Measurements of $b ightarrow s \mu^+ \mu^-$ transitions at LHCb

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$b ightarrow s \mu^+ \mu^-$ decays as a probe for New Physics

SM:



- $b \rightarrow s\mu^+\mu^-$ transitions occur via FCNC \rightarrow cannot occur at tree level in SM
 - New particles:
 - ◊ enhance/suppress decay rates
 - modify angular distribution of final state particles
 - introduce new sources of CP violation

Possible NP contributions:



Leptoquarks (tree-level)



Supersymmetry (loop-level)

Heavy Quark Effective Field Theory (HQEFT) for $b ightarrow s \mu^+ \mu^-$ decays

- Search for BSM physics in a model independent way
- Integrate out interesting heavy physics (at m_W):



Effective Hamiltonian

$$\mathcal{H}_{eff} = -rac{4G_F}{\sqrt{2}}V_{tb}V_{ts}^*\sum_i rac{\mathcal{C}_i^{(\prime)}\mathcal{O}_i^{(\prime)}}{i}$$

- Wilson Coefficients (Effective Coupling)
- Local operators



b-hadron physics at LHCb

Optimised for *b*-hadron physics Forward spectrometer (where most $b\bar{b}$ is produced)



- Vertex Locator
 - Separate b and c hadron production and decay vertices at high precision
- Ring Imaging Cherenkov (RICH) Detectors
 - ♦ PID of K, p, π
 - $\diamond~$ High K PID efficiency: $\sim 95\%$
 - \diamond Low hadron mis-ID: 5% ($\pi \rightarrow K$)
- Muon System
 - \diamond High μ PID efficiency: \sim 97%
 - \diamond Low hadron mis-ID: 1 3% ($\pi \rightarrow \mu$)

Deviations from SM in $b \rightarrow s\mu^+\mu^-$ decays at LHCb (Branching Fraction Measurements)



- Measurements **below SM** by $1-3\sigma$ levels
- Sizeable hadronic uncertainties ($\sim 20 30\%$) in SM calculations \rightarrow need for improved theory predictions

Deviations from SM in $b \rightarrow s \mu^+ \mu^-$ decays at LHCb (Angular Analyses)





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- Measurements in tension with SM predictions (1-3 σ levels)
- Sizeable hadronic uncertainties ($\sim 20-30\%)$ in SM calculations
 - \rightarrow need for improved theory predictions

Improved Theory Predictions at Low q^2 (Branching Fraction Measurements)



Use of novel parameterisation of non-local QCD form factors





Use of form factors from $N_f = 2 + 1 + 1$ lattice QCD

• Tensions between SM and experiment are still observed in most cases (agreement in $B \rightarrow K^* \mu \mu$)

Theory Explanations for the $B ightarrow {\cal K}^* \mu \mu$ anomaly

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- Hadronic contributions could be severely underestimated (e.g. $B^0 \rightarrow D^*D_s \rightarrow K^{*0}\mu\mu$: Phys.Rev.Lett. 125(2020) 1,011802)
- Results can be explained by an apparent shift in C_9 (charm loop)

Current Strategy

• Extraction of a limited set of observables in **bins of** q^2



Example: Angular analysis $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (Phys. Rev. Lett. 126, 161802 (2021), JHEP 1308 (2013) 131, JHEP 06 (2015) 084)

Explore Additional Strategies

Increase in data and theory developments allow:

- New approach to determine $B \to K^* \mu \mu$ amplitudes as continuous distributions in q^2
 - Able to exploit relations between observables that are inaccessible in binned fits to observables
 - \diamond Able to exploit q^2 shape information via unbinned fits
 - $\diamond\,$ Eliminates the need to correct theory predictions for q^2 averaging effects

Increases sensitivity to NP!

• more work is still required to fully account for $B\to D^*D_s\to K^{(*)}\mu\mu$ rescattering amplitudes

Direct measurements of Wilson Coefficients

• Unbinned fits allow for direct extraction of Wilson Coefficients



Direct measurements of Wilson Coefficients (Form Factors)



• q^2 spectrum has theory uncertainties both **local** and **non-local** contributions:

Local:

 Form-factors well described by: Lattice QCD (Phys. Rev. D 107 (2023) 014510, Phys. Rev. D 93, 025026 (2016))
 Light Cone Sum rules (JHEP 01 (2019) 150)

Non-Local:

 Far from resonances: estimations are made using perturbative bounds (Nucl.Phys.B612:25-58,2001, JHEP 1009 (2010) 089)

