

Gantries

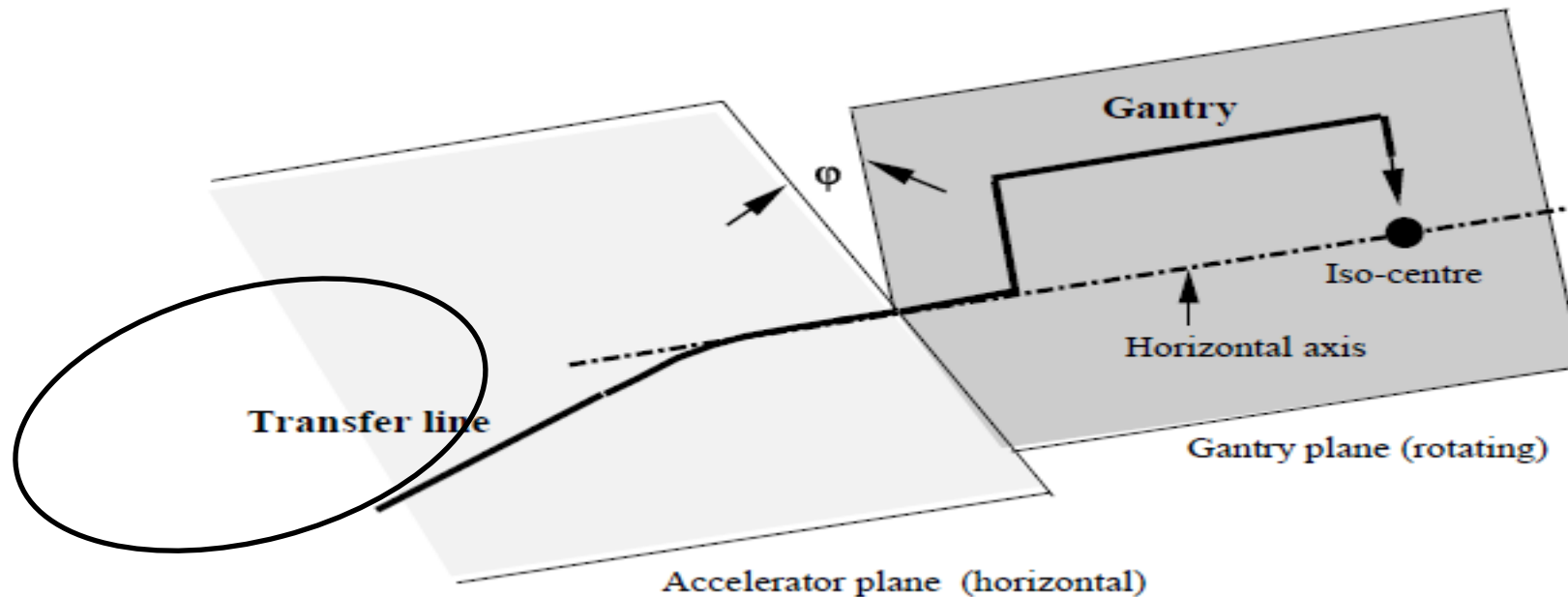
MARCO PULLIA



This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

What is a gantry?

A gantry is a section of beamline that can rotate around the isocenter in order to direct the beam onto the patient from any direction



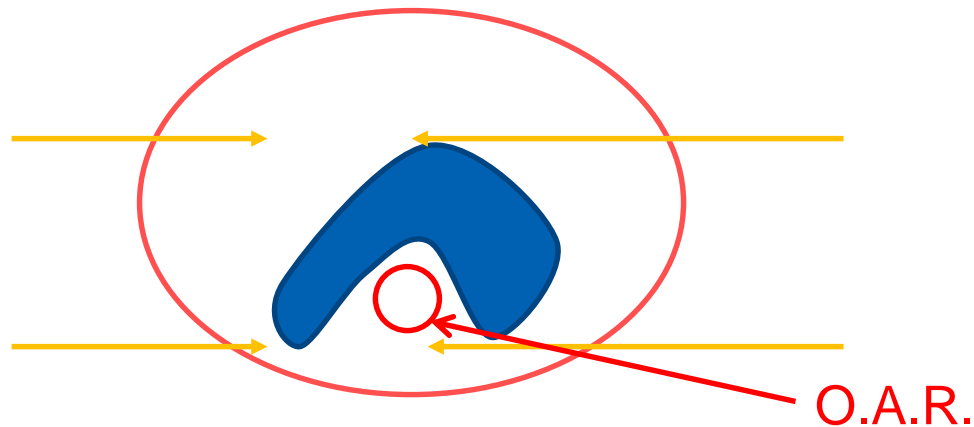
Why a gantry?

- To treat patients in supine position (eventually prone) in the same position in which CT, PET and MRI were acquired. Patient rotation only around gravity to preserve internal organs and soft tissue geometry.
- To provide the maximum flexibility in selecting the irradiation direction when optimising the dose delivery.
- To allow a “robust” treatment planning. Exploiting the sharp distal fall off can be risky in some cases and a gantry helps in avoiding fields directed towards an Organ At Risk (OAR).
- Avoid density heterogeneities
- Minimize SOBP extension (less energies required and better peak to plateau ratio)

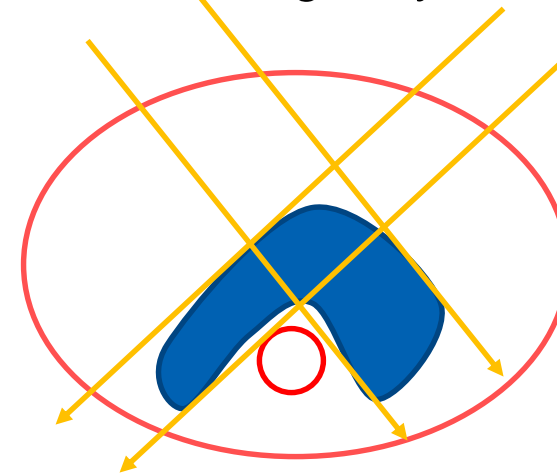
Why a gantry

Allows better, more robust planning:
e.g. minimize fields pointing towards OAR (Organ At Risk)

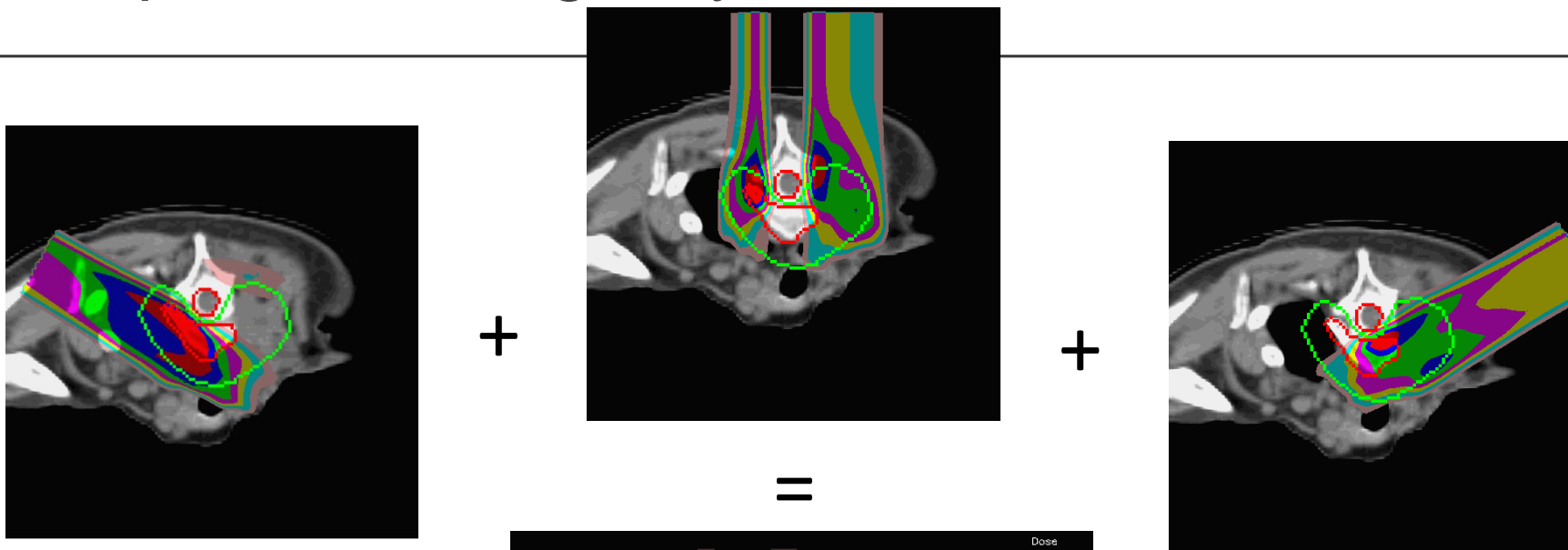
With horizontal line only



With gantry



Treatment planned with gantry

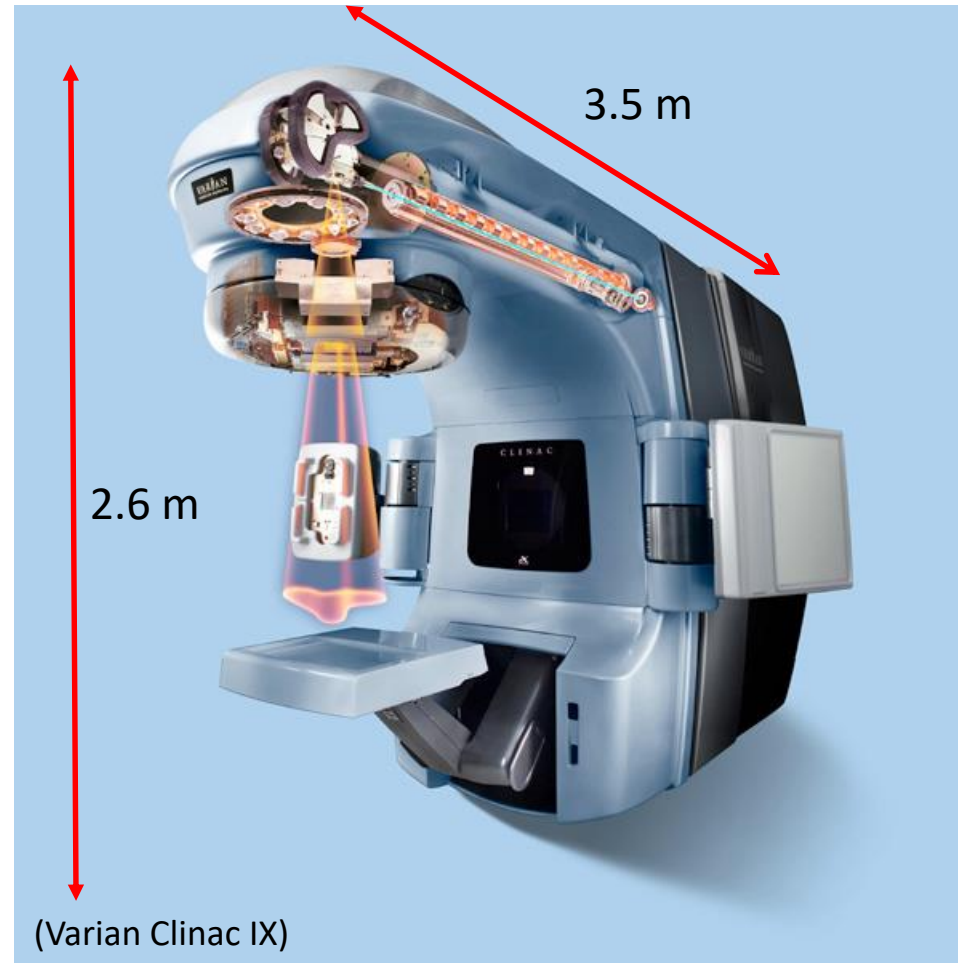


IMPT: each spot has an individually specified number of particles. The sum of the various fields is flat (or as required by clinics).

