

# Searching for the Stochastic Gravitational-Wave Background with Ground-Based Detectors

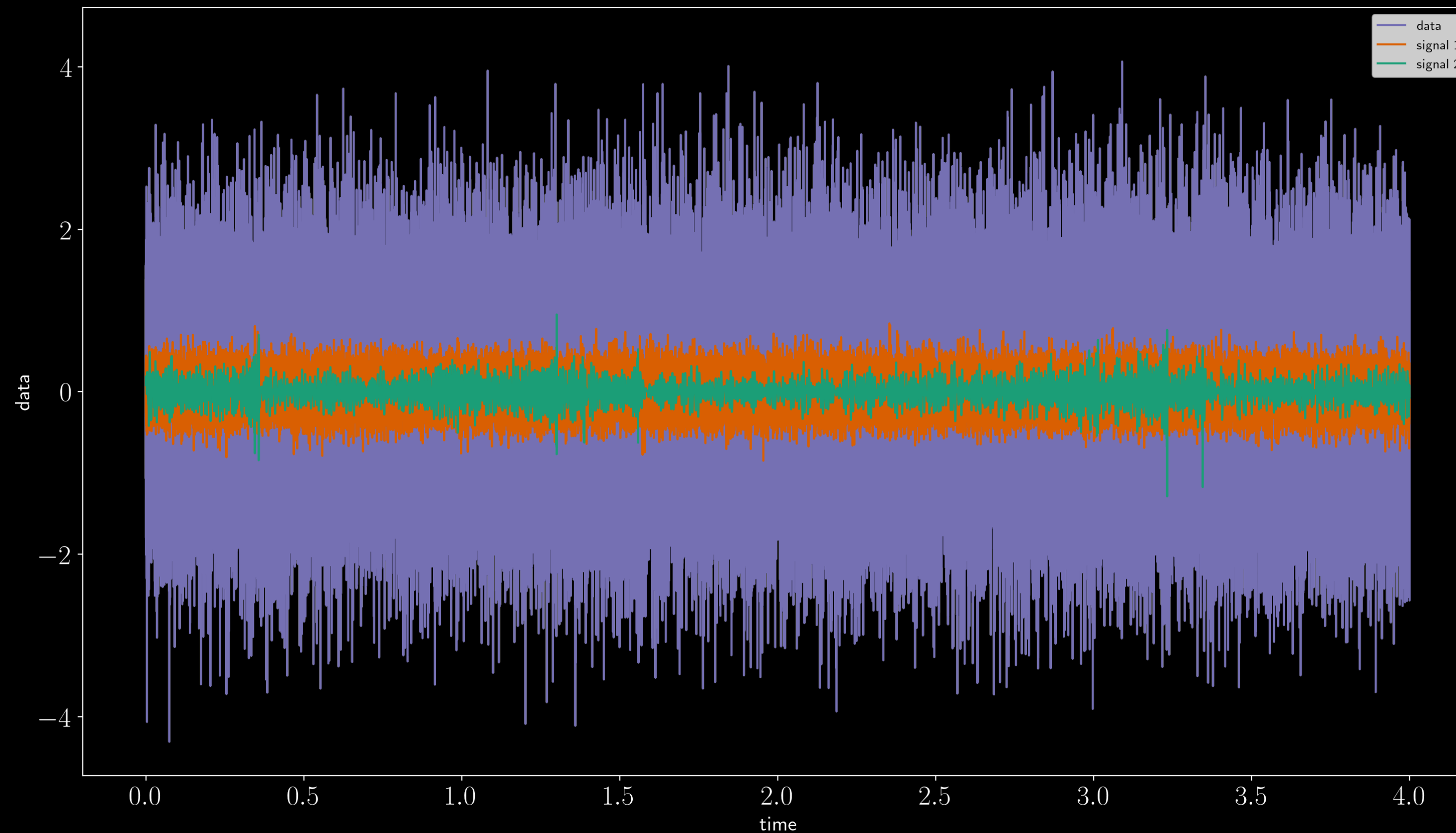
Jishnu Suresh  
Université catholique de Louvain

# STOCHASTIC GRAVITATIONAL WAVE BACKGROUND

Superposition of signals **too weak** or **too numerous** to individually detect

Looks **like noise** in a single detector

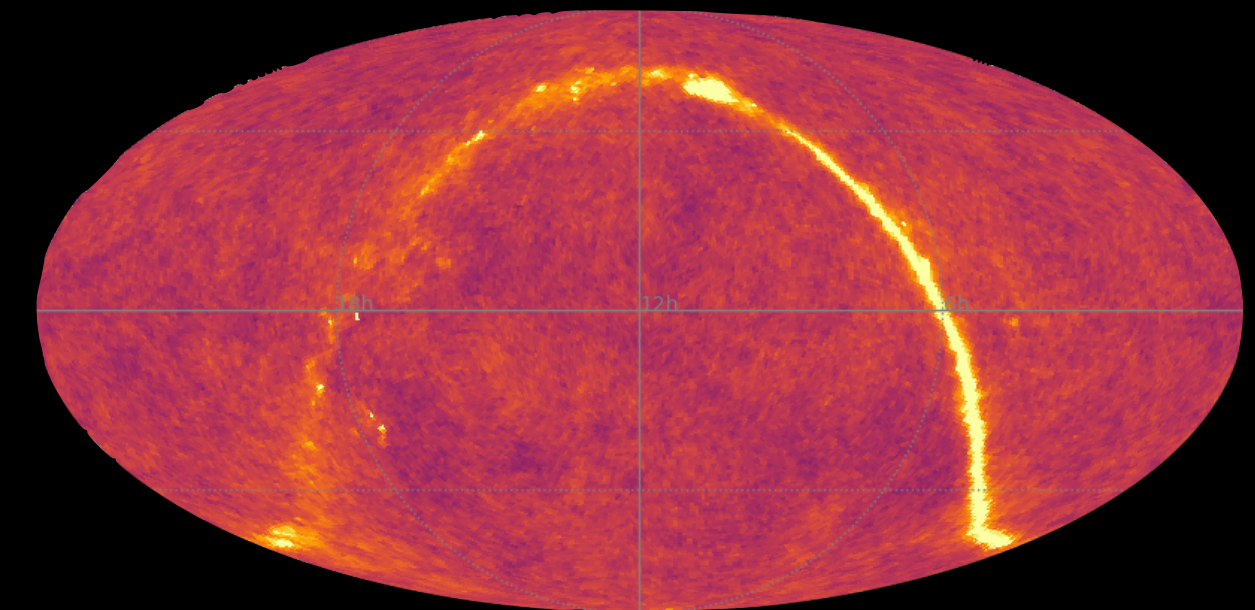
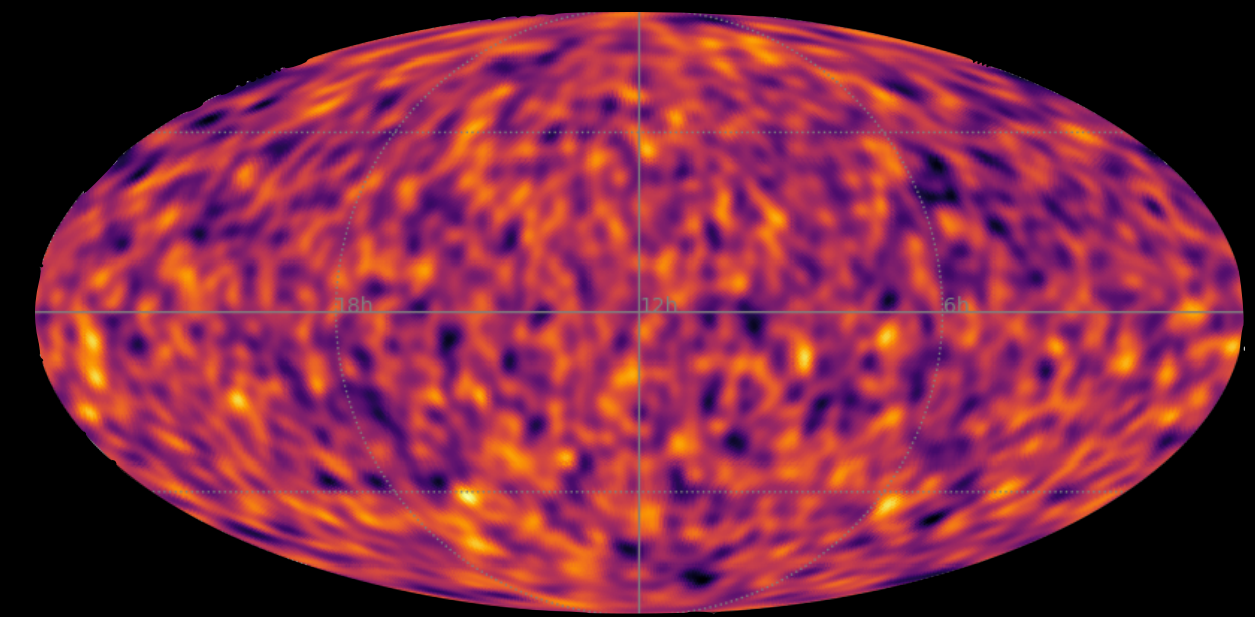
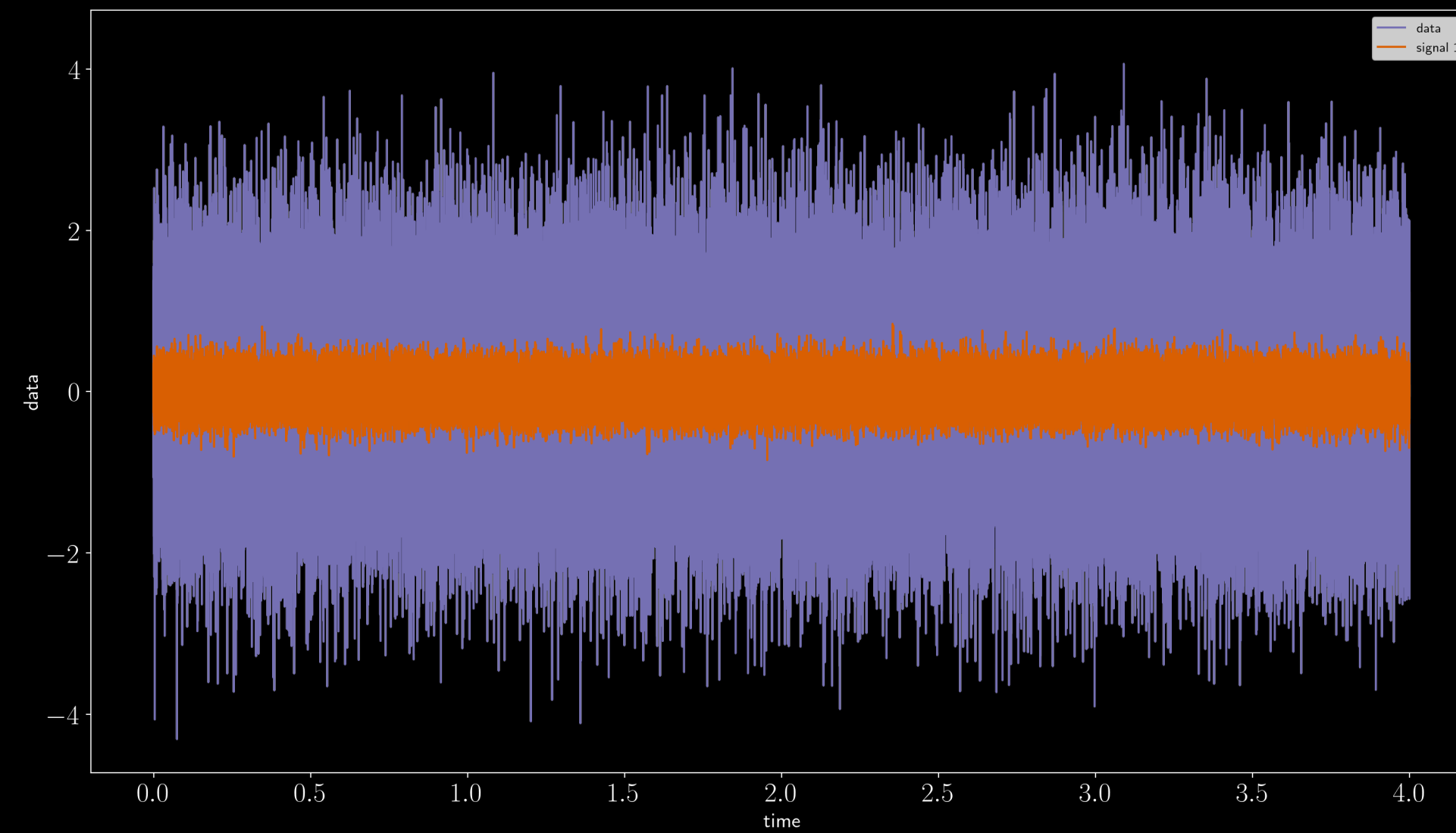
Characterized **statistically** in terms of moments (ensemble averages) of the metric perturbations



# STOCHASTIC GRAVITATIONAL WAVE BACKGROUND

In this talk, we will only consider the following cases

Unpolarized,



# WHAT DETECTION METHODS CAN WE USE?

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Problem: The stochastic signal looks more like noise in a single detector.

Solutions:

- Identify features that distinguish between the expected signal and noise.
  - Know our GW detector's noise sources well enough in amplitude and spectral shape.
- For multiple detectors having uncorrelated noise: cross-correlation separates the signal from the noise.

