

Improving Bayesian parameter estimation with the latest RHIC and LHC data including a new initial conditions model

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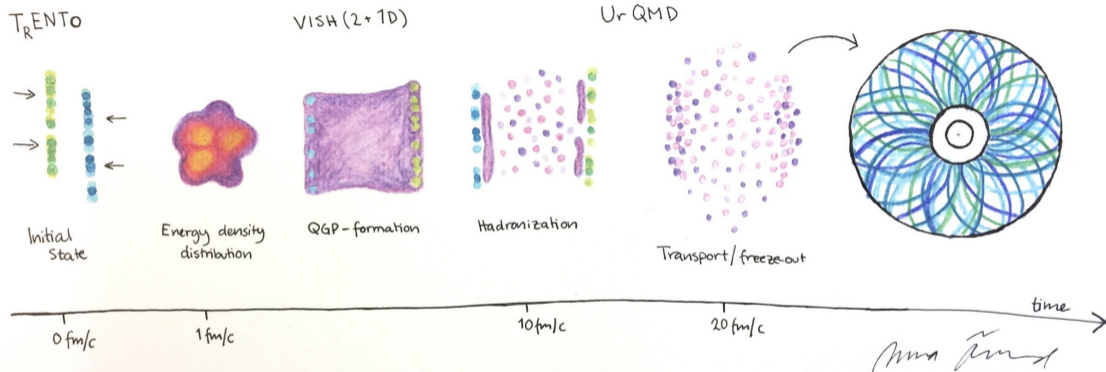
1. University of Jyväskylä, Finland

Wednesday 4th January, 2023

Spåtind 2023, Vinstra, Norway



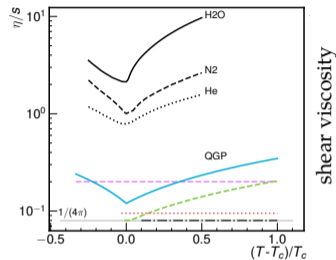
THE DIFFERENT STAGES OF HEAVY-ION COLLISIONS



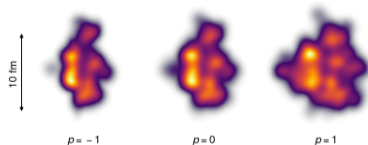
$$T^{\mu\nu} = e u^\mu u^\nu - (P + \Pi) \Delta_{\mu\nu} + \pi^{\mu\nu}, \quad \delta_\mu T^{\mu\nu} = 0$$

COLLECTION OF PARAMETERS

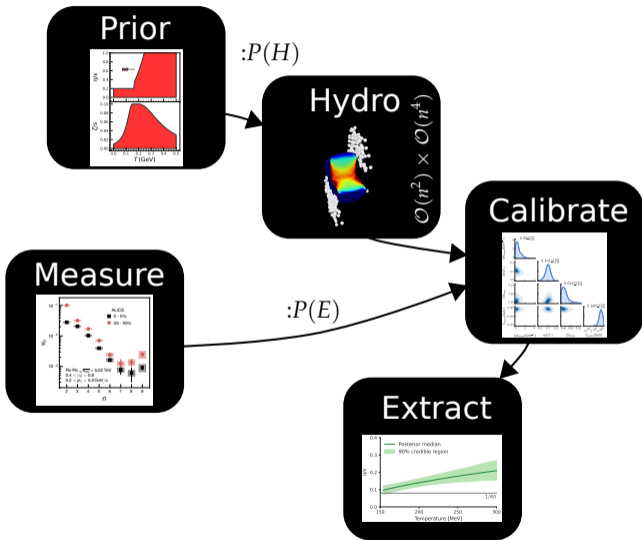
Parameter	Description
T_c	Temperature of const. $\eta/s(T)$, $T < T_c$
$\eta/s(T_c)$	Minimum $\eta/s(T)$
$(\eta/s)_{\text{slope}}$	Slope of $\eta/s(T)$ above T_c
$(\eta/s)_{\text{curve}}$	Curvature of $\eta/s(T)$ above T_c
$(\zeta/s)_{\text{peak}}$	Temperature of $\zeta/s(T)$ maximum
$(\zeta/s)_{\text{max}}$	Maximum $\zeta/s(T)$
$(\zeta/s)_{\text{width}}$	Width of $\zeta/s(T)$ peak
T_{switch}	Switching / particlization temperature
N(2.76 TeV)	Overall normalization (2.76 TeV)
N(5.02 TeV)	Overall normalization (5.02 TeV)
p	Entropy deposition parameter
w	Nucleon width
σ_k	Std. dev. of nucleon multiplicity fluctuations
d_{min}^3	Minimum volume per nucleon
τ_{fs}	Free-streaming time



