



Outline

Recent BaBar results on searches for low mass New Physics

- Physics motivation and experimental frame
- **Search for B Mesogenesis:**

 $B^0 \rightarrow \Lambda$ ψ_D Phys. Rev. D 107, 092001 – Published May 3 2023 **Dark Baryon Preliminary:** First presentation at Moriond EW 2023

- **Recently published other searches for New Physics**
 - Search for an **Axion Like Particle (ALP)**:

PhysRevLett.128.131802 (2022)

Search for **Darkonium** Υ_D : PhysRevLett.128.021802 (2022)

In B decays

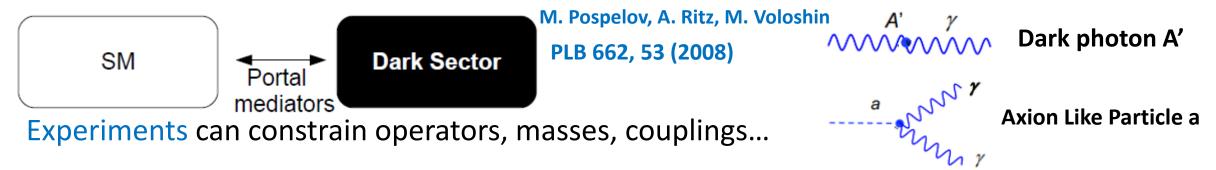
$$e^+e^- \rightarrow \gamma_{ISR} \Upsilon_D$$

Physics motivation

Understand what Dark Matter (DM) is made of and its mass scale

Plethora of models!

Effective theory approach provides different « portals » to access the dark sector



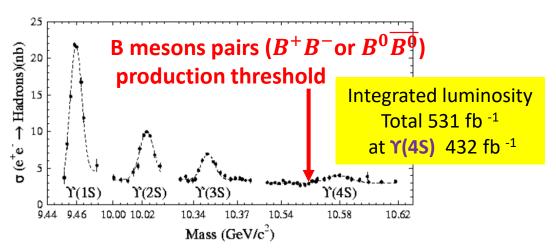
B Factories : great window to search for low mass New Physics produced through mixing between portals mediators and SM particles in :

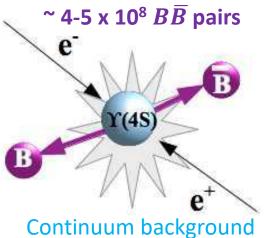
- B mesons decays
- Dark matter produced in e⁺ e⁻ interactions

B mesogenesis: recent approach to explain both

Dark Matter and Baryon Asymmetry in the Universe

Production mechanisms at BaBar — PEP II

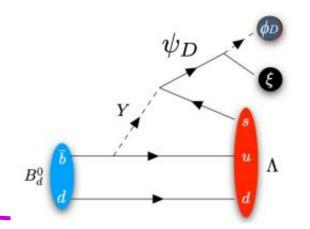




for B decays analyses

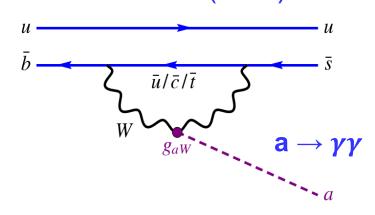
Process	σ (nb)
bb	1.1
cc	1.3
Light quark qq	~2.1
$ au^+ au^-$	0.9
e⁺e⁻	~40

B Mesogenesis dark anti-baryon $\psi_{\scriptscriptstyle extsf{D}}$

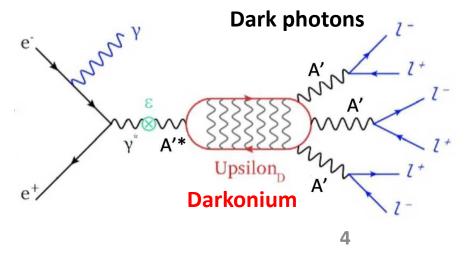


Axion Like Particle

 $B \rightarrow K + a (ALP)$



Darkonium



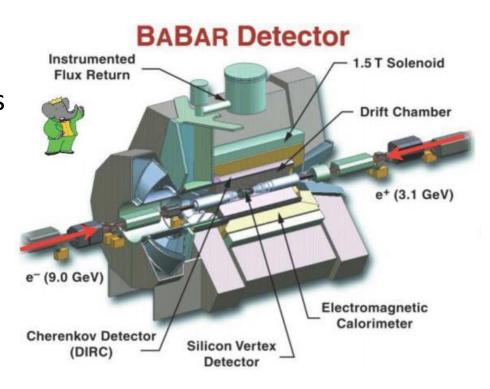
Detector performances for DM searches

B factories very well suited to dark sector searches:

- Clean e⁺e⁻ environment with quasi hermetic detector coverage; good missing energy (potential dark matter) reconstruction
- Reconstruct displaced vertices from long-lived particles
- Precise Particle Identification and reconstruction
- High statistics "precision frontier" data samples

Long history: more than 10 years of DM searches even if different from initial (CP violation) goal

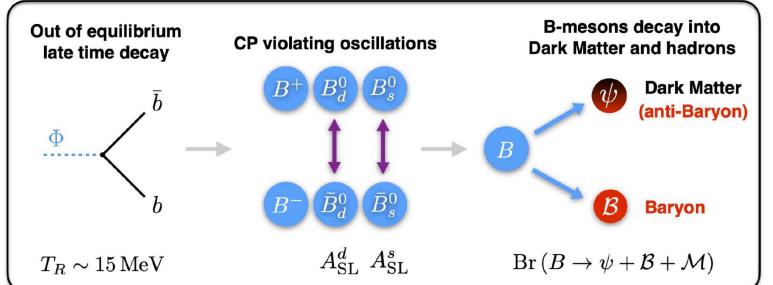
Analyses optimized and validated on <10% of data set, discarded before search performed



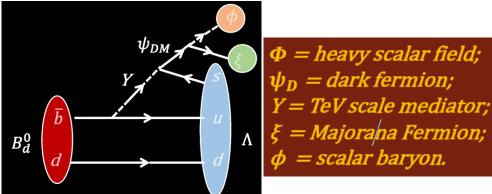
Search for B Mesogenesis (1)

Recently proposed mechanism to explain dark matter abundance and

Baryon Asymmetry of the Universe (BAU)



- G. Elor, M. Escudero and A. E. Nelson, Phys. Rev. D 99, 035031 (2019).
- G. Alonso-Alvarez, G. Elor and M. Escudero, Phys. Rev. D 104, 035028 (2021).



