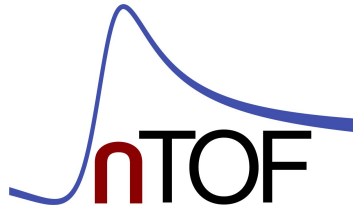


Analysis status of $^{94}\text{Nb}(n,\gamma)$ cross section measurement

J. Balibrea-Correa, V. Babiano-Suarez, J. Lerendegui-Marco, C. Domingo-Pardo, I. Ladarescu,
A. Tarifeño-Saldivia, V. Alcayne, D. Cano-Ott, E. González, T. Martínez, E. Mendoza, F. Calviño. A. Casanovas, C. Guerrero,
S. Heinitz, U. Köster, E. A. Maugeri, R. Dressler, D. Schumann, I. Mönch



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DE VALÈNCIA



CSIC

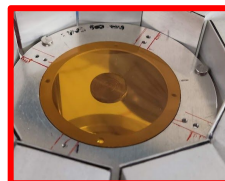
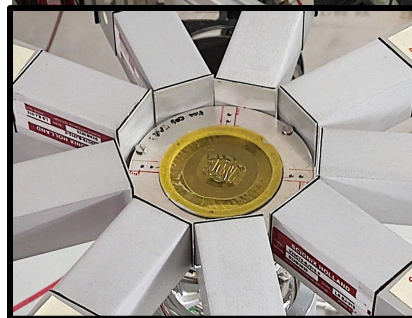
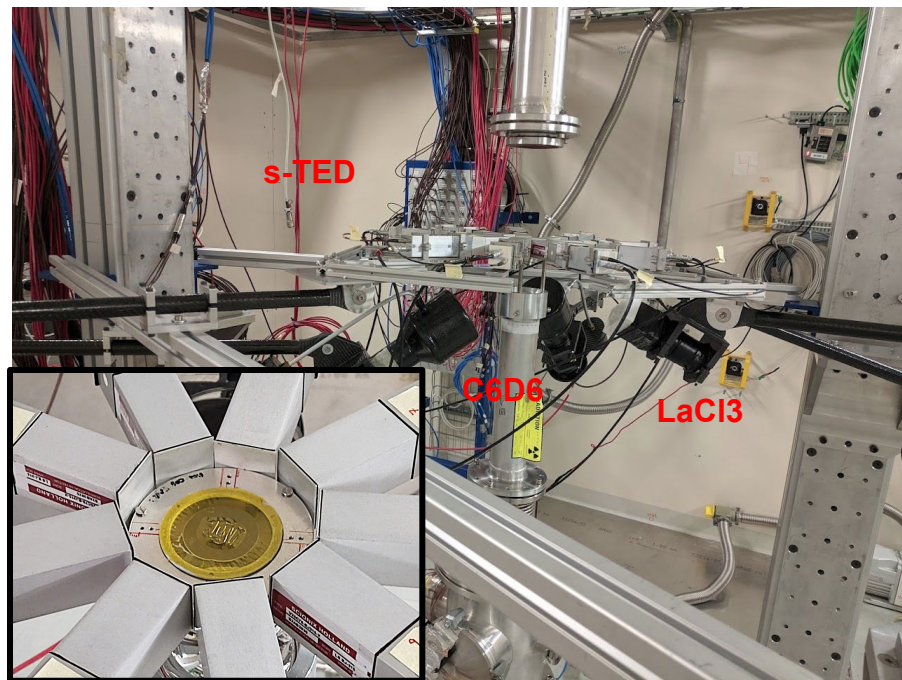
CONSEJO SUPERIOR DE INVESTIGACIONES CIENTÍFICAS



- Brief summary of prev. analysis status
- Monte Carlo simulations & dead time model
- Results of Maximum Likelihood estimation for exp. yield
- Self-shielding corrections
- Take home messages

First (n, γ) measurement in **EAR2** for **2022** campaign:

- Experiment performed between end of **March** and **April** of this year.
- Experimental setup:
 - **9 s-TEDs** in a ring configuration @ **4.5 cm**.
→Main detectors for (n, γ) (~1 L of C6D6).
 - **2 C6D6** @ **17.5 cm** with the **new PMT+VD**
→Validation.
 - **1 LaCl3** @ **9 cm**
→Spectroscopic inf. & angular distribution.
- A total of **3.2×10^{18}** protons / **3.0×10^{18}** INTC distributed in **several configurations** devoted to:
 - **Isotope** of interest
 - **Bkg** estimation
 - **Normalization** with a controlled geometry



In the last meeting:

- Pulse shape identification for all detectors



- Gain drift for all detectors along the measurement

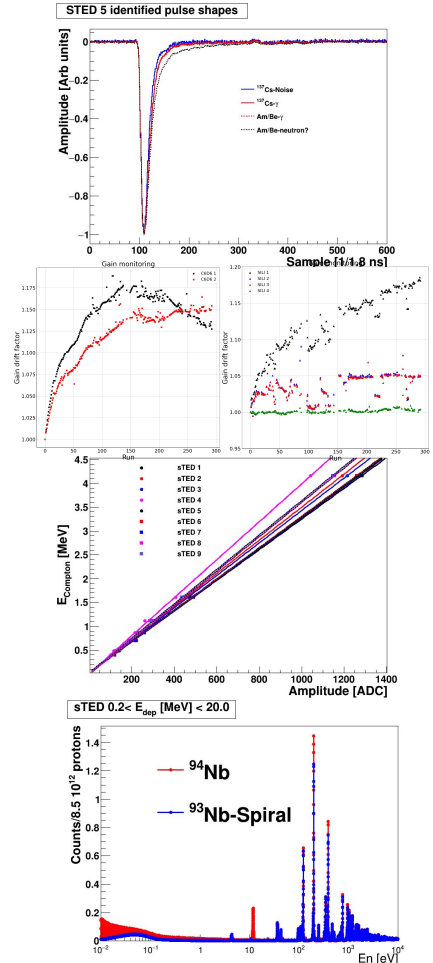


So, what is new since May?

- Detector individual energy calibration and t-flash

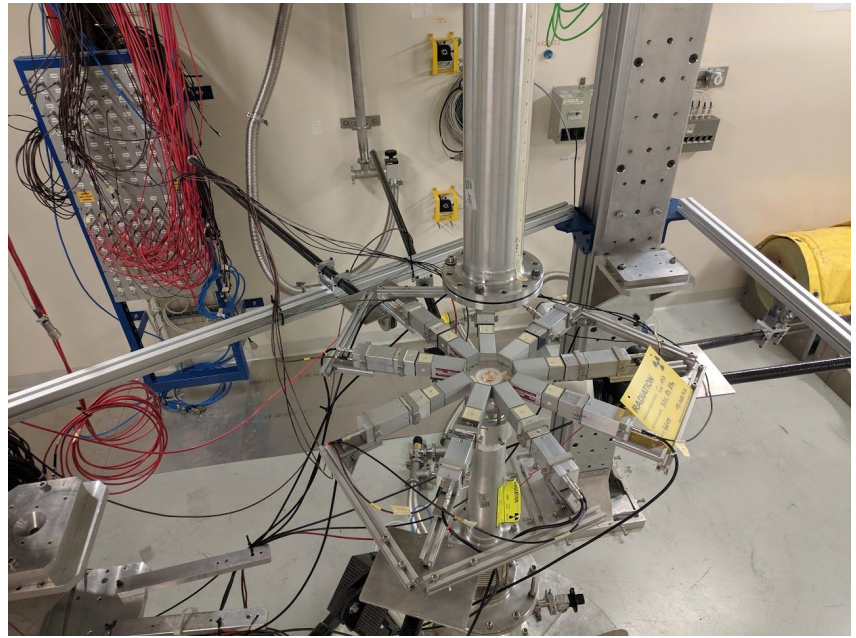
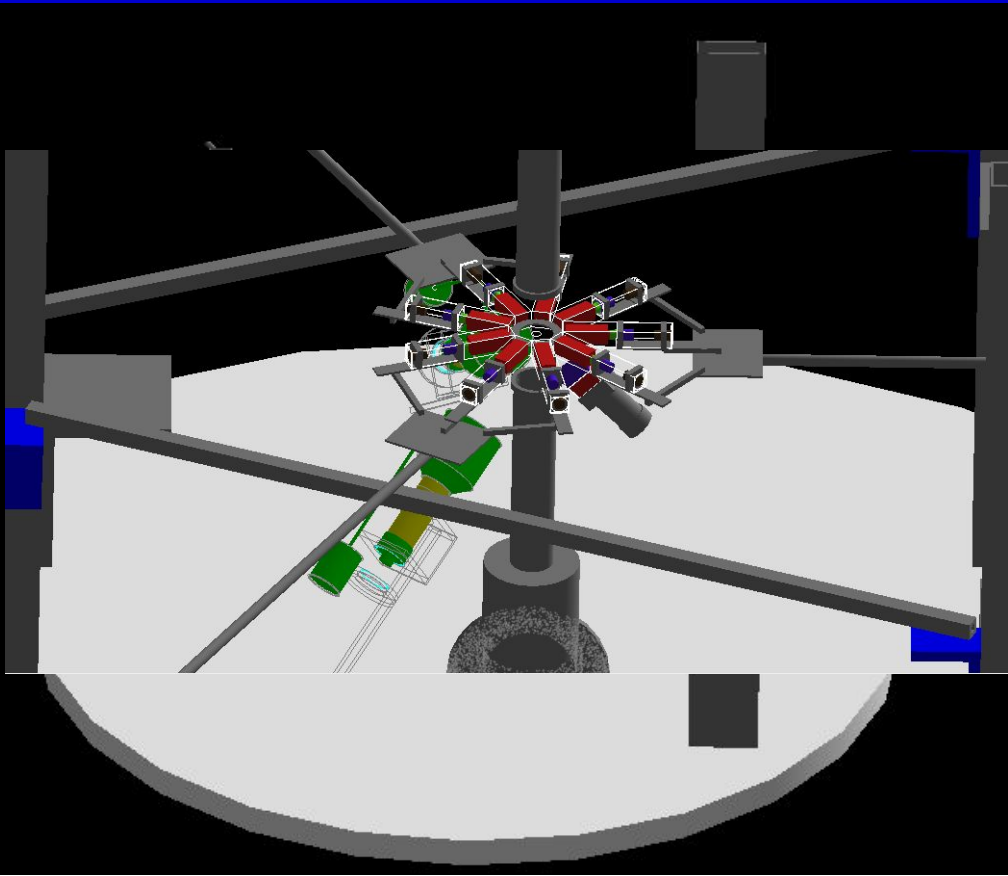


- Preliminary yield for all configurations



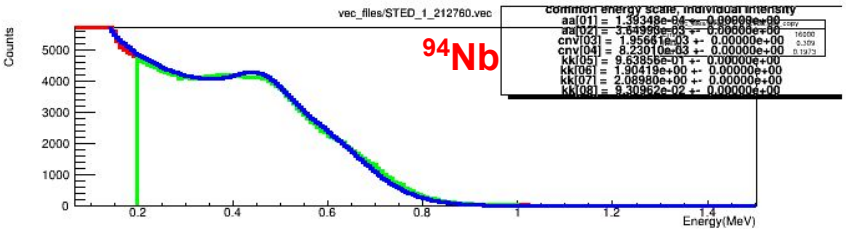
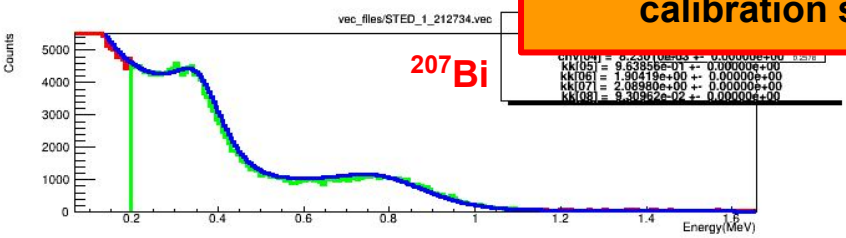
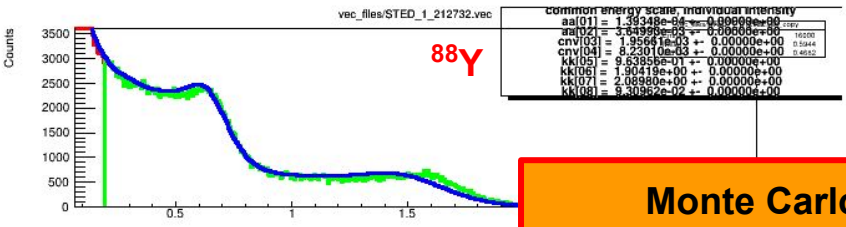
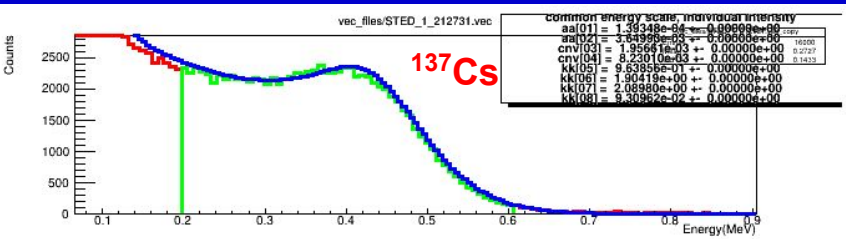
Monte Carlo & Dead time model