Trento p-value, <http://qcd.phy.duke.edu/trento/>



BAYESIAN PARAMETER ESTIMATION



Bayes' theorem:

$$P(H|E) = \frac{P(E|H) \cdot P(H)}{P(E)}$$

$$P(E) = \sum_{i=1}^n P(E|H_i)P(H_i)$$

- Find optimal set of model parameters that best reproduce the experimental data.
- Utilize constraints, such as flow observables, to help narrow down the $\eta/s(T)$ and such.

Testing a single set of parameters requires $\mathcal{O}(10^4)$ hydro events, and evaluating eight different parameters five times each requires $5^8 \times 10^4 \approx 10^9$ hydro events.

That's roughly 10^5 CPU years!

OUR ARSENAL OF OBSERVABLES - STOCHASTIC APPROACH

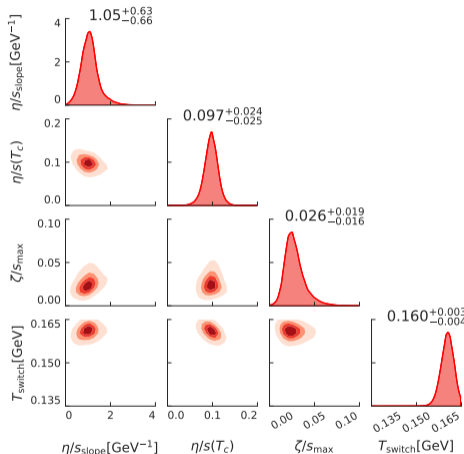
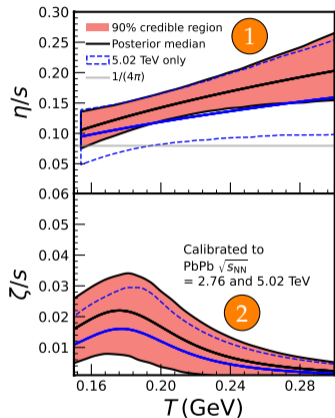
- Together, various flow observables cover the sensitivity for all components of transport properties.

Name	Symbol	Measure	Sensitivity-stochastic approach
Flow coefficients	v_n	System expansion and anisotropy of the flow	Average $\langle \eta/s \rangle$ and $\zeta/s(T)$ peak
(Normalized) Symmetric cumulants	(N)SC(k, l, m)	Correlation between magnitudes of flow harmonics	$\eta/s(T)$ temperature dependence
Non-linear flow mode coefficients	$\chi_{n,mk}$	Quantification of the non-linear response	$\eta/s(T)$ at the freeze-out
Symmetry-plane correlations	$\rho_{n,mk}$	Correlations between the directions of flow harmonics	$\eta/s(T)$

Thanks to excellent ALICE papers over years:

- Phys.Rev.Lett. 117 (2016) 182301, Phys.Lett. B773 (2017) 68, Phys.Rev. C 97 (2018) 024906, JHEP05 (2020) 085, Phys.Lett. B818 (2021) 136354, Phys.Rev.Lett. 127 (2021) 092302 - [flow](#)
- Phys.Rev.Lett. 106 (2011) 032301, Phys.Rev.C 88 (2013) 044910, Phys.Lett. B772 (2017) 567-577, Phys.Rev.C 101, 044907 (2020) - [N_{ch} and \$\langle p_T \rangle\$](#)

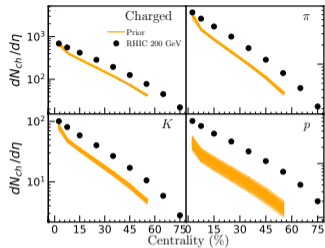
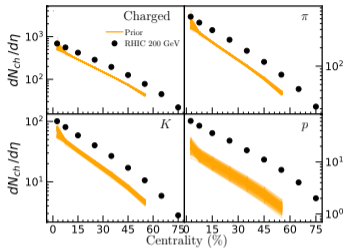
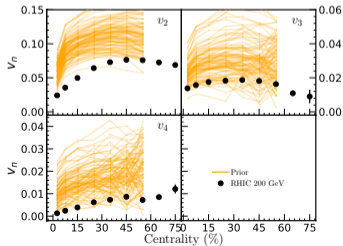
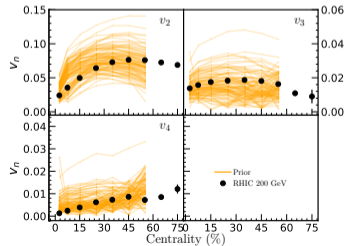
RESULTS: JYVASKYLA (2022) – COMBINED COLLISION ENERGY ANALYSIS (2.76 + 5.02 TeV)



