

# Higgs EFT interpretation of the search for $HH \rightarrow b\bar{b}\gamma\gamma$ in the ATLAS detector

Laura Pereira Sánchez

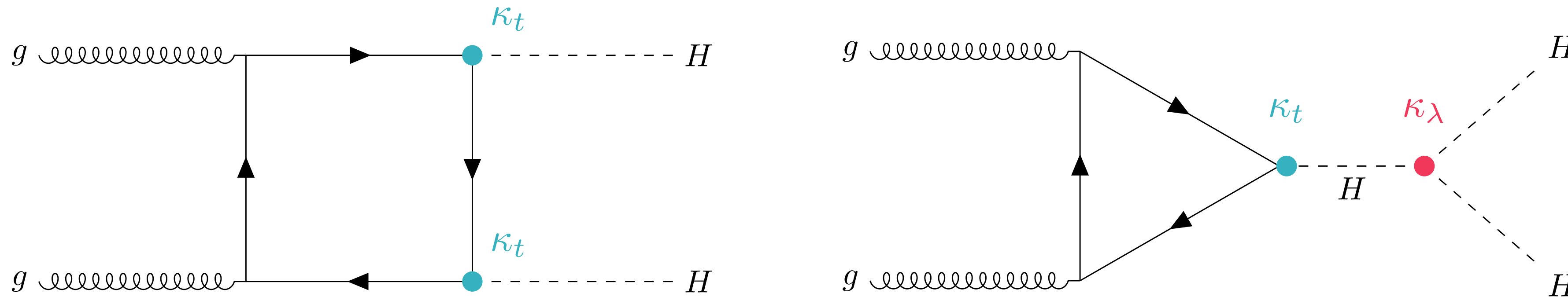
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Higgs EFT interpretation of the  
**Search for  $HH \rightarrow b\bar{b}\gamma\gamma$  in the ATLAS detector**

# $ggF$ $HH$ production

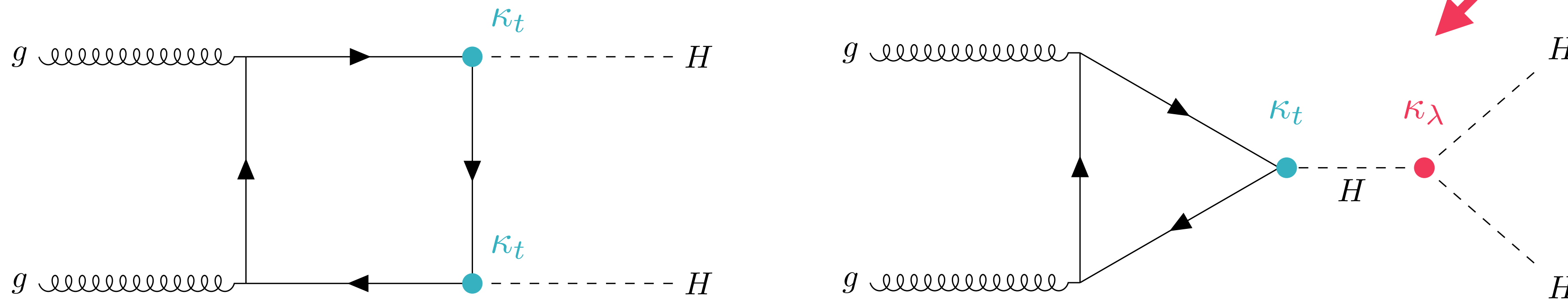
- Measuring  $HH$  production will give us access to the triple Higgs coupling (self coupling)  $\lambda_3$ , which gives information of the shape of the Higgs potential  $V(H) = \frac{1}{2}m_H^2 H^2 + \lambda_3 \nu H^3 + \frac{1}{4}\lambda_4 \nu H^4 + O(H^5)$ .
- The leading  $HH$  production mode is gluon gluon fusion ( $ggF$ ):



- The coupling modifier  $\kappa_\lambda$  controls the strength of the Higgs self coupling with respect to SM:  $\kappa_\lambda = \lambda_3 / \lambda_3^{SM}$
- Destructive interference between the two diagrams results in a very small SM cross section of  $\sigma_{ggF}^{HH} = 31.05$  fb at  $\sqrt{s} = 13$  TeV.

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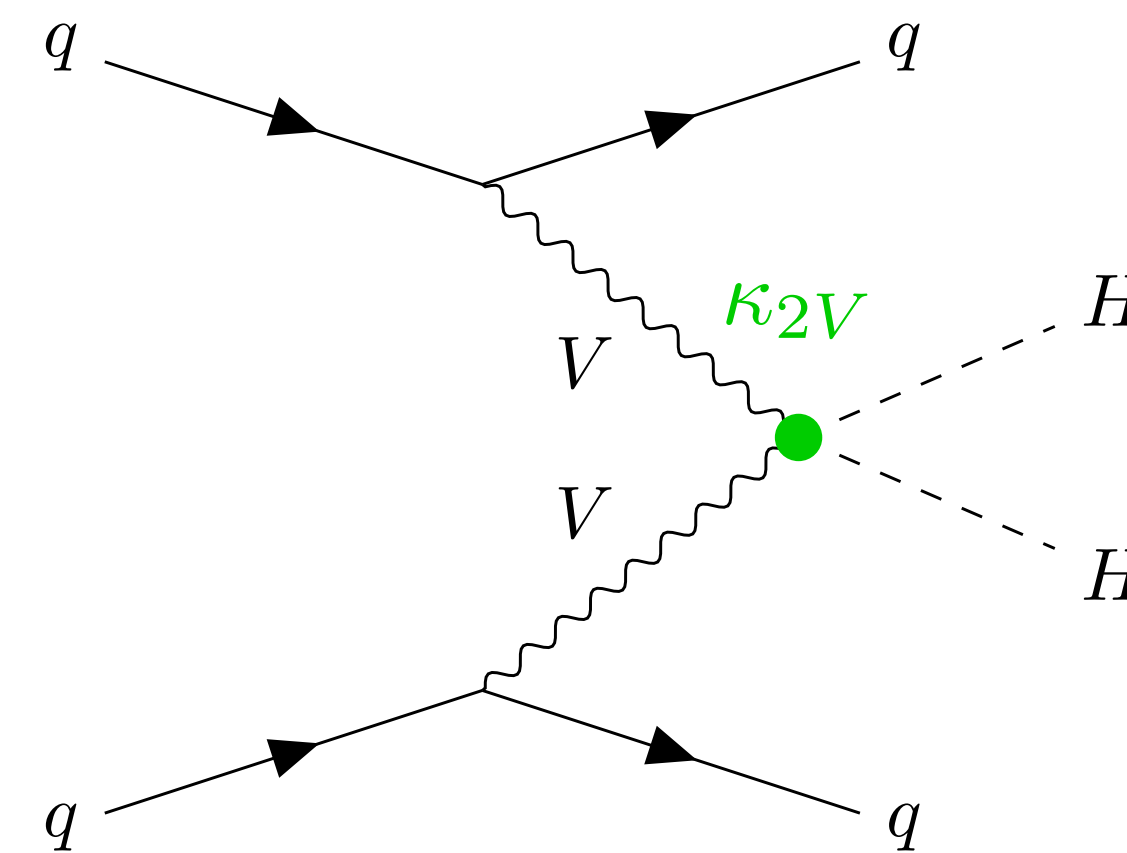
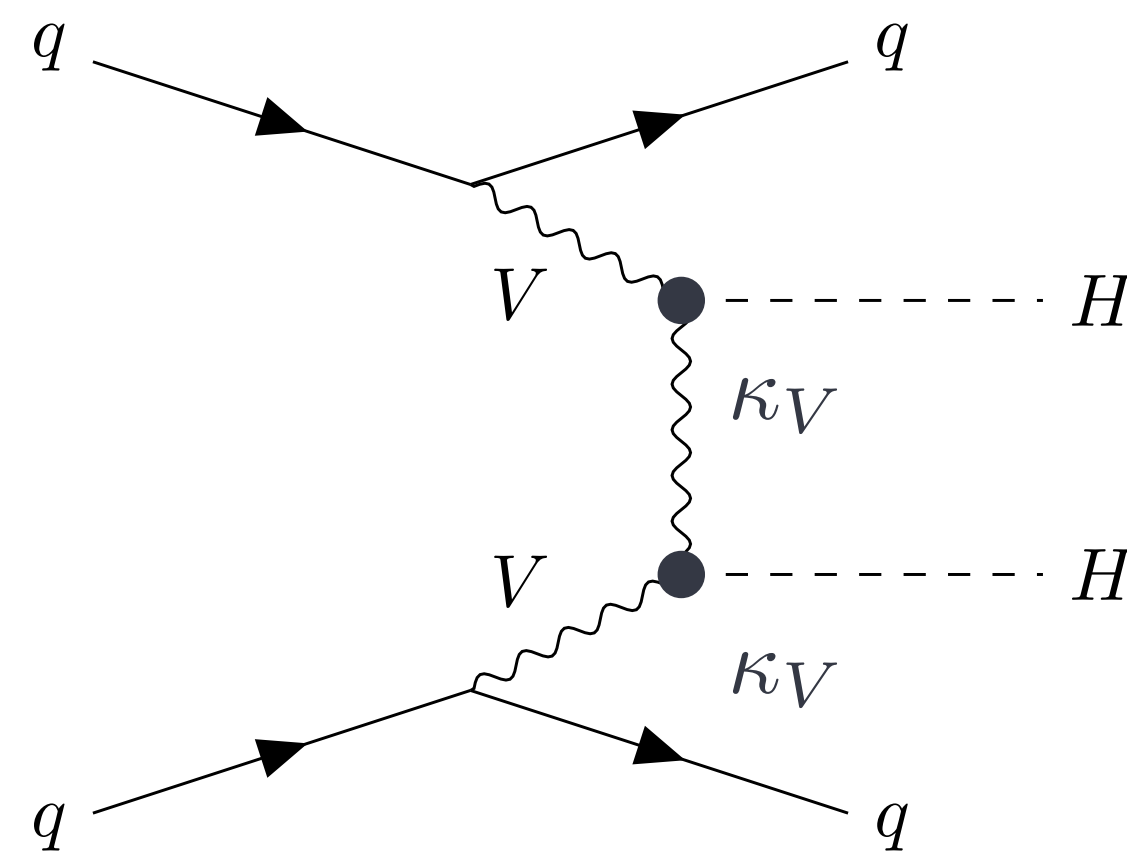
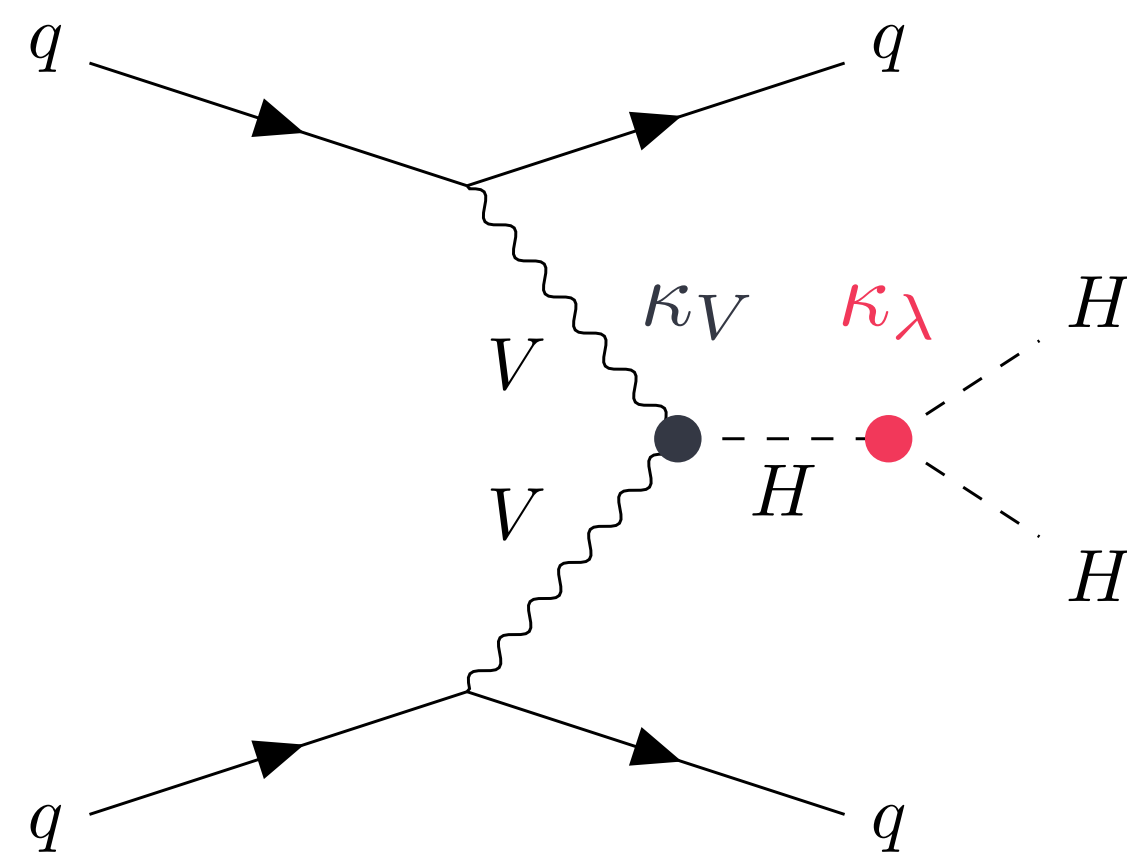


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# VBF HH production

- $HH$  production through  $VBF$  is the sub-leading  $HH$  production mode with a SM cross section of  $\sigma_{VBF}^{HH} = 1.73 \text{ fb}$  at  $\sqrt{s} = 13 \text{ TeV}$  (calculated at N3LO)

- The coupling modifiers  $\kappa_\lambda$ ,  $\kappa_V$  and  $\kappa_{2V}$  control the strength of the  $g_{HHH} = \frac{3m_H}{v^2}$ ,  $g_{VVH} = \frac{2m_V^2}{v}$  and  $g_{VVHH} = \frac{2m_V^2}{v^2}$  couplings with respect to the SM value.

