

Top Flavor Gauge Boson Physics (& Flavor Triviality)

Seung J. Lee

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Portoroz 2011

Based on works in collaboration with
Delaunay, Gedalia, Perez & Ponton (x2)
Csáki, Kagan, Perez & Weiler (in preparation)

Outline

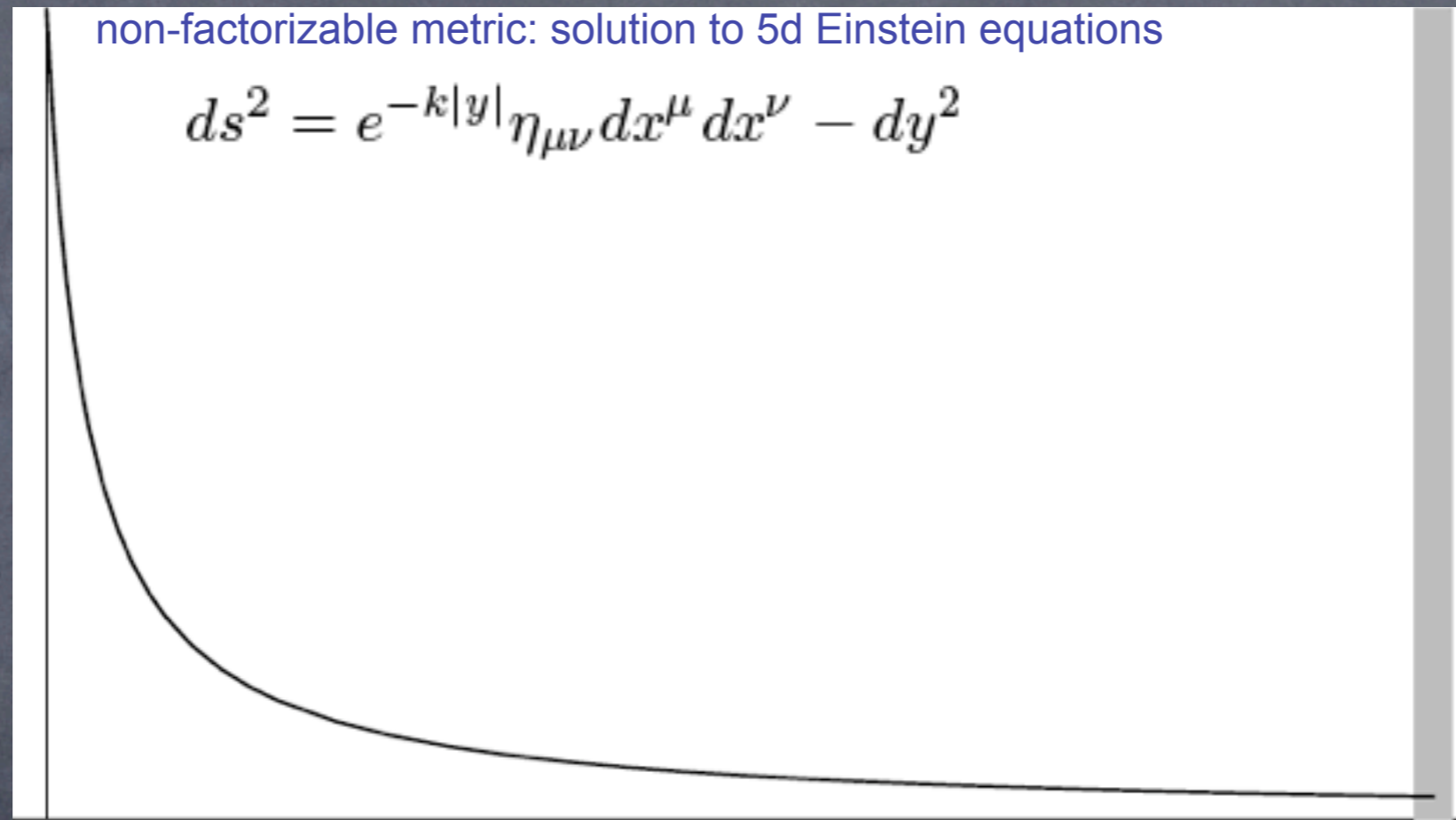
- Introduction: Warped XD & little CP problem
- Flavor symmetry (decoupling flavor and CP):
“Flavor triviality” and flavor gauge boson
- Improved EW fit: Global fit to the EW observables
- Signals from Flavor Trivial Warped Extra Dimension: KK gluon and flavor gauge boson
 - top physics at LHC and Tevatron
- Summary

RS & the hierarchy problem

Randall-Sundrum,
PRL (99)

- XD curved, but brane remains static and flat

SM+Higgs



UV (y=0)

IR (y=b)

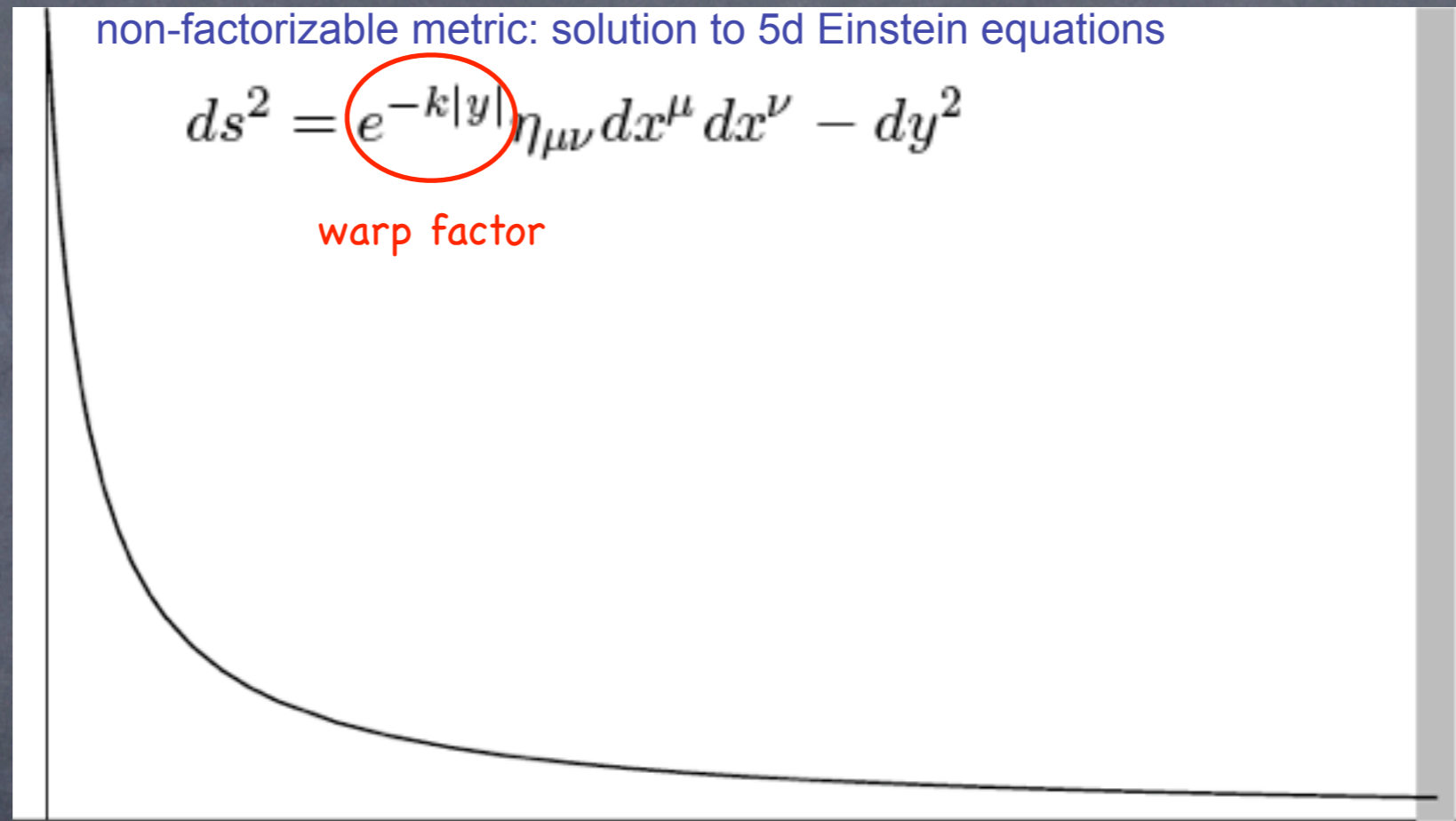
- address the hierarchy problem by providing concrete 5D realizations of various 4D EWSB scenarios: composite higgs, pseudo-Goldstone higgs, higgsless, etc.

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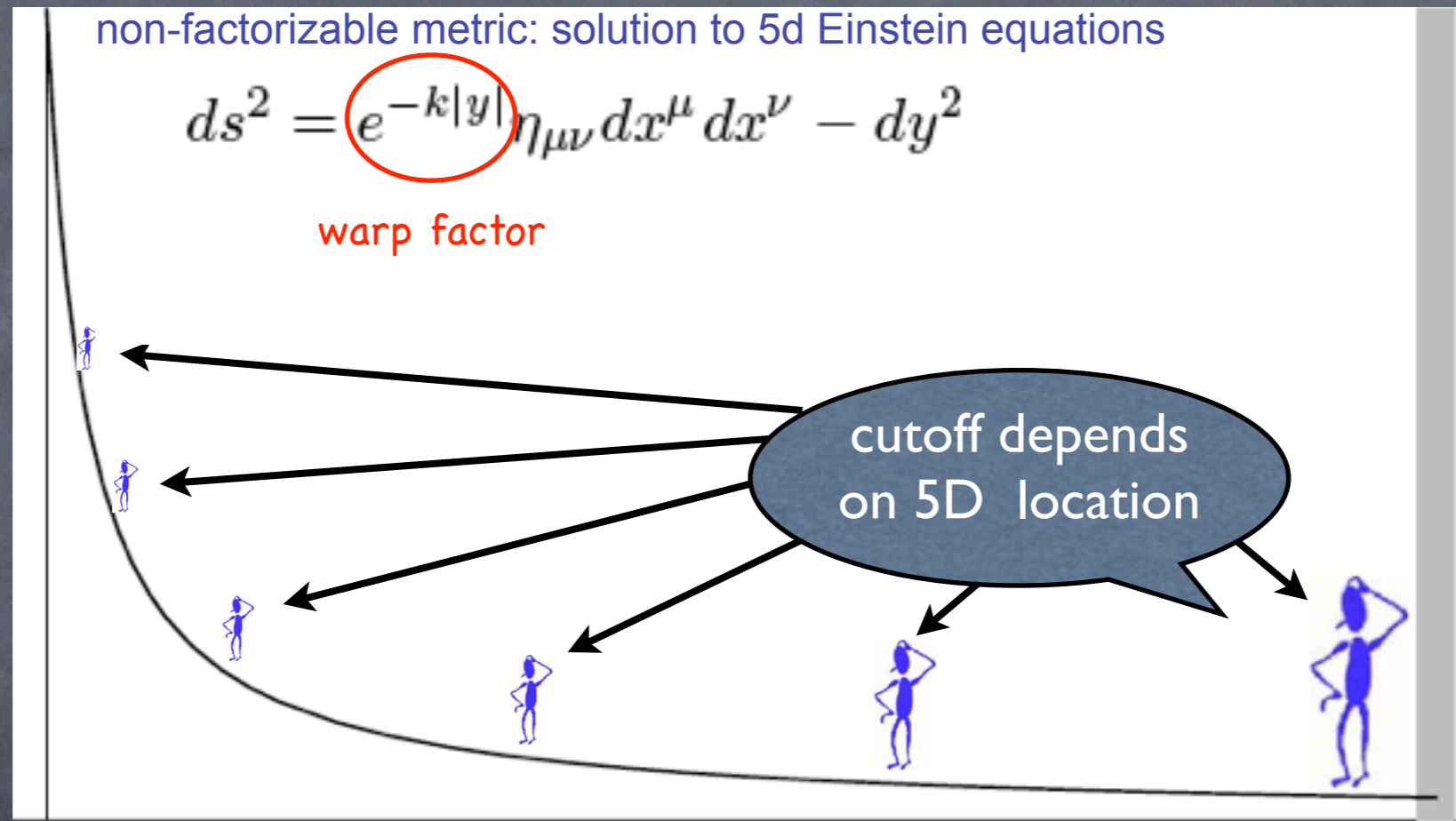
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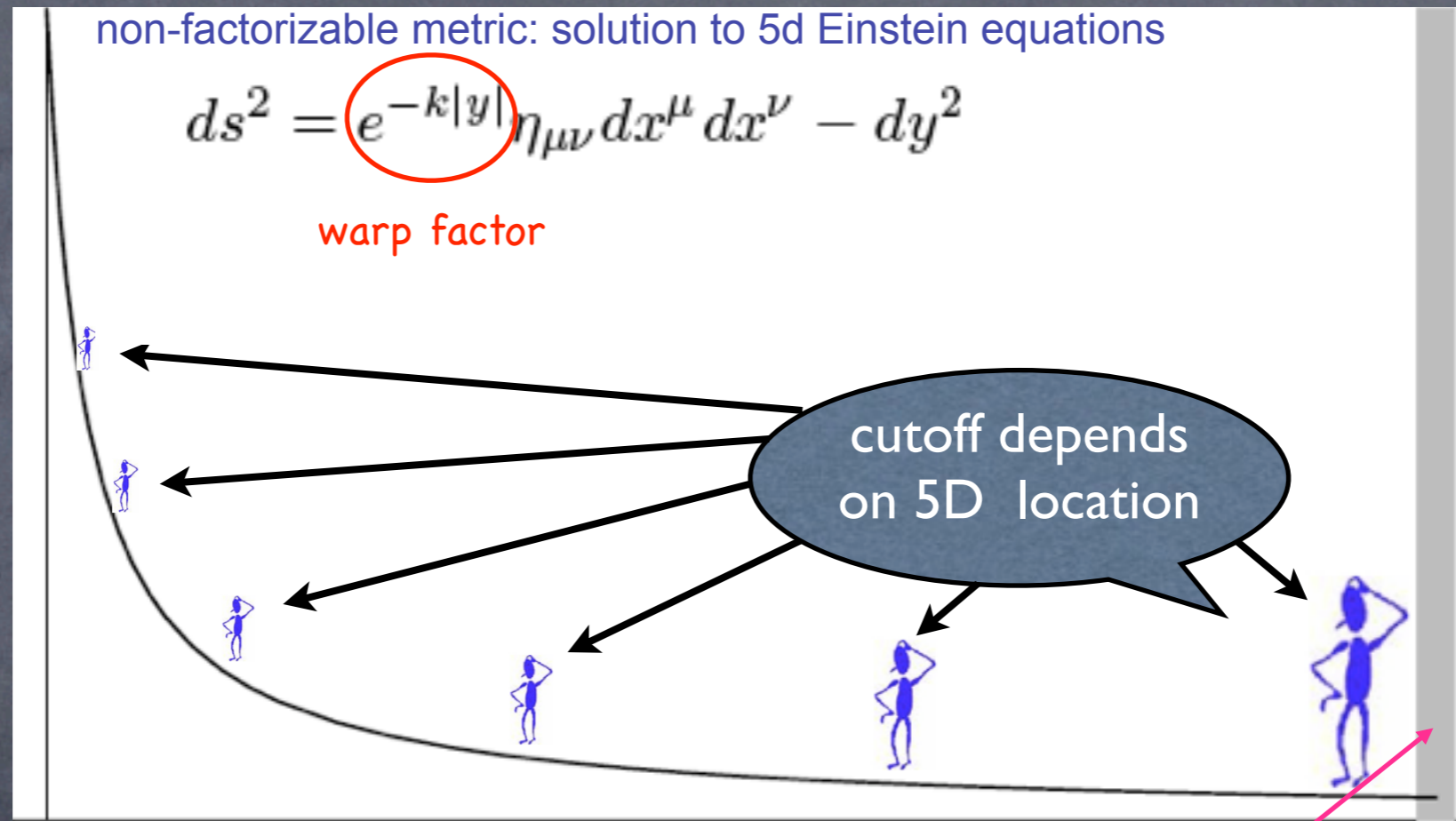
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UV ($y=0$)