Gantry in conventional radiotherapy

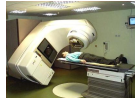
The whole linac is inside the gantry

The gantry head can pass between patient and floor for irradiation from below



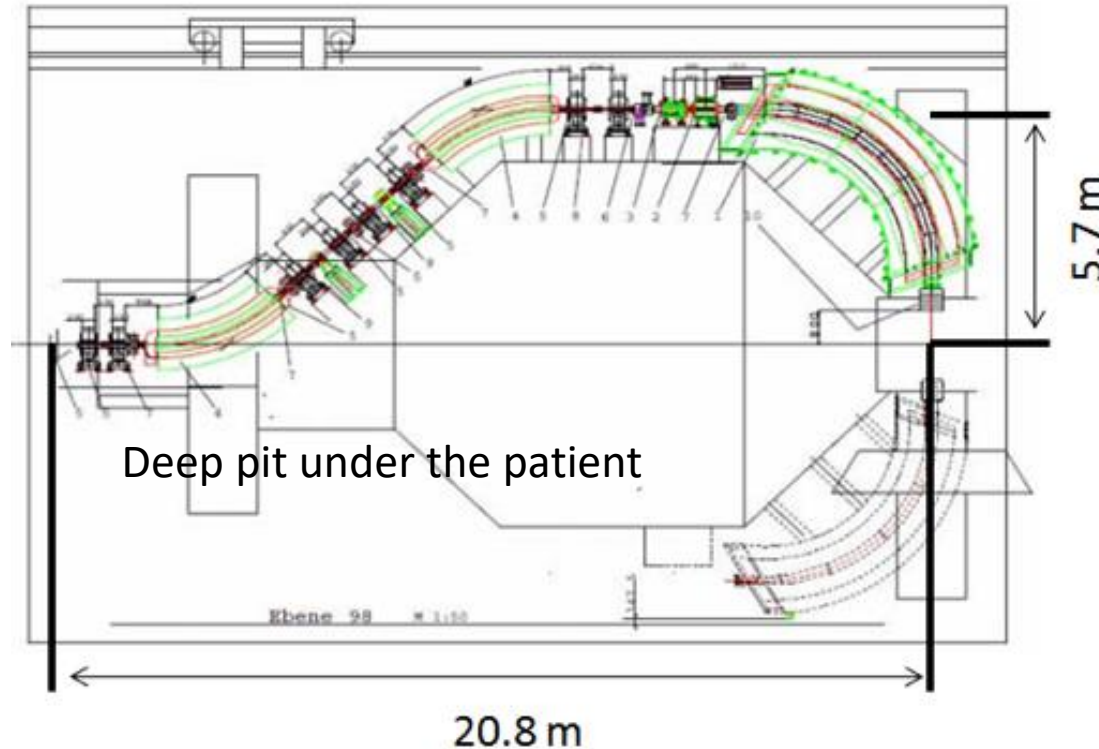
Gantries for particle therapy are large

Conventional RT

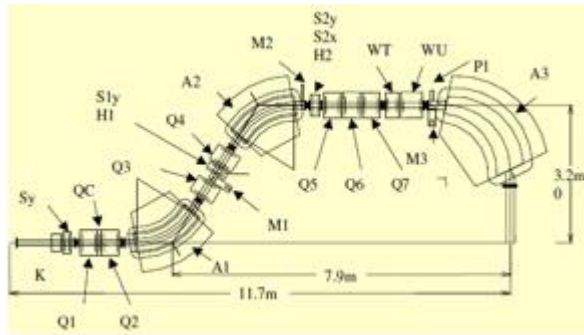


Carbon Ion Gantry

$B\rho < 6.6 \text{ Tm}$



The larger magnetic rigidity of carbon ions requires larger fields and larger bending radius

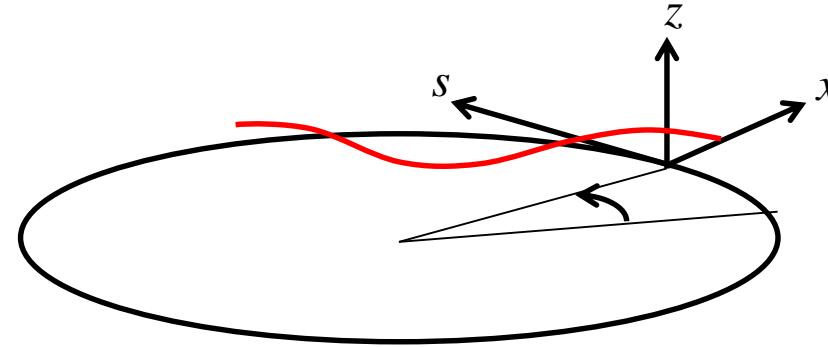


Proton Gantry
 $B\rho < 2.4 \text{ Tm}$

Betatron oscillations recap

With all the possible simplifications and linearizations, the motion of a particle with nominal energy along a magnetic lattice is described by the Hill's equation

$$\frac{d^2 y}{ds^2} + K(s)y = 0 \quad y = x \text{ or } z$$



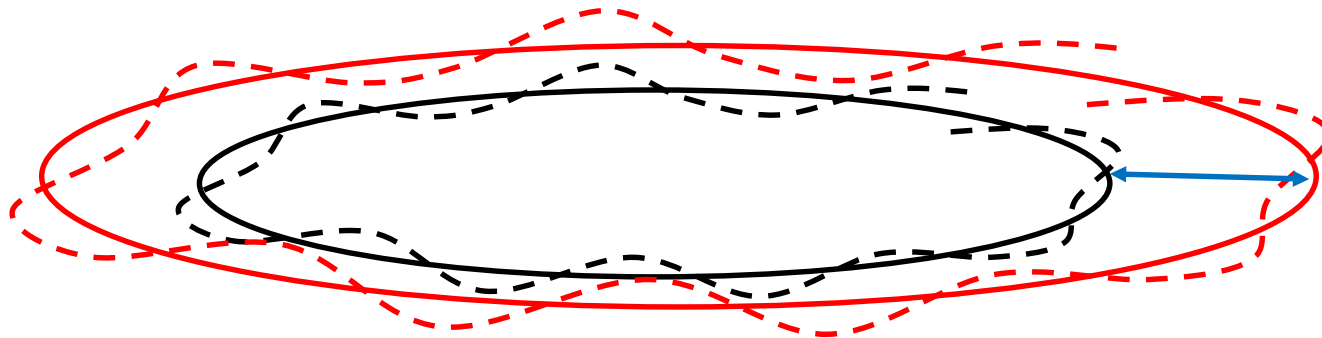
The Hill's equation looks like an harmonic oscillator, but $K(s)$ varies along the lattice depending on the magnetic element at position s . Assume it is constant inside each magnet and varies abruptly when passing from one element to the following (hard edge approximation)

For each element we can write a **transfer matrix** transporting the initial coordinates to the particle position at the element exit.

Dispersion

Particles with a (small) momentum deviation are bent differently wrt the nominal particle. Anyway a closed orbit/reference trajectory for particles with the considered momentum can be found and particles not moving on this new path orbit perform betatron oscillations around it.

The dispersion function expresses the closed orbit variation in terms of $\Delta p/p$



$$\begin{pmatrix} \Delta x(s) \\ \Delta x'(s) \end{pmatrix} = \begin{pmatrix} D_x(s) \\ D'_x(s) \end{pmatrix} \frac{\Delta p}{p_0}$$

The dispersion function originates from the dipoles and when no dipole is traversed the quantity

$$\gamma D^2 + 2\alpha D D' + \beta D'^2$$

stays constant .

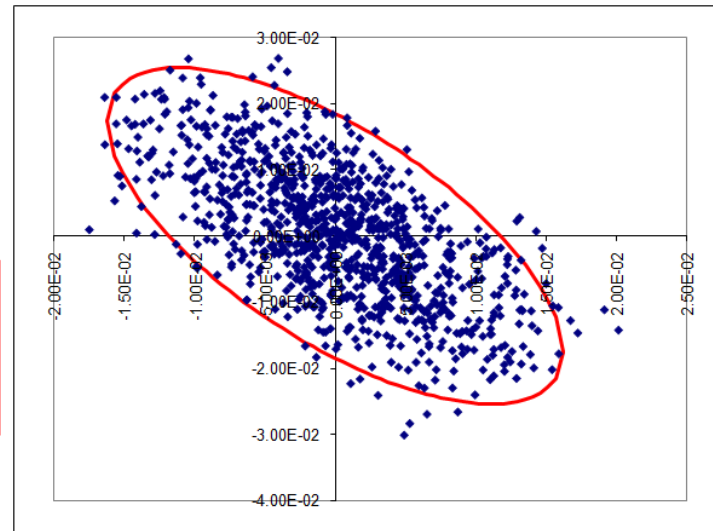
Circular accelerators and beamlines

The beta function in a ring is derived considering periodic conditions $K(s+L) = K(s)$. This defines clearly the meaning of the beta function, which describes the accelerator and the beam adapts to it.

In transfer lines the periodicity condition does not apply. One can choose the initial betas “freely”. Betas are useful if they describe the beam!

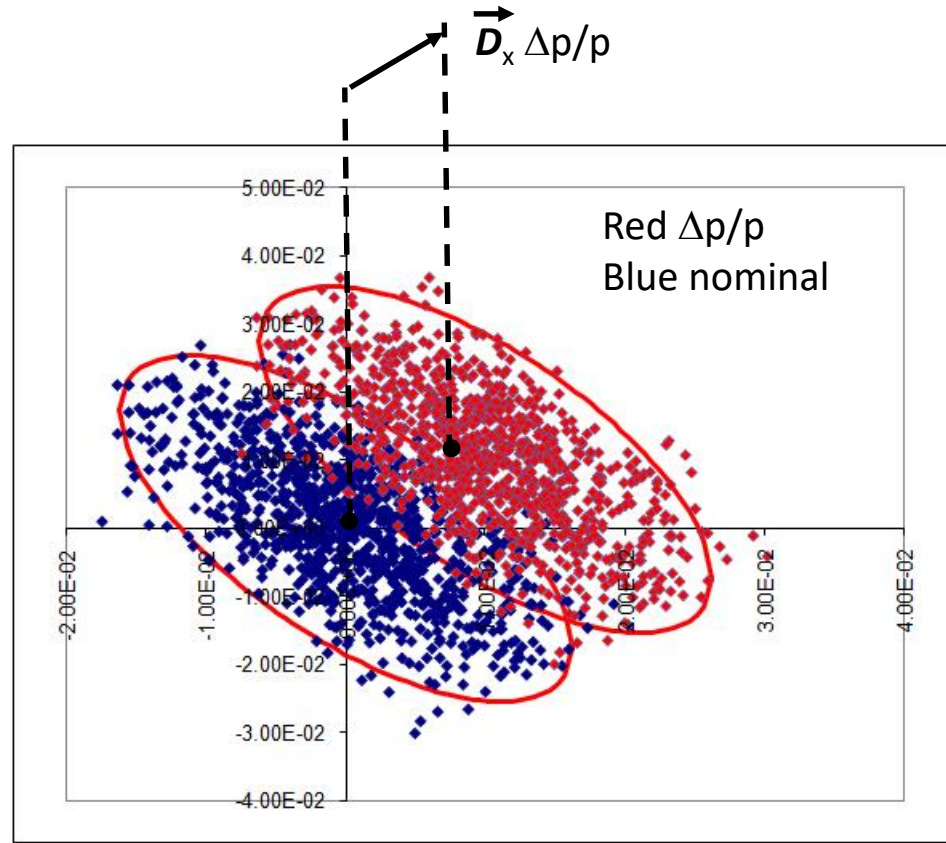
Area of ellipse = $\pi\varepsilon$

Now that it is defined by a beam ε is called **beam emittance**



Initial dispersion

As for betas, in a ring the periodic dispersion is clearly defined while in a transfer line the initial value of the dispersion function shall be based on the beam distribution



