

String Excitation by Initial Singularity of Inflation

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Introduction (1/2)

- Inflation models are the most vigorously studied model of the early universe.

ex. Starobinsky model, hill-top model, etc...

- A singularity occurs at very early stage of inflation due to a component of Ricci tensor diverges.

→ Does a singularity due to the divergence of a component of Ricci tensor cause problems?

Introduction (2/2)

- From what point of view do you find out “whether an initial singularity causes problems”?
- Effects of quantum gravity is important on the early stage of inflation.
 - Consider (bosonic) string theory to investigate the initial singularity of an inflationary universe.
- Focus on expectation values of mass of string.
- A singularity should be resolved by effects of quantum gravity if it causes a divergence of mass.

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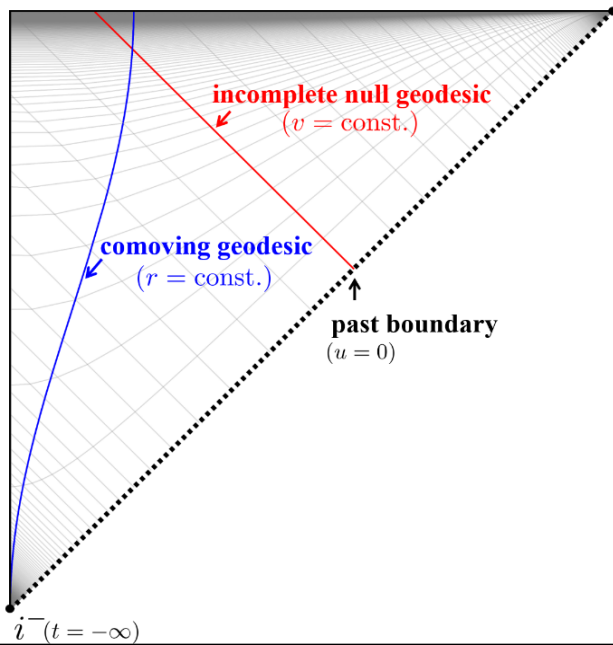
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Initial singularity in FLRW space-time

- Flat FLRW space-time in light cone coordinates

$$g_{\mu\nu} dx^\mu dx^\nu = -2dudv + a^2(u)dv^2 + a^2(u)(v - \eta(u))^2 d\Omega^2$$

- A singularity cause at the past $u = 0$ by divergence of uu component of Ricci tensor in inflationary universe.