Assume we live here

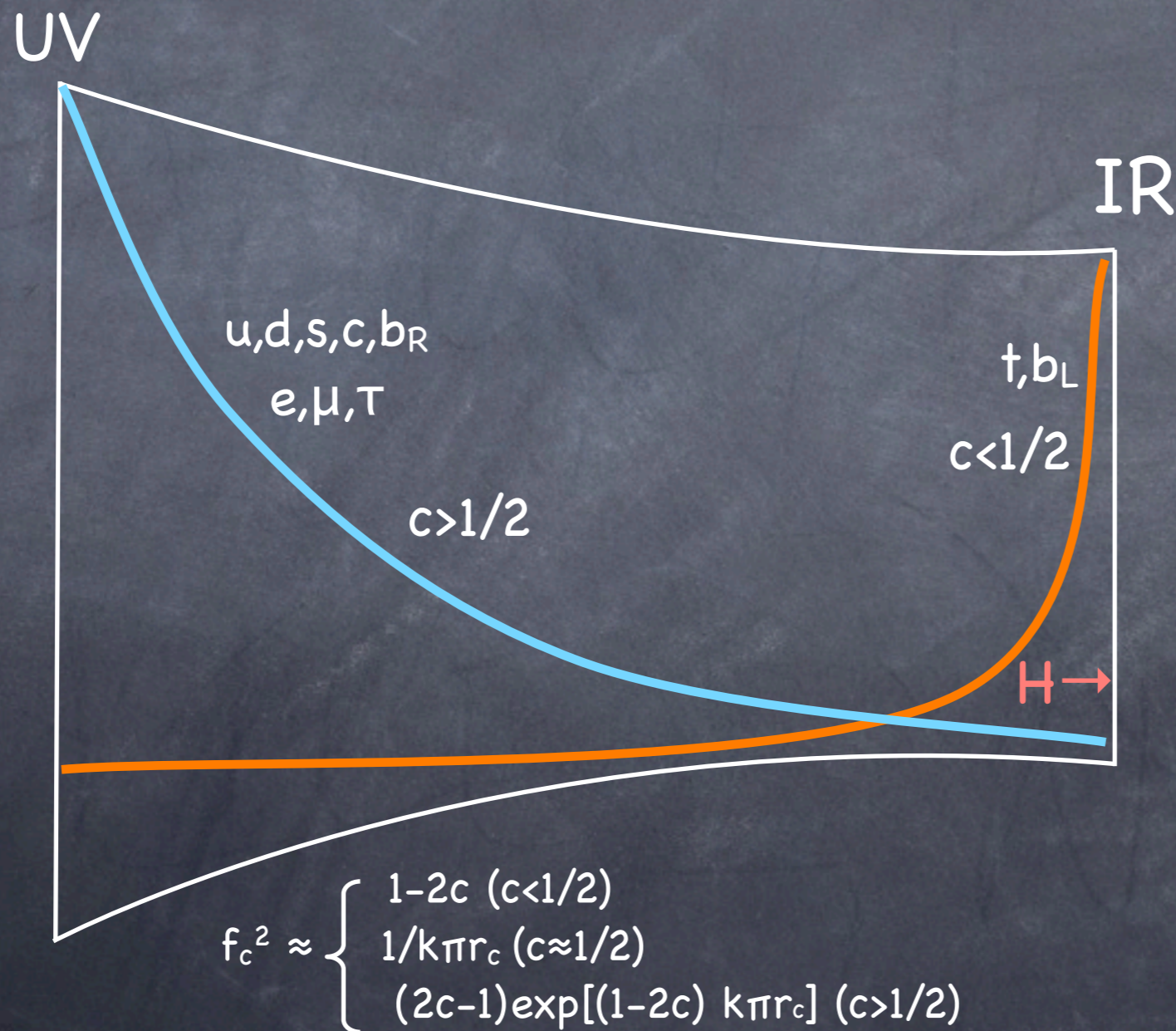
IR ($y=b$)

(red shift to TeV scale)
SM localized on IR brane

- address the hierarchy problem by providing concrete 5D realizations of various 4D EWSB scenarios: composite higgs, pseudo-Goldstone higgs, higgsless, etc.

Anarchic RS

- **Anarchic model:** Localization address flavor hierarchy and also give protection against flavor & CPV (with $C_{Q,U,D} \sim O(1)$)



$$Y_{U,D} \sim O(1)$$

$$M_{U,D} \sim v f_{c_Q} Y_{U,D} f_{c_{U,D}}$$

$$V_{CKM_{ij}} \approx f_{c_{Qi}} / f_{c_{Qj}}$$

Anarchic RS: Little CP Problem

- Stringent bound from flavor physics:

$$\epsilon_K = \frac{A(K_L \rightarrow (\pi\pi)_{I=0})}{A(K_S \rightarrow (\pi\pi)_{I=0})}$$

$$\epsilon'/\epsilon_K$$

Davidson, Isidori & Uhlig (07);
Gedalia, Isidori & Perez (09).

Csaki, Falkowski & Weiler;
Agashe, Azatov & Zhu (08)

$$\epsilon_{ij} \frac{\bar{Q}_i Q_j \bar{Q}_i Q_j}{\Lambda_F}$$

$$|\epsilon_{ij}| \sim O(1) \rightarrow \Lambda_F > O(6) \text{ TeV!}$$

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New Physics > 5–6 TeV: out of reach from LHC
for best case

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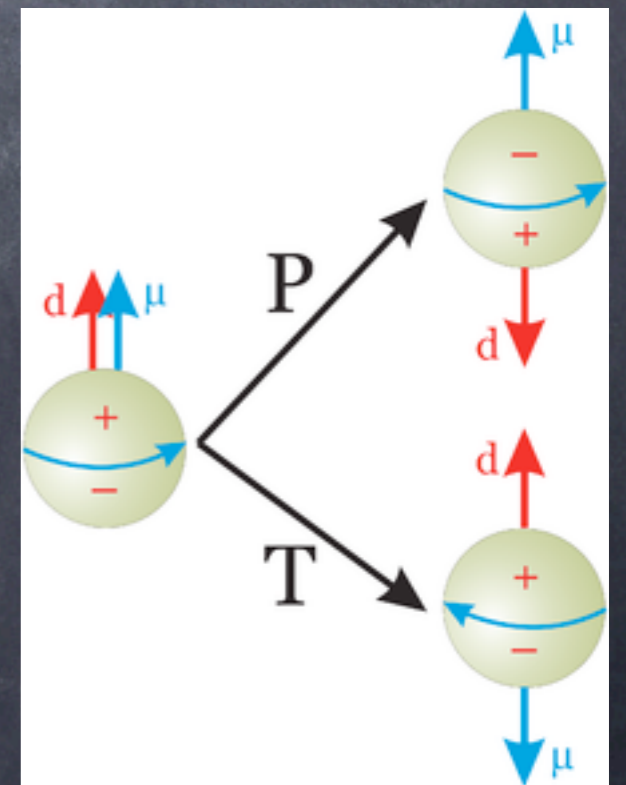
New Physics > 5-6 TeV: out of reach from LHC
for best case

- neutron Electric Dipole Moment
constraint: $O(20)$ larger than
bounds

measure for the distribution of positive and negative charge inside the neutron

Agashe, Perez & Soni (04)

Delaunay, Gedalia, SJL, Perez & Ponton (11)



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$$\epsilon_{ij} \frac{\bar{Q}_i Q_j \bar{Q}_i Q_j}{\Lambda_F^2}$$

New Physics

combining EDM and ϵ_K , $m_{KK} > 8.4$ TeV for anarchic RS

$\Lambda_F > O(6)$ TeV!

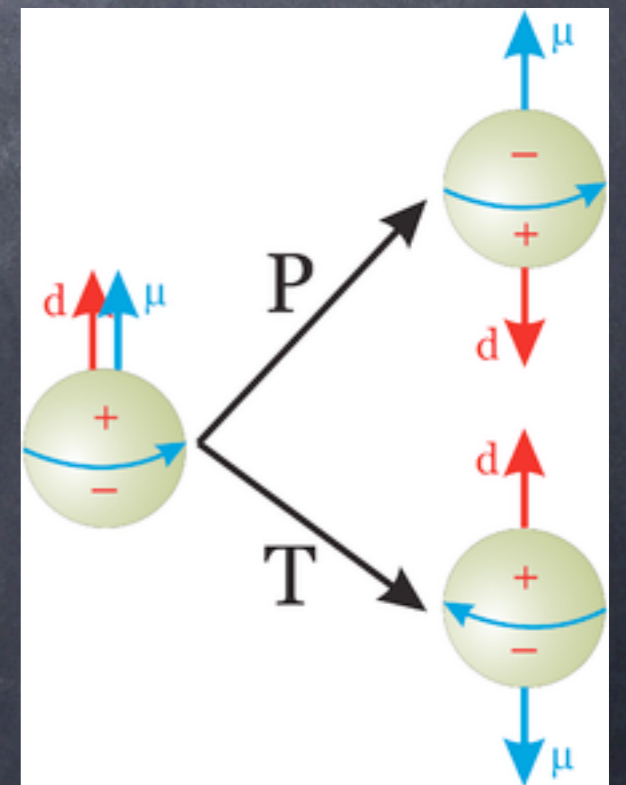
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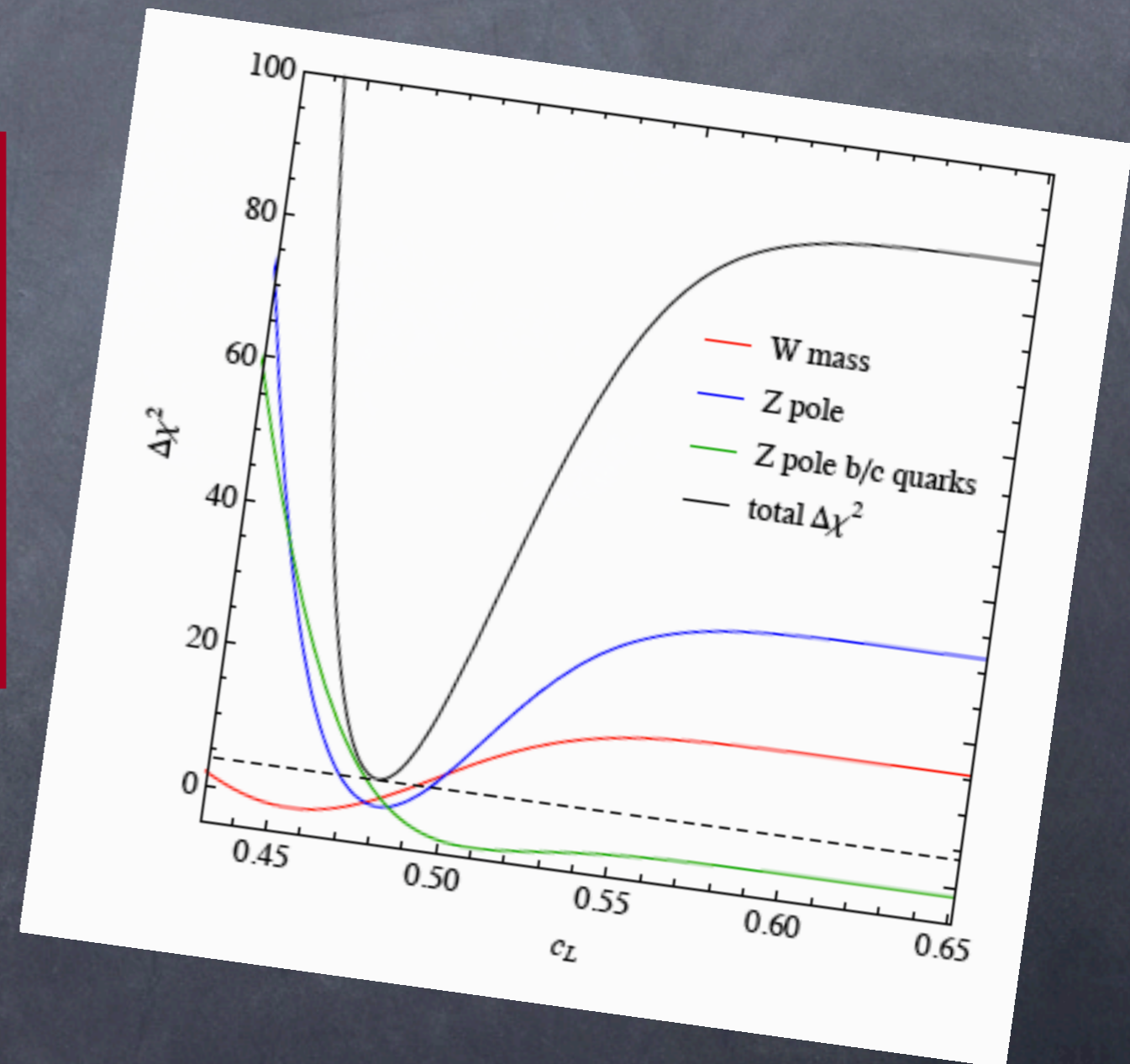
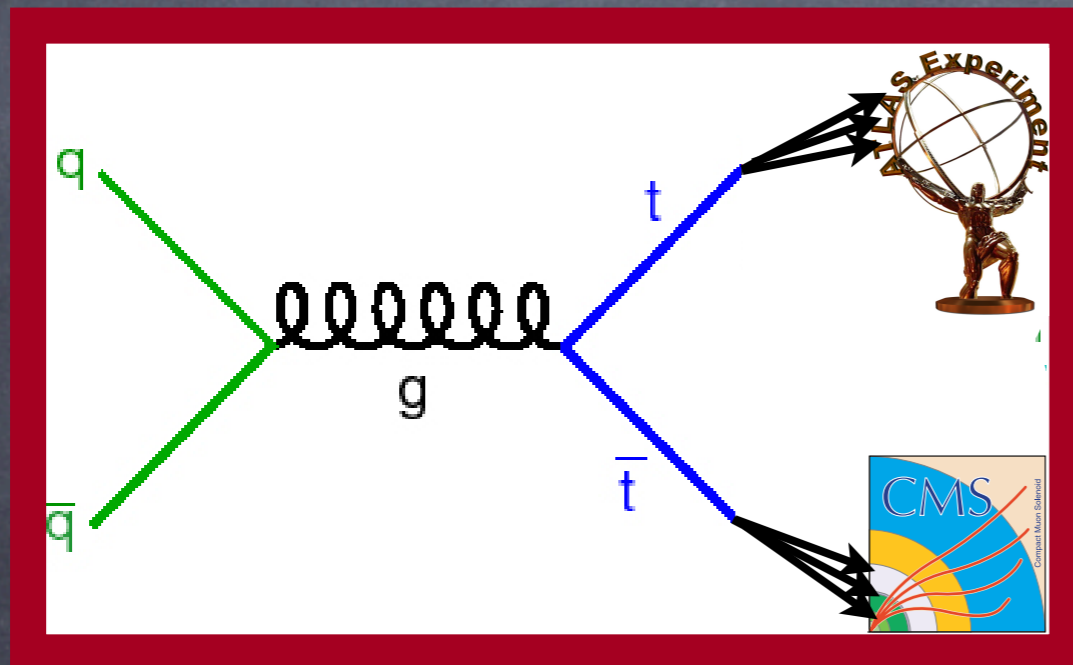
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extraordinary pheno' from "flavor triviality" & Improved Naturalness

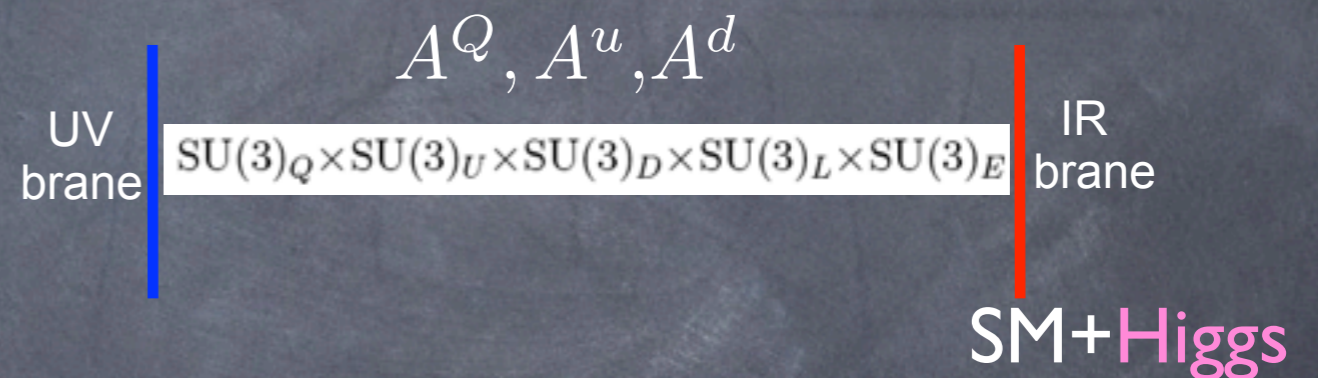
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Rattazzi-Zaffaroni model w/ original RS1 (dynamical MFV) Rattazzi & Zaffaroni (00)

- Basic idea: both the bulk and the IR are flavor conserving, while the UV is flavor violating: only two marginal, flavor violating operators (SM Yukawa couplings): \leftrightarrow (5D picture) bulk scalars:

gauging flavor symmetries
in the bulk of ED
 \Rightarrow Flavor gauge boson!