- BAU from B meson decays into a SM baryon (Λ , p, Λ_c ...) and a dark sector anti-
- **baryon** ψ_D + possible **SM** light mesons **M** (not in our analyses channels)
- Visible and dark sectors have equal but opposite matter-antimatter asymmetries (but total baryon number is conserved)

Preliminary $\Upsilon(4S)$

Search for B Mesogenesis (2)

Hadronic Recoil Method

« Signal » B meson:

- Reconstructed baryon (Λ, p)
- Missing ψ_D 4-momentum

(dark anti-baryon ψ_D escaping detection)

 $B^0 \rightarrow \Lambda \quad \psi_D \text{ PRD 107, 092001}$

 $B^+ \rightarrow p$ ψ_D Preliminary Presented at Moriond EW 2023

 $B^+ \rightarrow \Lambda_c^+ \psi_D$ will be out soon

- **Fully reconstructed** « Tag » B meson in known hadronic modes
 - Boosted Decision Tree (BDT) to suppress residual combinatorial backgrounds from $q\bar{q}$ and $B\bar{B}$ decays
 - Background estimated directly from $m(\psi_D)$ sideband data
 - Many $m(\psi_D)$ mass hypotheses tested

In $B^0 \rightarrow \Lambda \psi_D$ analysis, kinematic fit of $\Lambda \rightarrow p \pi$, including displaced vertex significance requirement

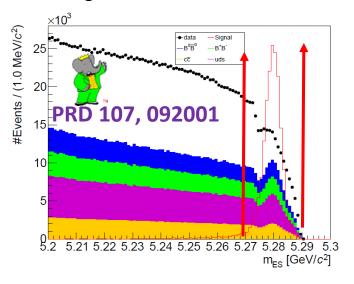
Background rejection

• Beam energy substituted mass m_{ES} : $\sqrt{E_{beam}^2 - p_{B_{Tag}}^2}$

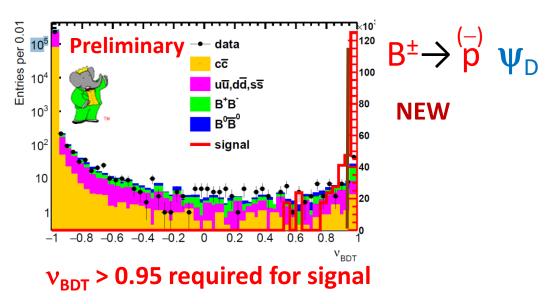
Peaks at B meson mass for signal

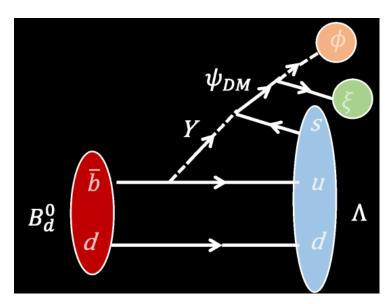
- Δ E relative to the beam energy : $E_{beam} E_{B_{Tag}}$ Centered at 0 for B signal events
- Event topology variables to separate $B\bar{B}$ more « spherical » events from more « jet-like » $q\bar{q}$ continuum events

m_{ES} for B_{Tag} in $B^0 \rightarrow \Lambda$ ψ_D analysis

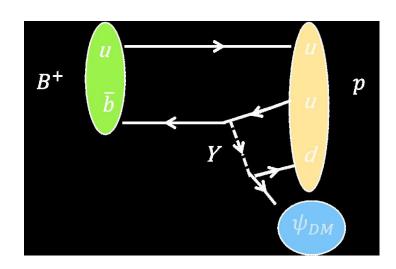


BDT output variable v_{BDT}





Probes O_{us}



Probes O_{ud}

G. Alonso-Alvarez, G. Elor, and M. Escudero Phys.Rev. D 104, 035028 (2021).

Dark SM SM

For the decay $B \to \psi_D \ \mathcal{B} \ \mathcal{M}$ to exist : new BSM TeV-scale color triplet scalar Y, coupling ψ_D and SM quarks

Effective low energy Lagrangian integrating out heavy mediator Y

$$\mathcal{L}_{eff} = \sum_{i,j} \mathfrak{O}_{u_i,d_j} \; \frac{y_{ij}^2}{M_Y^2}$$

 y_{ij}^2 Product of couplings

Four possible flavor combination Θ_{u_i,d_j} operators for B meson decays

Only one active in early universe Need to test four possibilities

$$\Theta_{\mathsf{ud}=\psi b u d}$$

$$\Theta_{us=\psi bus}$$

$$\mathfrak{O}_{\mathsf{cd}=\psi b c d}$$

$$\Theta_{cs=\psi bcs}$$

Search for B Mesogenesis (4)

Distribution of $m(\psi_D)$ (missing energy)



Upper Limit of Br(B⁰→ψ_bΛ)

 10^{-5}

 10^{-6}

PRD 107, 092001

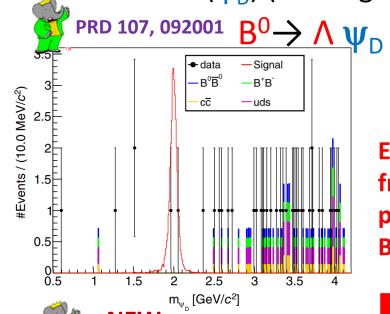
Belle previous limit

PRD 105. L051101 (2022)

Belle Experiment

3.5

BABAR Experiment



Exclude large fraction of parameter space for B mesogenesis

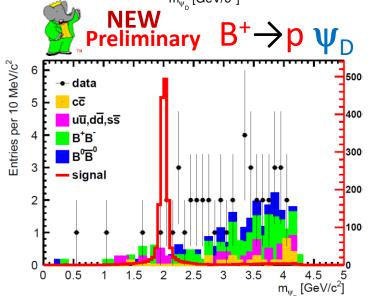


90% C.L. Upper limits on Br($B^+ \rightarrow p \psi_D$) vs m(ψ_D)

