

# *Setting limits on the toponium cross section*

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# *Toponium??*

According to a particle theory book from 2015:

The t-quark has, due to its large mass, only a fleeting lifetime. Thus no pronounced  $t\bar{t}$  states (*toponium*) are expected.

# About Toponium

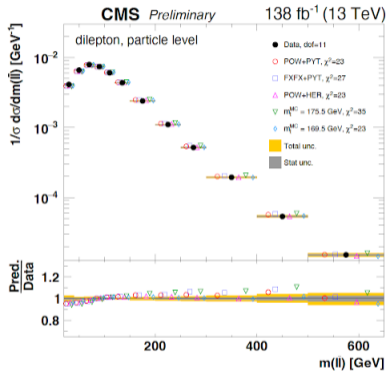
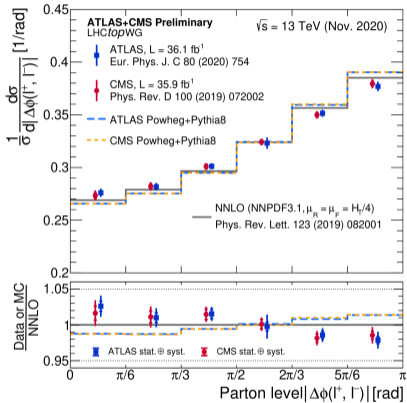
Top quarks decay into a W boson and b quark  $\approx 100\%$  of the time. Decays split into 3 branching ratios: dileptonic, semileptonic, full hadronic. The  $t\bar{t}$  pairs are produced dileptonically as:  $gg \rightarrow t\bar{t} \rightarrow (bW^+)(\bar{b}W^-) \rightarrow (bl^+\nu_l)(\bar{b}l^-\bar{\nu}_l)$

$$m_t \approx 172.44 \pm 0.13 \pm 0.47 \text{ GeV}^{[\text{arxiv:1403.4427}]}$$

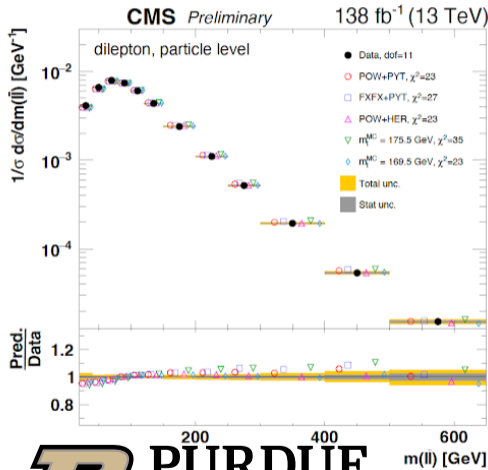
$$\underbrace{\frac{1}{m_t}}_{\text{production}} < \underbrace{\frac{1}{\Gamma_t}}_{\text{lifetime}} < \underbrace{\frac{1}{\Lambda_{\text{QCD}}}}_{\text{hadronization}} < \underbrace{\frac{m_t}{\Lambda^2}}_{\text{spin-flip}}$$

$10^{-27} \text{ s} \quad 10^{-25} \text{ s} \quad 10^{-24} \text{ s} \quad 10^{-21} \text{ s}$

# Toponium Evidence



# Evidence for Toponium?

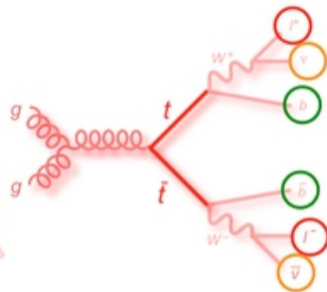
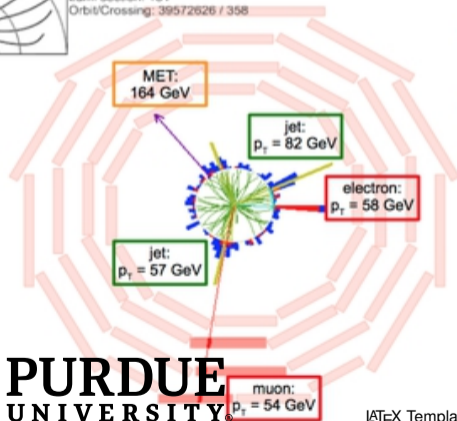