Monte Carlo geometry

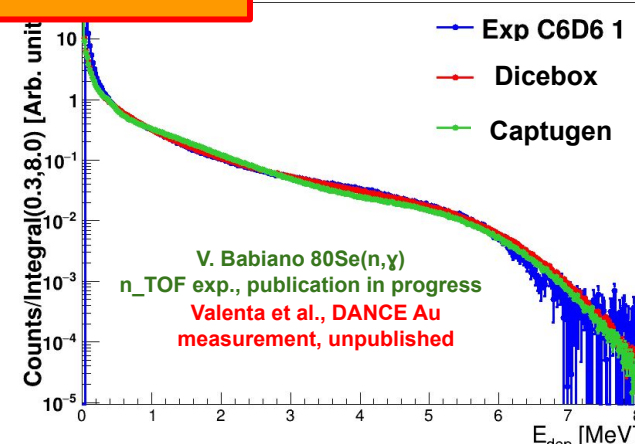
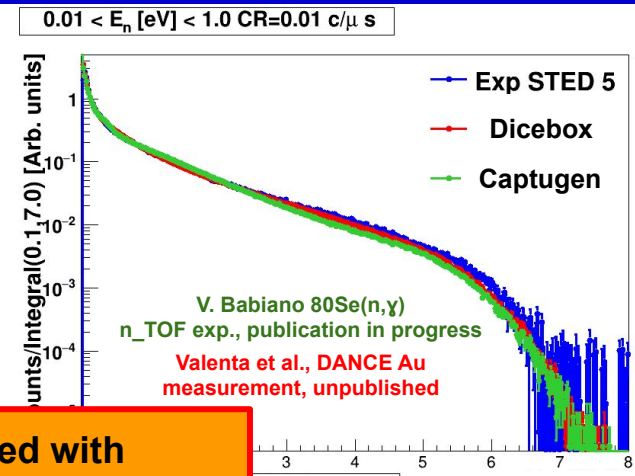


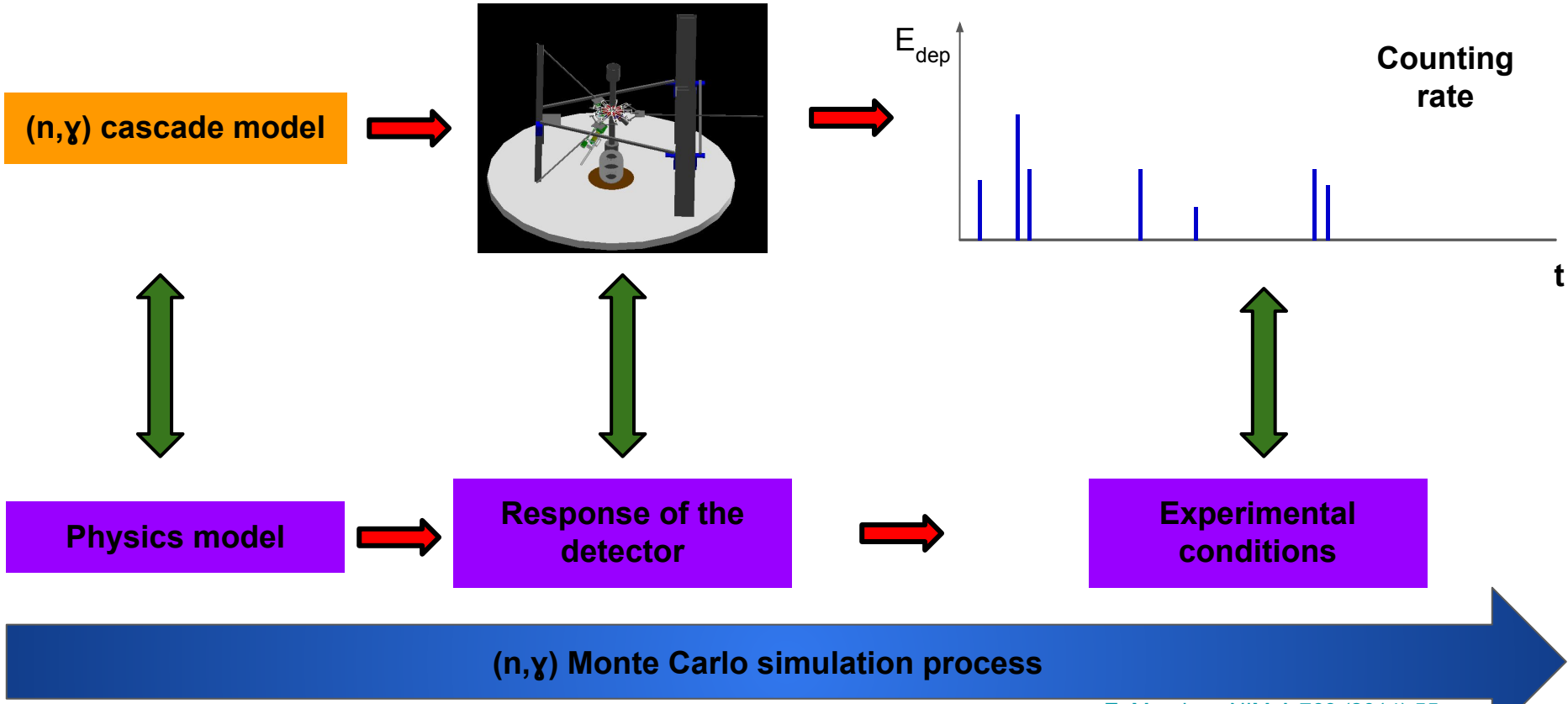
Detailed geometry of the setup implemented in GEANT4

Monte Carlo validation



Monte Carlo model validated with calibration samples and ¹⁹⁷Au(n,γ)



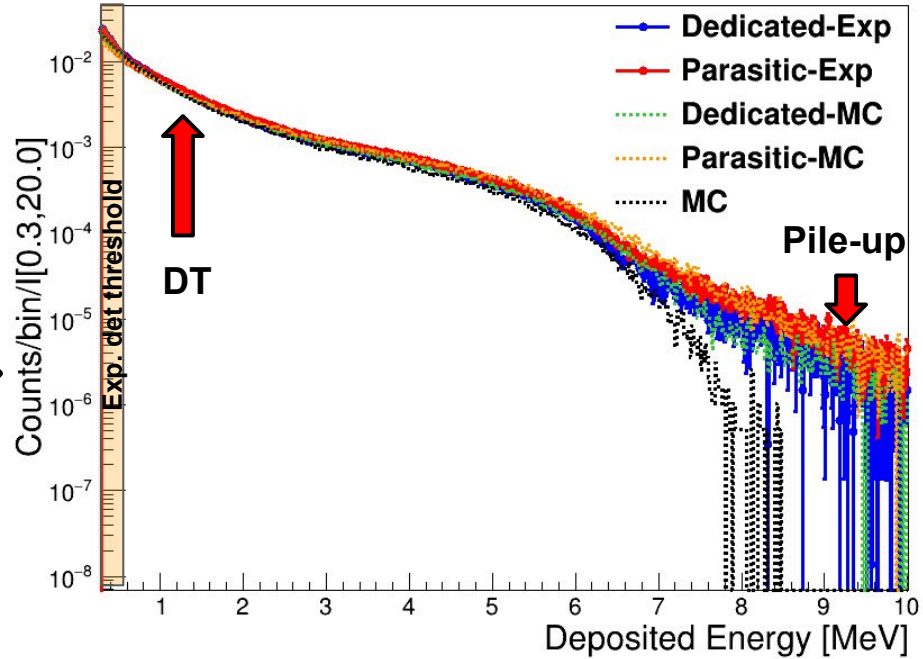
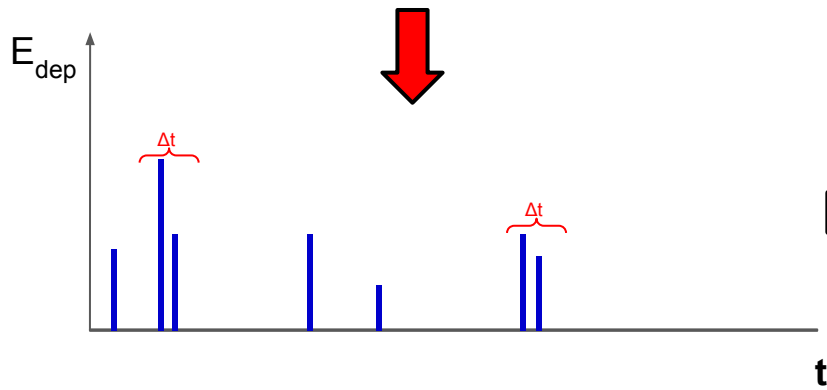


Dead time & pile-up modelization

C6D6 Deposited energy spectra at 4.8 $^{197}\text{Au}(n,\gamma)$ saturated resonance

^{79}Se -campaign

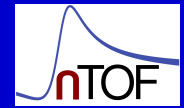
Apply dead time model



Comparing dead time & pile-up model spectra with ideal case we can get:

- Dead time corrected counting rates applying exp. detection threshold.
- Correction for pile-up+DT using a weighting function $Q=W(E_{\text{dep,DT}})/W(E_{\text{dep}})$

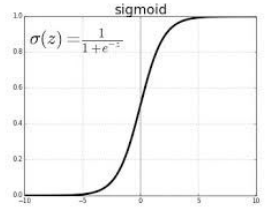
Dead time & pile-up modelization



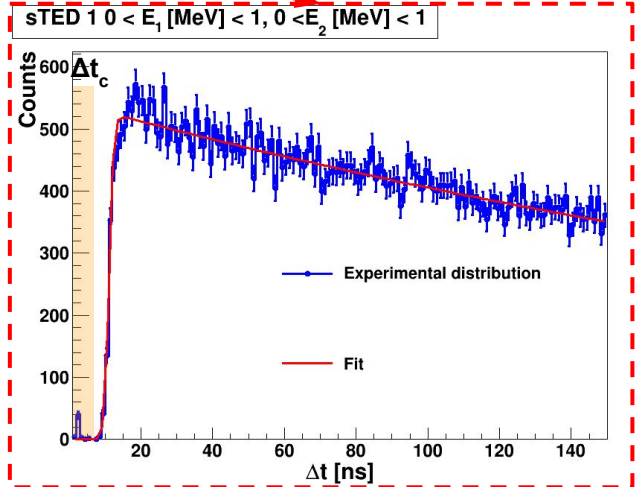
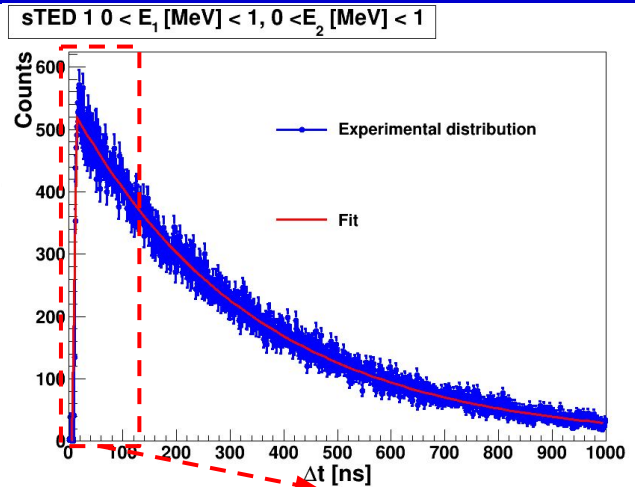
Small modification of [C. Guerrero et al.](#) dead time model:

- Dead time **depends** on **amplitude** of **consecutive** detected **signals**
- Dead-time is **characterized** by a **soft** detection probability **function**:

$$f(\Delta t) = \underbrace{\frac{1}{(1 + e^{-a(\Delta t - \Delta t_0)})}}_{\text{Detection probability}} \underbrace{Ae^{-CR\Delta t}}_{\text{Underlying distribution}}$$

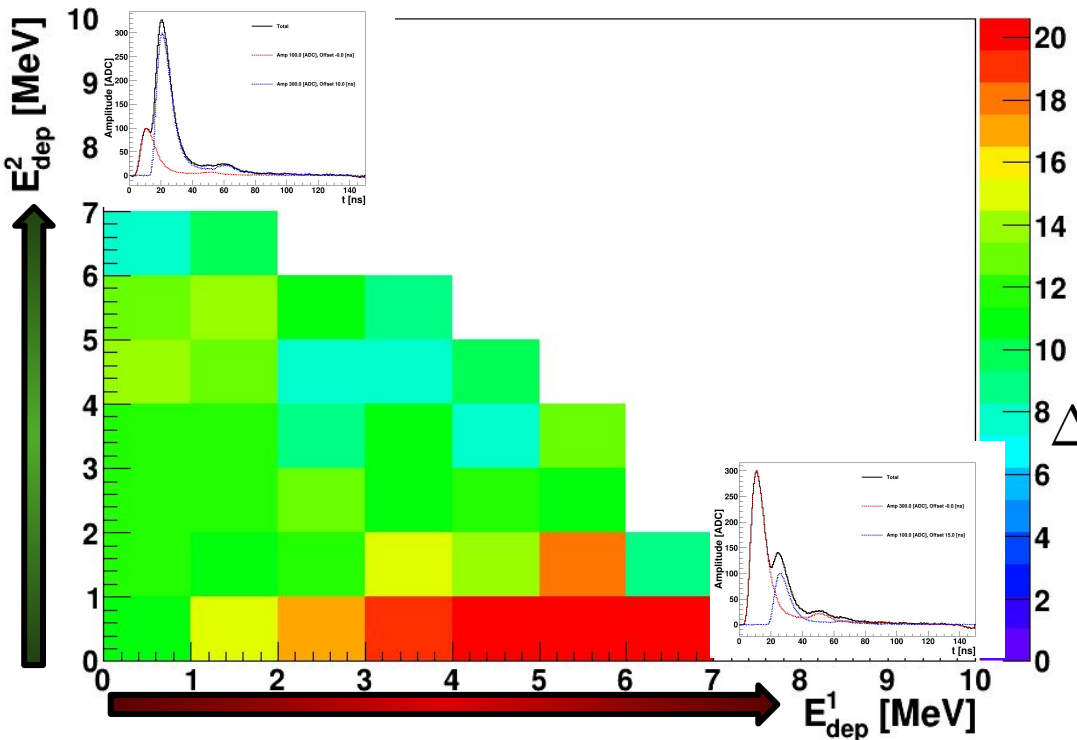


- Pile-up parameter Δt_c in the model is **not accessible** experimentally
 → **Estimated** matching experimental data in ^{197}Au **saturated resonance**.



Dead time function parametrization

Δt_0 [ns] values distribution for sTED 3 @ 5.0 eV



Dead time parameters for individual detectors fitted as a function of signals energy @ 4.9 ¹⁹⁷Au(n,γ) resonance:

$$f(\Delta t) = \frac{1}{(1 + e^{-a(\Delta t - \Delta t_0)})} A e^{-CR\Delta t}$$

$$\Delta t_0 = \Delta' t_0 + \Delta' t_1 E_{dep}^1 + \Delta' t_2 E_{dep}^2$$

50% detection chance

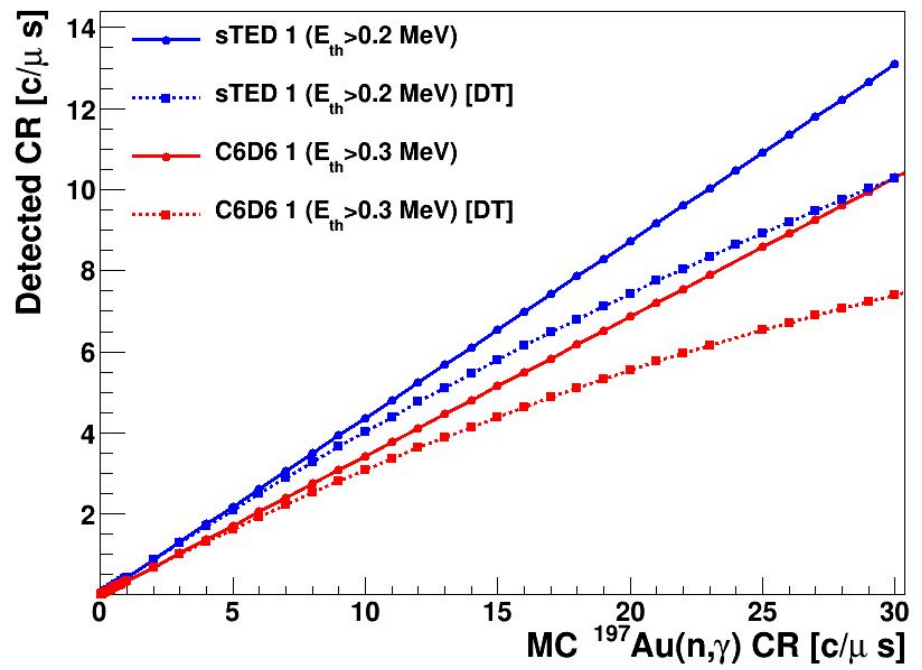
$$a = a_0 + a_1 E_{dep}^1 + a_2 E_{dep}^2$$

Probability change rate

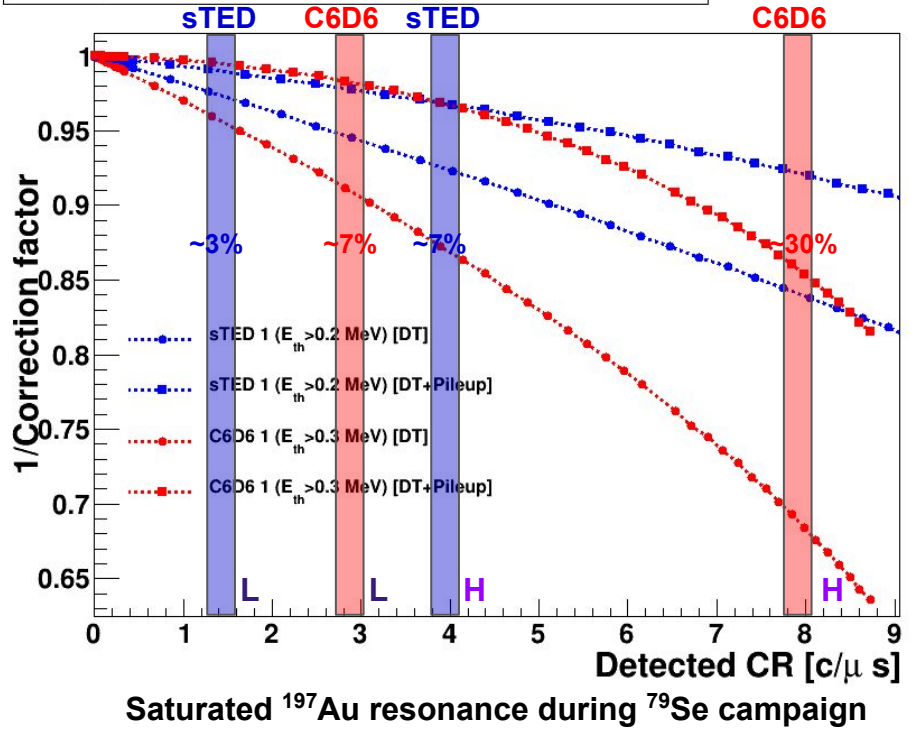
General rule: $E_{dep}^2 > E_{dep}^1 \rightarrow$ Easier to detect the second signal

Dead time & pile-up modelization

Comparison between sTED and C6D6 detectors



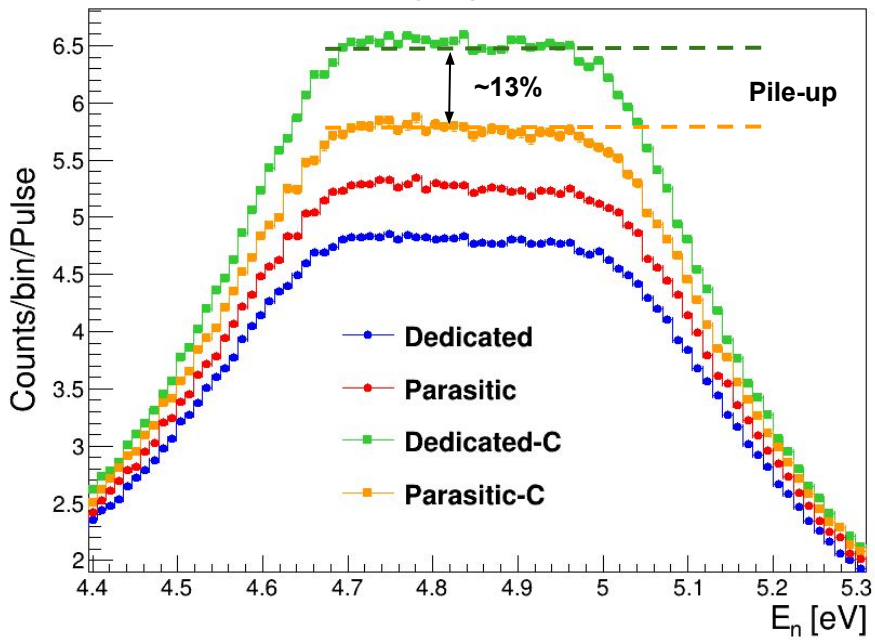
Comparison between sTED and C6D6 detectors



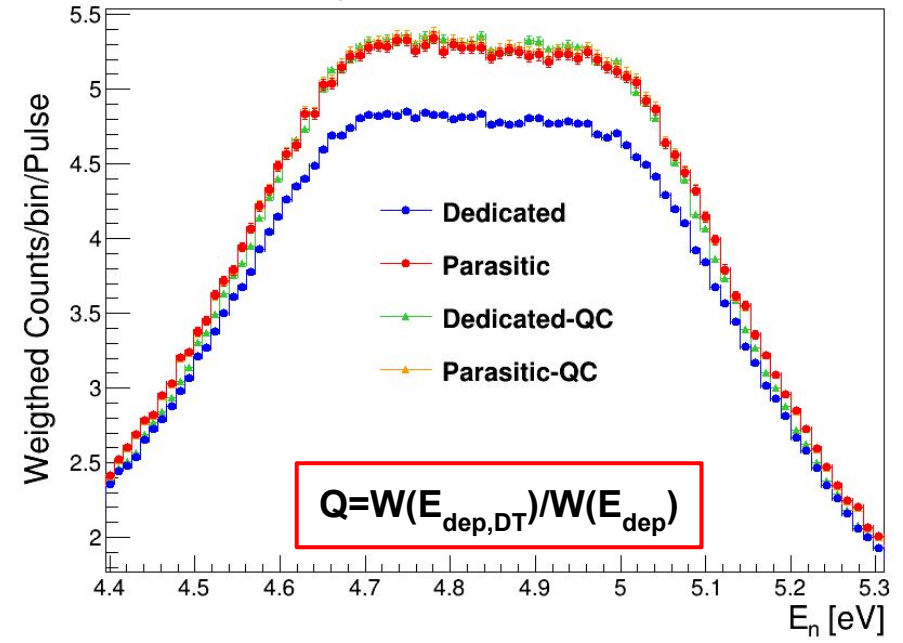
Dead time corrections can be large even for s-TEDs!

H: High intensity pulses
L: Low intensity pulses

Corrected only by dead time effect



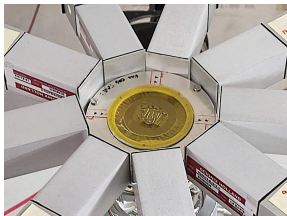
Corrected by dead time and pile-up effects



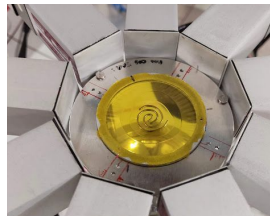
Correcting only by dead time is not enough: in addition on has a large pile-up effect which requires further corrections!

