

# Index theorem on $T^2/Z_N$ orbifold with magnetic flux

**Maki Takeuchi**(Kobe Univ, Japan)

Collaborator :Tatsuo Kobayashi(Hokkaido Univ),Hajime Otsuka(Kyusyu Univ), Makoto Sakamoto(Kobe Univ), Yoshiyuki Tatsuta(Scuola Normale Superiore),Hikaru Uchida(Hokkaido Univ)

(arXiv 2209.xxxxx)

Workshop on the Standard Model and Beyond  
Corfu, Greece, 7 September 2022

# Table of contents

1. Introduction
2. Purpose of my talk
3.  $T^2/Z_N$  orbifold model
4. Blow-up of  $T^2/Z_N$  orbifold
5. Conclusion

# 1. Introduction

**Question: Why are there 3 generations of quarks and leptons?**

# 1. Introduction

Question: Why are there 3 generations of quarks and leptons?

Accidental?



**Standard Model**

Has physical meaning ?



**Extra dimension model**

# 1. Introduction

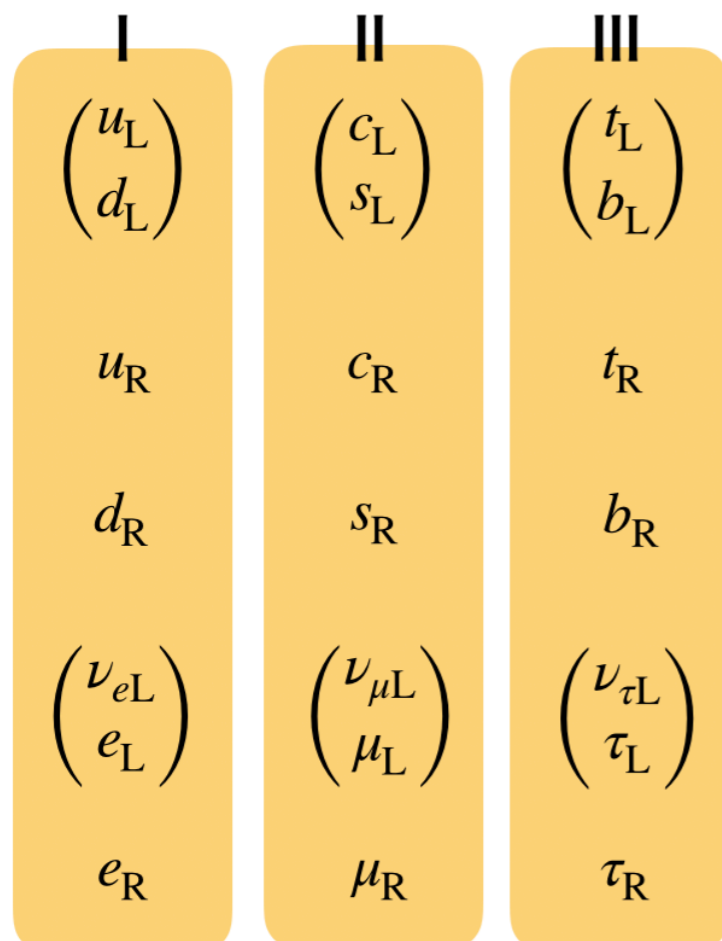
Question: Why are there 3 generations of quarks and leptons?

Accidental? 

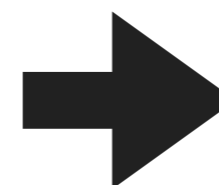
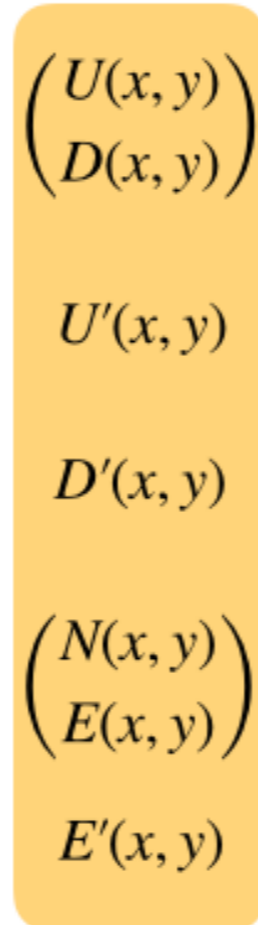
Has physical meaning? 

**Standard Model**

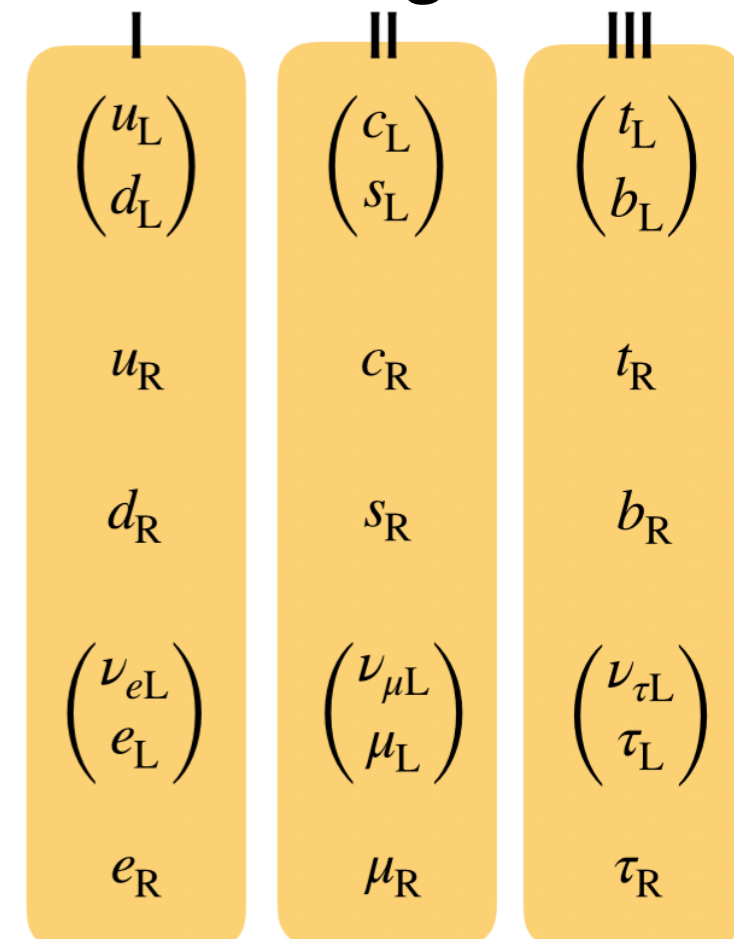
**Extra dimension model**



6d chiral fermion



Zero modes = generation



**Degenerate**

# 1. Introduction

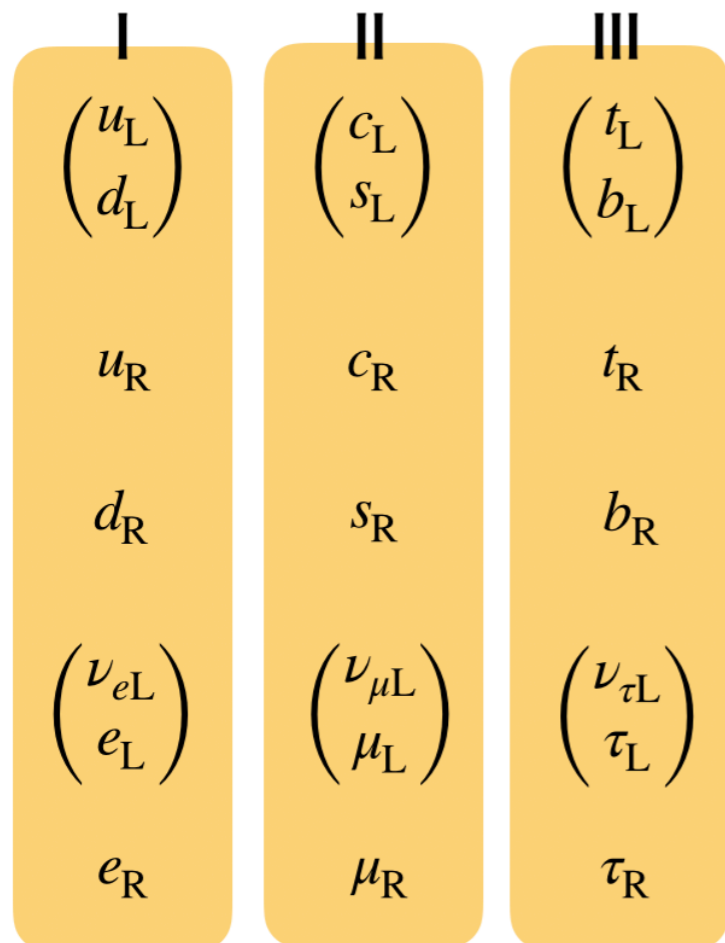
Question: Why are there 3 generations of quarks and leptons?

