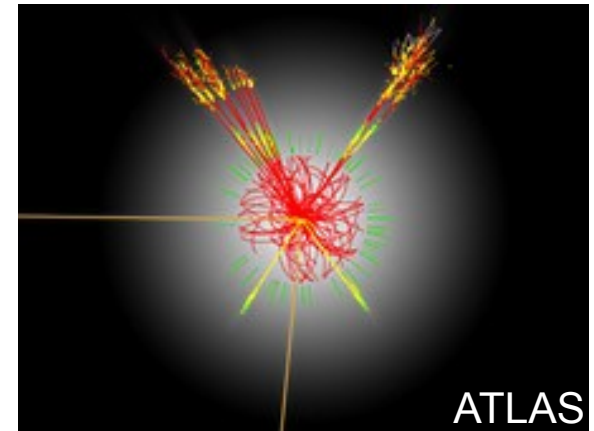
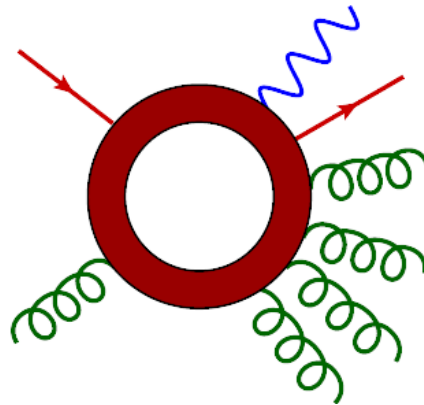
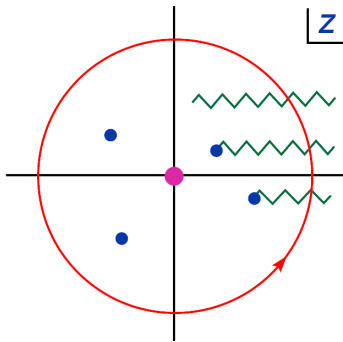


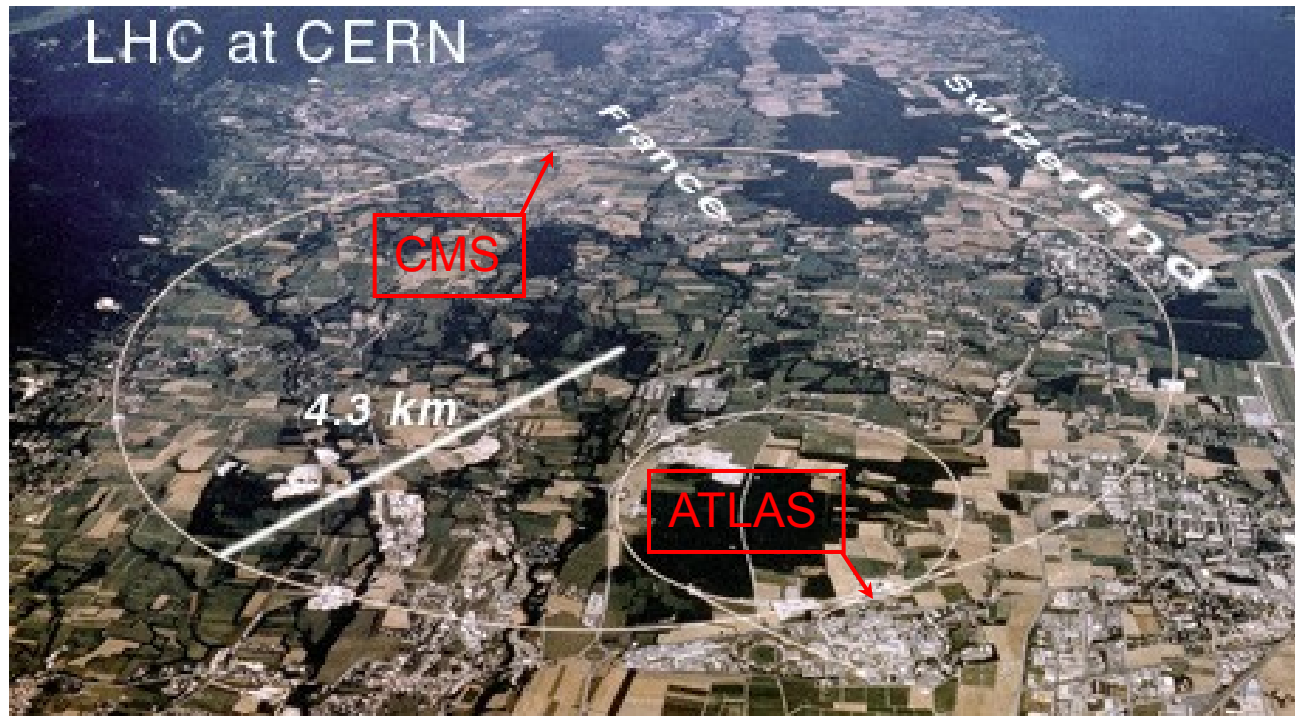
New Tools for Forecasting Old Physics at the LHC



Lance Dixon (CERN & SLAC)

CERN Colloquium
20 January, 2011

The Large Hadron Collider



- Proton-proton collisions at **7 → 14 TeV** center-of-mass energy, **3.5 → 7 times greater** than previous (Tevatron)
- Luminosity (collision rate) → **10—100 times greater**
- **New window** into physics at shortest distances – **opening now!**

New Physics around the Corner

Expect new physics at the **100 GeV – 1 TeV** mass scale, associated with **electroweak symmetry breaking**. At least, a **Higgs boson** (or similar)

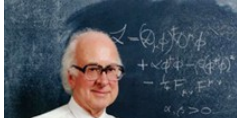
- Many theories predict a host of **new massive particles** in this mass range, including a **dark matter candidate**

- **supersymmetry**
- **new dimensions of space-time**
- **new forces**
- **etc.**

- Most new massive particles **decay rapidly** to old, ~massless particles: **quarks**, **gluons**, **charged leptons**, **neutrinos**, **photons**

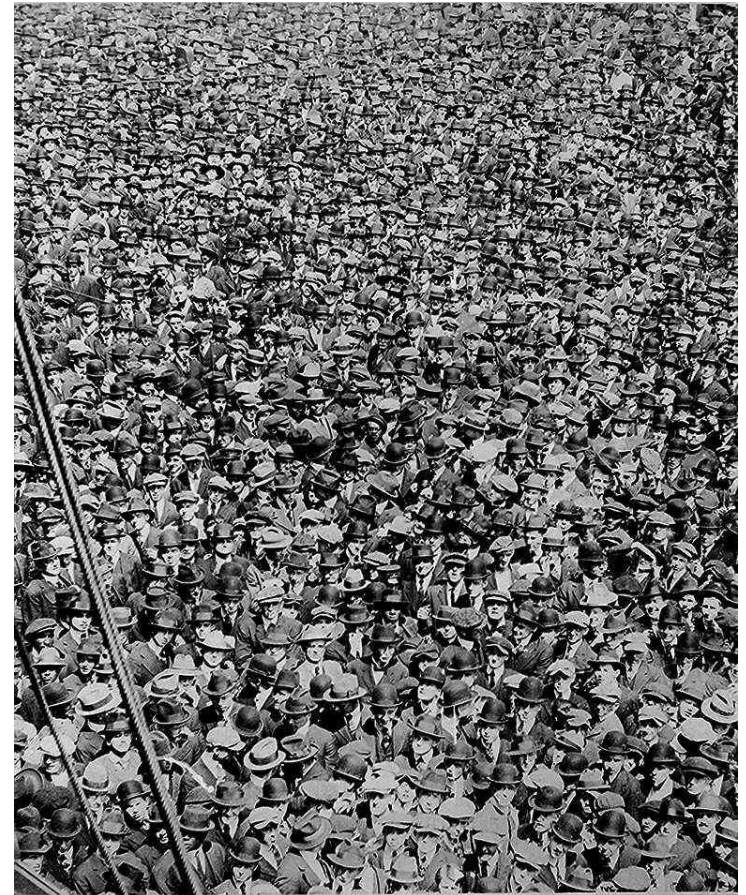
- **How to distinguish new physics from old (Standard Model)?**
- **From other types of new physics?**

Signals vs. Backgrounds



electron-positron colliders
– small backgrounds

vs.

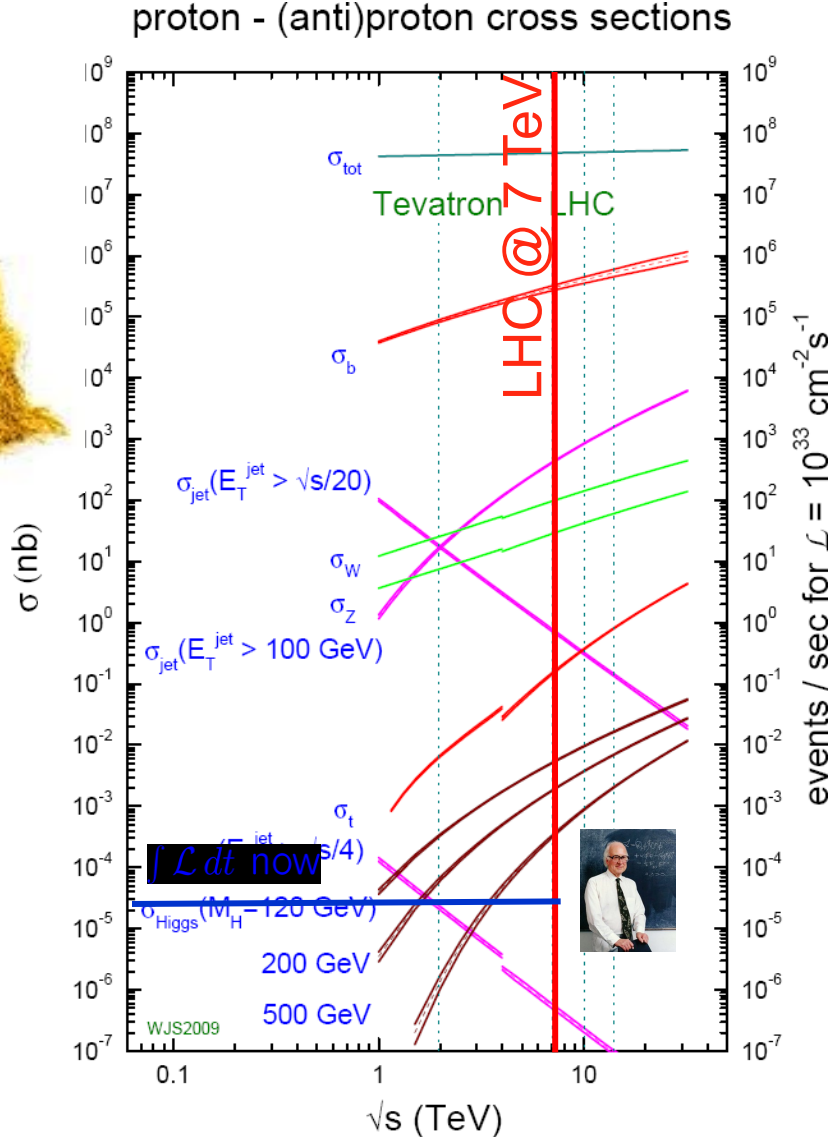


hadron colliders
– large backgrounds

LHC Data Dominated by Jets



new physics →



Jets come from quarks and gluons.

- q, g from decay of new particles?
- Or from old QCD?

