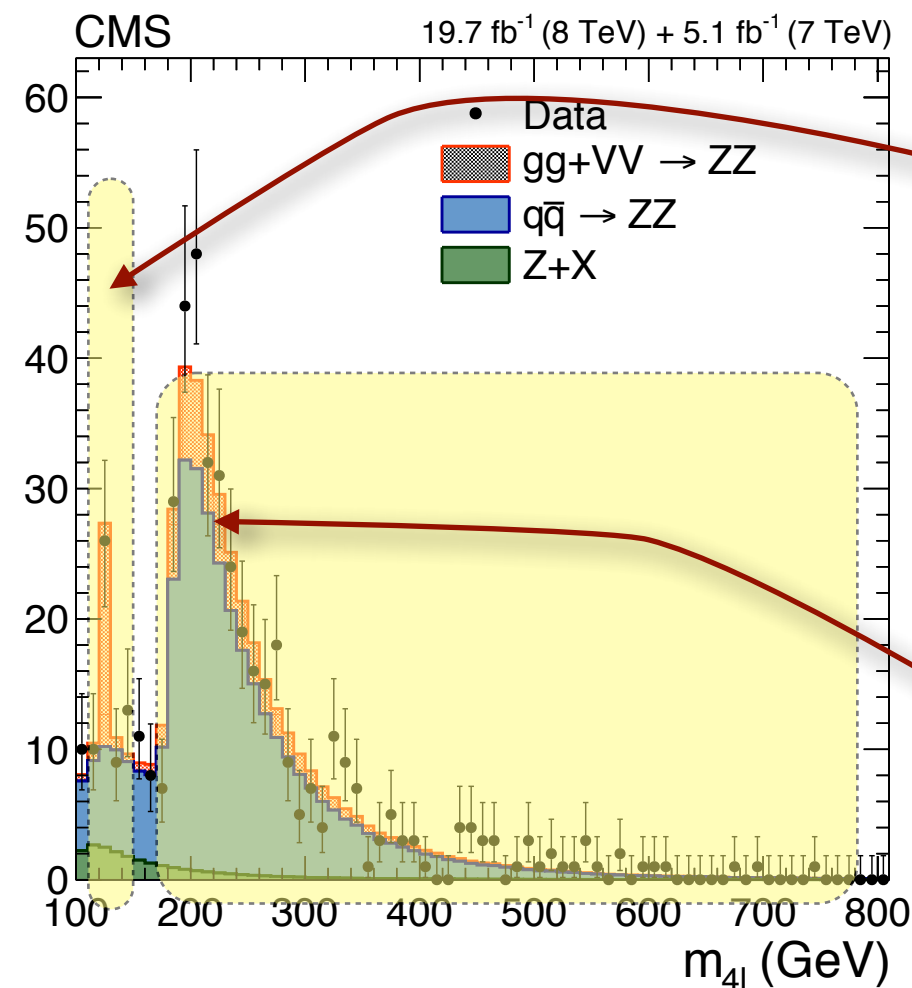
$$|\mathcal{A}_{H \rightarrow VV}|^2 \sim (q^2)^2$$

- Sizeable contribution from off-shell
- Enhances effect of interference



Width measurement from off-shell

$$gg \rightarrow H \rightarrow VV$$



$$\sigma^{\text{on}} \int_{M_H^2 - \Delta^2}^{M_H^2 + \Delta^2} dq^2 \frac{|A_{gg \rightarrow H \rightarrow VV}|^2}{(q^2 - M_H^2) + \Gamma_H^2 M_H^2} \sim \frac{g_{ggH}^2(M_H^2) g_{HVV}^2(M_H^2)}{\Gamma_H}$$

