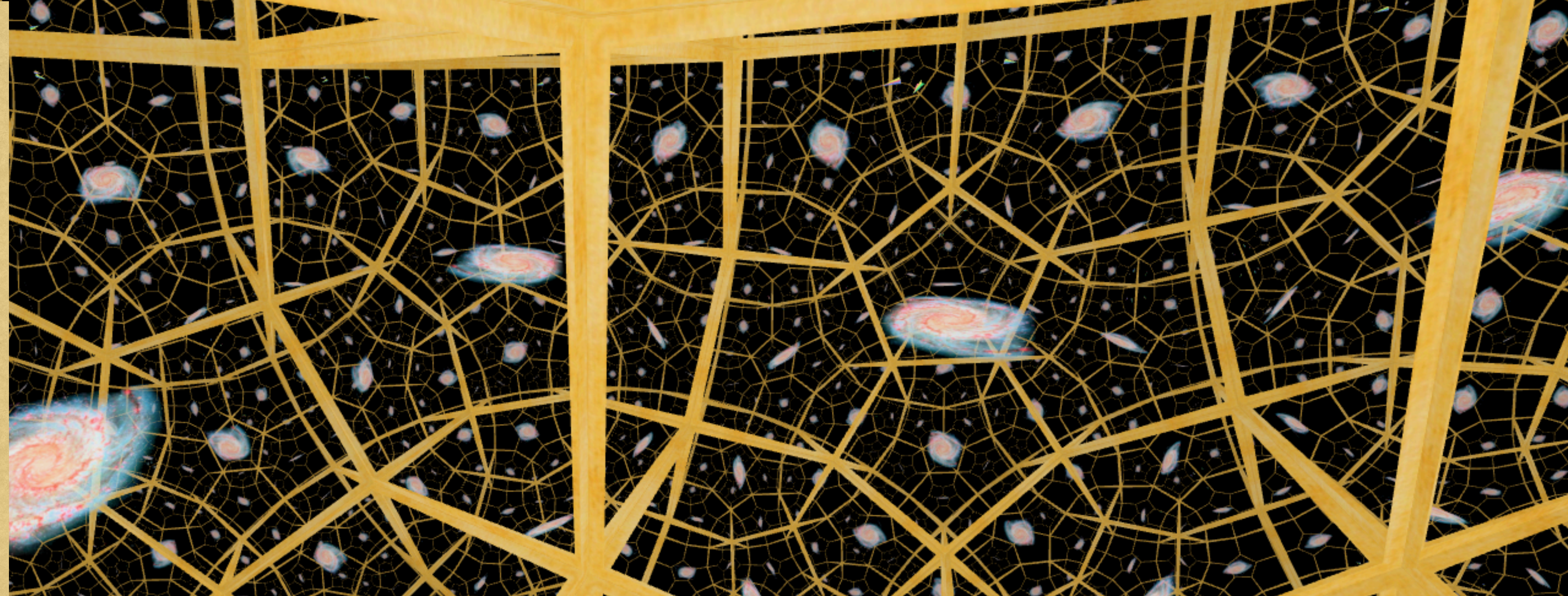




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Progress on Old and New
Themes in cosmology

3 May 2023



COSMOLOGICAL MODEL WITH EXPANSION BLIND TO THE SPATIAL CURVATURE

Quentin Vigneron

Nicolaus Copernicus University, Toruń, POLAND



Motivations:

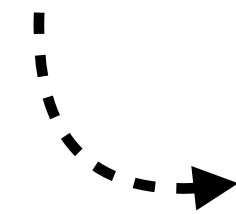
The theory:

Cosmology:

Topological term in the Einstein equation

(arXiv: [2204.13980](https://arxiv.org/abs/2204.13980))

$$R_{\mu\nu} - \bar{R}_{\mu\nu} = T_{\mu\nu} - T/2g_{\mu\nu} + \Lambda g_{\mu\nu}$$



Reference curvature related to
the spacetime topology

TOPOLOGICAL TERM IN THE EINSTEIN EQUATION

Rosen's bi-connection theory [e.g. Rosen (1980), *General relativity with a background metric*]

The theory is composed of:

- ▶ The **physical** Lorentzian structure $(g, {}^4\nabla)$ and its Riemann curvature tensor $R^\mu{}_{\alpha\beta\gamma}$
- ▶ A **reference, non-dynamical**, connection ${}^4\bar{\nabla}$ and its reference Riemann curvature tensor $\bar{R}^\mu{}_{\alpha\beta\gamma}$

Modified Einstein's equation:

$$R_{\mu\nu} - \bar{R}_{\mu\nu} = \kappa T_{\mu\nu}^* + \Lambda g_{\mu\nu} \quad ; \quad g^{\mu\nu} \nabla_\mu \bar{R}_{\nu\alpha} - \frac{1}{2} g^{\mu\nu} \nabla_\alpha \bar{R}_{\mu\nu} = 0$$

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My choice of reference metric: the reference connection determines the spacetime topology

Vigneron Q., 2022c, arXiv: [2204.13980](https://arxiv.org/abs/2204.13980)

$$\bar{R}_{\mu\nu} = \begin{pmatrix} 0 & 0 \\ 0 & \bar{\mathcal{R}}_{ij}(x^k) \end{pmatrix} \quad \text{where } \bar{\mathcal{R}}_{ij} = \text{(maximally symmetric) curvature of the spatial covering space } \tilde{\Sigma}.$$

$\tilde{\mathcal{M}} = \mathbb{R} \times$	\mathbb{R}^3	S^3	H^3	others
$\bar{R}_{\mu\nu}$	$\mathbf{0}$	$\text{diag}(0; 2\bar{h}_{ij})$	$-\text{diag}(0; 2\bar{h}_{ij})$	$\neq 0$

Motivations:

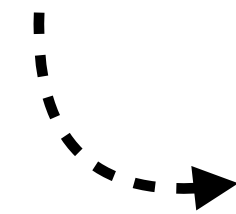
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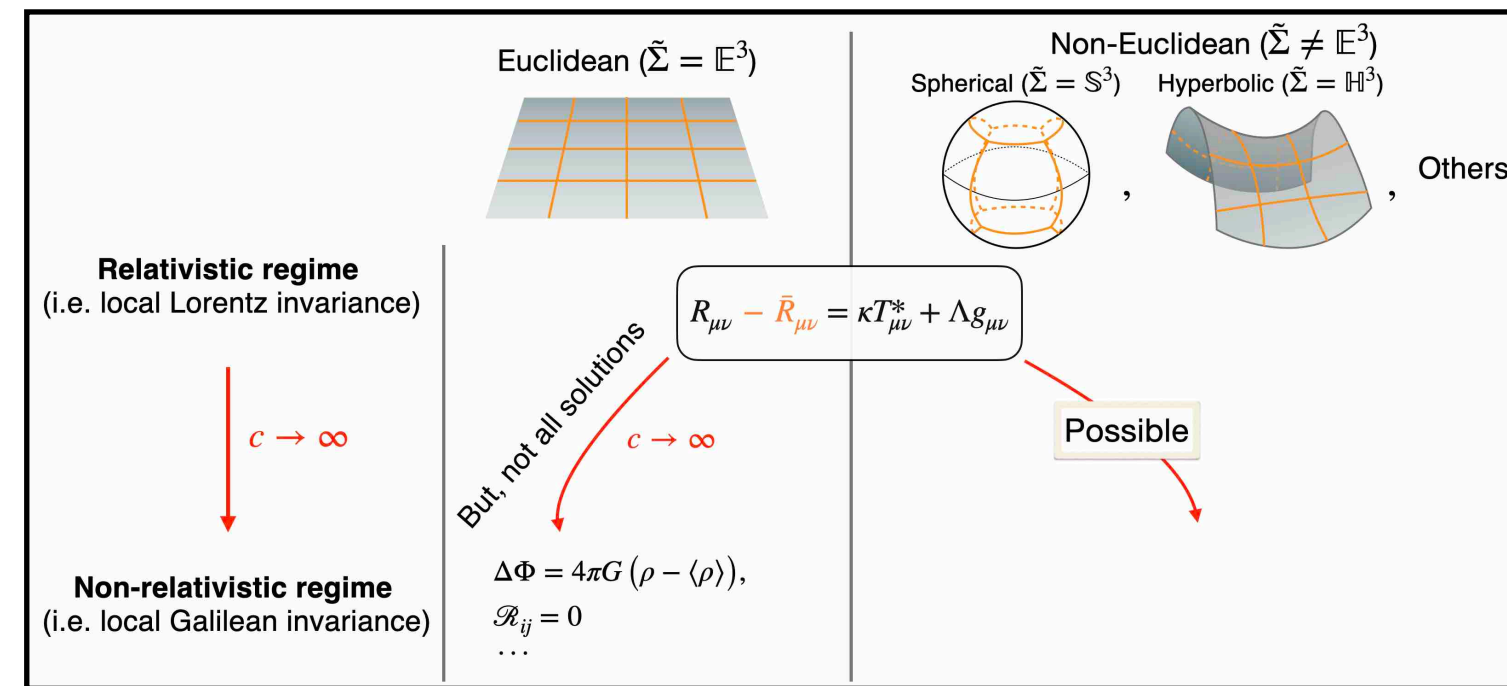
Reference curvature related to
the spacetime topology

