

CP violation for electroweak baryogenesis in SMEFT

Henning Bahl, EF, Sven Heinemeyer, Judith Katzy, Marco Menen, Krisztian Peters, Matthias Saimpert and Georg Weiglein
2202.11753

In collaboration with

EF, Marta Losada, Yehonatan Viernik, Yossi Nir
1911.08495 (μ) [PRL]
2002.00099 (τ, t, b) [JHEP]
2006.06940 (EWBG) [JHEP]

Elina Fuchs
CERN & LU Hannover & PTB

HEFT 2022, Granada
June 15, 2022



Leibniz
Universität
Hannover



QuantumFrontiers

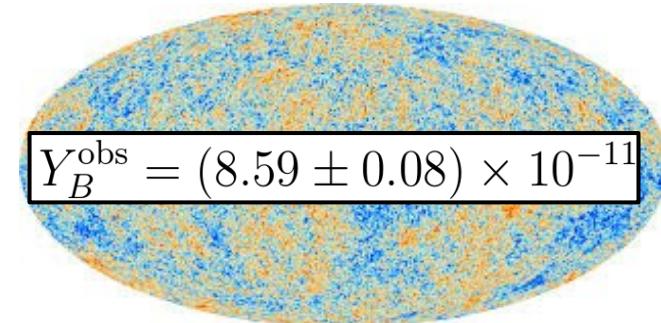
BSM CP violation for baryon asymmetry

Sakharov conditions for baryon asymmetry

I. B number violation

II. CP violation

III. Out of thermal equilibrium



[PLANCK/ESA 2013]

- Observed baryon asymmetry $Y_B^{\text{obs}} = \frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-10}$
- SM: δ_{CKM} and $\bar{\theta}_{\text{QCD}} < 10^{-10}$ by far **insufficient**

Gavela, Hernandez, Orloff, Pene '93
Huet, Sather '94

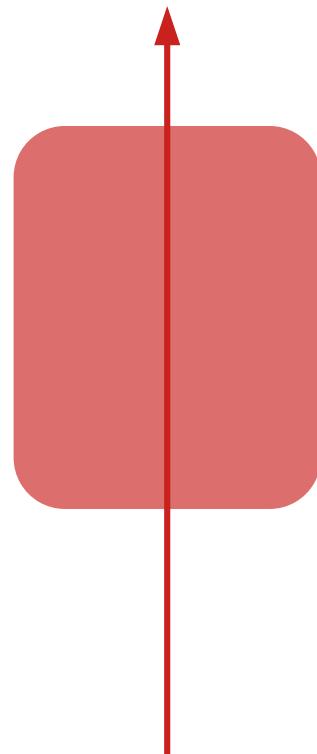
Electroweak baryogenesis:
during e.w. phase transition
→ connected to the Higgs
→ potentially testable
at colliders

Need CP violation beyond the SM

Constraints on CP violation - schematic

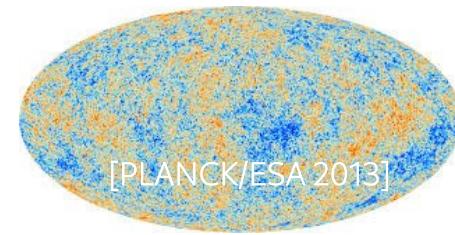


Constraints on CP violation - schematic



PLANCK: observed baryon asymmetry

$$Y_B = \frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-10}$$



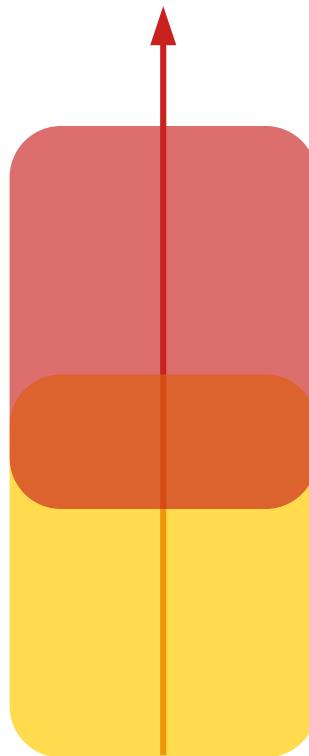
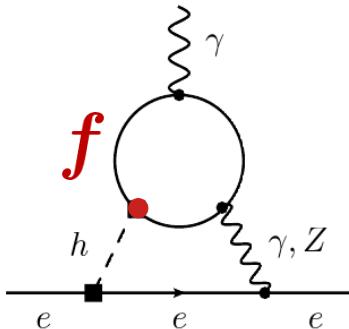
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Constraints on CP violation - schematic

ACME [Nature '18]: upper bound

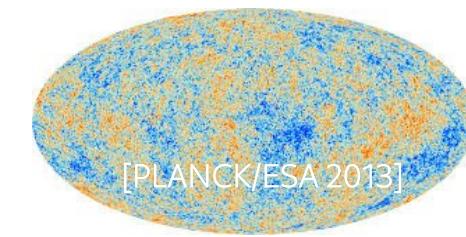
$$d_e \leq 1.1 \times 10^{-29} e \text{ cm at } 90\% \text{ CL}$$

Electric dipole moments restrict CPV



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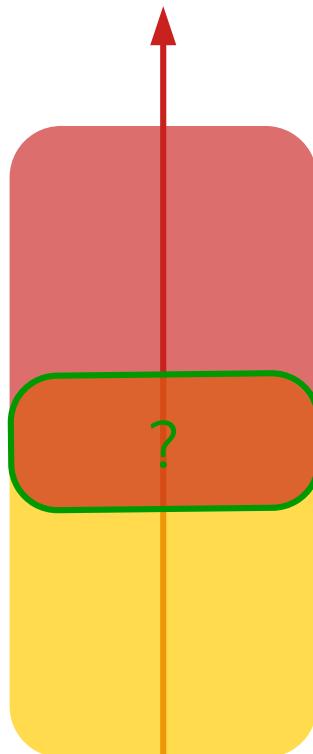
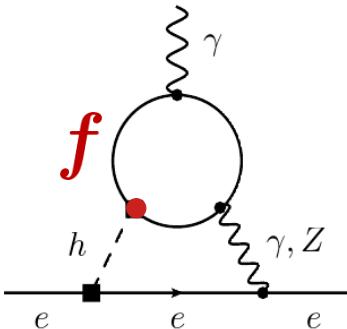


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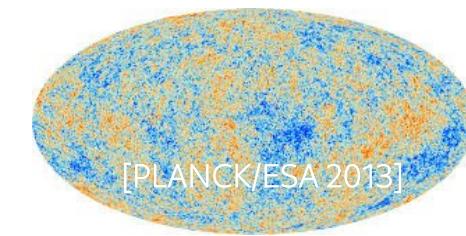
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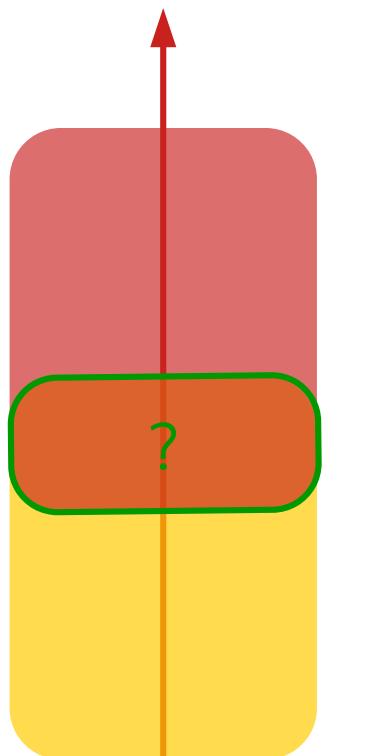
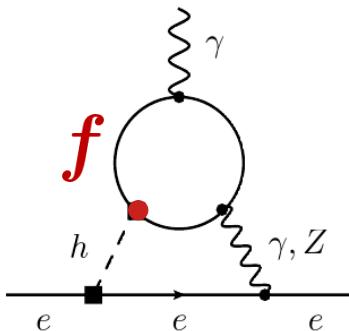


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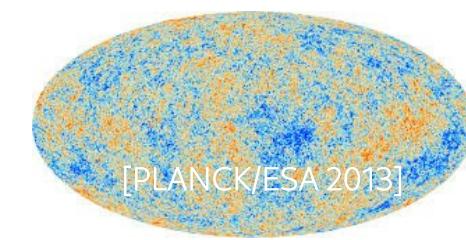
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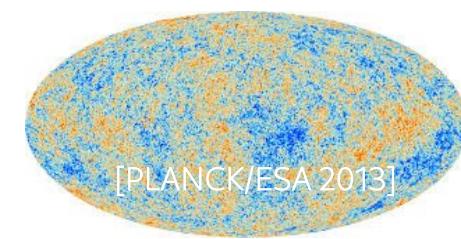
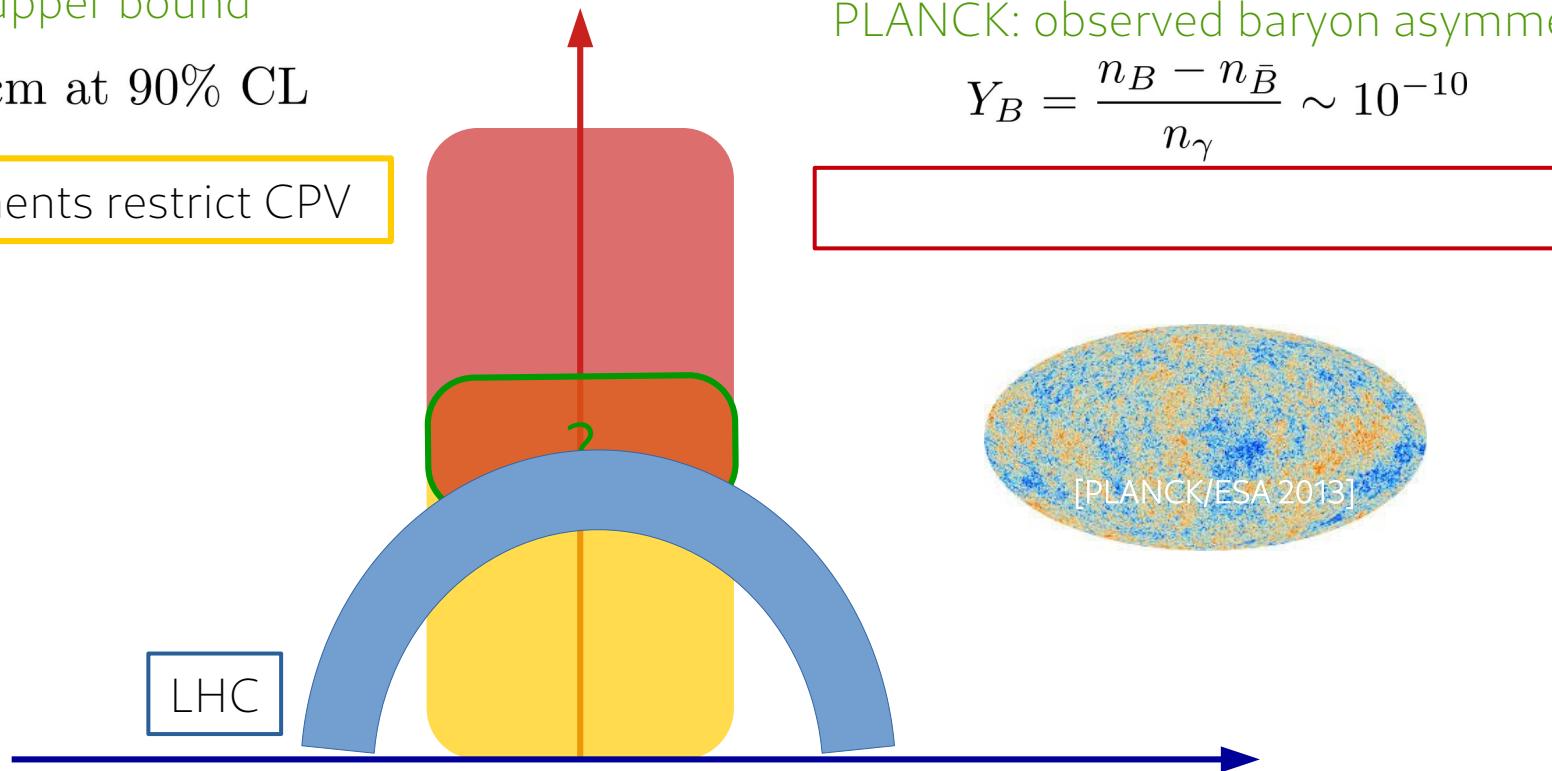
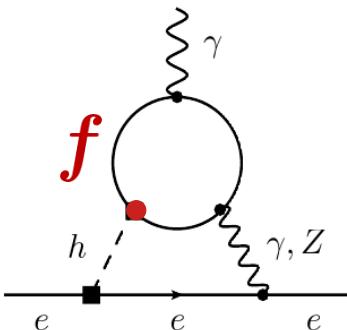


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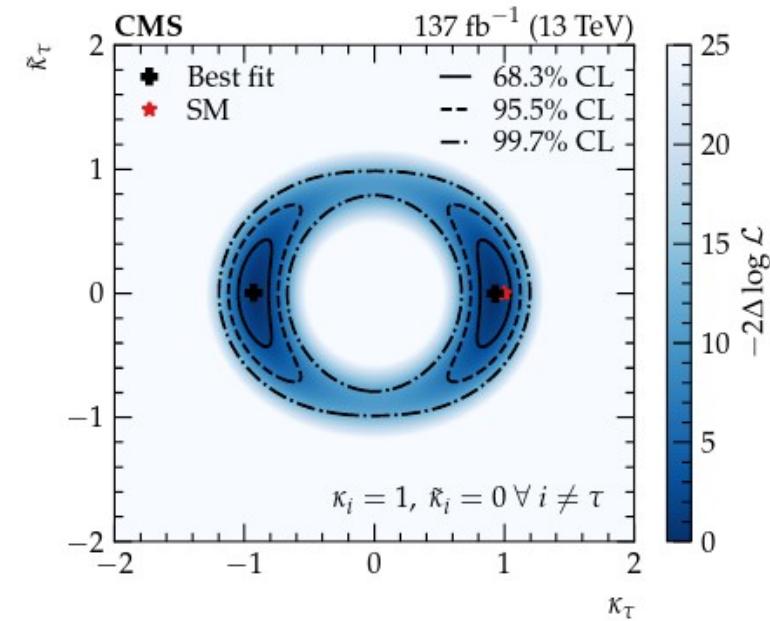
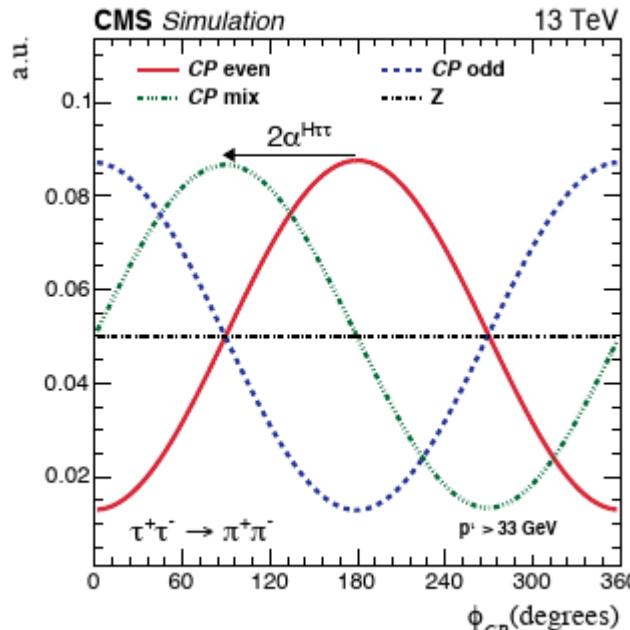
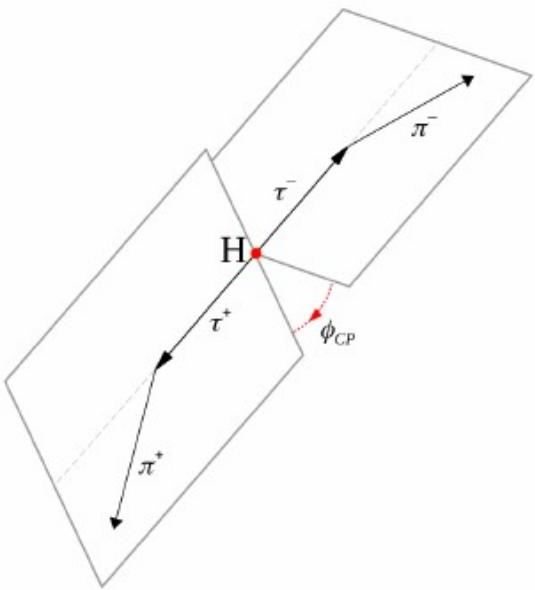
Electric dipole moments restrict CPV



Timely: CP-odd observables at LHC

CMS 2110.04836, CMS-HIG-20-006

CMS $H \rightarrow \tau\tau$ analysis



Outline

1.) Framework

2.) Baryogenesis

3.) Electric dipole moments

4.) Higgs signal strengths and angular observables at the LHC

5.) Complementarity

Complex Yukawa in SMEFT dim-6

- Consider dim-6 Yukawa with real and imaginary part

$$\mathcal{L}_{\text{Yuk}} = Y_f \overline{F_L} F_R H + \frac{1}{\Lambda^2} (X_R^f + i X_I^f) |H|^2 \overline{F_L} F_R H + \text{h.c.}$$

cf [de Vries, Postma, van de Vies '18] where $X \equiv \pm i Y_f$

- Relative size of dim-6 normalized to dim-4 $T = m_f^{(6)} / m_f^{(4)}$

Our coordinates →

$$T_{R,I}^f = \frac{v^2}{2\Lambda^2} \frac{X_{R,I}^f}{Y^f}$$

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$$\begin{aligned} \mathcal{L}_f = & \frac{y_f v}{\sqrt{2}} \left[1 + \frac{v^2}{2\Lambda^2} \frac{X_R^f + i X_I^f}{y_f} \right] \overline{f_L} f_R + \frac{y_f}{\sqrt{2}} \left[1 + \frac{3v^2}{2\Lambda^2} \frac{X_R^f + i X_I^f}{y_f} \right] \overline{f_L} f_R h \\ & + \frac{3v}{2\sqrt{2}\Lambda^2} (X_R^f + i X_I^f) \overline{f_L} f_R h h + \frac{1}{2\sqrt{2}\Lambda^2} (X_R^f + i X_I^f) \overline{f_L} f_R h h h. \end{aligned}$$