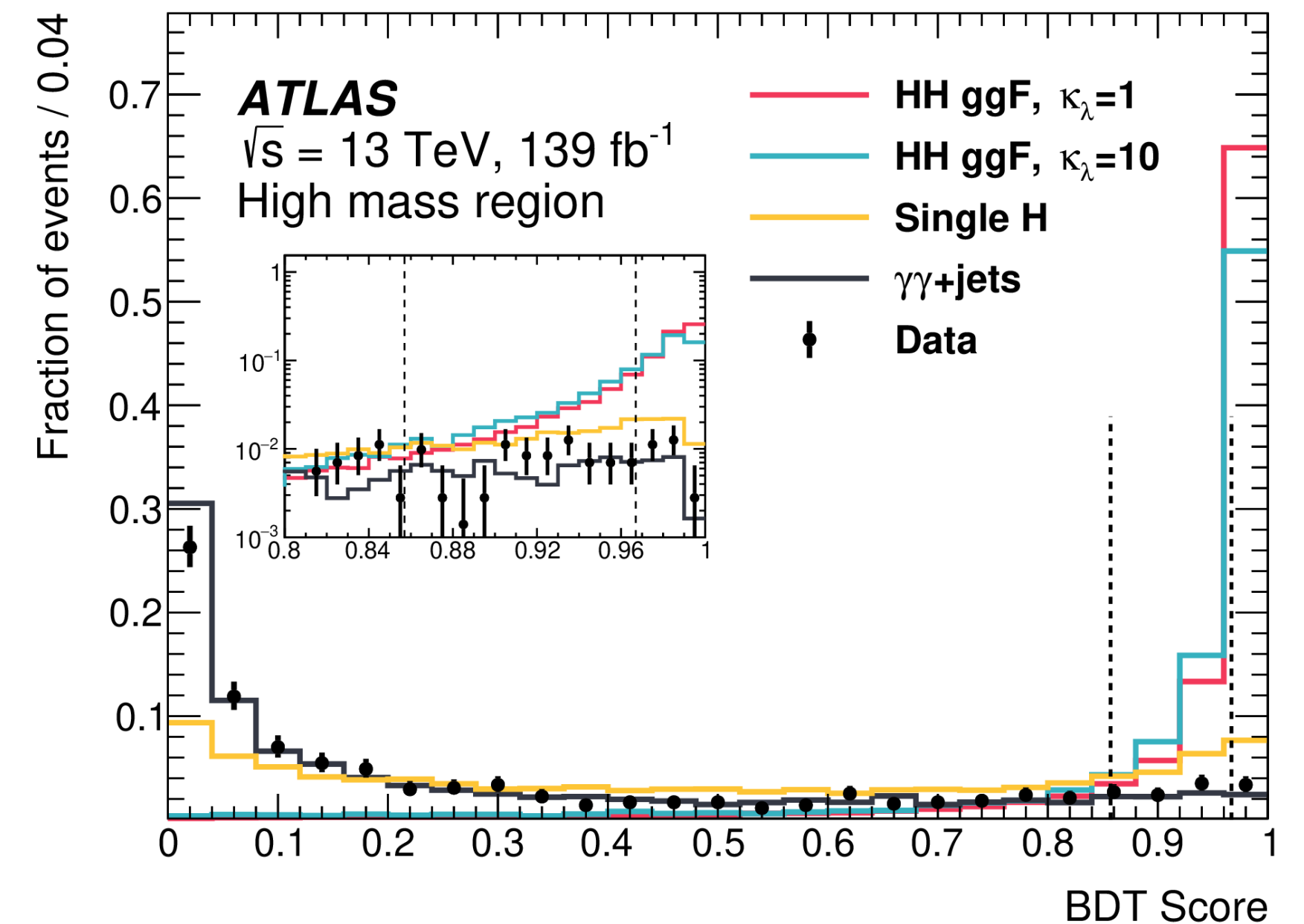
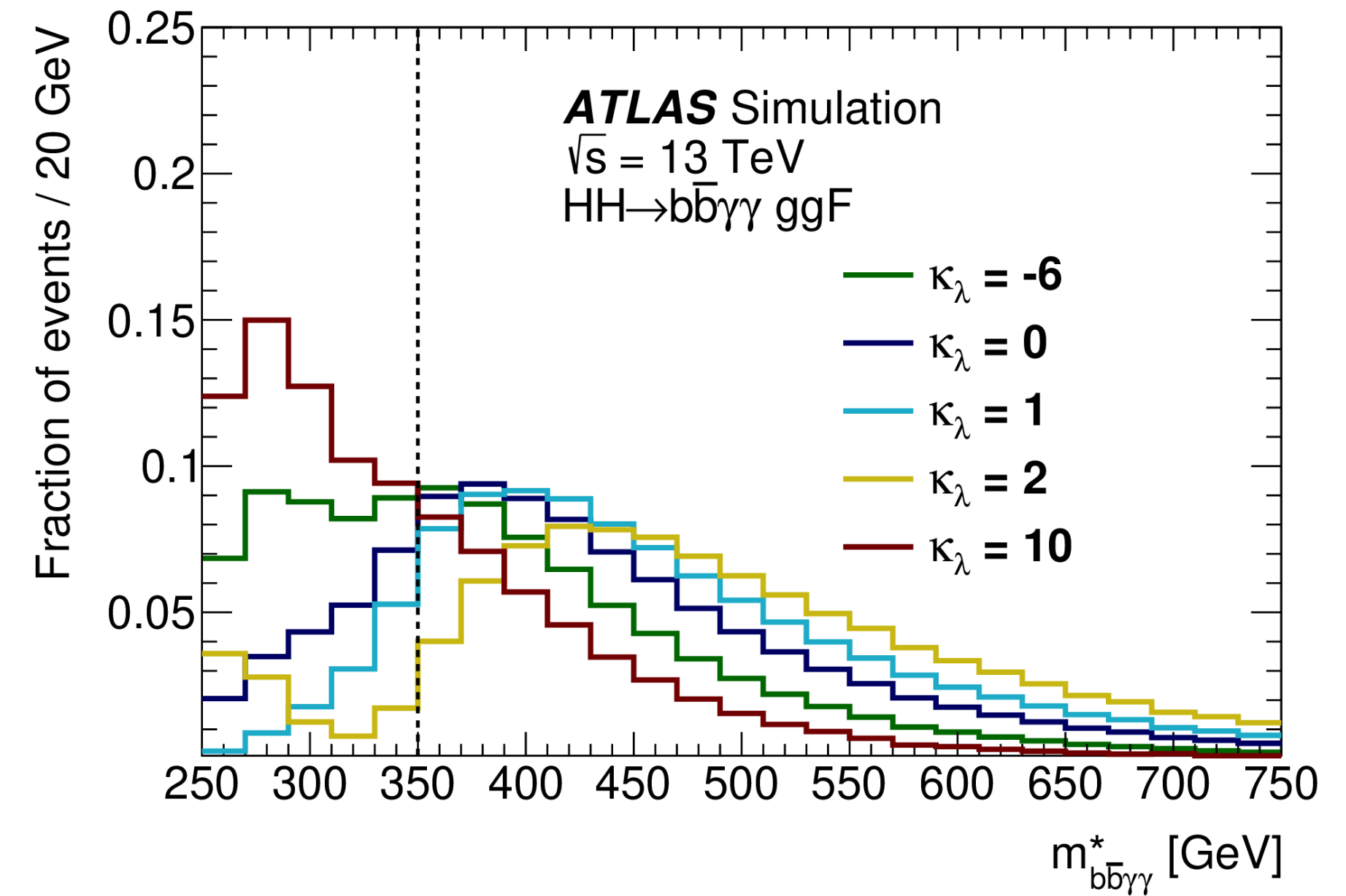


- Given the larger cross section, searches for  $ggF$   $HH$  production provide better sensitivity to  $\kappa_\lambda$  but the VBF topology has a unique sensitivity to  $\kappa_{2V}$ .

# $HH \rightarrow b\bar{b}\gamma\gamma$ event selection

- The paper (optimised for ggF  $HH$  production) has been published in [PhysRevD](#)
- Events with 2 isolated photons, 2  $b$ -jets at 70% WP and zero leptons are selected and divided in two regions:
  - High mass region:  $m_{b\bar{b}\gamma\gamma}^* > 350$  GeV  $\rightarrow$  sensitive to SM
  - Low mass region:  $m_{b\bar{b}\gamma\gamma}^* < 350$  GeV  $\rightarrow$  sensitive to BSM
- Independent BDTs are trained in each region to separate signal from background.
- Two regions are defined from each BDT resulting in 4 final categories optimised for ggF  $HH$   $\kappa_\lambda = 1$  and  $\kappa_\lambda = 10$ :

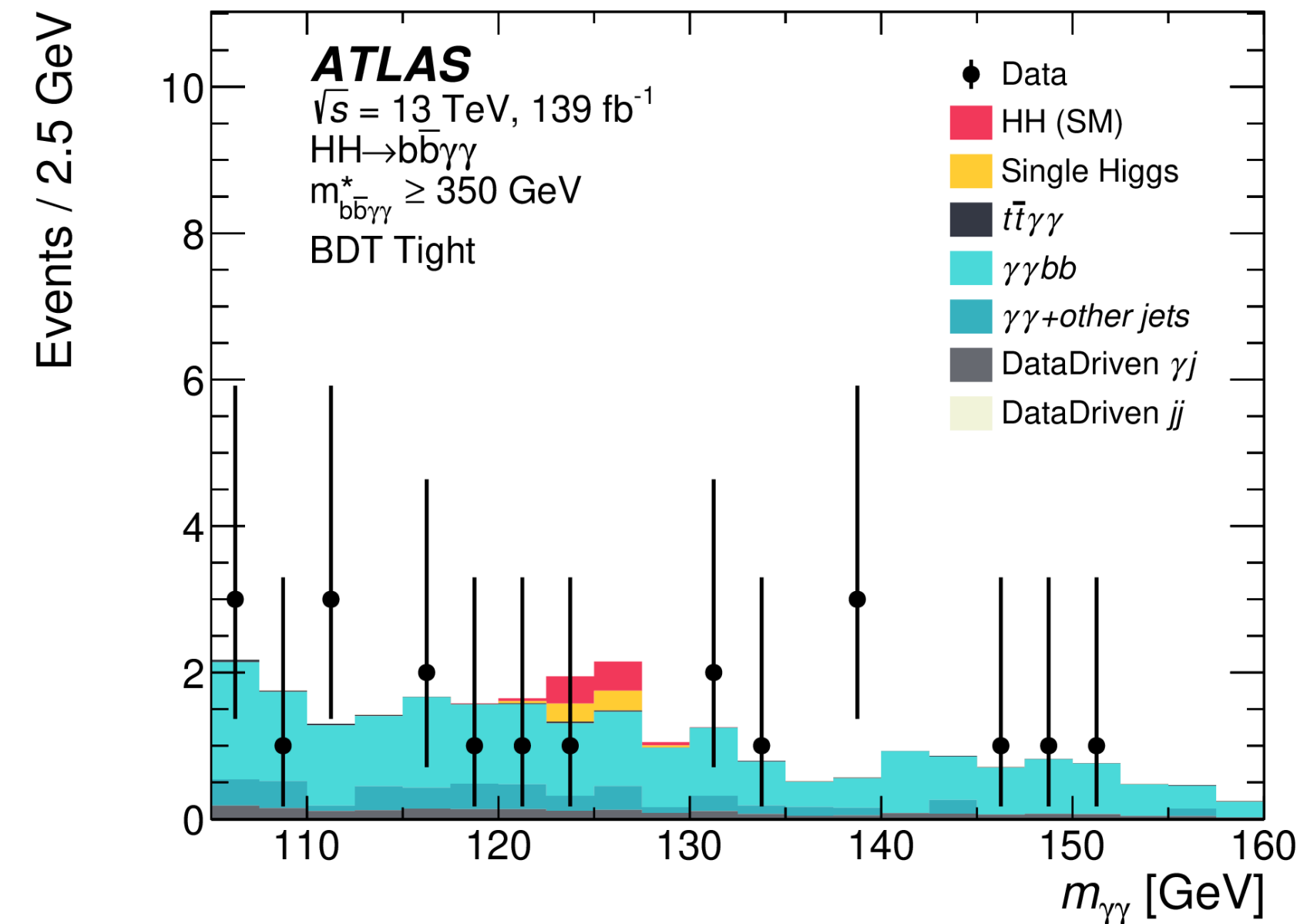
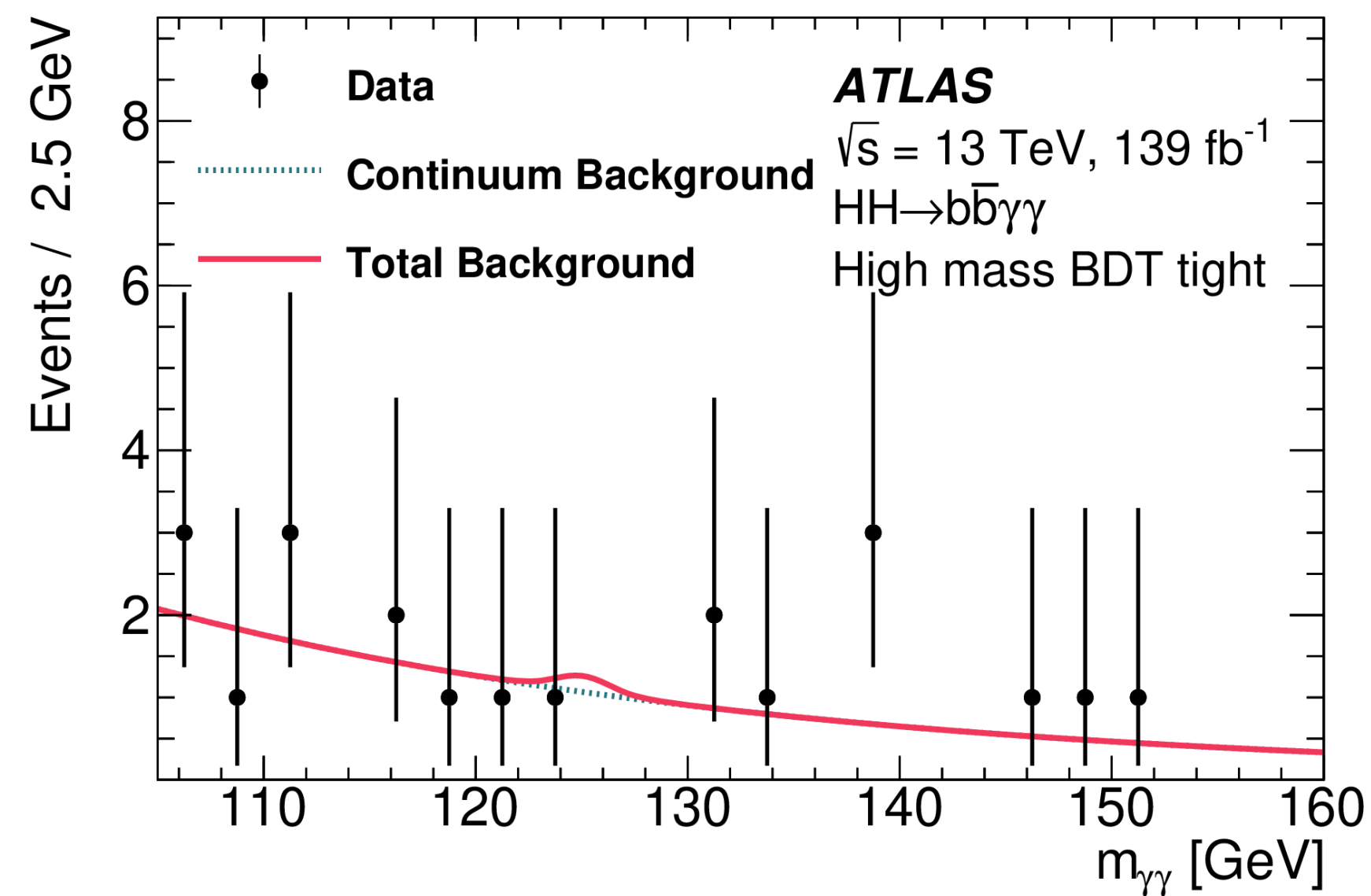
Category	Selection criteria
High mass BDT tight	$m_{b\bar{b}\gamma\gamma}^* \geq 350$ GeV, BDT score $\in [0.967, 1]$
High mass BDT loose	$m_{b\bar{b}\gamma\gamma}^* \geq 350$ GeV, BDT score $\in [0.857, 0.967]$
Low mass BDT tight	$m_{b\bar{b}\gamma\gamma}^* < 350$ GeV, BDT score $\in [0.966, 1]$
Low mass BDT loose	$m_{b\bar{b}\gamma\gamma}^* < 350$ GeV, BDT score $\in [0.881, 0.966]$