- The scalar VEVs give rise to the Yukawa couplings on the IR brane:

$$\phi_u = (3, \bar{3}, 1) \text{ and } \phi_d = (\bar{3}, 1, 3) \text{ of } SU(3)_Q \times SU(3)_u \times SU(3)_d$$

$$\bar{Q}H\phi_u u + (u \leftrightarrow d)$$

$$\phi_u \sim Y_u, \quad \phi_d \sim Y_d$$

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- Basic idea:
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violating
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=> Flavor

Simplest setup w/ SM on IR brane
(all of the SM being composite)

No explanation
for the absence of EWP corrections
(why S-parameter is small)
& little hierarchy

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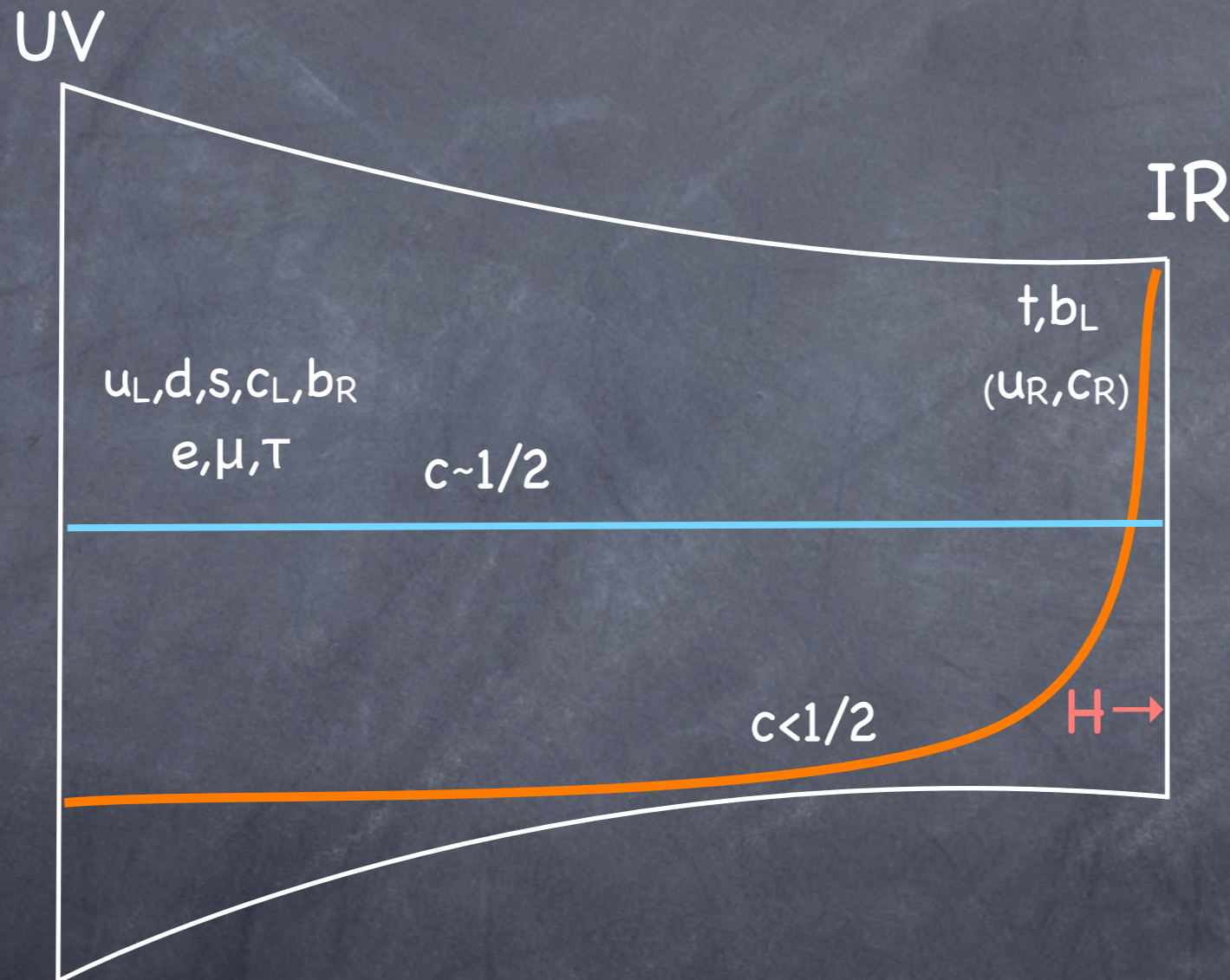
Flavor Triviality

Delaunay, Gedalia, SJJ,
Perez & Ponton (10,11)

- giving-up on flavor puzzle solution:

Generalization of RZ model w/ SM into 5D bulk

(anomaly mediation, gauge mediation)



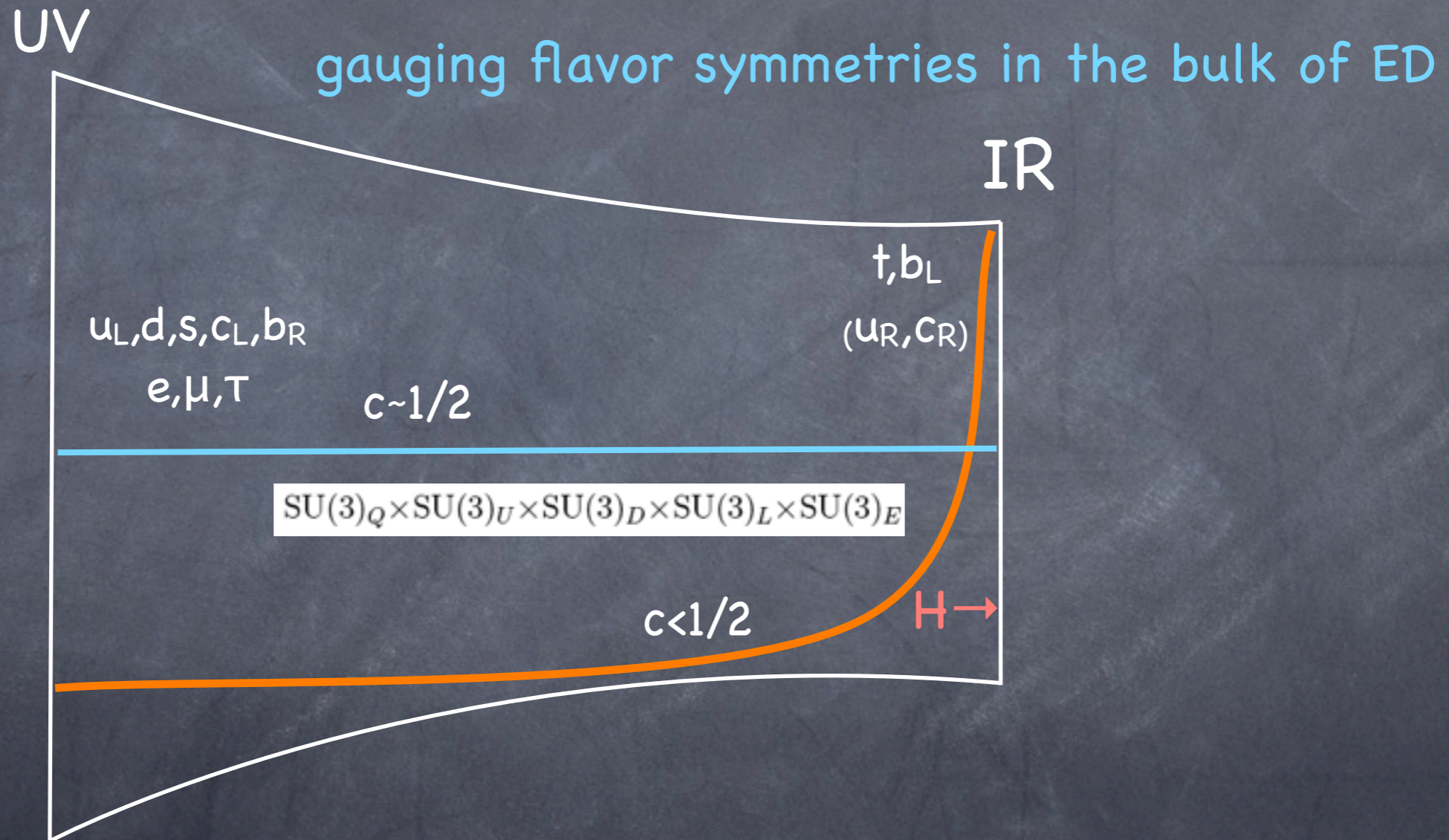
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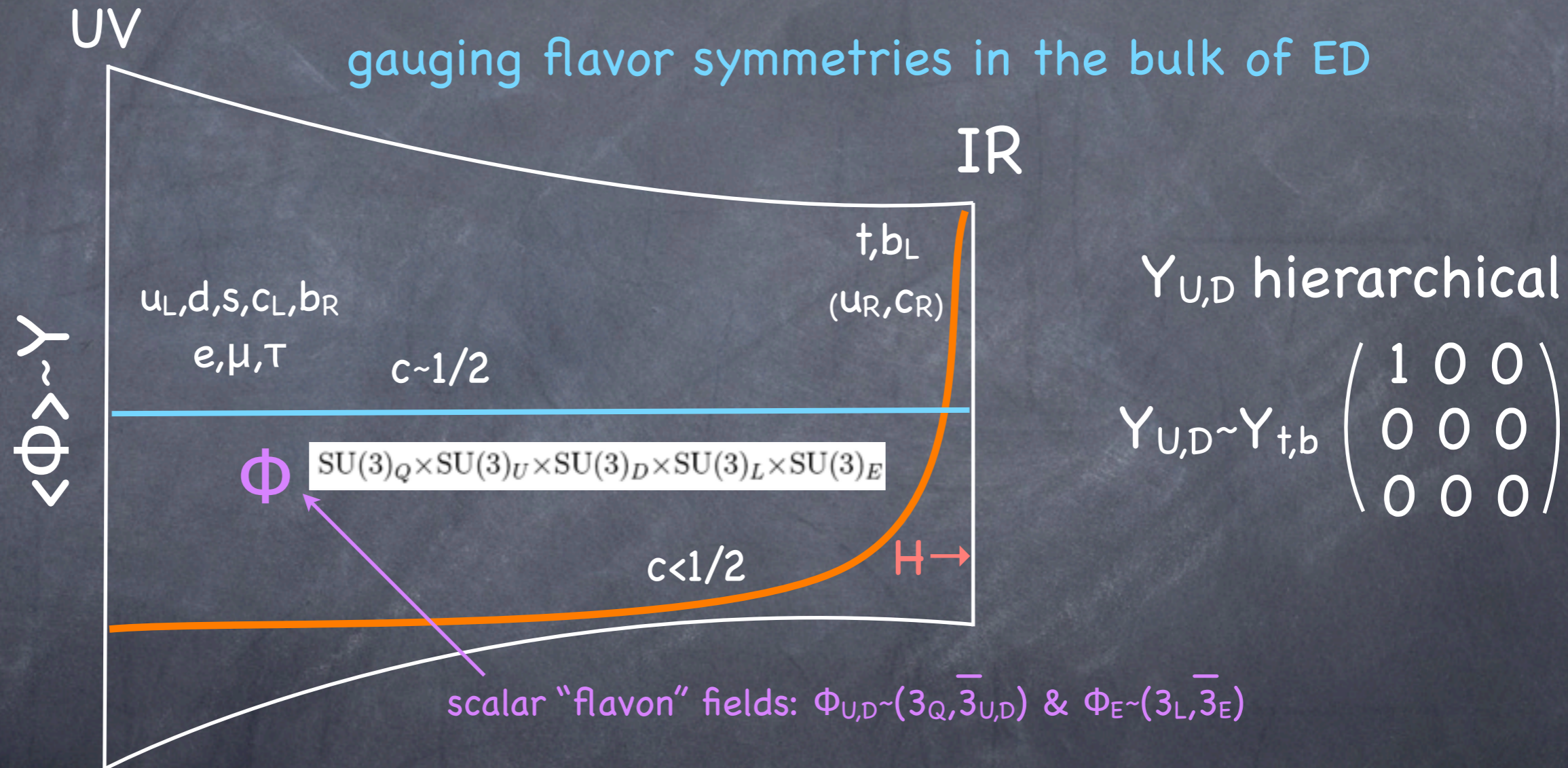
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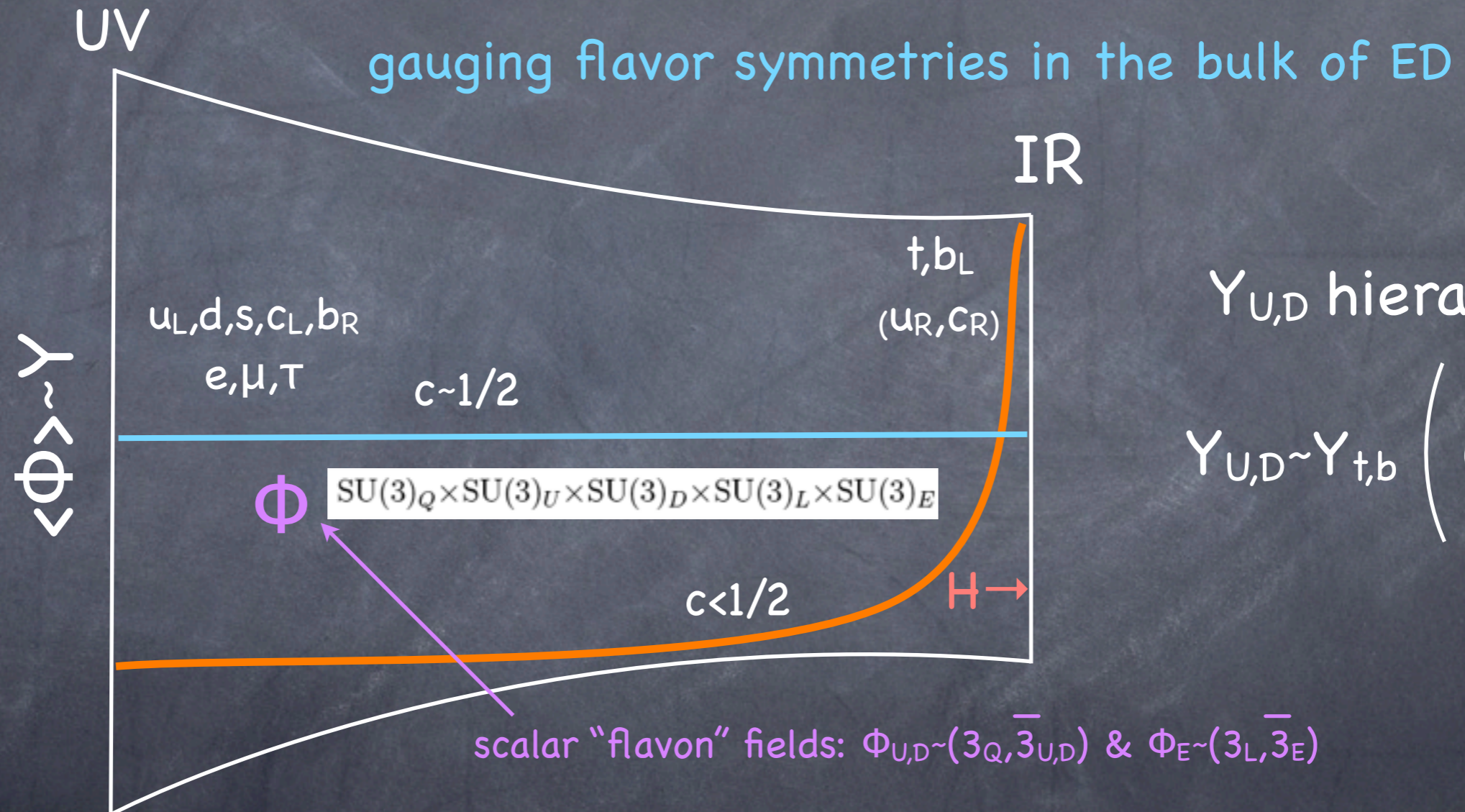
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- SM flavor puzzle solution postponed to the UV:
 $G_F \rightarrow$ approx. $SU(2)_Q \times SU(2)_U \times SU(2)_D \times SU(3)_L \times SU(3)_E$
- hierarchical Y "shined" to the IR by the Φ 's

