## Fusion Plasma Physics and ITER: An Introduction

## 1. Plasma Physics for Magnetic Fusion

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The views and opinions expressed herein do not necessarily reflect those of the ITER Organization.

#### **Overview of Lectures**

- 1. Fusion Plasma Physics in Magnetic Fusion DJ Campbell
- 2. Physics of Tokamak Plasmas DJ Campbell
- 3. Fusion Technology for ITER and the ITER Project G Janeschitz
- 4. Further development towards a Fusion Power Plant G Janeschitz

## Lecture 1 - Synopsis

- Introduction to thermonuclear fusion
  - some aspects of inertial confinement fusion
- Basics of magnetic confinement fusion the tokamak
- (Some) plasma physics for magnetic confinement fusion in tokamaks
- The ITER Project

## **An Introduction - ITER**

- 7 partners representing >50% of the world's population have embarked on the ITER project
- ITER is designed to produce 500MW of fusion power (tenfold power amplification) for extended periods of time (several 100s)
- 10 years construction
   20 years operation
   5 years de-activation
- These lectures will explain the background to ITER and how it fits into the development of fusion energy

#### **ITER tokamak**



## **ITER: Fusion Power Production**



## **Introduction to Thermonuclear Fusion**

## Why Fusion ?



- Fuel: abundant, world-wide distributed:
  - sufficient deuterium in seawater for millions of years
  - tritium is produced from lithium
     sufficient ore supplies for thousands of years (millions of years including seawater resources

#### Safety: no risk of major accidents:

- reactor contains fuel for only a few minutes burn
- Waste: no long-term burden:
  - low radio-toxicity after < 100 years</li>
  - no CO<sub>2</sub>

#### Fusion – the fundamental principle

 Energy gain from fusion, like fission, is based on Einstein's equation:

#### $E = \Delta mc^2$

- mass loss for DT reactions corresponds to ~ 0.4%
- As illustrated, energy gain per unit mass is greater for fusion
  - energy gain/ reaction:

DT fusion: 17.6 MeV U fission: ~200 MeV



## **Essential Fusion Reactions**



+ 20% of Energy (3.5 MeV)

+ 80% of Energy (14.1 MeV)

• The D-T fusion reaction is the simplest to achieve under terrestrial conditions

 $^{2}D + ^{3}T \Rightarrow ^{4}He (3.5 \text{ MeV}) + ^{1}n (14.1 \text{ MeV})$ 

Two other important reactions for DT fusion:

 $^{1}n + ^{6}Li \Rightarrow ^{4}He + ^{3}T + 4.8 MeV$ 

 $^{1}n + ^{7}Li \Rightarrow ^{3}He + ^{3}T + ^{1}n - 2.5 MeV$ 

these reactions will allow a fusion reactor to breed tritium



#### **Fusion Power Production**

- High temperatures (~10 keV) are required for significant thermonuclear fusion energy production (⇒ dealing with plasmas!):
  - nuclei must overcome Coulomb barrier to approach close enough to fuse

 $\Rightarrow$  T ~ 300 keV required to overcome Coulomb barrier of ~0.4 MeV

 but, reaction rate is dominated by a small population of high energy nuclei which react by quantum tunnelling

 $\Rightarrow$  for 100 keV nuclei, tunnelling probability, w = 3.2 × 10<sup>-2</sup>

• Fusion power density for an optimal 50:50 D-T mixture:

$$P_{F} = \frac{n^{2}}{4} \left\langle \sigma v \right\rangle E_{F} = \frac{p^{2}}{16} \frac{\left\langle \sigma v \right\rangle}{\left( kT \right)^{2}} E_{F} \quad (E_{F} = 17.6 MeV)$$

- e.g. for magnetic fusion reactor parameters, with p ~ 10 atm,  $P_F \sim 7.5 \text{ MWm}^{-3}$ 