Full Lagrangian
→ focus on Yukawa
& mass terms

Impact on fermion mass & Yukawa

$$m_f = \frac{Y_f v}{\sqrt{2}} \left(1 + T_R^f + iT_I^f \right), \quad \lambda_f = \frac{Y_f}{\sqrt{2}} \left(1 + 3T_R^f + 3iT_I^f \right)$$

rotate into basis where mass is real

$$m_f \overline{f_L} f_R$$

$$\tan \theta_f = \frac{T_I^f}{1 + T_R^f}$$

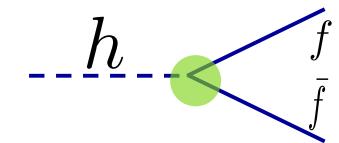
$$\frac{Y_f v}{\sqrt{2}} \left[1 + T_R^f + \mathcal{O}(T^{f2}) \right]$$

$$\frac{Y_f}{\sqrt{2}} \left[1 + 3T_R^f + 2iT_I^f + \mathcal{O}(T^{f2}) \right].$$

Higgs characterization model

Consider also simpler description of effective Higgs coupling modifiers (kappa framework)

$$\mathcal{L}_{\text{Yuk}} = - \sum_f \frac{y_f}{\sqrt{2}} \bar{f} (c_f + i\gamma_5 \tilde{c}_f) f h,$$



Translate kappa SMEFT:

$$g_f = c_f + i\tilde{c}_f = 3 - \frac{2}{1 + T_f^R + iT_f^I} \quad \text{with} \quad T_f^{R,I} \equiv \frac{v^2}{2\Lambda^2} \frac{X_f^{R,I}}{y_f}$$

Allow also modifications of real parts of HVV couplings

$$\mathcal{L}_V = c_V H \left(\frac{M_Z^2}{v} Z_\mu Z^\mu + 2 \frac{M_W^2}{v} W_\mu^+ W^{-\mu} \right)$$

Capture BSM effects in effective Hgg and Hy γ couplings: $c_g, \tilde{c}_g, c_\gamma, \tilde{c}_\gamma$

Limits on CPV in Higgs couplings

SMEFT of dim. 6 $\mathcal{L}_{\text{Yuk}} = -\sum_f y_f \overline{F_L} F_R H + \frac{1}{\Lambda^2} (X_R^f + i X_I^f) |H|^2 \overline{F_L} F_R H + \text{h.c.}$

used in EF, Losada, Nir, Viernik '19, '20, '20
see also de Vries, Postma, v. de Vis '19;
Brod, Cornell, Skodras, Stamou '22

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Inclusive LHC
Higgs rates

Angular
Higgs distributions

Electric dipole
moments

Remaining allowed regions for 1 or several complex Yukawas?

Goals:

Limits on CPV in Higgs couplings

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several f:
EDM cancellation?

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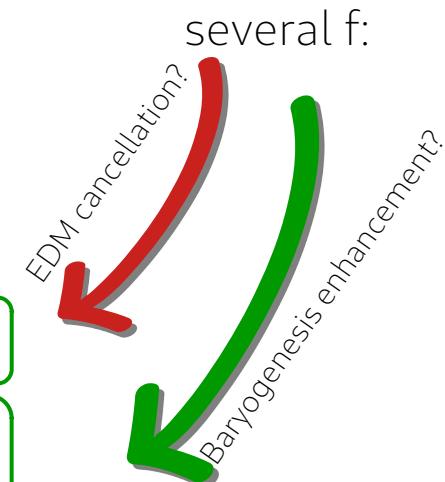
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Calculate baryon asymmetry
within experimental limits → Maximal possible asymmetry?

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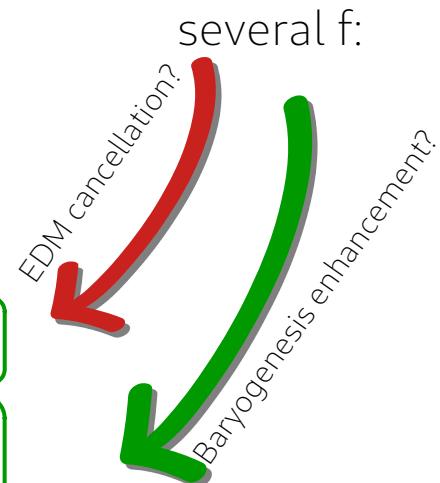
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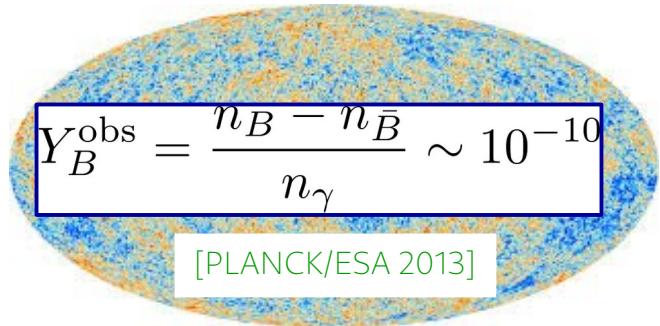
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BSM for baryogenesis: focus here on CPV, assume electroweak phase transition can be enhanced separately → later: models

Electroweak baryogenesis

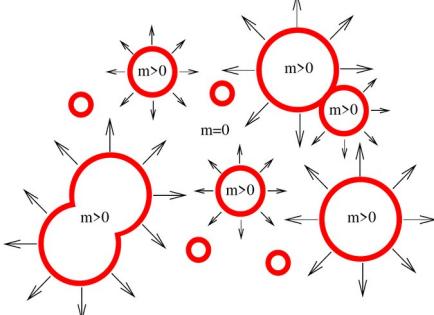


Lots of literature, e.g.

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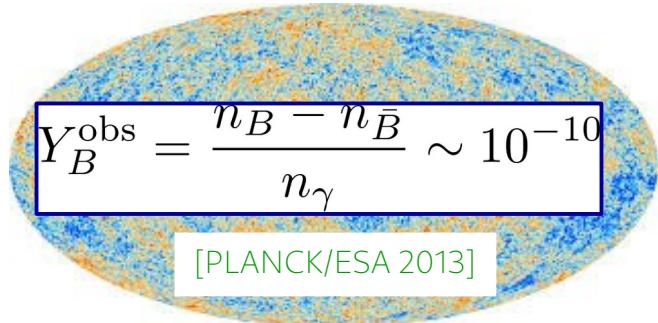
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- CP violation
- 1st order electroweak phase transition



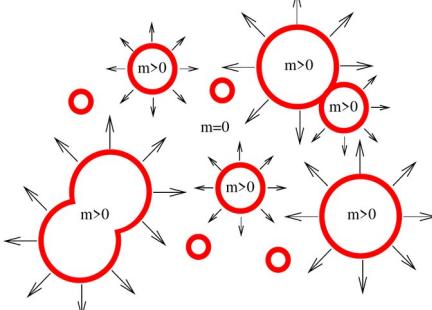
Bubbles of the
broken phase
expand

Electroweak baryogenesis



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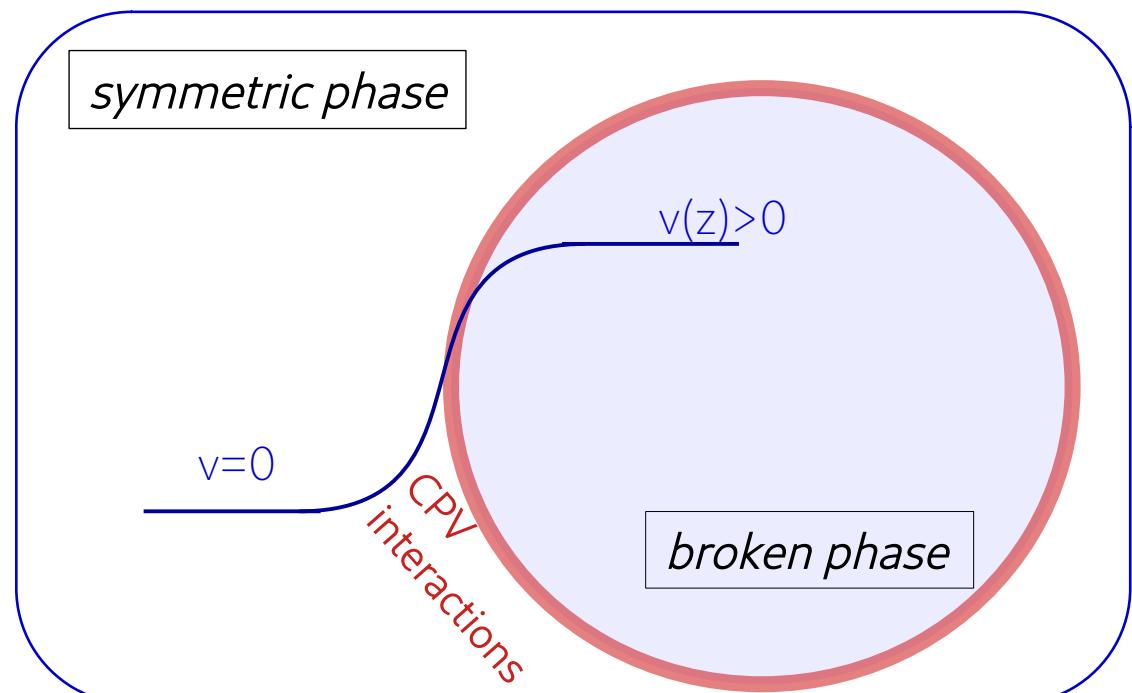
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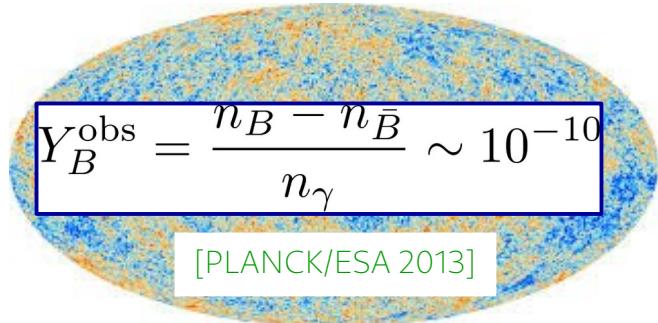
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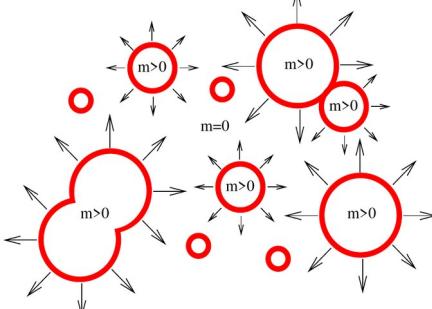


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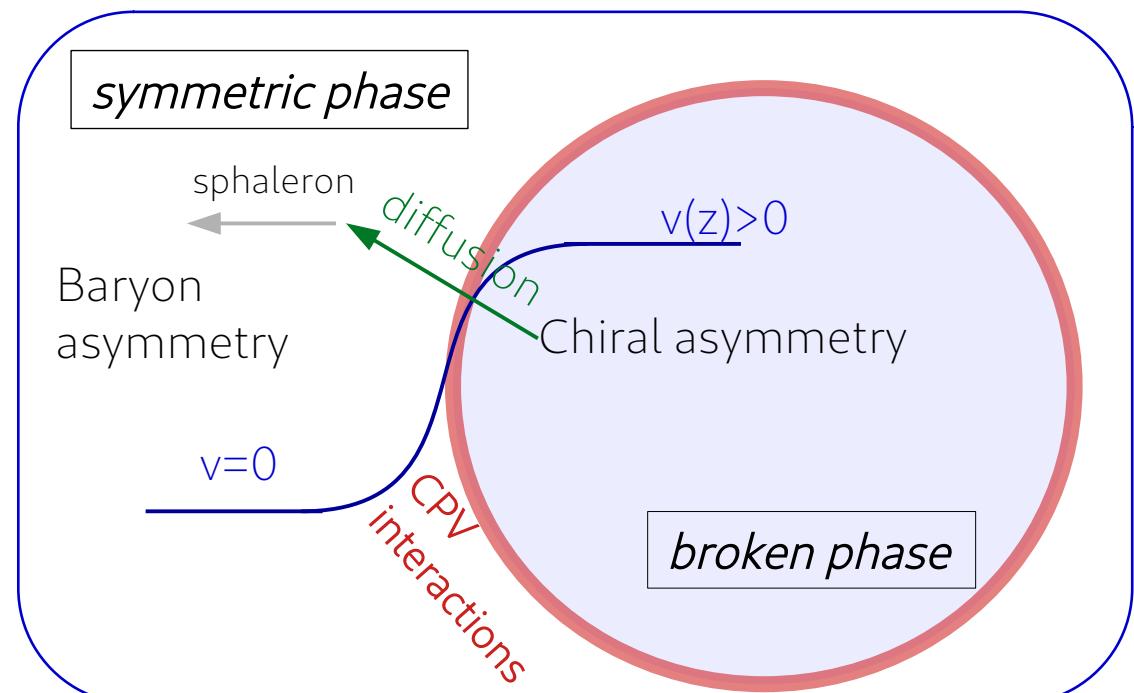
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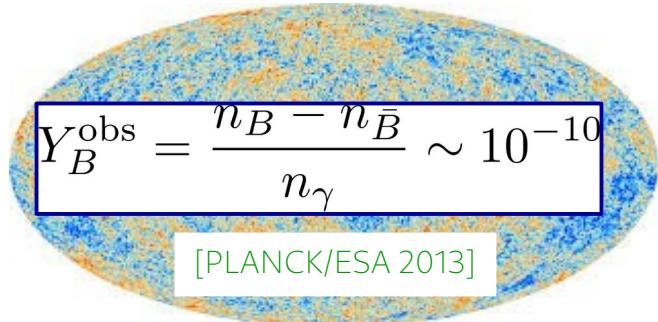
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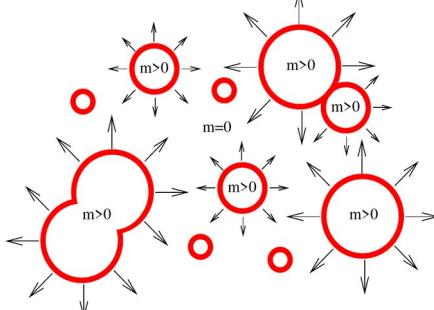


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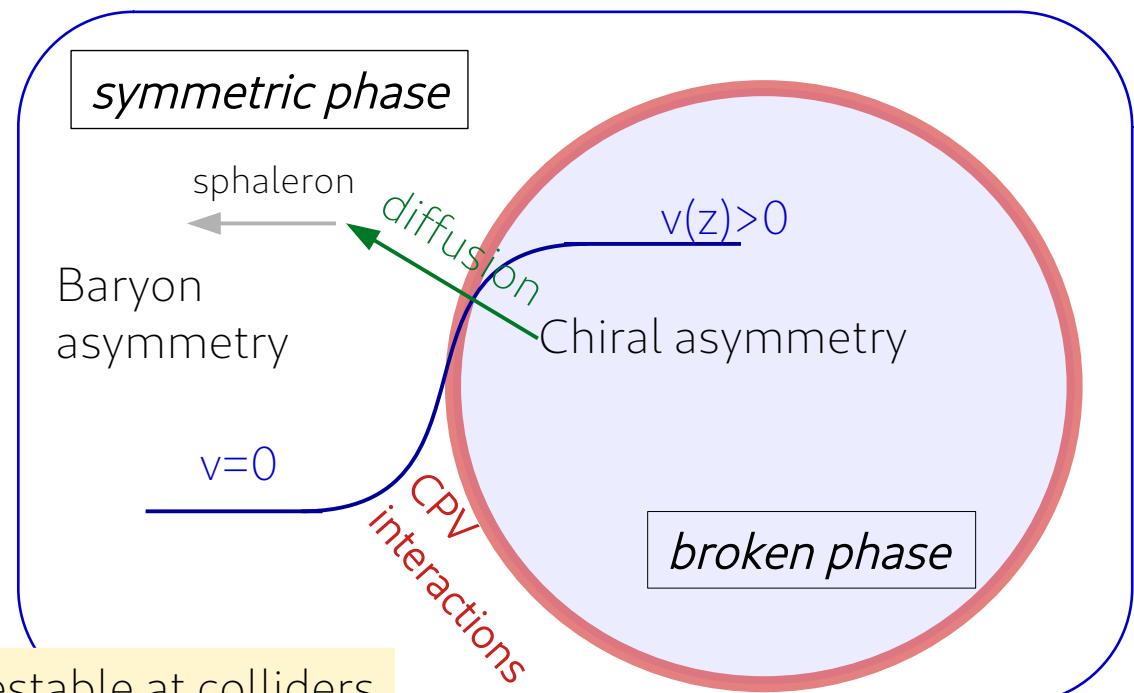


Bubbles of the broken phase expand

Higgs physics → testable at colliders

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Transport equations for baryogenesis

Transport equations for each fermion and Higgs, set of coupled differential equations

$$\partial_\mu f^\mu = -\Gamma_M^f \mu_M^f - \Gamma_Y^f \mu_Y^f + \Gamma_{ss}^f \mu_{ss} - \Gamma_{ws}^f \mu_{ws}^f + S_f$$

relaxation Yukawa Strong weak
 sphaleron CPV source

Transport equations for baryogenesis

Transport equations for each fermion and Higgs, set of coupled differential equations

$$\partial_\mu f^\mu = \text{CP-conserving interactions} + S_f$$

CPV source

Transport equations for baryogenesis

Transport equations for each fermion and Higgs, set of coupled differential equations

$$\partial_\mu f^\mu =$$

CP-conserving interactions

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CPV
source

Approximations

- vev-insertion (VIA)
caveat: VIA – WKB
discrepancy
- thin wall
- diffusion

$$Y_B \propto S_f \propto \text{Im} [m_f^* m'_f] \tilde{c}_f$$

Same scaling
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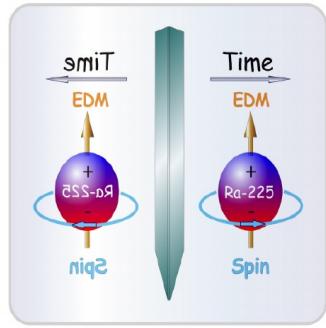
Benchmark
of wall
properties:

$$\frac{Y_B}{Y_B^{\text{obs}}} = 28\tilde{c}_t - 0.2\tilde{c}_b - 0.03\tilde{c}_c - 2 \cdot 10^{-4}\tilde{c}_s - 9 \cdot 10^{-8}\tilde{c}_u - 4 \cdot 10^{-7}\tilde{c}_d \\ - 11\tilde{c}_\tau - 0.1\tilde{c}_\mu - 3 \cdot 10^{-6}\tilde{c}_e$$

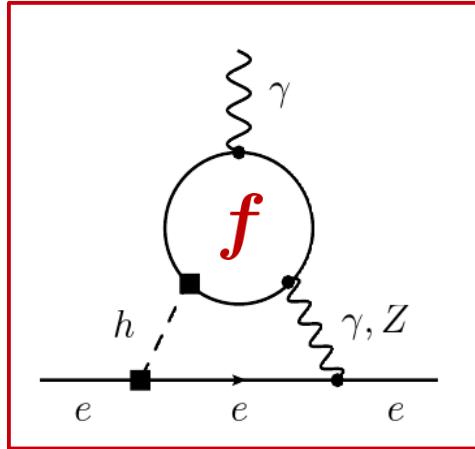
EF, Losada, Nir, Viernik '19, '20, '20
Aharony-Shapira '21

Electron's Electric Dipole Moment

[Hewett, Weerts et al '12]



EDM violates \mathcal{T} and \mathcal{P}
 $\Rightarrow \mathcal{CP}$



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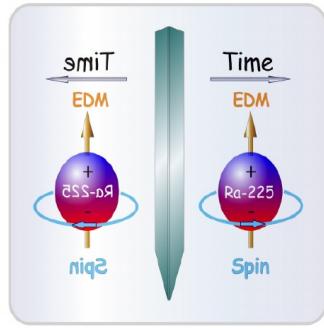
for t, b, c, τ , μ : electron EDM most sensitive

Using [Panico, Pomarol, Riembau '18], [Brod, Haisch, Zupan '13],
[Brod, Stamou '18],...

