## Precision Monte Carlos <br> for multi-boson processes

Jonas M. Lindert

Science \& Technology Facilities Council

UK Research and Innovation

Workshop on.... measurements and reinterpretations... at the LHC
U Sussex
15th June 2022

## Multibosons

## VV

VVV

VBS-VV

## Multibosons



VBF-V $\xrightarrow{\text { background }}$ VBS-VV

## Multibosons


forward-jet dynamics

## Multibosons


forward-jet dynamics

## Multibosons



## Multibosons: theory state-of-the art fixed-order

$V+$ jets<br>NLO QCD+NLO EW<br>NLO QCD+NLO EW

VV NNLO QCD+NLO EW VBF-V NLO QCD*+LO EW

VBS-VV<br>NLO QCD+NLO EW / NLO QCD*+LO EW

## Multibosons: theory state-of-the art fixed-order

V+jets<br>NLO QCD+NLO EW

VV<br>NNLO QCD+NLO EW<br>VBF-V<br>NLO QCD+NLO EW<br>[JML, Pozzorini, Schönherr, 2204.07652]

VVV
NLO QCD+NLO EW

VBS-VV
NLO QCD+NLO EW / NLO QCD*+LO EW

*: VBF approximation

## Multibosons: theory state-of-the art Monte Carlo

| V+jets | VV | VVV |
| :---: | :---: | :---: |
| NLOPS QCD $(0,1,2 \mathrm{j}) \times$ | NLOPS QCD $(0,1) \times$ | NLOPS QCD |
| NLO EWvirt | NLO EWvirt |  |
|  | NNLOPS QCD |  |
|  | NLOPS EW |  |
| VBF-V | NBS-VV |  |
|  |  |  |
|  |  | NLOPS* QCD |

*: VBF approximation

## Multibosons: theory state-of-the art Monte Carlo

V+jets<br>NLOPS QCD ( $0,1,2 \mathrm{j}) \times$ NLO EWvirt

| VV | VVV |
| :--- | :--- |
| NLOPS QCD $(0,1) x$ |  |
| NLO EWvirt |  |$\quad$ NLOPS QCD

*: VBF approximation

## Perturbative expansion:VV,VVV

(single perturbative order at LO)

$$
\begin{aligned}
\mathrm{d} \sigma=\mathrm{d} \sigma_{\mathrm{LO}}+ & \alpha_{S} \mathrm{~d} \sigma_{\mathrm{NLO}}+\alpha_{\mathrm{EW}} \mathrm{~d} \sigma_{\mathrm{NLO} \text { EW }} \\
& \mathrm{NLO} \mathrm{QCD} \quad \mathrm{NLO} \mathrm{EW} \\
& +\alpha_{S}^{2} \mathrm{~d} \sigma_{\mathrm{NNLO}}+\alpha_{\mathrm{EW}}^{2} \mathrm{~d} \sigma_{\mathrm{NNLO} \mathrm{EW}}+\alpha_{S} \alpha_{\mathrm{EW}} \mathrm{~d} \sigma_{\mathrm{NNLO} \text { QCDxEW }} \\
& \mathrm{NNLOQCD} \quad \mathrm{NNLOEW} \quad \mathrm{NNLO} \text { QCD-EW } \\
& +\alpha_{S}^{3} \mathrm{~d} \sigma_{\mathrm{NNLO}}+\ldots \\
& \mathrm{N} 3 \mathrm{LO} \mathrm{QCD}
\end{aligned}
$$



NNLO QCD


## Perturbative expansion:VV,VVV

(single perturbative order at LO)