 $\mathcal{O}_{us}^1 = (\psi_D b)(us)$

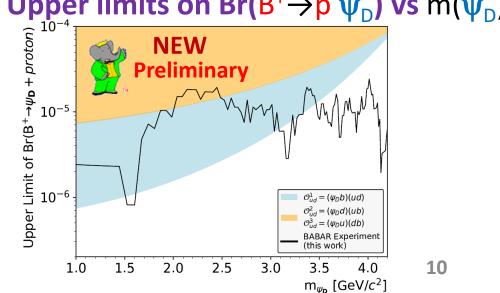
 $\mathcal{O}_{us}^2 = (\psi_D s)(ub)$

 $m_{\psi_{\mathbf{p}}}$ [GeV/ c^2]



BaBar analyses scan (many) $m\psi_D$ hypotheses As Belle looks at discrete $m\psi_D$ hypotheses

Result also useful to constrain SUSY model jhep02(2023)224 (see back-up)





Search for an Axion-Like Particle (ALP) in $B^+ \rightarrow K^+$ a (a $\rightarrow \gamma\gamma$)

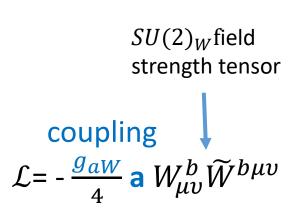
PhysRevLett.128.131802 (2022)

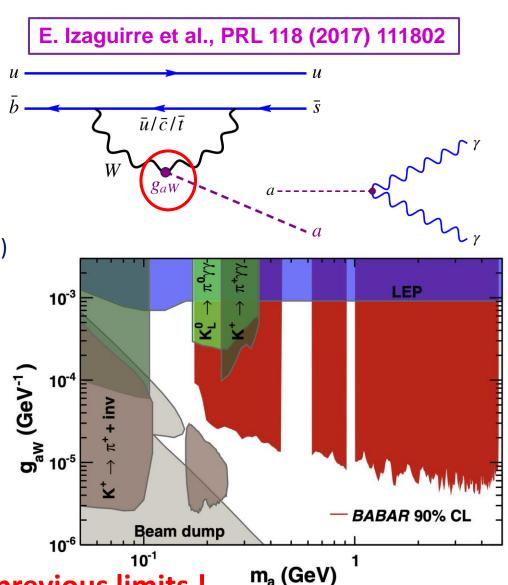
Pseudo-Goldstone bosons: Axion-Like Particles (ALPs) resulting from Spontaneously-broken global symmetries, Could help resolve naturalness issues of SM parameters (like CP strong Pb) and serve as mediators to dark sectors

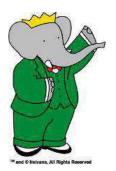
Can be produced in FCNC B decay processes, like $B \rightarrow K a$, with $a \rightarrow \gamma\gamma$ (SM BR prediction very small for $B \rightarrow K \gamma\gamma$)

Search for peak in reconstructed $\gamma\gamma$ mass

90% CL exclusion bounds on the ALP coupling g_{aW}







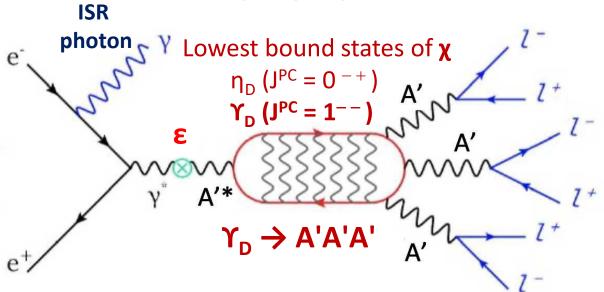
Search for **Darkonium**:

PhysRevLett.128.021802 (2022)

H. An et al., PRL 116 (2016) 151801

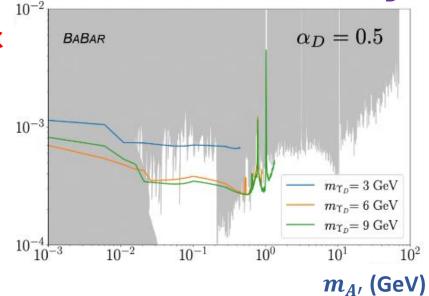
Dark photon A' couples strongly to the dark matter fermion (χ) via coupling α_D and A' mixes with SM photon via kinetic mixing with strength ϵ

Darkonium

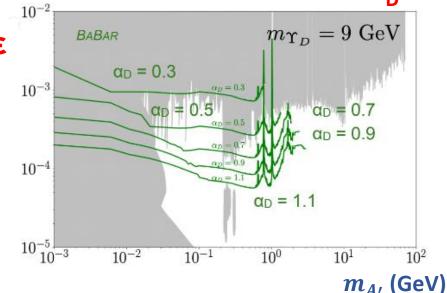


90% C.L. upper limits on the kinetic mixing parameter $oldsymbol{\epsilon}$ versus $m_{A'}$

For different values of $m(\Upsilon_D)$



For different values of $\alpha_{\rm D}$



3 dark photons A'

Six tracks final state : A' \rightarrow f +f - (f = e, μ , π)



CONCLUSION

BaBar data open an interesting and important window for searching for various physics beyond the Standard Model

Clean B factory environment: extremely well suited to searches for light dark sector Significant improvements in constraining dark sector

B mesogenesis, ALPs and **darkonium** searches are the most recent in a long, flourishing, and still developing history of dark sector and exotic searches.