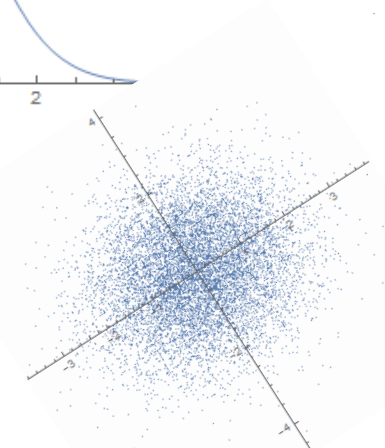
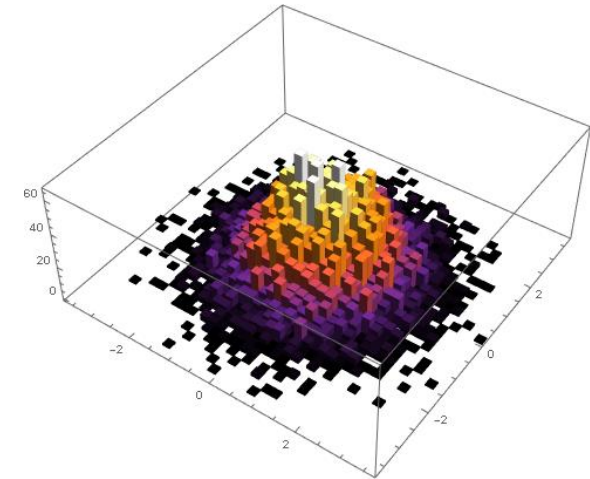
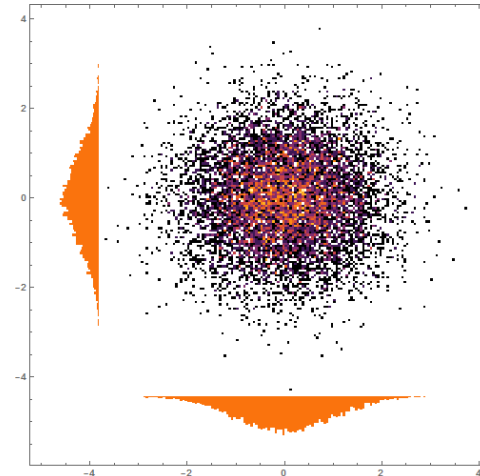
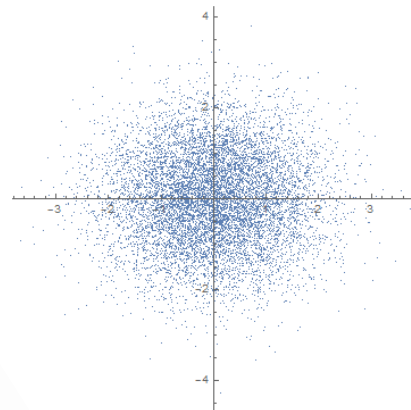
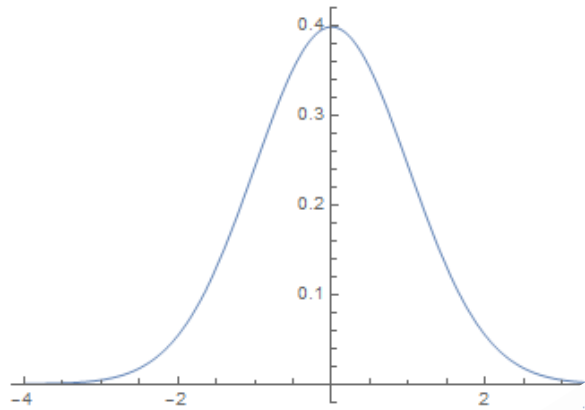
Matching to rotating gantries

Symmetric beam method with zero dispersion (exact)

- The beam at the entrance to the gantry must have zero dispersion and must be **rotationally symmetric** i.e. the same distribution (e.g. gaussian or KV) with equal Twiss functions and equal emittances in both planes at the entry to the gantry.
- The gantry must be designed to be a closed dispersion bump in the plane of bending (achromatic transport)

Matching to gantries

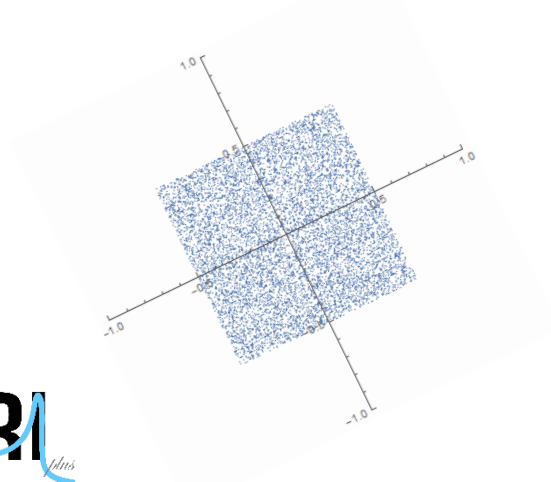
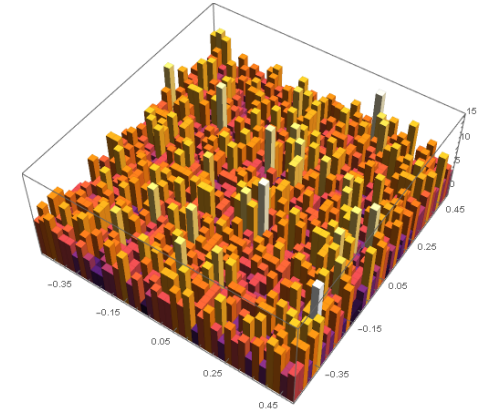
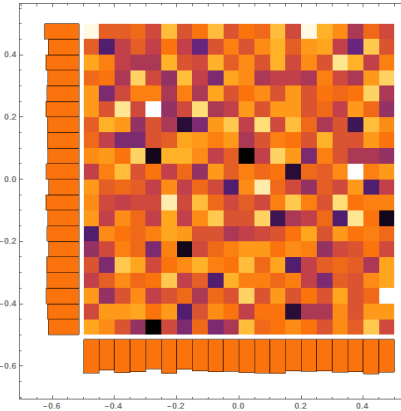
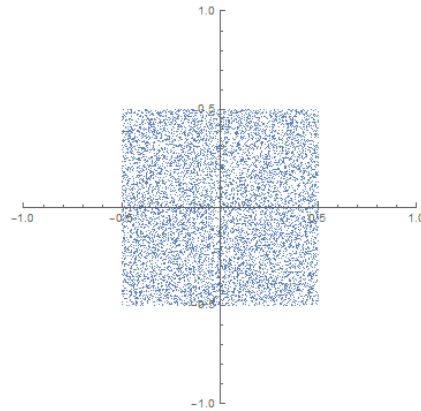
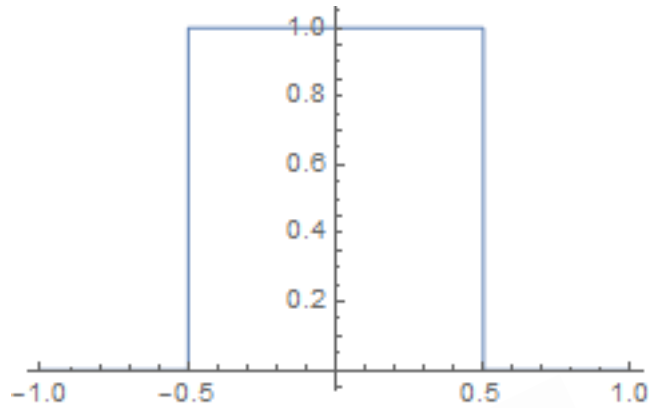
Factorised distribution $\rho_x(x) = \frac{e^{-\frac{x^2}{2}}}{\sqrt{2\pi}}$, $\rho_y(y) = \frac{e^{-\frac{y^2}{2}}}{\sqrt{2\pi}}$, $\rho(x,y) = \rho_x \cdot \rho_y$



Round distribution!
When rotated it's invariant!

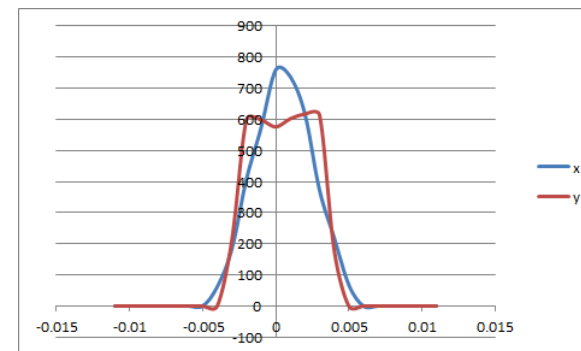
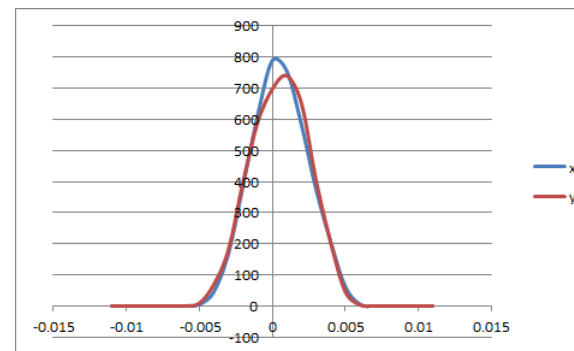
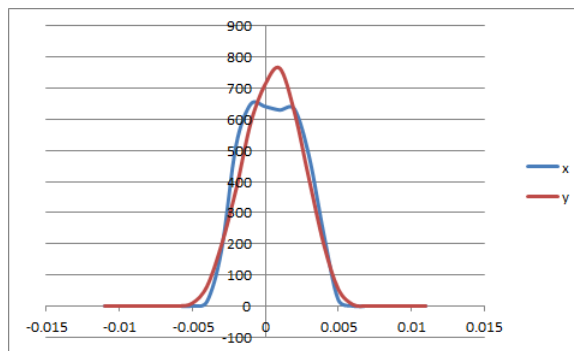
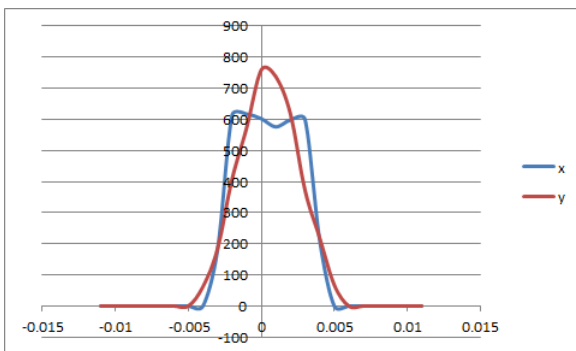
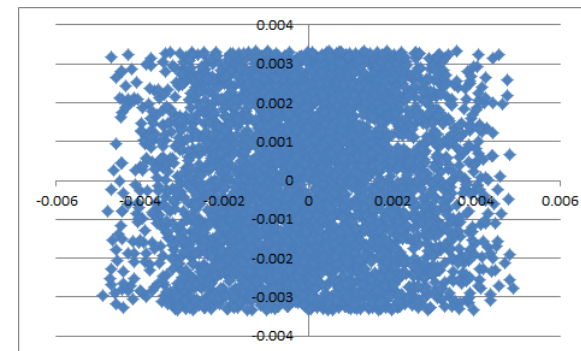
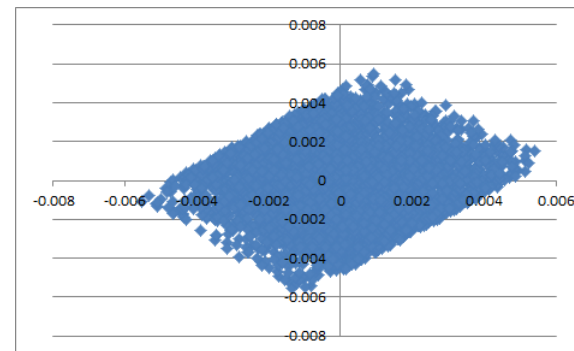
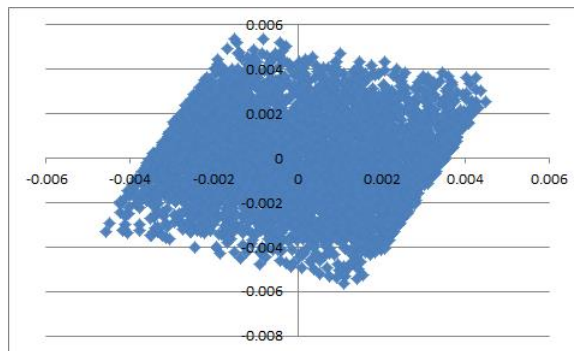
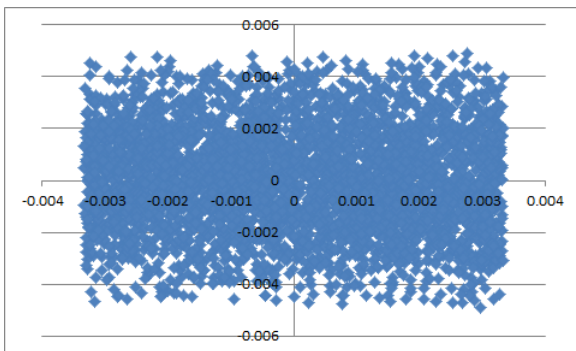
Matching to gantries

Factorised distribution $\rho_x(x) = 1/2A$, $\rho_y(y) = 1/2A$, $\rho(x,y) = \rho_x \cdot \rho_y$



Square distribution!
When rotated it's not invariant!