NLO QCD NLO EW
$+\alpha_{S}^{2} \mathrm{~d} \sigma_{\mathrm{NNLO}}+\alpha_{\mathrm{EW}}^{2} \mathrm{~d} \sigma_{\text {NNLO EW }}+\alpha_{S} \alpha_{\mathrm{EW}} \mathrm{d} \sigma_{\mathrm{NNLO}} \mathrm{QCDxEW}$

scheme variation, e.g. Gmu vs. $a(m Z)$
scale variation at NNLO


NLO QCD + EW VS.
NLO QCD x EW

## Perturbative expansion for VV

$$
\begin{aligned}
\mathrm{d} \sigma=\mathrm{d} \sigma_{\mathrm{LO}}+ & \alpha_{S} \mathrm{~d} \sigma_{\mathrm{NLO}}+\alpha_{\mathrm{EW}} \mathrm{~d} \sigma_{\mathrm{NLOEW}} \\
& \mathrm{NLO} \mathrm{QCD} \quad \mathrm{NLO} \mathrm{EW} \\
& +\alpha_{S}^{2} \mathrm{~d} \sigma_{\mathrm{NNLO}}+\alpha_{\mathrm{EW}}^{2} \mathrm{~d} \sigma_{\mathrm{NNLOEW}}+\alpha_{S} \alpha_{\mathrm{EW}} \mathrm{~d} \sigma_{\mathrm{NNLO} \text { QCDxEW }} \\
& \mathrm{NNLO} \mathrm{QCD} \text { ? NNLO EW } \quad \text { NNLO QCD-EW } \\
& +\alpha_{S}^{3} \mathrm{~d} \sigma_{\mathrm{NNLO}}+\ldots \\
& \text { N3LO QCD }
\end{aligned}
$$

## NNLO QCD + NLO EW

| 41-SF-ZZ | $p p \rightarrow \ell^{+} \ell^{-} \ell^{+} \ell^{-}$ | ZZ |
| :---: | :---: | :---: |
| 41-DF-ZZ | $p p \rightarrow \ell^{+} \ell^{-} \ell^{\prime+} \ell^{\prime-}$ | ZZ |
| 31-SF-WZ | $p p \rightarrow \ell^{+} \ell^{-} \nu_{\nu_{\ell}}$ | WZ |
| 3l-DF-WZ | $p p \rightarrow \ell^{+} \ell^{-} \ell^{\prime} \nu_{\ell^{\prime}}$ | WZ |
| 21-SF-ZZ | $p p \rightarrow \ell^{+} \ell^{-} \nu_{\ell^{\prime} \bar{\nu}_{\ell^{\prime}}}$ | ZZ |
| 2l-SF-ZZWW | $p p \rightarrow \ell^{+} \ell^{-} \bar{\nu}_{\ell} \bar{\nu}_{\ell}$ | $\mathrm{ZZ}, \mathrm{WW}$ |
| 2l-DF-WW | $p p \rightarrow \ell^{+} \ell^{-} \nu_{\ell} \bar{\nu}_{\ell^{\prime}}$ | WW |

In Matrix+OpenLoops all (massive) diboson processes are now available at NNLO QCD + NLO EW
[M. Grazzini, S. Kallweit, JML, S. Pozzorini, M.Wiesemann; I 9 | 2.00068]

# NNLO QCD + NLO EW for dibosons: pTV2 

 pTV2

- moderate QCD corrections
- NNLO/NLO QCD very small at large pTV2
$\rightarrow$ NNLO QCD uncertainty: few percent
- NLO EW/LO=-(50-60)\% @ I TeV
- difference very conservative upper bound on $\mathcal{O}\left(\alpha_{S} \alpha\right)$
- multiplicative/factorised combination clearly superior (EW Sudakov logs $\times$ soft QCD)
- dominant uncertainty at large pTV2: $\mathcal{O}\left(\alpha^{2}\right) \sim \alpha_{\mathrm{w}}^{2} \log ^{4}\left(Q^{2} / M_{W}^{2}\right)$

Estimate: $\frac{1}{2} \delta_{\mathrm{EW}}^{2}$

## NNLO QCD + NLO EW for dibosons: pTV2



## Giant QCD K-factors and EW corrections: pTVI



## MEPS @ NLO QCD + EW:WW(+jet)

[Bräuer, Denner, Pellen, Schönherr, Schumann; '20]
-More rigorous solution: merge VV incl. approx. EW corrections with VV with Sherpa's MEPS@NLO QCD + EWvirt