## **Power Gain – the Lawson Criterion**

- Obtaining the temperature required to produce fusion reactions involves heating the plasma
  - for a net power gain, fusion power out must exceed input heating power (including correction for loss processes, e.g. bremsstrahlung, synchrotron radiation ....)
- One can define a parameter,  $\tau_E$ , the energy confinement time, which characterizes the rate of power loss:

$$\tau_E = \frac{W_{th}}{P_{loss}} = 3 \frac{\int nkT \, dV}{P_{loss}}$$

• Then, the overall power balance can be written (P<sub>heat</sub> = ext power):

$$P_{heat} = \left(\frac{3\overline{n}k\overline{T}}{\tau_{E}} - \frac{n^{2}}{4}\langle\sigma v\rangle E_{\alpha}\right) V$$
Ignition (P<sub>heat</sub> = 0) implies:  $\overline{n}\tau_{E} > \frac{12}{\langle\sigma v\rangle}\frac{k\overline{T}}{E_{\alpha}} \Rightarrow Q_{DT} = \frac{P_{F}}{P_{heat}} \to \infty$ 

#### **Power Gain – the Lawson Criterion**



#### **Plasma Fusion Performance**

<i>Temperature (T<sub>i</sub>):</i>	1-2 × 10 <sup>8</sup> K (~10 × temperature of sun's core)		
Density (n <sub>i</sub> ):	1 × 10 <sup>20</sup> m <sup>-3</sup> (~10 <sup>-6</sup> of atmospheric particle density)		
<b>Energy confinement time (<math>\tau_E</math>):</b> few seconds ( $\propto$ current $\times$ radius <sup>2</sup> ) (plasma pulse duration ~1000s)			
Fusion power amplification	$Q = \frac{Fusion Power}{Input Power} \sim n_i T_i \tau_E$		
⇒ Present devices: Q ≤ 1			
<i>⇒ITER:</i> Q ≥ 10			
$\Rightarrow$ "Controlled ignition": Q $\geq$ 30			

#### Plasma Fusion Performance – Tokamaks Fusion Triple Product

- Existing experiments have achieved nTτ values
   ~ 1×10<sup>21</sup> m<sup>-3</sup>skeV
  - ~ Q<sub>DT</sub> = 1
- JET and TFTR have produced DT fusion powers of >10MW for ~1s
- ITER is designed to a scale which should yield
   Q<sub>DT</sub> ≥ 10 at a fusion power of 400 - 500MW for 300-500s



#### **Inertial Confinement Fusion**

- The US National Ignition Facility at Livermore is expected to achieve ignition within the next 2 years:
  - uses high power lasers to compress a small DT capsule embedded in a "hohlraum" (indirect drive)



WJ Hogan, Ch. 8 in Landolt-Börnstein-Handbook on Energy Technologies, Springer Verlag (2005)

#### **Inertial Confinement Fusion**

#### • Concept of ignition in MCF and ICF is somewhat different:

- in MCF, ignition criterion is based on power balance
- in ICF, ignition criterion is based on burn propagation in capsule and fuel burnup criterion

on:

 $nT\tau_{E} \sim 5 \times 10^{21} (m^{-3} keV.s)$ 

(fuel burn-up fraction ~ few %)

#### • ICF criterion:

 $\rho R \sim 2 (kg.m^{-2})$ 

(fuel burn-up fraction  $\sim 30\%$ )

	ITER	NIF
n <sub>i</sub> (m <sup>-3</sup> )	1×10 <sup>20</sup>	1.1×10 <sup>31</sup>
ρ (kgm <sup>-3</sup> )	4.2×10 <sup>-7</sup>	5.7×10 <sup>4</sup>
<t> (keV)</t>	~ 10	~ 10
(atm)	3.3	4.5×10 <sup>11</sup>
т <sub>Е</sub> (S)	~ 3.5	~ 10 <sup>-10</sup>
a (m)	2.0	3.5×10⁻⁵
V (m <sup>3</sup> )	830	1.8×10 <sup>-13</sup>
E <sub>plas</sub> (J)	3.5×10 <sup>8</sup>	9.3×10 <sup>3</sup>
Output	500 MW	10-20 MJ

# Basics of Magnetic Confinement Fusion: The Tokamak

#### Plasma Toroidal Magnetic Confinement

 Magnetic fields cause ions and electrons to spiral around the field lines:

$$F = q(E + v \times B)$$

 in a toroidal configuration plasma particles are lost to the vessel walls by relatively slow diffusion across the field lines



A special version of this torus is called a tokamak:

'toroidal chamber' and 'magnetic coil' (Russian)

## Magnetic Confinement in a Tokamak

#### The Tokamak:

#### External coils

- to produce a toroidal magnetic field

#### Transformer with primary winding

- to produce a toroidal current in the plasma
- this plasma current creates a poloidal magnetic field

