

State of the art Superconducting Magnets

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Content

- Why magnetic fields – basic equations – accelerators and magnets, detectors, confinement
- Type of magnets – longitudinal, transverse, toroidal; example with electromagnets
- Superconductivity and accelerators; type of superconductors
- Superconducting magnets for accelerators (and beam lines)
- From Tevatron to LHC; example of LHC magnets
- Beyond LHC: new magnets for HiLumi LHC
- Beyond Hilumi: HFM, Nb3Sn, HTS,
- New technologies: NI coils

Motion of a charged particle in an electric and a magnetic field

$$F = q(\vec{v} \times \vec{B} + \vec{E})$$

\vec{E} Electric field accelerates (along the motion), i.e., gives momentum, energy

$\vec{v} \times \vec{B}$ Magnetic field act as **perfect centripetal force**, i.e., bend trajectories of charged particles, without changing the modulus of the momentum, i.e., the energy.

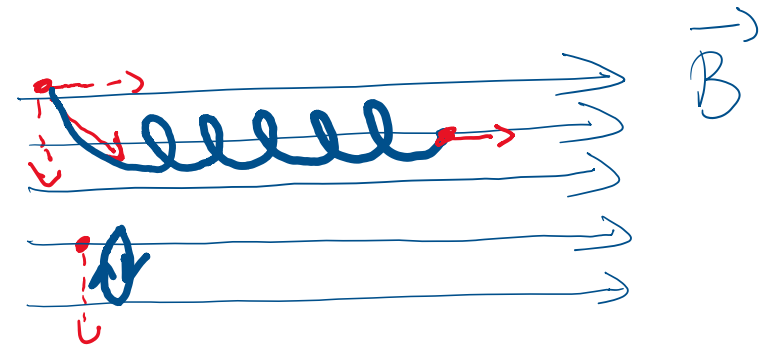
Considering B as the component perpendicular to the motion:

$$F_c = m \frac{v^2}{\rho} \quad F_m = q v B \quad F_m = F_c$$

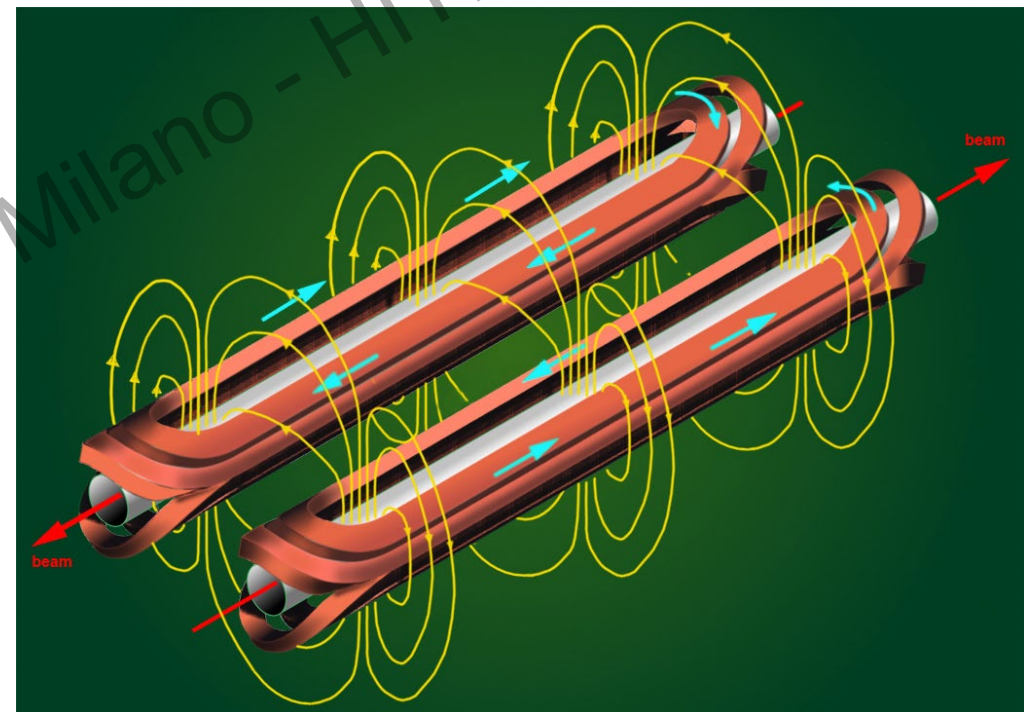
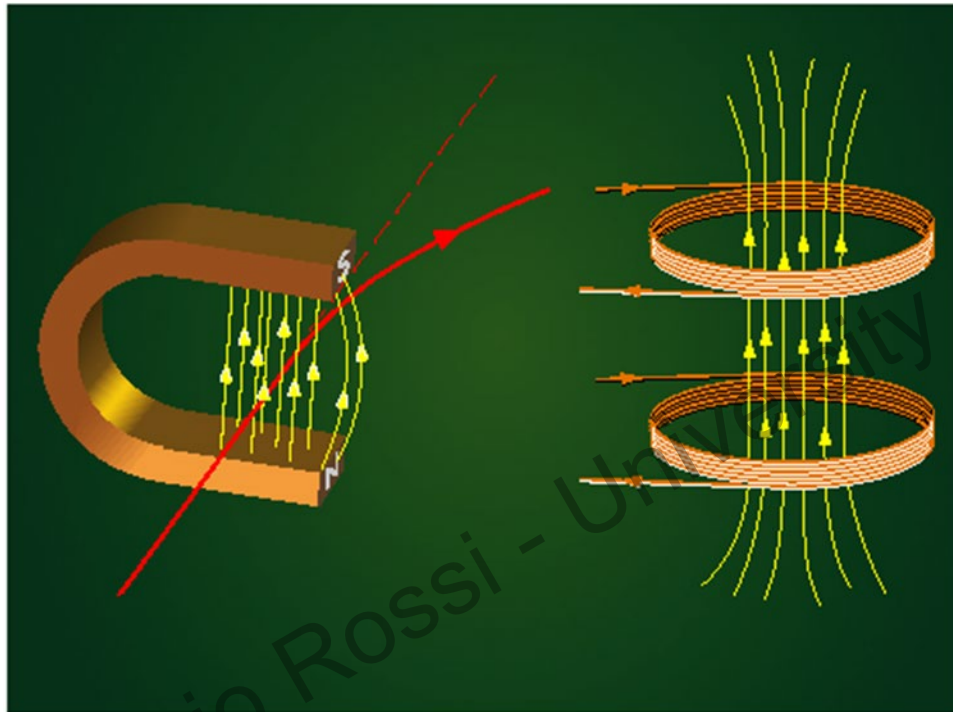
$$q v B = \frac{m v^2}{\rho} \Rightarrow B \rho = \frac{m v}{q}$$

$$B \rho = \frac{p}{q}$$

MAGNETIC RIGIDITY [T.m]



In accelerators we find two types of magnetic configuration
solenoidal fields to confine particle along trajectories or to bend particles trajectories in the midplane (in split coil configuration)
Transverse fields: long tube of fields perpendicular to trajectories (dipoles)



Field-energy expression for non-relativistic case: classical cyclotron

Particles curved by magnetic field use same voltage many times...

NON-relativist case

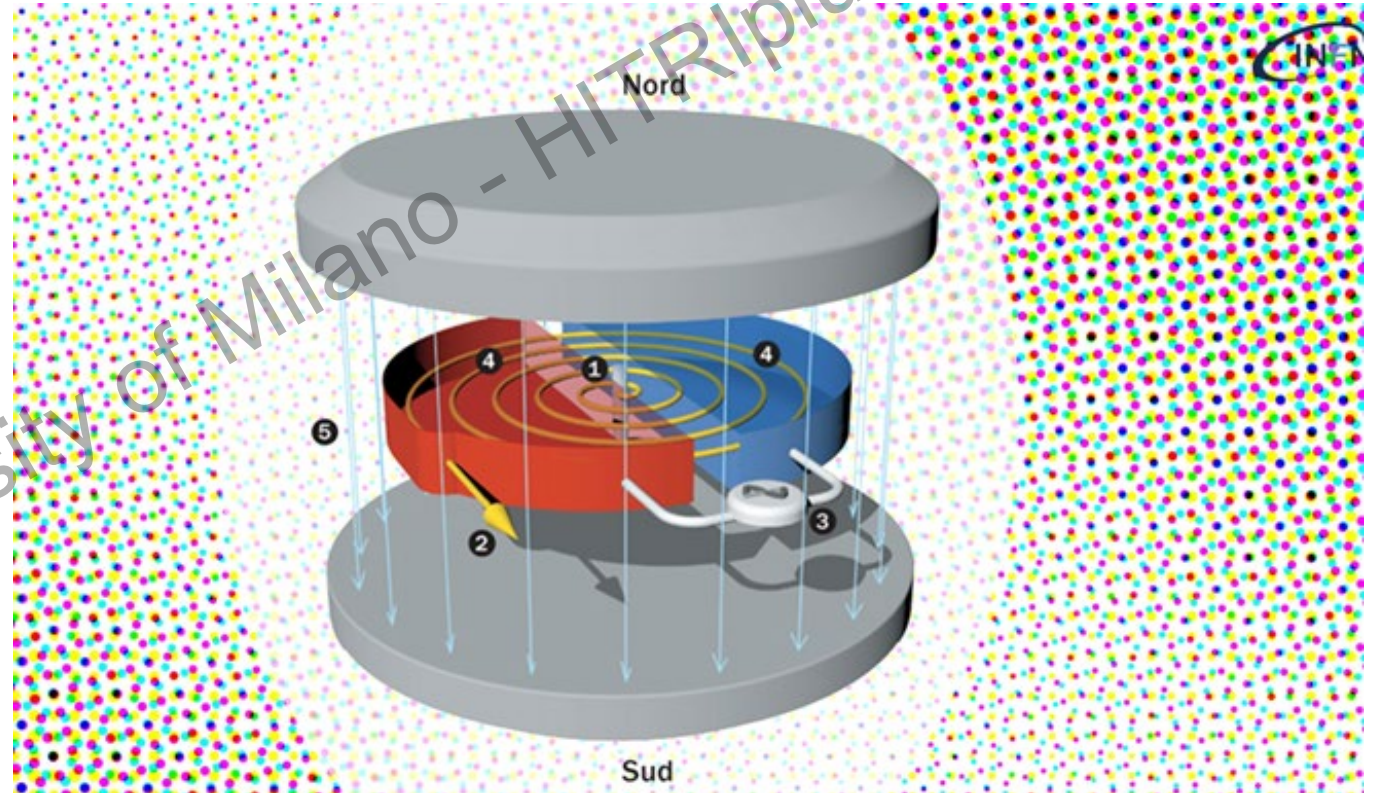
For proton or electron of charge e and mass m , kinetic energy E_k and field are related by:

$$E_k = \frac{1}{2} \frac{e}{m} (B\rho)^2$$

Where B is the magnetic field (perpend. To the motion) and ρ is the curvature radius.

For ions with charge Z and atomic no. A :

$$E_k/A = \frac{1}{2} \left(\frac{e}{m_p} \right) \left(\frac{Z}{A} \right)^2 (B\rho)^2$$



Relativistic case, most used for HEP synchrotron and colliders... and (more complex) general expression

c being the light velocity, the expression is (mass energy is negligible, $E_k \approx E$)

$$E = e c B \rho$$

Which makes the maximum energy in terms of practical units (GeV, tesla, m) very simple:

$$E_k \cong 0.3 B \rho$$

We here that beam energy scaled linearly with field and size of the accelerators.

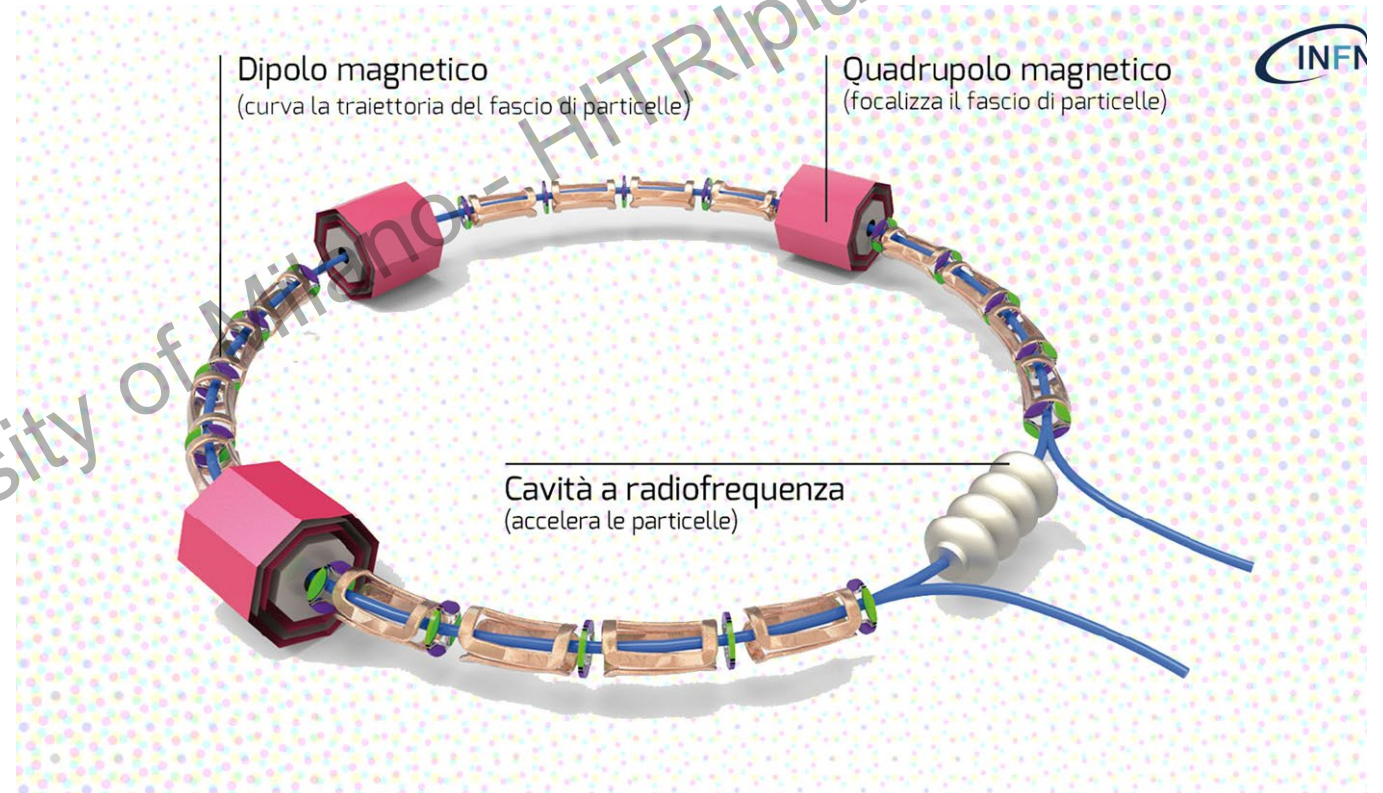
In intermediate energy we must use:

$$E^2 - E_0^2 = p^2 c^2$$

$$E^2 - E_0^2 = e^2 c^2 \rho^2 B^2$$

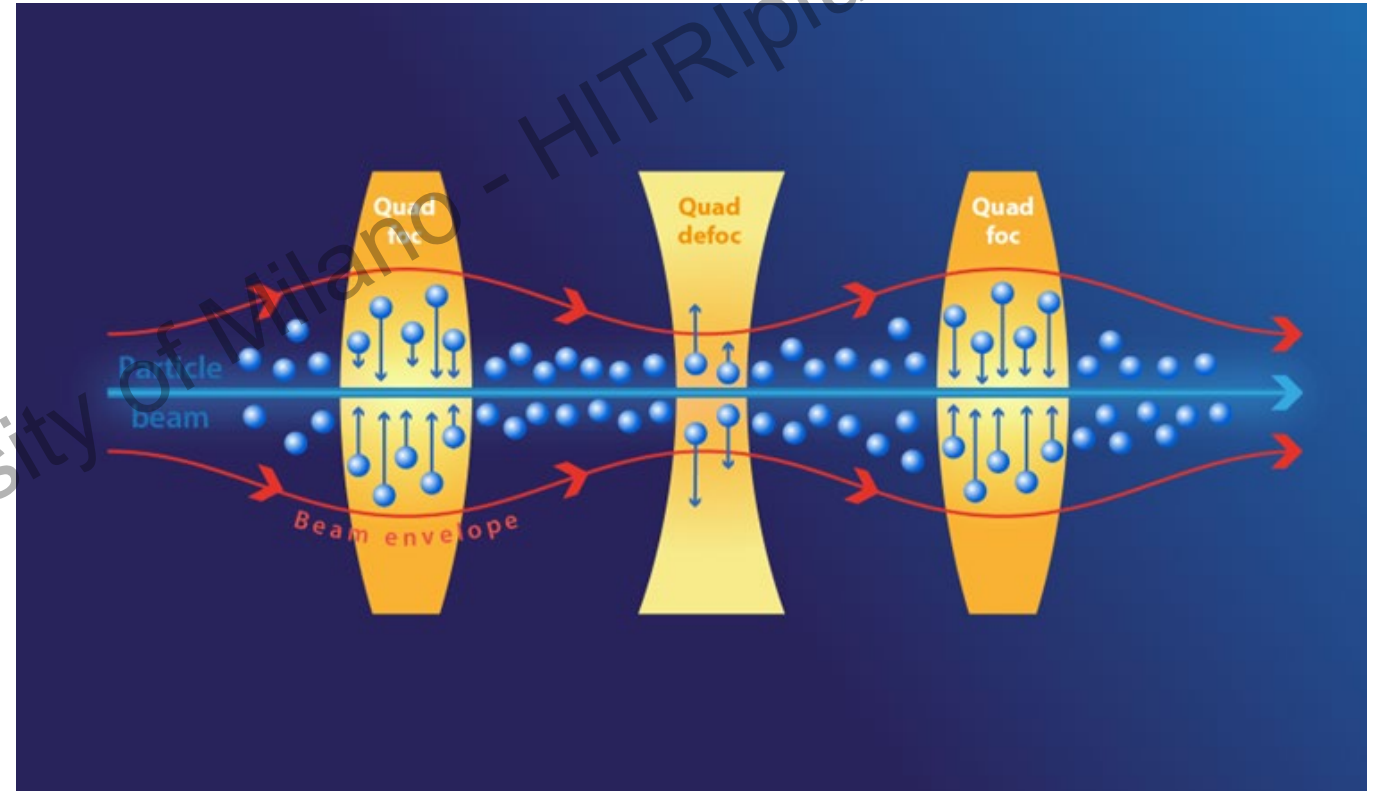
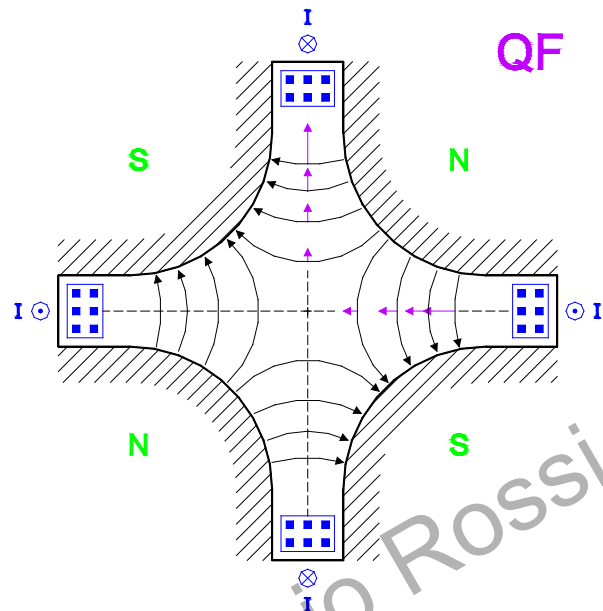
Again, dividing by e to have energy in eV:

$$E_k = \sqrt{E_0^2 + B^2 \rho^2 c^2} - E_0$$



Magnetic forces are necessary for focusing particle beams (quads) and higher order (sextupoles, octupoles, etc...) are needed, too.

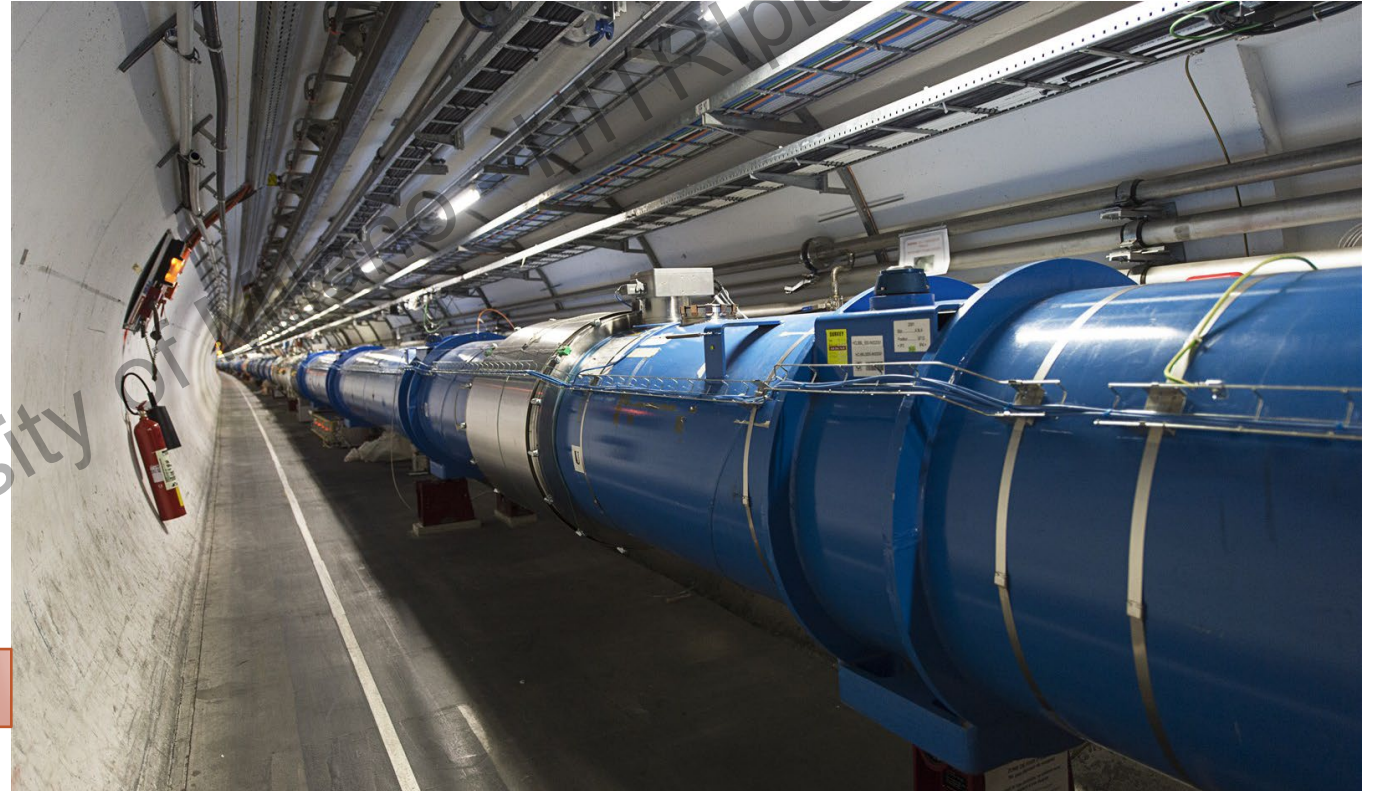
Quadrupoles are needed for focusing. When focusing in one plane (example: x-z) it defocuses in the other plane (y-z). So a doublet of FD or DF quads are needed; **total effect is focusing** →



Magnets are key components of any particle accelerator for circular accelerators fills most of the tunnel

About 80% of the 27 km long Large Hadron Collider of CERN, the superaccelerator that has allowed the discovery of the Higgs boson, is filled by a continuous lattice of 3 dipole (45 m) and one quadrupole – with HO magnets (about 5 m)

Series of dipoles in the LHC tunnel →



Comparison between needs for accelerator and detector magnets

- Magnetic fields needed for
 - electric charge identification
 - momentum spectrometry
 - $p = mv = q \rho B$; $\phi = q/p \quad B L$
 $\Rightarrow BL$ is often the comparison parameter
- If momentum analysis is done by tracking inside the field volume:
 - $\Delta p/p \propto 1/BL^2 \Rightarrow$ **large volume pays off better than high field**
 - Field homogeneity appreciated but NOT critical
(field knowledge of 0.1% usually suffices)



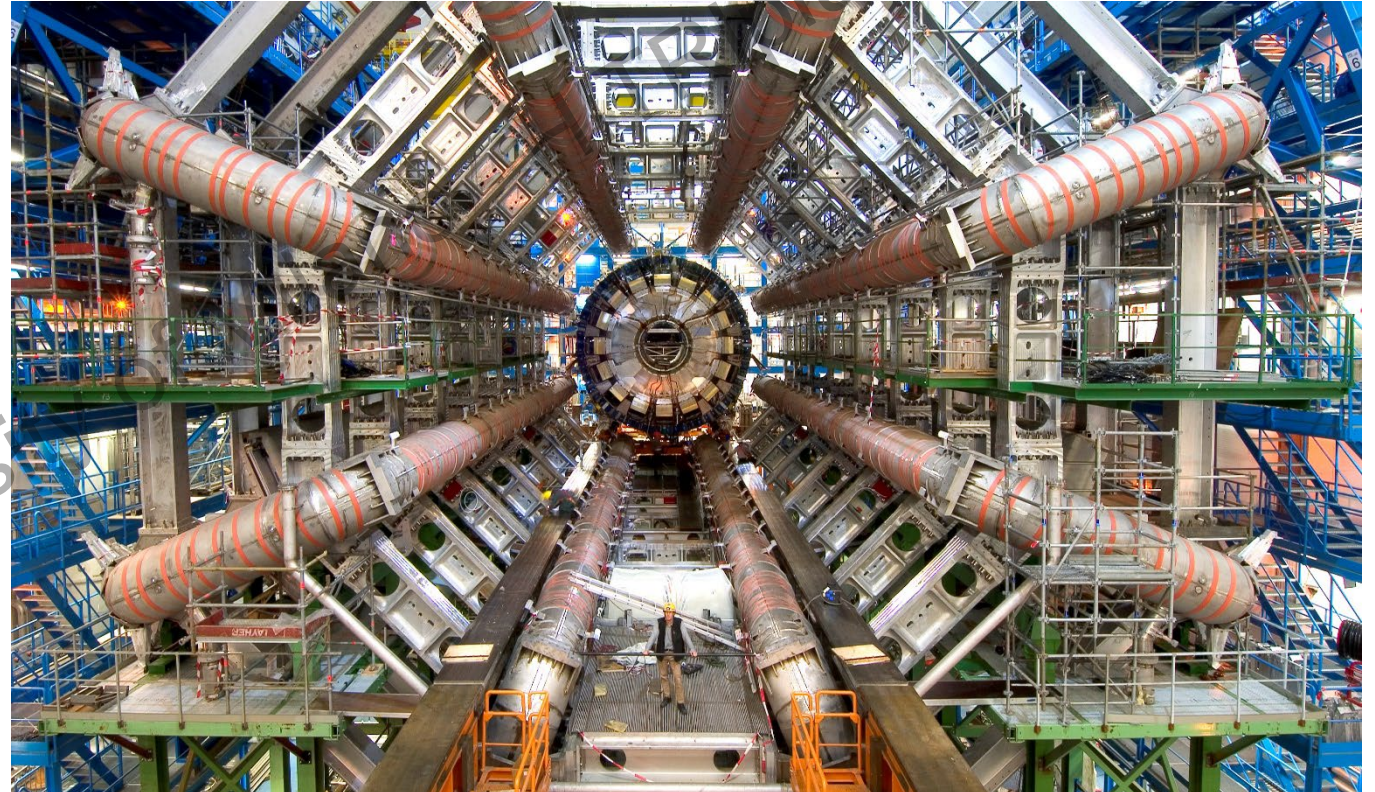
Toroids: a perfect transverse field with a heavy toll to efficiency

Toroidal field is always perpendicular to the particle trajectories, in any direction they go.

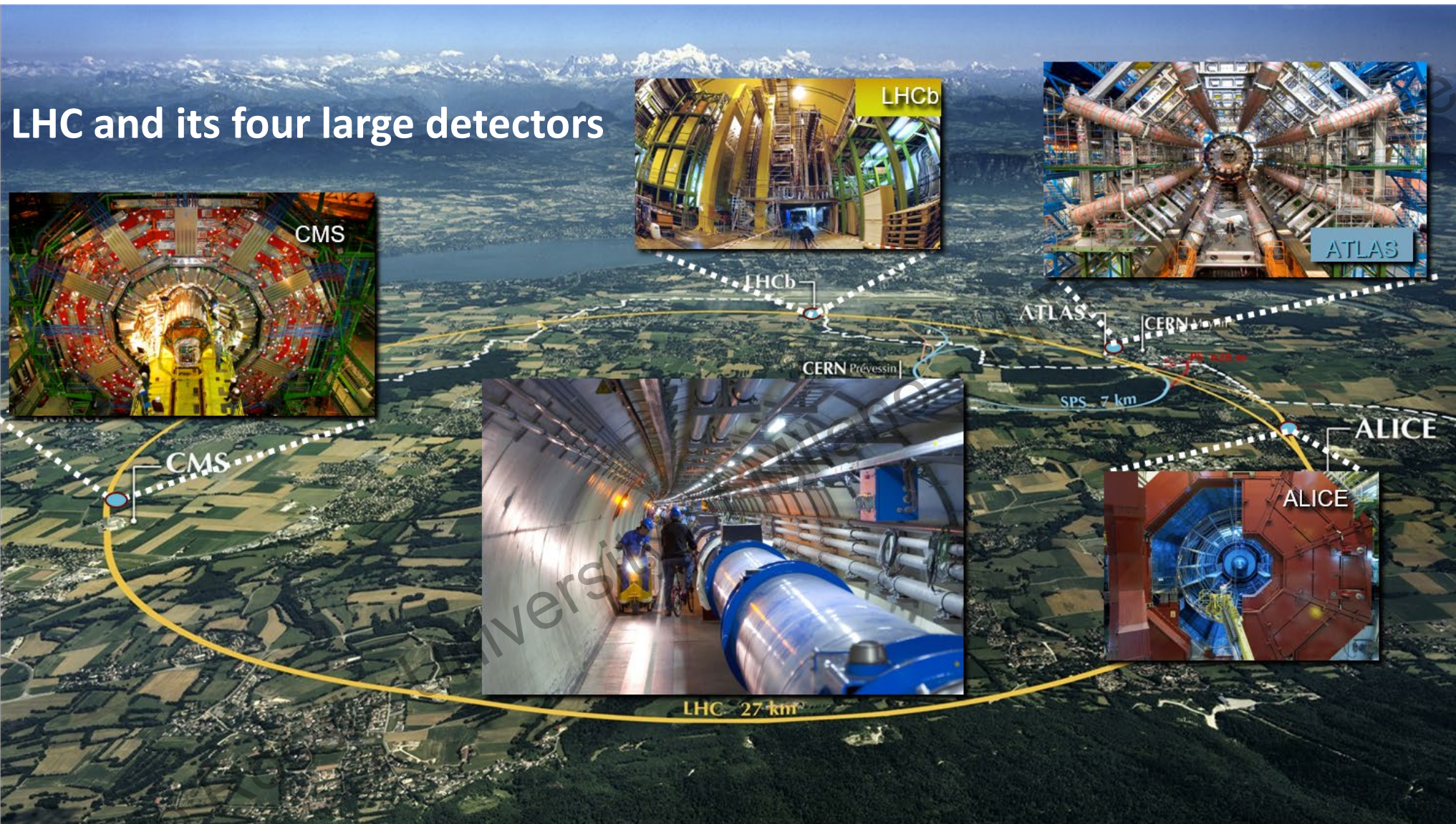
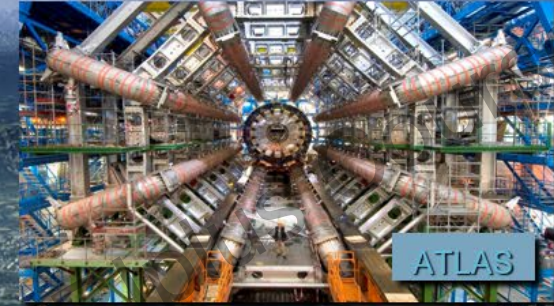
So they are used for spectroscopy → ATLAS

However the ratio between useful field and peak field on the coil is very bad. They are used only when really needed like detectors or magnetic confinement or when no magnetic moment is allowed (because of the closed field lines)

type	$B_{\text{useful}}/B_{\text{coil}}$
Solenoids	0.8-0.99
Dipoles	0.7-0.95
Toroids	0.25-0.6



LHC and its four large detectors



Field computation (no ferromagnetic material)

General formulation to compute field from current distribution

$$\mathbf{B} = \frac{\mu_0}{4\pi} I \int_c d\mathbf{l} \times \frac{\mathbf{r}}{r^3} dl$$

$$\mathbf{B} = \frac{\mu_0}{4\pi} \int_V \mathbf{J} \times \frac{\mathbf{r}}{r^3} dv$$

for infinitely long wire

$$\mathbf{B} = \frac{\mu_0}{2\pi} \frac{I}{d} \mathbf{e}_\theta$$

for finite straight wire, I
flowing from \mathbf{r}_1 to \mathbf{r}_2 (from
Wilson book)

$$\mathbf{B} = \frac{\mu_0}{4\pi} I \frac{(\mathbf{r}_1 \times \mathbf{r}_2)}{r_1 r_2} \frac{r_1 + r_2}{r_1 r_2 + \mathbf{r}_1 \cdot \mathbf{r}_2}$$



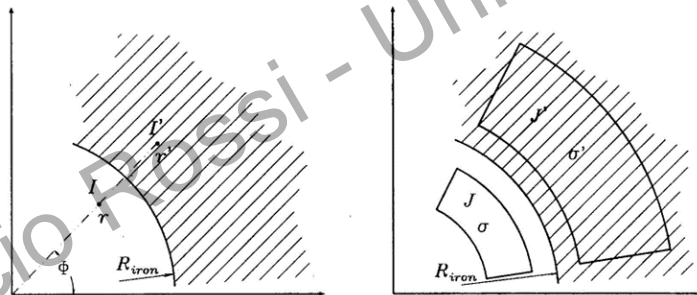
If there us iron → use FE codes Or a quick evaluation by IMAGE CURRENT methods

Use of iron yoke to contain the stray field is very common. Then image current methods (magnetic mirror) could be used, either with infinite (first guess) or finite μ_r .

- ❑ The method is actually applicable when:
 μ_r is constant and known

$$I_{im} = I \frac{\mu_r - 1}{\mu_r + 1}$$

- ❑ the iron has a shape that respects the symmetry of the problem (infinite iron plane, or circular iron fully surrounding the current region), see picture below.



Tip:

If iron is fully saturated, and the magnetization is known, the field is the sum of air core coil current plus the magnetization surface currents, $\mathbf{J}_s = \mathbf{M} \times \mathbf{n}$

Solenoids – field computation

J_s = linear density of the surface current (A/m)

J = current density, t = coil thickness

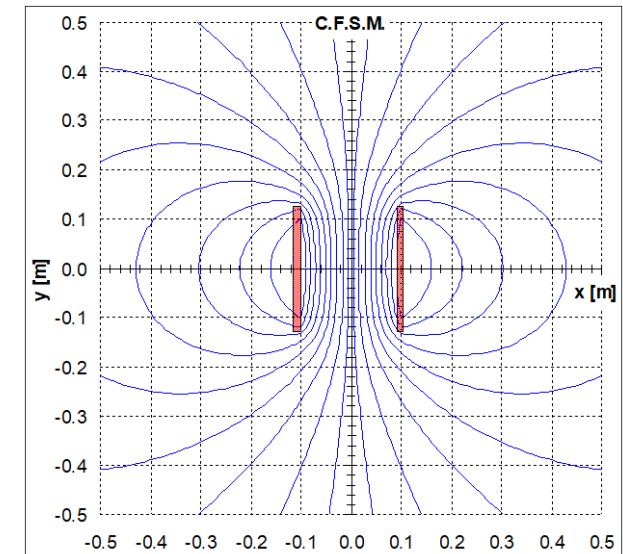
IN/L = ampere-turns (total current) per unit length

$$B = \mu_0 J_s = \mu_0 J \cdot t = \mu_0 I \frac{N}{L}$$

Most solenoids are not so far from this very easy computation.

When $L/\varnothing = 1.25$ the central field is about 15% less than the field infinite solenoid (with the same peak field on the coil).

