

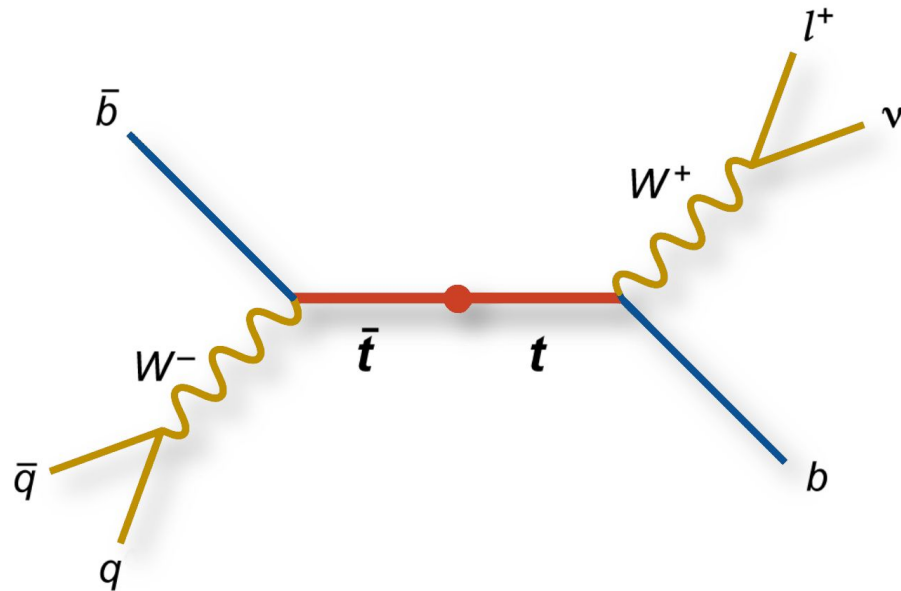


Top quark mass measurement in lepton and jets final state at the CMS experiment

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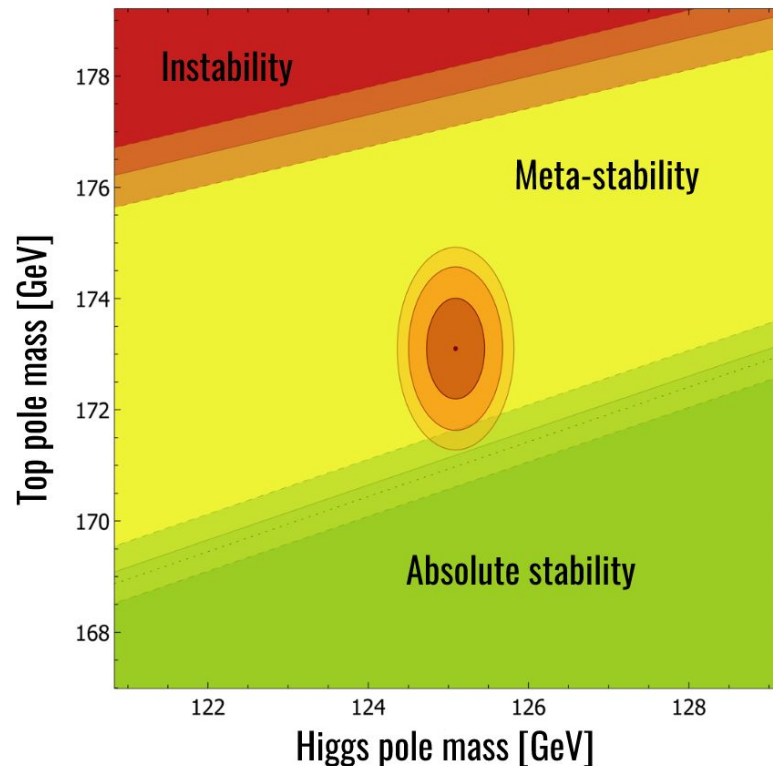
Overview

- Motivation
- Latest results
- Analysis strategy
 - event topology and the semileptonic channel
 - kinematic fit
 - binned profile likelihood method
- Future and current challenges
- Summary



Motivation

- top quark is the most massive particle in the SM
 - important probe for SM or ingredient for BSM models
- stability of the electroweak vacuum parameterized often as a function of the pole masses of Higgs and top
- “To rule out absolute stability to 3σ confidence, the uncertainty on the top quark pole mass would have to be pushed below 250 MeV” [1]



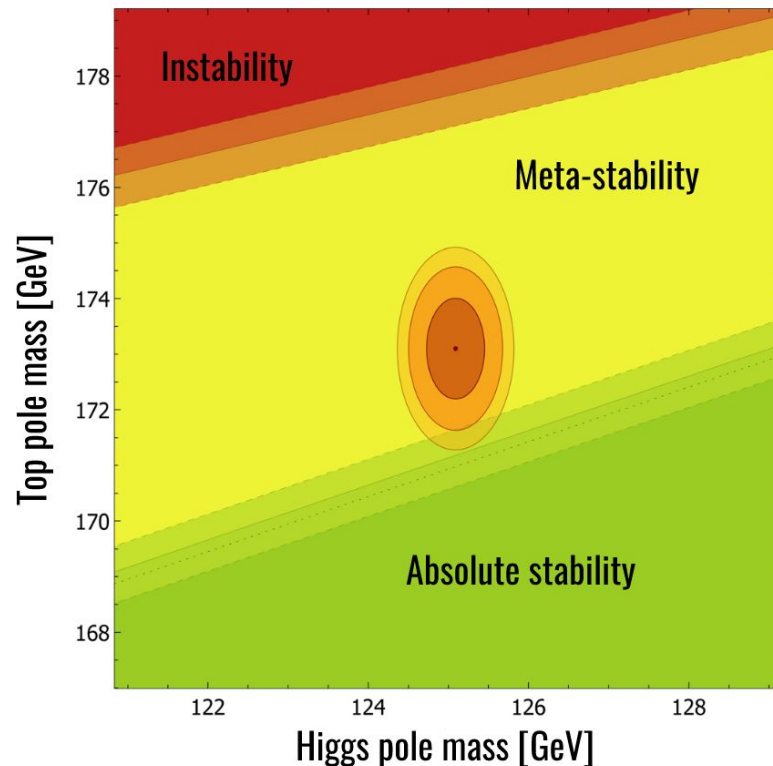
EW vacuum stability illustration [1]

[1] [arxiv:1707.08124](https://arxiv.org/abs/1707.08124)

Motivation

- top quark is the most massive particle in the SM
 - important probe for SM or ingredient for BSM models
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How close are we?