One of the recent *BaBar* results not covered in this talk:

Search for Heavy Neutral Leptons from Taus (see back-up)

Phys. Rev. D 107, 052009 - Published 23 March 2023



BACK-UP SLIDES



BABAR papers



Presentation at Lake Louise 2023

Extensive program of searches for physics beyond the Standard Model, and dark sector in particular:

- Search for heavy neutral leptons in τ decays arXiv 2207.09575 [hep-ex] (accepted to PRD)
- Lepton universality in Y(3S) decays
 Phys. Rev. Lett .125, 241801 (2020)
- Search for LFV in Y(3S) →e µ
 Phys. Rev. Lett. 128, 091804 (2022)
- Rare and forbidden D decays Phys. Rev. Lett. 124, 071802 (2020)
- Search for LFV in D⁰ → X⁰ e μ
 Phys. Rev. D 101, 112003 (2020)

Search for Darkonium
 Phys. Rev. Lett. 128 021802 (2022)

arXiv:2302.00208 [hep-ex]

B mesogenesis

- Axion like particle
 Phys. Rev. Lett. 128, 131802 (2022).
- Dark Leptophilic scalar Phys. Rev. Lett. 125,181801 (2020).
- Six quark dark matter
 Phys. Rev. Lett. 122, 072002 (2019).
- Dark photon
 Phys. Rev. Lett. 113, 201801 (2014);
 Phys. Rev. Lett. 119, 131804 (2017).
- Muonic dark force
 Phys. Rev. D 94, 011102 (2016).
- Dark Higgs bosons
 Phys. Rev. Lett. 108, 211801 (2012)



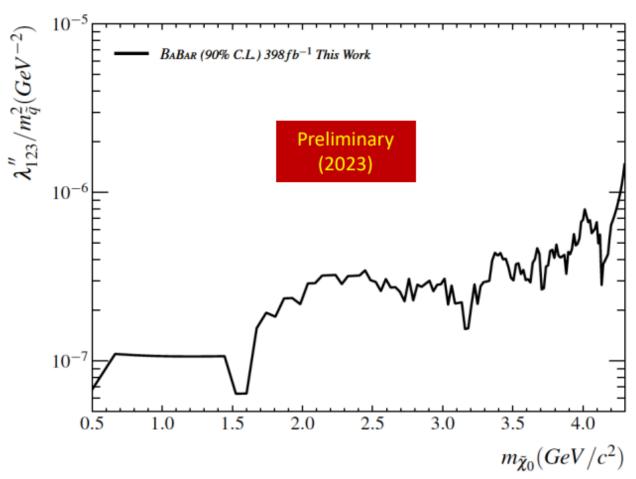
SUSY constraints from $B^{\pm} \rightarrow (p) + invisible$

C. O. Dib et al.

SUSY model: (Journal of High Energy Physics 2023 - **jhep02(2023)224**) also constrained by this result

Model with : R-parity violation (RPV) and Light neutralino $\tilde{\chi}_0$

Limits on the RPV coupling $\lambda_{123}^{\prime\prime}$ divided by the relevant squark mass squared



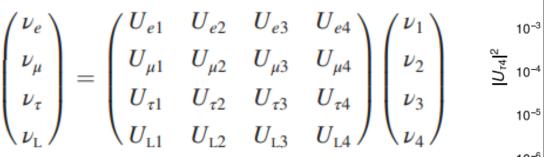
Search for Heavy Neutral Leptons from τ decays

Search for heavy neutral leptons using tau lepton decays at BABAR

Phys. Rev. D 107, 052009 – Published 23 March 2023 (Analysis covered at Moriond EW 2023)

Heavy neutral leptons : additional neutrino states v_n ("sterile neutrinos") mixing with active neutrinos v_l with mixing strength $|U_{ln}|^2$, U_{ln} element

of the extended PMNS matri $/\nu_e$ \



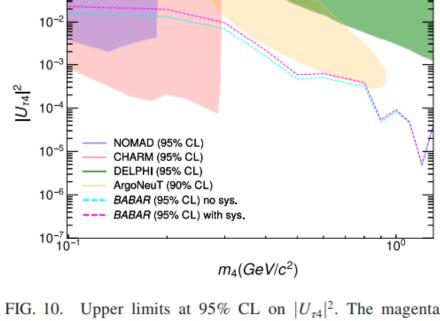
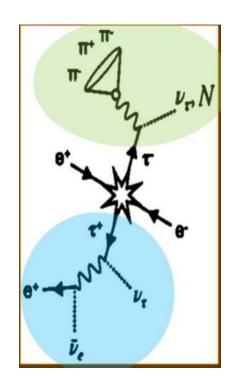


FIG. 10. Upper limits at 95% CL on $|U_{\tau 4}|^2$. The magenta line represents the result when uncertainties are included. The magenta line is expected to be a very conservative upper limit. Limits from Nomad [29], Charm [20], and Delphi [30] are also shown for reference. The recent ArgoNeuT result is also shown [31].

Model independent search:

- looking only at kinematics
- For a HNL $N=v_4$ mixing with v_{τ}
- Signal side : 3-prong pionic τ decay
- Tag side : leptonic τ decay



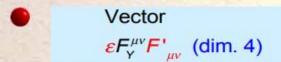


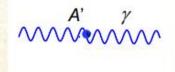
Link between SM and Dark Matter



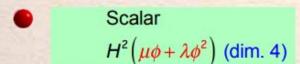
The SM may be connected to the dark sector through so-called portals, these links are the lowest-dimensional operators that may provide coupling of the dark sector to the SM (higher-dimensional operators are mass suppressed)

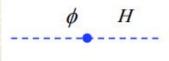
PLB 662, 53 (2008)



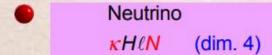


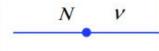
new U(1) symmetry \rightarrow dark photon, coupling to SM γ/Z via kinetic mixing ε



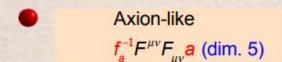


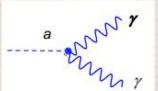
new dark scalar, mixing with Higgs





new heavy neutral lepton, mixing with left-handed SM doublet and Higgs

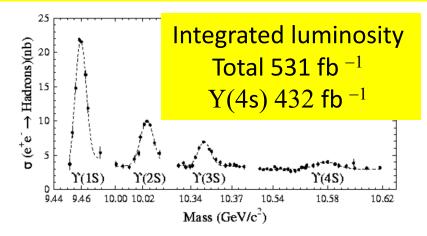




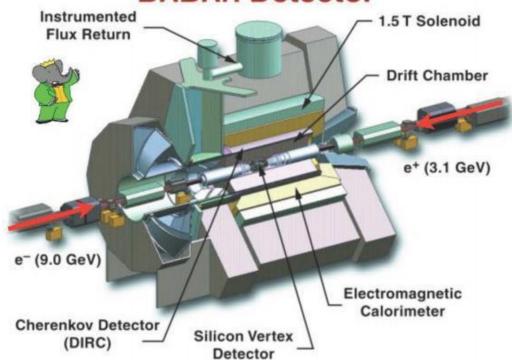
new axion/axion-like particle, coupling to SM gauge and fermion fields