# $HH \rightarrow b\bar{b}\gamma\gamma$ fit

- A likelihood fit to the  $m_{\gamma\gamma}$  is performed
- The non-resonant background is parametrised as an exponential function where the shape is estimated from simulation while the normalisation is estimated from data through a fit to the  $m_{\gamma\gamma}$  side band ( $120 \text{ GeV} < m_{\gamma\gamma} > 130 \text{ GeV}$ ).
- The resonant processes are parametrised from a double side crystal ball function in the  $120 > m_{\gamma\gamma} > 130 \text{ GeV}$  region, where the shape and normalisation is obtained from simulation.



# $HH \rightarrow b\bar{b}\gamma\gamma$ limits

Kappa framework

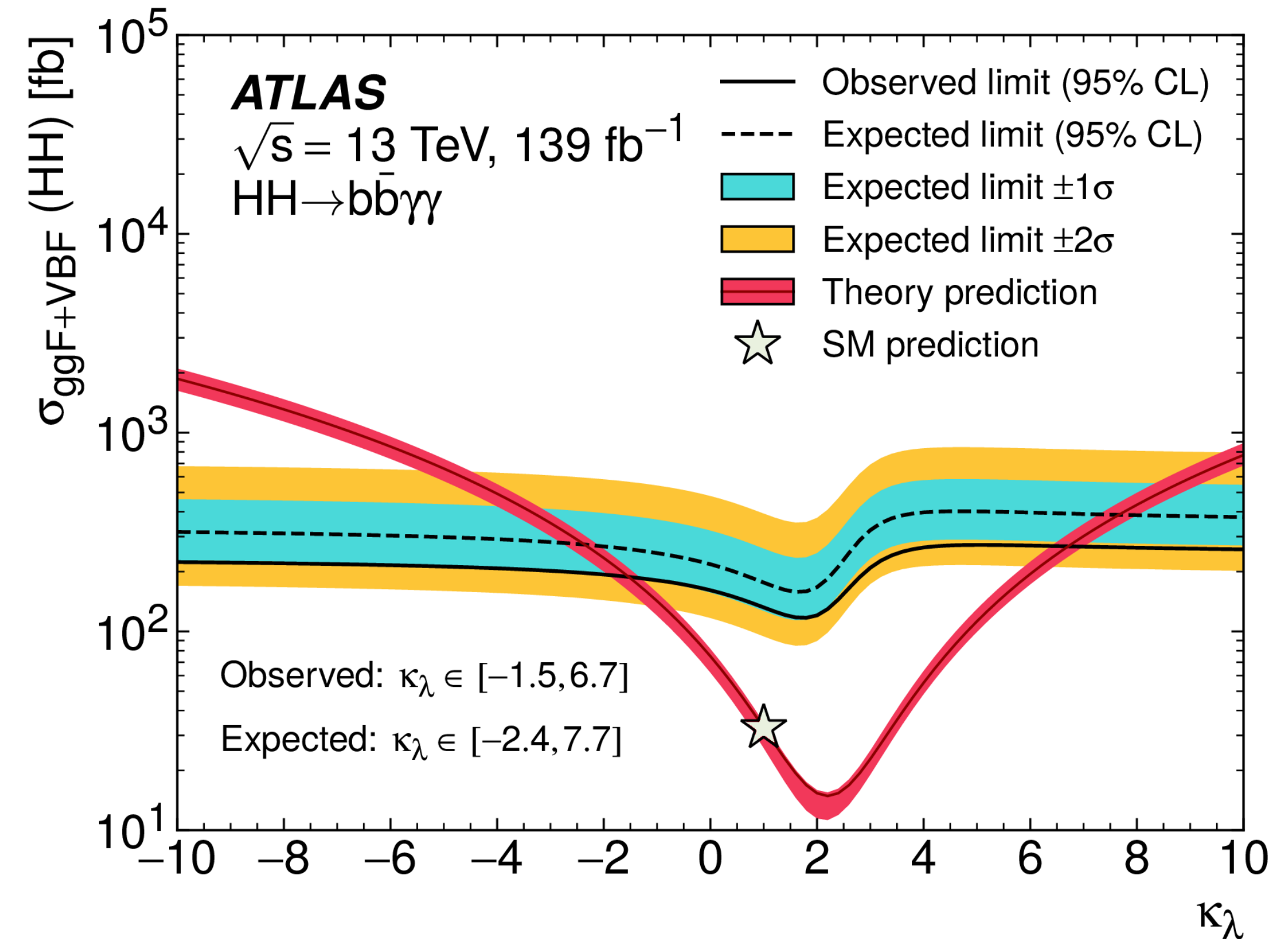
- The analysis is optimised for  $ggF$   $HH$ , however  $VBF$   $HH$  events are also considered as signal.



- Observed (expected) limits at 95% CL:

- $$\sigma_{ggF+VBF}^{HH} < 4.2 \text{ (5.7)} \times \sigma_{ggF+VBF}^{HH SM}$$

- $$-1.5 \text{ (-2.4)} < \kappa_\lambda < 6.7 \text{ (7.7)}$$





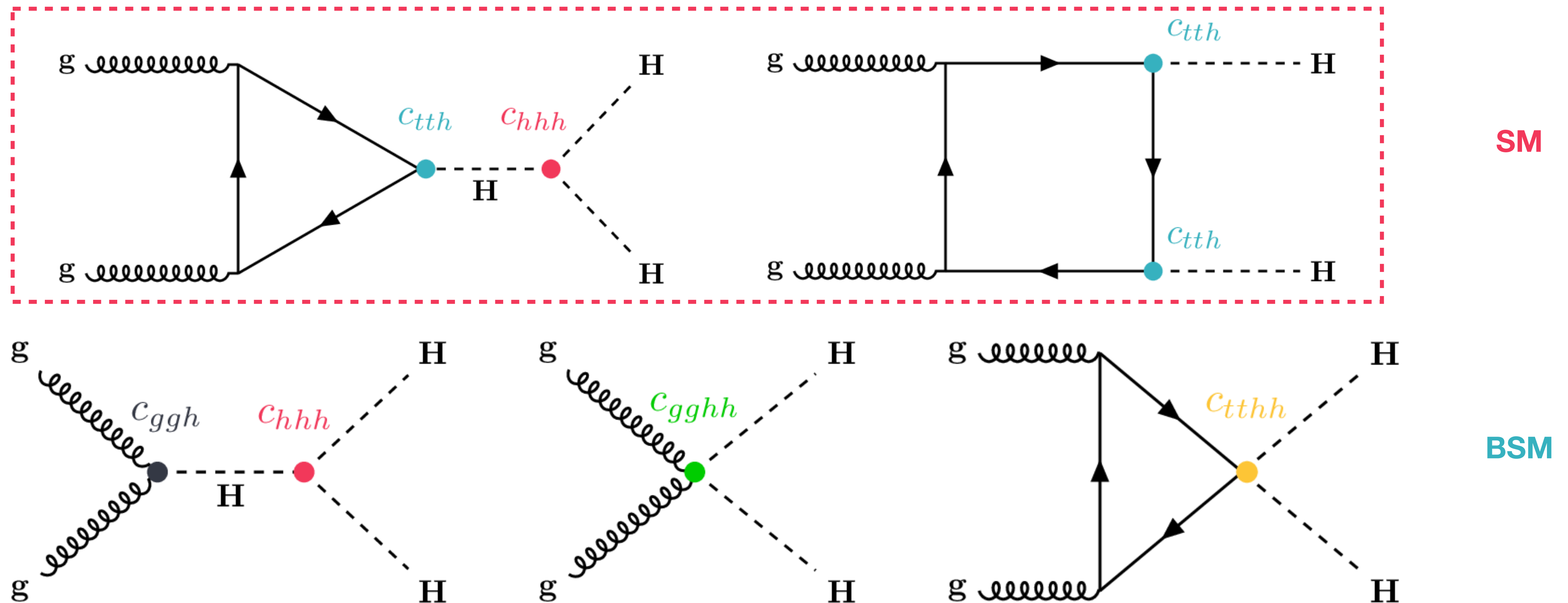
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# HEFT interpretation

- The Effective Field Theory (EFT) framework can be used as a tool to:

1. Make a more general measurement of the Higgs self-coupling → Natural extension of the kappa framework
2. Explore Beyond-the-Standard-Model scenarios produced a  $E \gg E_{LHC}$ .

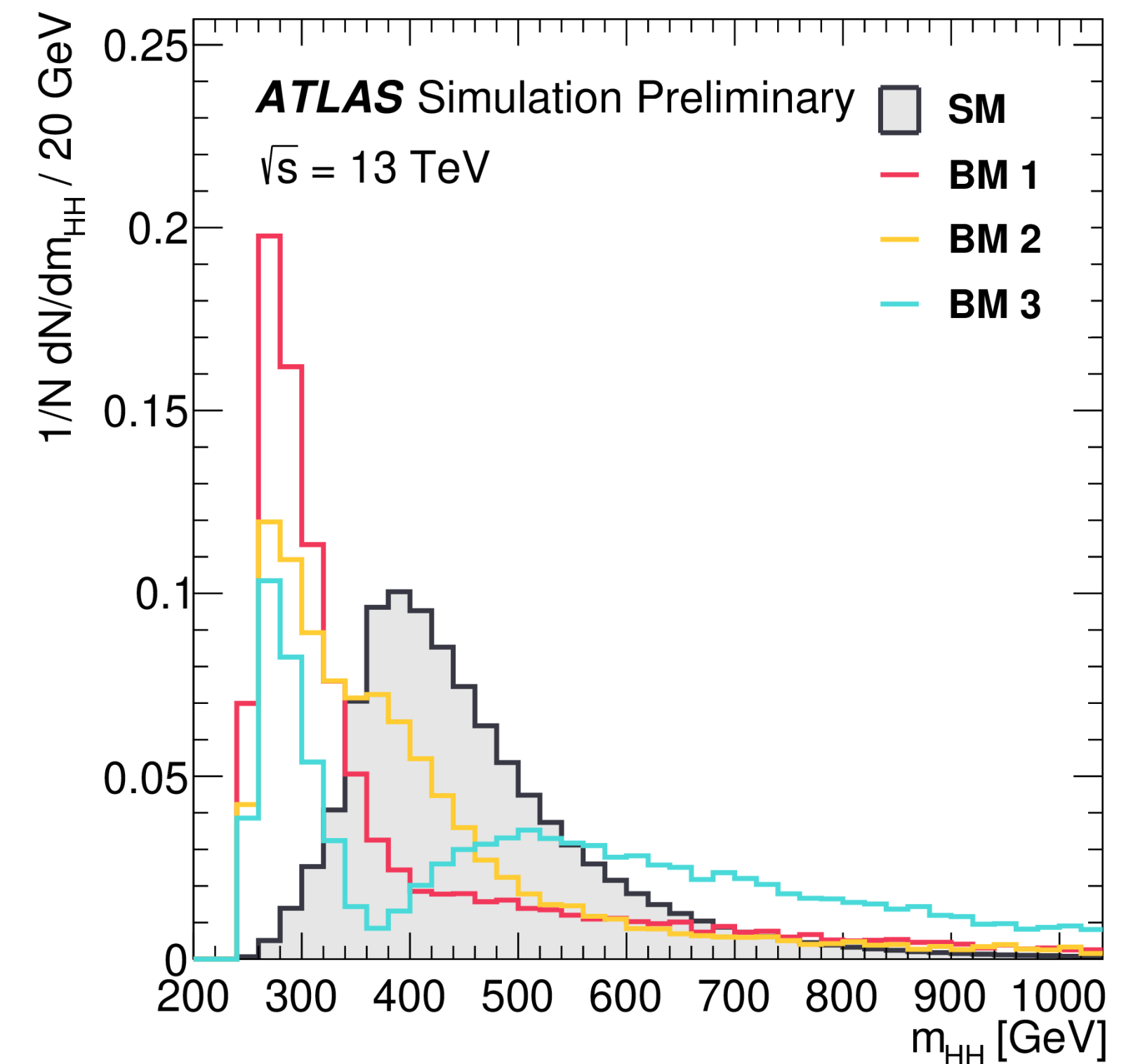
- Where  $c_{hhh} = \kappa_\lambda$  and  $c_{tth} = \kappa_t$
- In the SM:  $c_{hhh} = c_{tth} = 1$  and  $c_{ggh} = c_{tthh} = c_{gggh} = 0$