Note that whatever long is the solenoid, at the ends B_{axis} is about half B_{central} and that inevitably there is a large radial field (this is important for HTS tapes)



Solenoids – computation of field in any point

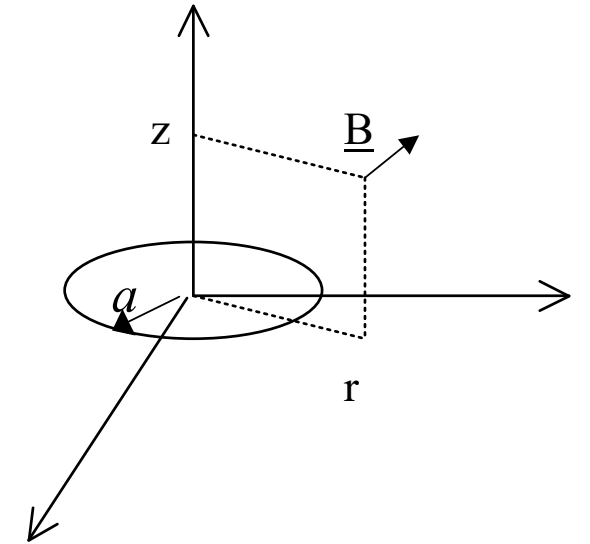
The formula of the loop can be analytically integrated only along the axis

$$\mathbf{B} = \frac{\mu_0}{2} I \frac{a^2}{(a^2 + z^2)^{3/2}}$$

In a generic point (r,z) the field components are:

$$\begin{cases} B_z(r, z) = \frac{\mu_0 I}{2\pi} \frac{1}{\sqrt{(r+a)^2 + z^2}} \left(J_1 + \frac{a^2 - r^2 - z^2}{(r-a)^2 + z^2} J_2 \right) \\ B_r(r, z) = \frac{\mu_0 I}{2\pi} \frac{z}{r\sqrt{(r+a)^2 + z^2}} \left(-J_1 + \frac{a^2 + r^2 + z^2}{(r-a)^2 + z^2} J_2 \right) \end{cases}$$

$$J_1 = \int_0^{\pi/2} \frac{1}{\sqrt{1 - k^2 \sin^2 \varphi}} d\varphi, \quad J_2 = \int_0^{\pi/2} d\varphi \sqrt{1 - k^2 \sin^2 \varphi} \quad ; \quad k = \sqrt{\frac{4R_c r}{(r + R_c)^2 + z^2}}$$

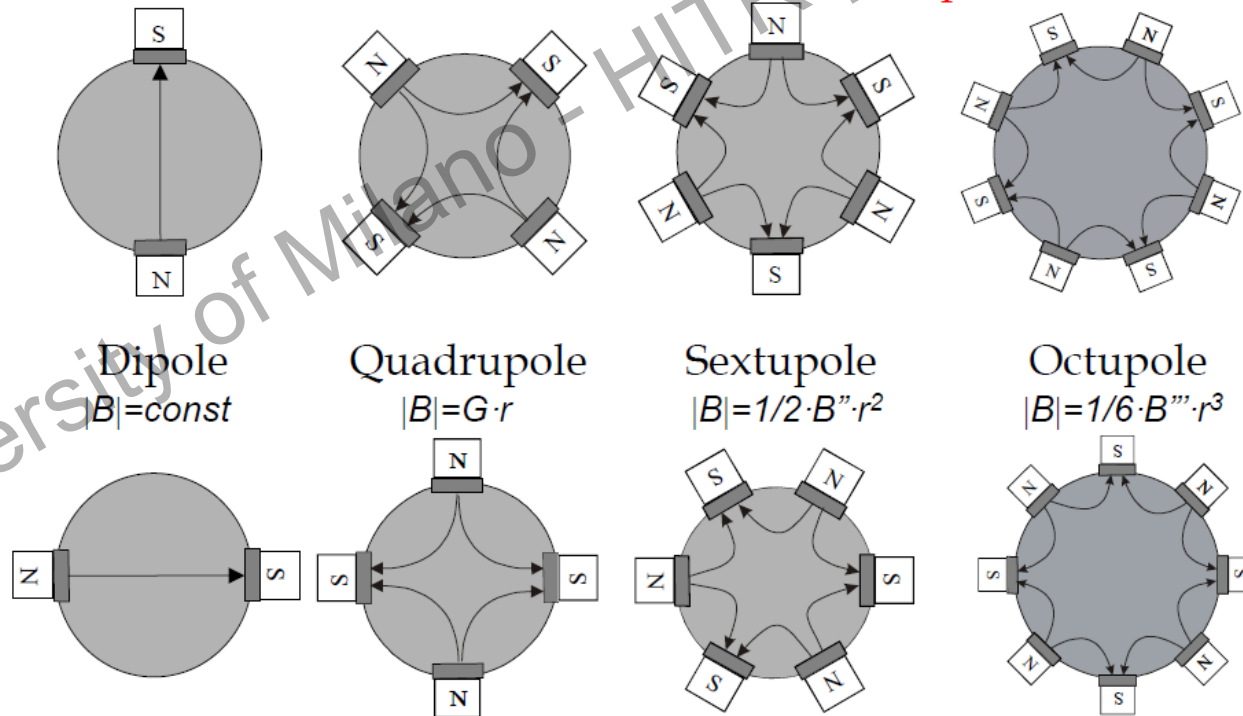


Transverse fields: any field can be represented as a sum of various multipoles

Courtesy of D. Tommasini, CERN, CAS Varna, 2010

Key point:
A dipole magnet generate "only" dipole field (the other components are inevitable errors... usually very small of the order of 10^{-5} to 10^{-2})

NORMAL : vertical field on mid-plane

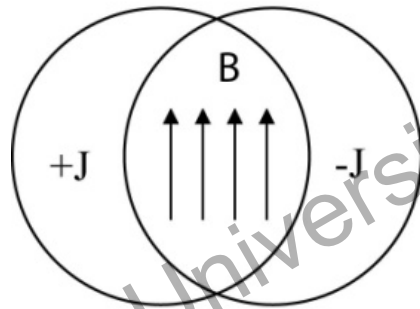


SKEW : horizontal field on mid-plane

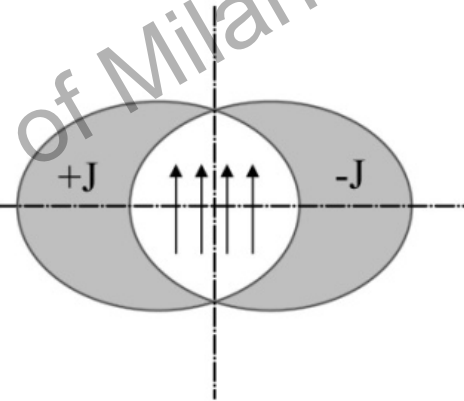
How to generate multipoles field with currents: overlapping cylinders or overlapping ellipses

- Field shape DIPOLE
 - Half the field of a solenoid for same J and coil thickness
 - $J = J_0 \cos\vartheta$
- Field shape QUADRUPOLES
 - $J = J_0 \cos 2\vartheta$

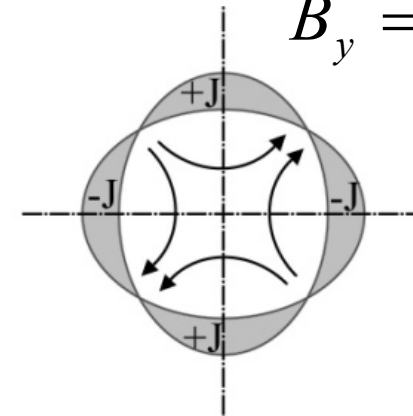
Field of a long solenoid with N turns and length L
 $B = \mu_0 NI/L$ or $B = \mu_0 J d$
 So, a dipole of same J and same coil thickness d gives half field of a solenoid (but is **transverse field**)



$$\mathbf{B} = \frac{\mu_0 J d}{2} \mathbf{e}_y$$



$$B_y = \frac{\mu_0 J b d}{a + b}$$



a, b ellipses parameters

$$B_x = \frac{\mu_0 J (a - b)}{a + b} y$$

$$B_y = \frac{\mu_0 J (a - b)}{a + b} x$$

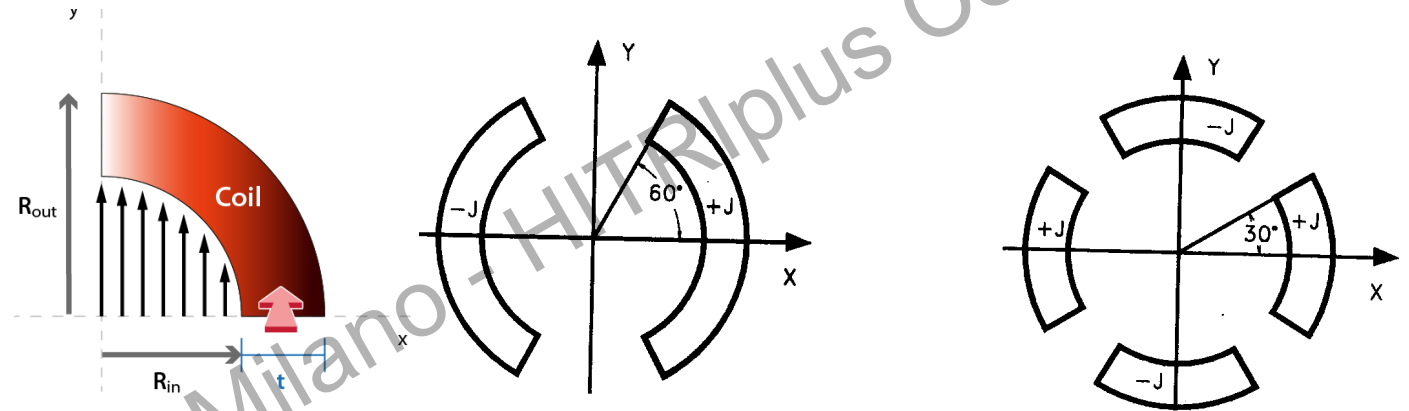
If the d (coil thickness) is constant ? (easy to manufacture) one can vary the J as $\cos\theta$

Constant internal radius, sector coils:

$$J_s = J \cdot d \cdot \cos(\theta), J_s = J \cdot d \cdot \cos(2\theta), \dots$$

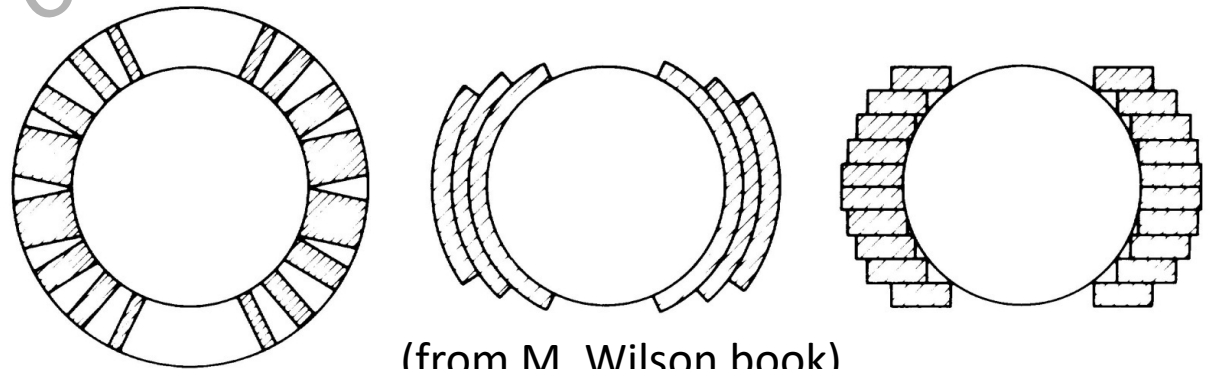
$d = t = \text{coil thickness}$

These generate perfect dipole
 (quadrupole) field inside a circle

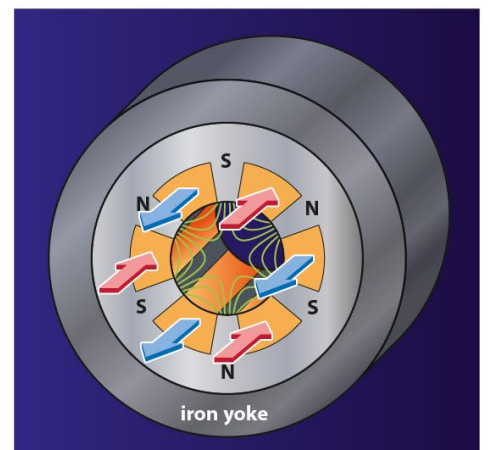
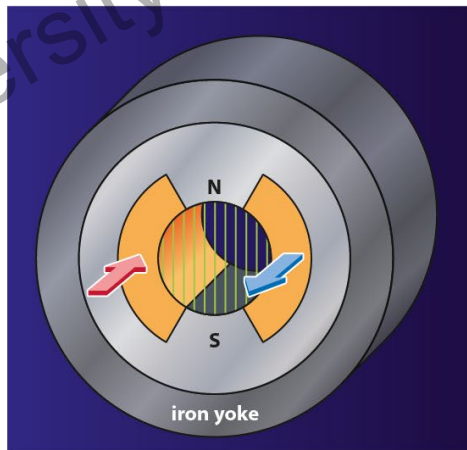
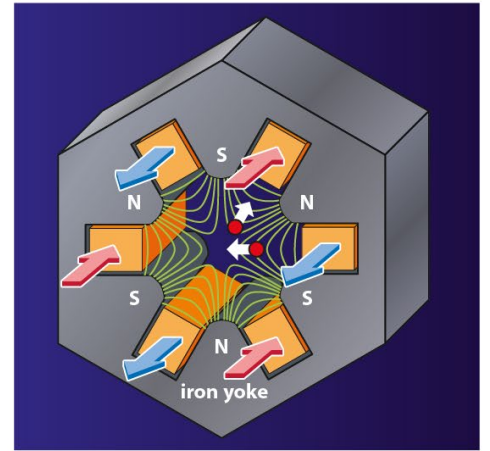
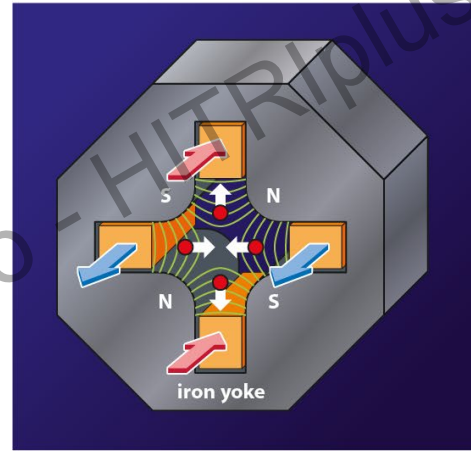
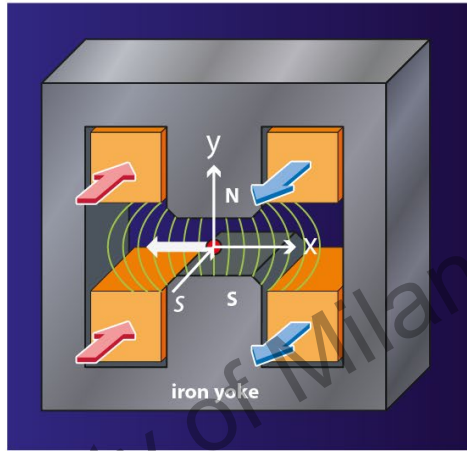
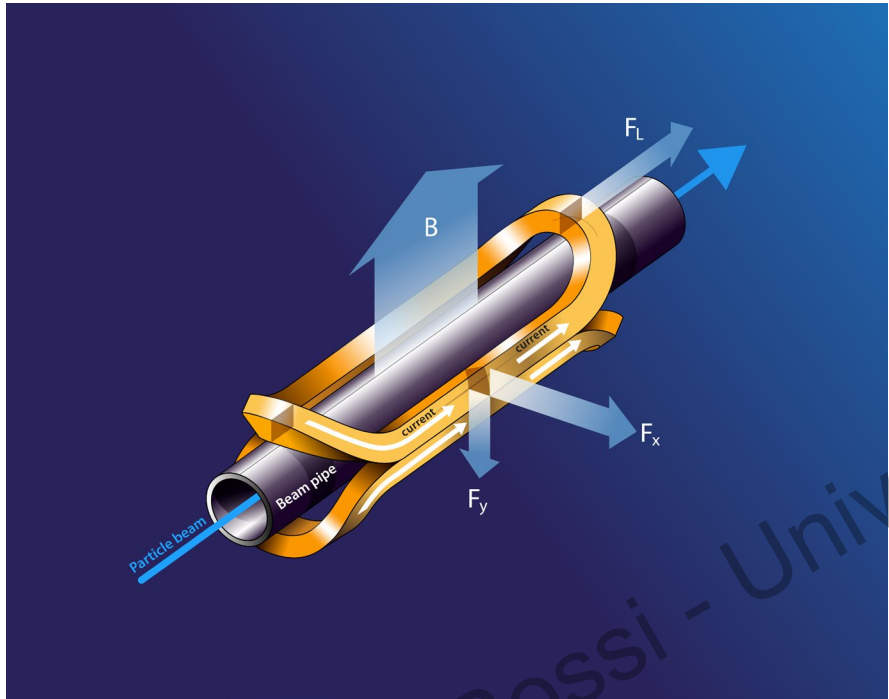


Approximation of $\cos\theta$ with coil blocks
 (left) and multiple shells (centre) and
 of intersecting ellipses.

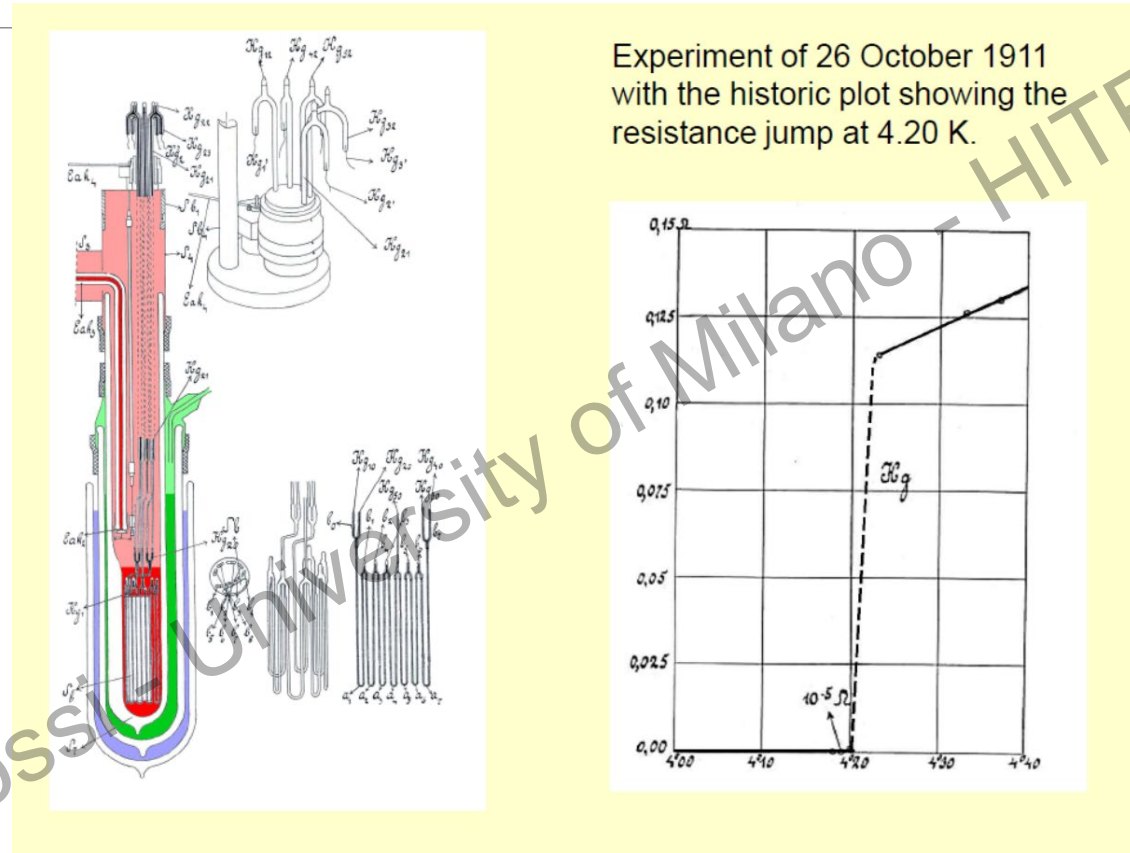
DEVIATION FROM PERFECT FIELD:
 10^{-4} (0.01%) that is taken as the **unit**.



Again: a dipole 3D and 2D scheme of dipole, quadrupole, sextupole, both for iron and coil dominated



But how to get much current? Superconductivity... but at low T



Why we need superconductivity

Superconducting LHC



Tunnel: 27 km

Field : 8.3 T

Cryoplant power at the plug: 40 MW: **always on**

~ 70 MW for LHC. 150 MW for the accelerator complex

180 MW for the whole CERN complex

Normalconducting LHC



Tunnel 120 km

Field : 1.8 T

Dissipated power at collision: ~ 2,200 MW

Average power (0.4 coefficient): 900 MW only for accelerator

However, power is required only when operating while power in SC magnets is needed to keep it cold... a factor 2 favorable in energy for NC

Zero resistance: really?

Gallop experiment on the current decay in a superconducting loop:

$$\rho < 10^{-26} \Omega\text{m}$$

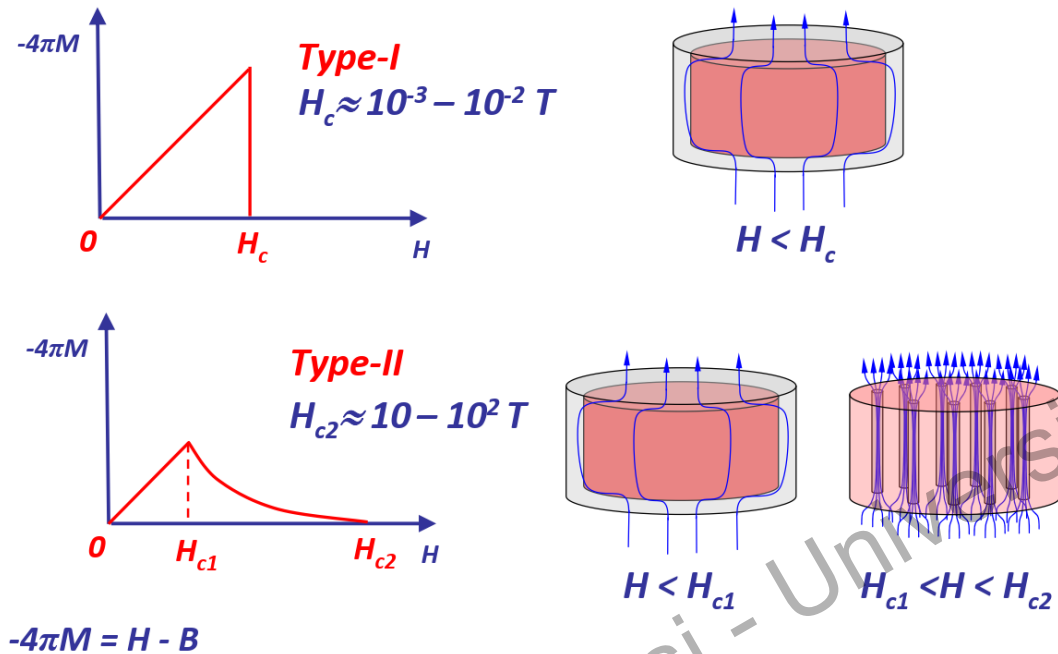
One experiment lasted two years, with no sign of current decay and had to stop because the supply of LHe was interrupted by transport strike !

Resistivity table

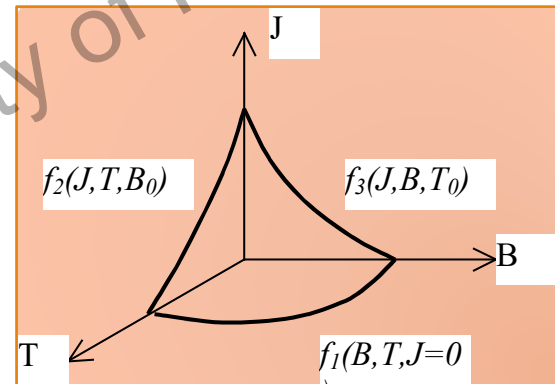
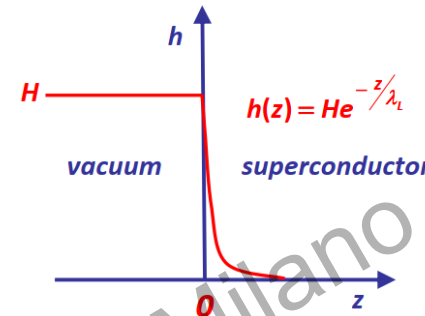
Materials	ρ (Ωm)
↓	↓
Vacuum	∞
Insulators	$10^{20} \div 10^{10}$
Semiconductors	$10^5 \div 10^{-3}$
Metals	$10^{-5} \div 10^{-10}$
Superconductors	≈ 0

Superconductivity: more than zero resistance

Type-I and Type-II superconductors



Meissner-Ochsenfeld effect



Superconductivity is a thermodynamic state defined by T, B, J similar to T, P, V for the ideal gas: $P \Leftrightarrow B, V \Leftrightarrow J$

Superconductivity exists only below the critical surface: increase one parameters depress possible value of others...

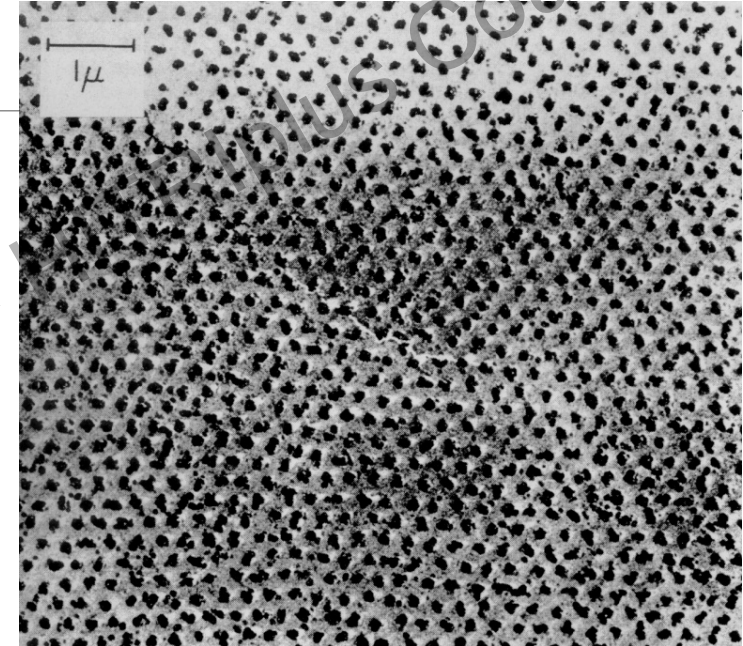
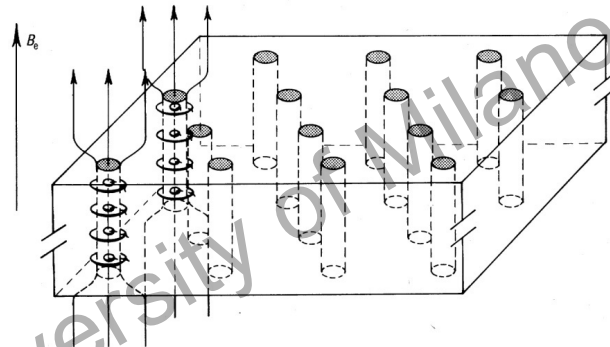
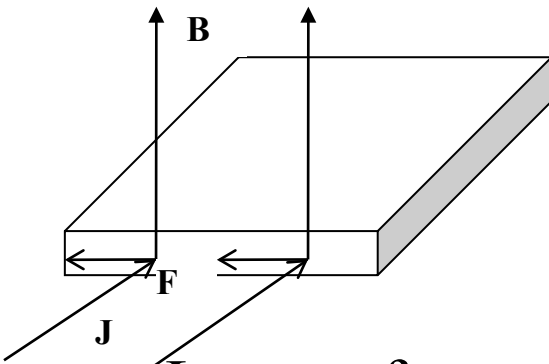
B and J are linked by consideration of free-energy and by Maxwell equations... So inevitably an increase in current brings also an increase in field...

Deception! B_c is very low: $\sim 1-10 \text{ mT}$!

Fluxoids or quantum magnetic flux

Flux penetration **in the material**
is in quanta:

$$\Phi = h/2e \cong 2 \cdot 10^{-15} \text{ Wb}$$



Lorentz force : $F_p = -J_c \times B$: to avoid movements and heating it is needed a **pinning**
given by defects.

NbTi: $F_{p \text{ max}} \approx 15 \text{ GN/m}^3$ (or 15 N/mm^3 !!) $\Rightarrow J_c \approx 3 \text{ GA/m}^2$ (3000 A/mm^2) at 5 T

Stability margins...

Superconductors are NOT stable against perturbation albeit very small. ΔE of μJ are enough to drive superconductor normal!

Heat capacity drops at low temperature ($T \ll T_{\text{Debye}}$):

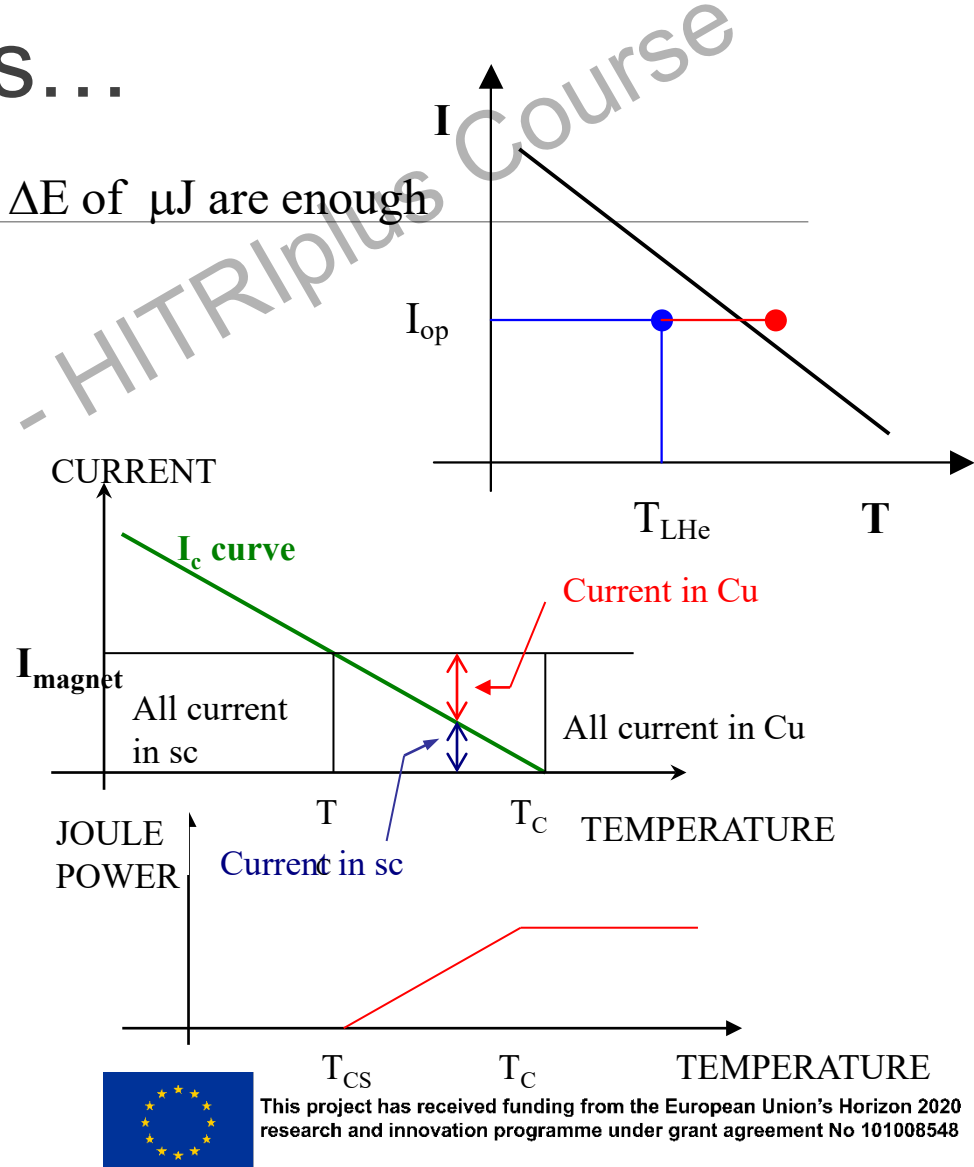
$C \propto T^3 \Rightarrow \Delta T = \Delta E / \gamma C$. So even small ΔE generates sensible ΔT

\Rightarrow operating point of the magnet beyond critical surface \Rightarrow **QUENCH**

Electrodynamic stability: intimate contact between the superconductor and a good conductivity material.

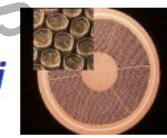
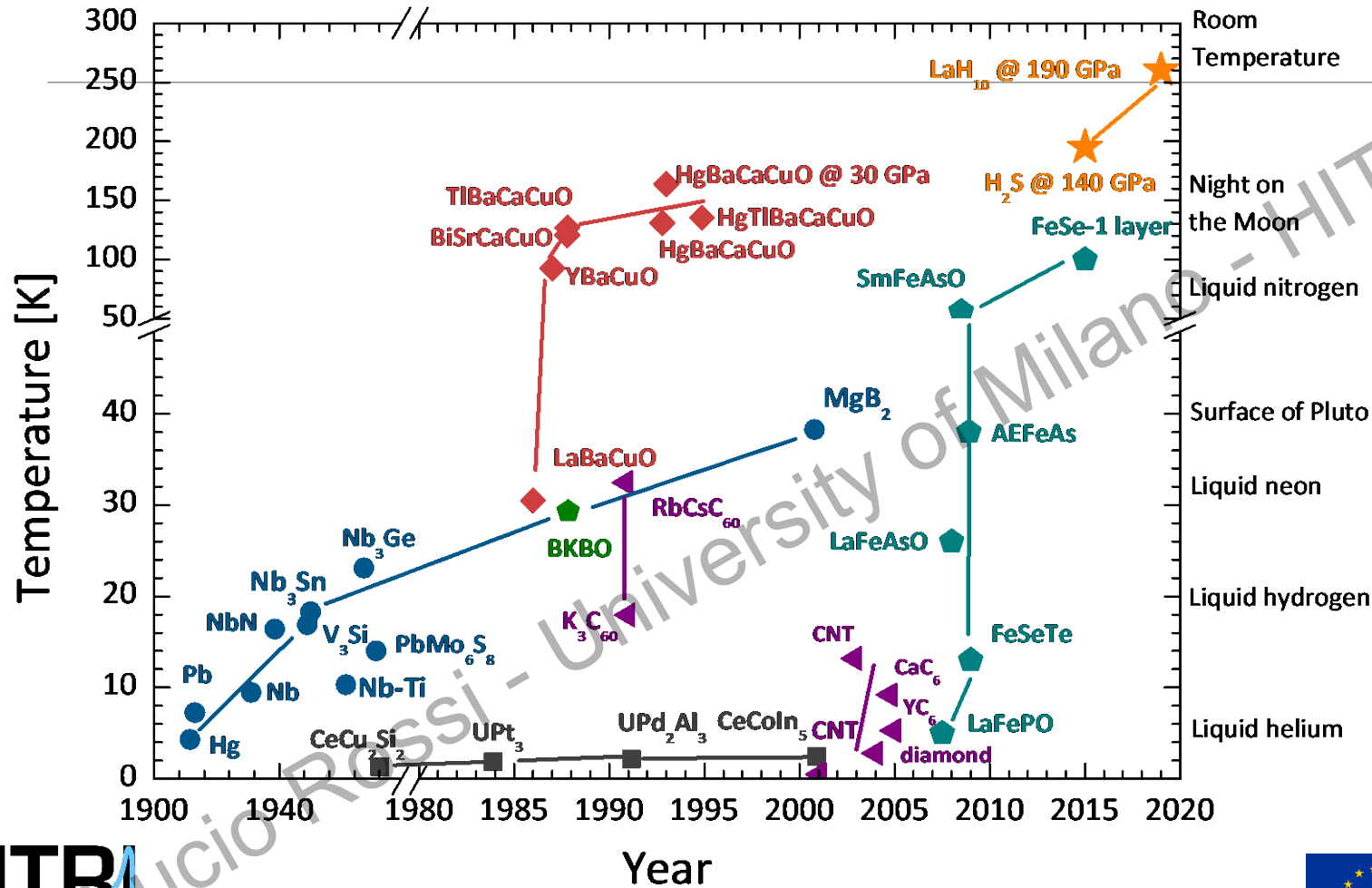
Adiabatic (or intrinsic) stability: to cure the flux rearrangement that generates heat

Direct cooling: LHe and more HEII are very good coolant, capable to remove heating in milliseconds! Latent heat 10-1000 times that of solid specific heat!



This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

Critical temperature vs time...



NbTi



Nb₃Sn



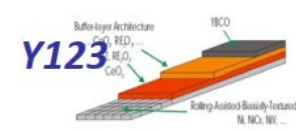
MgB₂



Bi2223



Bi2212



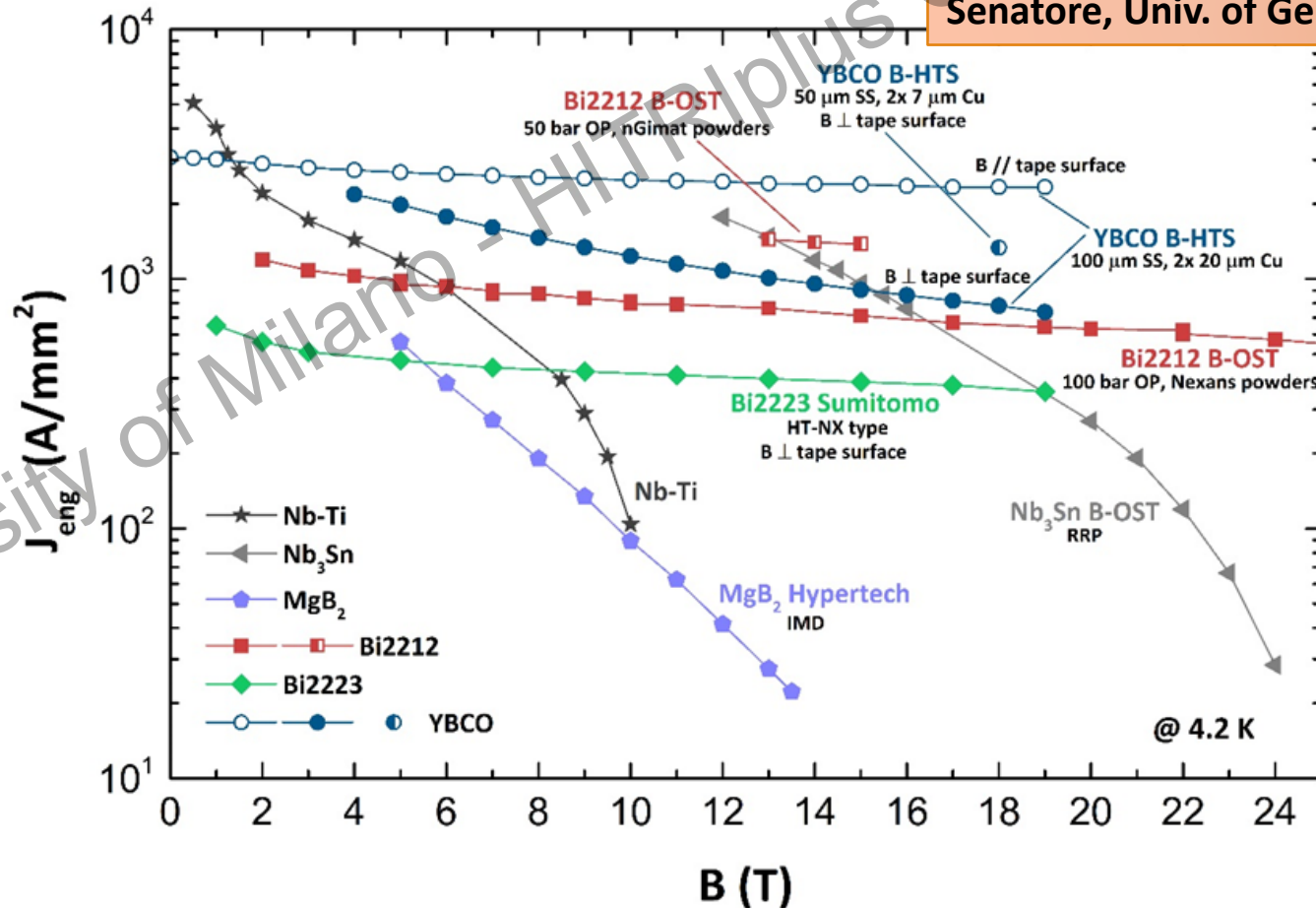
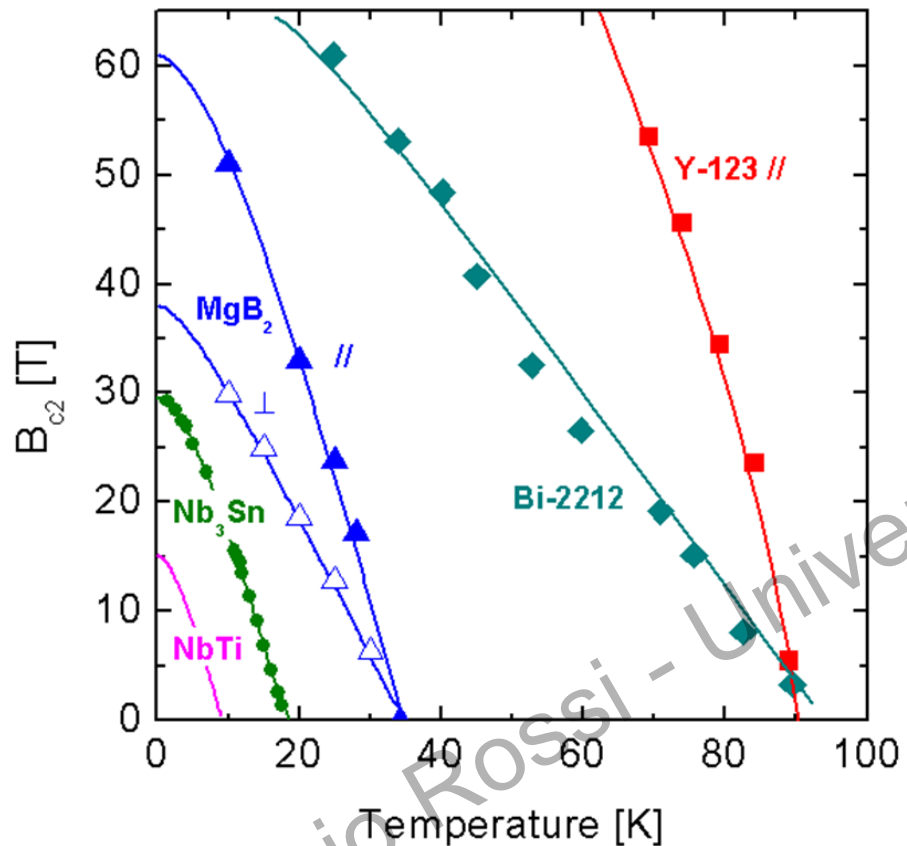
Y123

Thanks to C. Senatore, Univ of Geneva

Critical field vs. temperature (at zero current)

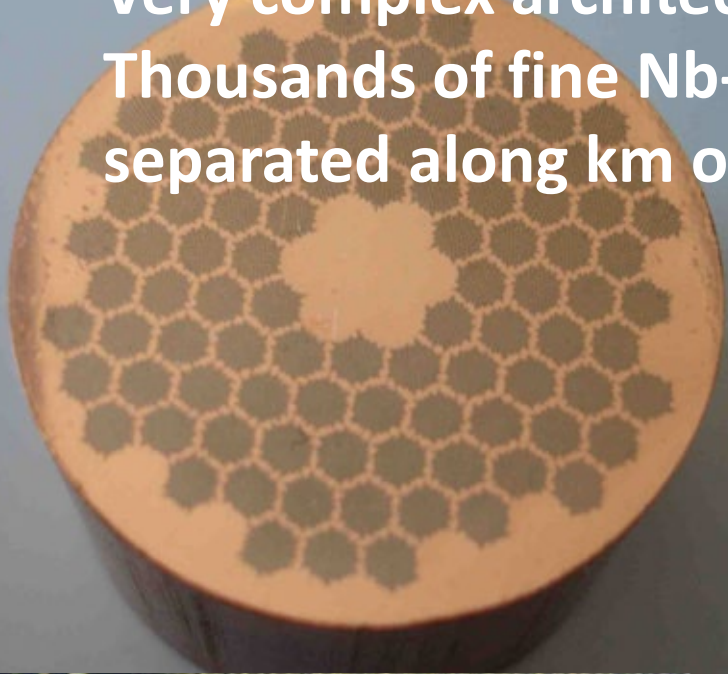
Critical current density (eng.) vs field at fixed T

Plot courtesy of C. Senatore, Univ. of Geneva

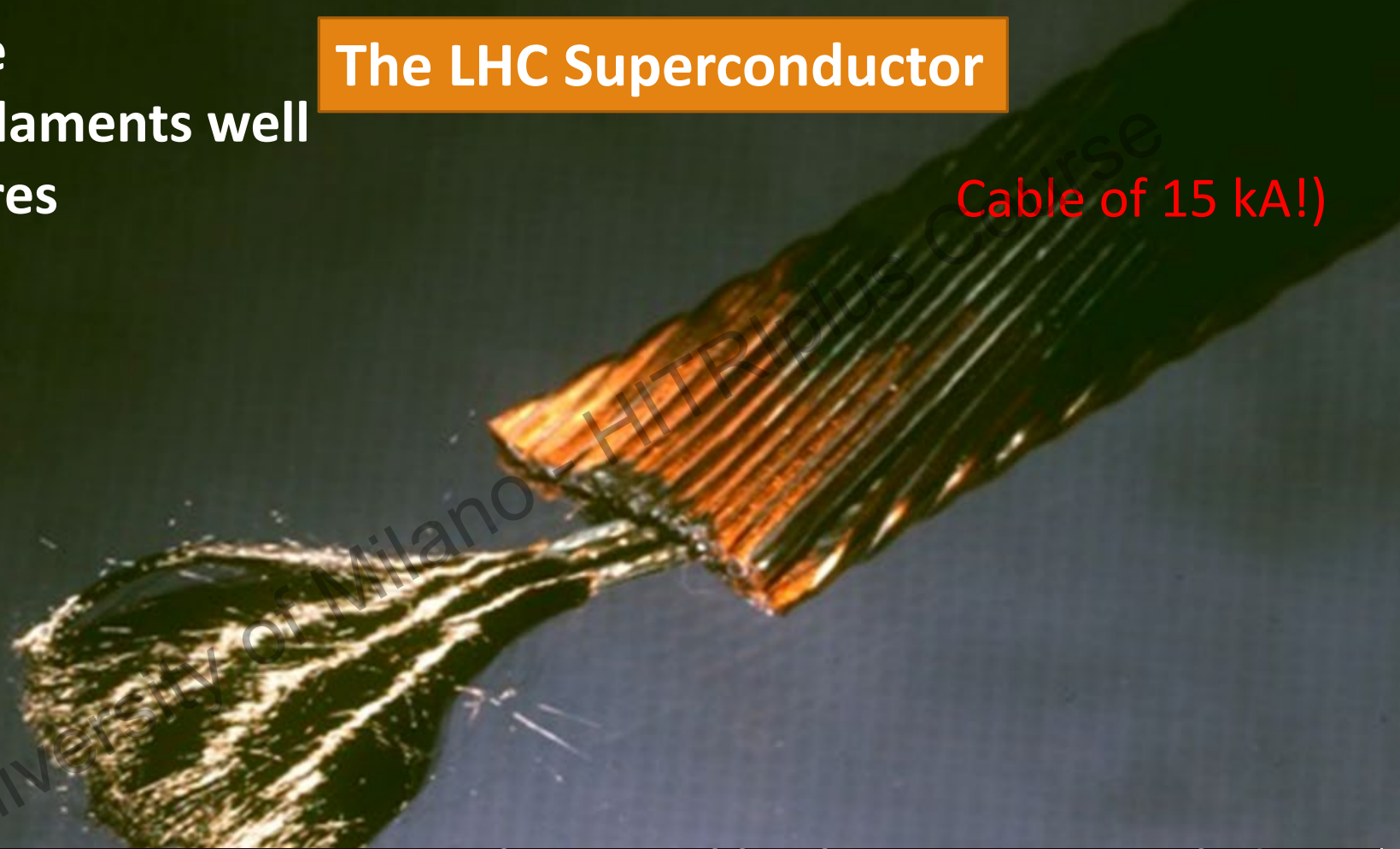


The LHC Superconductor

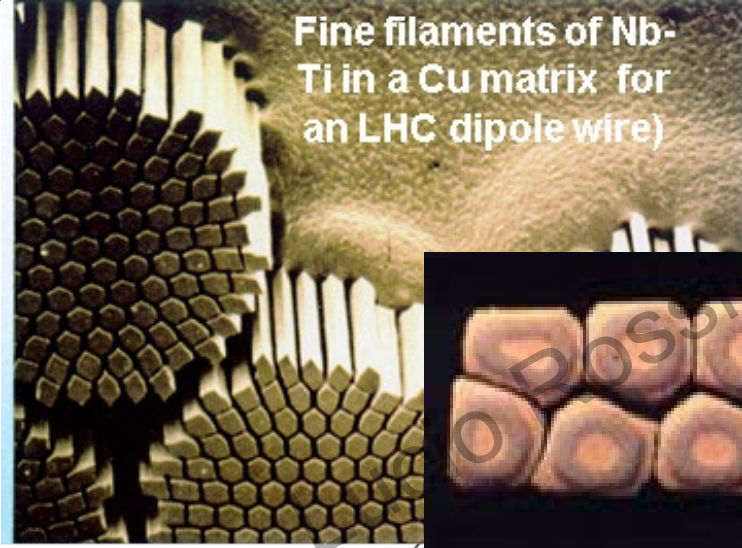
Very complex architecture
Thousands of fine Nb-Ti filaments well separated along km of wires



Cable of 15 kA!