#### • Finally, poloidal coils

- to control the position and shape of the plasma



## **JET: Joint European Torus**

- JET is currently the largest tokamak
  - Major/ minor radius: 3 m/ 1 m
  - Plasma volume ~100 m<sup>3</sup>
  - Toroidal field: 3.4 T
  - Plasma Current: 7 MA
- In DT experiments in 1997, a peak fusion power of 16 MW was produced (Q<sub>DT</sub> ~ 0.6)



#### **JET - Internal**



#### Magnetic Confinement in a Tokamak



- In configurations with only a toroidal field, ions and electrons drift vertically in opposite directions:
  - caused by field gradient and curvature
  - resultant electric field destroys plasma
- An additional poloidal field allows particles to follow helical paths, cancelling the drifts
- "Winding number" of helix is an important stability parameter for the system:

$$q_{c} = \frac{aB_{\phi}}{RB_{\theta}}$$

- q = safety factor
- R/a = aspect ratio

#### **Toroidal Magnetic Confinement Systems**



- Numerous toroidal confinement configurations are being studied:
  - Tokamak has progressed most rapidly and is ready for the "thermonuclear" step
- Note that in the stellarator, the helical magnetic surfaces are produced entirely by external coils:

#### **A Digression: Stellarators**





Wendelstein 7-X (Germany) Modular Stellarator Large Helical Device (Japan) Heliotron

- Operation without a plasma current has some advantages (eg steady-state operation), but coil configuration is more complicated
  - LHD is already in operation, while W7-X will enter operation in middle of decade
  - overall, stellarator energy confinement is similar to that in "equivalent current" tokamaks

## Plasma Equilibrium in a Tokamak



- Plasma is force-free, ie "in equilibrium":
  - implies both internal and external force balance
  - ignoring internal flows and electric fields, force balance equation takes form:

$$\mathbf{j} \times \mathbf{B} = \nabla p$$

- It follows that

 $\mathbf{B} \bullet \nabla p = \mathbf{j} \bullet \nabla p = 0$ 

- ⇒ pressure is constant on magnetic surfaces
- ⇒ current lines lie within magnetic surfaces
- Can define poloidal magnetic flux function,  $\Psi$ , satisfying,

 $\mathbf{B} \bullet \nabla \Psi = \mathbf{0}$ 

## Plasma Equilibrium in a Tokamak



• Formal definition of safety factor:

$$q = \frac{d\Phi}{d\Psi} \xleftarrow{\text{toroidal flux}} poloidal flux}$$

- absolute value of *q* and its variation across the plasma radius are important in plasma stability
- define magnetic shear as:

$$S = \frac{r}{q} \frac{dq}{dr}$$

 by elongating the plasma, more current can be squeezed into the plasma ring at fixed q:

$$\kappa = \frac{b}{a}$$

 Typically the pressure (temperature, density) and current profiles are peaked on the plasma axis:

- the profile of q is then the inverse, with  $q(0) \sim 1$ 

## Plasma Equilibrium in a Tokamak

 Since there internal force balance between the plasma pressure and the magnetic field, it is conventional to work with a normalized pressure, poloidal beta:

$$\beta_{p} = \frac{2\mu_{0} \int p.dV}{\left\langle B_{\theta}(a) \right\rangle_{line}^{2}.V}$$

- when  $\beta_p < 1$ , plasma is paramagnetic when  $\beta > 1$  plasma is diamagnetic
- when  $\beta_p > 1$ , plasma is diamagnetic
- equilibrium condition limits  $\beta_p$  to approximately  $\beta_p < R/a$
- The plasma beta, i.e. pressure normalized to the toroidal field, is an important measure of plasma stability and of efficient use of field:

$$\beta(\%) = 100 \times \frac{2\mu_0 \int p.dV}{B_{\phi}^2(0).V}$$

• The plasma internal inductance characterizes how peaked the current profile is and is also a significant factor in plasma stability:

$$\ell_{i} = \frac{\left\langle B_{\theta}^{2} \right\rangle_{vol}}{\left\langle B_{\theta}^{2}(a) \right\rangle_{flux}} = \frac{4U_{p}}{\mu_{0}R_{0}l_{p}^{2}} \quad \text{stored poloidal field energy inside plasma}$$

