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## Theory overview: high-pt searches for flavor

## Plan of this talk



Assuming that well-known B anomalies are indeed due to new physics, I will discuss possible implications of these deviations in high-p searches @ LHC

## B anomalies in a nutshell



Both sets of $B$ anomalies challenge assumption of lepton flavor universality (LFU), which is usually taken for granted in high-energy physics

## B anomalies in a nutshell



Mass/scale suppression of effective operators suggests that explanations of $b \rightarrow c$ anomalies should lead to testable high- $p$ s signatures, while $b \rightarrow s$ case looks much less promising

## Simplified models for B anomalies

$$
\lambda_{i j}^{q} \lambda_{\alpha \beta}^{l}\left(C_{T}\left(\bar{Q}_{L}^{i} \gamma_{\mu} \sigma^{a} Q_{L}^{j}\right)\left(\bar{L}_{L}^{\alpha} \gamma^{\mu} \sigma^{a} L_{L}^{\beta}\right)+C_{S}\left(\bar{Q}_{L}^{i} \gamma_{\mu} Q_{L}^{j}\right)\left(\bar{L}_{L}^{\alpha} \gamma^{\mu} L_{L}^{\beta}\right)\right)
$$



| Model | Mediator | $b \rightarrow s$ | $b \rightarrow c$ |
| :---: | :---: | :---: | :---: |
| Colorless vectors | $B^{\prime}=(1,1,0)$ | $\checkmark$ | $\mathbf{X}$ |
|  | $W^{\prime}=(1,3,0)$ | $X$ | $\checkmark$ |
| Scalar leptoquarks | $S_{1}=(\overline{3}, 1,1 / 3)$ | $X$ | $\checkmark$ |
|  | $S_{3}=(\overline{3}, 3,1 / 3)$ | $\checkmark$ | $\mathbf{X}$ |
| Vector leptoquarks | $U_{1}=(3,1,2 / 3)$ | $\checkmark$ | $\checkmark$ |
|  | $U_{3}=(3,3,2 / 3)$ | $\checkmark$ | $\mathbf{X}$ |

$b \rightarrow s(b \rightarrow c)$ anomalies alone can be accommodated by several simple single-mediator models

## Simplified models for B anomalies

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\lambda_{i j}^{q} \lambda_{\alpha \beta}^{l}\left(C_{T}\left(\bar{Q}_{L}^{i} \gamma_{\mu} \sigma^{a} Q_{L}^{j}\right)\left(\bar{L}_{L}^{\alpha} \gamma^{\mu} \sigma^{a} L_{L}^{\beta}\right)+C_{S}\left(\bar{Q}_{L}^{i} \gamma_{\mu} Q_{L}^{j}\right)\left(\bar{L}_{L}^{\alpha} \gamma^{\mu} L_{L}^{\beta}\right)\right)
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| Scalar leptoquarks | $S_{1}=(\overline{3}, 1,1 / 3)$ | $\mathbf{X}$ | $\checkmark$ |
|  | $S_{3}=(\overline{3}, 3,1 / 3)$ | $\checkmark$ | $\mathbf{X}$ |
| Vector leptoquarks | $U_{1}=(3,1,2 / 3)$ | $\checkmark$ | $\checkmark$ |
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$\mathrm{U}_{1}$ singlet vector leptoquark (LQ) is only single-mediator model that can explain both sets of anomalies

## A digression on LQs



Both scalar \& vector LQ have important advantage with respect to other tree-level mediators that they do not induce tree-level contributions to $B$ mixing \& $\tau \rightarrow \mu v v$

## Well-known LQ search strategies @ LHC


pair production

t-channel Drell-Yan (DY)

## Photon \& lepton content of proton



## Resonant LQ production @ LHC



Non-zero lepton parton distribution functions allow for resonant LQ production @ LHC, but single lepton-jet final states are not part of exotics search canon of ATLAS \& CMS