Motivations:

The theory:

Cosmology:

Non-relativistic limit



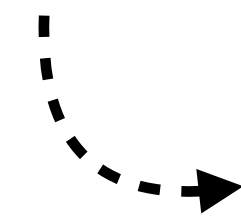
First order covariant formalism

$$S_{\text{EH}} = \int_{\mathcal{M}} \sqrt{-g} g^{\mu\nu} R_{\mu\nu} dx^4$$

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Reference curvature related to the spacetime topology

MOTIVATION 1: EXISTENCE OF A NON-RELATIVISTIC LIMIT

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$c \rightarrow \infty$

$$\Delta\Phi = 4\pi G (\rho - \langle\rho\rangle),$$

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...

MOTIVATION 1: EXISTENCE OF A NON-RELATIVISTIC LIMIT

But, not all solutions

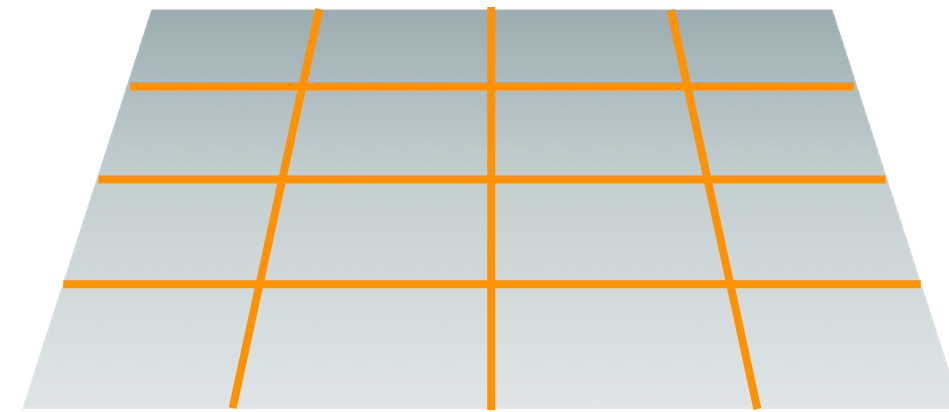
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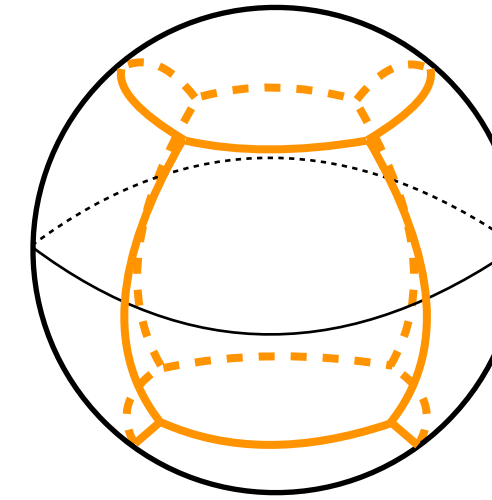
MOTIVATION 1: EXISTENCE OF A NON-RELATIVISTIC LIMIT

Euclidean ($\tilde{\Sigma} = \mathbb{E}^3$)

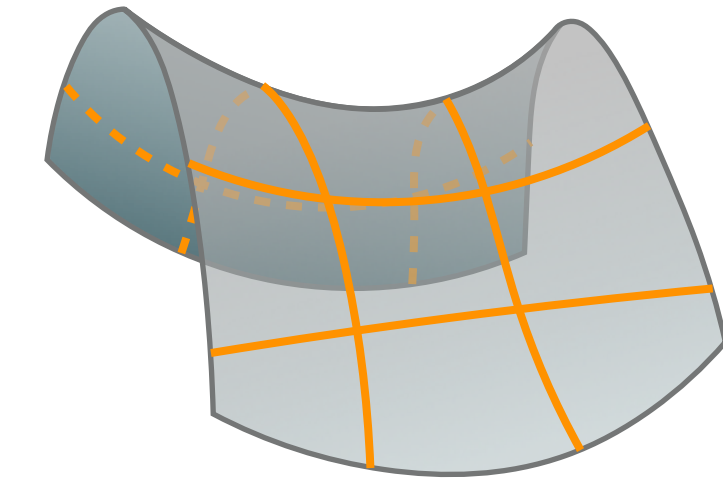


Non-Euclidean ($\tilde{\Sigma} \neq \mathbb{E}^3$)

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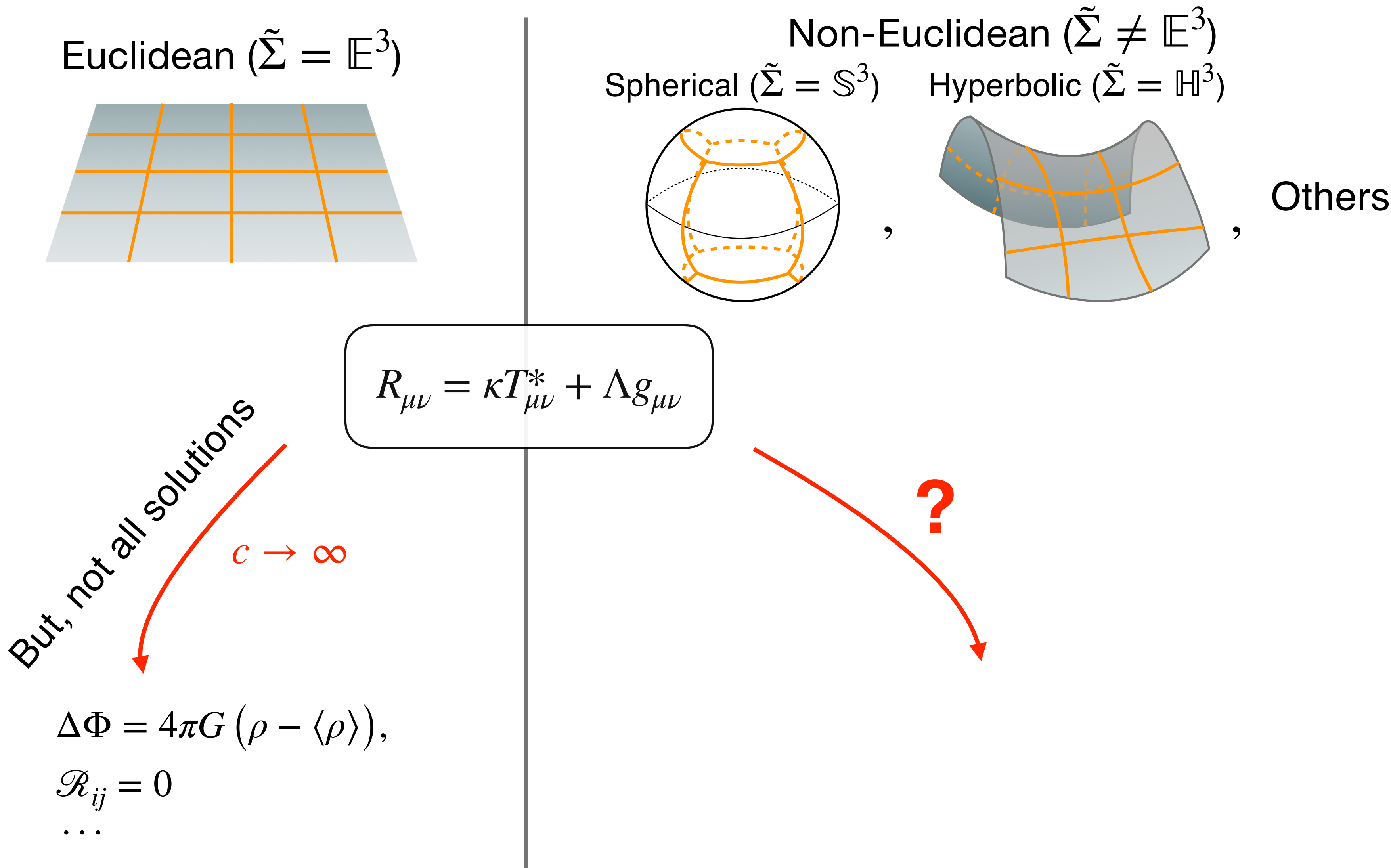
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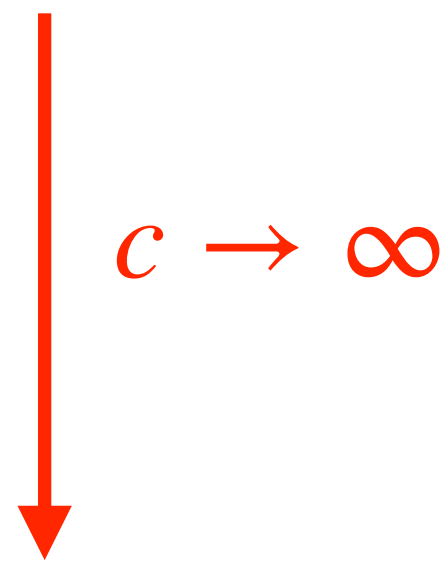
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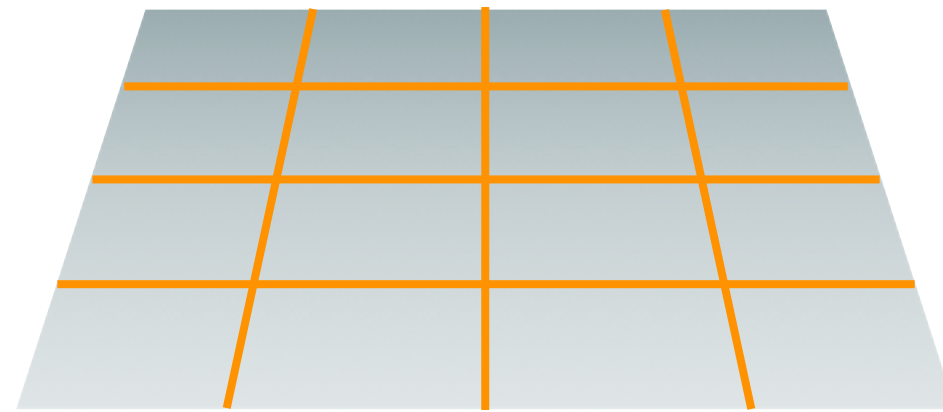
Limit defined for any topology