SM assumptions on couplings (running)

$$g = \xi g^{SM}$$

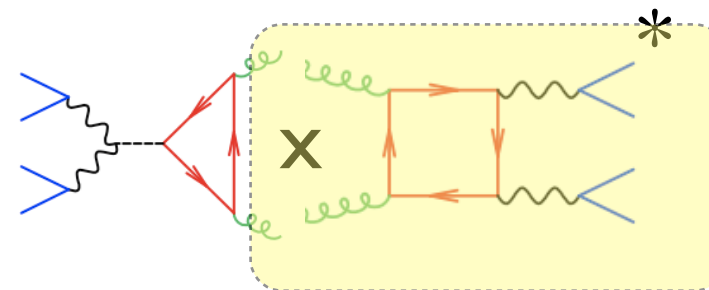
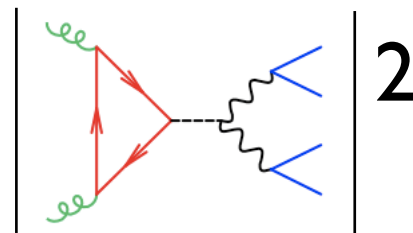
$$\Gamma_H = \xi^4 \Gamma_H^{SM}$$

$$\sigma^{\text{off}} \int_{q^2 \gg M_H^2} dq^2 \frac{|A_{gg \rightarrow H \rightarrow VV}|^2}{(q^2 - M_H^2) + \Gamma_H^2 M_H^2} \sim \int dq^2 g_{ggH}^2(q^2) g_{HVV}^2(q^2)$$

$$\sigma^{\text{exp}} = \sigma^{\text{back}} + \sigma^{\text{on}} + \sigma^{\text{off}} \times \frac{\Gamma_H}{\Gamma_H^{SM}} + \sigma^{\text{int}} \times \sqrt{\frac{\Gamma_H}{\Gamma_H^{SM}}}$$

CMS

$$3.2_{-2.2}^{+2.8} \text{ MeV}$$



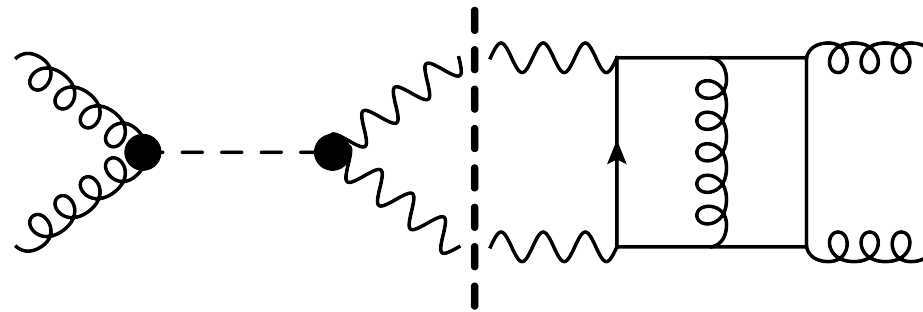
need precision on $gg \rightarrow VV$

NNLO for background but sizeable

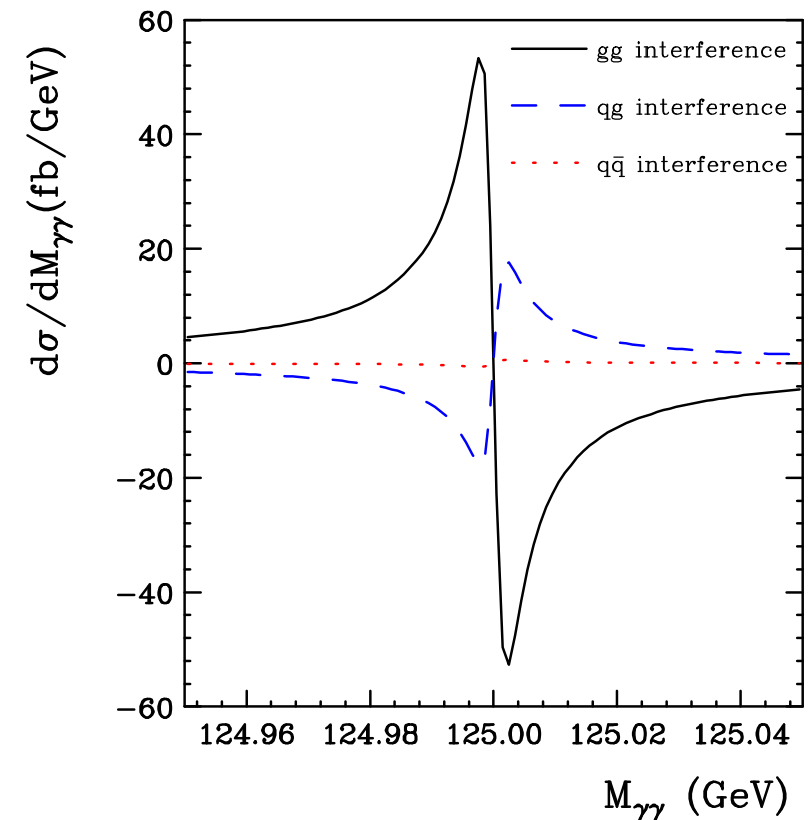
Width from interference $H \rightarrow \gamma\gamma$