## Dilepton searches @ ATLAS: 13*

## Dijet searches @ ATLAS: 12*

## Lepton-jet final state searches @ ATLAS: 1*



## LHC limits on $1^{\text {st }} \& 2^{\text {nd }}$ generation LQs


pair production (PP)

## Singlet vector LQ models for B anomalies

$$
\begin{aligned}
& \mathcal{L} \supset \frac{g_{U}}{\sqrt{2}}\left[\beta_{L}^{i j} \bar{Q}_{L}^{i, a} \gamma_{\mu} L_{L}^{j}+\beta_{R}^{i j} \bar{d}_{R}^{i, a} \gamma_{\mu} \mu_{R}^{j}\right] U^{\mu, a}+\text { h.c. }, \\
& \left|\beta_{L}^{22}\right| \lesssim\left|\beta_{L}^{32}\right| \ll\left(\beta_{L}^{23}|\lesssim| \beta_{L}^{33} \mid=\mathcal{O}(1)\right.
\end{aligned}
$$

| Parameters |  | Branching ratios |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\beta_{L}^{33}$ | $\beta_{L}^{23}$ | $\mathrm{BR}\left(U \rightarrow b \tau^{+}\right)$ | $\mathrm{BR}\left(U \rightarrow t \bar{\nu}_{\tau}\right)$ | $\mathrm{BR}\left(U \rightarrow s \tau^{+}\right)$ | $\mathrm{BR}\left(U \rightarrow c \bar{\nu}_{\tau}\right)$ |
| 1 | 0 | $51 \%$ | $49 \%$ | $0 \%$ | $0 \%$ |
| 1 | 1 | $25 \%$ | $22 \%$ | $25 \%$ | $27 \%$ |

## Possible singlet vector LQ signatures



Flavor structure as suggested by $b \rightarrow c$ anomalies singles out $p p \rightarrow \tau^{+} \tau^{-}$, bt as most interesting channels $-\mathrm{pp} \rightarrow \tau \mu, \mathrm{\tau v}, \mu^{+} \mu^{-}, \mathrm{tv}, \mathrm{cv}$ may be important as well in case of discovery or if $b \rightarrow c$ anomalies disappear

## Ditau searches @ LHC Run II




Three different ditau LHC Run II analyses, all considering events without \& with an extra b-jet

## Ditau searches @ LHC Run II




ATLAS data agrees with background predictions but both CMS analyses see a $3 \sigma$ excess

## Ditau searches @ LHC Run II




Non-resonant (resonant) excess in b-tag (b-veto) sample fits (does not fit) LQ explanation

## ATLAS ditau limits on singlet vector LQs



## ATLAS ditau limits on singlet vector LQs



## HL-LHC projections for singlet vector LQs



## HL-LHC projections for singlet vector LQs

weaker but complementary information provided by searches for resonant 3rd_ generation LQ signatures


95\% CL region favoured by b c anomalies
probably all singlet vector LQ explanations of $b \rightarrow c$ anomalies can be tested via ditau searches @ HL-LHC

## Beyond simplified LQ models



Ultraviolet complete LQ models typically contain new degrees of freedom besides LQ such as a heavy gluon $\mathrm{G}^{\prime}$, a $Z^{\prime}$, vector-like leptons (VLLs) L, additional Higgses, etc. New states cannot be arbitrarily heavy in models that address $b \rightarrow c$ anomalies

## Bounds on $\mathbf{G}^{\prime}$ motivated by b $\boldsymbol{\rightarrow} \mathbf{c}$ anomalies



## VLLs in gauged vector LQ models




Curbing LQ contributions to $B_{s}$ mixing requires VLLs with mass $M\llcorner$ not far from 1 TeV

## VLLs in gauged vector LQ models



VLL production in context of gauged vector LQ models addressing $b \rightarrow c$ anomalies is expected to lead to high-multiplicity final states with $\tau, b, t \& E_{T, \text { miss }}$

