

Hints of New Physics and EFT: Flavor

Anders Eller Thomsen

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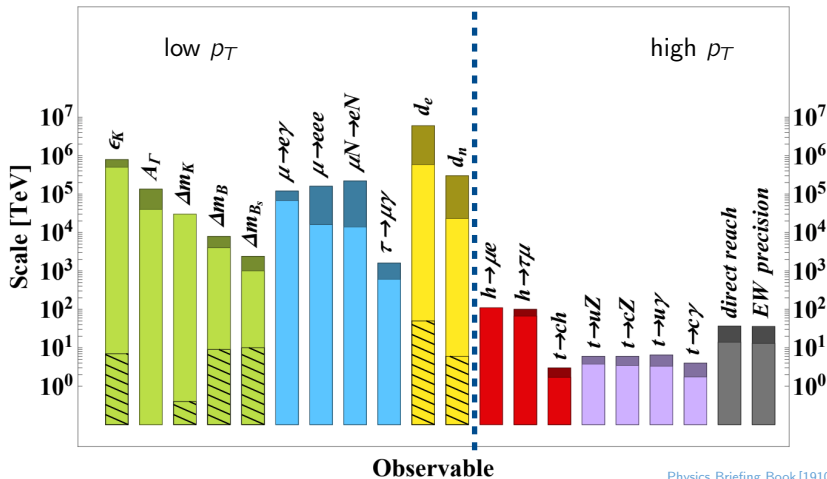
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ALBERT EINSTEIN CENTER
FOR FUNDAMENTAL PHYSICS

Based on work with A. Greljo and A. Palavrić

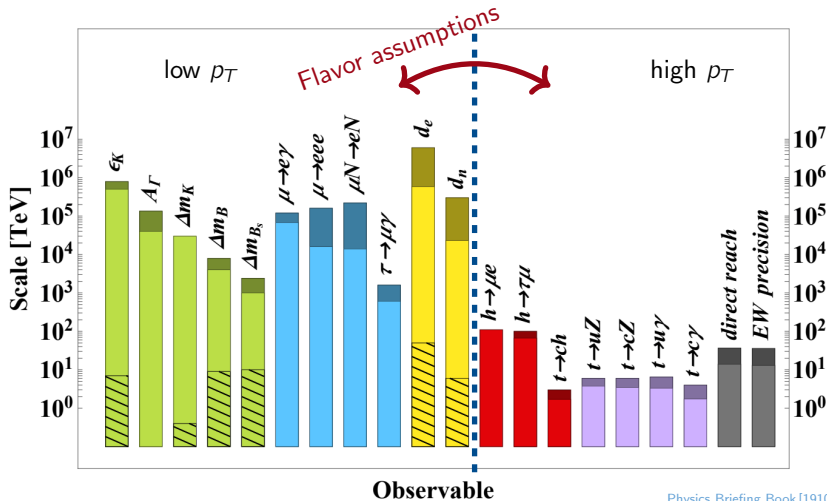
Workgroup 1 meeting, 1 July 2022

Probing high-scale new physics



Physics Briefing Book [1910.11775]

Probing high-scale new physics



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Flavor assumptions in EFT fits discussed at [LHC EFT topical meeting on flavor assumptions](#) (January, 2022)

The Standard Model

Symmetries:

$$SU(3)_c \times SU(2)_L \times U(1)_Y \times \text{Poincaré}$$

Matter fields:

$$q_i, u_i, d_i, \ell_i, e_i, H$$

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The Yukawa couplings break $G_F \rightarrow U(1)_B \times U(1)_{L_e} \times U(1)_{L_\mu} \times U(1)_{L_\tau}$:

$$\mathcal{L}_{\text{yuk}} = -y_u \bar{q} \tilde{H} u - y_d \bar{q} H d - y_e \bar{\ell} H e + \text{H.c.}$$

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But G_F identifies equivalent theories

$$\{y_u, y_d, y_e\} \sim \{U_q y_u U_u^\dagger, U_q y_d U_d^\dagger, U_\ell y_e U_e^\dagger\}, \quad \text{for } U \in G_F$$