Beam Rotation



0°

20°

45°

90°

Sigma matching

Given a particle distribution we will call sigma matrix the covariance or better correlation matrix

Sigma Matrix

$$\sigma M = \begin{pmatrix} \langle x x \rangle & \langle x x' \rangle & \langle x y \rangle & \langle x y' \rangle \\ \langle x' x \rangle & \langle x' x' \rangle & \langle x' y \rangle & \langle x' y' \rangle \\ \langle y x \rangle & \langle y x' \rangle & \langle y y \rangle & \langle y y' \rangle \\ \langle y' x \rangle & \langle y' x' \rangle & \langle y' y \rangle & \langle y' y' \rangle \end{pmatrix}$$

When the beam is transported through (a system described by) a transfer matrix TM, then the sigma matrix transforms according to

$$\sigma M_f = TM \cdot \sigma M_i \cdot TM^T$$

Assume uncorrelated beam

$$\sigma M = \begin{pmatrix} \langle x x \rangle & \langle x x' \rangle & 0 & 0 \\ \langle x' x \rangle & \langle x' x' \rangle & 0 & 0 \\ 0 & 0 & \langle y y \rangle & \langle y y' \rangle \\ 0 & 0 & \langle y' y \rangle & \langle y' y' \rangle \end{pmatrix}$$

$$= \begin{pmatrix} \sigma_{11} & \sigma_{12} & 0 & 0 \\ \sigma_{21} & \sigma_{22} & 0 & 0 \\ 0 & 0 & \sigma_{33} & \sigma_{34} \\ 0 & 0 & \sigma_{43} & \sigma_{44} \end{pmatrix} = \begin{pmatrix} \sigma_{11} & \sigma_{12} & 0 & 0 \\ \sigma_{12} & \sigma_{22} & 0 & 0 \\ 0 & 0 & \sigma_{33} & \sigma_{34} \\ 0 & 0 & \sigma_{34} & \sigma_{44} \end{pmatrix}$$

$$\sigma M = \begin{pmatrix} \epsilon_{\text{rms},x} \beta_x & -\epsilon_{\text{rms},x} \alpha_x & 0 & 0 \\ -\epsilon_{\text{rms},x} \alpha_x & \epsilon_{\text{rms},x} \gamma_x & 0 & 0 \\ 0 & 0 & \epsilon_{\text{rms},y} \beta_y & -\epsilon_{\text{rms},y} \alpha_y \\ 0 & 0 & -\epsilon_{\text{rms},y} \alpha_y & \epsilon_{\text{rms},y} \gamma_y \end{pmatrix}$$

Transport through a rotated gantry

Rotation matrix

$$RM = \begin{pmatrix} \cos[\theta] & 0 & \sin[\theta] & 0 \\ 0 & \cos[\theta] & 0 & \sin[\theta] \\ -\sin[\theta] & 0 & \cos[\theta] & 0 \\ 0 & -\sin[\theta] & 0 & \cos[\theta] \end{pmatrix}$$

Gantry transfer matrix

$$GM = \begin{pmatrix} g_{11} & g_{12} & 0 & 0 \\ g_{21} & g_{22} & 0 & 0 \\ 0 & 0 & g_{33} & g_{34} \\ 0 & 0 & g_{43} & g_{44} \end{pmatrix}$$

$$\sigma M_f = GM \cdot RM \cdot \sigma M_i \cdot RM^T \cdot GM^T$$

$$\sigma M_f = GM \cdot RM \cdot \sigma M_i \cdot RM^T \cdot GM^T$$

$$\begin{pmatrix} g_{11} & g_{12} & 0 & 0 \\ g_{21} & g_{22} & 0 & 0 \\ 0 & 0 & g_{33} & g_{34} \\ 0 & 0 & g_{43} & g_{44} \end{pmatrix} \cdot \begin{pmatrix} \cos[\theta] & 0 & \sin[\theta] & 0 \\ 0 & \cos[\theta] & 0 & \sin[\theta] \\ -\sin[\theta] & 0 & \cos[\theta] & 0 \\ 0 & -\sin[\theta] & 0 & \cos[\theta] \end{pmatrix} \cdot \begin{pmatrix} \sigma_{11} & \sigma_{12} & 0 & 0 \\ \sigma_{21} & \sigma_{22} & 0 & 0 \\ 0 & 0 & \sigma_{33} & \sigma_{34} \\ 0 & 0 & \sigma_{43} & \sigma_{44} \end{pmatrix} \cdot \begin{pmatrix} \cos[\theta] & 0 & -\sin[\theta] & 0 \\ 0 & \cos[\theta] & 0 & -\sin[\theta] \\ \sin[\theta] & 0 & \cos[\theta] & 0 \\ 0 & \sin[\theta] & 0 & \cos[\theta] \end{pmatrix} \cdot \begin{pmatrix} g_{11} & g_{21} & 0 & 0 \\ g_{12} & g_{22} & 0 & 0 \\ 0 & 0 & g_{33} & g_{43} \\ 0 & 0 & g_{34} & g_{44} \end{pmatrix}$$

Let's evaluate some matrix element

$$sM_f(1,1) = g_{11}^2 (\sigma_{11} \cos^2[\theta] + \sigma_{33} \sin^2[\theta]) + g_{11} g_{12} (\sigma_{12} \cos^2[\theta] + \sigma_{21} \cos^2[\theta] + \sigma_{34} \sin^2[\theta] + \sigma_{43} \sin^2[\theta]) + g_{12}^2 (\sigma_{22} \cos^2[\theta] + \sigma_{44} \sin^2[\theta])$$

$$sM_f(3,3) = g_{33}^2 (\sigma_{33} \cos^2[\theta] + \sigma_{11} \sin^2[\theta]) + g_{33} g_{34} (\sigma_{34} \cos^2[\theta] + \sigma_{43} \cos^2[\theta] + \sigma_{12} \sin^2[\theta] + \sigma_{21} \sin^2[\theta]) + g_{34}^2 (\sigma_{44} \cos^2[\theta] + \sigma_{22} \sin^2[\theta])$$

Beam size independent of rotation angle

$$sM_f(1,1) = g_{11}^2 (\sigma_{11} \cos^2[\theta] + \sigma_{33} \sin^2[\theta]) + g_{11} g_{12} (\sigma_{12} \cos^2[\theta] + \sigma_{21} \cos^2[\theta] + \sigma_{34} \sin^2[\theta] + \sigma_{43} \sin^2[\theta]) + g_{12}^2 (\sigma_{22} \cos^2[\theta] + \sigma_{44} \sin^2[\theta])$$

$$sM_f(3,3) = g_{33}^2 (\sigma_{33} \cos^2[\theta] + \sigma_{11} \sin^2[\theta]) + g_{33} g_{34} (\sigma_{34} \cos^2[\theta] + \sigma_{43} \cos^2[\theta] + \sigma_{12} \sin^2[\theta] + \sigma_{21} \sin^2[\theta]) + g_{34}^2 (\sigma_{44} \cos^2[\theta] + \sigma_{22} \sin^2[\theta])$$

Complete symmetry: $\sigma_{11} = \sigma_{33}$; $\sigma_{12} = \sigma_{34}$; $\sigma_{22} = \sigma_{44}$ that is same emittance and same Twiss parameters in the horizontal and vertical plane; in this case **the beam is in itself rotationally symmetric** and any gantry transfer matrix gives an angle independent result;

Beam size independent of rotation angle

$$sM_f(1,1)=g_{11}^2 (\sigma_{11} \cos[\theta]^2+\sigma_{33} \sin[\theta]^2)+g_{11} g_{12} (\sigma_{12} \cos[\theta]^2+\sigma_{21} \cos[\theta]^2+\sigma_{34} \sin[\theta]^2+\sigma_{43} \sin[\theta]^2)+g_{12}^2 (\sigma_{22} \cos[\theta]^2+\sigma_{44} \sin[\theta]^2)$$

$$sM_f(3,3)=g_{33}^2 (\sigma_{33} \cos[\theta]^2+\sigma_{11} \sin[\theta]^2)+g_{33} g_{34} (\sigma_{34} \cos[\theta]^2+\sigma_{43} \cos[\theta]^2+\sigma_{12} \sin[\theta]^2+\sigma_{21} \sin[\theta]^2)+g_{34}^2 (\sigma_{44} \cos[\theta]^2+\sigma_{22} \sin[\theta]^2)$$

Round beam in (x, y): $\sigma_{11} = \sigma_{33}$; then the beam size at the isocenter independent of the gantry angle is obtained imposing **$g_{12} = 0$ and $g_{34} = 0$** . The gantry matrix then looks like

$$GM = \begin{pmatrix} g_{11} & 0 & 0 & 0 \\ g_{21} & g_{22} & 0 & 0 \\ 0 & 0 & g_{33} & 0 \\ 0 & 0 & g_{43} & g_{44} \end{pmatrix} = \begin{pmatrix} g_{11} & 0 & 0 & 0 \\ g_{21} & 1/g_{11} & 0 & 0 \\ 0 & 0 & g_{33} & 0 \\ 0 & 0 & g_{43} & 1/g_{33} \end{pmatrix}$$

This is referred to as “**point-to-point**” optics, since particles with the same initial position in (x, y) end up in the same position independently of the initial divergence (x', y')

Beam size independent of rotation angle

$$sM_f(1,1)=g_{11}^2 (\sigma_{11} \cos[\theta]^2+\sigma_{33} \sin[\theta]^2)+g_{11} g_{12} (\sigma_{12} \cos[\theta]^2+\sigma_{21} \cos[\theta]^2+\sigma_{34} \sin[\theta]^2+\sigma_{43} \sin[\theta]^2)+g_{12}^2 (\sigma_{22} \cos[\theta]^2+\sigma_{44} \sin[\theta]^2)$$

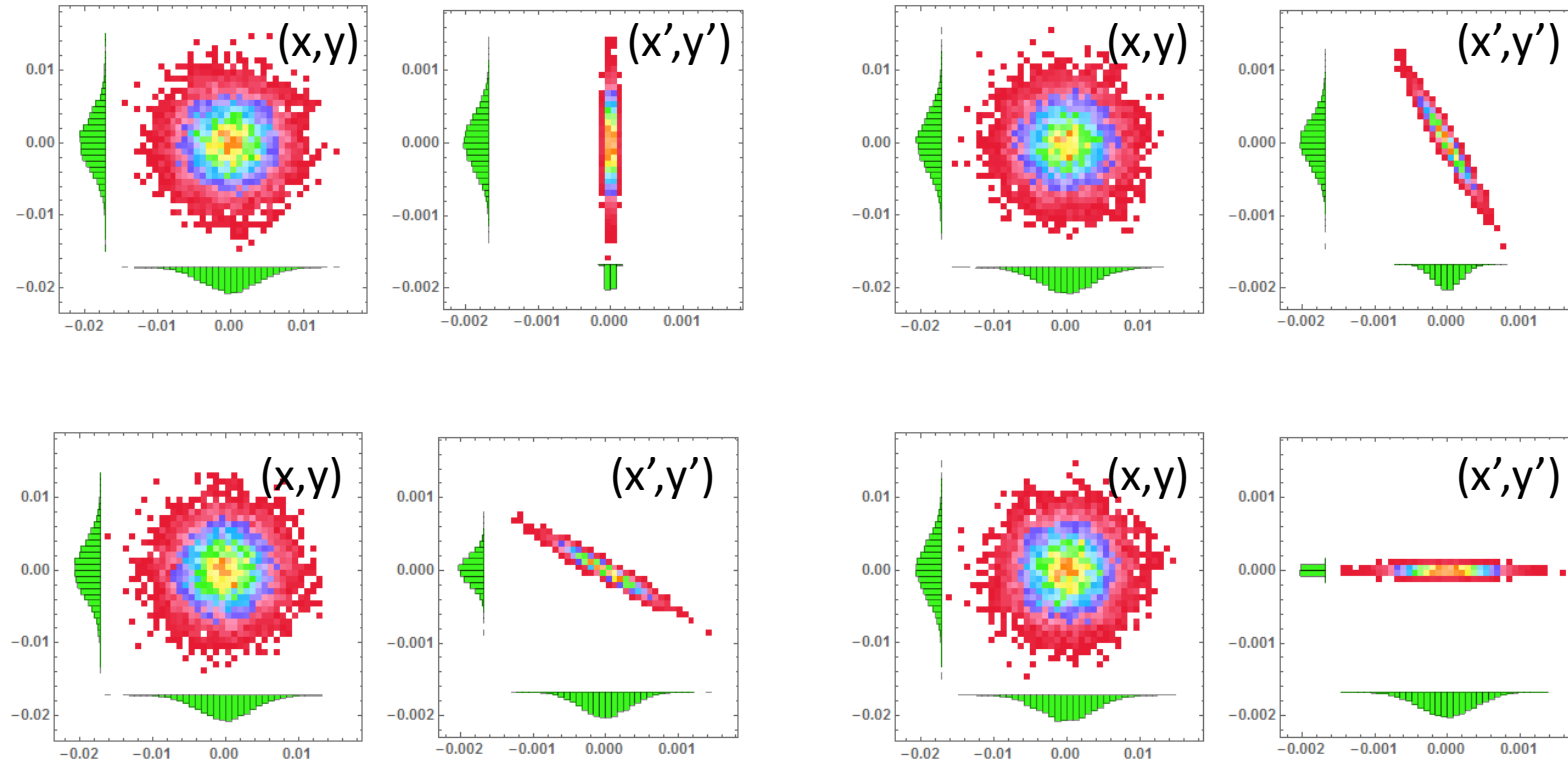
$$sM_f(3,3)=g_{33}^2 (\sigma_{33} \cos[\theta]^2+\sigma_{11} \sin[\theta]^2)+g_{33} g_{34} (\sigma_{34} \cos[\theta]^2+\sigma_{43} \cos[\theta]^2+\sigma_{12} \sin[\theta]^2+\sigma_{21} \sin[\theta]^2)+g_{34}^2 (\sigma_{44} \cos[\theta]^2+\sigma_{22} \sin[\theta]^2)$$