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Data from two detectors:

$$d_1 = h + n_1 \quad d_2 = h + n_2 \quad h \rightarrow \text{common GW signal component}$$

Cross-correlation:

$$\langle d_1 d_2 \rangle = \langle h^2 \rangle + \langle n_1 n_2 \rangle + \langle h n_2 \rangle + \langle n_1 h \rangle = \langle h^2 \rangle + \langle n_1 n_2 \rangle$$

*(Note: Blue arrows point from  $\langle h n_2 \rangle$  and  $\langle n_1 h \rangle$  to 0)*

Assuming detector noise is uncorrelated:

$$\langle d_1 d_2 \rangle = \langle h^2 \rangle + \langle n_1 n_2 \rangle$$

*(Note: Blue arrow points from  $\langle n_1 n_2 \rangle$  to 0)*

$$\langle d_1 d_2 \rangle = \langle h^2 \rangle \equiv S_h$$

**Cross-correlation separates the signal from the noise**

Intensity of the background

# OPTIMAL FILTERING

What is the optimal way to correlate data from two physically separated and misaligned detectors to search for a SGWB

Cross-correlation estimator  $\hat{S}_h \simeq \int_{-\infty}^{\infty} df \int_{-\infty}^{\infty} df' \delta_T(f-f') \tilde{d}_1(f) \tilde{d}_2^*(f') \tilde{Q}^*(f')$

Variance  $\sigma^2 \simeq \frac{T}{2} \int_0^{\infty} df P_1(f) P_2(f) |\tilde{Q}(f)|^2$

What we meant by optimal: Choose  $Q$  to maximize SNR for fixed spectral shape

$$\tilde{Q}(f) \propto \frac{\Gamma_{12}(f) \Omega_t(f)}{P_1(f) P_2(f)}$$

Overlap reduction function  $\leftarrow$

expected signal spectrum  $\leftarrow$

de-weight correlation when noise is large  $\leftarrow$

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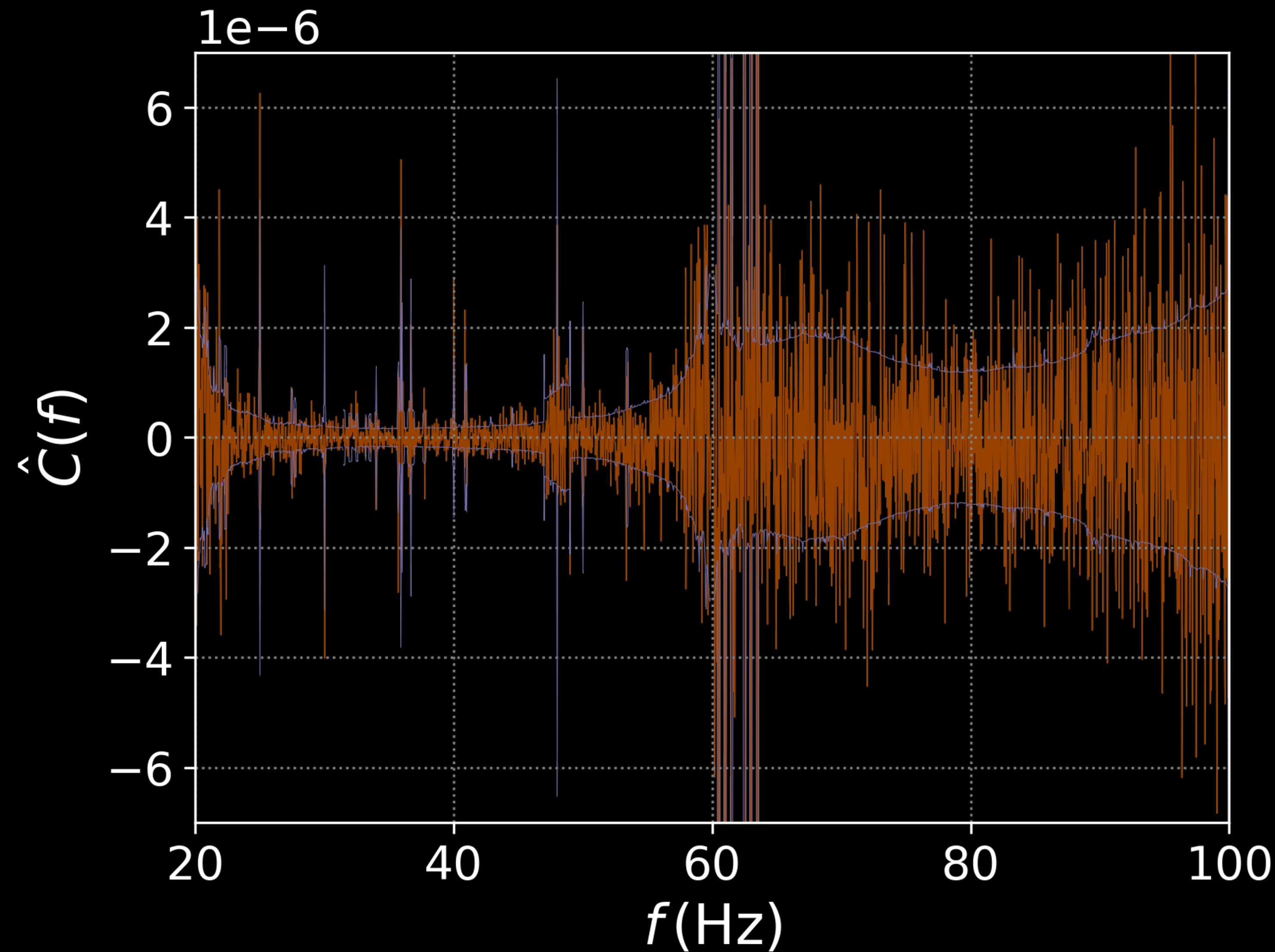
de-weight correlation when noise is large

We often choose a power-law functional form for the SGWB template spectrum

$$\Omega_t(f) = \Omega_{\text{ref}} \left( \frac{f}{f_{\text{ref}}} \right)^\alpha$$

## O1+O2+O3 RESULTS

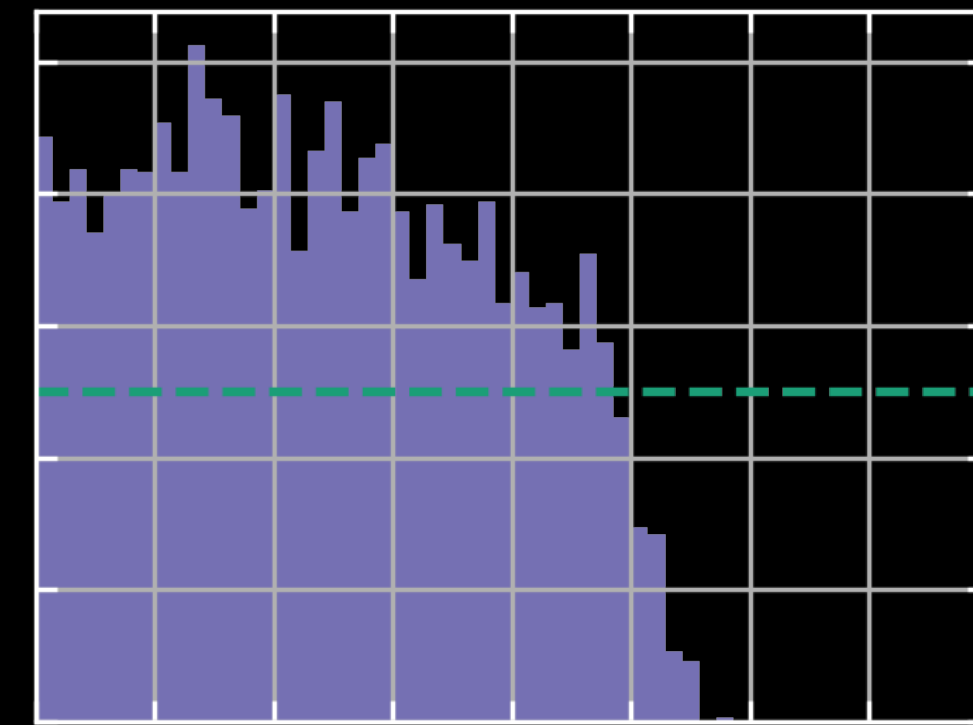
The observed cross-correlation spectra combining data from all three baselines in O3, as well as the HL baseline in O1 and O2. The spectrum is consistent with expectations from uncorrelated, Gaussian noise.



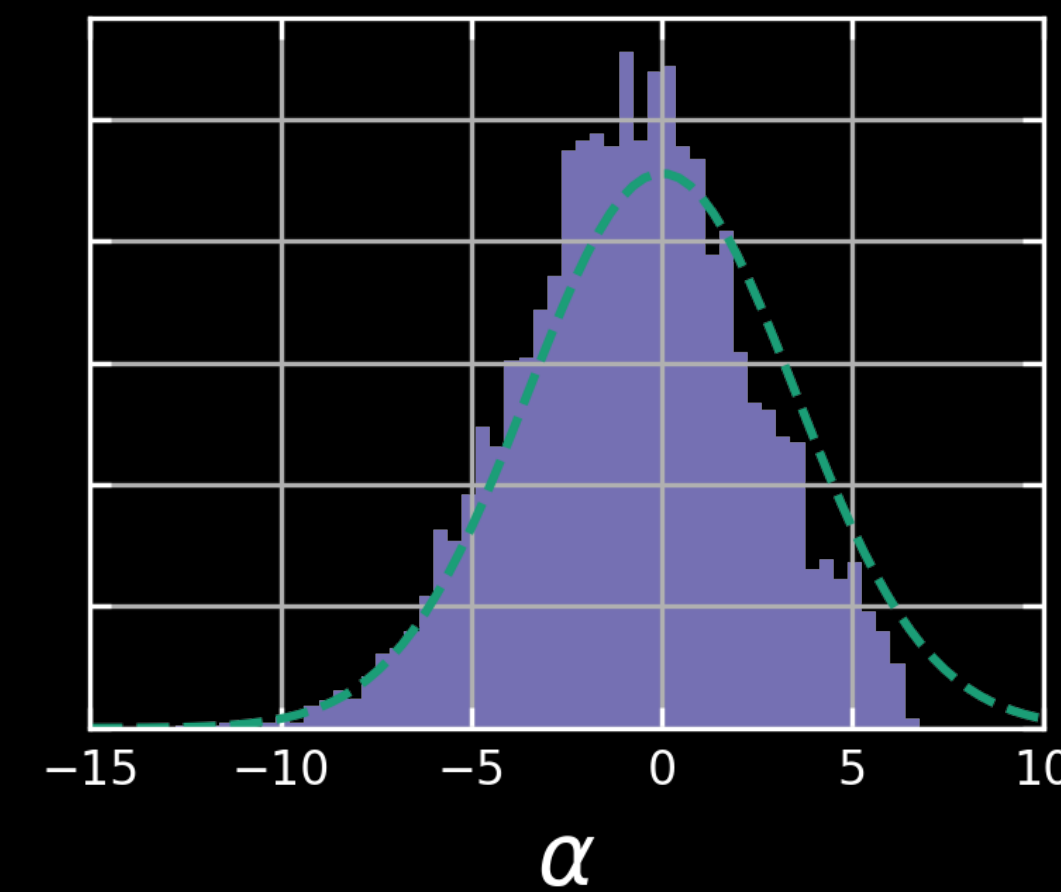
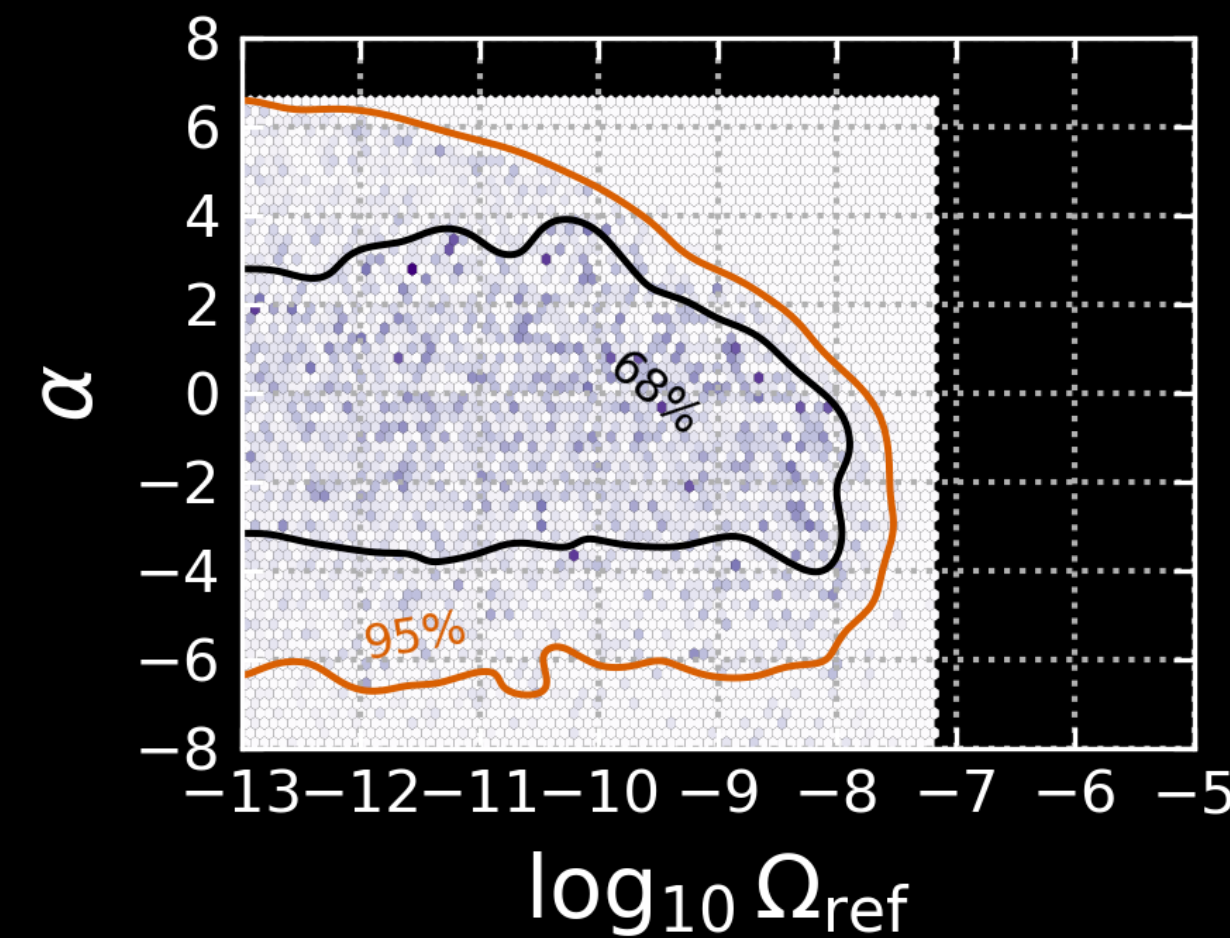


Since there was no evidence of an isotropic signal, we placed upper limits on  $\Omega_\alpha$  for different power-law indices  $\alpha$ .