# BSM exploration through HEFT

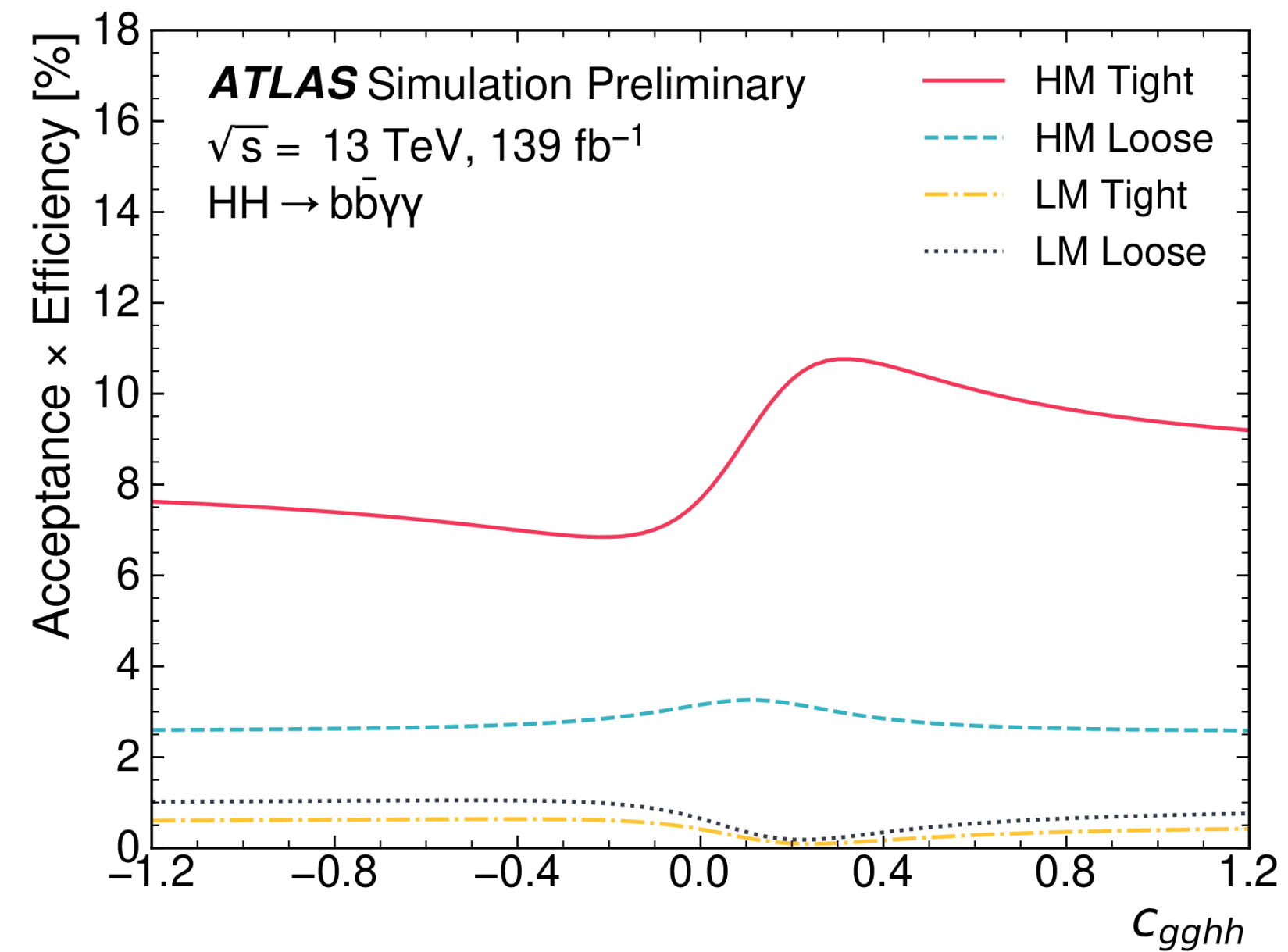
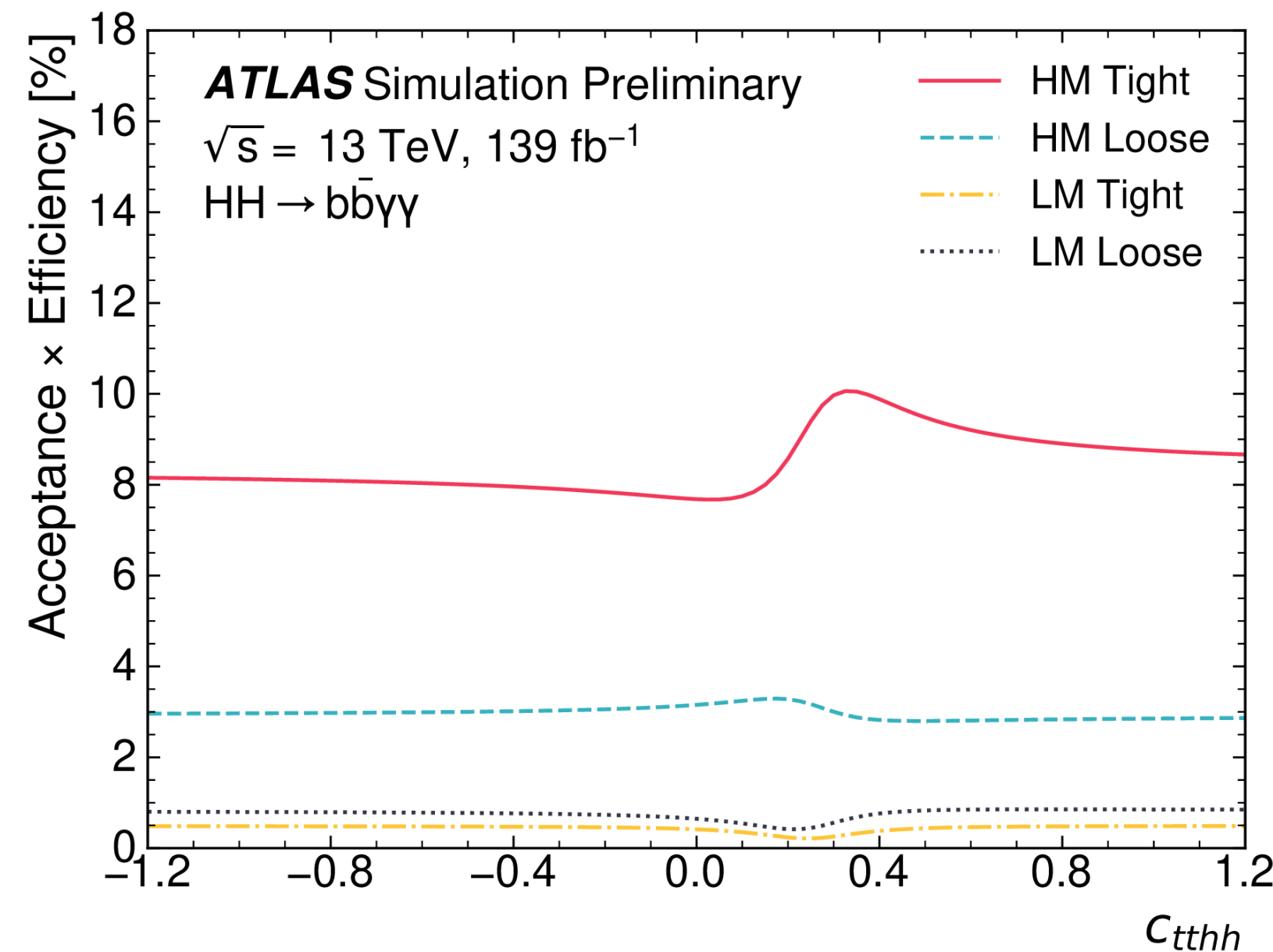
- HEFT shape benchmarks can be used to explore BSM scenarios which are uniquely sensitive to HH production.
- By varying the EFT couplings we are also modifying the shape of  $m_{hh}$ .
- Theorists have suggested a set of 7 benchmarks that fairly **represent the different shapes** obtained by the coupling variations in HEFT at NLO.

Benchmark model	$c_{hhh}$	$c_{tth}$	$c_{ggh}$	$c_{gggh}$	$c_{tthh}$
SM	1	1	0	0	0
BM 1	3.94	0.94	1/2	1/3	-1/3
BM 2	6.84	0.61	0.0	-1/3	1/3
BM 3	2.21	1.05	1/2	1/2	-1/3
BM 4	2.79	0.61	-1/2	1/6	1/3
BM 5	3.95	1.17	1/6	-1/2	-1/3
BM 6	5.68	0.83	-1/2	1/3	1/3
BM 7	-0.10	0.94	1/6	-1/6	1



# Scans to the HEFT Wilson coefficients

- Limits to  $c_{hhh} = \kappa_\lambda$  are already set by the standard  $HH \rightarrow b\bar{b}\gamma\gamma$  analysis.
- In HEFT, the Higgs field is a singlet and therefore  $c_{ggh} - c_{gghh}$  and  $c_{thh} - c_{tthh}$  are independent:
  - Alternatively, in SMEFT, the Higgs field is a doublet (SM like) and  $c_{ggh} - c_{gghh}$  and  $c_{thh} - c_{tthh}$  are dependent.
- The  $c_{thh} = \kappa_t$  and  $c_{ggh}$  Wilson coefficients are better constrained by single Higgs analyses.
- $HH$  is uniquely sensitive to the  $c_{tthh}$  and  $c_{gghh}$  Wilson coefficients.

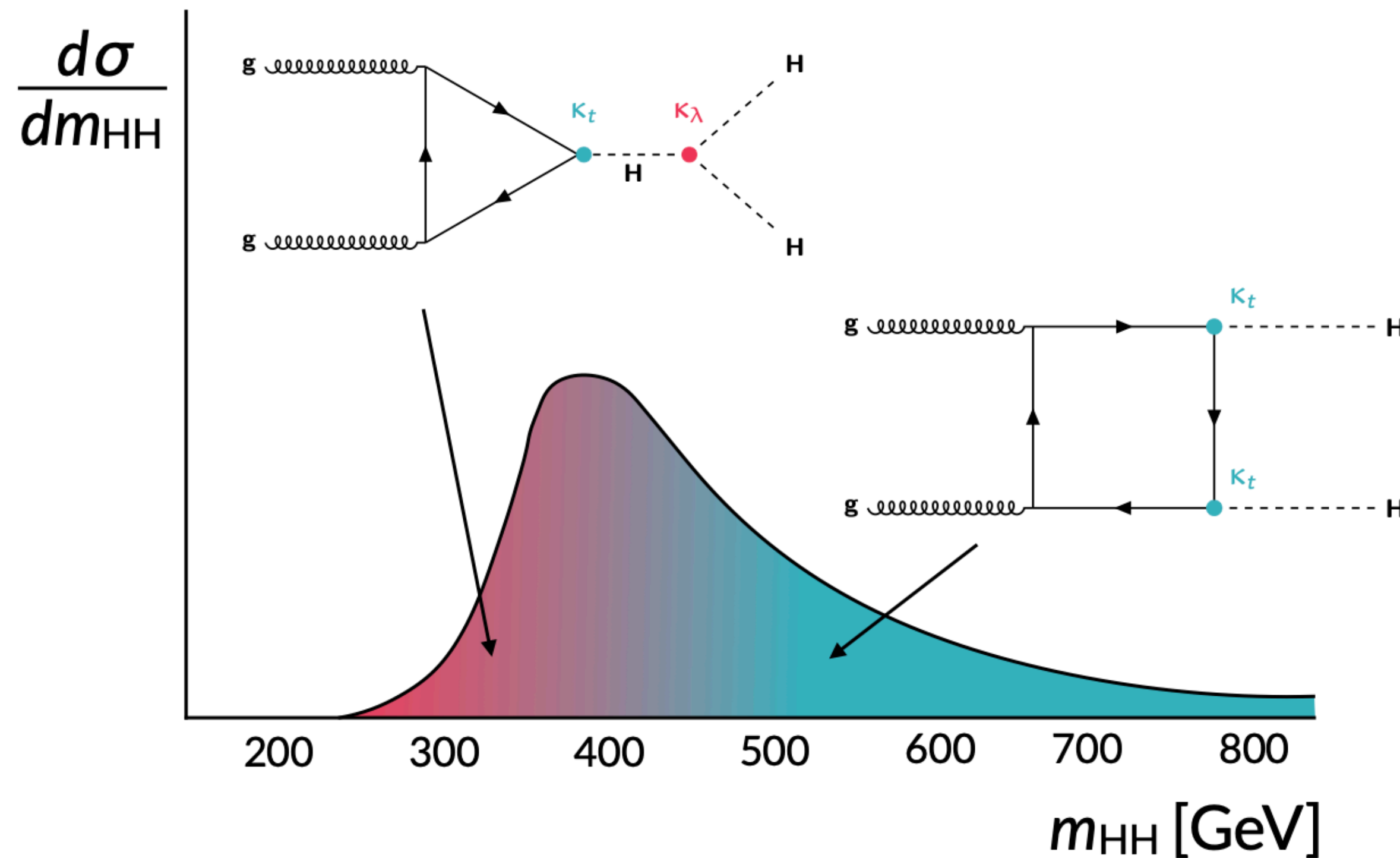


# Signal reweighting and uncertainty

- A reweighting of the SM ggF  $HH$  signal is used to generate the ggF HEFT signals ([arxiv:1806.05162](https://arxiv.org/abs/1806.05162)).

$$w_{EFT} = \frac{\sigma(m_{hh})}{\sigma^{SM}(m_{hh})} = \frac{\text{Poly}(m_{hh} | c_{hhh}, c_t, c_{tt}, c_{ggh}, c_{gghh})}{\text{Poly}(m_{hh} | 1, 1, 0, 0, 0)}$$

- An uncertainty on the selection efficiency is estimated from simulated events at truth level.