Round beam in (x', y'): $\sigma_{22} = \sigma_{44}$; then a beam size at the isocenter independent of the gantry angle is obtained imposing $g_{11} = 0$ and $g_{33} = 0$. The gantry matrix then looks like

$$GM = \begin{pmatrix} 0 & g_{12} & 0 & 0 \\ g_{21} & g_{22} & 0 & 0 \\ 0 & 0 & 0 & g_{34} \\ 0 & 0 & g_{43} & g_{44} \end{pmatrix} = \begin{pmatrix} 0 & g_{12} & 0 & 0 \\ -1/g_{12} & g_{22} & 0 & 0 \\ 0 & 0 & 0 & g_{34} \\ 0 & 0 & -1/g_{34} & g_{44} \end{pmatrix}$$

This is referred to as “**parallel-to-point**” optics, since particles with the same initial divergence (x', y') end up in the same position independently of the initial position (x, y)

Beam Rotation (point to point gantry)



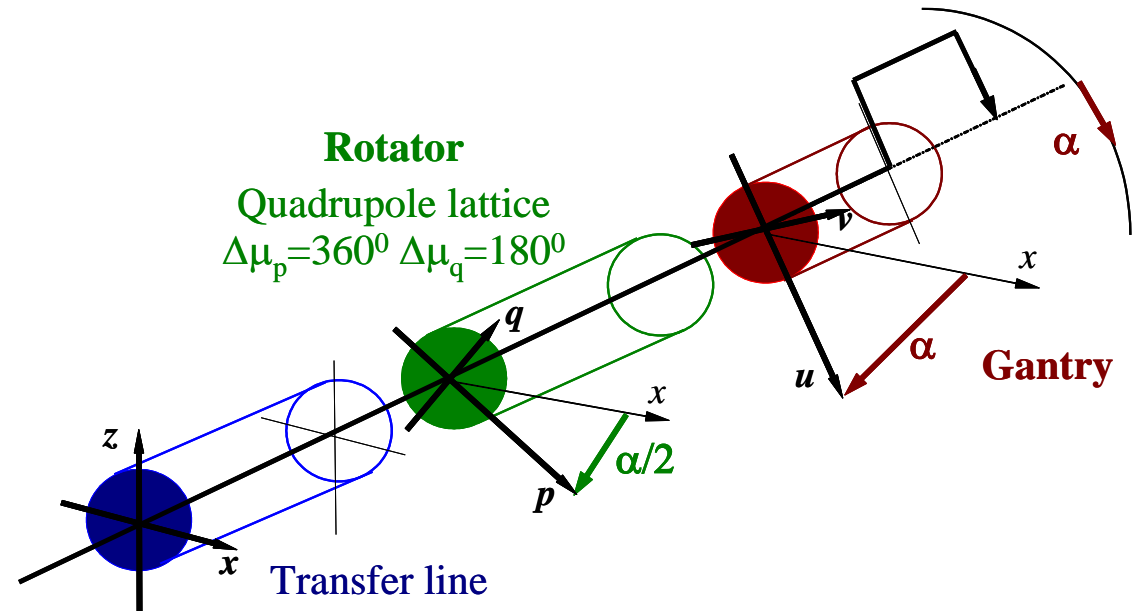
Rotator

The rotator is a lattice with transfer function
 It is rotated by half the gantry angle

$$RTM = \begin{pmatrix} r_{11} & r_{12} & 0 & 0 \\ r_{21} & r_{22} & 0 & 0 \\ 0 & 0 & -r_{11} & -r_{12} \\ 0 & 0 & -r_{21} & -r_{22} \end{pmatrix}$$

Historically/often

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$



Rotator

When the rotator is rotated by half of the gantry angle, **it maps the incoming beam to the gantry rotated axes**, as shown by the matrix multiplication of: rotation (fixed line to rotator) times rotator times rotation (rotator to gantry). If the rotator is rotated by θ and the gantry by 2θ , then

$$\begin{pmatrix} \cos[\theta] & 0 & \sin[\theta] & 0 \\ 0 & \cos[\theta] & 0 & \sin[\theta] \\ -\sin[\theta] & 0 & \cos[\theta] & 0 \\ 0 & -\sin[\theta] & 0 & \cos[\theta] \end{pmatrix} \cdot RTM \cdot \begin{pmatrix} \cos[\theta] & 0 & \sin[\theta] & 0 \\ 0 & \cos[\theta] & 0 & \sin[\theta] \\ -\sin[\theta] & 0 & \cos[\theta] & 0 \\ 0 & -\sin[\theta] & 0 & \cos[\theta] \end{pmatrix} = RTM$$

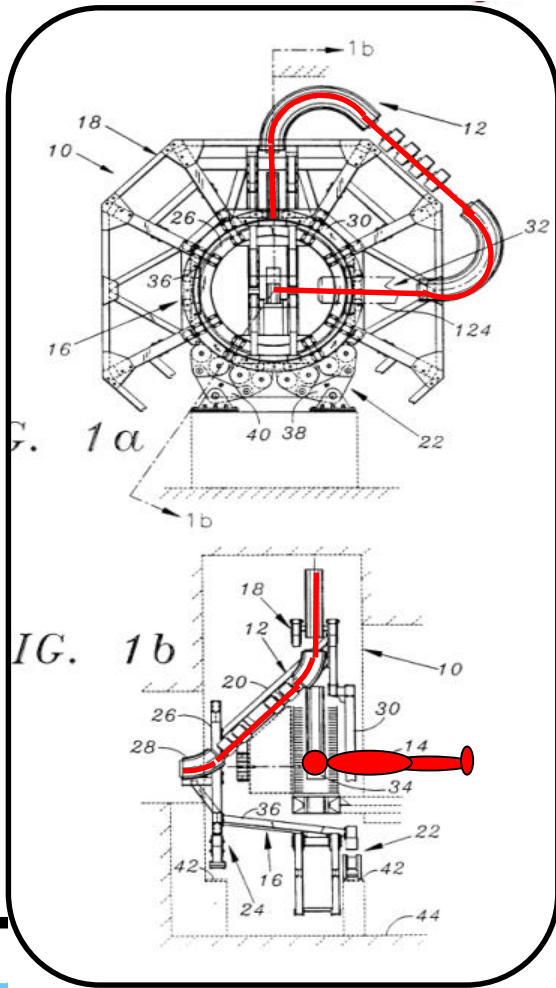
Which is independent of θ , showing that an incoming particle enters into the gantry always in the same way (with the same initial coordinates in the gantry reference system) independently of the rotation angle.

The rotator **rotates also the dispersion** function and allows therefore to close the dispersion in the fixed part of the beamline.

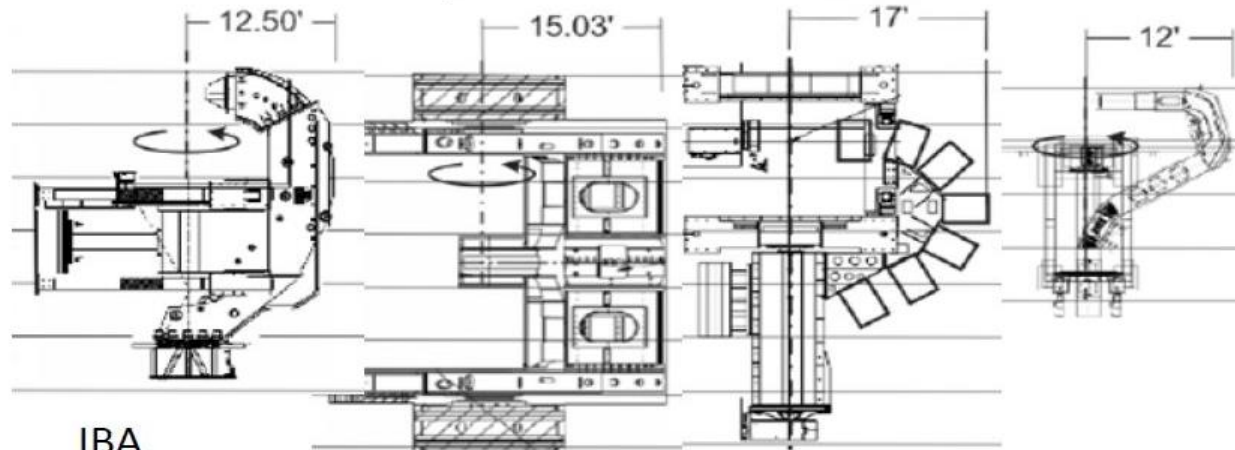
Gantry overview



Many geometries used



Corkscrew gantry at LLUMC

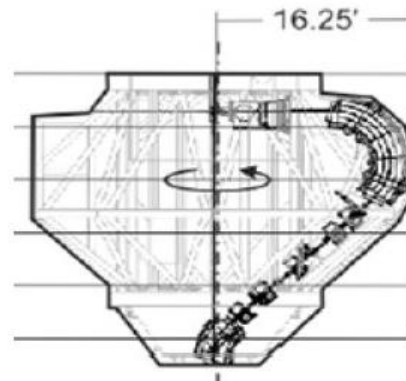


IBA Proteus One

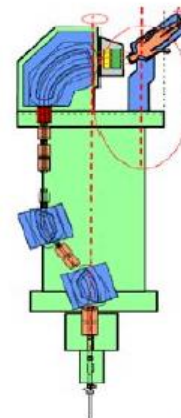
MeVlon

ProTom

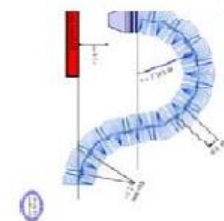
ProNova



IBA/SHI/Hitachi In-Plane



PSI 1 shortest radius



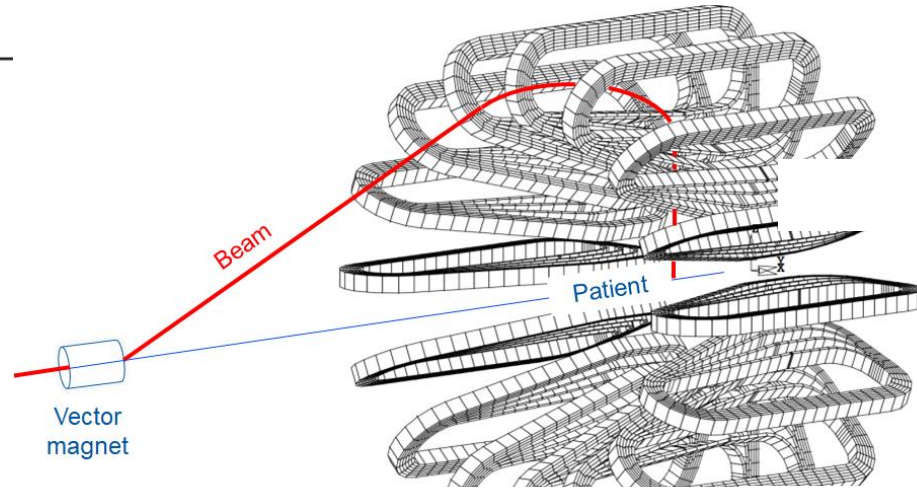
FFAG is the lightest

(Adapted from a slide of J. Flanz)

