

Massive Colour-Octet Bosons in Top-Quark Pair Production

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JOHANNES GUTENBERG
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Portorož 2011, April 11-14, Portorož, Slovenia

Three playgrounds for massive colour-octet bosons

1.) Theory

- axigluons in chiral colour models
- Kaluza-Klein gluons in warped extra dimensions
- colorons in technicolour models ...

2.) Phenomenology

- generically have strong couplings to quarks
 - interfere with standard-model gluon amplitude
- large effects in LHC observables

3.) Anomaly

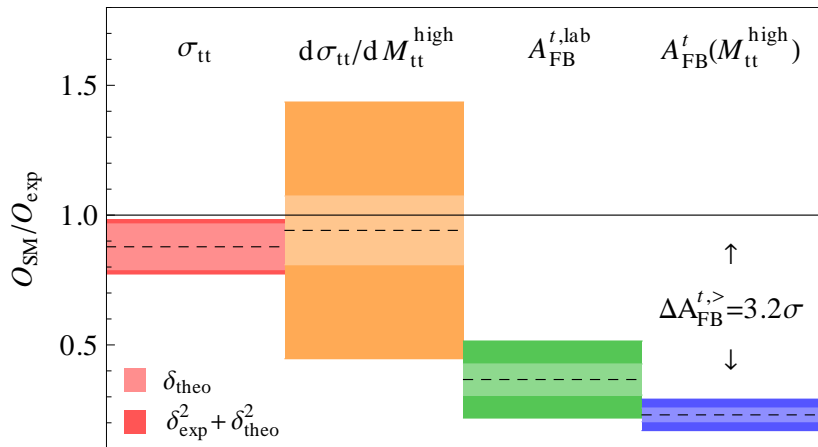
- may exhibit flavour non-universal couplings
- explain top-quark forward-backward asymmetry at Tevatron

This talk: constraints and effects in $t\bar{t}$ production

Top-quark pair production at the Tevatron

Standard-model predictions O_{SM} versus measurements O_{exp}

[CDF '11]

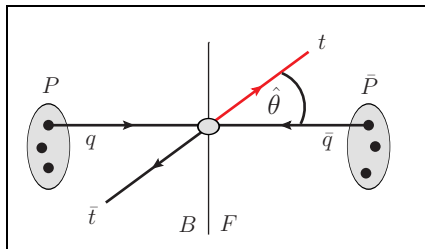


Asymmetric observables lie significantly below the measurement.

Forward-backward asymmetry in $t\bar{t}$ production

Charge-(a)symmetric cross section

$$\sigma_a(s) = \int_0^1 \cos \hat{\theta} \left[\frac{d\sigma(p\bar{p} \rightarrow t\bar{t}X)}{d \cos \hat{\theta}} - (+) \frac{d\sigma(p\bar{p} \rightarrow \bar{t}tX)}{d \cos \hat{\theta}} \right]$$



$$A_{\text{FB}}^t = \frac{N_t(F) - N_t(B)}{N_t(F) + N_t(B)} = \frac{\sigma_a}{\sigma_s}$$

Measurement at Tevatron: inclusive and in bins of invariant mass $M_{t\bar{t}}$

$$(A_{\text{FB}}^{t,\text{lab}})_{\text{exp}} = (15.0 \pm 5.5)\%$$

$$(A_{\text{FB}}^t)_{\text{exp}}^{M_{t\bar{t}} > 450 \text{ GeV}} \equiv (A_{\text{FB}}^{t,>})_{\text{exp}} = (47.5 \pm 11.2)\% \quad [\text{CDF '11}]$$

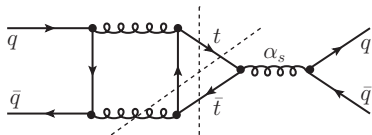
$$(A_{\text{FB}}^t)_{\text{exp}}^{M_{t\bar{t}} < 450 \text{ GeV}} \equiv (A_{\text{FB}}^{t,<})_{\text{exp}} = (-11.6 \pm 15.3)\%$$

Large charge asymmetry from massive colour octets

Need amplitude antisymmetric under $t \leftrightarrow \bar{t}$ interchange.