EW vacuum stability illustration [1]

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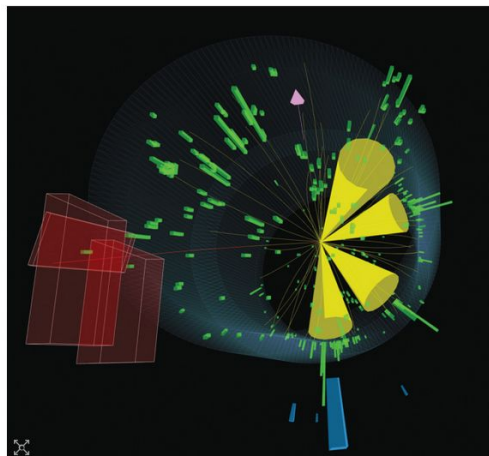
Previous results at LHC

- latest public result from 2022 with lepton+jets channel using 2016 data performed by the CMS Collaboration
 - $m_t = 171.77 \pm 0.38$ GeV [1]

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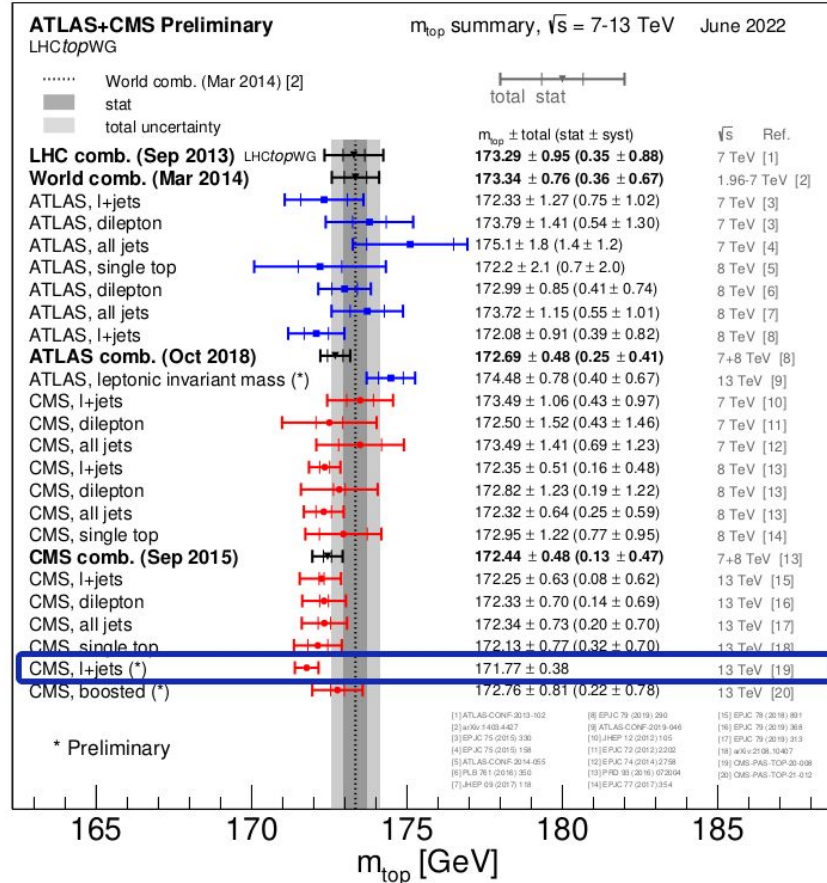
Top quark weighs in with unparalleled precision

1 July 2022



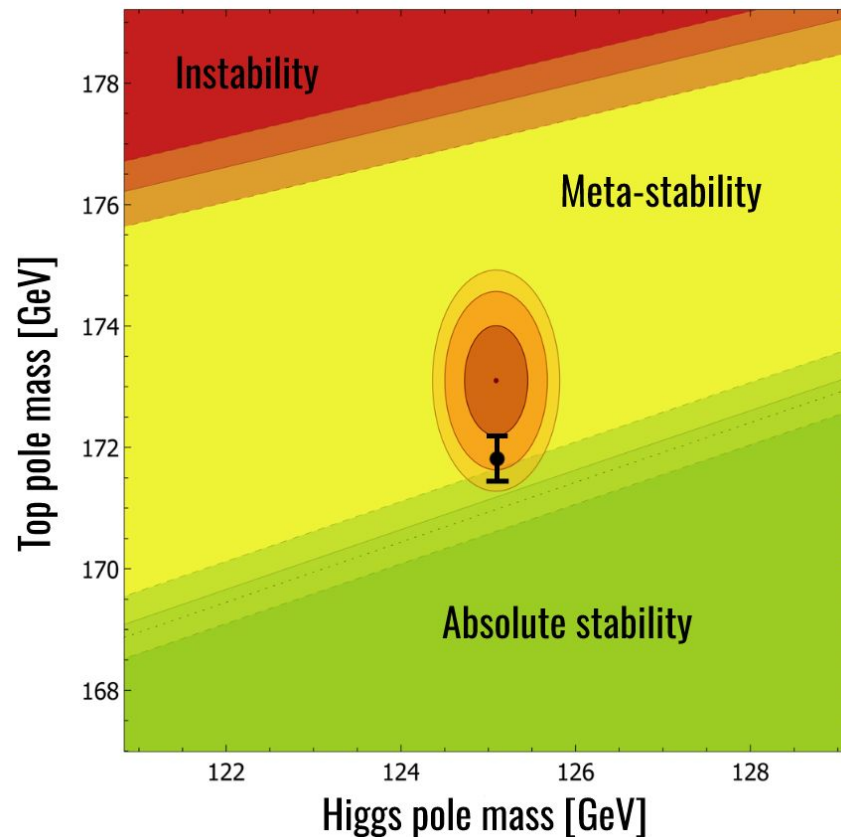
Top marks The classic signature of a top-quark pair at the LHC is four jets (yellow cones), one muon (red line and boxes) and missing energy from a neutrino (pink arrow). Credit: CMS

[1] [CMS-PAS-TOP-20-008](#)



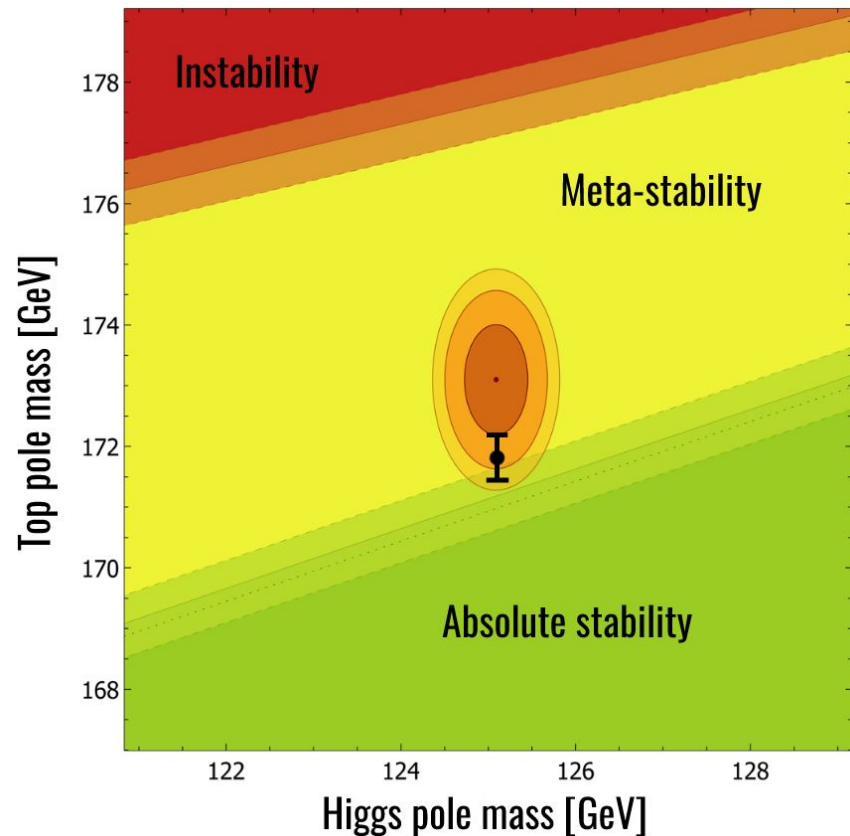
Lepton + jets 2016 results

- 171.77 ± 0.38 GeV in the stability figure



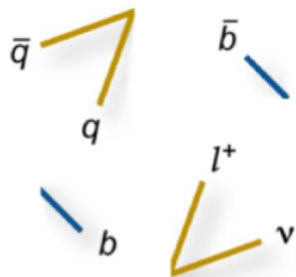
Lepton + jets 2016 results

- 171.77 ± 0.38 GeV in the stability figure
- our analysis in Helsinki extends this using data from 2017-2018
 - $\sim 100 \text{ fb}^{-1}$ \rightarrow around 3 times more than 2016
 - systematic uncertainties can be reduced if well constrained by the data using **profile likelihood methods**