## VLLs searches triggered by B anomalies

| Tau multiplicity | VLL production <br> + decay mode | Final state |
| :---: | :---: | :---: |
| $0 \tau$ | $\mathrm{EE} \rightarrow \mathrm{b}\left(\mathrm{t} v_{\tau}\right) \mathrm{b}\left(\mathrm{t} \nu_{\tau}\right)$ | $4 \mathrm{~b}+4 \mathrm{j}+2 v_{\tau}$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}\left(\mathrm{t} v_{\tau}\right) \mathbf{t}\left(\mathrm{t} \nu_{\tau}\right)$ | $4 \mathrm{~b}+6 \mathrm{j}+2 \nu_{\tau}$ |
|  | $\mathrm{NN} \rightarrow \mathrm{t}\left(\mathrm{t} v_{\tau}\right) \mathrm{t}\left(\mathrm{t} v_{\tau}\right)$ | $4 \mathrm{~b}+8 \mathrm{j}+2 \nu_{\tau}$ |
| $1 \tau$ | $\mathrm{EE} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{b}\left(\mathrm{t} \nu_{\tau}\right)$ | $4 \mathrm{~b}+2 \mathrm{j}+\tau+v_{\tau}$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}\left(\mathrm{t} v_{\tau}\right) \mathrm{t}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+4 \mathrm{j}+\tau+\nu_{\tau}$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{t}\left(\mathrm{t} v_{\tau}\right)$ | $4 \mathrm{~b}+4 \mathrm{j}+\tau+v_{\tau}$ |
|  | $\mathrm{NN} \rightarrow \mathrm{t}(\mathrm{b} \tau) \mathrm{t}\left(\mathrm{t} v_{\tau}\right)$ | $4 \mathrm{~b}+6 \mathrm{j}+\tau+\nu_{\tau}$ |
| $2 \tau$ | $\mathrm{EE} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{b}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+2 \tau$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{t}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+2 \mathrm{j}+2 \tau$ |
|  | $\mathrm{NN} \rightarrow \mathrm{t}(\mathrm{b} \tau) \mathrm{t}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+4 \mathrm{j}+2 \tau$ |



Recently CMS performed first dedicated search for VLLs in gauged vector LQ model, exploring final states with at least three b-jets \& two 3rd-generation leptons

## VLLs searches triggered by B anomalies

| Tau multiplicity | VLL production + decay mode | Final state |
| :---: | :---: | :---: |
| $0 \tau$ | $\mathrm{EE} \rightarrow \mathrm{b}\left(\mathrm{t} v_{\tau}\right) \mathrm{b}\left(\mathrm{t} \nu_{\tau}\right)$ | $4 \mathrm{~b}+4 \mathrm{j}+2 \nu_{\tau}$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}\left(\mathrm{t} v_{\tau}\right) \mathrm{t}\left(\mathrm{t} v_{\tau}\right)$ | $4 \mathrm{~b}+6 \mathrm{j}+2 \nu_{\tau}$ |
|  | $\mathrm{NN} \rightarrow \mathrm{t}\left(\mathrm{t} \nu_{\tau}\right) \mathrm{t}\left(\mathrm{t} v_{\tau}\right)$ | $4 \mathrm{~b}+8 \mathrm{j}+2 \nu_{\tau}$ |
| $1 \tau$ | $\mathrm{EE} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{b}\left(\mathrm{t} v_{\tau}\right)$ | $4 \mathrm{~b}+2 \mathrm{j}+\tau+v_{\tau}$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}\left(\mathrm{t} v_{\tau}\right) \mathrm{t}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+4 \mathrm{j}+\tau+\nu_{\tau}$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{t}\left(\mathrm{t} v_{\tau}\right)$ | $4 \mathrm{~b}+4 \mathrm{j}+\tau+\nu_{\tau}$ |
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| $2 \tau$ | $\mathrm{EE} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{b}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+2 \tau$ |
|  | $\mathrm{EN} \rightarrow \mathrm{b}(\mathrm{b} \tau) \mathrm{t}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+2 \mathrm{j}+2 \tau$ |
|  | $\mathrm{NN} \rightarrow \mathrm{t}(\mathrm{b} \tau) \mathrm{t}(\mathrm{b} \tau)$ | $4 \mathrm{~b}+4 \mathrm{j}+2 \tau$ |