Standard Model: A_{FB}^t at NLO

$$\sigma_a \sim \alpha_s^3$$



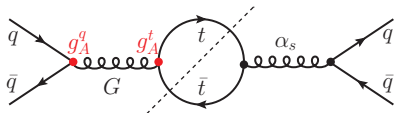
+ 10% electroweak corrections

NLO+NNLL: [Ahrens et al. '11]

$$(A_{\text{FB}}^{t,\text{lab}})_{\text{SM}} = 4\% - 5.6\%$$

New physics: A_{FB}^t at tree level

$$\sigma_a \sim -\alpha_s g_A^q g_A^t$$



+ t - and u -channel new physics

Tree-level effects in cross section:

$$\sigma_s \sim \alpha_s g_V^q g_V^t.$$

Good fit to $t\bar{t}$ observables expected for $g_A^q g_A^t < 0$ and $|g_V/g_A| < 1$.

Axigluons from chiral colour

Extended colour gauge group

[Frampton, Glashow '87][Frampton, Shu, Wang '09]

$$SU(3)_A \times SU(3)_B \xrightarrow{\langle \Sigma \rangle} SU(3)_C$$

with associated gauge fields A_μ , B_μ and couplings g_A , g_B .

Mass eigenstates: **axigluon** G with $M_G = \mathcal{O}(1 \text{ TeV})$ and massless gluon g

$$\begin{pmatrix} G_\mu \\ g_\mu \end{pmatrix} = \begin{pmatrix} \sin \theta & -\cos \theta \\ \cos \theta & \sin \theta \end{pmatrix} \begin{pmatrix} A_\mu \\ B_\mu \end{pmatrix}, \quad \tan \theta = \frac{g_A}{g_B},$$

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Assign charges to light ($\ell = q_{1,2}$) and heavy ($h = q_{3,4}$) quarks

	$Q_{1,2}$	$u_{1,2}^c$	$d_{1,2}^c$	$Q_{3,4}$	$u_{3,4}^c$	$d_{3,4}^c$
$SU(3)_A$	3	1	1	1	$\bar{3}$	$\bar{3}$
$SU(3)_B$	1	$\bar{3}$	$\bar{3}$	3	1	1

$$g_V^\ell = g_V^h$$

$$g_A^\ell = -g_A^h$$

Important for $t\bar{t}$ observables:

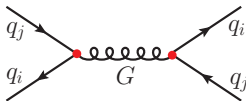
$$g_A^\ell g_A^h < 0, \quad \left| \frac{g_V}{g_A} \right| = \cos(2\theta) < 1$$

Constraints on massive colour octets

[Bai, Hewett, Kaplan, Rizzo '11, Haisch, SW '11, et al.]

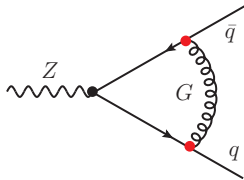
- Flavour-changing neutral currents at tree level

Neutral meson mixing



- Electroweak precision observables

$Zq\bar{q}$ coupling, oblique corr. S, T



- Dijet production at the LHC

[ATLAS, CMS '10,'11]

Resonances in $pp \rightarrow G \rightarrow jj$ spectrum and angular distribution

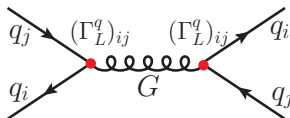
Flavour-changing neutral currents

$$\mathcal{L} \supset - (Y_u)_{ij} \bar{Q}_L^i H u_R^j + g_s \left[(\Gamma_L^u)_{ij} \bar{u}_L^i \not{G} u_L^j + (\Gamma_R^u)_{ij} \bar{u}_R^i \not{G} u_R^j + (u \leftrightarrow d) \right]$$

Flavour violation due to **misalignment** of axigluon- with Yukawa-couplings,

$$\Gamma_L^{u,d} = U_{u,d}^\dagger \hat{\Gamma}_L U_{u,d} \quad \text{with } \hat{\Gamma}_L = \text{diag}(g_L^\ell, g_L^\ell, g_L^h), \quad (\Gamma_R^{u,d} = \hat{\Gamma}_R).$$

Neutral meson mixing at tree level:



$$\sim \frac{(g_L^h - g_L^\ell)^2}{M_G^2} = \frac{4}{M_G^2 \sin^2(2\theta)}$$

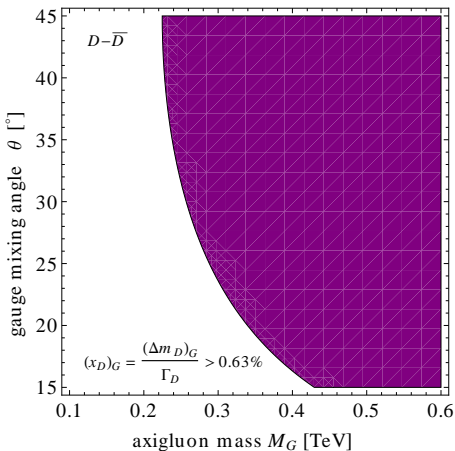
SM weak charged currents require $U_u = U_d V_{CKM}^\dagger$.