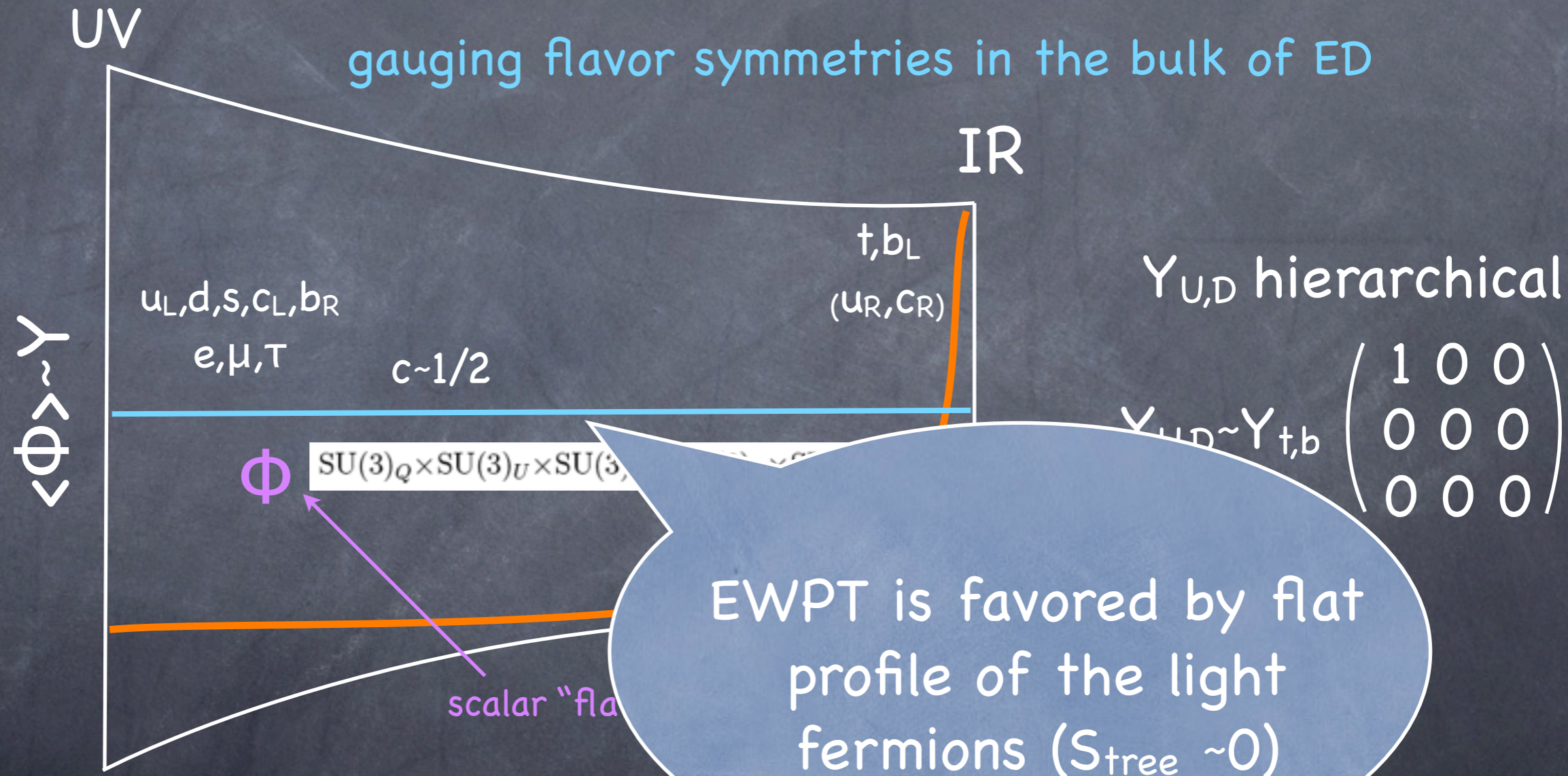
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Electroweak Precision Test

• $O_{\text{LEP}} \approx \langle O \rangle_{\text{SM}} (1 + \delta_{\text{NP}}) \rightarrow |\delta_{\text{NP}}| < 10^{-3}$ (precision tests!)

• Naturalness: $\Lambda_{\text{NP}} \sim \text{TeV} \rightarrow \text{EFT}$

• 20 (D=6) operators "measured" @LEP1+2:

(i) gauge prop. : $B_{\mu\nu} H^\dagger W_{\mu\nu} H, |H^\dagger D_\mu H|$ ($\rightarrow S, T$)

(ii) vertex corr. : $(\bar{f} \gamma^\mu f) H^\dagger D_\mu H, (\bar{f} \sigma^a \gamma^\mu f) H^\dagger \sigma^a D_\mu H$ ($Z \rightarrow ff$)

(iii) 4ψ : $(\bar{f}_1 \gamma^\mu f_1) (\bar{f}_2 \gamma^\mu f_2), (\bar{f}_1 \sigma^a \gamma^\mu f_1) (\bar{f}_2 \sigma^a \gamma_\mu f_2)$ ($ee \rightarrow ff$)

(+3rd gen. quark effects: $U(3)_3 \rightarrow U(2)_3 \times U(1)_3$)

Electroweak Precision Test

EWPT fitting procedure:

i) matching UV physics to L_{eff}

ii) constraints fundamental parameters ($\Lambda_{\text{NP}}, \dots$) with χ^2 -like analysis

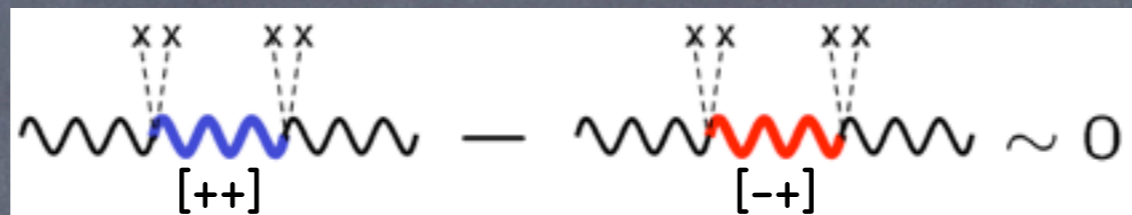
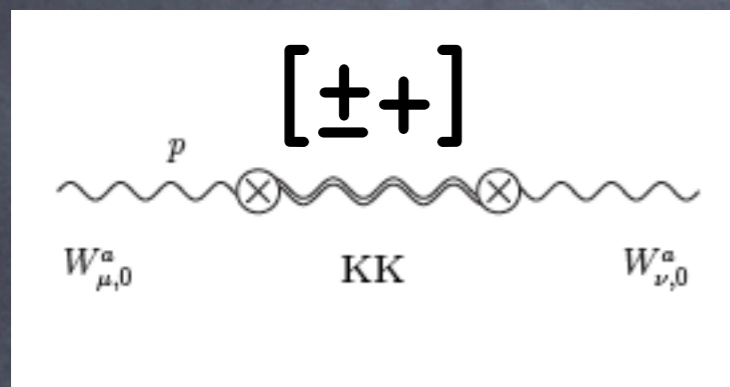
$$\chi^2 = \sum_{ij} (\langle O \rangle - O_{\text{LEP}})_i (\sigma^2)^{-1}_{ij} (\langle O \rangle - O_{\text{LEP}})_j$$



Tree level EWPOs

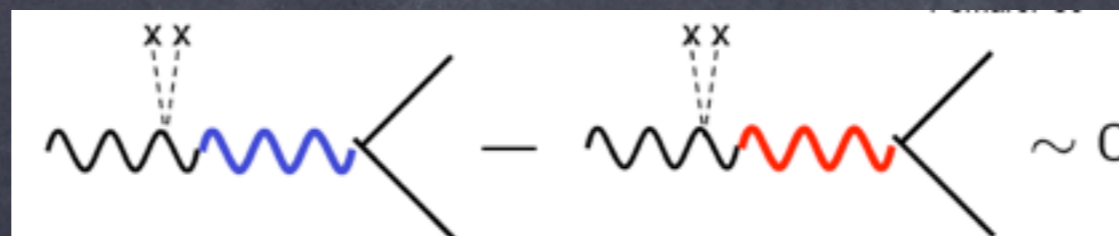
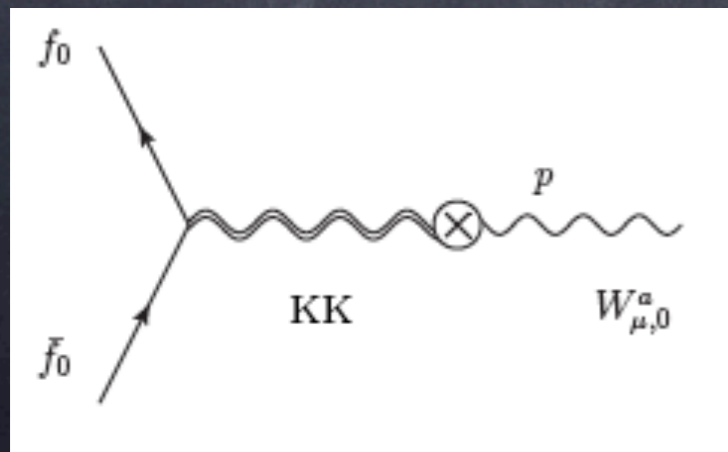
Agashe, Delgado, May, Sundrum

- Custodial symmetry: Avoid large corrections to the M_z/M_w ratio (T parameter): $SU(2)_L \times SU(2)_R$



Agashe, Contino, DaRold, Pomarol

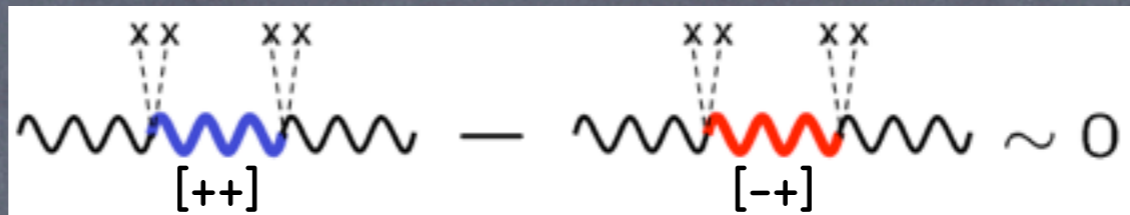
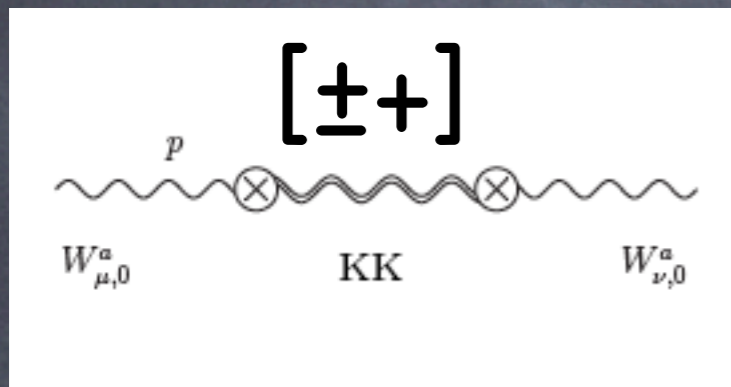
- To avoid Large corrections to Zbb coupling, δg_{bL} : impose discrete LR symmetry with $SU(2)_L \times SU(2)_R$



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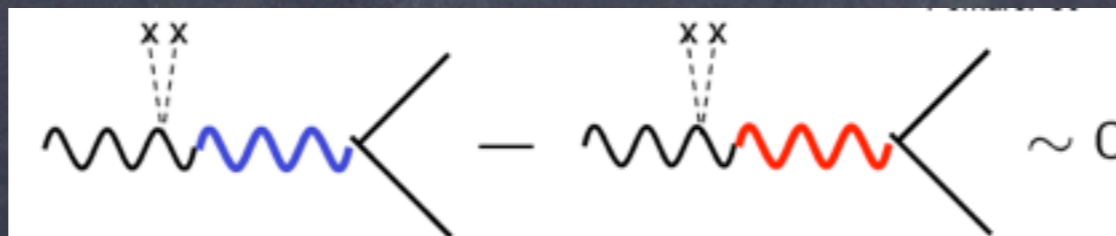
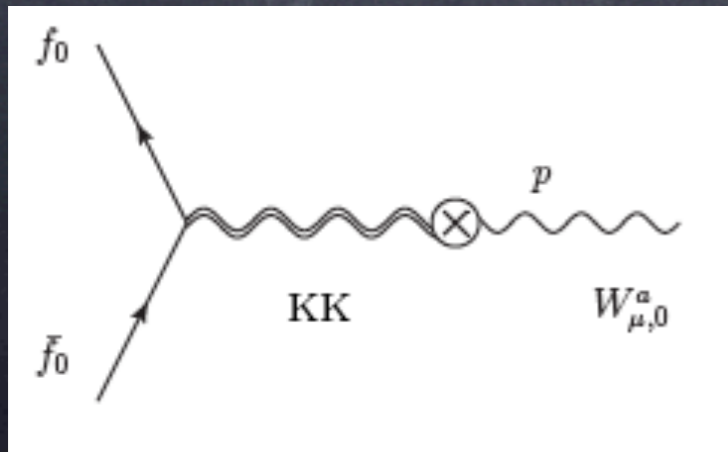