Accidental? 

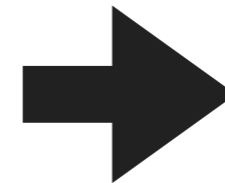
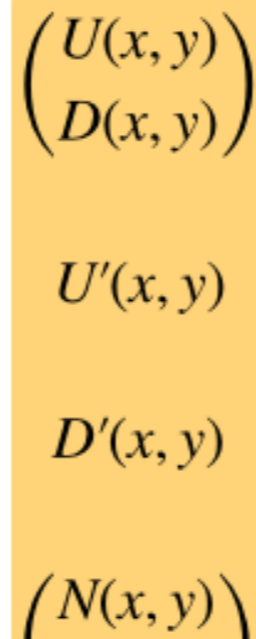
Has physical meaning? 

**Standard Model**

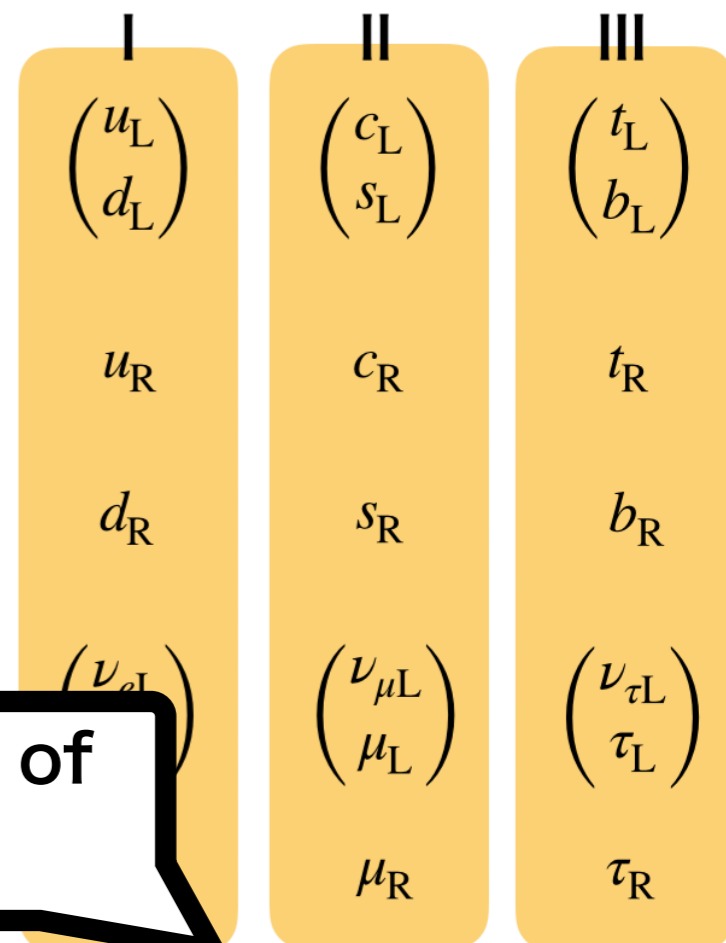
**Extra dimension model**



6d chiral fermion



Zero modes



Decided by topology of extra dimension

Degenerate

# 1. Introduction

Atiyah-Singer index theorem

$$n_+ - n_- \propto \int_{\mathcal{M}} F^{D/2} + \int_{\mathcal{M}} R^{D/4} + (F, R \text{ mixing term})$$

$n_{\pm}$ : chiral zero modes #     $D$ : dimension     $\mathcal{M}$ : a smooth manifold

# 1. Introduction

Atiyah-Singer index theorem

$$n_+ - n_- \propto \int_{\mathcal{M}} F^{D/2} + \int_{\mathcal{M}} R^{D/4} + (F, R \text{ mixing term})$$

Contribution of flux  
 $D=2n$

Contribution of curvature  
 $D=4n$





## 2. Purpose of my talk

In previous paper, we obtain the following formula on  $T^2/Z_N$  orbifolds

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

$M$  : flux quanta,  $V_+$  : sum of winding number at fixed points

Makoto Sakamoto, **Maki Takeuchi**, Yoshiyuki Tatsuta,  
**Phys. Rev. D** 102 (2020) 025008

## 2. Purpose of my talk

In previous paper, we obtain the following formula on  $T^2/Z_N$  orbifolds

Contribution of flux

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

$M$  : flux quanta,  $V_+$  : sum of winding numbers of first twisted sector

What's physical meaning?

# 2. Purpose of my talk

## Result of my talk

Contribution of flux

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

$M$  : flux quanta,  $V_+$  : sum of winding number at fixed points

Winding number contains **localized flux** and **curvature**.  
+1 comes from the sum of **localized curvatures**.

# 2. Purpose of my talk

## Result of my talk

Contribution of flux

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

$M$  : flux quanta,  $V_+$  : sum of winding number at fixed points

+1 removes contribution of localized curvatures in  $V_+$ .

# 2. Purpose of my talk

## Result of my talk

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

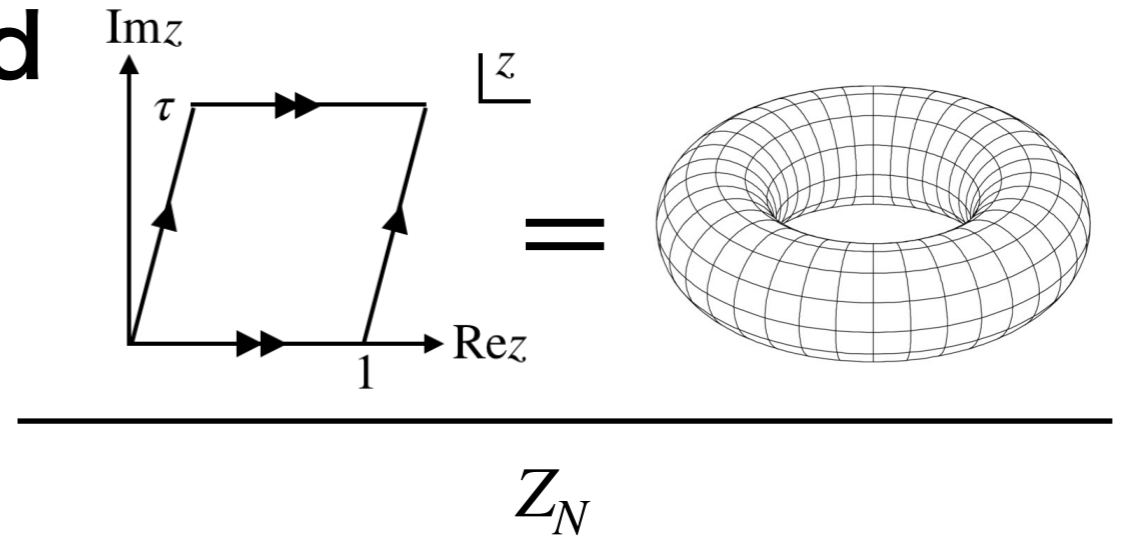
$M$  : flux quanta,  $V_+$  : sum of winding number at fixed points

**Only contribution of flux!**

# 3. $T^2/Z_N$ orbifold model

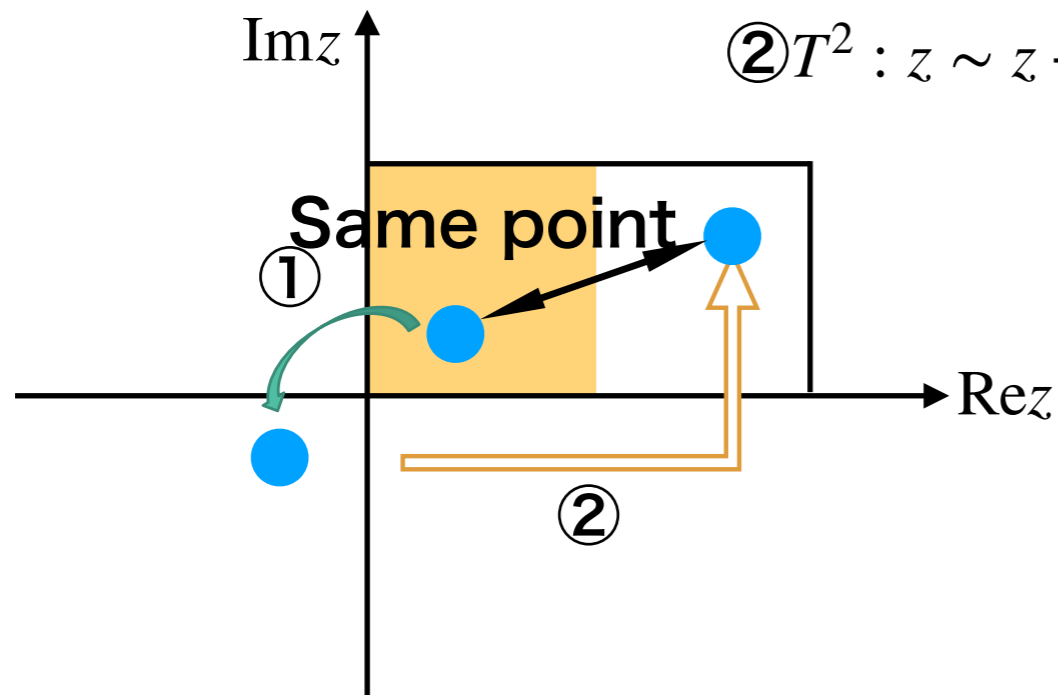
- $\mathcal{M}^4 \times T^2/Z_N$  in flux background