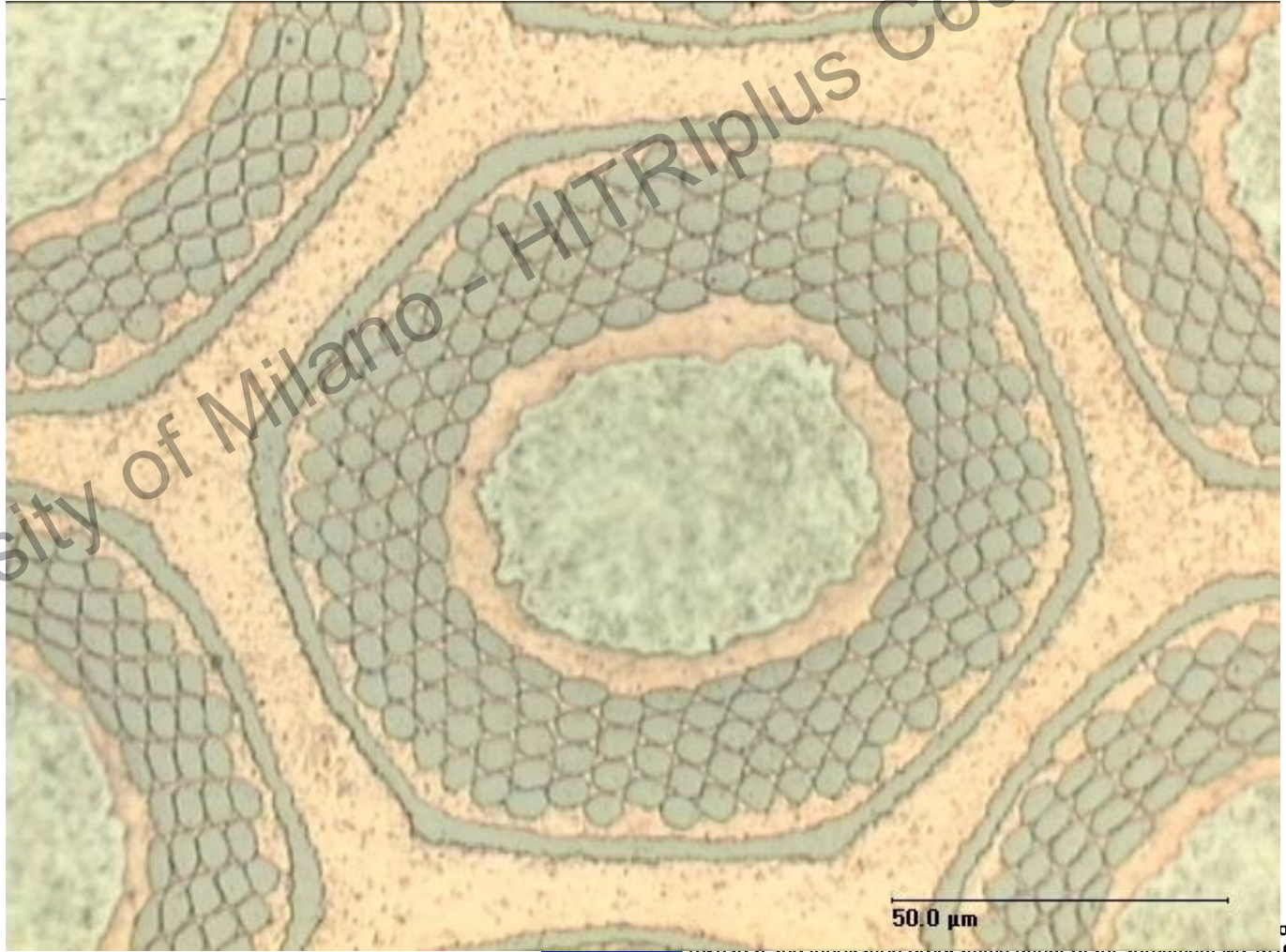
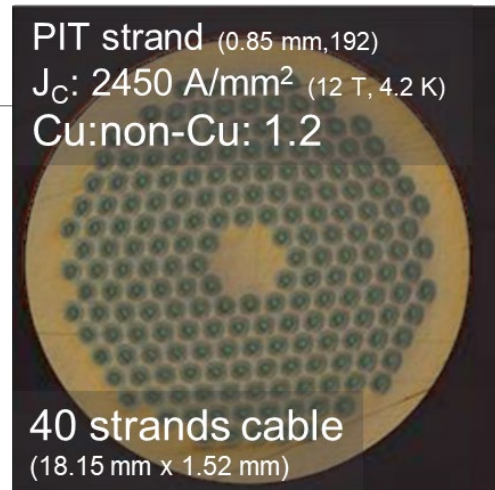
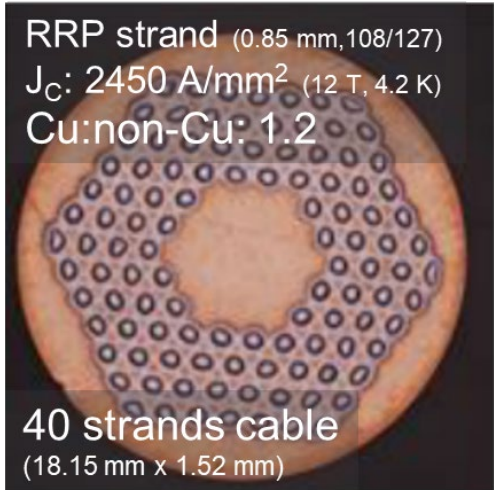
Fine filaments of Nb-Ti in a Cu matrix for an LHC dipole wire)



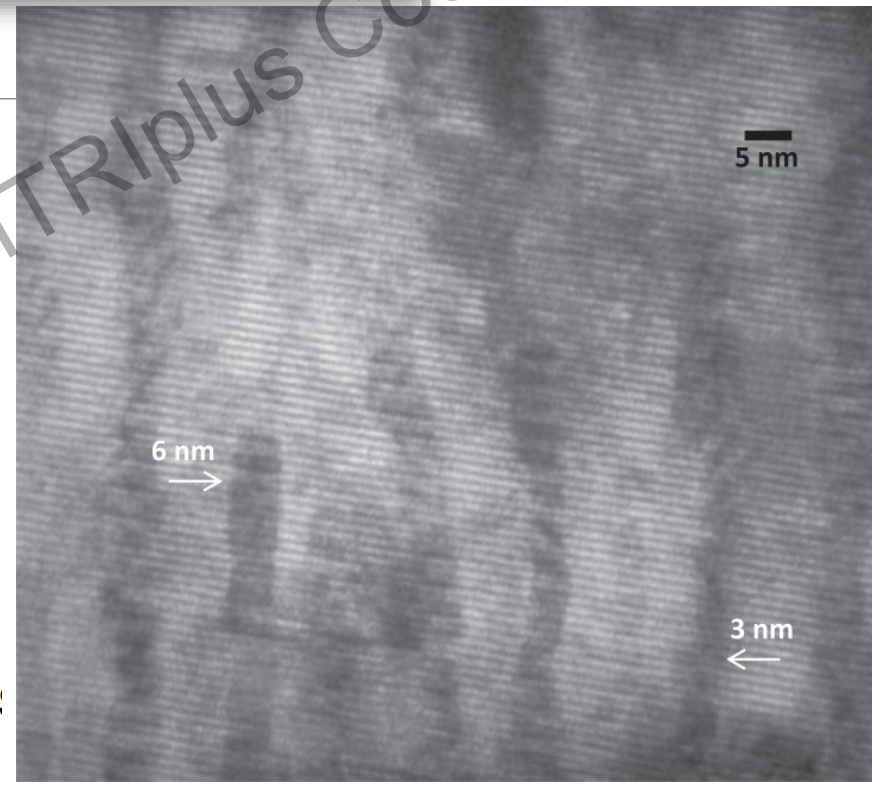
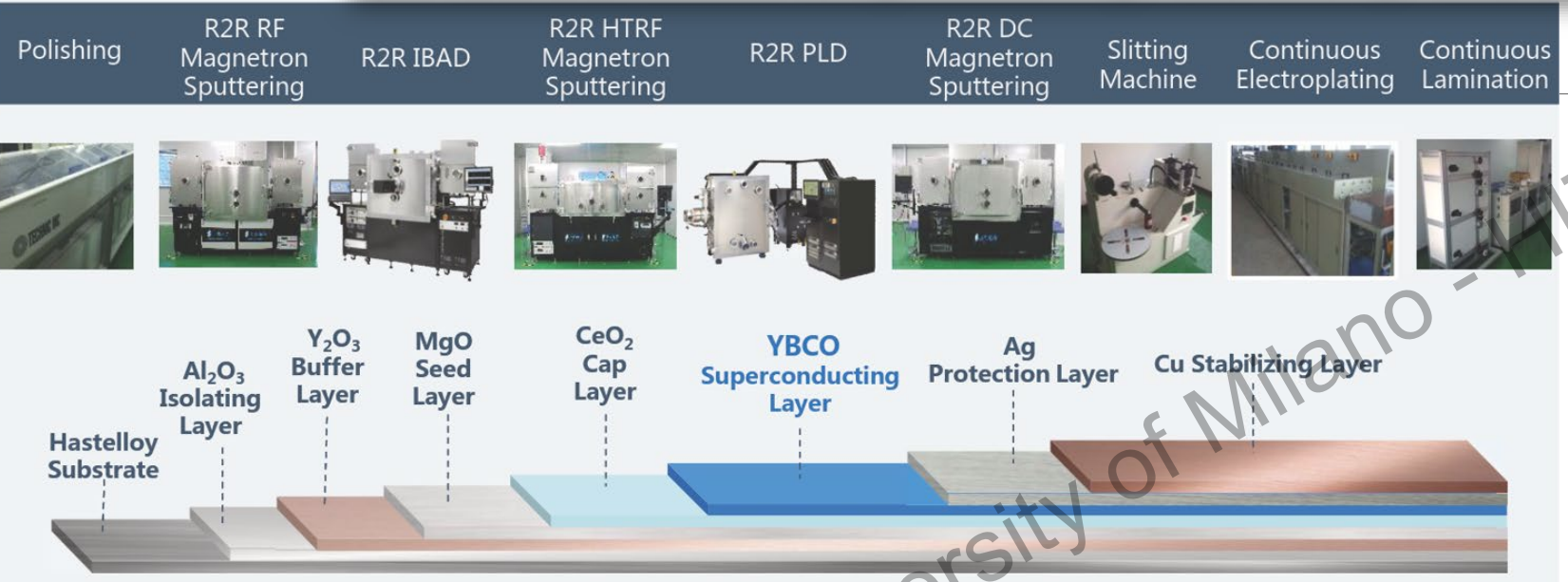
Multi-wire cable: the way to 10-100 kA!



Innovative materials: Nb₃Sn with high J_c



Inovative Materilas: HTS Bi-2212 and especially YBCO / REBCO



Double disordered YBCO coated conductor of industrial scale: high currents in high magnetic field



D Abramov¹, A Ballarino², C Barth³, L Bottura², R Dietrich⁴, A Francis¹,
J Jaroszynski¹, G S Majkic⁵, J McCallister¹, A Polyanskii¹, L Rossi²,
A Rutt⁴, M Santos¹, K Schlenga⁴, V Selvamanickam⁵, C Senatore³,
A Usoskin⁴ and Y L Viouchkov¹



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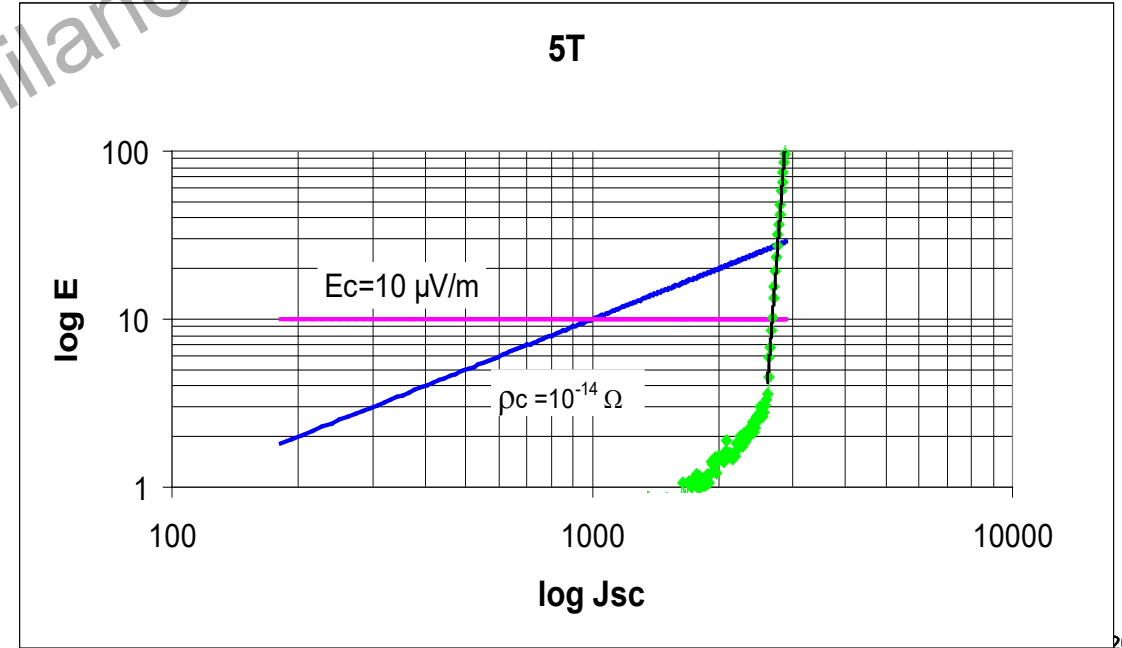
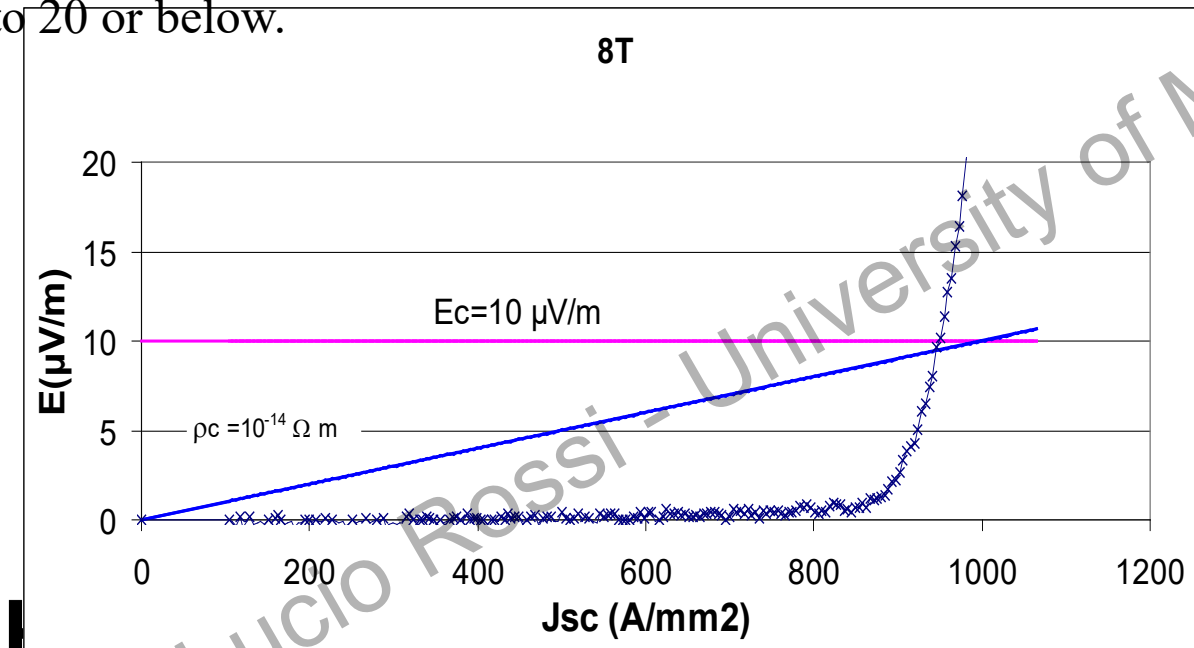
Max current in SC: I_c , E-J curve

Transition at fixed temperature: $V = k I^n$, so we have to adopt a criterion to define I_c .

Electric field. I_c is the current generating an electric field $E_c = 10^{-5}$ V/m $\Rightarrow E = E_c (J/J_c)^n$

Resistivity. I_c is the current showing an apparent resistivity of $\rho_c = 10^{-14}$ Ω m.

The exponent n , called also n -value or n -index, is related to the homogeneity of the material or of the superconducting properties. For good superconductors $n \sim 30 - 60$ or more. Near critical surface, $B > 0.9 B_{c2}$ the n -values drops down to 20 or below.



Carrying a lot of current: what a difference for magnets!



Resistive magnets of PS accelerator at CERN (1.5 tesla)

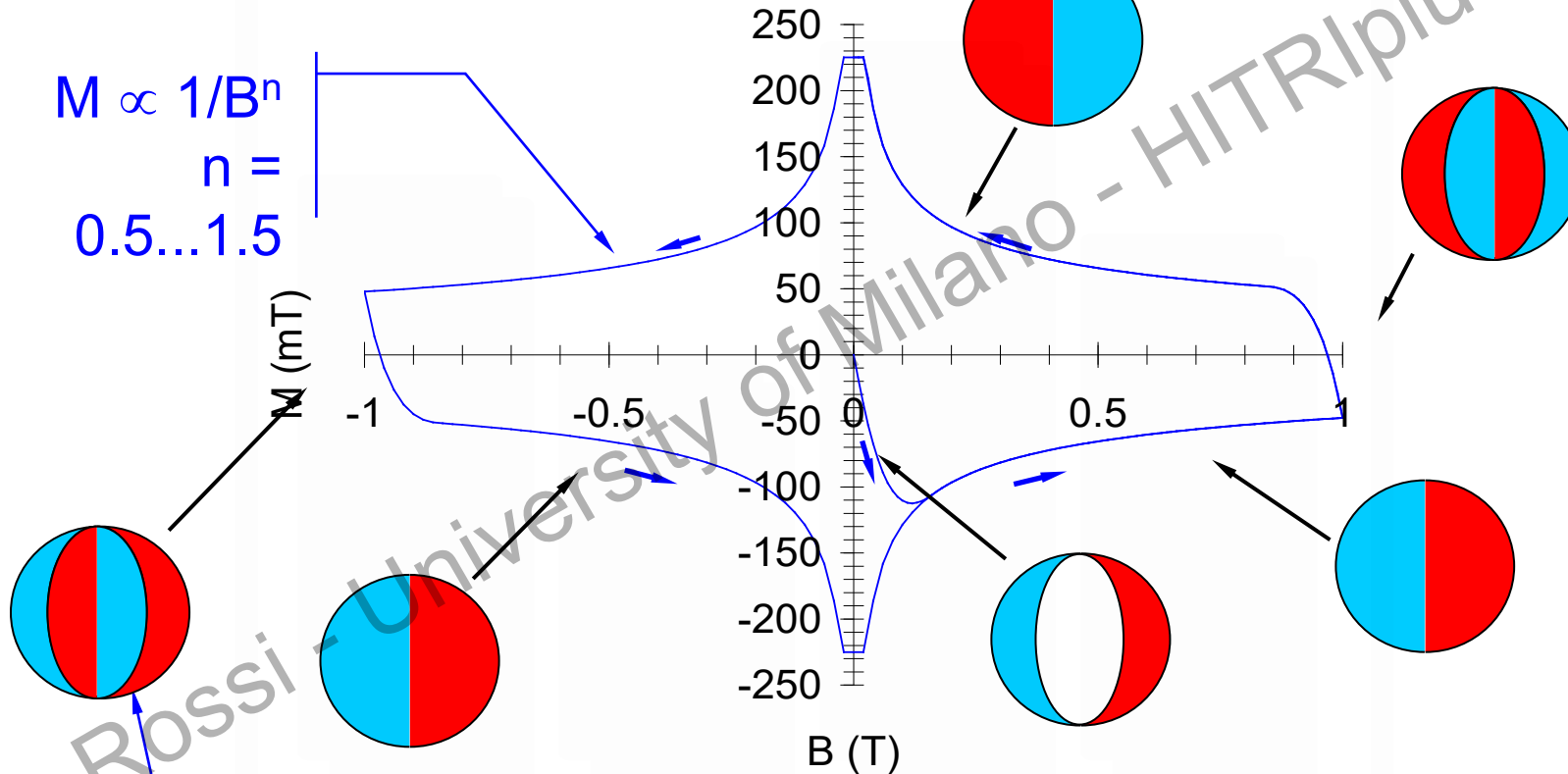
SC magnets at Tevatron at Fermilab (USA) 3 times more powerful!



(Dia)Magnetization and fine filaments

$$M \cong J_c D_{\text{eff}}$$

$$M \propto 1/B^n$$
$$n = 0.5 \dots 1.5$$

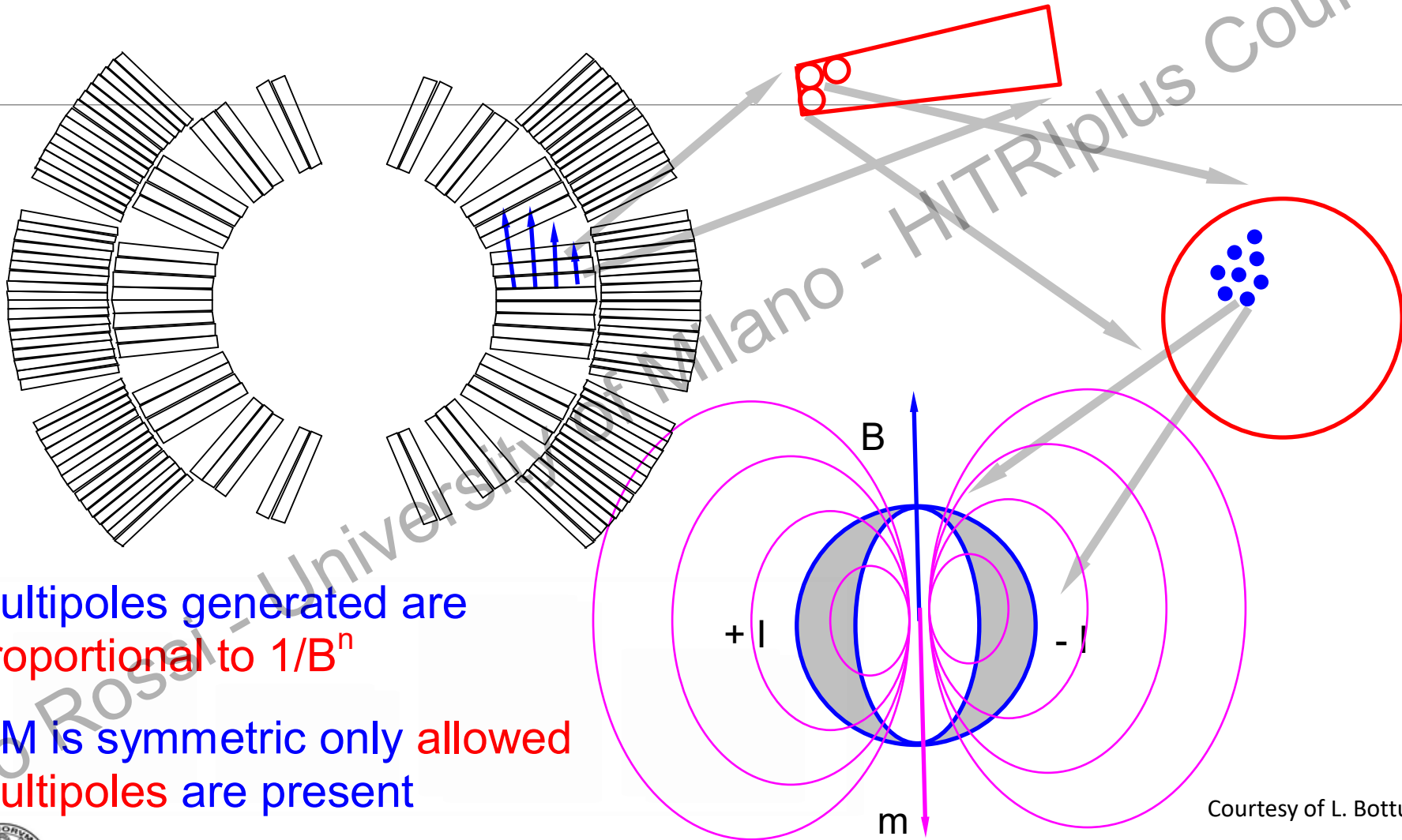


shielding current layer in a superconducting filament (persistent currents)

Courtesy of L. Bottura



Effect of magnetization



Multipoles generated are
proportional to $1/B^n$

If M is symmetric only **allowed**
multipoles are present

Courtesy of L. Bottura

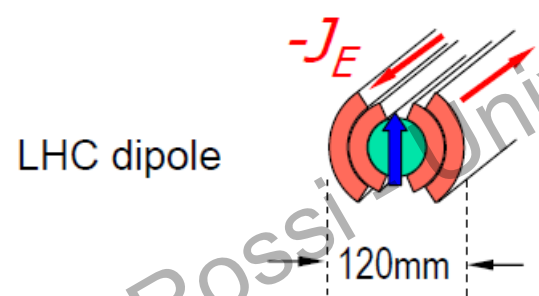


The advantage of high current density

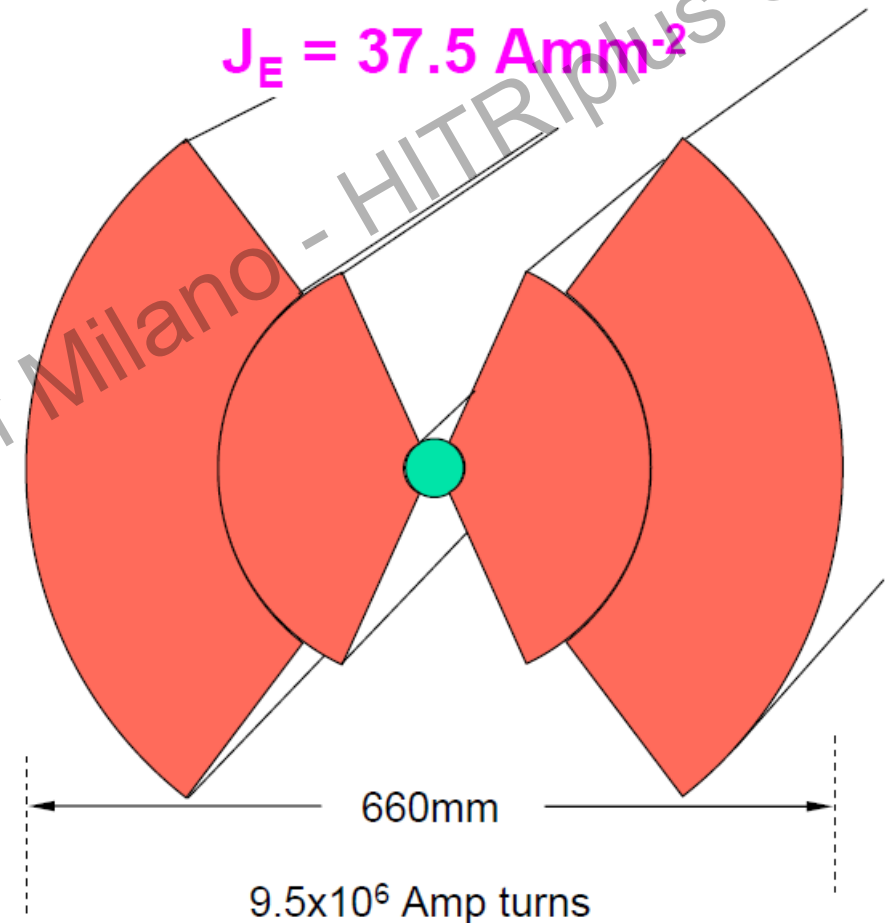
- The field produced by an ideal dipole is:

$$B \approx 1/\pi \mu_0 J_E t$$

$$J_E = 375 \text{ Amm}^{-2}$$



9.5×10^5 Amp turns
 $= 1.9 \times 10^6$ A.m per m



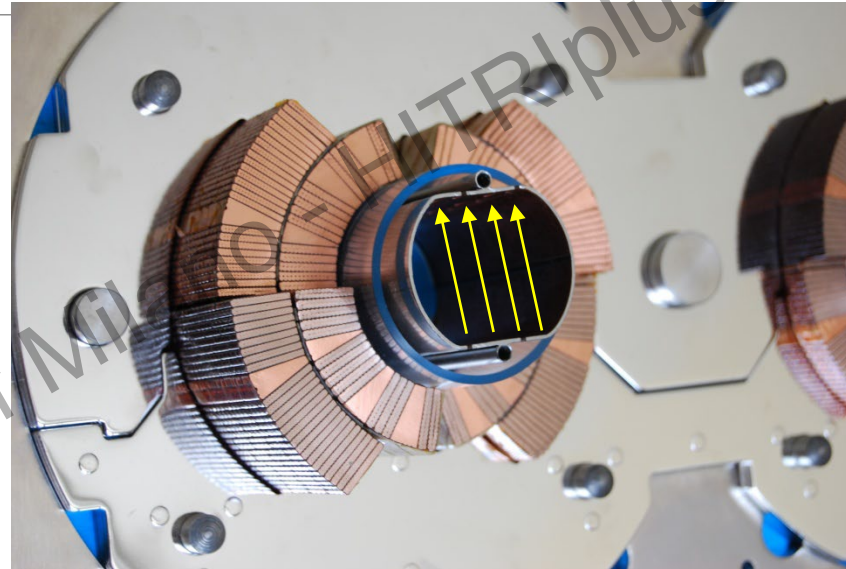
9.5×10^6 Amp turns
 $= 1.9 \times 10^7$ A.m per m

Accelerator magnets : basic - field

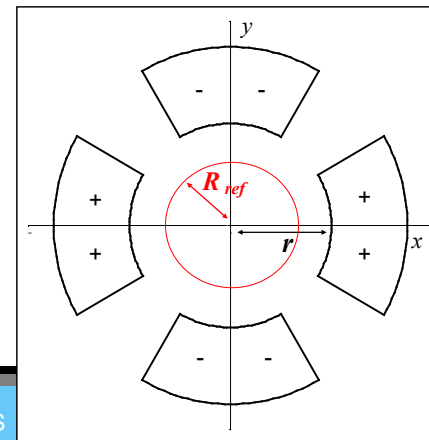
The basic shape : mix between cos ϑ and shell

Shells with const J is a very good approximation

Field expansion



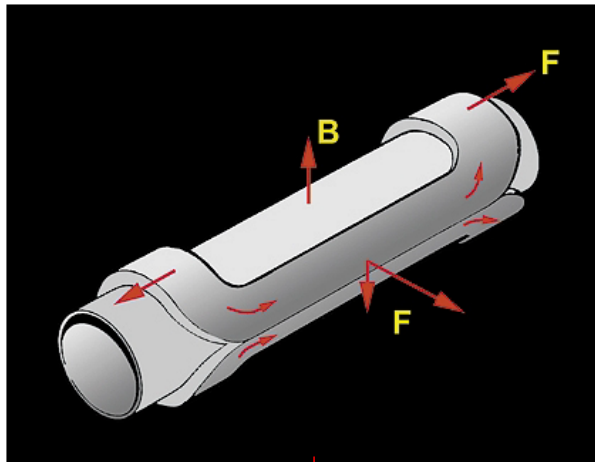
$$B_y + iB_x = 10^{-4} B_1 \sum_{n=1}^{\infty} (b_n + ia_n) \left(\frac{x + iy}{R_{ref}} \right)^{n-1}$$



Accelerator Magnets: Basic Design - forces

$J_{\text{overall}} \approx 500 \text{ A/mm}^2$! e.m. forces are not kept by conductors but tend to tear apart the winding.

Principle

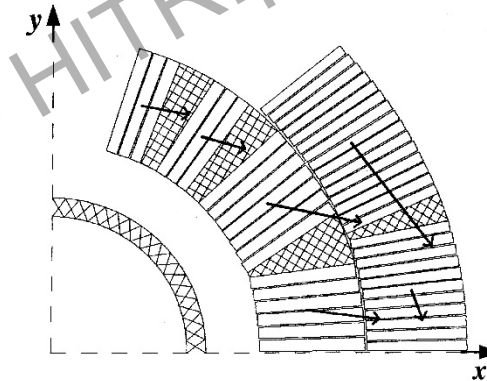


Reality



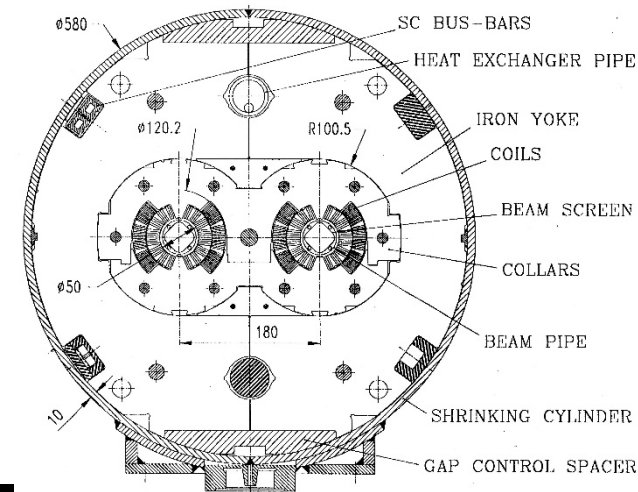
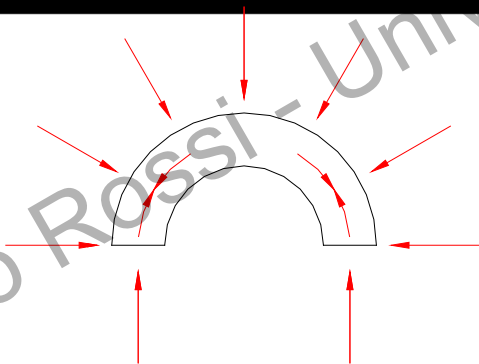
e.m. forces

NOT SELF-SUPPORTING



How to contain them

More difficult in twin magnets!



European Union's Horizon 2020
grant agreement No 101008548

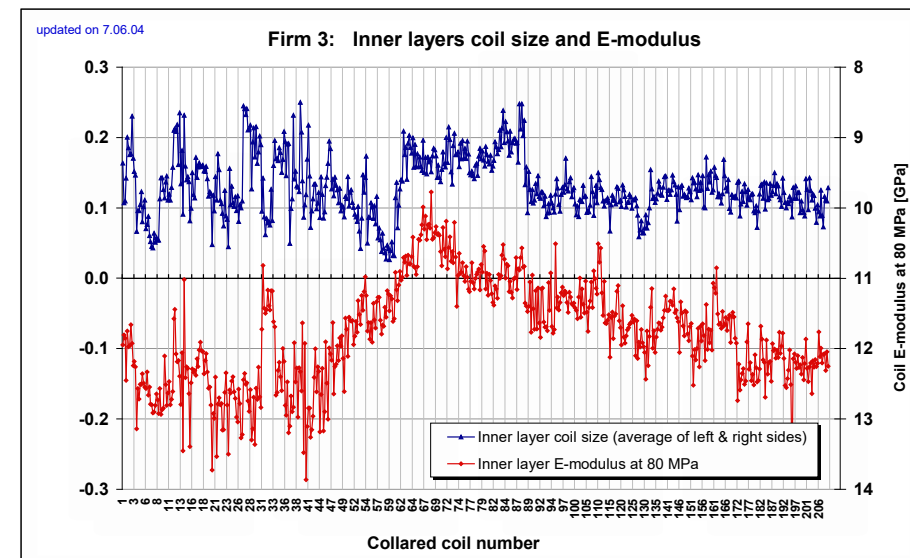
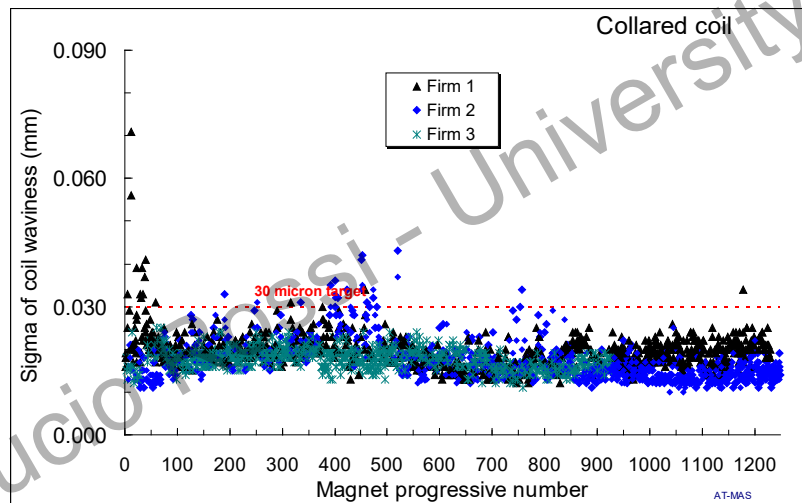
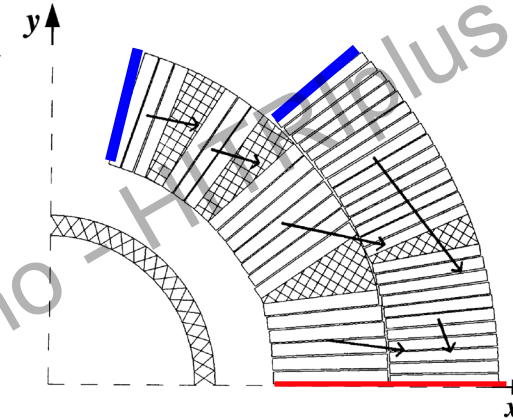
Accelerator magnets : total quality

All in series

- The worst performing determines the accelerator energy
- They must be all equal (within few units, $1 \mu = 10^{-4}$)

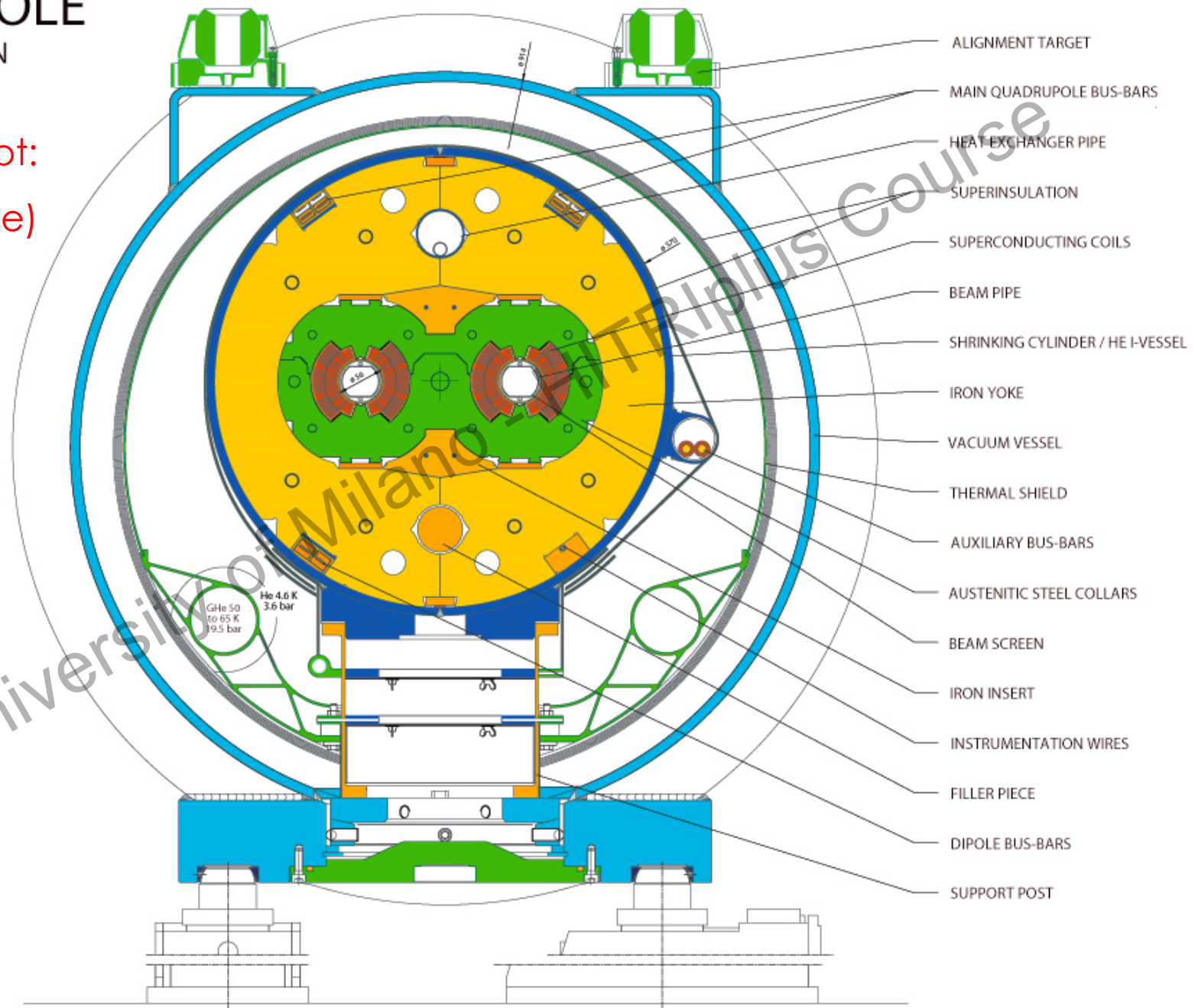
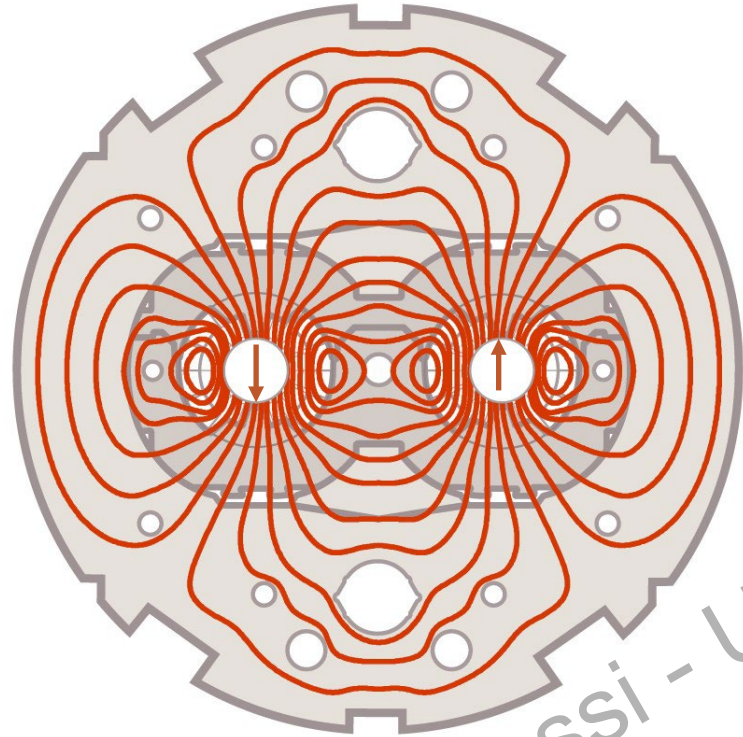
The field quality is typically controlled at a fraction of unit over a wide range.