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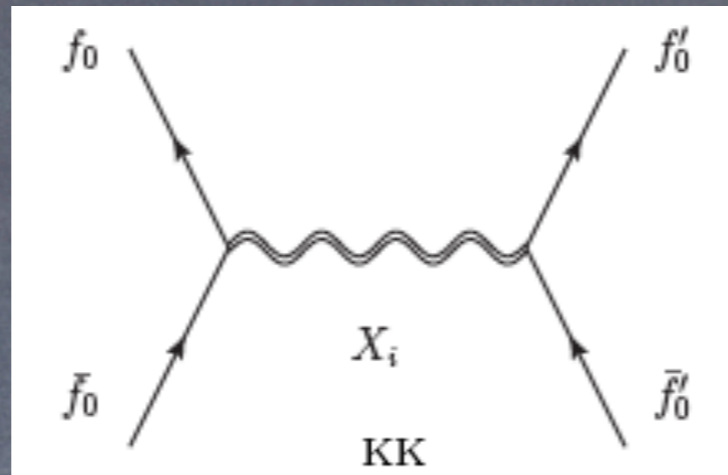
symmetry works for both anarchic and flavor trivial models

$SU(2)_L$: impose $SU(2)_R$



Tree level EWPOs

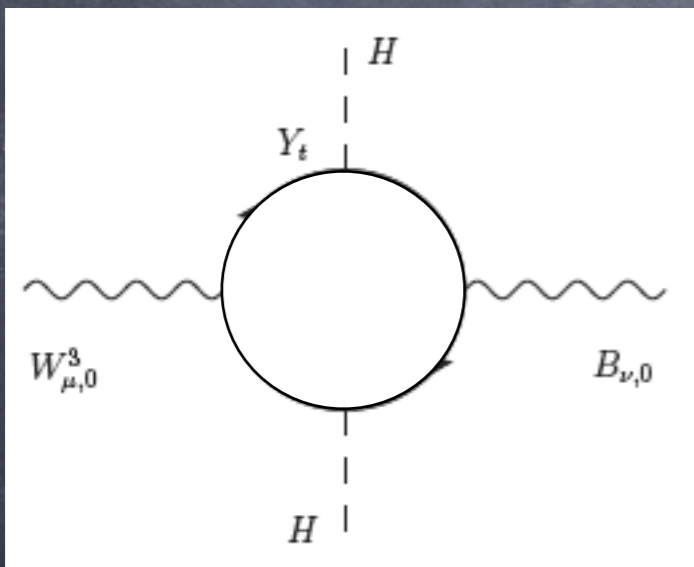
- 4-fermion operators involving light fermions



- Anarchic RS: negligible due to small wave function overlap between KK GB and light elementary fermions
=> fit for "oblique parameters" (S,T,U) + δg_{bL}
- FT RS: 4-fermion operators become important (large non-oblique corrections) => one must look at more than $O(35)$ EWPOs in order to assess whether EWPTs are passed

Radiative corrections to EWPOs

- Loop corrections are important:
 - T & δg_{bL} are finite to all orders (custodial $SU(2)_V \times P_{LR}$)
 - S is finite at 1-loop in 5D for bulk Higgs



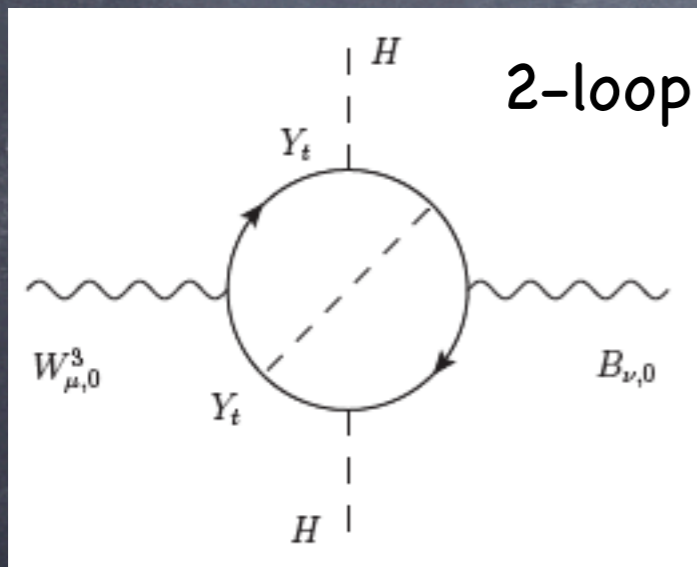
5D NDA

$$g_5^2 Y_5^2 B_{\mu\nu} W^{\mu\nu} H^2 \sim \text{finite}$$

$$(-1) + (-1) + 5/2 + 5/2 + 3 = 6$$

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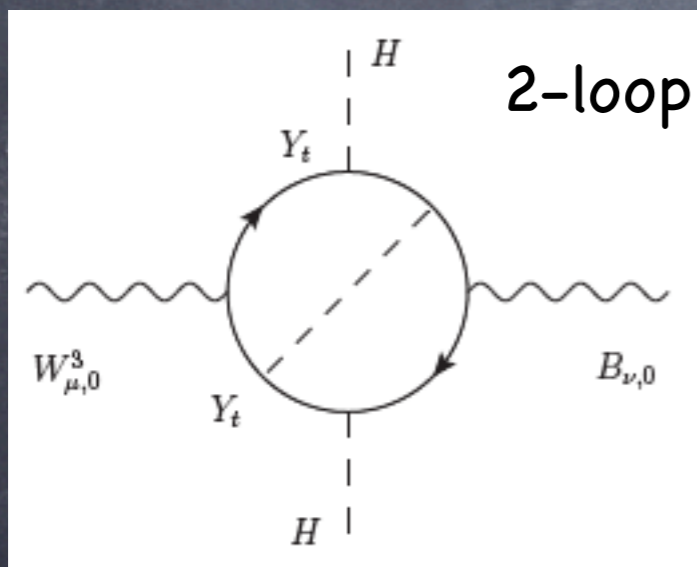
5D NDA

$$g_5^2 Y_5^4 B_{\mu\nu} W^{\mu\nu} H^2 \sim \text{Log} \Lambda_5$$

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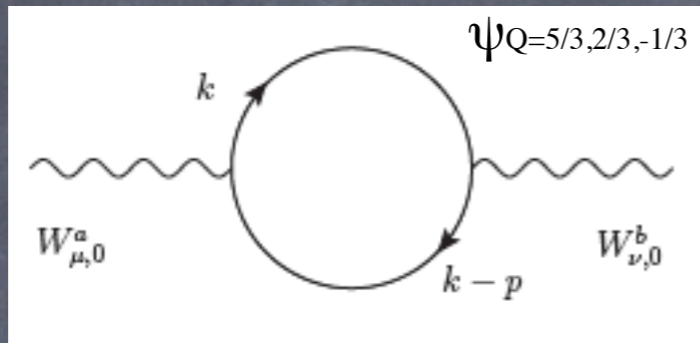
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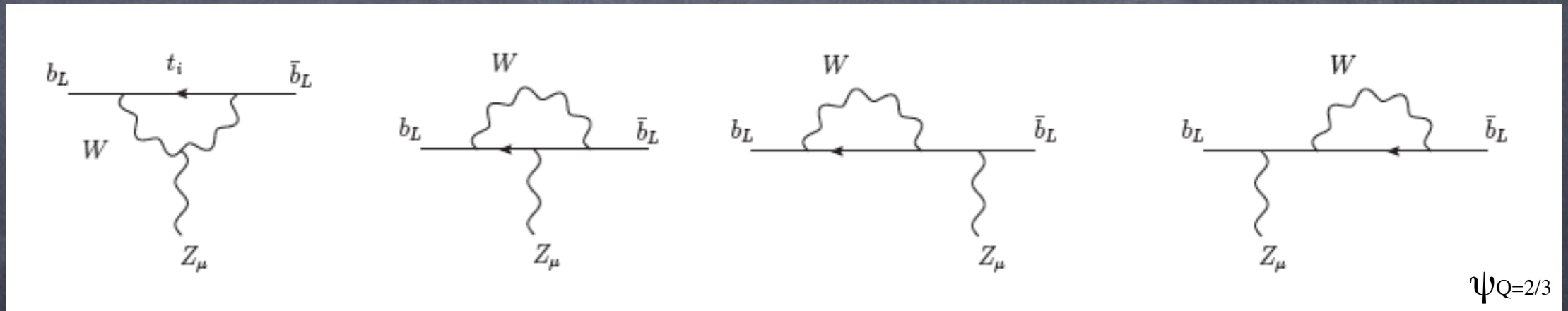
- 2-loop corrections are suppressed for perturbative Y_t and Y_b (less than 20%)
- Higher-loop corrections are more suppressed by NDA expansion $(Y^2 \Lambda_5 / 16\pi^2) \ll 1$

Radiative corrections to EWPOs

(i) S,T



(ii) $Z \rightarrow b_L b_L$



χ^2 global fit results

- we (re)did the EWPT fit for anarchic RS (most precise to date): @ 2σ : $m_{KK} > 4.6$ (3.9) TeV for 1(3) d.o.f.

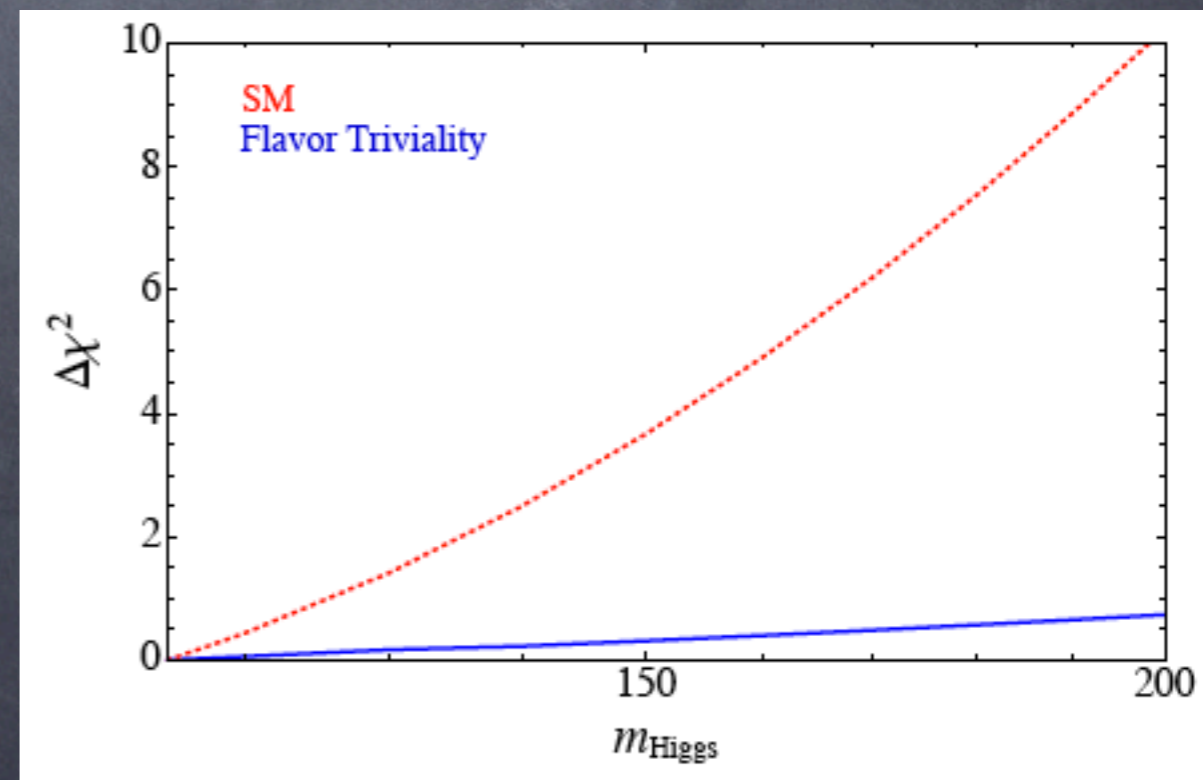
is the celebrated anarchic RS model visible @ LHC?

- Flavor Triviality fit: m_{KK} , C_t , C_b , C_{Q3} , C_{Qi} , C_{ui} , C_{di} , C_L , C_e : @ 2σ : $m_{KK} > 2.1$ (1.7) TeV for 1(6 important d.o.f.s) d.o.f.

$(\chi^2/\text{dof})_{SM} = 219.9/232 \sim 0.95 < (\chi^2/\text{dof})_{FT RS} = 217.3/223 \sim 0.97$

- We find FTRS provides a better fit than SM (with $m_{KK} = 3.9$ TeV)

- We find fit depends only mildly on the Higgs mass \Rightarrow large Higgs mass values are still compatible with the model, without spoiling the EW fit



χ^2 global fit results

- slight preference of u_R being composite (localize towards IR brane)

- w/ sweet-spot parameters:

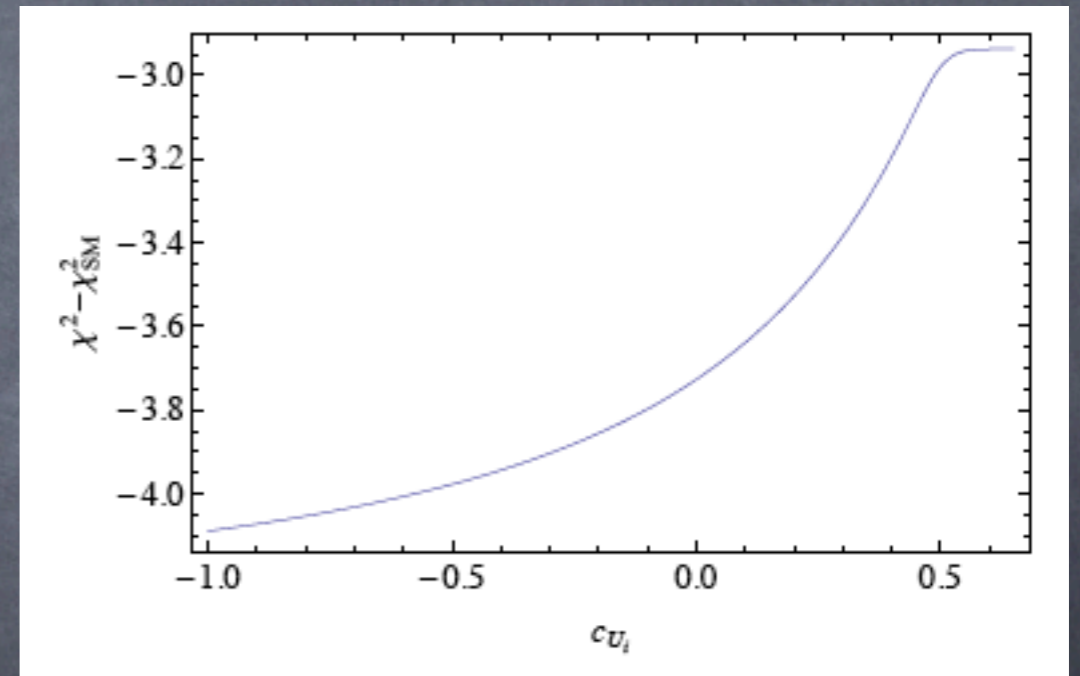
- $CQ \approx (0.50, 0.50, 0.02)$

- $CD \approx (0.63, 0.63, 0.57)$

- $CU \approx (0.15, 0.15, 0.48)$

- $\alpha_U Y_U \approx (4.3 \times 10^{-5}, 0.021, 4.2)$

- $\alpha_D Y_D \approx (0.01, 0.19, 0.45)$.



with $\alpha_{U,D} \approx 4, 0.12,$

$$Y_b^2 \gg Y_t^2$$

down alignment: $[m_D, Y_D] \approx 0$

Flavor physics is decoupled!

