





The Multipole Expansion of the Expansion Rate in the Local Universe

Based on: Phys. Rev. D 107, 023507 (2023)

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Progress on Old and New Themes in cosmology (PONT) 04/05/2023

Introduction

The Standard Model of Cosmology

☐ The standard cosmological model relies on the theory of general relativity, with the cosmological principle (CP).

 \square In the standard model of cosmology, all the cosmological parameters (density parameters Ω , Hubble parameter H ...) depend only on time.

☐ Hubble constant tension problem is one of the main problems in standard cosmology.

In the local universe:

By CMB:
$$H_0 = 67.4 \pm 0.5 \, km/s/Mpc$$

(Riess et al. 2021)

 $H_0 = 73.2 \pm 1.3 \, km/s/Mpc$

(Planck Collaboration et al. 2020)

Introduction

□ One of the possible solution for Hubble constant tension is that the local universe is not homogeneous and isotropic.

□ Traditionally, deviations from the cosmological principle (CP) predictions are treated perturbatively by expanding the cosmological parameters into a smooth background component and a fluctuating part.

☐ We developed a completely non-perturbative and model-independent approach to investigate the local inhomogeneities.

 \square An observable that captures deviations from the linear redshift-distance relation is the expansion rate fluctuation (natural units c = 1),

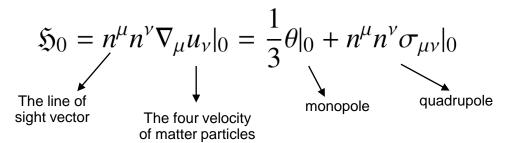
$$\eta \equiv \log \frac{z}{H_0 d_L}$$

☐ It is linear in the distance modulus

$$\hat{\eta} = \log \left[\frac{z}{H_0} \right] + 5 - \frac{\mu}{5}$$

 $lue{}$ If we ignore the error in the redshift, then $\hat{\eta}$ is Gaussian and unbiased estimator.

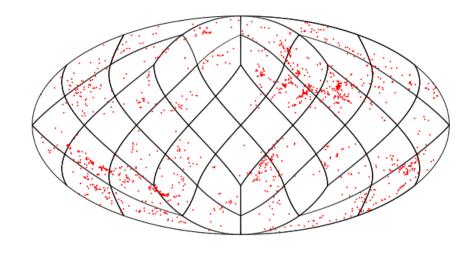
☐ If the observer is comoving with matter flow (\approx zLG), then $\hat{\eta}$ will converge to $\log(\mathfrak{H}_0/H_0)$ (for z→0), where \mathfrak{H}_0 is the effective Hubble parameter (Kristian & Sachs 1966, Clarkson & Maartens 2010. Heinesen 2021):



However, if we choose an observer that does not see the CMB dipole (z_{CMB} in FRW background), then η provides insights into the radial component of peculiar velocities (if v<<1)

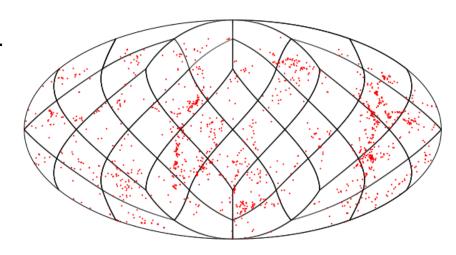
$$\eta \approx \frac{v(1+z)}{z \ln 10}$$

☐ By using **HEALPix** (Hierarchical Equal Area isoLatitude Pixelisation of a 2-sphere).



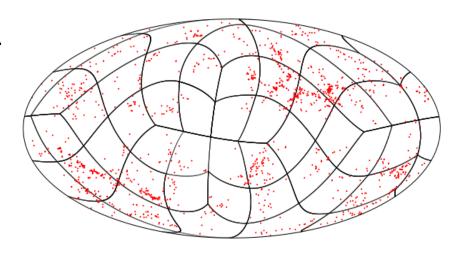
Number of pixels = 48

- ☐ By using **HEALPix** (Hierarchical Equal Area isoLatitude Pixelisation of a 2-sphere).
- ☐ Rotation to handle angular incompleteness.