- Coil positioning and e.m. forces



LHC DIPOLE CROSS SECTION

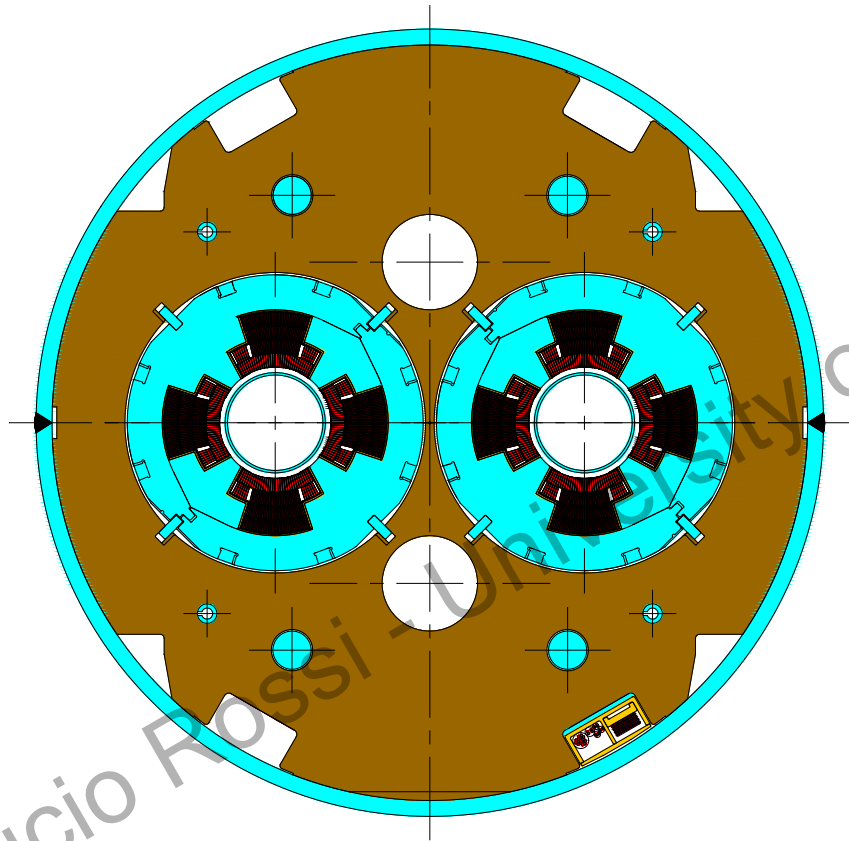
Twin Concept:
MB Main Bend (Dipole)



- ALIGNMENT TARGET
- MAIN QUADRUPOLE BUS-BARS
- HEAT EXCHANGER PIPE
- SUPERINSULATION
- SUPERCONDUCTING COILS
- BEAM PIPE
- SHRINKING CYLINDER / HE I-VESSEL
- IRON YOKE
- VACUUM VESSEL
- THERMAL SHIELD
- AUXILIARY BUS-BARS
- AUSTENITIC STEEL COLLARS
- BEAM SCREEN
- IRON INSERT
- INSTRUMENTATION WIRES
- FILLER PIECE
- DIPOLE BUS-BARS
- SUPPORT POST

Lucio Rossi - Università del Piemonte Orientale

MQY wide aperture quadrupole



70 mm ID coil

$G = 160 \text{ T/m}$ at 4.5 K

$I = 3620 \text{ A}$

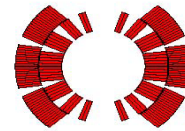
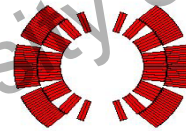
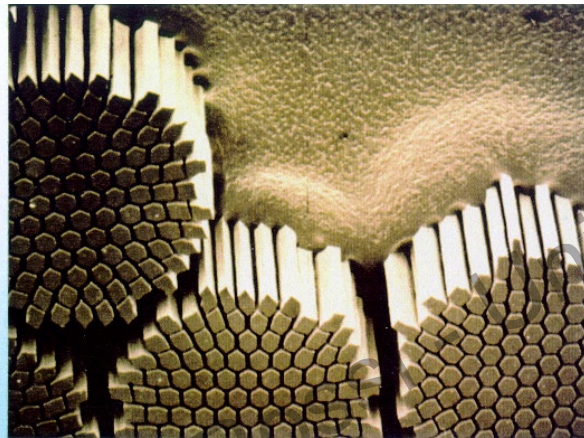
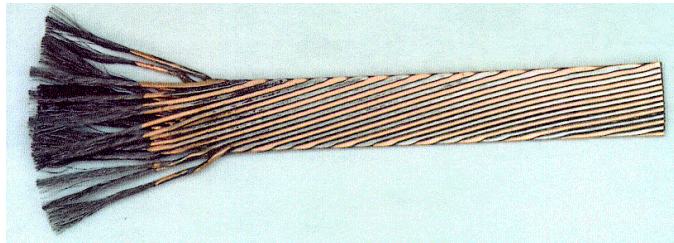
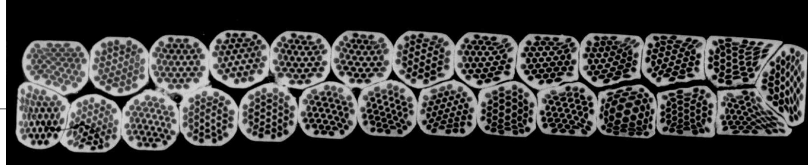
$E = 141 \text{ kJ/m/aperture}$

$L_{\text{mag}} = 3.4$

- Four-layer, graded shell coil.
- Free standing collars, fully supporting the forces.
- Two-in-one iron yoke.

This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

LHC MB X-sect: conductor (Rutherford cable)



designed with the ROXIE code developed at CERN for the LHC (S. Russenschuck)

Conductor position optimization:

Control of harmonics
Balance of margin among blocks

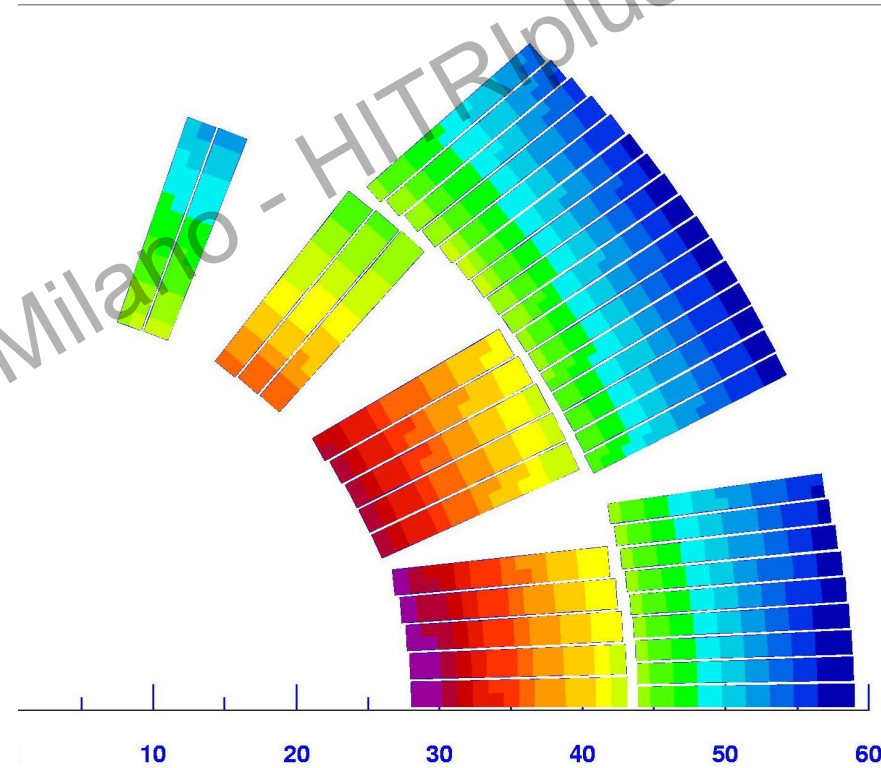
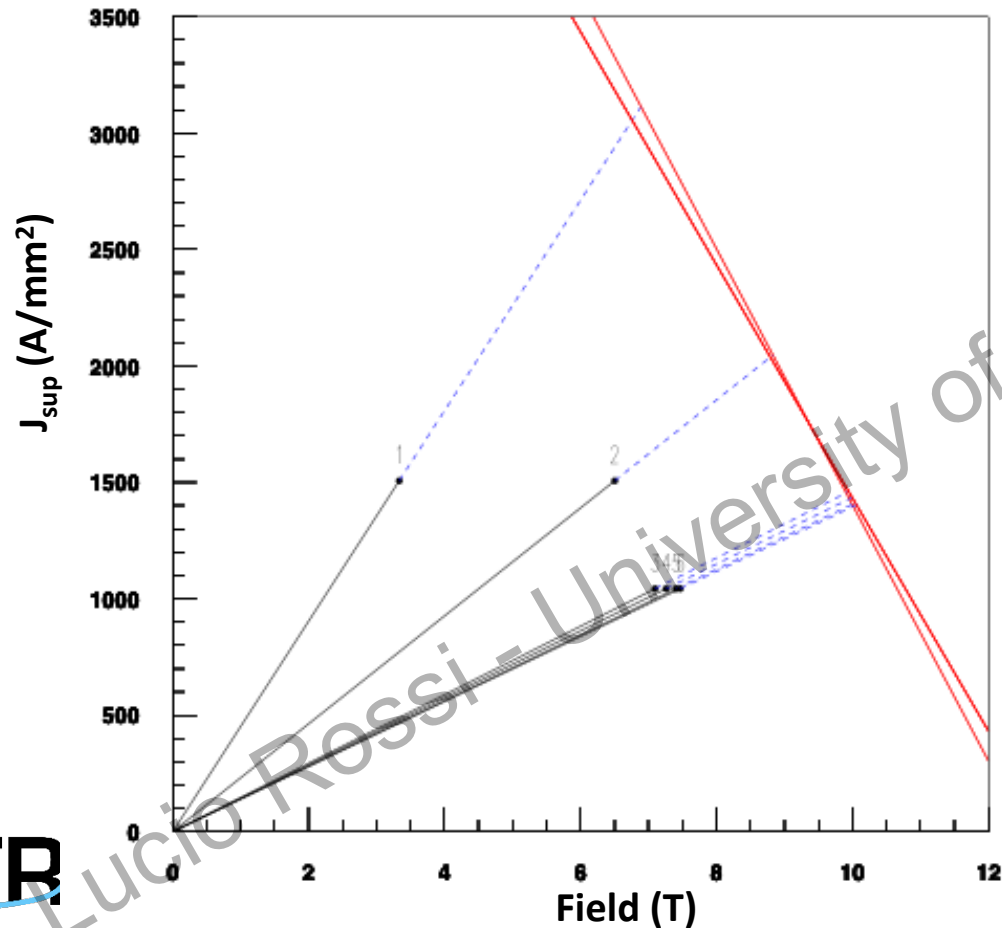
Stable against inevitable errors

Minimum shear among conductors

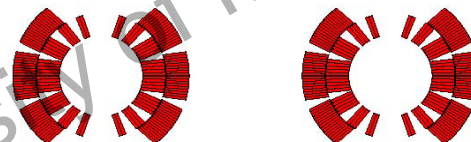
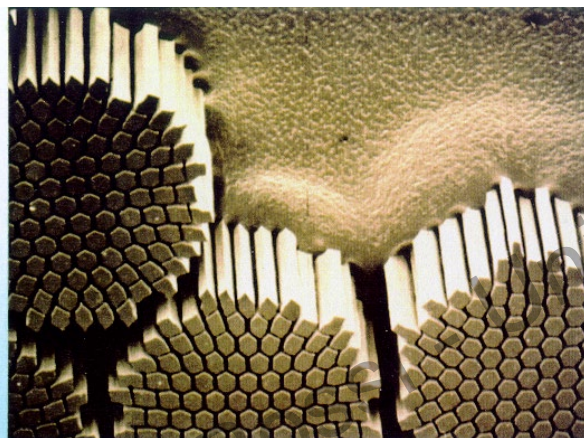
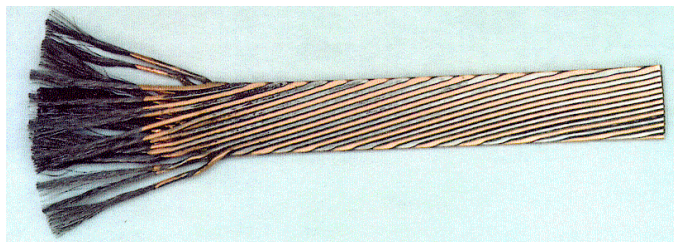
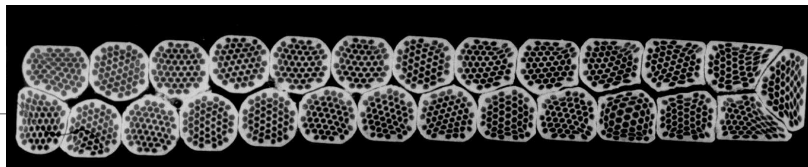
Balance between T margin of inner/outer

No quench anymore in straight part

The basic relation: B vs I



LHC MB X-sect: conductor (Rutherford cable)



Conductor position optimization:

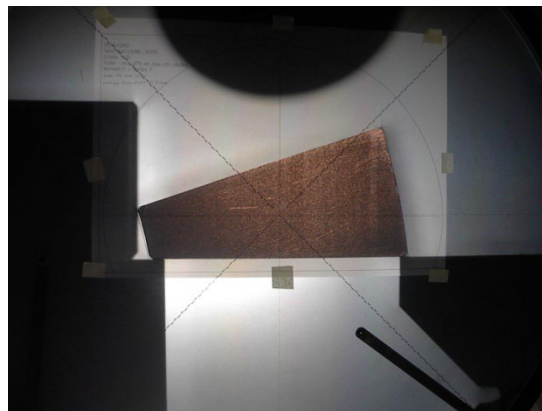
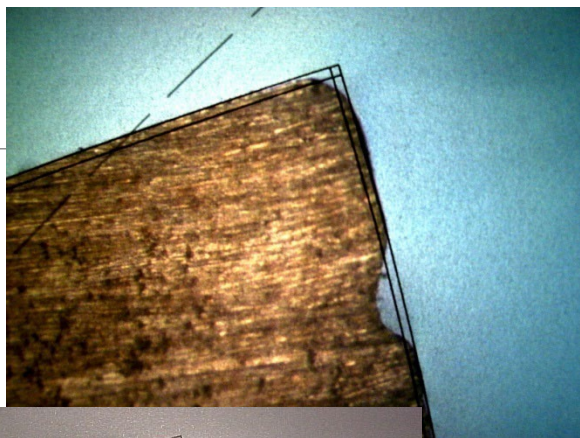
Control of harmonics
Balance of margin among blocks

Stable against inevitable errors

Minimum shear among conductors

Balance between T margin of inner/outer

LHC MB X-sect: copper wedges



Cu Spacer

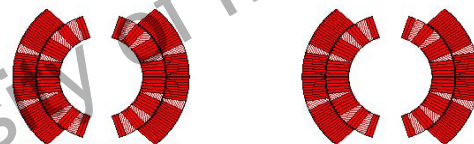
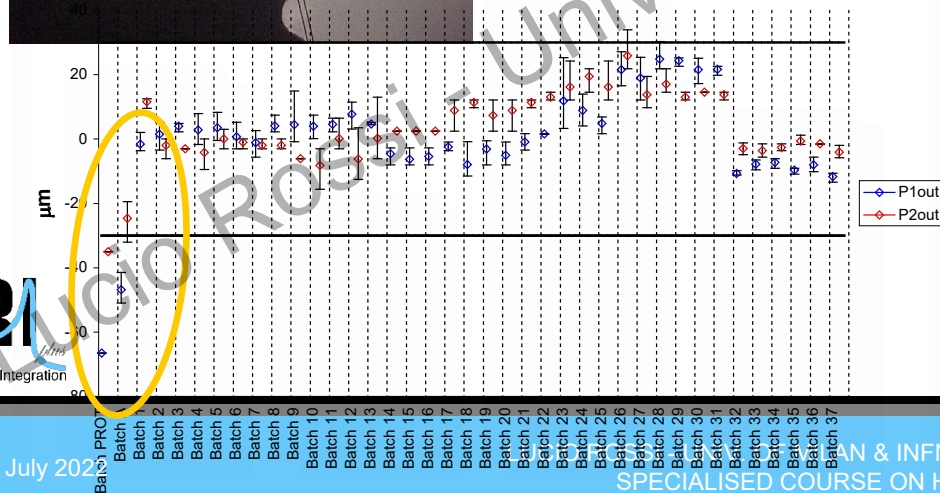
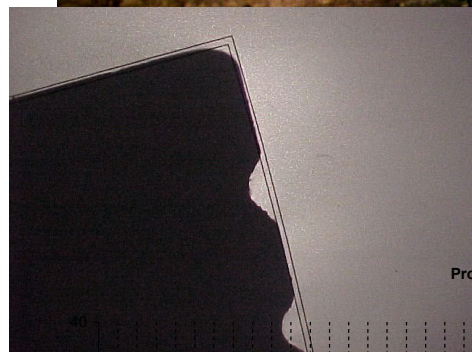
Precise at $\pm 20 \mu\text{m}$

Used to steer production

Change of Cu wedge
0.2-0.5 mm of inner
wedges in July 2001 (3
CM, 15 CC).

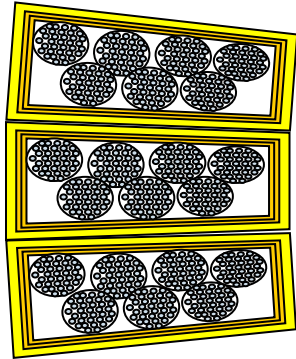
Effect in 2002.

~35 CM old Xsect.



LHC MB X-sect: conductor and ground Insulation, Interlayer

Rutherford Cables Insulation



-2 layers of Apical 200 AV insulation

-1 layer Pixeo to glue cables together at 185°C (-0, +5 critical)

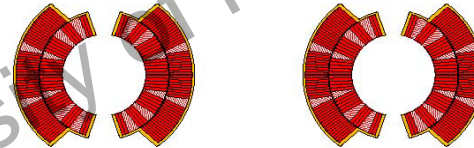
Ground isolation

Four layers 125AH

Polyimide insulation

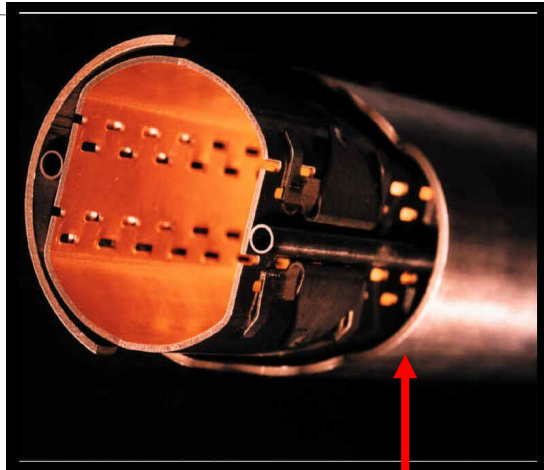
Around cable and around coils

Important elements are dimensions, $\pm 3\%$ of thickness, and creep (Apical creeps less than kapton)



Inter layer
To allow HEII to flow

LHC MB X-sect: insulated CBT

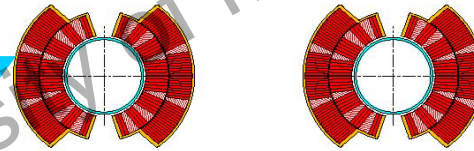


Cold bore tube

StSt tubes

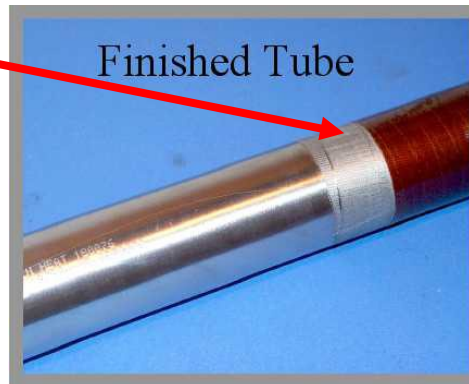
Insulation done at CERN
Special Insulation
technique > 20 kV

Clearance between coils and
insulated CBT is about 0.5 mm
over the 15 m length

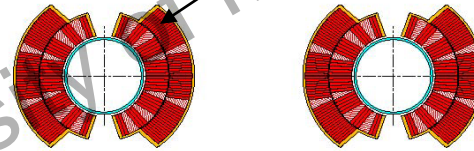


CBT

Finished Tube



LHC MB X-sect: Quench Heater

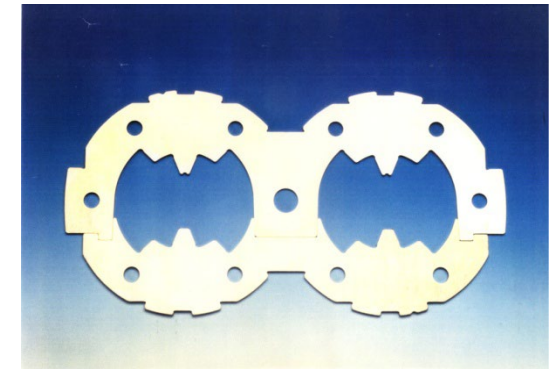
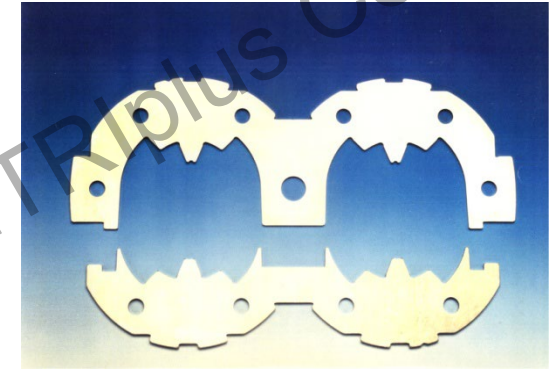
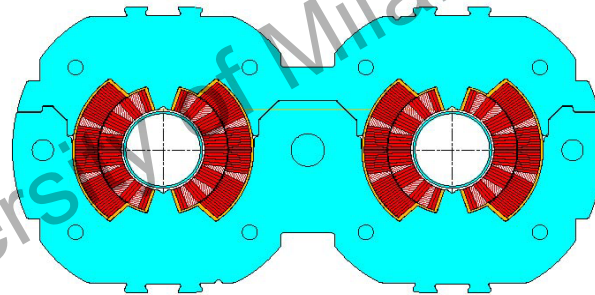
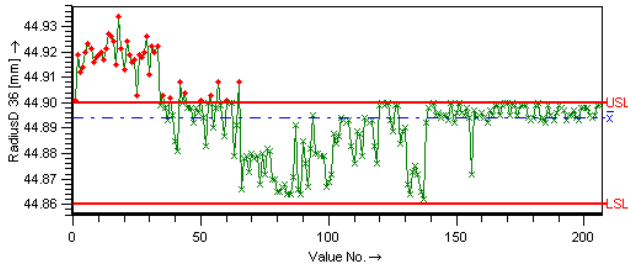
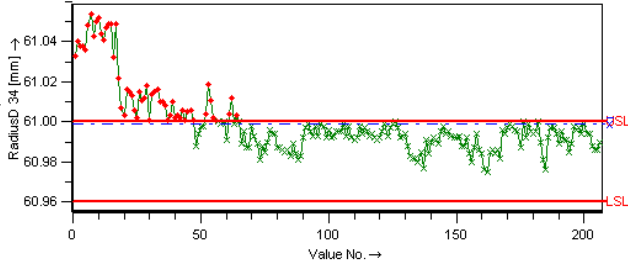


Strips of stainless steels partially coated with copper to adjust resistance

Encapsulated like a sandwich in two foils of 75 μm of polyimide

Fired by current pulse, heat must diffuse from strip to coils in 20-50 ms !!

LHC MB X-sect: Collars

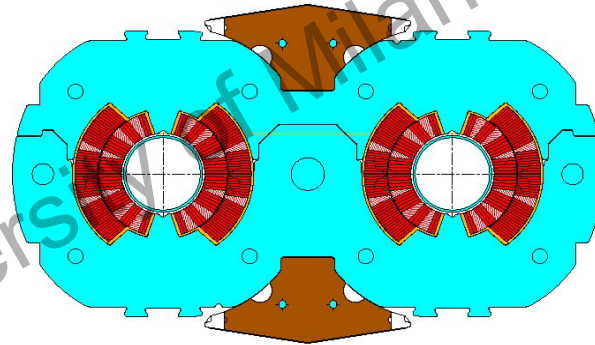


Collars and collaring are the main controllers of the final coil shape

LHC MB X-sect: magnetic insert

The iron Insert

Punched together
with the yoke
lamination



Introduced to ease the
mechanical assembly

It serves for FQ

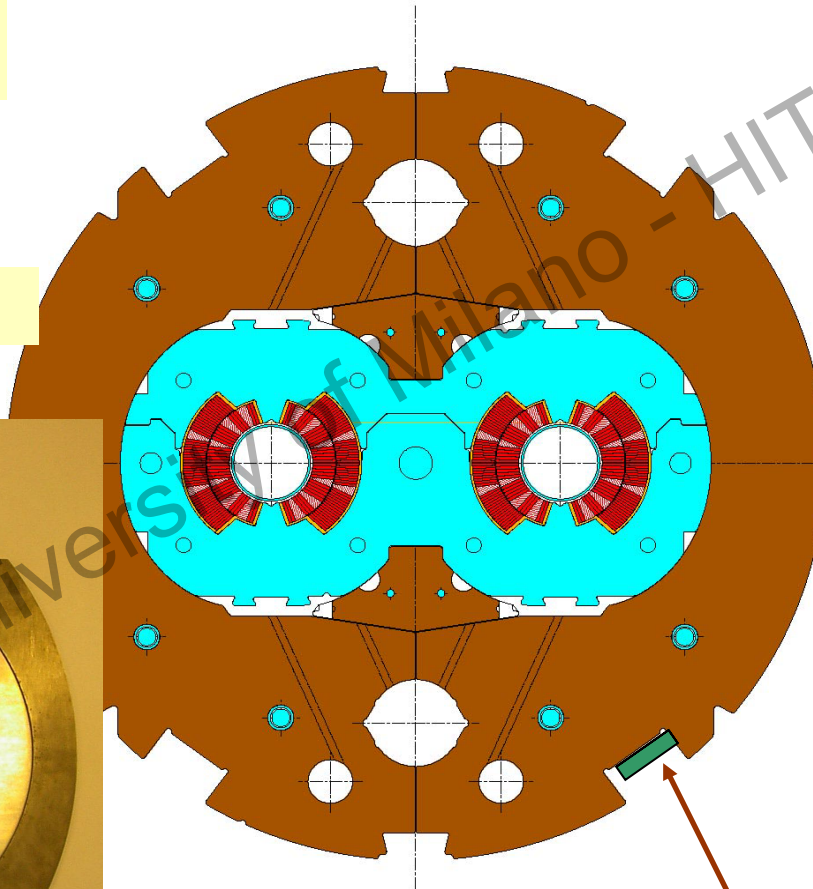
By tapering we cured
unwanted quadrupole
and octupole
components

LHC MB X-sect: yoke laminations

One supplier for the steel
45,000 tons

Precise vertical gap

Regular *Nested*



The iron yoke:

Stray field

15% field increase
(but big gain in
protection)

If saturates affect FQ
(sextupole)

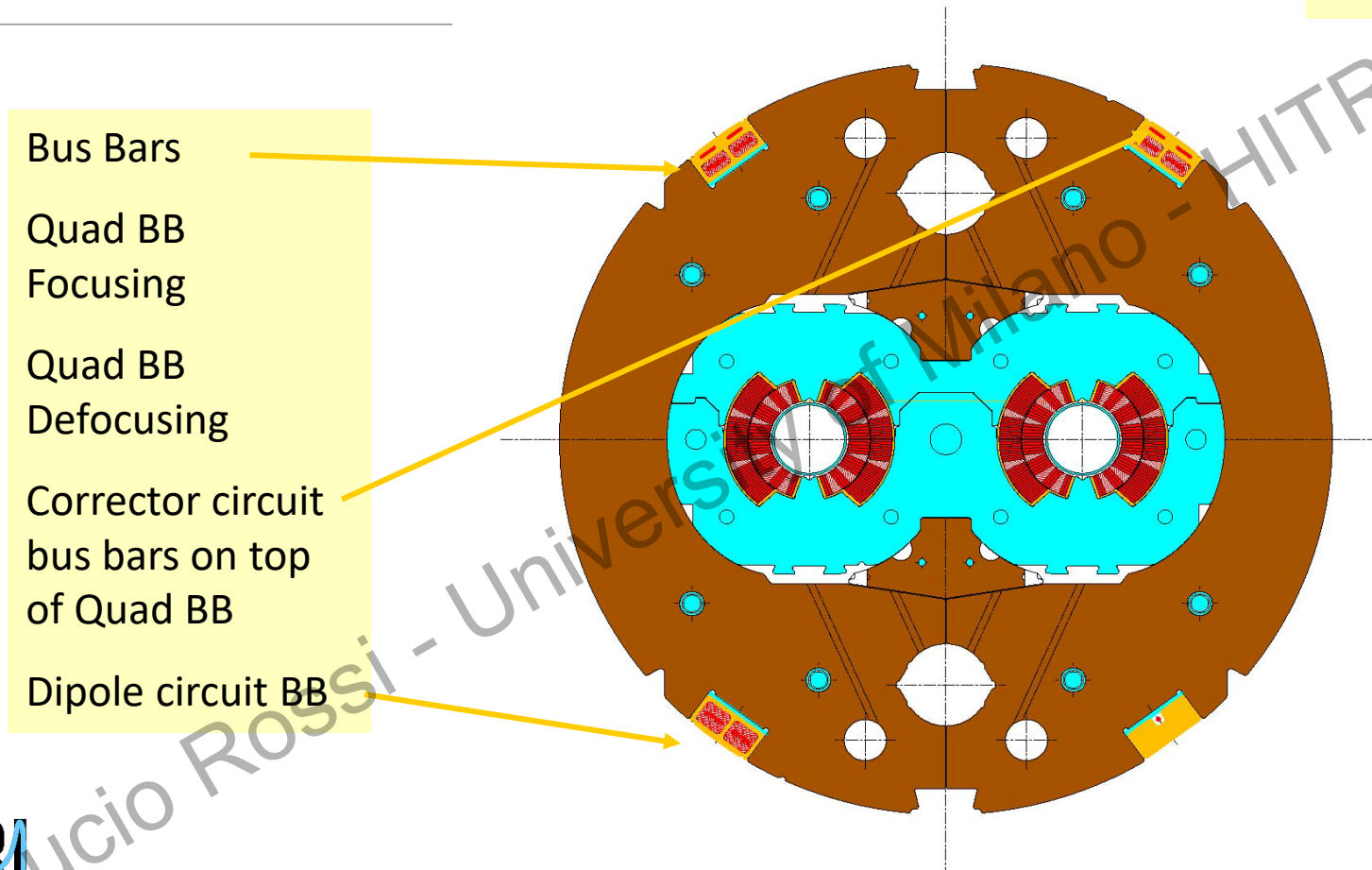
Trim of magnetic
length

Temperature probe

It has received funding from the European Union's Horizon 2020
and innovation programme under grant agreement No 101008548

LHC MB X-sect: Bus Bars & fillers

160 km of main BusBars!!



LHC MB X-sect: Shrinking cylinder and support

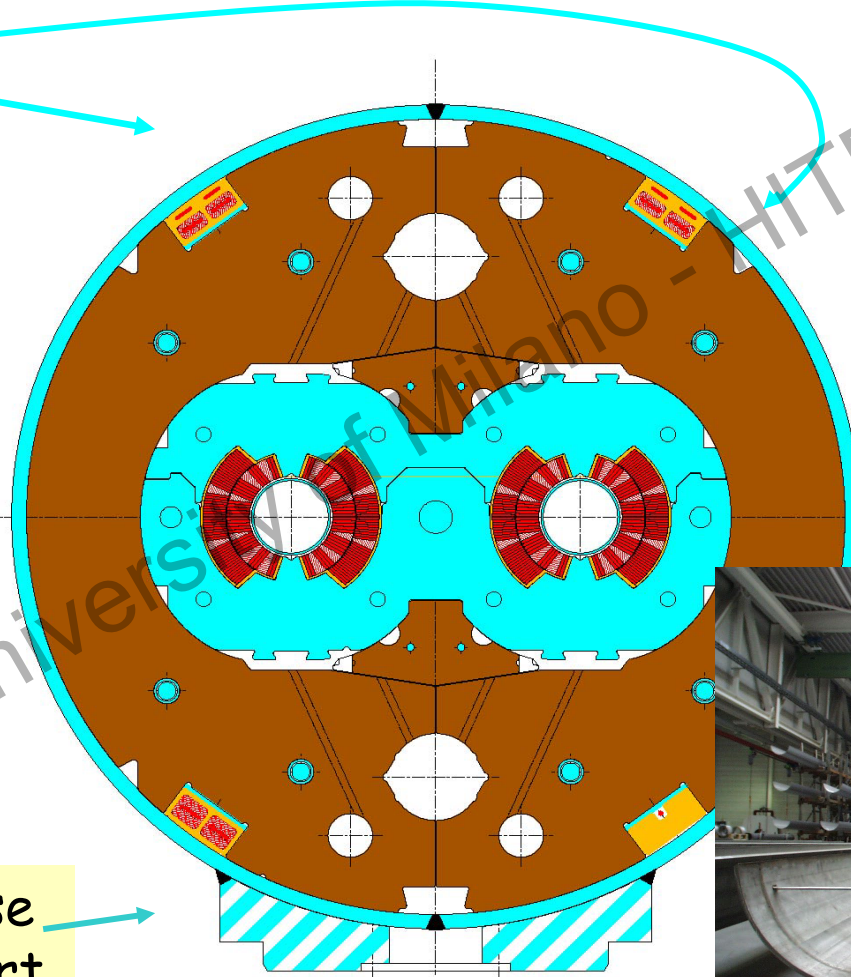
Two half shells, welded on the magnet

Many difficulties

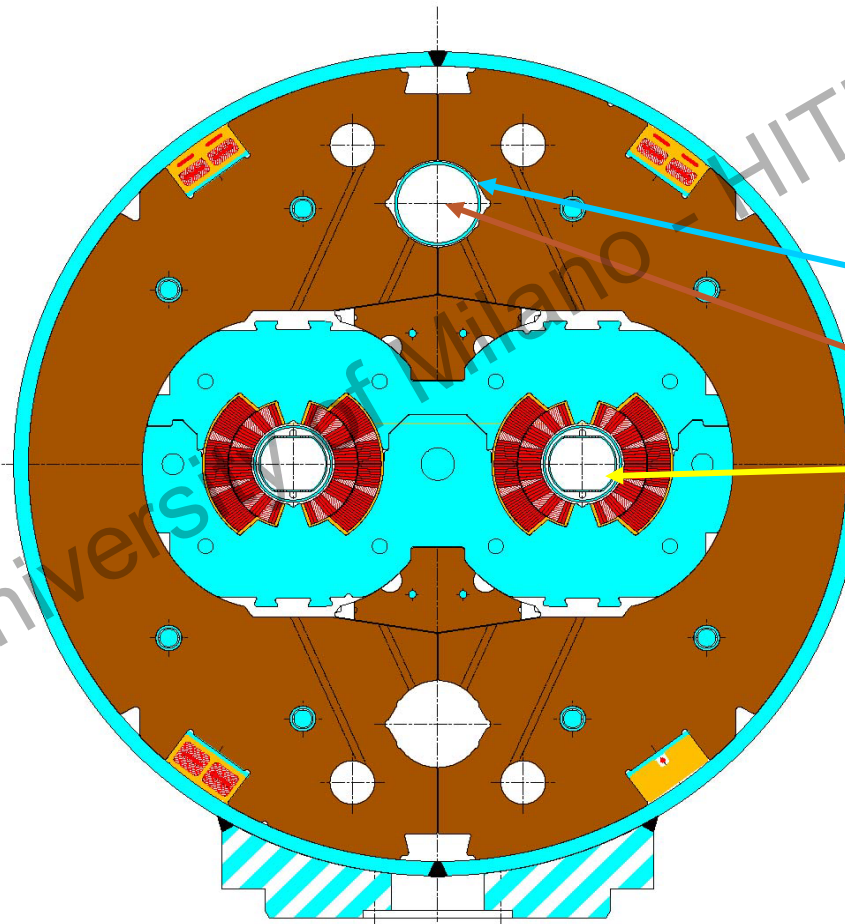
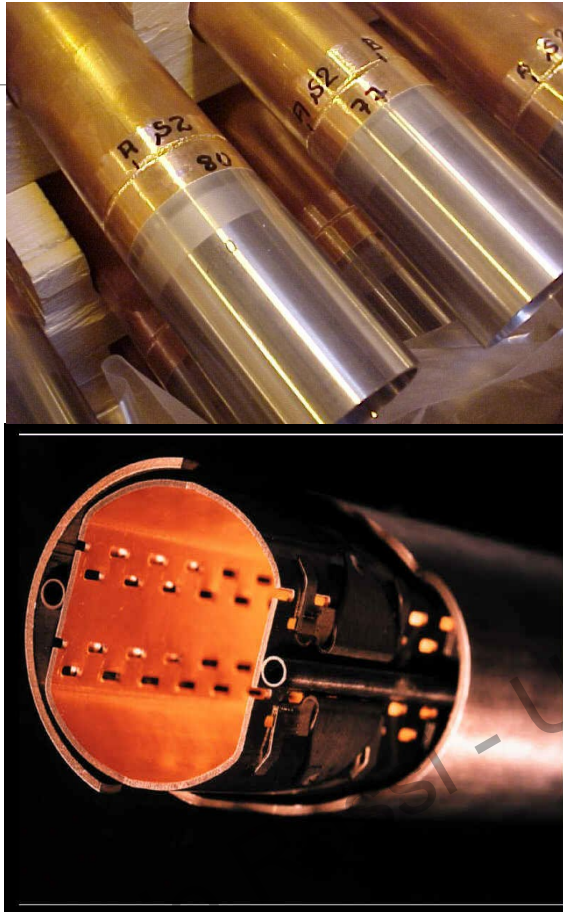
Curvature released from ± 1 to ± 2.5 mm: still not OK.

To cure this we went to sorting

Precise support



LHC MB X-sect: beam screen and HXT



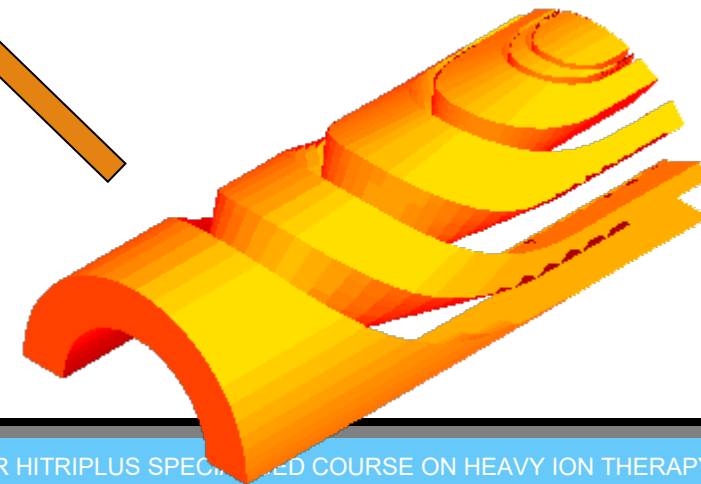
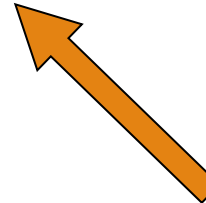
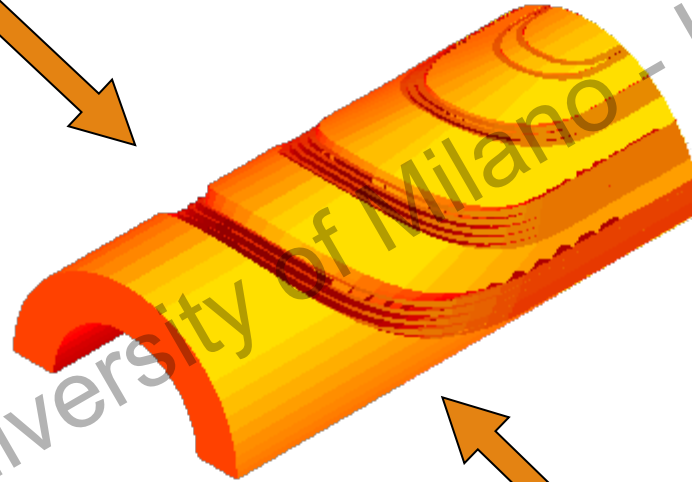
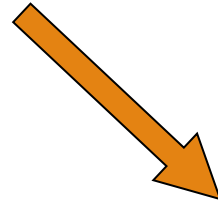
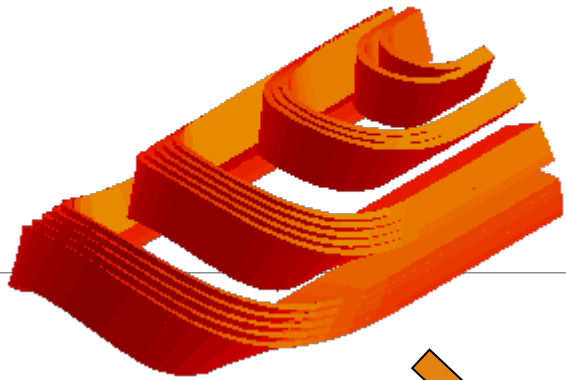
Copper Heat Exchange Tubes

HEII satur.