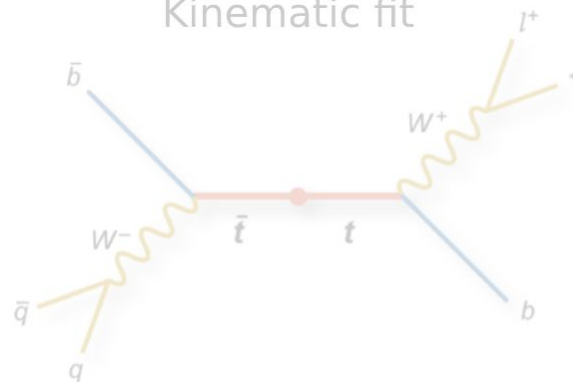


Analysis strategy

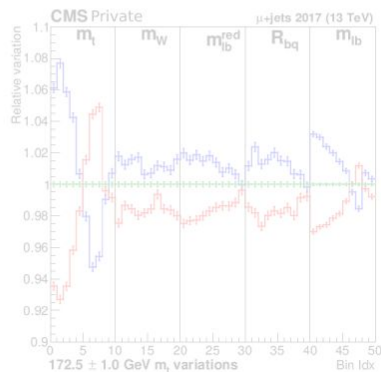
Object selections



Kinematic fit

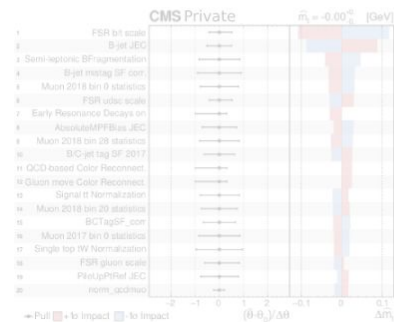


Binned profile likelihood fit



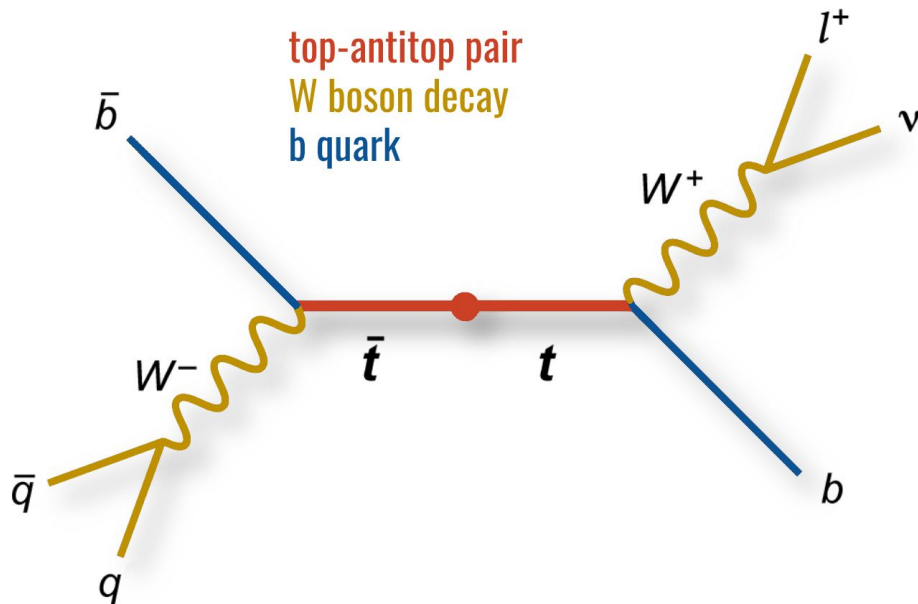
Results

17X.XX ± 0.XX GeV



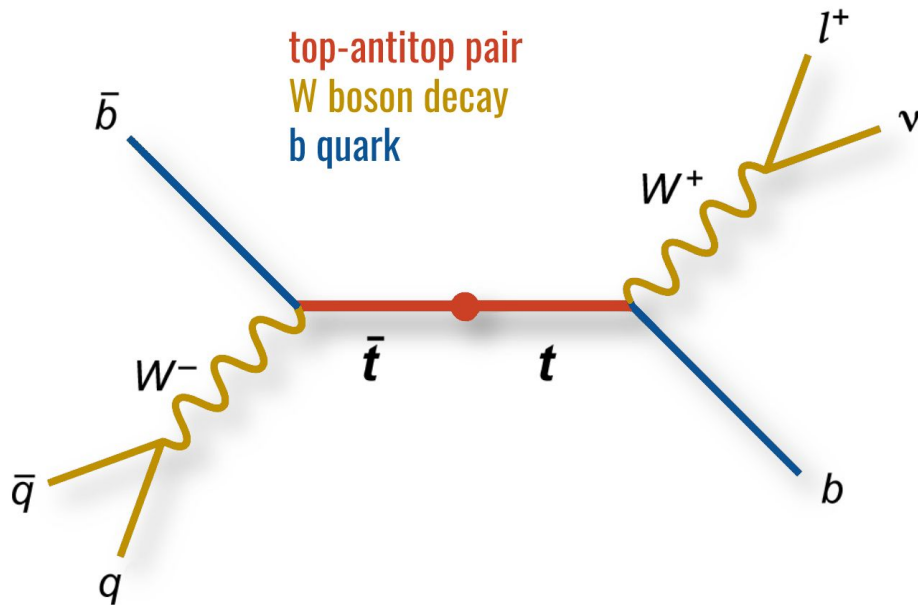
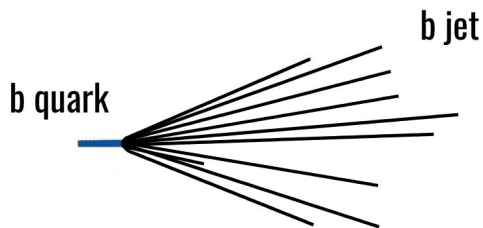
Mass measurement in $t\bar{t}$ topology

- direct mass measurement
 - mass from the reconstructed decay products
 - most precise method at LHC at the moment
- $t\bar{t}$ at 13 TeV LHC offers largest statistics
- top decays almost exclusively to W and b
- channels categorized by the W boson decays
 - dileptonic ($W \rightarrow l + \nu$, $W \rightarrow l + \nu$)
 - + clear signal
 - missing momentum from neutrinos
 - hadronic ($W \rightarrow qq$, $W \rightarrow qq$)
 - + largest statistics
 - difficult to distinguish jets
 - **semileptonic channel** ($W \rightarrow qq$, $W \rightarrow l + \nu$)
 - ± signal from the single lepton
 - + only one neutrino
 - ± decent statistics

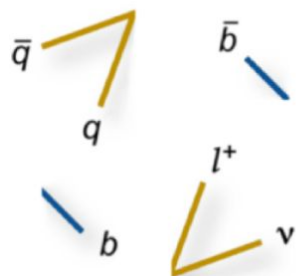


Semileptonic channel selections

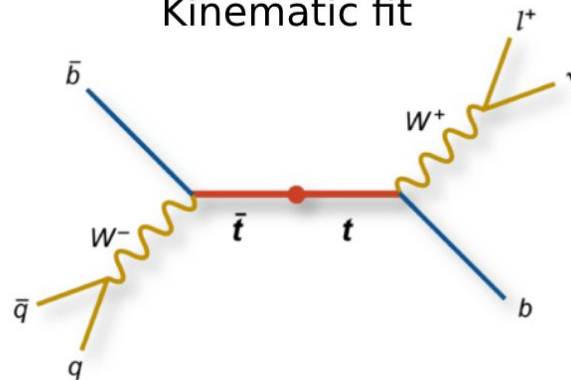
- exactly one muon/electron
- at least 4 jets
 - of which 2 are b-jets
 - challenging to tag
 - of which 2 are light quark jets from W
 - boosted W complicates things
- neutrino inferred using missing momentum