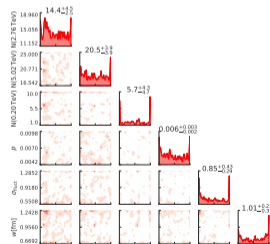
- 1 Significantly improved $\eta/s(T)$ uncertainty
- 2 Non-zero $\zeta/s(T)$
- 3 Overall better convergence for parameter components

Together with two collision energies and added observables, the uncertainty has reduced!

INCLUSION OF RHIC DATA

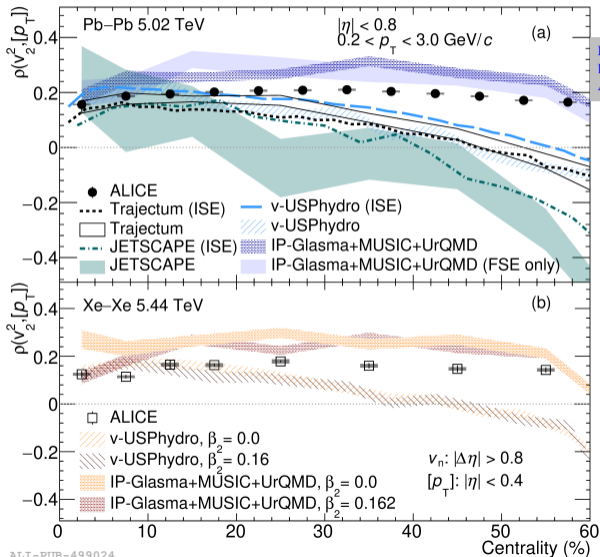
Fixed $w_{nucl} = 0.5$  $w_{nucl} = [0.67-1.24]$ 

- Included observables for RHIC 200 GeV data: $v_2 - v_4$, $\langle p_T \rangle$ and N_{ch} for charged particles and PID
- Fixed nuclen width $w = 0.5$ and now relaxed $[0.67-1.24]$
- small improvements for N_{ch} and better for v_n
- Posterior distributions don't converge



Preliminary

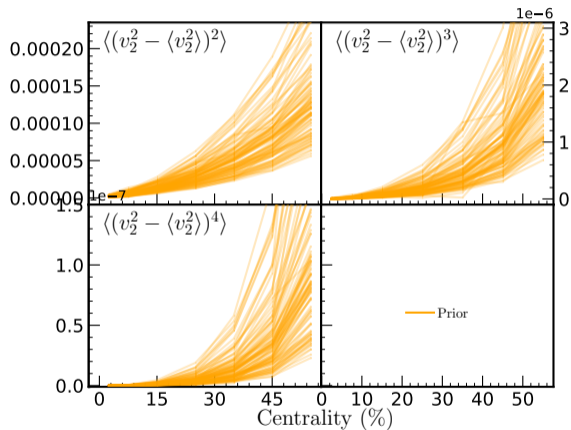
$v_n, [p_T]$ CORRELATION - SHORTAGE OF T_{RENT0} MODEL



P. Bozek, R. Samanta, Phys. Rev. C 102, 034905
 B. Schenke, C. Shen, D. Teaney, Phys. Rev. C 102, 034905
 ALICE, arXiv:2111.06106

$$\rho(v_2^2, [p_T]) = \frac{\langle \delta v_2^2 \delta [p_T] \rangle}{\sqrt{\langle (\delta v_2^2)^2 \rangle \langle (\delta [p_T])^2 \rangle}} \quad (1)$$