χ^2 global fit results

slight
toward

composite $u_R \Rightarrow$ exciting
collider phenomenology,
with no flavor constraint

composite (localize

w/ sweet-spot parameters:

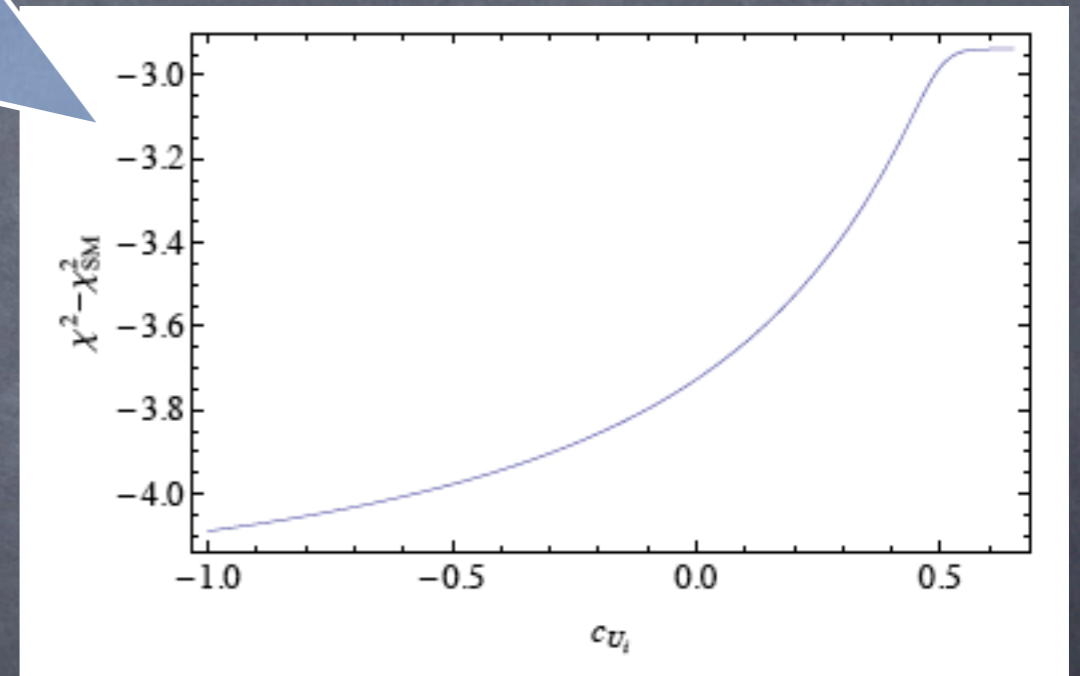
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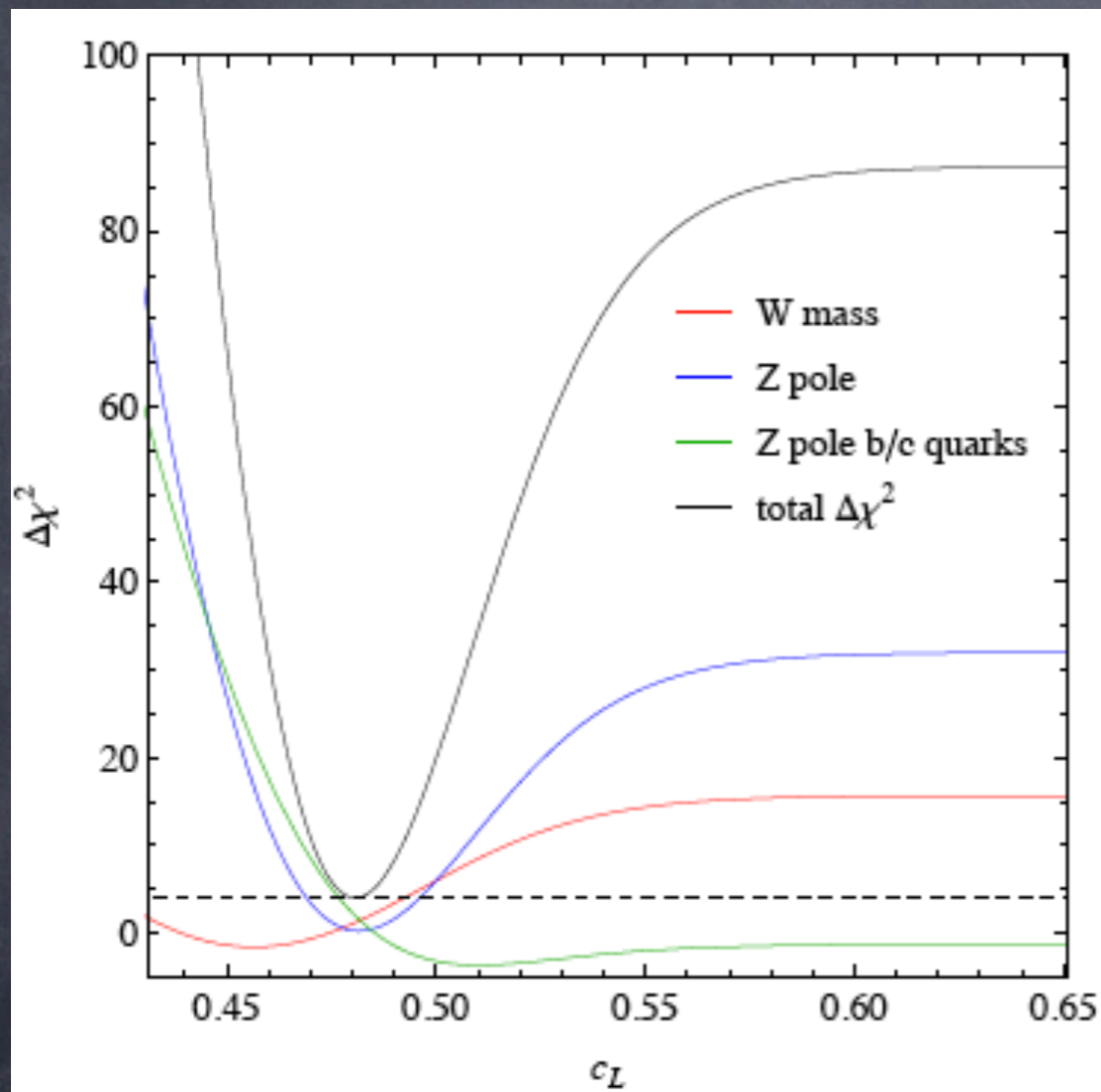
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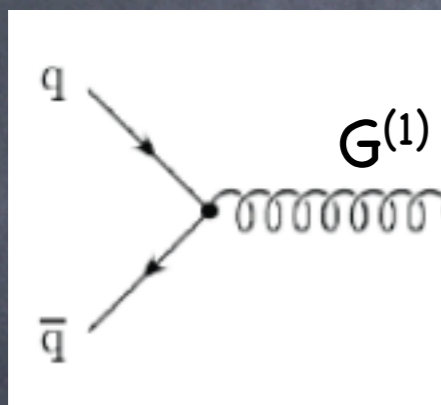
	$\chi^2 - \chi_{\text{SM}}^2$	$\chi_{\text{min}}^2 - \chi_{\text{SM}}^2$	$\Delta\chi^2$
W mass	1.37	0.12	1.25
Z line shape & lepton A_{FB}	-2.47	-2.74	0.27
Z pole b&c quarks	5.30	3.45	1.85
Leptonic polarization asymmetries	0.20	0.23	-0.04
s_W^2 hadronic charge asymmetry	-1.95	-2.18	0.23
Deep Inelastic Scattering	-0.13	-0.12	$< 10^{-2}$
Atomic Parity Violation	3.23	0.11	3.12
LEP2 hadronic cross-section	-1.97	-0.97	-1.00
LEP2 muon pair	$< 10^{-2}$	0.03	-0.03
LEP2 tau pair	-0.04	-0.03	-0.01
OPAL electron pair	-0.02	-0.02	$< 10^{-2}$
L3 W pair	-0.17	-0.11	-0.06
Z pole s quark	0.07	0.09	-0.02
LEP2 $ee \rightarrow bb$	-3.22	-1.75	-1.48
LEP2 $ee \rightarrow cc$	-0.18	-0.08	-0.10
total	0.02	-3.98	4

Interplay between Flavor Physics & Collider (top) Physics



FT RS is ultra visible

- generic RS signature: KK gluon production (& FGB for FT model)



$$= g^* \left(\frac{1}{k\pi r_c} - f_{c_q}^2 \right)$$

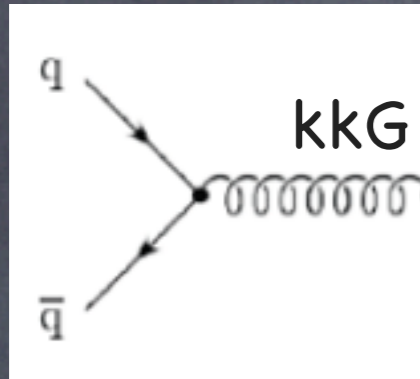
$$f_c^2 \approx \begin{cases} 1-2c & (c < 1/2) \\ 1/k\pi r_c & (c \approx 1/2) \\ (2c-1)\exp[-(2c-1)k\pi r_c] & (c > 1/2) \end{cases}$$

suppressed in anarchic RS
(only through "γ-ρ mixing")

O(1) × g* for FT RS

- 7(14)TeV LHC : $\sigma_{KK} \approx 0.04(6)\text{fb}$ in anarchic RS w/ $m_{KK}=4.6\text{TeV}$
- 7(14)TeV LHC : $\sigma_{KK} \approx 3(500)\text{pb}$ in FT RS w/ $m_{KK}=2.1\text{TeV}$ (mainly due to composite u_R , and also due to low m_{KK})

RS KK_G and forward-backward asymmetry



$$g^* \gamma^\mu (v_q + \gamma_5 a_q) \begin{cases} a_{u,d} \approx 0 \text{ in anarchic RS} \\ a_u \approx 0(1) \text{ in FT RS} \end{cases}$$

- Anarchic models \Rightarrow tiny effect (vectorial & suppressed production couplings).

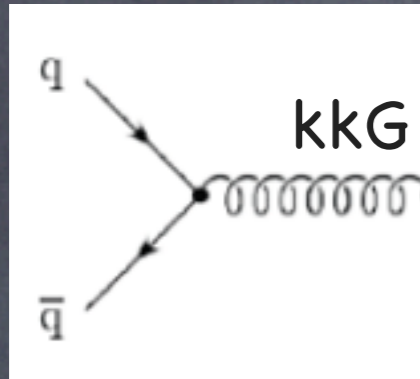
Agashe, Belyaev, Krupovnickas, Peres & Virzi (06);
Bauer, Goertz, Haisch, Pfoh & Westhoff (10).

$$\sigma^{\text{NP}} / \sigma^{\text{SM}}(p_T > 400 \text{ GeV}): 2 \sim 3$$

$$v_q \equiv -\frac{1}{\xi} + \frac{1}{2}(f_{qL}^2 + f_{qR}^2), \quad a_q \equiv \frac{1}{2}(f_{qL}^2 - f_{qR}^2)$$

\downarrow
 $k\pi R \sim \log(M_{\text{Pl}}/\text{TeV})$

RS KKG and forward-backward asymmetry

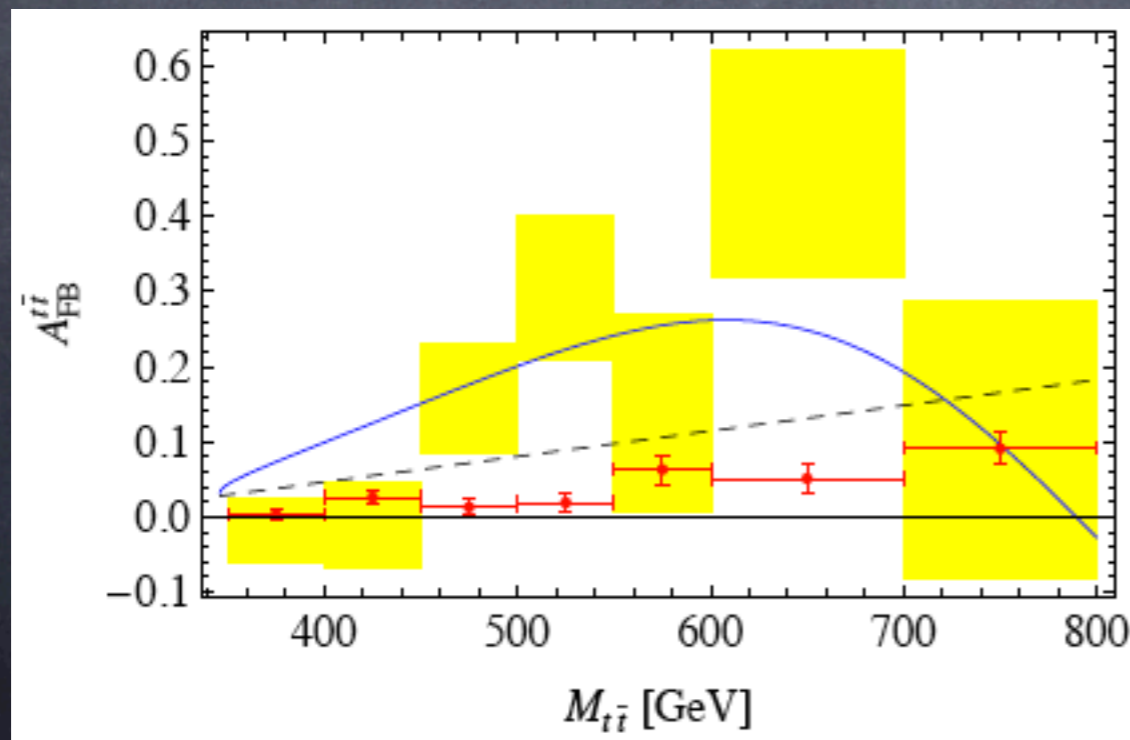


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- Anarchic models => tiny effect (vectorial & suppressed production couplings).

Agashe, Belyaev, Krupovnickas, Peres & Virzi (06);
Bauer, Goertz, Haisch, Pfoh & Westhoff (10).