	Uniform prior			Log-uniform prior		
$\alpha$	O3	O2	Improv.	O3	O2	Improv.
0	$1.7 \times 10^{-8}$	$6.0 \times 10^{-8}$	3.6	$5.8 \times 10^{-9}$	$3.5 \times 10^{-8}$	6.0
2/3	$1.7 \times 10^{-8}$	$4.8 \times 10^{-8}$	4.0	$3.4 \times 10^{-9}$	$3.0 \times 10^{-8}$	8.8
3	$1.3 \times 10^{-9}$	$7.9 \times 10^{-9}$	5.9	$3.9 \times 10^{-10}$	$5.1 \times 10^{-9}$	13.1



posteriors for  $\alpha$  and  $\Omega_{\text{ref}}$



# HOW DO WE MAP THE SGWB SKY?

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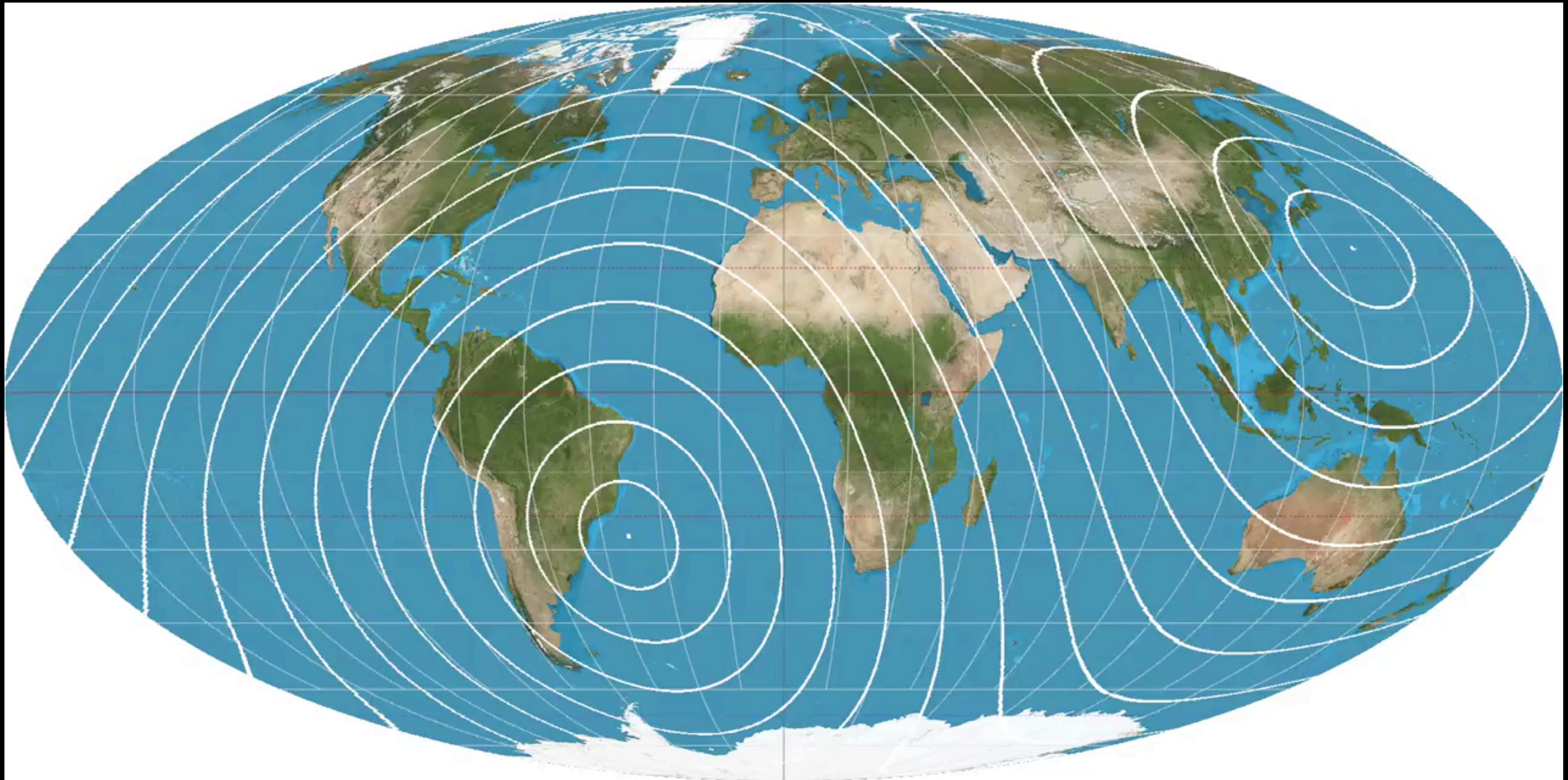
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# HOW DO WE MAP THE SGWB SKY?

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**Cross-correlation is essentially  
a one-dimensional map of the sky**

# HOW DO WE MAP THE SGWB SKY?

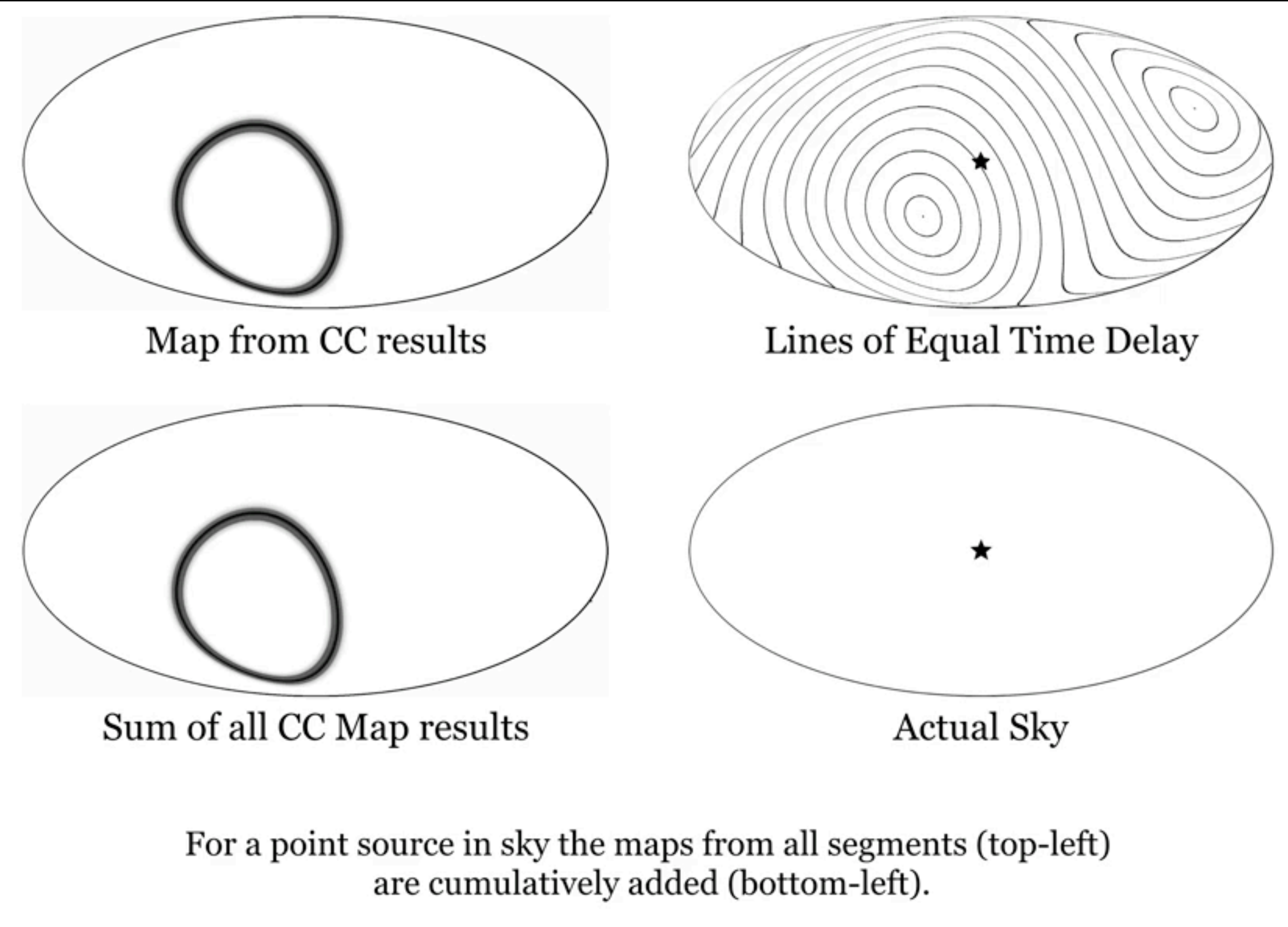


The white circles indicate positions in the sky map that will have equal time/phase delay when the signal from that part of the sky arrives in the LIGO Hanford and Livingston detectors.

# SGWB MAPPING

Cross-correlation is essentially a one-dimensional map of the sky.

When we consider the time delay between two detectors and the Rotation of the earth.



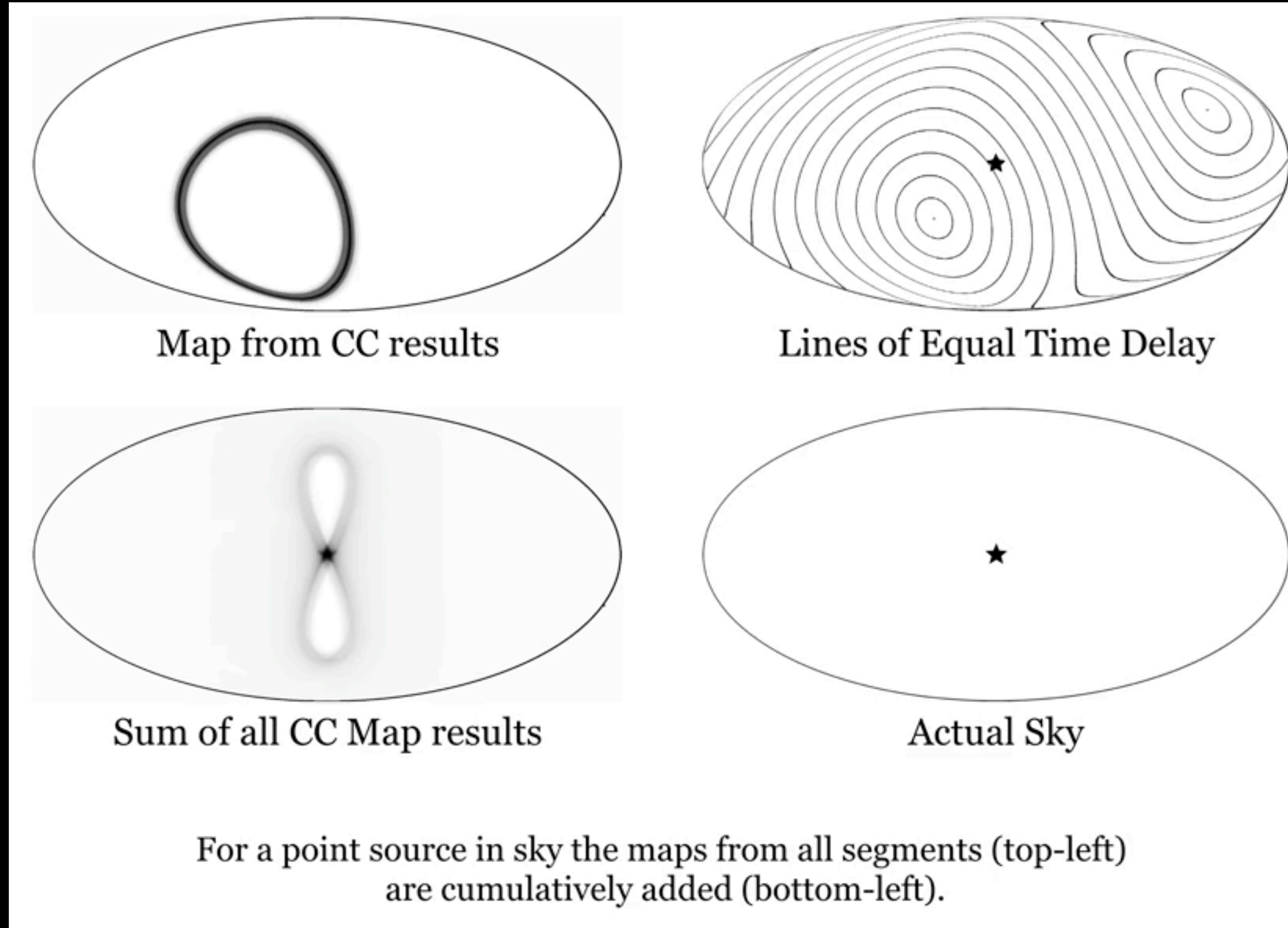
- LIGO Hanford-Livingston baseline.
- Assume that the detector has perfect sky coverage.
- All the processes are in the time domain.
- No detector noise.
- Strong monochromatic signal.