Align Γ_L^d with down-quark Yukawa coupling, $U_d = 1$:

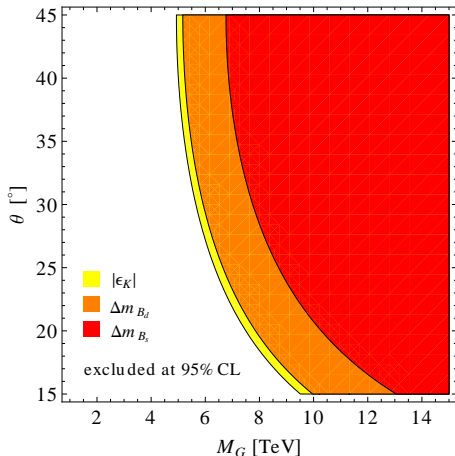
- no flavour violation in $K - \bar{K}$, $B_d - \bar{B}_d$, $B_s - \bar{B}_s$ mixing
- minimal flavour violation in $D - \bar{D}$ mixing $\sim (V_{ub}^* V_{cb})^2$

Constraints on axiglucos from meson mixing

Align down-quark sector, $U_d = 1$.



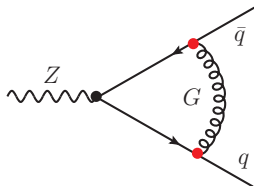
Align up-quark sector, $U_u = 1$.



Minimal constraint from $D - \bar{D}$ mixing: $M_G \sin(2\theta) > 225 \text{ GeV}$

Corrections to electroweak couplings

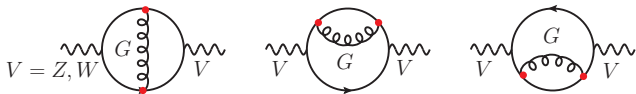
$Zq\bar{q}$ vertex



Effective coupling $Zq\bar{q}$ ($m_q = 0$)

$$\delta \mathcal{G}_{L,R}^{l,h} \sim (g_{L,R}^{l,h})^2 \frac{M_Z^2}{M_G^2} \log \left(\frac{M_Z^2}{M_G^2} \right)$$

Oblique corrections

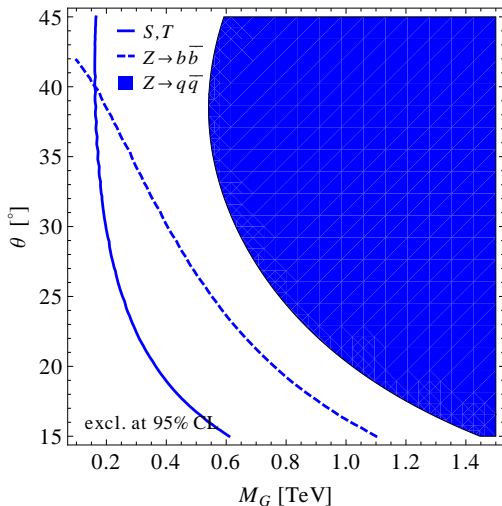


SM: Peskin-Takeuchi parameters $S, T, U = 0$ at reference point.

Axigluon: $\{S, T\} \sim \left\{ 1, \frac{m_t^2}{M_Z^2} \right\} \left[(g_L^h)^2 + 2 (g_R^h)^2 \right] \frac{m_t^2}{M_G^2} \log^2 \left(\frac{m_t^2}{M_G^2} \right)$

Electroweak precision observables

Measurements at LEP: $e^+e^- \rightarrow Z \rightarrow q\bar{q}$ at $\sqrt{S} = M_Z$ [LEP EWWG '05]



$Z \rightarrow q\bar{q}$ ($\Gamma_Z, \sigma_{\text{had}}, Zb\bar{b}$):

$$M_G \gtrsim 500 \text{ GeV}$$

$Z \rightarrow b\bar{b}$:

combined fit to

- partial decay rate R_b
- asymmetry parameters A_b, A_{FB}^b

Oblique corrections S, T :

require fine-tuning of
4th quark generation

Dijet production at the LHC

Consider $pp \rightarrow G \rightarrow jj$ at $\sqrt{S} = 7$ TeV.

[ATLAS, CMS '10,'11]

- Dijet invariant mass spectrum

constrains $\sigma(pp \rightarrow G) \mathcal{B}(G \rightarrow jj)$
for resonances with $\Gamma_G/M_G \lesssim 15\%$.