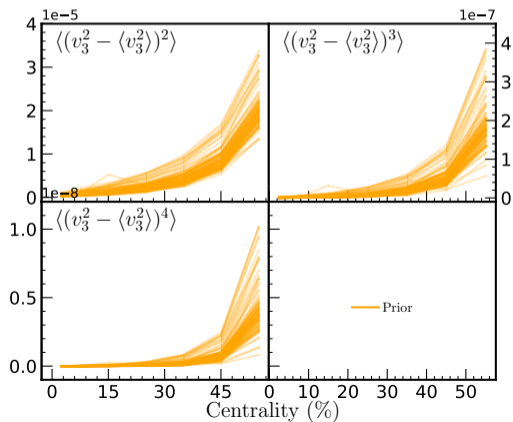
- Correlation between $[p_T]$ and v_2 :
 - can be used to differentiate initial state models
 - More peripheral \rightarrow best described by models with IP-Glasma
 - strong centrality dependence on the models with Trento
- Ongoing progress:
 - Calculate sensitivity
 - Adapt it to the Bayesian Analysis

MOMENTS OF δv_n 

- Characterizes the fluctuation of different order
- $|\eta| < 0.8$ and $0.2 \text{ GeV} < p_T < 5.0 \text{ GeV}$

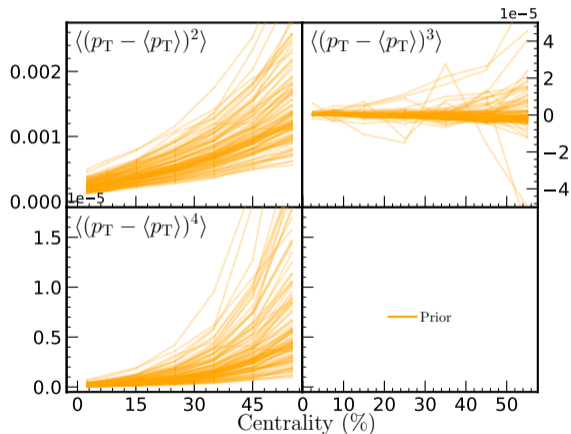
Moments of a distribution

Variance: $(X - \mu)^2$, **skewness:** $(X - \mu)^3$
and kurtosis: $(X - \mu)^4$



PRIOR DISTRIBUTION OF $\rho(v_n^2, [p_T])$

Moments of $[p_T]$ distribution

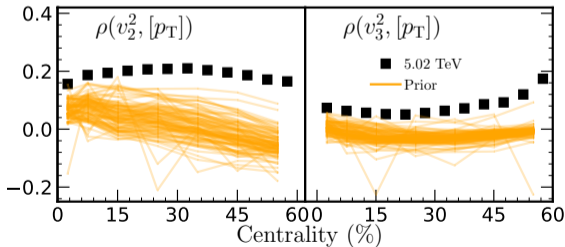


■ $|\eta| < 0.8$ and $0.2 \text{ GeV} < p_T < 5.0 \text{ GeV}$

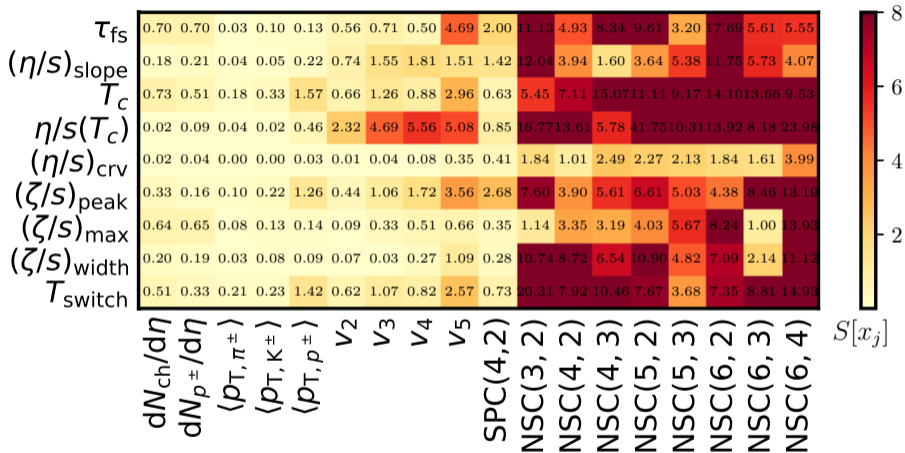
$v_n^2, [p_T]$ correlation

$$\rho(v_2^2, [p_T]) = \frac{\langle (v_2^2 - \langle v_2^2 \rangle) ([p_T] - \langle [p_T] \rangle) \rangle}{\sqrt{\langle (v_2^2 - \langle v_2^2 \rangle)^2 \rangle \langle ([p_T] - \langle [p_T] \rangle)^2 \rangle}}$$