Animation Credit: A. Ain

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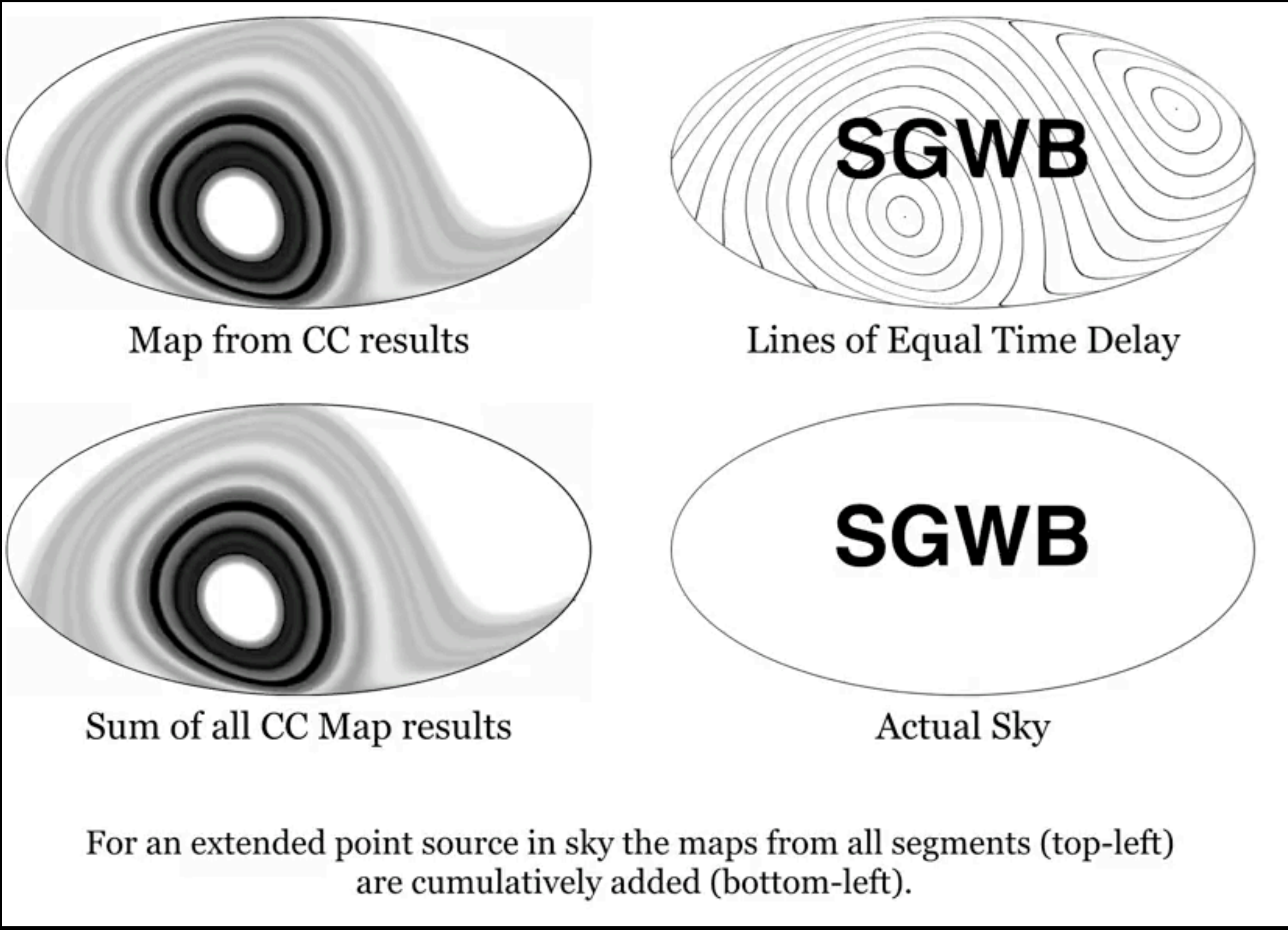
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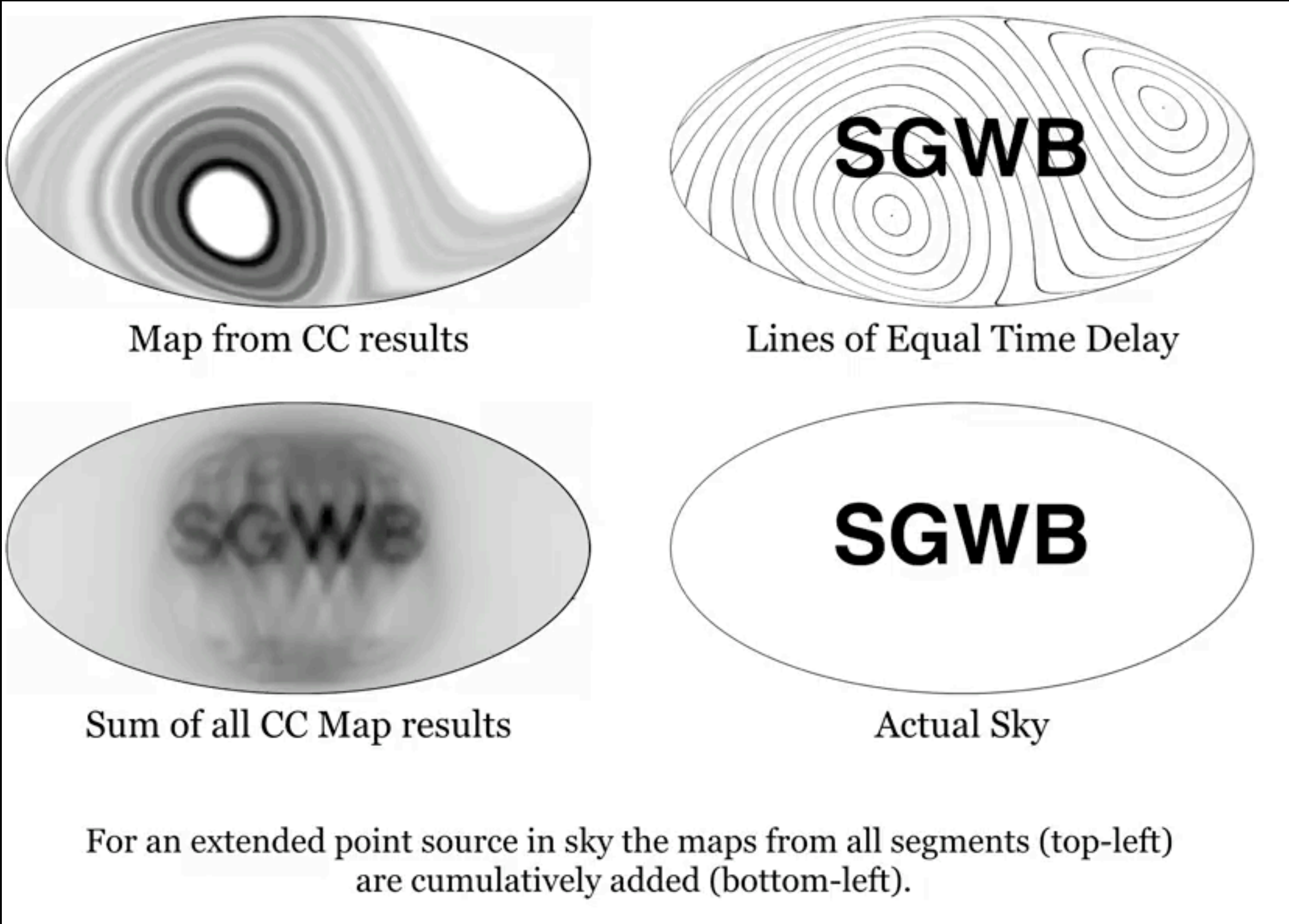
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# ANISOTROPIC SEARCH

Anisotropic search tries to measure the direction of the sky from where the signal comes.  
In this mapping process, we consider:

- The time delay between two detectors
- Rotation of the earth.

Recall: the SGWB energy density  $\Omega_{\text{GW}}(f, \hat{\mathbf{n}}) \equiv \frac{f}{\rho_c} \frac{d\rho_{\text{GW}}}{df} = \frac{2\pi^2}{3H_0^2} f^3 \mathcal{P}(f, \hat{\mathbf{n}})$

Cross-correlation is essentially a one-dimensional map of the sky.

Anisotropy can be expanded in pixel or spherical harmonic basis

# HOW DO WE MAP THE SGWB SKY?

The anisotropy of the SGWB can be characterized using the dimensional energy density parameter

$$\Omega_{\text{GW}}(f, \hat{\mathbf{n}}) \equiv \frac{f}{\rho_c} \frac{d\rho_{\text{GW}}}{df} = \frac{2\pi^2}{3H_0^2} f^3 \mathcal{P}(f, \hat{\mathbf{n}})$$

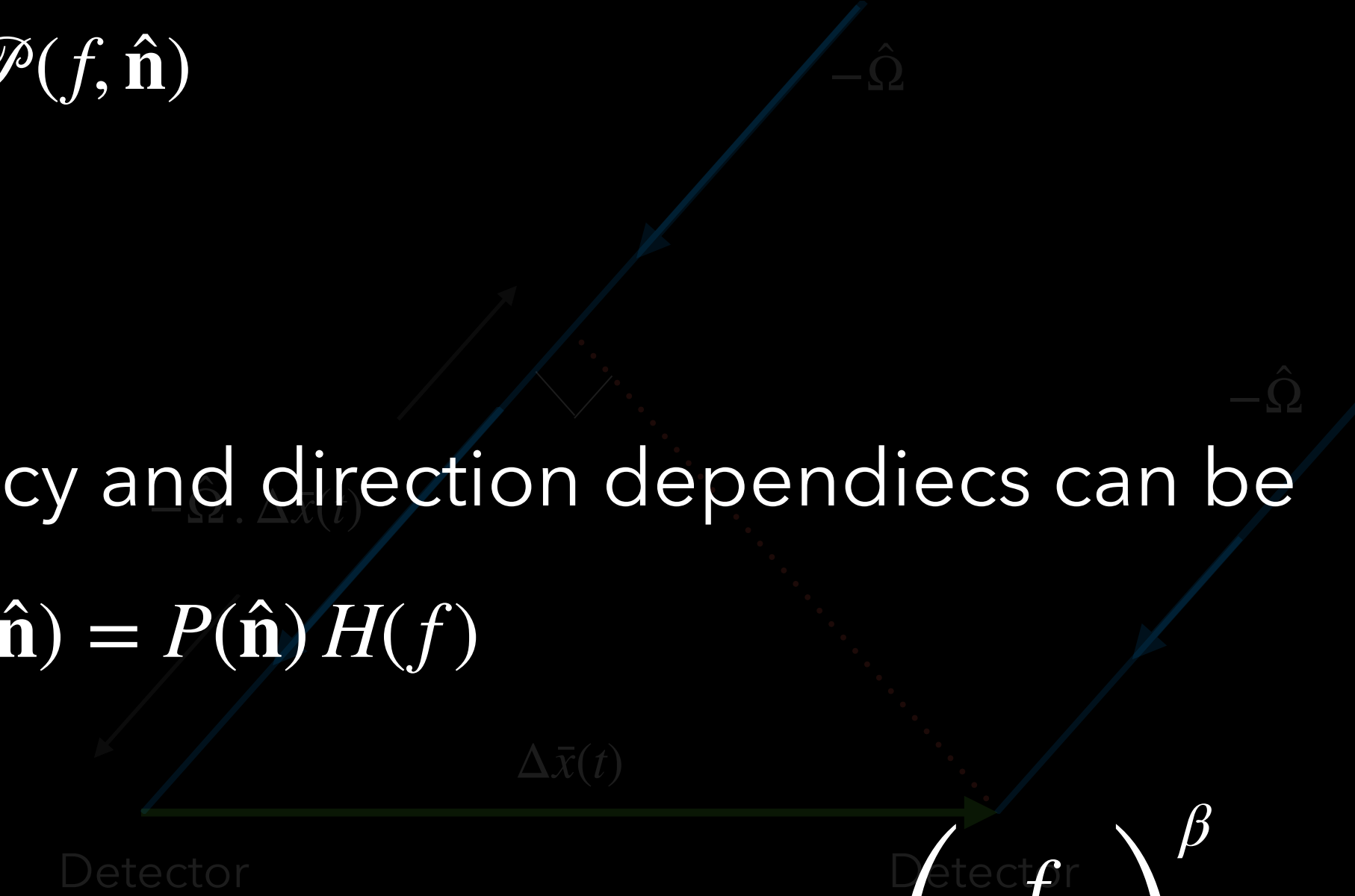
- Essentially Earth Rotation Synthesis Imaging

Most of the analysis performed so far assumes that the frequency and direction dependencies can be separated:

$$\mathcal{P}(f, \hat{\mathbf{n}}) = P(\hat{\mathbf{n}}) H(f)$$

- map making: use time-dependent phase delay
- Use spectral filters
  - to enhance signal power
  - to reduce noise power

Where the common choice of spectral shape is  $H(f) = \left(\frac{f}{f_{\text{ref}}}\right)^\beta$



Here  $\Delta \vec{x}(t)$  is the separation or baseline vector between the two detectors; as the Earth rotates, its direction changes, but its magnitude remains fixed. The direction to the source  $\hat{\Omega}$  is also fixed in the barycentric frame. The phase difference between signals arriving at two detector sites from the same direction is also shown.