• Every process shown also comes with one more jet at $\sim 1/5$ the rate

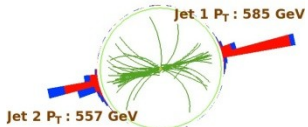
• Should understand Standard Model production of

$X + 1, 2, 3, \dots$ jets
where

$X = W, Z, tt, WW, H, \dots$

A Few Postcards from the Frontier

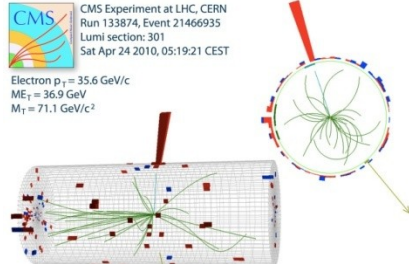
Run : 138919
Event : 32253996
Dijet Mass : 2.130 TeV



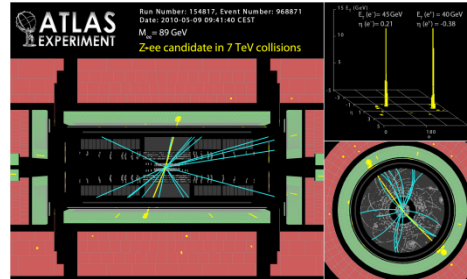
2 jets

CMS Experiment at LHC, CERN
Run 133874, Event 21466935
Lumi section: 301
Sat Apr 24 2010, 05:19:21 CEST

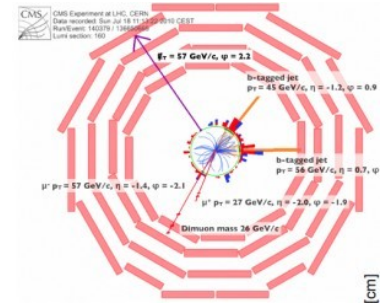
Electron $p_T = 35.6$ GeV/c
 $ME_T = 36.9$ GeV
 $M_T = 71.1$ GeV/c²



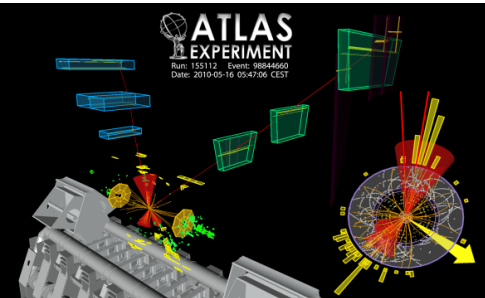
W



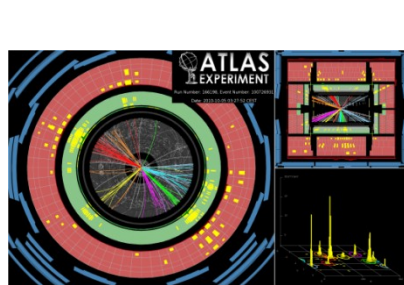
Z



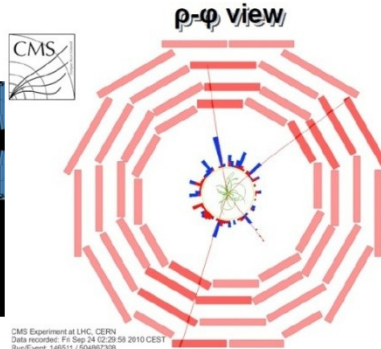
top



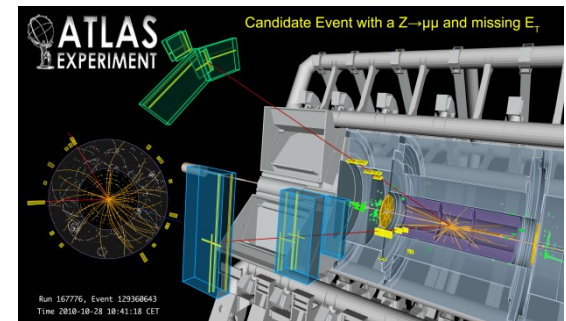
W + 3 jets



8 jets

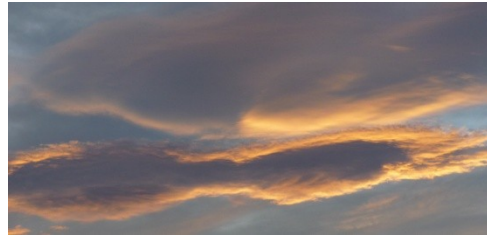
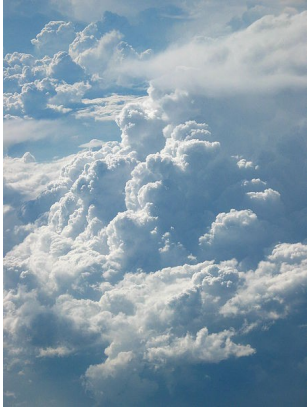


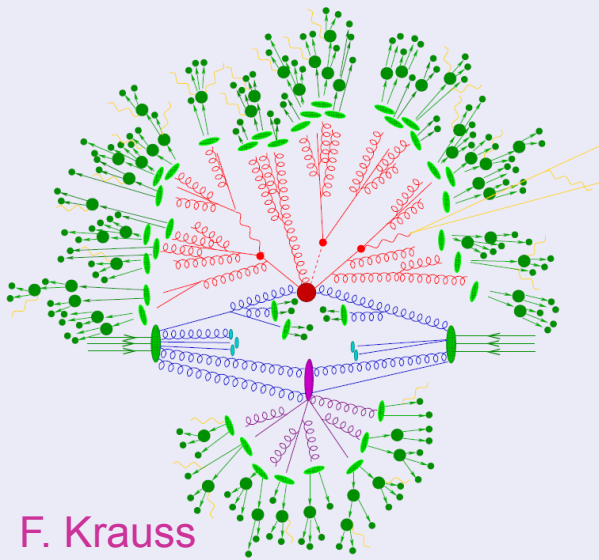
ZZ



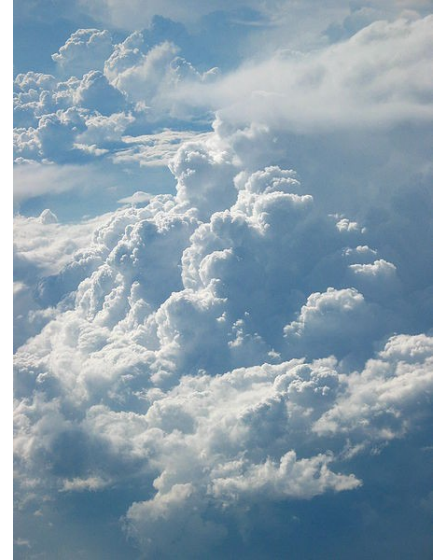
ZZ?

Now let's talk about the weather...





LHC events and clouds



F. Krauss

- **Both have fractal properties**
- **Cannot predict individual events**
 - Clouds: unpredictable turbulence, etc., on **small** distance scales
 - LHC: quantum mechanics, plus unpredictable QCD (strongly coupled) at **long** distances
- **All about predicting suitable ensembles**
 - weather, or climate, using global circulation models (**large** distances)
 - cross sections or probabilities that (ideally) are only sensitive to **short** distances (**infrared safe**) using perturbative QCD
 - **new physics ~ climate change**

Asymptotic Freedom

Gross, Wilczek, Politzer (1973)

Glueon self-interactions make quarks almost free, and make QCD calculable at short distances (high energies)

Quantum fluctuations of massless virtual particles polarize vacuum

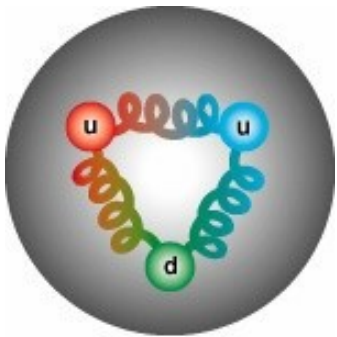
QED: electrons screen charge (e larger at short distances)

QCD: gluons anti-screen charge (g_s smaller at short distances)

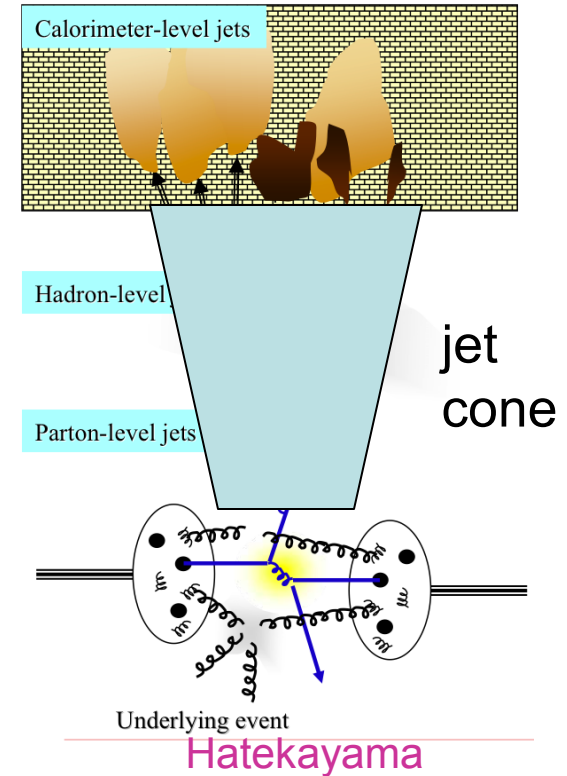
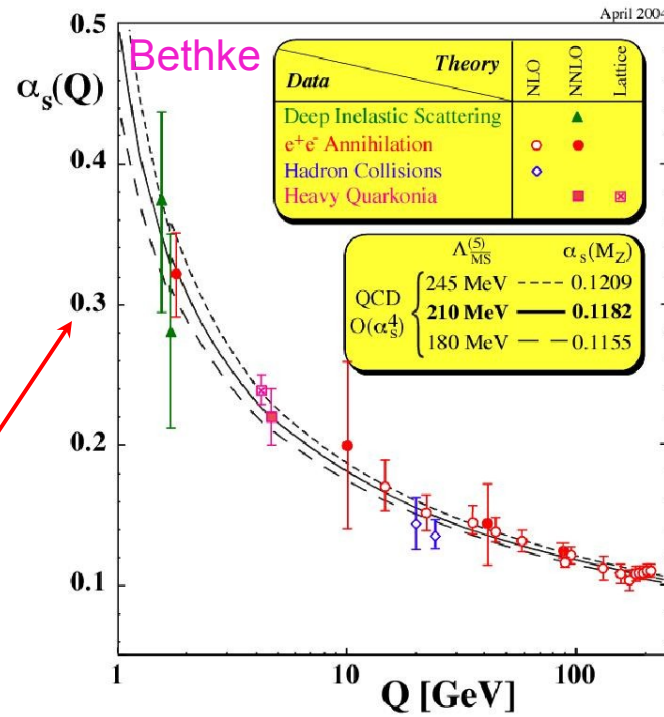


Asymptotic Freedom (cont.)

Running of α_s is *logarithmic*, *slow* at short distances (large Q)



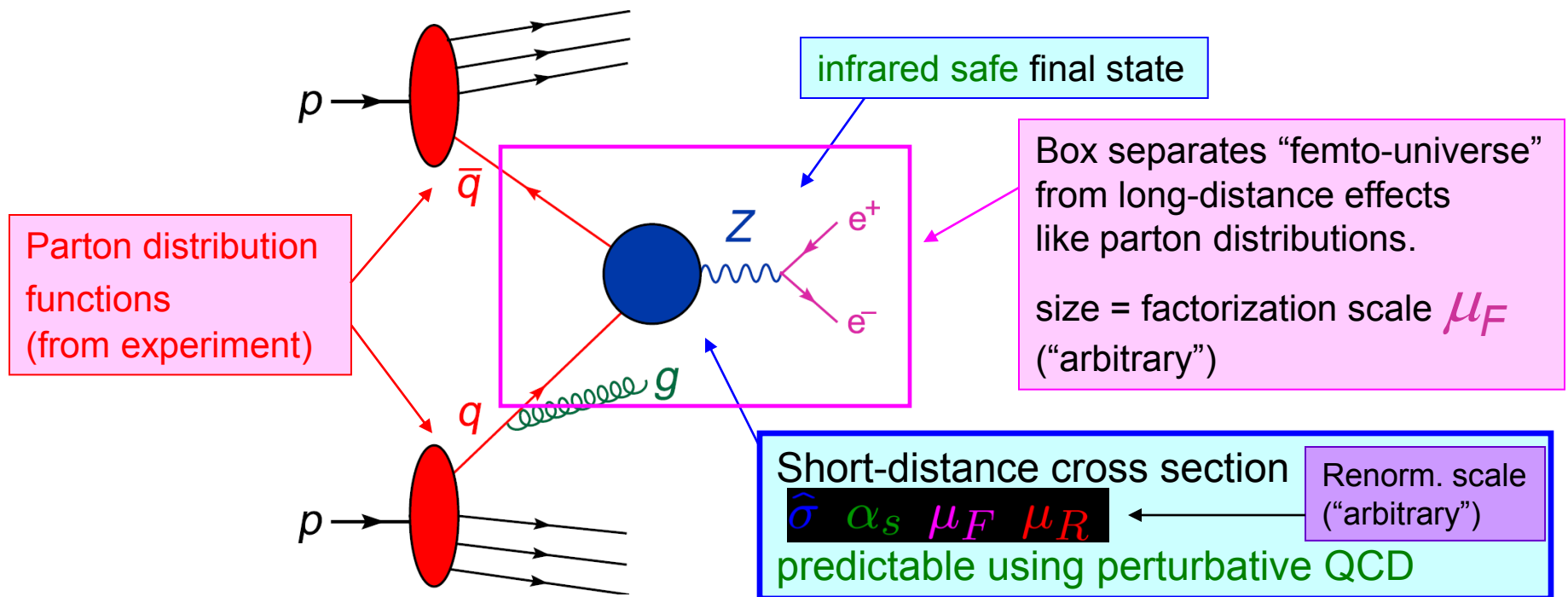
confining



calculable

QCD Factorization & Parton Model

Asymptotic freedom: At short distances, **quarks** and **gluons** (**partons**) in proton are **almost free**, and are sampled “one at a time”