54 parameters \longrightarrow 13 *physical* parameters

The flavor puzzle

Flavor of the SM:

$$y_u \sim \begin{pmatrix} \text{light blue} & & \\ & \text{medium blue} & \\ & & \text{dark blue} \end{pmatrix} \quad V_{\text{CKM}} \sim \begin{pmatrix} \text{dark blue} & \text{light blue} & \\ \text{light blue} & \text{dark blue} & \\ & & \text{dark blue} \end{pmatrix}$$

Not visible in colliders

$$y_{d,e} \sim \begin{pmatrix} \text{light blue} & & \\ & \text{medium blue} & \\ & & \text{medium blue} \end{pmatrix} \quad V_{\text{PMNS}} \sim \begin{pmatrix} \text{dark blue} & \text{medium blue} & \text{light blue} \\ \text{medium blue} & \text{medium blue} & \text{medium blue} \\ \text{medium blue} & \text{medium blue} & \text{medium blue} \end{pmatrix}$$

The flavor puzzle

Flavor of the SM:

$$y_u \sim \begin{pmatrix} \square & & \\ & \square & \\ & & \blacksquare \end{pmatrix} \quad V_{\text{CKM}} \sim \begin{pmatrix} \blacksquare & \square & \\ \square & \blacksquare & \square \\ & \square & \square & \blacksquare \end{pmatrix}$$

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- Is the structure in the flavor sector meaningful?
- How does potential new physics couple to flavor?
- What is (if any) the flavor symmetry of the SM?

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- How does potential new physics couple to flavor?
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y_t is the leading (only non-perturbative) breaking of G_F in the SM:

$$y_u \sim \begin{pmatrix} & & \\ & & \\ & & \text{dark blue} \end{pmatrix} : \quad G_F \rightarrow U(2)_q \times U(2)_u \times U(3)_d \times U(3)_\ell \times U(3)_e \times U(1)_B$$

Flavor in the SMEFT

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_O \frac{C_O}{\Lambda^{\dim O - 4}} O$$

Constructed from the same fields and symmetries as \mathcal{L}_{SM}

dim $O = 5$: 1 classes, 12 operators

dim $O = 6$: 59 classes, 2499 operators ($\Delta B = 0$)

e.g. $C_{uu}^{psrt} (\bar{u}_p \gamma_\mu u_r) (\bar{u}_s \gamma^\mu u_t)$
45 parameters

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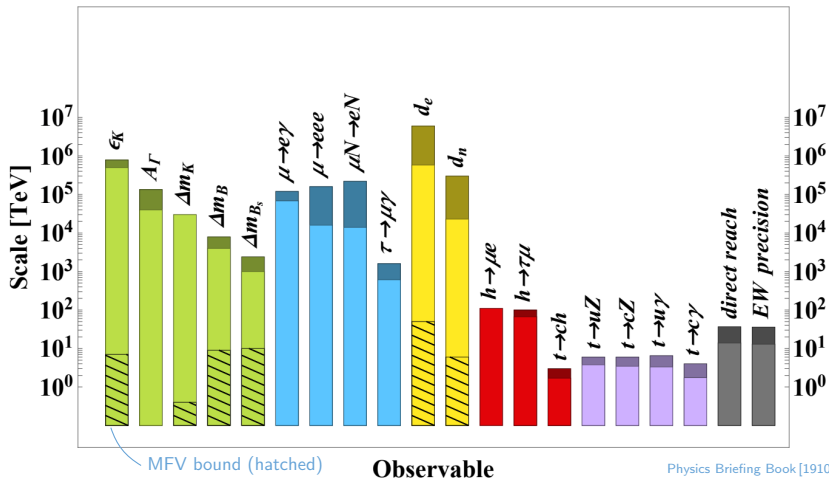
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Chart the space of SMEFT with flavor symmetries

Greljo, Palavrić, AET [2203.09561]
and also Faroughy *et al.* [2005.05366]

- Organization in perturbative spurion expansion, reduces the number of (relevant) operators
- Reduction in number of parameters makes global SMEFT fits possible
- Differentiation of SMEFT in universality classes pointing to different NP

NP with MFV



NP with MFV assumption can reside at the TeV-scale, but it is very (too?) restrictive

D'Ambrosio *et al.* [hep-ph/0207036]