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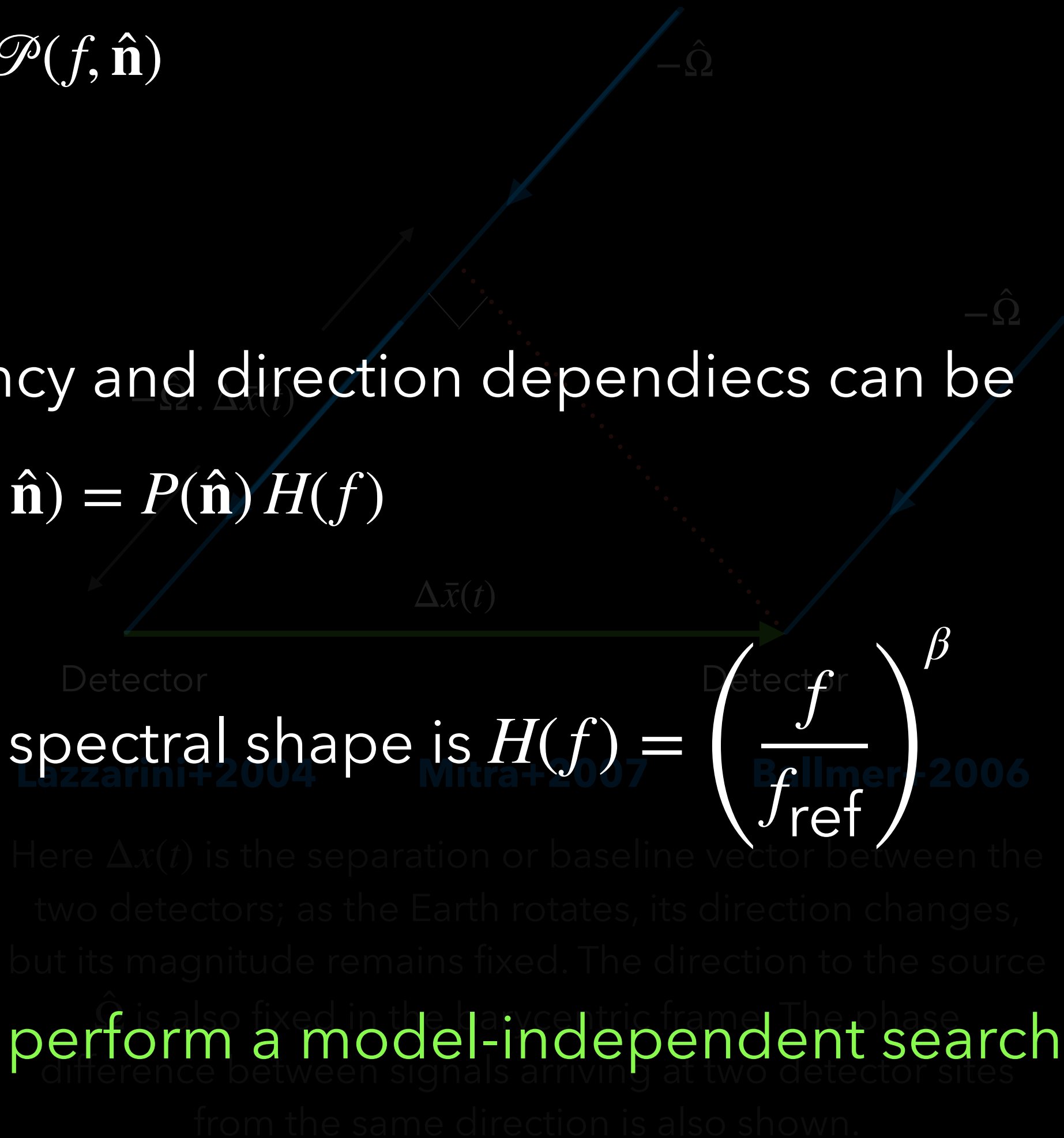
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We will perform a model-independent search

PyStoch : fast HEALPix based SGWB mapmaking

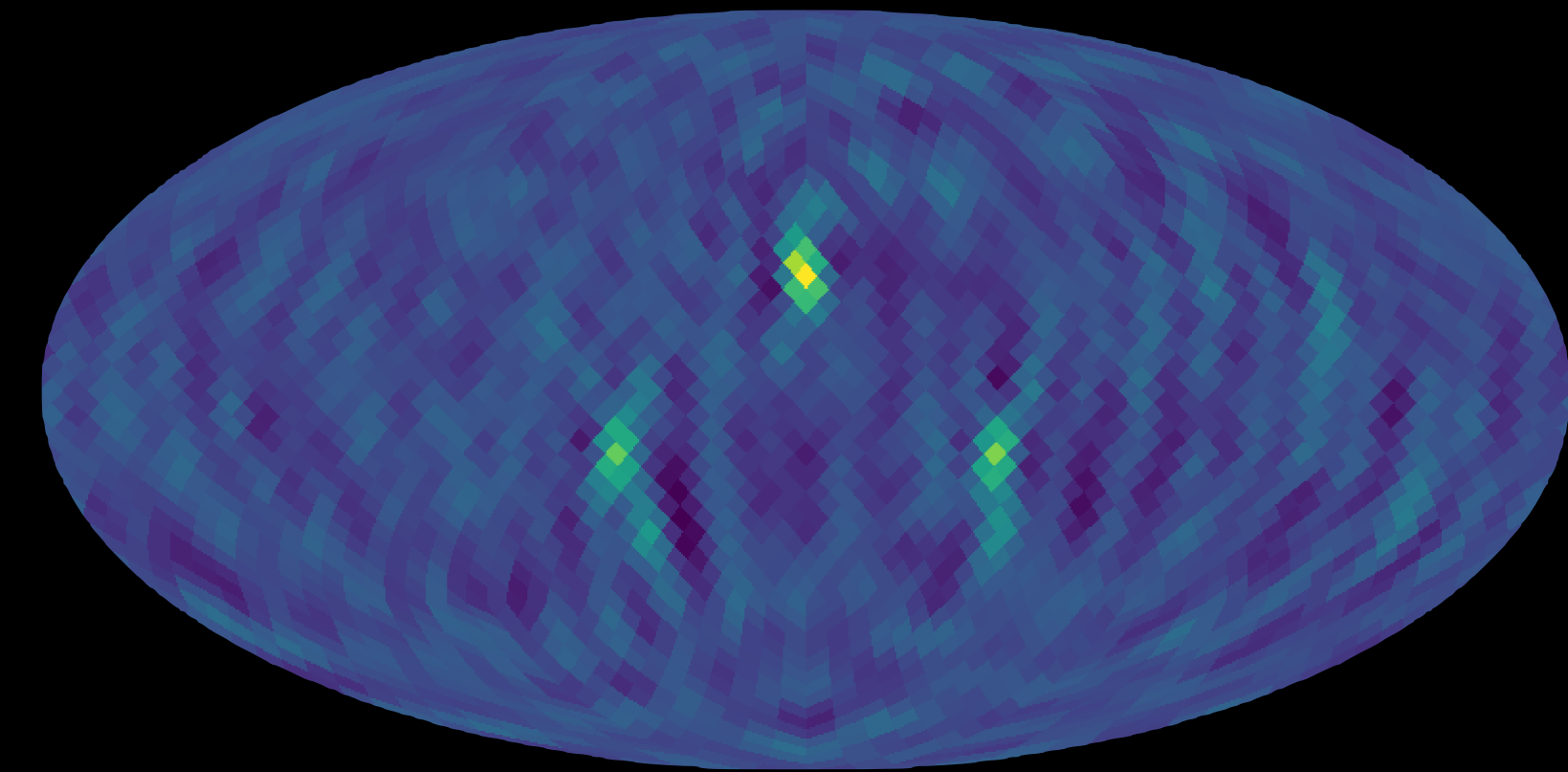
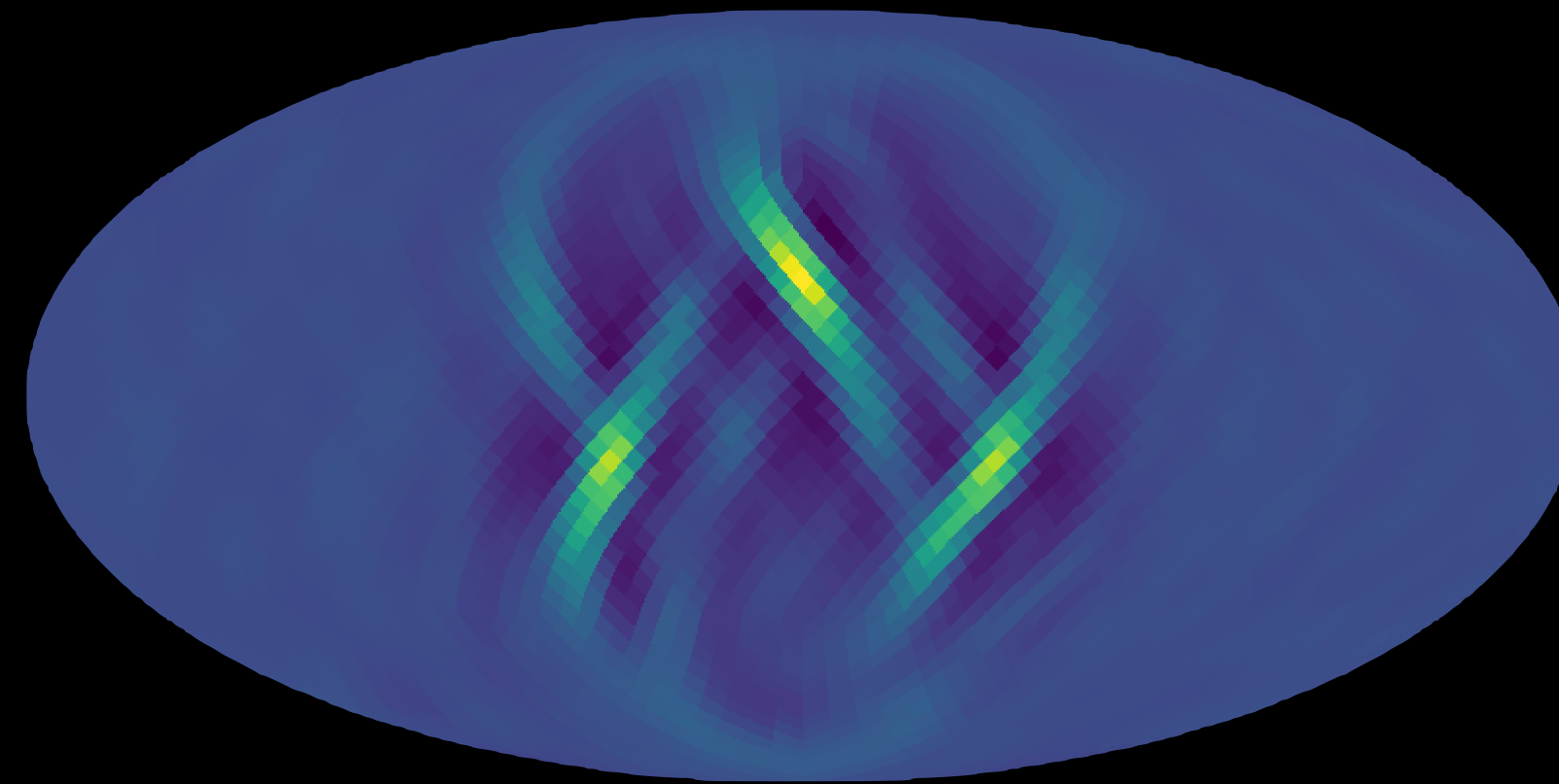
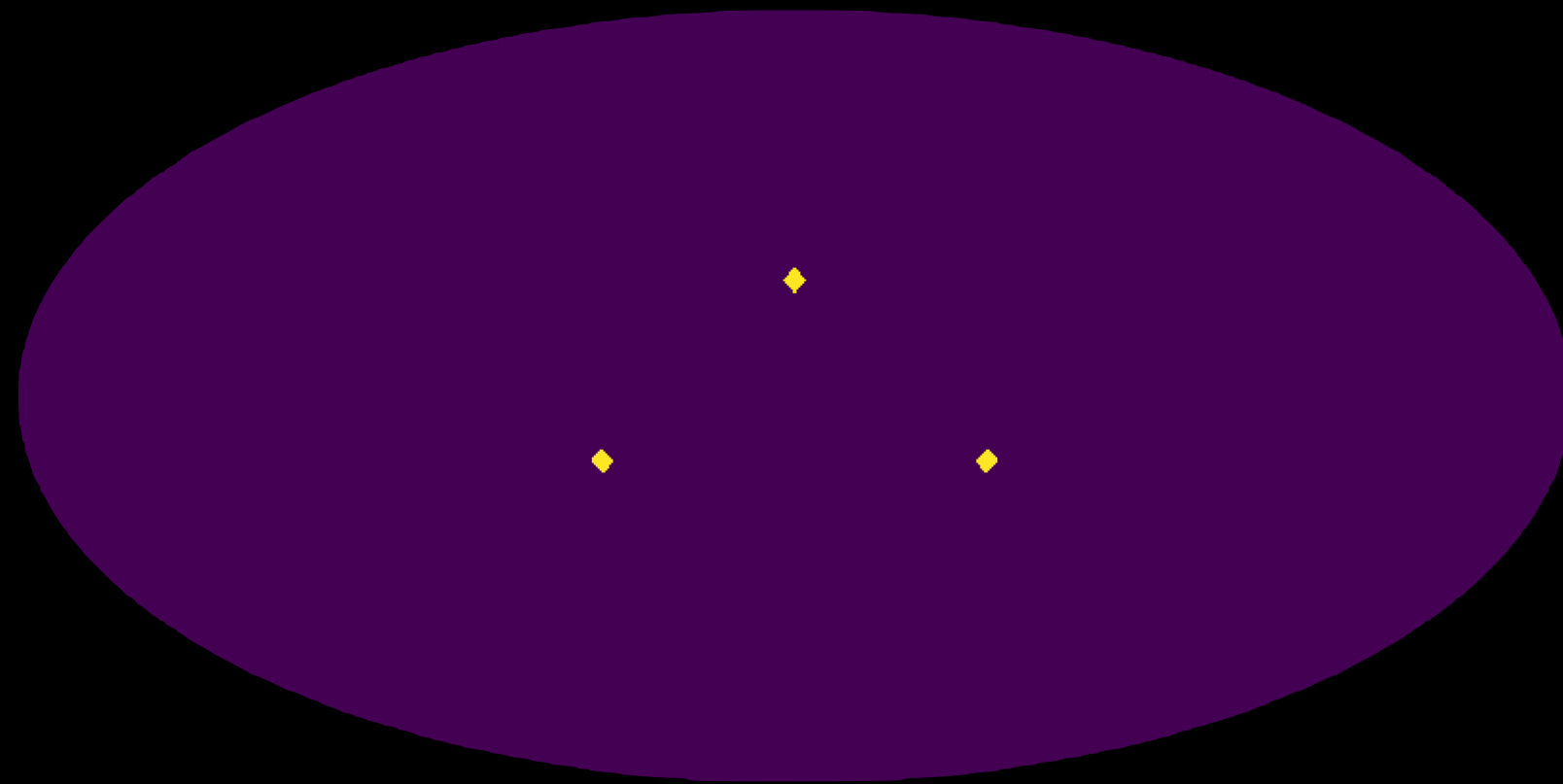