Number of pixels = 48

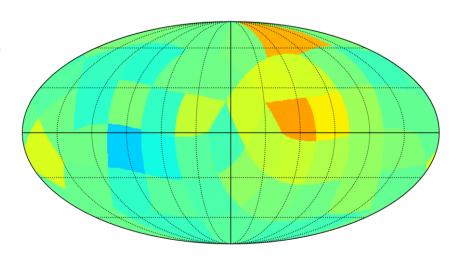
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Number of pixels = 48

- ☐ By using **HEALPix** (Hierarchical Equal Area isoLatitude Pixelisation of a 2-sphere).
- ☐ Rotation to handle angular incompleteness.
- ☐ Define a piece-wise function.

$$\eta = \left\langle \log \frac{z}{H_0 d_L} \right\rangle_{\Omega}$$



☐ It is then decomposed into the spherical harmonic basis

$$\eta = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\theta, \phi)$$

☐ The spherical harmonic coefficients

$$a_{\ell m} \equiv \int_{0}^{2\pi} \int_{0}^{\pi} \eta(\theta, \phi) Y_{\ell m}^{*}(\theta, \phi) \sin \theta \ d\theta d\phi$$

The zeroth order estimator is

$$\hat{a}_{\ell m}^{(0)} = \frac{4\pi}{N_{pix}} \sum_{p=1}^{N_{pix}} \eta(\Omega_p) Y_{\ell m}^*(\theta_p, \phi_p)$$



$$\hat{a}_{\ell m}^{(k+1)} = \hat{a}_{\ell m}^{(k)} + \frac{4\pi}{N_{pix}} \sum_{p=1}^{N_{pix}} \left(\eta(\Omega_p) - \eta^{(k)}(\theta_p, \phi_p) \right) Y_{\ell m}^*(\theta_p, \phi_p) \quad , \quad \eta^{(k)}(\theta_p, \phi_p) = \sum_{\ell=0}^{\ell_{max}} \sum_{m=-\ell}^{\ell} \hat{a}_{\ell m}^{(k)} Y_{\ell m}(\theta_p, \phi_p)$$

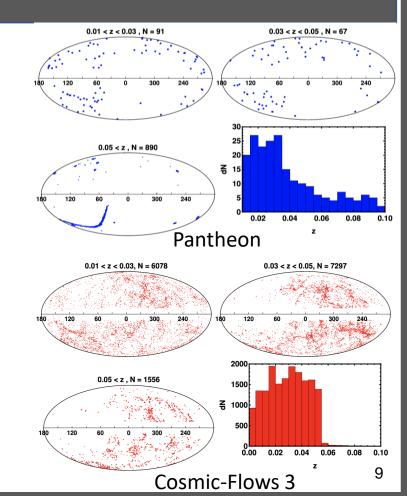
$$\Box$$
 The power spectrum $\hat{C}_{\ell} = \frac{1}{2\ell+1} \sum_{\ell=0}^{\ell} |a_{\ell m}^{(\infty)}|^2$

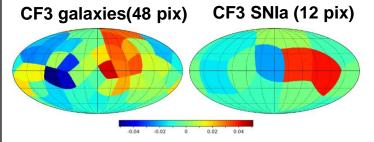
Data

- 1) <u>Pantheon</u> SNIa sample (Scolnic et al. 2018)
- Almost isotropic for z < 0.05 with (158 objects).
- The typical error in the distance modulus ($\Delta\mu=0.15$).