- ATLAS [arxiv:1910.08819] and CMS see an excess of low  $m_{\ell\bar{\ell}}$  events (dilepton events)
- Fuks et. al. showed that this can be explained via a  $t\bar{t}$  scalar resonance [arxiv:2102.11281]
- Using an event selection procedure biased for toponium production Fuks et. al. achieved a S/N ratio of 7% when considering only  $t\bar{t}$  events

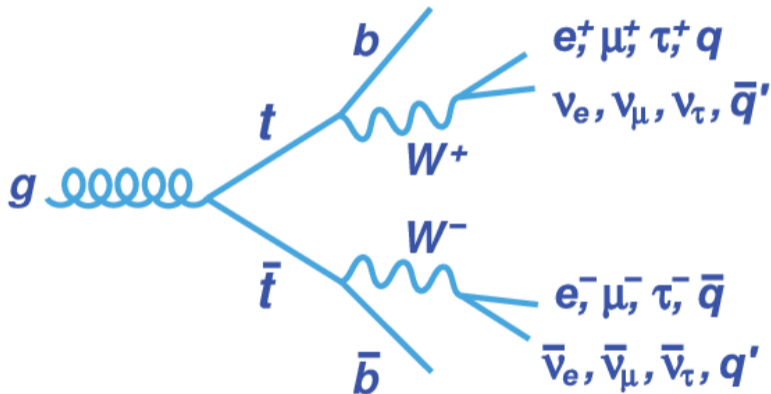
# Top Quark Physics



CMS Experiment at LHC, CERN  
Data recorded: Wed Jul 8 19:26:24 2015 CEST  
Run/Event: 251244 / 83494441  
Lumi section: 151  
Orbit/Crossing: 39572626 / 358



# Top Quark Physics



# Spin Correlation Analysis

$$\underbrace{\frac{1}{m_t}}_{\text{production}} < \underbrace{\frac{1}{\Gamma_t}}_{\text{lifetime}} < \underbrace{\frac{1}{\Lambda_{\text{QCD}}}}_{\text{hadronization}} < \underbrace{\frac{m_t}{\Lambda^2}}_{\text{spin-flip}}$$

$10^{-27} \text{ s}$        $10^{-25} \text{ s}$        $10^{-24} \text{ s}$        $10^{-21} \text{ s}$



## Basic of Spin-corr

For measurements in pp collisions, using the  $\{\hat{k}, \hat{r}, \hat{n}\}$  basis vectors, the angular distribution for the two leptons (each measured in its parent top CM frame) can be stated precisely. For each of the 15 coefficients that make up  $\mathbf{B}_{1/2}$  and  $C$ , a change of variables can be made to obtain a single-differential cross section dependent only on that coefficient. For example, after integrating out the azimuthal angles,  $\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_1 d \cos \theta_2} = \frac{1}{4} (1 + B_1(i) \cos \theta_1 + B_2(i) \cos \theta_2 - C(i, i) \cos \theta_1 \cos \theta_2)$ , where

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_1^i d \cos \theta_2^j} = \frac{1}{4} \left( 1 + B_1(i) \cos \theta_1^i + B_2(j) \cos \theta_2^j - C(i, j) \cos \theta_1^i \cos \theta_2^j \right),$$

## Cont.

from which we can derive single-differential cross sections wrt  $\cos \theta_1^i$ ,  $\cos \theta_2^i$ , and  $\cos \theta_1^i \cos \theta_2^j$ :

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_1^i} = \frac{1}{2} \left( 1 + B_1(i) \cos \theta_1^i \right)$$

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_2^i} = \frac{1}{2} \left( 1 + B_2(i) \cos \theta_2^i \right)$$

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_1^i \cos \theta_2^j} = \frac{1}{2} \left( 1 - C(i, j) \cos \theta_1^i \cos \theta_2^j \right) \log \left( \frac{1}{|\cos \theta_1^i \cos \theta_2^j|} \right).$$

15 variables at parton level:

- The 3  $\cos \theta_1^i \cos \theta_2^i$ , to measure  $C(i, i)$ , the "diagonal" spin correlation coefficient for each axis  $i$ .
- The 6 sums and differences  $\cos \theta_1^i \cos \theta_2^j \pm \cos \theta_1^j \cos \theta_2^i$ , to measure the sums and differences of the "cross" spin correlation coefficients  $C(i, j) \pm C(j, i)$  for each pair of axes  $i, j$ .
- The 6  $\cos \theta_1^i$  and  $\cos \theta_2^i$ , to measure  $B_1(i)$  and  $B_2(i)$ , the top and anti-top polarization coefficients w.r.t each reference axis  $i$ .
- In addition, we measure 4 distributions based on modified axes  $\hat{r}^*$  and  $\hat{k}^*$ , equal to  $\pm \hat{r}$  or  $\pm \hat{k}$  depending on  $\text{sign}(|y_t| - |y_{\bar{t}}|)$ .

- The distribution  $\cos \varphi_{\text{lab}} = \hat{\ell}_1^{\text{lab}} \cdot \hat{\ell}_2^{\text{lab}}$ , as above but using lepton directions measured in the laboratory frame (which have excellent experimental resolution).
- The distribution  $|\Delta\phi_{\ell\ell}|$ , the difference in azimuthal angle  $\phi$  between the two leptons in the lab frame (whose shape is partly dependent on the above coefficients, particularly the  $C(i, i)$ , and has excellent experimental resolution).

Observable	Measured coefficient	Coefficient function
$\cos \theta_1^k$	$B_1^k$	$b_k^+$
$\cos \theta_2^k$	$B_2^k$	$b_k^-$
$\cos \theta_1^r$	$B_1^r$	$b_r^+$
$\cos \theta_2^r$	$B_2^r$	$b_r^-$
$\cos \theta_1^n$	$B_1^n$	$b_n^+$
$\cos \theta_2^n$	$B_2^n$	$b_n^-$

**Table:** (Partial table of) Observables and corresponding measured coefficients and production spin density matrix coefficient functions. 15 independent spin parameters and 25 direct observables/distributions and indirect observables as well.\*\*\*

# Statement of Classifier Problem in $n = 7$

From this we want to map this vector:

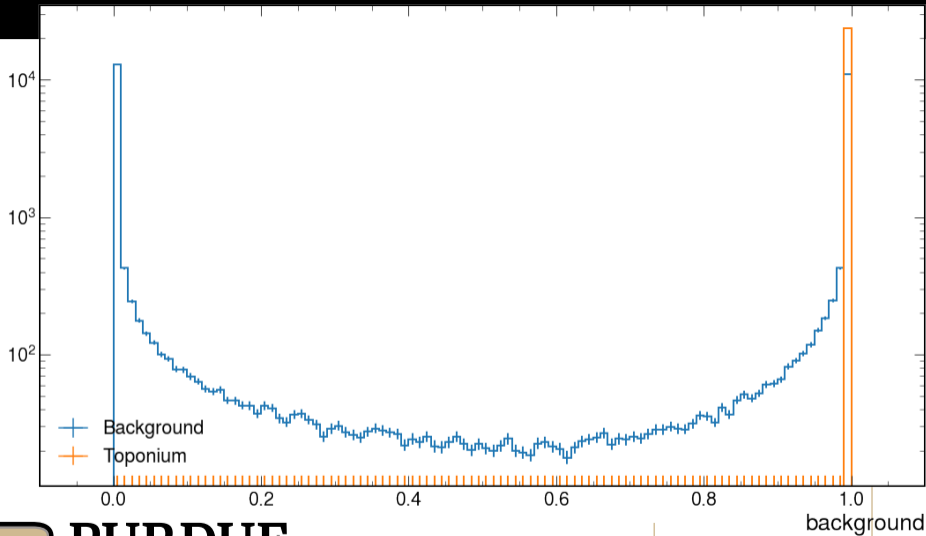
$$\mathbf{v} = \langle \rho_t, \eta_t, \phi_t, \rho_{\bar{t}}, \eta_{\bar{t}}, \phi_{\bar{t}}, m_{t\bar{t}} \rangle \quad (1)$$