$$\frac{T^2 : z \sim z + 1 \sim z + \tau}{Z_N : z \sim \rho z \quad (\rho = e^{i\frac{2\pi}{N}})}$$

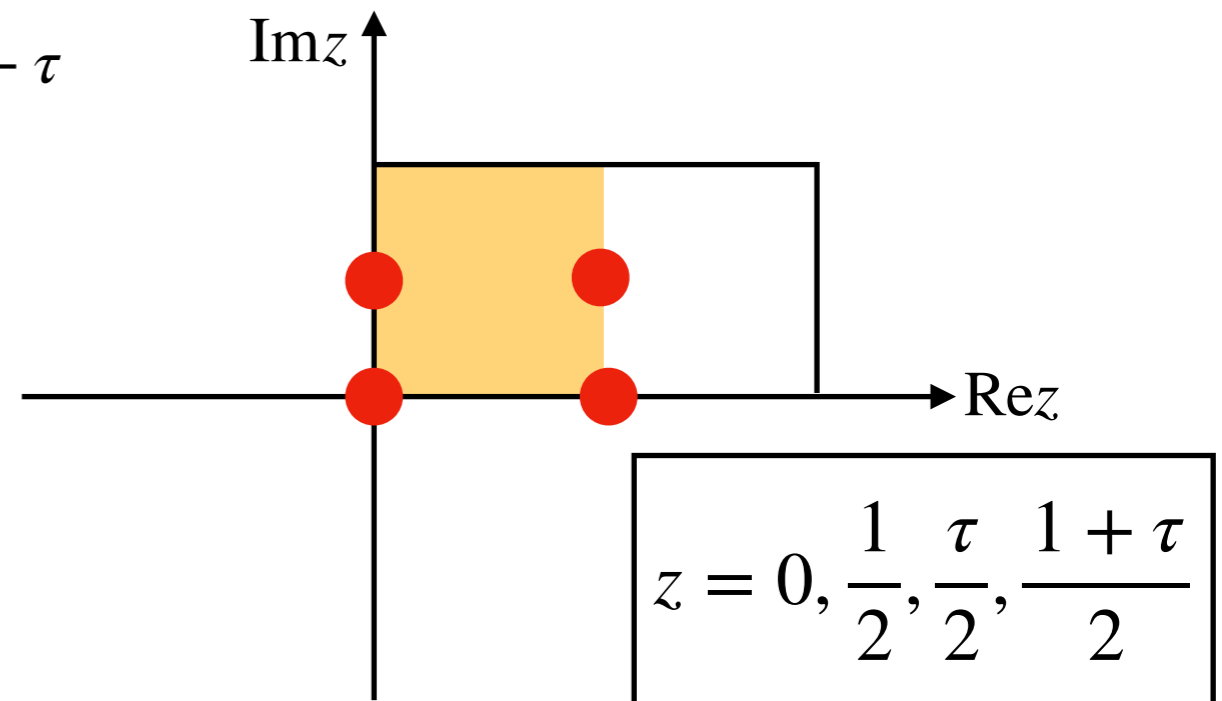


**Ex)**  $T^2/Z_2$  ( $\tau = i$ )    ①  $Z_2 : z \sim -z$

②  $T^2 : z \sim z + 1 \sim z + \tau$



Independent region is  $1/N$  of  $T^2$



Fixed point

# 3. $T^2/Z_N$ orbifold model

$Z_N$  eigen function  $\psi_{T^2/Z_N^{\pm,n,j}}^m(z)$

## Boundary condition

$$\psi_{T^2/Z_{N^+,n,j}}^m(z+1) = U_1(z) \psi_{T^2/Z_{N^+,n,j}}^m(z)$$

$$\psi_{T^2/Z_{N^+,n,j}}^m(\rho z) = \rho^m \psi_{T^2/Z_{N^+,n,j}}^m(z)$$

$$\psi_{T^2/Z_{N^+,n,j}}^m(z+\tau) = U_2(z) \psi_{T^2/Z_{N^+,n,j}}^m(z)$$

$$\rho = e^{i\frac{2\pi}{N}}$$

$$U_1(z) = e^{iq\Lambda_1(z)}, U_2(z) = e^{iq\Lambda_2(z)}$$

## Winding number

Define winding number of  $Z_N$  eigen function  $\psi_{T^2/Z_{N^+,n,j}}^m(z)$  at fixed point  $z_j^f$

$$\psi_{T^2/Z_{N^+,n,j}}^m(\rho z + z_j^f) = \rho^{\chi_+} \psi_{T^2/Z_{N^+,n,j}}^m(z + z_j^f) \Rightarrow \text{winding \# : } \chi_+$$

In case of  $z_j^f = 0$ ,  $\chi_+ = m$



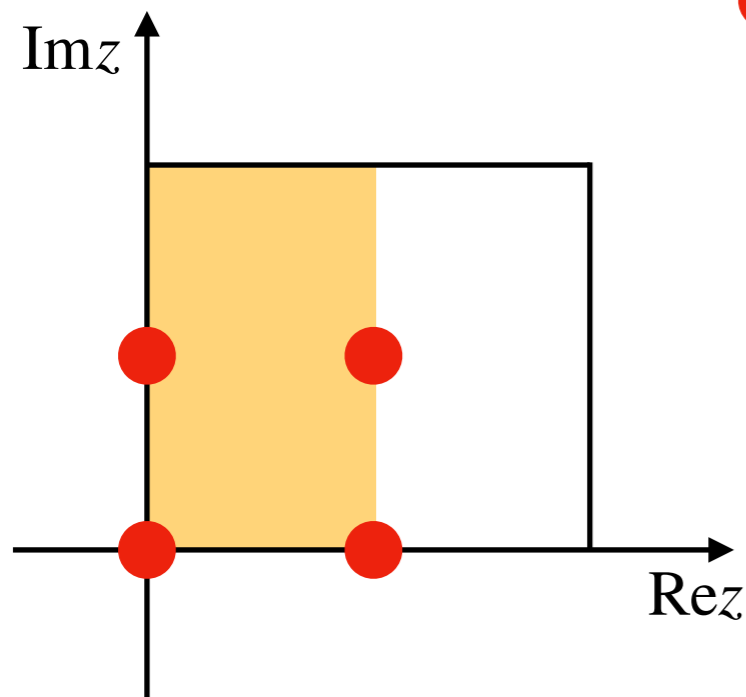
# 3. $T^2/Z_N$ orbifold model

## Atiyah-Singer index theorem on $T^2/Z_N$ orbifold

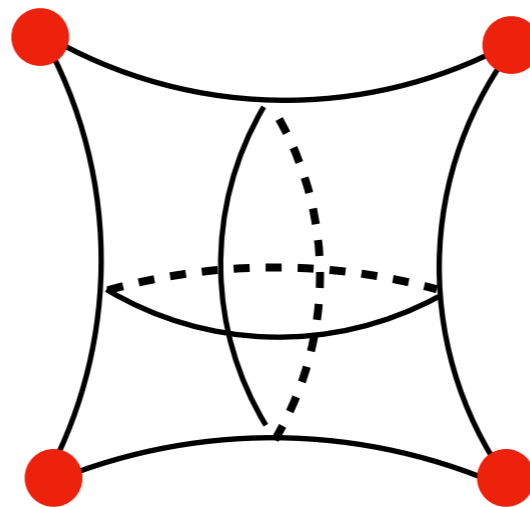
It is difficult to apply AS index theorem because fixed points are singular points.

➔ We consider blow-up manifold without singular points.