Expected limit on VLL mass of 650 GeV but CMS observes $2.8 \sigma$ excess for VLL mass hypothesis of 600 GeV \& as a result no VLL masses are excluded at 95\% CL

## $Z^{\prime}$ for $b \rightarrow s$ anomalies: $L_{\mu}-L_{\tau}$ models



b $\rightarrow$ s anomalies explained

excluded by
$B_{s}$ mixing

## $Z^{\prime}$ for $b \rightarrow s$ anomalies: $L_{\mu}-L_{\tau}$ models



## Dilepton searches in $L_{\mu}-L_{\tau}$ models



Z' couplings that follow minimal flavor violating (MFV) pattern excluded by dilepton searches

## Dilepton searches in $L_{\mu}-L_{\tau}$ models



Gauging $L_{\mu}-L_{\tau}$ gives gives $Z^{\prime}$ with vectorial couplings to $\mu, \tau \&$ corresponding v. Introduce vector-like quarks $Q$ to generate bsZ' coupling \& suppress $Z^{\prime}$ couplings to light quarks

## Searches for bs $\mu \mu$ contact interactions


[ATLAS-CONF-2021-012]


First search for bs $\mu \mu$ four-Fermi operator by ATLAS, but bounds on suppression scale are a factor of $\mathrm{O}(20)$ below sensitivity needed to test $\mathrm{b} \rightarrow \mathrm{s}$ anomalies model independently

## Testing LFU with dilepton events @ LHC



CMS observes good agreement with LFU up to masses of 1.5 TeV , but above 1.8 TeV there is slight excess in dielectron channel leading to a deviation of LFU ratio from 1

## Testing LFU with dilepton events @ LHC



CMS recently also measured difference between dimuon \& dielectron forwardbackward asymmetry ( $\mathrm{A}_{\text {FB }}$ ). Result is found to agree with zero within 2.4б. Like rate measurement, also AFB results show a slight dielectron excess

## Conclusions \& outlook

- Beyond SM models that explain all B-physics anomalies generically lead to signatures (e.g. pp $\rightarrow \tau^{+} \tau^{-}, b \tau, t_{t} \&$ high-multiplicity final states with $\left.\tau, b, t \& E_{T, \text { miss }}\right)$ testable @ LHC. If b $\rightarrow$ c anomalies persist, IMHO likely that LHC sees something
- BSM models that explain only $b \rightarrow s$ anomalies can be easily hidden from leaving imprint on high-pT LHC physics. Still, searches for bs $\mu \mu$ contact interactions, LFU violation in dilepton production, etc. may shed light on origin of anomalies
- Signals in Higgs \& diboson physics connected to anomalies possible (e.g. $h \rightarrow \tau \mu$ \& exotics decays of heavy Higgses) but model dependent - cf. backup for details


## Backup



## A digression on LFU

| Decay | Precision | Channels | Deviation |
| :---: | :---: | :---: | :---: |
| $Z$ | $0.3 \%$ | $e, \mu, \tau$ | - |
| $W$ | $0.8 \%$ | $e, \mu$ | - |
| $W$ | $3 \%$ | $\tau$ | $2.8 \sigma$ |
| $\mu, \tau$ | $0.15 \%$ | $e, \mu$ | - |
| $\pi$ | $0.3 \%$ | $e, \mu$ | - |
| $K$ | $0.4 \%$ | $e, \mu$ | - |
| $J / \psi$ | $0.65 \%$ | $e, \mu$ | - |
| $D_{s}$ | $6 \%$ | $\mu, \tau$ | - |

Before 2012, stringent experimental test of LFU in B-meson decays did not exist

Combined LEP results hint towards LFU violation in W-boson decay with significance of $2.8 \sigma$
[LEPEWWG, hep-ex/0511027]