perform the whole analysis on a laptop in a few minutes\*



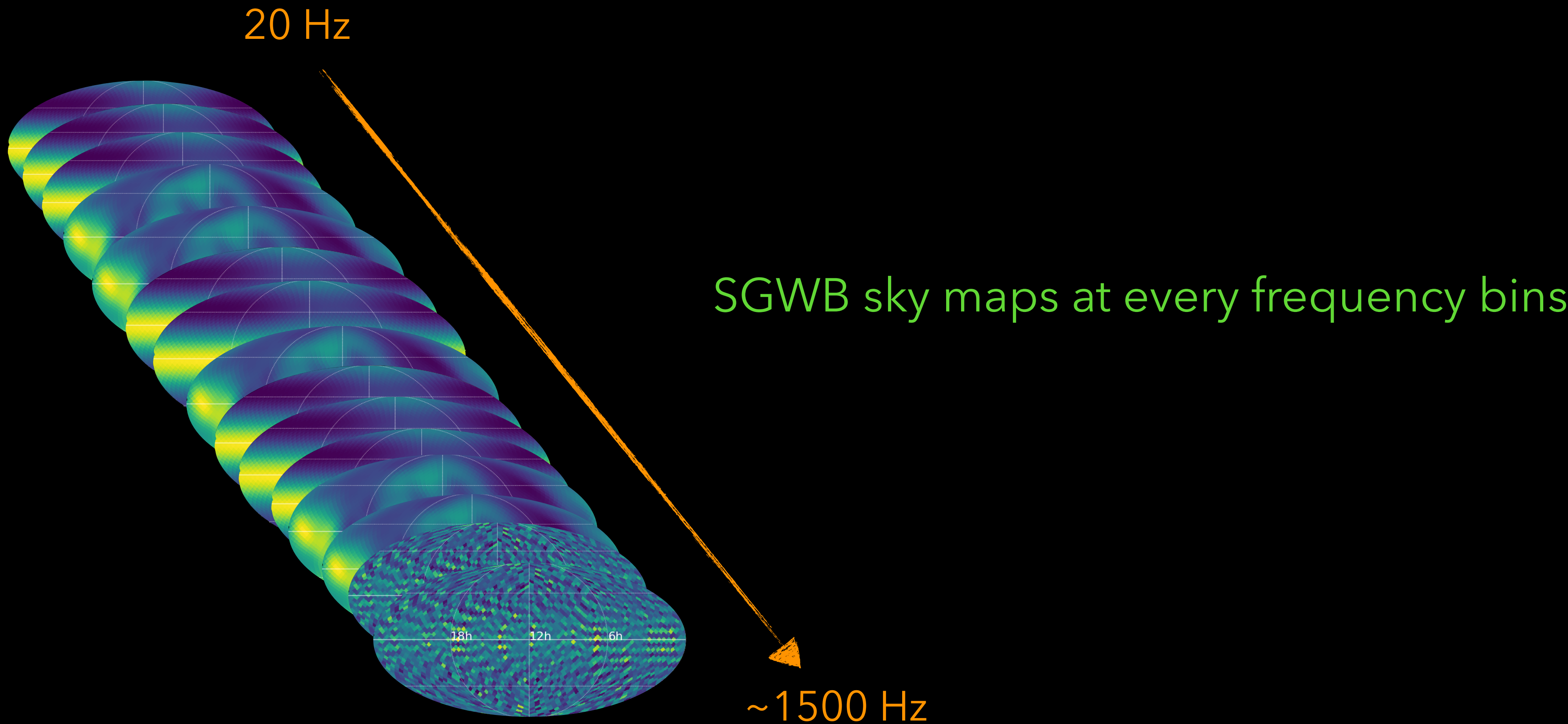
Produces the narrowband maps as an intermediate result

so separate search for different frequency spectra becomes redundant



# ALL-SKY ALL-FREQUENCY SEARCH

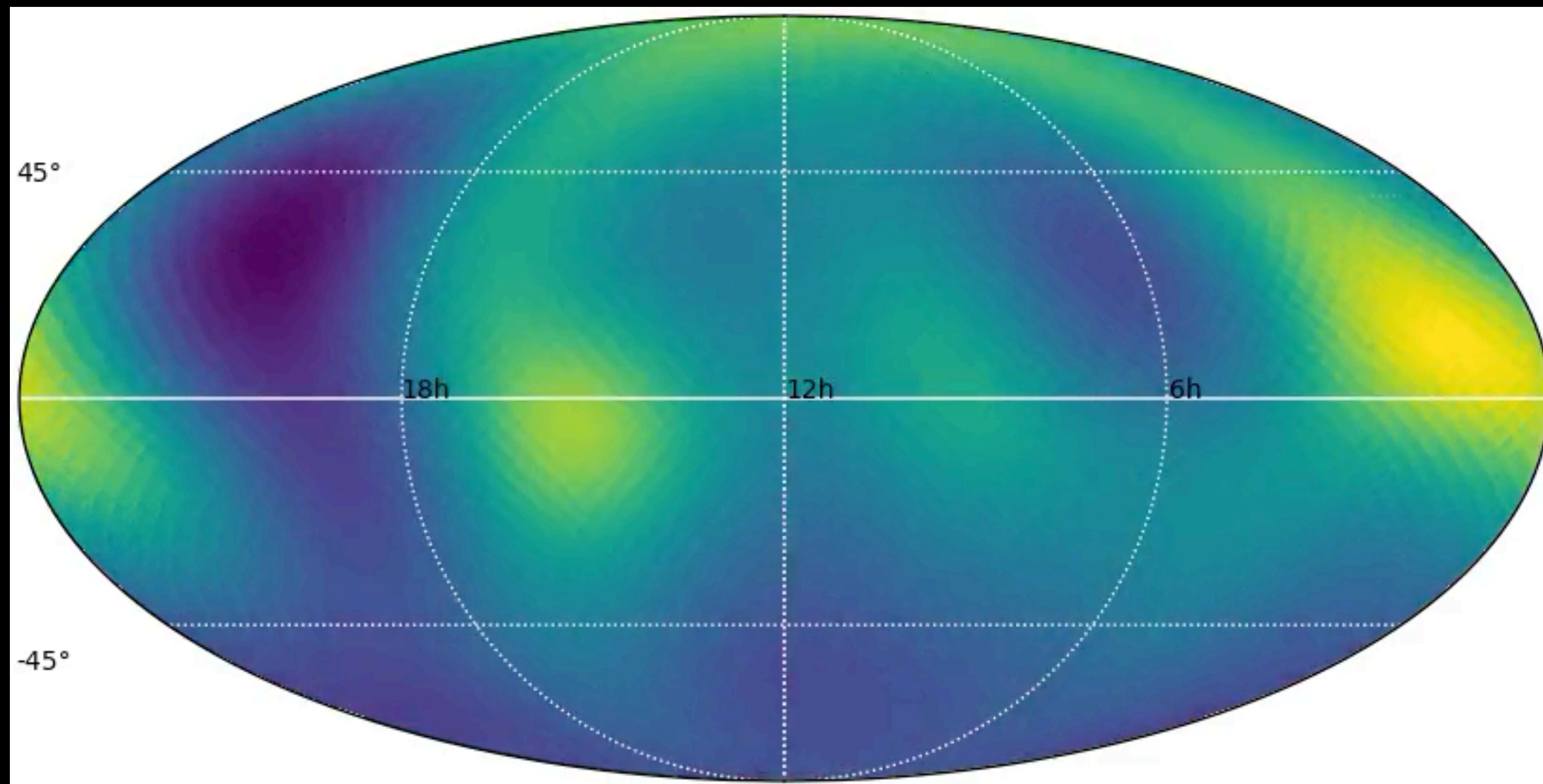
Now we have all the ingredients to perform an all-sky, all-frequency search, which assumes **no** specific power-law model for the SGWB



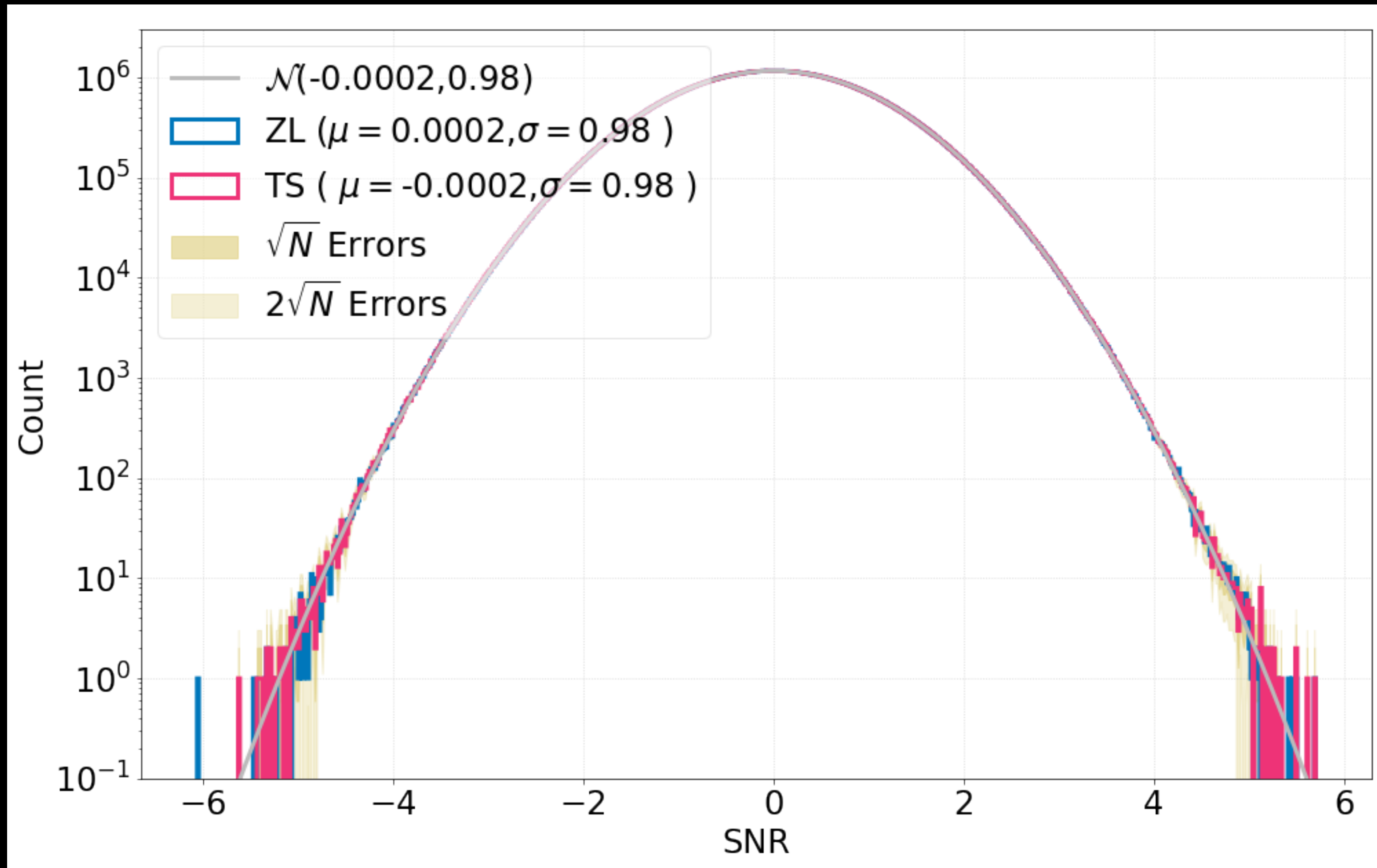
# ALL-SKY ALL-FREQUENCY SEARCH

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We presented the **first atlas** of SGWB sky from this analysis.



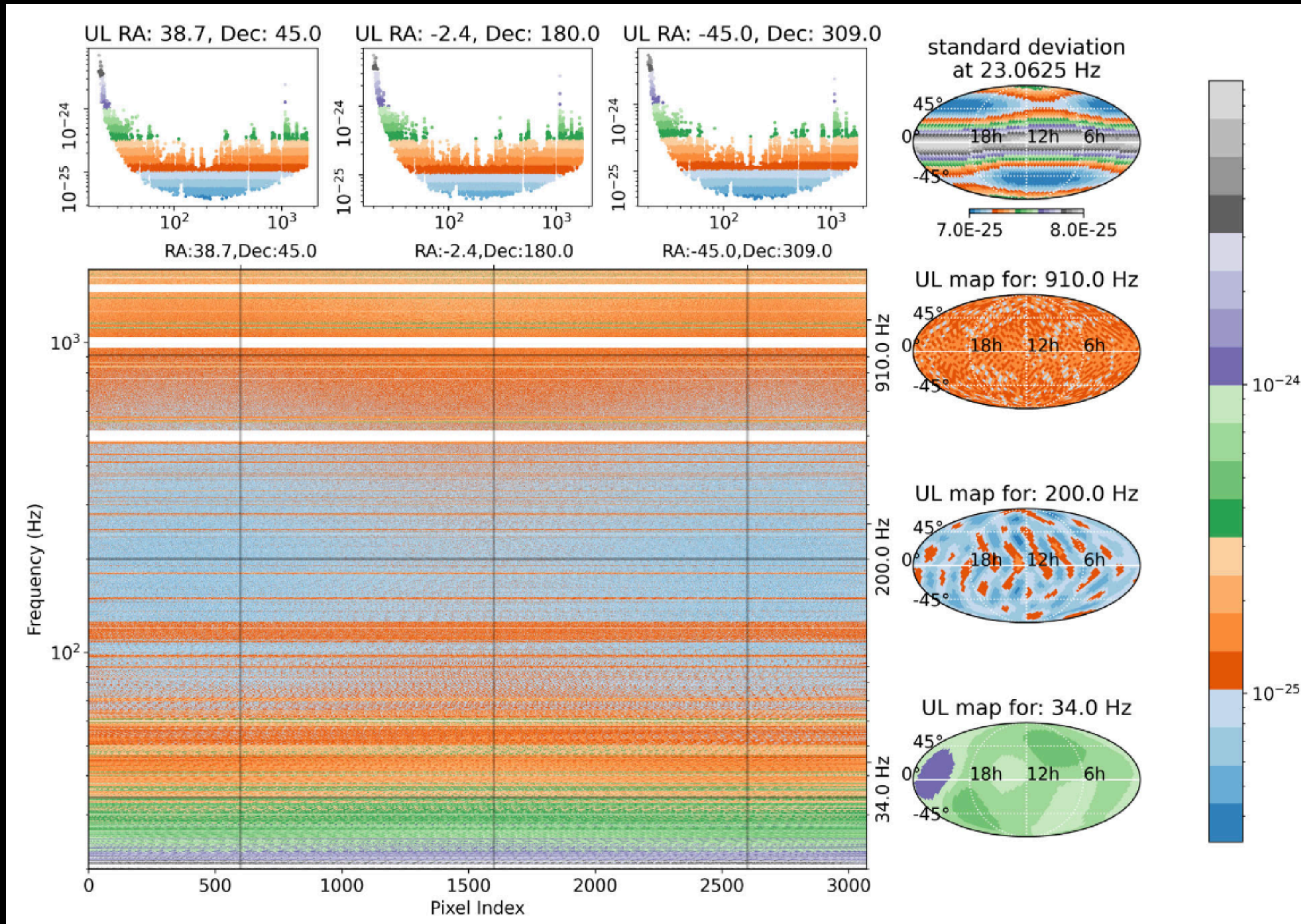
GW data from LIGO-Virgo-KAGRA's first three observing runs (O1 + O2 + O3)



The zero-lag (ZL) data is consistent with the time-shifted (TS) data within 2-sigma error bars.