- FT models: [Delaunay, Gedalia, SJL, Perez & Ponton \(11\)](#)



$$\hat{A} \propto \beta_t \hat{s} |\mathcal{D}|^2 a_t a_q g_{s*}^2 [g_s^2 (\hat{s} - m_{\text{KK}}^2) + 2g_{s*}^2 \hat{s} v_t v_q]$$

$\beta_t \equiv \sqrt{1 - 4m_t^2/\hat{s}}$

$$A^{t\bar{t}}(M_{t\bar{t}} > 450 \text{ GeV}) = 0.19 \quad \sim 2\sigma$$

$$A^{t\bar{t}}(|\Delta y| > 1) = 0.24$$

$$A^{t\bar{t}} = 0.118$$

$$\sigma^{\text{NP}}/\sigma^{\text{SM}}(p_T > 400 \text{ GeV}): 2\sim 3$$

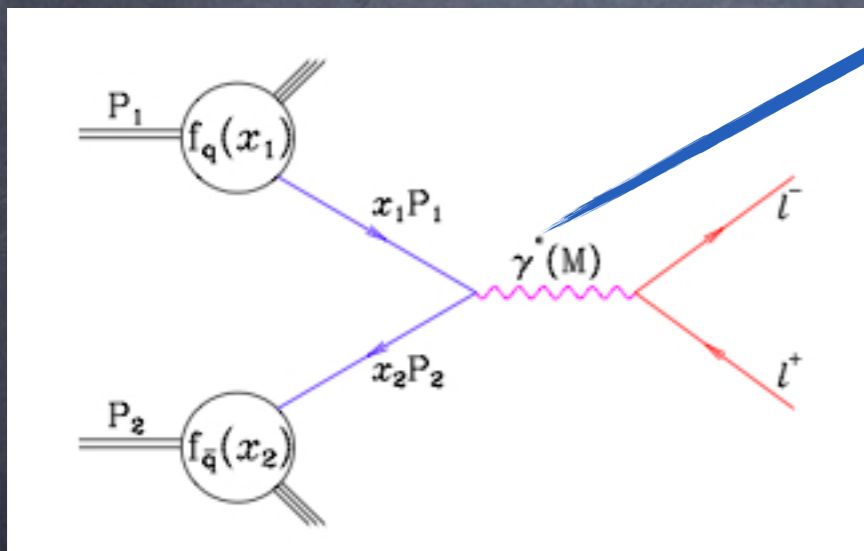
$$v_q \equiv -\frac{1}{\xi} + \frac{1}{2}(f_{qL}^2 + f_{qR}^2), \quad a_q \equiv \frac{1}{2}(f_{qL}^2 - f_{qR}^2)$$

$k\pi R \sim \log(M_{\text{Pl}}/\text{TeV})$

Flavor Gauge Bosons @ LHC

Csáki, Kagan, SJL, Perez & Weiler (in preparation)

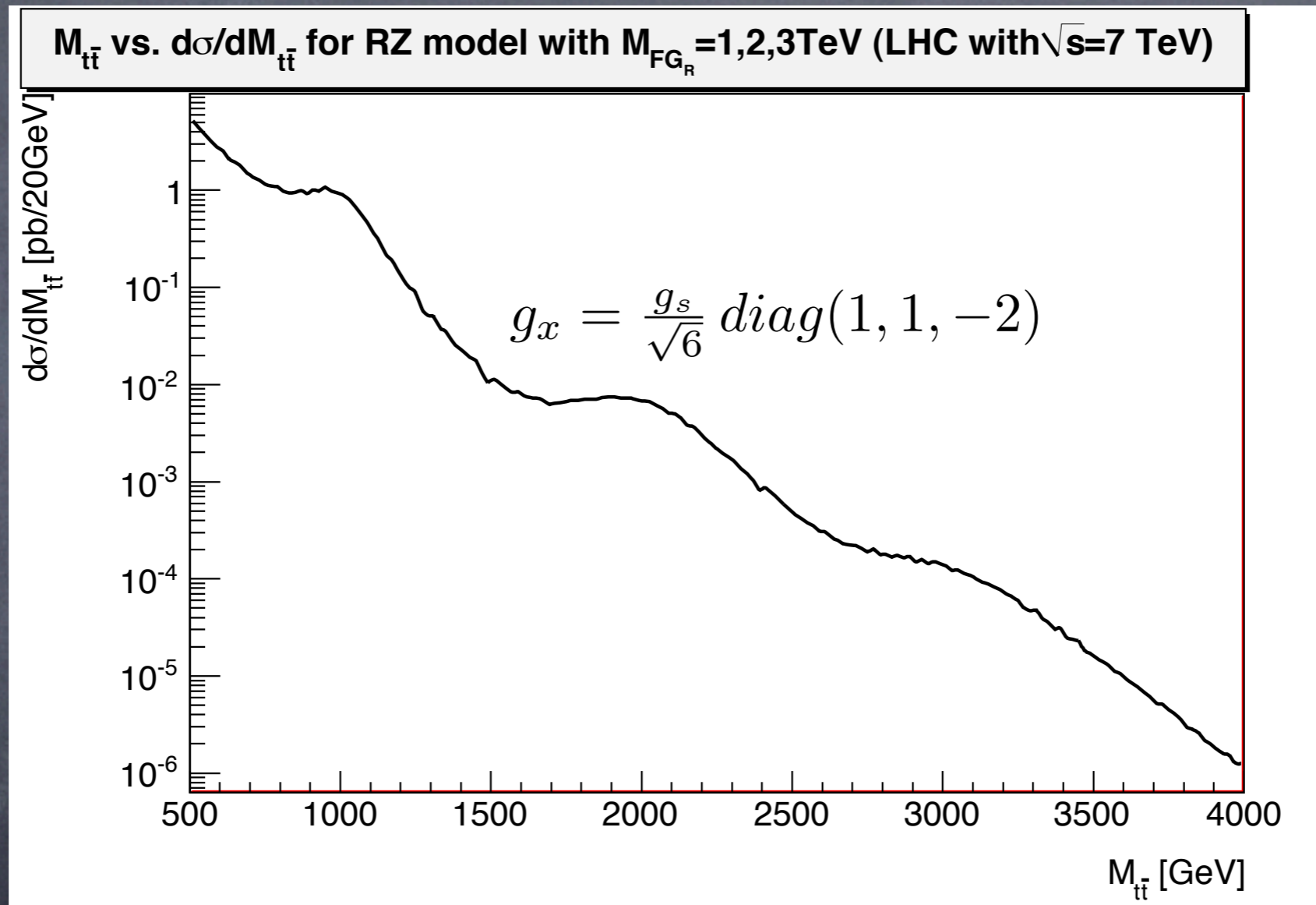
$$g_{eff} G_{\mu}^{(1)KK} \bar{\psi} \psi$$



- Flavor gauge bosons do not have massless modes (flavor is broken)
- no $\gamma - \rho$ mixing !
- But quark composite mixing can be flavor universal & large
- with BKT, it is possible to have $m_{FGB} < M_{KK}$

FGBs at the LHC (preliminary)

- The flavor gauge bosons & scalars might be observable.



$$\mathcal{L} = \frac{\lambda}{2} V_\mu^8 \left(\frac{1}{\sqrt{3}} \bar{u}_R \gamma^\mu u_R + \frac{1}{\sqrt{3}} \bar{c}_R \gamma^\mu c_R - \frac{2}{\sqrt{3}} \bar{t}_R \gamma^\mu t_R \right) + \frac{\lambda}{2} \left((V_\mu^4 - iV_\mu^5) \bar{u}_R \gamma^\mu t_R + (V_\mu^6 - iV_\mu^7) \bar{c}_R \gamma^\mu t_R + h.c. \right)$$

Flavor Gauge Boson @ Tevatron?

$$\mathcal{L} = g_{eff} \bar{u}_R V_\mu^A \frac{T^A}{2} \gamma_\mu u_R + h.c.$$

• Can partially explain A_{FB} with the usual constraints:

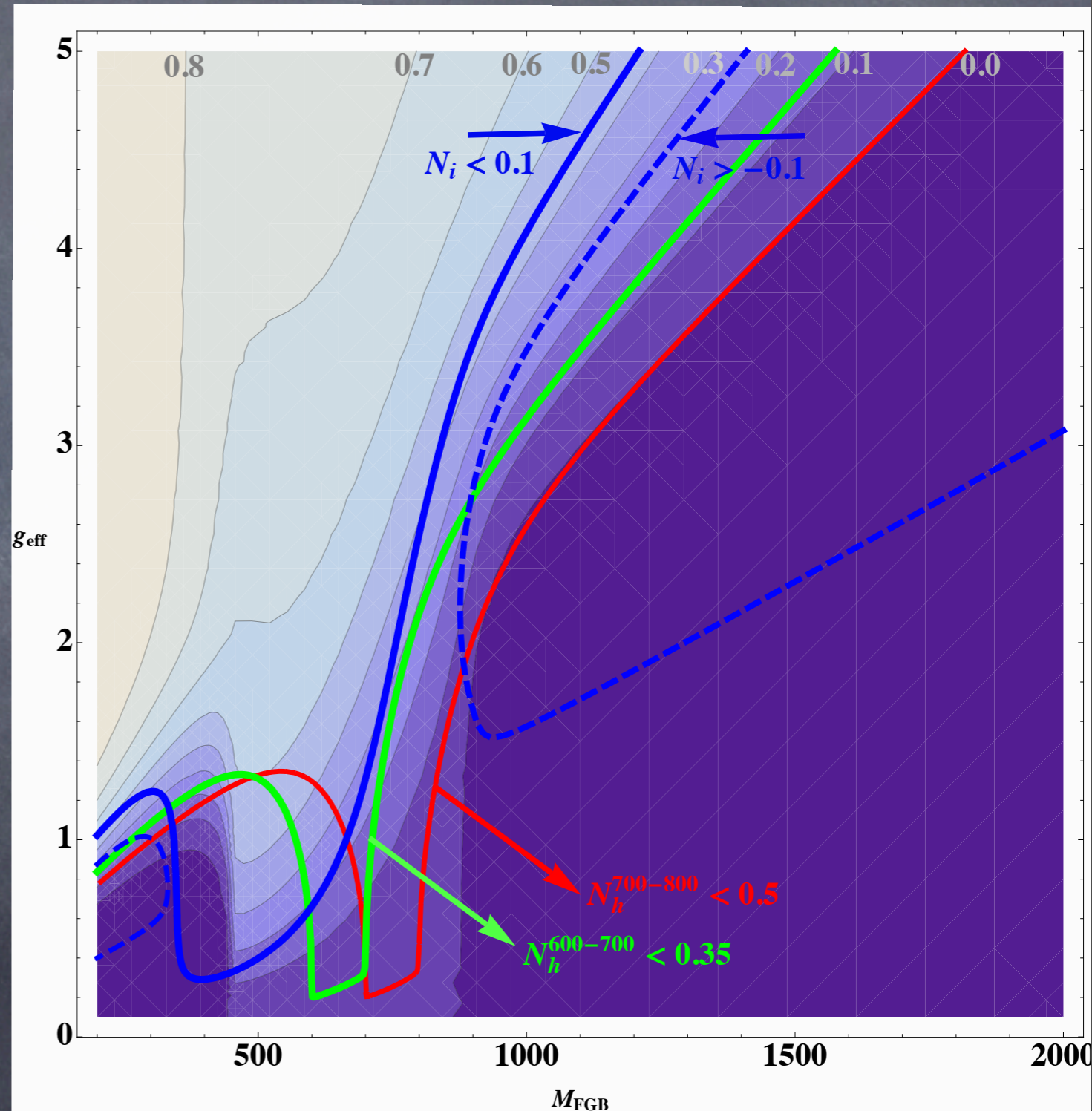
i) $\delta\sigma_{700-800\text{ GeV}}^{\text{NP}} / \sigma_{700-800\text{ GeV}}^{\text{SM}} \lesssim 47\%$

ii) $\delta\sigma_{t\bar{t}}^{\text{NP}} / \sigma_{t\bar{t}}^{\text{SM}} \lesssim 10\%$

• $M_{\text{FGB}} < 900\text{ GeV}$, $g_{\text{eff}} \sim O(1)$

$$A_{FB}^{t\bar{t}}(M_{\text{inv}} > 450\text{ GeV}) \lesssim 10\%$$

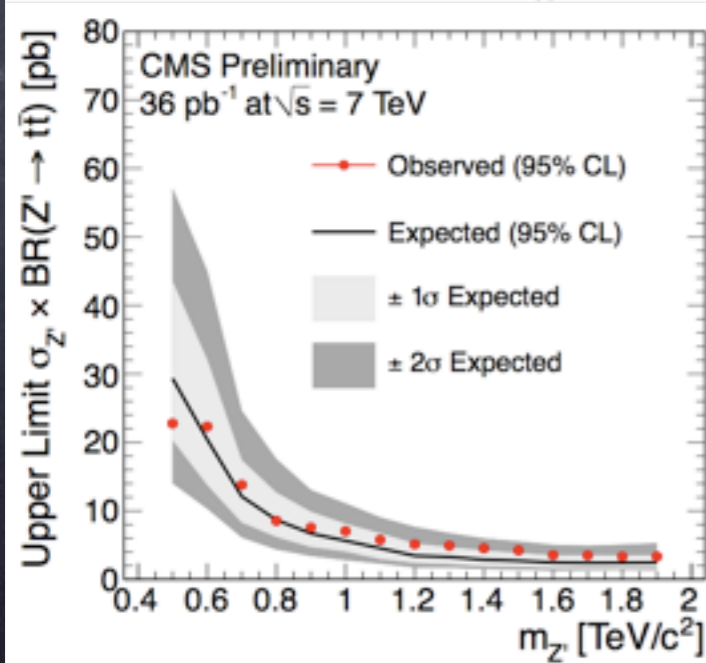
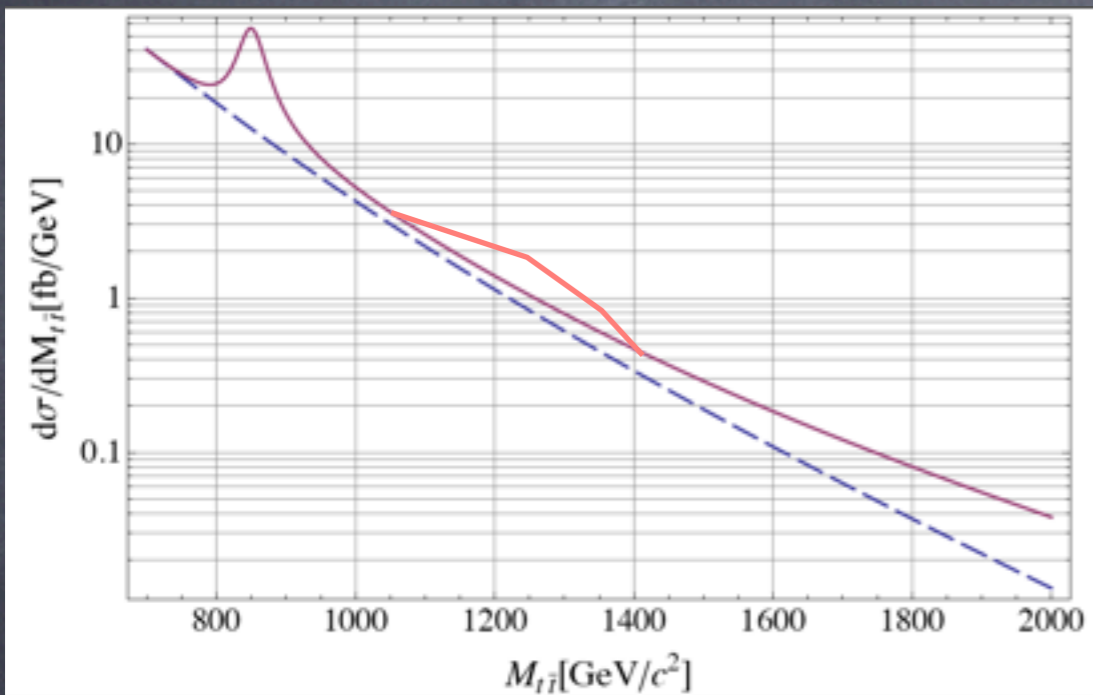
• $\sigma_{\text{NP}} / \sigma_{\text{SM}}(p_T > 400\text{ GeV})$: 2-3