Short-Distance Cross Section in Perturbation Theory

$$\hat{\sigma}(\alpha_s, \mu_F, \mu_R) = [\alpha_s(\mu_R)]^{n_\alpha} \left[\underset{\text{LO}}{\hat{\sigma}^{(0)}} + \frac{\alpha_s}{2\pi} \underset{\text{NLO}}{\hat{\sigma}^{(1)}}(\mu_F, \mu_R) + \left(\frac{\alpha_s}{2\pi}\right)^2 \underset{\text{NNLO}}{\hat{\sigma}^{(2)}}(\mu_F, \mu_R) + \dots \right]$$

Problem: Leading-order (LO) predictions only **qualitative**

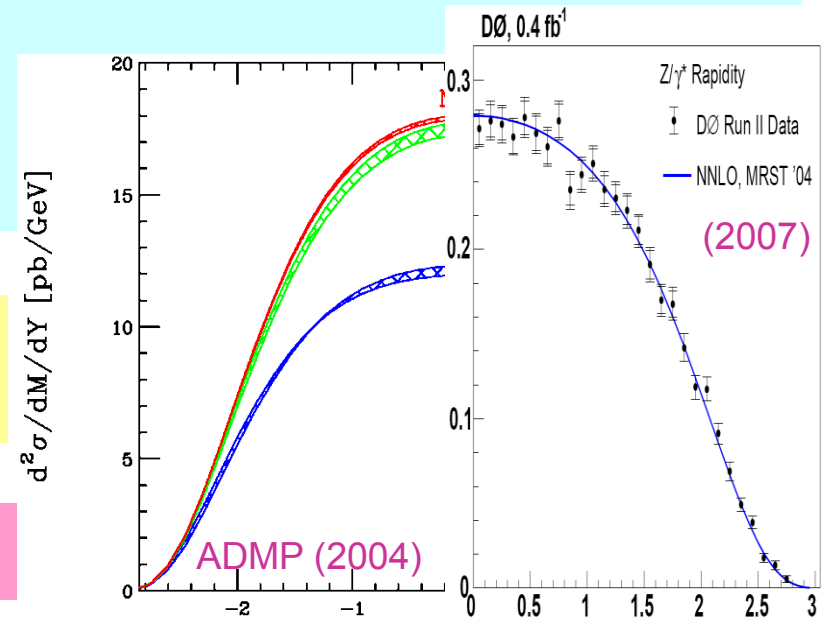
due to **poor convergence**

of expansion in α_s μ

(setting $\mu_R = \mu_F = \mu$)

Example: Z production at Tevatron
- Distribution in rapidity Y

~50% corrections, LO \rightarrow NLO

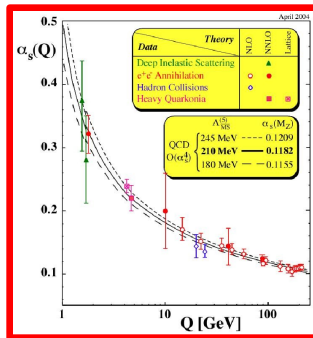
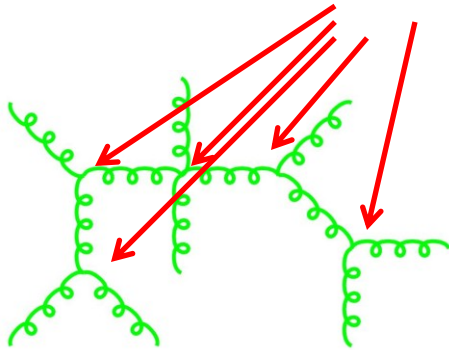


by NNLO, a precision observable

LO uncertainty increases with n_{jets}

$$\hat{\sigma}(\alpha_s, \mu_F, \mu_R) = [\alpha_s(\mu_R)]^{n_\alpha} \left[\hat{\sigma}^{(0)} + \frac{\alpha_s}{2\pi} \hat{\sigma}^{(1)}(\mu_F, \mu_R) + \left(\frac{\alpha_s}{2\pi}\right)^2 \hat{\sigma}^{(2)}(\mu_F, \mu_R) + \dots \right]$$

LO
NLO
NNLO

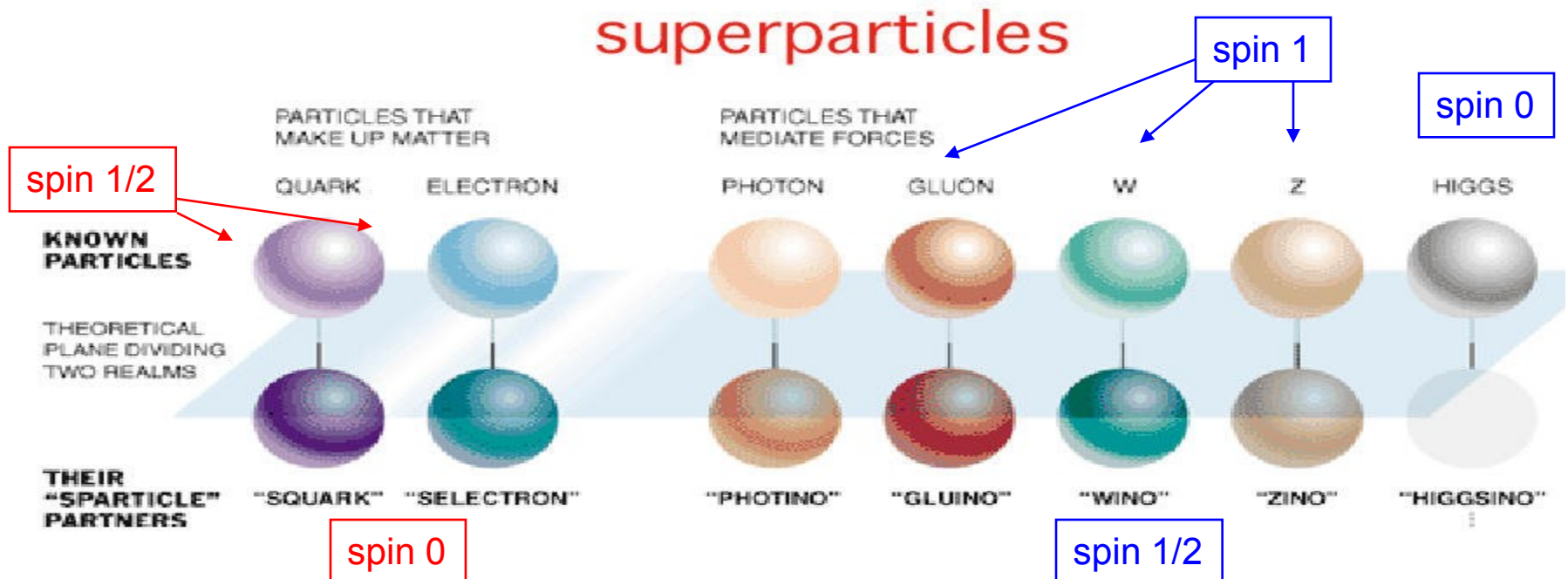


Uncertainty brought under much better control with NLO corrections: $\sim 50\%$ or more $\rightarrow \sim 15\text{-}20\%$

NLO really required for quantitative control of multi-jet final states

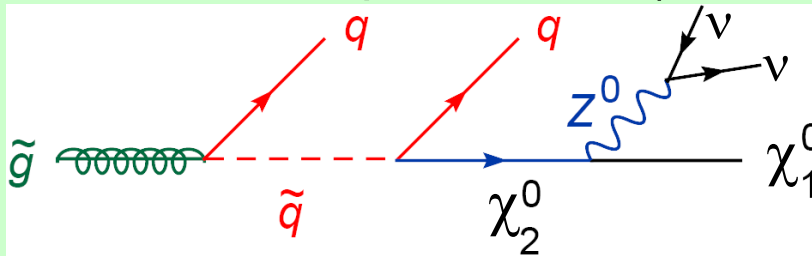
New Physics Example: Supersymmetry

- Symmetry between **fermions (matter)** and **bosons (forces)**
- Very elegant, also solves theoretical puzzles
- Lightest supersymmetric particle can be **dark matter**
- For **every** elementary particle already seen, another one should show up soon at LHC!

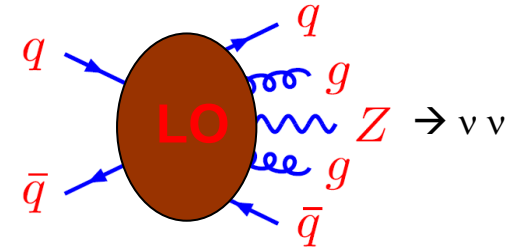


Backgrounds to Supersymmetry at LHC

- **Decay from gluino to neutralino**
(dark matter, escapes detector)



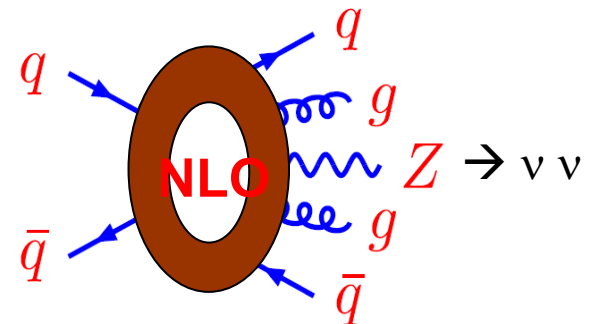
- **2 gluinos in event** \rightarrow
Signal: missing energy (MET) + 4 jets
- **SM background:** $Z + 4$ jets,
 $Z \rightarrow$ neutrinos



Current state of art for $Z + 4$ jets based on **LO approximation**
 \rightarrow normalization still quite uncertain

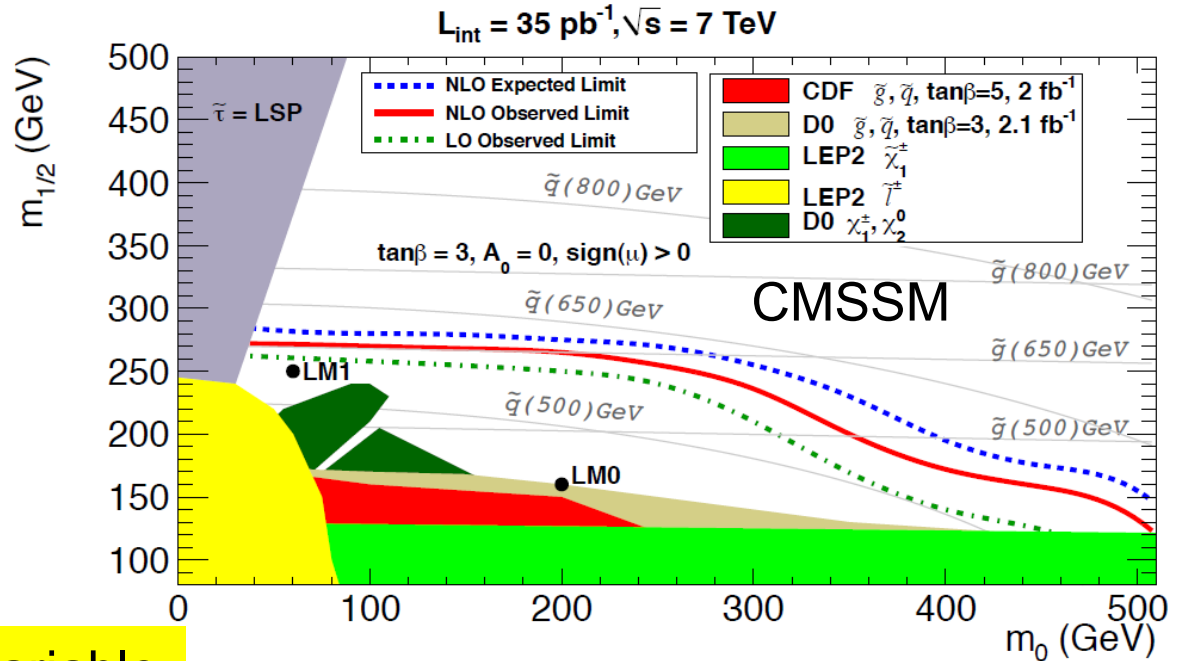
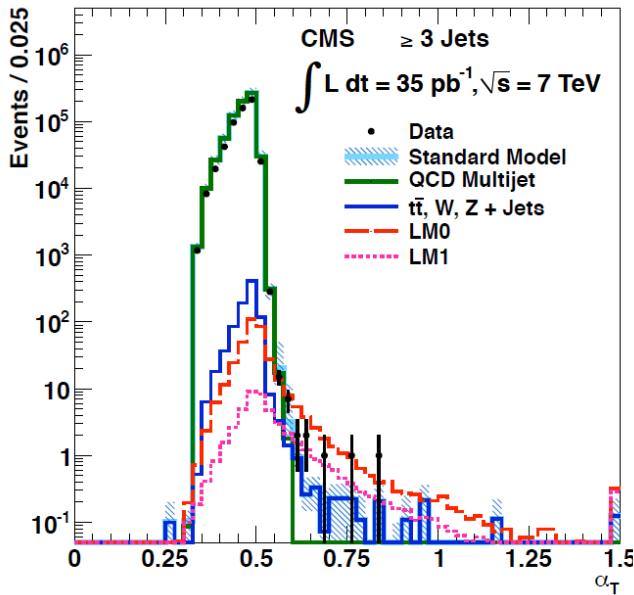
- Motivates goal of