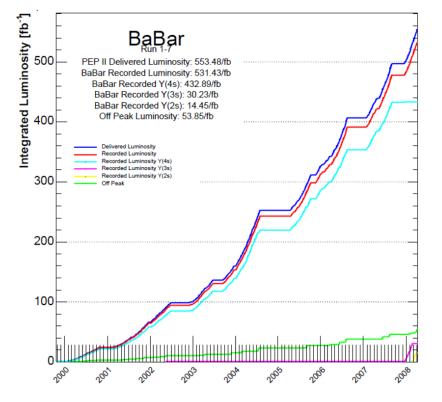
B factories provide tests of vector, scalar and axion-like portals in the lower-mass region

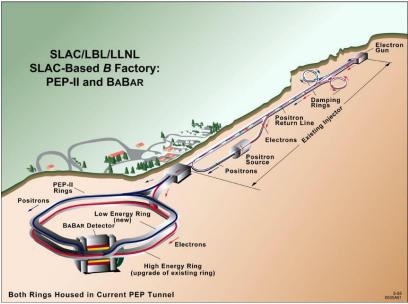
BaBar at SLAC PEPII: 1999 - 2008



BABAR Detector







Collider Signals of Baryogenesis and Dark Matter from B Mesons (B Mesogenesis)

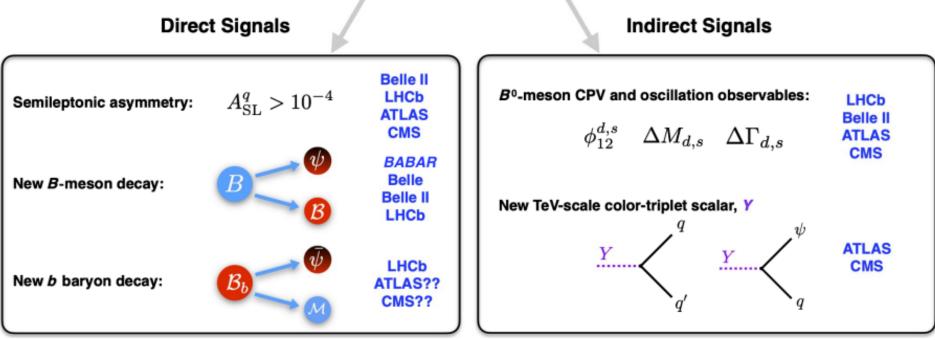
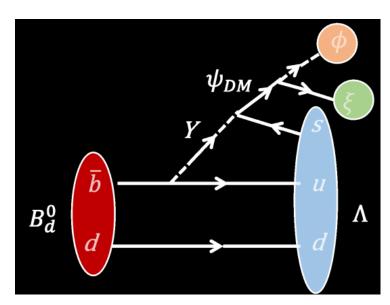
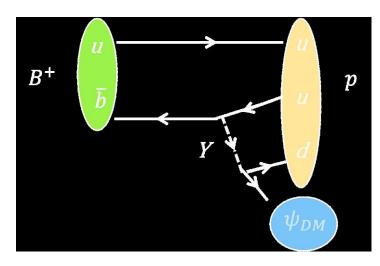


FIG. 1. Summary of the collider implications of baryogenesis and dark matter from B mesons [39], i.e., B-Mesogenesis. The distinctive signals of the mechanism are (i) the requirement that at least one of the semileptonic (CP) asymmetries in B_q^0 decays is $A_{SL}^q > 10^{-4}$, (ii) that both neutral and charged B mesons decay into a dark sector antibaryon (appearing as missing energy in the detector), a visible baryon, and any number of light mesons with $Br(B \to \psi BM) > 10^{-4}$, and (iii) that b-flavored baryons should decay into light mesons and missing energy at a rate $Br(B_b \to \bar{\psi}M) > 10^{-4}$. In addition, we include as indirect signals the various oscillation observables in the $B_q^0 - \bar{B}_q^0$ system as they are linked to A_{SL}^q , and the presence of a new TeV-scale color-triplet scalar Y that is needed to trigger the $B \to \psi BM$ decay. We also highlight the existing experiments that can probe each corresponding signal. Notation: B, B meson; B, B0 baryon; B1, any number of light mesons; B2, B3 dark sector antibaryon (ME in the detector).



Probes O_{us}



Probes O_{ud}

G. Alonso-Alvarez, G. Elorand, and M. Escudero, Phys.Rev. D 104, 035028 (2021).

In order for the $B \to \psi B \mathcal{M}$ decay to exist, a new BSM TeV-scale bosonic mediator is needed. In particular, this state should be a color-triplet scalar Y which couples to ψ and SM quarks. The LHC and flavor observables set relevant constraints on the mass and couplings of this color-triplet scalar which we discuss in detail in Sec. V. This heavy mediator can be integrated out to yield a low

energy Lagrangian of the form $\mathcal{L}_{\text{eff}} = \sum_{i,j} \mathcal{O}_{u_i d_j} \frac{y_{ij}^2}{M_Y^2}$, with

 y_{ij}^2 being the product of the two relevant dimensionless couplings. The four possible flavor combination operators \mathcal{O}_i of interest for *B*-meson decays are

$$\mathcal{O}_{ud} = \psi b u d,$$
 (15a) $\mathcal{O}_{us} = \psi b u s,$ (15b) $\mathcal{O}_{cd} = \psi b c d,$ (15c)

$$c_s = \psi b c s$$
,

(15d)