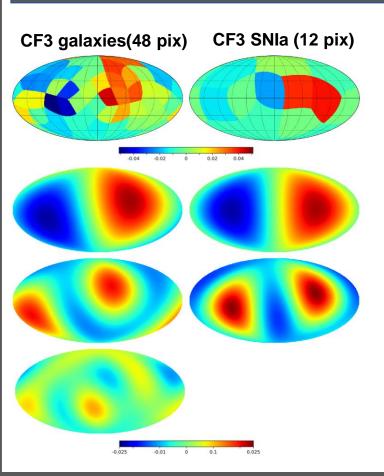
- 2) Cosmic-Flows 3 (CF3) galaxies (Tully et al. 2017)
- Almost isotropic for z < 0.05 with (13661 objects).
- We divided the sample into two independent subsamples:
- A) **CF3 SNIa**: 286 objects ($\Delta \mu = 0.18$).

B) **CF3 galaxies**: 13375 objects ($\Delta \mu = 0.46$).

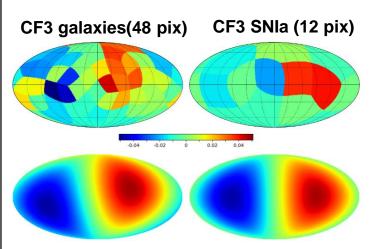




 η maps



 $\eta \; maps$

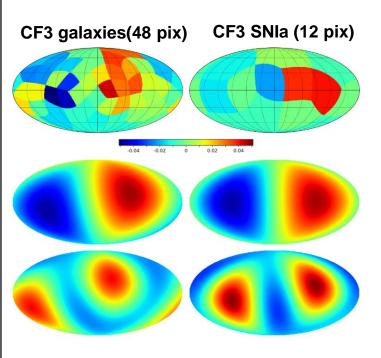


η maps

Same intensity and direction (CF3g, CF3sn and Pantheon). Same direction as bulk component of local group motion.

0.025 -0.01 0 0.1 0.025

12



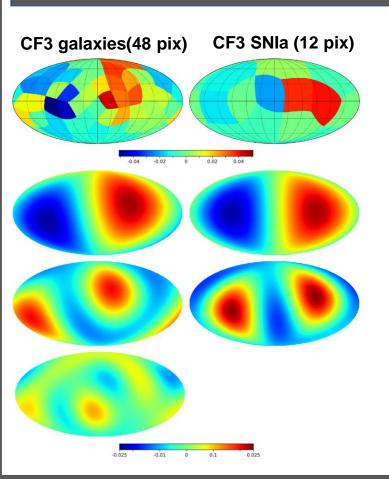
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Quadrupole intensity is half that of the dipole Alignement with the direction of the dipole No quadrupole component in Pantheon data!

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13

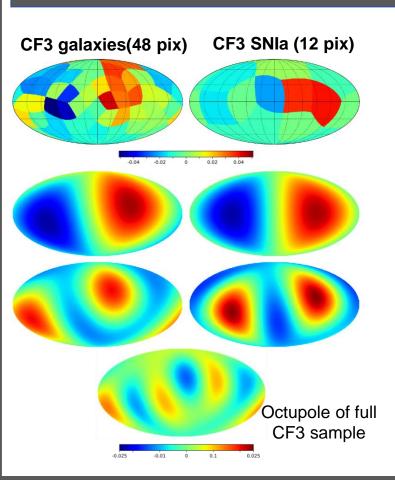


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Octupole: only for CF3 the S/N is acceptable (~3.1)



η maps

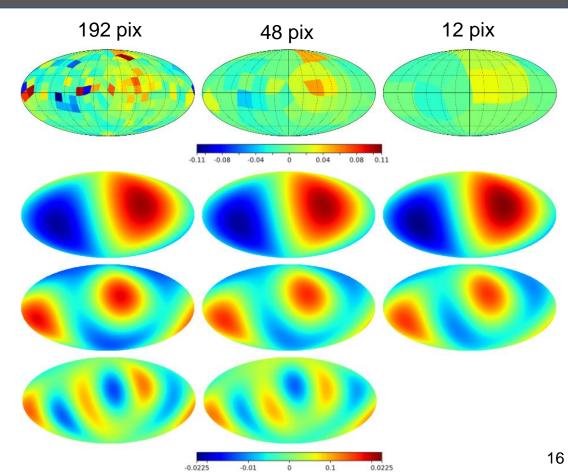
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The Robustness of the Results