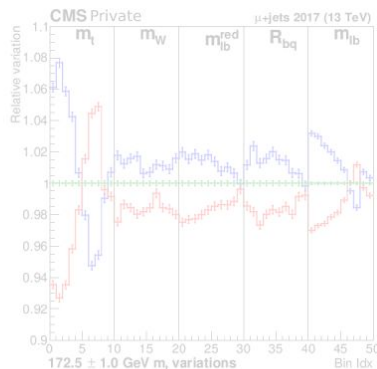
Object selections



Kinematic fit

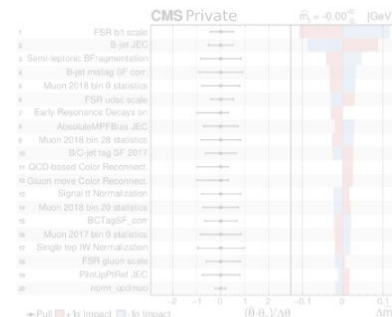


Binned profile likelihood fit



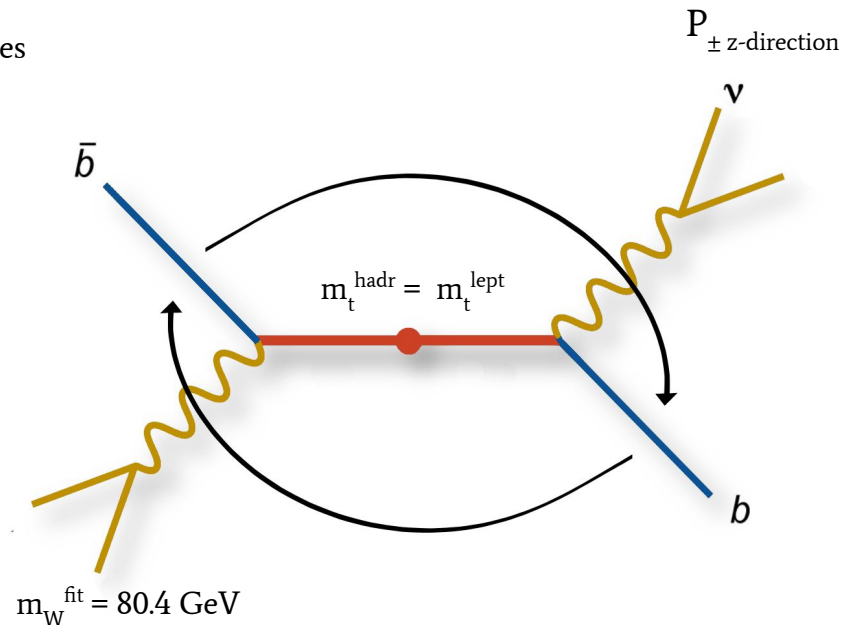
Results

$17X.XX \pm 0.XX \text{ GeV}$



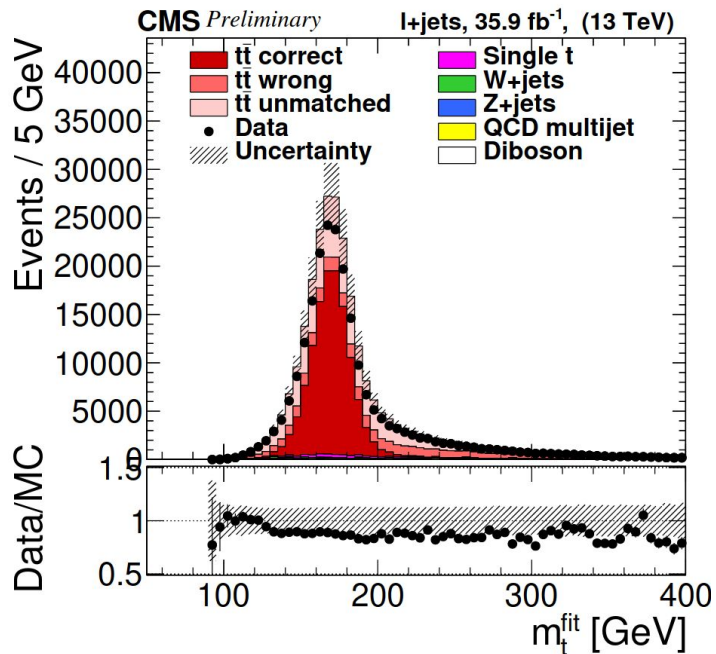
Kinematic fit

- semileptonic event hypothesis tested for combinations of selected objects
 - two possible combinations for b jets
 - neutrino momentum z-component has two possible values
- kinematic fit constraints:
 - $m_W^{\text{fit}} = 80.4 \text{ GeV}$
 - $m_t^{\text{hadr}} = m_t^{\text{lept}}$
- gives χ^2 using object-parton resolution
 - example: further the reconstructed qq system from W mass, larger the χ^2
- goodness-of-fit for each permutation determined
 - $P_{\text{gof}} = \exp(-1/2 \chi^2)$
 - hypothesis with the highest P_{gof} value is used

