

# Charm physics at BESIII

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FPCP Conference

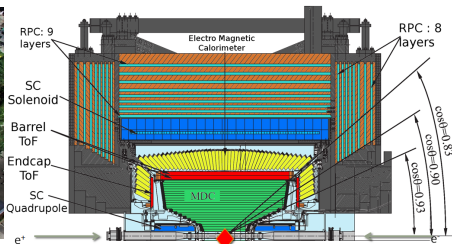
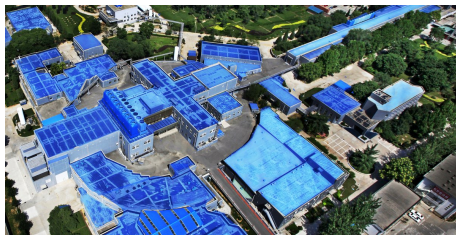
30th May 2023



- 1 Charm physics at the BESIII experiment
- 2  $D \rightarrow K^- \pi^+$
- 3  $D \rightarrow K^- \pi^+ \pi^- \pi^+$
- 4  $D \rightarrow K^+ K^- \pi^+ \pi^-$
- 5 Summary and conclusion

# The BESIII experiment

- BEPCII is a symmetric  $e^+e^-$  collider with a peak luminosity of  $1 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$  at  $\sqrt{s} = 3.773 \text{ GeV}$
- Tracking: Helium-based multilayer drift chamber (MDC)
- PID: Plastic scintillator TOF system and  $\frac{dE}{dx}$
- Magnet: 1.0 T superconducting solenoid
- Neutral particle tracking: CsI(Tl) electromagnetic calorimeter (EMC)

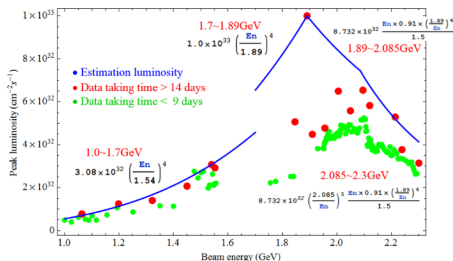


Overview of (left) BEPCII and (right) BESIII.

# The BESIII experiment

## Key datasets for charm physics:

- 2010-2011:  $2.9 \text{ fb}^{-1}$  at  $\psi(3770)$
- 2013-2019:  $7.3 \text{ fb}^{-1}$  of  $D_s \bar{D}_s^*$
- 2020:  $4.5 \text{ fb}^{-1}$  of  $\Lambda_c^+ \bar{\Lambda}_c^-$
- 2021-2022:  $5.0 \text{ fb}^{-1}$  at  $\psi(3770)$
- 2022-:  $\sim 8 \text{ fb}^{-1}$  at  $\psi(3770)$

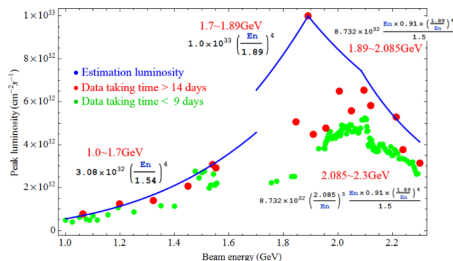


BEPCII peak luminosity.

# The BESIII experiment

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BEPCII peak luminosity.

Charm threshold data at  $\psi(3770) \rightarrow D \bar{D}$  provide a unique access to strong-phase information that is essential for charm mixing and  $\gamma$  measurements at  $B$  factories

# Recent charm results from BESIII

BESIII has a rich programme of charm physics:

① Strong-phase measurements

- Measurement of  $\delta_{K\pi}$  [EPJC 82 1009 \(2022\)](#)
- $D \rightarrow K^- \pi^+ \pi^- \pi^+$  strong-phase measurement [JHEP 5 \(2021\) 164](#)
- $D \rightarrow K^+ K^- \pi^+ \pi^-$   $F_+$  measurement [Phys. Rev. D 107 032009](#)

② Amplitude analysis

③ Semileptonic charm decays

④ Searches for rare decays

⑤ Branching fraction measurements

No time to cover all topics in this talk!  
I will mainly focus on strong-phase measurements in  
charm decays...