$pp \rightarrow Z + 4$ jets
at NLO in QCD



New Limits on Supersymmetry from LHC

CMS, 1101.1628



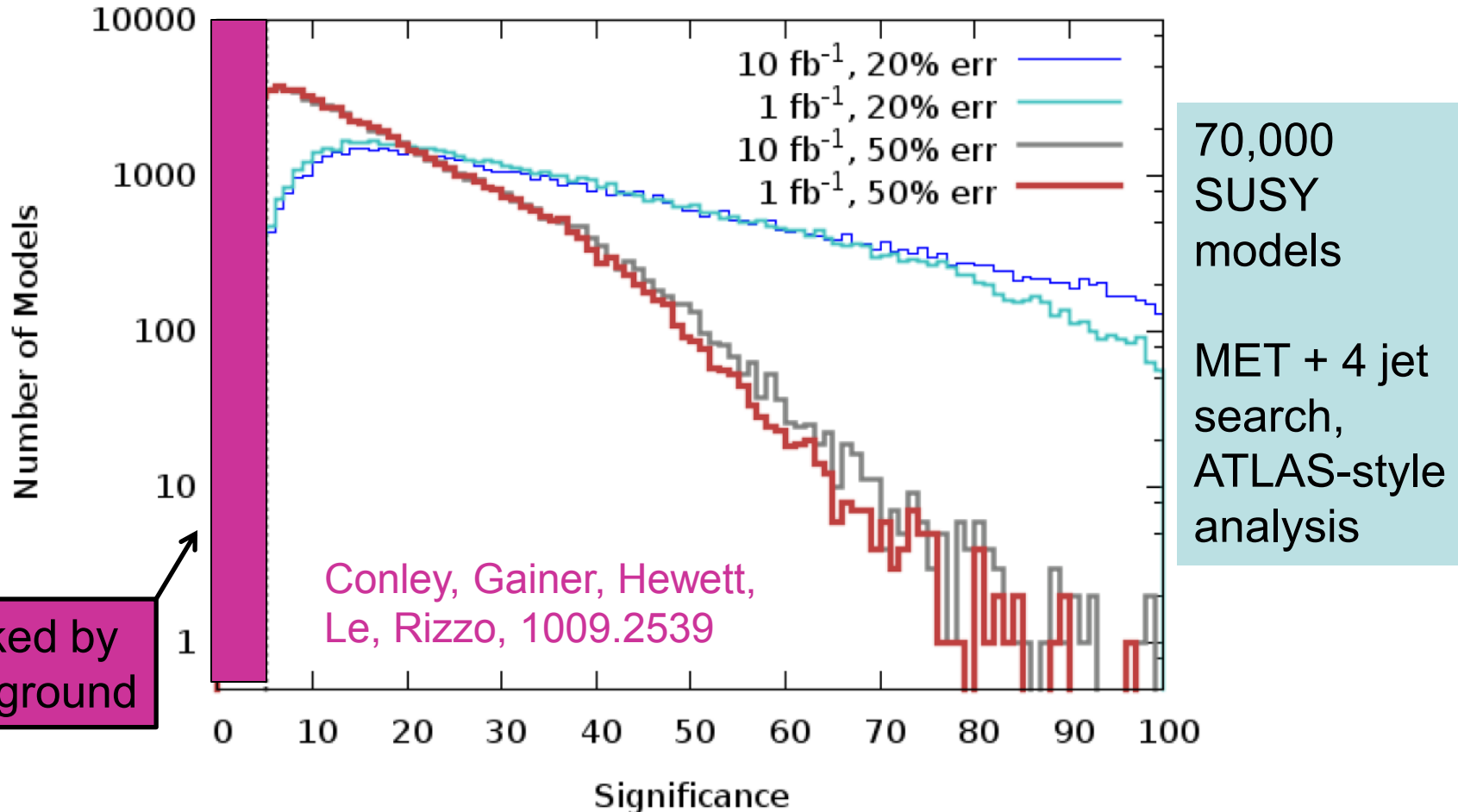
$\alpha_T = \text{MET} + \text{jets shape variable}$

- LHC off to an extremely promising start!

- As data increases rapidly this year, better SM theory can help

Reducing Background Systematics Improves SUSY Search Sensitivity

Significance for 4j0l, flat priors

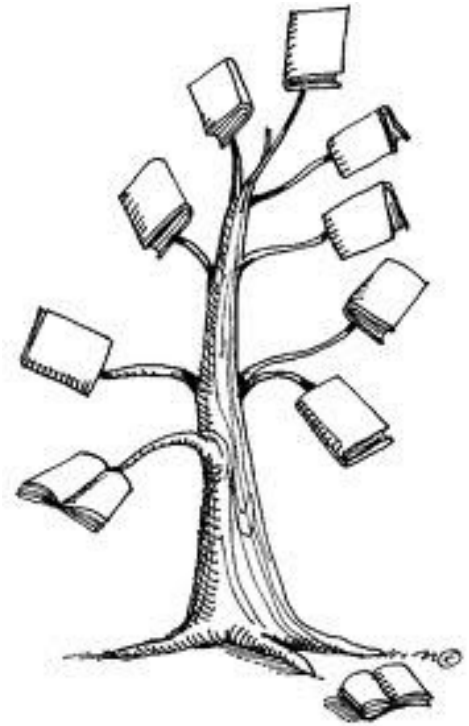
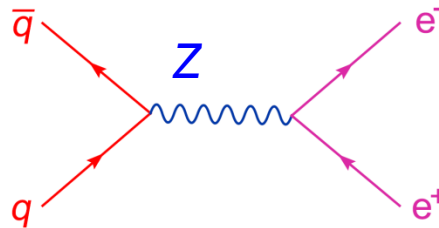


LO = Trees

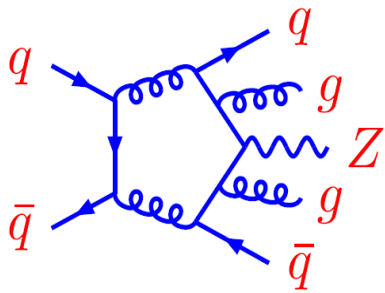
LO cross section can be computed using only Feynman diagrams with **no closed loops** – called **tree diagrams**. Here is a very simple one:

$$\hat{\sigma}(0)$$

LO



Although there are many kinds of trees, some harder than others, “textbook” methods often suffice



NLO = Loops



NLO cross section needs Feynman diagrams with **exactly one closed loop**

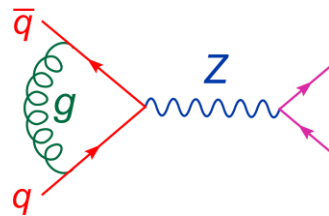
Where the **fun really starts** – textbook methods quickly fail, even with very powerful computers

- NLO also needs tree-level amplitudes with one more parton
- Both terms **infinite(!)** – combine them to get a finite result

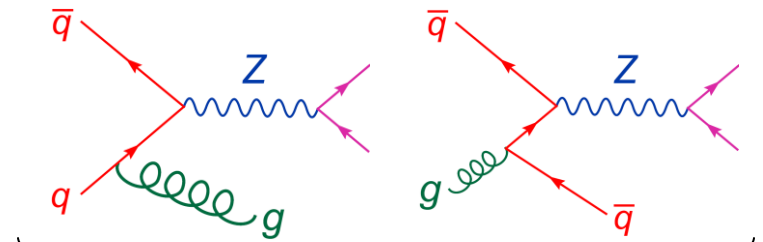


$$\hat{\sigma}(1)$$

NLO



1 loop



tree + 1 parton

- **One-loop amplitudes were the bottleneck for a long time** – focus today just on this part of the problem

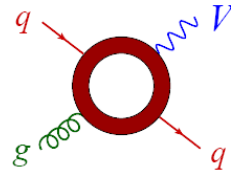
Loops get difficult quickly!

For $pp \rightarrow W + n \text{ jets}$ (maximum number of external gluons only)

of jets

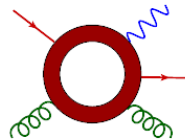
1-loop Feynman diagrams

1



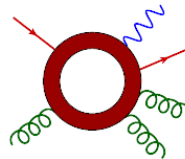
11

2

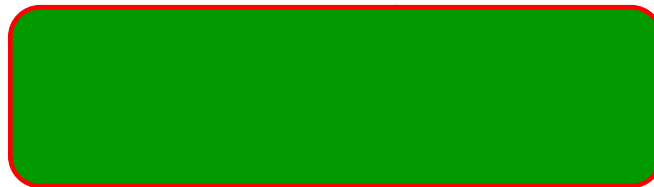


110

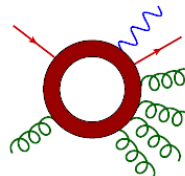
3



1,253



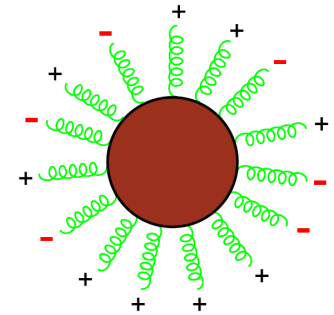
5



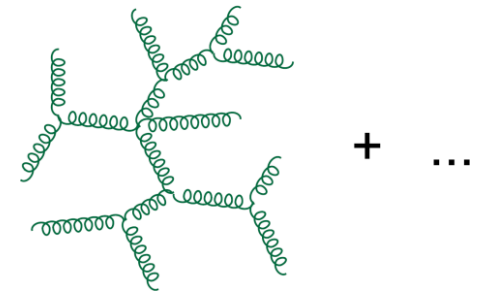
256,265

A Better Way to Compute?

- **Backgrounds** (and many **signals**) require detailed understanding of **scattering amplitudes** for many ultra-relativistic (“massless”) particles – especially **quarks** and **gluons** of **QCD**



- Long ago, **Feynman** told us how to do this – in principle



- However, **Feynman diagrams**, while **very general and powerful**, are **not optimized** for these processes
- There are more efficient methods for multi-jet processes!

Remembering a Simpler Time...

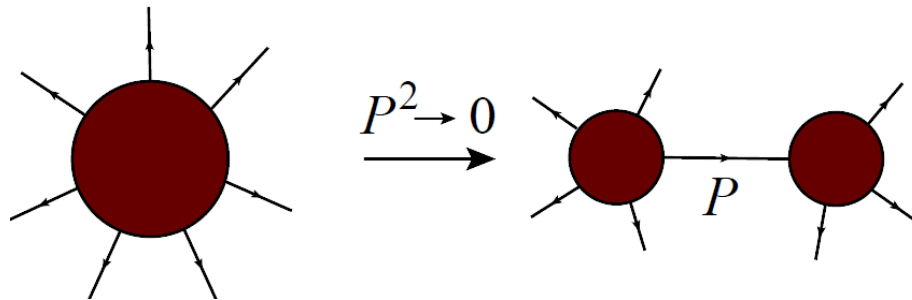


- In the 1960s there was no QCD, no Lagrangian or Feynman rules for the strong interactions

The Analytic S-Matrix

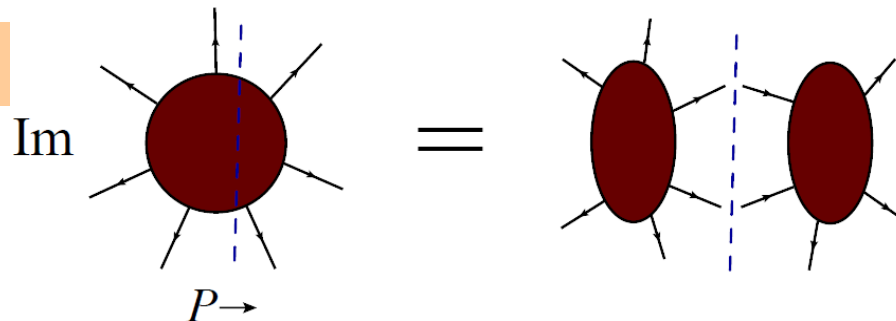
Bootstrap program for strong interactions: Reconstruct scattering amplitudes **directly** from **analytic properties**: “on-shell” information

• Poles



Landau; Cutkosky;
Chew, Mandelstam;
Eden, Landshoff,
Olive, Polkinghorne;
Veneziano;
Virasoro, Shapiro;
... (1960s)

• Branch cuts



Analyticity fell out of favor in 1970s with the rise of **QCD** & Feynman rules

Now **resurrected** for computing amplitudes in **perturbative QCD**
– as **alternative to Feynman diagrams!**

Perturbative information now assists analyticity.