Beam Screen

Inserted at CERN just before insertion in the tunnel

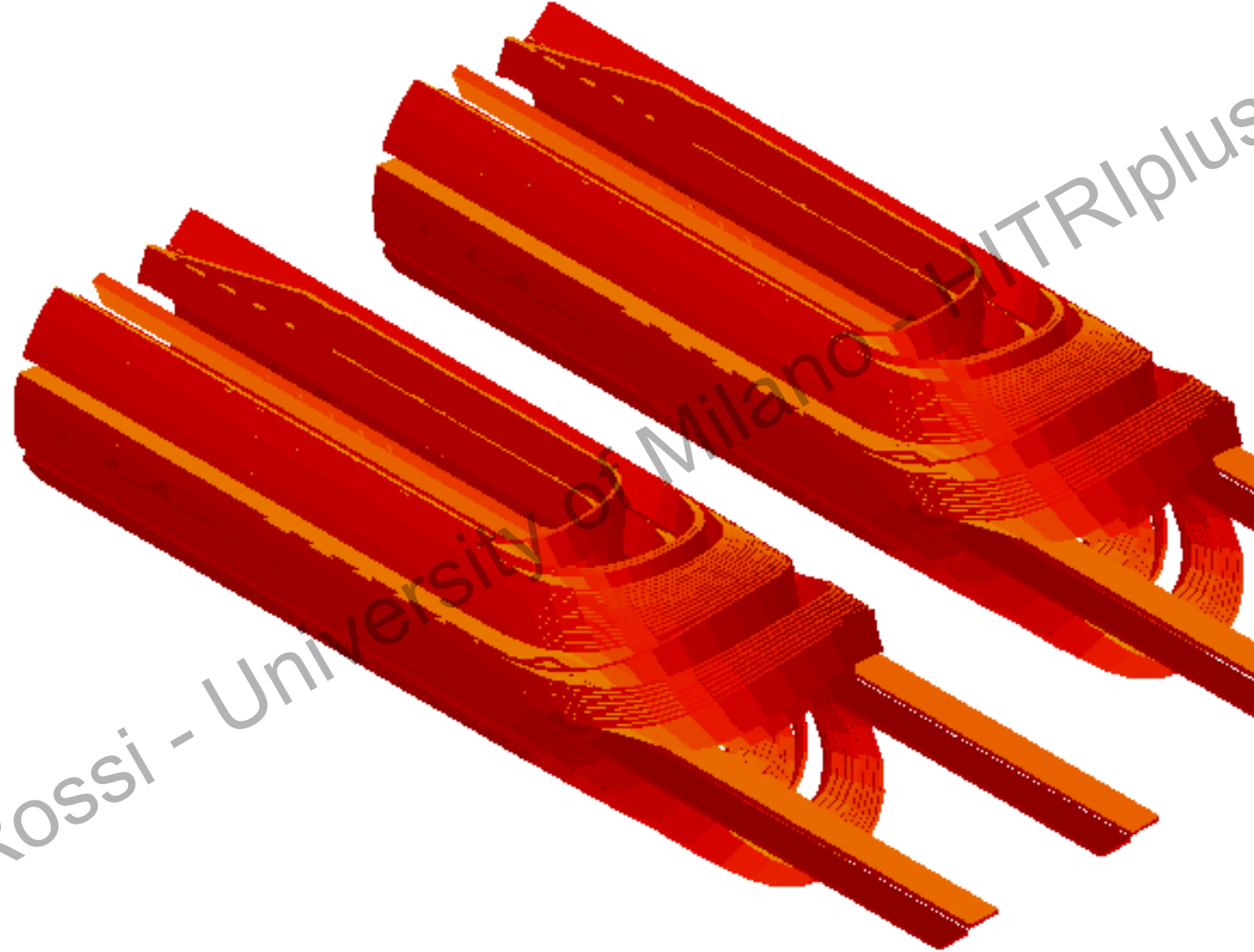
3D Inner layer lyre side



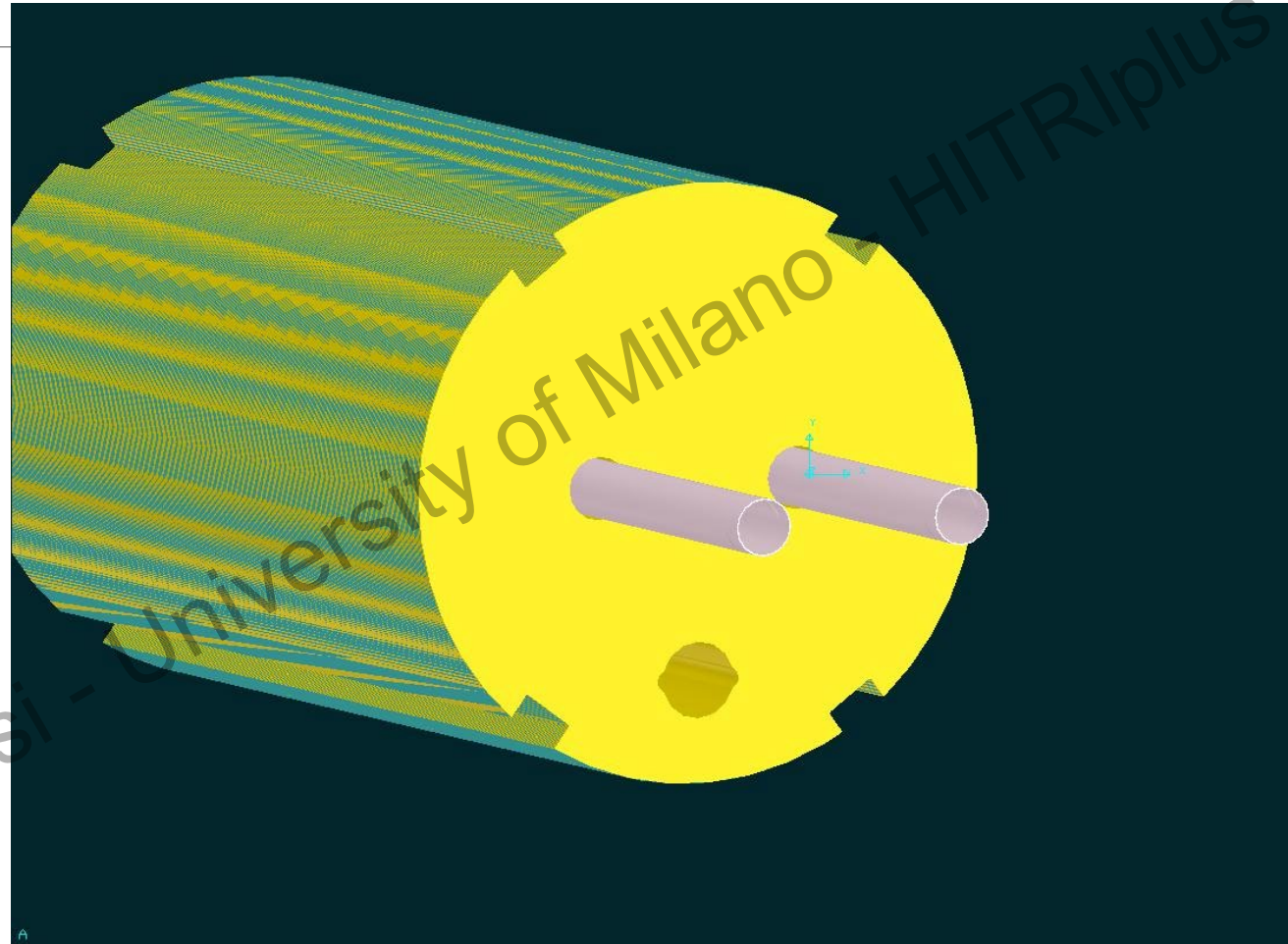
End Spacers: critical for Quench



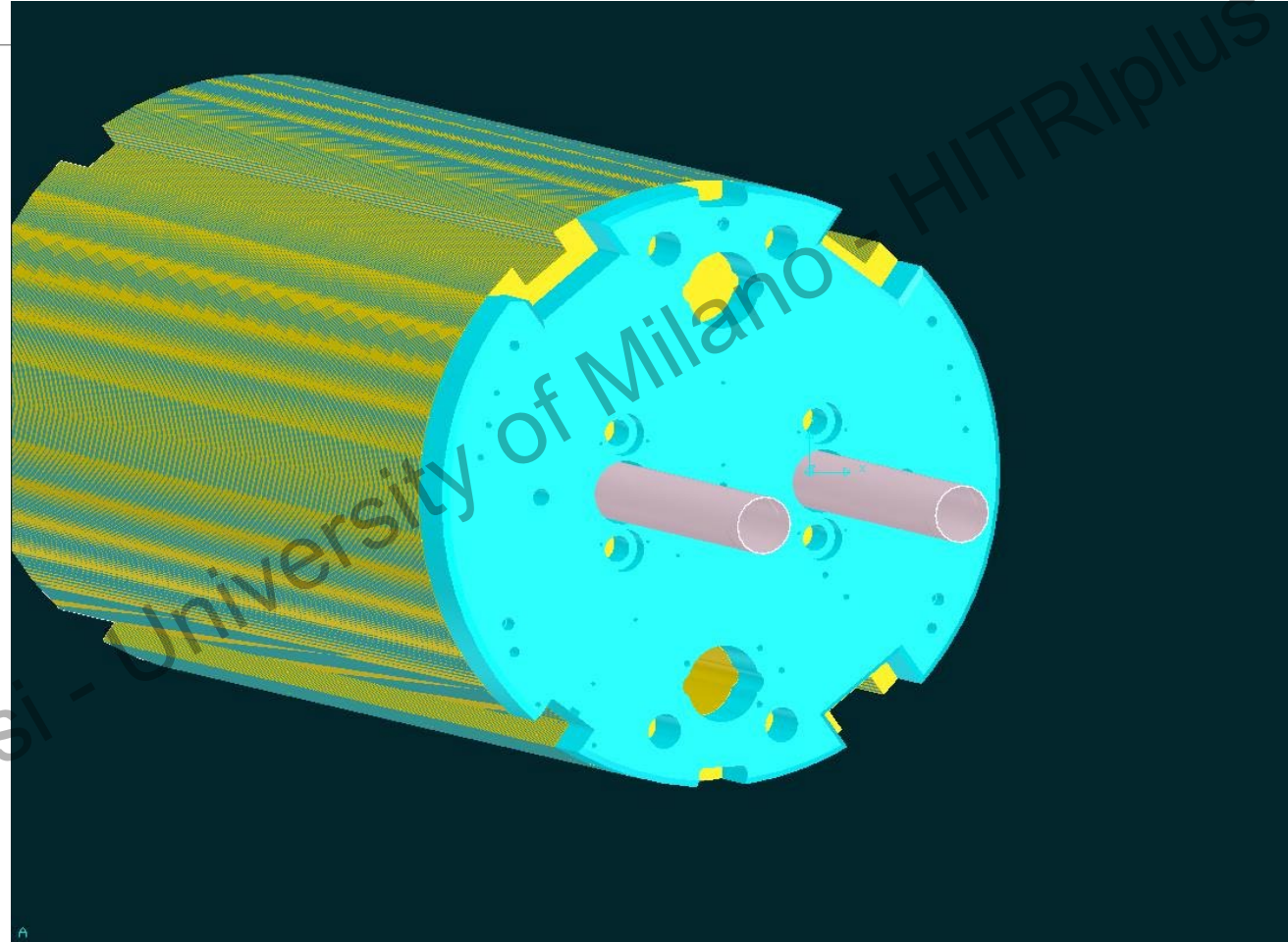
3D Connection side



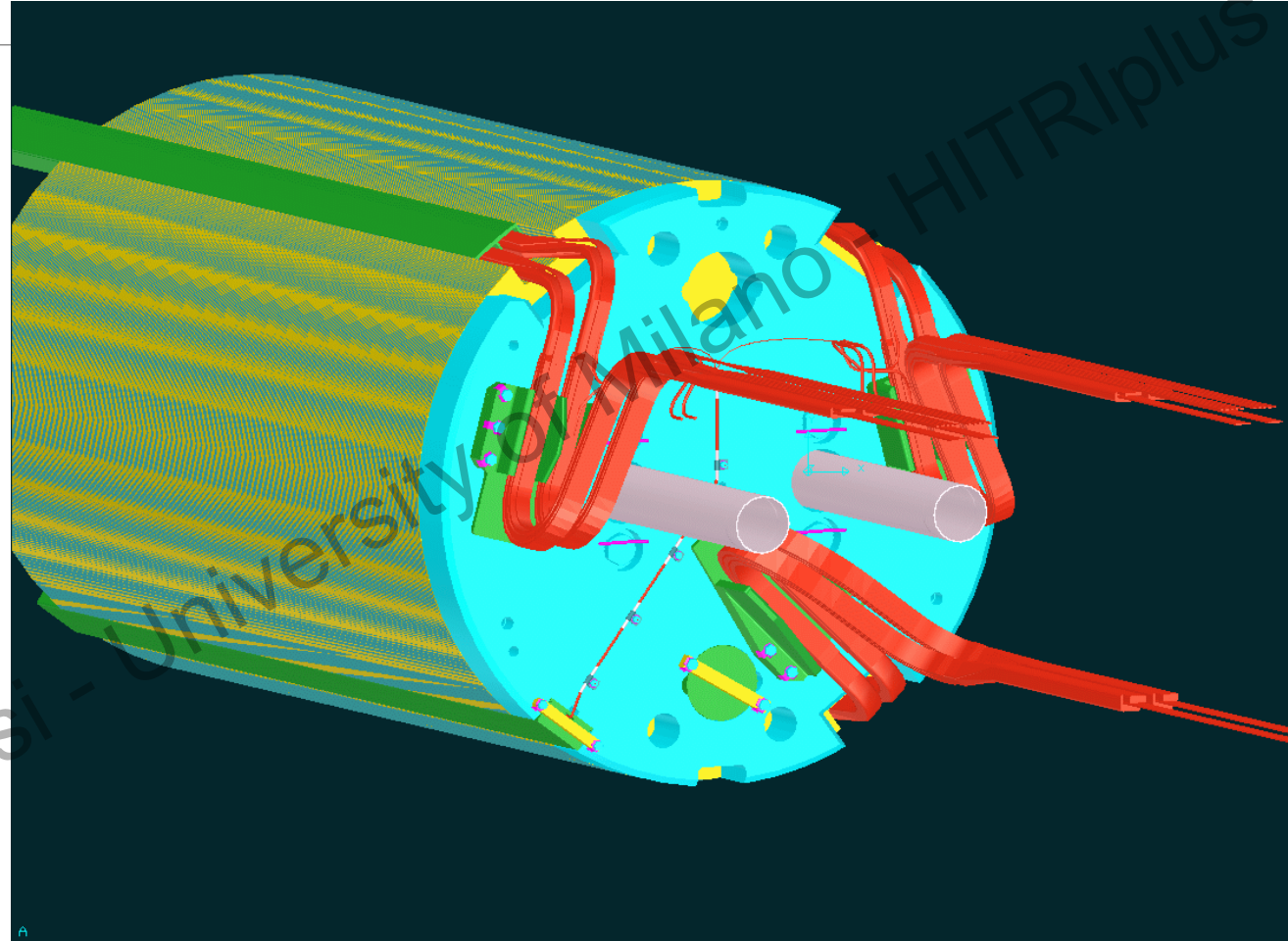
LHC MB - end part CBTs and Yoke



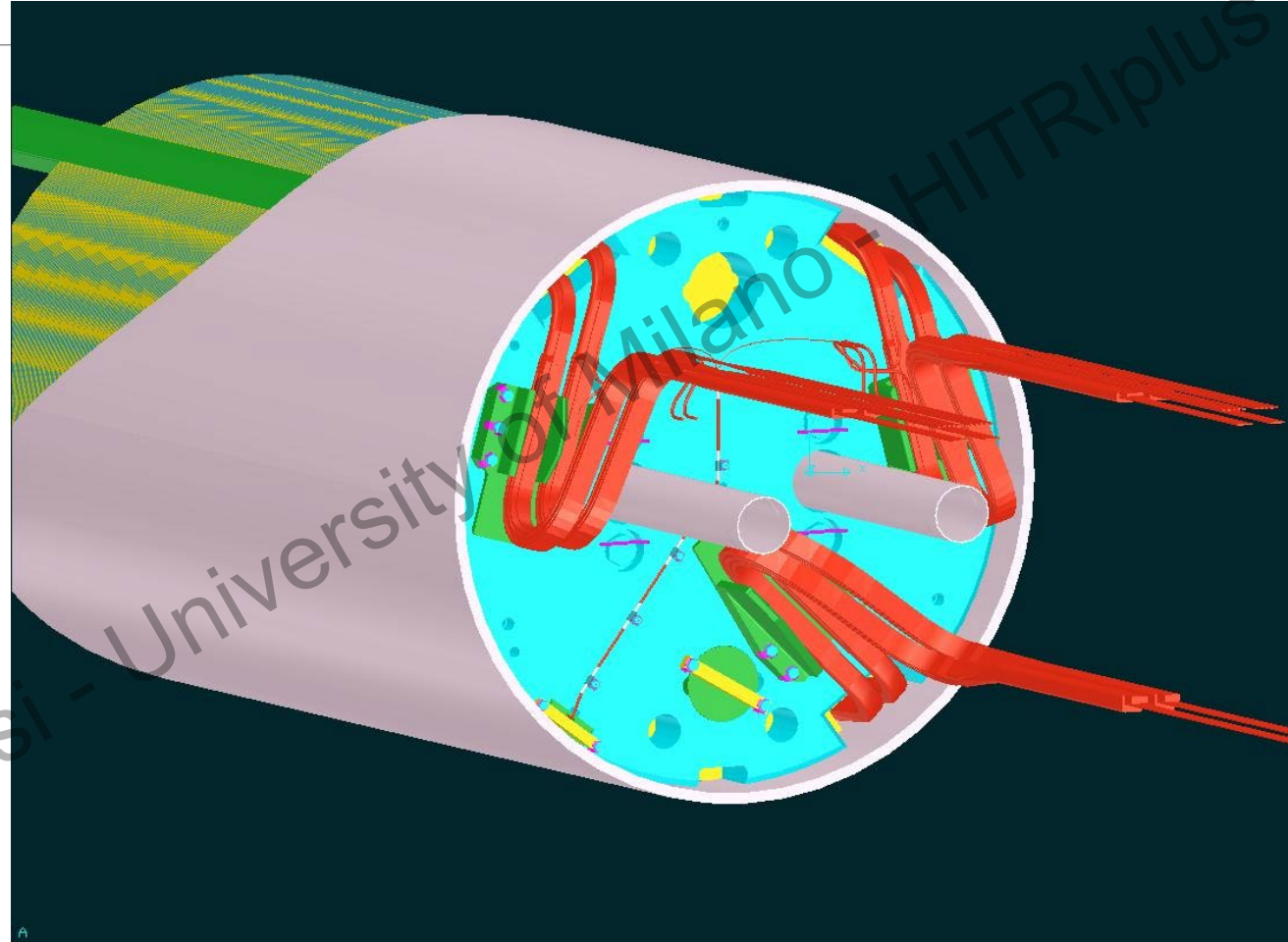
LHC MB -end part end plate



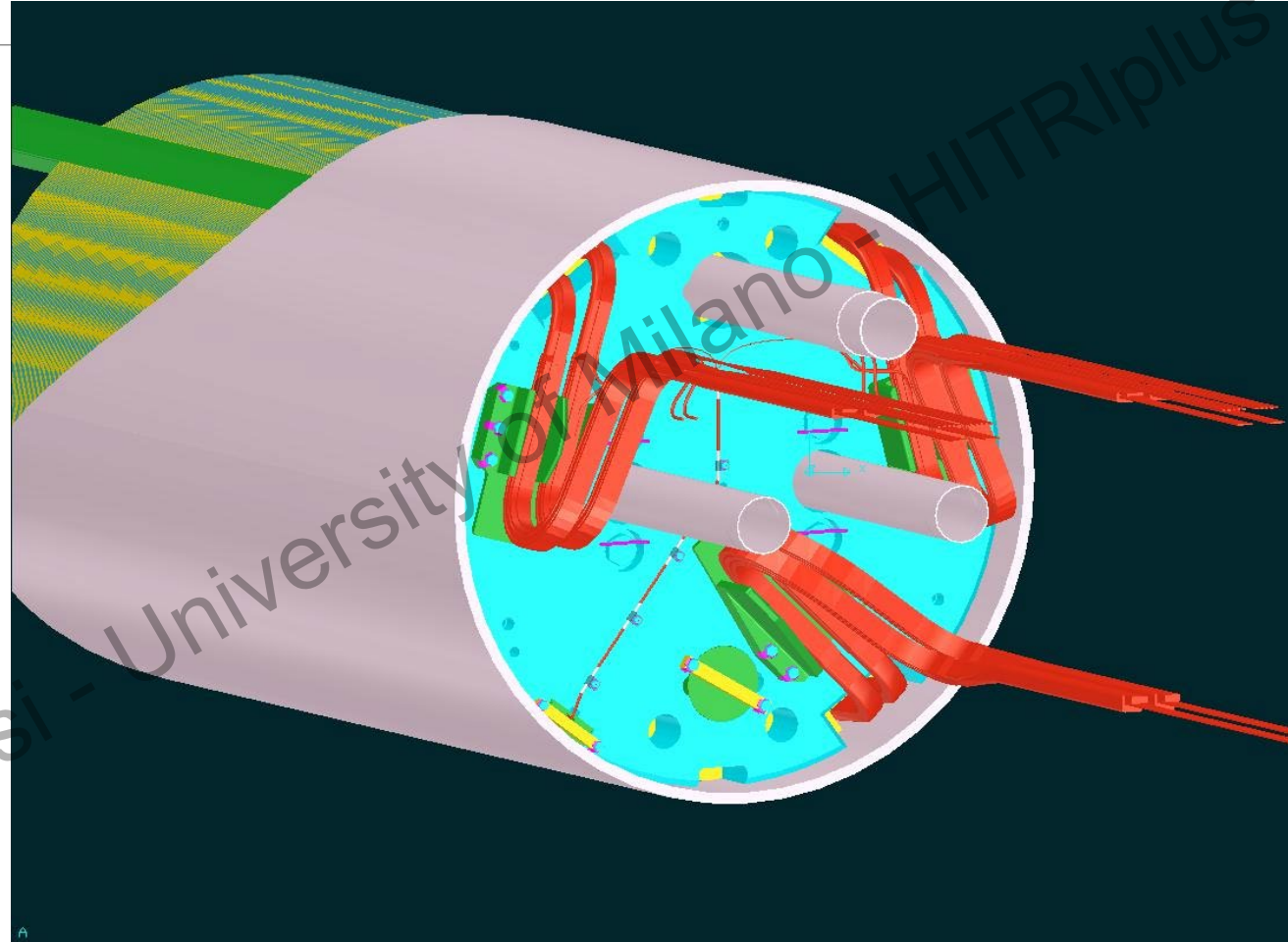
LHC MB-end part Bus Bars positioning



LHC MB -end part Shrinking cylinder



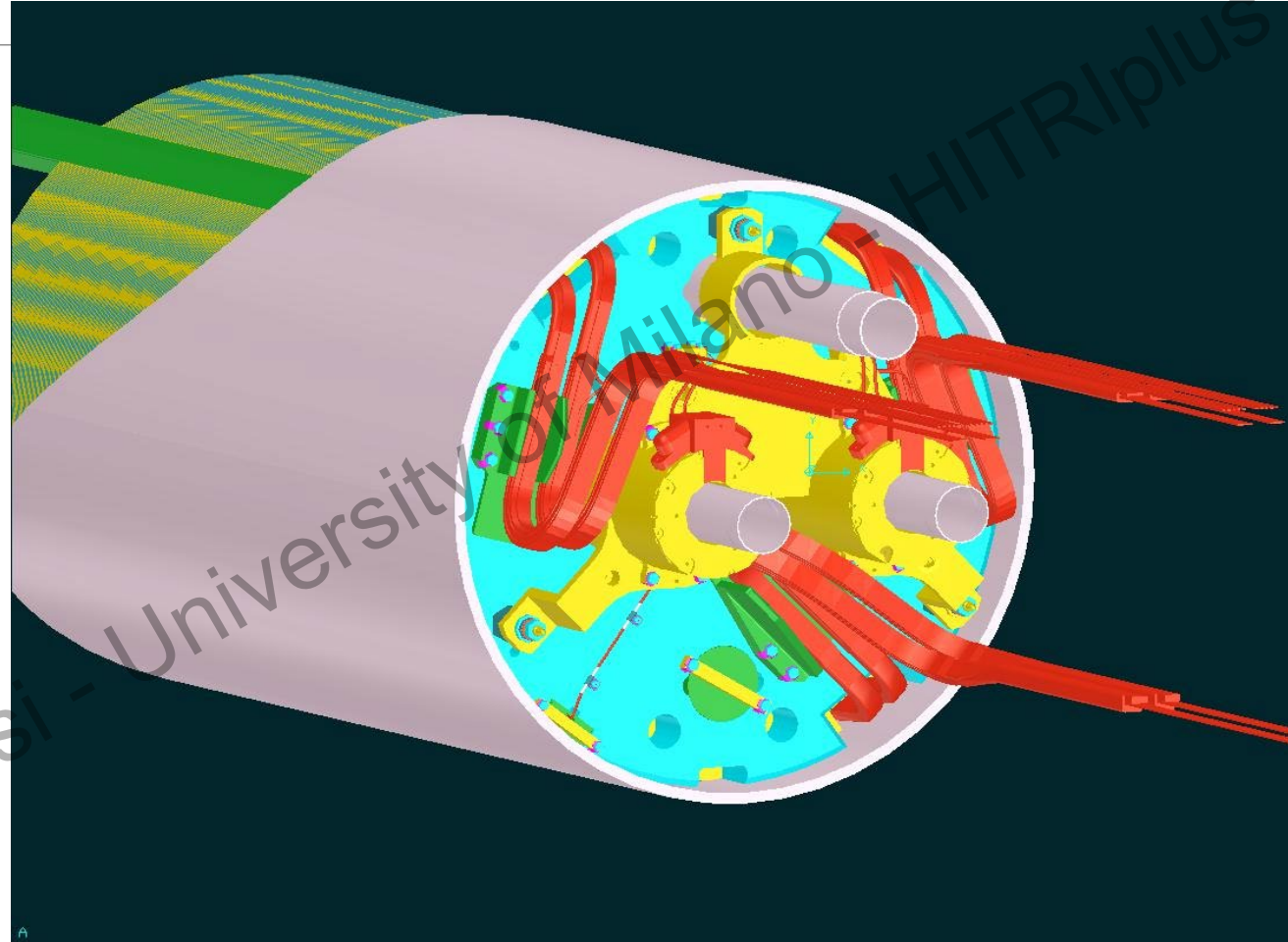
LHC Main Dipole -end part Cu HXT



LHC Main Dipole -end part Corrector Magnets (spool pieces)

Assembly in
CMAs is purely
mechanical

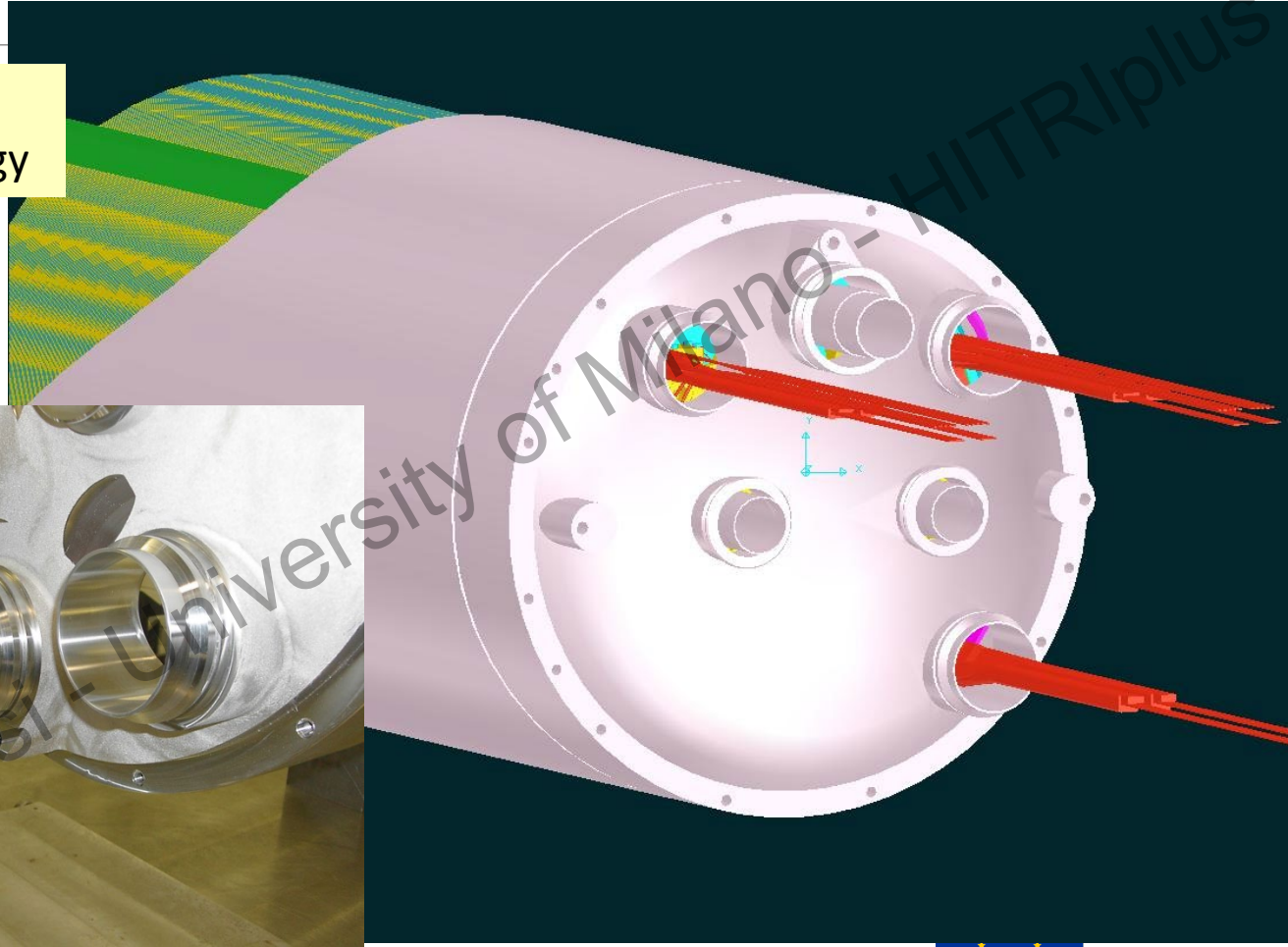
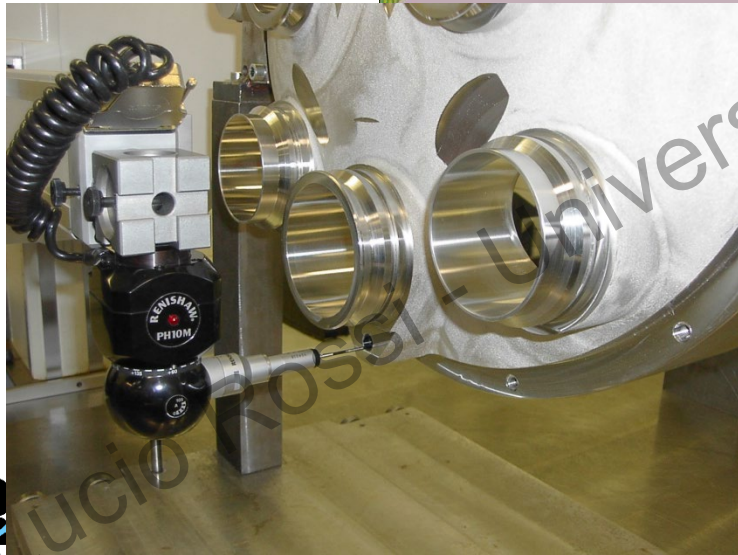
(tolerances of
B axis wrt
mech. frame
given by
supplier



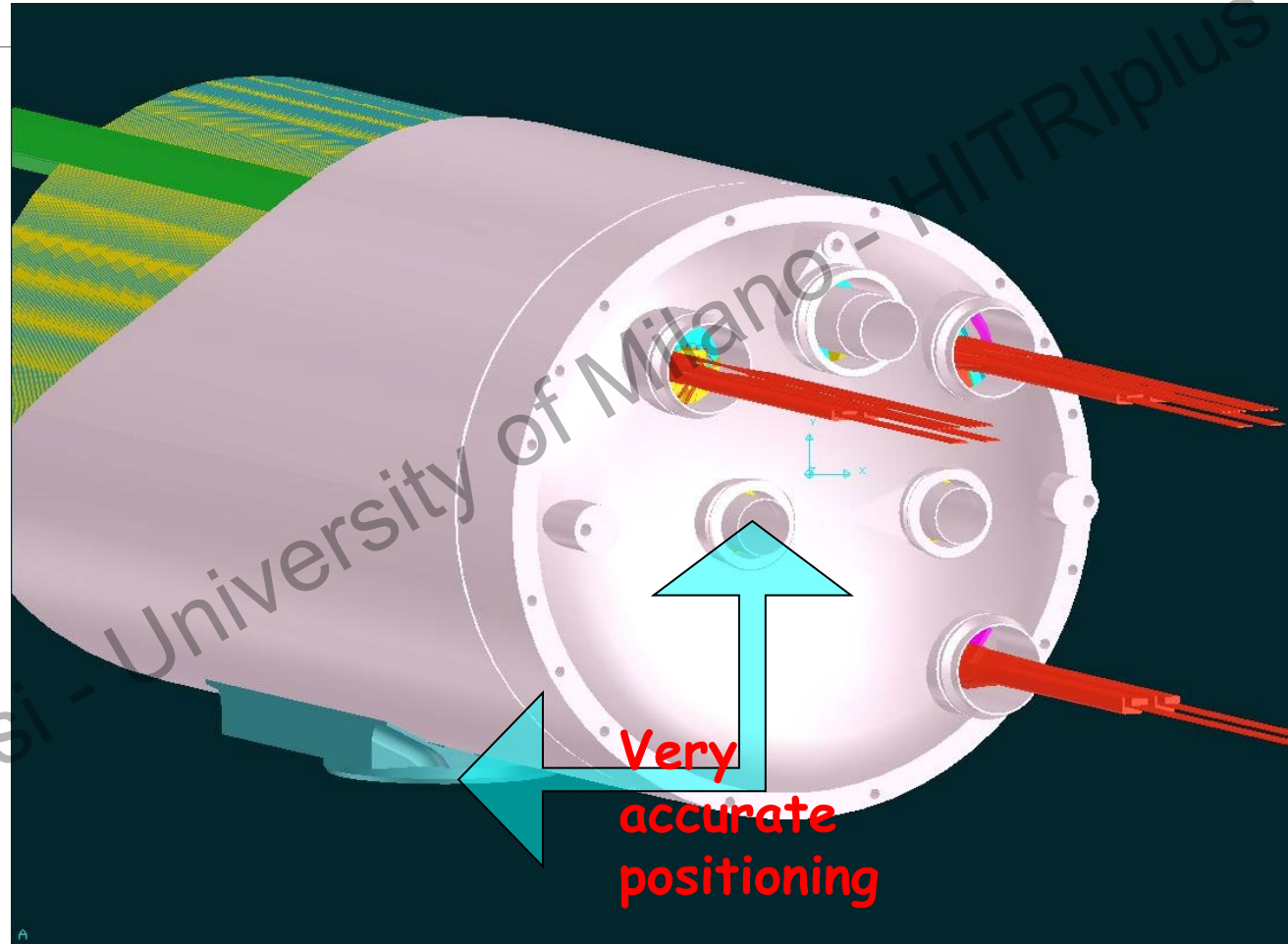
LHC Main Dipole -end part

End covers

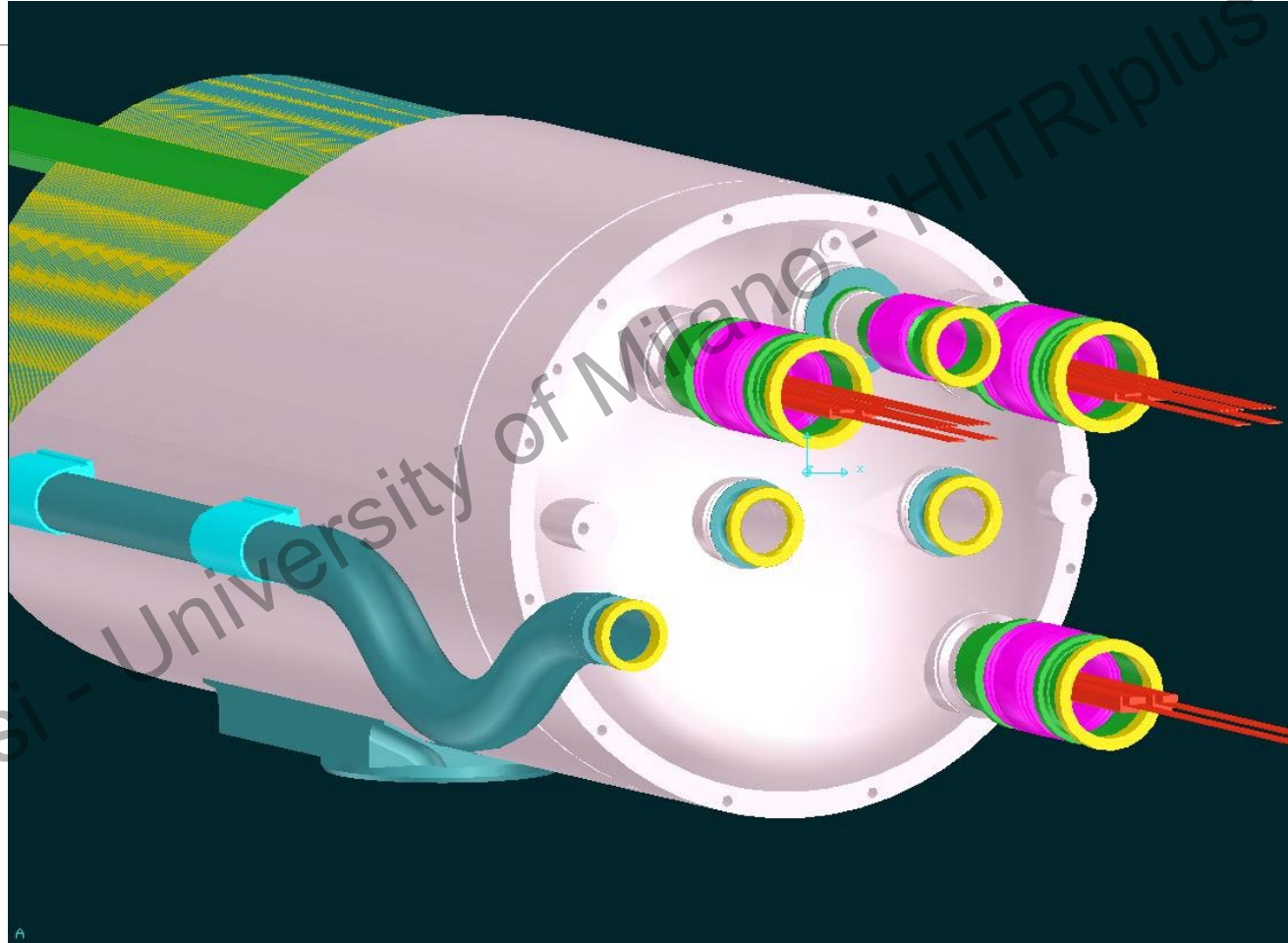
One supplier
Powder technology



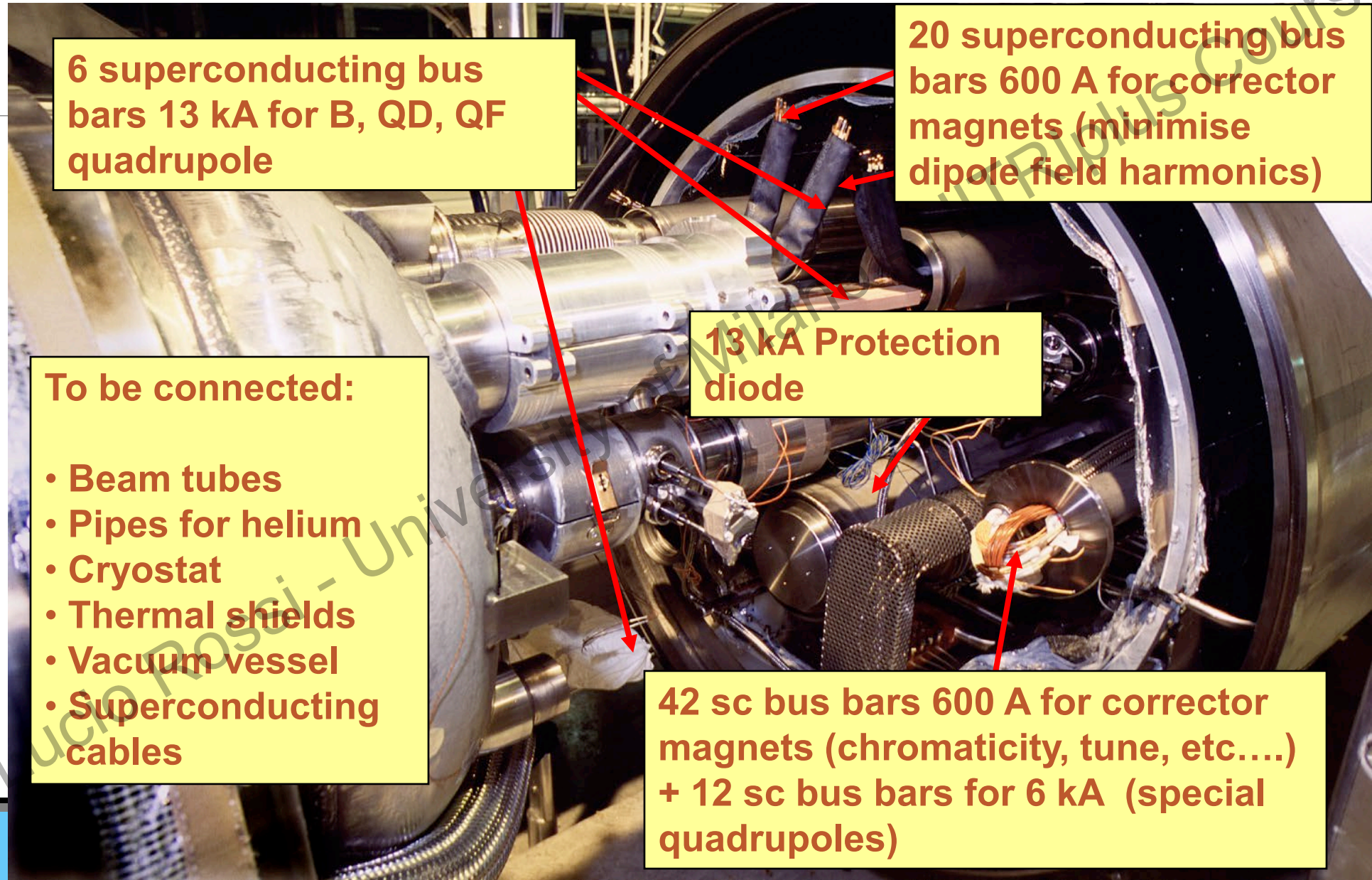
LHC Main Dipole -end part « Cold foot »



LHC Main Dipole -end part Bellows and N-line



Interconnection between SC LHC dipoles



6 superconducting bus bars 13 kA for B, QD, QF quadrupole

20 superconducting bus bars 600 A for corrector magnets (minimise dipole field harmonics)

13 kA Protection diode

To be connected:

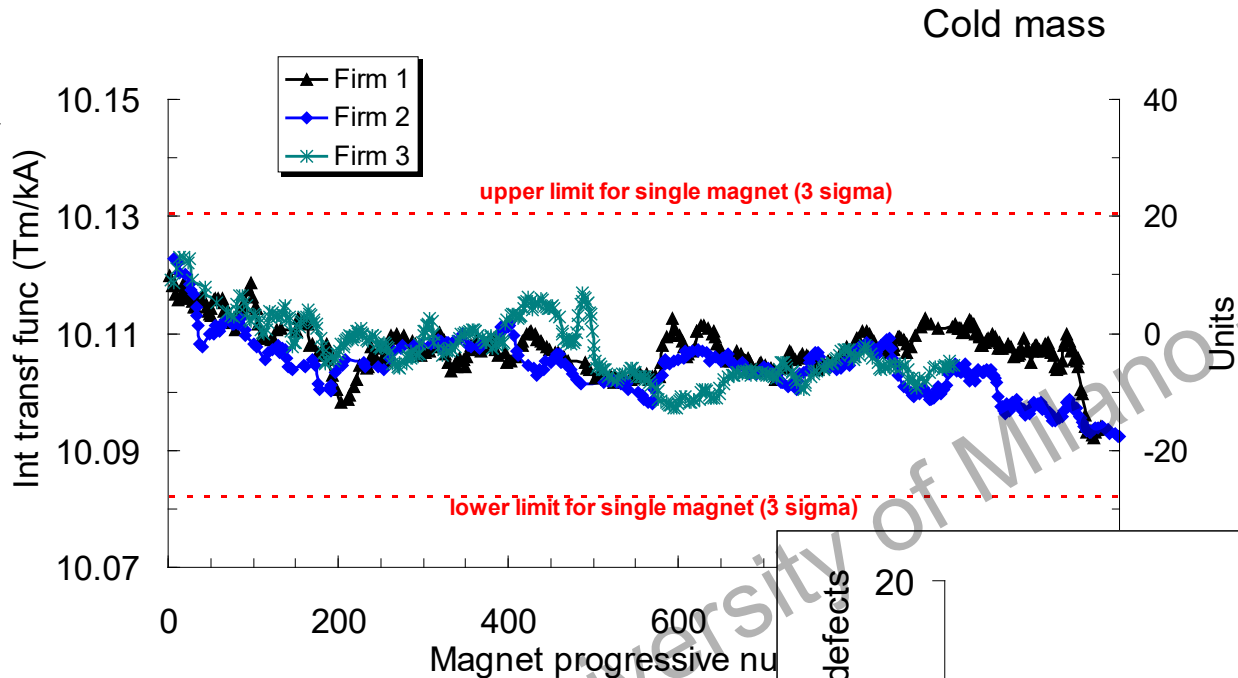
- Beam tubes
- Pipes for helium
- Cryostat
- Thermal shields
- Vacuum vessel
- Superconducting cables

42 sc bus bars 600 A for corrector magnets (chromaticity, tune, etc....) + 12 sc bus bars for 6 kA (special quadrupoles)

there is a series...