$B^+ \rightarrow K^+ \mu^+ \mu^-$ decay rate as a function of q^2 Obtain a model of the decay rate as a function of q^2 :

$$\frac{d\Gamma}{dq^{2}} = \frac{\alpha_{em}^{2} G_{F}^{2} | V_{tb} V_{ts}^{*} |^{2}}{128\pi^{5}} \kappa \left(q^{2}\right) \beta \left(q^{2}\right) \left\{\frac{2}{3} \kappa^{2} \left(q^{2}\right) \beta^{2} \left(q^{2}\right) | \mathcal{C}_{10}^{\mu} f_{+} \left(q^{2}\right) |^{2} + \frac{m_{\mu}^{2} \left(m_{B}^{2} - m_{K}^{2}\right)^{2}}{q^{2} m_{B}^{2}} | \mathcal{C}_{10}^{\mu} f_{0} \left(q^{2}\right) |^{2} + \kappa^{2} \left(q^{2}\right) \left[1 - \frac{1}{3} \beta^{2} \left(q^{2}\right)\right] \left|\mathcal{C}_{9}^{\mu} \text{, eff}_{f_{+}} \left(q^{2}\right) + 2\mathcal{C}_{7} \frac{m_{b} + m_{s}}{m_{B} + m_{K}} f_{7} \left(q^{2}\right) \right|^{2} \right\} \otimes \mathbb{R}(q^{2})$$

• Form Factors

• Wilson Coefficients
$$Improve mass resolution by performing a kinematic fit with m_{max}(B) a constraint$$

Structure of $C_9^{\mu, { m eff}}$ in $B^+ o K^+ \mu^+ \mu^-$ decay rate as a function of q^2



- Rely on once-subtracted dispersion relation that includes $D\bar{D} \to \mu\mu$ and $\tau\tau \to \mu\mu$ amplitudes
- $Y_{c\bar{c}}^{(0)}$ subtraction term to ensure convergence at large q^2

$$B^+
ightarrow K^+ \mu^+ \mu^-$$

(Example of earlier isobar naive approach which ignores 2 particle states)



Run 1 Analysis (EPJC (2017) 77: 161)

- Degeneracy of J/ψ and ψ_{2S} phases lead to 4 equivalent solutions
- Run 2 analysis (following the dispersion relation) currently in WG review

Extension to $B ightarrow K^* \mu \mu$

(based off EPJC (2018) 78: 453)

$$\mathcal{A}_{\lambda}^{L,R} = \mathcal{N}_{\lambda} \left\{ \left(\mathcal{C}_{9} \mp \mathcal{C}_{10} \right) \mathcal{F}_{\lambda} \left(\boldsymbol{q}^{2} \right) + \frac{2m_{b}M_{B}}{q^{2}} \left[\mathcal{C}_{7} \mathcal{F}_{\lambda}^{T} \left(\boldsymbol{q}^{2} \right) - 16\pi^{2} \frac{M_{B}}{m_{b}} \mathcal{H}_{\lambda} \left(\boldsymbol{q}^{2} \right) \right] \right\}$$

- Form Factors
- Wilson Coefficients
- Non-local hadronic matrix elements

Two approaches pursued at LHCb (currently ongoing):

- 1. Expand $\mathcal{H}_{\lambda}(q^2)$ as a polynomial in $z(q^2)^1$ and fitting simultaneously (Chrzaszcz et al., JHEP 10 (2019) 236) with:
 - External inputs coming from J/ψ and ψ_{2S} measurements
 - Theory points in negative q^2 region
- 2. Include all known contributions to C_9 (combine approaches of Egede et al., EPJC 78 (2018) 6, 453 and Cornella et al., EPJC 80 (2020) 12, 1095) \rightarrow fit to full q^2 spectrum

¹Conformal mapping of q^2 to the unit circle

Extension to $B \rightarrow K^* \mu \mu$ (based off EPJC (2018) 78: 453)

Sensitivity studies with pseudo-experiments



Left: Fits to z-expansion of $\mathcal{H}_{\lambda}\left(q^{2}\right)$ with negative q^{2} theory inputs. **Right:** 2D sensitivity scans for Wilson Coefficients. (approach 1)



(approach 2)

Future Prospects

- Tensions between SM theory and experiment **persist**, independent of recent status of LFU violation
- Model of the strong phase with q² allows for extra sensitivity of the imaginary parts of the Wilson Coefficients
 → work ongoing
- Continue with the robust approach of binned measurements
 - \rightarrow However, in order to take advantage of:
 - the increase in datasets
 - sensitivity to the tau loop (motivated by $R(D^{0(*)})$)

we employ the **new unbinned approach**