Relativistic regime
(i.e. local Lorentz invariance)



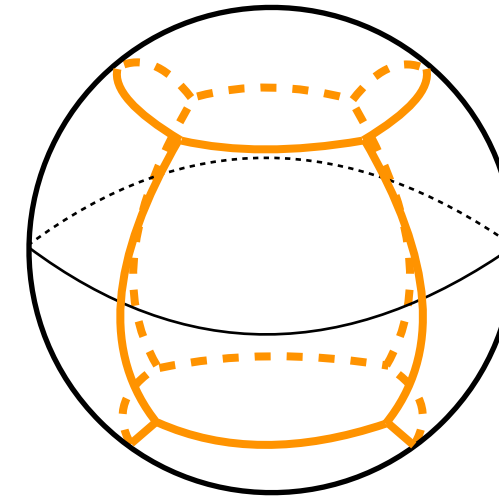
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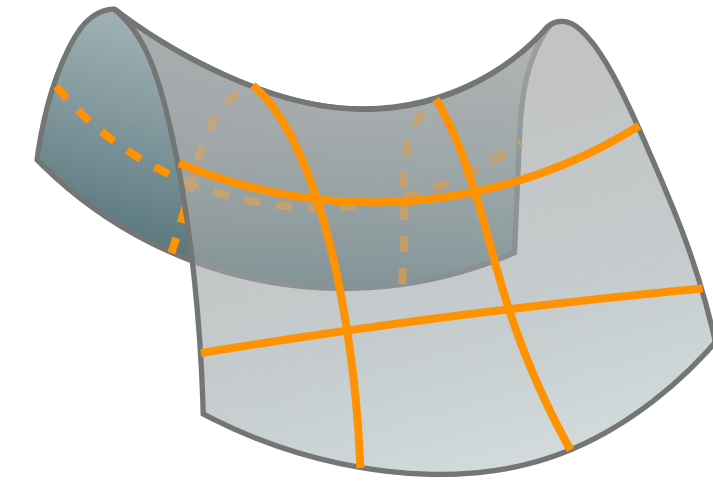


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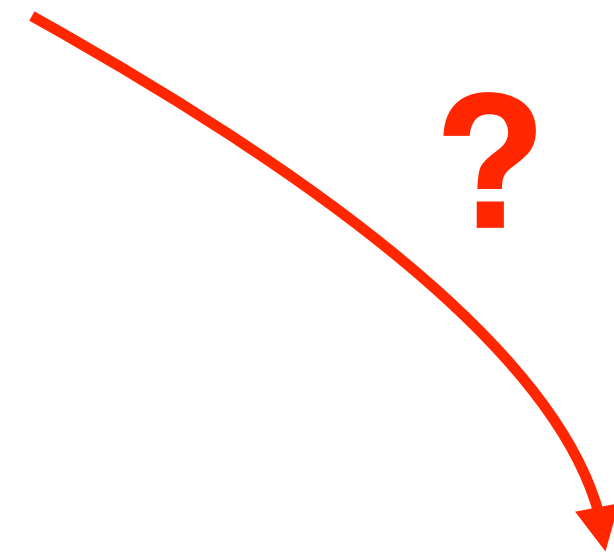
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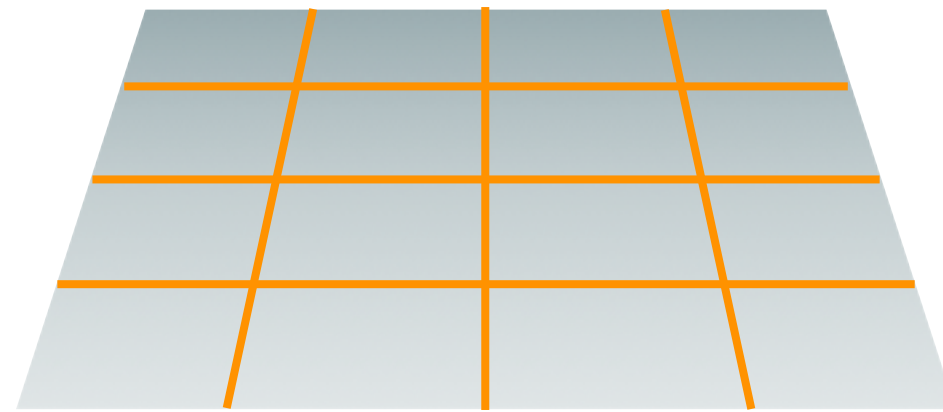
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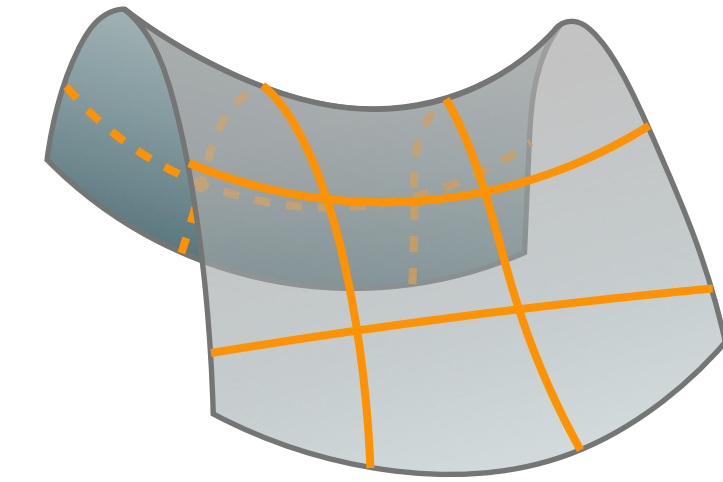
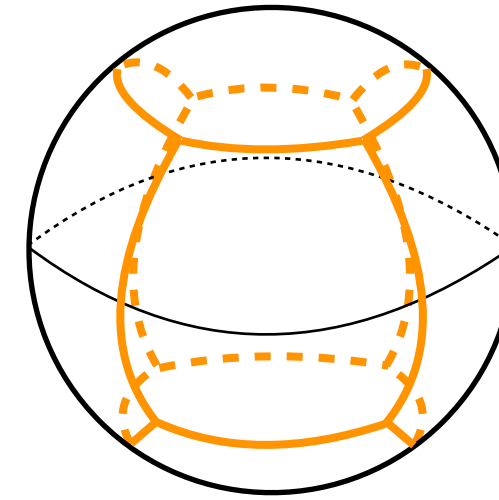
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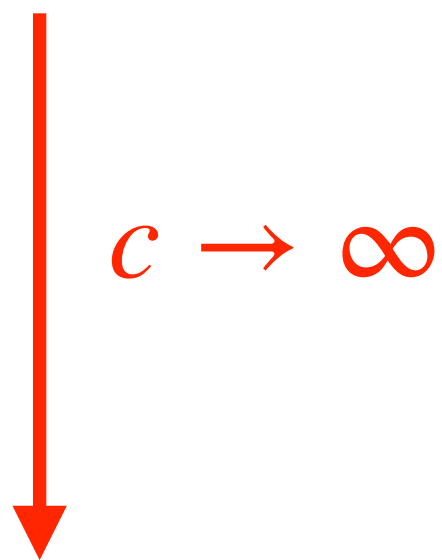
Limit does not exist