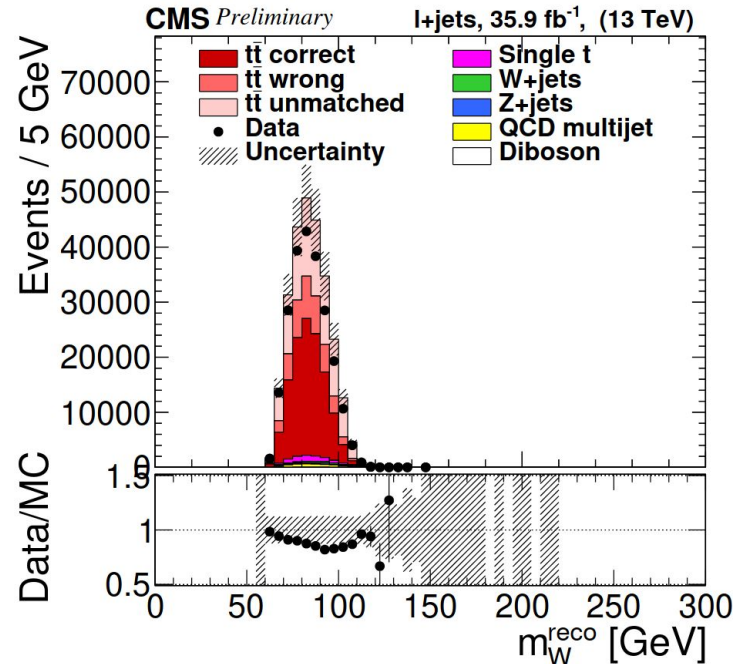


Control plots after kinematic fit

- data-simulation agreement can be confirmed

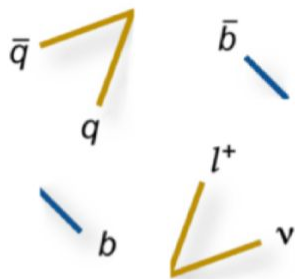


top mass from the fit

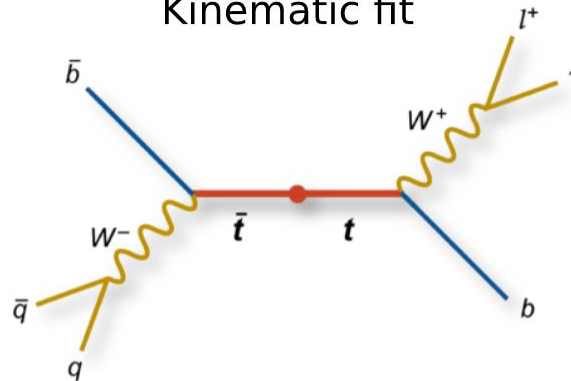


W mass

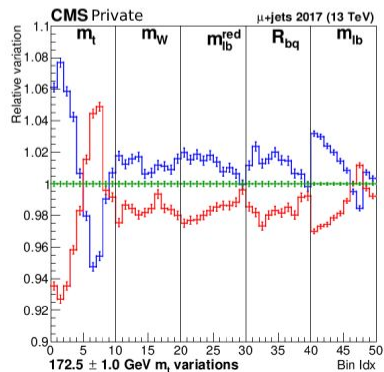
Object selections



Kinematic fit

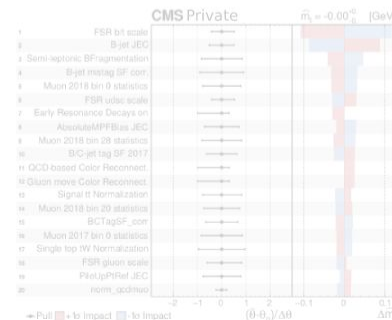


Binned profile likelihood fit



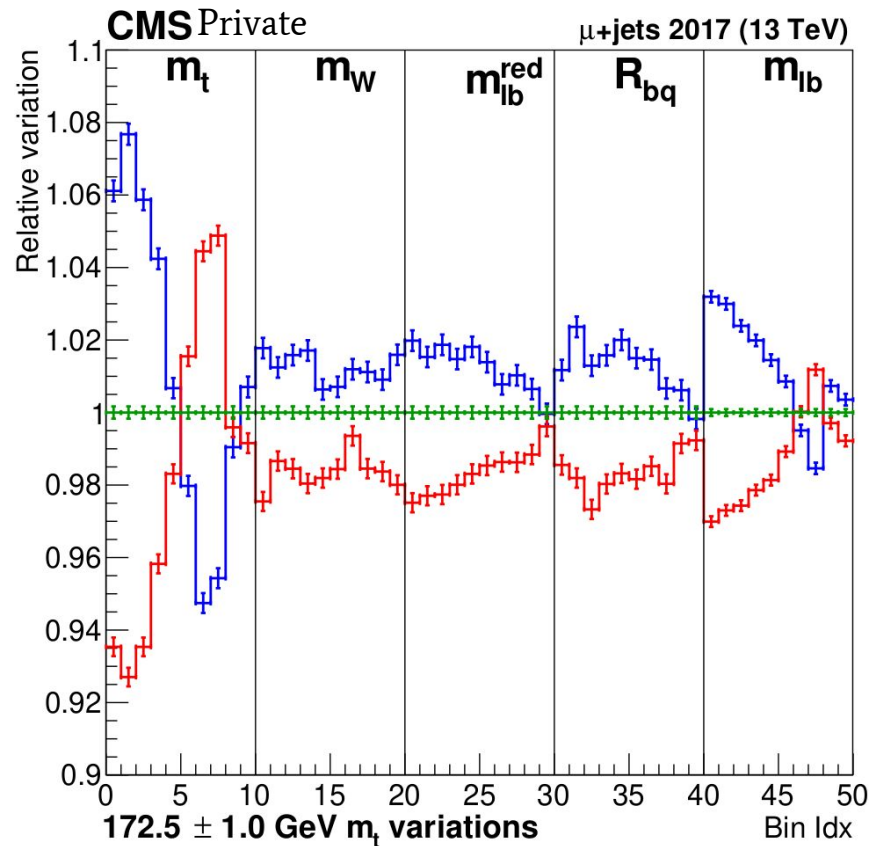
Results

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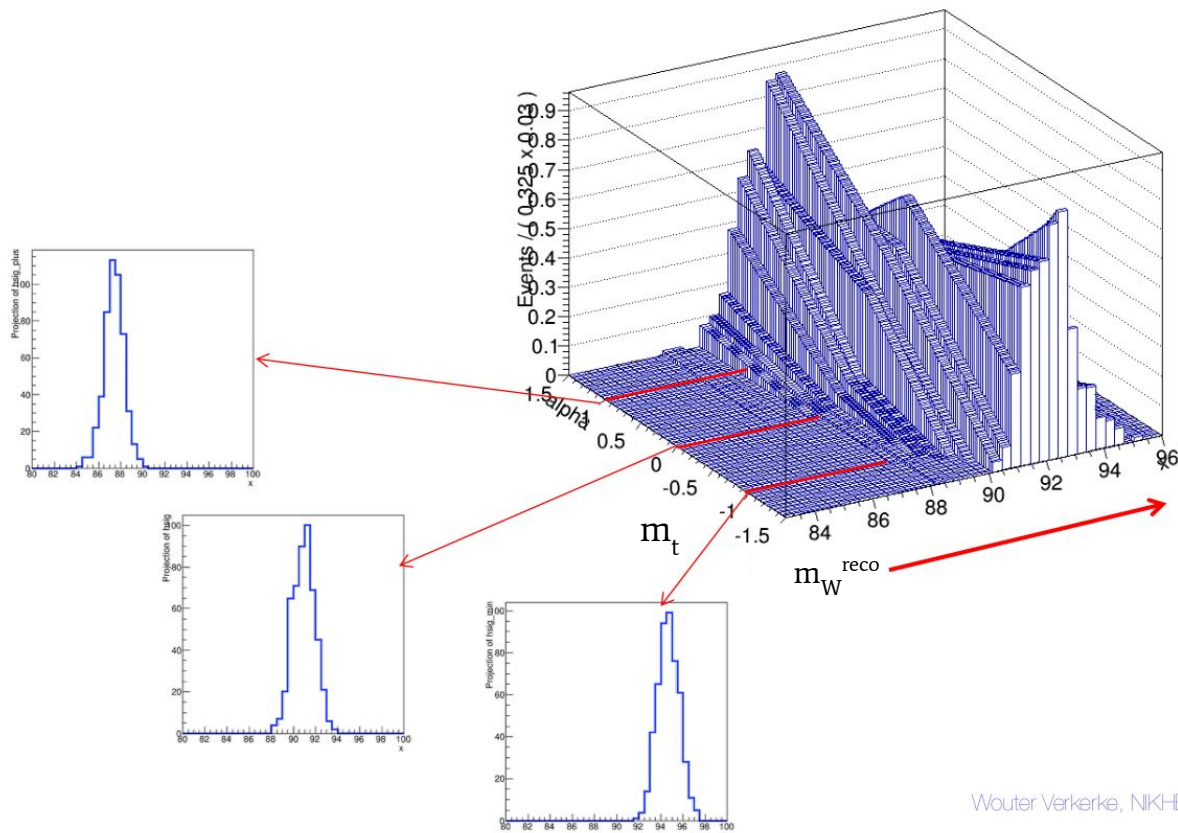
Binned profile likelihood fit

- m_t is the parameter of interest
- binned fit in five dimensions
 1. m_t^{fit} = obvious choice for m_t
 2. m_W^{reco} = jet energy scale of light quarks
 3. $m_{\text{lb}}^{\text{red}}$
 4. $R_{\text{bq}}^{\text{reco}}$ = b-jet energy scale (originally from ATLAS)
 5. $m_{\text{lb}}^{\text{reco}} (P_{\text{good}} < 0.2)$
- **Combine** tool used for the fitting
 - handles well year combinations
 - used extensively by Higgs community
 - nowadays also in precision measurements
- m_t inferred from the maximized likelihood