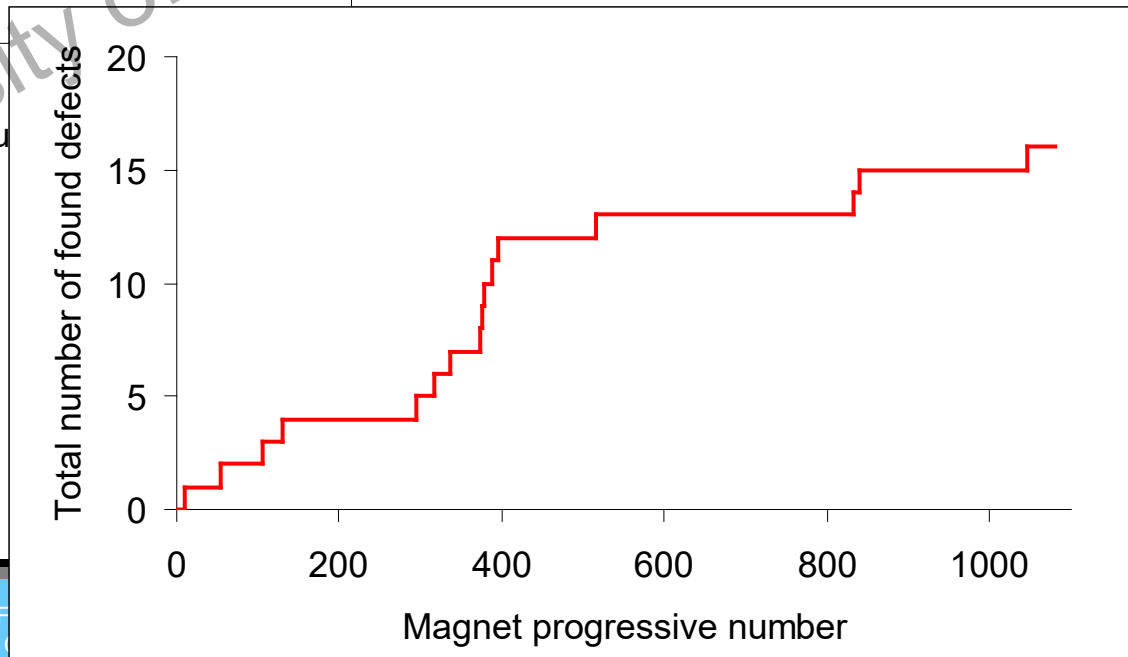


Steering production (and check assembly) Field Measurements - CERN supply

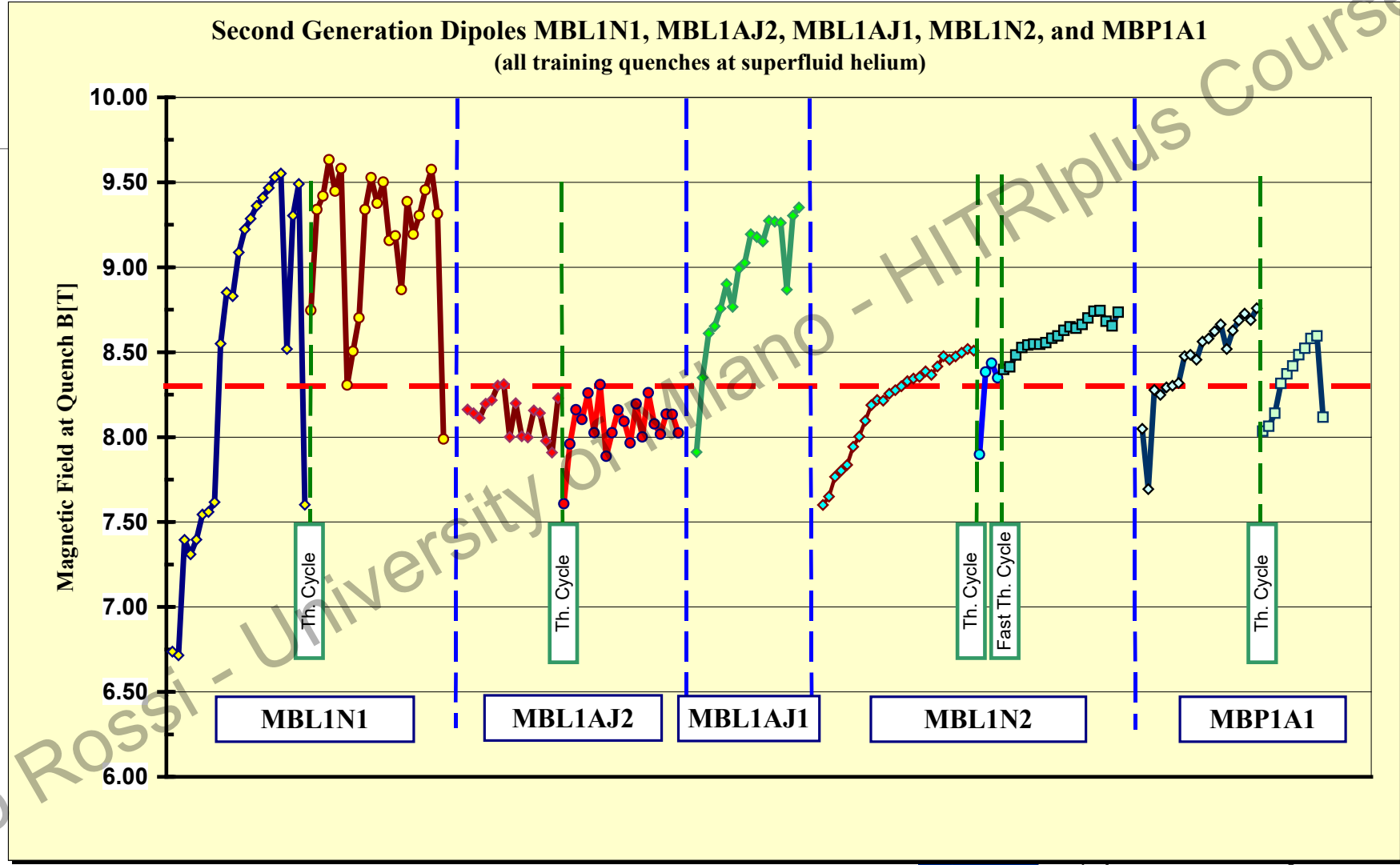


Introduced first to steer the FQ toward beam dynamics targets. Note the uniformity among different manufacturers

It has also helped to detect a number of defects.
It has also been used to detect subtle electrical shorts



Cold test at 4.2 K or 1.9 K: training curve...



Why SC Magnets quench?

Tiny energy release (disturbances) suffice to lose SC status

Time


Space

	Point	Distributed
<i>Transient</i>	Energy (J)	Energy/Vol. (J/m^3)
<i>Continuous</i>	Power (W)	Power/Vol. (W/m^3)

Continuous Distributed Perturbances:

- AC losses (hysteretic and coupling losses, eddy currents)
- Intrinsic dissipation due to smooth transition (I_{op} too near or above I_c !)
- Thermal load (vacuum degradation,...). This could be a serious effect in cryocooled system.

These perturbations are usually predictable and estimate must be done at design level

 coupling losses can depends on interstrand resistance, i.e. on manufacture technique and on prestress and e.m. forces \Rightarrow more difficult to evaluate

Continuous Point Perturbances:

- Joints inside coils
- Release of mechanical energy (hysteresis of the stress-strain relation)
- Localised heat input (suspension rods with bad thermal anchoring)

These effects are well understood and predictable (it does not mean easy to cure !)

The big enemy: The transient perturbances

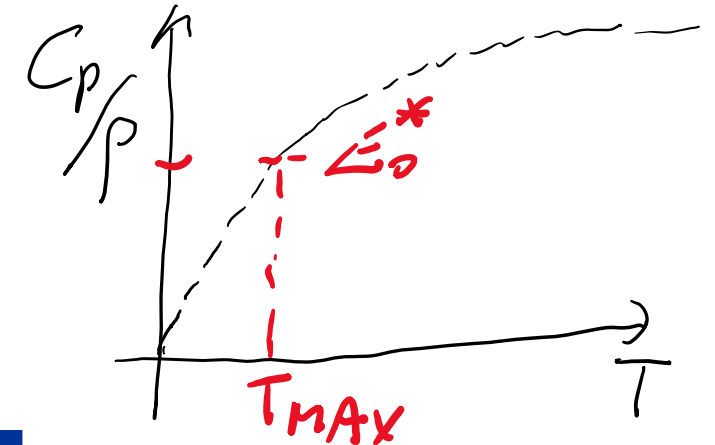
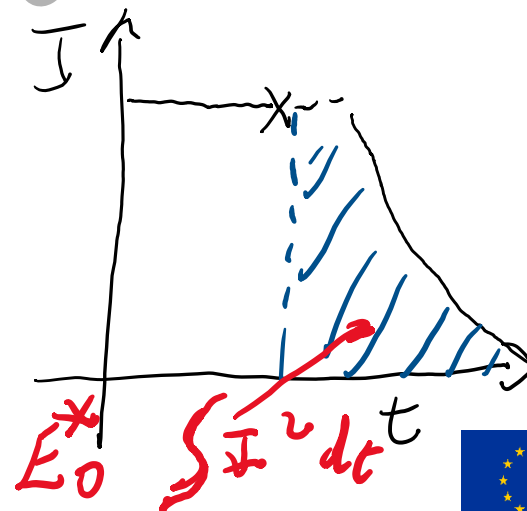
Transient Distributed/Point Perturbances

- Flux jumps. This effect is cured almost definitely for NbTi. Effects could be seen on NbSn with very high current density and very large effective filament diameter. This effect can be detected at low field, during current ramp.
- Mechanical origin: movements, friction, sudden release of elastic energy...
- crack in the resin

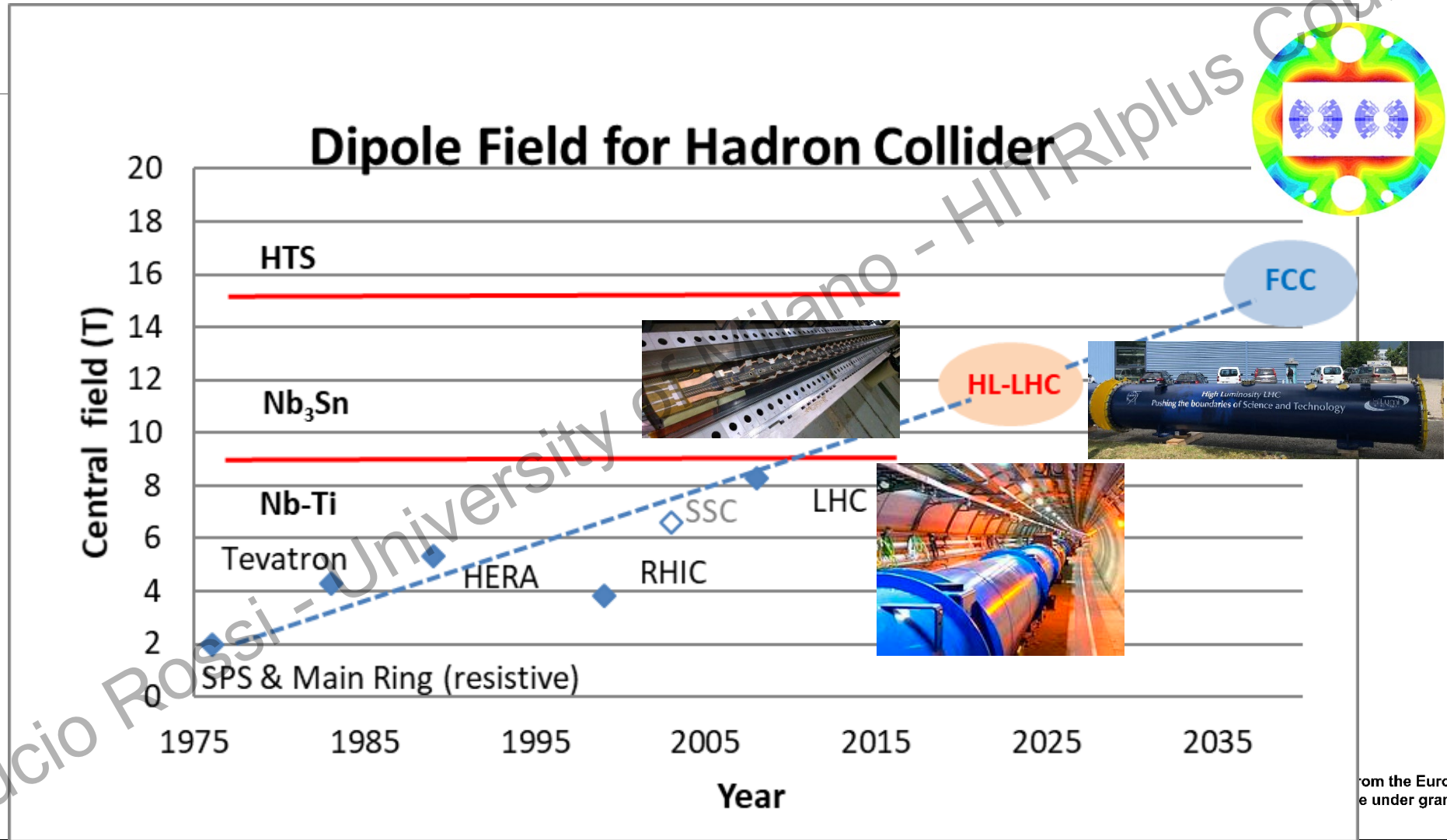
Basically these last two mechanism are now understood, in principle, and acoustic emission experiments did prove it almost visually.

Still they are less predictable and more difficult to avoid. They depend on magnet geometry, material properties, local conditions and on many details. They can depend on magnet history (previous quench, overheating, thermal induced stress, etc.)

$$\int_{T_{op}}^{T_{hot}} \frac{C}{\rho} dT = \frac{1}{A_{cu} A_{tot}} \int_0^{\infty} I_{op}^2 dt$$



HiLumi LHC: preparing technology for next big steps



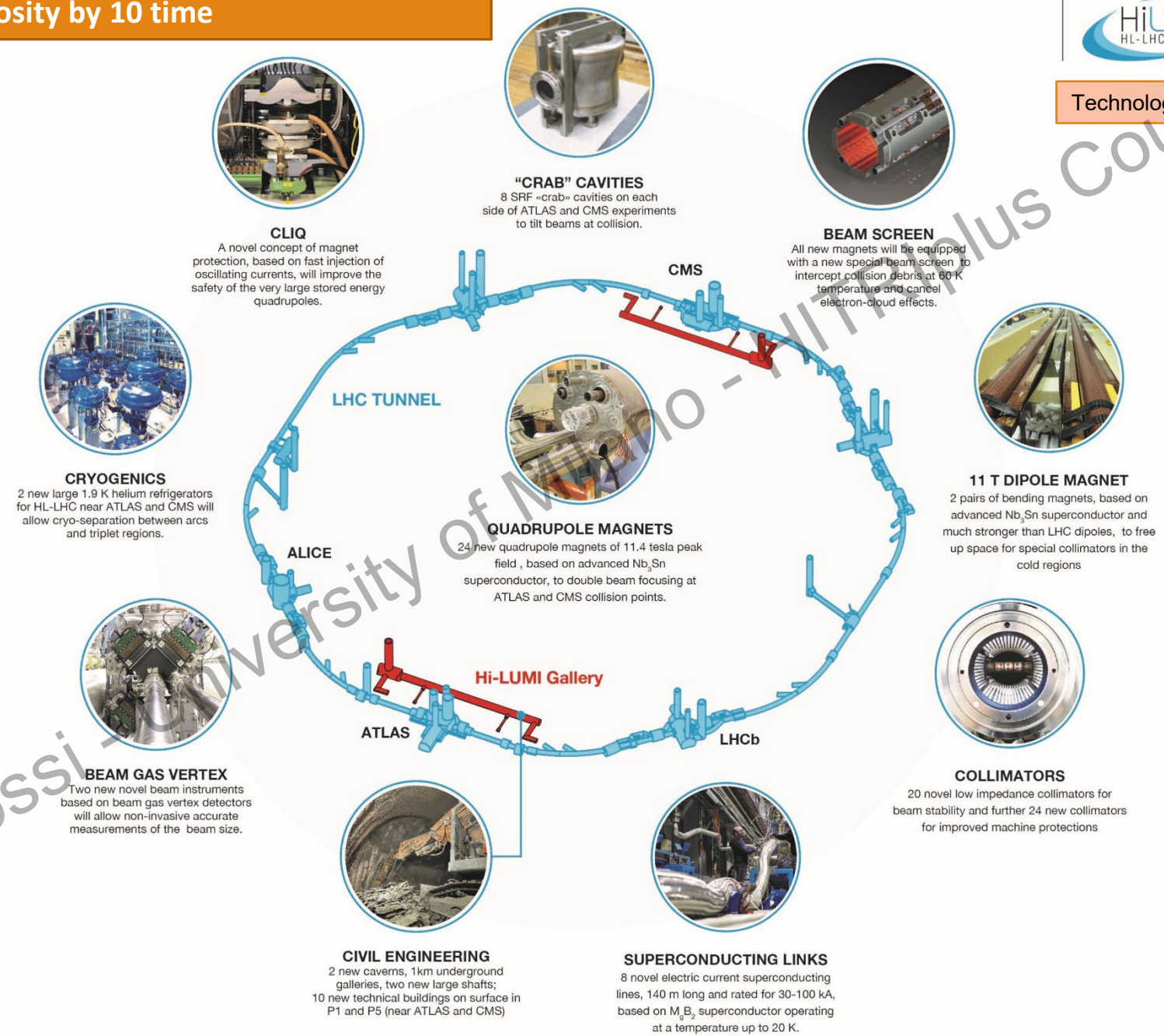
Actually the next step started in 2010...**The High Luminosity LHC project aims at increase the luminosity by 10 time**



CERN 2019

Technology landmarks

LHC is already highly optimized... an upgrade need a broad spectrum of new technologies, especially (not only) for magnets...
HL-LHC is a technology intensive project!



CLIQ
A novel concept of magnet protection, based on fast injection of oscillating currents, will improve the safety of the very large stored energy quadrupoles.



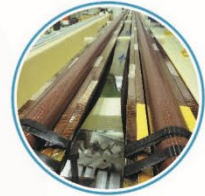
“CRAB” CAVITIES
8 SRF «crab» cavities on each side of ATLAS and CMS experiments to tilt beams at collision.



BEAM SCREEN
All new magnets will be equipped with a new special beam screen to intercept collision debris at 80 K temperature and cancel electron-cloud effects.



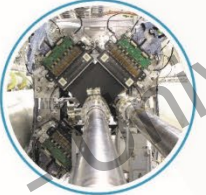
CRYOGENICS
2 new large 1.9 K helium refrigerators for HL-LHC near ATLAS and CMS will allow cryo-separation between arcs and triplet regions.



11 T DIPOLE MAGNET
2 pairs of bending magnets, based on advanced Nb₃Sn superconductor and much stronger than LHC dipoles, to free up space for special collimators in the cold regions



QUADRUPOLE MAGNETS
24 new quadrupole magnets of 11.4 tesla peak field, based on advanced Nb₃Sn superconductor, to double beam focusing at ATLAS and CMS collision points.



BEAM GAS VERTEX
Two new novel beam instruments based on beam gas vertex detectors will allow non-invasive accurate measurements of the beam size.



COLLIMATORS
20 novel low impedance collimators for beam stability and further 24 new collimators for improved machine protections

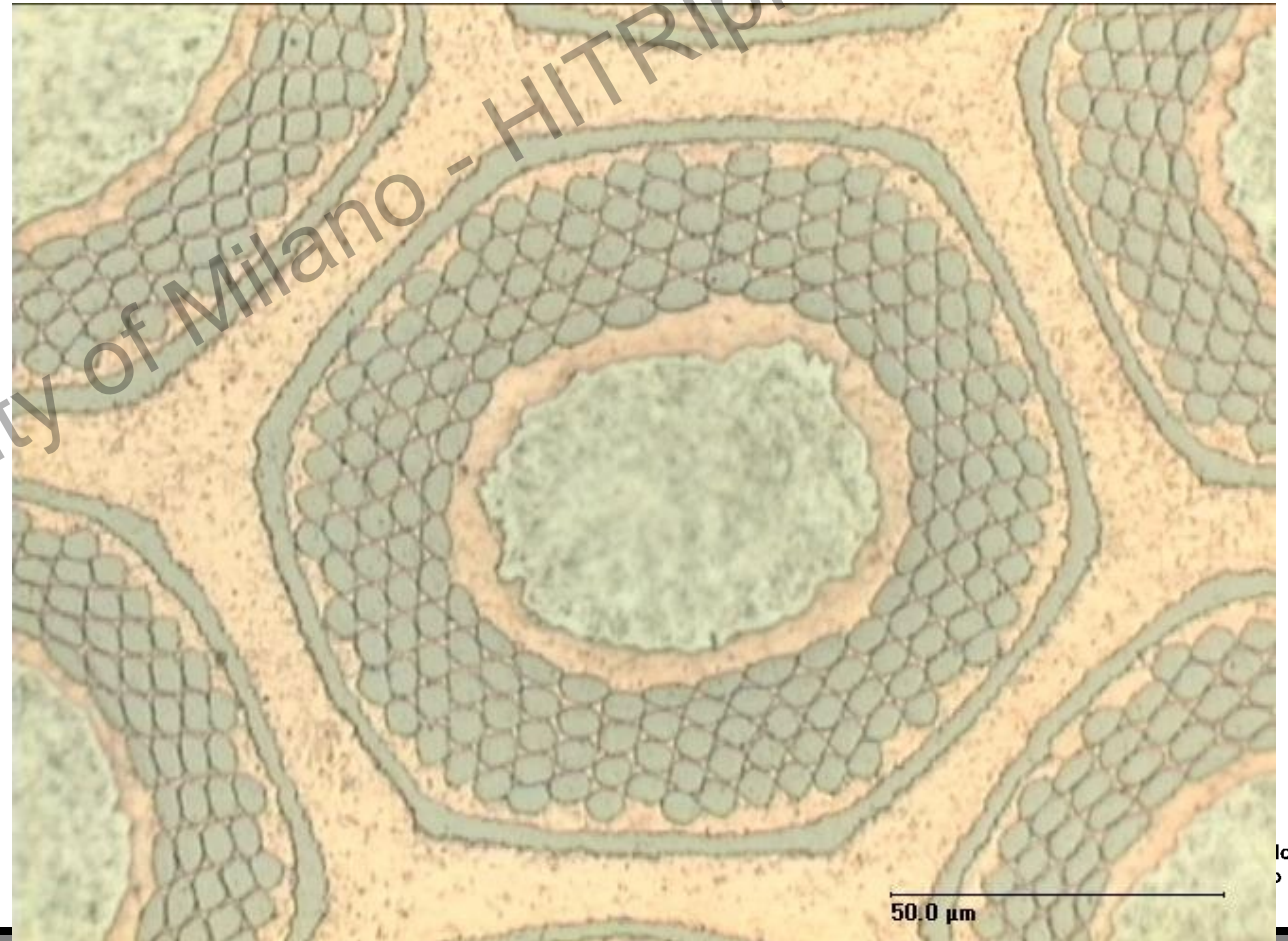
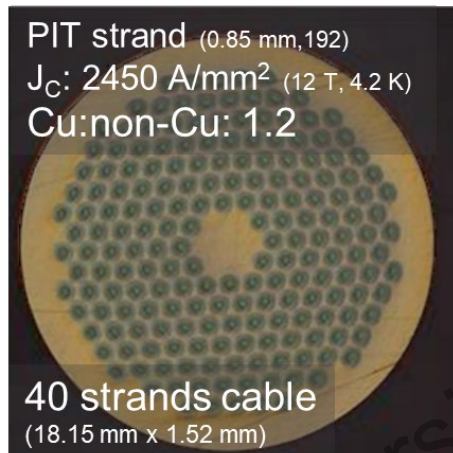
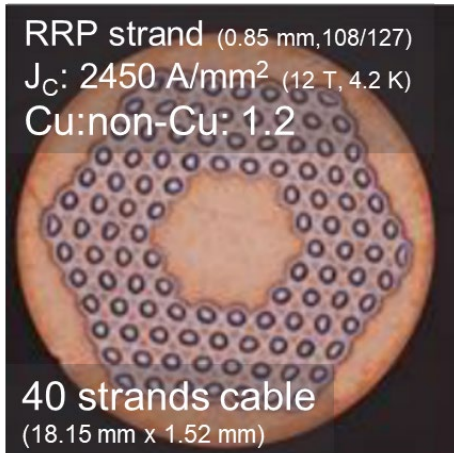


CIVIL ENGINEERING
2 new caverns, 1km underground galleries, two new large shafts; 10 new technical buildings on surface in P1 and P5 (near ATLAS and CMS)



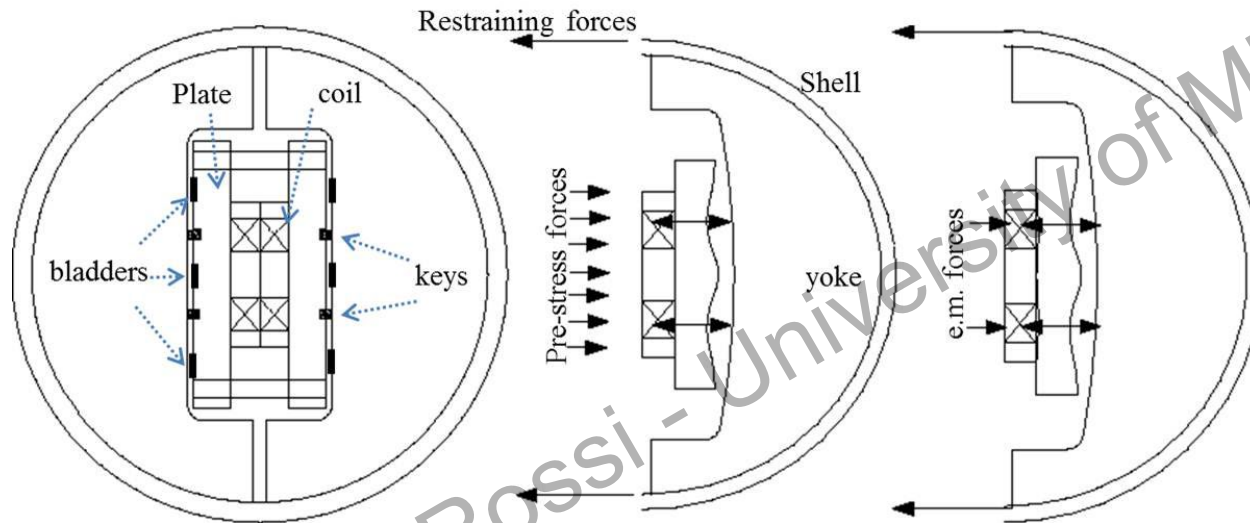
SUPERCONDUCTING LINKS
8 novel electric current superconducting lines, 140 m long and rated for 30-100 kA, based on M₀B₂ superconductor operating at a temperature up to 20 K.

Nb₃Sn : high **J_e** also at 12- 16 T but it is brittle and needs thermal treatment of whole coil @ 700 °C !

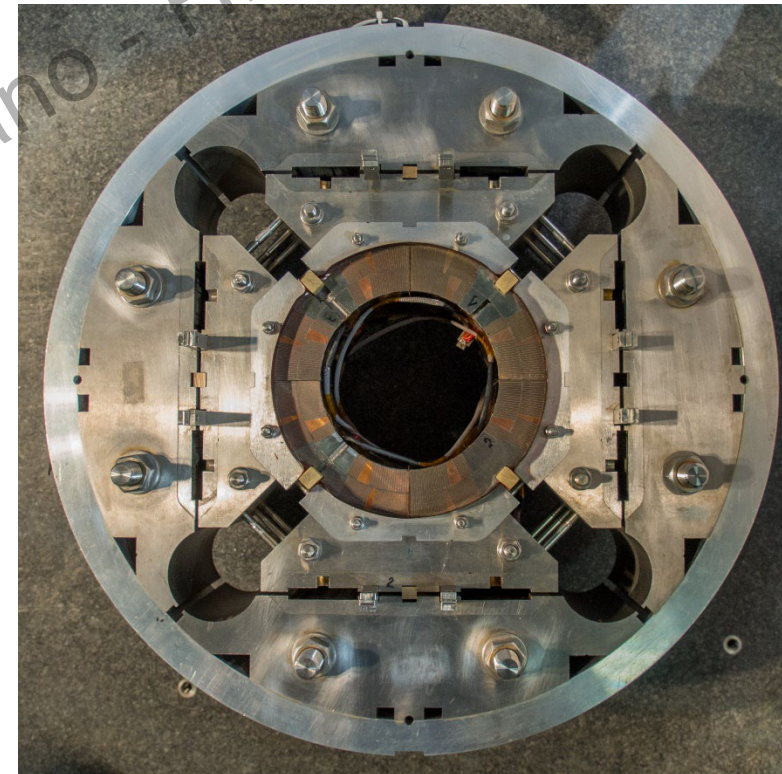


New way to keep the stress : HiLumi paves the way to reach a controlled precompression of ~ 130 MPa

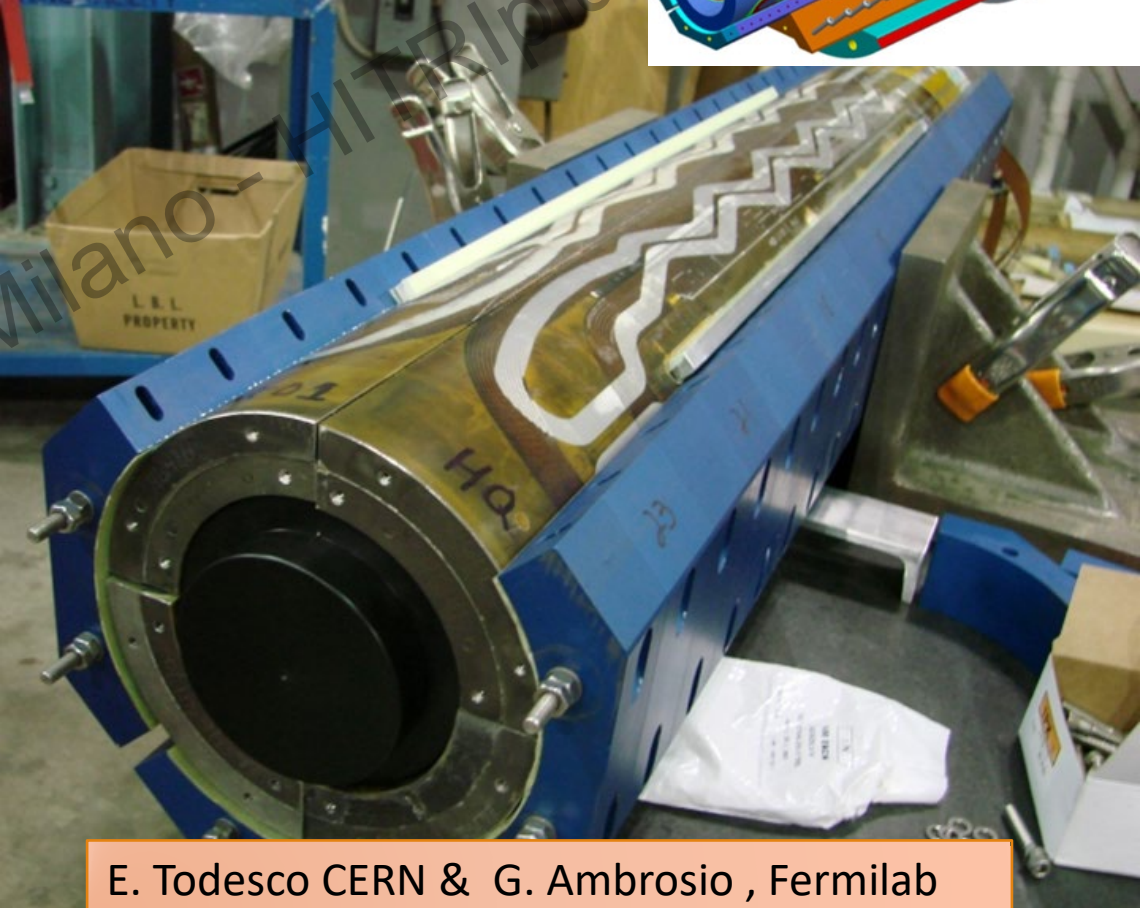
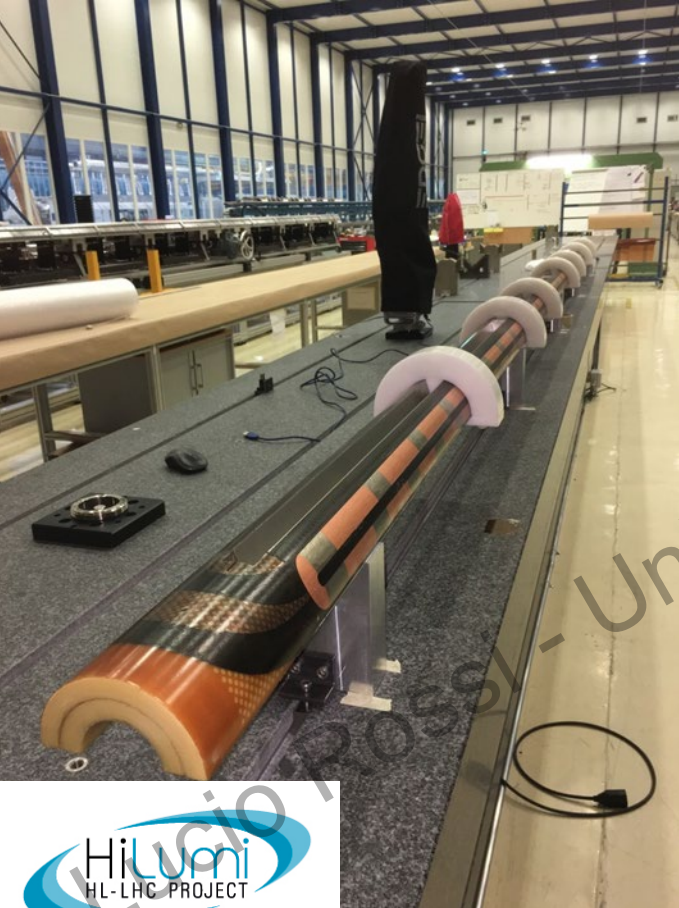
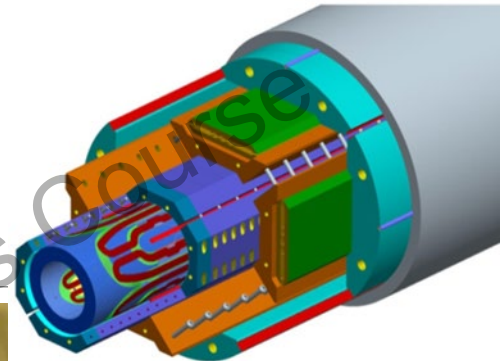
New concept: bladders and keys



Cross section of the Quad for HiLumi



The HiLumi Quad: 12 T in $\varnothing=150$ mm ~ equivalent to a dipole 15 T- 50 mm

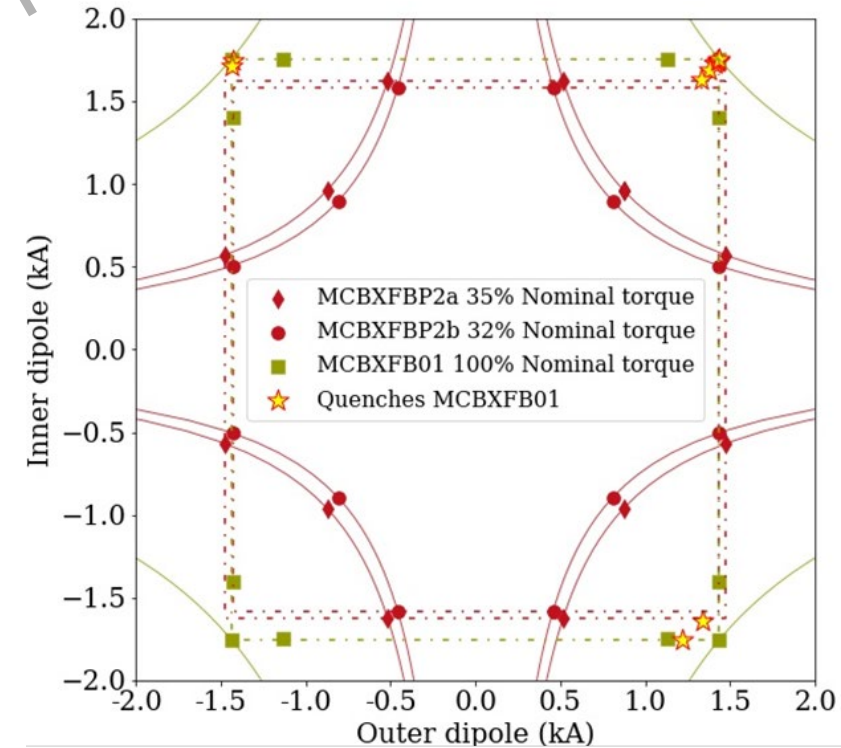
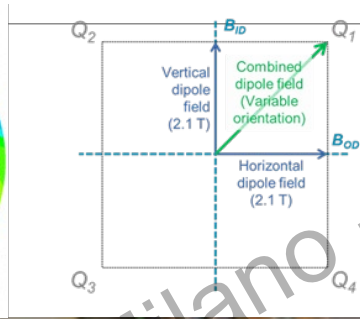
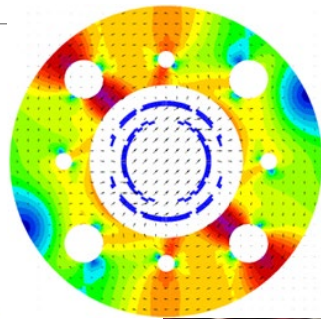
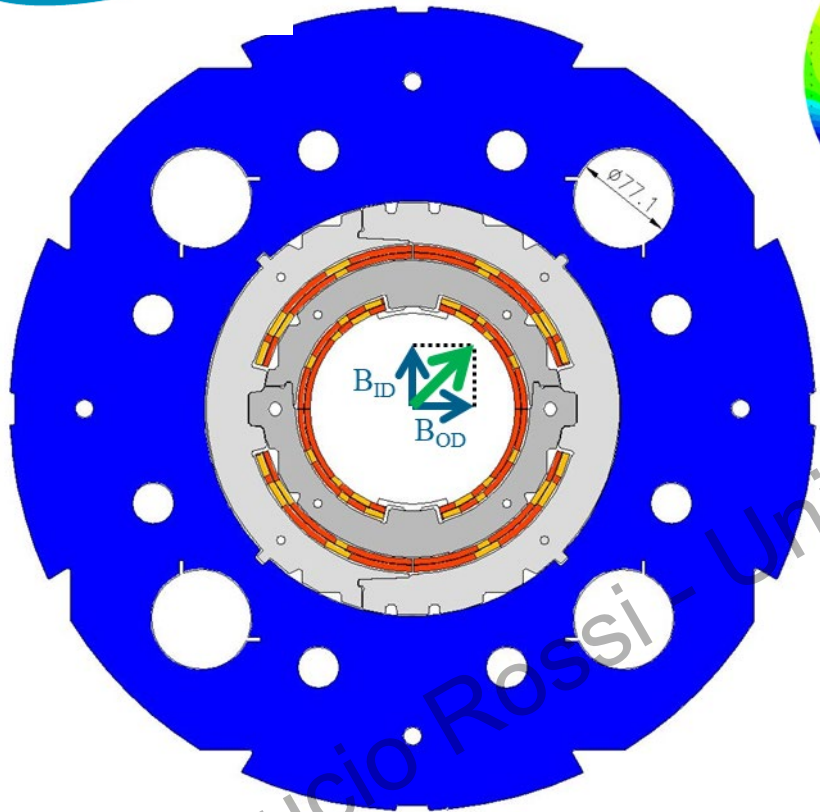


E. Todesco CERN & G. Ambrosio , Fermilab



2020
548

Other development with classical Nb-Ti: nested X-Y dipole Steering – Bending with compact magnet, space is precious...