Channel	$b\bar{b}\gamma\gamma$	
Region	LM	HM
BM 1	4.9%	0.4%
BM 2	2.4%	4.2%
BM 3	6.1%	4.7%
BM 4	5.3%	1.4%
BM 5	10.8%	2.5%
BM 6	2.8%	1.3%
BM 7	0.3%	2.3%

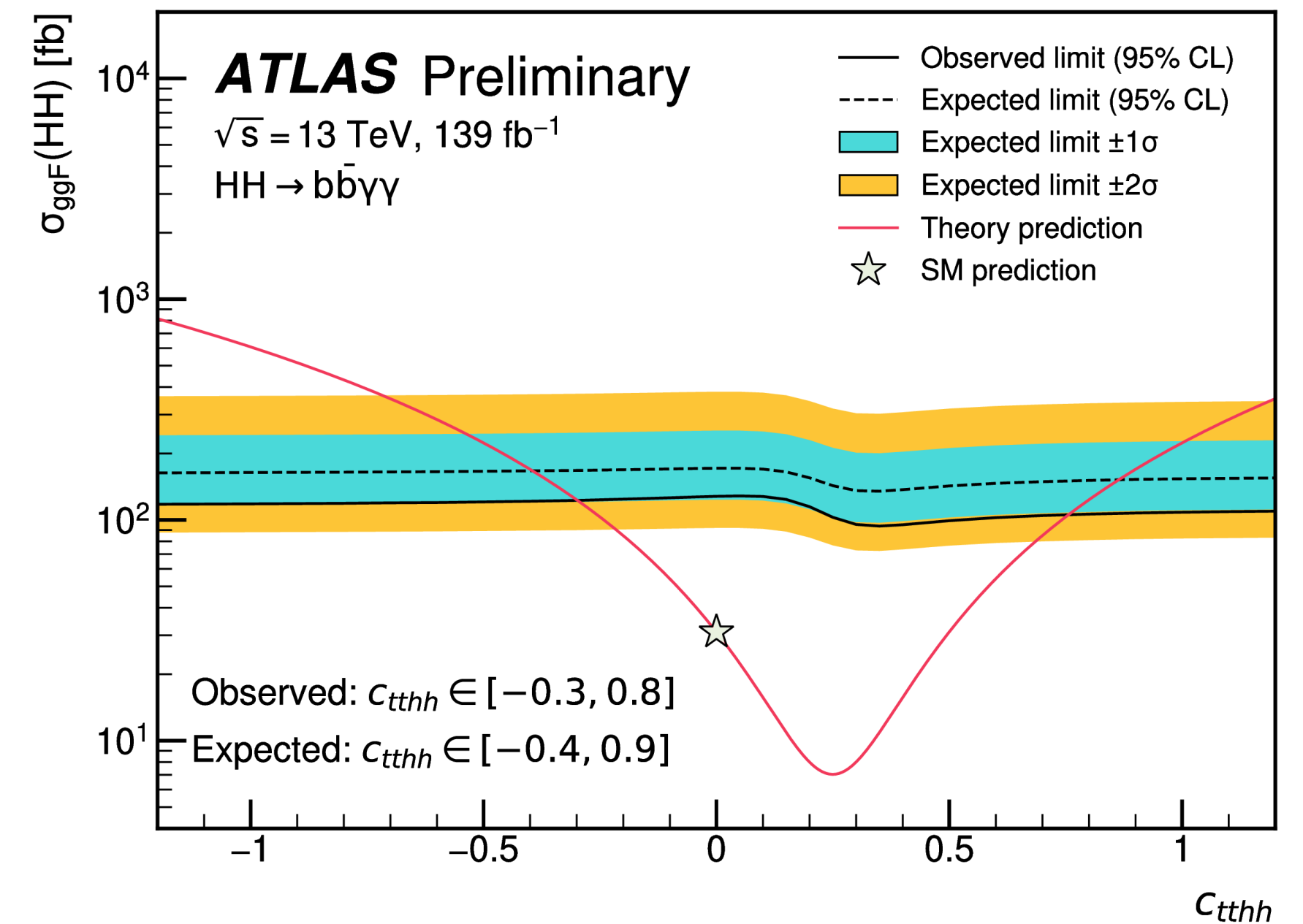
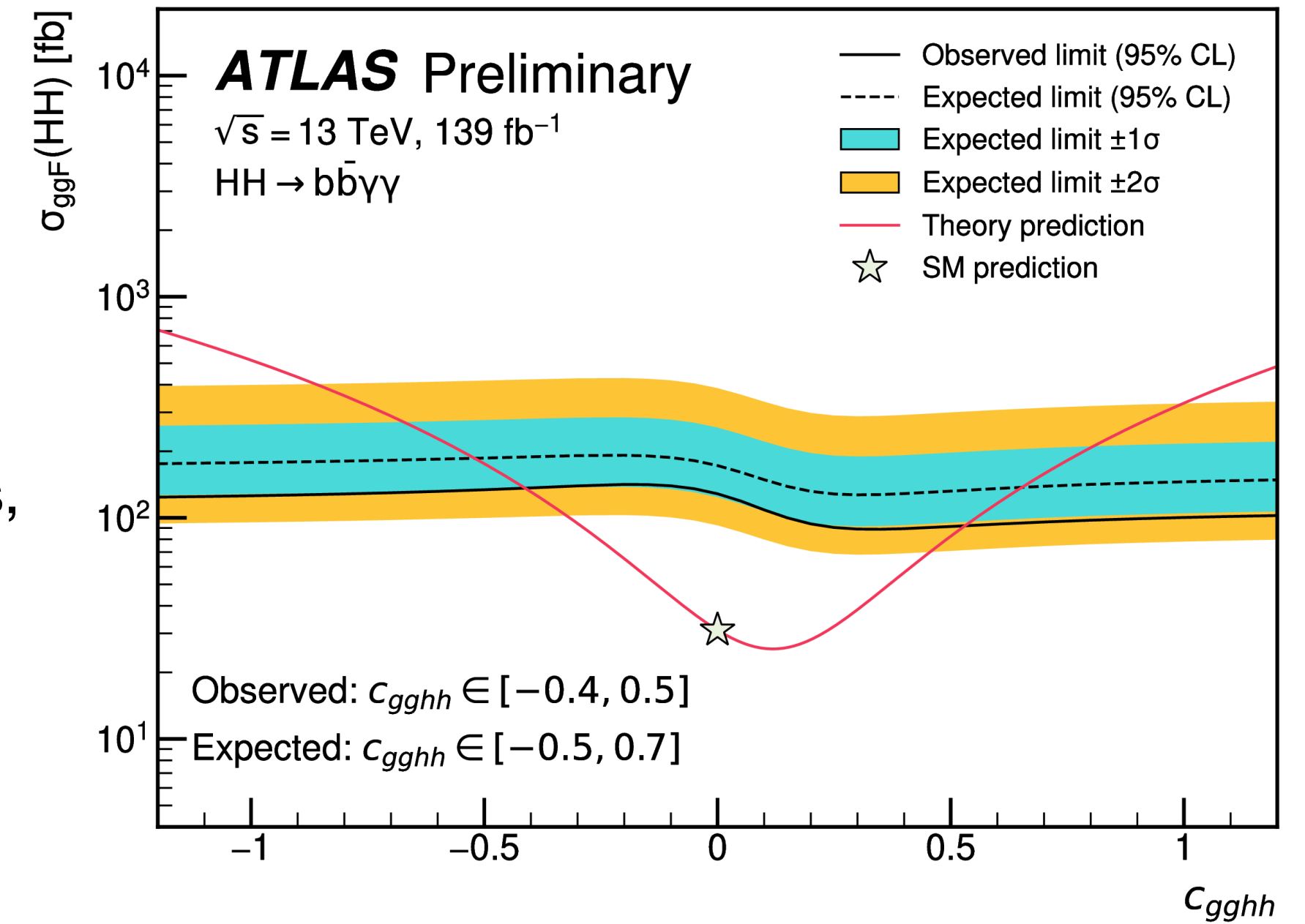
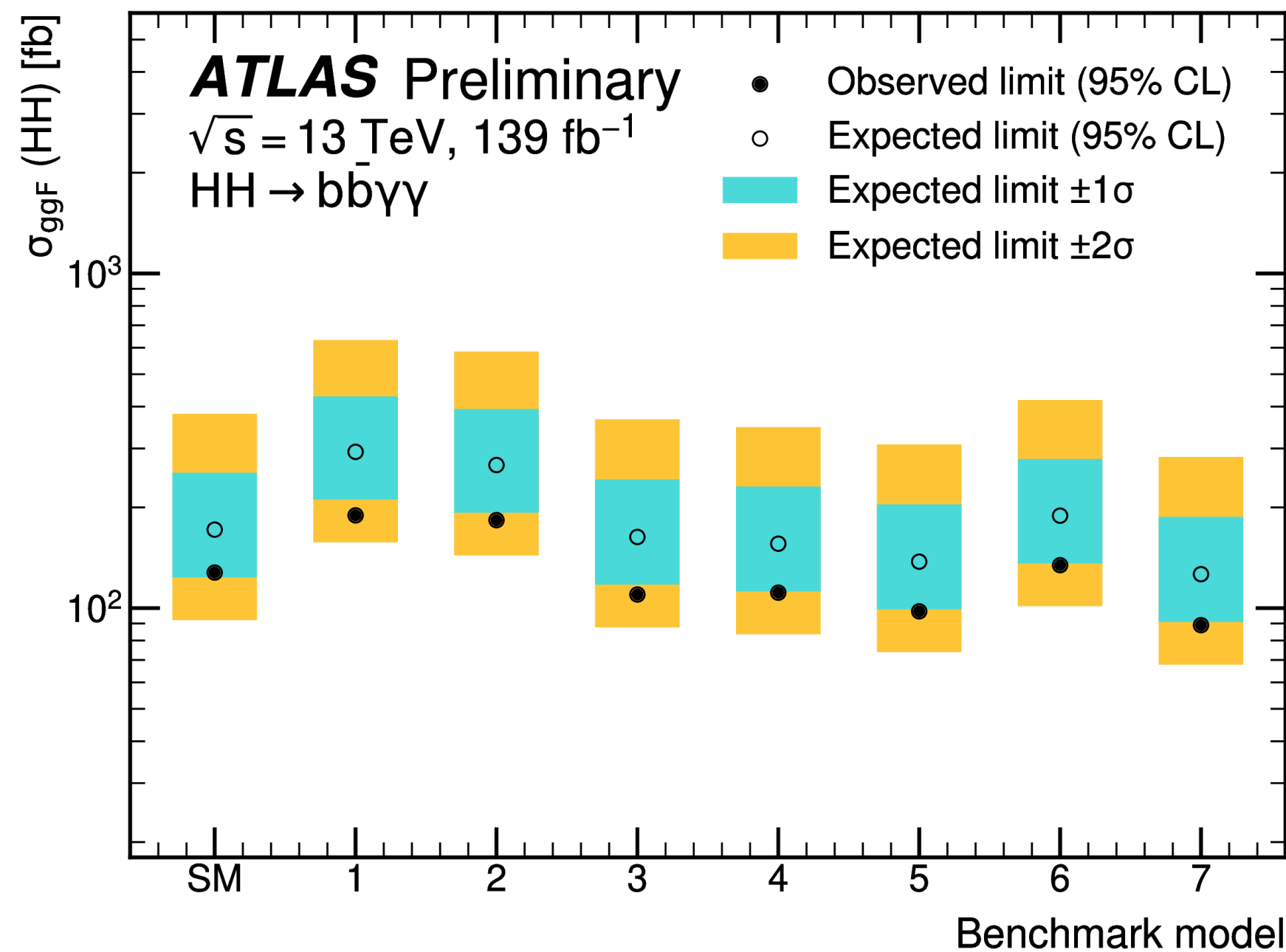
Channel	$b\bar{b}\gamma\gamma$	
Region	LM	HM
$c_{gghh}$	9.8%	6.2%
$c_{tthh}$	8.7%	1.6%

- The effect on the reweighting on the  $m_{\gamma\gamma}$  shape is found to be negligible.



# HEFT interpretations results

- The SM VBF  $HH$  process is treated as background.
- Limits on  $\sigma_{ggF}(HH)$  are set at 95% CL on the SM, the 7 shape benchmarks, the  $c_{tthh}$  and the  $c_{gghh}$  scans.