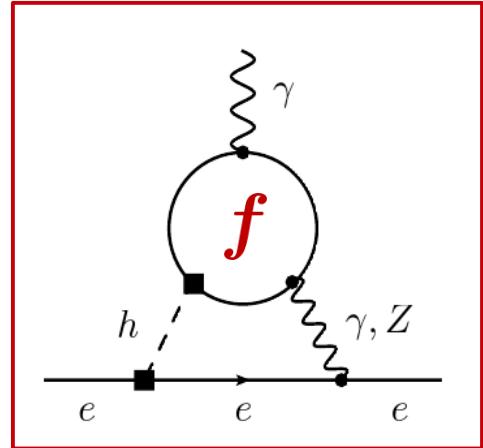
See also recent [Brod, Cornell, Skodras, Stamou '22]

Electron's Electric Dipole Moment

[Hewett, Weerts et al '12]



EDM violates \mathcal{T} and \mathcal{P}
 $\Rightarrow \mathcal{CP}$



ACME [Nature '18]:

$$d_e \leq 1.1 \times 10^{-29} e\text{ cm at 90\% CL}$$

for t, b, c, τ , μ : electron EDM most sensitive

Using [Panico, Pomarol, Riembau '18], [Brod, Haisch, Zupan '13], [Brod, Stamou '18],...

See also recent [Brod, Cornell, Skodras, Stamou '22]

$$\begin{aligned} \frac{d_e}{d_e^{\text{ACME}}} = & c_e (870.0 \tilde{c}_t + 3.9 \tilde{c}_b + 2.8 \tilde{c}_e + 0.01 \tilde{c}_s + 8 \cdot 10^{-5} \tilde{c}_u + 7 \cdot 10^{-5} \tilde{c}_d + 3.4 \tilde{c}_\tau + 0.03 \tilde{c}_\mu) \\ & + \tilde{c}_e (610.1 c_t + 3.1 c_b + 2.3 c_c + 0.01 c_s + 7 \cdot 10^{-5} c_u + 6 \cdot 10^{-5} c_d + 2.8 c_\tau + 0.02 c_\mu \\ & - 1082.6 c_V) \\ & + 2 \cdot 10^{-6} c_e \tilde{c}_e. \end{aligned}$$

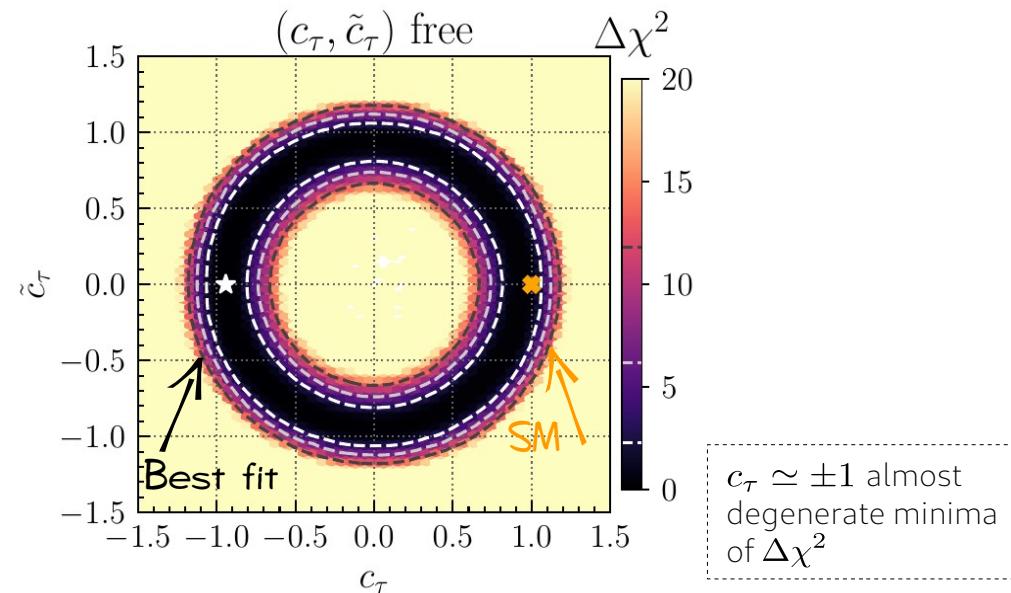
Cancellations possible

CP structure of Higgs couplings - τ

$$\mathcal{L}_{\text{Yuk}} = - \sum_f \frac{y_f}{\sqrt{2}} \bar{f} (c_f + i\gamma_5 \tilde{c}_f) f h,$$

Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22

Global fit using **HiggsSignals** + recent analyses



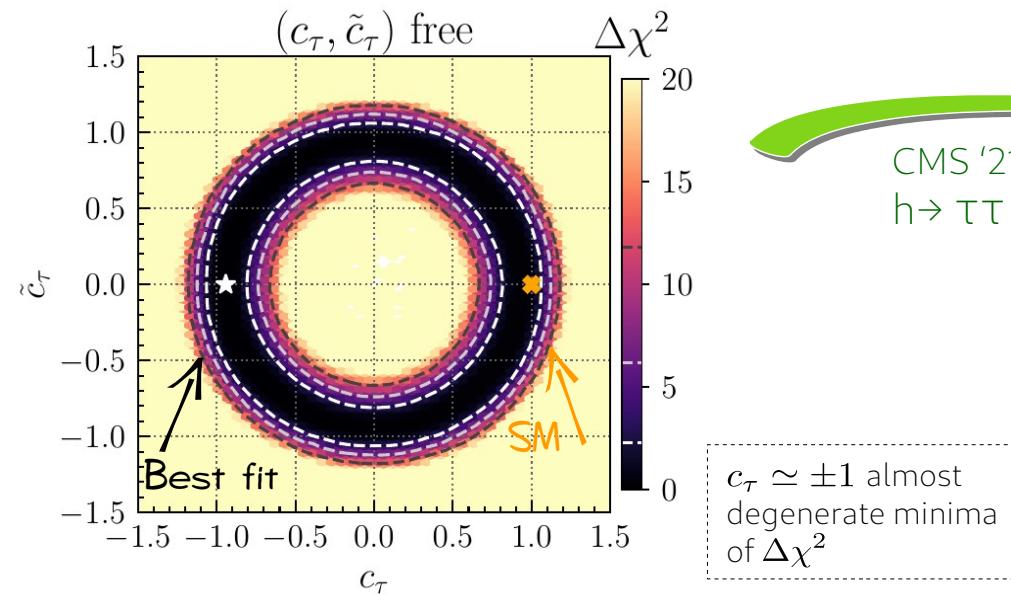
Ring-structure from upper/lower bound on BR

CP structure of Higgs couplings - τ

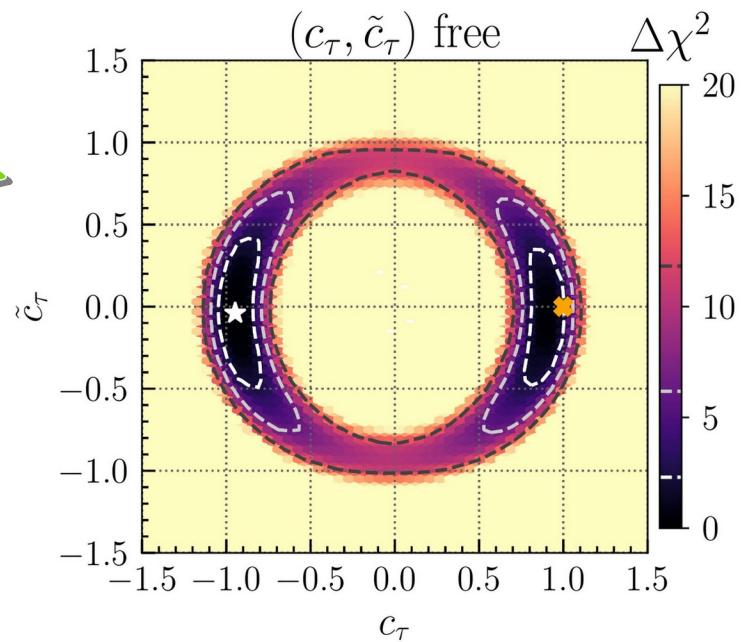
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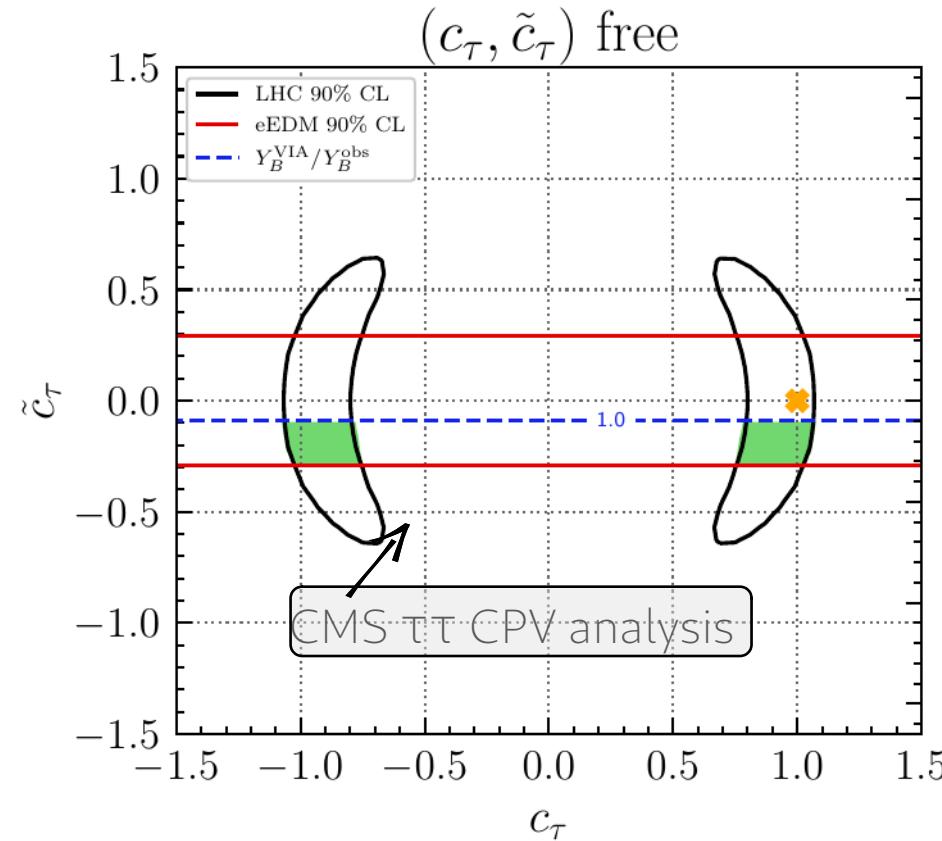


Ring-structure from upper/lower bound on BR



Complementary (τ): LHC, EDM, EWBG

Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22

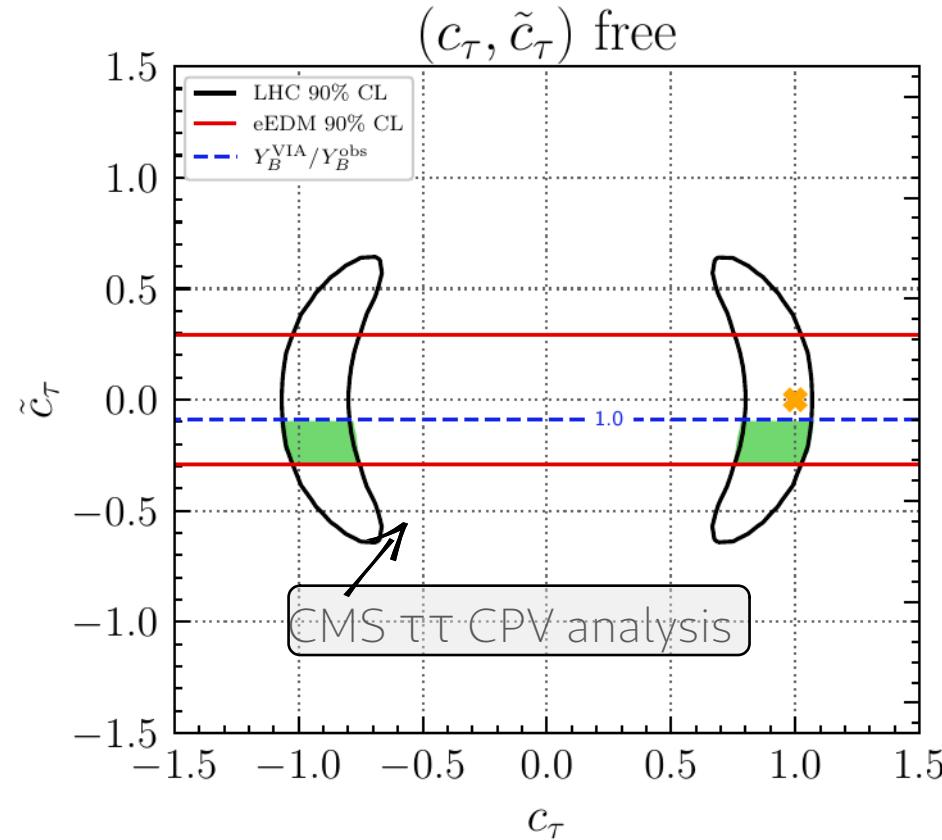


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Electron electric dipole moment

$$d_e \propto \tilde{c}_f$$

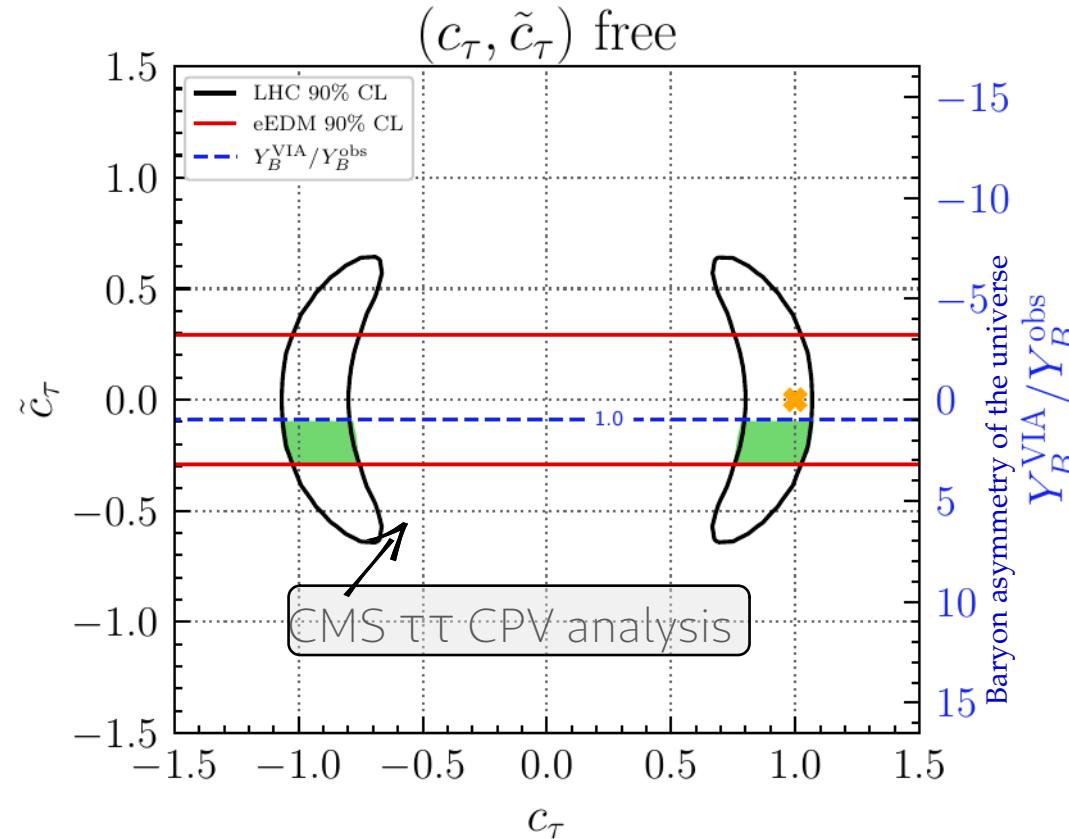


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See also

Brod, Haisch, Zupan '13

De Vries, Postma, van de Vis '18

EF, Losada, Nir, Viernik '19, '20, '20

Aharony-Shapira '21

Brod, Cornell, Skodras, Stamou '22

Electroweak baryogenesis
 $Y_B \propto \tilde{c}_f$

Caveat: "optimistic" scenario,
large uncertainty
(vev-insertion approximation)
→ almost **upper bound**