- ▶ In diphoton channel, interference small for total cross section but asymmetry produces shift in invariant mass : enhanced by detector resolution

Dicus, Willenbrock (1986)
 Dixon, Siu (2003)
 Martin (2012,2013)
 deF et al (2013)
 Dixon, Li (2013)



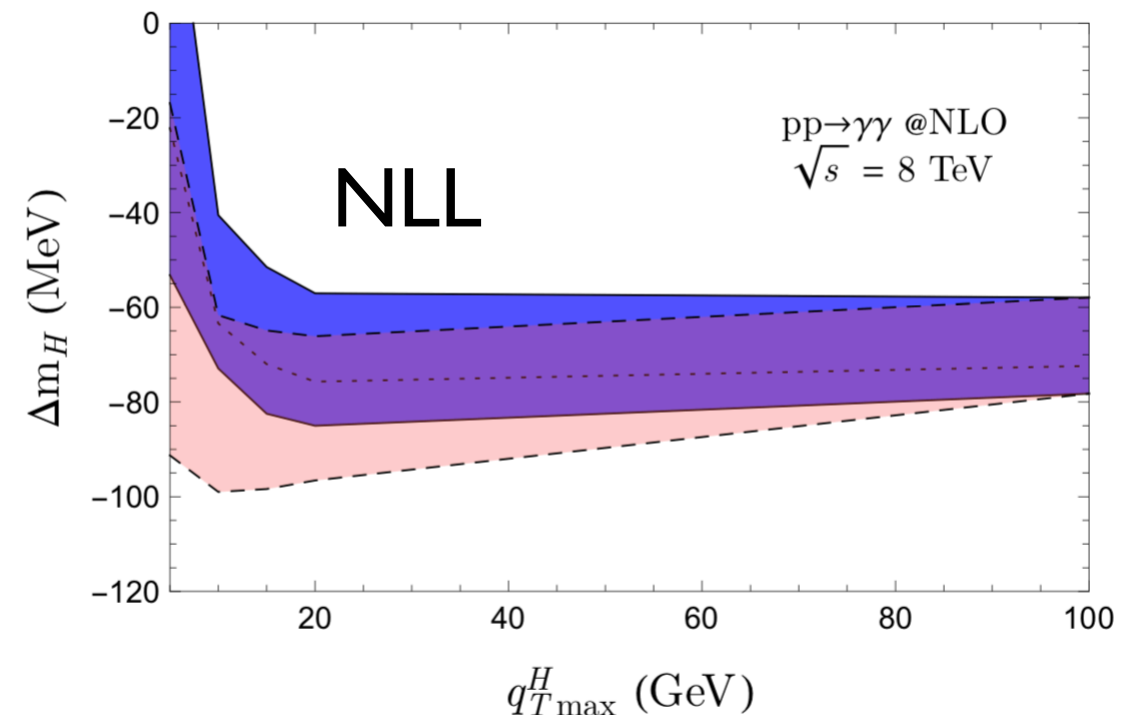
Known to $\mathcal{O}(\alpha_s^3)$



- ▶ Look at $\Delta M_H = M_H^{\gamma\gamma} - M_H^{ZZ}$

or precise mass measurements
 in different H q_T bins in $\gamma\gamma$

Cieri, Coradeschi, deF, Fidanza (2017)



- ▶ Search for -2% effect of interference in cross section
- ▶ All effects might be enhanced by BSM

Campbell, Carena,
 Harnik, Liu (2017)