under

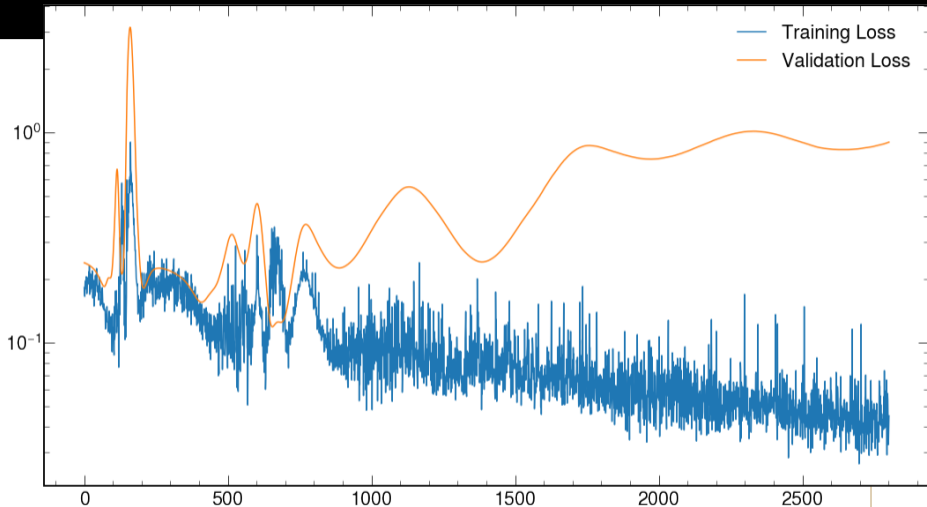
$$\Gamma : \mathbb{E}^7 \longrightarrow \mathbb{R}$$

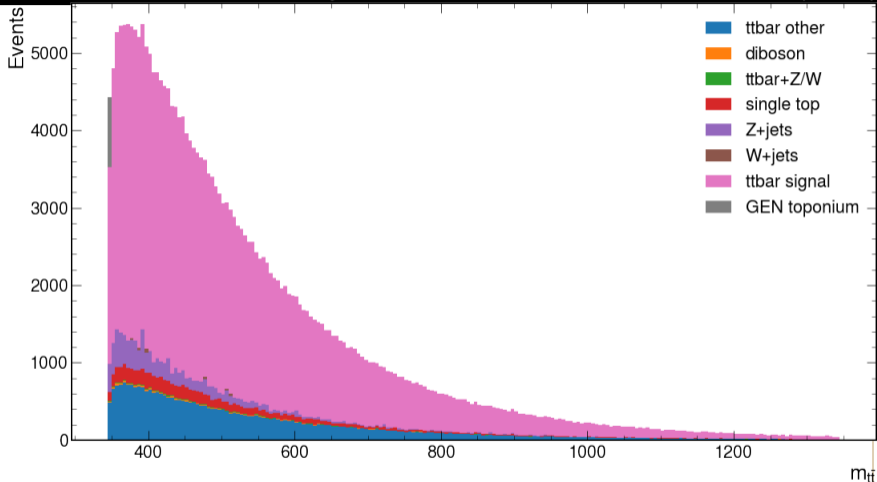
where  $\Gamma$  is a kernel map determined by the choice of classifier architecture. The output of  $\Gamma(\mathbf{V})$  is a probability that that  $\mathbf{V}$  is a signal event as determined by the classifier.