(arXiv: [2204.13980](https://arxiv.org/abs/2204.13980))

MOTIVATION 1: EXISTENCE OF A NON-RELATIVISTIC LIMIT

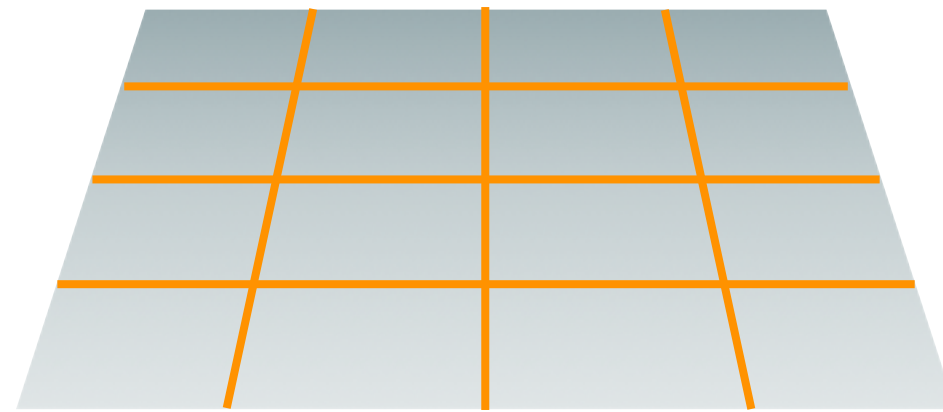
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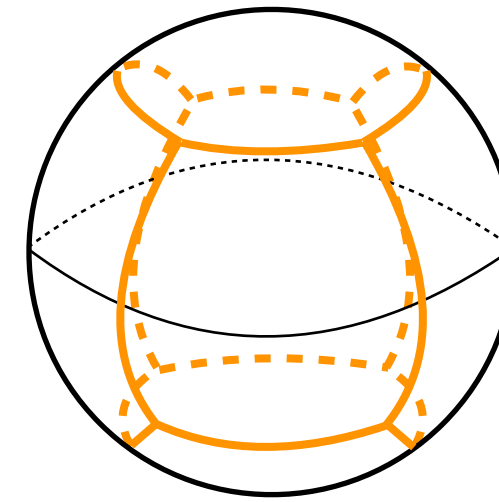
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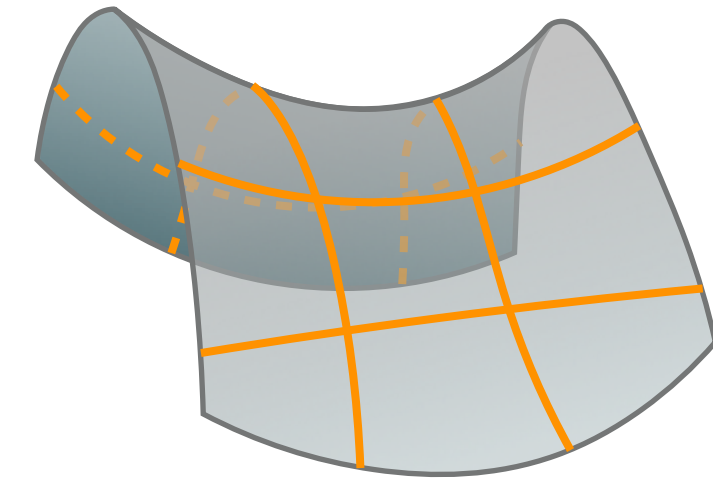


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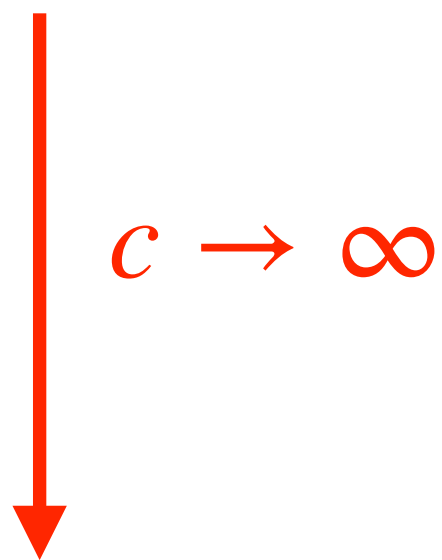
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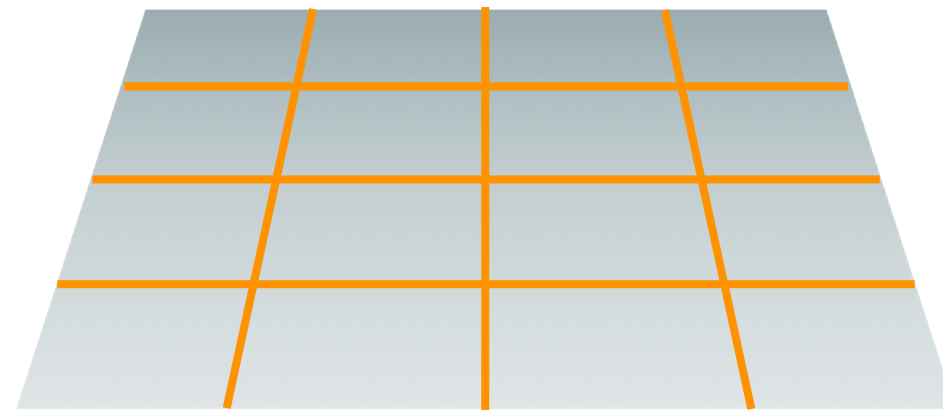
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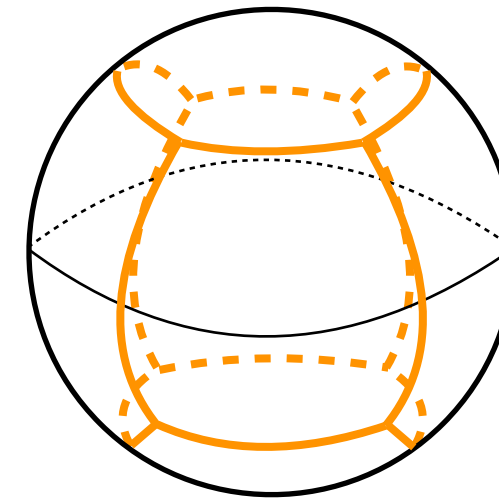
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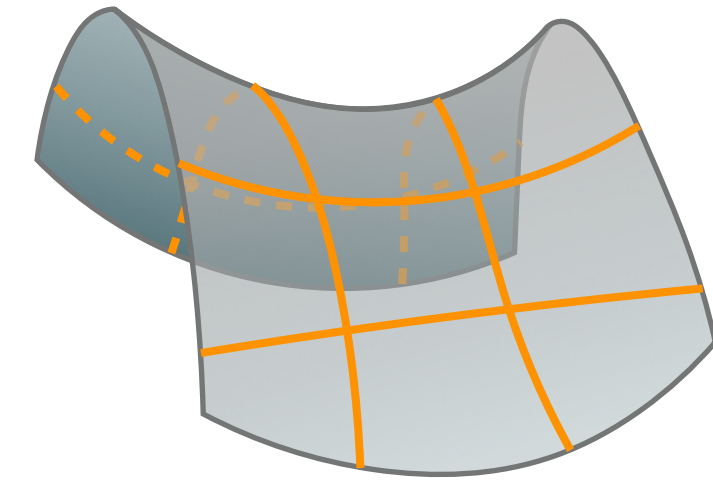