Interpolation example

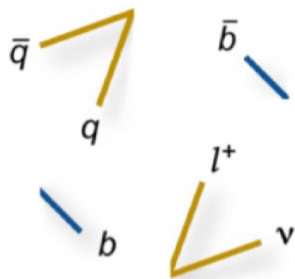
- interpolation between central simulation and variations by histogram morphing
- linear interpolation between -1,0,1 variations



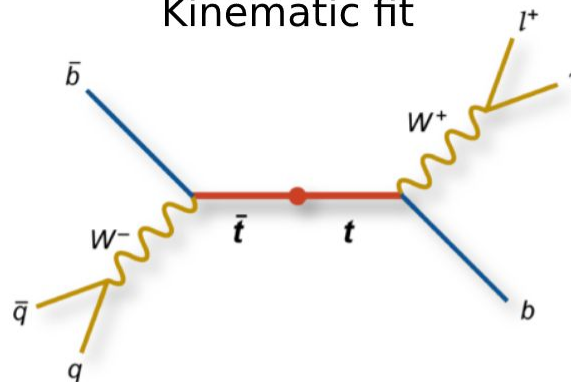
Wouter Verkeke, NIKHEF

Source: [Wouter Verkeke - Advanced Statistics. Systematic uncertainties and profiling](#)

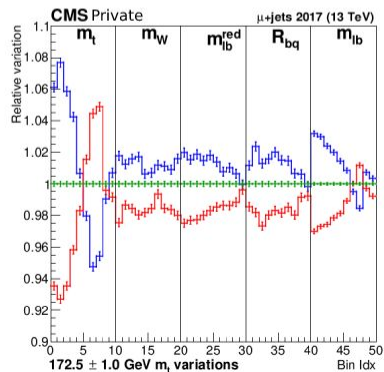
Object selections



Kinematic fit

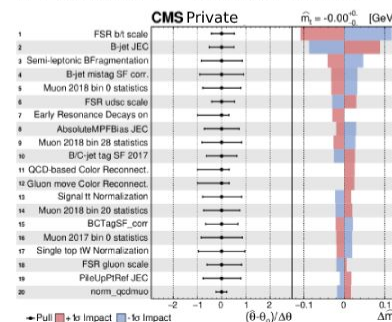


Binned profile likelihood fit



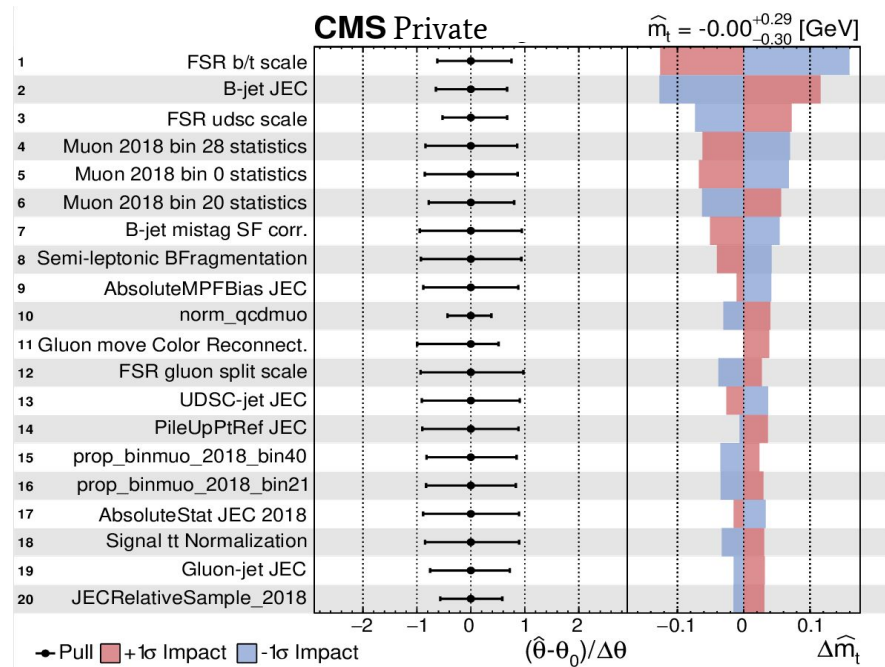
Results

$17X.XX \pm 0.XX \text{ GeV}$



Impact plot for 2018

- uncertainty on the m_t by each nuisance parameter
- impact plotting
 - fit all nuisance parameters together
 - fixing single nuisance to ± 1 value and refitting
 - impact = difference in the value of the m_t after fit
- in example plot we have blinded impact where the simulation is used as a target instead of data
 - impacts still estimate the uncertainty
- bJEC and b/t quark final state radiation (FSR) leading systematics

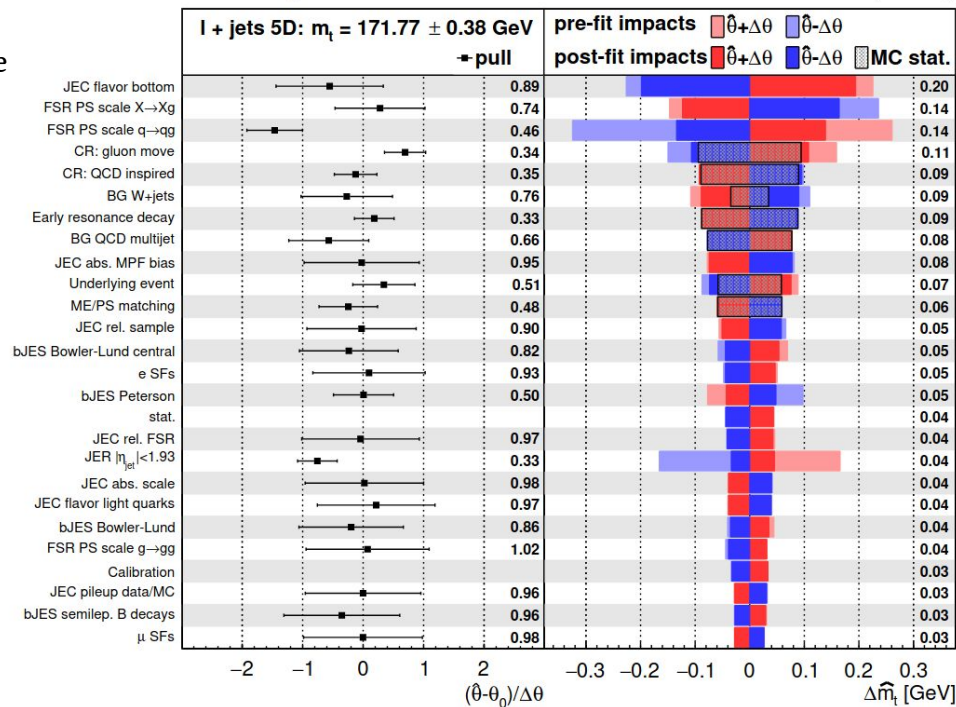


Impacts for 2016

- now fitted against real data
 - pulls indicate how different systematics should be shifted to best model the data
 - ideally pulls are below one sigma
- large pull in qFSR, opposite to bFSR
 - qFSR -1.5σ
 - bFSR $+0.3 \sigma$
- uses split FSR scheme
 - light quarks and heavy quarks not correlated
- however, there is indication that
 - α_s^{FSR} should be bigger in simulation
 - qFSR and bFSR are correlated
 - these are being studied for 2017-2018 analysis