Ex)  $T^2/Z_2$



● Singular point



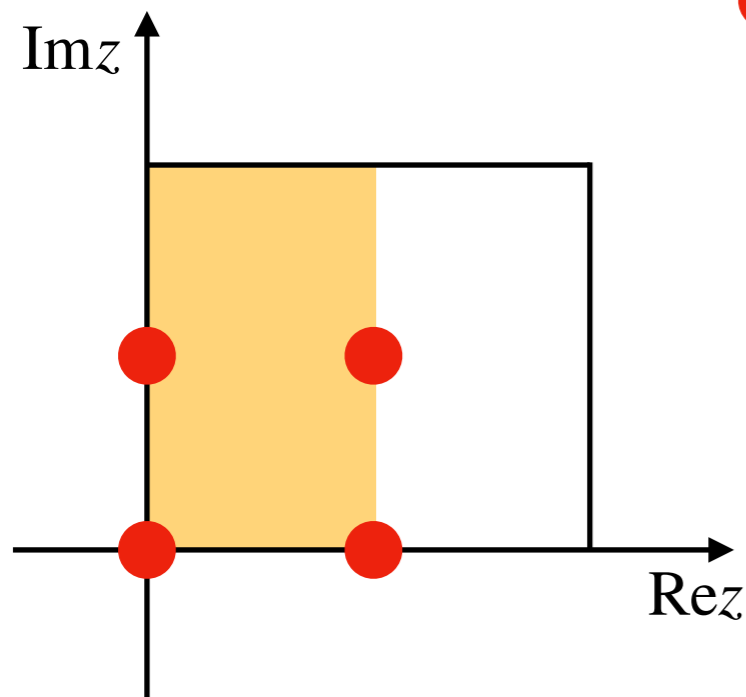
# 3. $T^2/Z_N$ orbifold model

## Atiyah-Singer index theorem on $T^2/Z_N$ orbifold

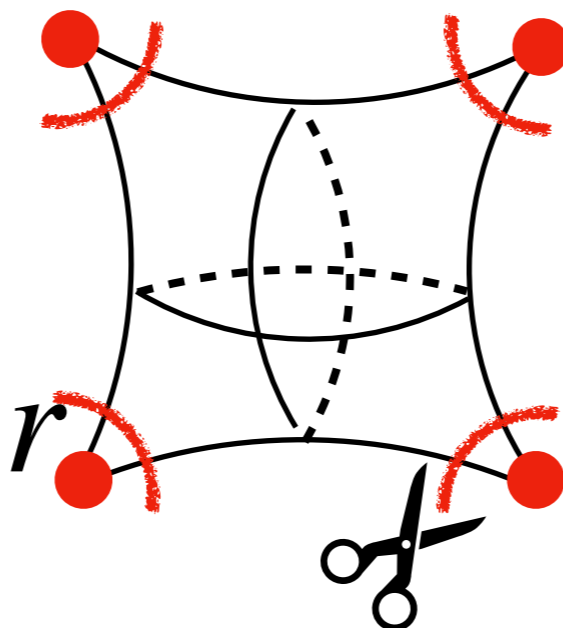
It is difficult to apply AS index theorem because fixed points are singular points.

➔ We consider blow-up manifold without singular points.

Ex)  $T^2/Z_2$



● Singular point



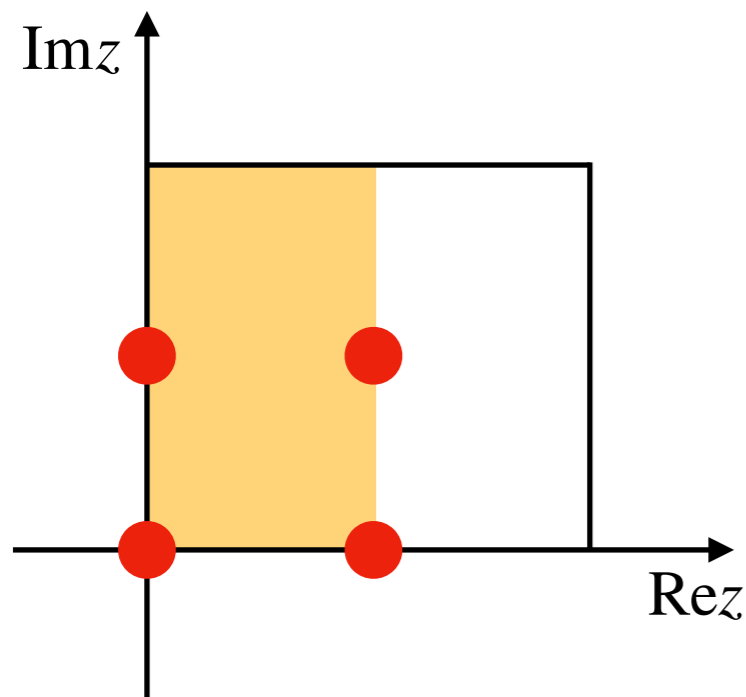
# 3. $T^2/Z_N$ orbifold model

## Atiyah-Singer index theorem on $T^2/Z_N$ orbifold

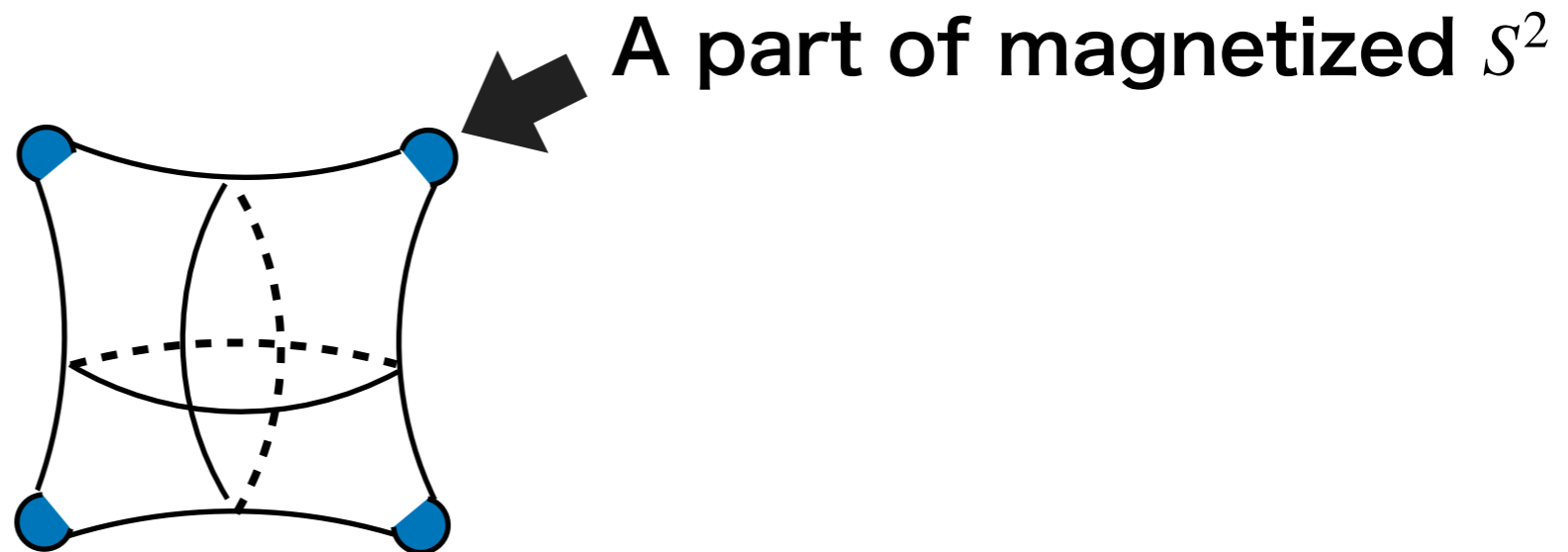
It is difficult to apply AS index theorem because fixed points are singular points.