- However, not NNLO QCD accurate

FO


FO

## MEPS @ NLO QCD + EW:WW(+jet)

[Bräuer, Denner, Pellen, Schönherr, Schumann; '20]
-More rigorous solution: merge $\mathrm{V} \mathrm{Vj}_{j}$ incl. approx. EW corrections with VV with Sherpa's MEPS@NLO QCD + EWvirt - However, not NNLO QCD accurate

FO


## MEPS@NLO QCD + EWvirt

$\mathrm{pp} \rightarrow \mu^{+} v_{\mu} \mathrm{e}^{-} \bar{v}_{\mathrm{e}}+$ jets @ 13 TeV


## MEPS @ NLO QCD + EW: ZZ (+jet)

[Bothmann, Napoletano, Schönherr, Schumann, Villani; '21]


- scheme variation: Gmu vs. a(mZ)
- EWsud based on [Bothmann, Napoletano, '20]:
process-independent implementation of Sudakov logs, see also [Pagani, Zaro '21]


## MEPS @ NLO QCD + EW: ZZ(+jet)

[Bothmann, Napoletano, Schönherr, Schumann, Villani; '2 I]



## PS MC: NNLO QCD + PS for WW via MiNNLOps

[Lombardi, Wiesemann; Zanderighi '2।]


- MiNNLOps physical down to pTVV=0
- Latest implementation does not require computationally expensive reweighting required earlier


## PS MC: NNLO QCD + PS for WW via MiNNLOps

[Lombardi, Wiesemann; Zanderighi '2।]



## PS MC: NNLO QCD + PS for ZZ via MiNNLOps

## Geneva: NNLO+PS

[Alioli, Broggio, Gavardi, Kallweit, Lim, Nagar, Napoletano, '2 I]

## MiNNLOps: nNNLO+PS

[Buonocore, Koole, Lombardi, Rottoli, Wiesemann, Zanderighi, '2 I]


## PS MC: NNLO QCD + PS for ZZ via MiNNLOps

[Buonocore, Koole, Lombardi, Rottoli, Wiesemann, Zanderighi, '2 I]


- nNNLO+PS predictions in good agreement with CMS results [arXiv:2009.0 I | 86]
- inclusion of EW corrections (through fixed-order NLO K-factor) required to describe tails of distributions


## PS MC: NLO QCD + NLO EW PS

[Chiesa, Re, Oleari '20]



# NLO (QCD + EW) PS (QCD + QED)/ NLO QCD PS (QCD + QED) 

## NLO (QCD + EW) PS (QCD + QED)/ NLO QCD PS QCD

- Note: resonance-aware NLO EW matching required (POWHEG-BOX-RES [Ježo, Nason, ' I 5])
- Missing: photon-induced channels
-Question: NLO (QCD + EW) PS (QCD + QED) / (NLO QCD PS QCD) × NLO EW


## PS MC: NNLO QCD $\times$ NLO EW PS forWZ

[MML, Lombardi, Wiesemann,, Zanderighi, Zanoli, to appear]


- NNLOPS QCD x NLOPS EW combination via reweighing (NLOPS EW resonance-aware)
- Next: combination at generator level


## gg-induced WW and ZZ production



- Formally same order as NNLO QCD
- Enhanced due to gg flux
- Interference with H->VV


- Sizeable QCD corrections (formally N3LO QCD)
- For m4I < 340 GeV 1/Mt expansion reliable



## NLO + PS for gg $\rightarrow \mathrm{VV} / \mathrm{H} \rightarrow 4 \mathrm{I}$

[Alioli, Ferrario Ravasio, JML, Röntsch, '2 1]

-ggWW/ggZZ @ NLO QCD + PS available!
-crucial for off-shell Higgs measurements

## Conclusions

- There is no clear scale/signature for new physics effects: Let's explore the unknown leaving no stone unturned!
- Precision is key for SM (QCD/EW/Higgs) measurements, SM parameter determination, as well as for BSM searches.