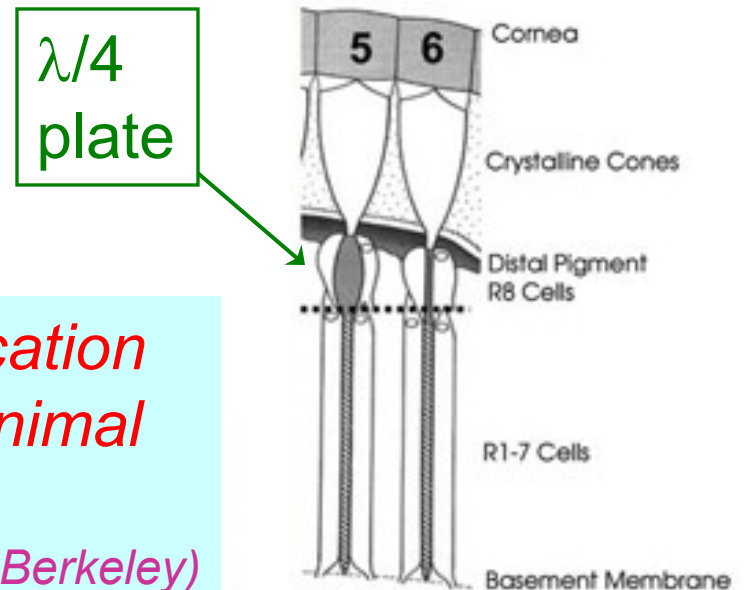
The Tail of the Mantis Shrimp

- Reflects left and right circularly polarized light differently

- Led biologists to discover that its eyes have differential sensitivity
- It communicates via the **helicity formalism**

“It's the most private communication system imaginable. No other animal can see it.”

- Roy Caldwell (U.C. Berkeley)

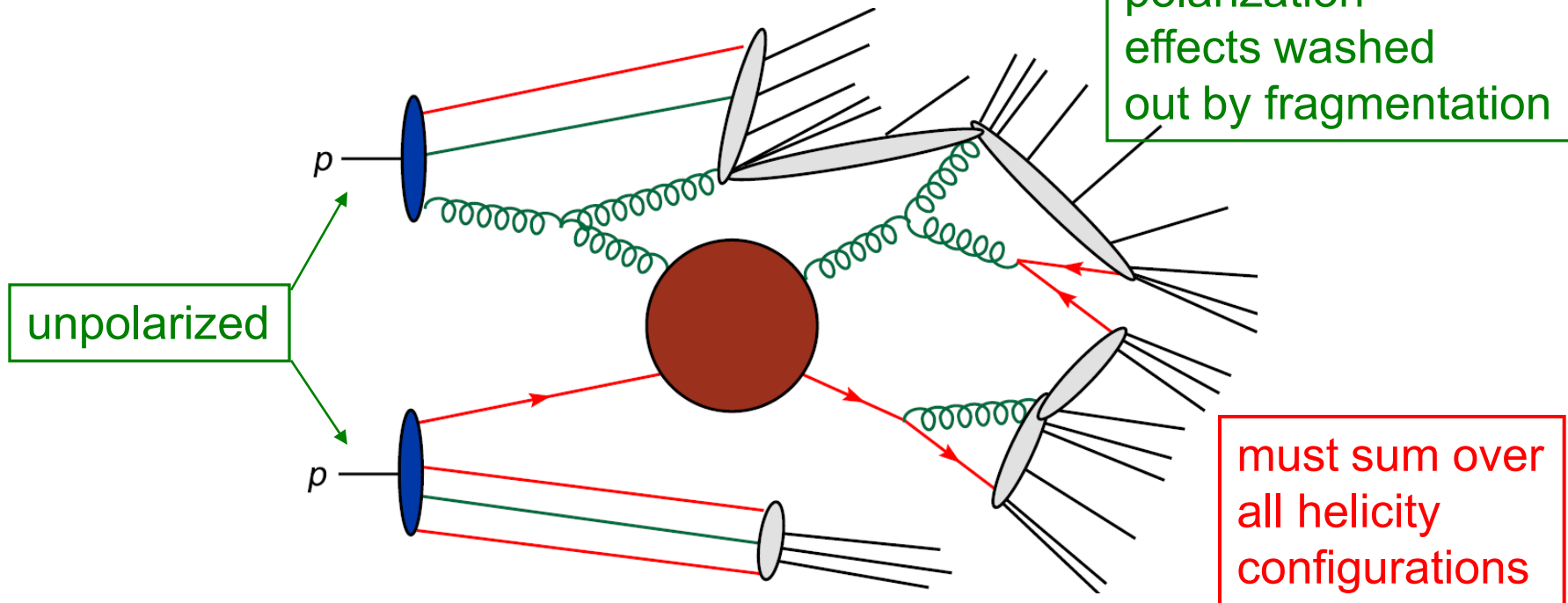


What the Biologists Didn't Know

Particle theorists have also evolved capability to communicate results via **helicity formalism**

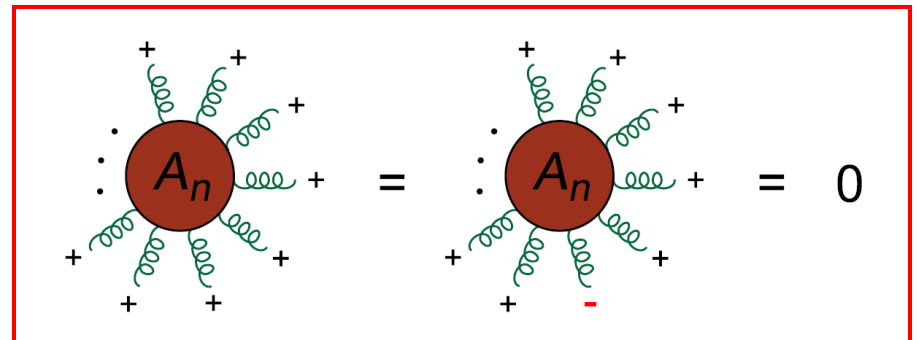
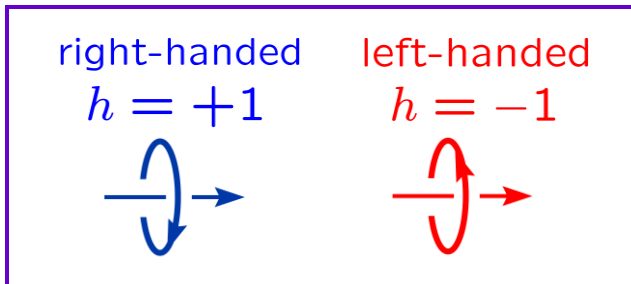
LHC experimentalists are blind to it

almost all final-state polarization effects washed out by fragmentation



Helicity Formalism Exposes Tree-Level Simplicity in QCD

Many helicity amplitudes either vanish or are very short



Analyticity makes it possible to recycle this simplicity into loop amplitudes

$$\frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$

Parke-Taylor formula (1986)

For Efficient Computation

Reduce

the number of “diagrams”

Reuse

building blocks over & over

Recycle

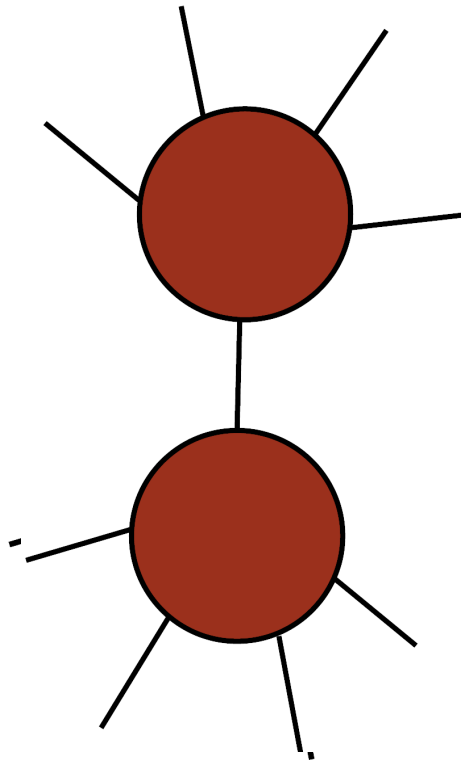
lower-point (1-loop) & lower-loop (tree)
on-shell amplitudes

Recurse



Recycling “Plastic” Amplitudes

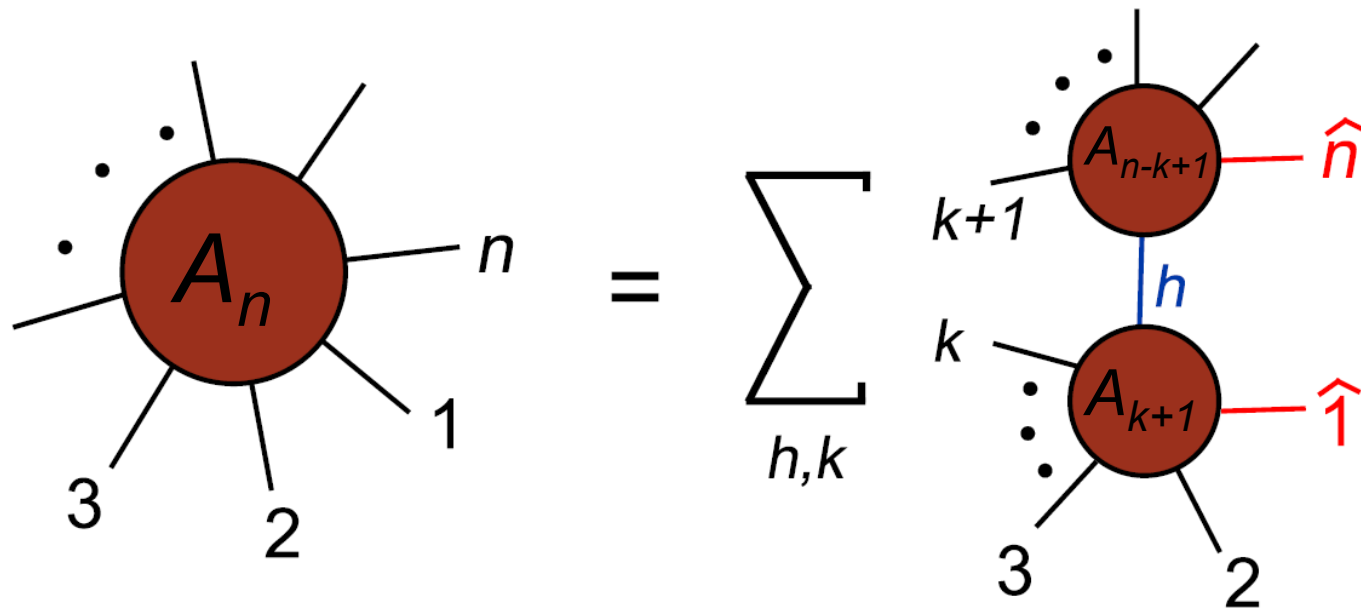
Amplitudes fall apart into simpler ones in special limits
– pole information



Pole information

→ BCFW (On-shell) Recursion Relations

Britto, Cachazo, Feng, Witten, hep-th/0501052

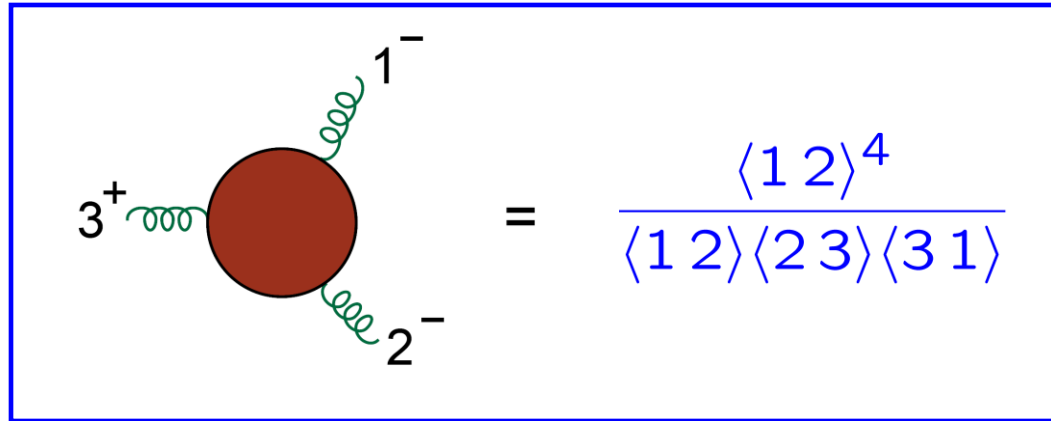


A_{k+1} and A_{n-k+1} are **on-shell** tree amplitudes with **fewer** legs, and with momenta **shifted** by a **complex** amount

Trees recycled into trees



All Gluon Tree Amplitudes Built From:



A diagram showing a central brown circle representing a three-gluon vertex. Three wavy green lines extend from the circle: one to the left labeled 3^+ , one to the top-right labeled 1^- , and one to the bottom-right labeled 2^- . To the right of the circle is an equals sign followed by a fraction: the numerator is $\langle 1 2 \rangle^4$ and the denominator is $\langle 1 2 \rangle \langle 2 3 \rangle \langle 3 1 \rangle$.