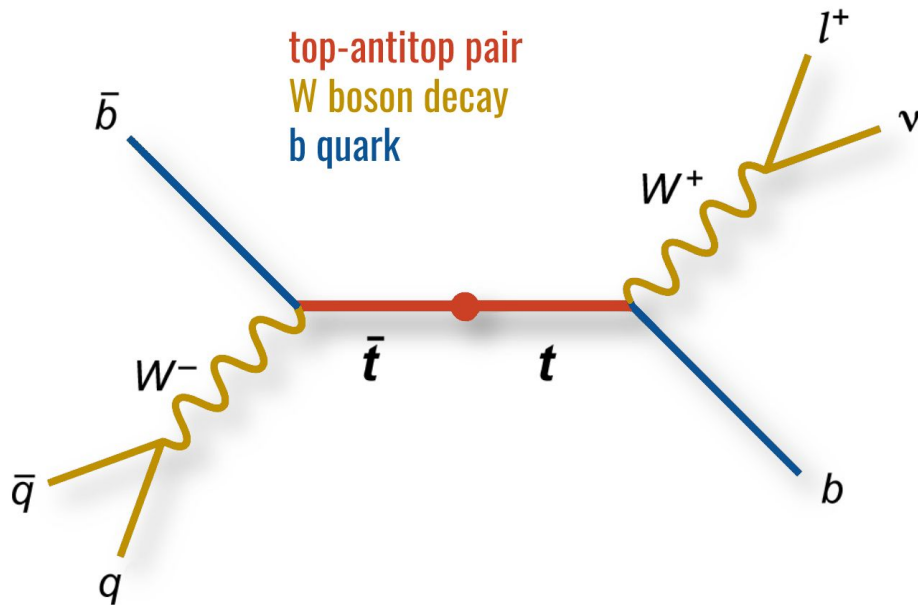
CMS Preliminary

36 fb⁻¹ (13 TeV)



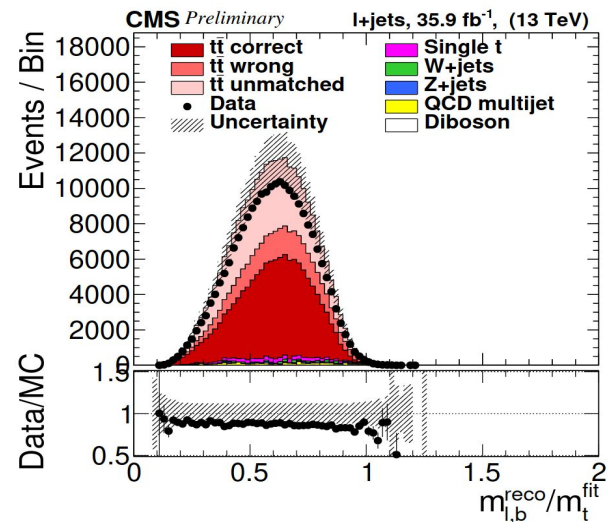
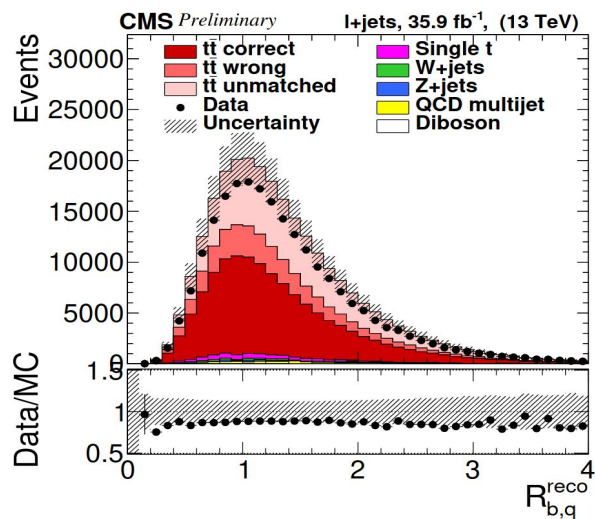
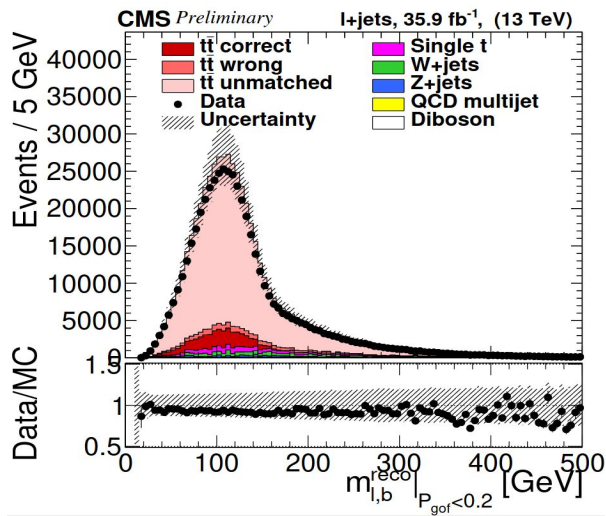
Summary

- latest public lepton + jets 2016 results presented
 - **most precise top mass measurement yet**
 - open questions related to FSR modeling
- we will aim to publish 2017-2018 results in 2023
 - around 3 times more data than 2016
 - some differences in the details of the analysis
- direct top mass measurement in $t\bar{t}b$ semileptonic channel
 - event selections
 - kinematic fit
 - binned profile likelihood method



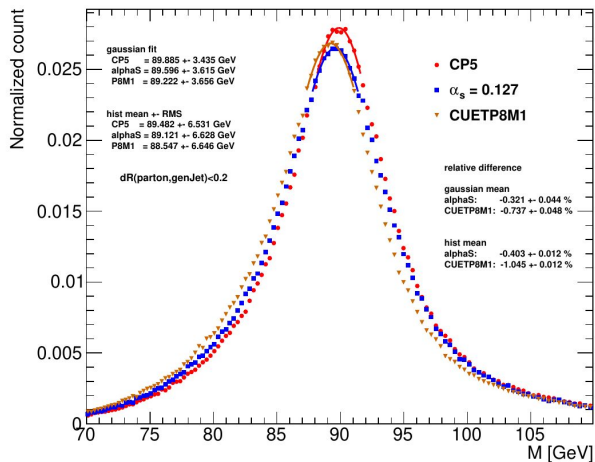
Backup

Backup: rest of the fit dimensions

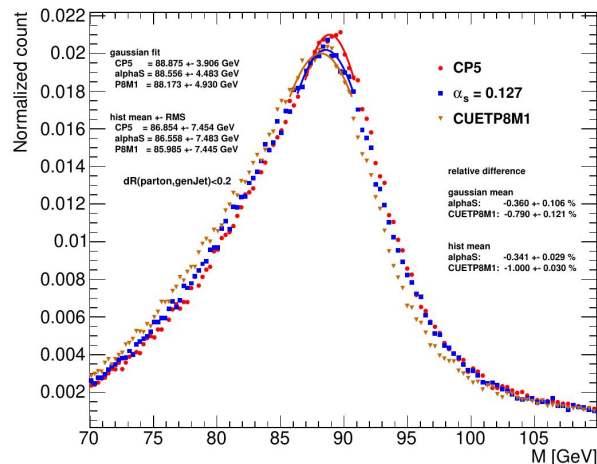


Backup: FSR correlation

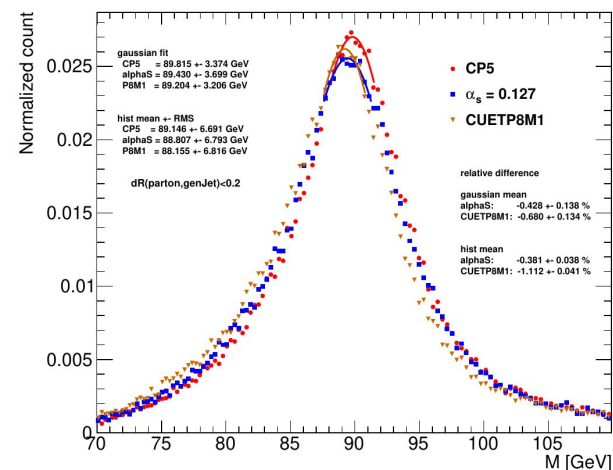
- new studies indicating correlation between light and b quark FSR in Pythia
- performed using $Z \rightarrow qq'$ decays and reconstructing Z-mass from the simulated quark jets