- Clear centrality dependence
- Gains negative values



OBSERVABLE SENSITIVITIES



Sensitivity

$$S[x_j] = \Delta/\delta., \text{ where}$$

$$\Delta = \frac{|\hat{O}(\vec{x}') - \hat{O}(\vec{x})|}{|\hat{O}(\vec{x})|}$$

- NSCs most sensitive to multiple different parameters
- v_n s show sensitivity to specific shear viscosity

$$(\eta/s)(T) = (\eta/s)(T_c) + (\eta/s)_{slope}(T - T_c) \left(\frac{T}{T_c}\right)^{(\eta/s)_{curve}}$$

$$(\zeta/s)(T) = \frac{(\zeta/s)_{max}}{1 + \left(\frac{T - (\zeta/s)_{peak}}{(\zeta/s)_{width}}\right)^2}$$

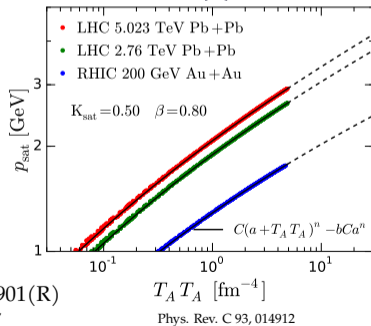
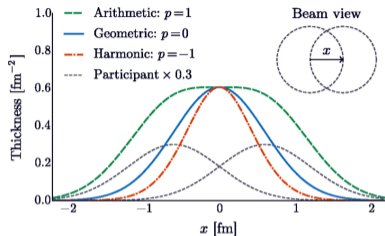
T_RENTo vs. EKRT: ONGOING

T_RENTo [1]

- Flexibility to produce some other models
- Unable to predict ($\sqrt{s_{NN}}$ - Cent) dependence
- Has six free parameters

EKRT [2]

- Only two free parameters, K_{sat} and β_{sat}
- ($\sqrt{s_{NN}}$ - Cent) dependence comes automatically from the gluon saturation and mini-jet production
- Computationally a bit heavier \rightarrow much improved (H. Hirvonen's talk)



[1]. J. S. Moreland, J. E. Bernhard, and S. A. Bass, PRC **92** 011901(R)

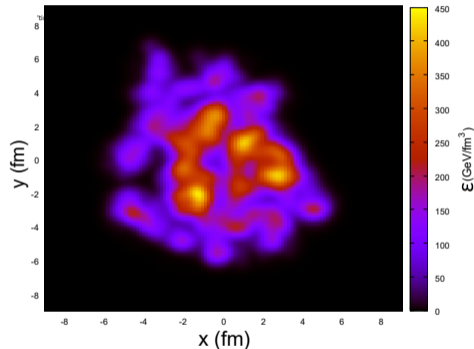
[2]. H. Niemi, K. J. Eskola, and R. Paatelainen, PRC **93** 024907

EKRT

- Low- p_T particles are produced by saturation
- Transverse energy profile is

$$\frac{dE_T}{d^2r}(p_{sat}, \sqrt{s_{NN}}, A, r, b) = \frac{K_{sat}}{\pi} p_{sat}^3 \Delta y$$
- The resulting energy density

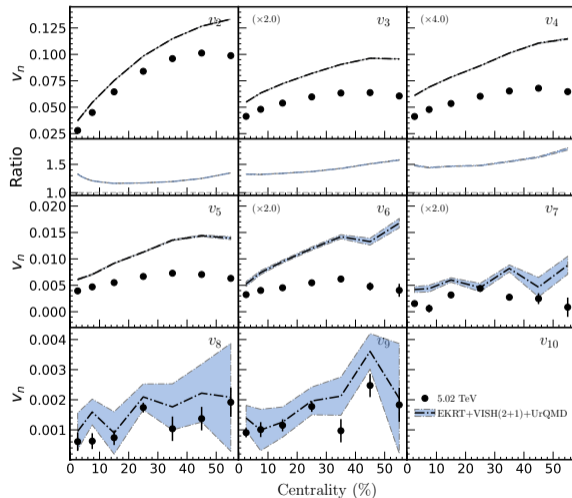
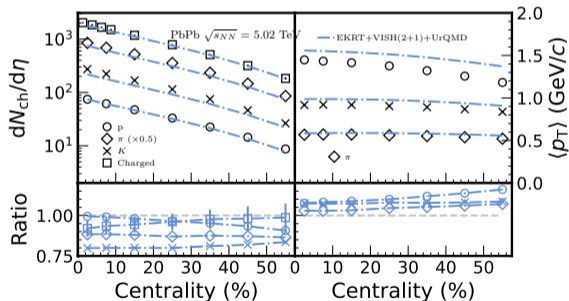
$$e(r, \tau_s(r)) = \frac{dE_T}{dV} = \frac{dE_T}{d^2r} \frac{1}{\tau_s(r) \Delta y} = \frac{K_{sat}}{\pi} p_{sat}^4$$
- Variable β describes how are soft particles from 3 \rightarrow 2 processes included
- K_{sat} gives momentum range when saturation is considered to take place