Minimal constraints on axigluon mass

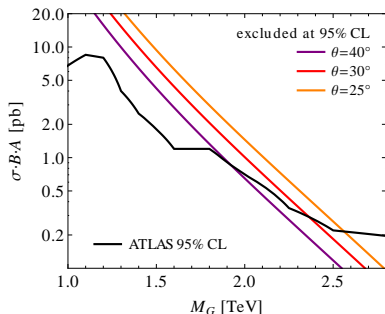
$$M_G \gtrsim 1.8 \text{ TeV} \quad \text{at 95\%CL.}$$

- Angular distribution of jets

constrains contact interactions $\mathcal{L}_{4q} = \frac{2\pi}{\Lambda^2} (\bar{q}_L \gamma^\mu q_L) (\bar{q}_L \gamma_\mu q_L),$

$$\Lambda \gtrsim 5.7 \text{ TeV} \quad \text{at 95\%CL.}$$

Translated into bound on axigluon: $M_G \gtrsim 1.9 \text{ TeV}.$



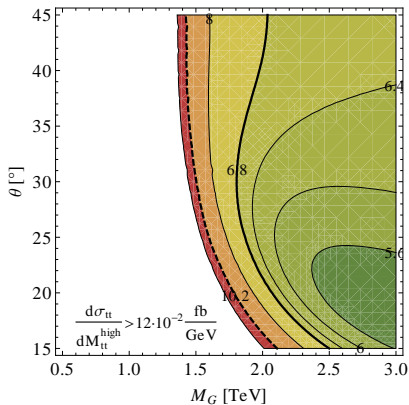
Axiglons in top-quark pair production

Consider symmetric and asymmetric cross sections at high $M_{t\bar{t}}$.

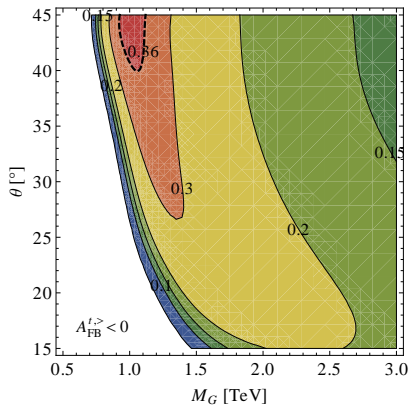
$$\left(\frac{d\sigma_{t\bar{t}}}{dM_{t\bar{t}}^{\text{high}}}\right) \text{ prefers } M_G \gtrsim 2 \text{ TeV}$$

\leftrightarrow

$$A_{\text{FB}}^{t,>} \text{ prefers } M_G \simeq 1 \text{ TeV}$$

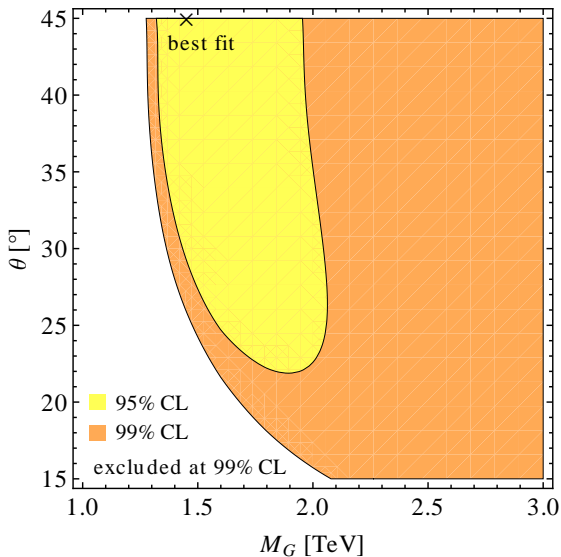


$$\begin{aligned} \left(\frac{d\sigma_{t\bar{t}}}{dM_{t\bar{t}}}\right)_{\text{exp}} M_{t\bar{t}} &\in [0.8, 1.4] \text{ TeV} \\ &= (6.8 \pm 3.4) \cdot 10^{-2} \frac{\text{fb}}{\text{GeV}} \end{aligned}$$



$$\begin{aligned} (A_{\text{FB}}^t)_{\text{exp}} M_{t\bar{t}} &> 450 \text{ GeV} \\ &= (47.5 \pm 11.2)\% \end{aligned}$$

Global fit to $t\bar{t}$ observables $\sigma_{t\bar{t}}$, $\frac{d\sigma_{t\bar{t}}}{dM_{t\bar{t}}^{\text{high}}}$, $A_{\text{FB}}^{t,<}$, and $A_{\text{FB}}^{t,>}$