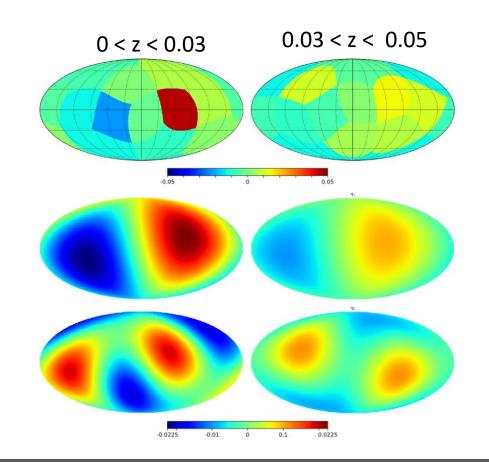
☐ Independent from pixelisation strategy.



The Robustness of the Results

☐ Independent from pixelisation strategy.

☐ Independent from the sample depth.



The Robustness of the Results

☐ Independent from pixelisation strategy.

Independent from the sample depth.

☐ Consistent with results of fitting the data with the spherical harmonics.

	Fourier	Fit	Fourier	Fit	Fourier	Fit	Fourier	Fit	
Sample	l	d	b	d	C_1	10^{4}	$C_2/10^4$		
CF3	285 ± 5	287 ± 5	11 ± 4	10 ± 4	5.1 ± 0.8	13±08	1 + 0 3	12+03	
[0.01, 0.05]	203 ± 3	267 ± 3	11 ± 4	10 ± 4	J.1 ± 0.6	4.5 ± 0.6	1 ± 0.3	1.2 ± 0.3	
CF3g	296 ± 6	290 ± 5	18 ± 5	5 ± 4	4.0 ± 0.6	53+00	13+04	13+03	
[0.01, 0.05]	290 ± 0	290 ± 3	10 ± 3	J ± 4	4.0 ± 0.0	3.3 ± 0.9	1.5 ± 0.4	1.5 ± 0.5	
CF3sn	222 + 22	293 ± 14	-8 ± 18	0 ± 10	3.7 ± 1.5	5.5 ± 2.6	1.7 ± 1.4	1 Q _ 1 1	
[0.01, 0.05]	322 ± 23							1.0 ± 1.1	
Pantheon	332 + 30	312 + 23	_0 + 10	10 + 26	3.9 ± 2.7	31+32	0.5 ± 1.4	06+10	
[0.01, 0.05]	332 ± 39	312 ± 23	7 1 19	17 ± 20	J.7 ± 2.1	J.T ± J.Z	0.5 I 1.4	0.0 I 1.0	

Errors and Biases

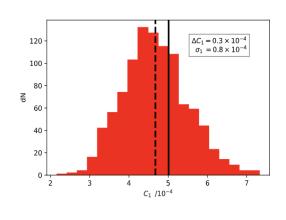
- The errors and biases are computed by 1000 Monte Carlo simulations.
- Variance of the power spectrum coefficients

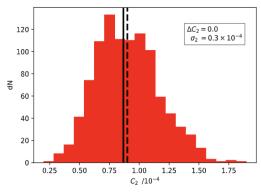
$$V\left[\hat{C}_{\ell}\right] = \left(\frac{1}{2\ell+1}\right)^{2} \sum_{n=0}^{2\ell} V[w_{n}^{(\ell)}]$$