 y_{ij}^2 being the product of the two relevant dimensionless couplings. The four possible flavor combination operators \mathcal{O}_i of interest for *B*-meson decays are

$$\mathcal{O}_{ud} = \psi b u d, \tag{15a}$$

$$\mathcal{O}_{us} = \psi b u s, \tag{15b}$$

$$\mathcal{O}_{cd} = \psi b c d,$$
 (15c)

$$\mathcal{O}_{cs} = \psi b c s,$$
 (15d)

where all fermions are assumed to be right-handed⁶ and color indices are contracted in a totally antisymmetric way. These operators can induce the decay of the \bar{b} quark within the B meson into two light quarks and a dark antibaryon ψ . The resulting possible hadronic processes are summarized in Table I for the different operators in Eq. (15). Matrix elements involving the operators in Eq. (15) depend on the precise pairing of the spinors. Each of the operators can come in three different versions: "type 1" $\mathcal{O}_{ij}^1 = (\psi b)(u_i d_j)$, "type 2" $\mathcal{O}_{ij}^2 = (\psi d_j)(u_i b)$, and "type 3" $\mathcal{O}_{ij}^3 = (\psi u_i)(d_j b)$. This distinction becomes relevant for some of the constraints discussed in the next sections.

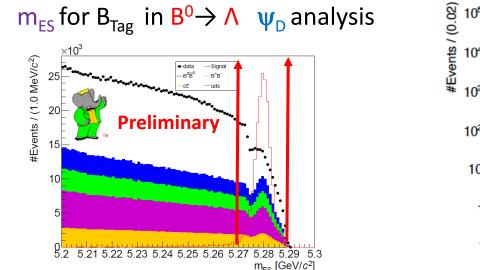
As we will see in Sec. V, <u>flavor constraints on the Y</u> triplet scalar imply that only one of these operators can be active in the early Universe. In practice, this means that we only expect one dominant flavor combination of these possible operators at collider experiments and not a combination of the above. Therefore, only one of the sets of decay channels listed in Table I is expected to have a sizeable branching ratio, while all others should be suppressed.

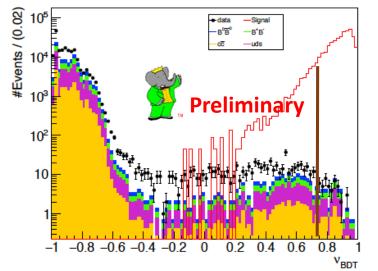
TABLE I. The lightest final state resulting from the new decay of b quarks as necessary to give rise to baryogenesis and dark matter production. We list each of the possible flavorful operators that can equally lead to B-Mesogenesis; see Eq. (15). For a given operator, the rate of each decay is fairly similar given that $m_{B^\pm} \simeq m_{B_d^0} \simeq m_{B_s^0} \sim m_{\Lambda_b}$. ΔM refers to the difference in mass between the initial and final SM hadron. Note that additional light mesons can be present in the final state, which act to decrease ΔM by their corresponding masses.

Operator and decay	Initial state	Final state	ΔM (MeV)	
$ \mathcal{O}_{ud} = \psi b u d \\ \bar{b} \to \psi u d $	B_d B_s B^+ Λ_b	$\psi + n(udd)$ $\psi + \Lambda(uds)$ $\psi + p(duu)$ $\bar{\psi} + \pi^{0}$	4340.1 4251.2 4341.0 5484.5	Recent BaBar
$\mathcal{O}_{us} = \psi bus$ $\bar{b} \to \psi us$	$egin{array}{c} B_d \ B_s \ B^+ \ \Lambda_b \end{array}$	$\psi + \Lambda(usd)$ $\psi + \Xi^{0}(uss)$ $\psi + \Sigma^{+}(uus)$ $\bar{\psi} + K^{0}$	4164.0 4025.0 4090.0 5121.9	results
$\mathcal{O}_{cd} = \psi b c d$ $\bar{b} \to \psi c d$	$egin{array}{c} B_d & B_s \ B^+ \ \Lambda_b \end{array}$	$\psi + \Lambda_c + \pi^-(cdd)$ $\psi + \Xi_c^0(cds)$ $\psi + \Lambda_c^+(dcu)$ $\bar{\psi} + \bar{D}^0$	2853.6 2895.0 2992.9 3754.7	BaBar Result Coming soon
$\mathcal{O}_{cs} = \psi b cs$ $\bar{b} \to \psi cs$	$egin{array}{c} B_d \ B_s \ B^+ \ \Lambda_b \end{array}$	$ \psi + \Xi_c^0(csd) $ $ \psi + \Omega_c(css) $ $ \psi + \Xi_c^+(csu) $ $ \bar{\psi} + D^- + K^+ $	2807.8 2671.7 2810.4 3256.2	

arXiv:2302.00208v1 (submitted to Phys. Rev. D)

Main systematics from Monte-Carlo-Data correction in BDT distributions using control region in data





 $B^0 \rightarrow \Lambda \psi_D$

 $v_{BDT} > 0.75$ required for signal

The inclusive background MC samples do not accurately reproduce the data. This discrepancy arises from a mis-modeling of several branching fractions used in the simulation, resulting in differences in Btag reconstruction efficiencies, as well as differences in charged and neutral particle reconstruction efficiencies, PID efficiencies, and the modeling of variables used in the BDT. We **correct the simulation in a two-step procedure**, **using sideband data** selected with the criteria, described above, applied before the BDT selection, except with the looser requirement $5.20 < m_{FS} < 5.29 \text{ GeV/c}^2$.

- The region -0.5 < BDT < 0.75, largely dominated by $e^+e^- \rightarrow q \overline{q}$ (q = u; d; s; c) events, is used to extract a correction factor for continuum production, f_{udsc} , by rescaling the corresponding MC predictions to the number of observed events.
- The correction factor for $B\overline{B}$ production, is determined from data in the complementary region BDT < 0.5, assuming equal contributions from $B^0\overline{B^0}$ and B+B-. We obtain $f_{udsc} = 1.34 \pm 0.10$ and $f_{udsc} = 1.06 \pm 0.08$. Under the assumption that the $B^0\overline{B^0}$ correction factor is independent of the signal B decay mode, we rescale the signal efficiency by $f_{udsc} = 1.06 \pm 0.08$. Under the assumption that the corresponding uncertainty as a systematic uncertainty.