# (Some) plasma physics for magnetic confinement fusion in tokamaks

#### **Some Basic Plasma Physics**

 Plasmas studied in fusion research are essentially, quasi-neutral, but there is a characteristic scale length for shielding of the potential due to individual charges, the Debye length:

$$\lambda_D = \left(\frac{\varepsilon_0 T}{ne^2}\right)^{1/2} = 2.35 \times 10^5 \left(\frac{T}{n}\right)^{1/2} (T \text{ in keV})$$

- the assumption of quasi-neutrality is satisfied if  $n\lambda_D^3 >> 1$  (~10<sup>8</sup> in tokamak)
- Characteristic plasma frequency:

$$\omega_{p,e} = \left(\frac{ne^2}{\varepsilon_0 m_e}\right)^{1/2} = 56.4n^{1/2} \text{ s}^{-1}$$

• Motion of charged particles in confining magnetic fields can be characterized as a gyro-motion around of Larmor radius,  $\rho_L$ , around a guiding centre:

$$\rho_{Lj} = \sqrt{2} \frac{m_j v_{Tj}}{\left| e_j \right| B}$$

- $\thicksim$  100 $\mu m$  for electrons at 10keV and 3T
- ~ 5mm for protons at 10keV and 3T

#### **Particle Orbits in the Tokamak**

JA Wesson, *Tokamaks*, 3<sup>rd</sup> edition, OUP (2004)



- Gradients and curvature in the magnetic field lead to modifications in the particle trajectories:
  - "Passing" particles orbit shift:  $\delta r_p \sim \varepsilon . \rho_{L\theta} = q . \rho_{L\phi}$  ( $\varepsilon = a / R$ )
  - "Trapped" particles "banana" width:  $\Delta r_t \sim \varepsilon^{0.5} \rho_{L\theta}$
  - Guiding centre orbit of trapped particles bounce back and forth on outer half of torus due to magnetic mirror formed by toroidal field ( $B_{\phi} = R_0 B_0/R$ )
  - At low collision frequencies, a fraction of particles are trapped:  $f = \sqrt{2r} / (R_0 + r)$

#### **Plasma Resistivity**

 How does one calculate the current achievable in a tokamak plasma (ignoring stability considerations)?

- Ohm's law for magnetized plasma:  

$$E + \mathbf{v} \times \mathbf{B} = \eta \mathbf{j}$$
inductive electric field  
fluid velocity magnetic field  
- average ionic charge number, "effective charge":  $Z_{eff} = \frac{\sum_{i=1}^{n} n_{i} Z_{i}^{2}}{n_{e}}$ 

Coulomb logarithm characterizes average over electron-ion interactions:

$$\ln \Lambda_{ei} = 15.2 - \frac{1}{2} \ln \left( \frac{n_e}{10^{20}} \right) + \ln (T_e) \quad (T_e \text{ in keV})$$

- "Classical" (or Spitzer) parallel resistivity:

$$\eta_{par} = 1.65 \times 10^{-9} f(Z_{eff}) Z_{eff} \frac{\ln \Lambda_{ei}}{T_e^{1.5}} \Omega m \quad (T_e \text{ in keV}) \quad (f(Z_{eff}) \sim 1)$$

 $-\,$  e.g.  $T_{e}$  = 1keV,  $\eta_{par}$  ~ 2 ×10^{-8}  $\Omega m$  – room temperature copper

#### "Neoclassical" Plasma Resistivity

 "Trapping" of particles in toroidal magnetic mirrors leads to an enhancement of the plasma resistivity:

$$\eta_{neo} \approx \frac{\eta_{par}}{(1 - \epsilon^{0.5})^2} f(v^*, \epsilon, Z_{eff})$$

- this effect first became detectable in the hot "collisionless" plasmas characteristic of JET scale devices
- Comparisons between resistivity profiles calculated from T<sub>e</sub> measurements and from resistive diffusion analysis of plasma current showed better agreement with the neoclassical resistivity





## "Bootstrap" Current

M Kikuchi, M Azumi, Plasma Phys Control Fusion 37 1215(1995)



- The "bootstrap" current is a further "neoclassical" consequence of the presence of trapped particles:
  - momentum exchange between trapped and passing particles, together with density and temperature gradients, lead to an additional component of current:

- locally: 
$$j_{bs}(\varepsilon \rightarrow 1) = -\frac{1}{B_{\theta}}\frac{dp}{dr}$$
 globally:  $I_{bs} = C\varepsilon^{1/2}\beta_{\rho}I_{\rho}$  (C ~ 1/3 – 2/3)

#### **Overview of a Tokamak Plasma Pulse**

#### Simulated ITER plasma pulse



#### **Tokamak Plasma Pulse – Flux Consumption**

• Flux consumption during an ITER plasma pulse:

$$\Psi_{\textit{tot}} = \Psi_{\textit{bd}} + \Psi_{\textit{ramp}} + \Psi_{\textit{ind}} + \Psi_{\textit{res}}$$

 $\Psi_{bd}$  = breakdown loss ~ 5 - 10 Wb

 $\Psi_{ramp} = C_E \mu_0 R_0 I_p \sim 25 \text{ Wb} (C_E \sim 0.4 - 0.5 \text{ is an empirical coefficient})$ 

$$\begin{split} \Psi_{ind} &= L_p I_p \sim 180 \text{ Wb} \\ \left[ L_p &= \mu_0 R_0 \left( \ln \frac{8R_0}{a} + \frac{\ell_i}{2} - 2 \right) \approx 2R_0 \ (\mu \text{H}) \right] \\ \Psi_{res} &= \text{resistive loss} = V_1 I_p \sim 30 - 40 \text{ Wb} \end{split}$$
$$\Psi_{tot} \sim 240 - 260 \text{ Wb}$$

- During the current flat-top at 15 MA, the single turn loop voltage, V<sub>I</sub> < 100 mV, due to:</li>
  - the high plasma temperature ( $T_e(0) \sim 25 \text{ keV}$ )
  - a bootstrap current contribution of ~10%
  - external "non-inductive" current drive of ~10%

## **The ITER Project**

## What is ITER?

ITER is a major international collaboration in fusion energy research involving the EU (plus Switzerland), China, India, Japan, the Russian Federation, South Korea and the United States

- The overall programmatic objective:
  - to demonstrate the scientific and technological feasibility of <u>fusion</u> <u>energy</u> for peaceful purposes
- The principal goal:
  - to design, construct and operate a <u>tokamak experiment</u> at a scale which satisfies this objective
- ITER is designed to confine a <u>Deuterium-Tritium plasma</u> in which <u> $\alpha$ -particle heating</u> dominates all other forms of plasma heating:

#### ⇒ a burning plasma experiment

## **ITER Scope - Mission Goals**

#### Physics:

- ITER is designed to produce a plasma dominated by  $\alpha\mbox{-particle}$  heating
- produce a significant fusion power amplification factor (Q ≥ 10) in long-pulse operation
- aim to achieve steady-state operation of a tokamak (Q = 5)
- retain the possibility of exploring 'controlled ignition' ( $Q \ge 30$ )

#### Technology:

- demonstrate integrated operation of technologies for a fusion power plant
- test components required for a fusion power plant
- test concepts for a tritium breeding module

#### **ITER - Major Components**



## **The ITER Project - Current Status**

- Spring 2006: ITER Joint Work Site established in Cadarache design teams arrive from Naka and Garching
- November 2006: ITER Agreement signed in Paris
- Late 2006: Design Review begins
- Early 2007: Construction activities launched
- October 2007: ITER Organization formally established
- July 2010: ITER Baseline (scope, schedule, cost) approved by ITER Council
- July 2010: New Director-General, Osamu Motojima appointed by ITER Council
- August 2010: Building construction begins on-site

## ITER Overall Project Cost (OPC)

- The total cost for the Construction Phase approved in July 2010 is 4584.7 kIUA
- Table shows total cost over lifetime of project:

Construction Phase	4584.7 kIUA
<b>Operation Phase</b>	188 kIUA per year
<b>Deactivation Phase</b>	EUR 281 Million
Decommissioning Phase	EUR 530 Million

#### NB: 1 kIUA = 1M \$US (1989) = 1.5M Euro (2010)

#### **ITER Construction Schedule**



### The ITER Project Team - Domestic Agencies



 90% of ITER components will be supplied "in-kind" by the Members through their Domestic Agencies

#### **Itinerary of ITER Components**



#### **ITER Construction at Cadarache**

#### ITER Site platform levelling complete and construction underway



#### **ITER Site after Construction**



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http://www.iter.org - and associated links