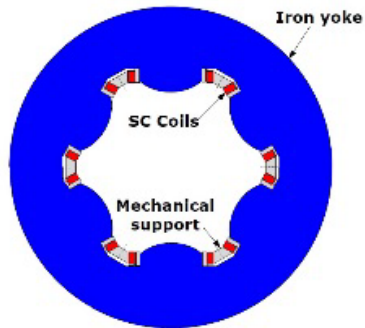
Courtesy S.Ferradas (CERN)

Thanks to F. Toral, Ciemat

Superferric magnet technology: 2-3 T iron driven by SC coils very convenient in terms of performance/cost

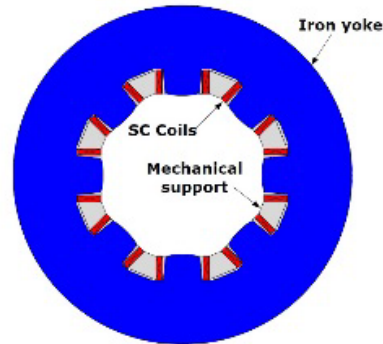
Thanks to M. Statera, INFN-LASA

OD=320 mm



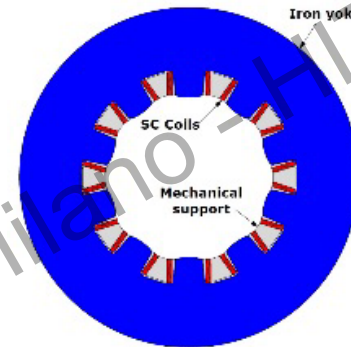
Sextupole

OD=320 mm



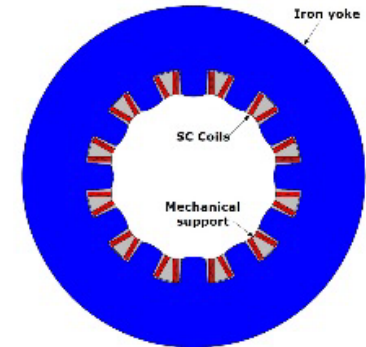
Octupole

OD=320 mm



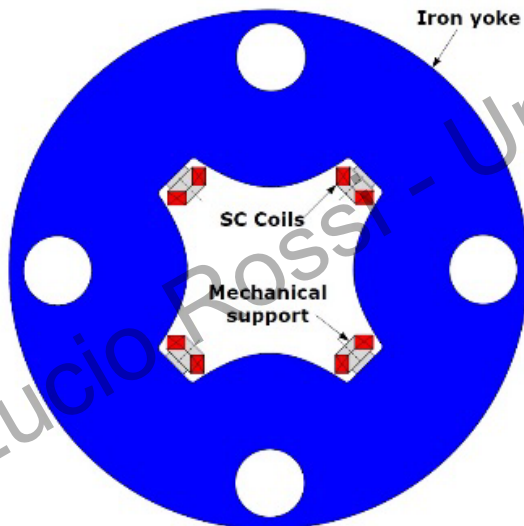
Decapole

OD=320 mm



Dodecapole

OD=460 mm



Skew quad

Physical length:

- 90-120 mm from 6-pole to 10-pole
- 430 mm 12-pole normal
- 840 mm 4-pole skew

Conductor type: NbTi

Peak field on cond.: 2-3 T

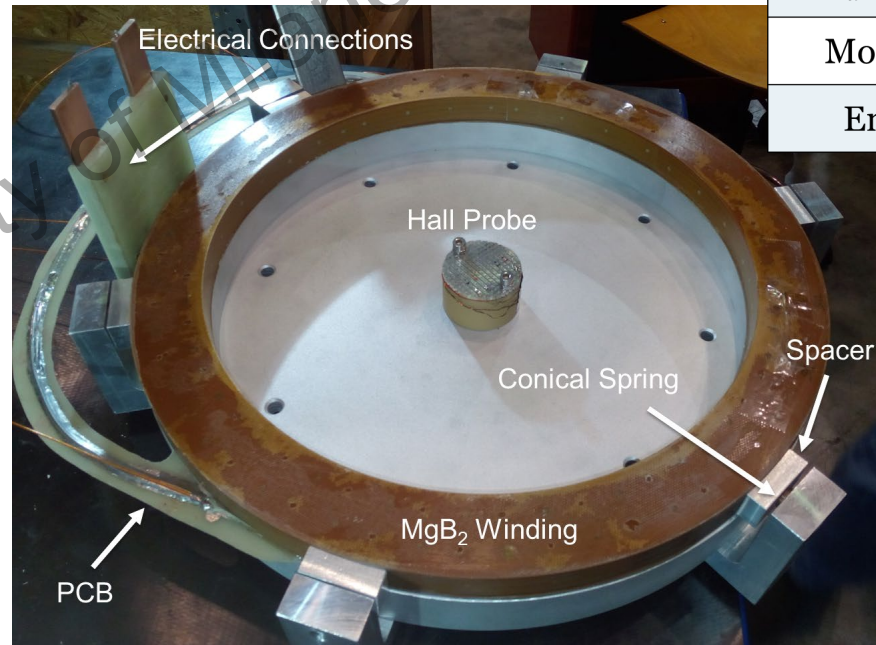
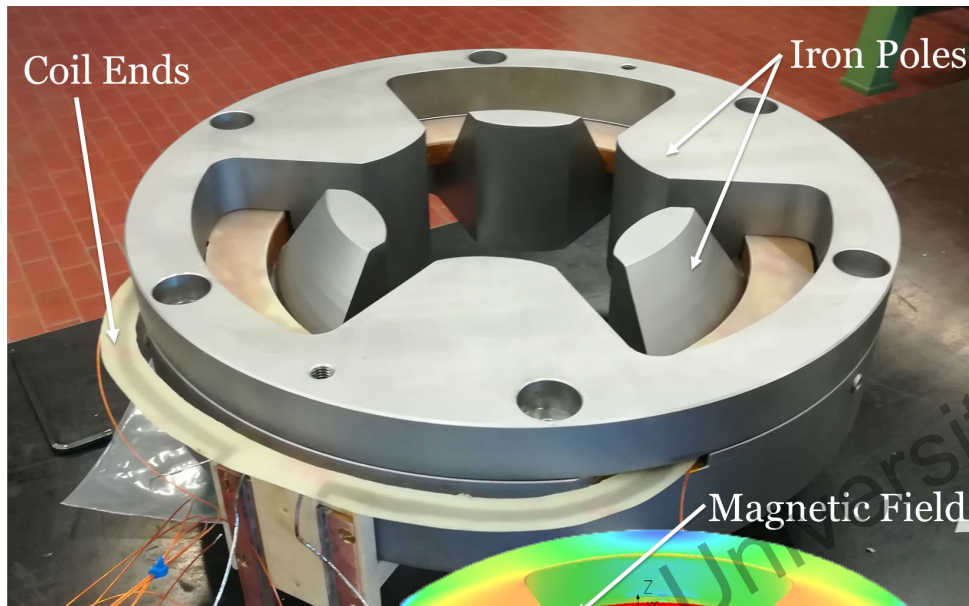
Operating current: 120-180 A

Margin on load line: 40%

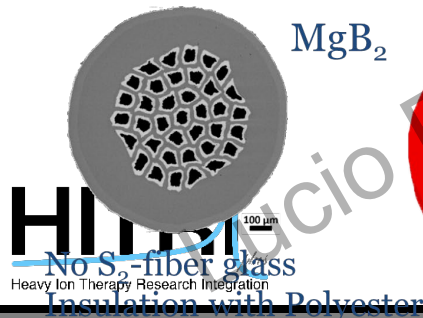


Next step for superferric: going cryogen-free HTS or RSCM design with MgB₂

S. Mariotto and R. Valente
INFN-MI-LASA



	Value
Integrated Field @ R=50 mm	70.06
B _{max} on SC @ I _{op}	1.37 T
B _{max} on Iron @ I _{op}	3.52 T
Module Length	384 mm
Energy @ I _{op}	3754 J

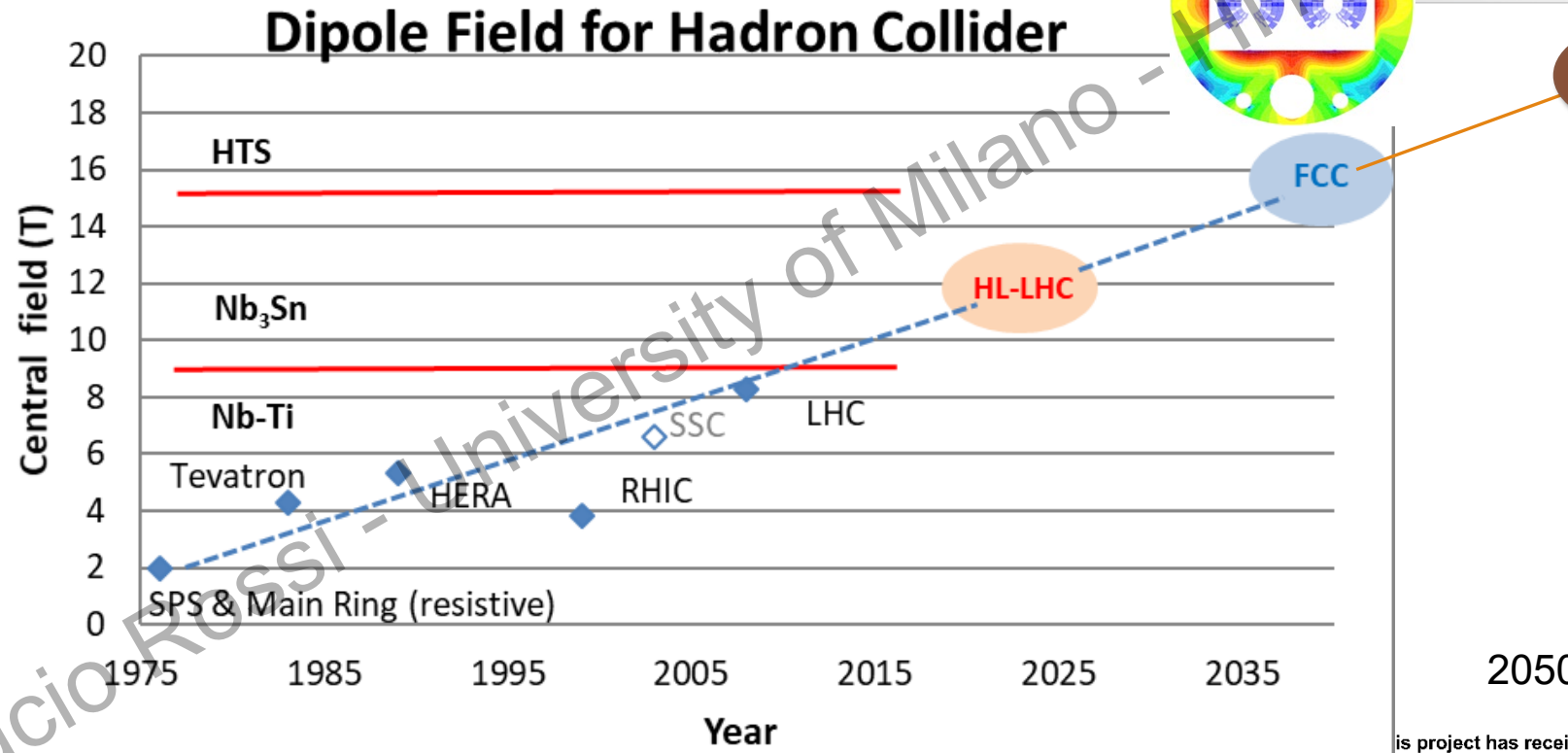
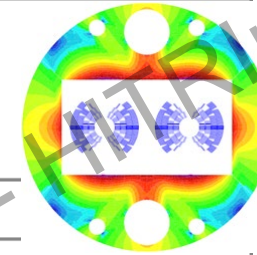


and funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

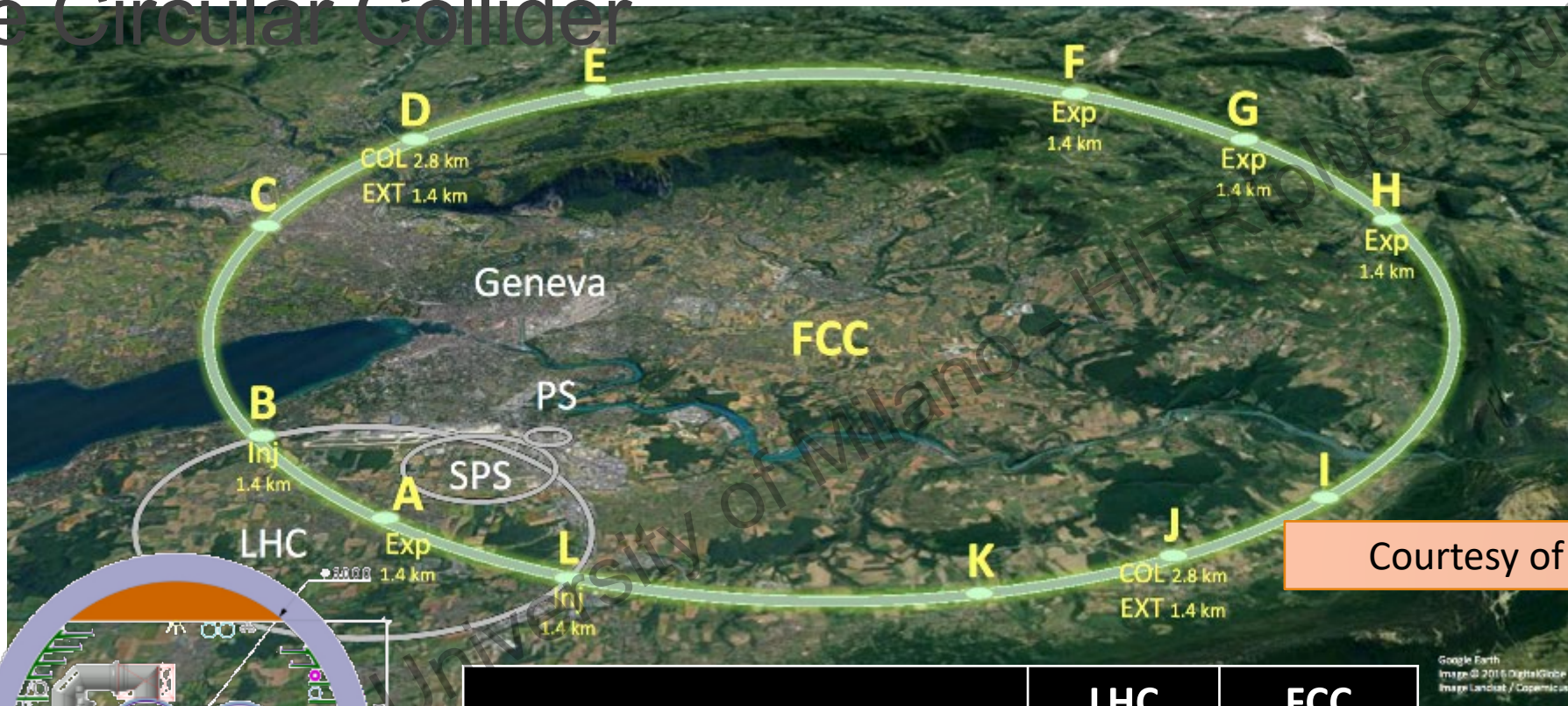
Can we extrapolate linearly from the past beyond HiLumi: 16-20 T FCC? \Rightarrow ELN



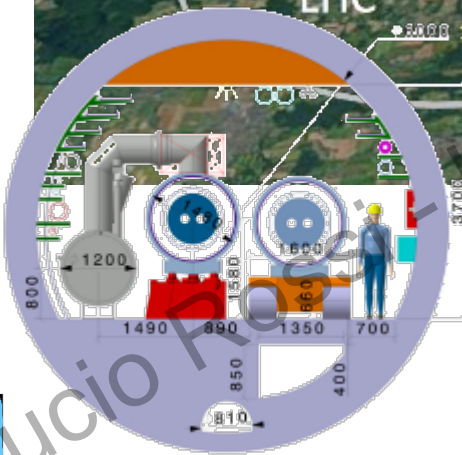
5T HTS demos



Future Circular Collider



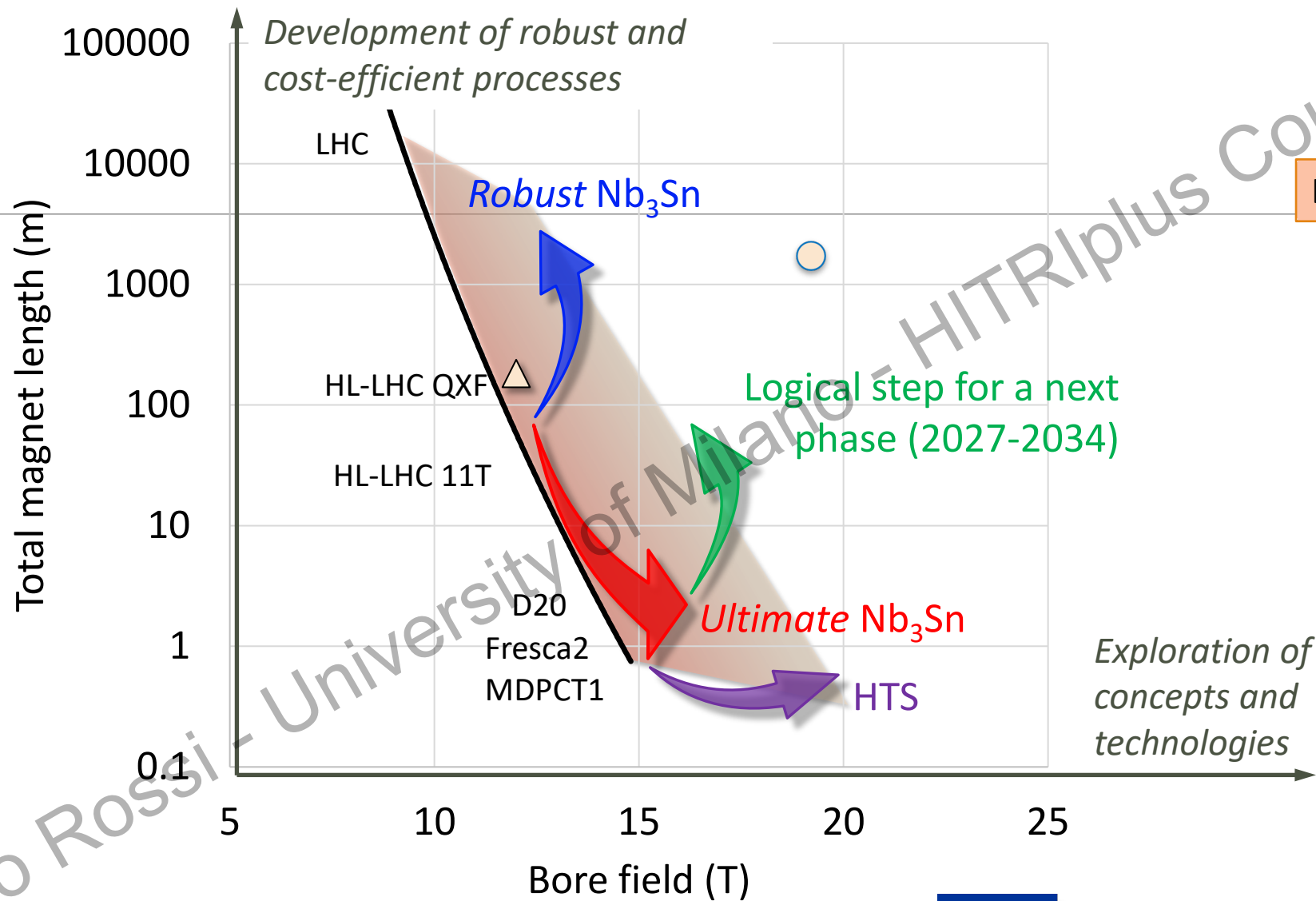
Courtesy of M. Benedikt, FCC



	LHC	FCC
Circumference (km)	26.7	97.5
Dipole field (T)	8.33	16
C.o.M. energy (TeV)	14	100

Received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

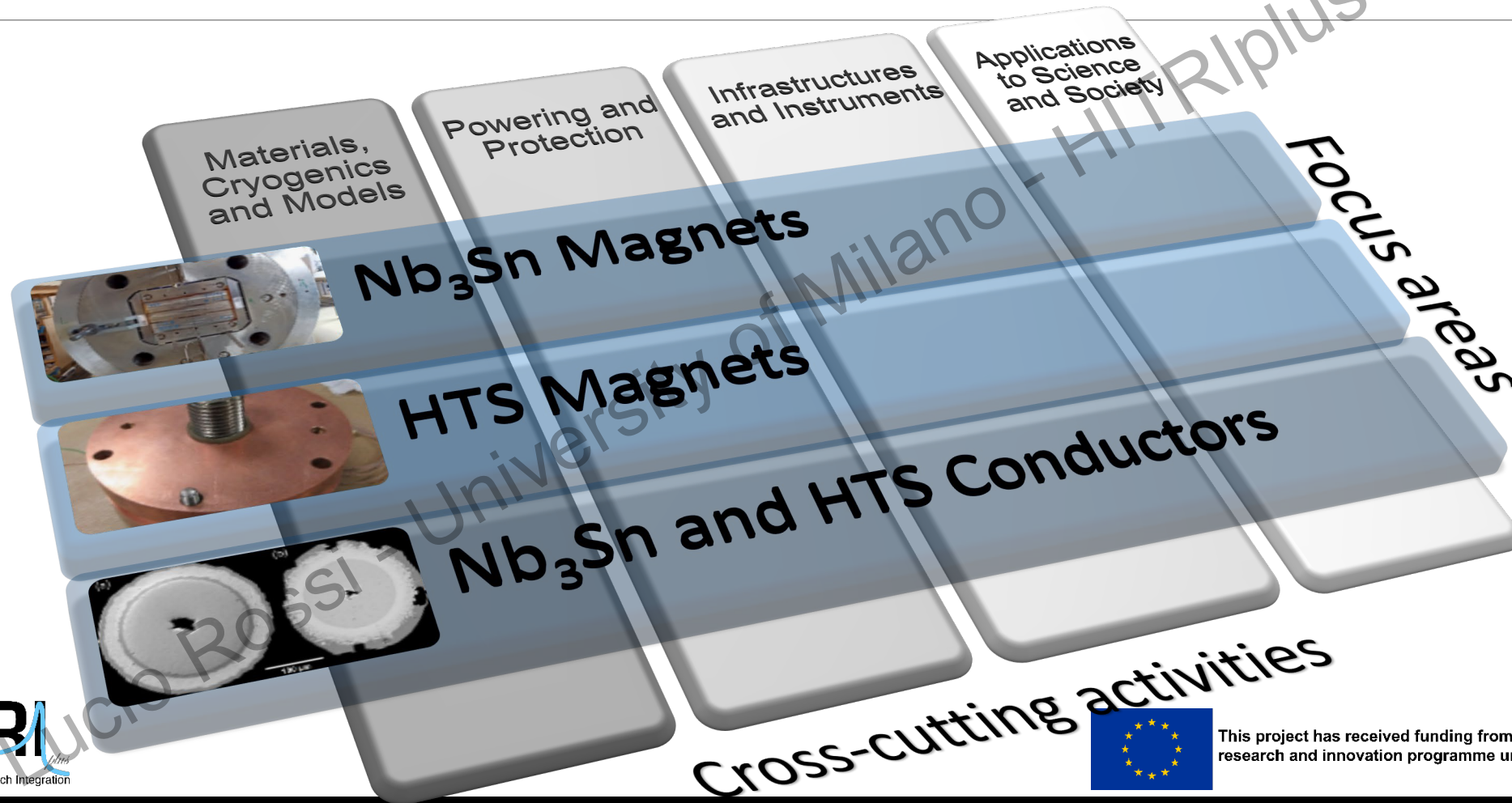




L. Bottura, CERN

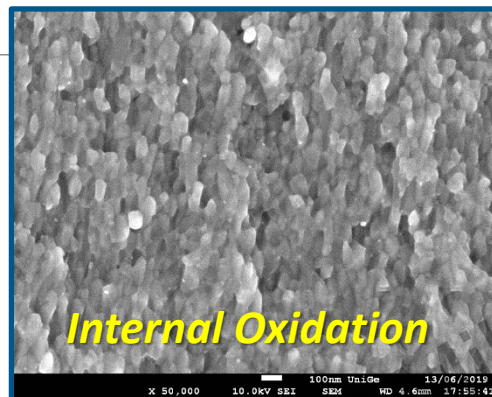
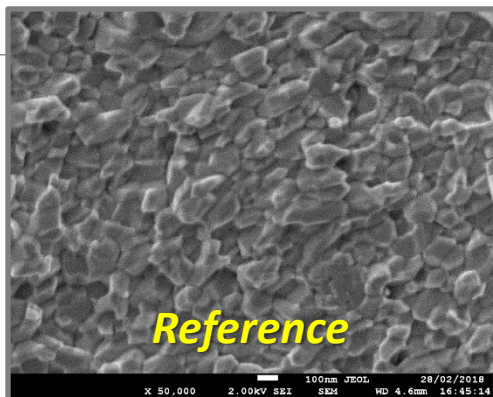


Activities – Topics matrix



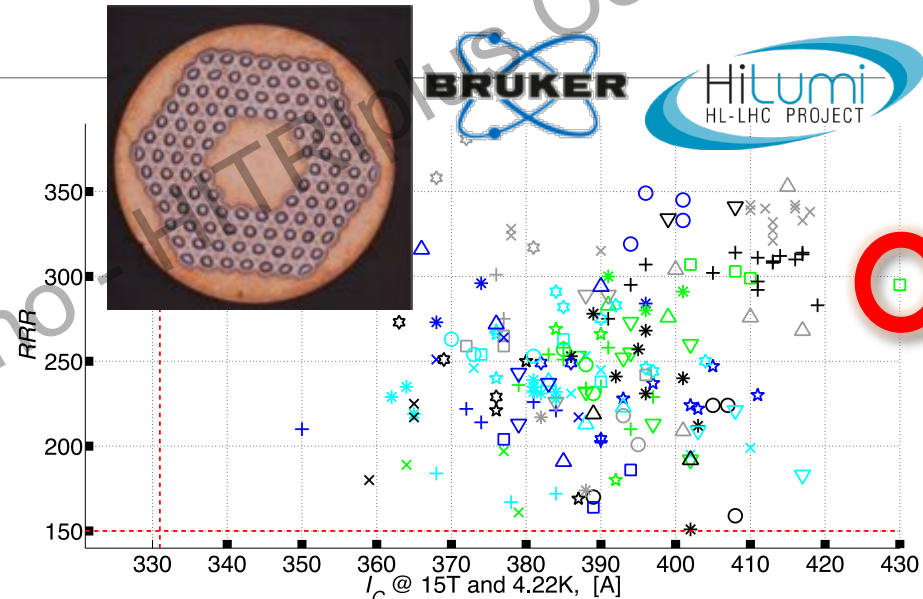
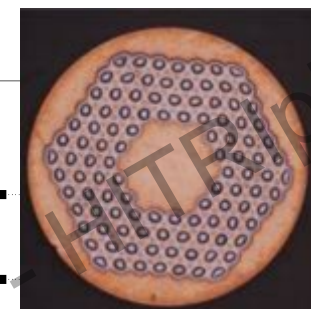
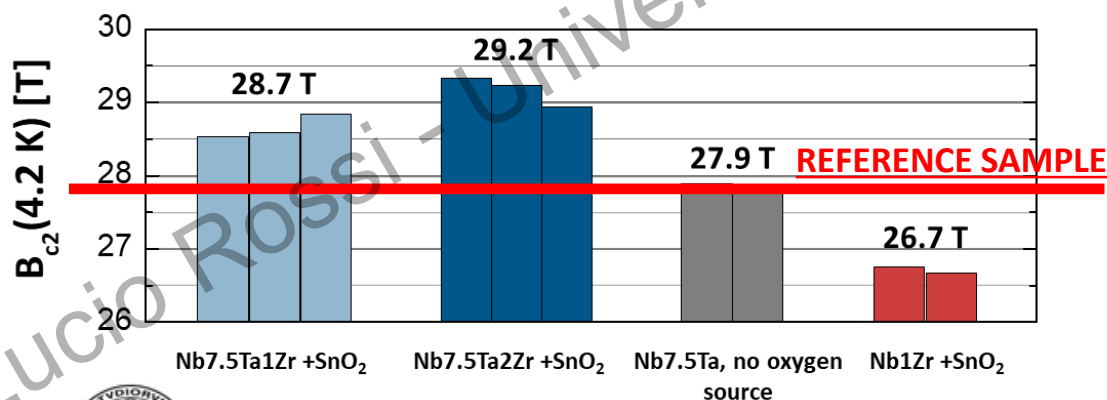
Towards the ultimate performance of Nb₃Sn

Performance target for FCC-hh $J_c(4.2K, 16 T) = 1'500 A/mm^2$



Grain refinement from 120 nm (left) to 60 nm (right)

Enhanced grain boundary pinning



1750 A/mm² @ 15T, 4.2K

≈ 1400 A/mm² @ 16T, 4.2K

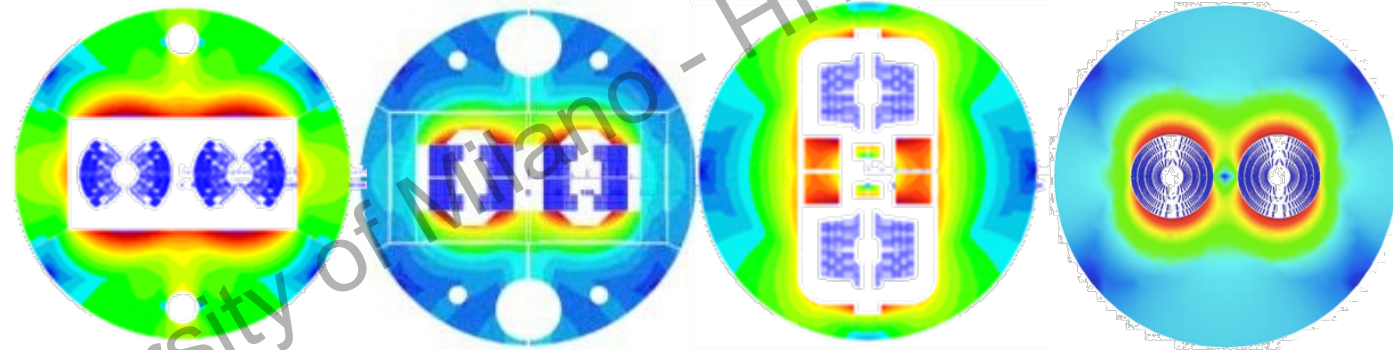
And record-high upper critical field !!

LNCMI-Grenoble

If advanced superconductors are there, magnets come



$OD = 600 \text{ mm}$
 $L = 2 \text{ m}$
 50 mm aperture
 $B_{\text{ultimate}} = 16 \text{ T}$



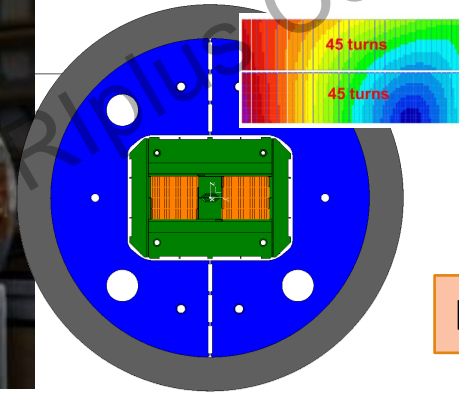
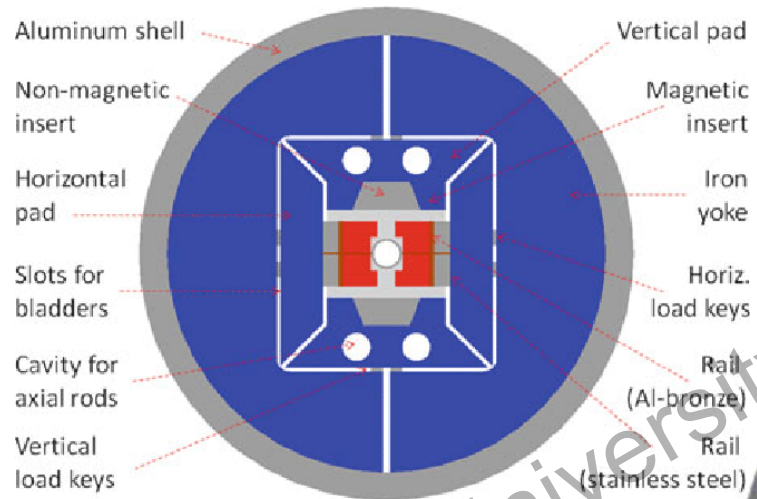
		$\cos(\theta)$	blocks	common coil	CCT
Current	(A)	10000	11230	16100	18055
Inductance	(mH/m)	50	40	19.2	19.2
Stored energy	(kJ/m)	2500	2520	2490	3200
Coil mass	(tons)	7400	7400	9200	9770



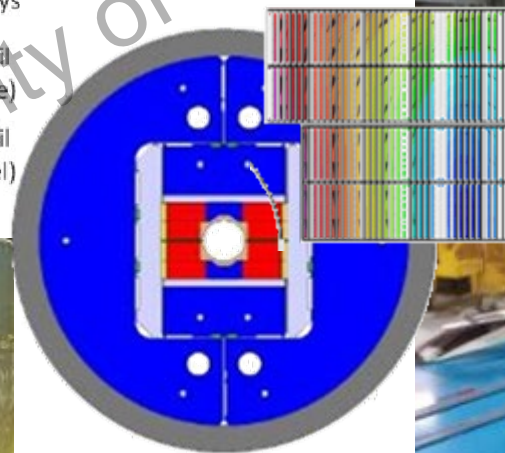
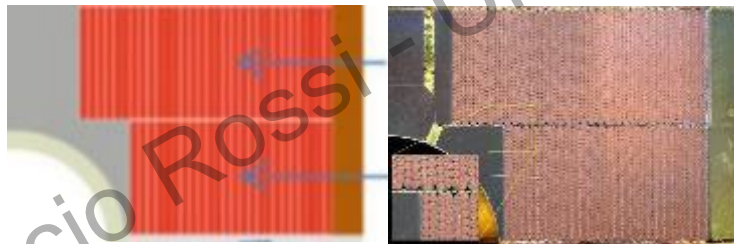
European Union's Horizon 2020 agreement No 101008548

Exploring new coil lay-out (again with bladder-key system)

RMC/eRMC (2-decks, no aperture), 16.5 T



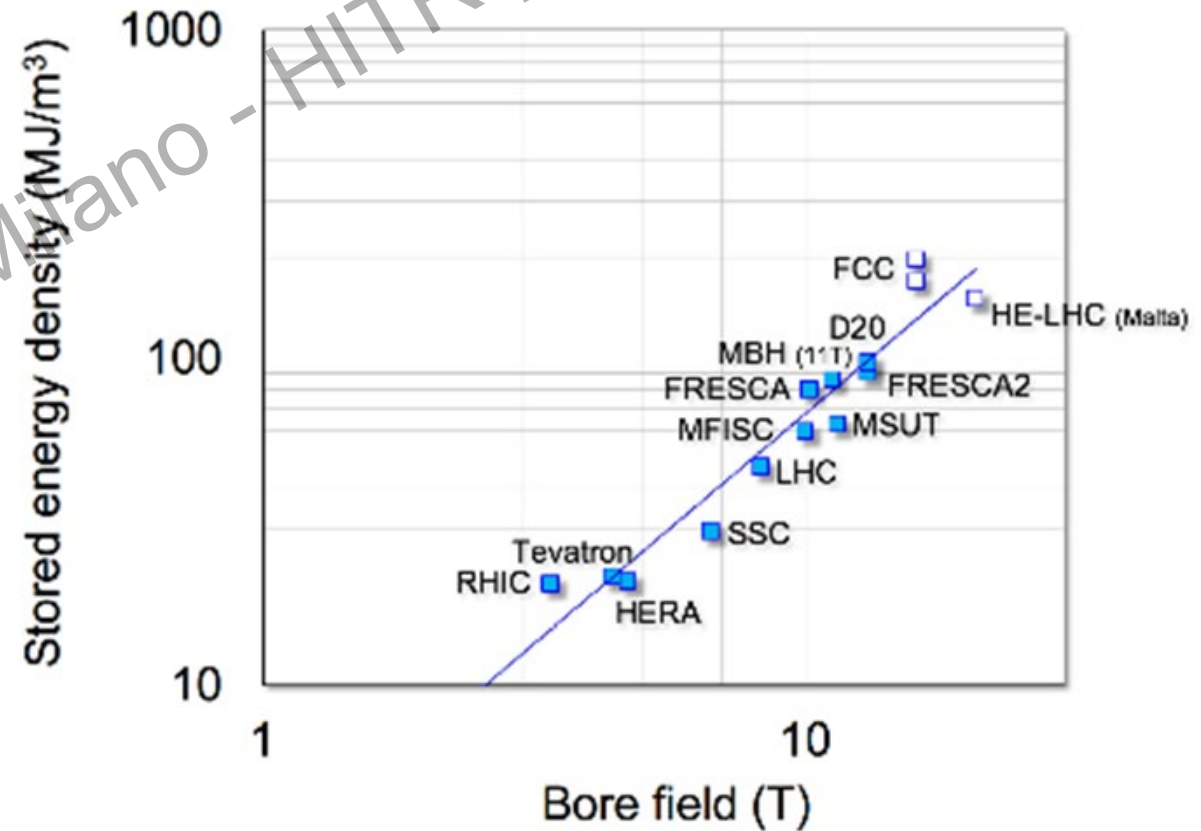
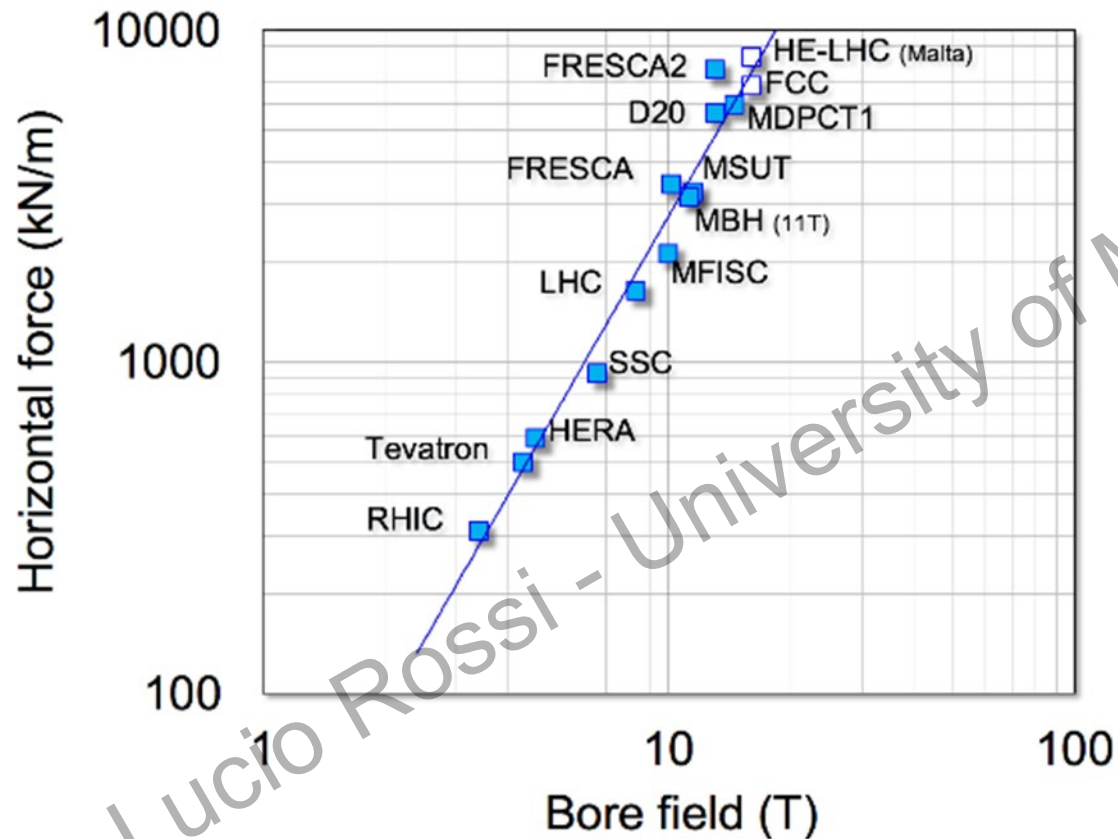
L. Bottura, CERN



