

Large N expansion of two dimensional QCD

Federico Ambrosino

Supervisors: Shota Komatsu, Joao Caetano

August 5, 2022

CERN - Department of Theoretical Physics, University of Oxford



Plan

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- Does $SU(3)$ YM exhibit confinement in **4** dimensions?
- If so, what is the mechanism responsible for this?
- What is the **full** spectrum of the theory?

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- What is the **mesons'** spectrum? And their dynamics?

Why QCD in 1+1 dimensions?

Old idea... but renewed interest

- Field Theory perspective:

- String theory perspective:

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 - Recent success in worldsheet constructions in lower dimensions (e.g. AdS_3 / CFT_2): can we do the same here?

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Overall: **simpler** and **insightful**

Quantum Chromodynamics on the line

An $SU(N)$ gauge theory coupled with N_f fundamental & massive fermions:

$$\mathcal{L} = -\frac{1}{4} \text{Tr} F \wedge \star F + \bar{\Psi}_A (i\not{D} - m_A) \Psi_A$$

Convenient to use **lightcone coordinates**: $x_{\pm} = \frac{1}{\sqrt{2}}(x_1 \pm x_0)$.

Massive simplification: fix gauge redundancy through **lightcone gauge**:
 $A_- = A^+ = 0$.

$$\mathcal{L}_{gf} = -\frac{1}{2} (\partial_- A_+)^2 + \bar{\Psi}_A (i\not{D} - m_A - g\gamma_- A_+) \Psi_A$$

A confining potential for quarks

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Solution:

$$A_+ = -\frac{g}{2} \int dy_- |x_- - y_-| j^+(x_+, y_-)$$

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Inserting A_+ on-shell:

$$\mathcal{L}_{gf} \supset \frac{g}{2N} \int dx_- dy_- \text{Tr} j_+(x_+, x_-) |x_- - y_-| j^+(x_+, y_-)$$

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A linear (confining) potential between quarks!

A large N expansion

A linear potential arises. . . but this is still not enough!

Goal: showing that there are no free quarks in the spectrum.

Further simplification: large N limit $N \rightarrow \infty$ with $g^2 N$ fixed

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- QCD has no *small parameter* in the IR phase
- W.r.t. $g^2 N$, diagrams organise in a **$1/N$ expansion** \Rightarrow **Perturbation theory!**
- $3 \neq \infty \dots$ but maybe $3 \gg 1!$ Again: **toy model** for **universal features**

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Large $N \longleftrightarrow$ matrix model: double line notation $\begin{array}{c} \longrightarrow \\ \longleftarrow \end{array} = \text{wavy line}$

In $N \rightarrow \infty$ only **planar diagrams** in double line **survive**



Non-perturbative quark propagator

In this approximation the theory is (almost) solvable

The **full, non-perturbative**, quark propagator:

$$\begin{array}{c} p \\ \rightarrow \end{array} \text{---} \text{---} \text{---} \text{---} \begin{array}{c} \text{S} \\ \text{---} \end{array} \begin{array}{c} p \\ \rightarrow \end{array} \text{---} \text{---} \text{---} \text{---} = \begin{array}{c} p \\ \rightarrow \end{array} \text{---} \text{---} \text{---} \text{---} + \begin{array}{c} p \\ \rightarrow \end{array} \text{---} \text{---} \begin{array}{c} \text{1PI} \\ \text{---} \end{array} \begin{array}{c} p \\ \rightarrow \end{array} \text{---} \text{---} \text{---} \text{---} + \begin{array}{c} p \\ \rightarrow \end{array} \text{---} \text{---} \begin{array}{c} \text{1PI} \\ \text{---} \end{array} \begin{array}{c} \text{1PI} \\ \text{---} \end{array} \begin{array}{c} p \\ \rightarrow \end{array} \text{---} \text{---} \text{---} \text{---} + \dots$$

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Is bootstrapped with Dyson Equation:

$$\text{---} \xrightarrow{p} \textcircled{1PI} \text{---} \xrightarrow{p} \text{---} = \text{---} \xrightarrow{p} \text{---} \textcircled{S} \text{---} \xrightarrow{p-k} \text{---} \textcircled{1PI} \text{---} \xrightarrow{p} \text{---}$$