where

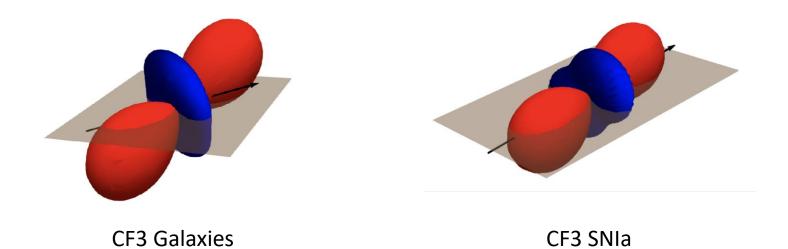
and

$$V[w_n^{(\ell)}] = \begin{cases} 2\sigma_{\ell 0}^4 + 4\sigma_{\ell 0}^2 a_{\ell 0}^2 & n = 0 \\ 8\sigma_{\ell n}^{(R)4} + 16\sigma_{\ell n}^{(R)2}\Re[a_{\ell n}]^2 & \ell \geq n > 0 \\ 8\sigma_{\ell (n-\ell)}^{(I)4} + 16\sigma_{\ell (n-\ell)}^{(I)2}\Im[a_{\ell (n-\ell)}]^2 & 2l \geq n > \ell \end{cases}$$
 and
$$\sigma_{\ell m}^{(R)2} = V[\Re[\hat{a}_{\ell m}]] \quad , \quad \sigma_{\ell m}^{(I)2} = V[\Im[\hat{a}_{\ell m}]]$$





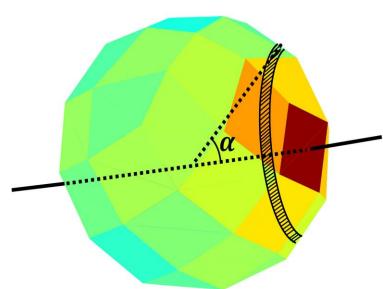
□ 3D structure of the quadrupole (Galactic plane shown for reference)

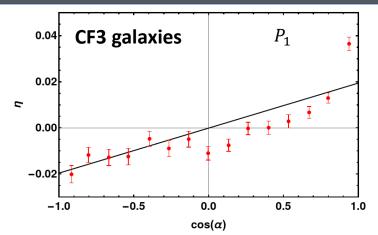


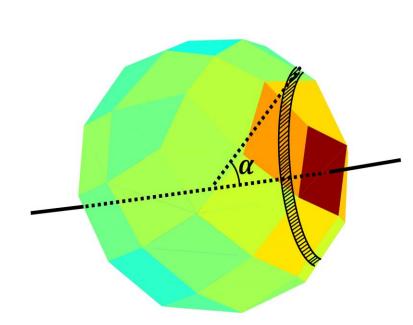
☐ The quadrupole is symmetric around the axis of the maximum ($l\sim285$, $b\sim11$).

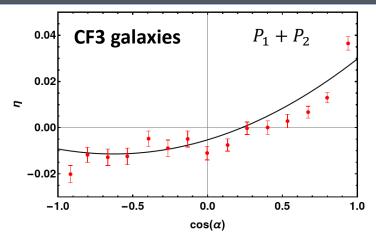
☐ Sufficient to expand using the Legendre basis about the axis of symmetry

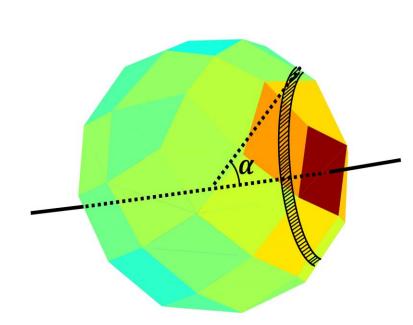
$$\eta(\alpha) = \sum_{\ell=1}^{\infty} a_{\ell} P_{\ell}(\cos \alpha)$$