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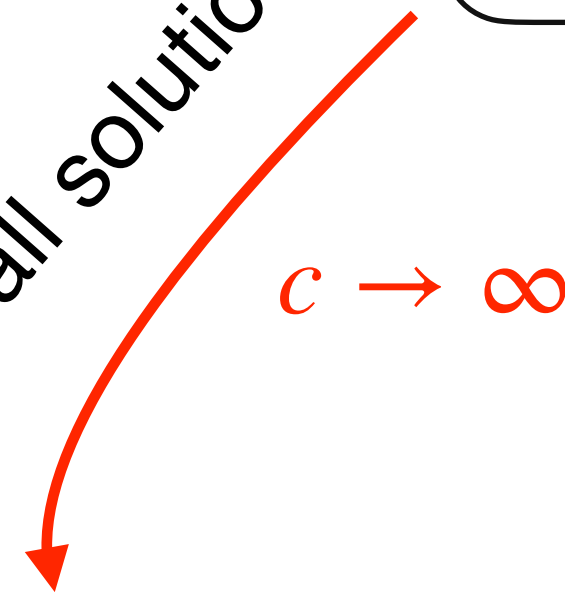
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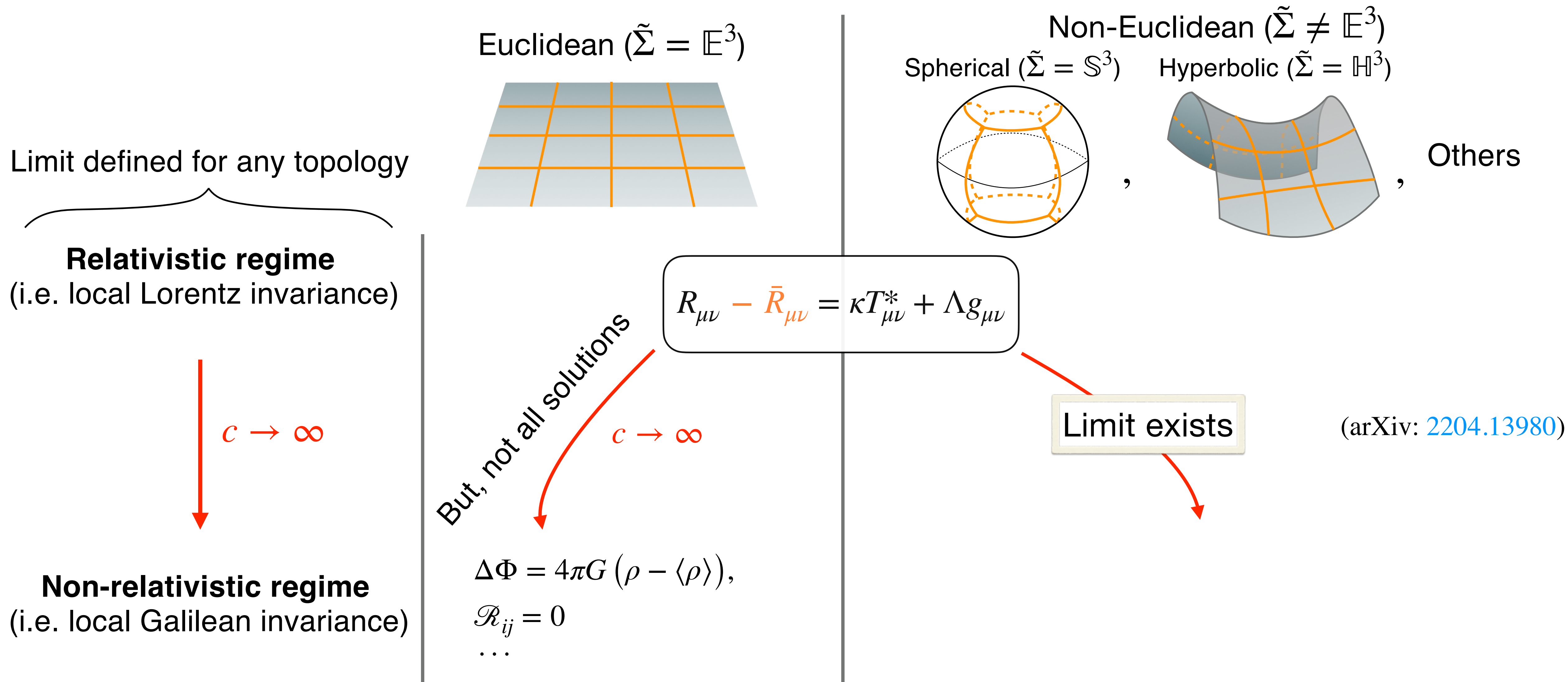
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MOTIVATION 1: EXISTENCE OF A NON-RELATIVISTIC LIMIT



Summary: Requiring compatibility between the relativistic and the non-relativistic regime in any topology.

MOTIVATION 2: FIRST ORDER COVARIANT FORMULATION

Action of general relativity:

- For a mathematically **well-posed** action
- For a **covariant first order** Lagrangian

$$\begin{aligned} S_{\text{EH} + \text{GHY}} &= \int_{\mathcal{M}} \sqrt{-g} g^{\mu\nu} R_{\mu\nu} dx^4 + 2 \int_{\partial\mathcal{M}} \sqrt{|h|} n^\mu h^{\alpha\beta} \left(\bar{\nabla}_\mu g_{\alpha\beta} - \bar{\nabla}_\alpha g_{\beta\mu} \right) dx^3 \\ &= \int_{\mathcal{M}} \sqrt{|g|} \mathcal{C}_{\mu[\beta}^\alpha \mathcal{C}_{\nu]\alpha}^\beta g^{\mu\nu} dx^4 \quad \text{where } \mathcal{C}_{\mu\nu}^\alpha := \Gamma_{\mu\nu}^\alpha - \bar{\Gamma}_{\mu\nu}^\alpha \text{ and } \bar{R}^\alpha_{\beta\mu\nu} = 0. \end{aligned}$$

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 &= \int_{\mathcal{M}} \sqrt{|g|} \mathcal{C}^\alpha_{\mu[\beta} \mathcal{C}^\beta_{\nu]\alpha} g^{\mu\nu} dx^4
 \end{aligned}$$

Einstein-Hilbert
Gibbons-Hawking-York term

where $\mathcal{C}^\alpha_{\mu\nu} := \Gamma^\alpha_{\mu\nu} - \bar{\Gamma}^\alpha_{\mu\nu}$ and $\bar{R}^\alpha_{\beta\mu\nu} = 0$.

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First order Lagrangian

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Hypothesis: $\bar{\Gamma}^{\alpha}_{\mu\nu}$ should be adapted as function of the spacetime topology, i.e. $\bar{R}^{\alpha}_{\beta\mu\nu} \neq 0$ for non-Euclidean topologies.

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$$\delta S_{\text{EH} + \text{GHY}} \left(\bar{R}_{\mu\nu} \neq 0 \right) = \int_{\mathcal{M}} \sqrt{-g} \left[R_{\mu\nu} - \bar{R}_{\mu\nu} - \frac{1}{2} \left(R_{\alpha\beta} - \bar{R}_{\alpha\beta} \right) g^{\alpha\beta} g_{\mu\nu} \right] \delta g^{\mu\nu} dx^4$$

Bi-connection equation : $R_{\mu\nu} - \bar{R}_{\mu\nu} = \kappa T^*_{\mu\nu} + \Lambda g_{\mu\nu}$

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Einstein-Hilbert Gibbons-Hawking-York term