**➔ We consider blow-up manifold without singular points.**

Ex)  $T^2/Z_2$



**Blow-up manifold**



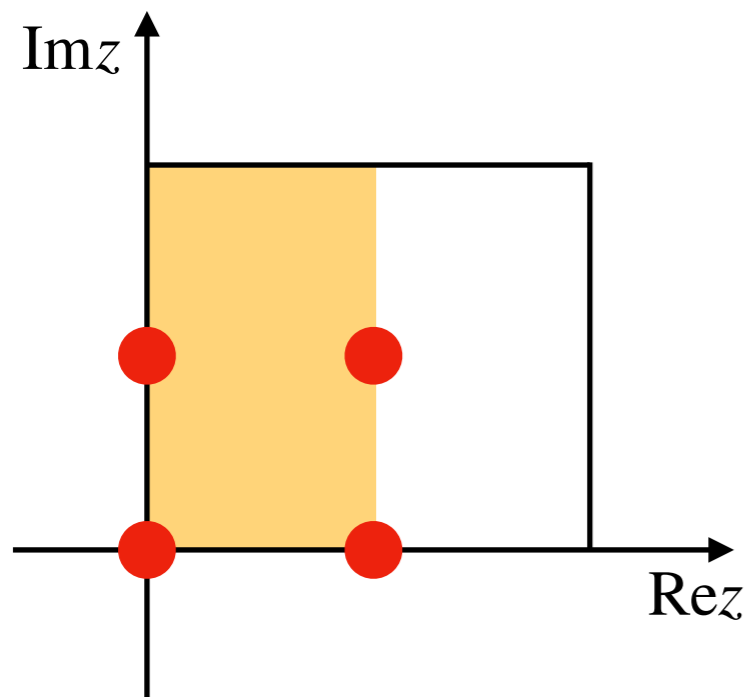
# 3. $T^2/Z_N$ orbifold model

## Atiyah-Singer index theorem on $T^2/Z_N$ orbifold

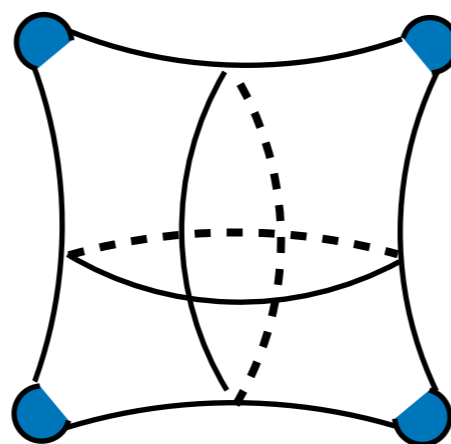
It is difficult to apply AS index theorem because fixed points are singular points.

**➔ We consider blow-up manifold without singular points.**

Ex)  $T^2/Z_2$



**Blow-up manifold**



**A part of magnetized  $S^2$**

**We can apply AS index theorem to blow-up manifold without singularity.**

# 4. Blow-up of $T^2/Z_N$ orbifold

$\psi_{T^2/Z_N,+}^0(z)$  with winding number cannot connect to  $\psi_{S^2,+}^0(z')$  ,  
because the boundary conditions are different at  $z \sim z_j^f$

# 4. Blow-up of $T^2/Z_N$ orbifold

**Purpose : To remove winding number**

We consider in case of fixed point  $z_j^f = 0$  ( $m \rightarrow \chi_+$  in case of  $z_j^f \neq 0$ )

$$\psi_{T^2/Z_N,+}^n(\rho z) = \rho^m \psi_{T^2/Z_N,+}^n(z) \quad \longrightarrow \quad \widetilde{\psi}_{T^2/Z_N,+}^n(\rho z) = \widetilde{\psi}_{T^2/Z_N,+}^n(z)$$

**“Singular” gauge transformation**

$$\widetilde{\psi}_{T^2/Z_N,\pm}(z) = U_{\xi^F} U_{\xi^R} \psi_{T^2/Z_N,\pm}(z) \quad U_{\xi^F} \propto \left( \frac{z}{\bar{z}} \right)^{\frac{\xi^F}{2}}, \quad U_{\xi^R} \propto \left( \frac{z}{\bar{z}} \right)^{\frac{\xi^R}{4}}$$

$\xi^F$  : localized flux at fixed point ,  $\xi^R$  : localized curvature at fixed point

$$\widetilde{\psi}_{T^2/Z_N,+}(\rho z) = \rho^{\xi^F - \frac{\xi^R}{2} + m} \widetilde{\psi}_{T^2/Z_N,+}(z)$$

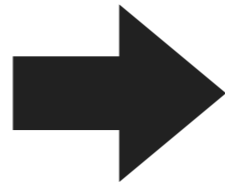
$$\rho^{\xi^F - \frac{\xi^R}{2} + m} = 1$$

However this gauge transformation changes flux, this is called “singular gauge transformation”.

# 4. Blow-up of $T^2/Z_N$ orbifold

Localized flux  $\xi^F$

$$\rho^{\xi^F - \frac{\xi^R}{2} + m} = 1$$



$$\xi^F = \frac{\xi^R}{2} - m \quad (z_j^f = 0)$$

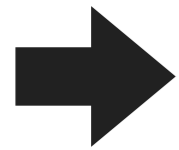
$$\xi^F = \frac{\xi^R}{2} - \chi_+ \quad (z_j^f \neq 0)$$

The information of winding numbers are replaced by localized flux  $\xi^F$  and localized curvature  $\xi^R$ .

Since  $\widetilde{\psi}_{T^2/Z_N,+}(z)$  have no winding numbers, these can be connected to  $\psi_{S^2,+}^0(z')$ .

# 4. Blow-up of $T^2/Z_N$ orbifold

Connection condition  $\tilde{\psi}_{T^2/Z_N, \pm}^0(z)$  and  $\psi_{S^2, +}^0(z')$

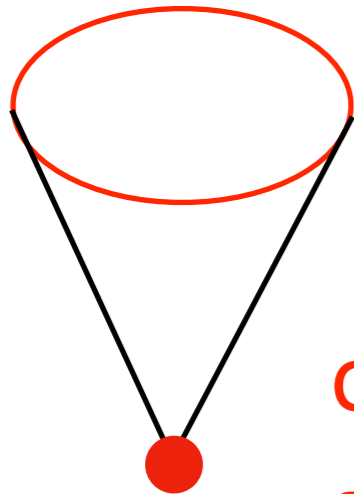


From  
non-  
holomorphic  
part

$$\frac{\pi r^2}{N \text{Im}\tau} M + \frac{\xi^F}{N} = \frac{N-1}{2N} M' \quad r \rightarrow 0 \quad \Rightarrow \quad \frac{\xi^F}{N} = \frac{N-1}{2N} M'$$

$\xi^F$  : localized flux.  $M'$  : Total flux of  $S^2$ ,  $\frac{N-1}{2N}$  : embedded area of  $S^2$

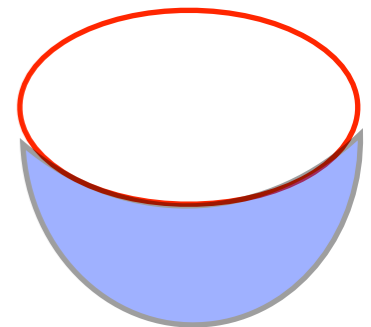
Physical meaning of connection condition



Cut out flux of  $T^2/Z_N$   
at fixed point

$$\frac{\xi^F}{N} = \frac{N-1}{2N} M'$$

Embedded flux of  $S^2$





# 5. Conclusion

Index theorem on blow-up manifold of  $T^2/Z_N$  orbifold

$$n_+ - n_- = \int_{\text{blow-up}} \frac{F}{2\pi} = \frac{M}{N} + \sum_{z_j^f} \frac{N-1}{2N} M' = \frac{M}{N} + \sum_{z_j^f} \frac{\xi^F}{N}$$

$n_{\pm}$  : chiral zero modes number ,  $M'$  : Total flux of  $S^2$  ,

$\frac{N-1}{2N}$  : embedded area of  $S^2$  ,  $\xi^F$  : localized flux at fixed point

# 5. Conclusion

Index theorem on blow-up manifold of  $T^2/Z_N$  orbifold

$$n_+ - n_- = \int_{\text{blow-up}} \frac{F}{2\pi} = \frac{M}{N} + \sum_{z_j^f} \frac{N-1}{2N} M' = \frac{M}{N} + \sum_{z_j^f} \frac{\xi^F}{N}$$

$n_{\pm}$  : chiral zero modes number ,  $M'$  : Total flux of  $S^2$  ,

$\frac{N-1}{2N}$  : embedded area of  $S^2$  ,  $\xi^F$  : localized flux at fixed point

**Only contribution of flux!**

# 5. Conclusion

Reinterpretation of index formula on  $T^2/Z_N$  orbifold

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

$M$  : flux quanta,  $V_+$  : sum of winding number  $\chi_+$  at fixed points

Winding number has contributions of localized flux and curvature.

$$\xi^F = \frac{\xi^R}{2} - m \quad (z_j^f = 0) \Rightarrow \xi^F = \frac{\xi^R}{2} - \chi_+ \quad (z_j^f \neq 0)$$

$$\sum_{z_j^f} \frac{\xi^F}{N} = - \sum_{z_j^f} \frac{\chi_+}{N} + \sum_{z_j^f} \frac{\xi^R}{2N} = - \frac{V_+}{N} + 1$$

# 5. Conclusion

Reinterpretation of index formula on  $T^2/Z_N$  orbifold

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

$M$  : flux quanta,  $V_+$  : sum of winding number  $\chi_+$  at fixed points

Winding number has contributions of localized flux and curvature.

$$\xi^F = \frac{\xi^R}{N} - \sum_{z_j^f} \frac{\chi_+}{N} + \sum_{z_j^f} \frac{\xi^R}{2N} = -\frac{V_+}{N} + 1$$

+1 removes contribution of localized curvatures in  $V_+$ .

$$\sum_{z_j^f} \frac{\xi^F}{N} = - \sum_{z_j^f} \frac{\chi_+}{N} + \sum_{z_j^f} \frac{\xi^R}{2N} = -\frac{V_+}{N} + 1$$

Thank you!