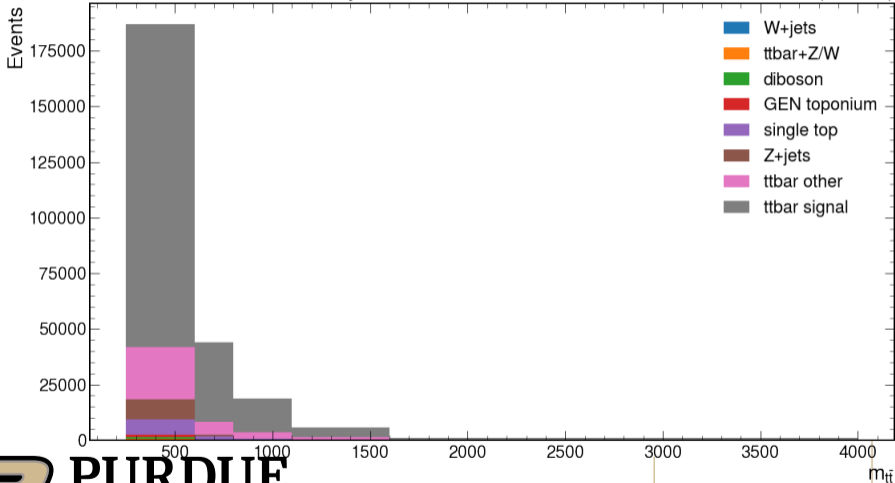
# *Results*

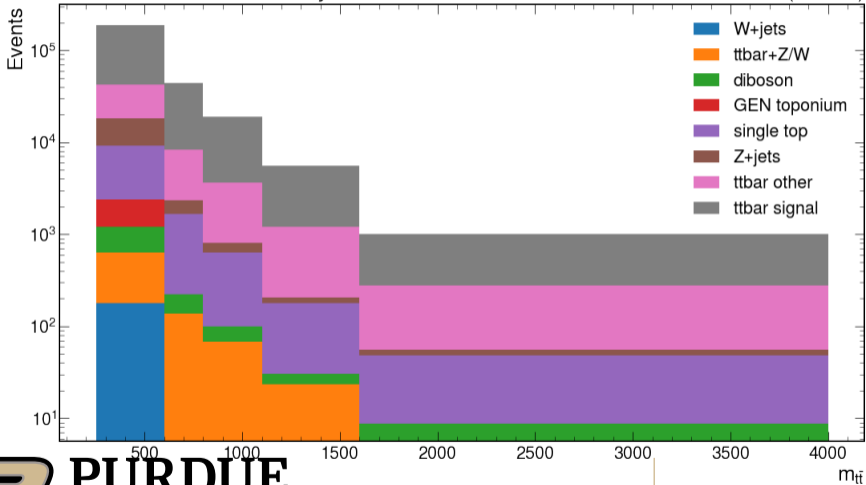


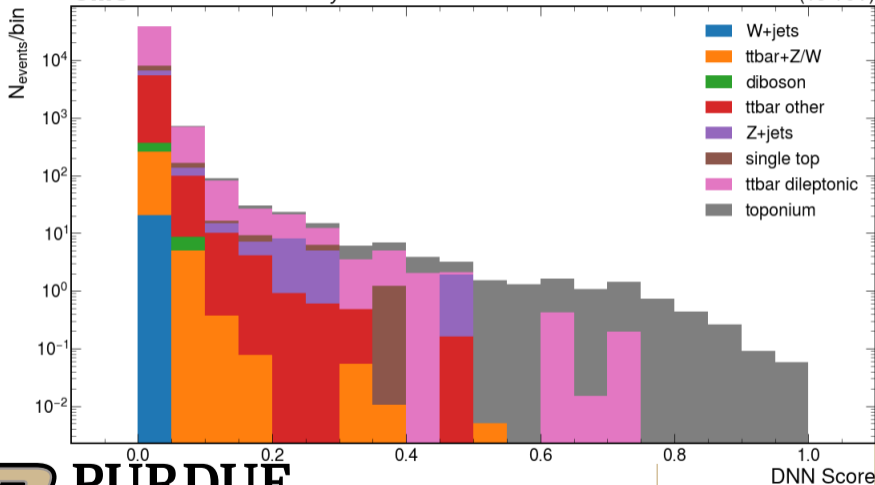












# *Questions and Feedback*

Thank you for your time.

Thank you for allowing me to take part in this internship.

Special thanks to AJ and Jason T. of the Jung Group, and Dr. Jung, Dr. Pando, Dr. Geller, Dr. Dahl, Dr. Fong, Dr. York, and Dr. Swordy.