# HEFT interpretation of $bb\gamma\gamma$ , $bb\tau\tau$ and their combination

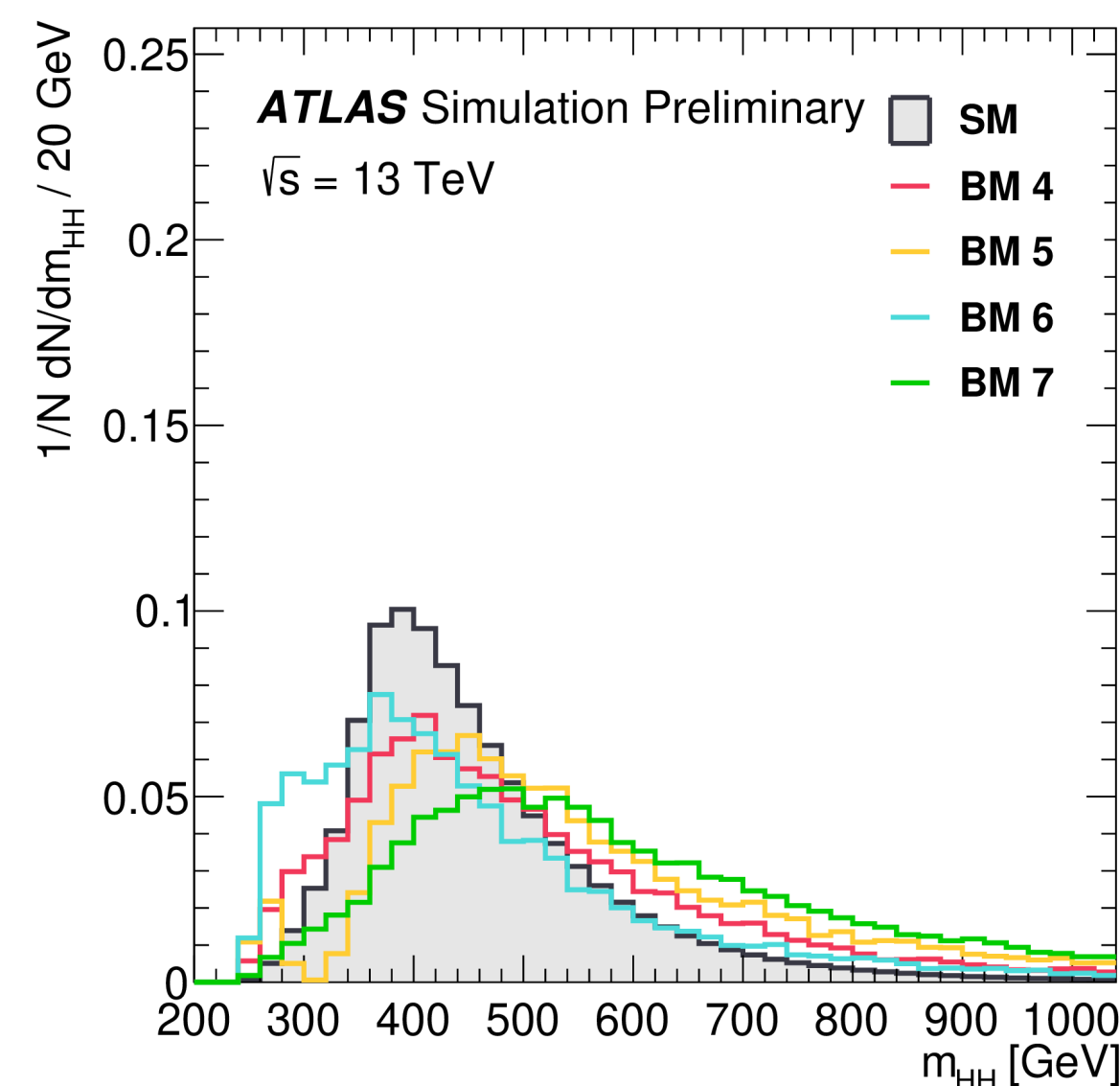
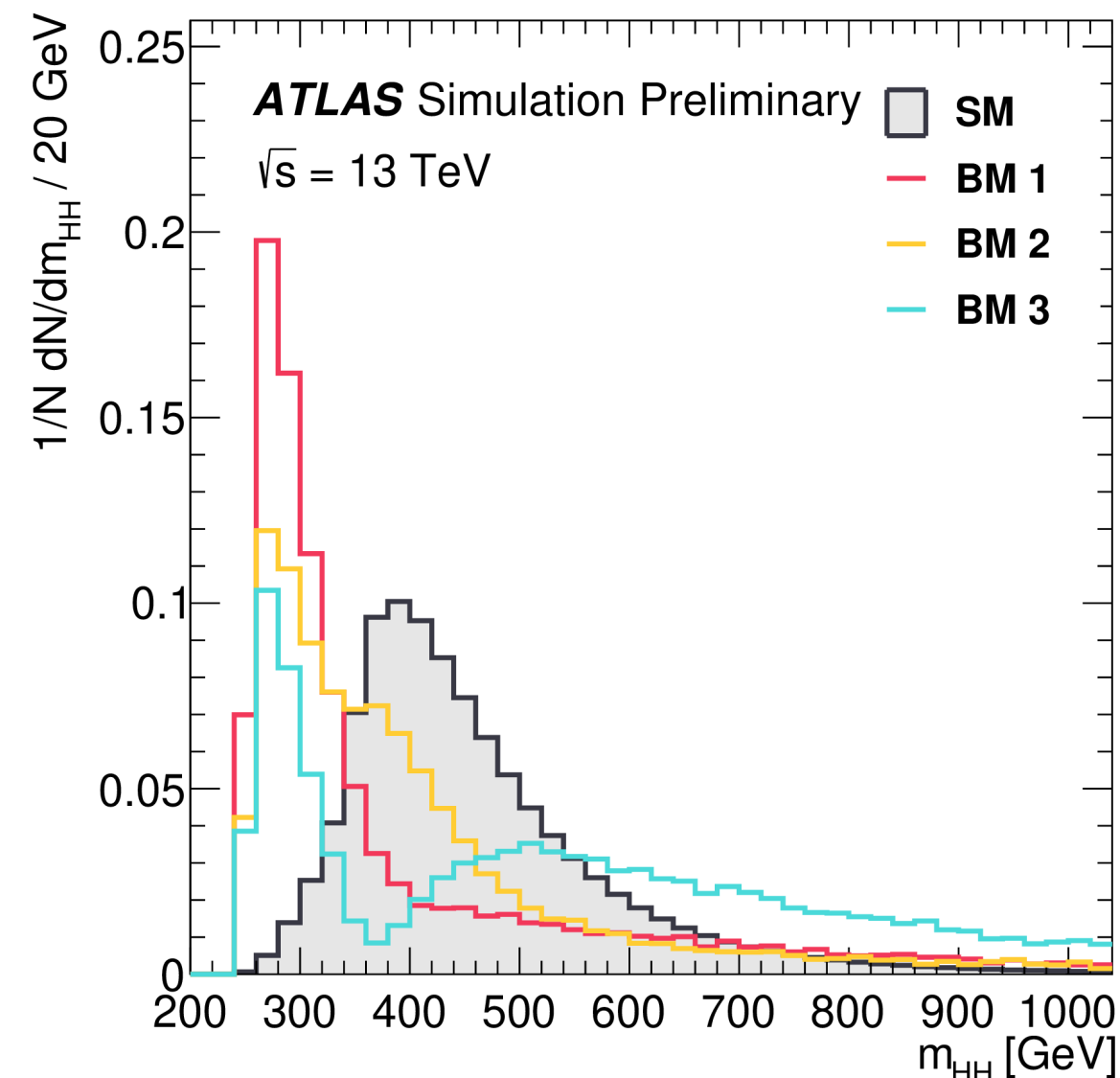
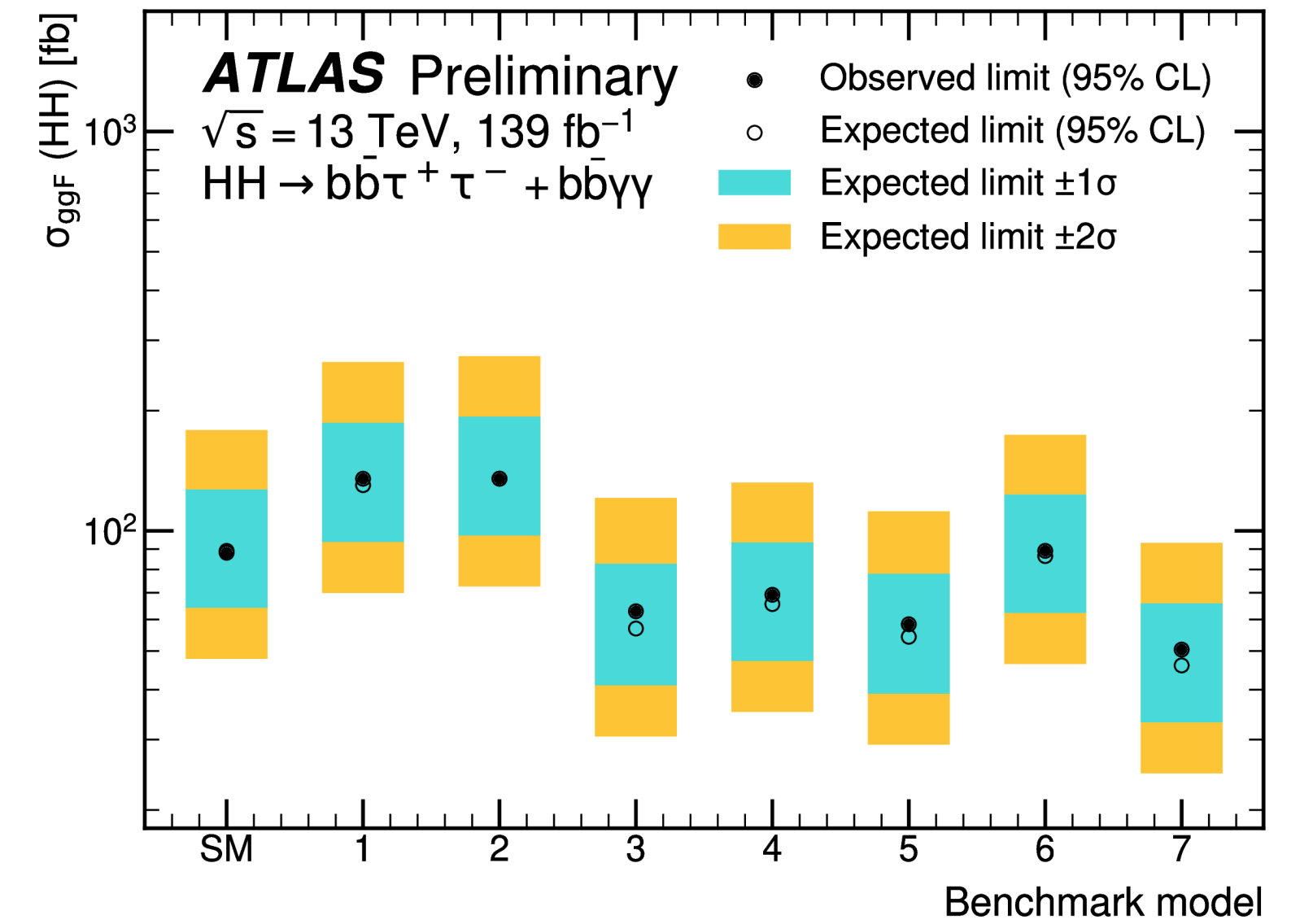
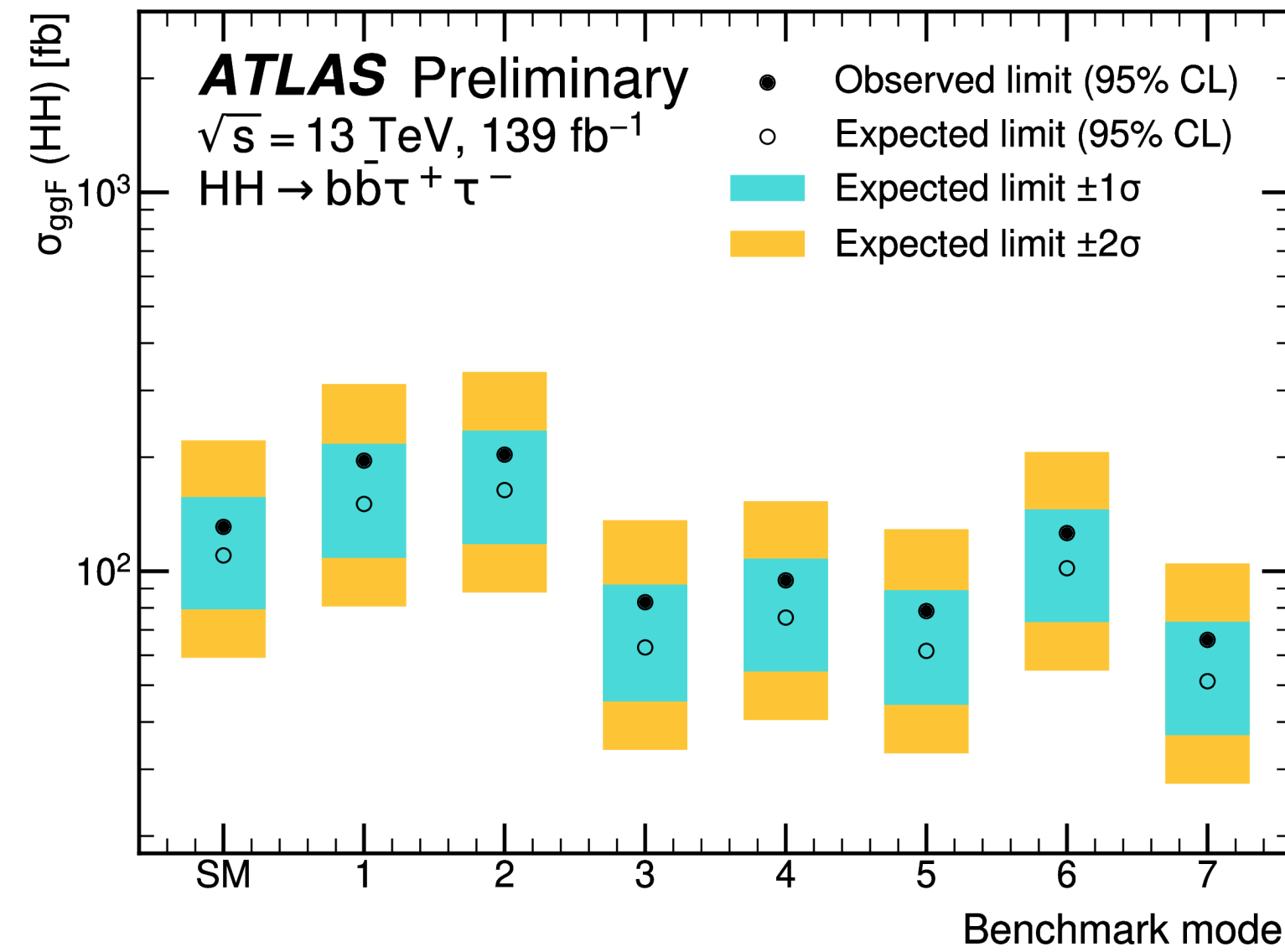
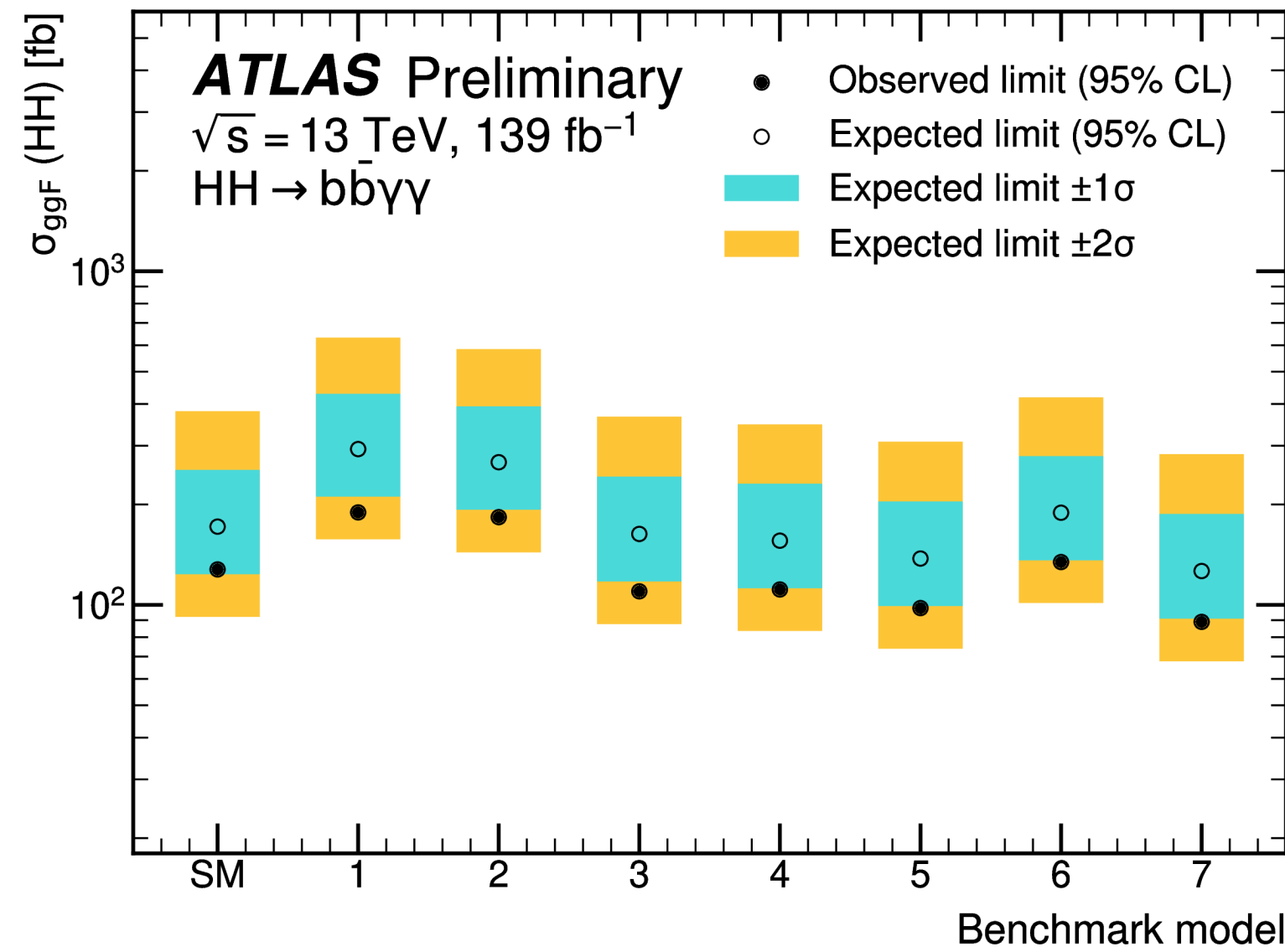