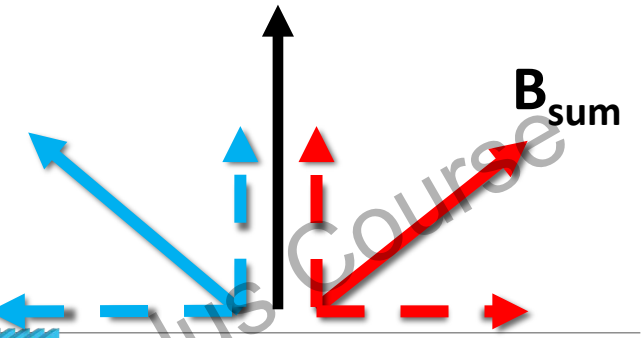
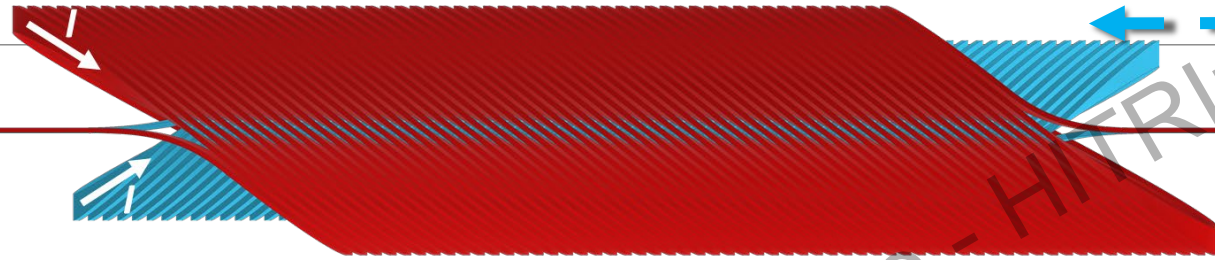
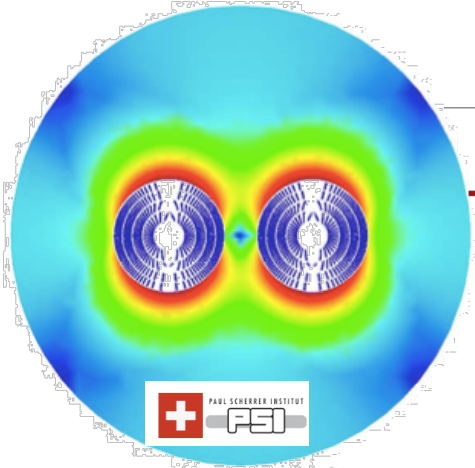
New structure to support forces/stresses

Effort in making magnets with huge energy safe

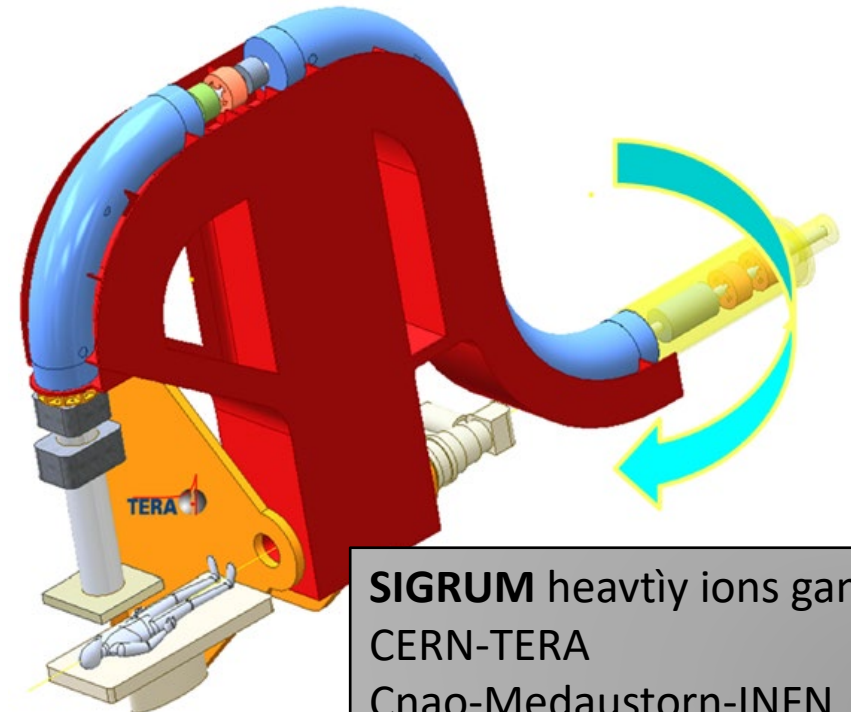
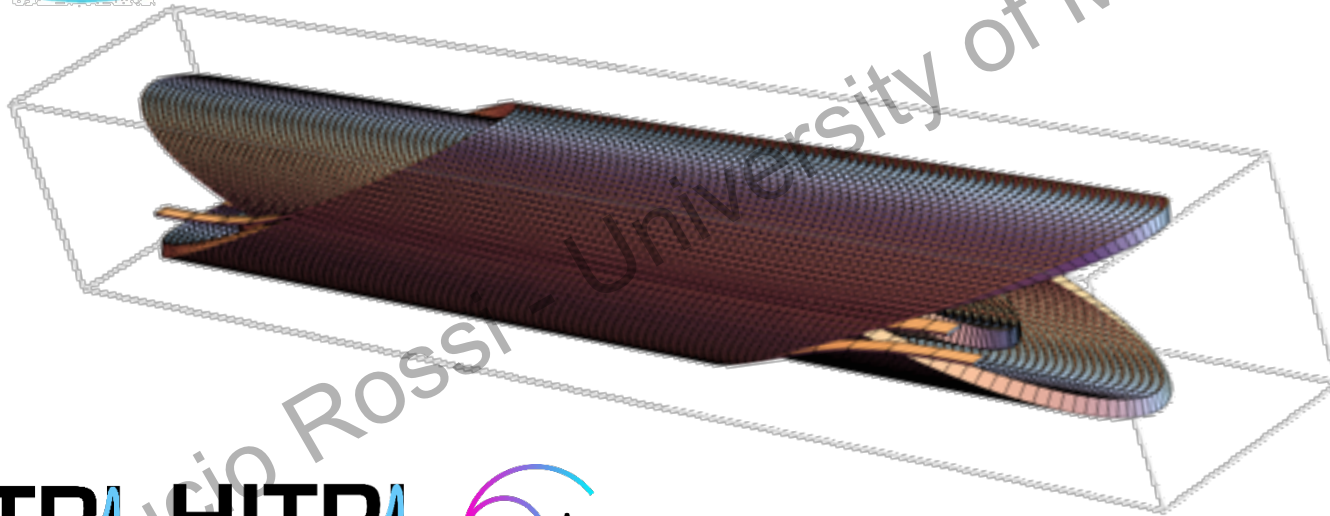
L. Bottura, CERN



CCT Canted CosTheta design : from FCC → Hadron Therapy



G. Ceruti & E. DeMatteis,
INFN-MI-LASA



SIGRUM heavy ions gantry
CERN-TERA
Cnao-Medaustorn-INFN

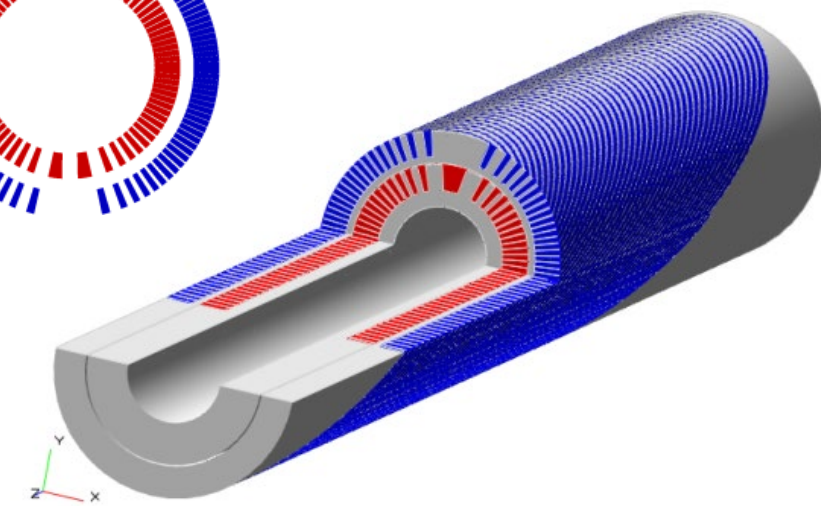
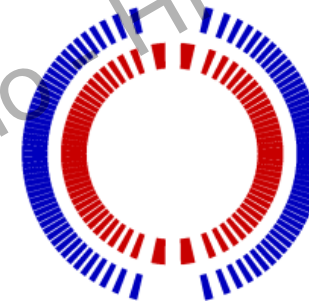
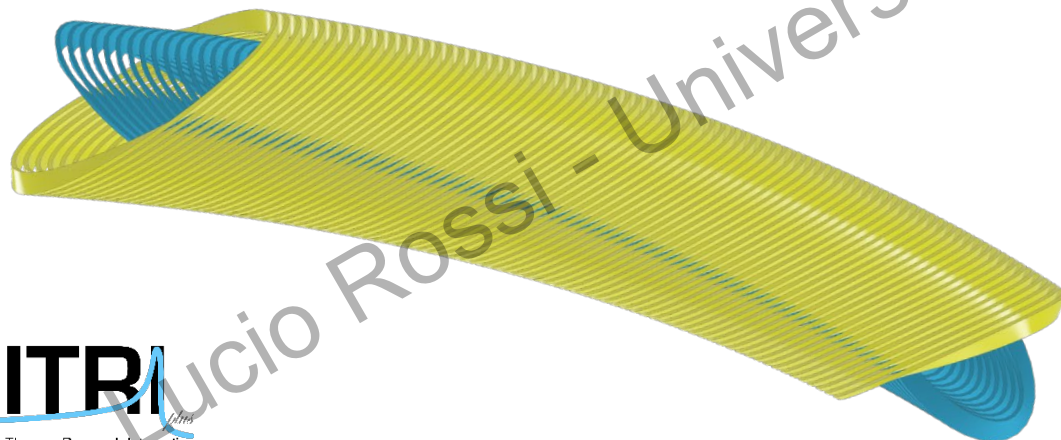
HITRI HITRI
Heavy Ion Therapy Research Integration Heavy Ion Therapy Research Integration

IFAST

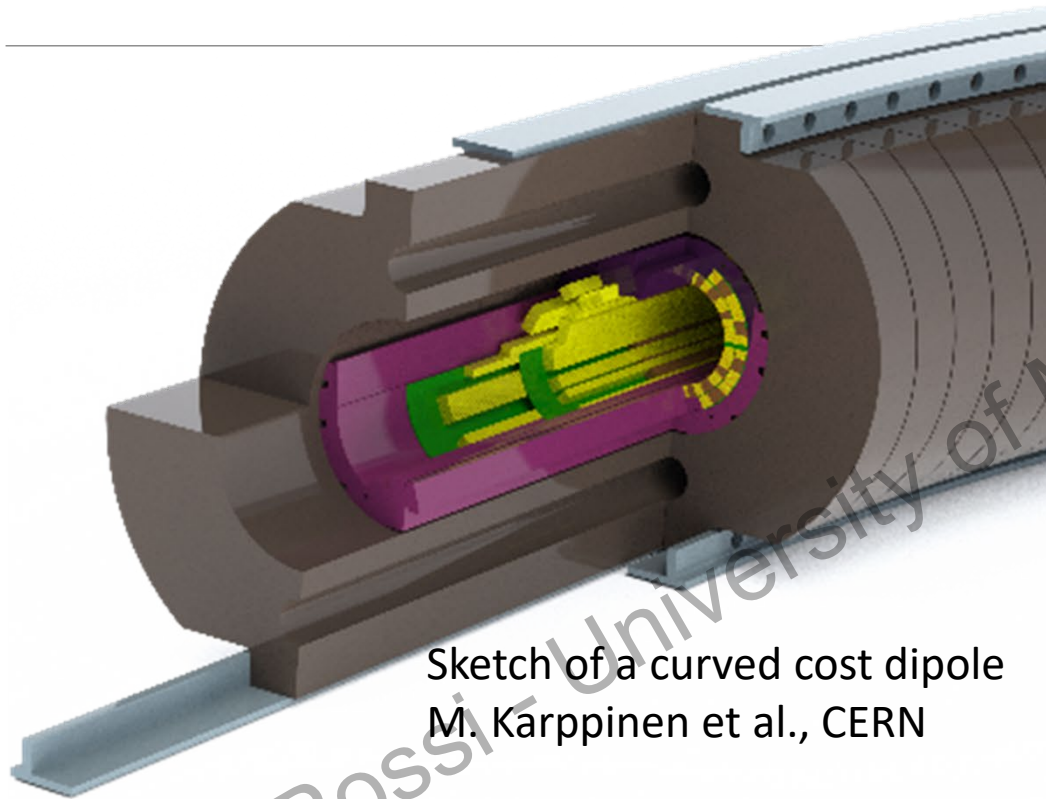
The CCT is perfect Costheta (in the bore)

Combined dip-quad

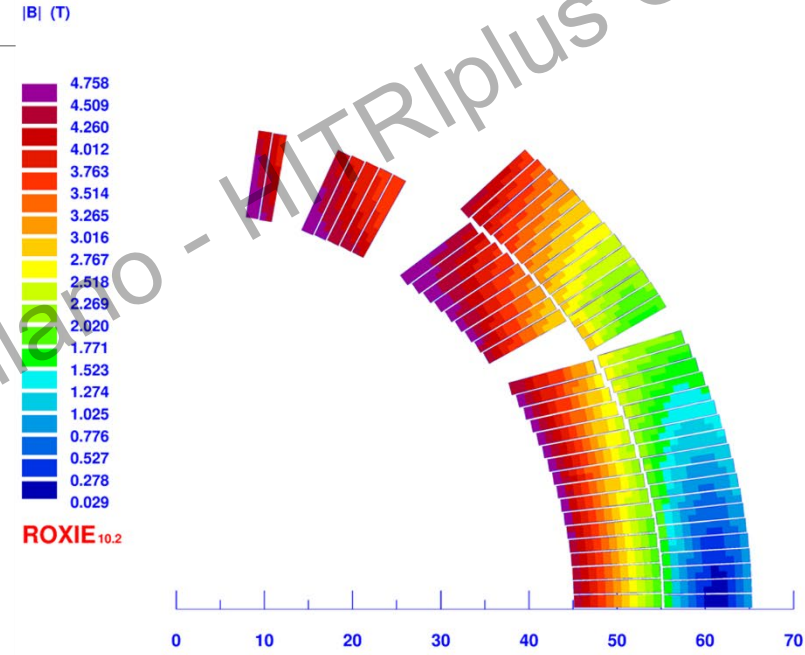
Curved 4 T – 0.4T/s



But also the “classical” CosTheta is being considered for curved dipole for gantry (INFN-SIG project inside the SIGRUM project)



Sketch of a curved cost dipole
M. Karppinen et al., CERN

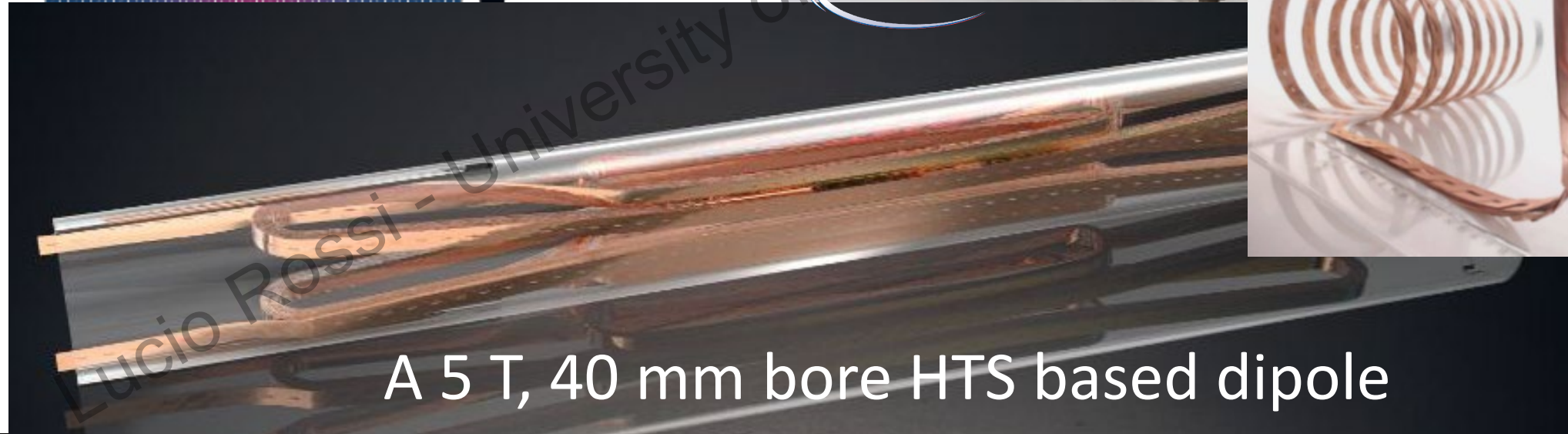
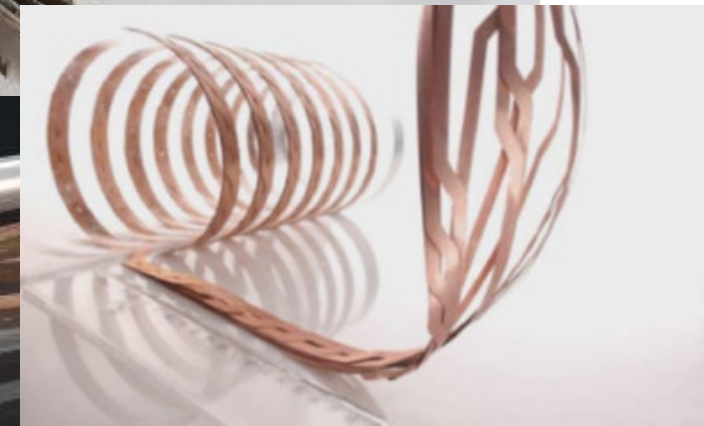
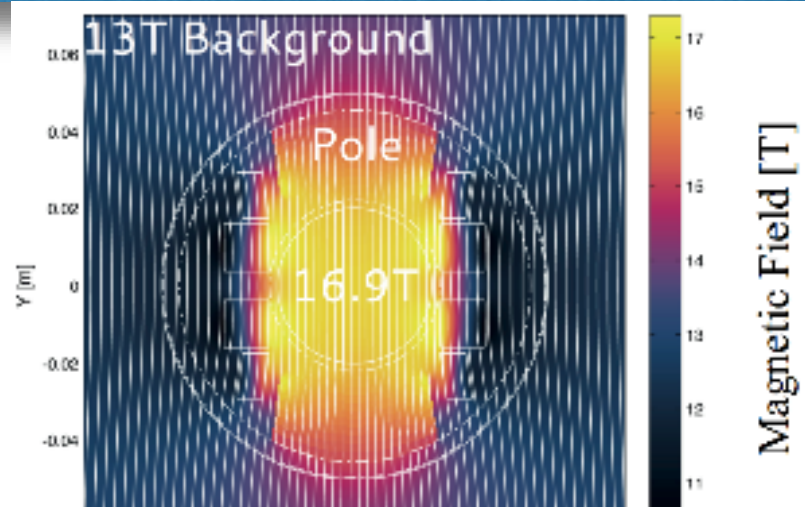


SIG project
INFN-LASA



High Temperature Superconductors – HTS

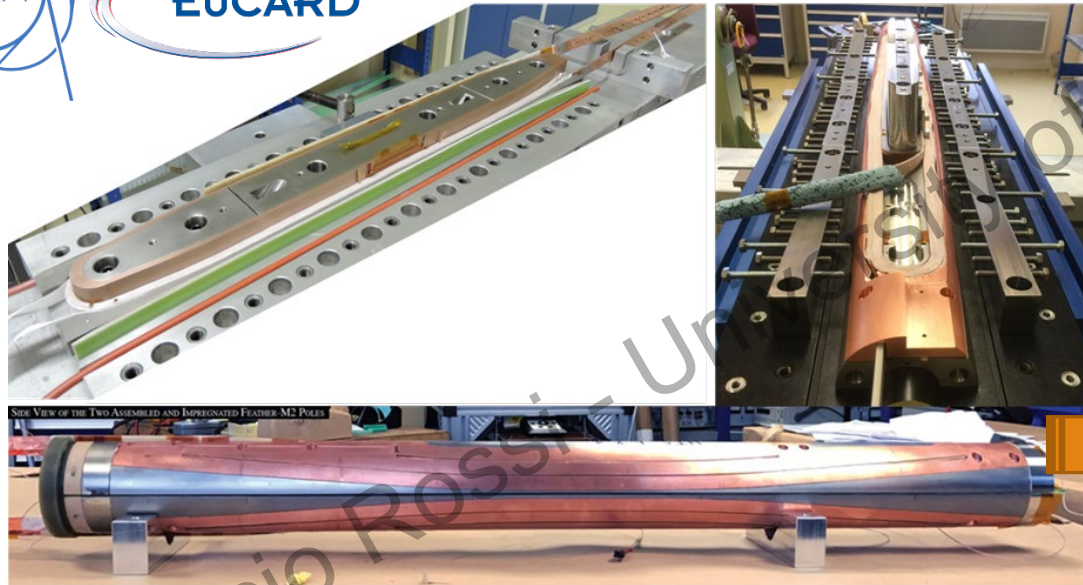
The dream of 20-25 tesla! (2 x HilumiLHC!)



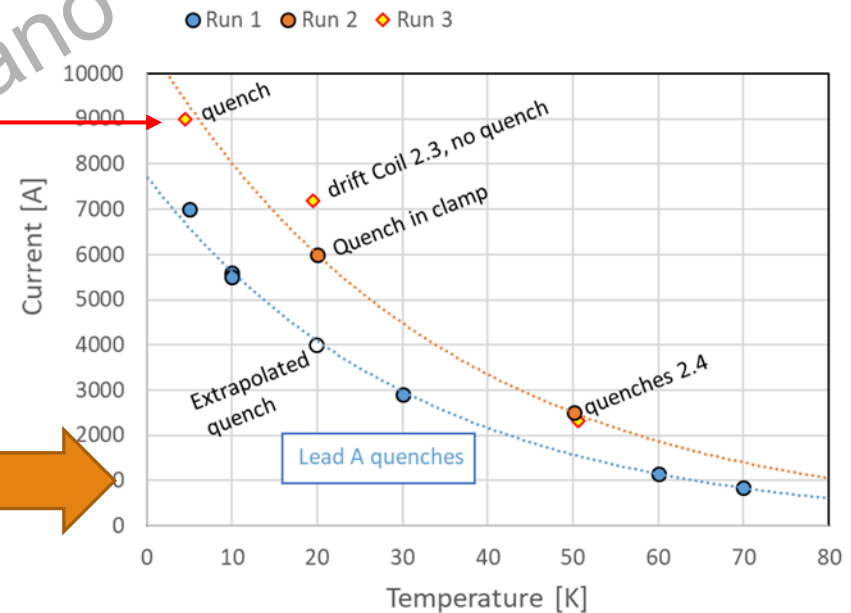
A 5 T, 40 mm bore HTS based dipole demonstrator

HTS for accelerator magnets: Eucard2 results

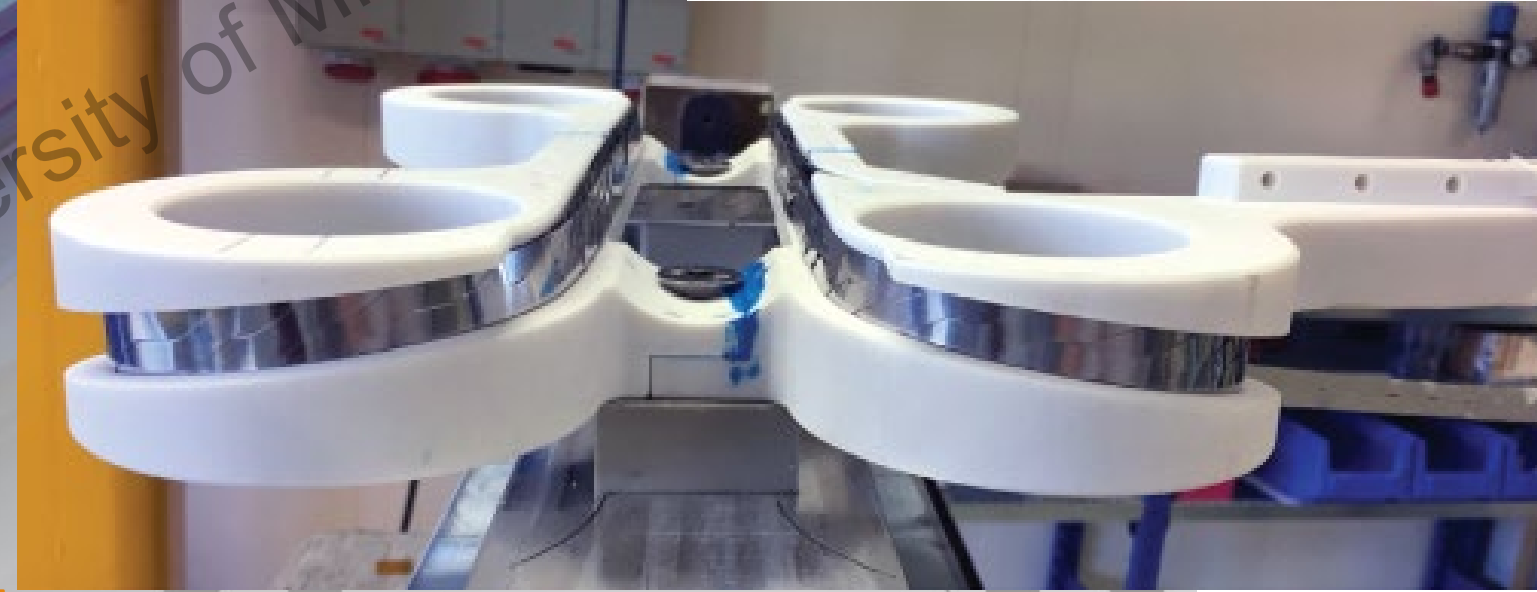
CERN Feather_M2: flare end race track,
REBCO Roebel cable, **40 mm bore**



> 3 T in 2017; 1st magnet
> 4.5 T in 2019; 2nd magnet



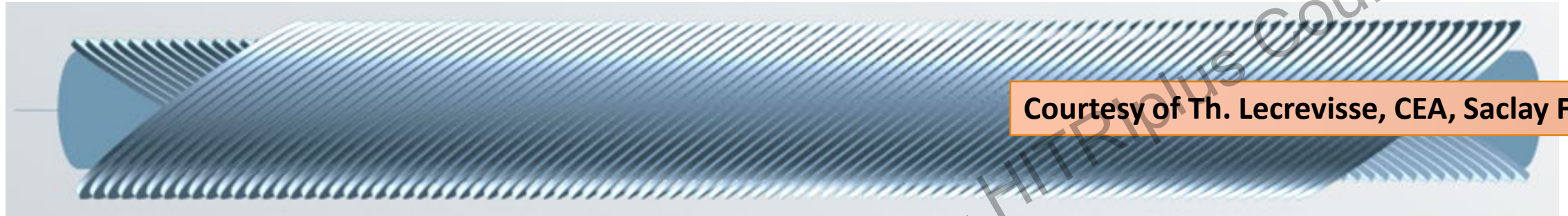
Trying the magnets of the future... 20 tesla or more...



J. Van Nugeteren – Little Beast Engineering,
Glyn Kirby CERN

& INFN-MI-LASA - SC MAGNETS FOR HITRIPLUS SPECIALISED COURSE ON HEAVY ION THERAPY RESEARCH

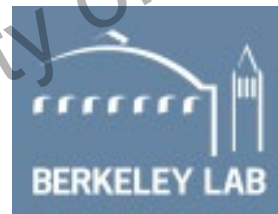
HTS dipoles with HTS superconductor



Courtesy of Th. Lecomte, CEA, Saclay FR

2-layer insulated CCT dipole of 88 Turns considering Frenet-Serret path without hardway bending, suitable for using HTS tape, under study in H202-I.FAST by CEA and INFN-MI-LASA.

CORC[®] By Advanced Conductor Technologies

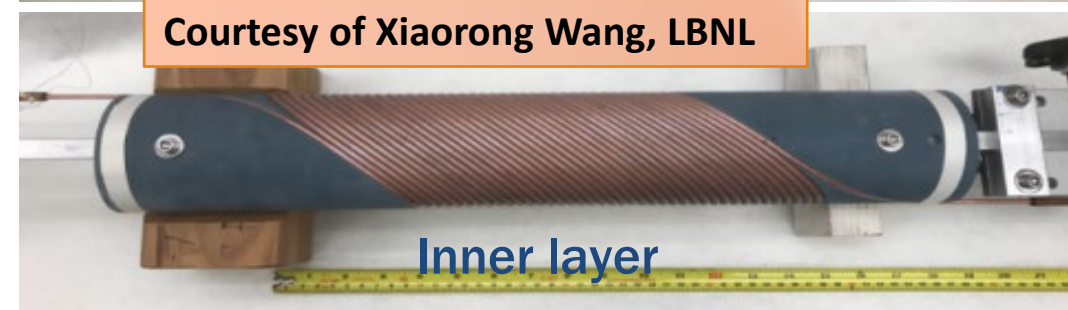


CCT program at LBNL for various applications: HEP, Hadron therapy...



Outer layer

Courtesy of Xiaorong Wang, LBNL

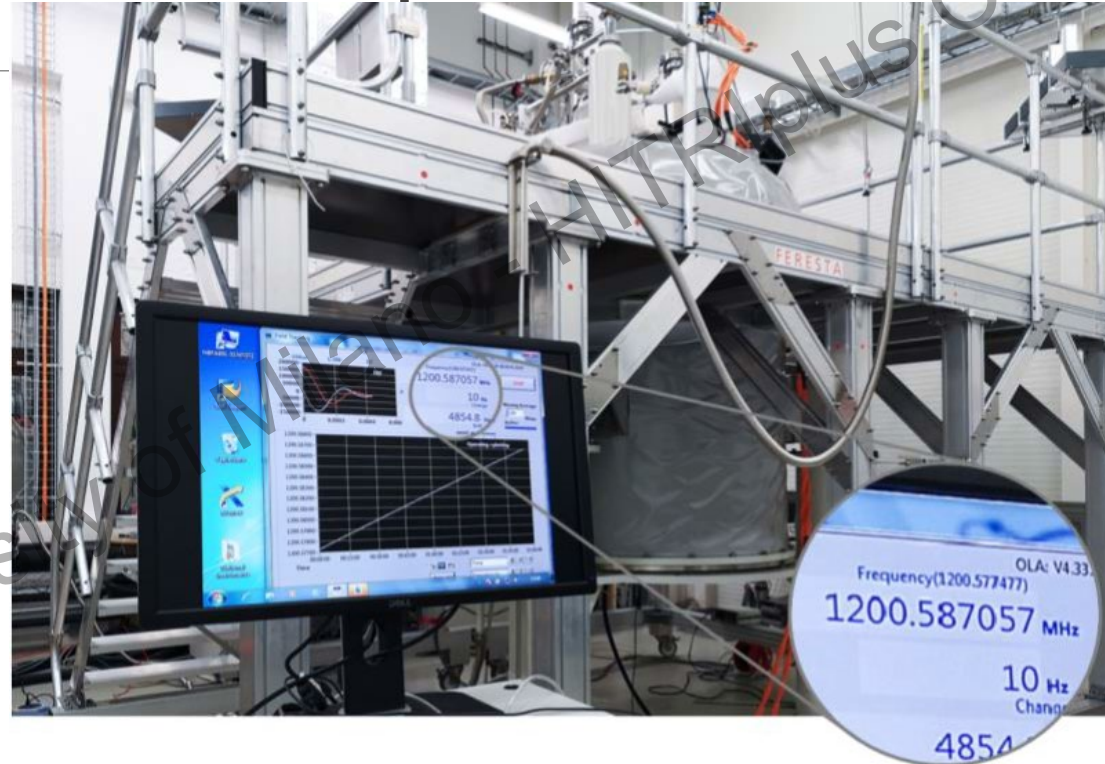
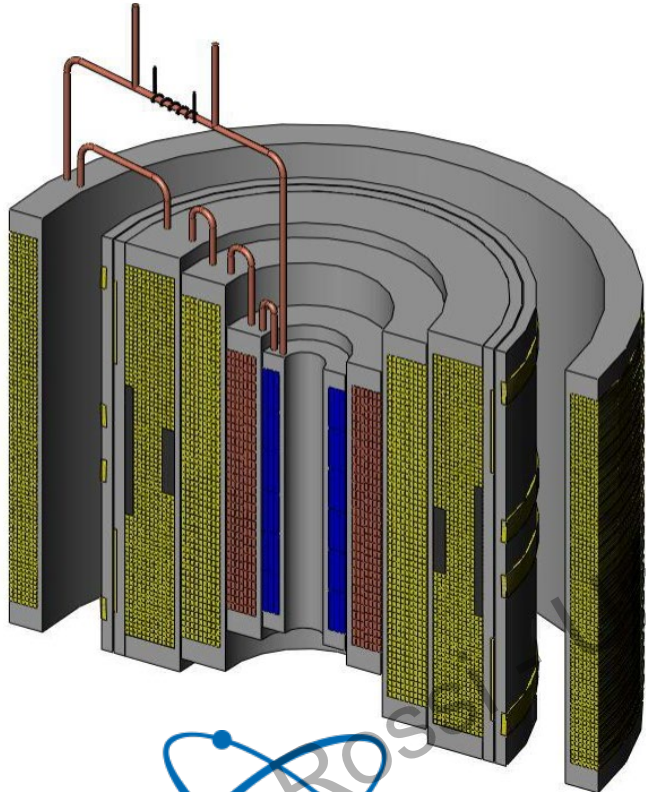


Inner layer



HITRI
Heavy Ion Therapy Research Integration

A big leap forward by a private company... Bruker Biospin



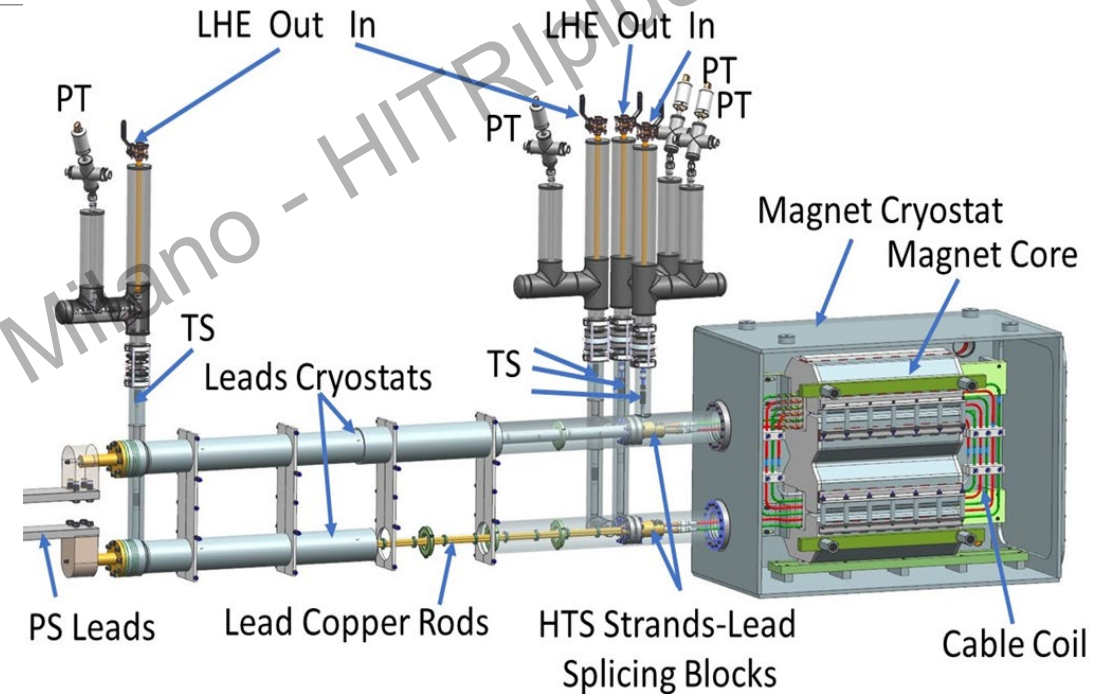
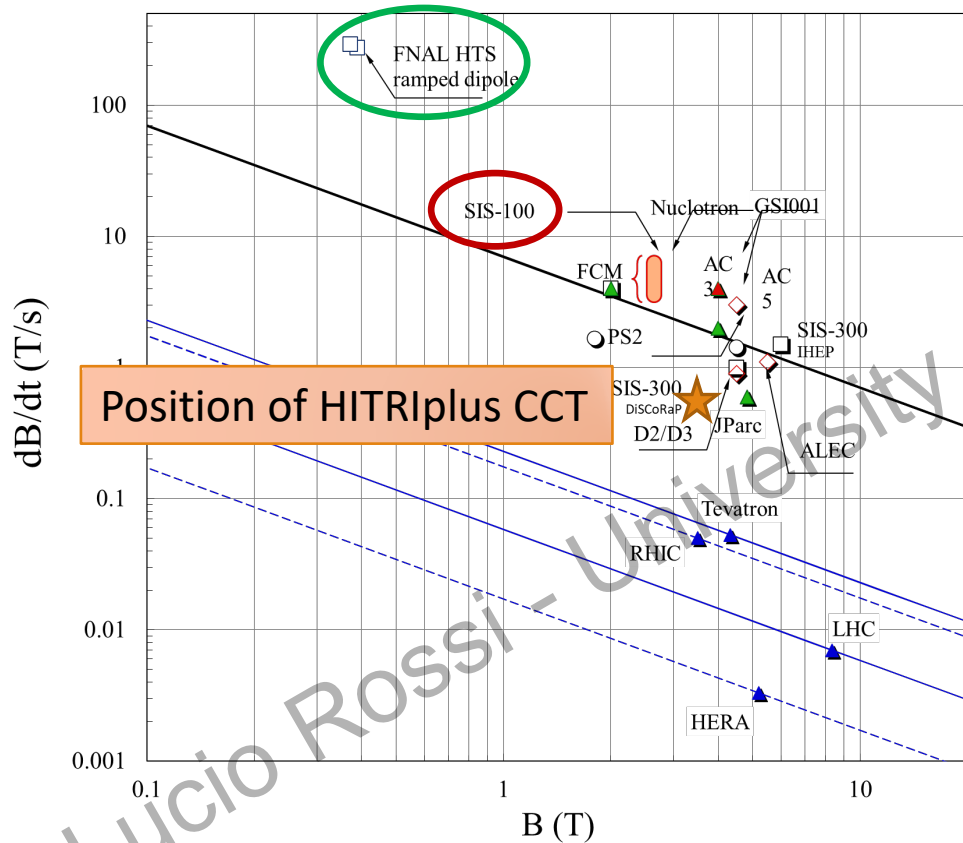
The First 1.2 GHz (28.2 T) NMR Magnet Reached Full Field in 2019, with HTS insert coil



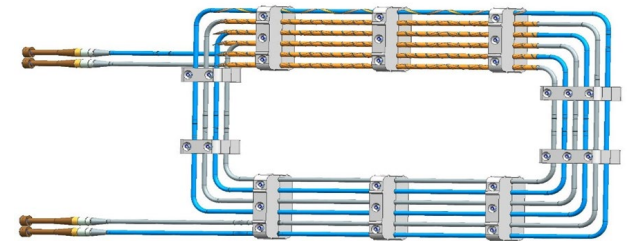
This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 101008548

SC Magnets for ramped operation

RCS for hadrons, Muon-C acceleration, Hadron Therapy



B = 0.3 T
dB/dt ~ 300 T/s
RepRate: 10 Hz
For RCS, for Muon-C...

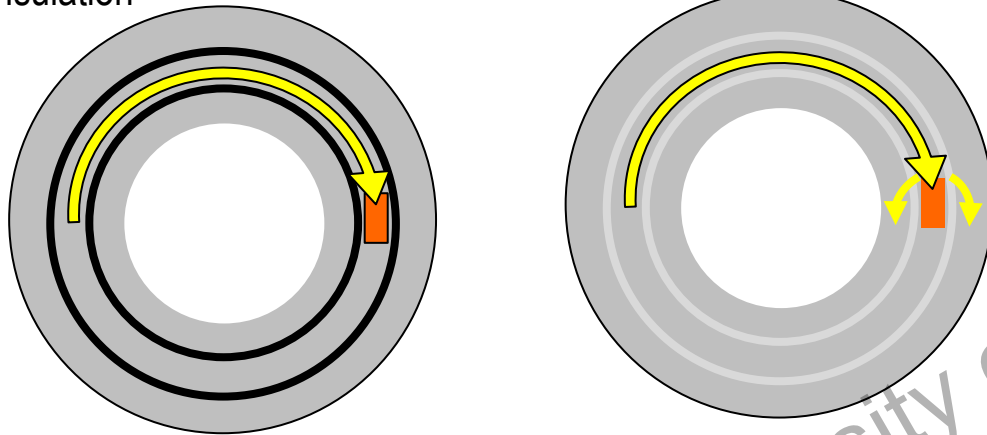


The new frontier: why to insulate a coils?

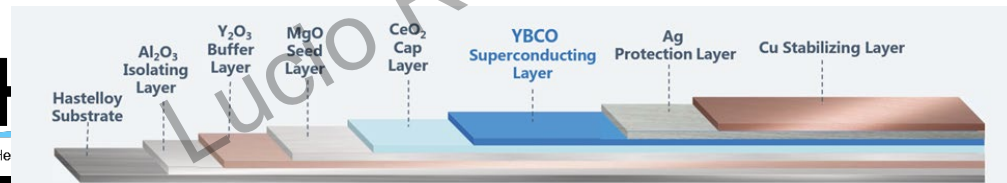
The raise of NI (non-insulated) coil...

Courtesy Prof. S. Hahn, NHMFL

Insulation



RRP strand (0.85 mm, 108/127)
 J_c : 2450 A/mm² (12 T, 4.2 K)
 Cu:non-Cu: 1.2
 40 strands cable
 (18.15 mm x 1.52 mm)



LNCMI/CEA Nougat
 HTS insert
 32.5 T in 50 mm
 (12.5 T HTS + 20 T resistive)



Sunam NI one-body
 HTS magnet
 26.4 T in 35 mm
 (26.4 T HTS multi-width)



Need to understand and control transient effects and control Field Quality during ramping up

Horizon 2020
 No 101008548

A few references (only -old- books)

for papers you can google the name reported in the paper or few other colleagues (E. Todesco, CERN, Soren Prestemon & P. Ferracin, LBNL, A. Ballarino (CERN), ecc...) or the CAS school on Superconducting magnets (2013)

- ✓ **M.N. Wilson, *Superconducting Magnets*, Clarendon Press Oxford**
- ✓ H.A. Brechna, *Superconducting Magnet Systems*, Springer Verlag
- ✓ K.-H. Mess, P. Schmüser, S. Wolff, *Superconducting Accelerator Magnets*, World Scientific
- ✓ E.W. Collings, *Applied Superconductivity*, Plenum Press
- ✓ B. Seeber (editor), *Handbook of Applied Superconductivity*, IoP Publishing
- ✓ L. Dresner, *Stability of Superconductors*, Plenum Publ. Corp.
- ✓ Y. Iwasa, *Case Studies in Superconducting Magnets*, Plenum Publ. Corp.

Thank you!