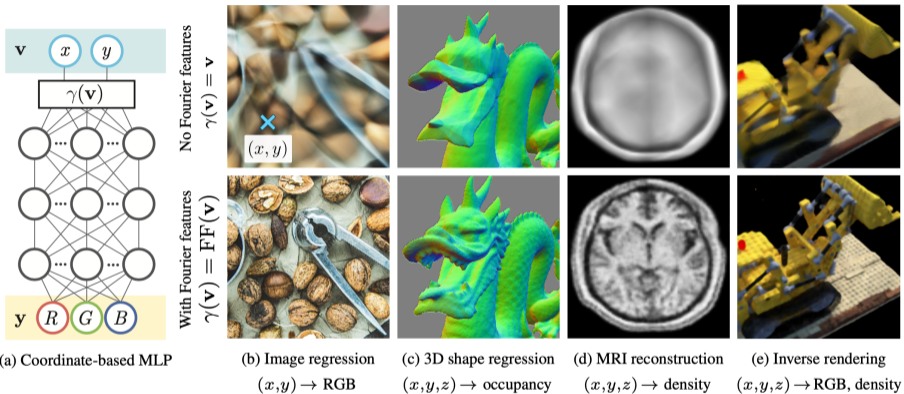
# References

- "Measurement of the top quark polarization and  $t\bar{t}$  spin correlations using dilepton final states in proton-proton collisions at  $\sqrt{s}=13$  TeV" (2019) [arXiv:1907.03729v2]
- "Measurement of the  $t\bar{t}$  production cross-section and lepton differential distributions in e dilepton events from pp collisions at  $\sqrt{s}=13$  TeV with the ATLAS detector" (2020) [arxiv:1910.08819]
- "Signatures of toponium formation in LHC run 2 data" (2021) [arXiv:2102.11281v2]

# *Mathematical Statement of Classifier Problem*

For each event, there is a list of cms2016 data and delphes gen-level data for toponium.\*\* How to overcome class imbalance when training the classifier? SMOTE? No. Upsampling? Yes.





$width = 400$

$dropout\_probability = 0.05$

$best\_model = nn.Sequential(nn.Dropout(dropout\_probability),$

$nn.Linear(7, width), nn.BatchNorm1d(width),$

$nn.LeakyReLU(0.01), nn.Linear(width, width),$

$nn.BatchNorm1d(width), nn.LeakyReLU(0.01),$

$nn.Dropout(dropout\_probability),$

$nn.Linear(width, width),$

$nn.BatchNorm1d(width), nn.LeakyReLU(0.01), nn.Dropout(dropout\_probability)$

$nn.Linear(width, width),$

$nn.BatchNorm1d(width), nn.LeakyReLU(0.01), nn.Dropout(dropout\_probability)$

$nn.Linear(width, 1), nn.Sigmoid())$  pass output through sigmoid to get predictions in

# Goal of Toponium analysis

A lot of work done so far:

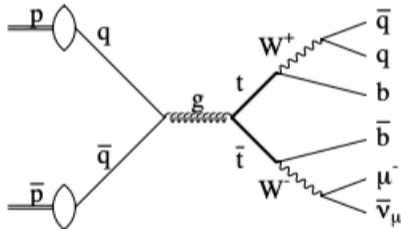
- MadGraph
- PYTHIA
- delphes
- BiLSTMs
- transformers??

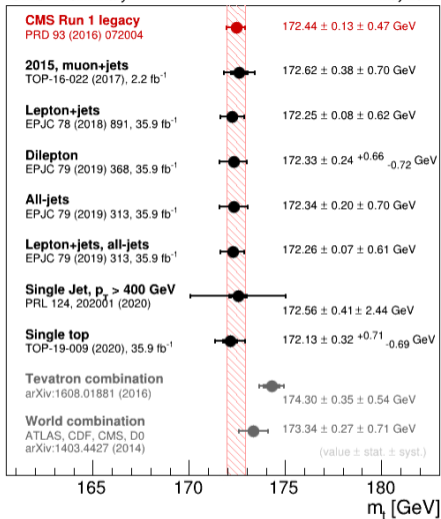
**Goal:** PUT A LIMIT ON THE CROSS SECTION FOR TOPONIUM USING HIGGS PAG - COMBINE AND MACHINE LEARNING TOOLS.