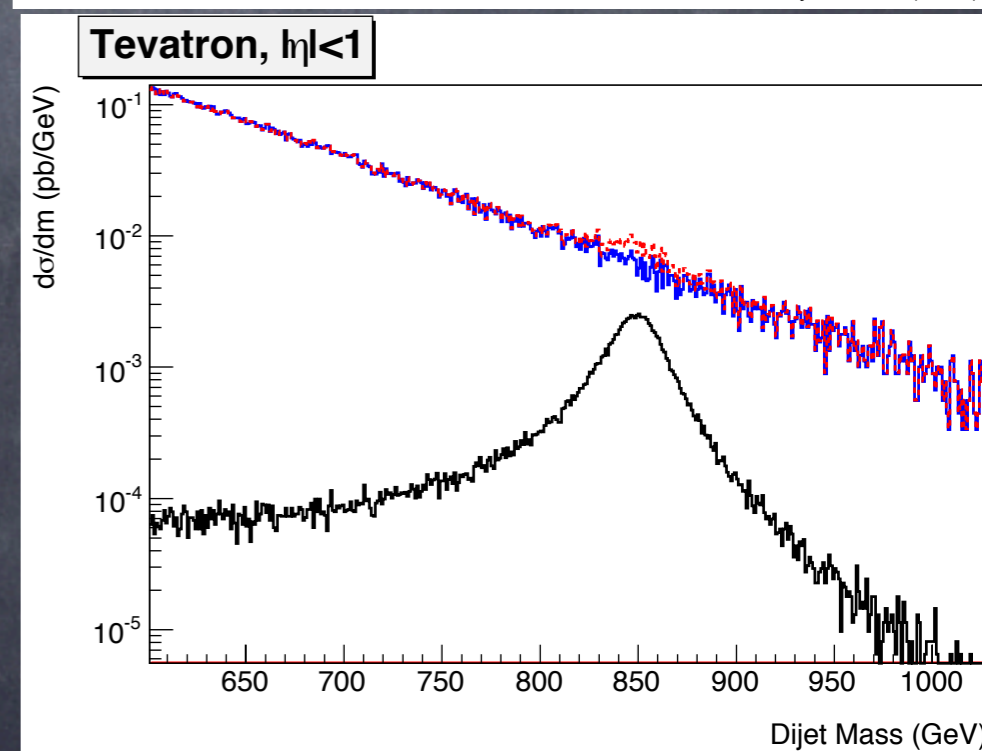
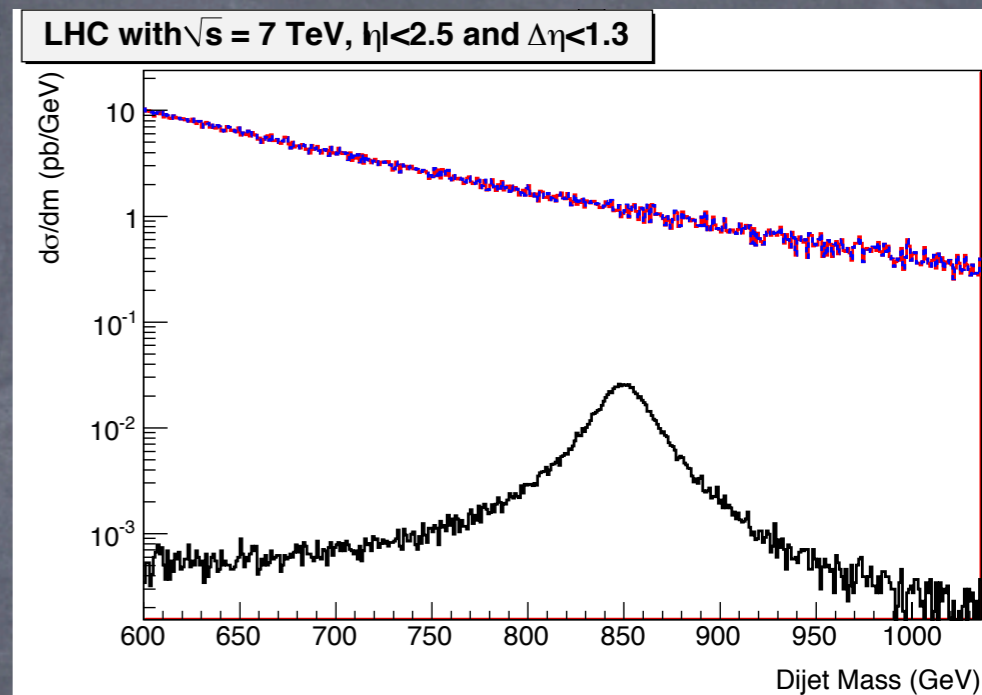
Signals and Constraint for A_{FB} sweet spot (Preliminary)

$$M_{FGB} < 900 \text{ GeV}, g_{eff} \sim O(1)$$

$t\bar{t}$



Dijet



Summary

- Warped XD provide an effective theory frame work for probing strong dynamics via AdS/CFT @ LHC
- Little CP Problem \Rightarrow flavor symmetry for NP scale relevant for EWSB and LHC \Rightarrow presence of flavor gauge boson
- Flavor trivial Warped Extra Dimension model can be tested @ LHC by 2012 (unlike popular anarchic RS, which we shows that $m_{KK} > 4.6$ TeV); it can't explain Tevatron hints entirely, but improvement is in the right direction
- The flavor gauge boson production cross section can be sizable and lead to observable effects at the Tevatron and at the LHC
- Top physics might be crucial for NP discovery @ LHC

hvala

Backup Slide

$$\begin{array}{c}
 Q(2,2)_{2/3} \\
 \leftarrow \text{SU}(2)_R \rightarrow \\
 \text{SU}(2)_L \updownarrow \left(\begin{array}{cc} X_L(-,+) & t_L(+,+) \\ T_L(-,+) & b_L(+,+) \end{array} \right)_X \sim (2,2)_{2/3} \sim \begin{pmatrix} 5/3 & 2/3 \\ 2/3 & -1/3 \end{pmatrix} \cdot \\
 \text{U}(1)_Q
 \end{array}$$

$$U(+,+) \sim (1,1)_{2/3}$$

$$D \sim (1,3)_{2/3} + (3,1)_{2/3}$$

$$L(+,+) \sim (2,2)_{-1}$$

$$E \sim (1,3)_0 + (3,1)_0$$

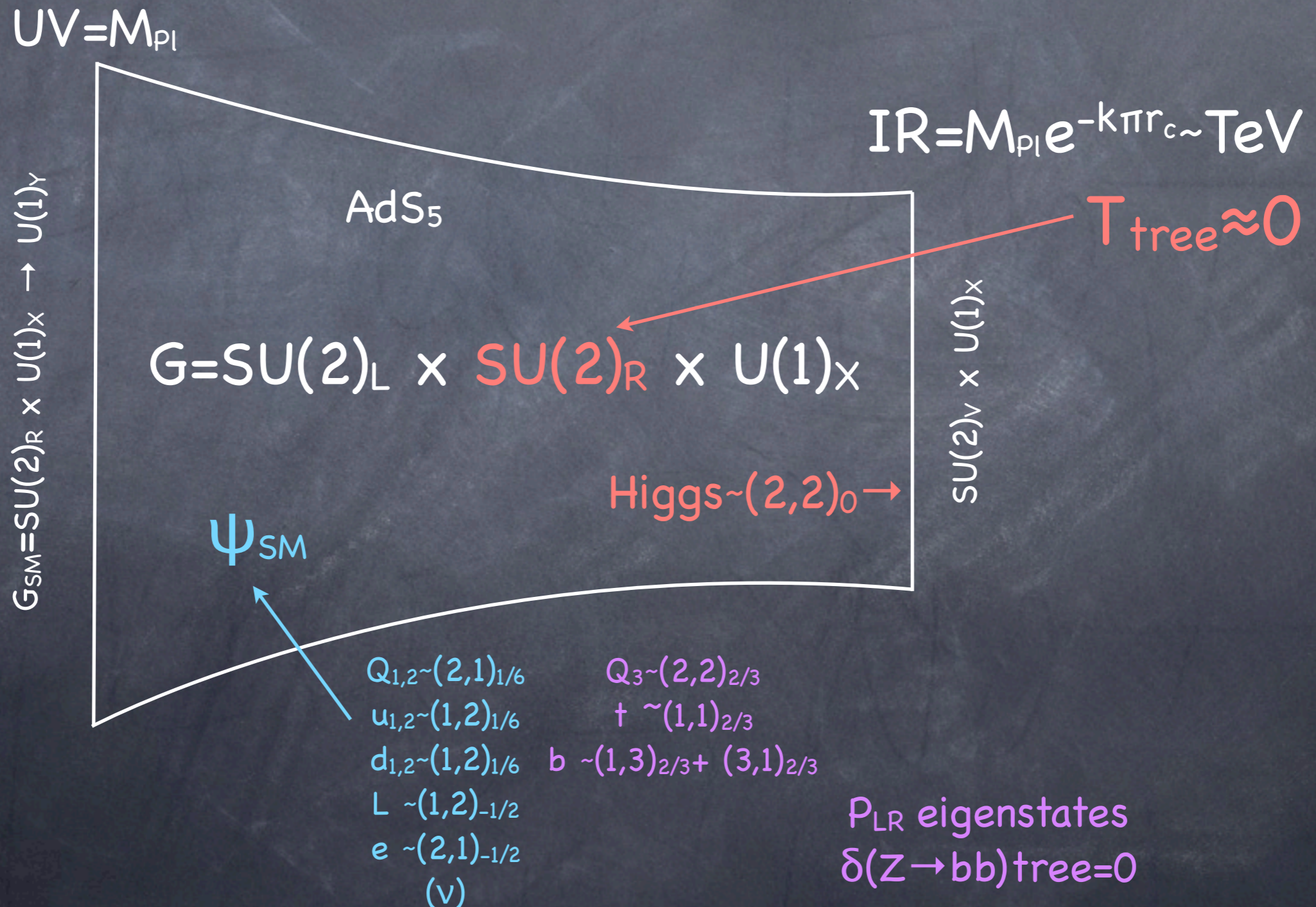
Realistic Warped Extra dimension (RS)

Randall, Sundrum

Agashe, Delgado, May, Sundrum

Agashe, Contino, DaRold, Pomarol

...



FCNC's from FGB

- When diagonalizing the Yukawa interaction, we can also diagonalize some of the GB interactions by redefining a basis of the gauge fields.

$$\phi_d = U_d \phi_d^{(diag)} V_d^\dagger$$

$$\phi_u = U_u \phi_u^{(diag)} V_u^\dagger$$

- However, there are four rotations needed, one has only three sets of flavor gauge bosons, thus the effect of one combination of these rotation matrices can not be eliminated.

$$A^u \rightarrow V_u A^u V_u^\dagger, \quad A^d \rightarrow V_d A^d V_d^\dagger, \quad A^Q \rightarrow U_d A^Q U_d^\dagger.$$

couplings of the u_R, d_R, d_L

$$g_* \bar{q}_i T_{ij}^a A_\mu^a \gamma^\mu q_j$$

couplings of u_L

$$g_* \bar{u}_L V A_\mu^Q V^\dagger \gamma^\mu u_L$$

- Similarly, the effect of the flavor rotations on the scalar VEVs will yield the following set of GB mass matrices

$$M_{KK}^2 \left(\alpha_u \text{Tr}[Y_u A_u A_u Y_u] + \beta_u \text{Tr}[A_Q V Y_u Y_u V^\dagger A_Q] + \gamma_u \text{Tr}[V^\dagger A_Q V Y_u A_u Y_u] \right. \\ \left. \alpha_d \text{Tr}[Y_d A_d A_d Y_d] + \beta_d \text{Tr}[A_Q Y_d Y_d A_Q] + \gamma_d \text{Tr}[A_Q Y_d A_d Y_d] \right)$$

diagonal real Yukawa couplings

FCNC's from FGB

- When diagonalizing the Yukawa interaction, we can also diagonalize some of the GB interactions by redefining a basis of the gauge fields.

$$\phi_d = U_d \phi_d^{(diag)} V_d^\dagger$$

$$\phi_u = U_u \phi_u^{(diag)} V_u^\dagger$$

- However, there are four rotations needed, one has only three sets of flavor gauge bosons, thus one of these rotation matrices can be absorbed into the Yukawa interaction

$$A^u \rightarrow V_u A^u V_u^\dagger, \quad A^d \rightarrow V_d A^d V_d^\dagger,$$

couplings of the u_R, d_R, d_L

$$g_* \bar{q}_i T_{ij}^a A_\mu^a \gamma^\mu q_j$$

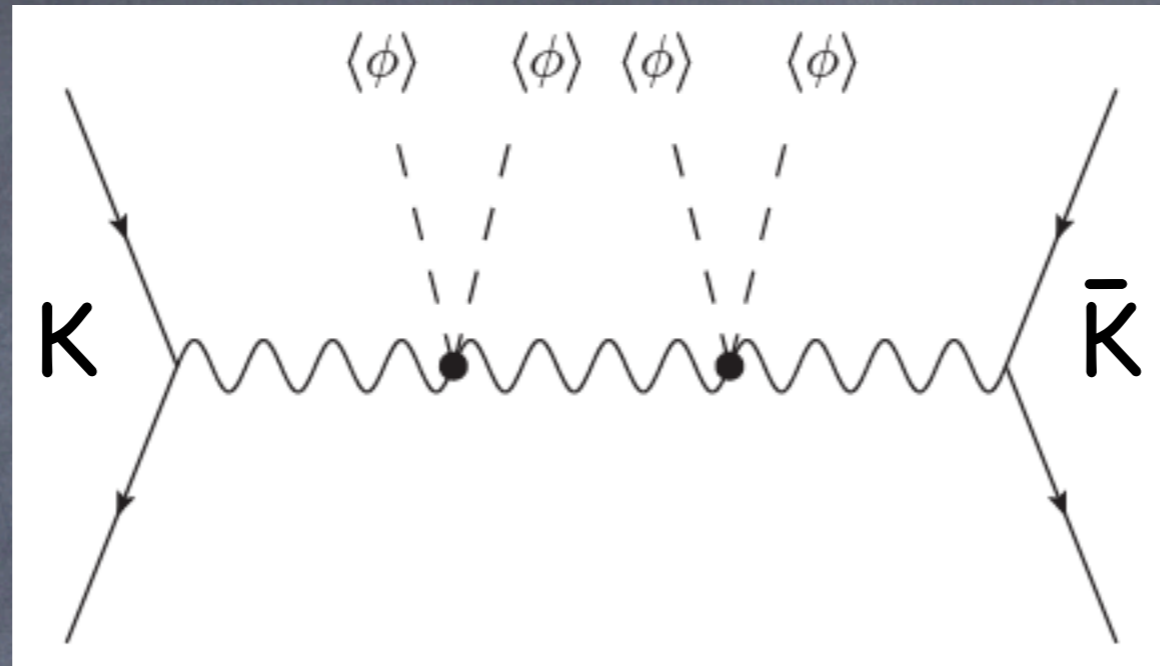
Clearly the only sources of flavor violation in this basis are explicitly proportional to V_{CKM}
 \Rightarrow MFV

- Similarly, the effect of the flavor VEVs will yield the following set of GB mass matrices

$$M_{KK}^2 (\alpha_u \text{Tr}[Y_u A_u A_u Y_u] + \beta_u \text{Tr}[A_Q V Y_u Y_u V^\dagger A_Q] + \gamma_u \text{Tr}[V^\dagger A_Q V Y_u A_u Y_u] + \alpha_d \text{Tr}[Y_d A_d A_d Y_d] + \beta_d \text{Tr}[A_Q Y_d Y_d A_Q] + \gamma_d \text{Tr}[A_Q Y_d A_d Y_d])$$

diagonal real Yukawa couplings

FCNC's from FGB



$$\mathcal{L}_{mass} = M_{KK}^2 \text{Tr}[A_Q A_Q] + \frac{4g_{Q*}^2 R^3}{3R'^2} \text{Tr}[\phi_u A_Q A_Q \phi_u] \quad g_* \sim 4 \text{ allowed}$$

$$\frac{1}{M_{KK}^2} \frac{2g_{Q*}^6 \alpha_u^4 y_t^4}{27(M_{KK} R')^4 + 42(g_{Q*} \alpha_u y_t)^2 (M_{KK} R')^2 + 16(g_{Q*} \alpha_u y_t)^4} ((V^\dagger Y_u^2 V)_{ij})^2 (\bar{d}_i \gamma_L^\mu d_j)(\bar{d}_i \gamma_L^\mu d_j)$$