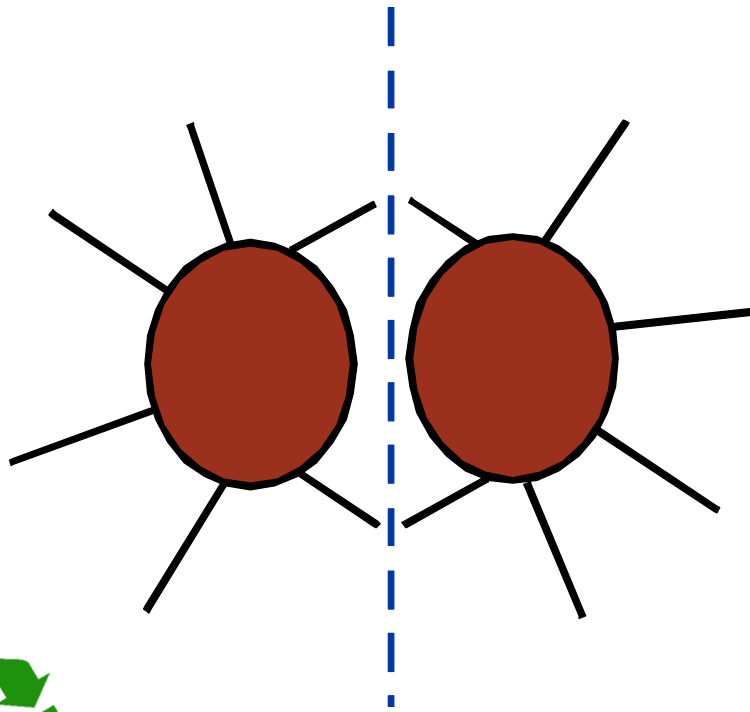
In contrast to Feynman vertices, it is on-shell, completely physical



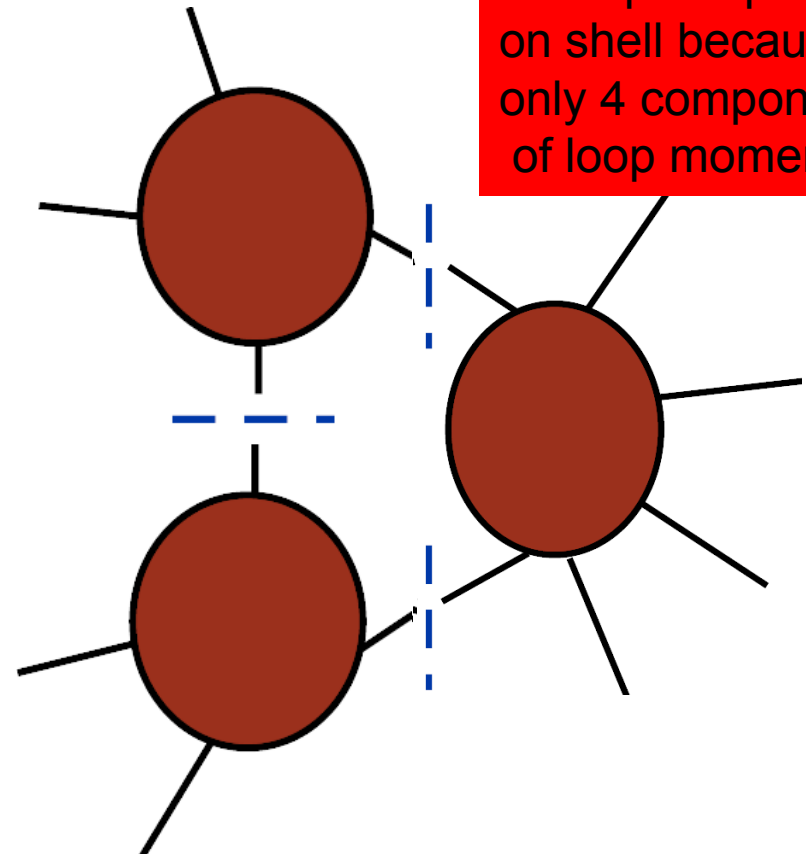
- On-shell recursion leads to very compact **analytic** formulae, and fast **numerical** implementation.
- Can do same sort of thing at **loop level**.

Branch cut information → Generalized Unitarity (One-loop Plasticity)

Ordinary unitarity:
put 2 particles on shell



Generalized unitarity:
put 3 or 4 particles on shell



Can't put 5 particles on shell because only 4 components of loop momentum



Trees recycled into loops!

One-Loop Amplitude Decomposition

Bern, LD, Dunbar, Kosower (1994)

- Missing from the old, nonperturbative analytic S-matrix



coefficients can be determined from products of **trees** using (generalized) unitarity

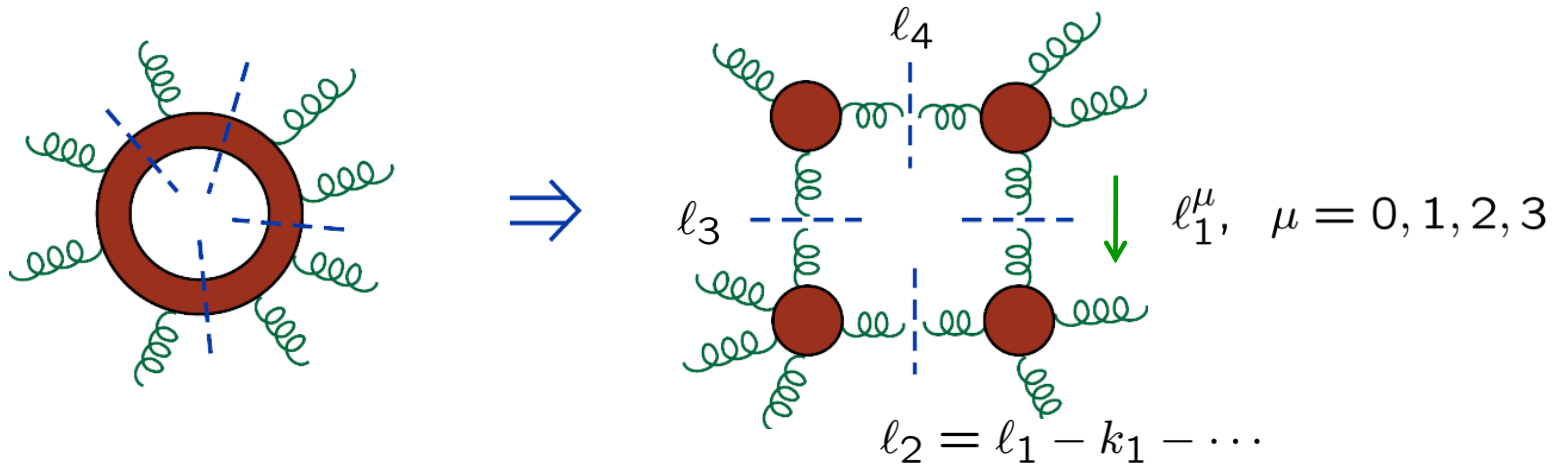
$$A^{1\text{-loop}} = \sum_i d_i \text{[box diagram]} + \sum_i c_i \text{[triangle diagram]} + \sum_i b_i \text{[bubble diagram]} + R + \mathcal{O}(\epsilon)$$

Known functions (integrals), same for all amplitudes

rational part; from D-dimensional trees, or recursively

Generalized Unitarity for Box Coefficients d_i

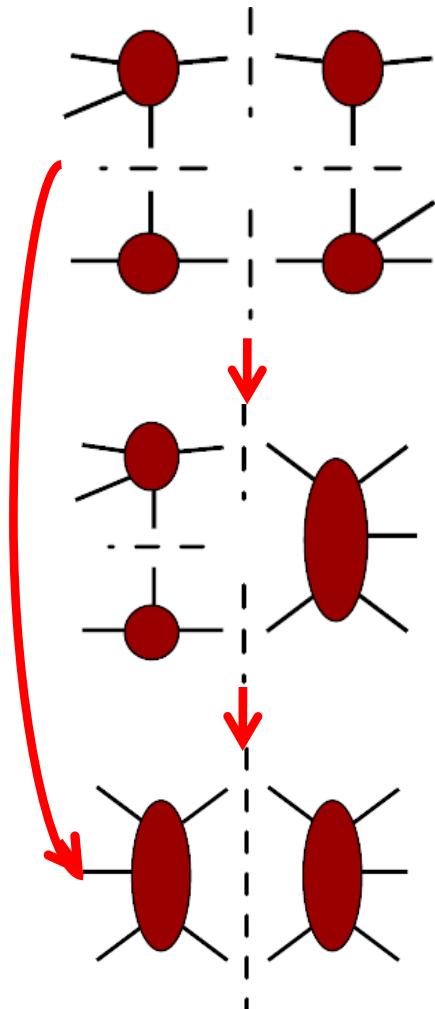
Britto, Cachazo, Feng, hep-th/0412103



Just multiply together 4 different tree amplitudes, evaluated at 2 different loop momenta that solve simple “quadruple cut” equations:

$$\begin{aligned}
 d_i &= A^{1\text{-loop}}(l_i)|_{l_i^2=m_i^2, i=1,2,3,4} \\
 &= \sum_{\pm} A_1^{\text{tree}}(l_0^{\pm}) A_2^{\text{tree}}(l_0^{\pm}) A_3^{\text{tree}}(l_0^{\pm}) A_4^{\text{tree}}(l_0^{\pm}) \\
 &= d_i^+ + d_i^-
 \end{aligned}$$

Rest of amplitude determined hierarchically



Each **box** coefficient comes **uniquely** from 1 “quadruple cut”

Ossola, Papadopolous, Pittau, hep-ph/0609007;
Mastrolia, hep-th/0611091; Forde, 0704.1835;
Ellis, Giele, Kunszt, 0708.2398; Berger et al., 0803.4180;...

Each **triangle** coefficient from 1 triple cut,
but “**contaminated**” by **boxes**

Each **bubble** coefficient from 1 double cut,
removing contamination by boxes and triangles

Bottom Line:

Trees recycled into loops!



Similar methods work for **multiple loops**
– especially in theories with lots of supersymmetry
like $N=4$ super-Yang-Mills and $N=8$ supergravity

Automated On-Shell Programs at One Loop

CutTools:

Ossola, Papadopolous, Pittau, 0711.3596

NLO WWW, WWZ, \dots

Binoth+OPP, 0804.0350

NLO $t\bar{t}b\bar{b}, t\bar{t} + 2 \text{ jets}, \dots$

Bevilacqua, Czakon, Papadopoulos,

Pittau, Worek, 0907.4723; 1002.4009; now going into MadGraph (Frederix, Frixione, ...)

Blackhat:

Berger, Bern, LD, Febres Cordero, Forde, H. Ita, D. Kosower, D. Maître; T. Gleisberg, 0803.4180, 0808.0941, 0907.1984, 1004.1659, 1009.2338

+ Sherpa \rightarrow NLO $W, Z + 3, 4 \text{ jets}$

Rocket:

Giele, Zanderighi, 0805.2152

Ellis, Giele, Kunstz, Melnikov, Zanderighi, 0810.2762

NLO $W + 3 \text{ jets}$ (large N_c), $W^+W^+ + 2 \text{ jets}$

EMZ, 0901.4101, 0906.1445; Melia, Melnikov, Rontsch, Zanderighi, 1007.5313

SAMURAI:

Mastrolia, Ossola, Reiter, Tramontano, 1006.0710

NGluon:

Badger, Biedermann, Uwer, 1011.2900

As a result...