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First order Lagrangian

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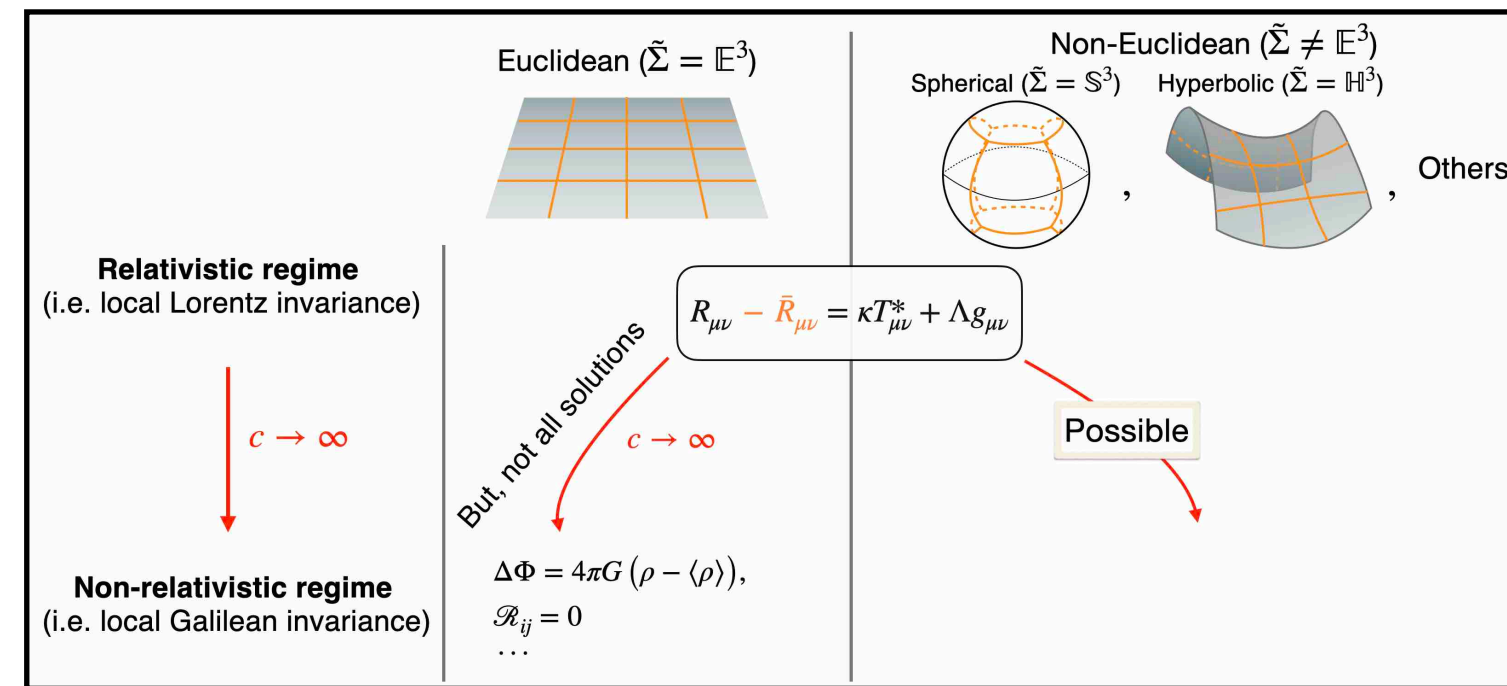
Summary: Having a well-posed and first order action of gravitation, for any topology.

Motivations:

The theory:

Cosmology:

Non-relativistic limit



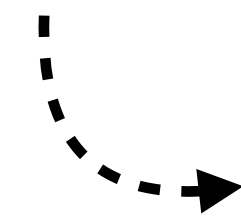
First order covariant formalism

$$S_{\text{EH} + \text{GHY}} = \int_{\mathcal{M}} \sqrt{|g|} \mathcal{E}^{\alpha}_{\mu[\beta} \mathcal{E}^{\beta}_{\nu]\alpha} g^{\mu\nu}$$

Topological term in the Einstein equation

(arXiv: [2204.13980](#))

$$R_{\mu\nu} - \bar{R}_{\mu\nu} = T_{\mu\nu} - T/2 g_{\mu\nu} + \Lambda g_{\mu\nu}$$



Reference curvature related to the spacetime topology

Expansion blind to the curvature (arXiv: [2212.00675](#))

$$\Omega_{\neq K} = 1, \quad \forall \Omega_K$$

CONSEQUENCE FOR THE ROLE OF SPATIAL CURVATURE

1. Exact homogeneous and isotropic solution:

$$\Lambda\text{CDM: } \begin{cases} \Omega_{\neq K} + \Omega_K = 1, \\ q = \Omega_m/2 + \Omega_{\text{rad}} - \Omega_\Lambda \end{cases}$$

2. Weak field limit: for $\Pi = 0 = \Pi_i$

$$(\Delta + 3K) \Psi = 4\pi G a^2 \delta\rho_\Delta,$$

$$\Lambda\text{CDM: } (\Delta + 2K) Q_i = -16\pi G a^2 q_i,$$

$$f''_{ij} + 2\mathcal{H}f'_{ij} + (2K - \Delta)f_{ij} = 8\pi G a^2 \Pi_{ij},$$

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—> Dynamical and geometrical effects of spatial curvature

$$\text{Bi-connection: } \begin{cases} \Omega_{\neq K} = 1, \quad \forall \Omega_K \\ q = \Omega_m/2 + \Omega_{\text{rad}} - \Omega_\Lambda \end{cases}$$

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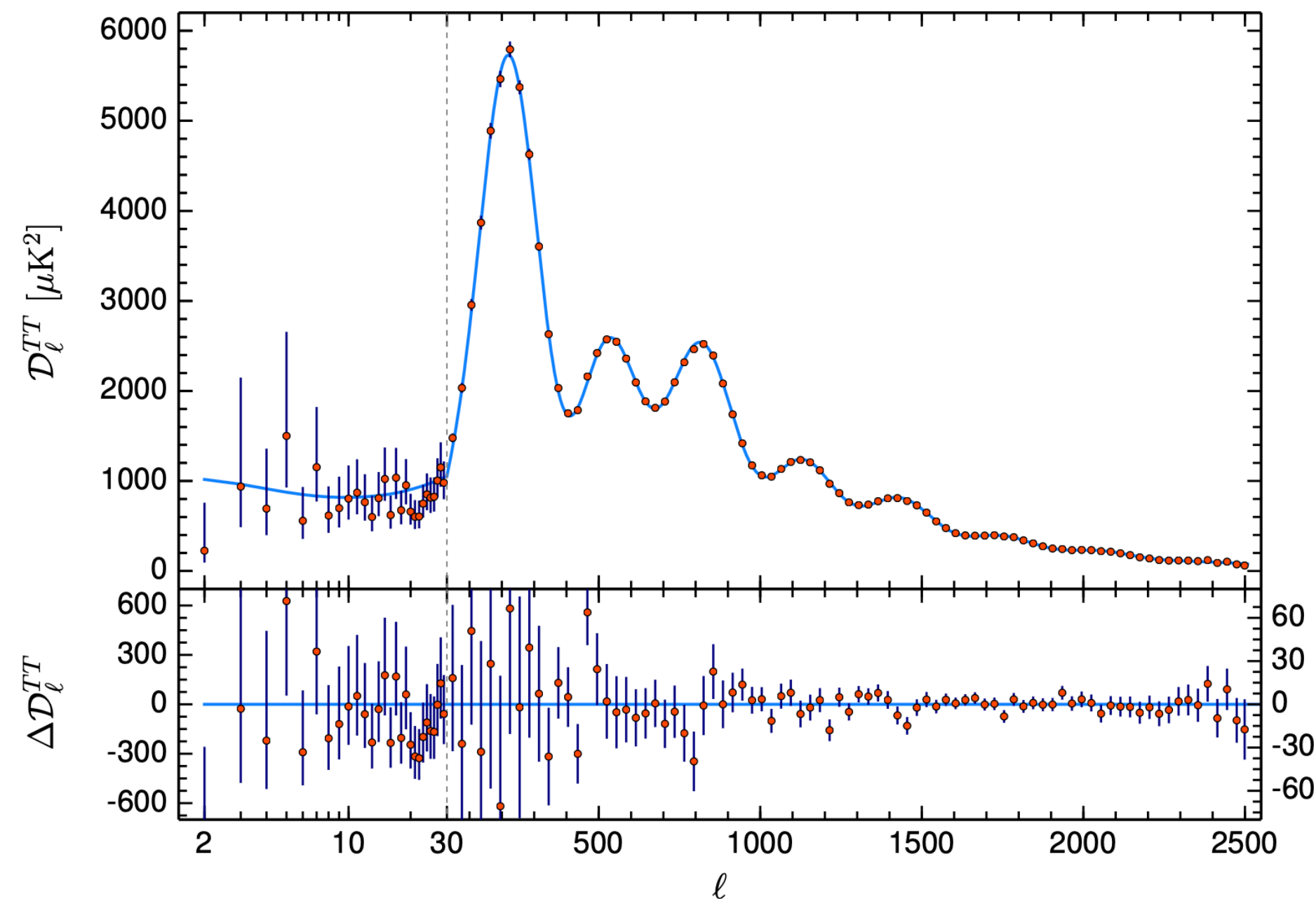
Strengths of this cosmological model:

1. Origin of the modifications not related to cosmology or tensions of ΛCDM .
2. Has the same number of free-parameters than the ΛCDM model.

