# Measurements of $b \rightarrow s \mu^{+} \mu^{-}$transitions at LHCb 

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30 May 2023<br>FPCP - Lyon, France



## $b \rightarrow s \mu^{+} \mu^{-}$decays as a probe for New Physics

## SM:



## Possible NP contributions:



Leptoquarks (tree-level)
$\rightarrow$ cannot occur at tree level in SM

- New particles:
$\diamond$ enhance/suppress decay rates
$\diamond$ modify angular distribution of final state particles
$\diamond$ introduce new sources of CP violation


## Heavy Quark Effective Field Theory (HQEFT) for

 $b \rightarrow s \mu^{+} \mu^{-}$decays- Search for BSM physics in a model independent way
- Integrate out interesting heavy physics (at $m_{W}$ ):


Full Theory


Effective Theory


## Effective Hamiltonian

$$
\mathcal{H}_{e f f}=-\frac{4 G_{F}}{\sqrt{2}} V_{t b} V_{t s}^{*} \sum_{i} \mathcal{C}_{i}^{(\prime)} \mathcal{O}_{i}^{\left({ }^{\prime}\right)}
$$

- Wilson Coefficients (Effective Coupling)
- Local operators



## b-hadron physics at LHCb

Optimised for $b$-hadron physics
Forward spectrometer (where most $b \bar{b}$ is produced)

- Vertex Locator
$\diamond$ Separate $b$ and $c$ hadron production and decay vertices at high precision
- Ring Imaging Cherenkov (RICH) Detectors
$\diamond$ PID of $K, p, \pi$
$\diamond$ High K PID efficiency: $\sim 95 \%$
$\diamond$ Low hadron mis-ID: $5 \%(\pi \rightarrow K)$
- Muon System
$\diamond$ High $\mu$ PID efficiency: $\sim 97 \%$
$\diamond$ Low hadron mis-ID: $1-3 \%(\pi \rightarrow \mu)$


## Deviations from SM in $b \rightarrow s \mu^{+} \mu^{-}$decays at LHCb (Branching Fraction Measurements)



JHEP 04 (2017) 142


JHEP 06 (2014) 133
JHEP 06 (2014) 133


JHEP 09 (2018) 145



Phys. Lett. B 753 (2016) 424


- Measurements below SM by $1-3 \sigma$ levels
- Sizeable hadronic uncertainties ( $\sim 20-30 \%$ ) in SM calculations $\rightarrow$ need for improved theory predictions


## Deviations from SM in $b \rightarrow s \mu^{+} \mu^{-}$decays at LHCb (Angular Analyses)



LHCb Coll. JHEP 09 (2015) 179


- Measurements in tension with SM predictions ( $1-3 \sigma$ levels)
- Sizeable hadronic uncertainties ( $\sim 20-30 \%$ ) in SM calculations $\rightarrow$ need for improved theory predictions


## Improved Theory Predictions at Low $q^{2}$ (Branching Fraction Measurements)

JHEP 09 (2022) 133


Use of novel parameterisation of non-local QCD form factors


Use of form factors from $N_{f}=2+1+1$ lattice QCD

- Tensions between SM and experiment are still observed in most cases (agreement in $B \rightarrow K^{*} \mu \mu$ )


## Theory Explanations for the $B \rightarrow K^{*} \mu \mu$ anomaly



- Hadronic contributions could be severely underestimated ( e.g. $B^{0} \rightarrow D^{*} D_{s} \rightarrow K^{* 0} \mu \mu$ : Phys.Rev.Lett. 125(2020) 1,011802)
- Results can be explained by an apparent shift in $C_{9}$ (charm loop)


## Current Strategy

- Extraction of a limited set of observables in bins of $\boldsymbol{q}^{\mathbf{2}}$


Example: Angular analysis $B^{0} \rightarrow K^{* 0} \mu^{+} \mu^{-}$
(Phys. Rev. Lett. 126, 161802 (2021), JHEP 1308 (2013) 131, JHEP 06 (2015) 084)

## Explore Additional Strategies

Increase in data and theory developments allow:

- New approach to determine $B \rightarrow K^{*} \mu \mu$ amplitudes as continuous distributions in $\boldsymbol{q}^{\mathbf{2}}$
$\diamond$ Able to exploit relations between observables that are inaccessible in binned fits to observables
$\diamond$ Able to exploit $q^{2}$ shape information via unbinned fits
$\diamond$ Eliminates the need to correct theory predictions for $q^{2}$ averaging effects


## Increases sensitivity to NP!

- more work is still required to fully account for $B \rightarrow D^{*} D_{s} \rightarrow K^{(*)} \mu \mu$ rescattering amplitudes


## Direct measurements of Wilson Coefficients

- Unbinned fits allow for direct extraction of Wilson Coefficients

An example: $B^{+} \rightarrow K^{+} \mu^{+} \mu^{-}$

## Short distance contributions <br> (Sensitive to NP)

Long distance contributions
(Resonances and $D D$ rescattering contributions)


- Able to simultaneously extract $C_{9}$ and $C_{10}$


## Direct measurements of Wilson Coefficients (Form Factors)

## An example: $B^{+} \rightarrow K^{+} \mu^{+} \mu^{-}$

## Short distance contributions (Sensitive to NP)

Long distance contributions (Resonances and DD rescattering contributions)


- $q^{2}$ spectrum has theory uncertainties both local and non-local contributions:

Local:
$\diamond$ Form-factors well described by:
Lattice QCD (Phys. Rev. D 107 (2023) 014510, Phys. Rev. D 93, 025026 (2016))
Light Cone Sum rules (JHEP 01 (2019) 150)
Non-Local:
$\diamond$ Far from resonances: estimations are made using perturbative bounds (Nucl.Phys.B612:25-58,2001, JHEP 1009 (2010) 089)

## $B^{+} \rightarrow K^{+} \mu^{+} \mu^{-}$decay rate as a function of $q^{2}$

## Obtain a model of the decay rate as a function of $q^{2}$ :

$$
\begin{aligned}
\frac{d \Gamma}{d q^{2}}=\frac{\alpha_{\text {em }}^{2} G_{F}^{2}\left|V_{t b} V_{t s}^{*}\right|^{2}}{128 \pi^{5}} & \kappa\left(q^{2}\right) \beta\left(q^{2}\right)\left\{\frac{2}{3} \kappa^{2}\left(q^{2}\right) \beta^{2}\left(q^{2}\right)\left|\mathcal{C}_{10}^{\mu} f_{+}\left(q^{2}\right)\right|^{2}+\frac{m_{\mu}^{2}\left(m_{B}^{2}-m_{K}^{2}\right)^{2}}{q^{2} m_{B}^{2}}\left|\mathcal{C}_{10}^{\mu} f_{0}\left(q^{2}\right)\right|^{2}\right. \\
& \left.+\kappa^{2}\left(q^{2}\right)\left[1-\frac{1}{3} \beta^{2}\left(q^{2}\right)\right]\left|\mathcal{C}_{9}^{\mu, \text { eff }} f_{+}\left(q^{2}\right)+2 \mathcal{C}_{7} \frac{m_{b}+m_{s}}{m_{B}+m_{K}} f_{T}\left(q^{2}\right)\right|^{2}\right\} \otimes \mathcal{R}\left(q^{2}\right)
\end{aligned}
$$

- Form Factors
- Wilson Coefficients



## Structure of $C_{9}^{\mu \text {,eff }}$ in $B^{+} \rightarrow K^{+} \mu^{+} \mu^{-}$decay rate as a function of $q^{2}$

e.g. Cornella et al., EPJC 80 (2020) 12. 1095


- Rely on once-subtracted dispersion relation that includes $D \bar{D} \rightarrow \mu \mu$ and $\tau \tau \rightarrow \mu \mu$ amplitudes
- $Y_{c \bar{c}}^{(0)}$ subtraction term to ensure convergence at large $q^{2}$

$$
B^{+} \rightarrow K^{+} \mu^{+} \mu^{-}
$$

(Example of earlier isobar naive approach which ignores 2 particle states)


- Degeneracy of $J / \psi$ and $\psi_{2 S}$ phases lead to 4 equivalent solutions
- Run 2 analysis (following the dispersion relation) currently in WG review


## Extension to $B \rightarrow K^{*} \mu \mu$

$$
\left.\mathcal{A}_{\lambda}^{L, R}=N_{\lambda}\left\{\left(C_{9} \mp C_{10}\right) \mathcal{F}_{\lambda}\left(q^{2}\right)+\frac{2 m_{b} M_{B}}{q^{2}}\left[C_{7} \mathcal{F}_{\lambda}^{T}\left(q^{2}\right)-16 \pi^{2} \frac{M_{B}}{m_{b}} \mathcal{H}_{\lambda}\left(q^{2}\right)\right]\right\} \right\rvert\,
$$

- Form Factors
- Wilson Coefficients
- Non-local hadronic matrix elements

Two approaches pursued at LHCb (currently ongoing):

1. Expand $\mathcal{H}_{\lambda}\left(q^{2}\right)$ as a polynomial in $z\left(q^{2}\right)^{1}$ and fitting simultaneously (Chrzaszcz et al., JHEP 10 (2019) 236) with:

- External inputs coming from $J / \psi$ and $\psi_{2 S}$ measurements
- Theory points in negative $q^{2}$ region

2. Include all known contributions to $C_{9}$ (combine approaches of Egede et al., EPJC 78 (2018) 6, 453 and Cornella et al., EPJC 80 (2020) 12, 1095) $\rightarrow$ fit to full $q^{2}$ spectrum
[^0]
## Extension to $B \rightarrow K^{*} \mu \mu$

## (based off EPJC (2018) 78: 453)

- Sensitivity studies with pseudo-experiments



Left: Fits to z-expansion of $\mathcal{H}_{\lambda}\left(q^{2}\right)$ with negative $q^{2}$ theory inputs.
Right: 2D sensitivity scans for Wilson Coefficients.
(approach 1)


Angular observables as a function of $q^{2}$ (approach 2)

## Future Prospects

- Tensions between SM theory and experiment persist, independent of recent status of LFU violation
- Model of the strong phase with $q^{2}$ allows for extra sensitivity of the imaginary parts of the Wilson Coefficients
$\rightarrow$ work ongoing
- Continue with the robust approach of binned measurements
$\rightarrow$ However, in order to take advantage of:
- the increase in datasets
- sensitivity to the tau loop (motivated by $R\left(D^{0(*)}\right)$ ) we employ the new unbinned approach


[^0]:    ${ }^{1}$ Conformal mapping of $q^{2}$ to the unit circle