Seasoning the SMEFT

SMEFT $\mathcal{O}(1)$ terms (dim-6, $\Delta B = 0$)		Lepton sector							
		MFV _L	U(3) _V	U(2) ² × U(1) ²	U(2) ²	U(2) _V	U(1) ⁶	U(1) ³	No symm.
Quark sector	MFV _Q								
	U(2) ² × U(3) _d								
	U(2) ³ × U(1) _{d₃}								
	U(2) ³								
	No symmetry								

Greljo, Palavric, AET [2203.09561]

- 28 hypothesis for the flavor symmetries
- Systematic chart from MFV to anarchy

$$U(3) \supset U(2) \times U(1) \supset U(2) \supset U(1)$$

- Spurions (minimal set) introduced to reproduce SM masses and mixings

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Greljo, Palavric, AET [2203.09561]

- Work from the Warsaw basis
- Construct the operators order by order in the spurions
- Generally:
 - Flavor-symmetric operators are relevant for EW/top/Higgs physics
 - Spurions are relevant for flavor physics (CKM is a result of spurions)

But also examples of non-trivial interplay [Bruggisser et al. \[2101.07273\]](#)

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Quark sector	MFV _Q	41	6	45	9	59	6	62	9	67	13	81	6	93	18	207	132
	U(2) ² × U(3) _d	72	10	78	15	95	10	100	15	107	21	122	10	140	28	281	169
	U(2) ³ × U(1) _{d₃}	86	10	92	15	111	10	116	12	123	21	140	10	158	28	305	175
	U(2) ³	93	17	100	23	118	17	124	23	132	30	147	17	168	38	321	191
	No symmetry	703	570	734	600	756	591	786	621	818	652	813	612	906	705	1350	1149

Greljo, Palavric, AET [2203.09561]

The $\mathcal{O}(1)$ terms in the spurion expansion are expected to be the relevant ones for EW/top/Higgs fits.

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The $\mathcal{O}(1)$ terms in the spurion expansion are expected to be the relevant ones for EW/top/Higgs fits.

- The symmetry of the SM broken by y_t (also in suggestion by LHC EFT WG)
- Discriminates third family. Good description of e.g. b anomalies
- Allows generic LFUV. No cLFV

Example: SMEFT with $U(2)^3 \times U(1)_{d_3}$ symmetry

Flavor symmetry $G = U(2)^3 \times U(1)_{d_3}$ with quark fields

$$q = \begin{bmatrix} q^a \sim (\mathbf{2}, \mathbf{1}, \mathbf{1})_0 \\ q_3 \sim (\mathbf{1}, \mathbf{1}, \mathbf{1})_0 \end{bmatrix}, \quad u = \begin{bmatrix} u^a \sim (\mathbf{1}, \mathbf{2}, \mathbf{1})_0 \\ u_3 \sim (\mathbf{1}, \mathbf{1}, \mathbf{1})_0 \end{bmatrix}, \quad d = \begin{bmatrix} d^a \sim (\mathbf{1}, \mathbf{1}, \mathbf{2})_0 \\ d_3 \sim (\mathbf{1}, \mathbf{1}, \mathbf{1})_1 \end{bmatrix}.$$

The spurions break G completely, and contains 9 physical parameters (incl. 1 phase)

$$V_q \sim (\mathbf{2}, \mathbf{1}, \mathbf{1})_0, \quad \Delta_u \sim (\mathbf{2}, \bar{\mathbf{2}}, \mathbf{1})_0, \quad \Delta_d \sim (\mathbf{2}, \mathbf{1}, \bar{\mathbf{2}})_0, \quad X_b \sim (\mathbf{1}, \mathbf{1}, \mathbf{1})_{-1}.$$

*Similar analyses available for all the symmetries

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$(\bar{q}q)(\bar{q}q)$

Examples from operator basis

$$\begin{aligned} \mathcal{O}(1) &: (\bar{q}_a q^b)(\bar{q}_b q^a), \quad (\bar{q}_a q_3)(\bar{q}_3 q^a), \quad \mathcal{O}(V) : (\bar{q}_a q_3)(\bar{q} V_q q^a), \quad \text{H.c.}, \\ \mathcal{O}(V^2) &: (\bar{q}_a V_q^\dagger q)(\bar{q} V_q q^a). \end{aligned} \quad (2.33)$$