Important Swedish contribution from Uppsala and Stockholm University

Upper limits at 95% CL on the ggF $HH$ cross-section [fb]										
Benchmark	$b\bar{b}\gamma\gamma$		$b\bar{b}\tau^+\tau^-$		Combination					
	Obs.	Exp.	Obs.	Exp.	Obs.	$-2\sigma$	$-1\sigma$	Exp.	$+1\sigma$	$+2\sigma$
SM	127.7	171.6	130.9	110.1	88.1	47.8	64.2	89.1	126.9	178.8
BM 1	189.3	293.0	195.8	150.5	135.0	69.9	93.8	130.2	186.3	264.7
BM 2	183.2	267.7	203.1	163.7	135.1	72.5	97.4	135.1	193.1	273.8
BM 3	109.8	163.2	82.8	62.9	62.9	30.6	41.0	56.9	82.7	120.8
BM 4	111.2	155.8	94.6	75.4	69.2	35.2	47.2	65.6	93.5	132.0
BM 5	97.8	137.7	78.5	61.6	58.4	29.1	39.1	54.3	78.1	111.9
BM 6	134.4	188.9	126.1	101.8	89.1	46.4	62.3	86.5	123.3	174.0
BM 7	88.9	126.3	65.9	51.2	50.4	24.7	33.2	46.0	65.8	93.3

Wilson coefficient	$b\bar{b}\gamma\gamma$		$b\bar{b}\tau^+\tau^-$		Combination	
	Obs.	Exp.	Obs.	Exp.	Obs.	Exp.
$c_{gghh}$	$[-0.4, 0.5]$	$[-0.5, 0.7]$	$[-0.4, 0.4]$	$[-0.4, 0.4]$	$[-0.3, 0.4]$	$[-0.3, 0.3]$
$c_{tthh}$	$[-0.3, 0.8]$	$[-0.4, 0.9]$	$[-0.3, 0.7]$	$[-0.2, 0.6]$	$[-0.2, 0.6]$	$[-0.2, 0.6]$

# HEFT interpretation of $bb\gamma\gamma$ , $bb\tau\tau$ and their combination



- The tightest limits are set for BM 7 (high  $m_{HH}$ ) while the less stringent limits are set to BM 1 and 2 (low  $m_{HH}$ ).
- The  $bb\tau\tau$  channel has better sensitivity, specially to high  $m_{HH}$ .

# Conclusions

- **These results were presented at [Moriond 22!](#)**
- This is the first public result of  $HH$  EFTs in ATLAS
- The pub note can be found at [ATL-PHYS-PUB-2022-019](#)



**ATLAS PUB Note**  
ATL-PHYS-PUB-2022-019  
18th March 2022



**HEFT interpretations of Higgs boson pair searches  
in  $b\bar{b}\gamma\gamma$  and  $b\bar{b}\tau\tau$  final states and of their  
combination in ATLAS**

The ATLAS Collaboration

# Backup

# $HH$ decay modes

- Due to the large branching ratio (BR), most searches require at least one  $H \rightarrow b\bar{b}$ . Different decay modes of the second Higgs are considered.

Targeted  $HH$  decays shown today

	$b\bar{b}$	$WW$	$\tau\tau$	$ZZ$	$\gamma\gamma$
$b\bar{b}$	33%				
$WW$	25%	4.6%			
$\tau\tau$	7.4%	2.5%	0.39%		
$ZZ$	3.1%	1.2%	0.34%	0.076%	
$\gamma\gamma$	0.26%	0.10%	0.029%	0.013%	0.0005%

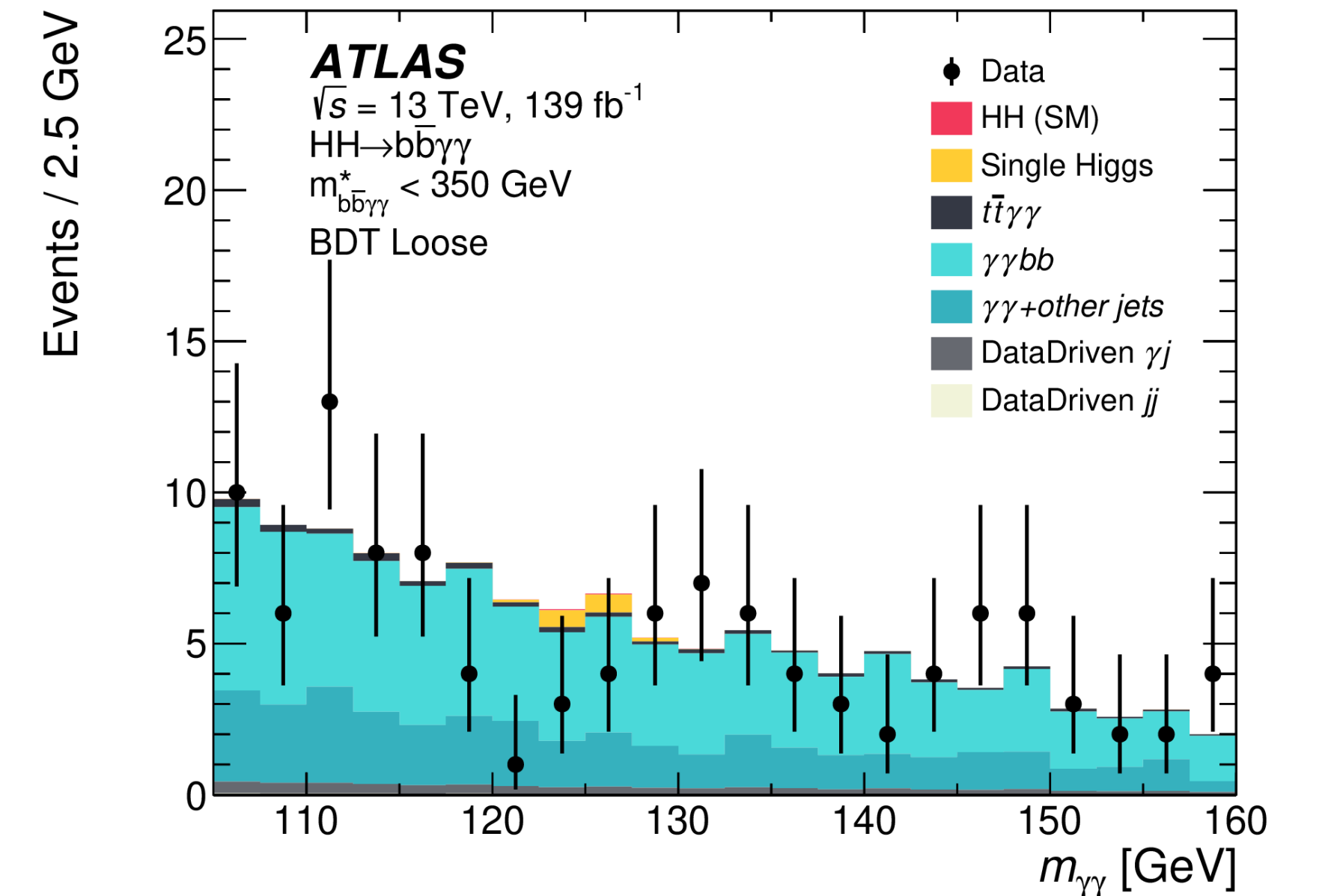
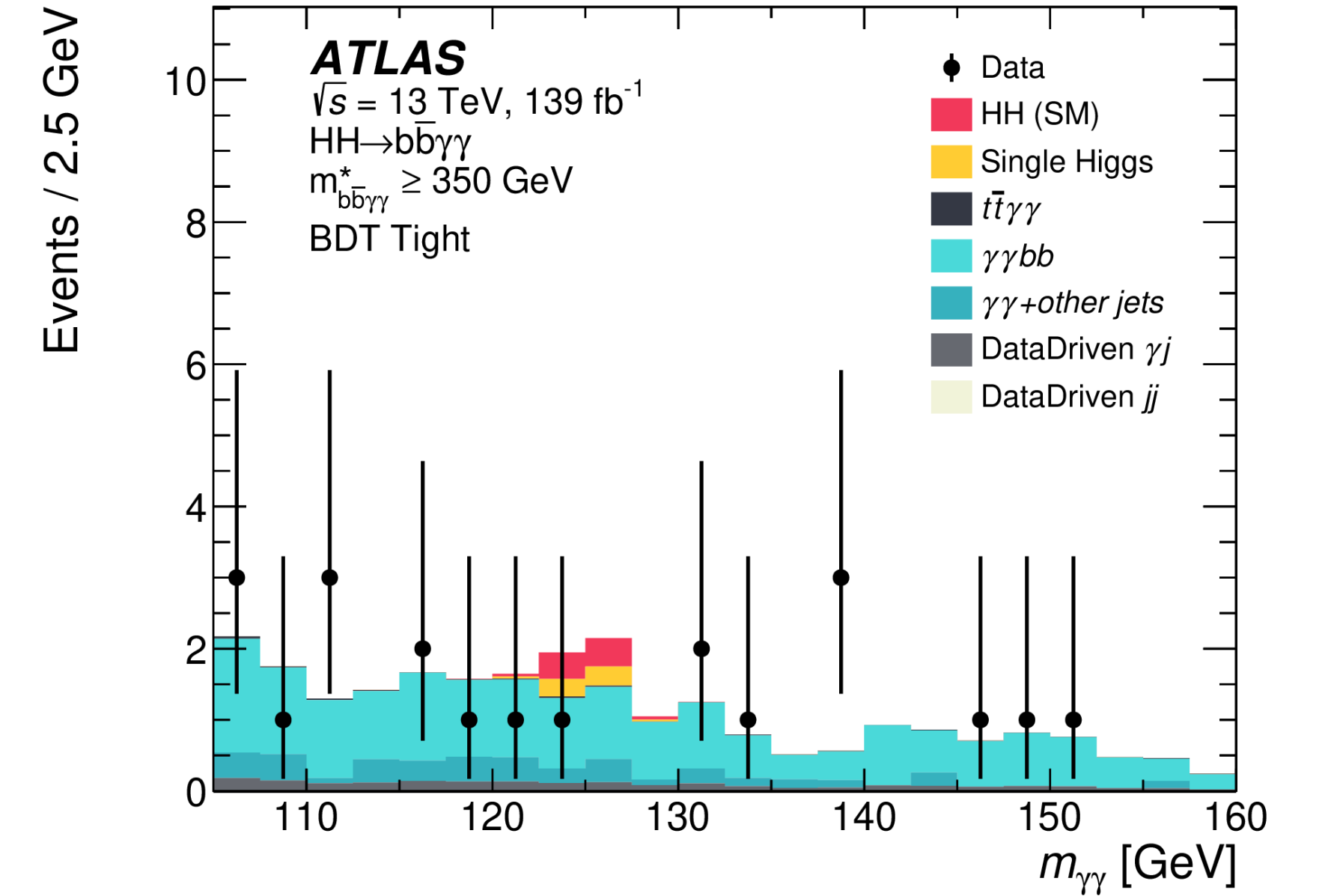
- ATLAS and CMS searches with full run 2 data for the following decay modes are presented:
  - $b\bar{b}b\bar{b}$  has the largest BR but large backgrounds arising from multijet production are challenging.
  - $b\bar{b}WW$ ,  $b\bar{b}ZZ$  and  $b\bar{b}\tau\tau$  have smaller BRs and can benefit from using leptons for triggering (hadronic  $b\bar{b}\tau\tau$  searches won't be presented).
  - $b\bar{b}\gamma\gamma$  has the smallest BR but it's a very sensitive analysis thanks to the clean  $m_{\gamma\gamma}$  resolution.
- Other final states without any  $H \rightarrow b\bar{b}$  are also included in the combinations with partial run 2 data.



# $HH \rightarrow b\bar{b}\gamma\gamma$ table

Table 1: The number of data events observed in the  $120 < m_{\gamma\gamma} < 130$  GeV window, the number of  $HH$  signal events expected for  $\kappa_\lambda = 1$  and for  $\kappa_\lambda = 10$ , and events expected for single Higgs boson production (estimated using MC simulation), as well as for continuum background. For the single Higgs boson, “Rest” includes the following production modes: VBF,  $WH$ ,  $tHq$ , and  $tHW$ . The values are obtained from a fit of the Asimov data set generated under the SM signal-plus-background hypothesis,  $\kappa_\lambda = 1$ . The continuum background component of the Asimov data set is obtained from the fit of the data sideband. The uncertainties in  $HH$  signals and single Higgs boson background include the systematic uncertainties discussed in Section ???. The uncertainty in the continuum background is given by the sum in quadrature of the statistical uncertainty from the fit to the data and the spurious-signal uncertainty.