Results of ML estimation

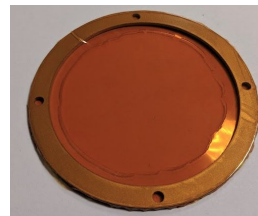
We have implemented a pulse by pulse maximum likelihood estimation for ^{93}Nb , ^{94}Nb and **Empty** contributions simultaneously using the three configurations:



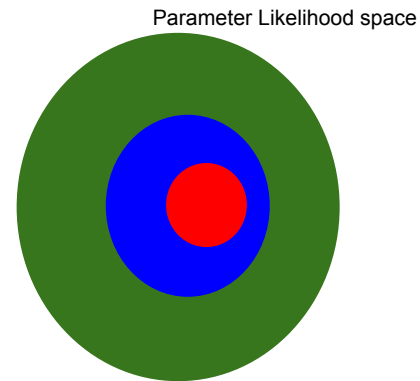
^{94}Nb , ^{93}Nb & **Empty**



^{93}Nb & **Empty**



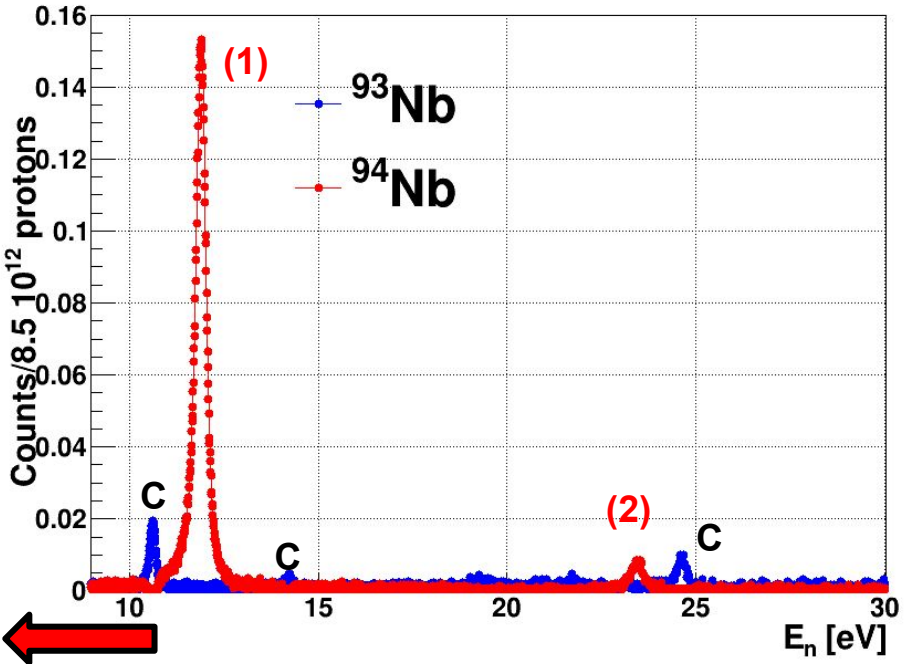
Empty



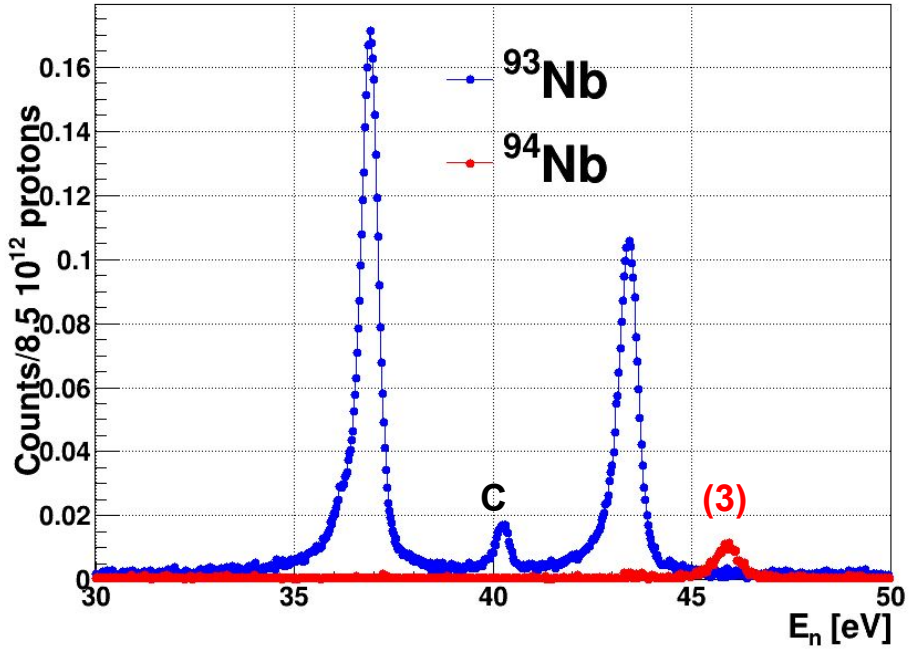
- **Maximum Likelihood estimation has smallest variance** → **Larger sensitivity** for small “signal”
- Allow **easily hypothesis contrast** via likelihood-ratio → Well behaved for **nested model** ($^{94}\text{Nb}=0$)
- **Registered counts** by detection systems in any ToF period **easily model** → $x \sim \frac{\lambda^x e^{-\lambda}}{x!}$
- **Systematics** (Norm. between configurations, Unc in no-beam components) included via **Monte-Carlo**

Next slides: Only s-TEDs included & confidence level >0.95

Maximum likelihood results

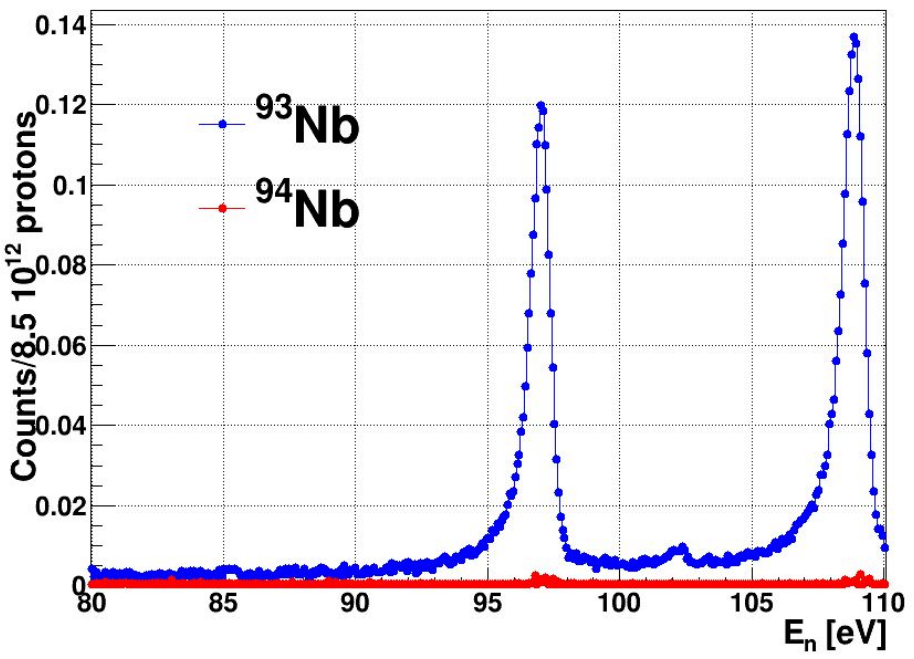


No enough sensitivity for ^{94}Nb thermal region because of large contribution from target decay

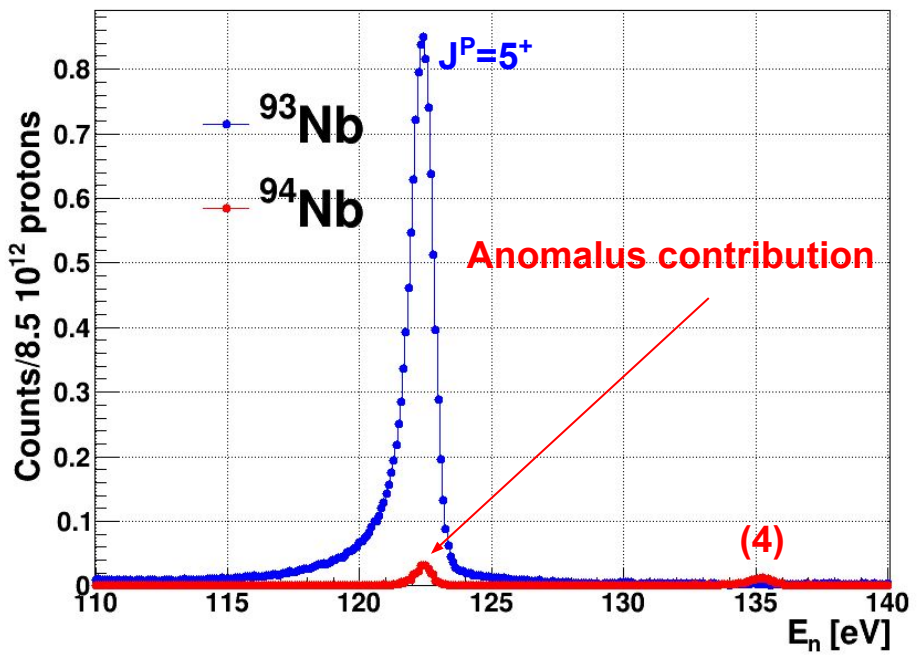


The contaminant resonances are only for ^{93}Nb spiral target and they do not overlap with “good resonances”

Maximum likelihood results

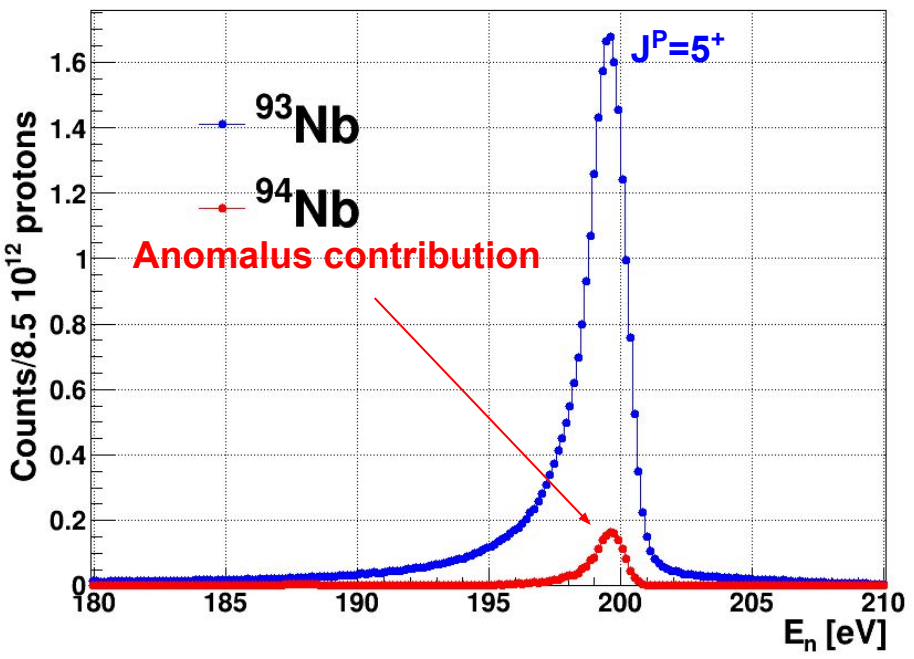


The calculated yields are not calibrated in ToF

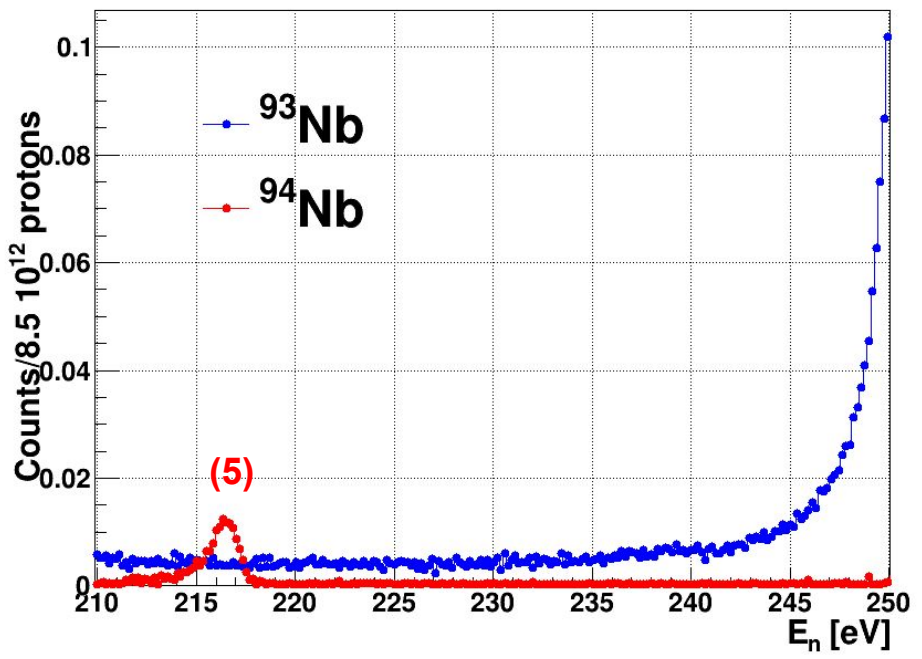


Anomalous contribution from " ^{94}Nb " to ^{93}Nb resonance
 → Related to sample geometry (Next slides)