That can be miraculously solved:

$$\text{---} \xrightarrow{p} \textcircled{1PI} \text{---} \xrightarrow{p} \text{---} = \frac{N_c g^2}{\pi} \left(\frac{\text{sgn}(p_-)}{\rho} - \frac{1}{p_-} \right)$$

Non-perturbative quark propagator: no free quarks

The resulting full propagator is:

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$\rho \rightarrow 0$ is a regulator to treat an IR divergence appearing

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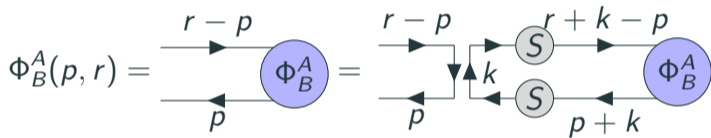
As $\rho \rightarrow 0$, poles are pushed to ∞

\implies **No free quark states in the spectrum!**

The model is indeed confining

Quark-antiquark

Not fully convinced? let's compute the scattering between two quarks.
 Boostroppable though Bethe-Salpeter:



From which we obtain the eigenproblem:

$$\mu^2 \phi_B^A(x) = \left(\frac{\alpha_A}{x} + \frac{\alpha_B}{1-x} \right) \phi_B^A(x) - \mathcal{P} \int_0^1 dy \frac{\phi_B^A(y)}{(y-x)^2}, \quad \phi_B^A(x) = \int dp_+ \Phi_B^A(p, r)$$

$$\alpha_{A,B} = \frac{\pi m_{A,B}^2}{g^2} - 1, \quad 2r_+ r_- = \frac{g^2}{\pi} \mu^2, \quad x = \frac{p_-}{r_-}, \quad x \in [0, 1]$$

Mesons spectrum

$$\mu^2 \phi_B^A(x) = \left(\frac{\alpha_A}{x} + \frac{\alpha_B}{1-x} \right) \phi_B^A(x) - \mathcal{P} \int_0^1 dy \frac{\phi_B^A(y)}{(y-x)^2}$$

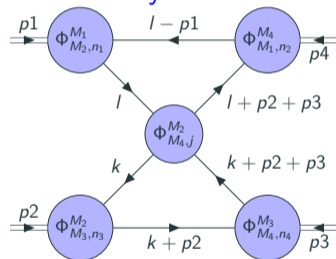
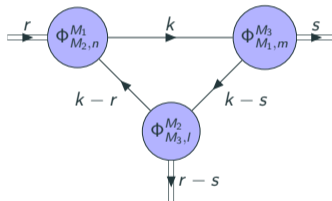
- There is no analytical expression for ϕ
- The spectrum is **discrete** \Rightarrow No free quarks final states
- **Only mesons states** with M^2 determined by μ^2
- Resemble QM of pair of particles in a box with linear potential
$$\mathcal{P} \int_0^1 dy \frac{\phi_B^A(y)}{(y-x)^2} = \mathfrak{F}^{-1} \left[(|p|) \tilde{\phi}(p) \right]$$
- High energy limit \rightarrow neglect linear potential \sim two particle confined in a box $\mu^2 \sim k^2$ (Regge trajectories)

What are we currently doing?

We are studying the scattering of mesons in this model.

Qualitative behaviour: in mesons correlation functions we have **only**:

- Local interactions
- Mesons exchange



The theory is confining \rightarrow No long range interactions mediated by gluons/quarks

Conclusions and outlook

Summarising:

- QCD₂ is a confining gauge theory
- No propagating dof from gluons: they provide linear potential
- No free quarks in the spectrum
- Pair of quarks can only exist in bound mesons states
- Mesons correlation functions respect the confining structure

Outlook:

- String Theory worldsheet construction?
- Beyond numerics for mesons: any analytical structure?
- Higher orders in $1/N$ expansion
- What about Baryons?

Thank you for you attention!

Questions?