Emulating the non-linear power spectrum in beyond-LCDM models

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Based on: 2203.11120 + 2303.09549

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Introduction

• The era of Stage-IV LSS surveys is finally arriving:







• And with it, the information on <u>small scales</u> will drastically increase

Introduction

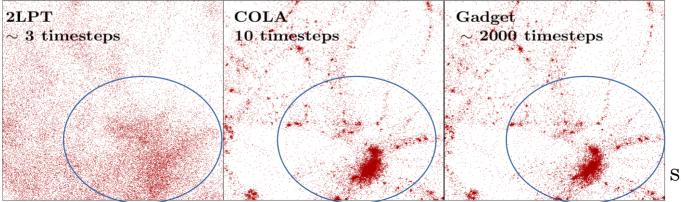
• The main quantity we want to get predictions on non-linear scales is the <u>matter power spectrum</u>

• One way of modelling this quantity in these scales is using <u>N-Body simulations</u>

- However:
 - 1) N-Body simulations are time consuming and computationally expensive
 - 2) Performing a full MCMC parameter estimation \rightarrow need an order of 10⁴-10⁵ simulations
 - 3) The power spectrum measured from these simulations is noisy \rightarrow affected by resolution issues
 - 4) All of this is worse in beyond-LCDM models

COLA

- <u>Issue 1</u>) To bypass the cost of running one N-Body sim we can use the COmoving Lagrangian Approximation method:
 - 2nd Order Lagrangian Perturbation Theory + Particle-Mesh Algorithm
 - Effective decoupling of large scale modes (LPT) and small scale modes (PM)
 - Reduced number of time-steps at the cost of losing resolution at small scales

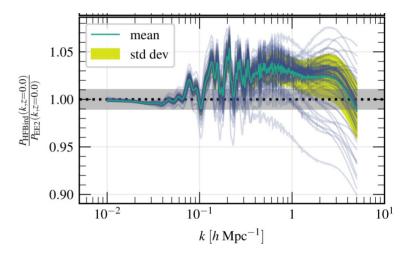


S. Tassev et al - 1301.0322

LCDM-Emulators

- <u>Issue 2</u>) Emulation techniques:
 - Instead of running 10⁴ simulations we run ~10² simulations
 - Emulation methods then interpolate the results of cosmological simulations using machine learning
 - Example of emulators for the non-linear power spectrum in LCDM:
 - Euclid Emulator 2 (EE2): M. Knabenhans et al 2010.11288
 - N-Body simulations PKDGRAV3
 - Approx 250 simulations

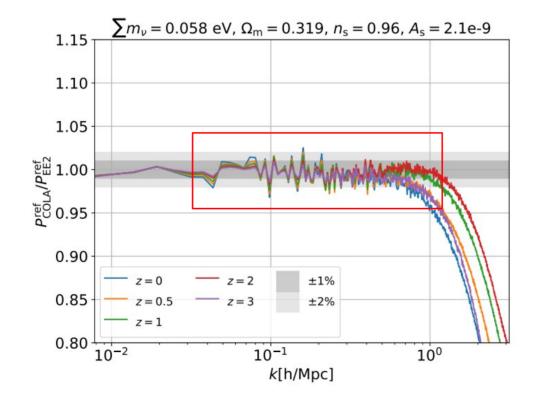
- Emulation of the boost:
$$B(k, z) = \frac{P_{\text{non}}(k, z)}{P_{\text{lin}}(k, z)}$$



Parameter	min.	max.	Reference
Ω_m	0.24	0.40	0.319
Ω_b	0.04	0.06	0.049
n_s	0.92	1.0	0.96
A_s	$1.7 imes 10^{-9}$	$2.5 imes 10^{-9}$	2.1×10^{-9}
h	0.61	0.73	0.67

Table 1: Parameter space EE2

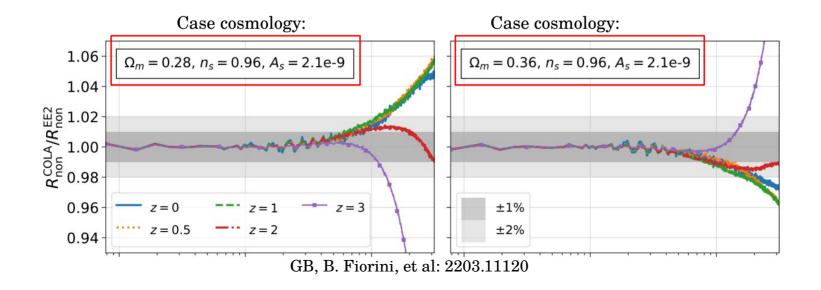
• Issue 3) The issue:



GB, B. Fiorini, et al: 2203.11120

• Issue 3) To reduce the noise coming from resolution effects we can instead compare the non-linear response function, which is defined as:

$$R_{\rm non} = \frac{P_{\rm non}^{\rm case}}{P_{\rm non}^{\rm ref}}$$



• Additionally, from the definition of the boost factor of a "case" cosmology (which can be any cosmology other than the reference one):

$$B^{\text{case}}(k,z) = \frac{P_{\text{non}}^{\text{case}}(k,z)}{P_{\text{lin}}^{\text{case}}(k,z)}$$