Basler, Mühlleitner, Müller '20

Cline, Kainulainen '20

Cline, Laurent '21, Postma '21

Kainulainen '21

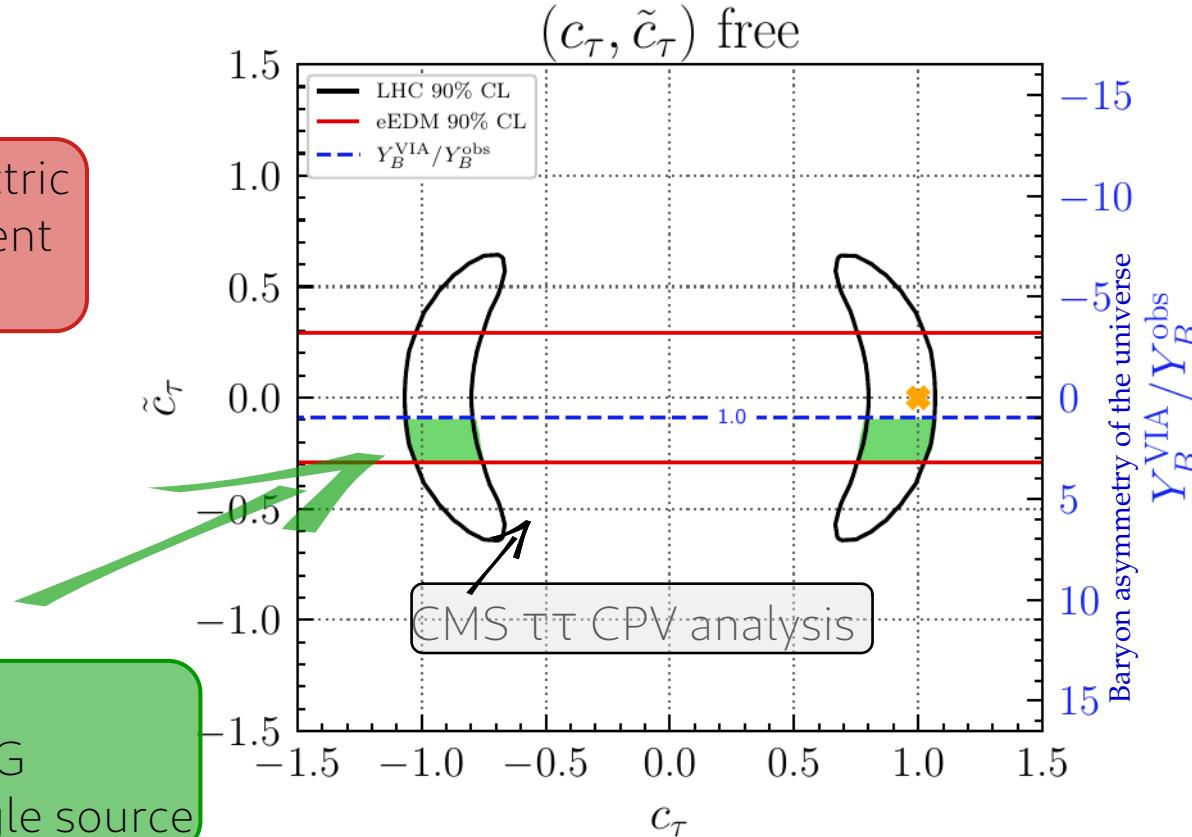
Postma, van de Vis, White '22

Complementary (τ): LHC, EDM, EWBG

Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22

Electron electric dipole moment
 $d_e \propto \tilde{c}_f$

Allowed by
LHC, EDM, EWBG
→ τ may be single source



See also

Brod, Haisch, Zupan '13

De Vries, Postma, van de Vis '18

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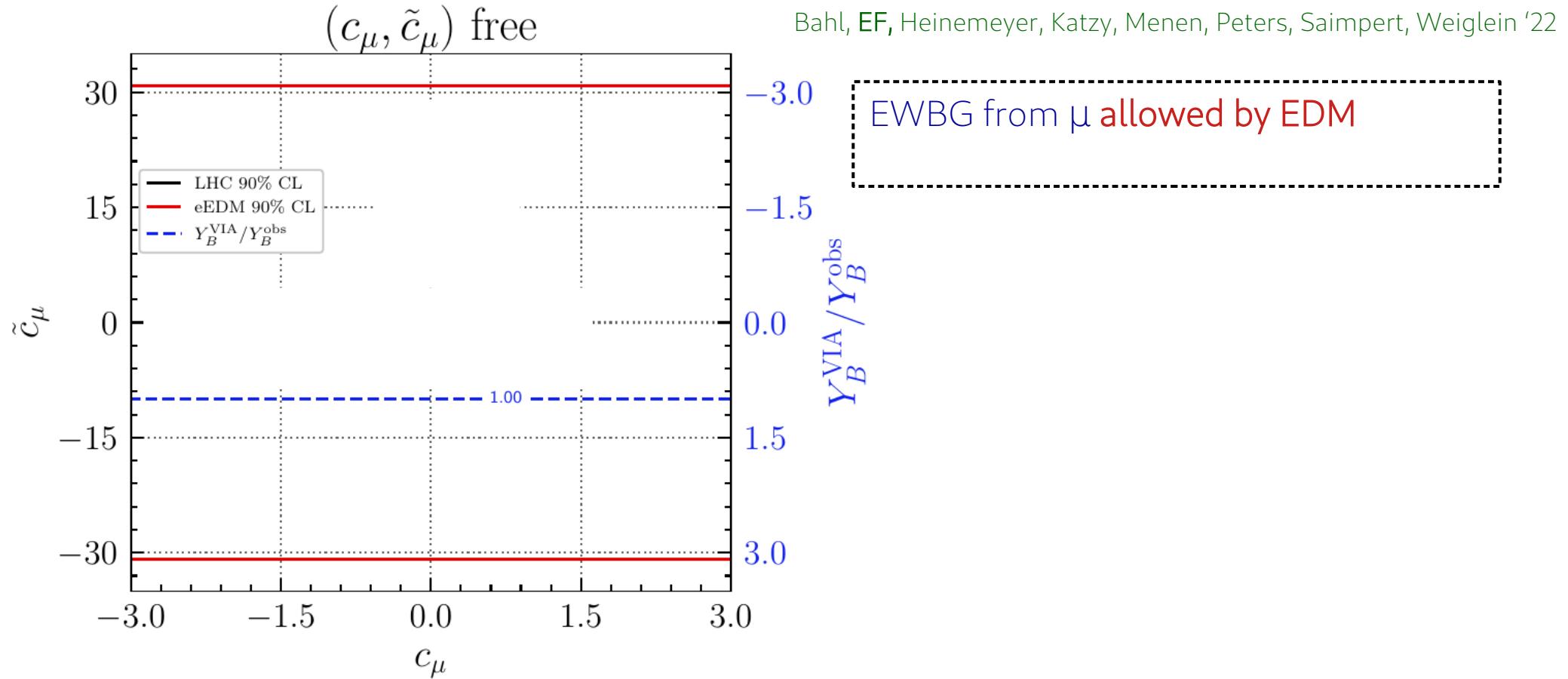
Cline, Kainulainen '20

Cline, Laurent '21, Postma '21

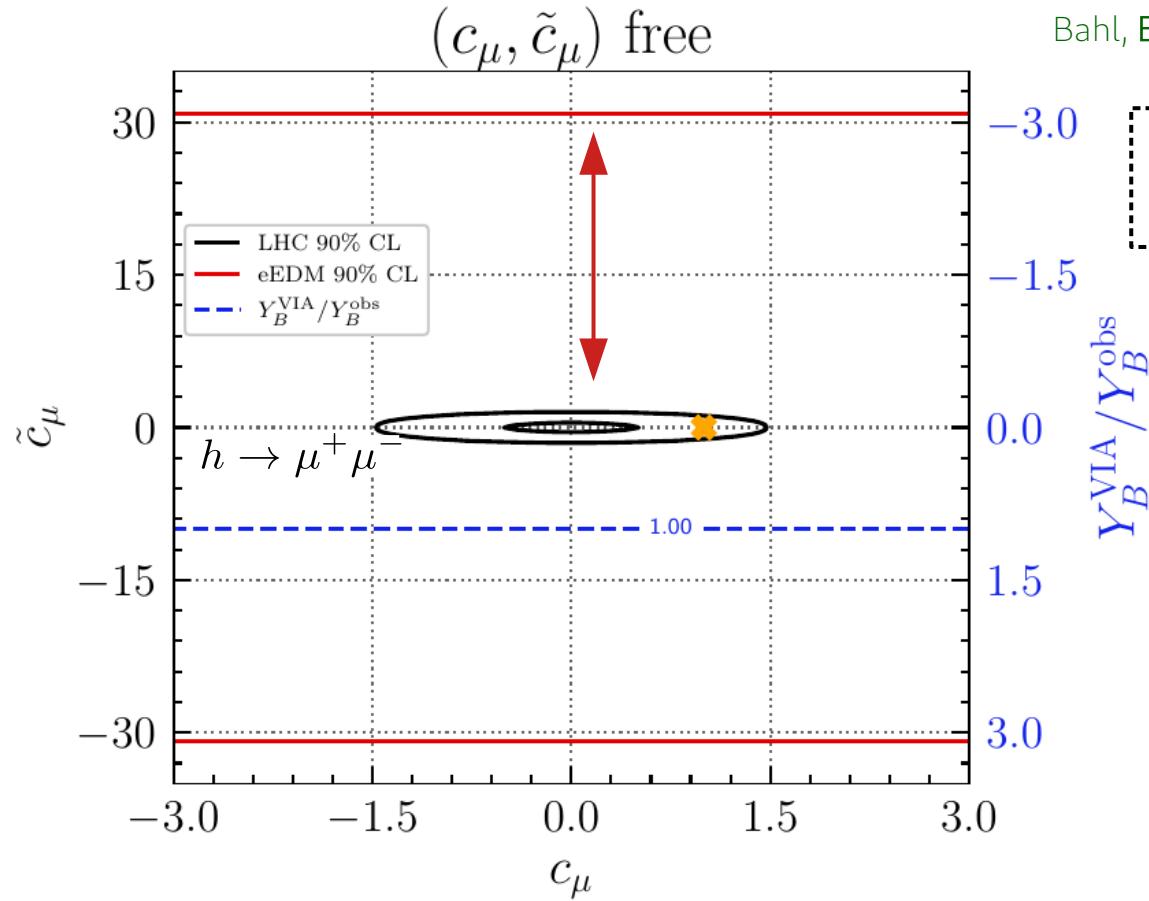
Kainulainen '21

Postma, van de Vis, White '22

Role of muon



Role of muon



Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22

EWBG from μ allowed by EDM

Excluded by LHC, but 17% contribution

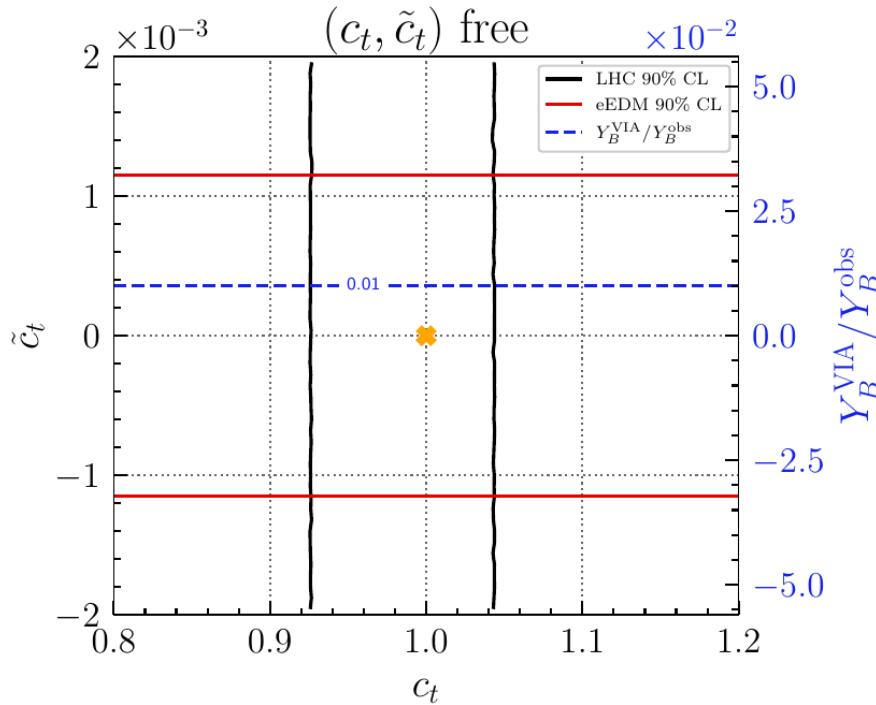
Confirmation of EF, Losada, Nir, Viernik [PRL 2020]

LHC stronger than EDM

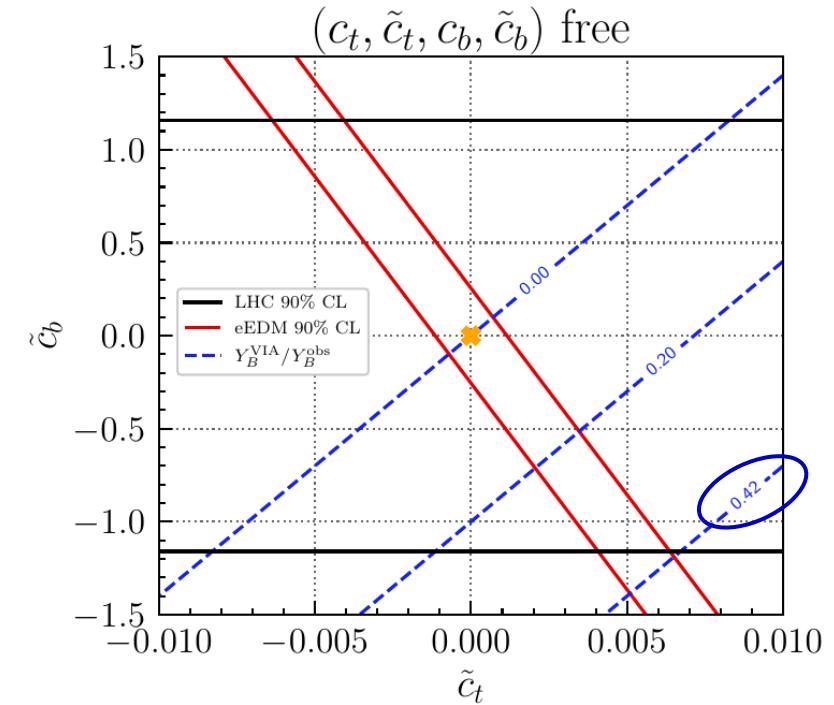
LHC probes cosmology

Complementarity

Top: EDM very constraining



t, b: cancellations of EDM allow larger CPV



t, b: each only 3-5 % of observed BAU

Combined: max. 42% of observed BAU

Summary of baryogenesis outcome

Maximal Y_B/Y_B^{obs} within LHC and EDM limits

| | t | b | c | τ | μ |
|--------|------|------|------|--------|-------|
| t | 0.03 | | | | |
| b | 0.42 | 0.05 | | | |
| c | 0.37 | 0.19 | 0.01 | | |
| τ | 6.9 | 6.9 | 6.9 | 3.2 | |
| μ | 0.18 | 0.19 | 0.16 | 3.2 | 0.16 |

- Calculated in VIA approach
 - In near-optimal benchmark scenario
- Robust upper bound

$Y_B/Y_B^{\text{obs}} < 1$

→ Disfavored by EWBG/
additional CPV needed

Summary of baryogenesis outcome

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$$Y_B/Y_B^{\text{obs}} \geq 1$$

→ EWBG possible in VIA,
but WKB might need additional CPV

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$$Y_B/Y_B^{\text{obs}} \geq 1$$

→ EWBG possible in VIA,
but WKB might need additional CPV

Also evaluated models with universal fermion coupling modifiers, and with vector coupling modifiers; investigated also complex electron Yukawa

SMEFT: Cut-off scales

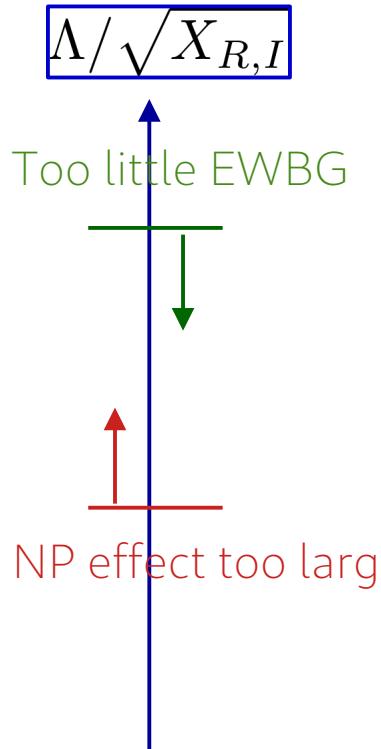
$$\Lambda/\sqrt{X_{R,I}}$$

Too little EWBG

Maximal scales for minimally required T_I (EWBG)

- $\tau: \Lambda/\sqrt{X_I^\tau} \lesssim 18 \text{ TeV } (0.01/T_I^\tau)^{1/2}$

SMEFT: Cut-off scales



Maximal scales for minimally required T_I (EWBG)