$(\bar{u}u)(\bar{u}u)$

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$(\bar{d}d)(\bar{d}d)$


$$\mathcal{O}(1) : (\bar{d}_a d^b)(\bar{d}_b d^a), \quad (\bar{d}_a d_3)(\bar{d}_3 d^a), \quad \mathcal{O}(\Delta V X) : (\bar{d}_a X_b d_3)(\bar{d} \Delta_d^\dagger V_q d^a), \quad \text{H.c.} \quad (2.35)$$

*Similar analyses available for all the symmetries

The SMEFTflavor package

In[14]:= `CountingTable["lep:U2diag"]`

Out[14]=

- SMEFTflavor is a Mathematica package to derive operator bases for flavor symmetries
- It is possible to implement custom flavor scenarios
- Example: $U(2)_V$ lepton symmetry with $\Delta \sim \mathbf{3}$ spurion 

lep:U2diag		$O[1]$		$O[\Delta\mathbf{1}]$	
$\psi^2 H^3$	OeH	2	2	1	1
$\psi^2 XH$	Oe(B,W)	4	4	2	2
$\psi^2 H^2 D$	OHL(1,3)	4		2	
	OHe	2		1	
(LL)(LL)	Oll	5		3	
(RR)(RR)	Oee	3		2	
(LL)(RR)	Ole	6	1	5	2
Total		26	7	16	5

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Explicit operator bases, e.g. $O_{le} = (\bar{l}\gamma_\mu l)(\bar{e}\gamma_\mu e)$:

In[13]:= `OperatorBasis["lep:U2diag"]["Ole", Spur["Δl"]] // TableForm // OpForm`

Out[13]/OpForm=

`(Operator[{l12a, l12b, e12c}, Spur[Δld], CGs[Tcda]] + H.c.)`

`Operator[{l12a, l12b, e12c}, Spur[Δld], CGs[Tbda]]`

`Operator[{l12a, l12b, e3, e3}, Spur[Δlc], CGs[Tbca]]`

`(Operator[{l3, l12a, e12b, e3}, Spur[Δlc], CGs[Tacb]] + H.c.)`

`Operator[{l3, l3, e12a, e12b}, Spur[Δlc], CGs[Tbca]]`

Removing redundancies: $U(2)^3$

Problem: Generic Yukawas seem to contain many parameters in, e.g. $U(2)^3$:

$$Y_{u,d} = \begin{bmatrix} a_1^{u,d} \Delta_{u,d} + a_2^{u,d} \Delta_u \Delta_u^\dagger \Delta_{u,d} + \dots & b_1^{u,d} V_q + b_2^{u,d} \Delta_u \Delta_u^\dagger V_q + \dots \\ c_1^{u,d} V_q^\dagger \Delta_{u,d} + \dots & d_1^{u,d} + d_2^{u,d} V_q^\dagger V_q + \dots \end{bmatrix}$$

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Solution: Additional freedom in the $G_F/G_{U(2)^3}$ flavor space:

$$U_q = \exp \begin{bmatrix} 0 & \lambda_1^q V_q + \lambda_2^q \Delta_u \Delta_u^\dagger V_q + \dots \\ -(\lambda_1^q)^* V_q^\dagger - (\lambda_2^q)^* V_q^\dagger \Delta_u \Delta_u^\dagger - \dots & 0 \end{bmatrix}$$
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Going to e.g. down-aligned Yukawas (10 parameters)

$$Y_{u,d} \xrightarrow{U_{q,u,d}} Y'_u = \begin{bmatrix} \Delta'_u & V'_q \\ 0 & y'_t \end{bmatrix}, \quad Y'_d = \begin{bmatrix} \Delta'_d & 0 \\ 0 & y'_b \end{bmatrix}$$

Redefined spurions

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Redefined spurions

No redundancy in the SMEFT basis once minimal Yukawas are chosen

- Flavor structure of NP will leave imprints in the SMEFT
- TeV scale NP must possess flavor structure to remain viable
- Operator bases made available for 28 flavor scenarios: ready-for-use in phenomenological studies
- Operator bases for other scenarios can be determined with

<https://github.com/aethomsen/SMEFTflavor>

Greljo, Palavric, AET [2203.09561]