# Backup 0

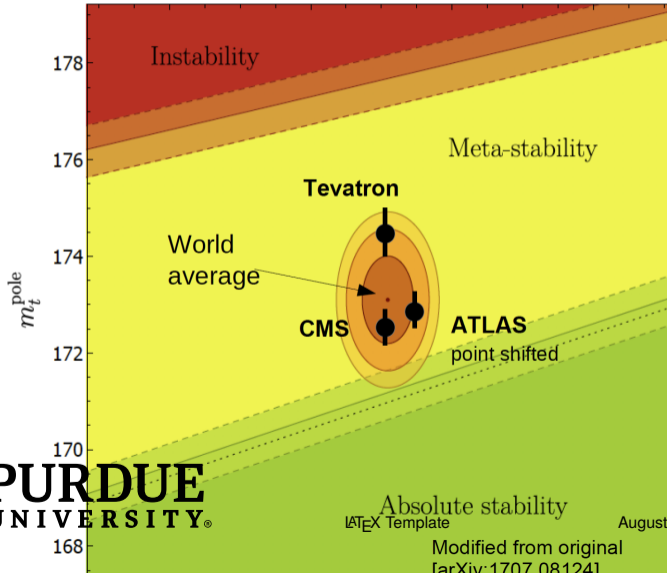
Top quark discovery (Tevatron, 1995)

$p\bar{p} \rightarrow t\bar{t} \rightarrow (bW^+)(\bar{b}W^-) \rightarrow (bq\bar{q})(\bar{b}\mu^-\bar{\nu}_\mu) \rightarrow 4\text{jets} + \mu^- + \bar{\nu}_\mu$





# Backup 1



# Backup 2

## $t\bar{t}$ Kinematic Reconstruction

➤ The four momenta of the neutrinos in the decay are needed to fully reconstruct the top quark four momenta

➤ Analytical solutions are obtained by applying **six kinematic constraints**:

$$E_x^{miss} = p_{x,\nu} + p_{x,\bar{\nu}}$$

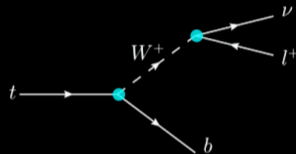
$$E_y^{miss} = p_{y,\nu} + p_{y,\bar{\nu}}$$

$$m_{W^+}^2 = (E_{l^+} + E_\nu)^2 - (p_{x,l^+} + p_{x,\nu})^2 - (p_{y,l^+} + p_{y,\nu})^2 - (p_{z,l^+} + p_{z,\nu})^2$$

$$m_{W^-}^2 = (E_{l^-} + E_{\bar{\nu}})^2 - (p_{x,l^-} + p_{x,\bar{\nu}})^2 - (p_{y,l^-} + p_{y,\bar{\nu}})^2 - (p_{z,l^-} + p_{z,\bar{\nu}})^2$$

$$m_{\bar{t}}^2 = (E_{l^+} + E_\nu + E_b)^2 - (p_{x,l^+} + p_{x,\nu} + p_{x,b})^2 - (p_{y,l^+} + p_{y,\nu} + p_{y,b})^2 - (p_{z,l^+} + p_{z,\nu} + p_{z,b})^2$$

$$m_t^2 = (E_{l^-} + E_{\bar{\nu}} + E_{\bar{b}})^2 - (p_{x,l^-} + p_{x,\bar{\nu}} + p_{x,\bar{b}})^2 - (p_{y,l^-} + p_{y,\bar{\nu}} + p_{y,\bar{b}})^2 - (p_{z,l^-} + p_{z,\bar{\nu}} + p_{z,\bar{b}})^2$$



➤ Choose the solution that gives lowest  $t\bar{t}$  invariant mass

➤ Reconstruct each event 100 times smearing  $W$  mass, jet and lepton directions and energies by their resolutions

➤ The solution to the top quark three momenta is the weighted average of the solutions corresponding to all lepton-jet

pairs (excluding the ones with most  $b$ -tags);

## Backup slide 3

*"Owing to its large mass, the lifetime of the top quark is very short. Consequently, the top pairs produced in the process  $q\bar{q} \rightarrow t\bar{t}$  do not have time to form bound states, such as those observed in the resonant production of charmonium ( $c\bar{c}$ ) and bottomonium ( $b\bar{b}$ ) states"... "Hence the top quark lifetime is of order  $t = 1/\Gamma \approx 5 \times 10^{-25}$  s. This is sufficiently short that the top quarks produced at the Tevatron decay in a distance of order  $10^{-16}$  m. This is small compared to the typical length scale for the hadronisation process, and therefore the  $t\bar{t}$  pairs produced at the Tevatron not only decay before forming a bound state, but also decay before hadronising." (Thomson)*