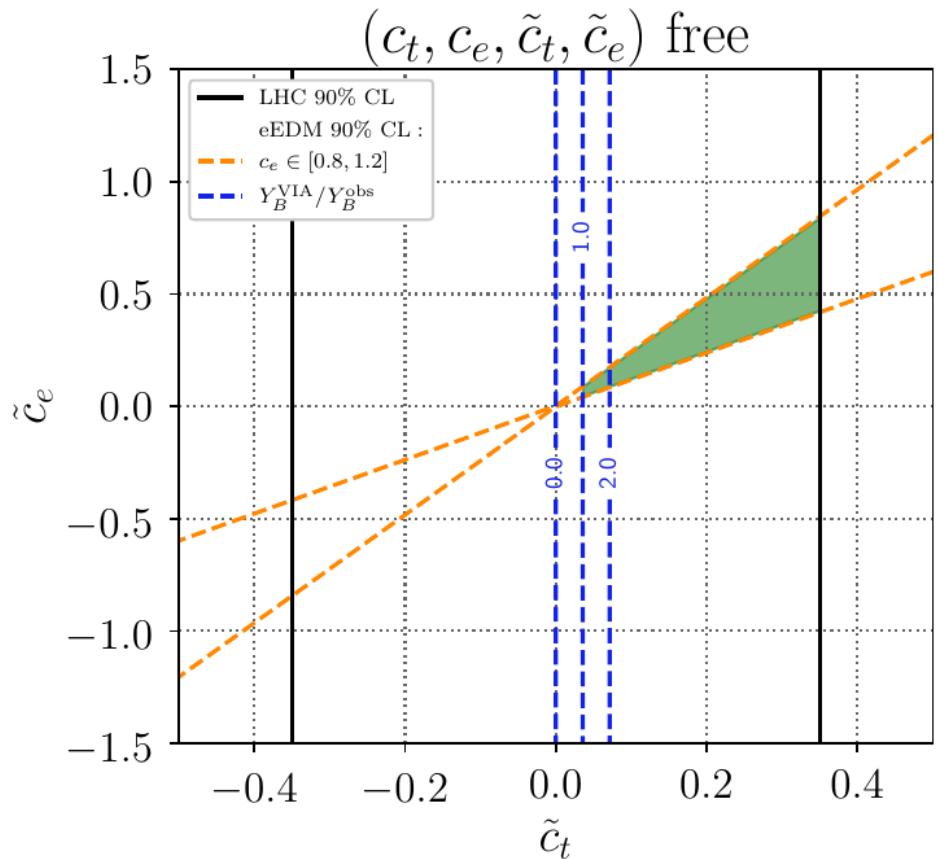
- $\tau: \Lambda/\sqrt{X_I^\tau} \lesssim 18 \text{ TeV } (0.01/T_I^\tau)^{1/2}$

Minimal scales for maximally allowed T (collider, EDM)

$$\Lambda/\sqrt{X_R^f}, \quad \Lambda/\sqrt{X_I^f} \gtrsim$$

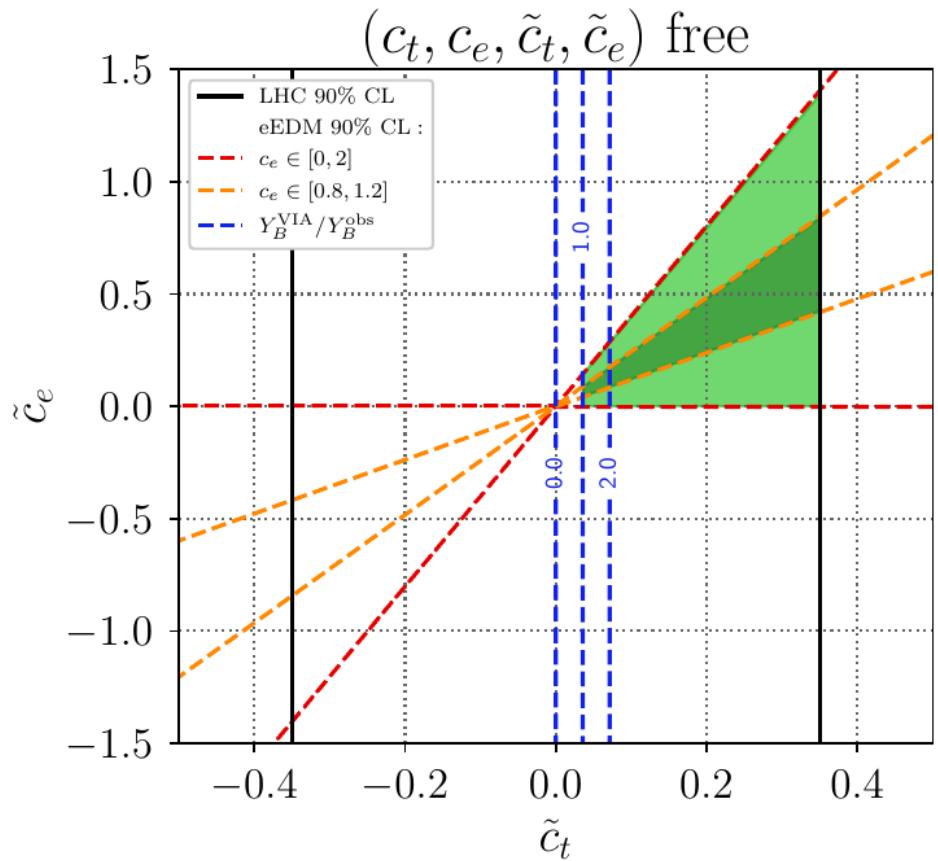
- $\tau: 2.4 \text{ TeV}, 3.1 \text{ TeV}$
- $b: 1.5 \text{ TeV}, 1.7 \text{ TeV}$
- $t: 8.7 \text{ TeV}$ from EDM
- $\mu: 10 \text{ TeV}, 12 \text{ TeV}$

Role of the electron



Interpretation of eEDM depends strongly on c_e .
If c_e small \rightarrow bound on other \tilde{c}_f much weakened

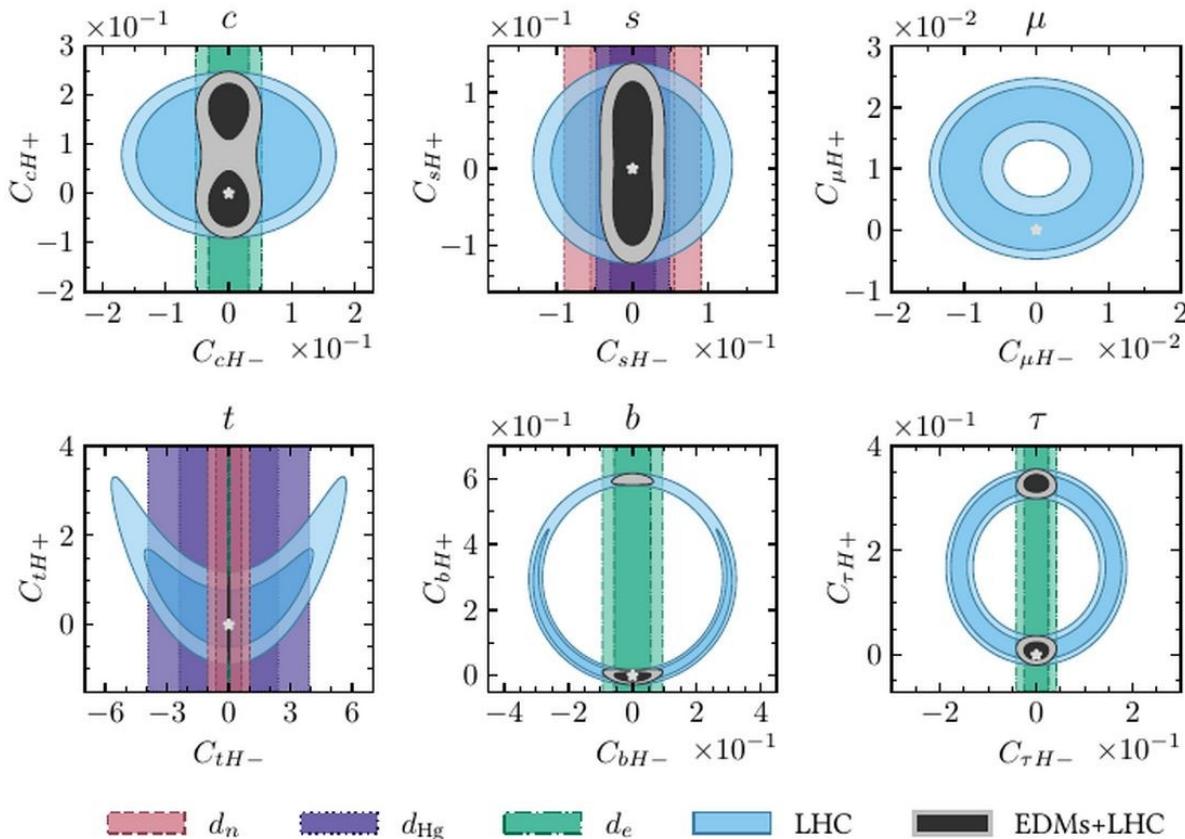
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EDMs and CPC LHC Higgs rates

Brod, Cornell, Skodras, Stamou 2203.03736



Global fit in SMEFT
in mass eigenstate basis

n, Hg, e EDMs

- RG evolution
- d_e most sensitive to c and 3rd gen.
- From 90% upper limit to likelihood: assuming Gaussian distribution of exp. uncertainty

LHC Higgs rates

- CP-conserving

$\Lambda=1$ TeV

Directions to improve tests of CPV

- Long-standing discrepancy in EWBG calculation
 - Perturbative VIA gives much larger prediction of Y_B than WKB, up to orders of magnitude
- Need likelihood from EDM bounds for global fit
- Improve (HL-)LHC studies of CPV in Higgs couplings
 - CP-odd observables
 - Machine Learning
- LHC HXSWG CPV subgroup in WG2, e.g.
 - CPV Benchmarks for UV models and EFT
 - Interplay of LHC and EDMs

Investigate further to which extent CPV in Higgs couplings can account for EWBG

Conclusions

- Complementarity of EDM, EWBG and LHC Higgs physics
- $H \rightarrow \tau\tau$ CP analysis excludes large \tilde{c}_τ , but τ remains viable EWBG source (VIA LO)
- LHC constrains cosmological scenarios, separates flavors; now also 2nd gen.
- Cancellations and enhancements with 2 fermions, e.g. t+b: few % $\rightarrow \sim 40\%$ of obs. Y_B
- Electron Yukawa has big impact on interpretation of electron EDM
- SMEFT generates Yukawa modifications, preferred scale $\Lambda / \sqrt{X_I} \sim \text{few-10-20 TeV}$

THANK YOU!

BACKUP

$$T_R, T_I, Y_f$$

Relation between SM mass and Yukawa fixes Y_f (a priori free coefficient of dim-4 term)

$T_R, T_I, Y_f \rightarrow 2$ free parameters per fermion: T_R, T_I

Modification of each vertex w.r.t. SM

$$r_f(T_R^f, T_I^f) \equiv \frac{|\lambda_f|^2 / |\lambda_f^{\text{SM}}|^2}{|m_f|^2 / |m_f^{\text{SM}}|^2} = \frac{(1 + 3T_R^f)^2 + 9T_I^{f2}}{(1 + T_R^f)^2 + T_I^{f2}}$$

production,
decay



Total Higgs width

$$\Gamma_h / \Gamma_h^{\text{SM}} = 1 + \text{BR}_f^{\text{SM}}(r_f - 1)$$

Transport equations

$$\partial f \equiv \partial_\mu f^\mu \approx v_w f' - D_f f'' \quad \text{Diffusion approximation}$$

$$\partial t = -\Gamma_M^t \mu_M^t - \Gamma_Y^t \mu_Y^t + \Gamma_{ss} \mu_{ss} + S_t$$

$$\partial b = -\Gamma_M^b \mu_M^b - \Gamma_Y^b \mu_Y^b + \Gamma_{ss} \mu_{ss} + S_b$$

$$\partial q = -\partial t - \partial b$$

$$\partial \tau = -\Gamma_M^\tau \mu_M^\tau - \Gamma_Y^\tau \mu_Y^\tau + S_\tau$$

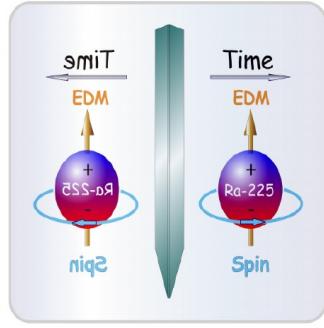
$$\partial l = -\partial \tau$$

$$\partial h = +\Gamma_Y^t \mu_Y^t - \Gamma_Y^b \mu_Y^b - \Gamma_Y^\tau \mu_Y^\tau$$

$$\partial u = +\Gamma_{ss} \mu_{ss} .$$

Electron's Electric Dipole Moment

[Hewett, Weerts et al '12]



EDM violates \mathcal{T} and \mathcal{P}
 $\Rightarrow \mathcal{CP}$

$$\frac{d_e^f}{e} \propto \left(\frac{Y^f}{Y_{\text{SM}}^f} \right)^2 T_I^f$$

ACME [Nature '18]:

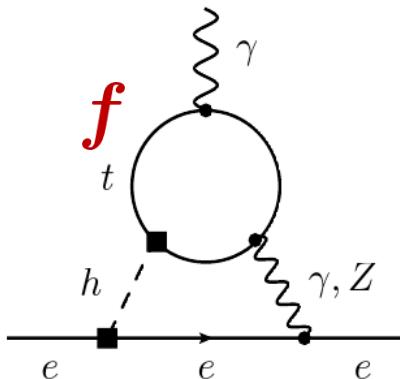
$$d_e \leq 1.1 \times 10^{-29} e \text{ cm at 90\% CL}$$

Using [Panico, Pomarol, Riembau '18],
 see also [Brod, Haisch, Zupan '13], [Brod, Stamou '18], ...

$$\frac{d_e^{(\ell)}}{e} \simeq -4Q_\ell^2 \frac{e^2}{(16\pi^2)^2} \frac{m_e m_\ell}{m_h^2} \frac{v}{\Lambda^2} \mathbf{X}_I^\ell \left(\frac{\pi^2}{3} + \ln^2 \frac{m_\ell^2}{m_h^2} \right), \quad \ell = \tau, \mu$$

$$\frac{d_e^{(\mathbf{b})}}{e} \simeq -4N_c Q_b^2 \frac{e^2}{(16\pi^2)^2} \frac{m_e m_b}{m_h^2} \frac{v}{\Lambda^2} \mathbf{X}_I^{\mathbf{b}} \left(\frac{\pi^2}{3} + \ln^2 \frac{m_b^2}{m_h^2} \right)$$

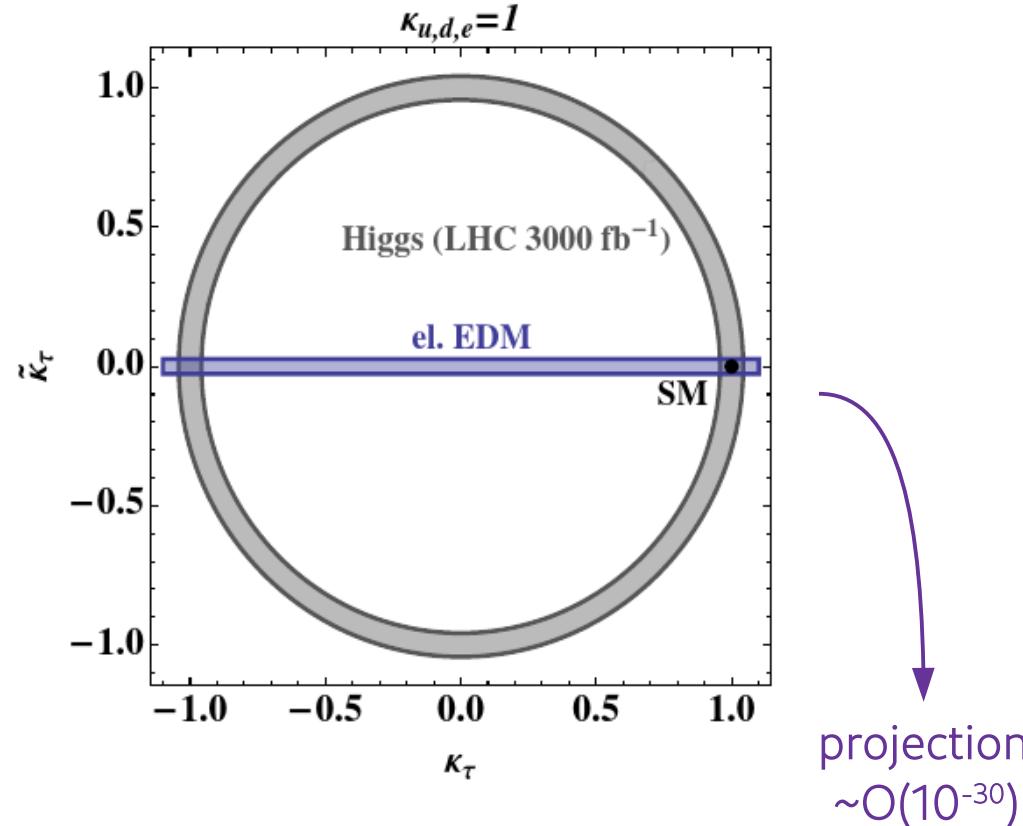
$$\frac{d_e^{(\mathbf{t})}}{e} \simeq -\frac{16}{3} \frac{e^2}{(16\pi^2)^2} \frac{m_e}{m_t} \frac{v}{\Lambda^2} \mathbf{X}_I^{\mathbf{t}} \left(2 + \ln \frac{m_t^2}{m_h^2} \right)$$



EDM & LHC limits CPV Yukawas

tau

Brod, Haisch, Zupan '13 (also for t, b)