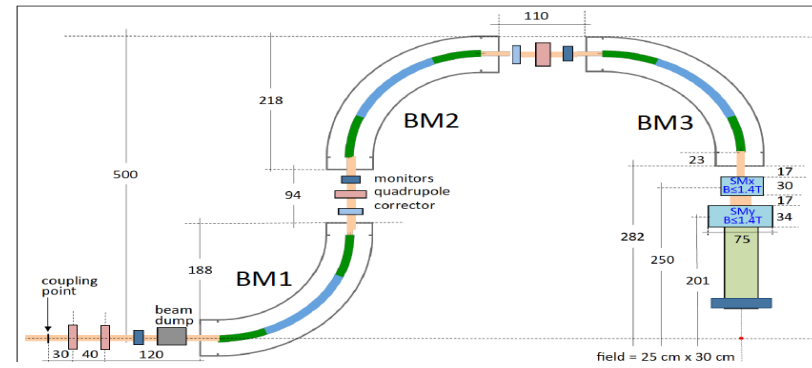
Alternative designs and large momentum acceptance



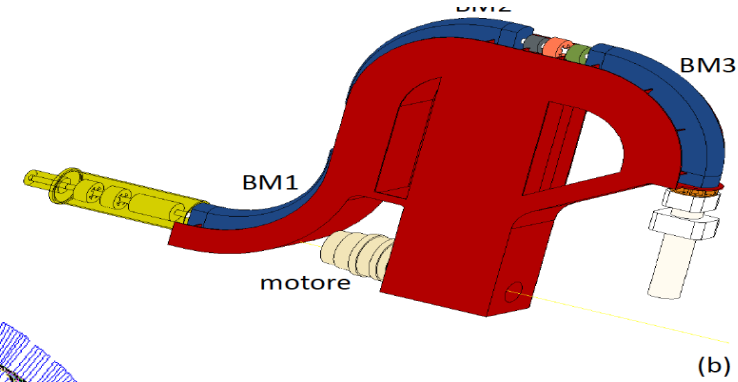
Leo cancer care



Toroidal magnet SC design (L. Bottura)

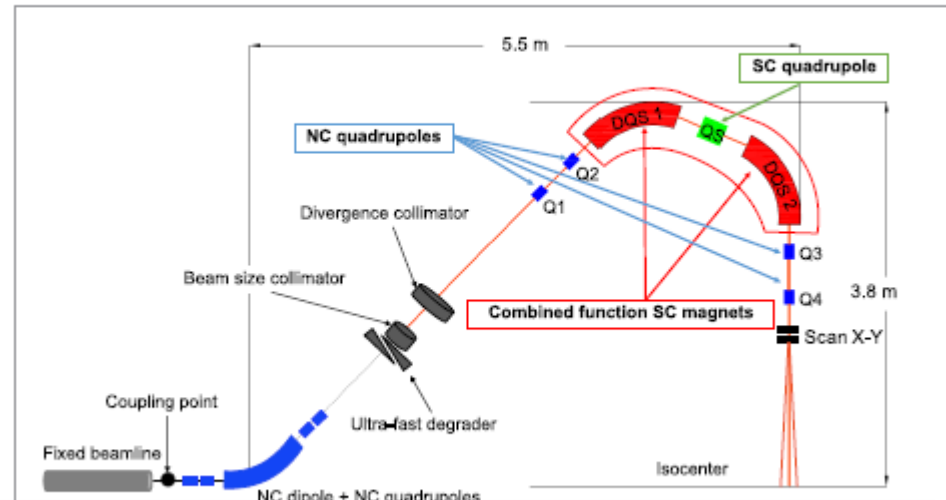


TERA-CERN-LBNL (SC canted cosine theta)

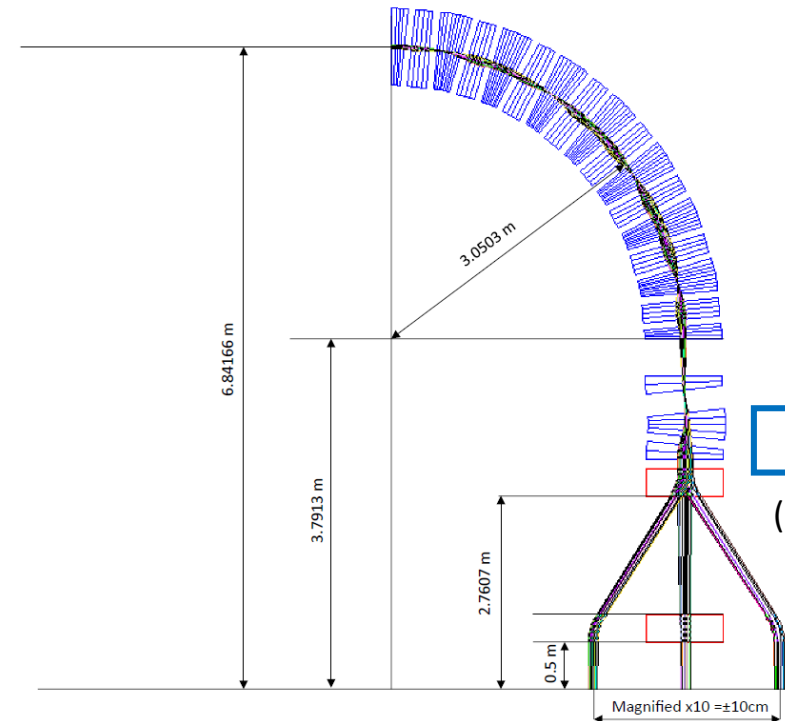


Phys. Med. Biol. 64 (2019) 175007 (13pp)

K P Nesteruk et al



PSI (large acceptance SC magnet)



FFAG

(Courtesy of Dejan Trbojevic)

financing from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

Proton gantries

Mitsubishi



Hitachi

IBA

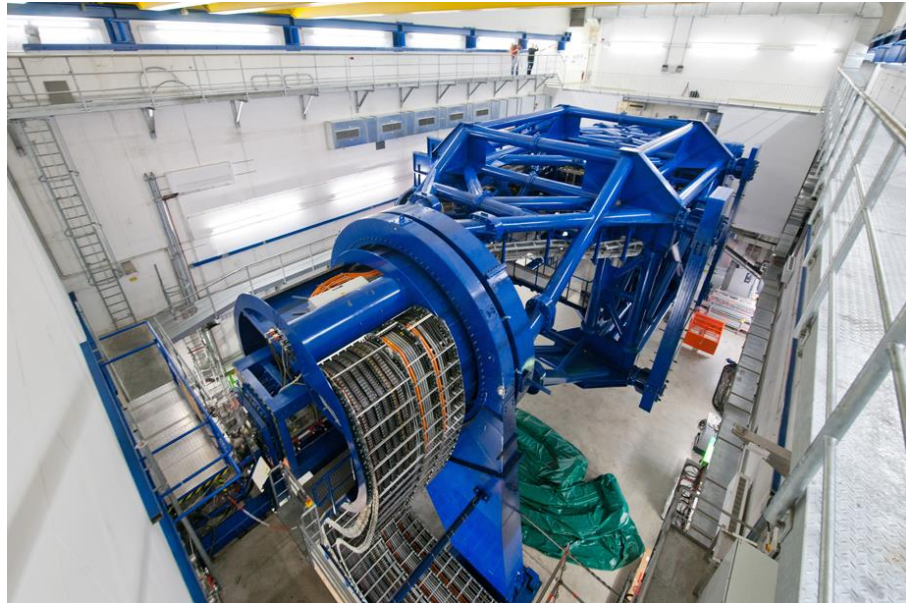


Horizon 2020
No 101008548

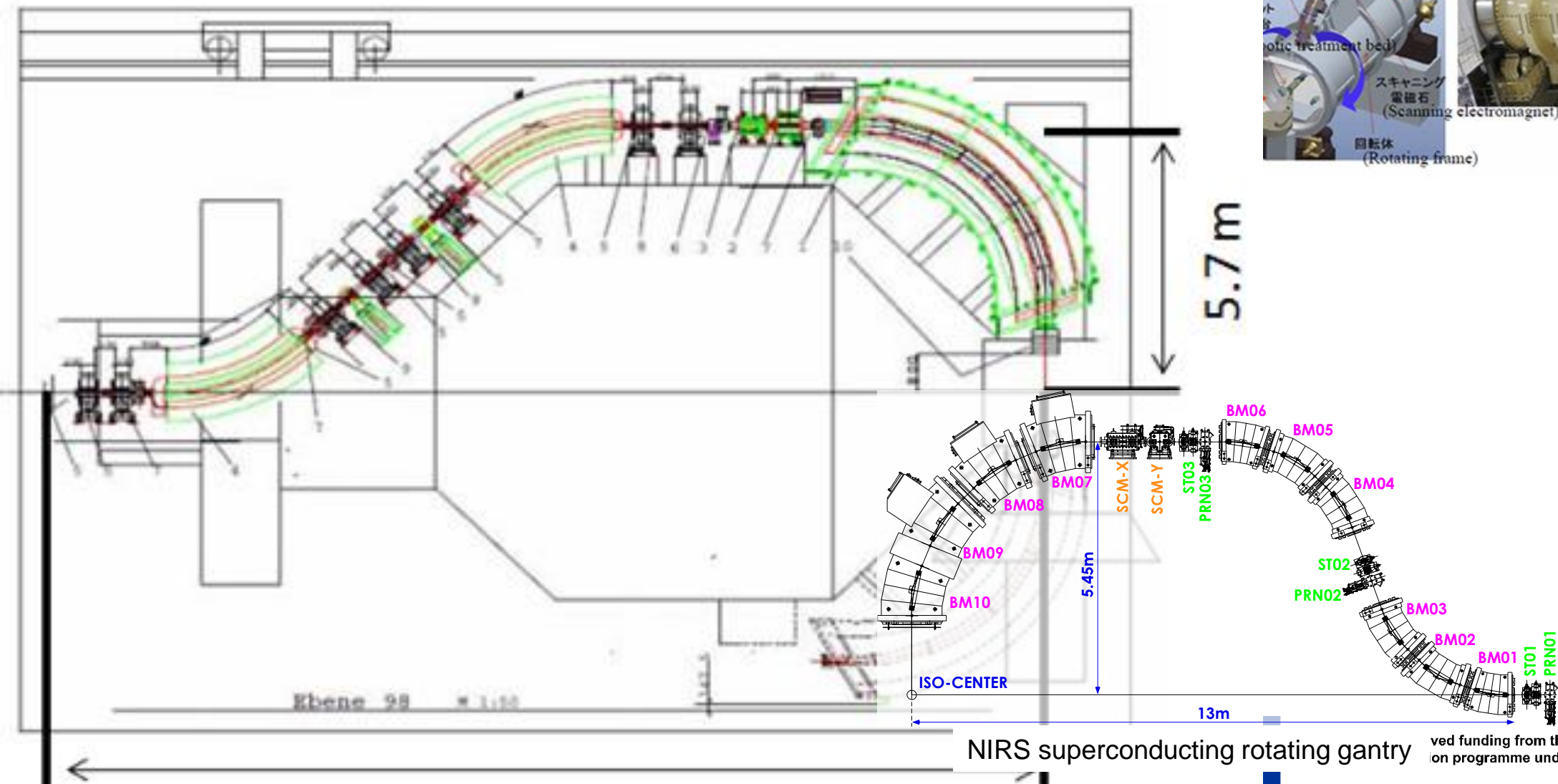
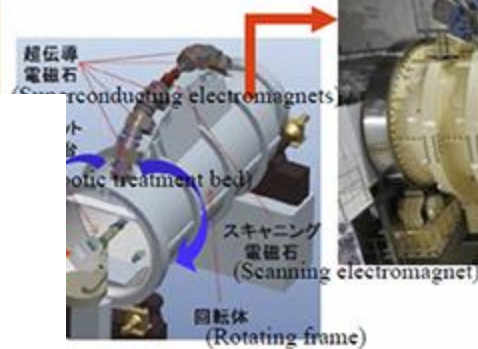
Carbon Ion Gantries