Maximum likelihood results

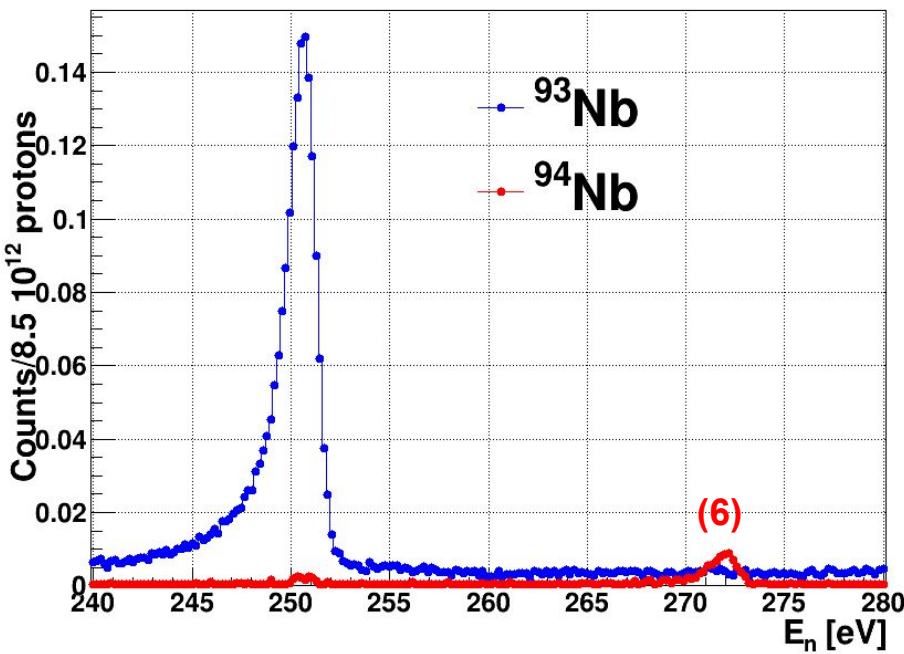


Another contribution from “ ^{94}Nb ” to same spin-parity ^{93}Nb

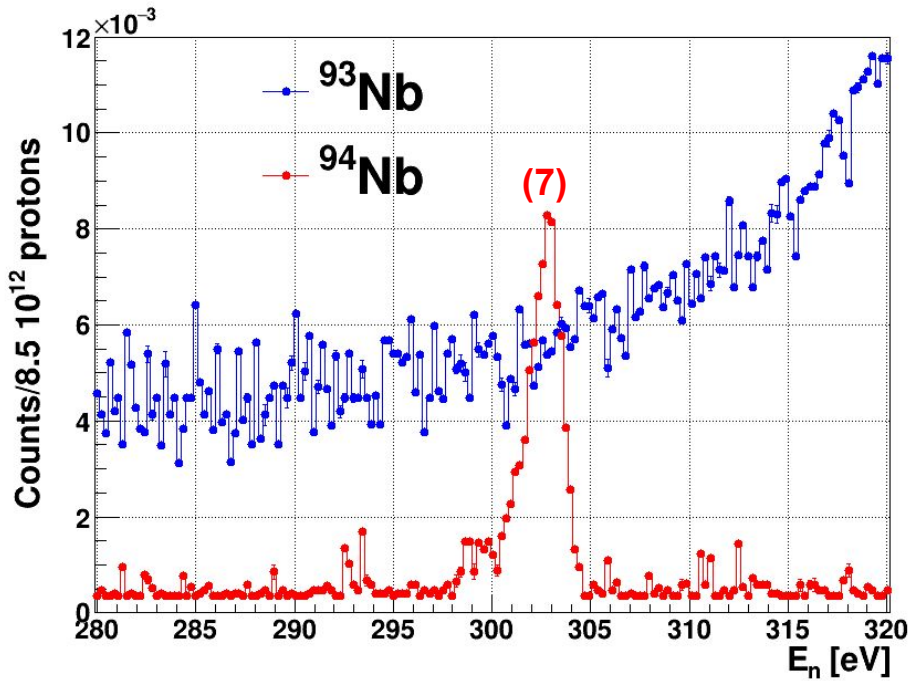


Good ^{94}Nb resonance definition

Maximum likelihood results

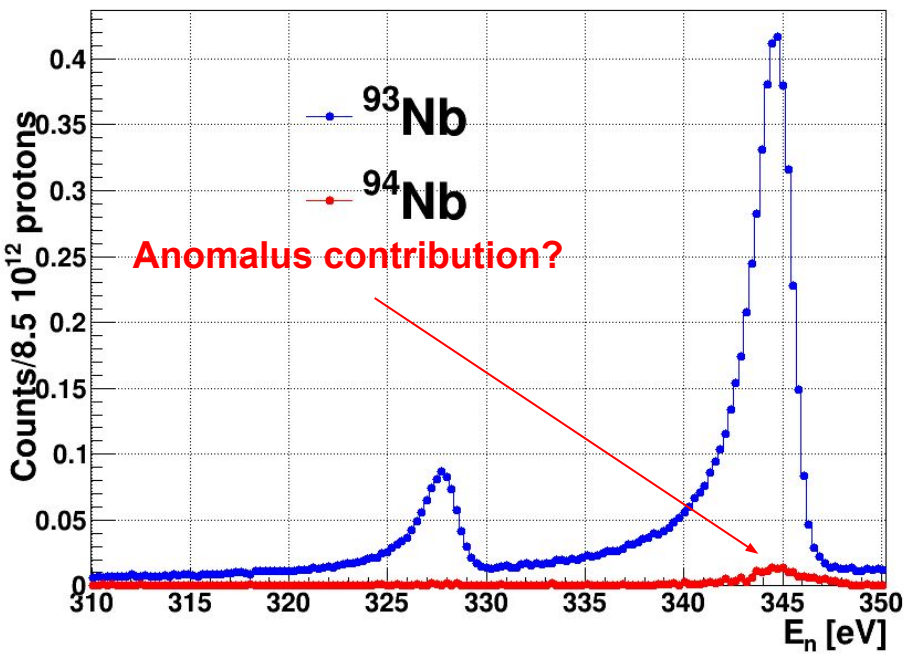


Resonances of ^{94}Nb appears to be so “small” because of the $^{94}\text{Nb}/^{93}\text{Nb}$ ratio $\sim 3\%$.

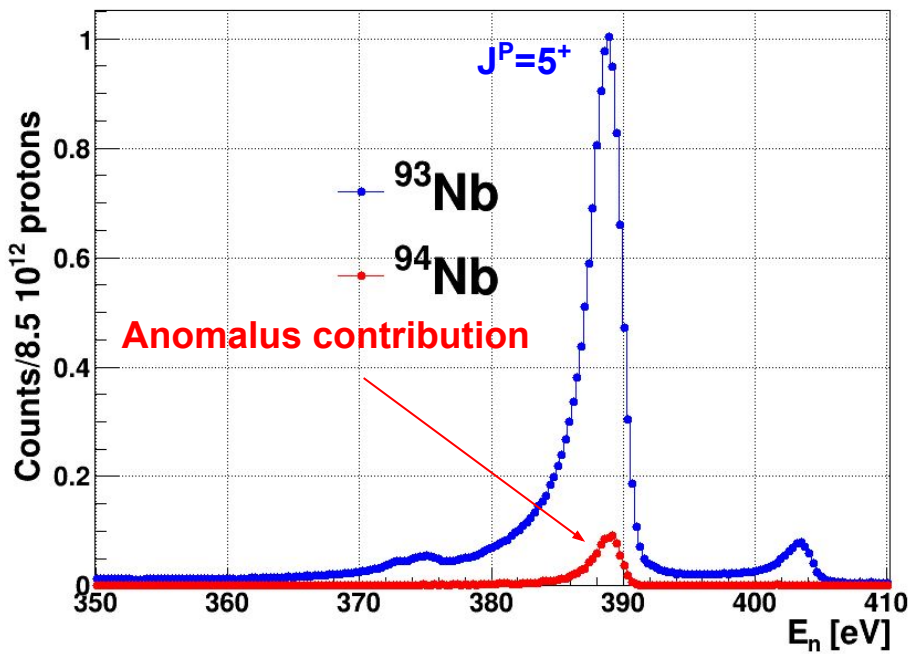


^{94}Nb resonance in a tail of a “big” ^{93}Nb resonance

Maximum likelihood results

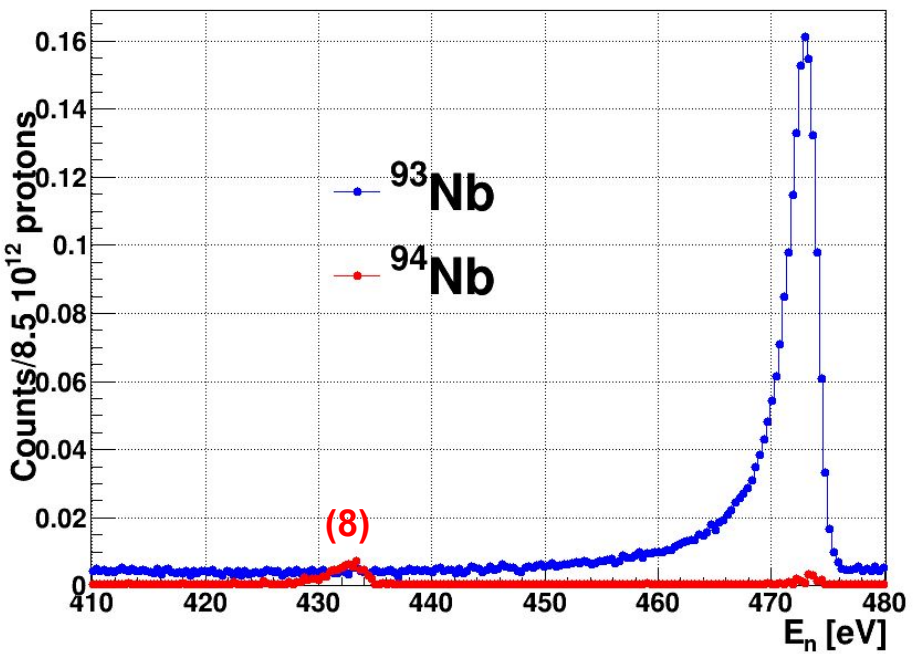


In this case, the anomalous contribution has a weird shape compared to other well-defined shapes