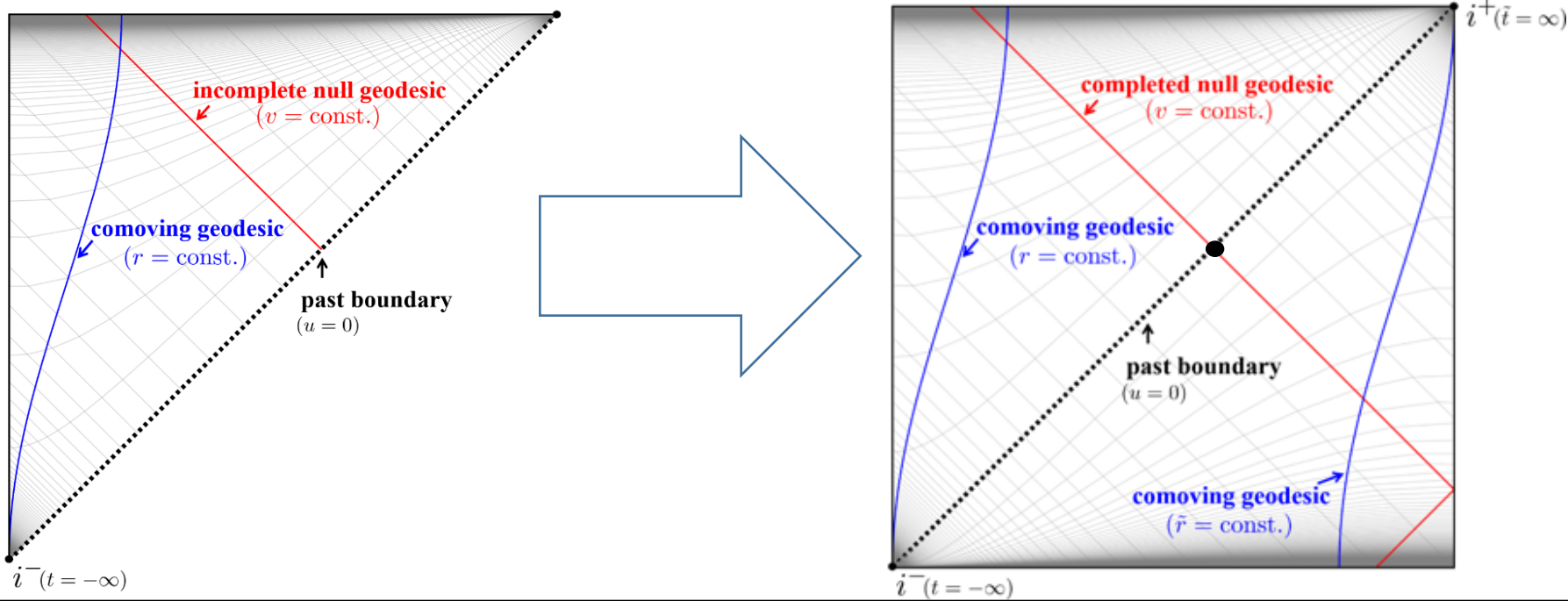
$$R_{uu} = -2A(u)$$

$$A(u) := \frac{\dot{H}}{a^2}, \quad a(u) \stackrel{u \rightarrow 0}{\sim} 0$$

- How do expectation values of mass behave when a string passes through the singularity?

Contracting universe

- Construct a contracting universe by using a continuity of the metric.
- The Bogoliubov transformation allows us to compare the masses before and after passing through the singularity.



EOM of string

- Consider a quantization of string and EOM in simplified FLRW space-time by taking the Penrose limit.

- $X_n^i(\tau)$ determine the motion of strings.

$$-\frac{d^2 X_n^i}{d\tau^2} + \mathcal{V}(\tau) X_n^i = E_n X_n^i, \quad \mathcal{V}(\tau) := \alpha'^2 p^2 A(\alpha' p \tau), \quad E_n := n^2$$

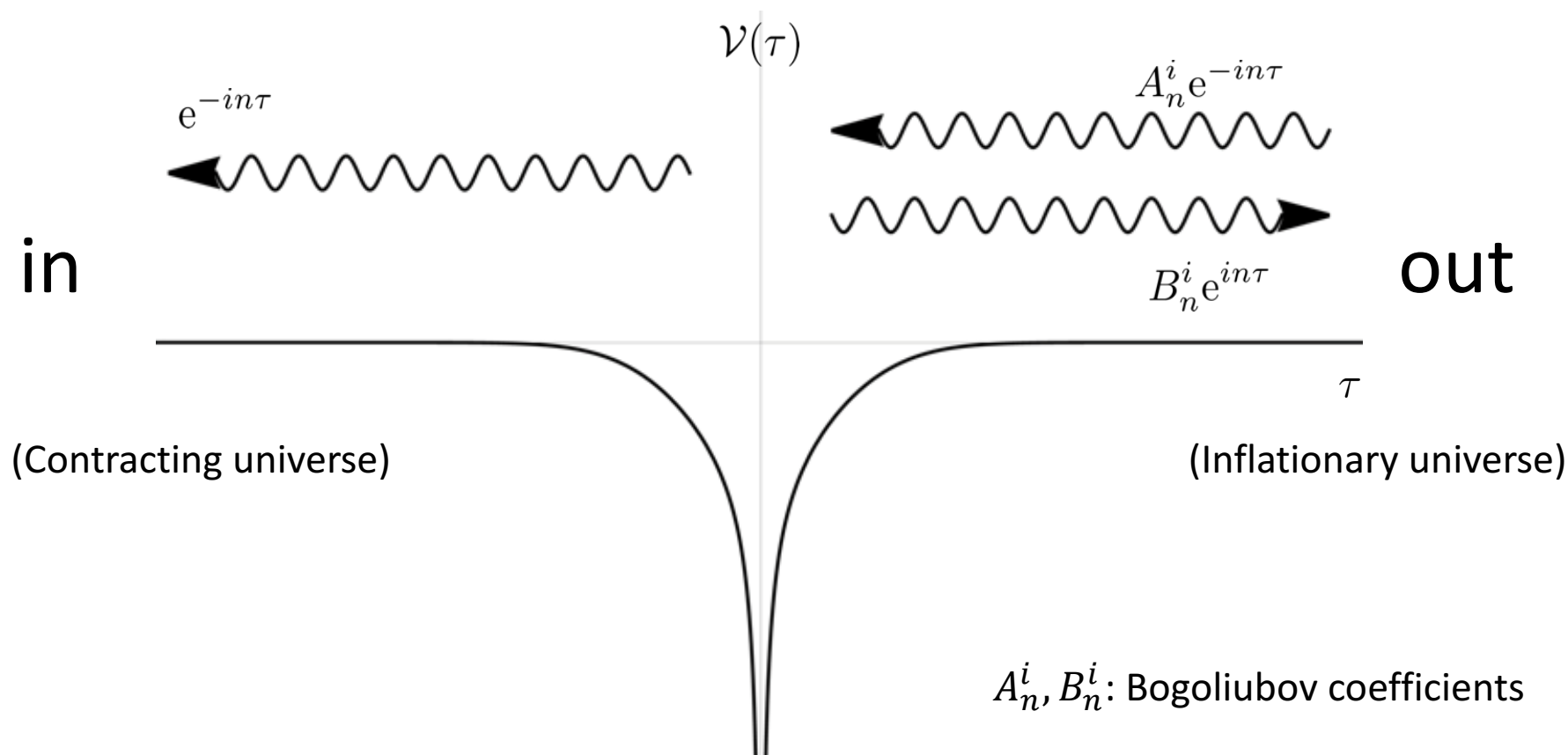
- This is Schrödinger-like equation has a potential determined by $A(u = \alpha' p \tau)$.