Incredible progress in theory predictions for multibosons V:


AND

- NNLO QCD + NLO EW available in MATRIX+OpenLoops
- MEPS @ NLO (QCD + EWapprox) available in Sherpa
- NLO (QCD + EW) + PS (QCD + QED) available in POWHEG
- NLO QCDgg PS available in POWHEG
- NNLO QCD PS via MiNNLO available (combined with NLOPS EW)

Remaining theory uncertainties: mixed QCD-EW, NNLO EW

Backup

## Giant K-factors and effect of jet veto



- at r2l $\rightarrow$ |: hard-VV topologies
- at $\mathrm{r} 2 \mathrm{I} \rightarrow 0$ : hard-Vj topologies



- for pTVI > $\mid$ TeV: hard-Vj topologies dominate over hard-VV
- Jet veto $H_{\mathrm{T}}^{\mathrm{jet}}<\xi_{\text {veto }} H_{\mathrm{T}}^{\text {lep }}$ corresponds to

$$
p_{\mathrm{T}, V_{2}} \geq \frac{1-\xi_{\text {veto }}}{1+\xi_{\text {veto }}} p_{\mathrm{T}, V_{1}}=\frac{2}{3} p_{\mathrm{T}, V_{1}} \quad \text { for } \quad \xi_{\text {veto }}=0.2
$$

(violated by off-shell topologies)

- Jet veto results in phase-space dominated by hard-VV


## Theory status for Tribosons

[Slide thanks to M. Schönherr]
NLO QCD corrections trivial, known for on-shell and o -shell processes.
NLO EW on-shell corrections calculated by Hefei group '14-'17, WWW also by Dittmaier, Huss, Knippen '17.

NLO EW off-shell corrections more involved, up to $2 \rightarrow 6$ complexity (like VBS, just with more and competing resonances)

$$
\begin{aligned}
& \text { - } p p \rightarrow \gamma \gamma \gamma / \gamma \gamma \ell \nu / \gamma \gamma \ell \ell \quad \text { Greiner, Schönherr '17 } \\
& \text { - } p p \rightarrow 3 \ell 3 \nu\left(\ell=e^{ \pm}, \mu^{ \pm}, 0 / 1 / 2\right. \text { SFOS channels, Schönherr '18 } \\
& \text { incl. WWW and WZZ topologies) } \\
& p p \rightarrow e^{\mp} \nu_{e} \mu^{ \pm} \nu_{\mu} \tau^{ \pm} \nu_{\tau}(W W W \text { only } \quad \text { Dittmaier, Knippen, Schwan '19 } \\
& \text { - } p p \rightarrow \gamma 2 \ell 2 \nu\left(\ell=e^{ \pm}, \mu^{ \pm}, 0,1\right. \text { SFOS channels, Ju, Lindert, Schönherr tbp } \\
& \text { incl. } \gamma W W \text { and } \gamma Z Z \text { topologies) }
\end{aligned}
$$

Generically, large contribution from photon-induced processes.

## Triboson production @ NLO QCD



- QCD correction driven by additional jet activity:VV+jet topologies with soft $V$
$\rightarrow$ 'giant K-factors'
$\rightarrow$ strong observable dependence
$\rightarrow$ NLO mandatory
- jet veto ( $\mathrm{p} \mathrm{T}_{\text {cut }}=50 \mathrm{GeV}$ ) reduces size and phase space dependence
$\rightarrow$ better: multi-jet merging

Triboson production: on-shell vs. off-shell



- at large mill and $p T_{\text {II }}$ large interference with other resonance structures


VS.


## Off-shell VVV production @ NLO EW



- Very large cancellations of EW corr. in qq and qY channels / highly observable dependent


## Interplay of WWW and Wh[ $\rightarrow$ WW* $]$

[Slide thanks to M. Schönherr]

- due to interference, Wh cannot be treated as independent background, but is part of the signal
$\rightarrow$ should not be subtracted

$\Rightarrow$ measure signature (e.g. $3 \ell+\mathrm{MET}$ ) in fiducial volume
$\rightarrow$ for limits on, e.g., AGCs: define fiducial region that has large WWW component, still measure signature, interferences can be as important as sought-after signal