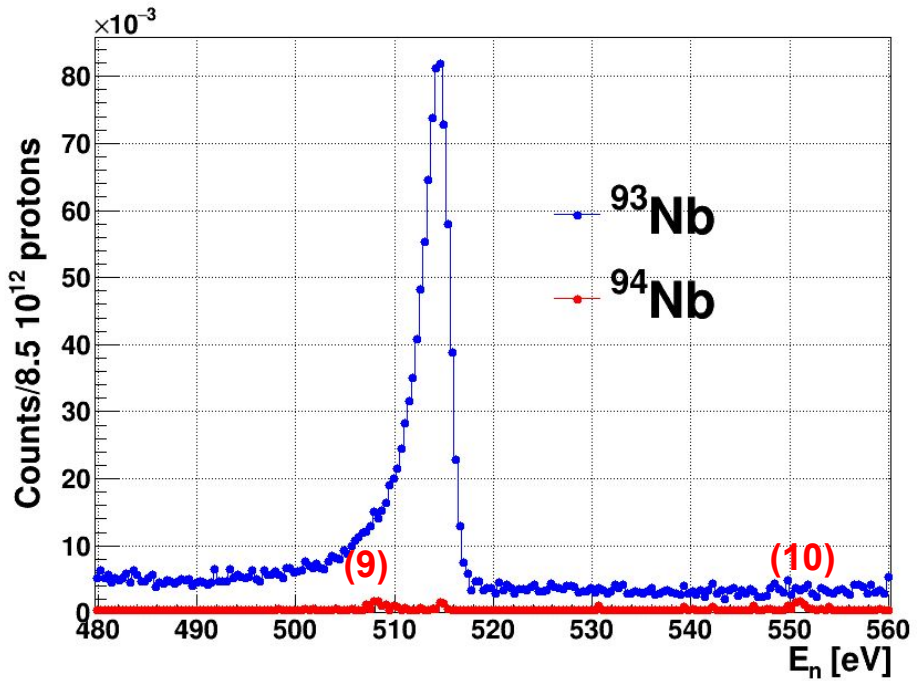


Another contribution to the same spin-parity resonance

Maximum likelihood results

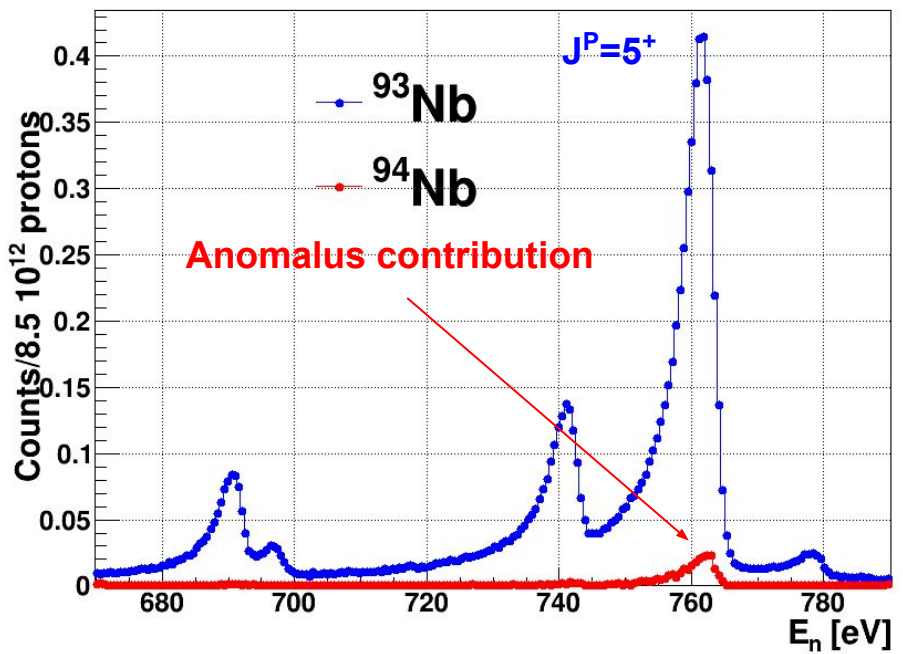
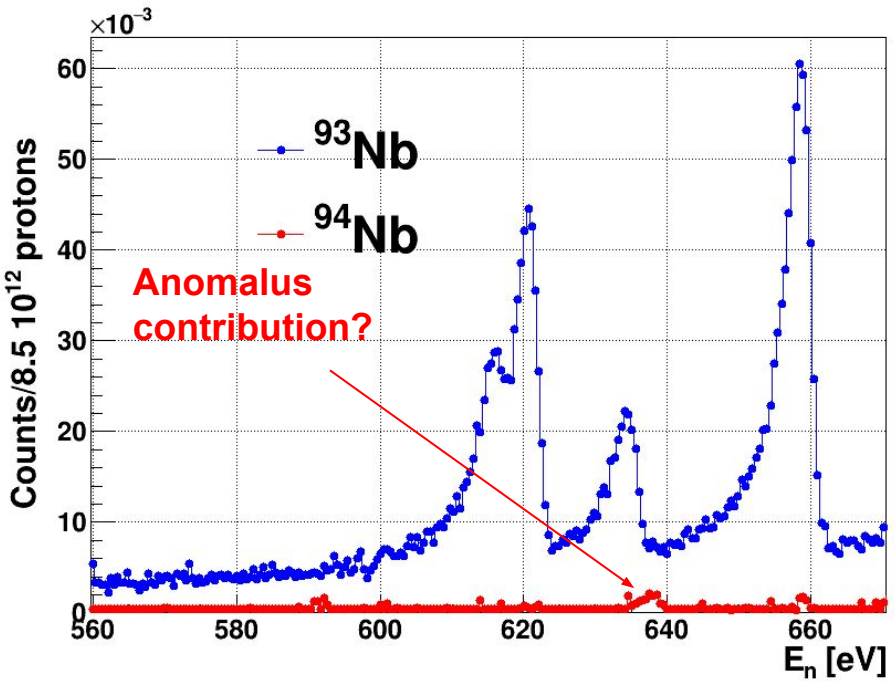


Small, but well located resonance

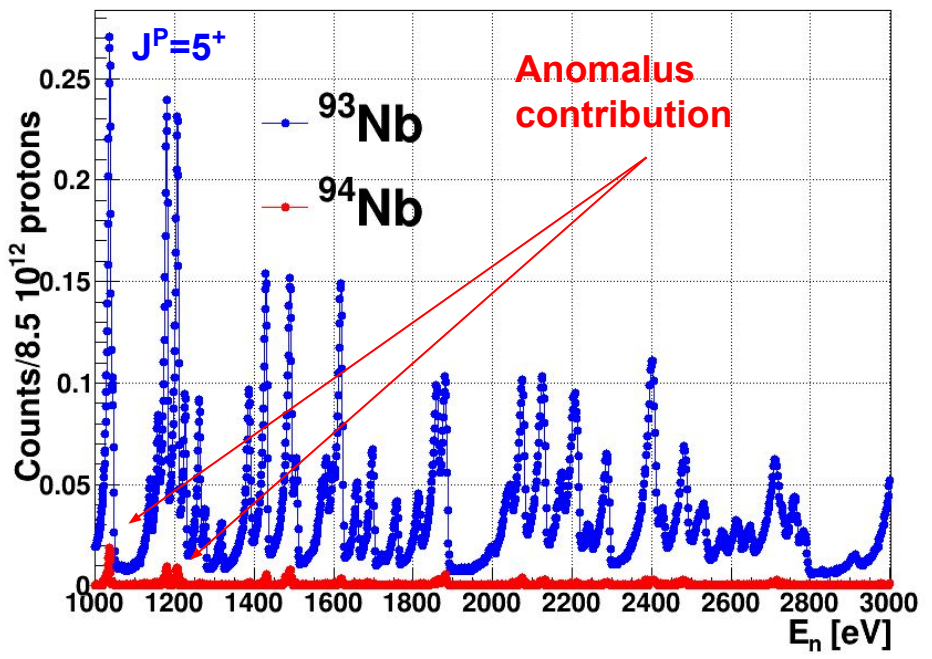
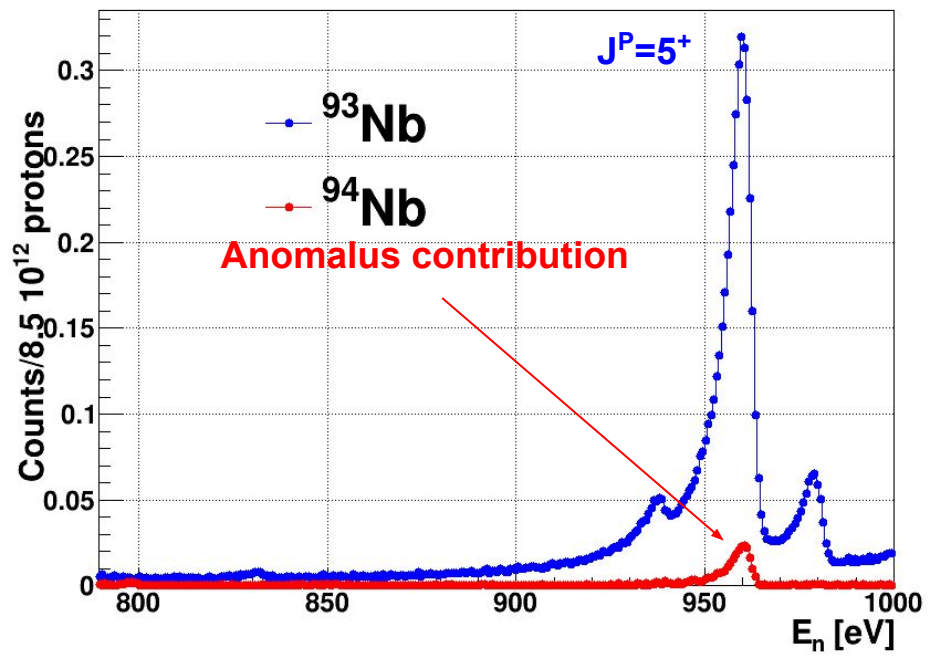


(9) and (10) are not a fluctuation to a confidence level > 95%

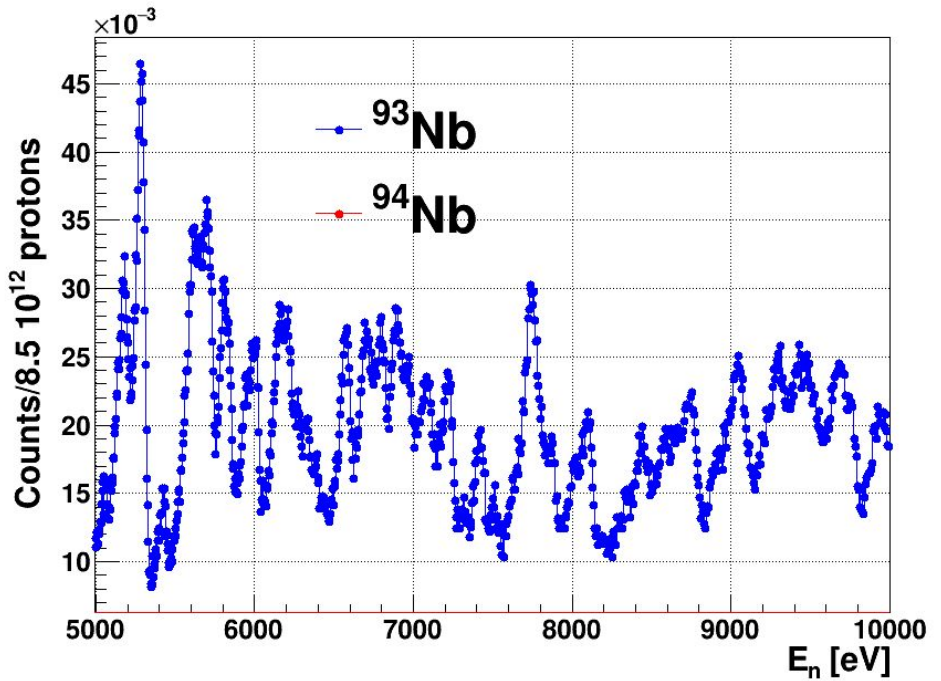
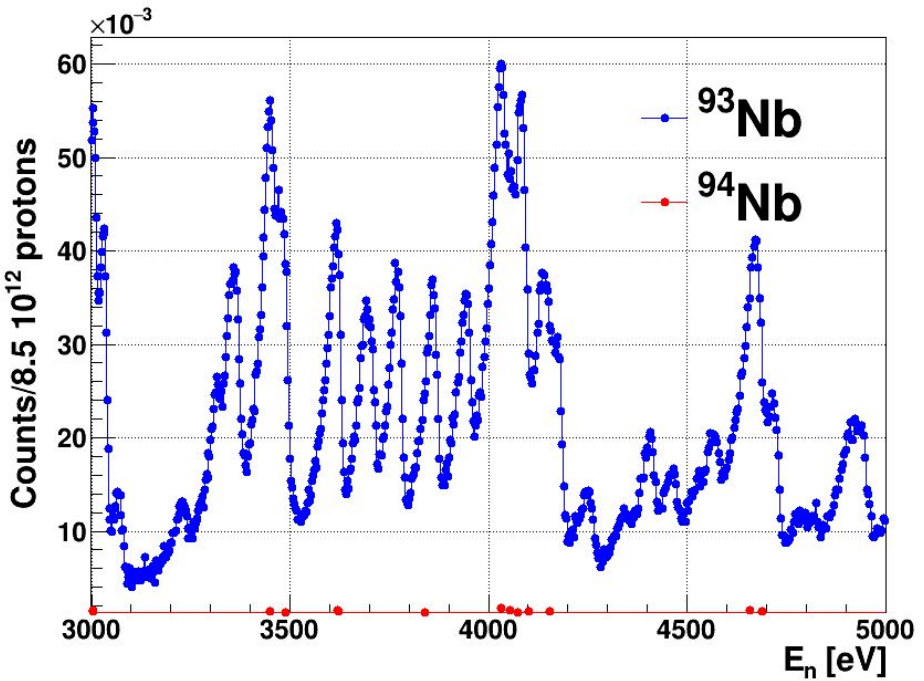
Maximum likelihood results