Back up

# 3. $T^2$ model

- $\mathcal{M}^4 \times T^2$  in flux background

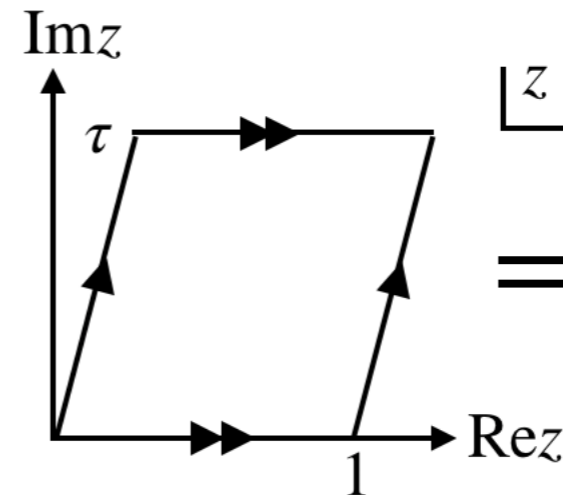
$$T^2 : z \sim z + 1 \sim z + \tau$$

## 6d Dirac action

$$S = \int d^4x \int d^2z \bar{\Psi}(x, z) i\Gamma^M D_M \Psi(x, z)$$

$$D_M = \partial_M - iqA_M$$

6d chiral fermion



## Mode expansion

$$\Psi(x, z) = \sum_n \psi_{R,n}^{(4)}(x) \otimes \psi_{T^2+,n}^{(2)}(z) + \psi_{L,n}^{(4)}(x) \otimes \psi_{T^2-,n}^{(2)}(z)$$

2d chiral fermion

4d chiral fermion

## 2d chiral fermion

$$\psi_{T^2+,n}^{(2)}(z) = \begin{pmatrix} \psi_{T^2+,n} \\ 0 \end{pmatrix}, \quad \psi_{T^2-,n}^{(2)}(z) = \begin{pmatrix} 0 \\ \psi_{T^2-,n} \end{pmatrix}$$

# 3. $T^2$ model

- $\mathcal{M}^4 \times T^2$  in flux background

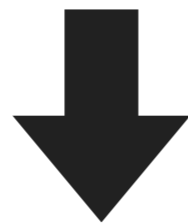
$T^2$  shift corresponds to gauge transformation.

## Boundary condition

$$\psi_{T^2 \pm, n, j}(z+1) = U_1(z) \psi_{T^2 \pm, n, j}(z)$$

$$\psi_{T^2 \pm, n, j}(z+\tau) = U_2(z) \psi_{T^2 \pm, n, j}(z)$$

$$U_1(z) = e^{iq\Lambda_1(z)}, U_2(z) = e^{iq\Lambda_2(z)}$$



**Flux quantization**  $\frac{qf}{2\pi} \equiv M \in \mathbb{Z}$

$M$ : flux quantization number  $f$ : flux



# 3. $T^2$ model

- $\mathcal{M}^4 \times T^2$  in flux background

## Eigenvalue equation

$$\left. \begin{aligned} 2D_{\bar{z}}\psi_{T^2+,n,j} &= 2(\partial_{\bar{z}} - iqA_{\bar{z}})\psi_{T^2+,n,j} = m_n\psi_{T^2-,n,j} \\ -2D_z\psi_{T^2-,n,j} &= -2(\partial_z - iqA_z)\psi_{T^2-,n,j} = m_n\psi_{T^2+,n,j} \end{aligned} \right\} \begin{aligned} -4D_zD_{\bar{z}}\psi_{T^2+,n,j} &= m_n^2\psi_{T^2+,n,j} \\ -4D_{\bar{z}}D_z\psi_{T^2-,n,j} &= m_n^2\psi_{T^2-,n,j} \end{aligned}$$

$\Rightarrow$  Zero mode has the eigenvalue  $m_n = 0$ .

## AS index theorem on $T^2$

$$n_+ - n_- = \int_{T^2} \frac{F}{2\pi} = M$$

↑ Chiral zero modes #
 ↑ flux quantization#

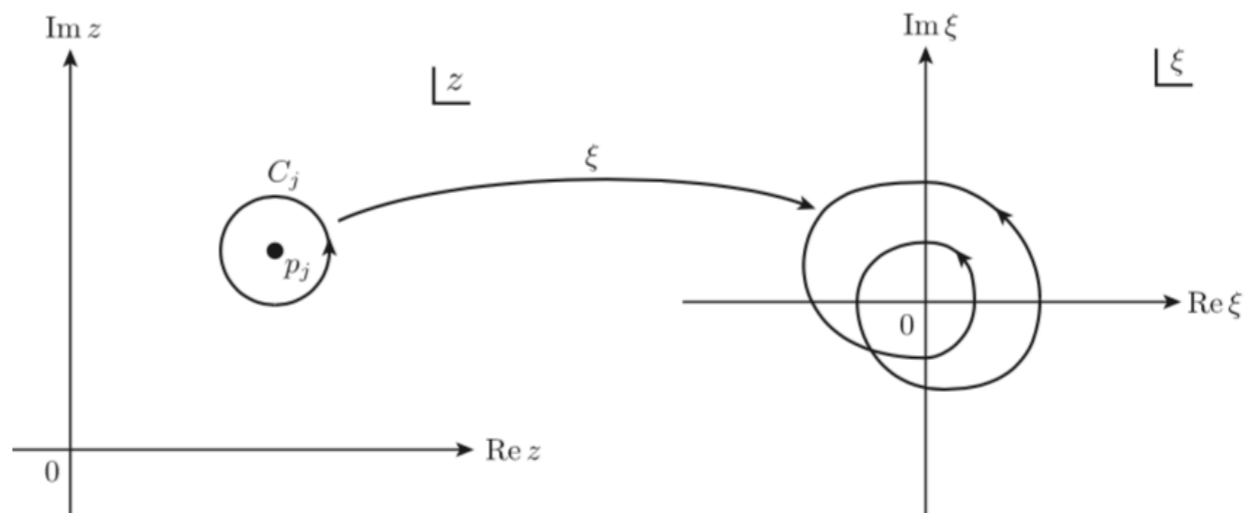
# 4. $T^2/Z_N$ orbifold model

## Winding number

Define winding number of  $Z_N$  eigen function  $\psi_{T^2/Z_{N^+},n,j}^m(z)$  at fixed point  $z_j^f$

$$\psi_{T^2/Z_{N^+},n,j}^m(\rho z + z_j^f) = \rho^{\chi_+} \psi_{T^2/Z_{N^+},n,j}^m(z + z_j^f) \Rightarrow \text{winding \#} : \chi_+$$