# Recent charm results from BESIII

## ① Strong-phase measurements

- $D \rightarrow K_S^0 \pi^+ \pi^- \pi^0$   $F_+$  measurement [arXiv:2305.03975](#)

## ② Amplitude analysis

- Amplitude analysis of  $D^0 \rightarrow K_L^0 \pi^+ \pi^-$  [arXiv:2212.0904](#)
- Observation of an  $a_0(980)$ -like state [Phys. Rev. Lett. 129, 182001](#)

## ③ Semileptonic charm decays

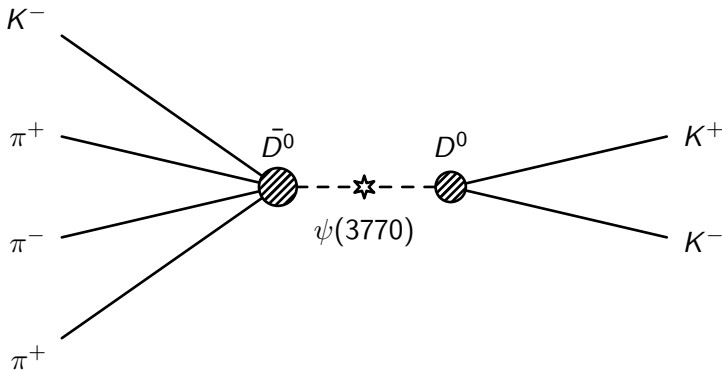
- First study of  $D_s^{*+} \rightarrow e^+ \nu_e$  [arXiv:2304.12159](#)
- $D_s^+ \rightarrow \tau^+ \nu_\tau$  [arXiv:2303.12600](#), [arXiv:2303.12600](#), [arXiv:2303.12468](#)
- Study of  $D_s^+ \rightarrow \pi^+ \pi^- e^+ \nu_e$  [arXiv:2303.12927](#)

## ④ Searches for rare decays

- See [Liang's talk](#)

... but here I provide references to some recent interesting charm results

# Double-tag analysis



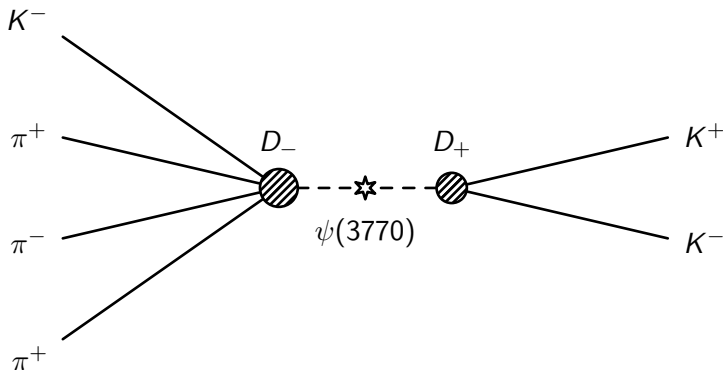
Double-tag method

The  $D$  mesons are produced in a quantum correlated state:

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|D^0\rangle|\bar{D}^0\rangle - |\bar{D}^0\rangle|D^0\rangle)$$



# Double-tag analysis



Double-tag method

Equivalently, we can consider the CP even (odd) eigenstates  $D_+$  ( $D_-$ ):

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|D_+\rangle|D_-\rangle - |D_-\rangle|D_+\rangle)$$

# Double-tag analysis

Double-tag analysis has many advantages:

- 1  $D\bar{D}$  pairs are quantum correlated, which provide direct access to the  $D^0\text{-}\bar{D}^0$  strong-phase difference
- 2 Measurements are, to first order, free from systematic uncertainties due to efficiencies and branching fractions
- 3 Full reconstruction ensures that the environment is extremely clean

Only one minor drawback:

- 1 Lower statistics

$$D \rightarrow K^- \pi^+$$

EPJC **82** 1009 (2022)

Improved measurement of the strong-phase difference  $\delta_D^{K\pi}$  in quantum-correlated  $D\bar{D}$  decays

What is measured:

- Strong-phase difference between CF and DCS  $D \rightarrow K^\mp \pi^\pm$  decays

Significance:

- Most precise measurement of  $\delta_D^{K\pi}$  in quantum-correlated  $D\bar{D}$  decays

$$D \rightarrow K^- \pi^+$$

The strong-phase difference  $\delta_D^{K\pi}$  between CF and DCS  $D \rightarrow K^- \pi^+$  is a key parameter in charm physics:

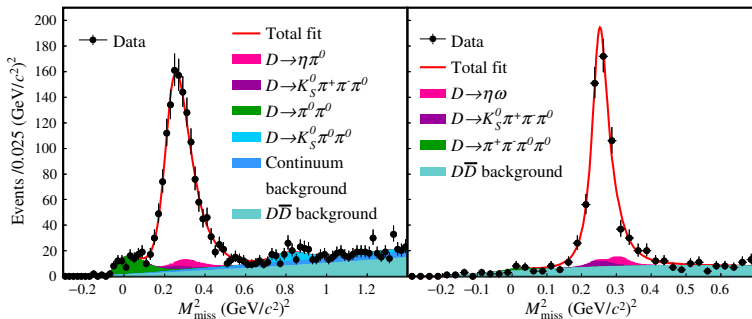
$$r_D^{K\pi} \exp(-i\delta_D^{K\pi}) = \frac{\langle K^+ \pi^- | D^0 \rangle}{\langle K^+ \pi^- | \bar{D}^0 \rangle}$$

Analysis is split into three main sections:

- 1 Determination of  $D \rightarrow K_L^0 X$  branching fractions
- 2 Measurement of the asymmetry  $\mathcal{A}_{K\pi}$  using  $CP$  tags
- 3 Measurement of  $r_D^{K\pi} \cos(\delta_D^{K\pi})$  and  $r_D^{K\pi} \sin(\delta_D^{K\pi})$  with  $K_{S,L}^0 \pi^+ \pi^-$  tags

$$D \rightarrow K^- \pi^+$$

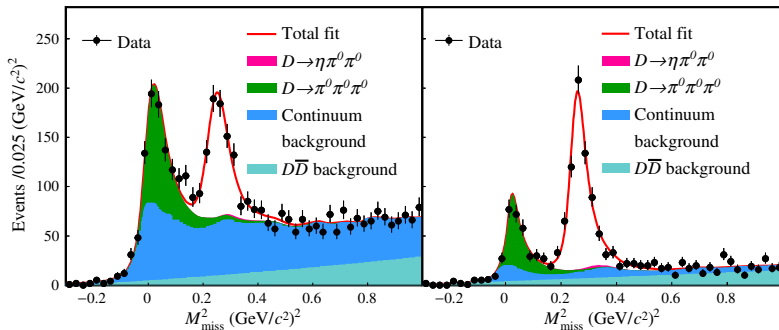
- Branching fractions must be determined independently of  $D \rightarrow K^- \pi^+ \Rightarrow$  Measure using pure  $CP$  tags
- $K_L^0 \pi^0$  and  $K_L^0 \omega$  are  $CP$ -even decays, and must therefore be tagged by  $CP$ -odd decays
  - $K_S^0 \pi^0$ ,  $K_S^0 \eta$ ,  $K_S^0 \eta' (\pi^+ \pi^- \eta, \pi^+ \pi^- \gamma)$ ,  $K_S^0 \omega$



Missing mass squared in  $CP$ -tagged (left)  $D \rightarrow K_L^0 \pi^0$  and (right)  $D \rightarrow K_L^0 \omega$ .

$$D \rightarrow K^- \pi^+$$

- Branching fractions must be determined independently of  $D \rightarrow K^- \pi^+ \Rightarrow$  Measure using pure  $CP$  tags
- Similarly,  $K_L^0 \pi^0 \pi^0$  is a  $CP$ -odd decay and is tagged by  $CP$ -even decays
  - $K^+ K^-$ ,  $\pi^+ \pi^-$ ,  $K_S^0 \pi^0 \pi^0$ ,  $\pi^+ \pi^- \pi^0$