CONSEQUENCE FOR THE ROLE OF SPATIAL CURVATURE

A possibility:

- ▶ A positive curvature ($\Omega_K \sim -0.05$) gives a better fit of the CMB power spectrum alone.
[e.g. Di-Valentino et al. (2020): [1911.02087](#)]
- ▶ But, it increases the Hubble tension: ΔH_0 from $67 \leftrightarrow 73$ to $55 \leftrightarrow 73$ km/s/Mpc: "**Curvature tension**"



Planck VI (2018)

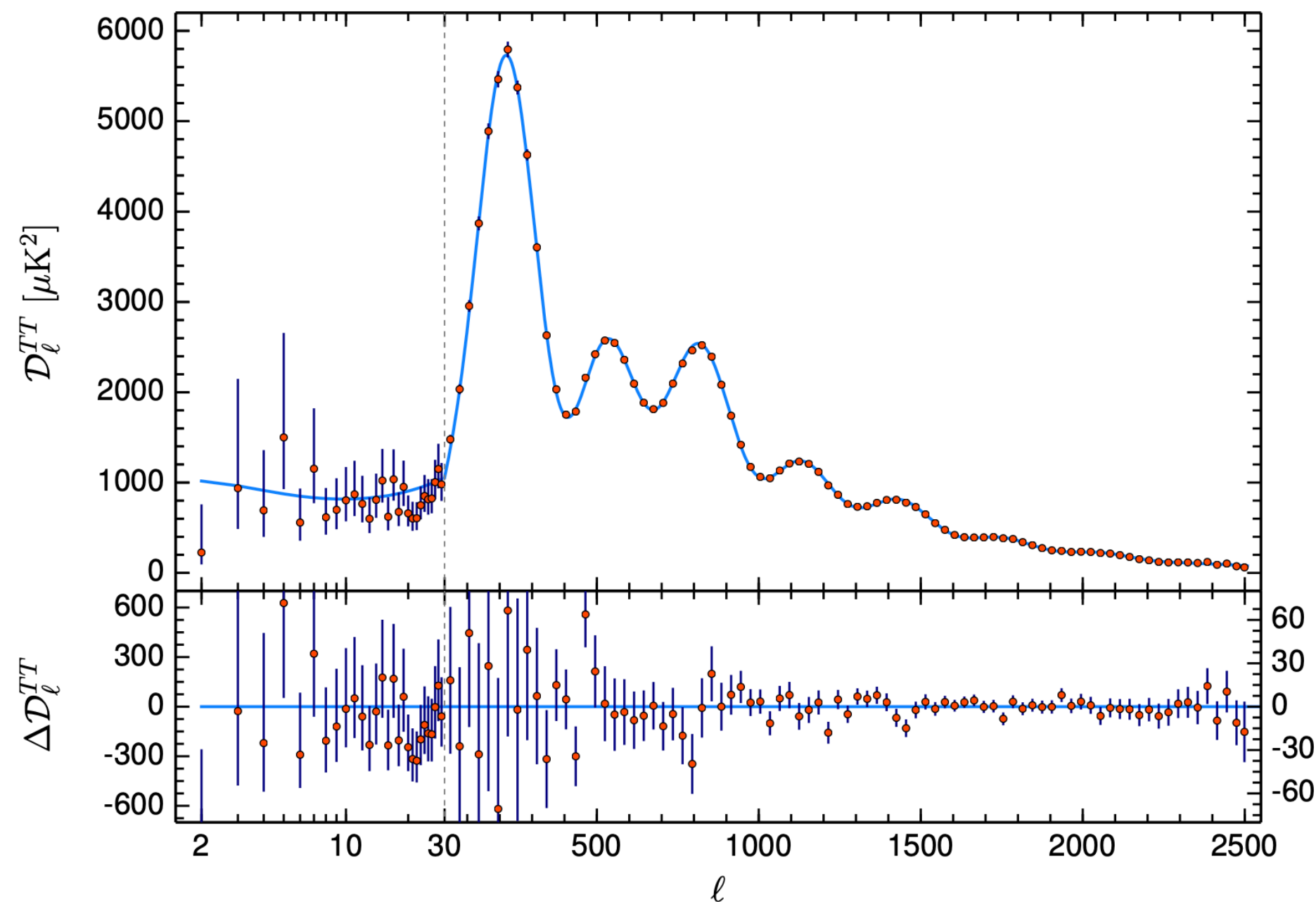
$$S_\chi(x) = \begin{cases} \sin x, & \Omega_K < 0 \\ x, & \Omega_K = 0 \\ \sinh x, & \Omega_K > 0 \end{cases}$$

$$d_A(z) = \frac{1}{\sqrt{|\Omega_K|} H_0 (1+z)} S_\chi \left(\sqrt{|\Omega_K|} H_0 \int_0^z \frac{dz'}{H(z')} \right)$$