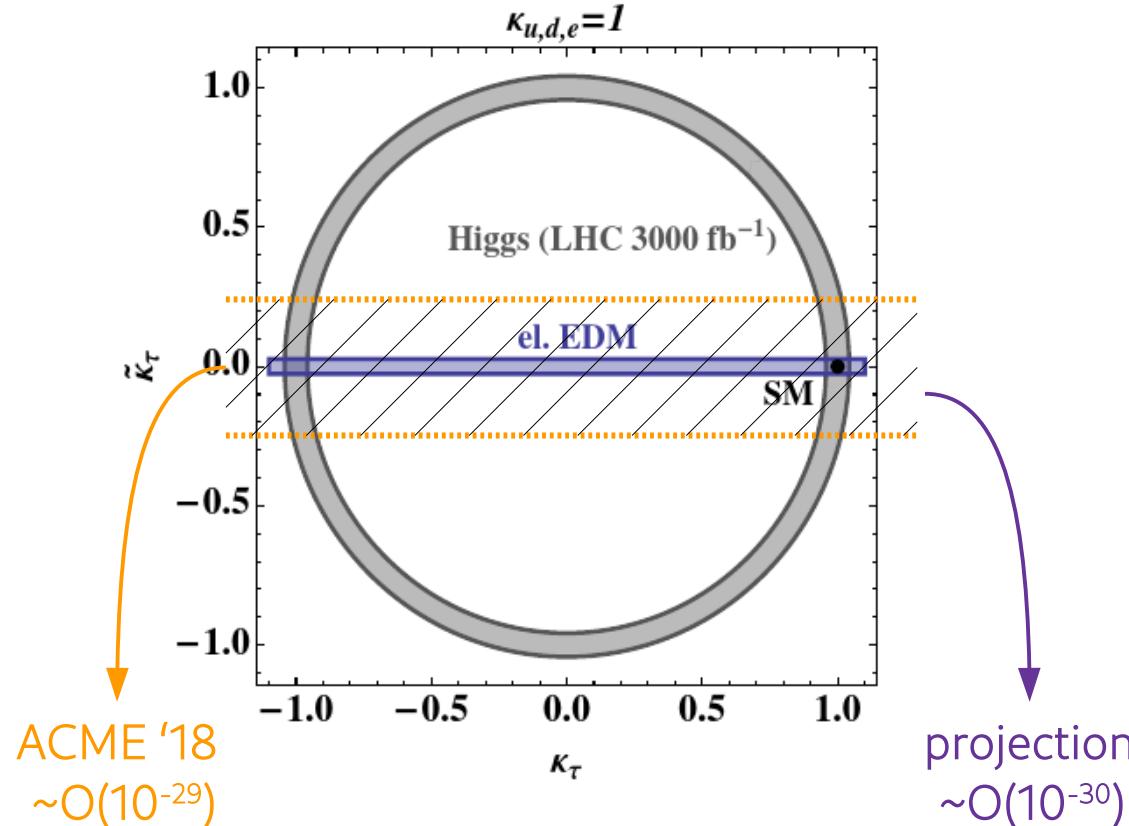


$d_e [\text{e cm}]$:

EDM & LHC limits CPV Yukawas

tau

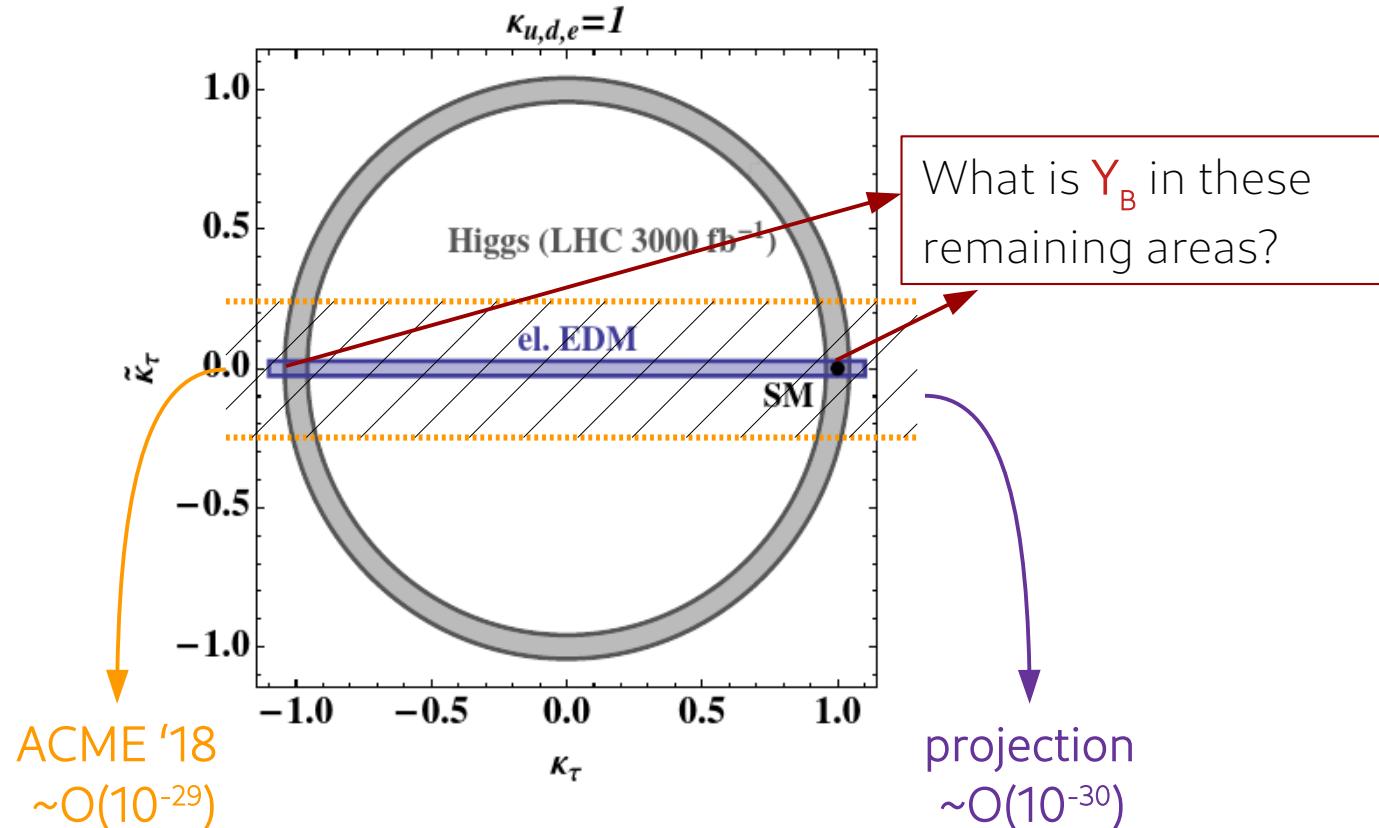
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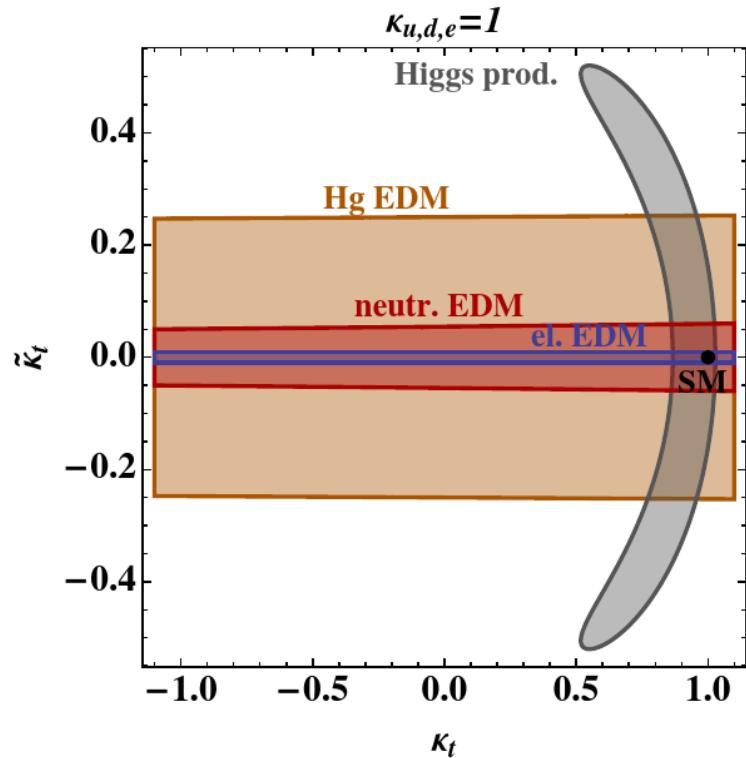
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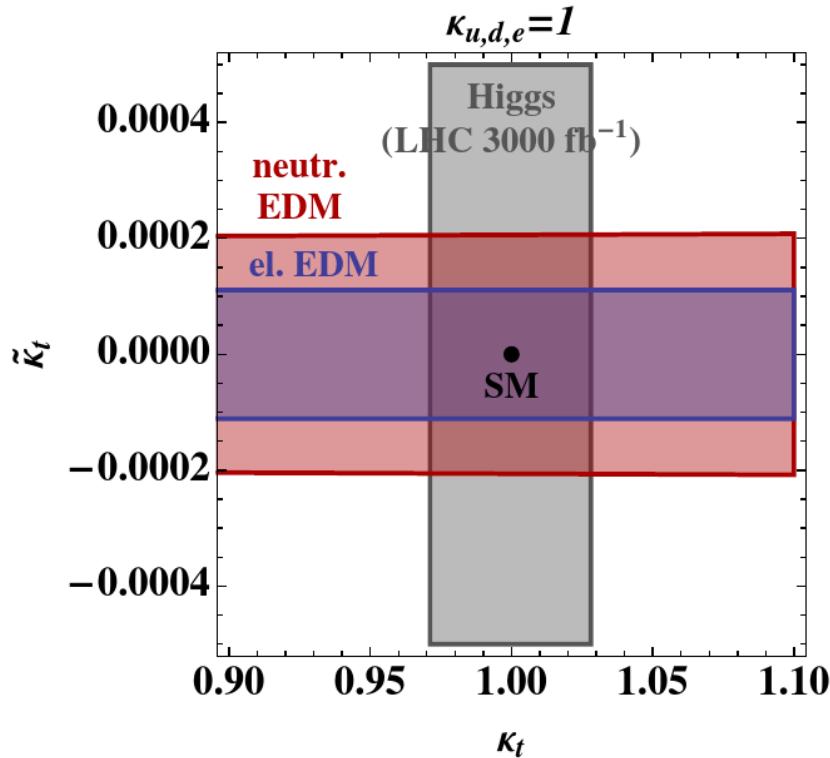
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top

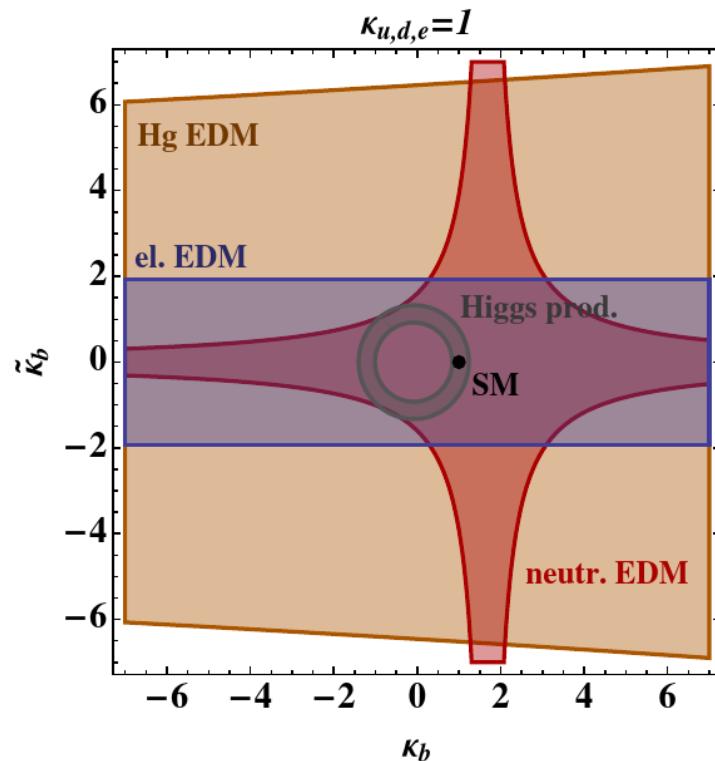


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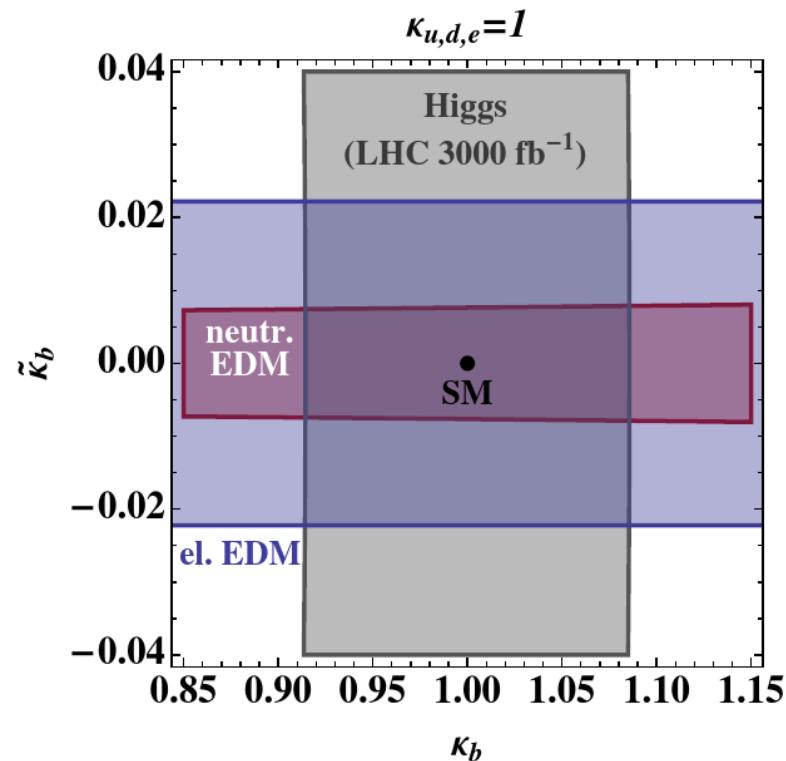


EDM & LHC limits CPV Yukawas

bottom

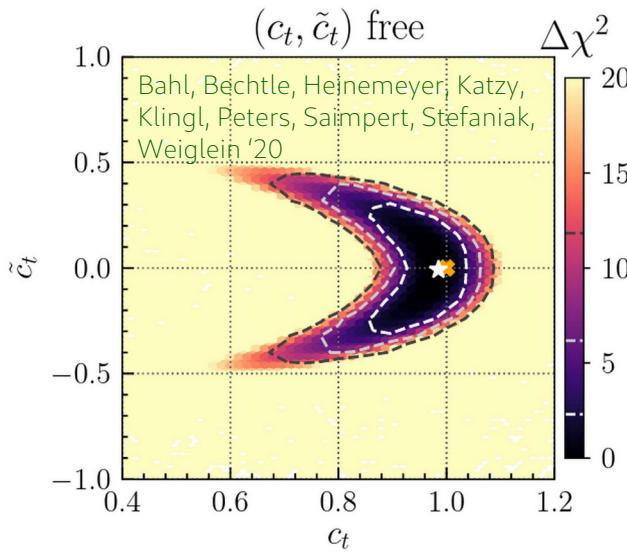


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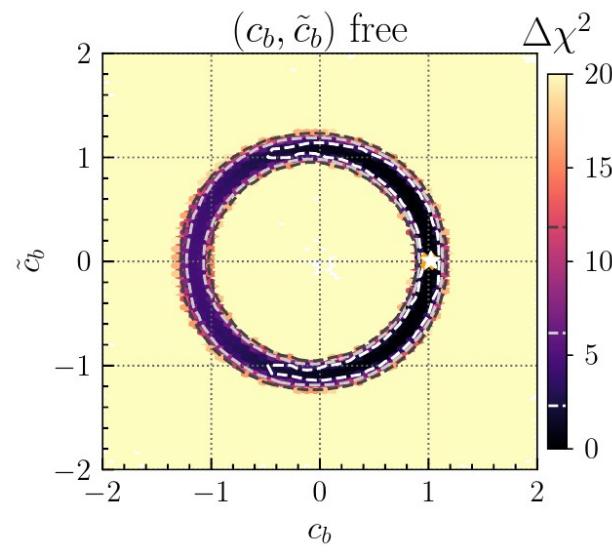


Top, bottom, and their combination

Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22



Top: ellipse (ggF) cut
off by $h \rightarrow \gamma\gamma$

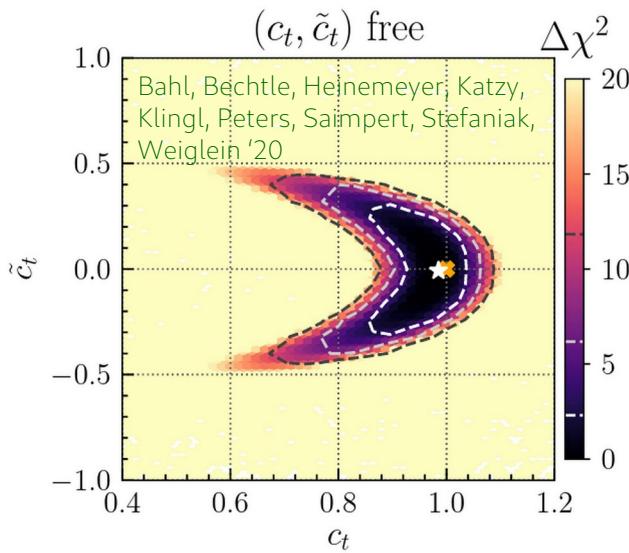


Bottom: ring ($h \rightarrow bb$)
reduced by ggF
(positive interference with t)

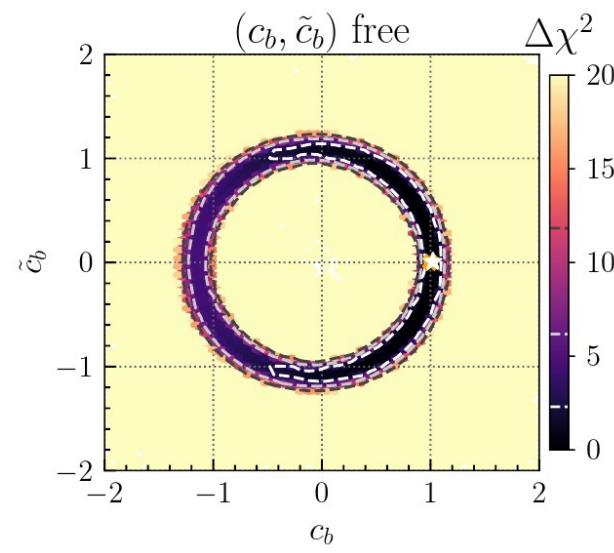
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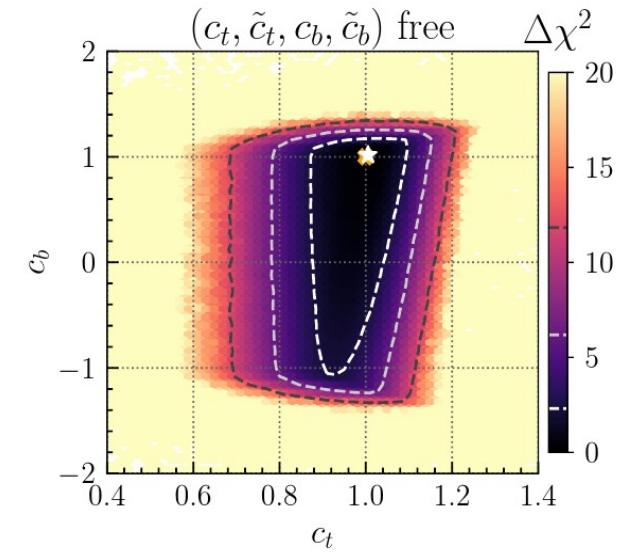
Floating several coupling modifiers simultaneously