Courtesy of Henry Hirvonen

EKRT PbPb 5.02 TeV RESULTS

- Particle distributions relatively well described
- $\langle p_T \rangle$ a bit overestimated
- Large overestimation in flow harmonics



OUTLOOK

Experiments

- RHIC data (AuAu collisions) - Energy and system size dependence (ongoing)
- LHC pPb and pp data - System size dependence
- Use new observables in BA
 - Higher order ($n > 5$) Symmetric cumulants (Anna)
 - Improved Symmetric Plane Correlation (SPC) : independent from flow magnitude correlations (Maxim)
 - Asymmetric Cumulants (AC) (Cindy)
 - $\rho(v_n, p_T) = \langle \delta v_n^2 \delta [p_T] \rangle / \sqrt{\langle (\delta v_n^2)^2 \rangle \langle (\delta [p_T])^2 \rangle}$

Theory

- Improving the initial conditions with
 - EKRT (ongoing)
 - IP+Glasma
- Testing hydro limit with small systems?
- Study the substructure

Thank you for your attention!

Acknowledgments:

- CSC for providing the ~24 million CPU hours
- Harri Niemi, Kari Eskola, Jonah E. Bernhard, J. Scott Moreland and Steffen A. Bass for their useful comments

REQUIREMENT OF $\langle p_T \rangle$ EVENT BY EVENT

Full equation

$$\rho(v_2^2, [p_T]) = \frac{\langle (v_2^2 - \langle v_2^2 \rangle) ([p_T] - \langle [p_T] \rangle) \rangle}{\sqrt{\langle (v_2^2 - \langle v_2^2 \rangle)^2 \rangle \langle ([p_T] - \langle [p_T] \rangle)^2 \rangle}}$$

- With one set of hydro runs?
 - Store all particles from hydro
 - Uses too much space
- Use mean values from previous runs!

SPC SENSITIVITY COMPARISON

- Previous biased scalar product method

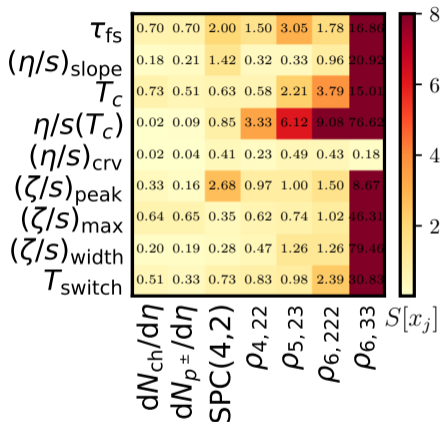
$$\langle \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle_{SP} = \frac{\langle v_{n_1}^{a_1} \dots v_{n_k}^{a_k} \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle}{\sqrt{\langle v_{n_1}^{2a_1} \rangle \dots \langle v_{n_k}^{2a_k} \rangle}}$$

- New Gaussian distribution method

$$\begin{aligned} \langle \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle_{GE} &= \int d\Theta N_{\Theta}(\Theta) \cos \Theta \\ &\approx \sqrt{\frac{\pi}{4}} \frac{\langle v_{n_1}^{a_1} \dots v_{n_k}^{a_k} \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle}{\sqrt{\langle v_{n_1}^{2a_1} \rangle \dots \langle v_{n_k}^{2a_k} \rangle}} \end{aligned}$$

- Way to probe the non-linear response, e.g.

$$v_2^2 v_4 e^{i4(\Psi_4 - \Psi_2)} = \omega_2 \omega_4 c_2^2 c_4 e^{i4(\Phi_4 - \Phi_2)} + \omega_{422} \omega_2^2 c_2^2$$



- The new SPC shows higher dependence on the specific bulk viscosity rather than the specific shear viscosity.