Search for Axion-Like Particles (ALP)- (1)

E. Izaguirre et al., PRL 118 (2017) 111802

Many extensions of SM include spontaneously-broken global symmetries, resulting in pseudo-Goldstone bosons : **Axion-Like Particles (ALPs)**

Could help resolve issues of naturalness of SM parameters (like CP strong Pb) and serve as mediators to dark sectors

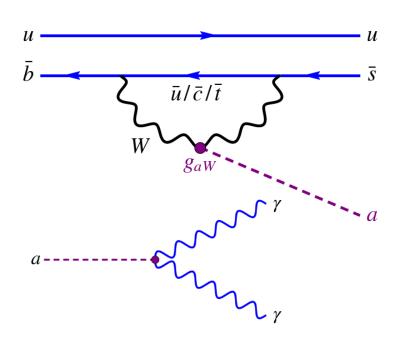
ALPs (a) couple primarily to pairs of SM gauge bosons (coupling g_{aw}).

Can be produced in FCNC B decay processes, like $B \rightarrow K a$, with $a \rightarrow \gamma\gamma$ (nearly 100% BF for m(a) < m(W))

For low axion mass and small coupling g_{aW} , the axion lifetime can become "long", i.e. non-prompt $\tau \sim 1 / m_a^3 g_{aW}^2$

Very Small SM BR prediction

for: $B \rightarrow K \gamma \gamma$





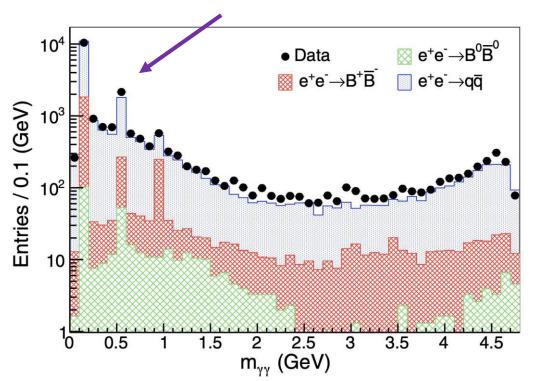
Search for Axion-Like Particles (ALP)- (2)

Searches for ALPs in $B^+ \rightarrow K^+$ a (a $\rightarrow \gamma\gamma$)

Phys. Rev. Lett. 128, 131802 (2022)

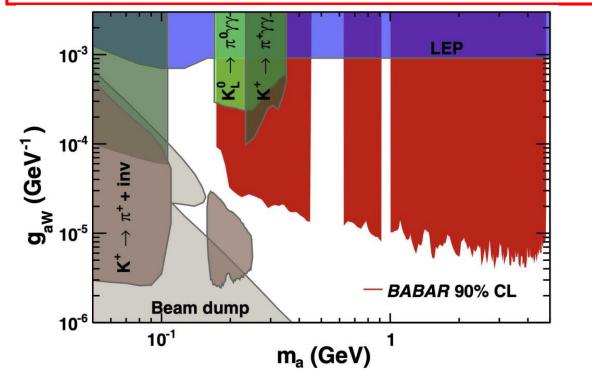
Search for peak in the reconstructed yy mass

peaking contributions from $\pi 0$, η , η' ηc





90% CL exclusion bounds on the ALP coupling g_{aW}



Up to two orders of magnitude improvements over previous limits



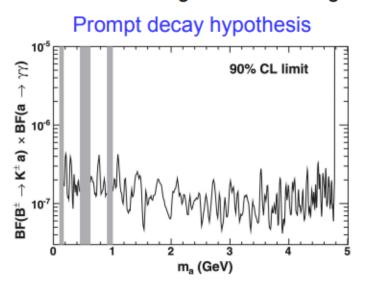
Axion-Like Particles

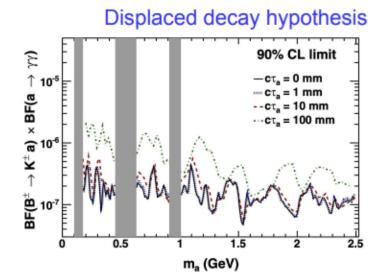


Phys. Rev. Lett. 128, 131802 (2022)

Scan $m_{\gamma\gamma}$ with steps equal to the signal mass resolution (~ 8 – 14 MeV)

461 signal mass hypotheses fit with unbinned ML fits to a hypothetical signal peak
 + smooth background over range of ~24 – 60 σ around each hypothesis





In low mass region ($m_{\gamma\gamma}$ < 2.5 GeV) the signal sensitivity is also assessed for non-prompt signal hypotheses: $c\tau$ = 1, 10, 100mm

- · displaced vertex not reconstructed, but ALP resolution degraded
- No significant excess observed

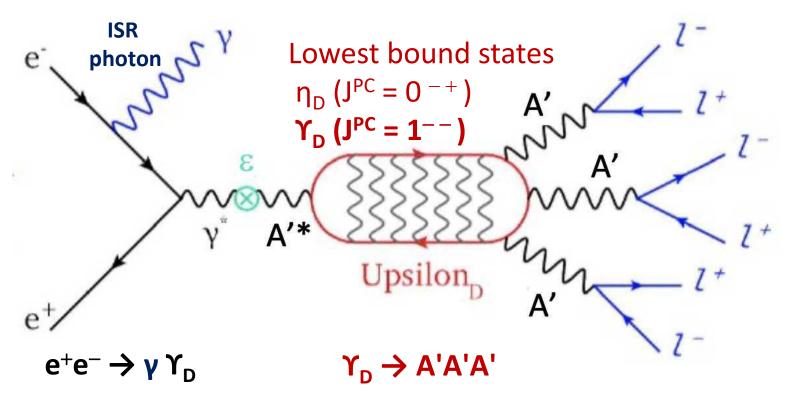
Presentation at Lake Louise 2023



Search for Darkonium (1)

PRL 128, 021802 (2022)