Top: ellipse (ggF) cut off by $h \rightarrow \gamma\gamma$



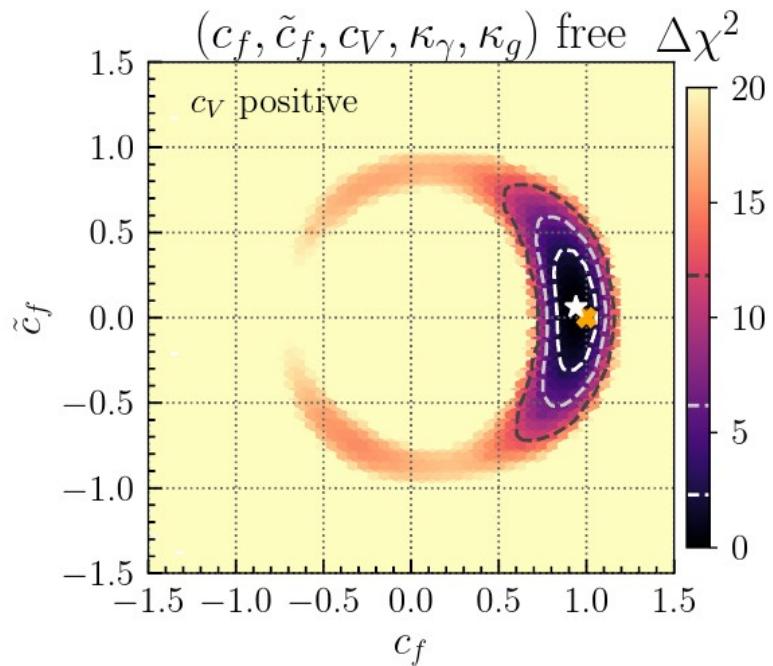
Bottom: ring ($h \rightarrow bb$) reduced by ggF
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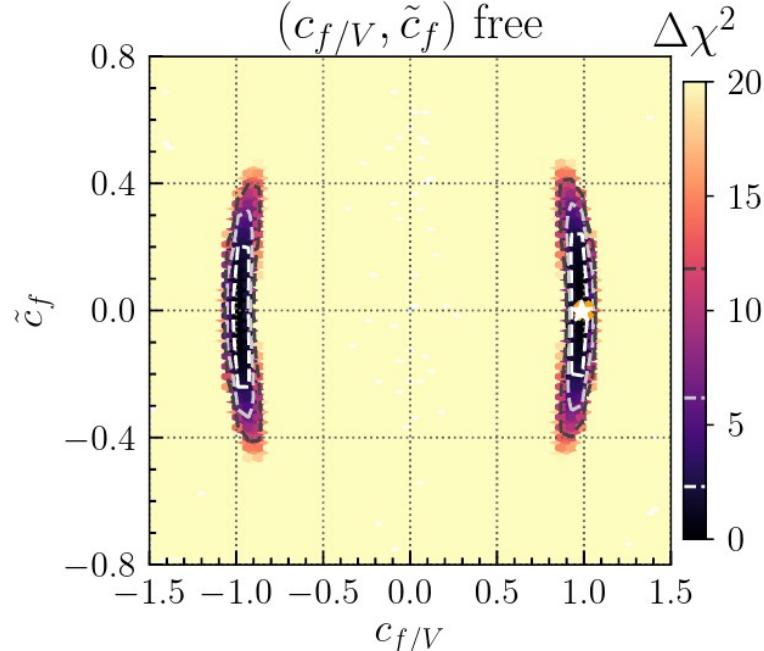
t+b: small c_b can be compensated by \tilde{c}_b

Varying vector couplings

Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22



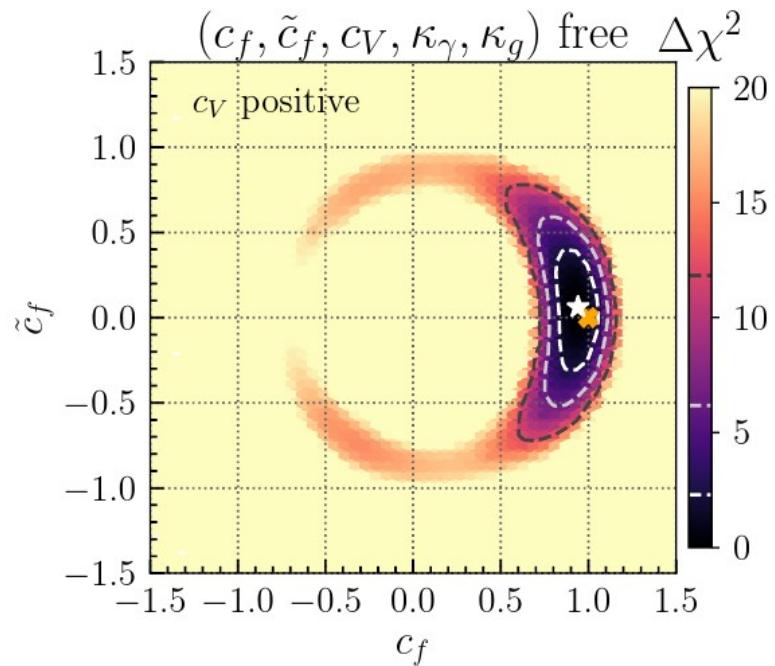
Universal fermion
coupling modifier:
Dominated by top



General mixing scenario: $c_f = c_V$
No CPV included in vector couplings

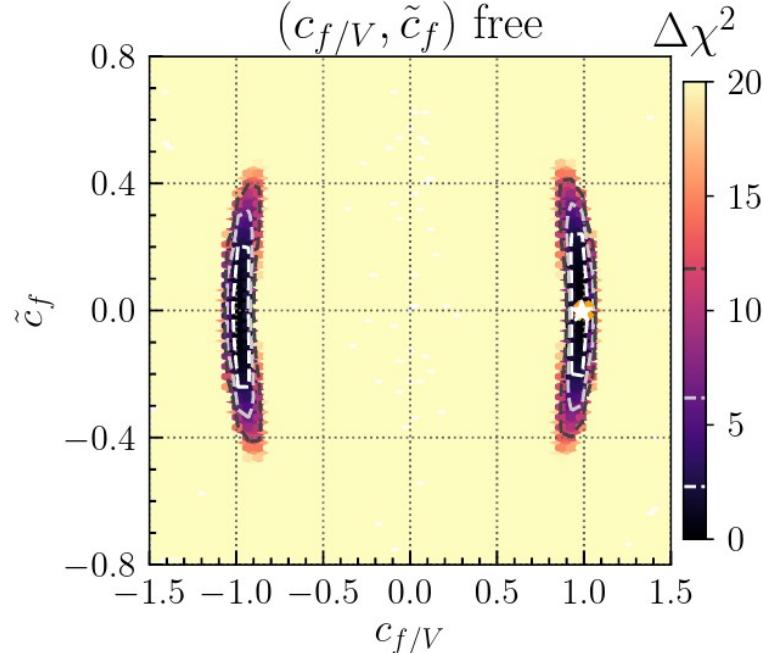
Varying vector couplings

Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22



Universal fermion coupling modifier:
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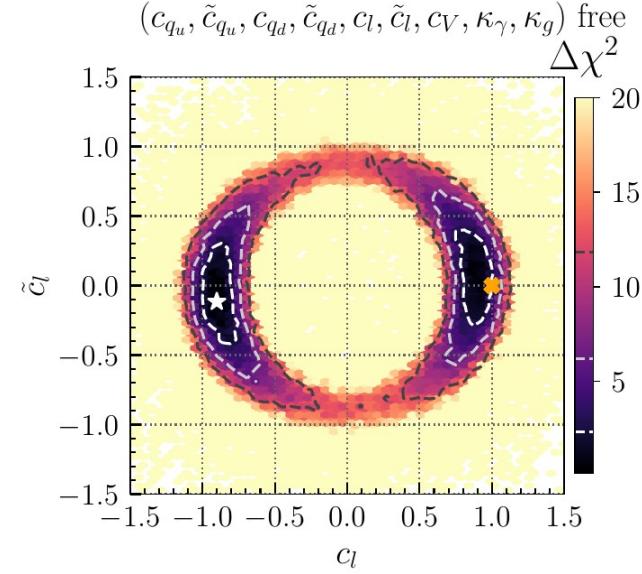
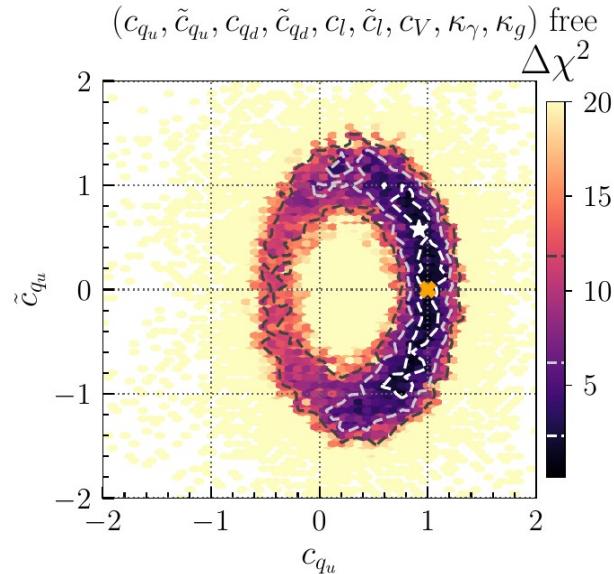
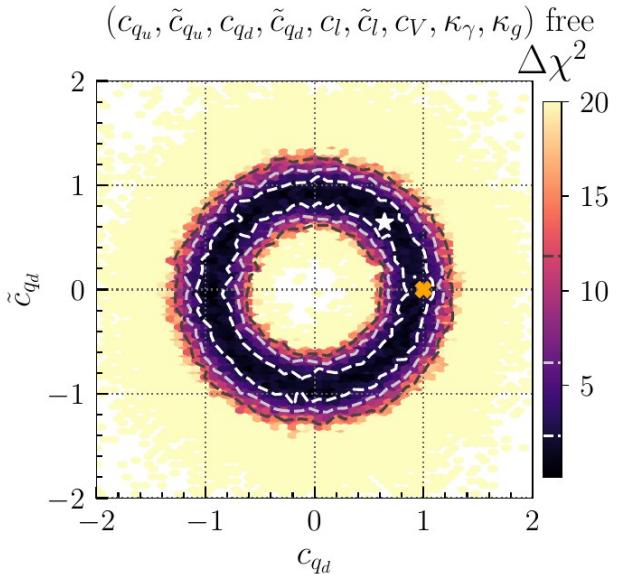
+ fitted more models
with up to 9 free parameters



General mixing scenario: $c_f = c_V$
No CPV included in vector couplings

General model: 9-parameter fit

Bahl, EF, Heinemeyer, Katzy, Menen, Peters, Saimpert, Weiglein '22



Lepton vs Quark Source

Lepton vs Quark Source

- Lepton advantages:
 - No strong sphaleron washout
 - Large diffusion
 - τ : still sizeable Yukawa
 - μ : weak EDM bound

Lepton vs Quark Source

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 - Large diffusion
 - τ : still sizeable Yukawa
 - μ : weak EDM bound
- Robustness: τ overshoots Y_b^{obs}
 - $O(1)$ uncertainties do not change conclusion
 - Quarks larger uncertainties

Lepton vs Quark Source

- Lepton advantages:
 - No strong sphaleron washout
 - Large diffusion
 - τ : still sizeable Yukawa
 - μ : weak EDM bound
- Robustness: τ overshoots Y_b^{obs}
 - $O(1)$ uncertainties do not change conclusion
 - Quarks larger uncertainties
- Benchmark choices:
 - Wall velocity, thickness, ...
 - → investigated impact in 2007.06940, see also Postma, van de Vis, White '16; de Vries, Postma, van de Vis '18;

2-step: baryon density from L density

Solve for left-handed particle density or directly baryon density

$$n_b''(z) - \frac{v_w}{D_q} n_b'(z) = \frac{\Gamma_{ws}(z)}{D_q} \left(\mathcal{R}n_b(z) + \frac{3}{2} n_L(z) \right) \equiv \frac{\Gamma_{ws}(z)}{D_q} \mathcal{R}n_b + f(z)$$

$$\begin{aligned} Y_B &= \frac{n_b(z > 0)}{s} = \frac{A_1}{s} = \frac{1}{s} \left(1 - \frac{\alpha_-}{\alpha_+} \right) B_1 = \frac{k}{D_q \alpha_+ s} B_1 \\ &= \frac{3\Gamma_{ws}}{2D_q \alpha_+ s} \int_0^{-\infty} e^{-\alpha_- x} n_L(x) dx . \end{aligned}$$

Chemical potentials

$$\mu_M^t = \frac{t}{k_t} - \frac{q}{k_q},$$

$$\mu_M^b = \frac{b}{k_b} - \frac{q}{k_q},$$

$$\mu_M^\tau = \frac{\tau}{k_\tau} - \frac{l}{k_l},$$

$$\mu_Y^t = \frac{t}{k_t} - \frac{q}{k_q} - \frac{h}{k_h},$$

$$\mu_Y^b = \frac{b}{k_b} - \frac{q}{k_q} + \frac{h}{k_h},$$

$$\mu_Y^\tau = \frac{\tau}{k_\tau} - \frac{l}{k_l} + \frac{h}{k_h},$$

$$\mu_{ss} = \sum_{i=1}^3 \frac{2q_i}{k_{q_i}} - \frac{u_i}{k_{u_i}} - \frac{d_i}{k_{d_i}}.$$

Matrix formalism: o.d.e. 1st order

$$\begin{pmatrix} t' \\ b' \\ \vdots \\ g'_t \\ g'_b \\ \vdots \end{pmatrix} - \begin{pmatrix} 0 & 1 & & & \\ & 0 & 1 & & \\ & & \ddots & \ddots & \\ & & & \frac{v_w}{D_t} & \\ & & & & \frac{v_w}{D_b} \\ & & & & \ddots \end{pmatrix} \begin{pmatrix} t \\ b \\ \vdots \\ g_t \\ g_b \\ \vdots \end{pmatrix} = \begin{pmatrix} 0 & & & & \\ & \ddots & & & \\ & & \frac{\Gamma_t}{D_t k_t} & & \\ & & & \frac{\Gamma_b}{D_b k_b} & \\ & & & & \ddots \\ & & & & & 0 \end{pmatrix} \begin{pmatrix} t \\ b \\ \vdots \\ g_t \\ g_b \\ \vdots \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ \vdots \\ S_t/D_t \\ S_b/D_b \\ \vdots \end{pmatrix}$$

$$\iff \bar{\chi}' - K\bar{\chi} = \bar{S},$$

$$K \equiv \begin{pmatrix} 0 & I_N \\ \Gamma & V \end{pmatrix}.$$

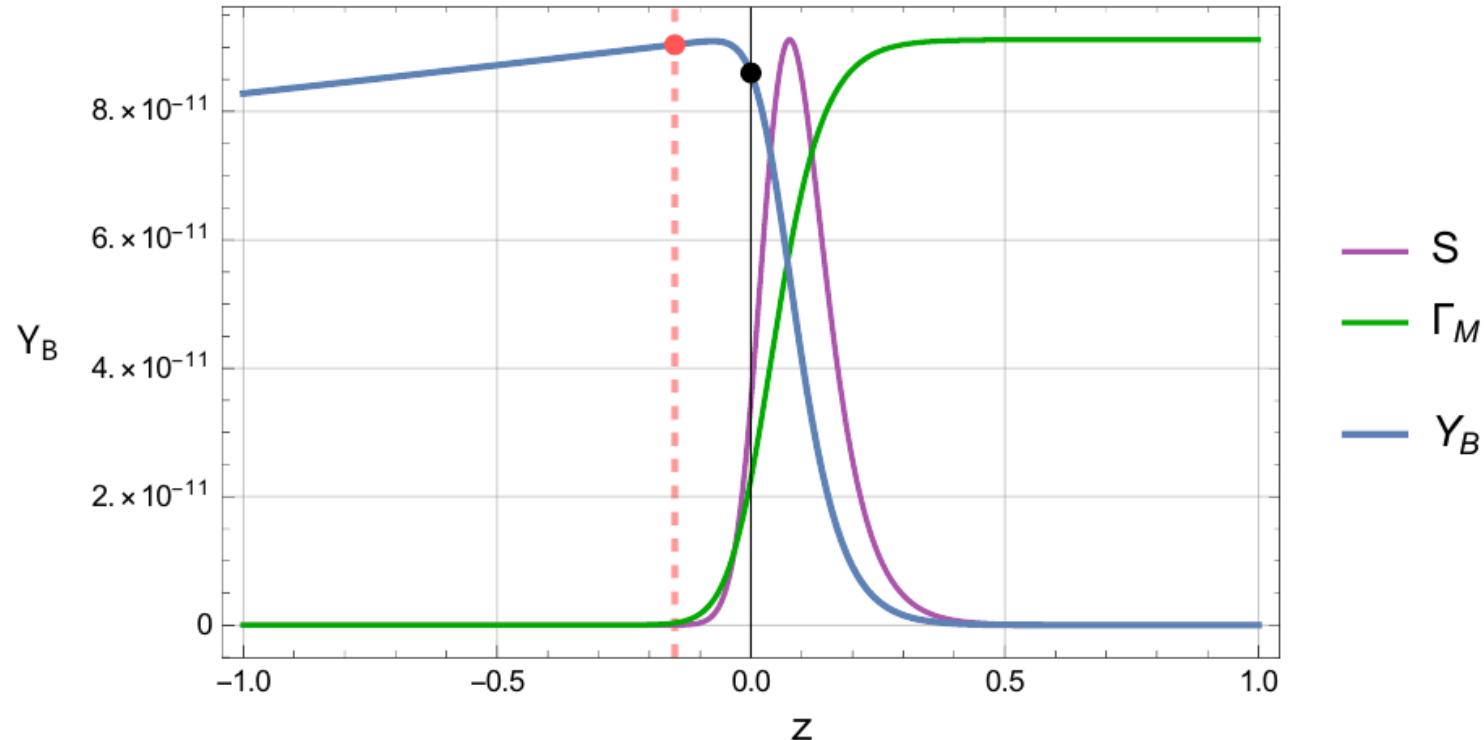
$$\det(K) = \det(-\Gamma) = (-1)^N \det(\Gamma)$$

Fuchs, Losada, Nir, Viernik '20

Particle dynamics

- CPV interactions across the expanding bubble wall **generate a chiral asymmetry**
- CPC interactions **wash out** the generated asymmetry
- Strong sphaleron process produces further washout in the quark sector
- Some of the remaining asymmetry **diffuses** into the symmetric phase; more efficient for leptons than quarks.
- Weak sphaleron process is efficient only in the symmetric phase, acting on left-handed multiplets and changing baryon number.
- Finally, the bubble wall catches up and freezes in the resulting baryon number density in the **broken phase**.