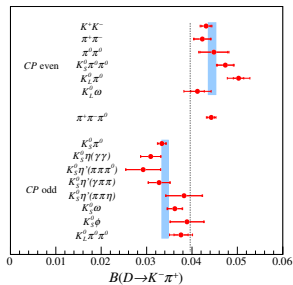
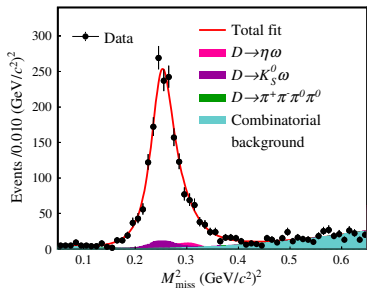
Missing mass squared in (left)  $\pi^+ \pi^- \pi^0$  tags and all other tags (left).

$$D \rightarrow K^- \pi^+$$

Asymmetry  $\mathcal{A}_{K\pi}$  of the branching fraction, tagged with  $CP$ -even and  $CP$ -odd decays, is sensitive to  $\delta_D^{K\pi}$ :

$$\mathcal{A}_{K\pi} = \frac{-2r_D^{K\pi} \cos(\delta_{K\pi}) + y}{1 + (r_D^{K\pi})^2}$$

Using external inputs for  $y$  and  $r_D^{K\pi}$ ,  $r_D^{K\pi} \cos(\delta_D^{K\pi})$  is calculated from  $\mathcal{A}_{K\pi}$



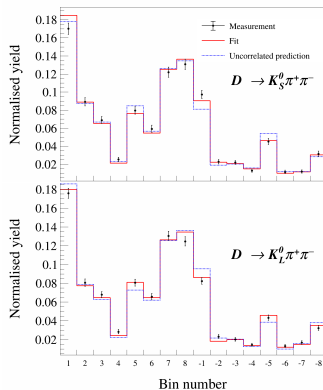
Left: Fit of the  $K^- \pi^+$  vs  $K_L^0 \omega$  double-tag yield. Right: Branching fraction of  $D \rightarrow K^- \pi^+$  measured using  $CP$  tags.

$$D \rightarrow K^- \pi^+$$

Double-tag yields using the  $K_{S,L}^0 \pi^+ \pi^-$  tags, in bins of phase space, are also sensitive to  $\delta_D^{K\pi}$ :

$$Y_i \propto \left( K_i + (r_D^{K\pi})^2 K_{-i} - 2r_D^{K\pi} \sqrt{K_i K_{-i}} \left[ c_i \cos(\delta_D^{K\pi}) - s_i \sin(\delta_D^{K\pi}) \right] \right)$$

- $\delta_D^{K\pi}$  is close to  $\pi \implies \sin(\delta_D^{K\pi})$  is much more sensitive to  $\delta_D^{K\pi}$
- Unique determination of  $\delta_D^{K\pi}$  without ambiguity
- External inputs for  $K_i$ ,  $c_i$  and  $s_i$  are recalculated without inputs from  $D \rightarrow K^- \pi^+$



Bin yields of  $D \rightarrow K_{S,L}^0 \pi^+ \pi^-$



$$D \rightarrow K^- \pi^+$$

Putting this all together:

$$\delta_D^{K\pi} = (187.6^{+8.9+5.4}_{-9.7-6.4})^\circ$$

Furthermore, the following branching fractions will be valuable additions to the PDG:

$$\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0) = (0.97 \pm 0.03 \pm 0.02) \times 10^{-2}$$

$$\mathcal{B}(D^0 \rightarrow K_L^0 \omega) = (1.09 \pm 0.06 \pm 0.03) \times 10^{-2}$$

$$\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0 \pi^0) = (1.26 \pm 0.05 \pm 0.03) \times 10^{-2}$$

$$D \rightarrow K^- \pi^+ \pi^- \pi^+$$

JHEP 5 (2021) 164

Measurement of the  $D \rightarrow K^- \pi^+ \pi^- \pi^-$  and  $D \rightarrow K^- \pi^+ \pi^0$  coherence factors and average strong-phase differences in quantum-correlated  $D\bar{D}$  decays

What is measured:

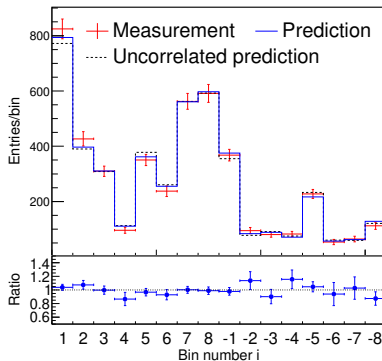
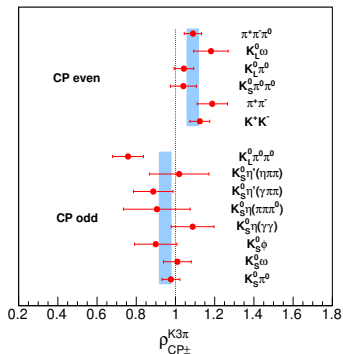
- Strong-phase difference and coherence factors between CF and DCS  $D \rightarrow K^\mp \pi^\pm \pi^\mp \pi^\pm$  decays in phase space bins

Significance:

- Crucial input to one of the most precise measurements of  $\gamma$

$$D \rightarrow K^- \pi^+ \pi^- \pi^+$$

The average strong-phase difference of multi-body decays may be determined analogously using  $CP$  and multi-body tags

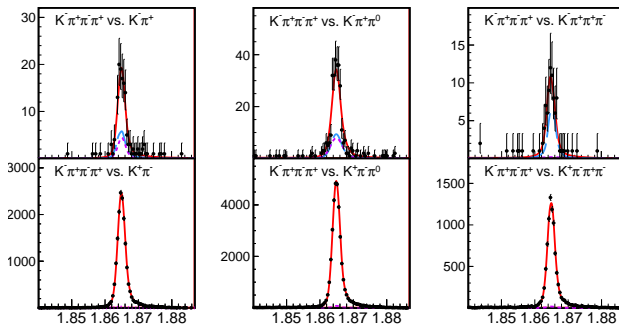


Left: Asymmetry of  $D \rightarrow K^- \pi^+ \pi^- \pi^+$  branching fraction with  $CP$ -even and  $CP$ -odd tags. Right: Double-tag yields of  $K_S^0 \pi^+ \pi^-$  vs  $K^- \pi^+ \pi^- \pi^+$ .

$$D \rightarrow K^- \pi^+ \pi^- \pi^+$$

However, averaging over phase space dilutes interference effects, which is parameterised in terms of the coherence factor:

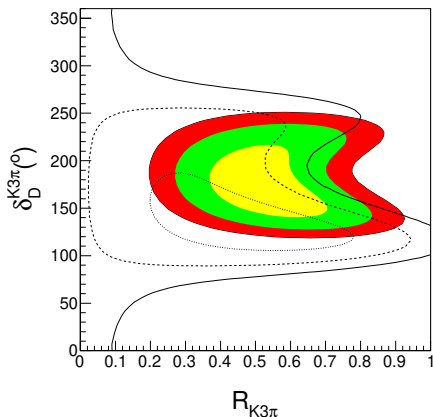
$$R_{K3\pi} \exp(i\delta_{K3\pi}) \equiv \frac{\int d\vec{x} \mathcal{A}_{\bar{D}^0 \rightarrow K^+ \pi^- \pi^+ \pi^-}(\vec{x}) \mathcal{A}_{D^0 \rightarrow K^+ \pi^- \pi^+ \pi^-}^*(\vec{x})}{A_{\bar{D}^0 \rightarrow K^+ \pi^- \pi^+ \pi^-} A_{D^0 \rightarrow K^+ \pi^- \pi^+ \pi^-}}$$



Like-sign yields are proportional to  $(1 - R^2)$  and brings sensitivity to  $R_{K3\pi}$ .

$$D \rightarrow K^- \pi^+ \pi^- \pi^+$$

Fit  $R_{K3\pi}$  and  $\delta_{K3\pi}$ : Huge improvements from CLEO-c analysis

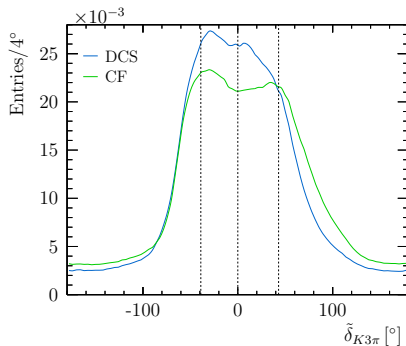


But:  $R_{K3\pi} = 0.52_{-0.10}^{+0.12}$ , so interference effects are very diluted