$$\text{※ } R_{uu} = -2A(u)$$

$$A(u) := \frac{\dot{H}}{a^2}, \quad a(u) \stackrel{u \rightarrow 0}{\sim} 0$$

String excitation by singularity (1/3)

- The Bogoliubov transformation from in-state (contracting universe) to out-state (inflationary universe)



String excitation by singularity (2/3)

- Calculate an expectation value of mass of out-state looked from in-state by using the Bogoliubov coefficients.

$$\langle M_{\text{out}}^2 \rangle = \frac{\langle 0_{\text{in}} | M_{\text{out}}^2 | 0_{\text{in}} \rangle}{\langle 0_{\text{in}} | 0_{\text{in}} \rangle} \sim \frac{2}{\alpha'} \sum_{n=1}^{\infty} \sum_{i=1}^2 n |B_n^i|^2$$

- The Bogoliubov coefficient B_n^i is the following as,

$$B_n^i \simeq \frac{p^2 \alpha'^2}{2in} \int_{-\infty}^{\infty} d\tau e^{-2in\tau} A(\alpha' p\tau)$$

so we can calculate the mass when we obtain uu components of Ricci tensor $A(u = \alpha' p\tau)$.

String excitation by singularity (3/3)

- In general, $A(u)$ takes the following form in an inflation models. (The values of β are different for each models)

$$A(u) = -\frac{\kappa}{|u|^\beta}, \quad \kappa > 0 : \text{const}, \quad 0 < \beta < 2$$

- The expectation value of mass converges for $0 < \beta < 1$ and diverges for $1 \leq \beta < 2$.

$$\langle M_{\text{out}}^2 \rangle = \begin{cases} \text{converge} & 0 < \beta < 1, \\ \text{diverge} & 1 \leq \beta < 2 \end{cases}$$

→ We can distinguish models whether break down or not by calculating the value of β .

Starobinsky model

- The potential of Starobinsky model is

$$V(\phi) = 3M_{pl}^2 \bar{H}^2 \left(1 - e^{-\frac{\phi}{\mu}} + \dots \right), \quad \mu : \text{const}$$

- We can calculate $A(u)$ easily as

$$A(u) \sim -\frac{\kappa}{u^2 (\log u)^2}$$

This corresponds to $1 \leq \beta < 2$, so the mass diverges.

→ The initial singularity in Starobinsky model should be removed by effects of quantum gravity.

Cosine type hill-top model

- The potential of Cosine type hill-top model is

$$V(\phi) = \frac{3M_{pl}^2 \bar{H}^2}{2} \left(1 + \cos \frac{\phi}{f} \right) = 3M_{pl}^2 \bar{H}^2 \left(1 - \frac{1}{2} \frac{M_{pl}^2}{2f^2} \frac{\phi^2}{M_{pl}^2} + \dots \right)$$

Then, we obtain

$$A(u) = \frac{-\kappa}{u^{2(1-\gamma)}}, \quad \gamma = \frac{-3 + \sqrt{9 + \frac{6M_{pl}^2}{f^2}}}{2}$$

- If we ignore the slow-roll condition,

$$\langle M_{\text{out}}^2 \rangle = \begin{cases} \text{converge} & f < \sqrt{\frac{6}{7}} M_{pl}, \\ \text{diverge} & f \geq \sqrt{\frac{6}{7}} M_{pl} \end{cases}$$

Summery

- The initial singularity causes at early stage of inflation.

→ Do inflation models break down due to the singularity as EFT of string theory?

	Is it OK?	Value of β
Starobinsky model	×	$1 \leq \beta < 2$
Cosine type hill-top model	△	$0 < \beta < 1$ or $1 \leq \beta < 2$

- We suggested the method to distinguish inflation models whether they break down or not due to the initial singularity.

- Stringy corrections should resolve the initial singularity even though the value of the Hubble parameter remains smaller than the string scale.

Outlook

- Calculate the time-dependence of the expectation value of mass (string excitation) without the contracting universe.
- Take into account a back reaction of the excitation.

Calculations

- $A(u)$ is calculated by potential $V(\phi)$.

$$V(\phi) = 3M_{pl}^2 \bar{H}^2 (1 + \delta V(\phi))$$

$$A(u) = \frac{1}{a^2} \frac{dH}{dt} \simeq -\frac{M_{pl}^2}{2u^2} \times \delta V'(\phi(u))^2$$

- EOMs of inflaton ϕ are

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0, \quad H^2 = \frac{1}{3M_{pl}^2} \left(\frac{1}{2} \dot{\phi}^2 + V(\phi) \right)$$