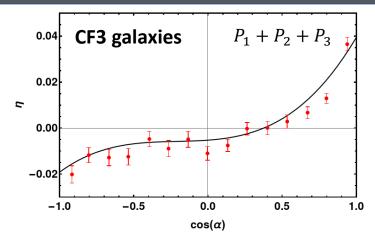


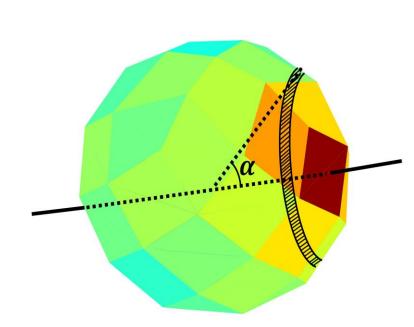


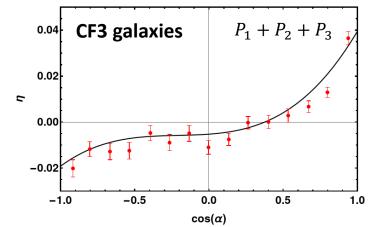


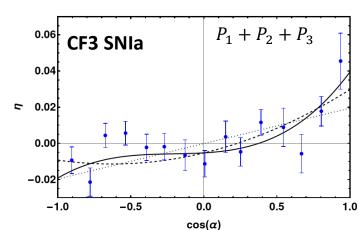








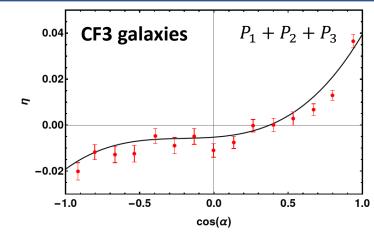




The same 3 Legendre coefficients can describe the measurements for both samples.

$$a_1 = (1.9 \pm 0.1) \times 10^{-2}$$

 $a_2 = (1.1 \pm 0.1) \times 10^{-2}$
 $a_3 = (1.1 \pm 0.1) \times 10^{-2}$

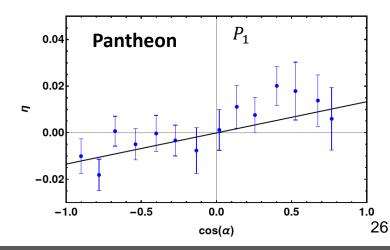


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 $a_3 = (1.1 \pm 0.1) \times 10^{-2}$

□ For Pantheon sample, the dipole is enough, since the error is larger (due to the small number of objects), and there is no objects in the interesting direction.

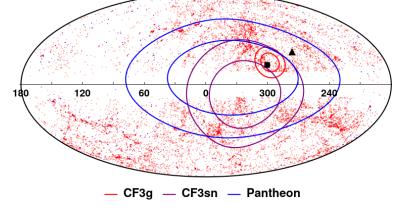


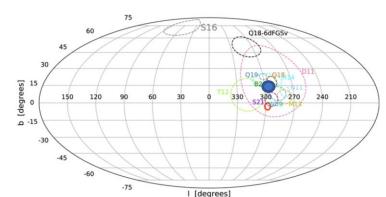
Bulk Motion

☐ The bulk velocity is related to the dipole.

$$v_b = \frac{a_1 \ln 10}{\langle (1+z)/z \rangle} = 318 \pm 22 \text{ km/s}$$

(S. S. Boruah, M. J. Hudson, and G. Lavaux. 2020)





(T. Hong et al. 2014)

 $v_b = 252 \pm 11 \,\mathrm{km/s}$

 $v_b = 292 \pm 28 \text{ km/s}$

Bulk flow direction

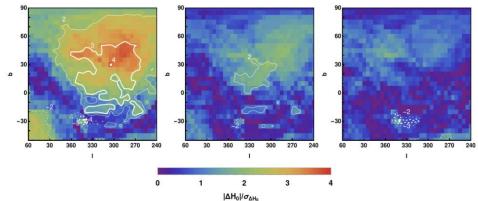
(Qin et al. 2021)

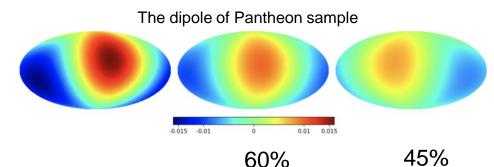
Example of Practical Use of η: subtract peculiar velocity fields

lacktriangle Since (η) is related to the radial peculiar velocity.