$$D \rightarrow K^- \pi^+ \pi^- \pi^+$$

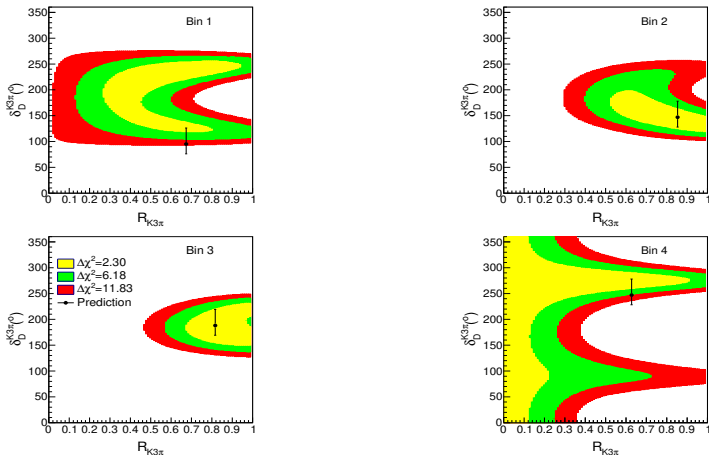
Can improve coherence factors by defining a binning scheme in terms of the normalised phase difference

$$\delta_{K3\pi}(\vec{x}) = \arg(\mathcal{A}_{\bar{D}^0}(\vec{x})\mathcal{A}_{D^0}^*(\vec{x})) - \arg\left(\int d\vec{x}' \mathcal{A}_{\bar{D}^0}(\vec{x}')\mathcal{A}_{D^0}^*(\vec{x}')\right)$$



$$D \rightarrow K^- \pi^+ \pi^- \pi^+$$

And indeed, the coherence factor is greatly improved in bin 2 and 3!



Binned fit of  $\delta_{K3\pi}$  and  $R_{K3\pi}$ .

$$D \rightarrow K^+ K^- \pi^+ \pi^-$$

Phys. Rev. D **107** 032009

Measurement of the  $CP$ -even fraction of  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

What is measured:

- Phase-space integrated strong-phase analysis of  $D \rightarrow K^+ K^- \pi^+ \pi^-$

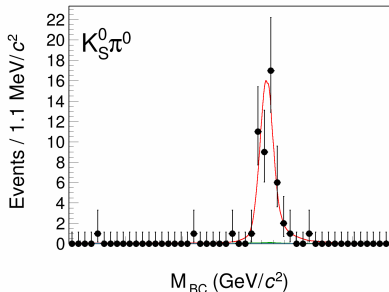
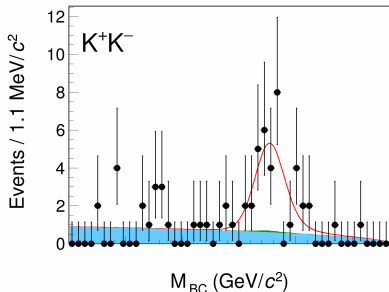
Significance:

- First model-independent study of the  $CP$  content of this decay



$$D \rightarrow K^+ K^- \pi^+ \pi^-$$

This decay suffers from both low branching fraction and low efficiency, as seen in the double-tag fits

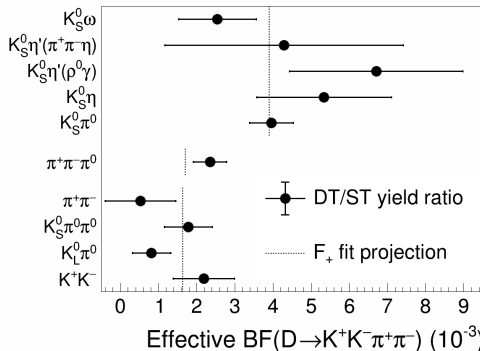


Double-tag fits of  $K^+ K^- \pi^+ \pi^-$  vs (left)  $K^+ K^-$  and (right)  $K_S^0 \pi^0$ .