The first one:
The HIT Gantry

$L = 25 \text{ m} \times \phi = 13 \text{ m}$,
600 t rotating mass



HIMAC



NIRS superconducting rotating gantry funded from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

Yamagata

Toshiba's Heavy Ion Therapy System

3: Compact Rotating Gantry



Compact rotating gantry at Yamagata University

▶ ▶ 🔊 3:07 / 5:36

Scorri per i dettagli

© 2021 Toshiba Energy Systems & Solutions Corporation



07/07/2022



2:55 / 5:36

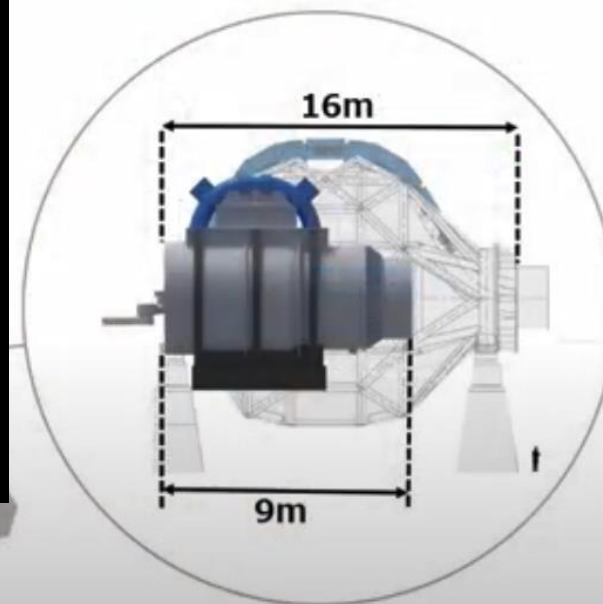
Scorri per i dettagli

© 2021 Toshiba Energy Systems & Solutions Corporation



<https://www.youtube.com/watch?v=8mAkVcW2LtA>
<https://www.youtube.com/watch?v=iVTBzCBCY90>

from our NIRS/QST prototype



*: compared to NIRS/QST gantry

Space around the isocenter

Patient size
Walk around patient
Imaging in situ
Couch rotation
Additional, future
instrumentation

Typical
~ 45 – 65 cm
~ 2 m opposite to
nozzle

Scattering, air and
distance degrade
beam quality



(Photon gantry used for illustration only, text refers to particles)

g from the European Union's Horizon 2020
nme under grant agreement No 101008548

Field size

Area that can be irradiated

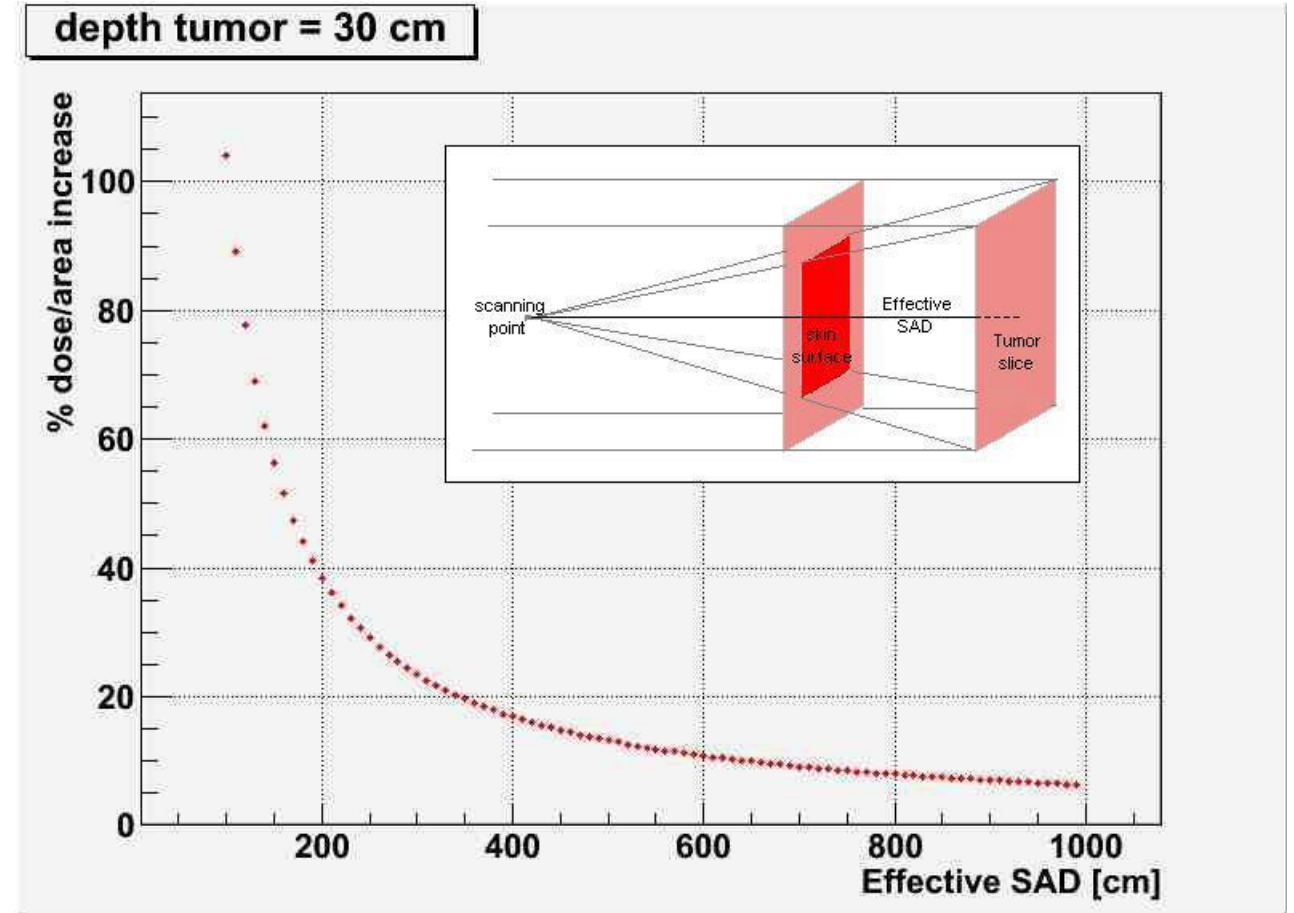
Trade off between performance and gantry cost/size



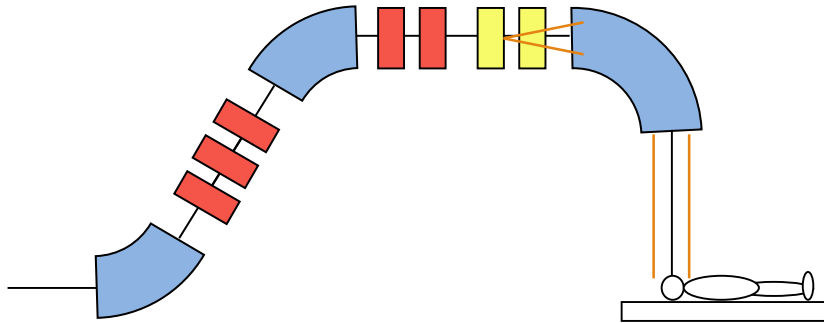
SAD

SAD - Source to Axis Distance

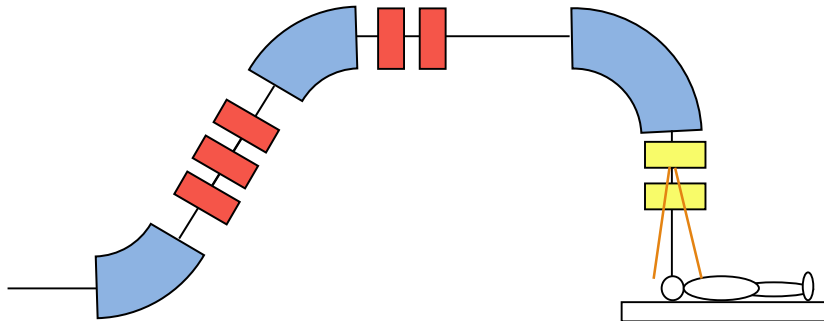
Dose increase to the skin
(which is a radiosensitive organ)



Scanning magnets position



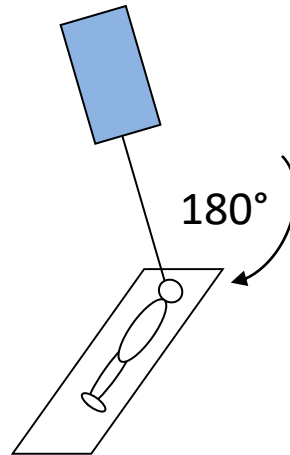
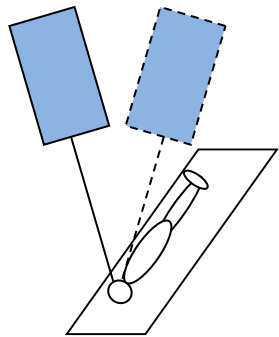
- Large aperture dipole: weight and power consumption
- Parallel scanning



- Large gantry radius and large room size

180 vs 360

By rotating the couch by 180°, all the beam directions are possible also with only 180° of rotation of the gantry

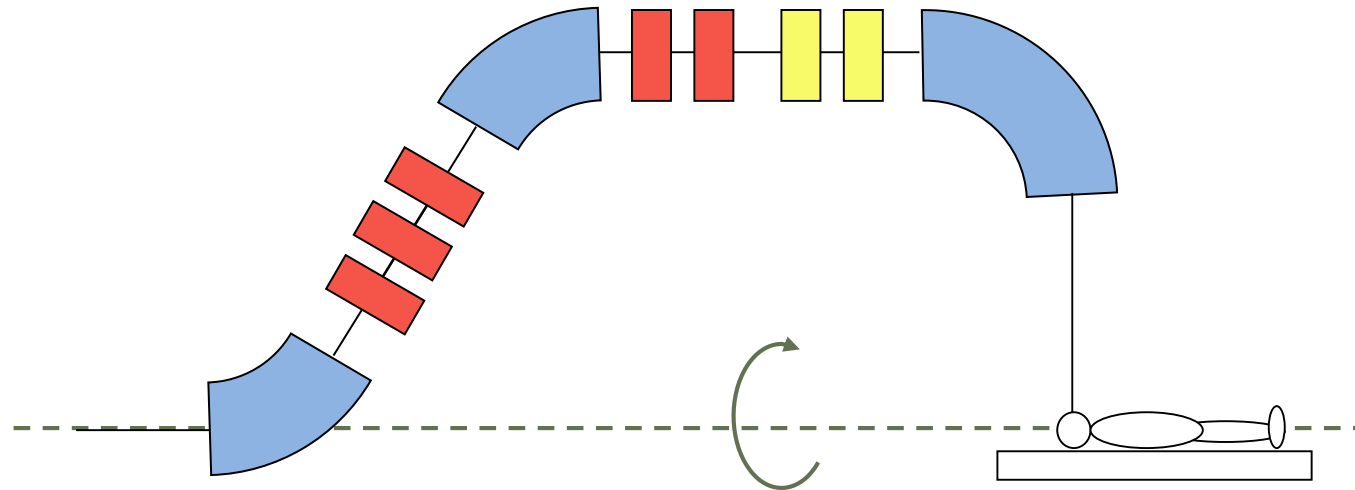


Rotation of the couch may require position verification (time and X Rays),
But it saves space and requires less shielding on the wall “not irradiated”.