## Perturbative expansion:VBF-V,VBS-VV

Example:WW+2jets


VS.

$$
\begin{equation*}
\mathrm{d} \sigma=\mathrm{d} \sigma\left(\alpha_{S}^{2} \alpha^{4}\right)+\mathrm{d} \sigma\left(\alpha_{S} \alpha^{5}\right)+\mathrm{d} \sigma\left(\alpha^{6}\right)+\ldots \tag{LO}
\end{equation*}
$$

QCD-background
interference
VBS-signal

## Perturbative expansion:VBF-V,VBS-VV

Example:WW+2jets


VS.


$$
\begin{equation*}
\mathrm{d} \sigma=\mathrm{d} \sigma\left(\alpha_{S}^{2} \alpha^{4}\right)+\mathrm{d} \sigma\left(\alpha_{S} \alpha^{5}\right)+\mathrm{d} \sigma\left(\alpha^{6}\right)+\ldots \tag{LO}
\end{equation*}
$$

QCD-background

$\cdots+\mathrm{d} \sigma\left(\alpha_{S}^{3} \alpha^{4}\right)+\mathrm{d} \sigma\left(\alpha_{S}^{2} \alpha^{5}\right)+\mathrm{d} \sigma\left(\alpha_{S} \alpha^{6}\right)+\sigma\left(\alpha^{7}\right)$
$\Rightarrow$ separation formally meaningless at NLO
$\Rightarrow$ strictly well defined measurements: fiducial cross sections

## QCD \& EW ZZ+2jets @ NLO QCD + EW

long-term program for VBS@NLO

- QCD and EW ss-WWjj at NLO QCD+EW: [Biedermann, Denner, Pellen '|6+'|7]
- EW WZjj at NLO QCD+EW: [Denner, Dittmaier, Maierhöfer, Pellen, Schwan, '|9]
- QCD and EW ZZZj at NLO QCD+EW: [Denner, Franken, Pellen, Schmidt, '20+'21]

EW ZZ+2jets @ NLO QCD + EW

$\cdot 2 \rightarrow 6$ particles at NLO EW!

| Order | $\mathcal{O}\left(\alpha^{6}\right)+\mathcal{O}\left(\alpha^{7}\right)$ | $\mathcal{O}\left(\alpha^{6}\right)+\mathcal{O}\left(\alpha_{\mathrm{s}} \alpha^{6}\right)$ | $\mathcal{O}\left(\alpha^{6}\right)+\mathcal{O}\left(\alpha^{7}\right)+\mathcal{O}\left(\alpha_{\mathrm{s}} \alpha^{6}\right)$ |
| :---: | :---: | :---: | :---: |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>100 \mathrm{GeV}$ |  |  |  |
| $\sigma_{\mathrm{NLO}}[\mathrm{fb}]$ | $0.08211(4)$ | $0.12078(11)$ | $0.10521(11)$ |
| $\delta[\%]$ | -15.9 | 23.6 | 7.7 |
| $M_{\mathrm{j}_{1} \mathrm{j}_{2}}>500 \mathrm{GeV}$ |  | $0.06077(25)$ |  |
| $\sigma_{\mathrm{NLO}}[\mathrm{fb}]$ | $0.06069(4)$ | $0.07375(25)$ | -17.5 |
| $\delta[\%]$ | -17.6 | 0.1 |  |

- In the VBS phase-space EW mode receives:
-very small QCD corrections (percent level)
$\rightarrow$ O(20\%) EW corrections


## QCD \& EW ZZ+2jets @ NLO QCD + EW

[Denner, Franken, Pellen, Schmidt; '21]



# QCD and EW V+2jets @ NLO QCD + EW 

[JML, S. Pozzorini, M. Schönherr; to appear]

## QCD-mode


-QCD: negative K-factor (increasing for large mjj), uncertainty ~20-25\%

- EW: up to $-10 \%$ in multi TeV


## EW-mode



- QCD: very small K-factor at large mjj,
uncertainty $\sim 10 \%$ (no VBF approximation)
- EW: up to $-20 \%$ in multi TeV