Maximum likelihood results

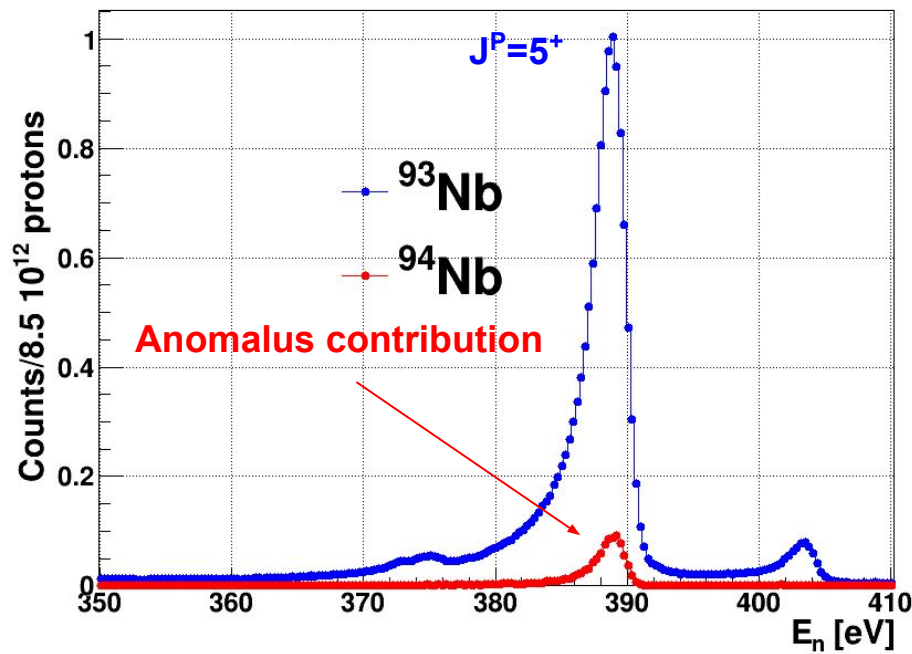


Maximum likelihood results



Beyond $E_n = 3$ keV: not enough sensitivity to extract any additional information from $^{94}\text{Nb}(n,\gamma)$

Self-shielding corrections



Anomalous contribution is clearly visible for a very well characteristic neutron resonances:

- $J^P=5^+$
- $\Gamma_n \gg \Gamma_\gamma \gg$
- $\Gamma_n/\Gamma_\gamma \gg$

Hypothesis we worked on:

- ^{93}Nb isomer contribution
 ^{94}Nb target: $^{93}\text{Nb} \rightarrow ^{94}\text{Nb} + ^{93*}\text{Nb} \rightarrow ^{94}\text{Nb}$
 ^{93}Nb target: $^{93}\text{Nb} \rightarrow ^{94}\text{Nb}$

Not possible due to angular momentum

- Direct neutron sensitivity contribution
 $\Gamma_n/\Gamma_\gamma \gg$

Not possible because of $\Gamma_n/\Gamma_\gamma \sim 3$

- Self-shielding \rightarrow Wires' diameter wrong?

Neutron flux simulation:

- $350 < E_n \text{ [eV]} < 400$ flat distribution
- $\sigma_x = 1.5 \text{ mm}, \sigma_y = 1.5 \text{ mm}$

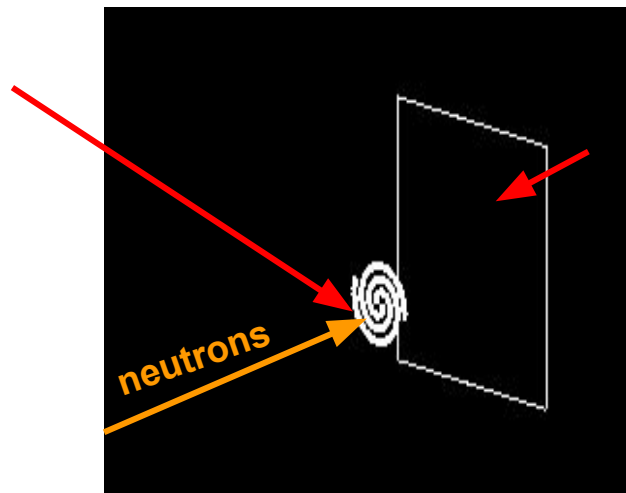
Targets:

- **geometries and thickness according to targets in the experiment**

Two sensitive detectors:

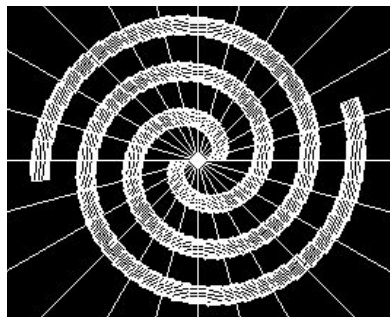
- **Sample** itself for (n,γ)
- neutron flux monitor for **transmission**

target

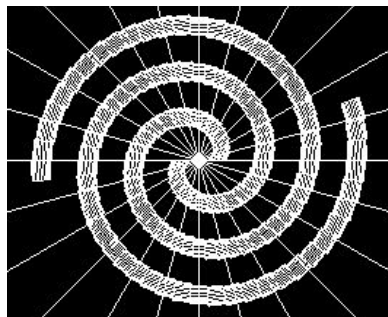


neutron flux monitor

diameter= 1.0 mm



diameter= 0.8 mm



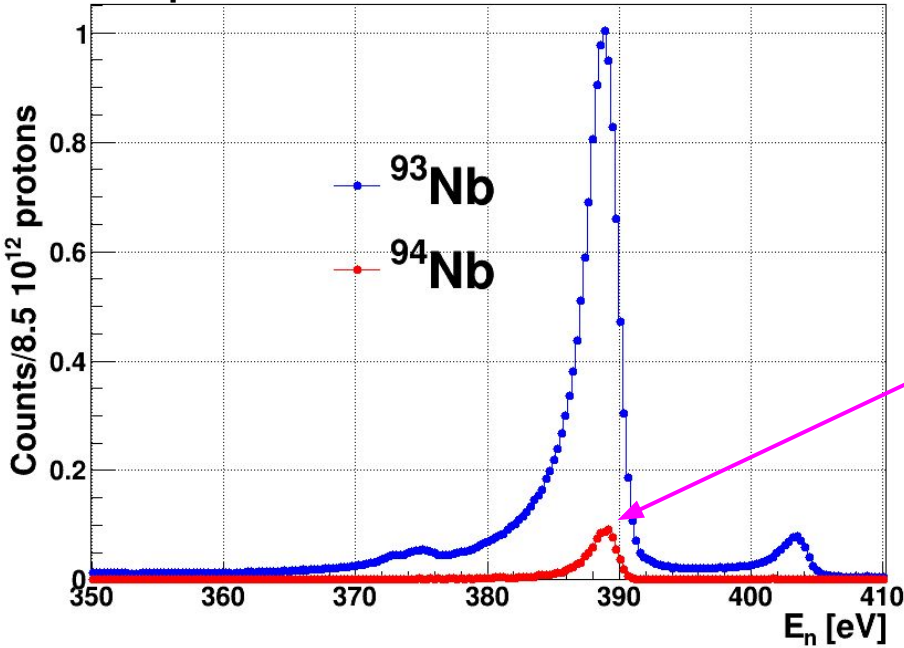
thickness= 2.1 mm



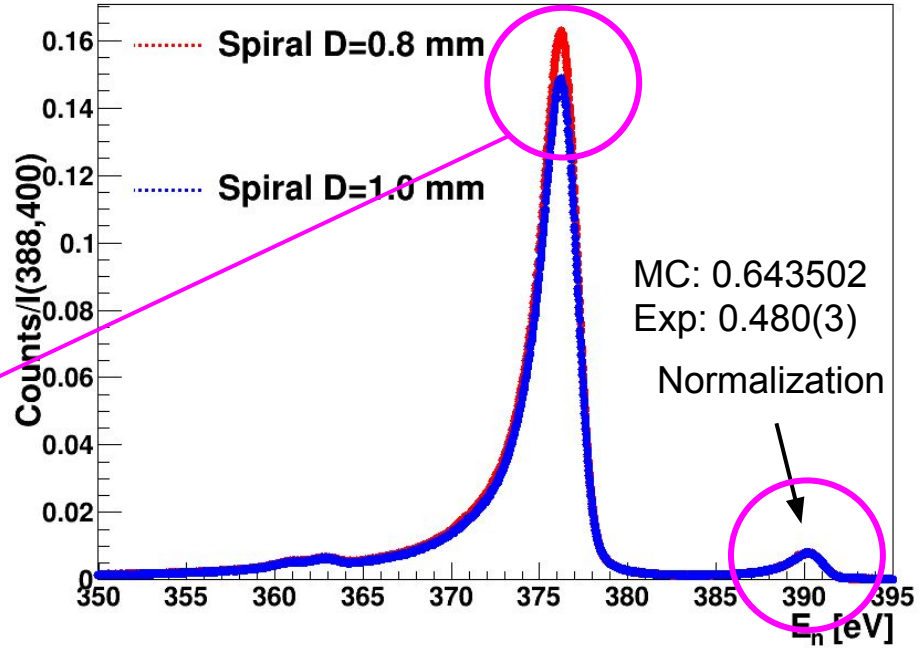
thickness= 0.02 mm



Experimental data obtained from ML



Monte Carlo data from simulations*RF

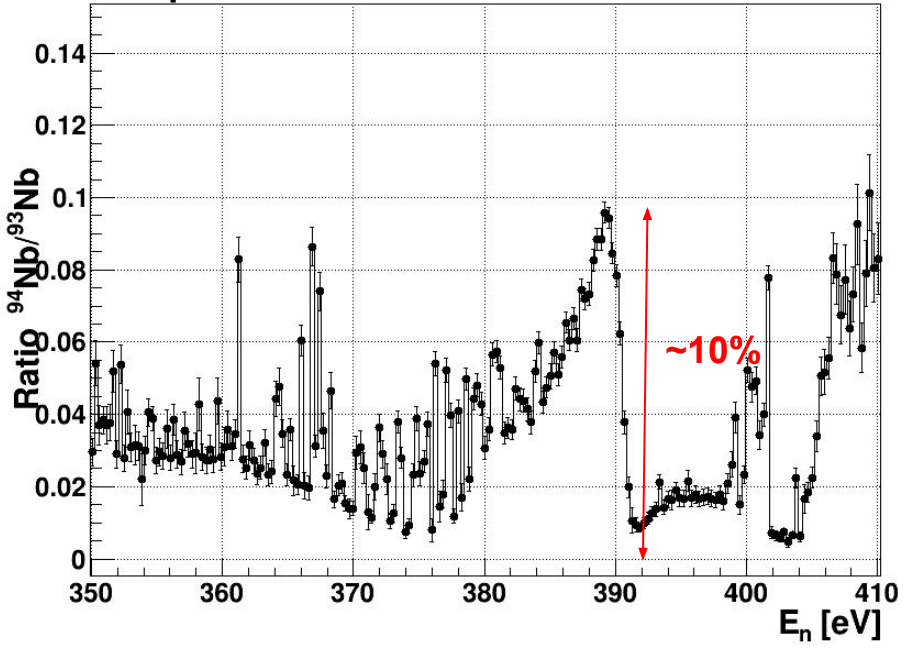


- The experimental data is not yet calibrated in distance of ToF
- **Monte Carlo** data shows the **same excess of counts** in the largest resonance, once normalized to **~400 eV** resonance