Dark matter bound states arise in some simple dark photon (A') models in which : the A' couples strongly to the dark matter fermion (χ) via coupling α_D and A' mixes with SM photon via kinetic mixing with strength ϵ



H. An et al., PRL 116 (2016) 151801

Final state : A' \rightarrow f +f - (f = e, μ , π)

Dark photons A' decay to pairs of leptons or pions BaBar search in **six-track final state** in 514 fb⁻¹

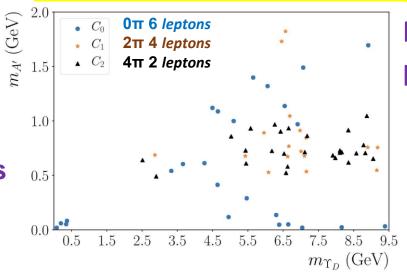
Dark photon lifetime can be long for small masses and small kinetic mixing ε hence prompt and displaced vertex signatures



Search for Darkonium (2)

PRL 128, 021802 (2022)

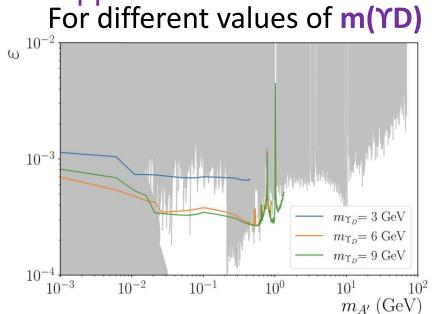
Search for 3 dark photons Prompt decays

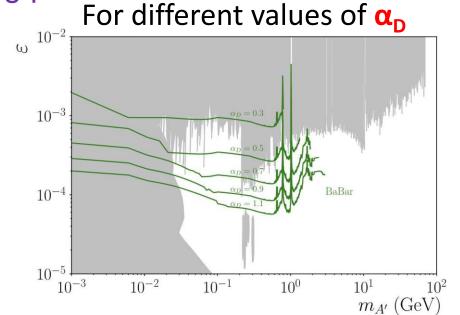


No significant signals observed in either prompt or displaced decay searches

Dark photon A' couples strongly to the dark matter fermion (χ) via coupling α_D and A' mixes with SM photon via kinetic mixing with strength ϵ

90% C.L. Upper limits on the kinetic mixing parameter ε as a function of m(A')





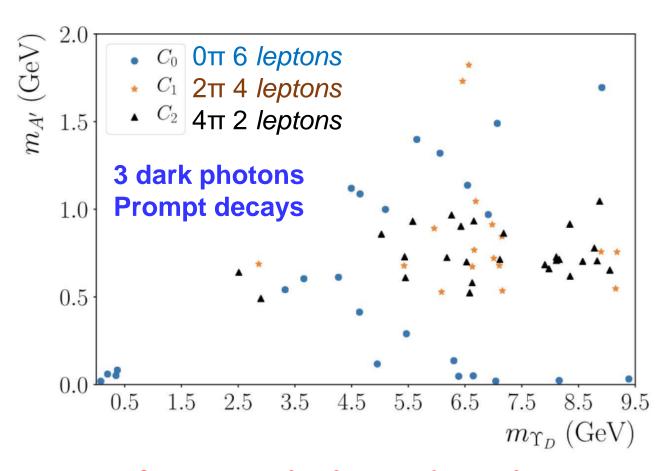


Search for **Darkonium**

Final state selection: 3 pairs of opposite-sign tracks (at least one lepton pair) which should all have same invariant mass

- Reconstruct Υ_D mass
- ISR photon may or may not be detected, but recoil mass against Υ_D should be consistent with zero
- MVA used to suppress backgrounds
- Scan m(YD) m(A') plane

PRL 128, 021802 (2022)



No significant signals observed in either prompt or displaced decay searches 2



Search for **Darkonium**

Final state selection: 3 pairs of opposite-sign tracks (at least one lepton pair) which should all have same invariant mass

- Reconstruct Υ_D mass
- ISR photon may or may not be detected, but recoil mass against Υ_D should be consistent with zero
- MVA used to suppress backgrounds
- Scan m(YD) m(A') plane

PRL 128, 021802 (2022)

3 dark photons Non prompt decays

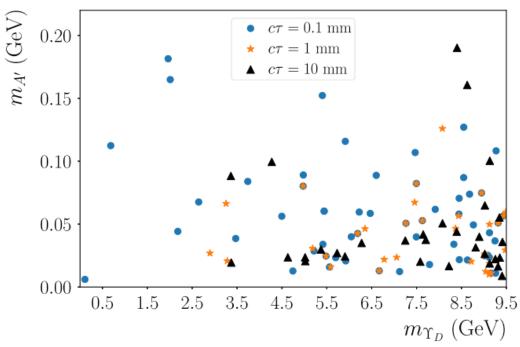


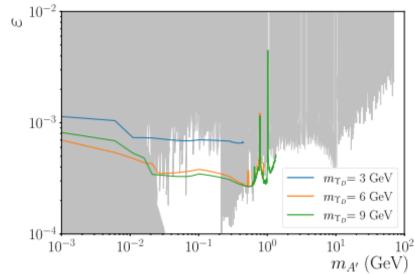
FIG. 4. The $(m_{\Upsilon_D}, m_{A'})$ mass distribution of event candidates passing all selection criteria for the datasets optimized for each dark photon lifetime.

Search for Darkonium

From PHYSICAL REVIEW LETTERS 128, 021802 (2022)

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Constraints other than Recent BaBar result (gray)



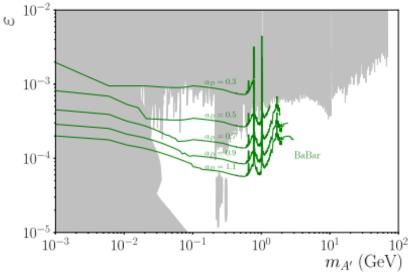


FIG. 6. The 90% C.L. upper limits on the kinetic mixing ε for (top) various Υ_D masses assuming $\alpha_D = 0.5$ and (bottom) various α_D values assuming $m_{\Upsilon_D} = 9$ GeV together with current constraints (gray area) [8–18].