Zqq (no bb)



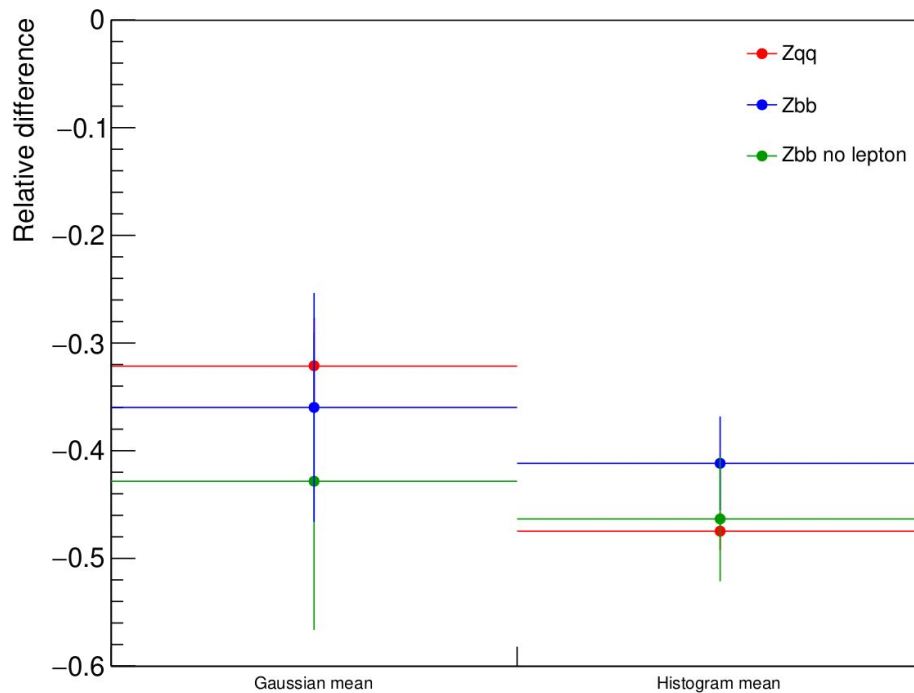
Zbb



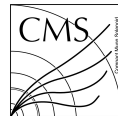
Zbb no lepton decays

Backup: FSR correlation

α_s variation from 0.118 to 0.127



Backup: Binned likelihood function

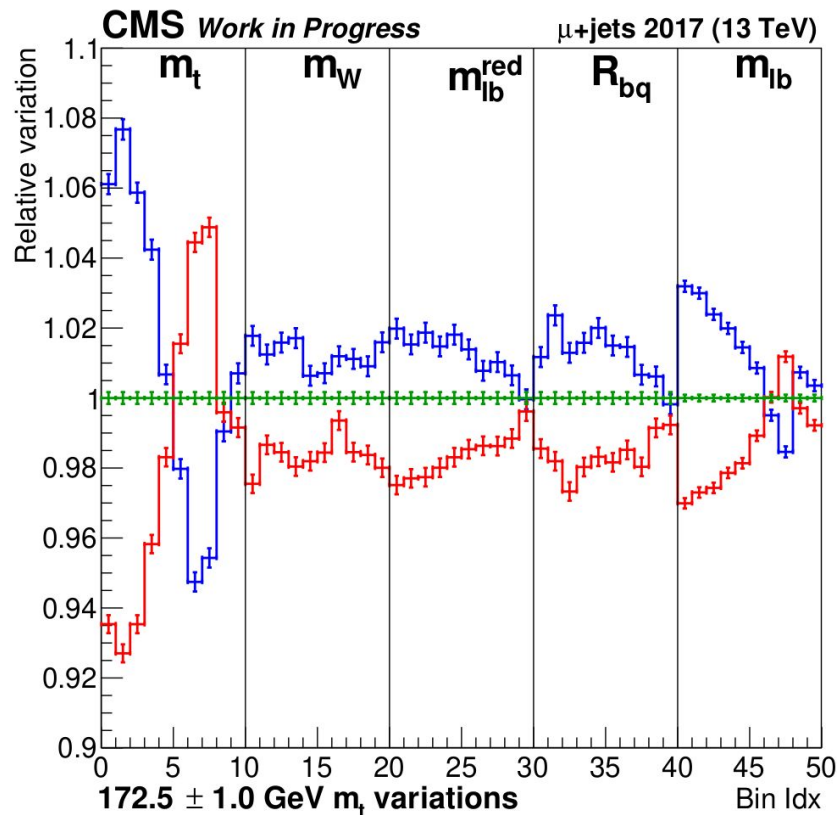


$$\mathcal{L}(\vec{n}) = \prod_{i \in \text{bins}} \mathcal{P} \left(n_i \mid \sum_{j \in \text{samples}} (1 + \kappa_j)^{\eta_j} \times \nu_i^j(\vec{\theta}, m_t) \right) \\ \times \prod_{k \in \text{nuisances}} \mathcal{G}(\theta_k) \times \prod_{j \in \text{samples}} \mathcal{G}(\eta_j).$$

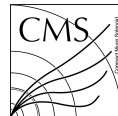
- n_i = number of events in data in bin i
- ν_i^j = expected number of events for simulated sample j
- $\mathcal{P}(n|\lambda)$ = Poisson probability to observe n events when λ predicted
- θ_k = nuisance parameter corresponding to a systematic uncertainty k
 - $\theta = 0$ is the central value and -1 and $+1$ are the down and up variations
- κ_j = scaling uncertainty of simulated sample j
 - minimally cross-section uncertainty
- systematics interpreted as nuisance parameters
- nuisances constrained by normalized Gaussians $\mathcal{G}(\theta_k)$
 - up/down variations shift the likelihood by one sigma

Backup: Likelihood fit dimensions

- binned fit in five dimensions
 - m_t^{fit}
 - m_W^{reco}
 - $m_{\text{lb}}^{\text{red}} = m_{\text{lb}}^{\text{reco}} / m_t^{\text{fit}}$
 - $R_{\text{bq}}^{\text{reco}} = (p_T^{\text{b-jet 1}} + p_T^{\text{b-jet 2}}) / (p_T^{\text{q-jet 1}} + p_T^{\text{q-jet 2}})$
 - $m_{\text{lb}}^{\text{reco}} (P_{\text{gof}} < 0.2)$



Backup: top mass definition



- top quark mass is theoretically complex to define
- what is measured in the direct measurements is the so-called Monte Carlo mass m_t^{MC}
- relationship with field theoretic mass definitions such as the pole mass m_t^{pole} is not straightforward

“top mass interpretation problem”

- theoretical work (A. Hoang [1]) has been conducted on this topic and as the precision of the measurements increases this additional uncertainty must be given serious thought

[1] <http://dx.doi.org/10.1146/annurev-nucl-101918-023530>