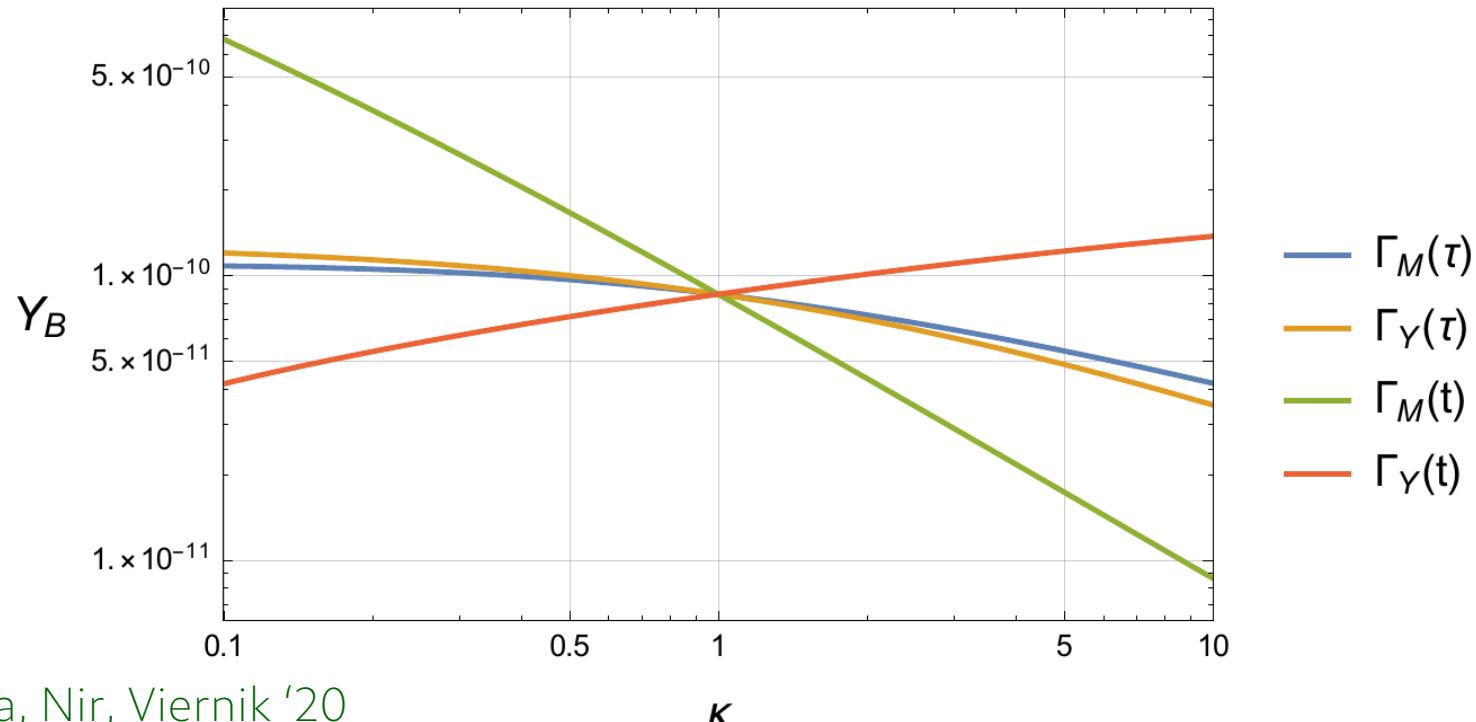
Thin-wall approximation



Fuchs, Losada, Nir, Viernik '20

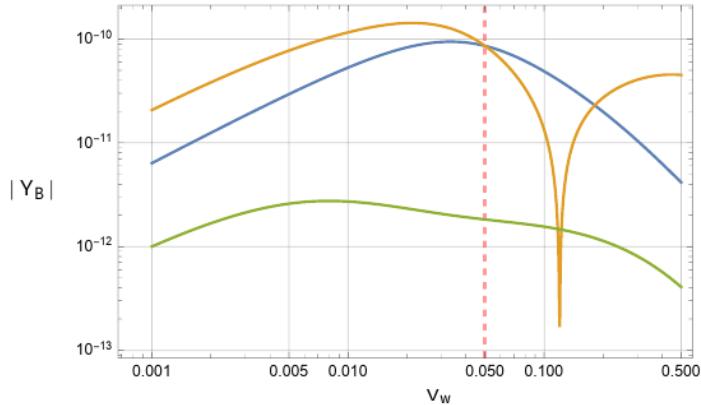
Uncertainty of Y_B from input rates

$$\Gamma_{M/Y}^f \rightarrow \kappa_{M/Y}^f \Gamma_{M/Y}^f$$

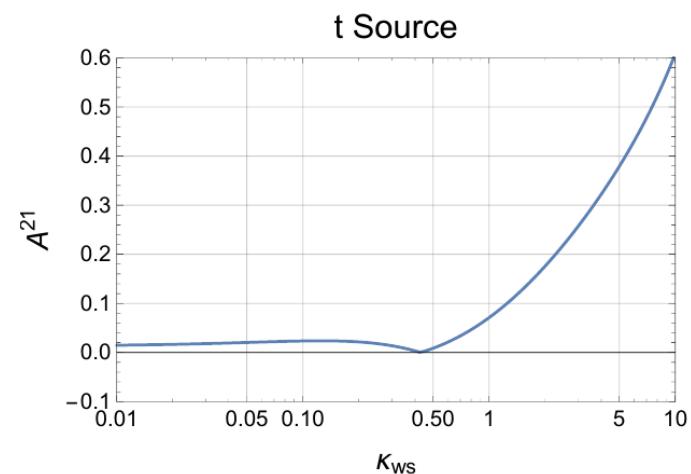
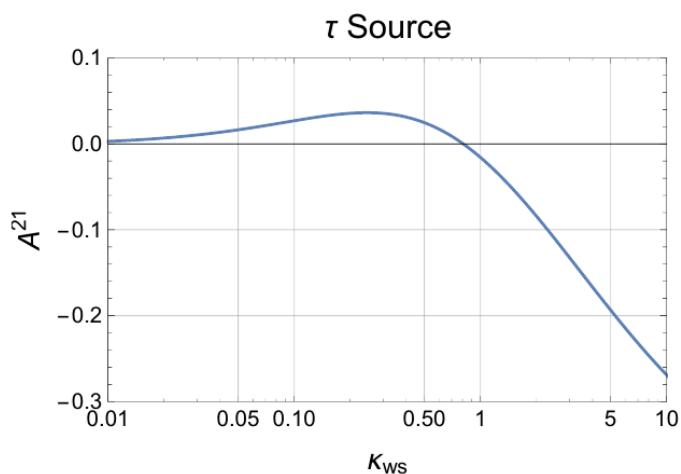


Fuchs, Losada, Nir, Viernik '20

Parameter dependence; 1-/2-step



Fuchs, Losada, Nir, Viernik '20



CP violation in the Higgs sector

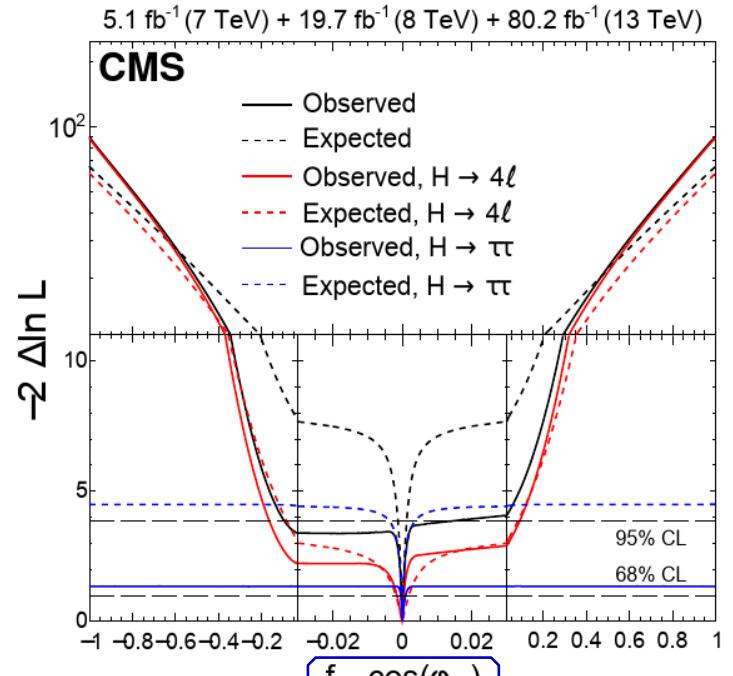
- Discovered Higgs compatible with $J^{PC}=0^{++}$
- Small CP-odd component possible

CP violation in the Higgs sector

- Discovered Higgs compatible with $J^{PC}=0^{++}$
- Small CP-odd component possible

Until recently:
mostly searches
for CPV in hVV

CMS 1903.06973: HVV anomalous couplings



Fraction of CP-odd cross section contribution

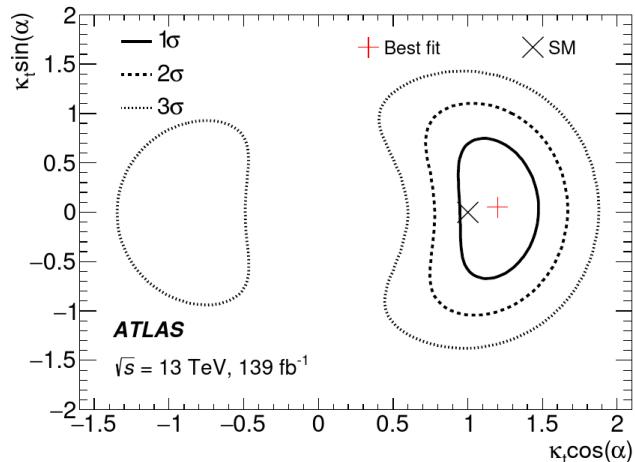
Phase of CP-odd coefficient

2020: LHC bounds on CPV in Yukawas

top $t\bar{t}h + t\bar{t}h \rightarrow \gamma\gamma$

ATLAS 2004.04545

See also
CMS 2003.10866



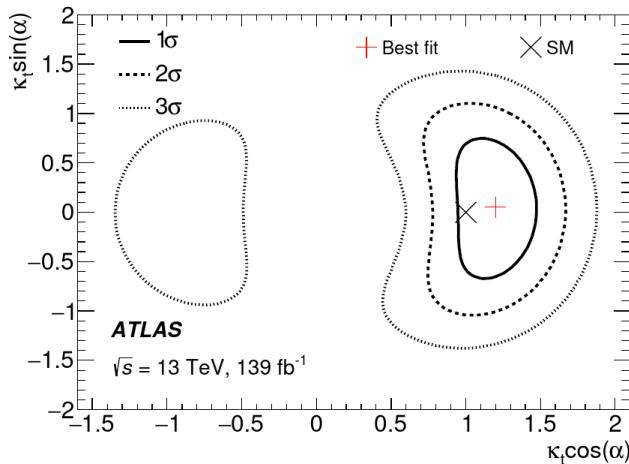
At 95% CL: $|\alpha| < 43^\circ$ (CP mixing angle)

Pure CP-odd excluded at: 3.9σ

2020: LHC bounds on CPV in Yukawas

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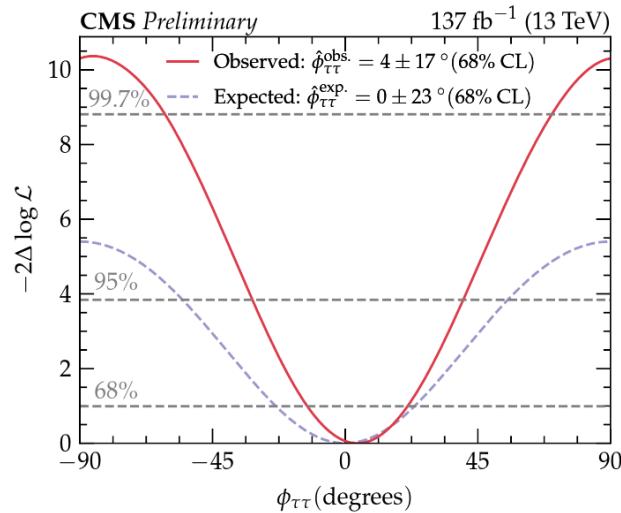


At 95% CL: $|\alpha| < 43^\circ$ (CP mixing angle)

Pure CP-odd excluded at: 3.9σ

See also
CMS 2003.10866

tau CMS HIG-20-006 $h \rightarrow \tau\tau$

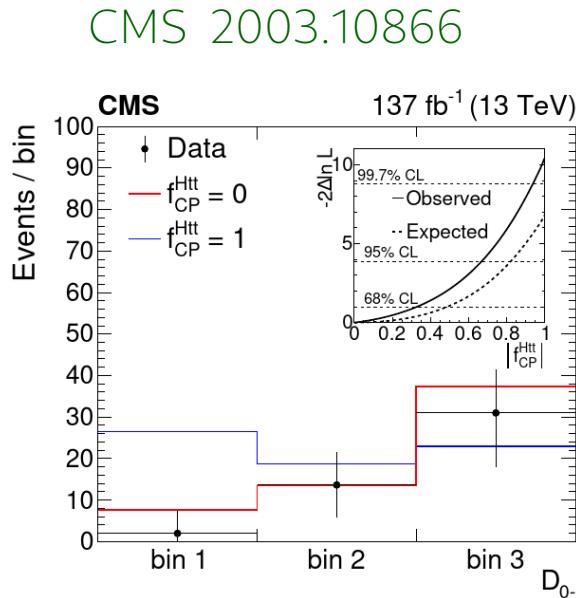


Sensitivity to $\phi_\tau \sim 4 \pm 17^\circ$

CMS htt CPV

$$f_{\text{CP}}^{\text{Htt}} = \frac{|\tilde{\kappa}_t|^2}{|\kappa_t|^2 + |\tilde{\kappa}_t|^2} \text{ sign}(\tilde{\kappa}_t / \kappa_t)$$

$$|f_{\text{CP}}^{\text{Htt}}| < 0.67$$



Many proposals

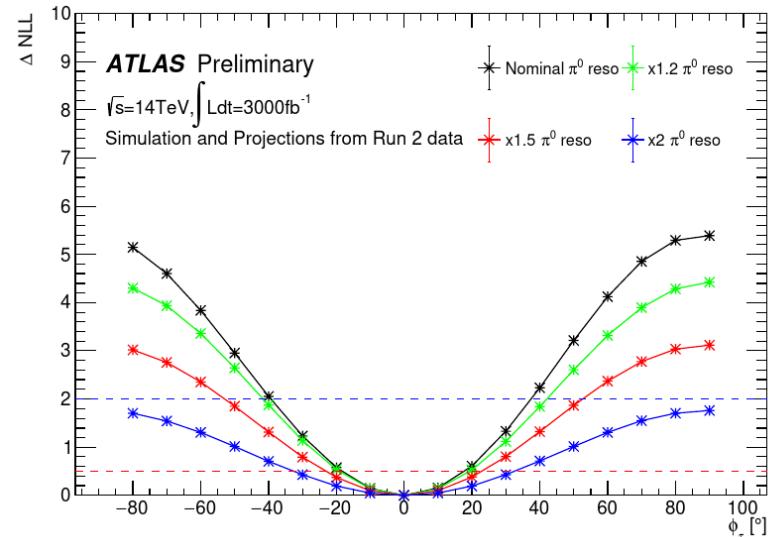
Bortolato, Kamenik, Košnik, Smolkovič '20

Pure CP-odd excluded at: 3.2σ

CPV ATLAS $h \rightarrow \tau\tau$: Prospects for HL

This note presents a study for the prospective measurement of the \mathcal{CP} quantum number of the Higgs boson coupling to τ leptons with 3000 fb^{-1} of proton–proton collisions at $\sqrt{s} = 14 \text{ TeV}$ using the ATLAS detector at the HL-LHC. Only $H \rightarrow \tau\tau$ events where both τ leptons decay via the $\tau^\pm \rightarrow \rho^\pm \nu_\tau \rightarrow \pi^0 \pi^\pm \nu_\tau$ chain are analysed and the acoplanarity angle $\varphi_{\mathcal{CP}}^*$, the angle between the planes spanned by the pion pairs, is used to determine the \mathcal{CP} -mixing angle. It is shown that considering only statistical uncertainties, a pseudoscalar Higgs boson can be excluded at 95% confidence level. The \mathcal{CP} -mixing angle can be measured with a statistical precision ranging between $\pm 18^\circ$ and $\pm 33^\circ$, depending on the precision of the π^0 reconstruction

$$\mathcal{L} = g_{\tau\tau} (\cos(\phi_\tau) \bar{\tau}\tau + \sin(\phi_\tau) \bar{\tau} i\gamma_5 \tau) h$$



ATLAS PHYS-PUB-2019-008

EDMs: e, n, Hg