We did follow-up studies on the outlier (SNR < -6) and found no astrophysical motivated channels. This outlier is also statistically insignificant, given the trial factors corrected p-value > 5%

Given no detection, we set the all-sky all-frequency upper limits on the SGWB strain

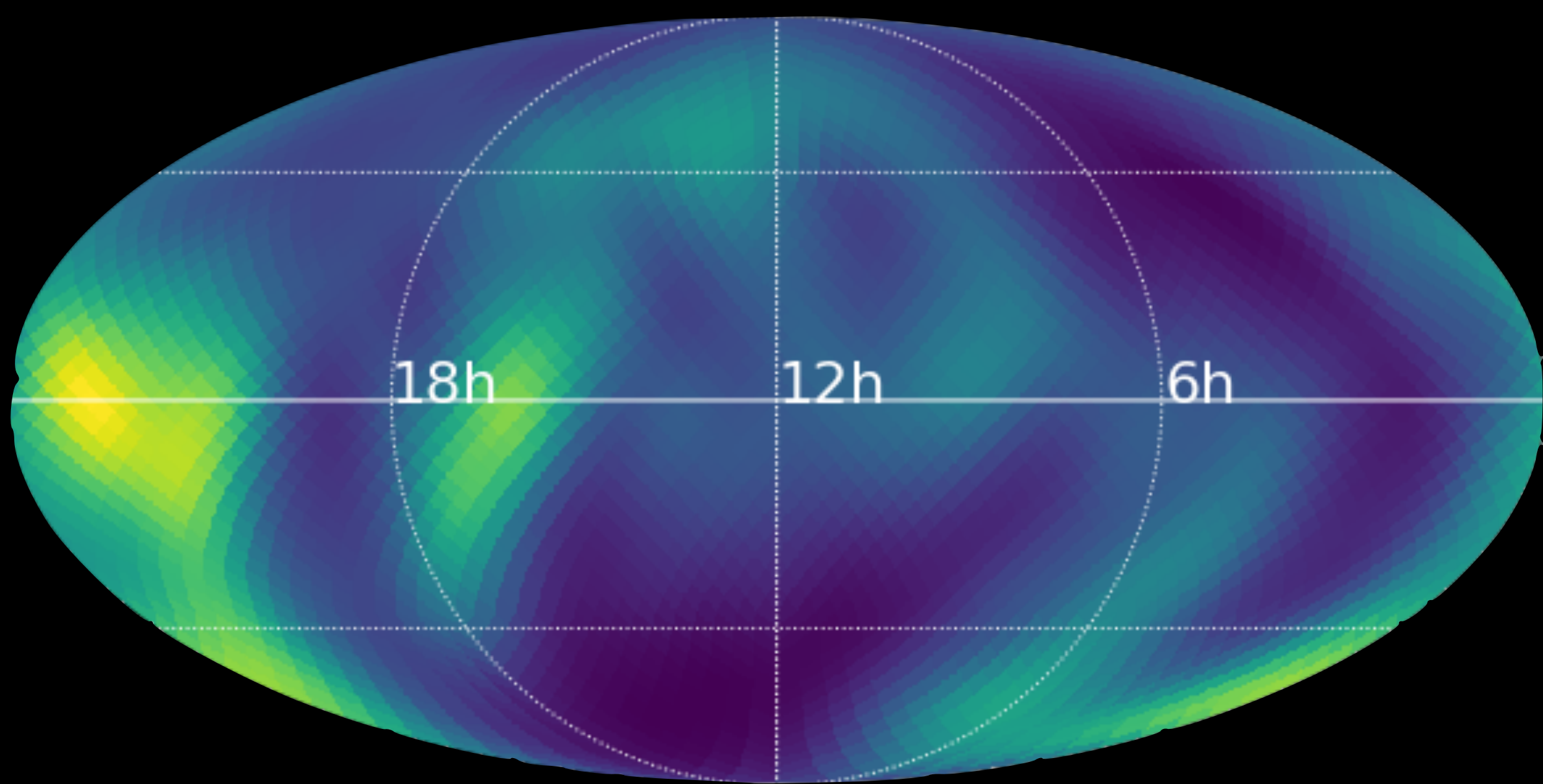


- The **colour bar** here denotes the range of **upper limit variations**.
- The **vertical** cross-section in this diagram shows the **frequency-dependent upper limit** in a particular direction.
- The **Horizontal** cross-sections form a **map of upper limits** in a particular frequency.
- **Notched frequencies** in a baseline appear as **horizontal white bands** in the plot.

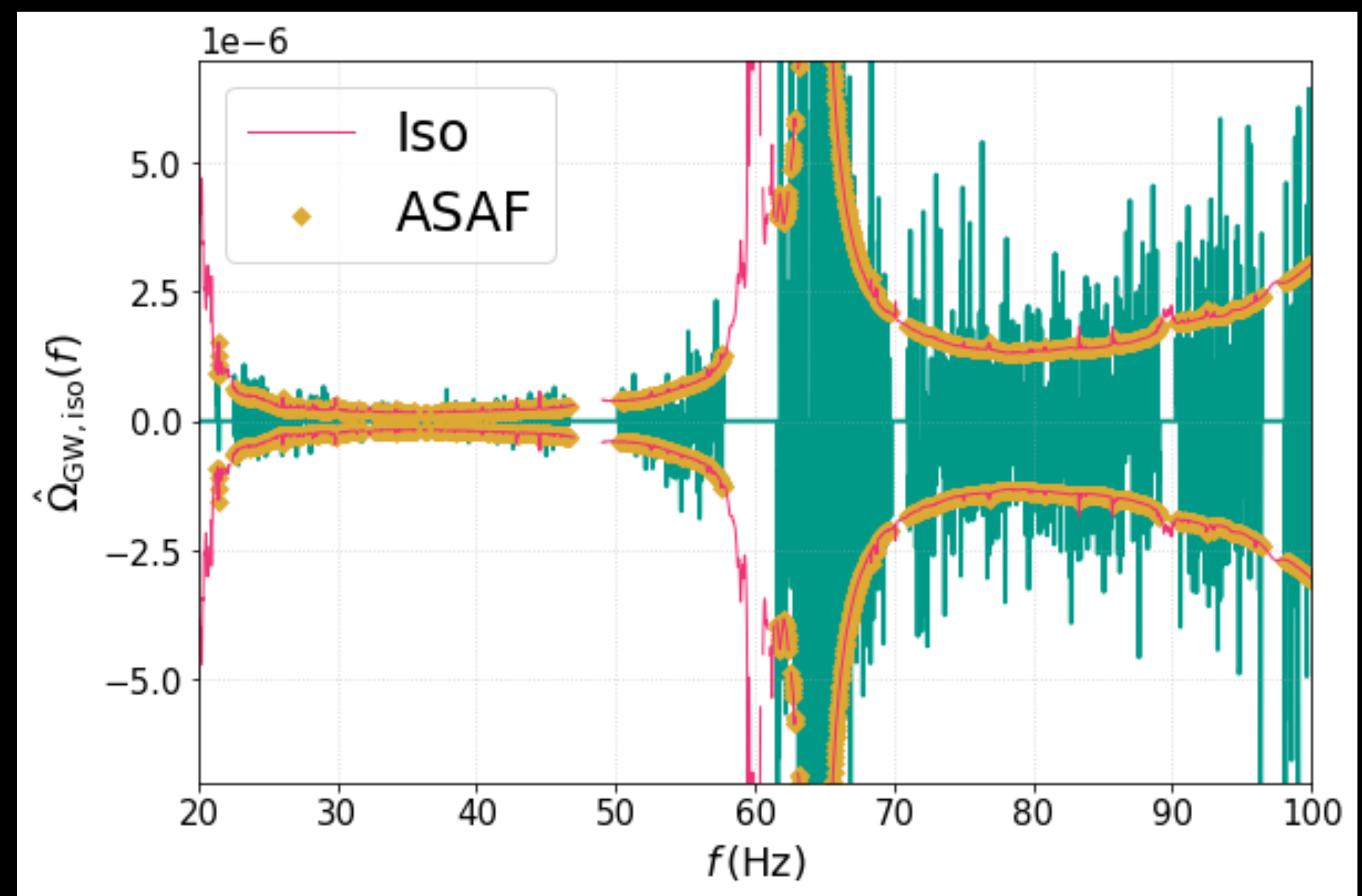


# ALL-SKY ALL-FREQUENCY RESULTS

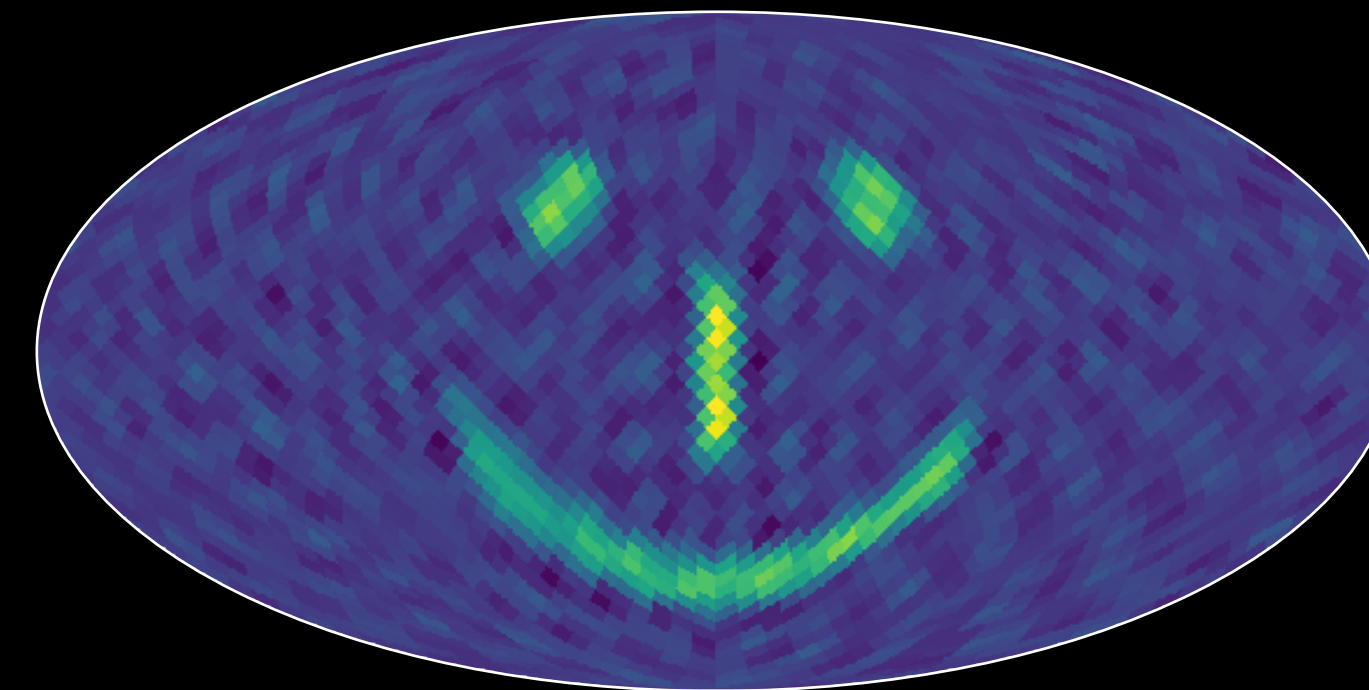
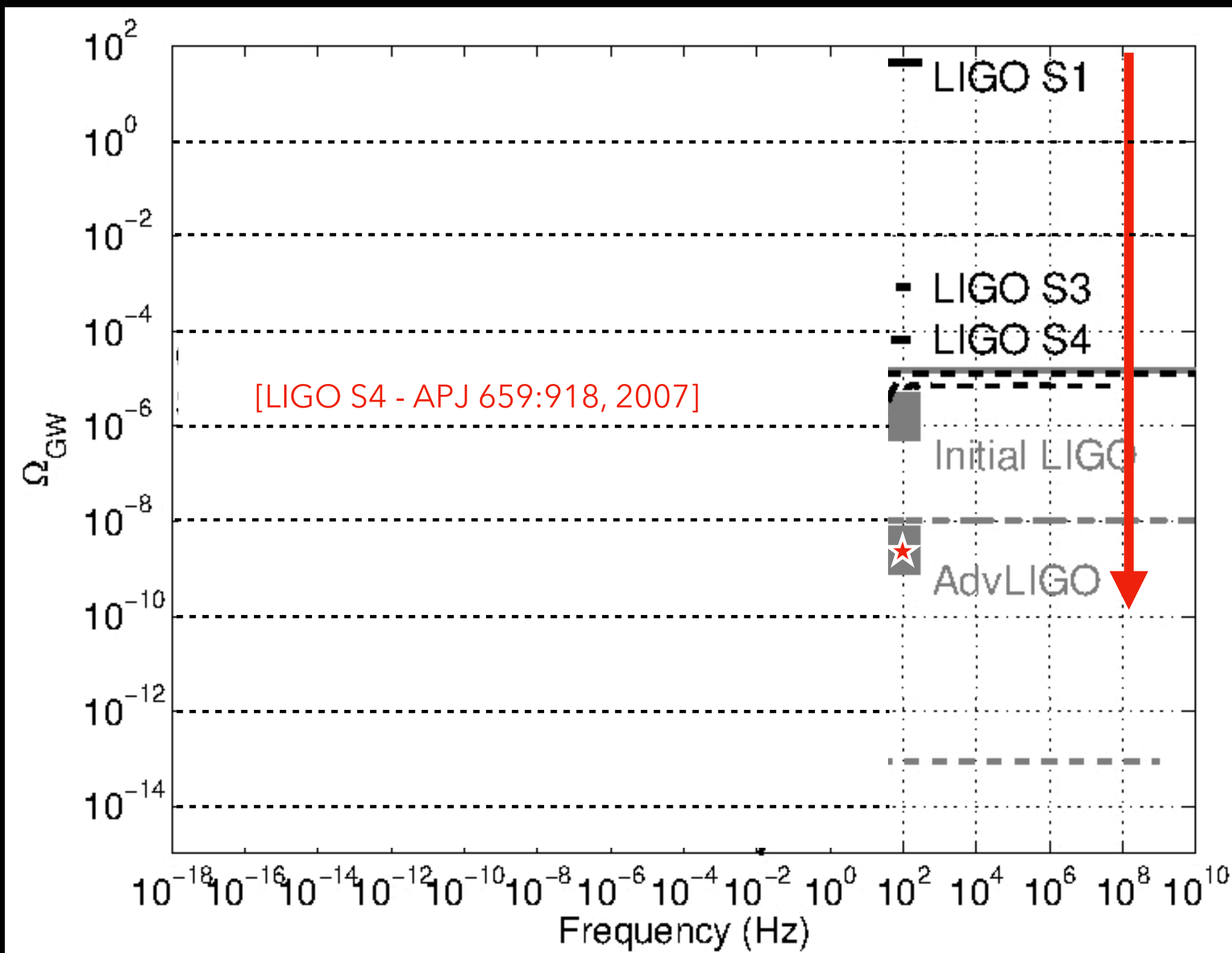
Assume a power law and combine these narrowband maps to obtain the 'usual' broadband results