ex) winding # = 2

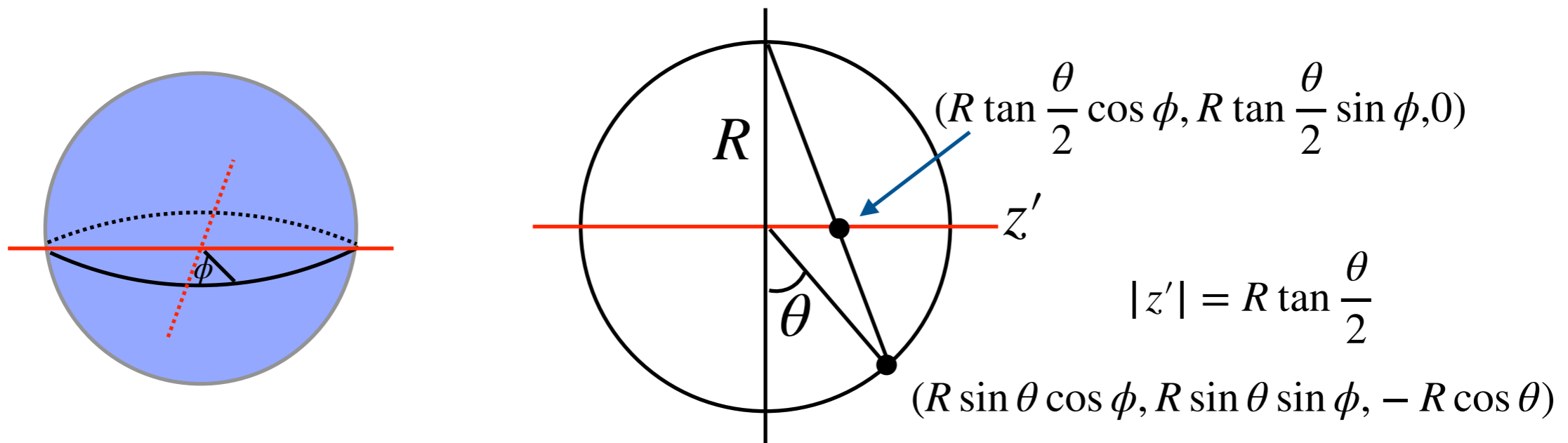


Winding number depends on eigen function & fixed point

# 5. Magnetized $S^2$

## Magnetized $S^2$

How to take the coordinate  $z'$  of  $S^2$



## AS index theorem on $S^2$

$$n_+ - n_- = \int_{S^2} \frac{F}{2\pi} = M'$$

Chiral zero modes # flux quantization #

# 5. Magnetized $S^2$

## Magnetized $S^2$

### Dirac equation

$$\frac{R^2 + |z'|^2}{R} i(\partial_{\bar{z}'} + i\frac{1}{2}\omega_{\bar{z}'} - iA_{\bar{z}'})\psi_{S^2,+}^n(z') = m_n \psi_{S^2,-}^n(z')$$

$$\frac{R^2 + |z'|^2}{R} i(\partial_{z'} - i\frac{1}{2}\omega_{z'} - iA_{z'})\psi_{S^2,-}^n(z') = m_n \psi_{S^2,+}^n(z')$$

$$\omega_{\bar{z}'} = \frac{i}{2} \frac{2}{R^2 + |z'|^2} z', \quad \omega_{z'} = -\frac{i}{2} \frac{2}{R^2 + |z'|^2} \bar{z}', \quad A_{\bar{z}'} = \frac{i}{2} \frac{M'}{R^2 + |z'|^2} z', \quad A_{z'} = -\frac{i}{2} \frac{M'}{R^2 + |z'|^2} \bar{z}'$$

$\Rightarrow$  Zero mode has the eigenvalue  $m_n = 0$ .