## LFU violation in W decays?



ATLAS LHC Run II measurement in full agreement with LFU as predicted in SM

## Ditau limits on singlet vector LQs from CMS




## LHC bounds: $p p \rightarrow \pi$ vs. $p p \rightarrow \pi v$



## $Z^{\prime}$ bounds in singlet vector LQ model



$Z^{\prime}$ searches in general not competitive with limits obtained from LQ or G' searches

## Another LQ search triggered by B anomalies

[ATLAS, 2101.11582]



## Testing LFU with dilepton events @ LHC




## Flavorful 2HDM with right-handed neutrinos



Box diagrams with a charged Higgs boson \& a right-handed neutrino are able to generate LFU violating effects needed to explain $b \rightarrow s$ anomalies

## Flavorful 2HDM with right-handed neutrinos



In 2016 explanation of muon anomalous magnetic moment possible without violating $h \rightarrow \tau \mu$ bound if Higgs sector close to alignment. Now possibility even stronger constrained

## Flavorful 2HDM with right-handed neutrinos



LHC phenomenology of model not worked out, but exotic decays such as $\mathrm{H}, \mathrm{A} \rightarrow \mathrm{tc}(\tau \mu)$ \& $H^{ \pm} \rightarrow \mathrm{cb}$ generically expected \& wait for interest of community. Challenging searches but may reveal first direct evidence of beyond SM physics \& unravel origin of flavor

## Resonant LQ production @ the LHC



At 13 TeV LHC, 9 events per $100 \mathrm{fb}^{-1}$ for minimal scalar LQ of $\mathrm{M}=3 \mathrm{TeV} \& \lambda_{\mathrm{eu}}=1$

## Resonant LQ production @ the LHC



Suppressed by $\mathrm{E}_{\mathrm{T}, \text { miss }}$ requirement \& jet veto

## Resonant LQ production @ the LHC



Suppressed by $\mathrm{E}_{\mathrm{T}, \text { miss }}$ requirement \& jet veto


## Resonant LQ production @ the LHC



Irreducible background particularly relevant @ high invariant lepton-jet mass


## Resonant LQ production @ the LHC



## Resonant LQ production @ the LHC



Suppressed by $\mathrm{E}_{\mathrm{T}, \text { miss }}$ requirement


## Resonant LQ production @ the LHC



Sum over backgrounds is a steeply falling distribution, while signal exhibits a narrow peak

## LHC limits on $\mathbf{1 s t}^{\text {st }} \boldsymbol{2} \mathbf{2}^{\text {nd }}$ generation LQs



## LHC limits on $\mathbf{1 s t}^{\text {st }} \boldsymbol{2} \mathbf{2}^{\text {nd }}$ generation LQs



## LHC limits on $1^{\text {st }} \boldsymbol{\&} \mathbf{2}^{\text {nd }}$ generation LQs



single LQ production (SP)

## LHC limits on $\mathbf{1 s t}^{\text {st }} \mathbf{2}^{\text {nd }}$ generation LQs



resonant LQ production

## LHC limits on $\mathbf{1 s t}^{\text {st }} \mathbf{2}^{\text {nd }}$ generation LQs



resonant LQ production

## LHC limits on $\mathbf{1 s t}^{\text {st }} \mathbf{2}^{\text {nd }}$ generation LQs



## LQ contributions to $\mathbf{b}+\mathbf{\tau}$ signature



For $\beta_{L}^{23}=0, b+\tau$ signal arises only from $2 \rightarrow 2$ process, while for $\beta_{L}^{23} \neq 0$ also $2 \rightarrow 3$ scattering is relevant. Since two topologies lead to final states with very different kinematic features, it is essential to develop two separate search strategies for them

## Kinematic distributions of b+t signal

LHC $14 \mathrm{TeV}, \mathrm{b}+\tau$


LHC $14 \mathrm{TeV}, \mathrm{b}+\tau$


## Kinematic distributions of b+t signal




## Kinematic distributions of b+t signal




## Mono-top \& mono-jet distributions




## $b+\tau$ constraints from $2 \rightarrow 2 \& 2 \rightarrow 3$ signal



## Constraints from new LQ search strategies



## Constraints from new LQ search strategies