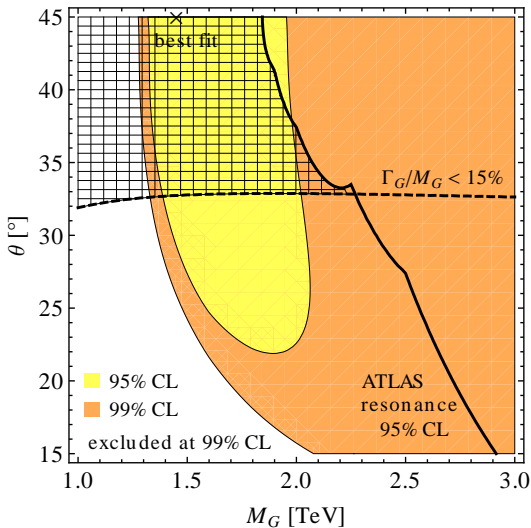


Best fit point x: $\chi^2 = 7.5$
 (cf. SM: $\chi^2 = 13.2$)

$$M_G = 1.45 \text{ TeV}, \theta = 45^{\circ}$$

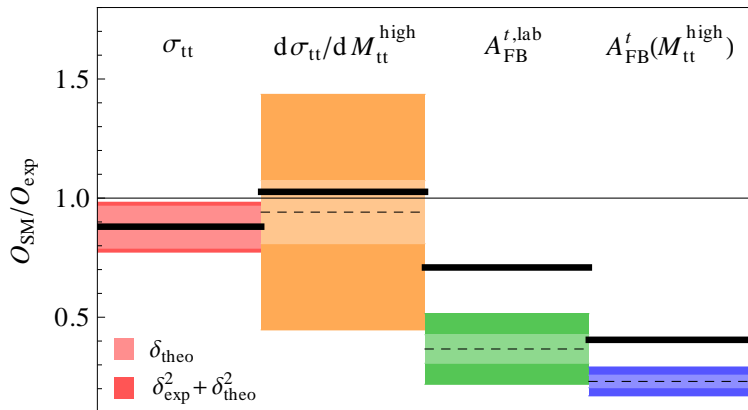
Constraints from resonance search in dijet production

$pp \rightarrow G \rightarrow jj$ @ $\sqrt{S} = 7 \text{ TeV}$, ATLAS 36pb^{-1}



Axigluons in $t\bar{t}$ observables

consistent w/ all constraints: $M_G = 1.9 \text{ TeV}$, $\theta = 45^\circ$



Tension between $\sigma_{t\bar{t}}$ and A_{FB}^t at high $M_{t\bar{t}}$:

Both increase as M_G decreases.

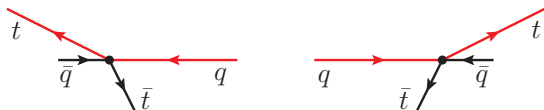
$$(A_{\text{FB}}^{t,>})_G = 19.2\%$$

$$(A_{\text{FB}}^{t,\text{lab}})_G = 11.2\%$$

Top-quark charge asymmetry at the LHC

The process $pp \rightarrow t\bar{t}$ is symmetric \Rightarrow no forward-backward asymmetry.

Top quarks are preferentially emitted along the beam axis at large rapidity:



Charge-asymmetric contributions to $q\bar{q} \rightarrow t\bar{t}$ can be probed by an asymmetry in pseudo-rapidities η ,

[Antunano, Kühn, Rodrigo '98,'08]

$$A_\eta = \frac{N(\Delta\eta > 0) - N(\Delta\eta < 0)}{N(\Delta\eta > 0) + N(\Delta\eta < 0)}, \quad \Delta\eta = |\eta_t| - |\eta_{\bar{t}}|,$$

$$= 0.060 \pm 0.134_{\text{stat}} \pm 0.026_{\text{syst}}.$$

[CMS '11]

To be taken home

Indirect constraints disfavour colour-octet bosons with masses at the weak scale.










For an axigluon with non-universal couplings to quarks:

- $D - \bar{D}$ meson mixing: $M_G \gtrsim 250 \text{ GeV}$
- $Z \rightarrow q\bar{q}$: $M_G \gtrsim 500 \text{ GeV}$

The top-quark charge asymmetry can be explained by such an axigluon with $M_G \sim 1.5 \text{ TeV}$.

Direct constraints from dijet production at the LHC challenge this solution by requiring $M_G \gtrsim 1.9 \text{ TeV}$.

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