Zero mode

$$\psi_{S^2,+}^0(z') = \frac{f_+(z')}{(R^2 + |z'|^2)^{\frac{M'-1}{2}}}$$

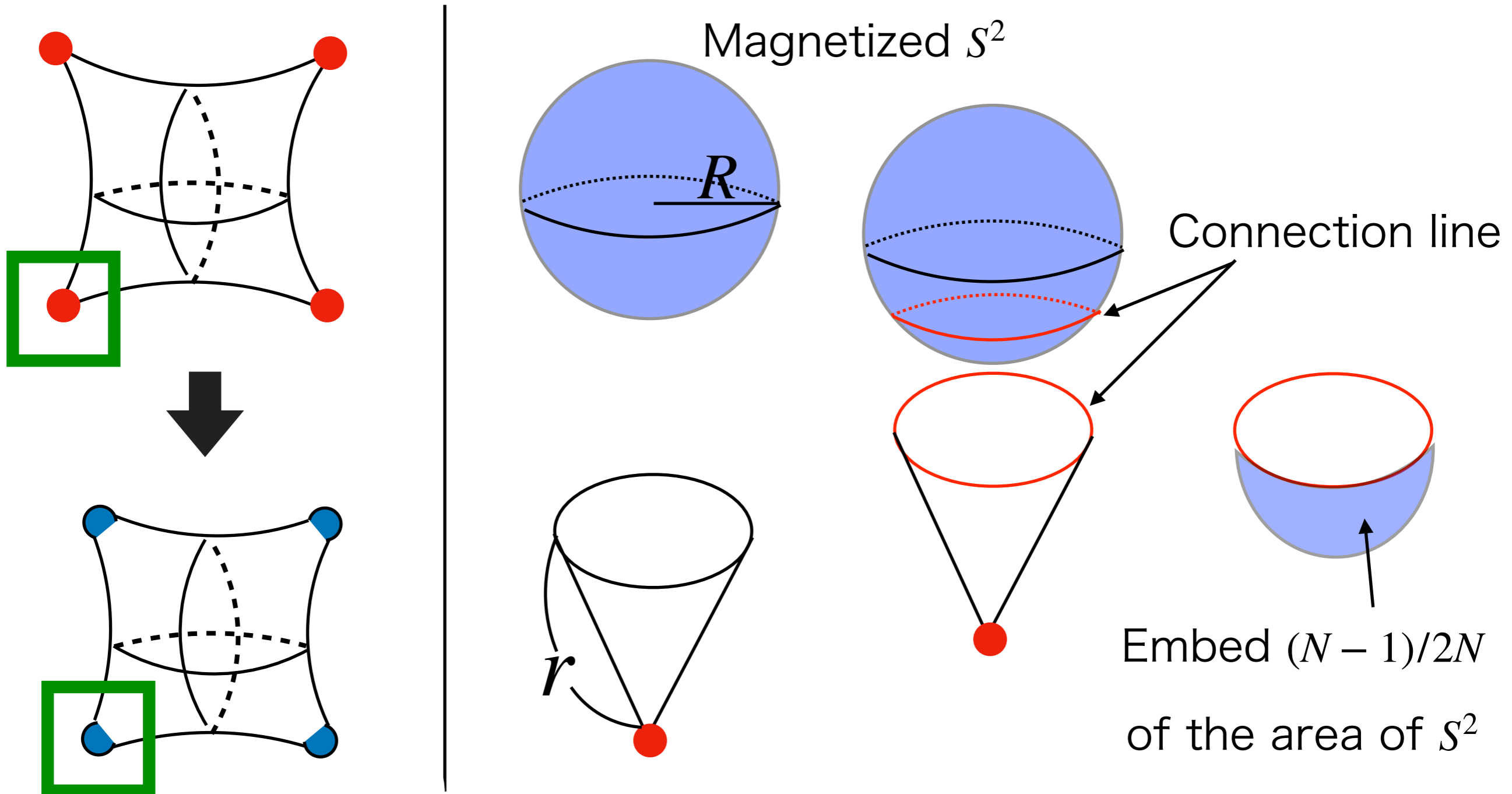
Lowest mode with - chirality

$$\psi_{S^2,-}^1(z') = \frac{f_-(z')}{(R^2 + |z'|^2)^{\frac{M'+1}{2}}}$$

$f_+(z'), f_-(z')$ : holomorphic function

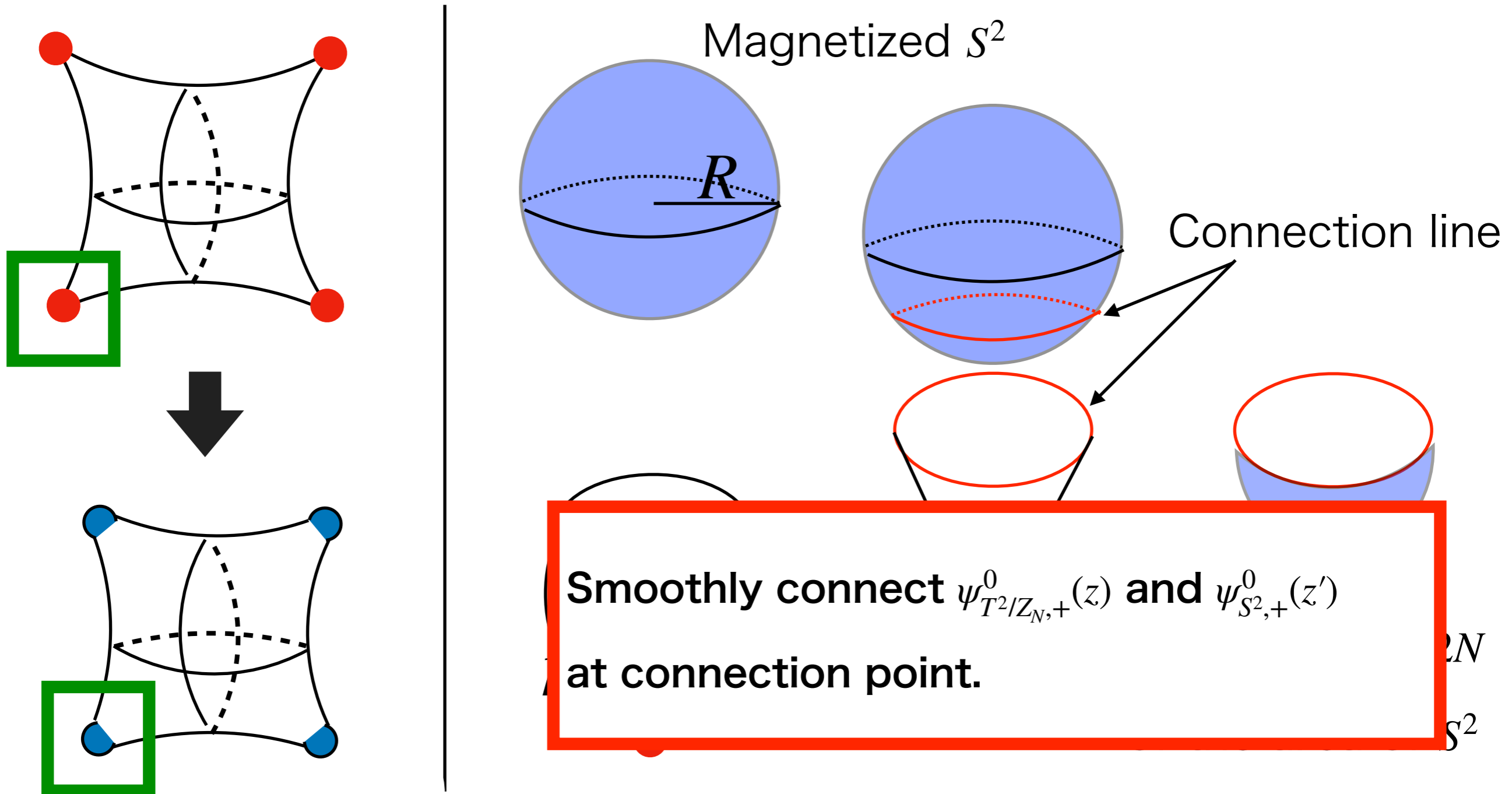
# 4. Blow-up of $T^2/Z_N$ orbifold

To apply AS index theorem, we introduce blow-up manifold



# 4. Blow-up of $T^2/\mathbb{Z}_N$ orbifold

To apply AS index theorem, we introduce blow-up manifold



# 4. Blow-up of $T^2/Z_N$ orbifold

## “Singular” gauge transformation

$$\widetilde{\Psi}_{T^2/Z_N, \pm}(z) = U_{\xi^F} U_{\xi^R} \Psi_{T^2/Z_N, \pm}(z) \quad U_{\xi^F} \propto \begin{pmatrix} z \\ \bar{z} \end{pmatrix}^{\frac{\xi^F}{2}}, \quad U_{\xi^R} \propto \begin{pmatrix} z \\ \bar{z} \end{pmatrix}^{\frac{\xi^R}{4}}$$

$$\Psi_{T^2/Z_N, +}(z)$$

Field strength  $F$

$$\int_{T^2/Z_N} \frac{F}{2\pi} = \frac{M}{N}$$

Gauge field  $A$

+

Winding number  $\chi_+$

$$\widetilde{\Psi}_{T^2/Z_N, +}(z)$$

Field strength  $F + \delta F$

$$\int_{T^2/Z_N} \frac{F + \delta F}{2\pi} = \frac{M}{N} + \frac{\xi^F}{N} \quad \xi^F : \text{localized flux at fixed point}$$

Gauge field  $A + \delta A$

$$\delta A = iU_{\xi^F} dU_{\xi^F}^{-1}$$

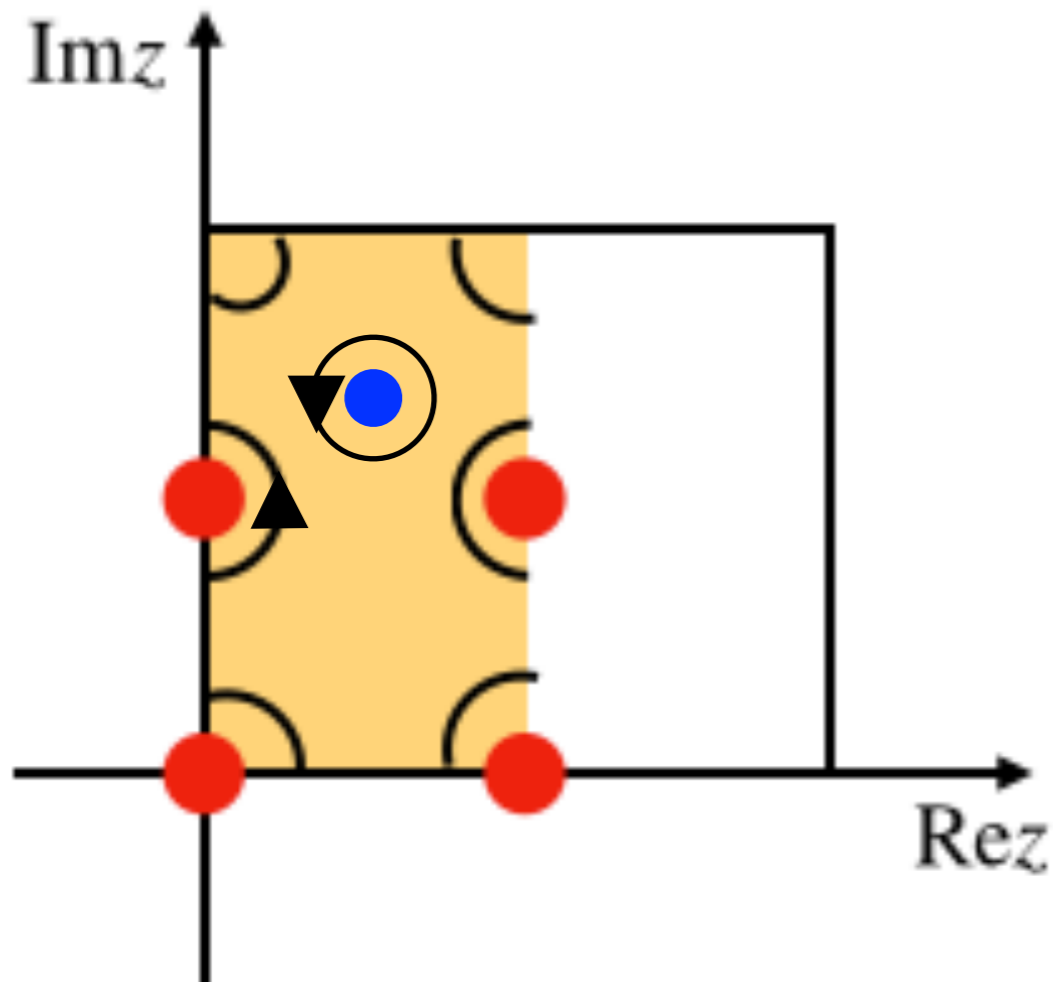
Localized curvature  $\xi^R$

$$\int_{T^2/Z_N \text{ fixed point}} \frac{\delta R}{2\pi} = \frac{\xi^R}{N}$$

+ No winding number

# 6. Blow-up of $T^2/Z_N$ orbifold

Localized curvature at fixed point  $\xi^R$



At fixed points

→ Winding is closed at  $180^\circ$

At the other points

→ Winding is closed at  $360^\circ$

At fixed points, there is

deflection angle  $2\pi - \frac{2\pi}{N} = 2\pi \frac{N-1}{N}$ .



# 6. Blow-up of $T^2/Z_N$ orbifold

Localized curvature at fixed point  $\xi^R$

Gauss-Bonnet theorem

$$\int_{\mathcal{M}} \frac{R}{2\pi} = \chi(\mathcal{M})$$

$R$  : Curvature

$\chi$  : Euler characteristic

+

$$(\text{deflection angle}) = 2\pi\chi$$

→ Localized curvature at fixed point  $\frac{\xi^R}{N} = \frac{N-1}{N}$

# 5. Conclusion

Reinterpretation of index formula on  $T^2/Z_N$  orbifold

$$(\text{generation \#}) = \frac{M}{N} - \frac{V_+}{N} + 1$$

$M$  : flux quanta,  $V_+$  : sum of winding number  $\chi_+$  at fixed points

Winding number has contributions of localized flux and curvature.

Index theorem implies the existence of  $l$  new zero modes = localized modes

$$\sum_{z_j^f} \frac{\xi^F}{N} = - \sum_{z_j^f} \frac{\chi_+}{N} + \sum_{z_j^f} \frac{\xi^R}{2N} + l = - \frac{V_+}{N} + 1 + l$$