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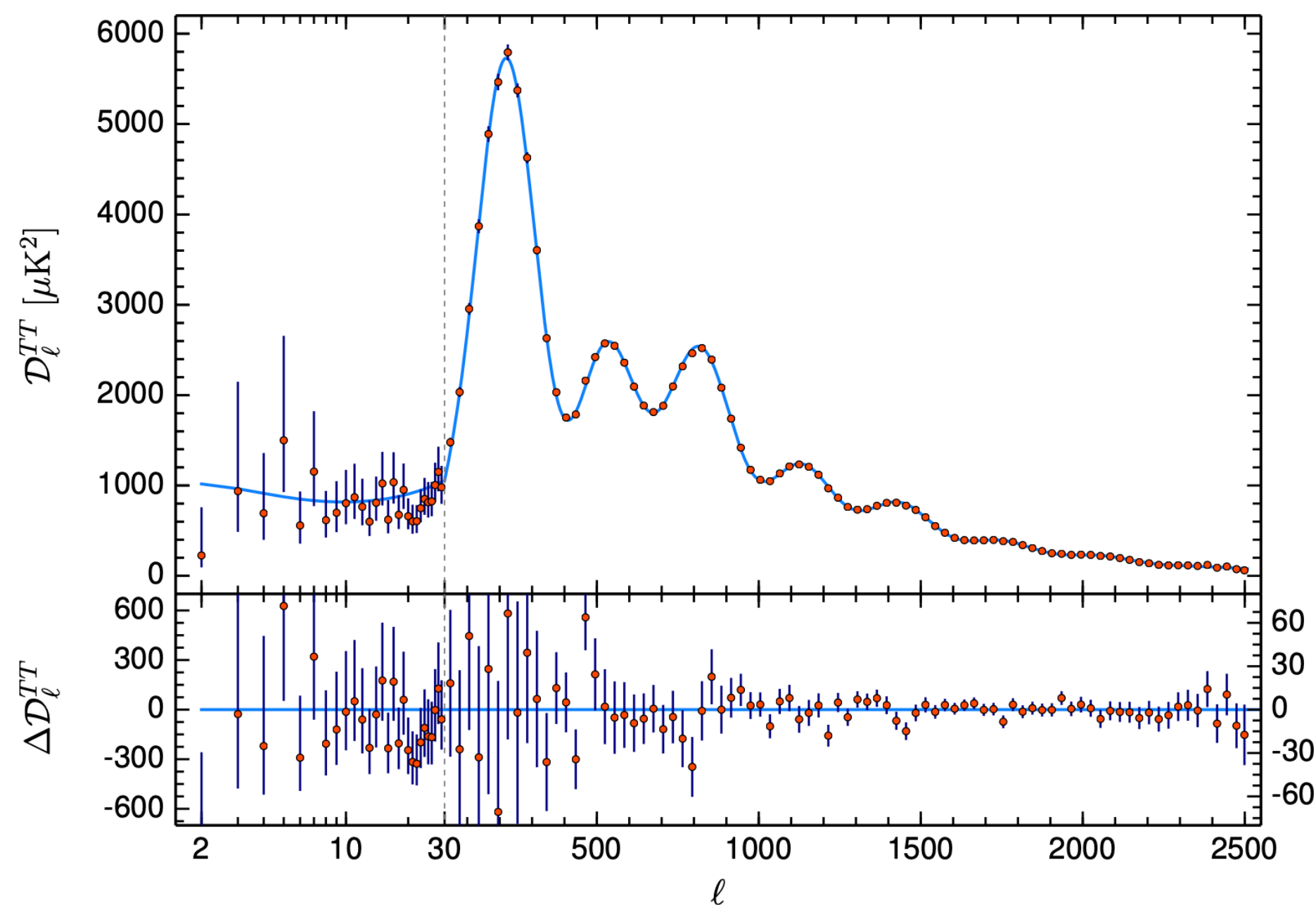
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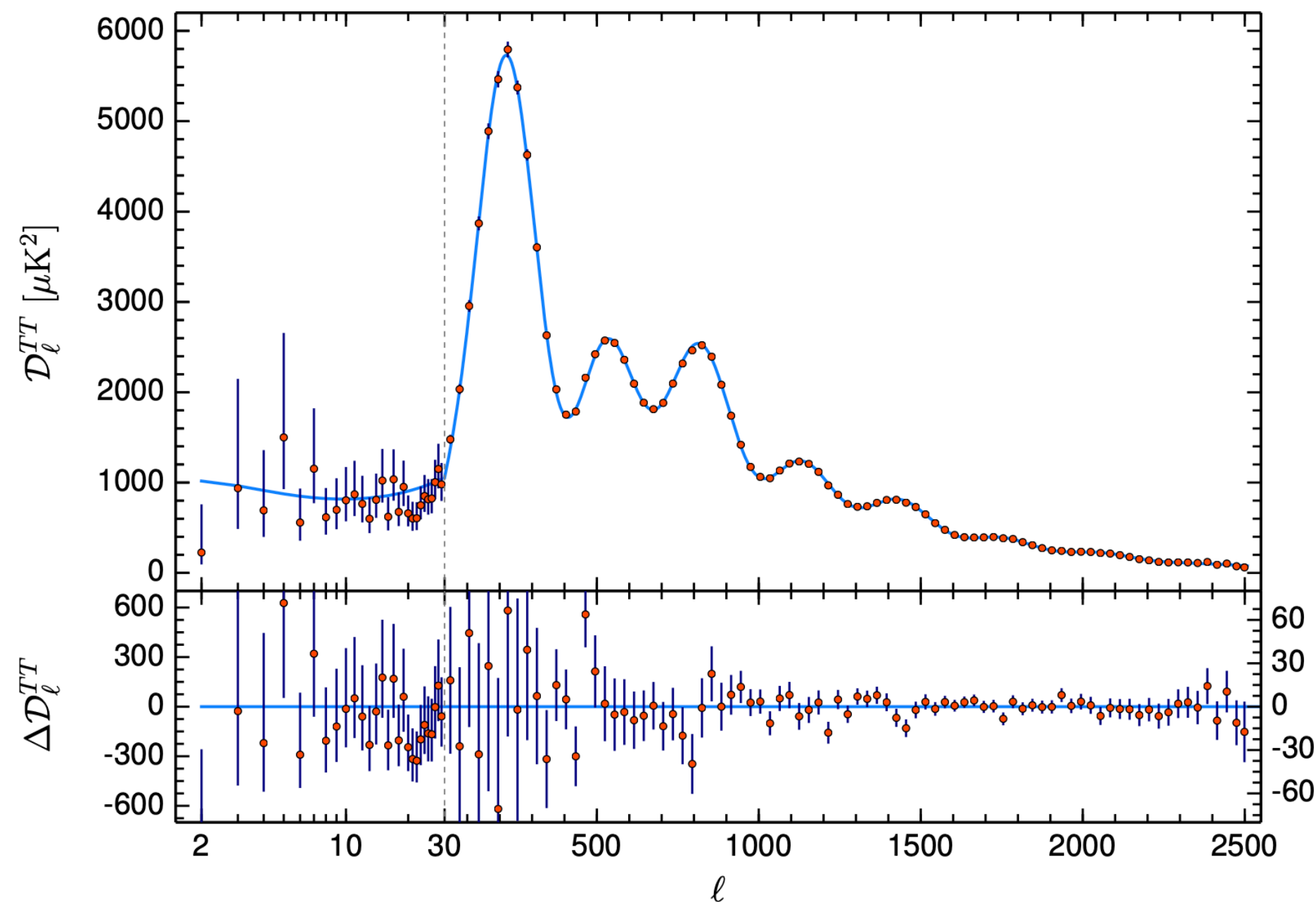
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Dynamical

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Geometrical Dynamical

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My guess: A priori NO, because the curvature tension is mainly due to geometrical effects of curvature.

CONSEQUENCE FOR EARLY UNIVERSE

$$R_{\mu\nu} - \bar{R}_{\mu\nu} = \kappa T_{\mu\nu}^* + \Lambda g_{\mu\nu}$$

- Two additional gauge invariant variables:
- "Spatial velocity" scalar mode: \mathcal{C}
 - "Spatial velocity" vector mode: \mathcal{C}^i

Wave equation sourced by anisotropic stress:

$$\begin{aligned}\mathcal{C}'' + 2\mathcal{H}\mathcal{C}' - \Delta\mathcal{C} &= -a^2\kappa\Pi, \\ \mathcal{C}_i'' + 2\mathcal{H}\mathcal{C}_i' - 2\Delta\mathcal{C}_i &= -a^2\kappa\Pi_i.\end{aligned}$$

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► Reference Bardeen potentials:

- Acceleration: $\tilde{\Phi} := \Phi - \mathcal{C}'' - \mathcal{H}\mathcal{C}'$
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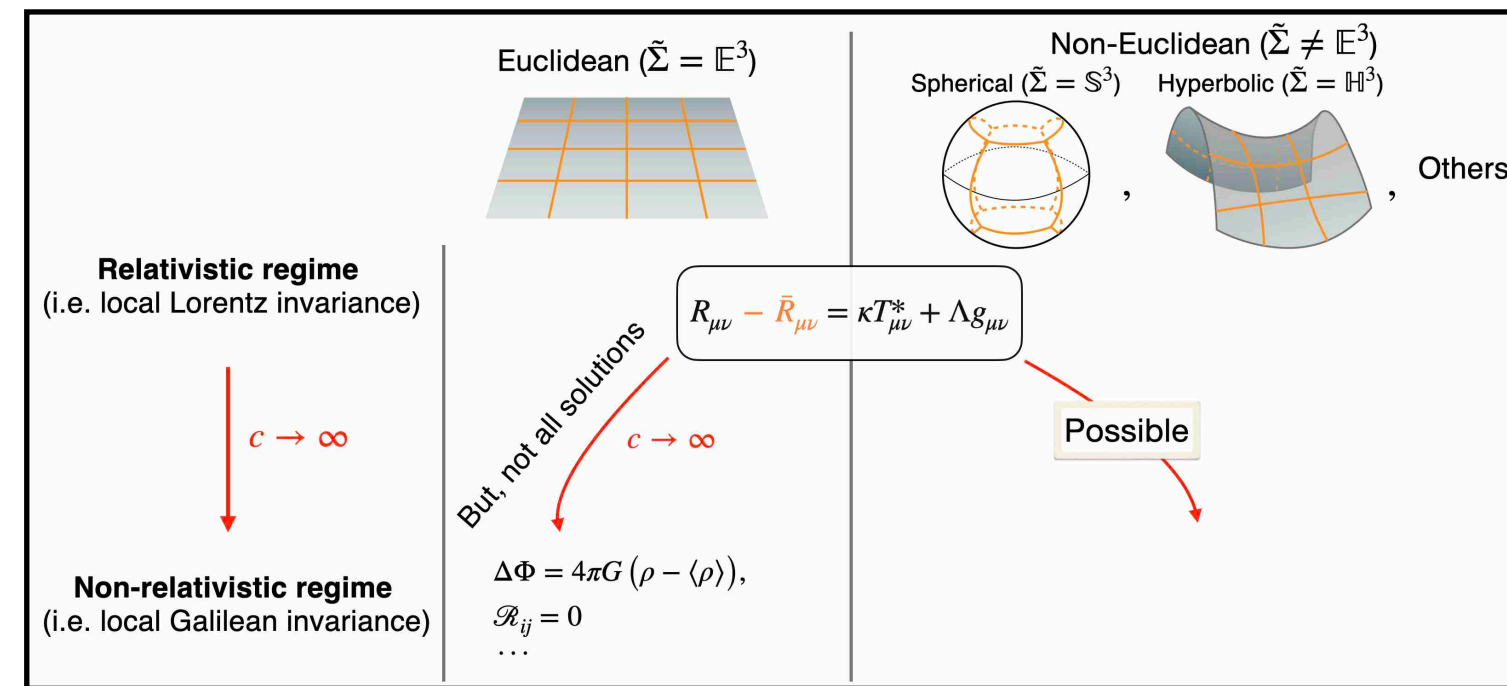
— — > **Early Universe analysis required.**

Motivations:

The theory:

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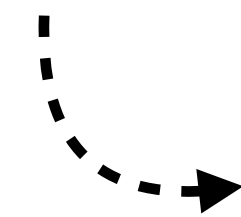
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