**Dramatic increase recently
in rate of NLO predictions
for new processes!**

Les Houches Experimenters' Wish List

2010

process wanted at NLO	background to
1. $pp \rightarrow VV + \text{jet}$	$t\bar{t}H$, new physics Dittmaier, Kallweit, Uwer; Campbell, Ellis, Zanderighi
2. $pp \rightarrow H + 2 \text{ jets}$	H in VBF BCDEGMRSW; Campbell, Ellis, Williams Campbell, Ellis, Zanderighi; Ciccolini, Denner Dittmaier
3. $pp \rightarrow t\bar{t}b\bar{b}$	$t\bar{t}H$ Bredenstein, Denner Dittmaier, Pozzorini; Bevilacqua, Czakon, Papadopoulos, Pittau, Worek
4. $pp \rightarrow t\bar{t} + 2 \text{ jets}$	$t\bar{t}H$ Bevilacqua, Czakon, Papadopoulos, Worek
5. $pp \rightarrow VVb\bar{b}$	VBF $\rightarrow H \rightarrow VV$, $t\bar{t}H$, new physics
6. $pp \rightarrow VV + 2 \text{ jets}$	VBF $\rightarrow H \rightarrow VV$ Melia, Melnikov, Rontsch, Zanderighi VBF: Bozzi, Jäger, Oleari, Zeppenfeld
7. $pp \rightarrow V + 3 \text{ jets}$	new physics Berger, Bern, Dixon, Febres Cordero, Forde, Gleisberg, Ita, Kosower, Maitre; Ellis, Melnikov, Zanderighi
8. $pp \rightarrow VVV$	SUSY trilepton Lazopoulos, Melnikov, Petriello; Hankele, Zeppenfeld; Binoth, Ossola, Papadopoulos, Pittau
9. $pp \rightarrow b\bar{b}b\bar{b}$	Higgs, new physics GOLEM

Feynman
diagram
methods

now joined
by

on-shell
methods
based on
analyticity
(unitarity)

table courtesy of
C. Berger

Top Quark Pairs + Jets

- Like (W,Z) + jets, a very important class of backgrounds
- Jets can boost the $t\bar{t}$ system, increasing missing E_T , and provide jets to pass various signal cuts.
- Cross sections large – no electroweak couplings
- State of art:
 - **NLO $t\bar{t}$ + 1 jet:** Dittmaier, Uwer, Weinzierl, hep-ph/0703120,...
 - **+ top decays:** Melnikov, Schulze, 1004.3284
 - **+ NLO parton shower:** Kardos, Papadopoulos, Trócsányi, 1101.2672
- **NLO $t\bar{t}$ + $b\bar{b}$:** Bredenstein, Denner, Dittmaier, Pozzorini, 0905.0110, 1001.4006; Bevilacqua, Czakon, Papadopoulos, Pittau, Worek, 0907.4723
- **NLO $t\bar{t}$ + 2 jets:** Bevilacqua, Czakon, Papadopoulos, Worek, 1002.4009

NLO $pp \rightarrow t\bar{t}b\bar{b}$ at LHC

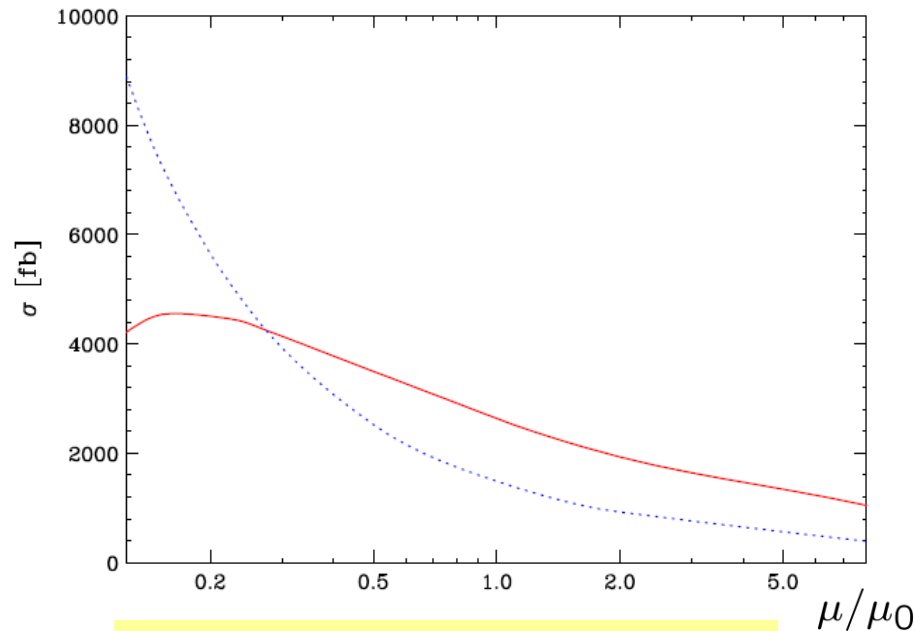
Background to $t\bar{t} + \text{Higgs}, H \rightarrow b\bar{b}$ at LHC (for λ_t)

First done using Feynman diagrams

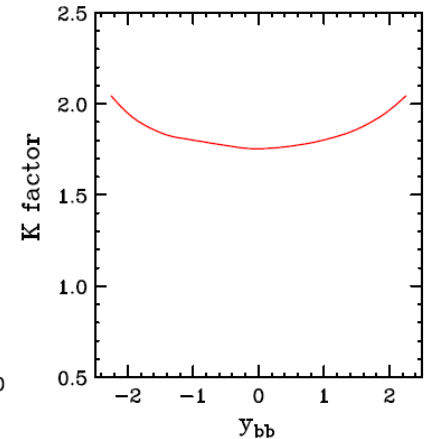
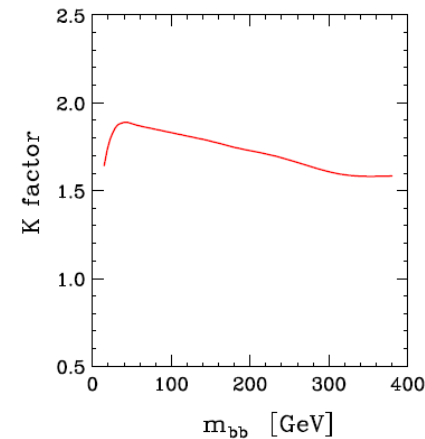
Recomputed using unitarity (**CutTools**)

Bredenstein et al.,
0807.1248, 0905.0110

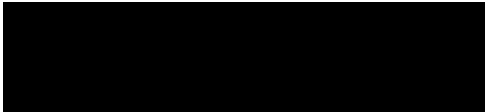
Bevilacqua et al., 0907.4723



much improved
scale uncertainties at NLO



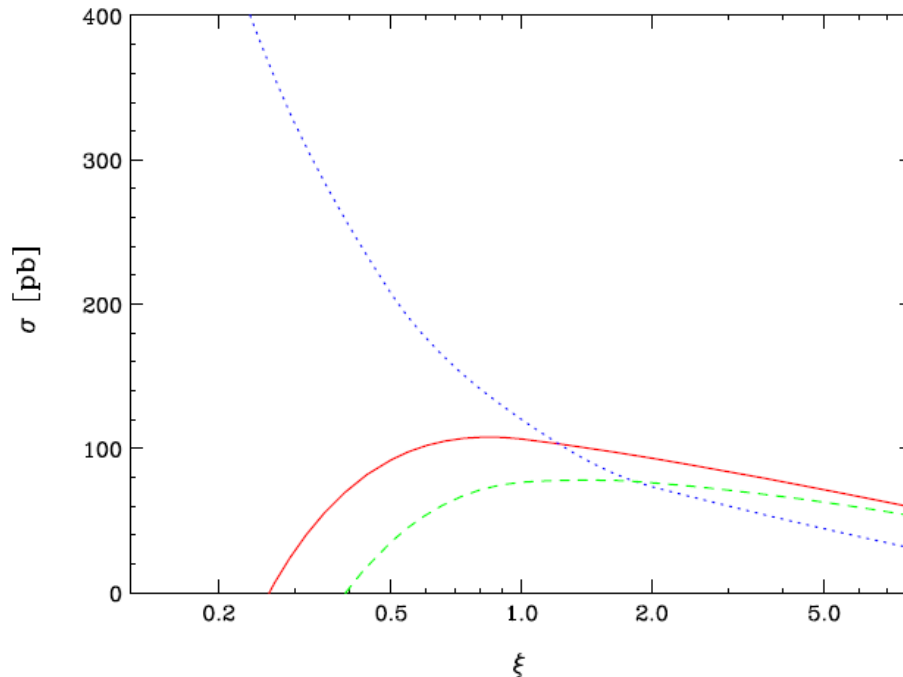
shape changes in bb distributions
from LO to NLO ($K = \text{NLO}/\text{LO}$)



Like , a background to 

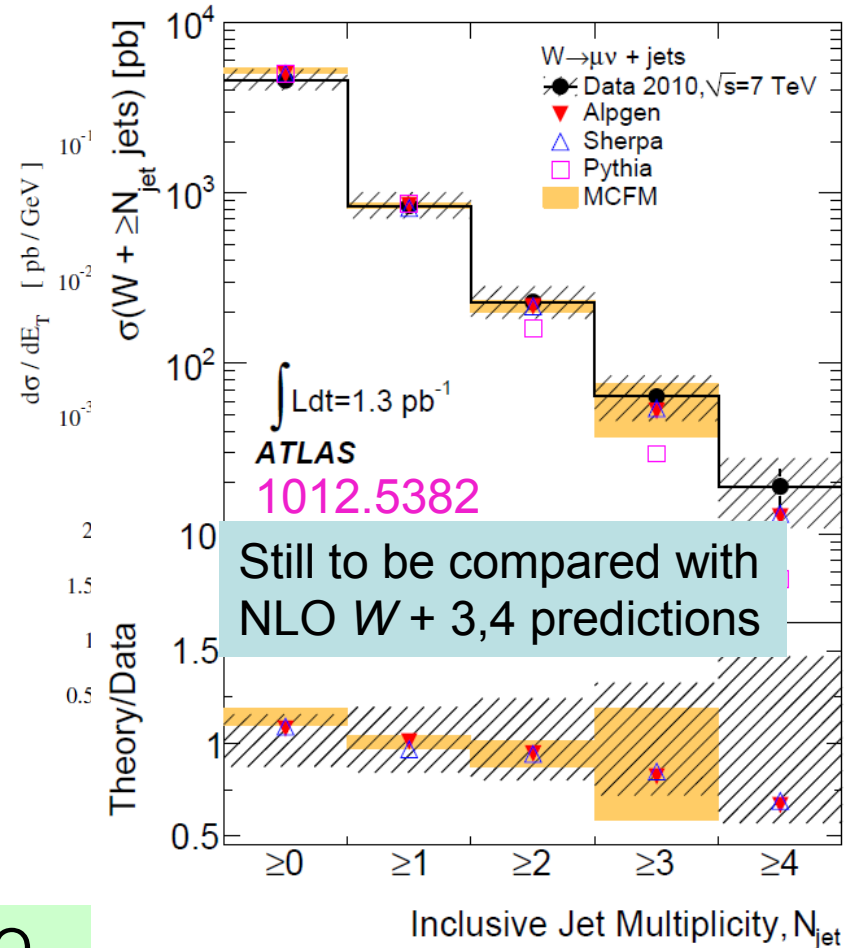
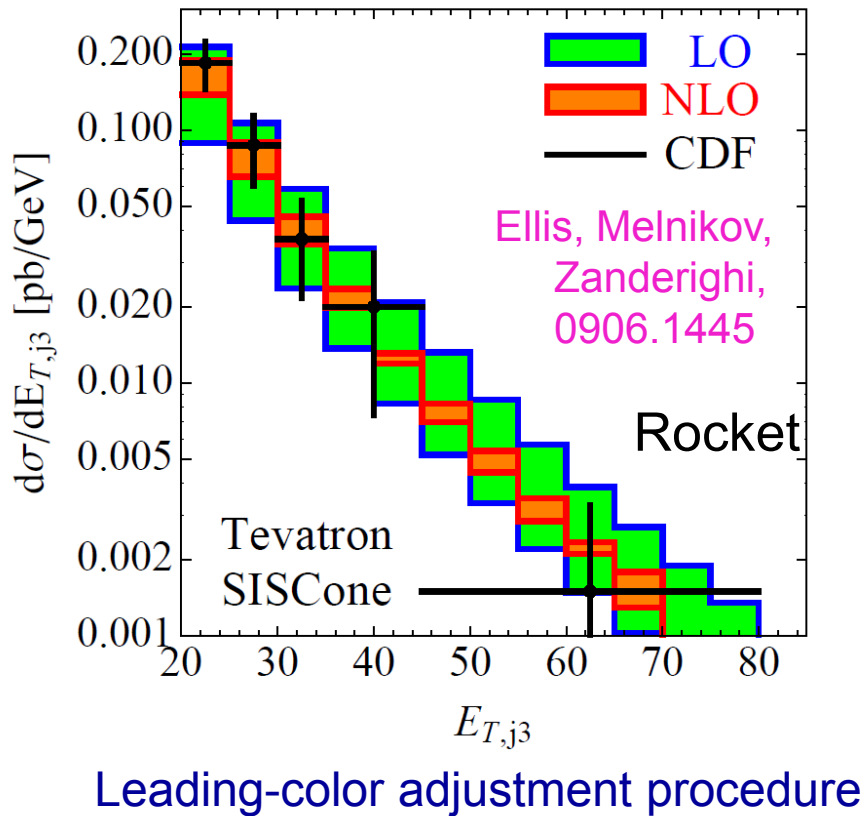
Only computed via unitarity (**CutTools**)