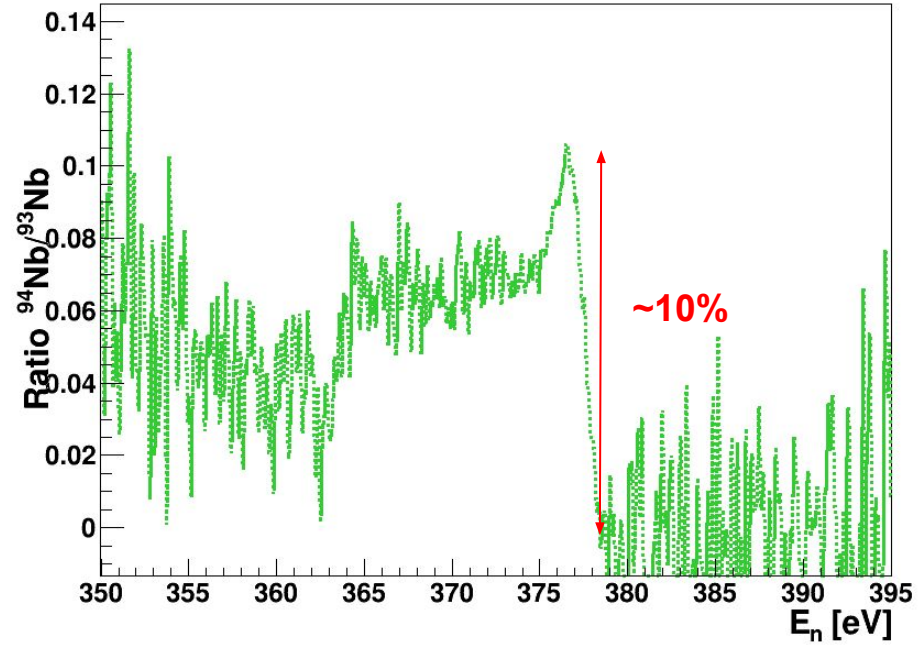
Excess of counts because of different wire diameter for ⁹³Nb and ⁹⁴Nb targets

Exp. & MC results

Experimental data obtained from ML



Monte Carlo data from simulations*RF



Reproduction of the effect is ok!

Work in progress for such complicated geometry!

EAR2 is not EAR1 !

DT and P-U corrections can be severe even for s-TEDs

The analysis of $^{94}\text{Nb}(n,\gamma)$ is in progress:

- Monte Carlo model implemented in GEANT4
- Dead time and pile model
- Maximum likelihood methodology applied to Exp. yield
- Self-shielding correction is ongoing

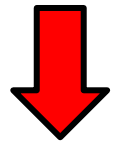
Thank you very much for your attention!



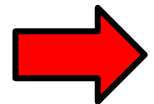
BACKUP

For each neutron bin energy (or ToF) we registered counts in the detector and proton intensity in a proton bunch:

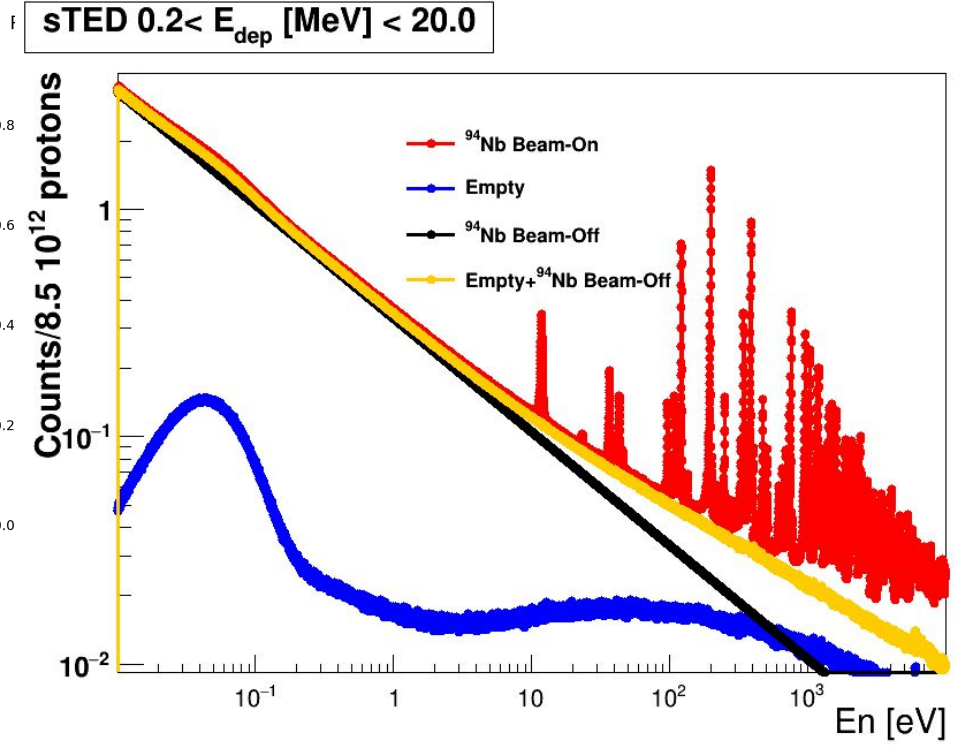
Protons 10^{-12}	Counts
8.308993	3
8.3541387	1
8.3149104	2
8.251564	3



Usually we accumulate statistics and histograms and use Gaussian statistics



In an effective way we are averaging the results:
 → Is there another way to do it?

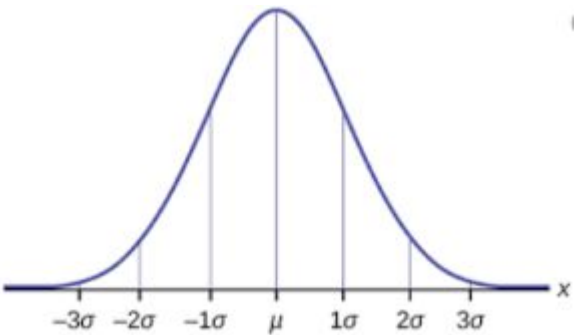


3
y
8.8

Bayes' theorem states (or conditional probability of the model constrained to experimental data):

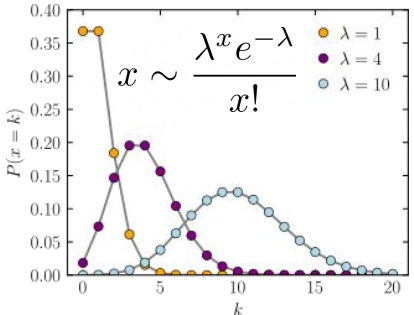
$$P(\text{theory} | \text{data}) \propto P(\text{data} | \text{theory}) \times P(\text{theory})$$

Posterior ↔ Conditional probability



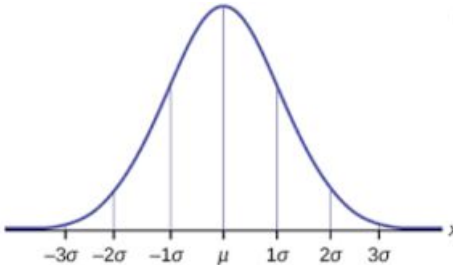
After experimental data

Likelihood ↔ Conditional probability



data modeling

Prior ↔ Probability



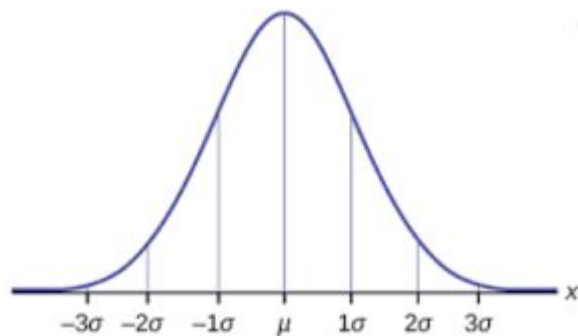
Prior experimental data



Bayes' theorem states (or conditional probability of the model constrained to experimental data):

$$P(\text{theory} | \text{data}) \propto P(\text{data} | \text{theory}) \times P(\text{theory})$$

Posterior \leftrightarrow Conditional probability



After experimental data

Asymptotic limit ($n \rightarrow \infty$):

- Prior does not matter \leftrightarrow Maximum Likelihood Estimation

$$l^* = \operatorname{argmax}_{\vec{\theta}} l(\vec{x} | \vec{\theta})$$

- Consistency: $\hat{\theta} \rightarrow \theta_*$
- Asymptotically Normal: $\frac{\hat{\theta} - \theta_*}{\hat{se}} \sim N(0, 1)$
- Asymptotically optimal of efficient ! \rightarrow Smallest variance
- Standard error can be computed using Fisher information

$$\hat{se} = \sqrt{1/I(\theta)} \quad I(\theta) = E_{\theta} \left(\frac{\partial^2 l}{\partial \theta^2} \right)_{\theta=\hat{\theta}}$$

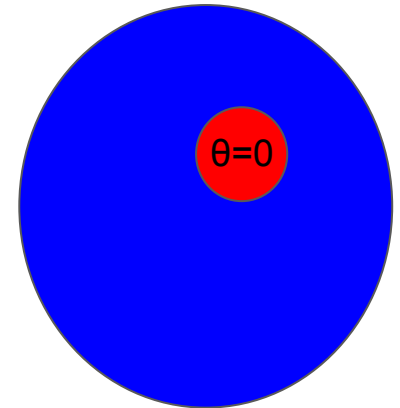
Together with the MLE it is a good good practice to include hypothesis contrast. The one we use in this work is the likelihood-ratio (Wilk-theorem):

$$\lambda_{LR} = -2 \ln \left[\frac{\sup_{\theta \in \Theta_0} \mathcal{L}(\theta)}{\sup_{\theta \in \Theta} \mathcal{L}(\theta)} \right]$$

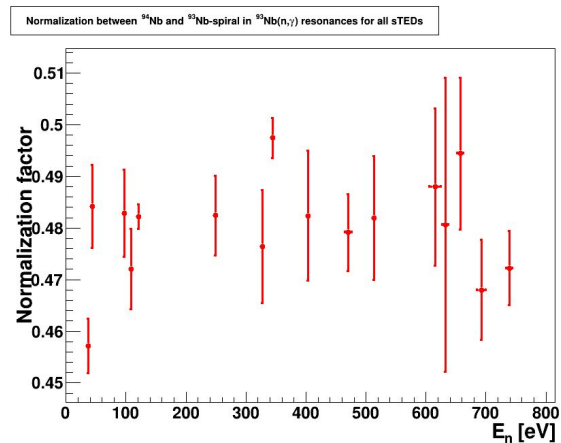
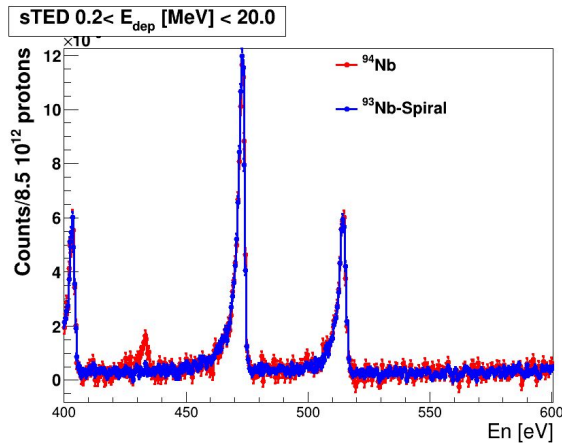
The likelihood-ratio assesses the goodness of fit of two competing statistical models based on the ratio of their likelihoods:

- It is especially well suited for nested hypothesis ($\theta=0$)
- It behaves as a χ_n where n is the number of parameters under test.
- It can be set a level of confidence α to contrast the hypothesis

Parameter space



- Normalization of Empty (γ): Shared in ^{94}Nb -target and ^{93}N -spiral, 1.10(5).
- Normalization of ^{93}N in ^{94}Nb -target because of different beam-intersection factor η :
 - Calculated as the ratio of integrals for 3⁻ and 4⁺ resonances 0.480(3).



- No beam background calculated from its configuration: 1.472(9) c/pulse
- No beam background with ^{94}Nb -target in place: 5785.5(2) c/pulse