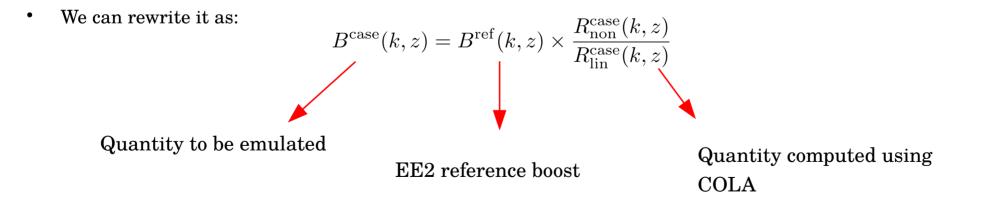
• We can rewrite it as:

$$B^{\text{case}}(k,z) = B^{\text{ref}}(k,z) \times \frac{R^{\text{case}}_{\text{non}}(k,z)}{R^{\text{case}}_{\text{lin}}(k,z)}$$

• Where we now see the Rnon function appearing explicitly

• Additionally, from the definition of the boost factor of a "case" cosmology (which can be any cosmology other than the reference one):

$$B^{\text{case}}(k,z) = \frac{P_{\text{non}}^{\text{case}}(k,z)}{P_{\text{lin}}^{\text{case}}(k,z)}$$



COLA & Modified Gravity

• To emulate the power spectrum for modified gravity we can use this prescription

• Where the <u>main quantity</u> we need to compute in COLA is the ratio (Rnon) of a case cosmology with respect to a reference one

• In MG theories this is simply:
$$\frac{P^{\mathrm{MG}}\left(k,z
ight)}{P^{\mathrm{ref}}\left(k,z
ight)}$$

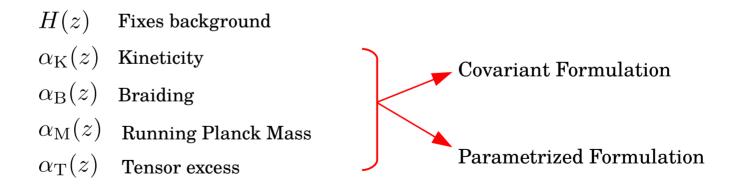
• So I will now present the MG theories we have created a fast way to run COLA simulations, and compare our results with available N-Body simulations in the literature

COLA & Modified Gravity

• We consider Hondeski gravity, whose action is big and messy:

$$S[g_{\mu\nu},\phi] = \int d^4x \,\sqrt{-g} \left[\sum_{i=2}^5 \frac{1}{8\pi G} \mathcal{L}_i[g_{\mu\nu},\phi] + \mathcal{L}_{\rm m}[g_{\mu\nu},\psi_M] \right]$$

• But can be made simpler:



E. Bellini and I. Sawicki 1404.3713

COLA & Modified Gravity

• MG N-Body simulations solve:

N-Body:

$$\nabla^{2} \Phi = 4\pi G a^{2} \delta \rho_{m} + \frac{1}{2} \nabla^{2} \delta \phi,$$

$$\partial_{\tau} \delta + k \vec{v} = 0,$$

$$(\partial_{\tau} + \mathcal{H}) \vec{v} = -\nabla \Phi,$$

$$L[\delta \phi] = S(\delta \rho_{m}, \delta \phi)$$

Needs different solver: Adds even more time and complexity

COLA:

$$\nabla^2 \Phi = 4\pi G_{\text{eff}} a^2 \bar{\rho}_{\text{m}} \delta_{\text{m}},$$
$$\partial_\tau \delta + k \vec{v} = 0,$$
$$[\partial_\tau + \mathcal{H}] v = -\nabla \Phi$$
$$G_{\text{eff}} = G_{\text{eff}} \left(\alpha_{\text{B}}, \alpha_{\text{M}}, \alpha_{\text{T}}\right)$$

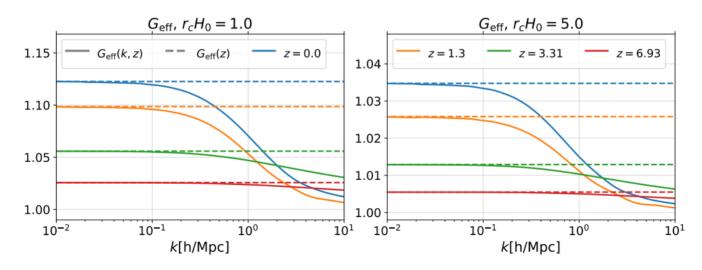
One single time-dependent function Linear theory

Screening

• In small scales, MG theories introduce an aditional force acting on particles, called 5th force. To shield this extra force we need to implement a screening mechanism in our simulations