Bevilacqua, Czakon,
Papadopoulos,
Worek, 1002.4009



Again large reduction in scale dependence from LO \rightarrow NLO

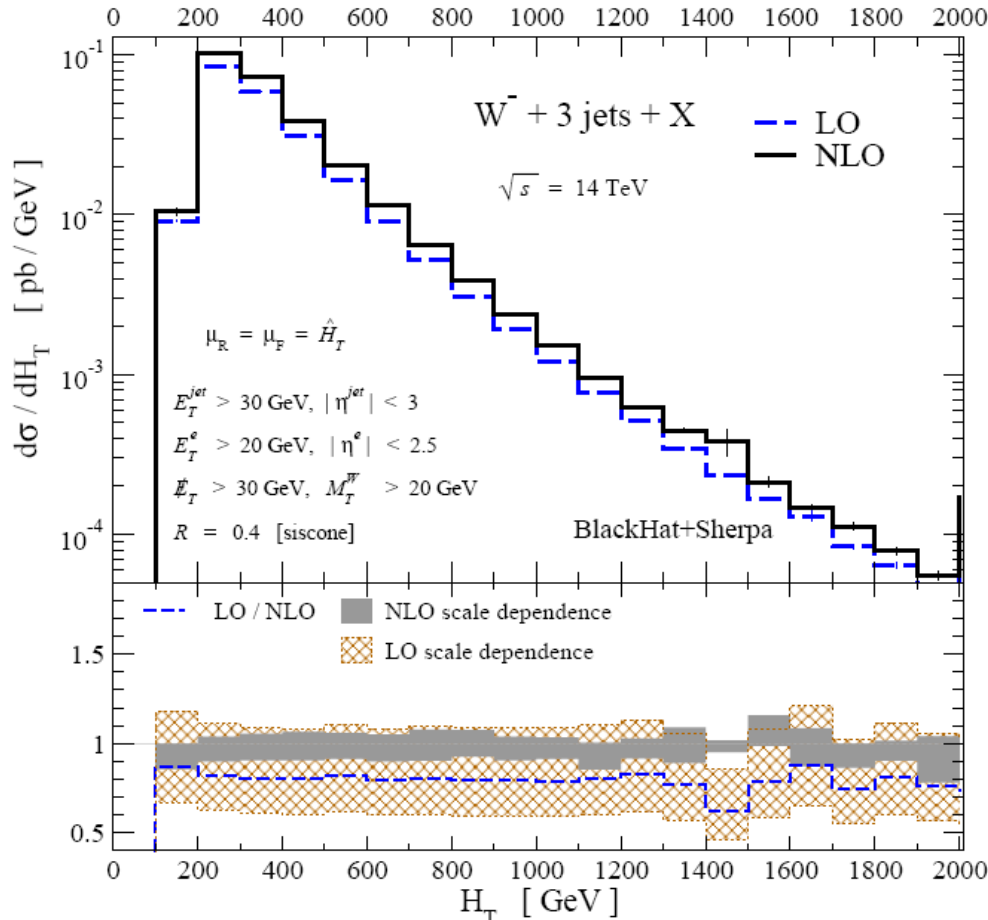
W + 3 jets at Tevatron → LHC



- Much smaller uncertainties than at LO.
- Agrees well with data; more data available now from Tevatron and LHC

Total Transverse Energy H_T at LHC

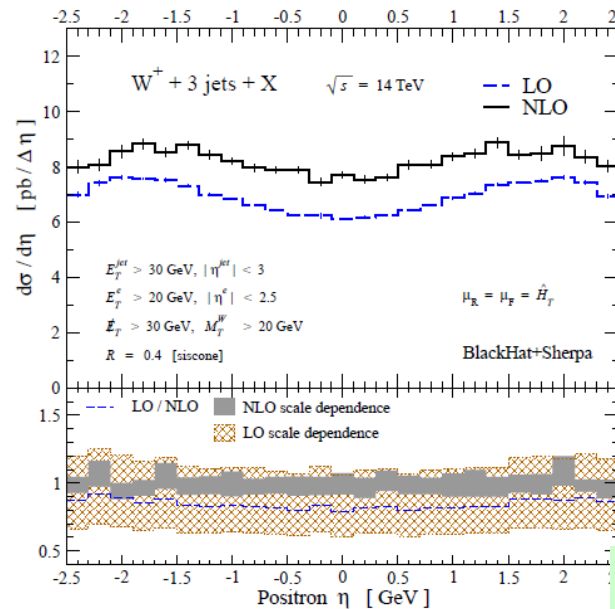
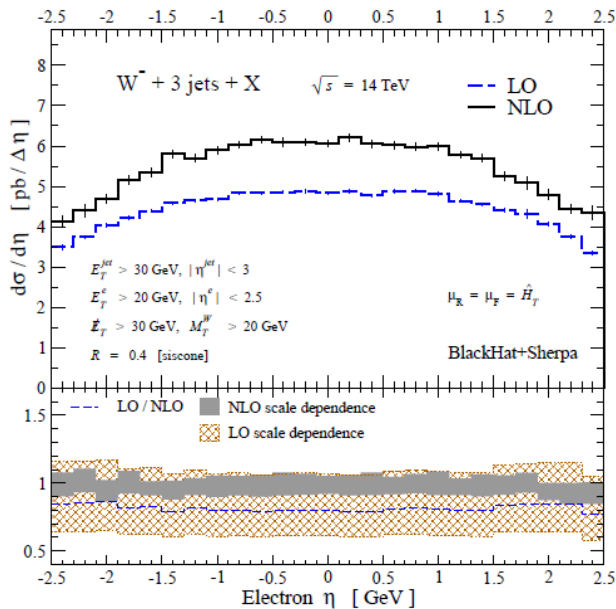
$$H_T = \sum_j E_{T,j}^{\text{jet}} + E_T^e + E_T^\nu \quad \text{often used in supersymmetry searches}$$



0907.1984

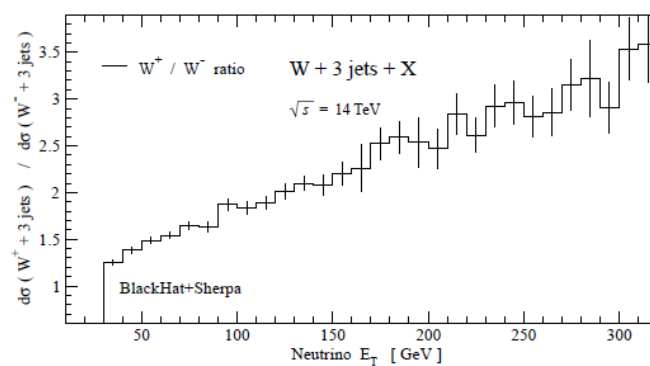
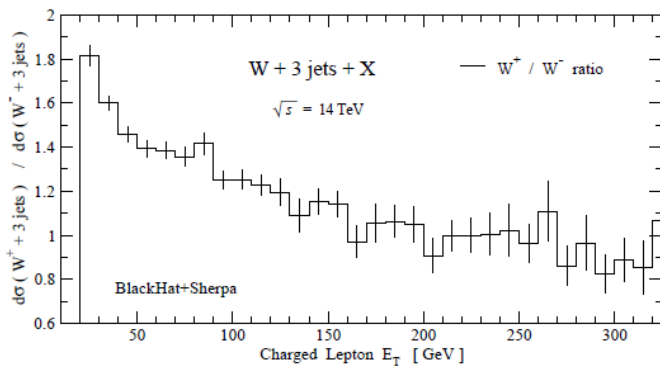
flat LO/NLO ratio
due to good
choice of
scale $\mu = H_T$

Positrons Differ from Electrons in $W^\pm + 3$ jets at LHC (pp)



rapidity
distributions
remember
 $u(x)/d(x)$ large
as $x \rightarrow 1$

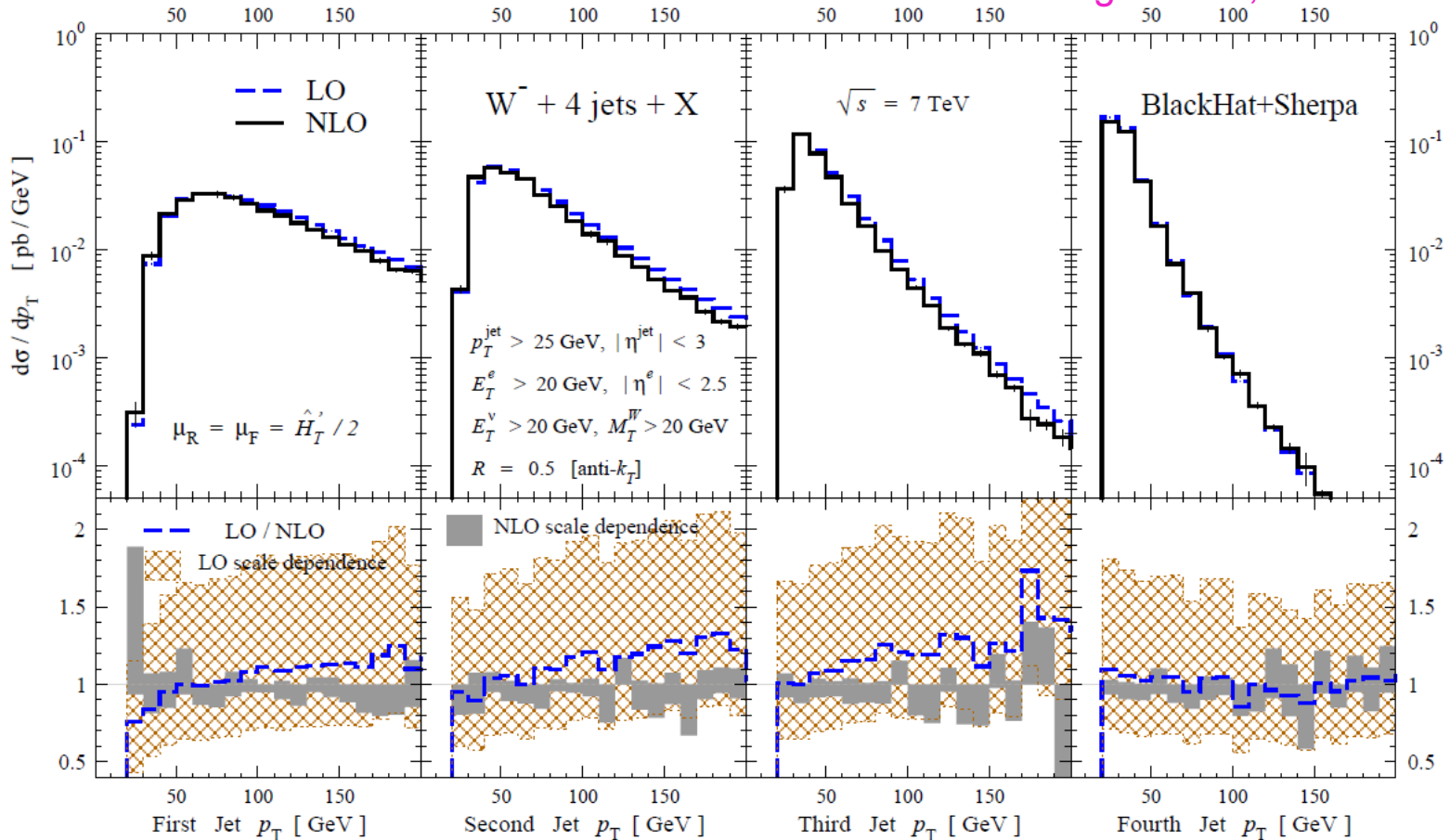
0907.1984



W^+/W^- transverse
ratios trace a
remarkably large
left-handed
 W polarization
– may be useful
to separate it from
top, new physics

NLO $pp \rightarrow W + 4 \text{ jets}$

Berger et al., 1009.2338



**First hadron collider process known at NLO with 5 objects in final state.
Also important SUSY background.**

One indicator of NLO progress

$pp \rightarrow W + 0 \text{ jet}$	1978	Altarelli, Ellis, Martinelli
$pp \rightarrow W + 1 \text{ jet}$	1989	Arnold, Ellis, Reno
$pp \rightarrow W + 2 \text{ jets}$	2002	Campbell, Ellis
$pp \rightarrow W + 3 \text{ jets}$	2009	BH+Sherpa Ellis, Melnikov, Zanderighi
$pp \rightarrow W + 4 \text{ jets}$	2010	BH+Sherpa

Conclusions

- **New and efficient** computational approaches to one-loop QCD amplitudes now used to forecast important Standard Model backgrounds at the LHC
 - exploit **analyticity/unitarity**: build loop amplitudes out of trees
 - implemented **numerically** in several programs:
BlackHat, CutTools, NGLuon, Rocket, Samurai, ...
- Long and growing list of complex processes computed at NLO with these techniques:
 - VVV ($V=W$ or Z)
 - $t\bar{t}bb$, $t\bar{t}j$, $t\bar{t}jj$
 - W^+W^+jj
 - $Wjjj$, $Zjjj$, $Wjjjj$
- Also very important to incorporate into NLO Monte Carlos, a la MC@NLO & POWHEG (no time to discuss here)
- Success will assist in optimal exploitation of LHC data!