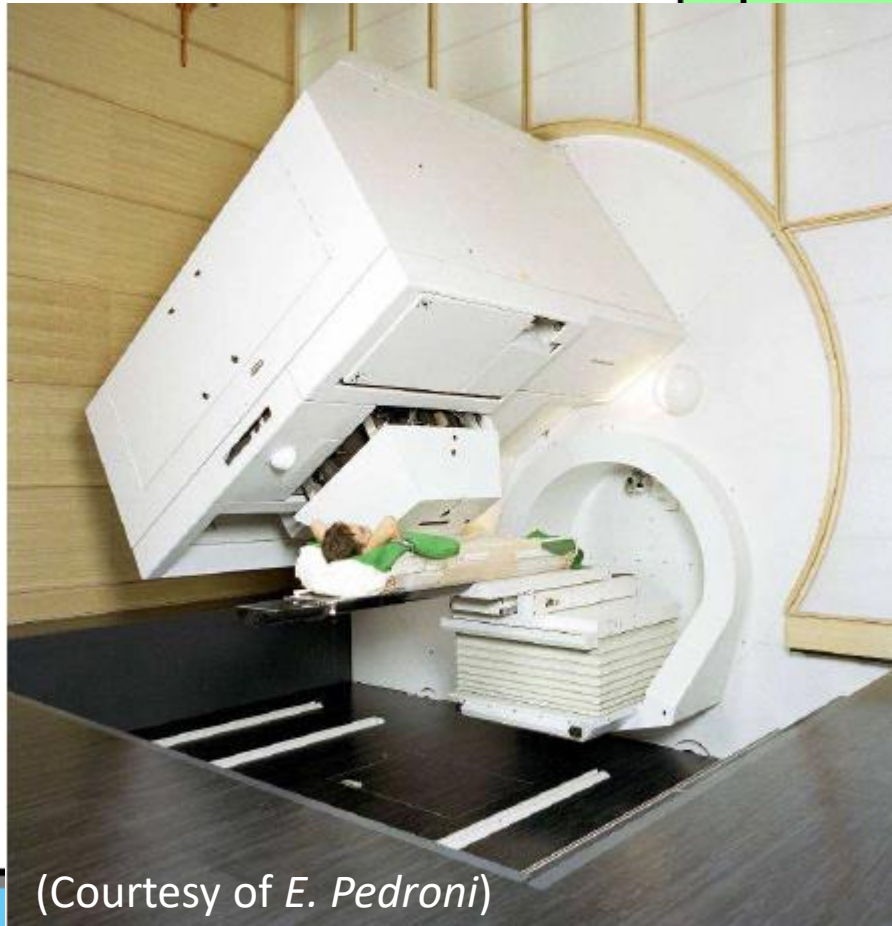
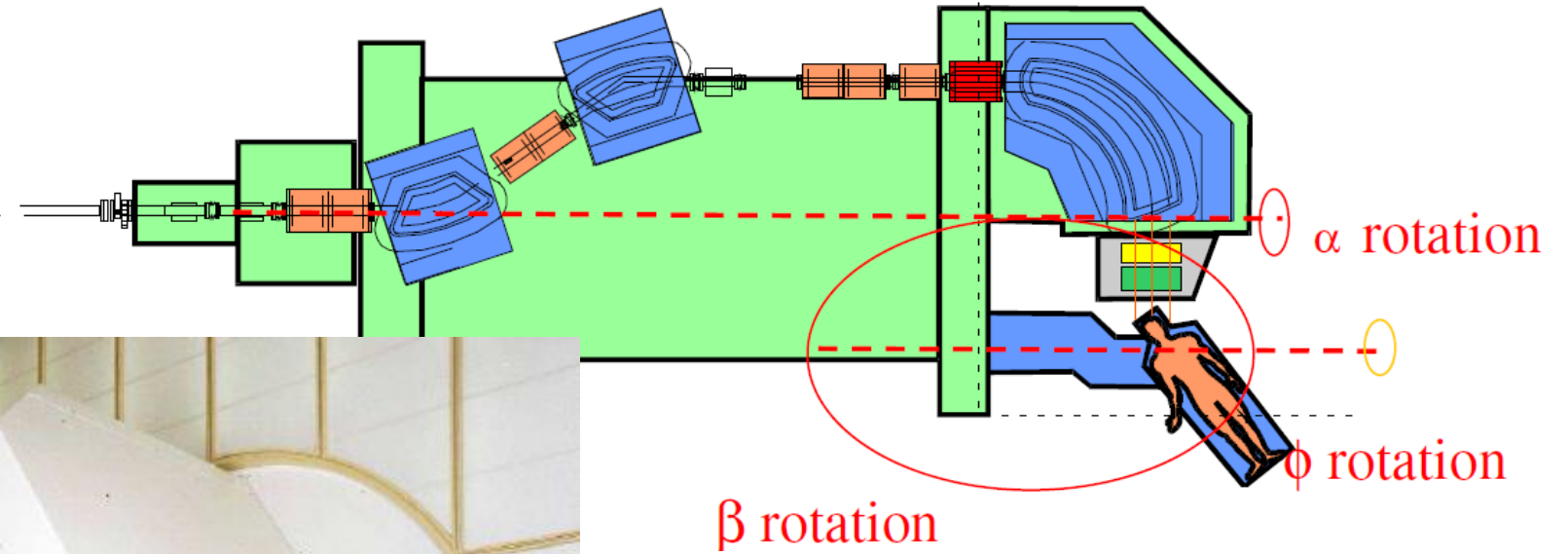


Fixed or mobile isocenter

Most of the existing gantries have a fixed isocenter on the rotation axis of the gantry. This implies large masses rotating at large radius.



PSI gantry 1



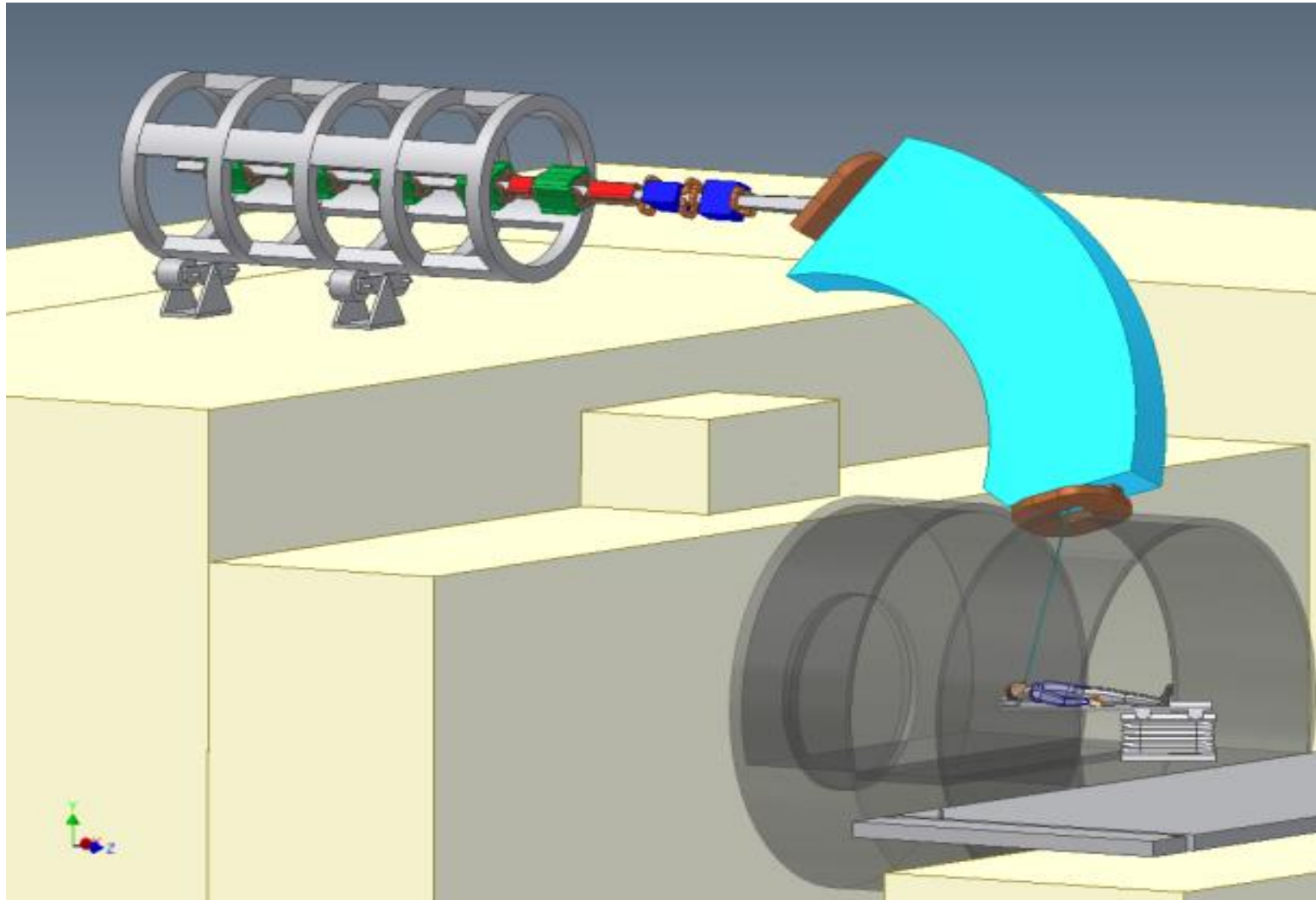
First scanning gantry worldwide

An isocenter, through which all the directions pass, exists but its position depends on gantry orientation.

“Riesenrad” gantry

ULICE Gantry

Dispersion is closed in the fixed line



90° bending magnet



Scanning magnets



Quadrupoles



Corrector magnets



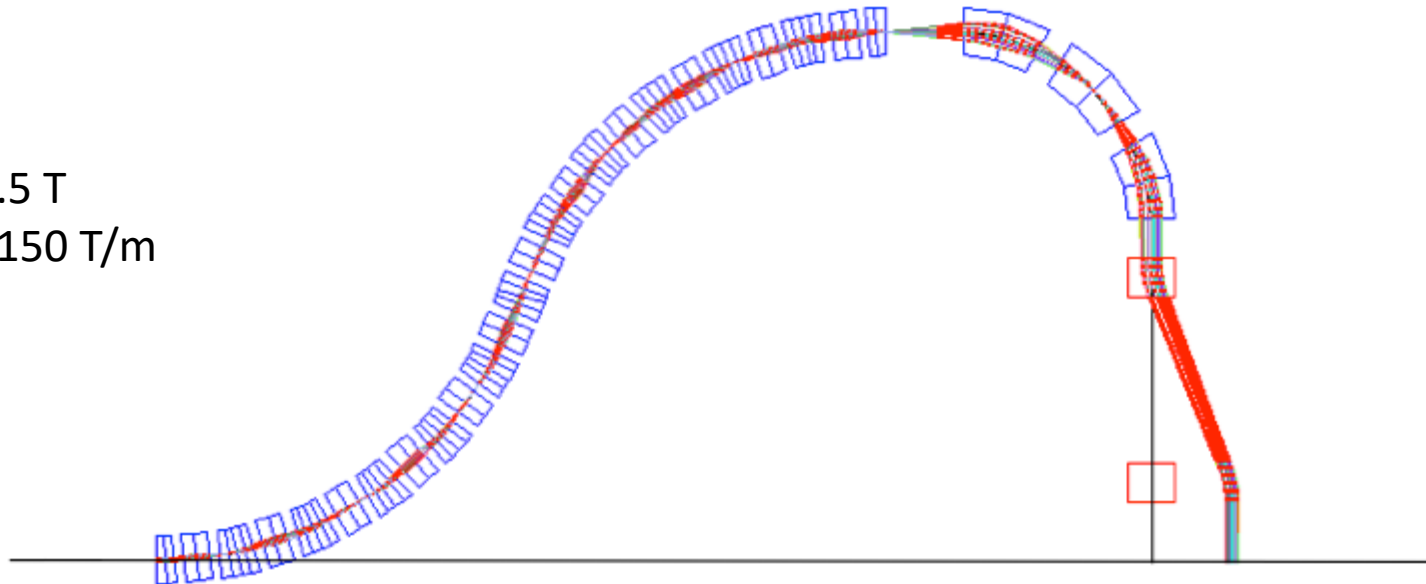
FFAG

What if dispersion is so small that $\Delta p/p = \pm 35\%$ goes through?

p 142 MeV
C 245 MeV

CARBON GANTRY height 4.091m

Bdip ~ 4.5 T
g up to 150 T/m



Plenty of other aspects to consider

Magnet misalignment and corrections

Mechanical structure deformations

Integration, floor, access to patient

Patient positioning

Position verification

“Range” verification, In room imaging

Dose delivery

Scanning speed

Maintenance

Safety and many other aspects to consider in gantry design...

...

Thank you for your attention

“Physics is like sex: sure, it may give some practical results, but that's not why we do it. ”

R. Feynmann