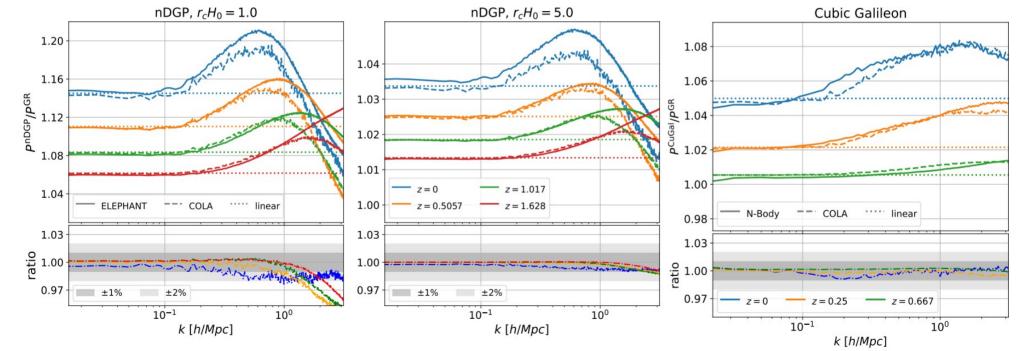
$$G_{\mathrm{eff}}\left(\alpha_{\mathrm{B}}\left(z\right),\alpha_{\mathrm{M}}\left(z\right),\alpha_{\mathrm{T}}\left(z\right)\right)\longrightarrow G_{\mathrm{eff}}\left(\mathbf{k},\alpha_{\mathrm{B}}\left(z\right),\alpha_{\mathrm{M}}\left(z\right),\alpha_{\mathrm{T}}\left(z\right)\right)$$

• In our COLA simulations we have implemented a numerical routine that numerically computes Geff as a function of scale, transitioning from its linear theory value to its GR one:



Screening

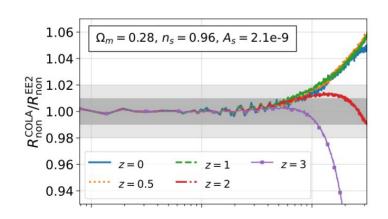
• Comparison with full N-Body simulations available in the literature

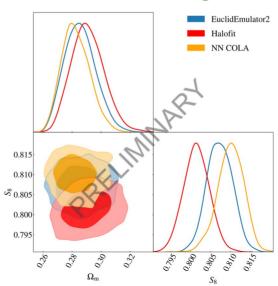


GB, K. Koyama, et al: 2303.09549

Conclusions

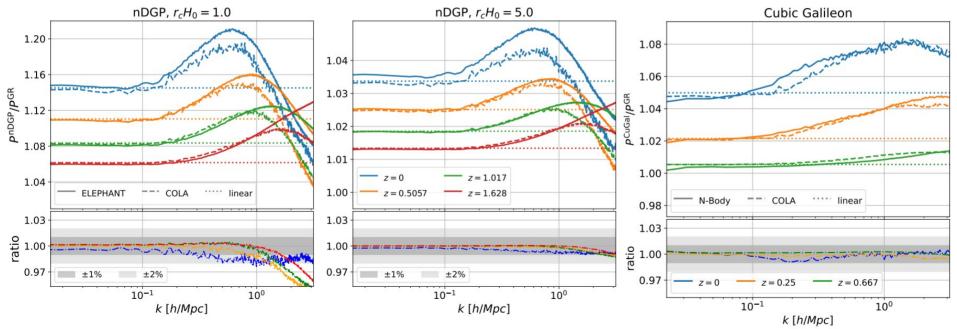
- Take aways:
 - COLA allows us to run faster simulations
 - Emulators reduce the number of required simulations
 - Emulating the Boost reduces resolution effects that worsen the agreement





Conclusions

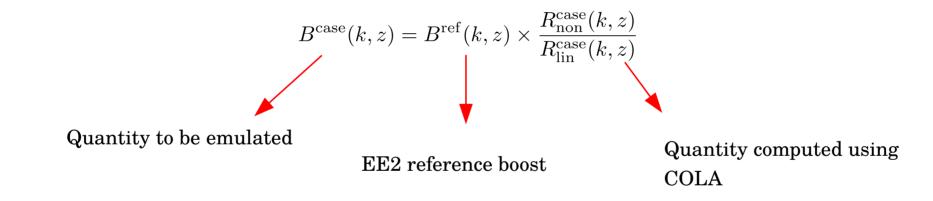
- Modified gravity:
 - Fast pipeline to generate simulations in Horndeski theories



Thank you

Backup

• Our methodology to investigate beyond-LCDM models is given by:



- As a proof of concept we created a Neural Network Emulator for:
 - LCDM with fixed sum of neutrino masses, 0.058 eV, and same parameter range as EE2
 - 400 COLA simulations

Backup

• We then performed a cosmic shear forecast analysis for an LSST-Y1 like survey