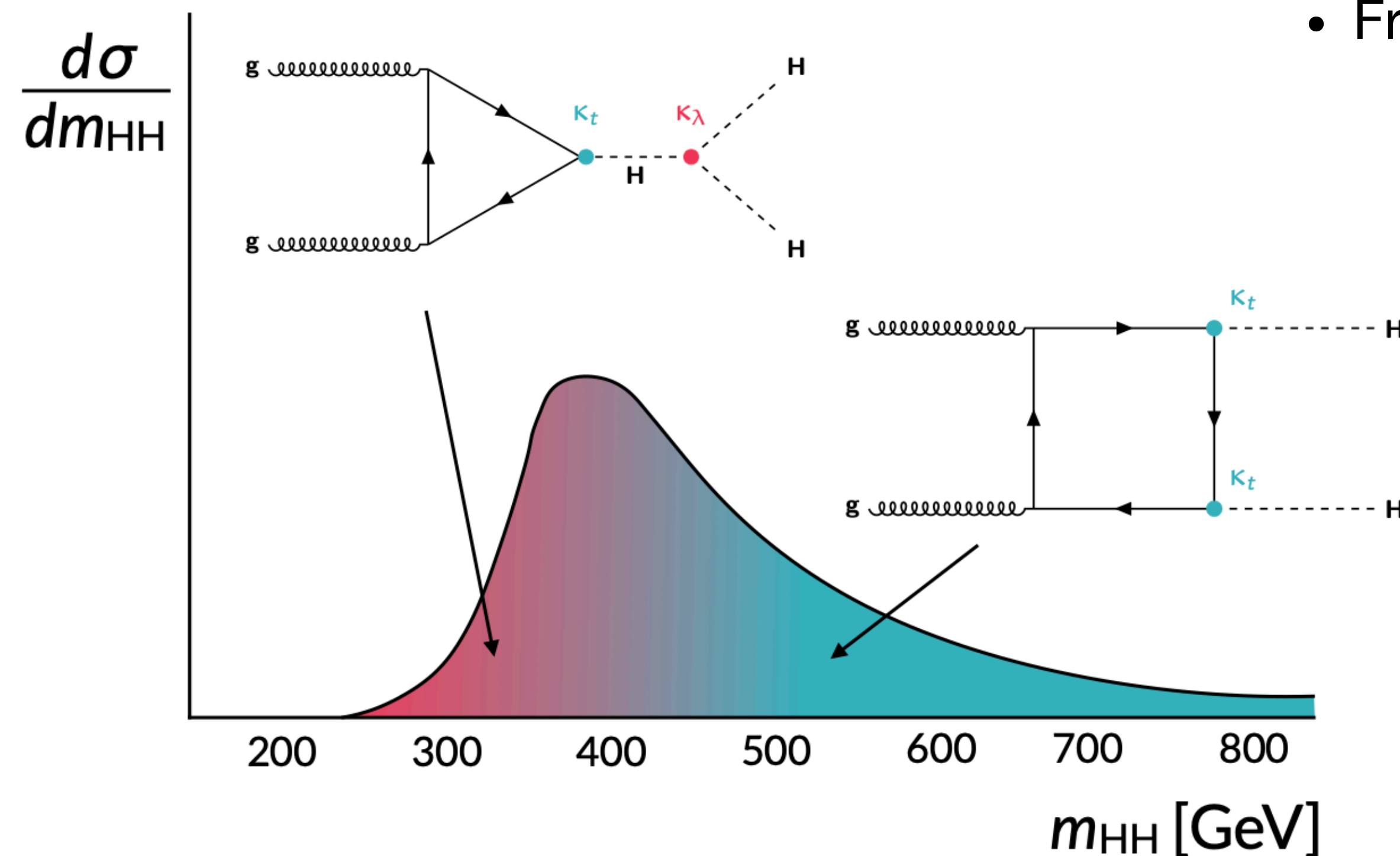
	High mass BDT tight	High mass BDT loose	Low mass BDT tight	Low mass BDT loose
Continuum background	$4.9^{+1.1}_{-1.3}$	$9.5^{+1.5}_{-1.7}$	$3.7^{+0.9}_{-1.1}$	$24.9^{+2.3}_{-2.5}$
Single Higgs boson background	$0.67^{+0.29}_{-0.13}$	$1.6^{+0.6}_{-0.2}$	$0.23^{+0.09}_{-0.03}$	$1.40^{+0.33}_{-0.16}$
$ggF+bbH$	$0.26^{+0.28}_{-0.16}$	$0.4^{+0.5}_{-0.2}$	$0.07^{+0.08}_{-0.04}$	$0.27^{+0.27}_{-0.16}$
$t\bar{t}H$	$0.19^{+0.03}_{-0.03}$	$0.49^{+0.09}_{-0.07}$	$0.107^{+0.022}_{-0.017}$	$0.75^{+0.13}_{-0.11}$
$ZH$	$0.142^{+0.035}_{-0.025}$	$0.48^{+0.09}_{-0.07}$	$0.040^{+0.020}_{-0.014}$	$0.27^{+0.06}_{-0.04}$
Rest	$0.074^{+0.032}_{-0.014}$	$0.16^{+0.07}_{-0.03}$	$0.012^{+0.008}_{-0.004}$	$0.111^{+0.030}_{-0.012}$
SM $HH(\kappa_\lambda = 1)$ signal	$0.87^{+0.10}_{-0.18}$	$0.37^{+0.04}_{-0.07}$	$0.049^{+0.006}_{-0.010}$	$0.078^{+0.008}_{-0.015}$
$ggF$	$0.86^{+0.10}_{-0.18}$	$0.35^{+0.04}_{-0.07}$	$0.046^{+0.006}_{-0.010}$	$0.072^{+0.008}_{-0.015}$
VBF	$(12.6^{+1.3}_{-1.2}) \cdot 10^{-3}$	$(16.1^{+1.4}_{-1.2}) \cdot 10^{-3}$	$(3.2^{+0.4}_{-0.4}) \cdot 10^{-3}$	$(6.9^{+0.5}_{-0.6}) \cdot 10^{-3}$
Alternative $HH(\kappa_\lambda = 10)$ signal	$6.5^{+1.0}_{-0.8}$	$3.6^{+0.6}_{-0.4}$	$4.5^{+0.7}_{-0.6}$	$8.5^{+1.3}_{-1.0}$
Data	2	17	5	14





# Signal reweighting

- The  $\sigma_{HH}$  depends only on  $m_{hh}$  and can be parametrised as a polynomial  $Poly(A) = \sum_{i=1}^N A_i * c_j$ , where  $A_i$  are some coefficients obtained from a fit and  $c_j$  are combinations of the 5 EFT couplings that play a role in HH production.

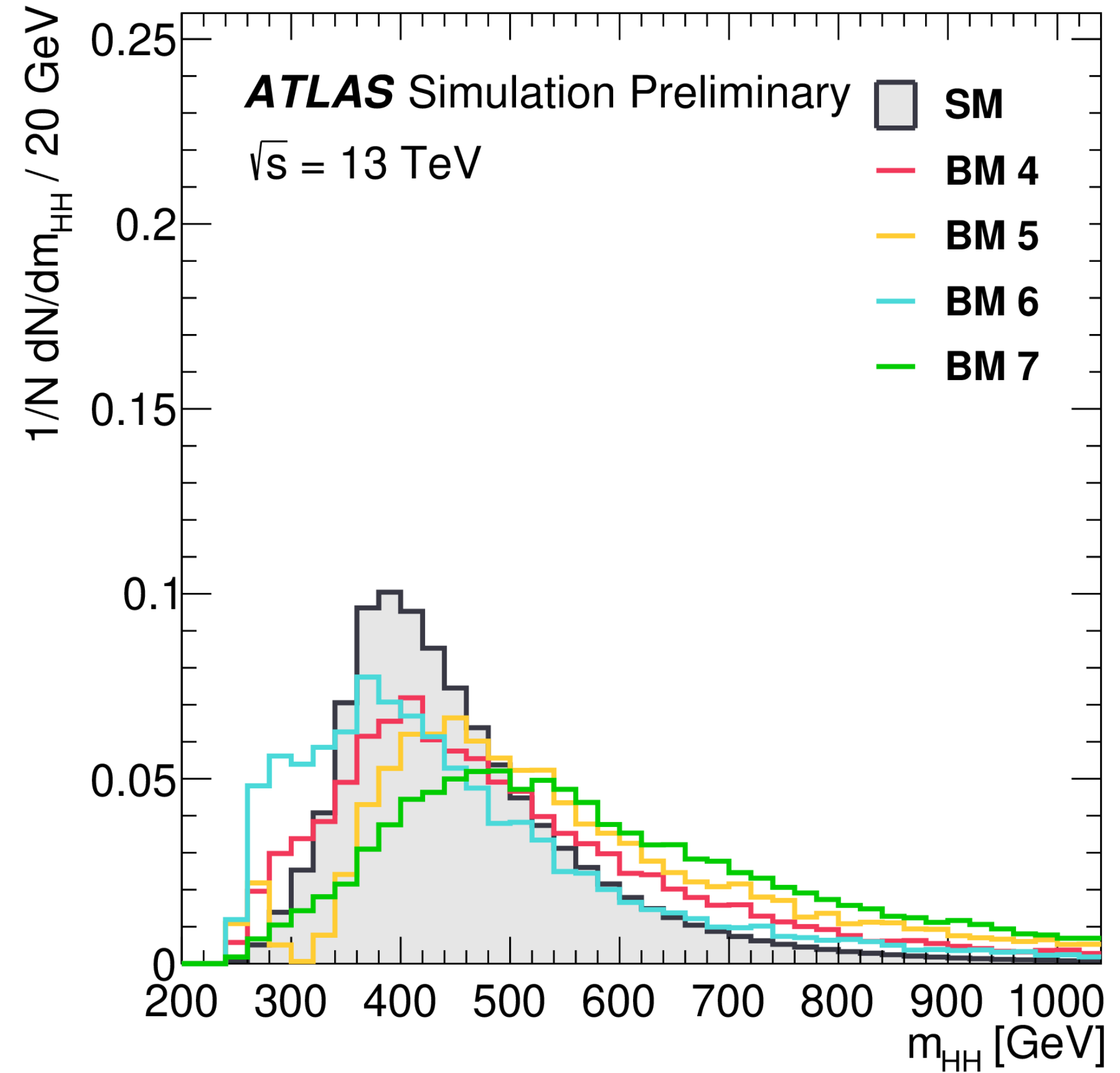
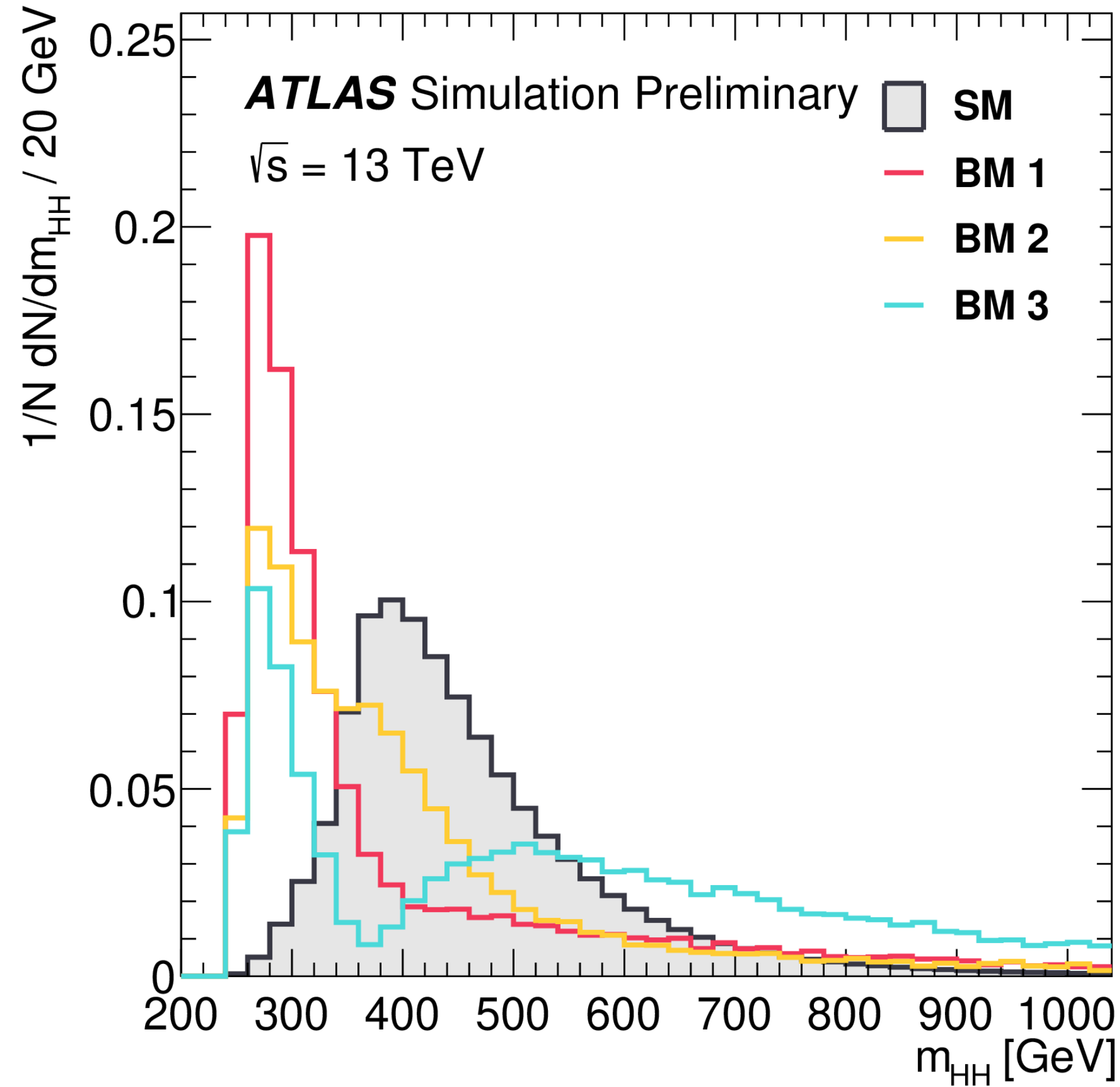


- From [arxiv:1806.05162](https://arxiv.org/abs/1806.05162):  $A_{1-23}$  coefficients to predict  $\sigma_{HH}^{NLO}$

$$w_{EFT} = \frac{\sigma(m_{hh})}{\sigma^{SM}(m_{hh})} = \frac{Poly(m_{hh} | c_{hhh}, c_t, c_{tt}, c_{ggh}, c_{gg hh})}{Poly(m_{hh} | 1, 1, 0, 0, 0)}$$

→ We can re-weight SM to any HEFT BSM point!

# Benchmarks



# Signal selection efficiency

Acceptance $\times$ Efficiency [%]	HM Loose	LM Loose	HM Tight	LM Tight	Total
SM	3.2	0.6	7.7	0.4	11.9
BM 1	1.3	2.9	3.8	1.5	9.5
BM 2	1.8	2.2	4.5	1.2	9.7
BM 3	2.2	1.3	8.3	0.6	12.4
BM 4	2.9	0.7	8.6	0.4	12.6
BM 5	3.1	0.3	9.8	0.1	13.3
BM 6	2.6	1.2	7.0	0.7	11.5
BM 7	3.1	0.3	10.8	0.2	14.4