Assume a power law and sum over all the directions of these narrowband maps to obtain the 'usual' isotropic results



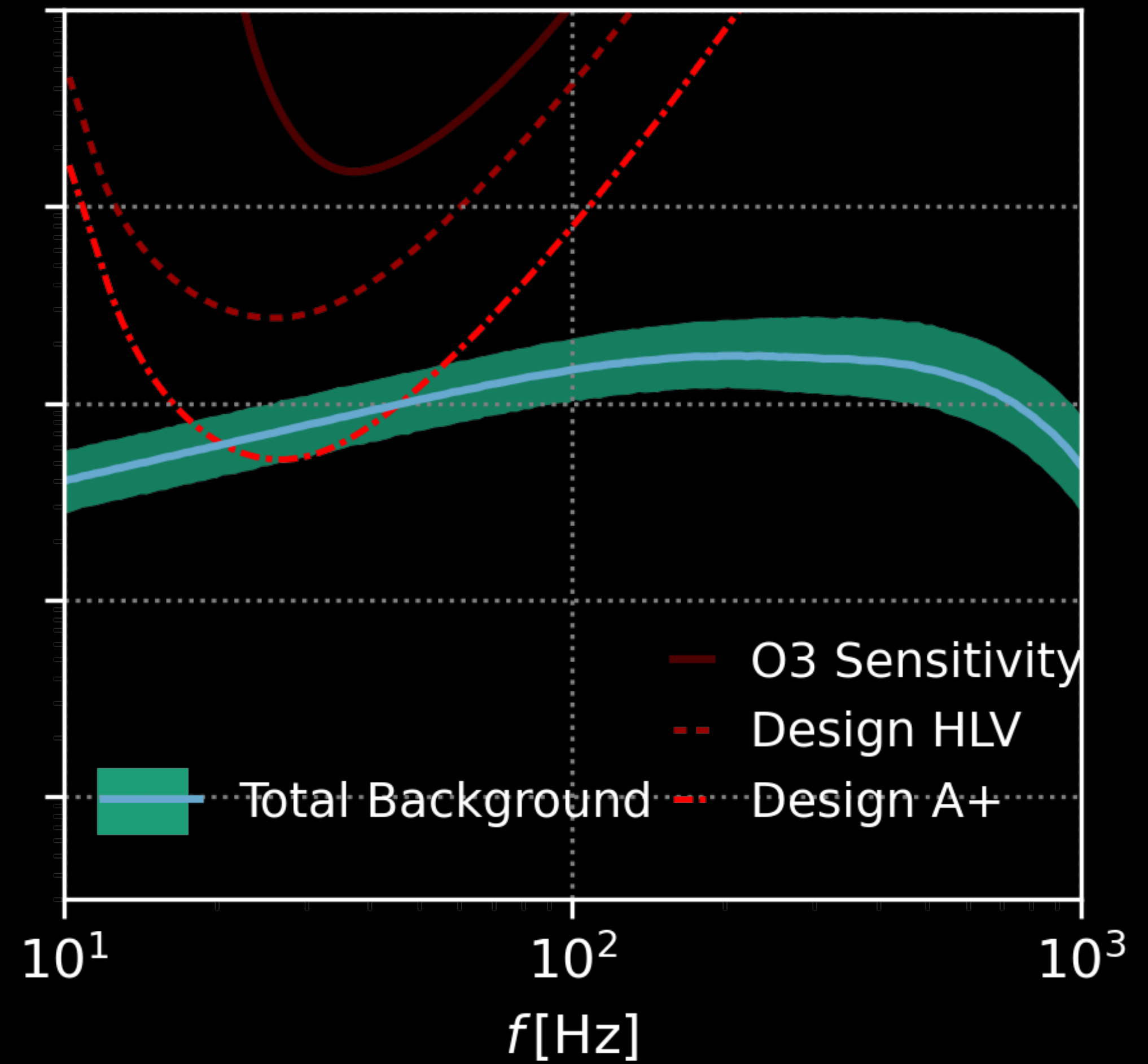
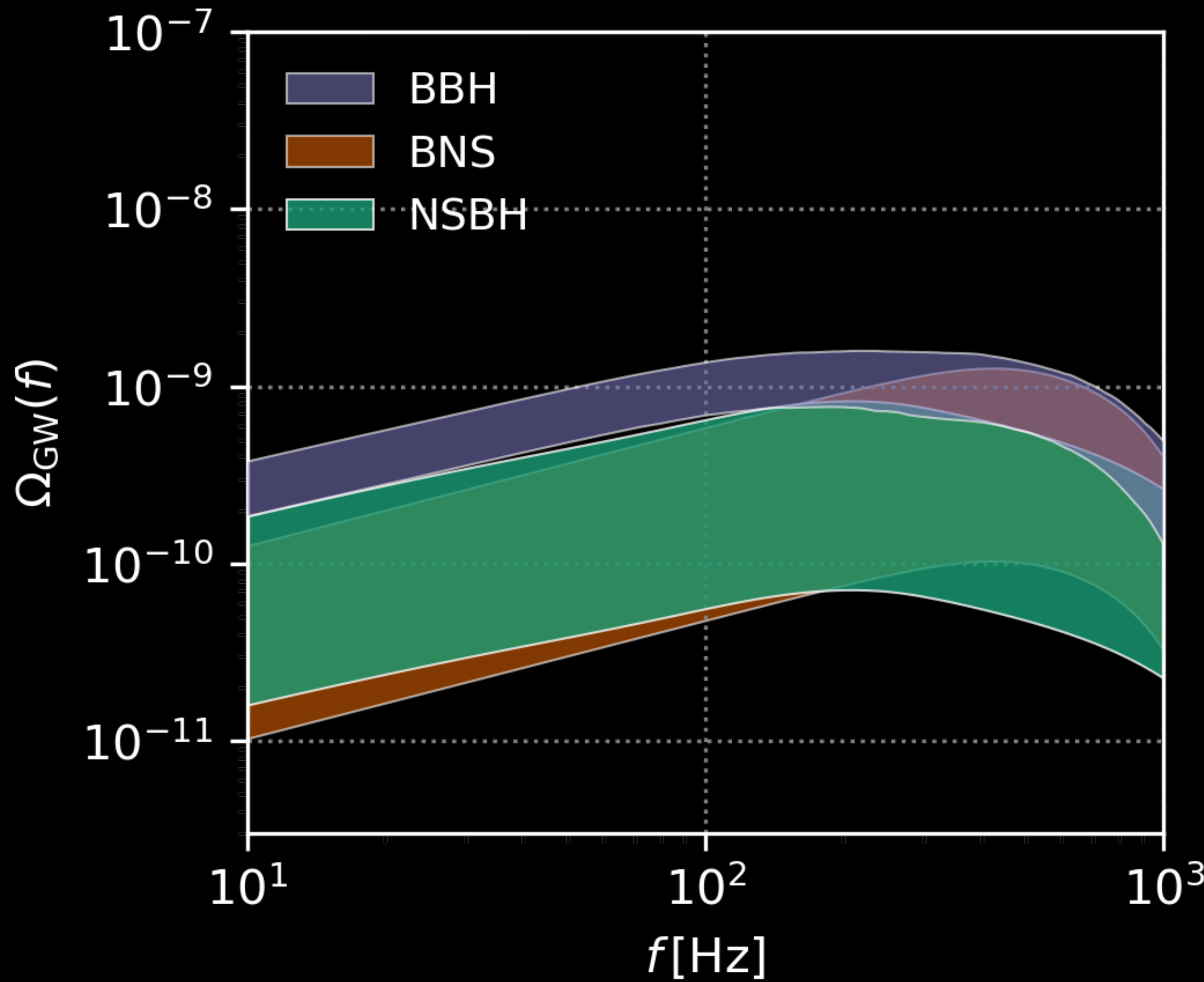
- New searches and techniques are opening up efficient ways to probe the dark universe.
- Plenty more work to do!
  - More detectors, More signals, More systems, and Dealing with real data.....



**Thank you!**

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BACKUP



The individual contributions expected from the collection of BNS, NSBH, and BBH mergers. While uncertainties on the energy density due to BNS and NSBH are due to Poisson uncertainties in their merger rates, our forecast for the SGWB due to BBHs includes systematic uncertainties associated with their imperfectly known mass distribution. (Right): Estimate of the total gravitational-wave background (green), as well as our current experimental sensitivity (red)

- Broadband: point sources with different power-law spectra.
- Narrowband: point sources having narrow GW frequency band (SN 1987A, ScoX-1, GC)
- Spherical harmonics search: Extended or diffuse sources - measure angular power spectra

<b>All-sky BBR Results</b>			<b>Max SNR (% <math>p</math>-value)</b>				<b>Upper limit ranges (<math>10^{-8}</math>)</b>	
$\alpha$	$\Omega_{\text{GW}}$	$H(f)$	<b>HL(O3)</b>	<b>HV(O3)</b>	<b>LV(O3)</b>	O1 + O2 + O3 <b>(HLV)</b>	O1 + O2 + O3 <b>(HLV)</b>	O1 + O2 <b>(HL)</b>
0	Constant	$\propto f^{-3}$	2.3 (66)	3.4 (24)	3.1 (51)	2.6 (23)	1.7–7.6	4.4–21
2/3	$\propto f^{2/3}$	$\propto f^{-7/3}$	2.5 (59)	3.7 (14)	3.1 (62)	2.7 (24)	0.85–4.1	2.3–12
3	$\propto f^3$	Constant	3.7 (32)	3.6 (47)	4.1 (12)	3.6 (20)	0.013–0.11	0.046–0.32

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<i>Narrow band Radiometer Results</i>					
<b>Direction</b>	<b>Max SNR</b>	<b><i>p</i>-value (%)</b>	<b>Frequency (Hz) (<math>\pm 0.016</math> Hz)</b>	<b>Best upper limit (<math>10^{-25}</math>)</b>	<b>Frequency band (Hz)</b>
Scorpius X-1	4.1	65.7	630.31	2.1	189.31–190.31
SN 1987A	4.9	1.8	414.0	1.7	185.13–186.13
Galactic Center	4.1	62.3	927.25	2.1	202.56–203.56

- Broadband: point sources with different power-law spectra.
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<b>SHD Results</b>								
$\alpha$	$\Omega_{GW}$	$H(f)$	Max SNR (% $p$ -value)				Upper limit range ( $10^{-9}$ )	
			HL(O3)	HV(O3)	LV(O3)	O1 + O2 + O3 (HLV)	O1 + O2 + O3 (HLV)	O1 + O2 (HL)
0	Constant	$\propto f^{-3}$	1.6 (78)	2.1 (40)	1.5 (83)	2.2 (43)	3.2–9.3	7.8–29
2/3	$\propto f^{2/3}$	$\propto f^{-7/3}$	3.0 (13)	3.9 (0.98)	1.9 (82)	2.9 (18)	2.4–9.3	6.4–25
3	$\propto f^3$	Constant	3.9 (12)	4.0 (10)	3.9 (11)	3.2 (60)	0.57–3.4	1.9–11