$$\mu = 5\log\left(\frac{z}{H_0}\right) + 25 - 5\eta$$

The map of the signal-to-noise ratio for ΔH_0 , between each direction and its opposite direction





Conclusion

- ☐ Designed a new observable measuring the *expansion rate fluctuations* and determined its multipolar structure.
 - A simple dipole term is a poor representation of the angular fluctuations in the local expansion rate.
- ☐ Find a signal for the quadrupolar component for both galaxies and SNIa of the CF3.
- ☐ The maximum of the quadrupole is aligned with dipole, and it is axially symmetric.
- ☐ The octupole (from galaxy sample only) has a maximum in the direction of the dipole.
- ☐ Need to build on the current work with updated and enlarged datasets: CosmicFlows-4, Pantheon+.

Spherical Harmonic analysis: parameters

Sample	N _{pix}	l_d	b_d	\hat{C}_1 (10 ⁻⁴)	$\frac{\frac{\eta_{1_{max}} - \eta_{1_{min}} }{2}}{(10^{-2})}$	p-value (%)	l_q	b_q	\hat{C}_2 (10 ⁻⁴)	$\frac{\eta_{2max} - \eta_{2min} }{2} $ (10 ⁻²)	p-value (%)	l_t	b_t	\hat{C}_3 (10 ⁻⁴)	$\frac{\frac{\eta_{3max} - \eta_{3_{min}} }{2}}{(10^{-2})}$	p-value (%)
CF3 [0.01, 0.05]	192	291 ± 15	12 ± 7	3.0 ± 1.6	1.6	0.03	323 ± 16	9 ± 4	2.4 ± 0.9	1.9	0.12	289 ± 24	16 ± 15	0.1 ± 0.4	1.2	30.44
CF3 [0.01, 0.05]	48	283 ± 6	12 ± 5	5.3 ± 0.8	1.9	< 0.01	310 ± 11	4 ± 8	0.9 ± 0.3	1.1	< 0.01	284 ± 7	12 ± 5	0.5 ± 0.2	1.3	0.01
CF3g [0.01, 0.05]	48	286 ± 7	4 ± 6	7.0 ± 1.0	2.0	< 0.01	338 ± 8	22 ± 5	1.1 ± 0.4	1.3	< 0.01	255 ± 9	11 ± 5	0.7 ± 0.2	1.5	0.01
CF3 [0.01, 0.05]	12	285 ± 5	11 ± 4	5.1 ± 0.8	1.9	< 0.01	308 ± 7	1 ± 7	1 ± 0.3	1.1	< 0.01	-	-	-	-	-
CF3g [0.01, 0.05]	12	296 ± 6	18 ± 5	4.0 ± 0.6	1.7	< 0.01	323 ± 34	2 ± 17	1.3 ± 0.4	1.4	< 0.01	ı	-	-	-	-
CF3sn [0.01, 0.05]	12	322 ± 23	-8 ± 18	3.7 ± 1.5	1.5	0.27	343 ± 15	-8 ± 10	1.7 ± 1.4	1.7	2.90	-	-	-	-	-
Pantheon [0.01, 0.05]	12	334 ± 42	6 ± 20	3.5 ± 2.7	1.6	4.37	337	-5	0.6 ± 1.9	1.6	33.33	-	-	-	-	-
CF3 [0.01, 0.03]	12	279 ± 5	12 ± 5	7.8 ± 1.0	2.3	< 0.01	310 ± 8	11 ± 6	2.9 ± 0.6	1.9	< 0.01	-	-	-	-	-
CF3 [0.03, 0.05]	12	301 ± 15	10 ± 14	1.1 ± 0.7	1.0	0.03	277 ± 28	-12 ± 11	0.9 ± 0.3	1.0	2.04	-	-	-	-	-