Clear quantum correlation: The yield of  $K^+ K^- \pi^+ \pi^-$  vs  $K^+ K^-$  ( $CP$ -even) is suppressed, while that of  $K_S^0 \pi^0$  ( $CP$ -odd) is enhanced

$$D \rightarrow K^+ K^- \pi^+ \pi^-$$

The strong-phase information is parameterised in terms of the  $CP$ -even fraction  $F_+$ :

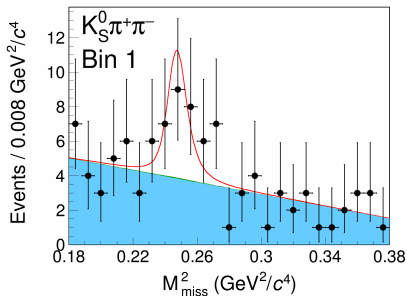
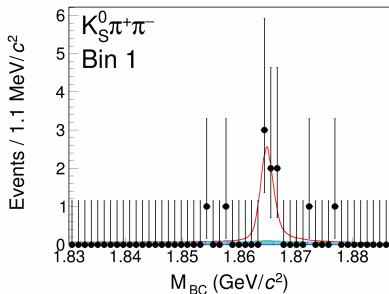
$$\frac{N_{DT}}{N_{ST}} \propto 1 \mp (2F_+ - 1)$$


Branching fractions of  $D \rightarrow K^+ K^- \pi^+ \pi^-$  measured against  $CP$ -even/odd tags.

Clearly,  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  is predominantly  $CP$ -even

$$D \rightarrow K^+ K^- \pi^+ \pi^-$$

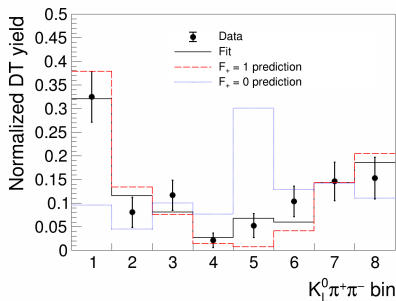
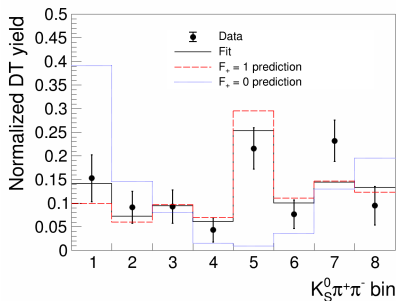
- To mitigate the low reconstruction efficiency, explore a novel partially reconstructed technique
- Poor low momentum kaon tracking efficiency  $\Rightarrow$   
Only reconstruct one kaon



Fits of (left) fully and (right) partially reconstructed  $K^+ K^- \pi^+ \pi^-$  vs  $K_S^0 \pi^+ \pi^-$ .

$$D \rightarrow K^+ K^- \pi^+ \pi^-$$

The multi-body tags  $K_{S,L}^0 \pi^+ \pi^-$ , which are split into phase-space bins, provide sensitivity to  $F_+$



Bin yields of  $K^+ K^- \pi^+ \pi^-$  vs  $K_{S,L}^0 \pi^+ \pi^-$ .

$$D \rightarrow K^+ K^- \pi^+ \pi^-$$

Combining the  $CP$  and  $K_{S,L}^0 \pi^+ \pi^-$  tags:

$$F_+ = 0.730 \pm 0.037 \pm 0.021$$

This is not the end of the story:

- ① A phase-space binned analysis is ongoing
- ② The statistical precision will improve greatly with more data
- ③ This fully charged mode will provide valuable inputs to  $\gamma$  and charm mixing at LHCb and Belle II

# Summary and conclusion

- 1 Many exciting measurements using quantum-correlated  $D\bar{D}$  pairs have been performed, using the  $3\text{ fb}^{-1}$  dataset
- 2 Previous analyses, such as  $\delta_{K\pi}$ , have been improved using more tags and more precise inputs
- 3 BESIII is now exploring four-body decays, and many binned analyses are in the pipeline using the larger  $8\text{ fb}^{-1}$  dataset
- 4 BESIII is expected to collect  $20\text{ fb}^{-1}$  of  $\psi(3770)$  data by 2024
  - Strong-phase measurements, which are currently statistics limited, will improve significantly in the next few years
  - This unique dataset will be essential for providing inputs to the foreseen datasets at LHCb and Belle II

Thanks for listening!